



Higgs Pheno Implications of the Aligned NMSSM

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In Collaboration with:

M. Carena, H. Haber, I. Low & C. Wagner arXiv:1510.09137

B. Bhattacharya, M. Carena & C. Wagner, In preparation

S. Baum, K. Freese & B. Shakya arXiv:1703.XXXXX



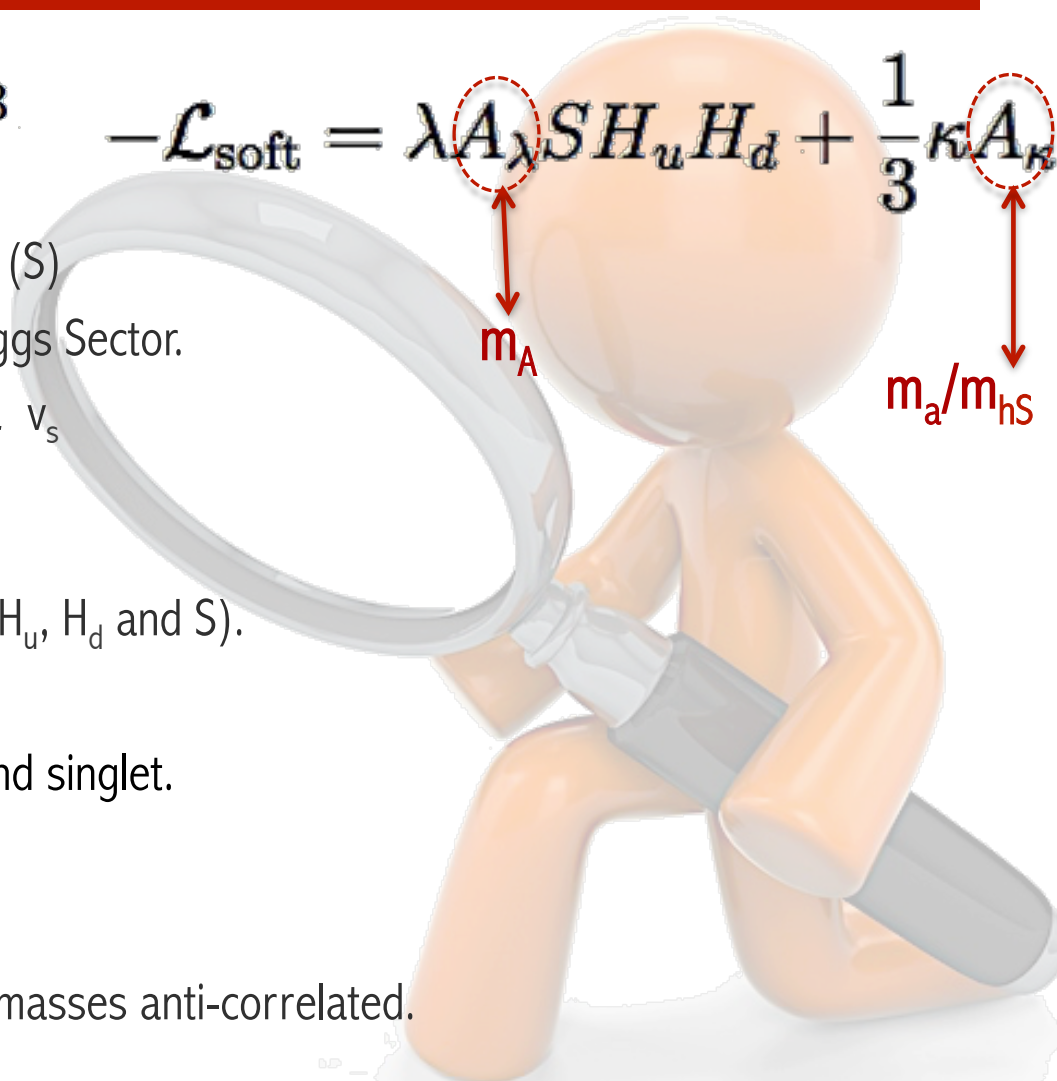
The Higgs Lamp Post:

$m_h \sim 125 \text{ GeV} + \text{SM-like}$



No New States at the LHC!!!

Alignment
SUSY \rightarrow NMSSM

$$W = \lambda S H_u H_d + \frac{\kappa}{3} S^3 \quad -\mathcal{L}_{\text{soft}} = \lambda A_\lambda S H_u H_d + \frac{1}{3} \kappa A_\kappa S^3$$


- 2 Doublets (H_u, H_d) + Singlet (S)
- Singlet couples only to the Higgs Sector.
- Singlet acquires vev: $\mu = \lambda v_s$
- **3 CP-Even Higgses:**
 - Mixing between all three (H_u, H_d and S).
- **2 CP-Odd Higgses:**
 - Mixtures of “MSSM” m_A and singlet.
- **Charged Higgs**
- Singlet-like CP-even and odd masses anti-correlated.
- **Singlino** mass: $2 \kappa \mu / \lambda$

NMSSM Higgs Sector

$$\tan \beta = \frac{v_1}{v_2}$$

$$H_{SM} = \sin \beta \Phi_1 + \cos \beta \Phi_2$$

$$H_{NSM} = -\cos \beta \Phi_1 + \sin \beta \Phi_2$$

$$\langle H_{SM} \rangle = v$$

$$\langle H_{NSM} \rangle = 0$$

- H_{SM} has completely SM-like couplings.
- Mixing between H_{SM} and H_{NSM} ,
“ $\cos(\beta - \alpha)$ ”, gives non-SM behavior of
observed h_{125} .

If

$$\cos(\beta - \alpha) = 0$$

SM-like HIGGS!!

ALIGNMENT

Higgs Basis

- Interaction basis: (H_u , H_d , S)
 - H_u : Couples only to up-type fermions
 - H_d : Couples only to down-type fermions
 - S : Only couples to Higgses

$$\begin{aligned} \langle H_u \rangle &= v_u \\ \langle H_d \rangle &= v_d \\ t_\beta &= v_u/v_d \\ \langle S \rangle &= \mu/\lambda \end{aligned}$$

CP-Even Higgs Bases

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- “Extended” Higgs basis: (H_{NSM}, H_{SM}, S)
 - H_{NSM} : (down, up, V) = ($y_d t_\beta, y_u/t_\beta, 0$)
 - H_{SM} : (down, up, V) = (y_d, y_u, g_{hVV})

$$\begin{aligned} \langle H_{NSM} \rangle &= 0 \\ \langle H_{SM} \rangle &= v \end{aligned}$$

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$$\begin{aligned} \langle H_{NSM} \rangle &= 0 \\ \langle H_{SM} \rangle &= v \end{aligned}$$

- Mass basis: (H^3, H^2, H^1)
 - $H^i = \kappa_{NSM}^i H_{NSM} + \kappa_{SM}^i H_{SM} + \kappa_S^i S$

CP-Even Higgs Bases

- Interaction basis: (H_u, H_d, S)
 - H_u : Couples only to up-type fermions
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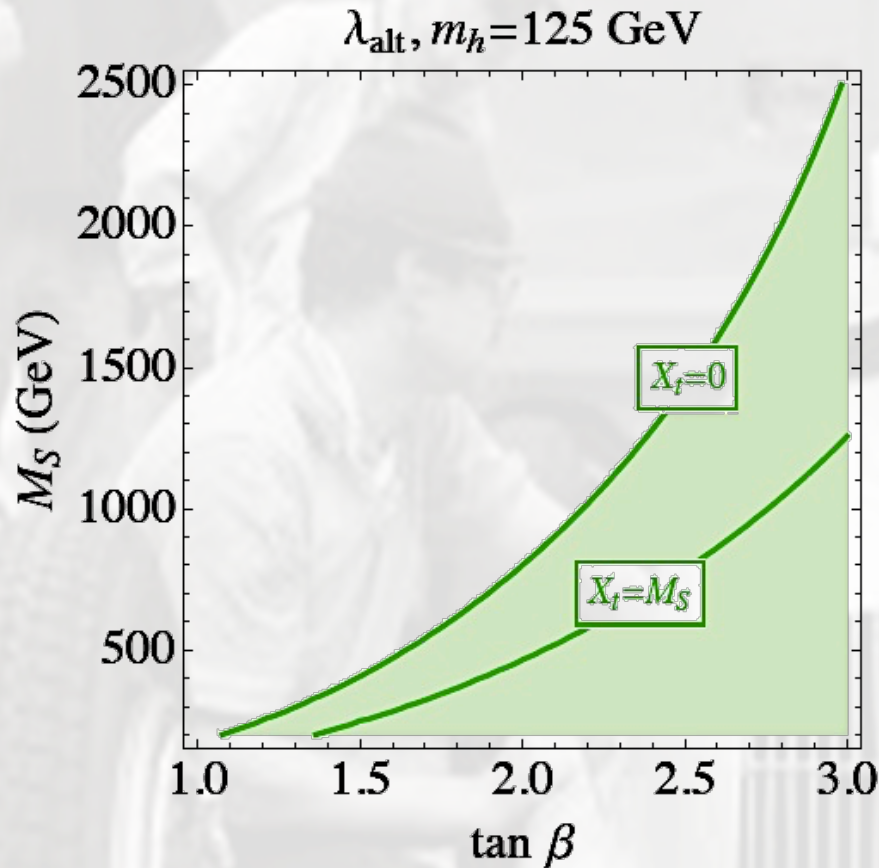
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 - H_{SM} : (down, up, V) = (y_d, y_u, g_{hW})

$$\begin{aligned} \langle H_{NSM} \rangle &= 0 \\ \langle H_{SM} \rangle &= v \end{aligned}$$

- Mass basis: $(H^3, H^2, H^1) \rightarrow (H, h_{125}, h_S)$
 - $H^i = \kappa_{NSM}^i H_{NSM} + \kappa_{SM}^i H_{SM} + \kappa_S^i S$

$$\begin{aligned} \text{Alignment:} \\ \kappa_{NSM}^{h_{125}} &= 0 \\ \kappa_S^{h_{125}} &= 0 \end{aligned}$$

CP-Even Higgs Bases



Alignment (No-Mixing):

$$m_h^2 \simeq \lambda^2 \frac{v^2}{2} \sin^2 2\beta + M_Z^2 \cos^2 2\beta + \Delta_{\tilde{t}}$$

$$\Delta_{\tilde{t}} = -\cos 2\beta (m_h^2 - M_Z^2)$$

Well Known

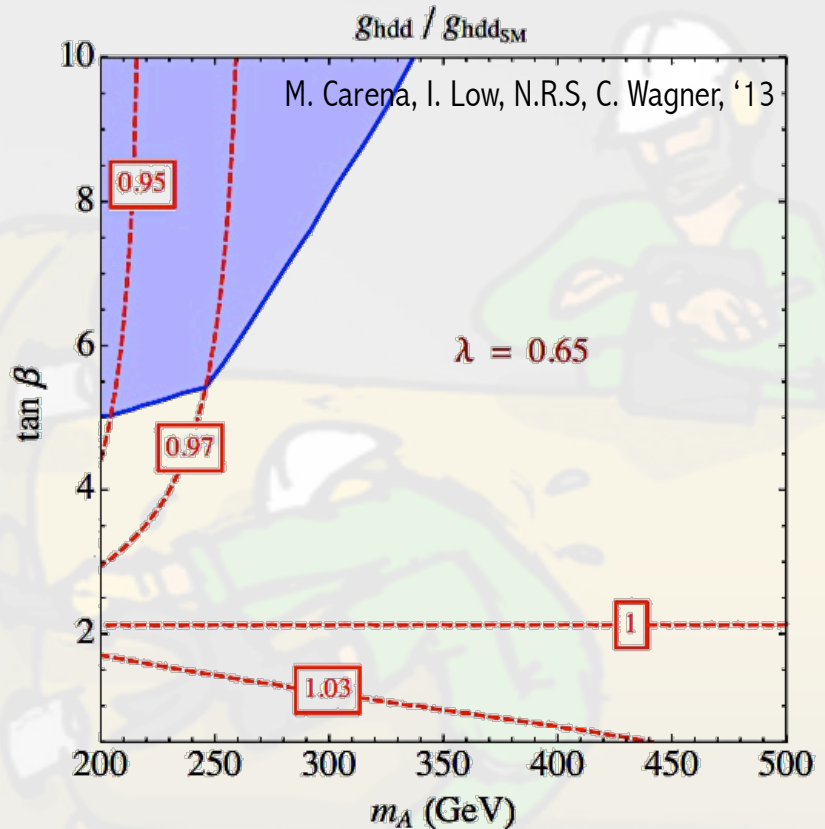
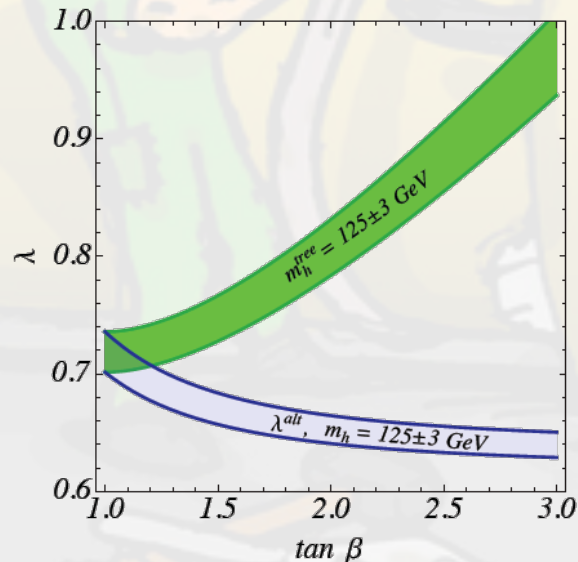
- 125 GeV Higgs
 - Tree-level contribution to Higgs mass from λ .
 - $\lambda \sim 0.65-0.7$
- Low $\tan \beta$
- Light Stops

125 GeV Higgs Naturally!

- Perturbative up to GUT scale.
 - $\lambda_{\max} \sim 0.7, \kappa_{\max} \sim \lambda/2$

Not so well known:

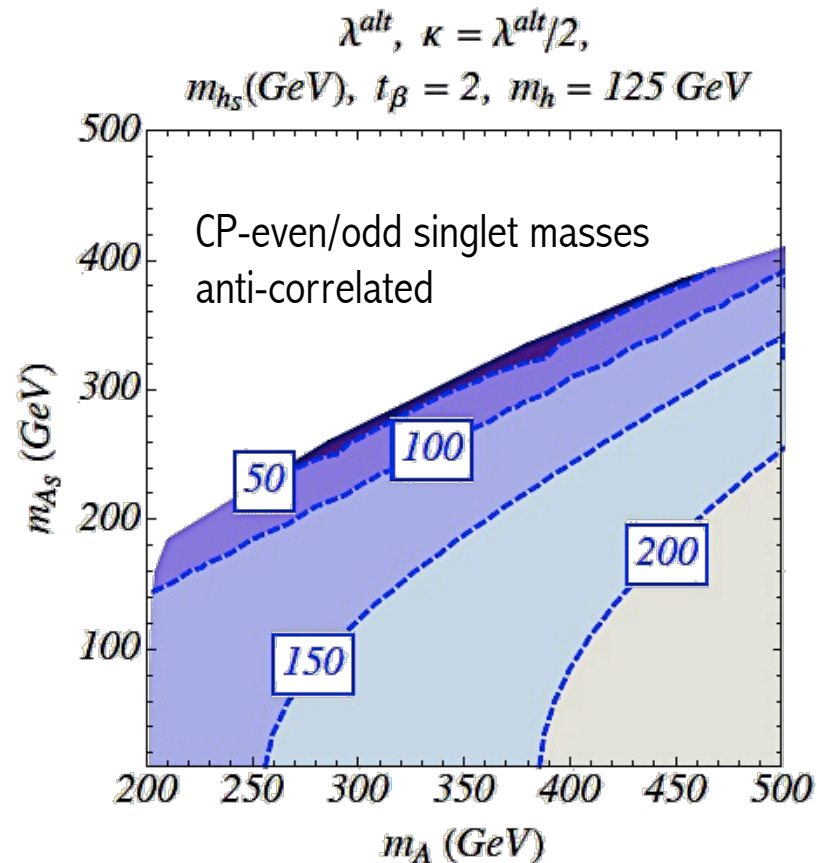
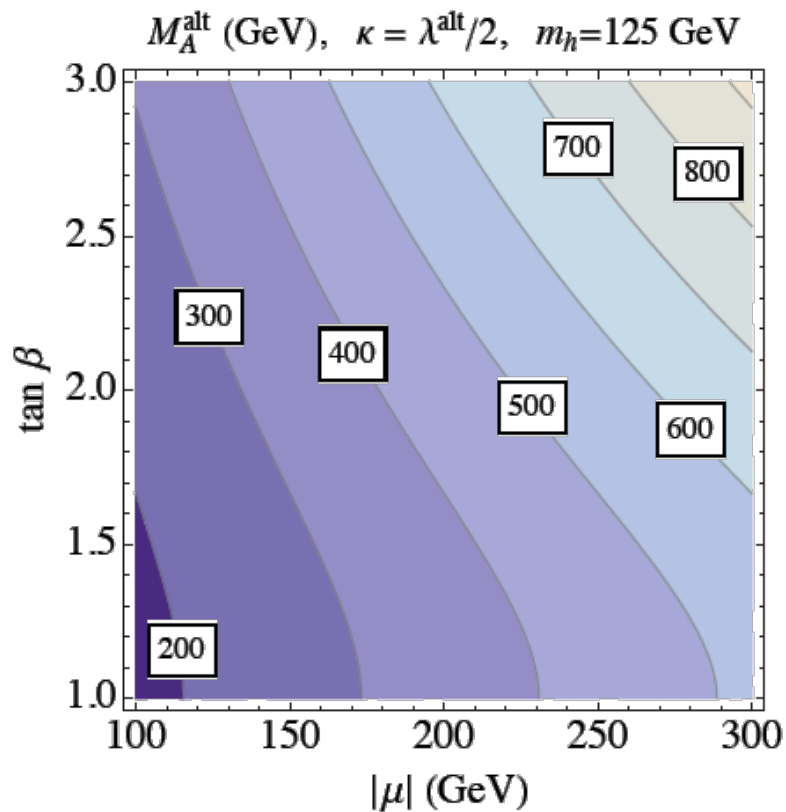
- Leads to excellent Alignment (very little mixing with Heavy Higgs) in the m_A - $\tan\beta$ plane.



$$\lambda_{\text{alt}}^2 = \frac{m_h^2 - M_Z^2 c_{2\beta}}{v^2 s_\beta^2}$$

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SM-Like Higgs Naturally!



$$1 - \frac{m_A^2}{4\mu^2} s_{2\beta}^2 - \frac{\kappa}{2\lambda} s_{2\beta} = 0$$

$h_{125} = H_{\text{SM}}$

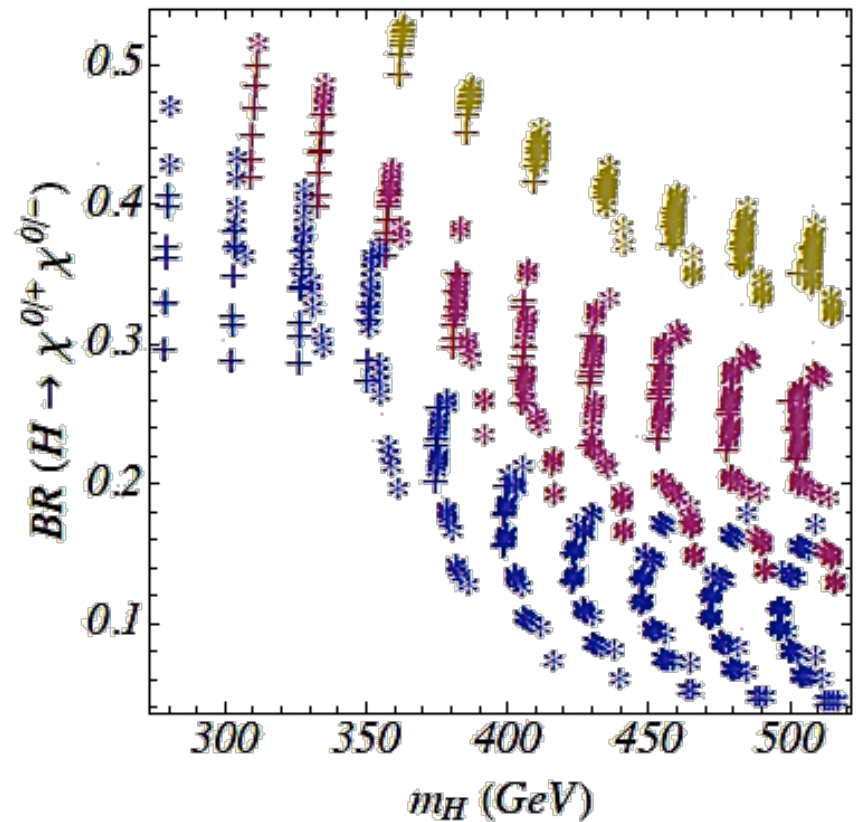
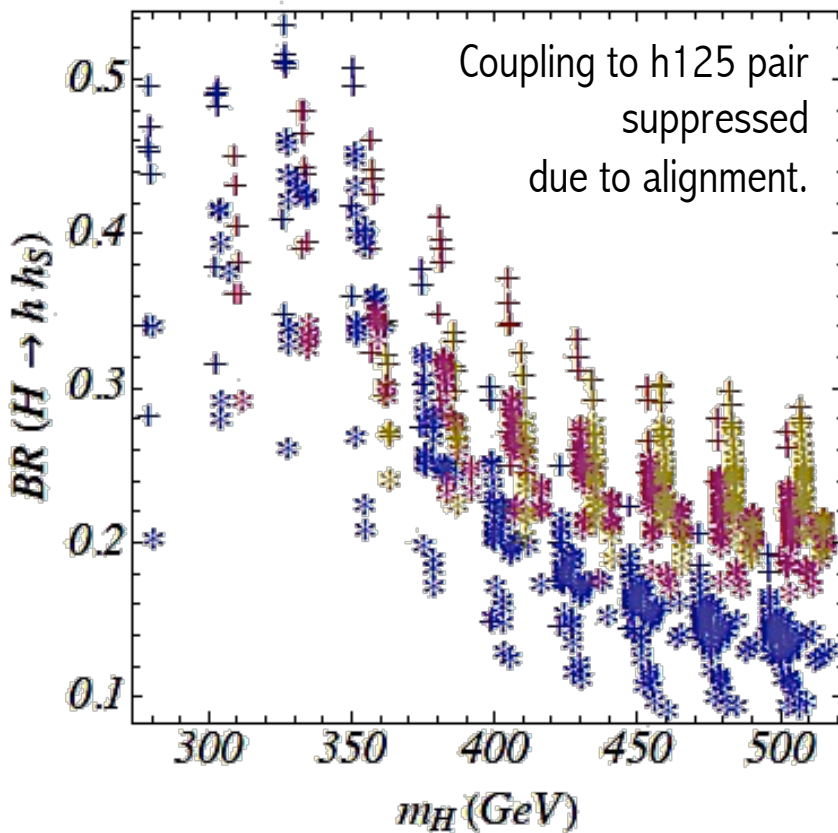
LIGHT SPECTRUM

Singlino: $2 \kappa \mu / \lambda \sim < \mu$

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Singlet Alignment

- Apart from tt , significant decays into H_2+H_1 and neutralino/charginos



. $\tan \beta = 2$. $\tan \beta = 2.5$. $\tan \beta = 3$

H_3 BR (A_2 BR similar)

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- Light Neutralinos:
 - Relic Density can be obtained via resonances or NMSSM “well-tempered” Neutralinos
 - Generically large direct detection cross-section via h125
 - Need “Blind-spot” cancellations. E.g:

$$\sigma_{SI} \propto \left\{ \left(\frac{2}{t_\beta} - \frac{m_\chi}{\mu} \right) \frac{2t_\beta}{m_h^2} + \frac{t_\beta}{m_H^2} + \frac{1}{m_{h_s}^2} \left(2S_{h,s} + \frac{\lambda v}{\mu} \right) \left[\frac{\lambda v}{\mu^2} m_\chi + S_{h,s} \left(\frac{2}{t_\beta} - \frac{m_\chi}{\mu} \right) + \frac{\kappa \mu}{\lambda^2 v} \right] \right\}^2.$$

C. Cheung, M. Papucci, D. Sanford, NRS, K. Zurek, '14

Conditions: $\mu > 0$

$\mu \sim m_\chi t_\beta / 2$ (h125 coupling reduced)

$h_s < h_{125}$ (need alignment for μ and m_A)

Cosmology?

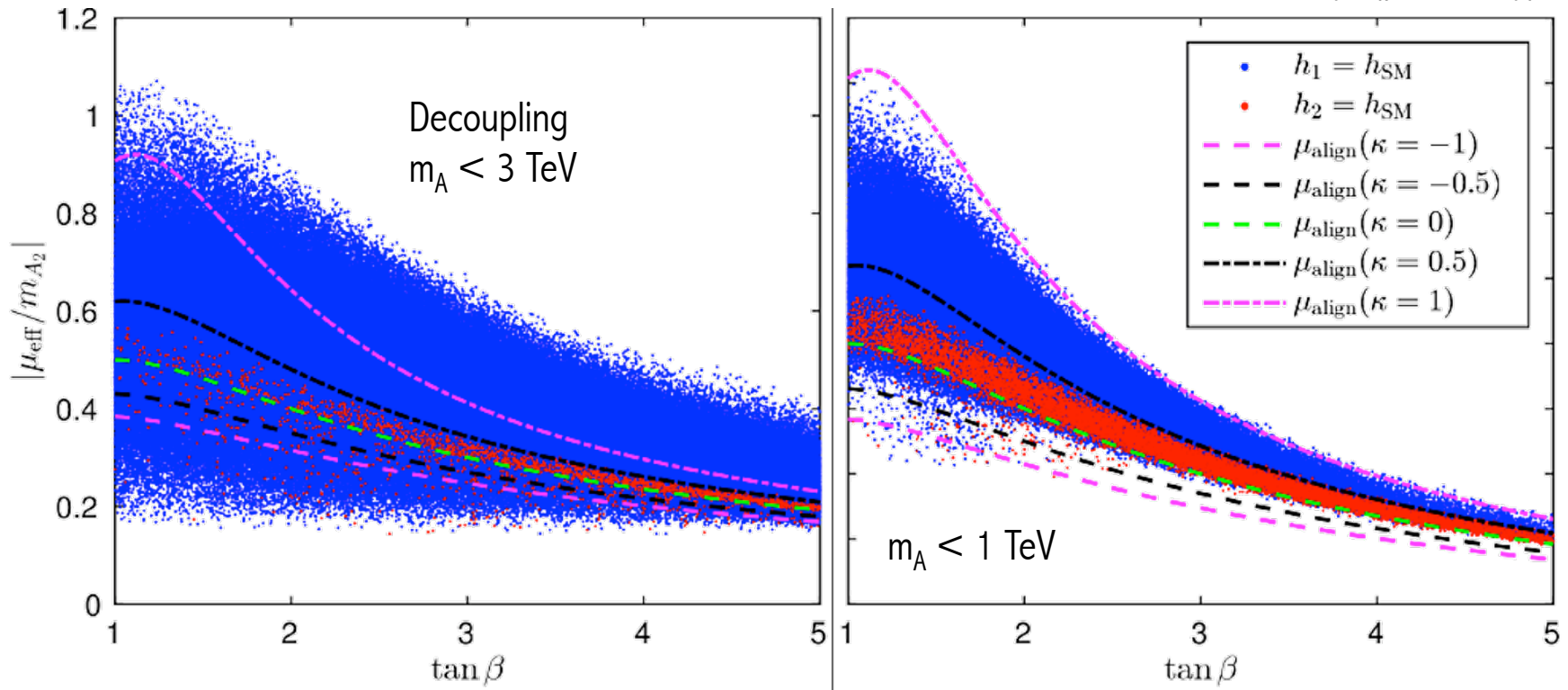
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- LSP could be $\sim O(10 \text{ GeV})$ or $\sim O(100 \text{ GeV})$
- $h_S < h_{125}$
 - Mixing can change total width of h_{125} – significant mixing still allowed
 - If $h_S < h_{125}/2$, $h_{125} \rightarrow h_S h_S$ – Exotic decays of h_{125}
 - h_S can have significant production cross-section $\sim \text{pb}$
 - if $h_S > 2\chi$, $\text{BR}(h_S \rightarrow \chi \chi) \sim 100\%$
 - Otherwise decays with SM-like BR
- For low $\tan\beta$, m_A can be few hundred GeV
 - Significant decays of both H and A into Higgses and Neutralinos
 - $H \rightarrow h_S h_S / h_S h_{125}$ etc.
 - Note: $H \rightarrow h_{125} h_{125} / WW/ZZ$ and $A \rightarrow Z h_{125}$ all suppressed by alignment/decoupling

Cosmology Benchmark

- LSP could be stable on detector time-scales but unstable on universe time-scale
- Consider collider phenomenology independent of cosmological constraints
- First question: Alignment with or without decoupling
 - Most interesting LHC phenomenology when all states light (no decoupling)
 - Pursue NMSSM parameter space using “alignment”
 - Pinpoint promising signatures
 - In particular will focus on **mono-Higgs**

Heavy Higgs Pheno

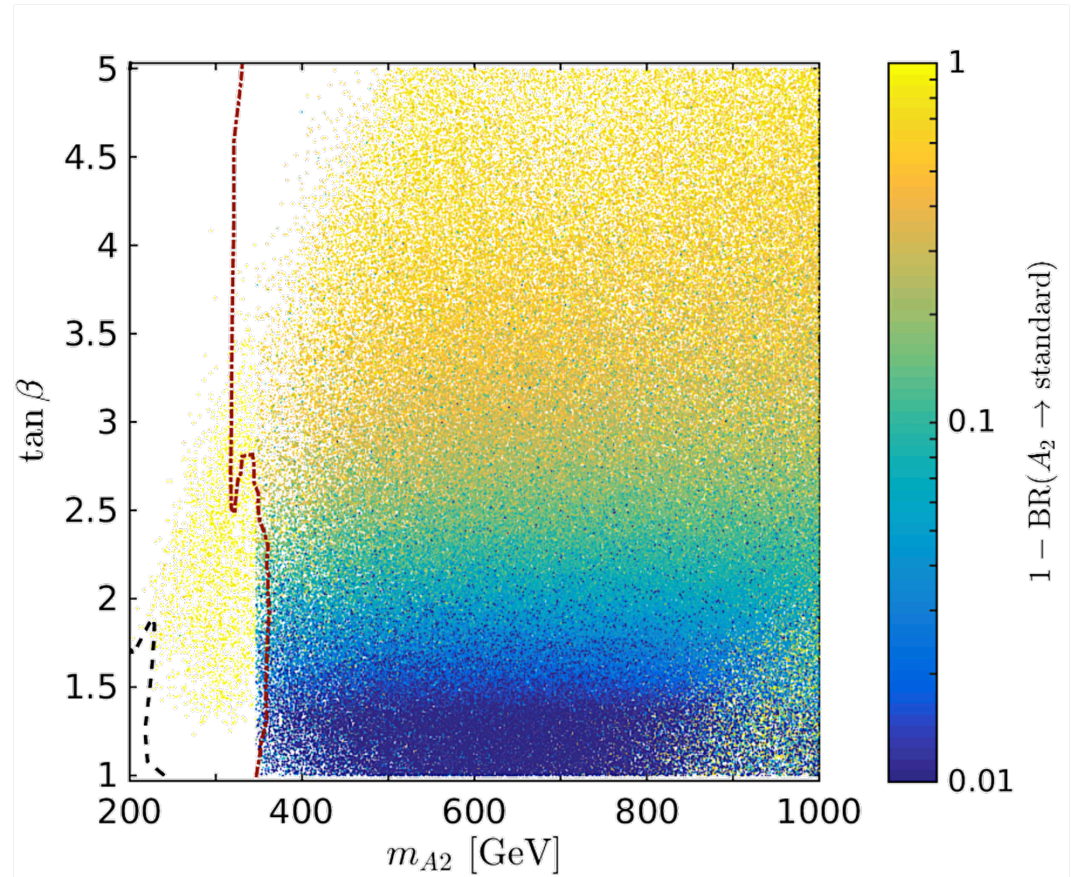
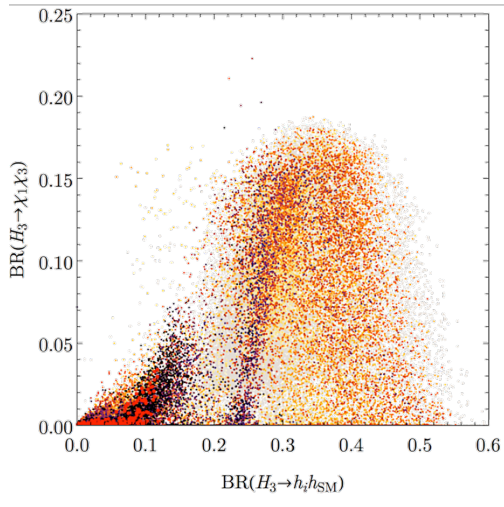
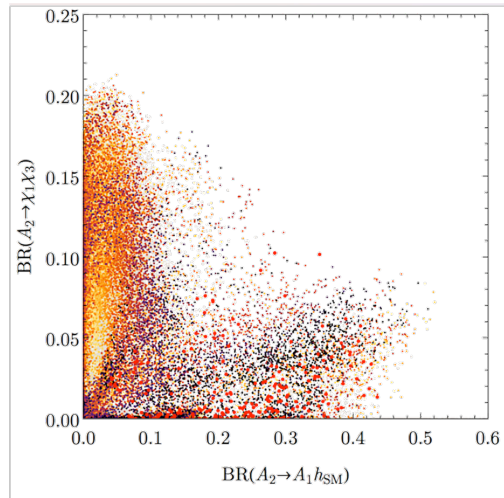


NMSSMTools Scans asking for SM-like h_{125}

Misalignment?

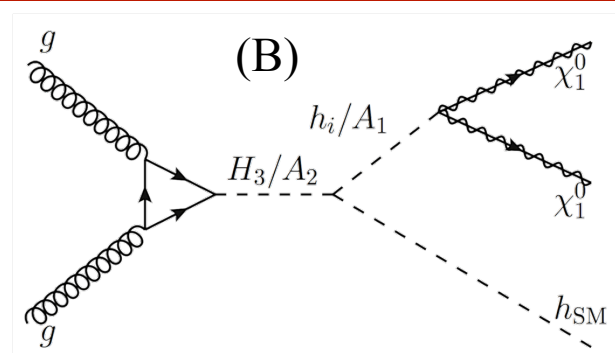
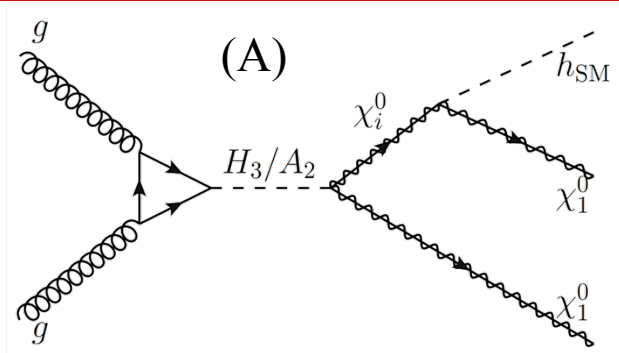
- $m_{A_2} \sim m_{H_3}$
- Completely “MSSM”-like A_2 has $x_{\text{sec}} \sim 2$ “MSSM”-like H_3
- Mixing with singlets will reduce these x_{sec}
- A_2 can mix significantly with A_1 consistent with alignment conditions
 - Still comparable x_{sec} to H_3
 - Significant BR into non-standards: $\chi_i \chi_1, A_1 h_{125}, h_S Z$
- H_3 mixes less with singlet, but enough to also have significant non-standard decays:
 - $\chi_i \chi_1, h_S h_{125}, A_1 Z$

A_2 and H_3



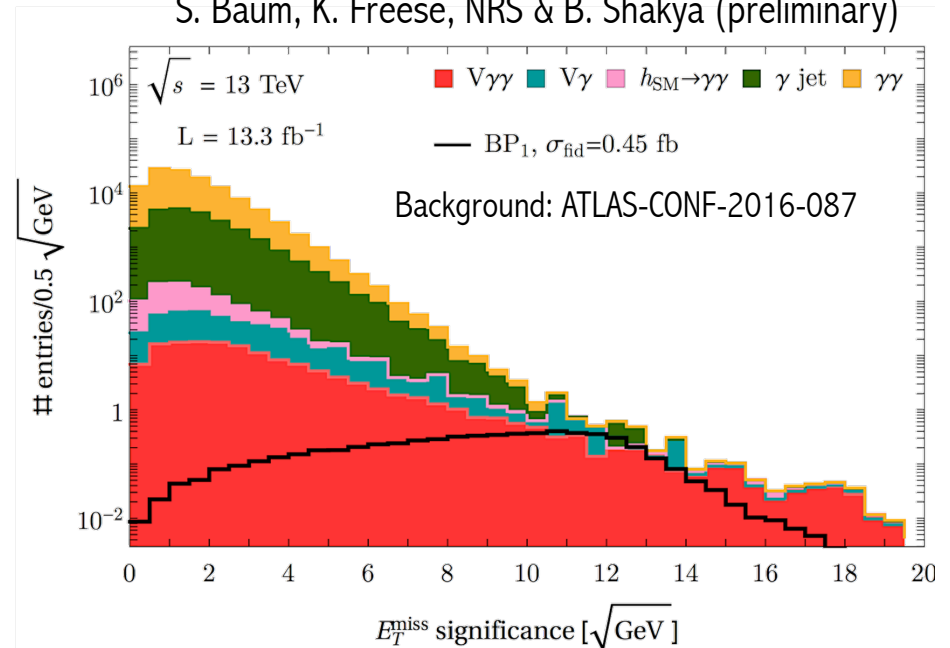
Non-Standard BR

If A_1 or $h_S > 2\chi$, $BR(\chi\chi)$ large
 Otherwise SM-like BR

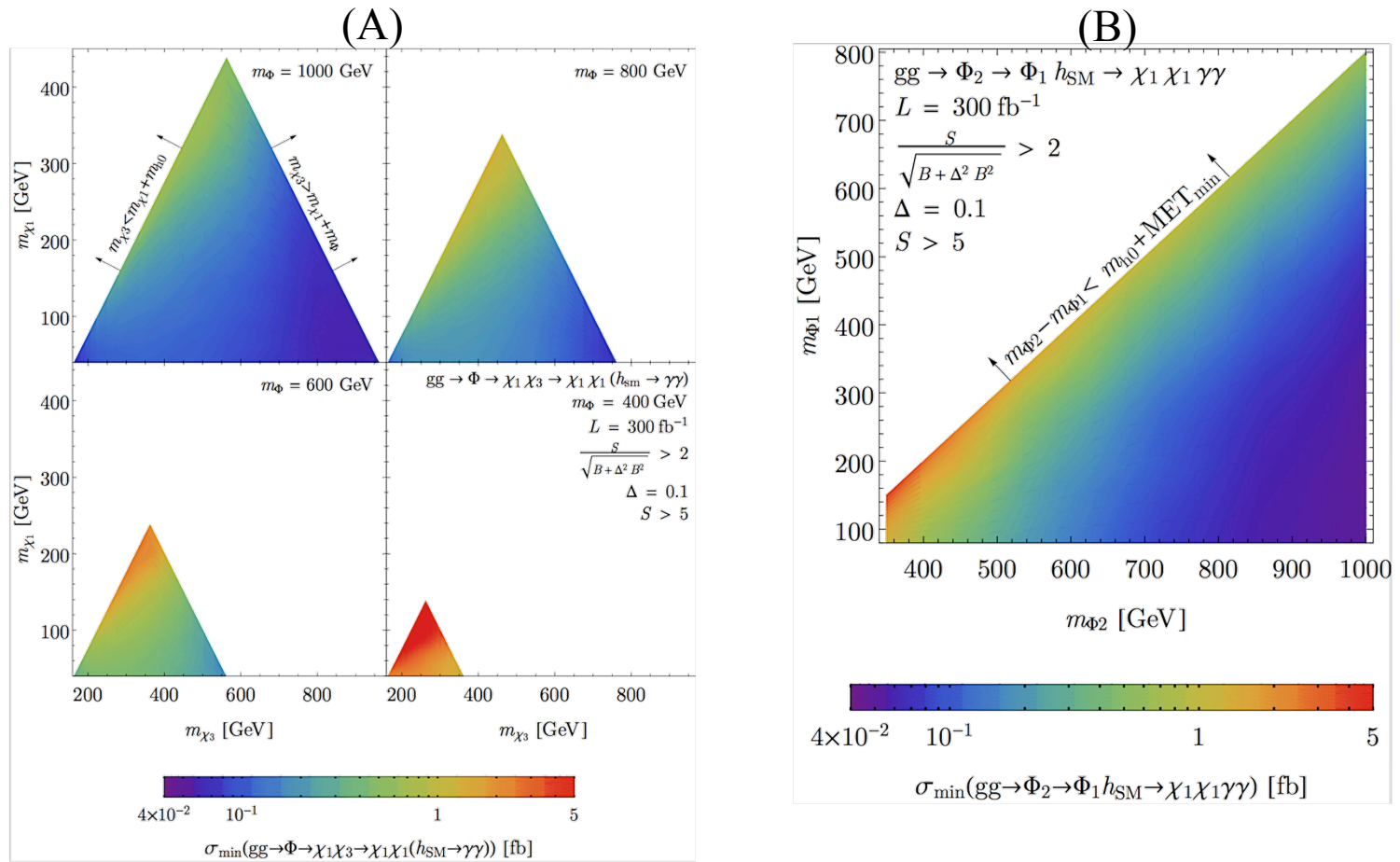


- A. Can add up contributions from both H_3 and A_2
- B. Expect stronger sensitivity –h125 back-to-back with MET

S. Baum, K. Freese, NRS & B. Shakya (preliminary)

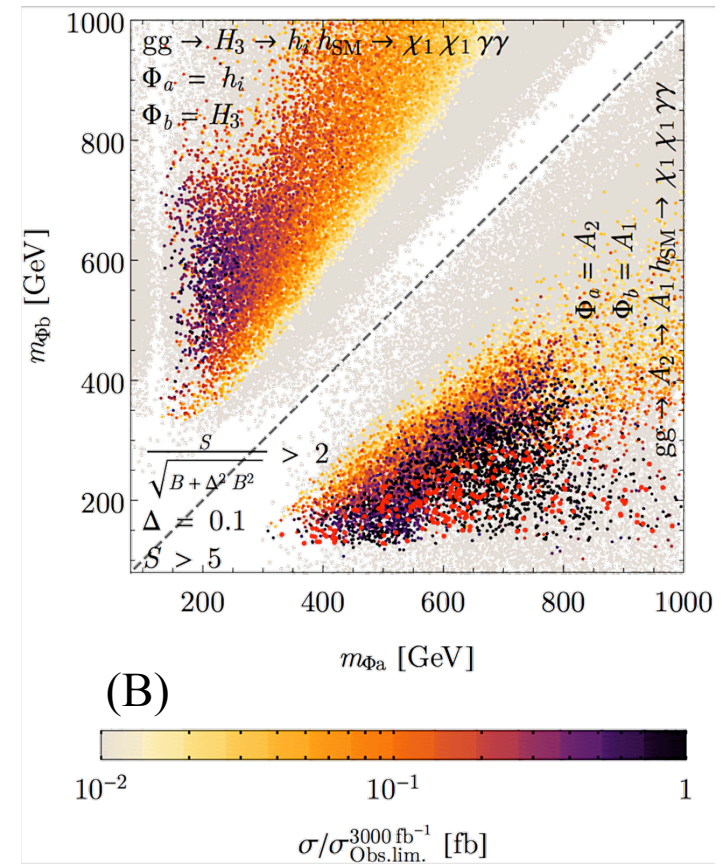
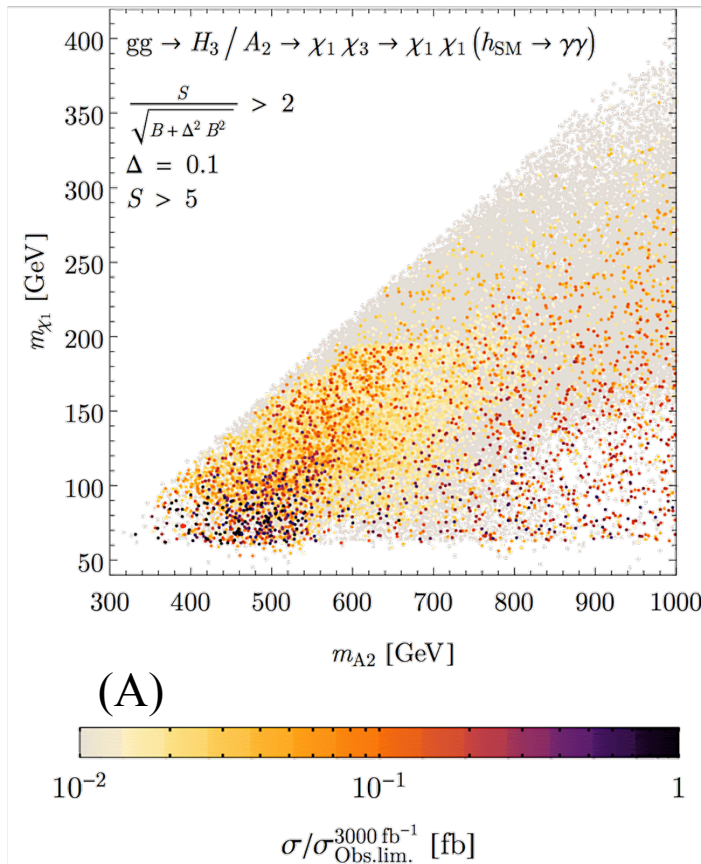


Mono-Higgs (to $\gamma\gamma$)



Depending on mass spectrum, can have sensitivity \sim sub-fb

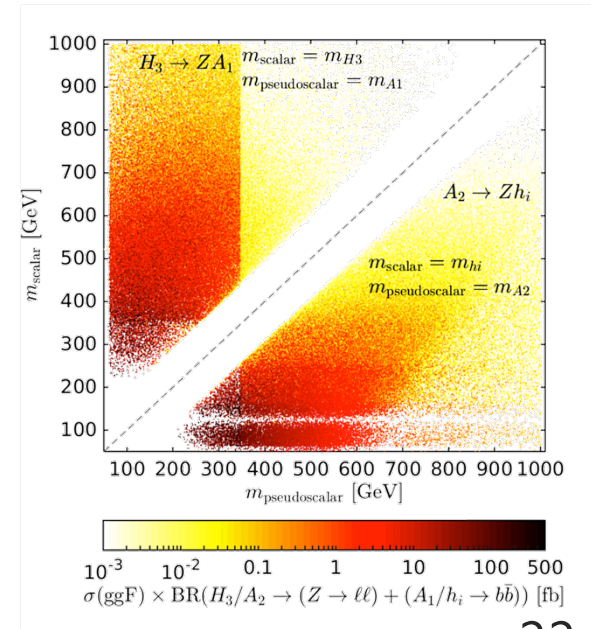
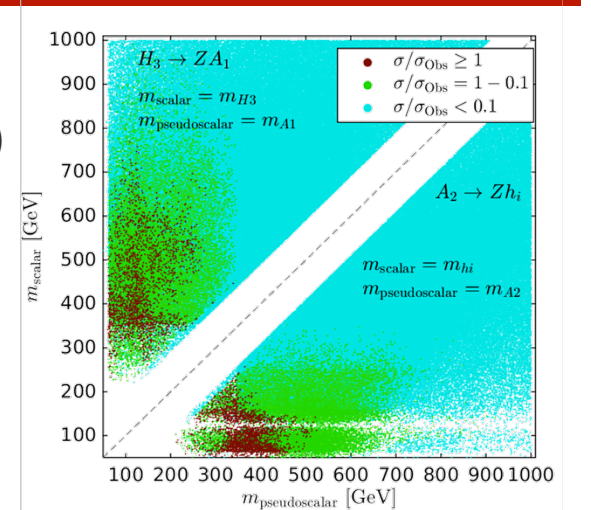
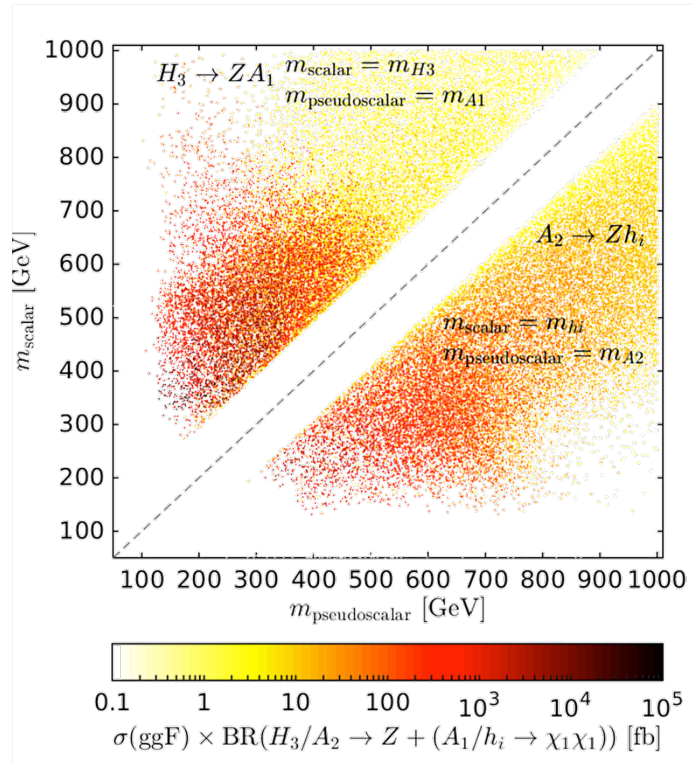
Reach?



As expected much better reach for channel (B) – even when tt is open

NMSSM

- $A/H \rightarrow Z h_s/a \rightarrow Z bb$ (CMS -- arXiv:1603.02991)
- $A/H \rightarrow Z h_s/a \rightarrow Z \chi \chi$ (mono-Z) (CMS -- arXiv:1701.02042)
- $A/H \rightarrow h_s/a$ $h_{125} \rightarrow bb \gamma \gamma, 4b?$



Other Signatures?

Conclusions and Outlook

- $m_h = 125 \text{ GeV} + \text{SM-like}$
 - Alignment: Decoupling or Prediction for parameters.
- NMSSM Higgs sector at low $\tan \beta$.
 - Perturbativity and the requirement of alignment with the singlet
 - light singlets (both CP-even and odd) and singlino + higgsinos (charged and neutral).
 - Consistency with Cosmology provides further guidance
 - singlet scalar lighter than h_{125} – not decoupled
 - Large BR of non-SM Higgs into singlet like states + neutralinos
 - Mono-Higgs reach studied in detail
 - Complimentary channels:
 - mono-Z
 - $2b 2 \gamma$, $4b$?



BACKUP SLIDES

$\tan \beta$	1 – 5	1 – 5
λ	0.5 – 1	0.5 – 2
κ	–0.5 – 0.5	–1 – 1
A_λ	–0.5 – 0.5 TeV	–1 – 1 TeV
A_κ	–0.5 – 0.5 TeV	–1 – 1 TeV
μ_{eff}	–0.5 – 0.5 TeV	–1 – 1 TeV
M_{Q_3}	1 – 10 TeV	1 – 10 TeV

TABLE I: NMSSM parameter ranges used in `NMSSMTools` scans.

- How much “non-standardness” is allowed by h125 measurements??

- $\kappa_{\text{NSM}} H_{\text{NSM}} + \kappa_{\text{SM}} H_{\text{SM}} + \kappa_S S$

- Singlet: Only couples to Higgses

- Ratios to SM:

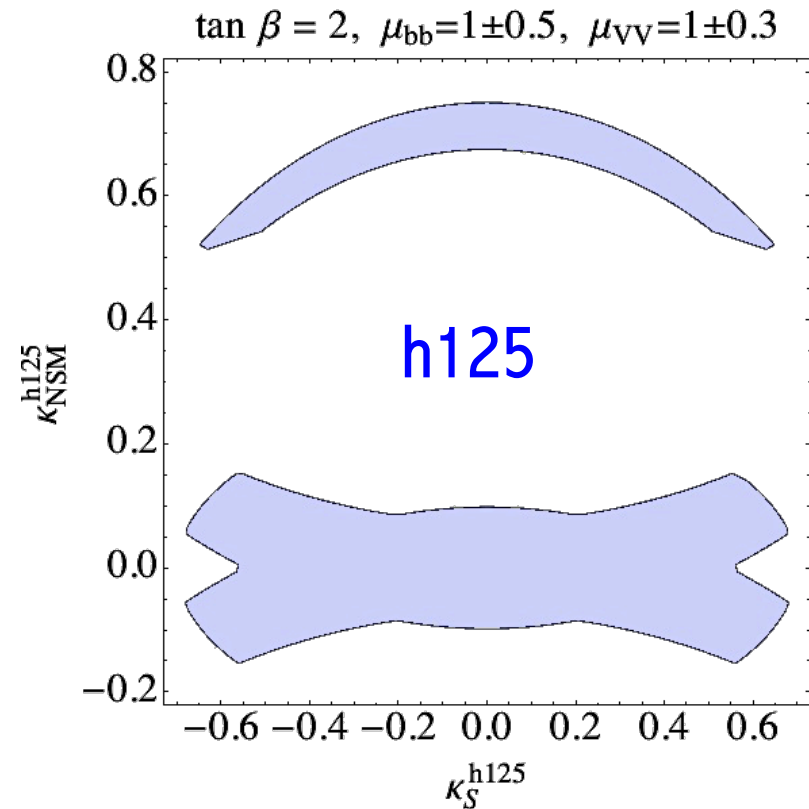
- $g_{\text{hgg}} = (\kappa_{\text{SM}} + \kappa_{\text{NSM}}/t_\beta)$

- $g_{\text{hdd}} = (\kappa_{\text{SM}} - \kappa_{\text{NSM}} t_\beta)$

- $g_{\text{hVV}} = \kappa_{\text{SM}}$

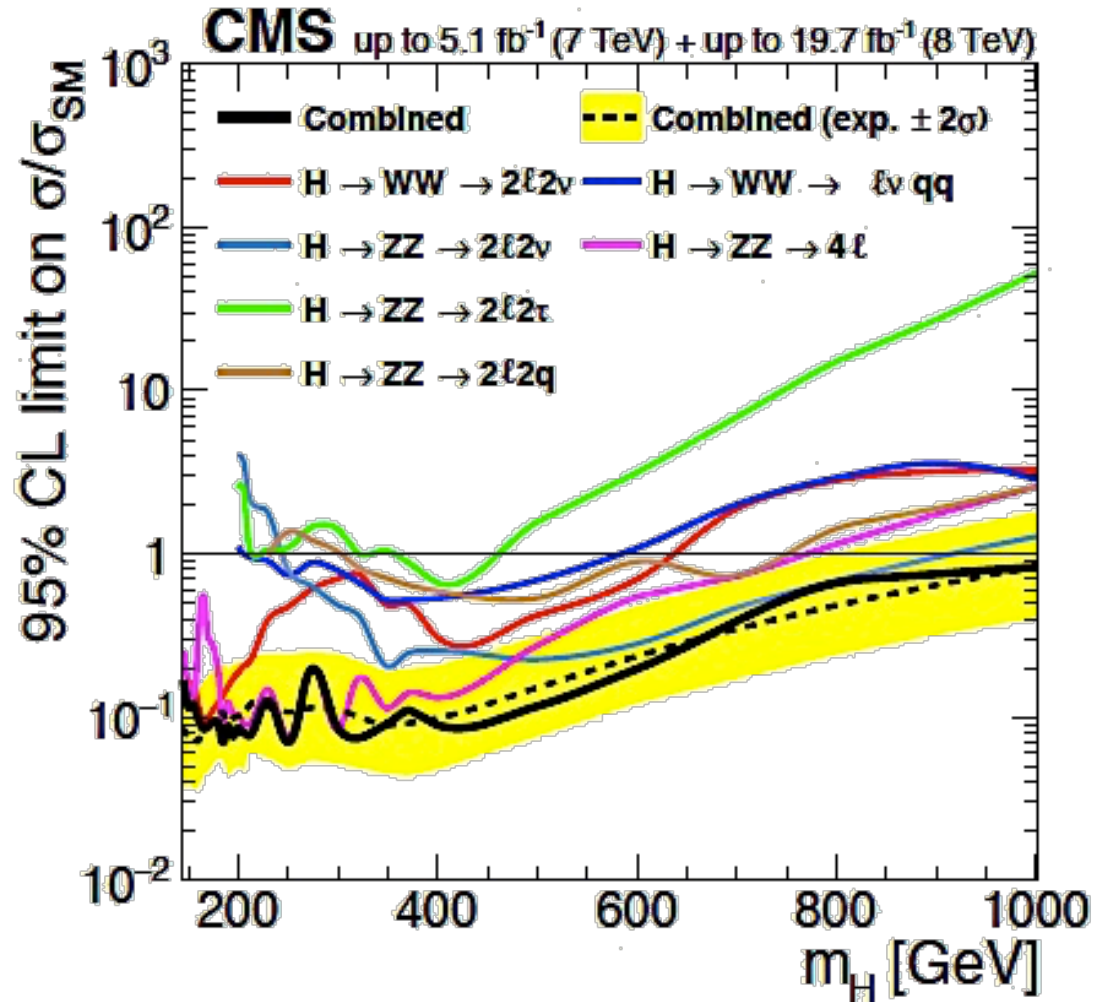
- Significant κ_S OK
- Large κ_{NSM} from sign change of g_{hdd}

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Contamination allowed in h125 ??

- CMS 1505.03831
- Strong constraints on SM-like Higgs decay to $VV \sim 12-6\%$
SM value for masses
160-500 GeV.



Heavy H to VV ?

- Strong constraints on SM-like Higgs decay to $VV \sim 12\text{-}6\%$ SM value for masses 160-500 GeV.

- Only κ_{SM}^i couples to VV

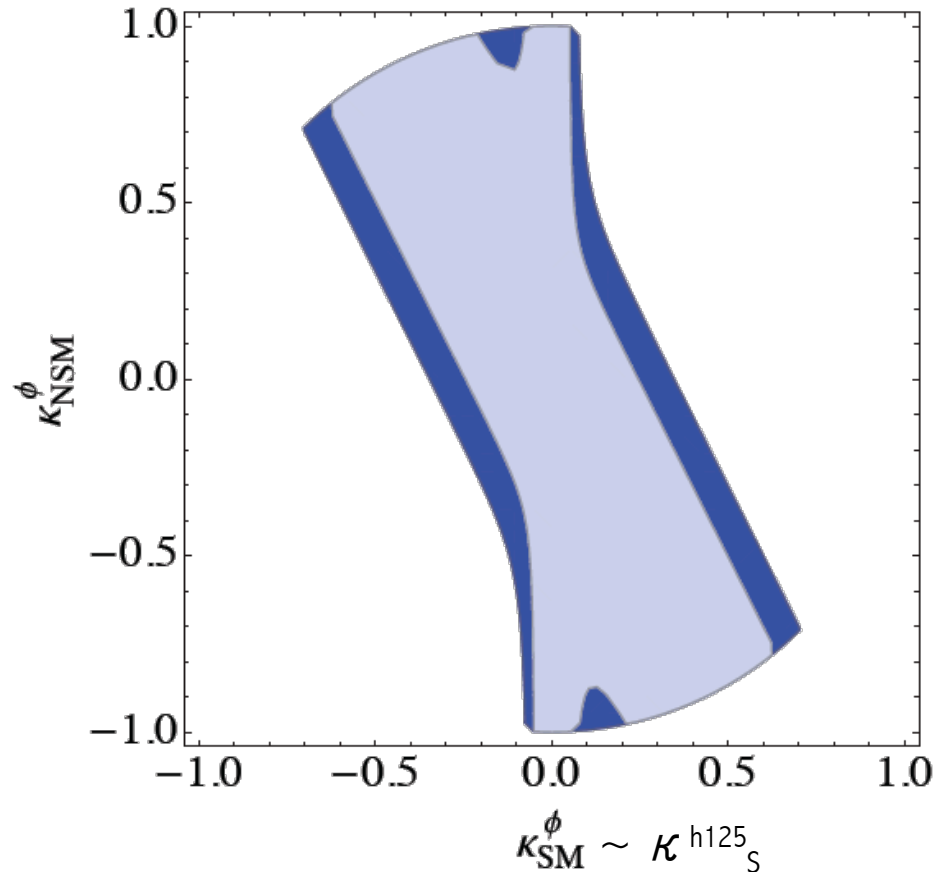
- What does this imply for SM and NSM components of extra Higgs??

- $160 \text{ GeV} < m_{h_S} < 350 \text{ GeV}$
- $BR(WW+ZZ) \sim 1$
- gF production XS impacted.

- With $\kappa_{NSM}^{h125} \sim 0$

- $\kappa_{SM}^{h_S} \sim \kappa^{h125}_S$
- κ^{h125}_S smaller than allowed by h125 measurements!

$\tan \beta = 2, \mu_{VV} < 0.12\text{--}0.06, m_\phi = 200 \text{ GeV}$



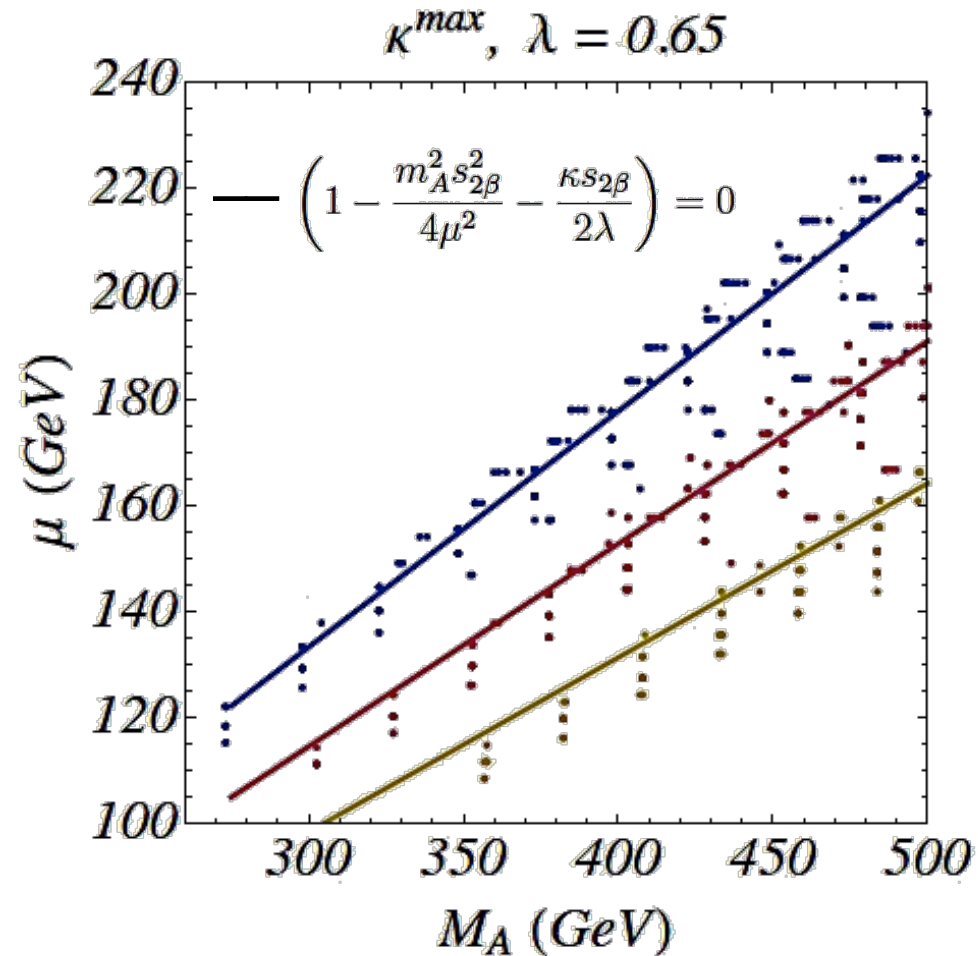
Direct Searches for heavy resonances?

- NMSSMTools + HiggsBounds/Signals
- Allowed “misalignment”

. $\tan \beta = 2$

. $\tan \beta = 2.5$

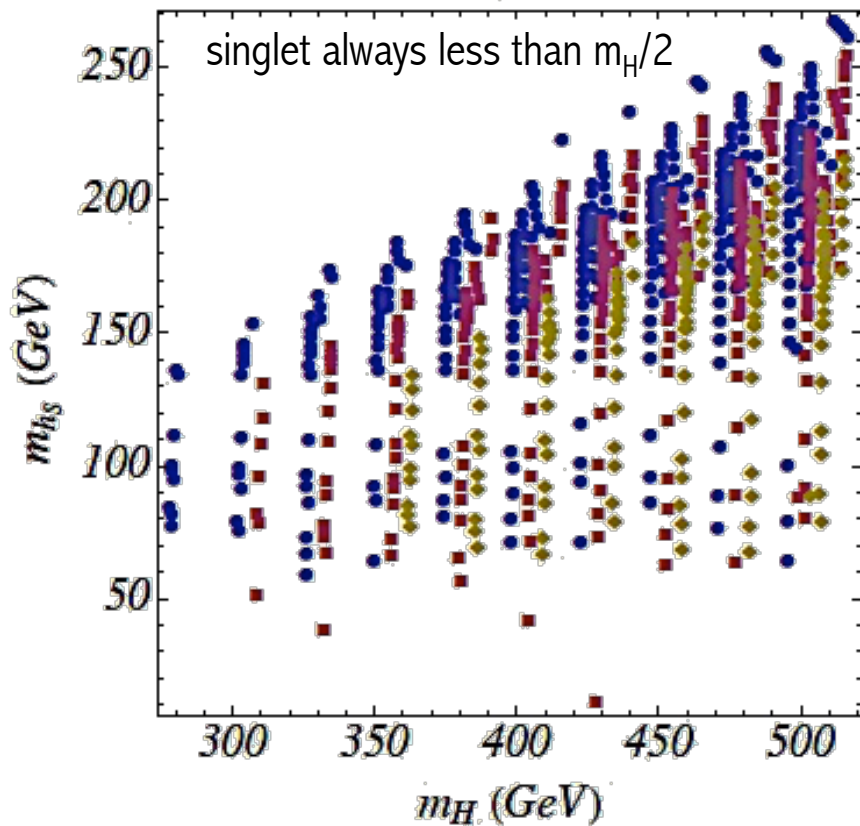
. $\tan \beta = 3$



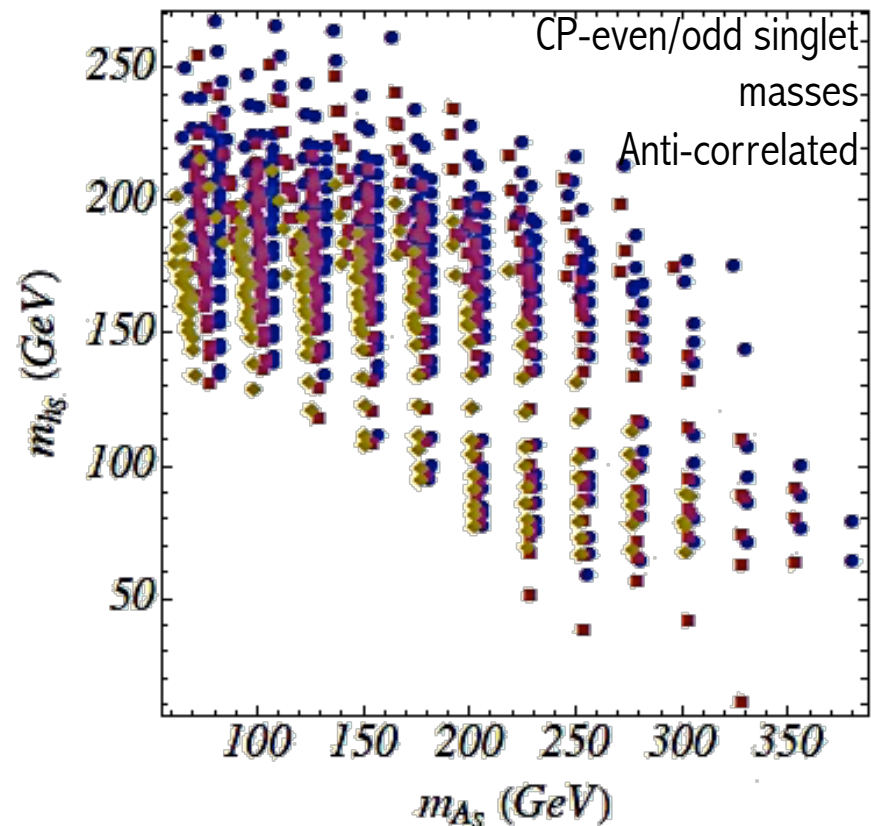
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Misalignment

$\kappa^{max}, \lambda = 0.65$



$\kappa^{max}, \lambda = 0.65$

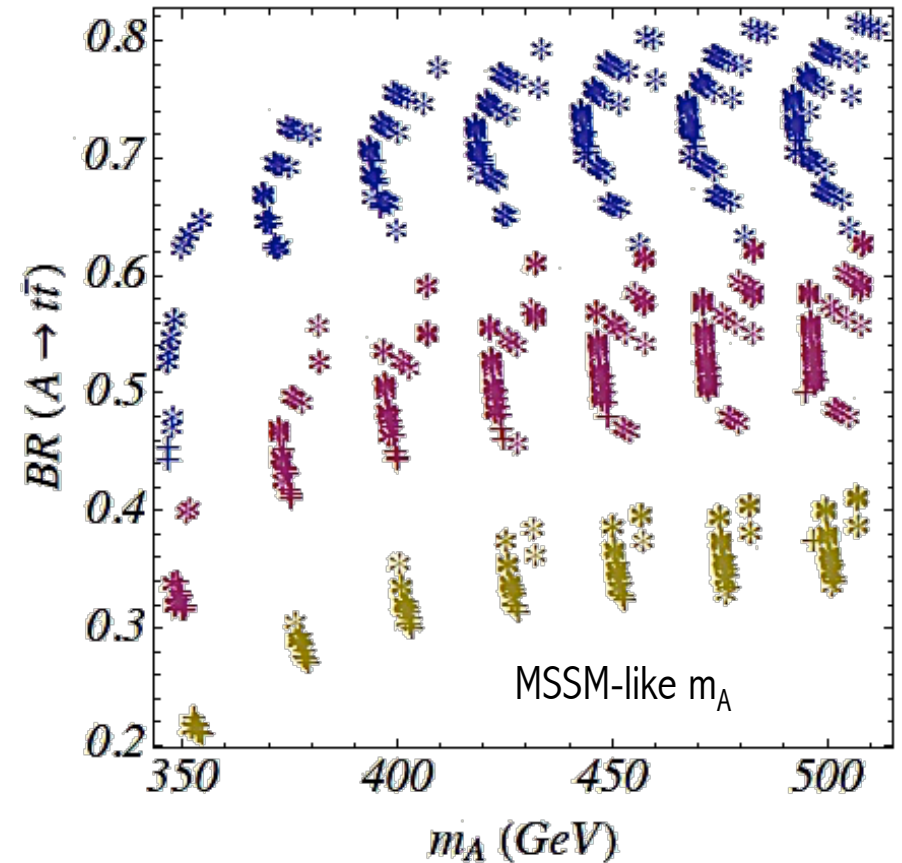
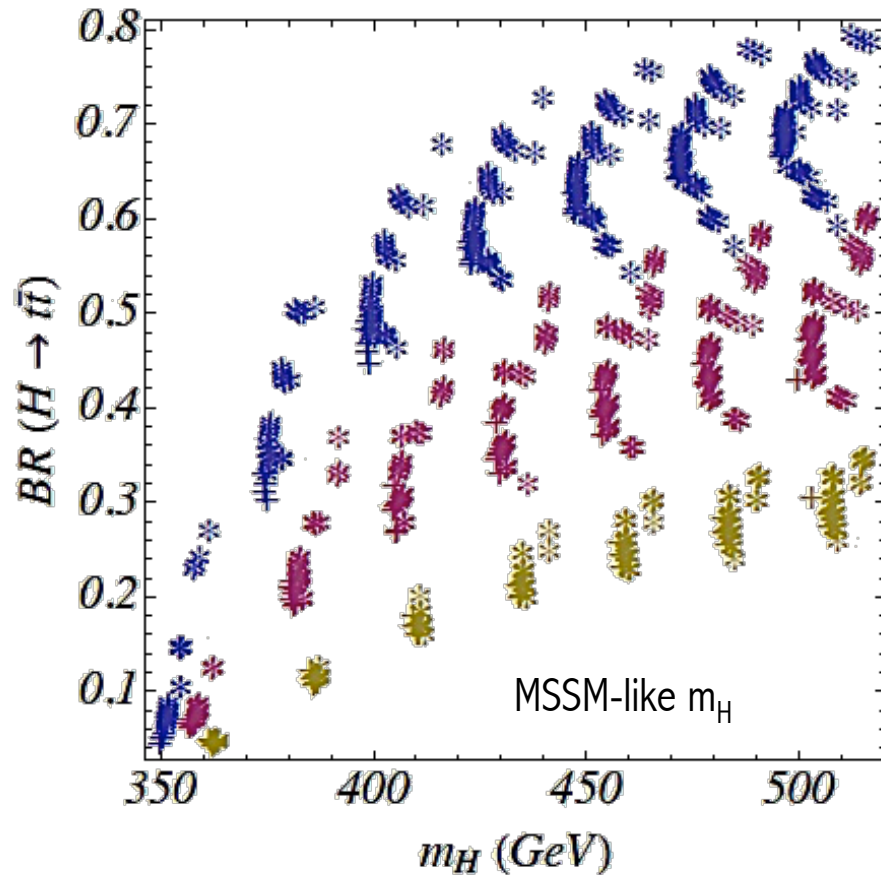


. tan β = 2 . tan β = 2.5 . tan β = 3

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Singlet Spectra

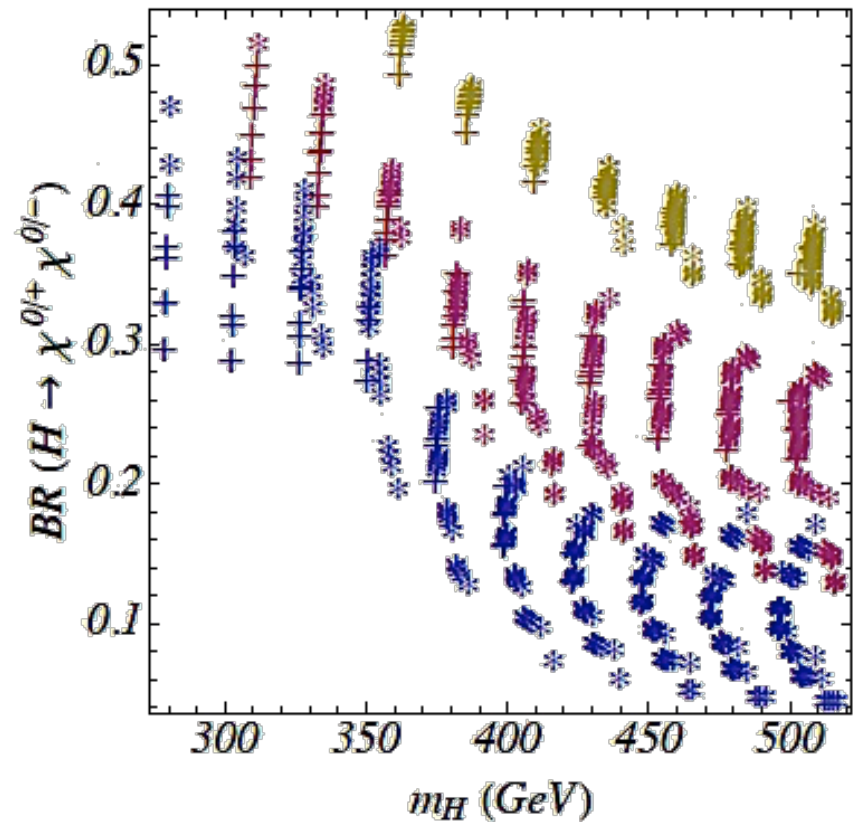
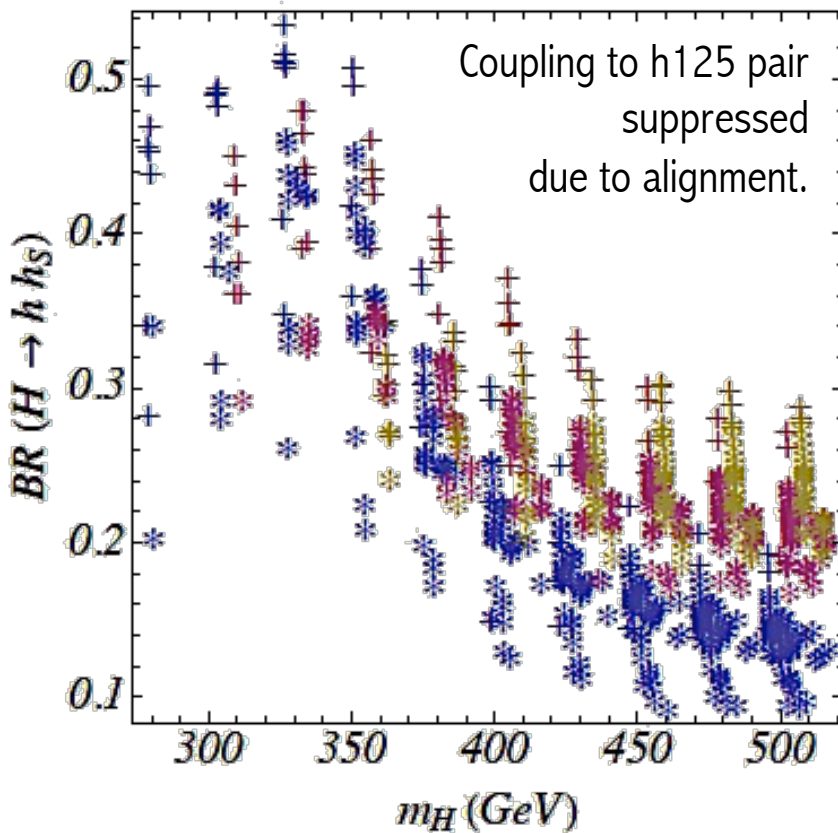
- Decay BR depends on $\tan \beta$



. $\tan \beta = 2$. $\tan \beta = 2.5$. $\tan \beta = 3$

MSSM-like A and H decays into $t\bar{t}$

- Apart from tt , significant decays into H_2+H_1 and neutralino/charginos



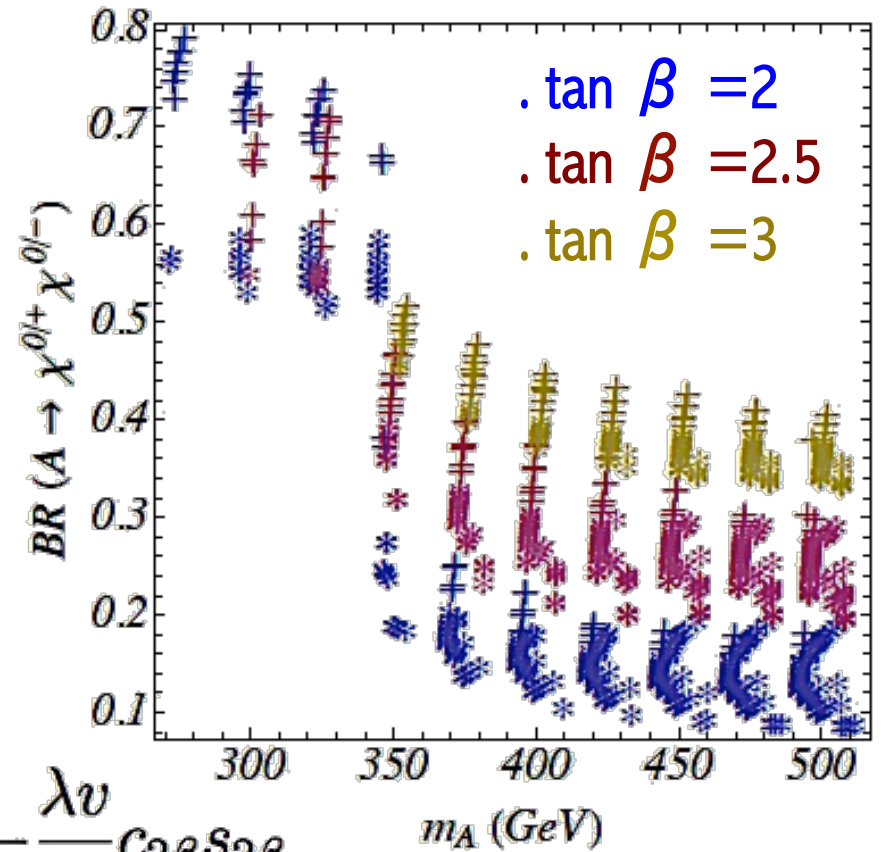
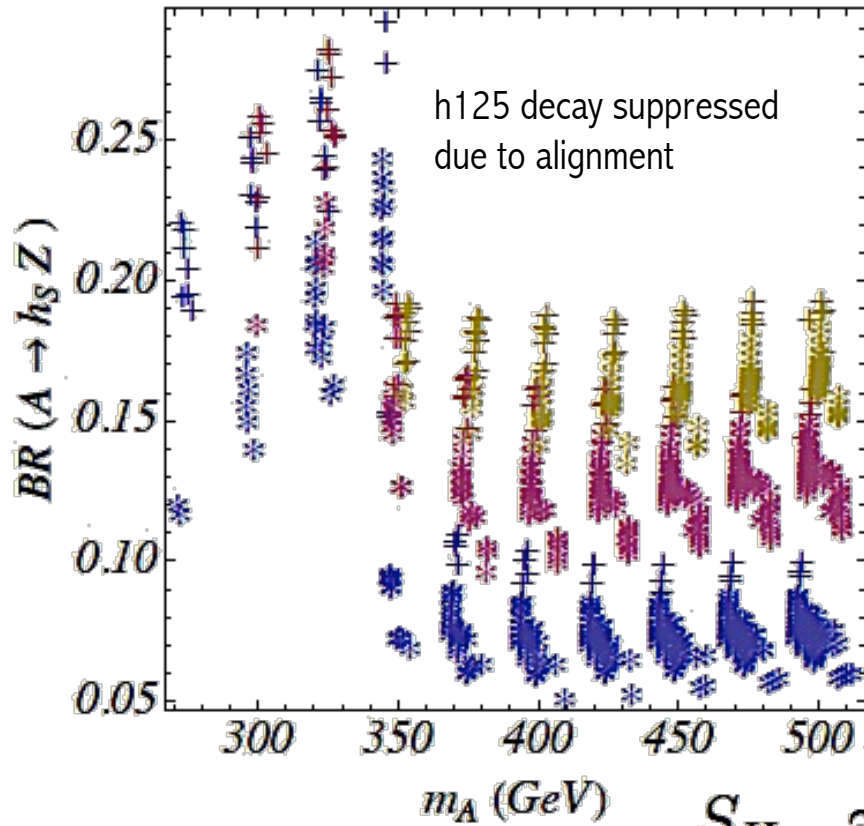
. tan β = 2 . tan β = 2.5 . tan β = 3

MSSM-like H BR

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- Into singlet-like H1/H2-Z and inos

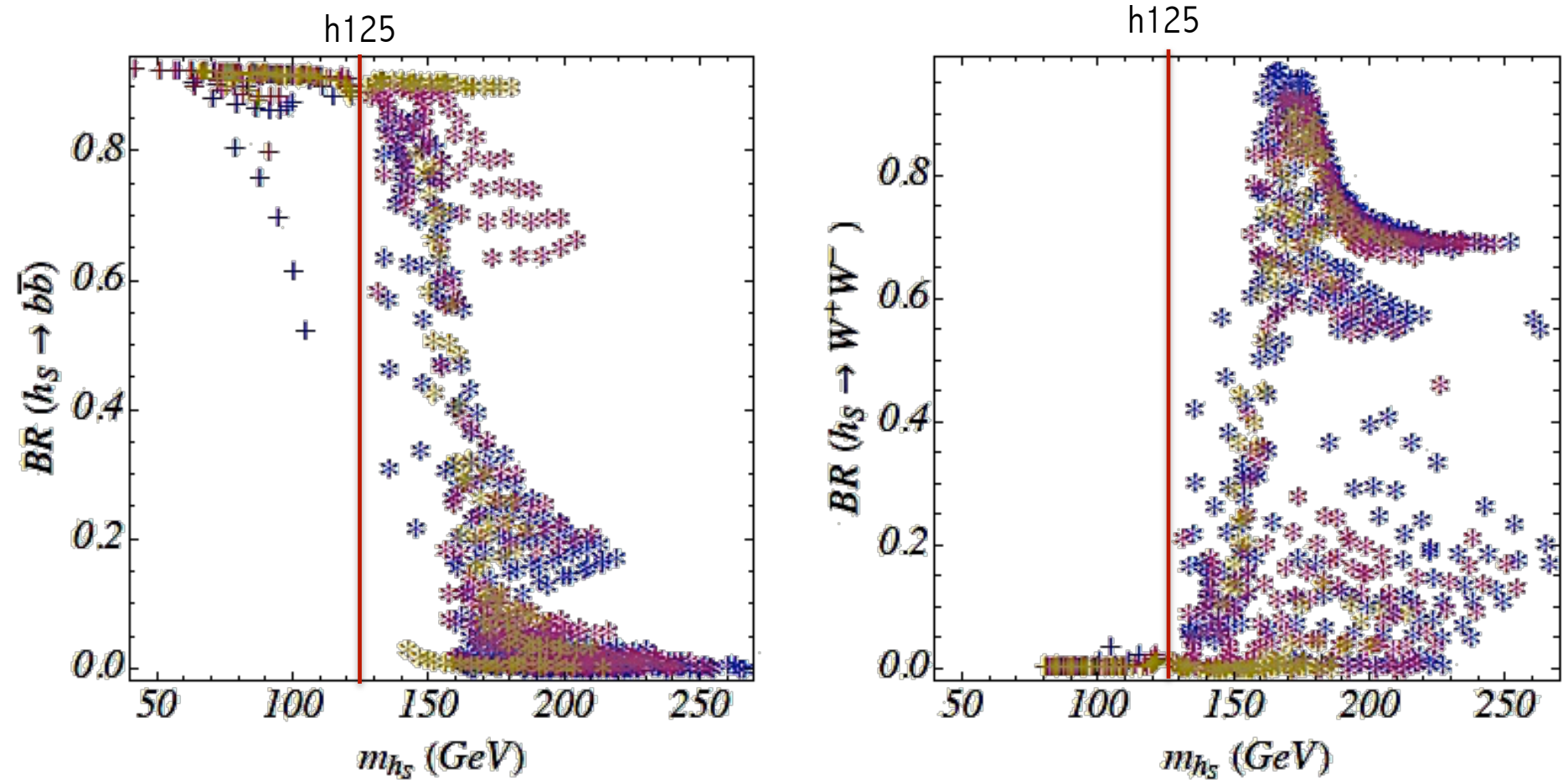
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$$S_{H,s} \approx -\frac{\lambda v}{2\mu} c_{2\beta} s_{2\beta}$$

MSSM-like A BR

- Singlet mainly decays to $b\bar{b}$ and WW

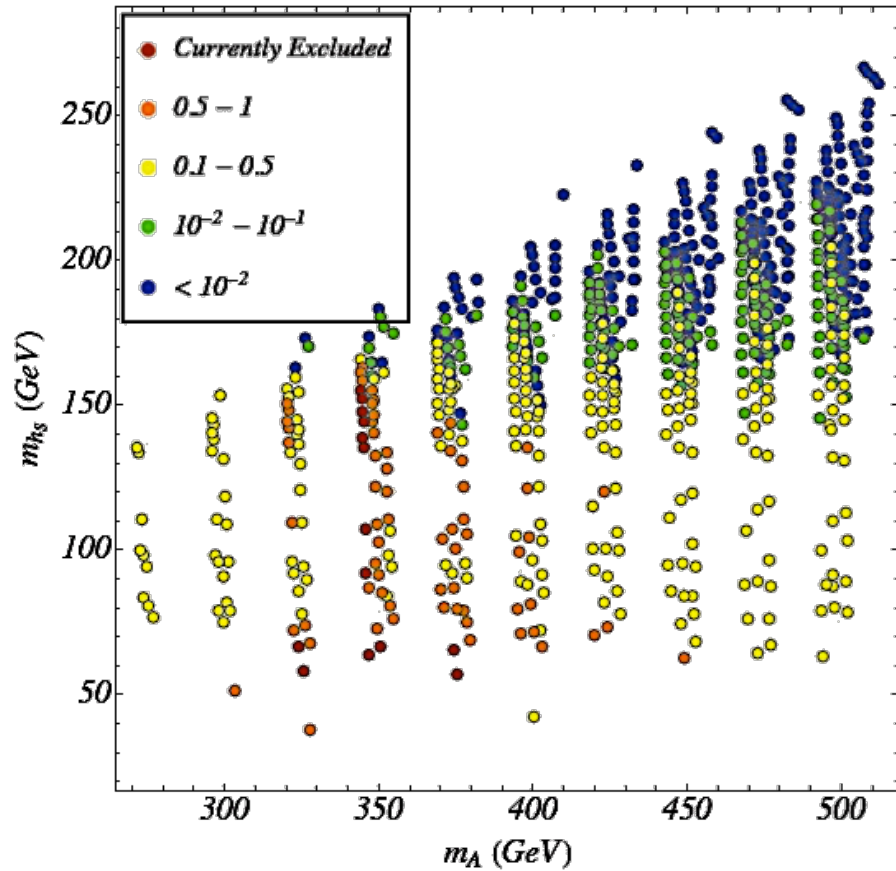


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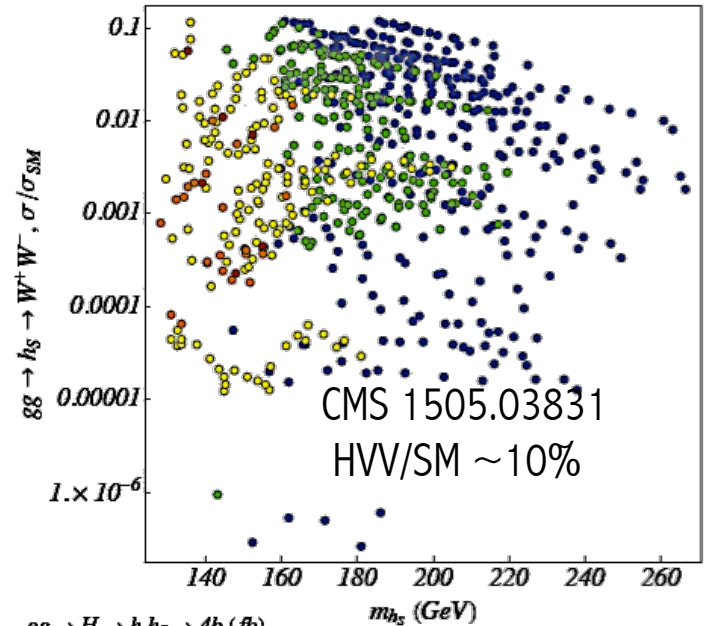
Singlet-like h_s BR

CMS PAS HIG-15-001

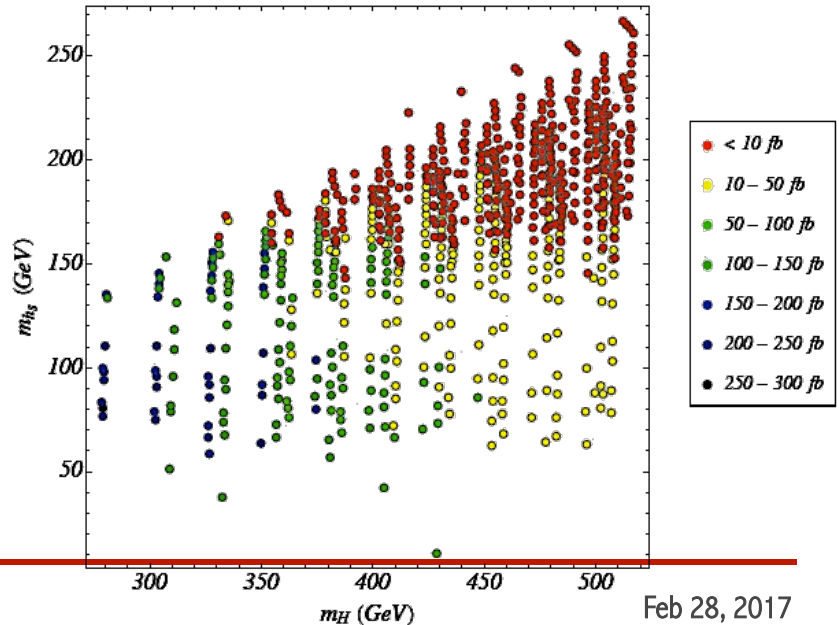
$gg \rightarrow A \rightarrow Z h_S \rightarrow ll bb, \sigma/\sigma_{Obs.Lim}$



8 TeV



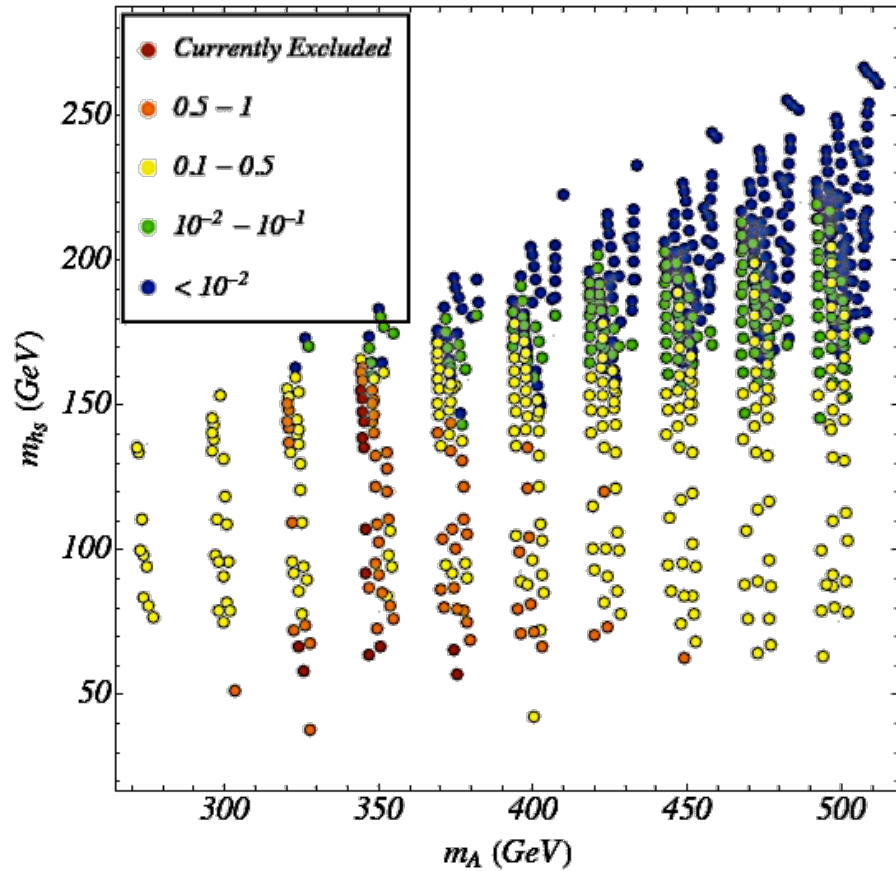
$gg \rightarrow H \rightarrow h h_S \rightarrow 4b (fb)$



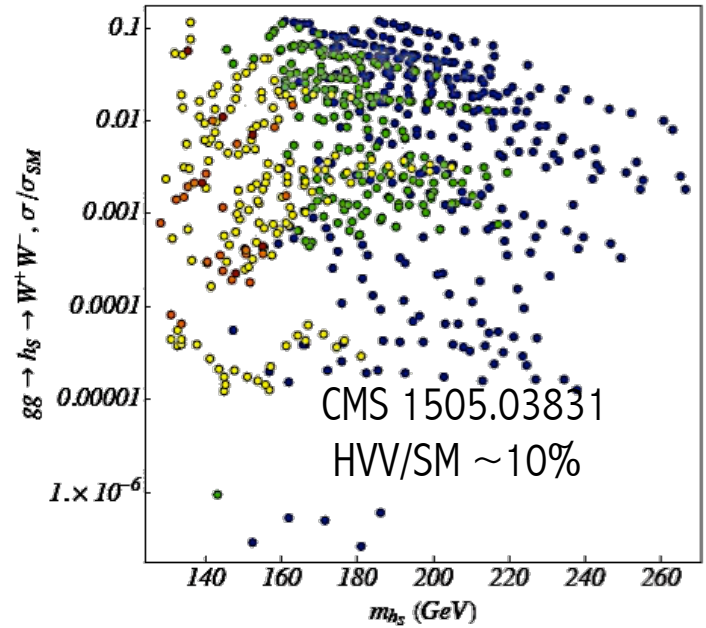
Now & Future

CMS PAS HIG-15-001

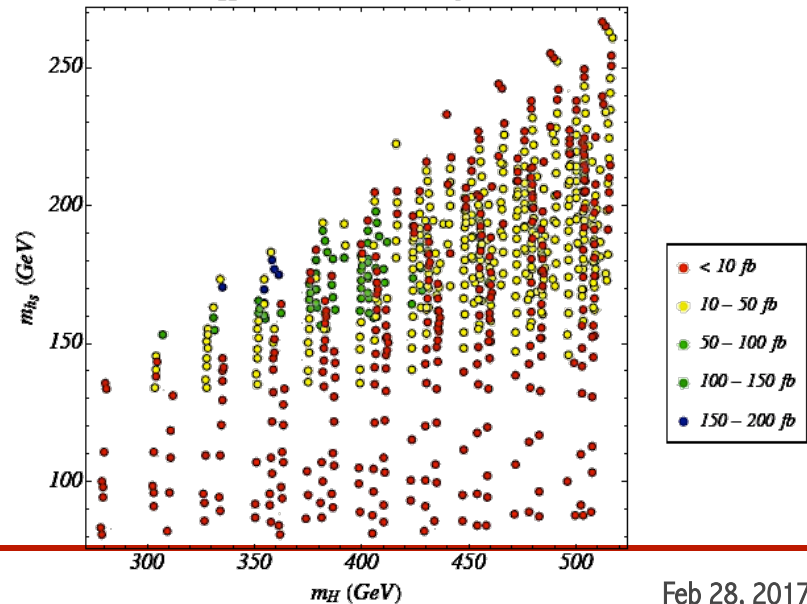
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8 TeV

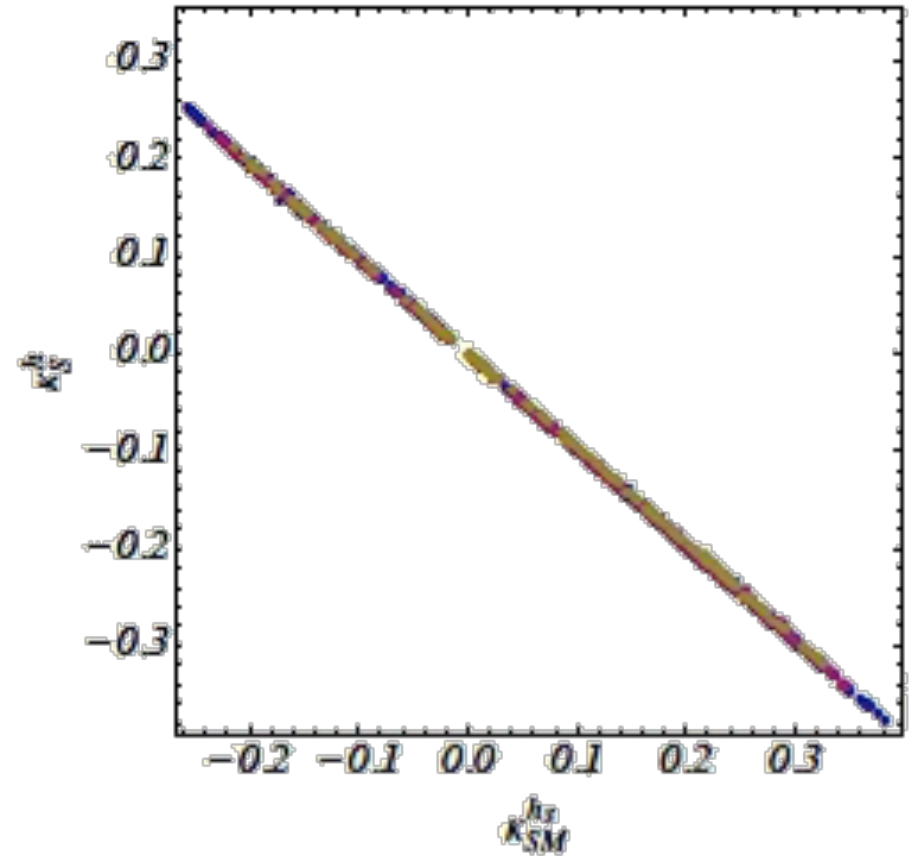
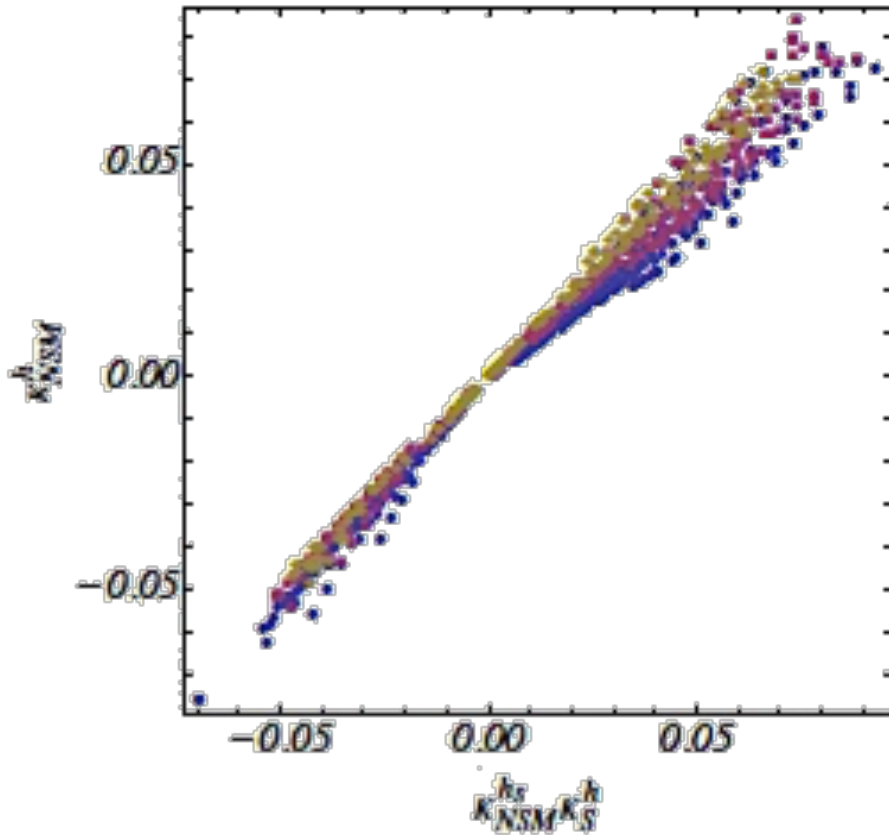


$gg \rightarrow H \rightarrow h h_S \rightarrow 2b 2W (fb)$



Now & Future

$$\begin{pmatrix} h \\ H \\ h_S \end{pmatrix} \simeq \begin{pmatrix} 1 & -\eta\eta' & \eta' \\ \mathcal{O}(\epsilon) & -1 & -\eta \\ -\eta' & -\eta & 1 \end{pmatrix} \begin{pmatrix} H^{\text{SM}} \\ H^{\text{NSM}} \\ H^{\text{S}} \end{pmatrix}$$



h125 Components

- Tree-level mass matrix in the $(H_{\text{NSM}}, H_{\text{SM}}, S)$ basis:

$$\begin{pmatrix} m_A^2 + s_{2\beta}^2 (m_Z^2 - \frac{1}{2}\lambda^2 v^2) & s_{2\beta} c_{2\beta} (m_Z^2 - \frac{1}{2}\lambda^2 v^2) & -\frac{\lambda v \mu}{\sqrt{2}} c_{2\beta} \left(\frac{m_A^2}{2\mu^2} s_{2\beta} + \frac{\kappa}{\lambda} \right) \\ c_{2\beta}^2 m_Z^2 + \frac{1}{2}\lambda^2 v^2 s_{2\beta}^2 & \sqrt{2}\lambda v \mu \left(1 - \frac{m_A^2}{4\mu^2} s_{2\beta}^2 - \frac{\kappa}{2\lambda} s_{2\beta} \right) & \\ \frac{\lambda^2 v^2 s_{2\beta}}{4} \left(\frac{m_A^2 s_{2\beta}}{2\mu^2} - \frac{\kappa}{\lambda} \right) + \frac{\kappa \mu}{\lambda} \left(A_\kappa + \frac{4\kappa \mu}{\lambda} \right) & & \end{pmatrix}$$

- **Alignment: Mixing between HNSM-SM=0 & SM-S =0**
- Alignment conditions, (+ stop corrections to always obtain h125):

$$\mathcal{M}_S^2(1, 2) = \frac{1}{t_\beta} \left[c_{2\beta} m_Z^2 - \mathcal{M}_S^2(2, 2) + \lambda^2 v^2 s_\beta^2 + \frac{3m_t^4 X_t (X_t - Y_t)}{4\pi^2 v^2 M_S^2} \left(1 - \frac{X_t^2}{6M_S^2} \right) \right] = 0$$

$$\mathcal{M}_S^2(2, 3) = 2\lambda v \mu \left(1 - \frac{m_A^2 s_{2\beta}^2}{4\mu^2} - \frac{\kappa s_{2\beta}}{2\lambda} \right) = 0$$

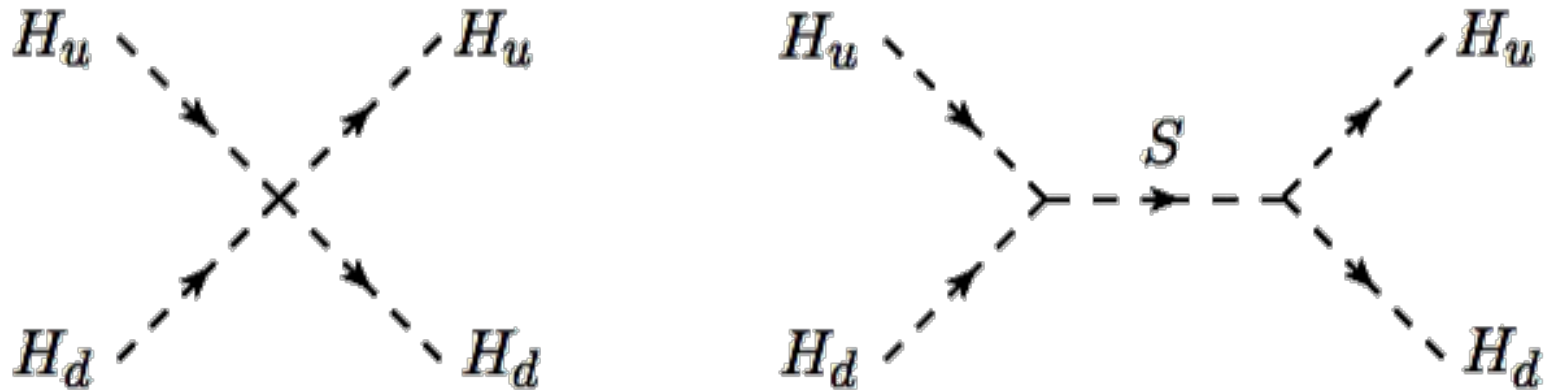
NSM-SM mixing cancels for

$$\lambda_{\text{alt}}^2 = \frac{m_h^2 - M_Z^2 c_{2\beta}}{v^2 s_\beta^2}$$

“Extended” Higgs Basis

$$W \supset \lambda S H_u H_d + \frac{M}{2} S^2 + \mu H_u H_d \quad V_{\text{soft}} \supset m_S^2 |S|^2$$

$$V \supset |F_S|^2 = |\lambda H_u H_d + M S|^2$$



$$\delta m_h^2 = \lambda^2 v^2 \sin^2 2\beta \left(1 - \frac{M^2}{M^2 + m_S^2} \right)$$

- It is well known that in the NMSSM there are new contributions to the lightest CP-even Higgs mass,

$$W = \lambda S H_u H_d + \frac{\kappa}{3} S^3$$

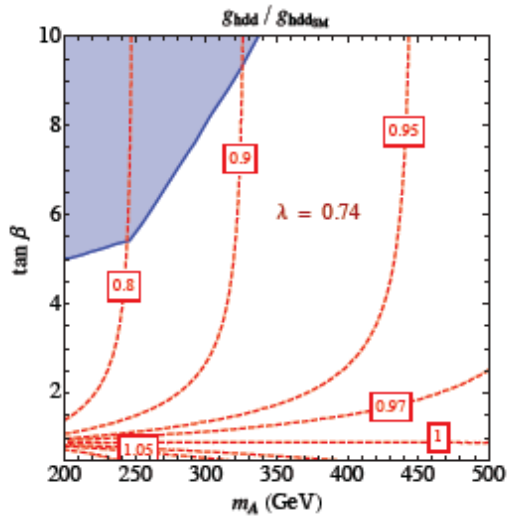
$$m_h^2 \simeq \lambda^2 \frac{v^2}{2} \sin^2 2\beta + M_Z^2 \cos^2 2\beta + \Delta_{\tilde{t}}$$

- It is perhaps less known that it leads to sizable corrections to the mixing between the MSSM like CP-even states. In the Higgs basis,

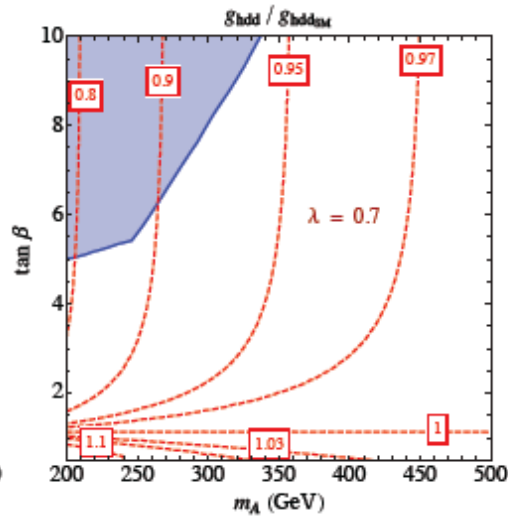
$$M_S^2(1,2) \simeq \frac{1}{\tan \beta} (m_h^2 - M_Z^2 \cos 2\beta - \lambda^2 v^2 \sin^2 \beta + \delta_{\tilde{t}})$$

- The last term is the one appearing in the MSSM, that are small for moderate mixing and small values of $\tan \beta$
- So, alignment leads to a determination of lambda,
- The values of lambda end up in a very narrow range, between 0.65 and 0.7 for all values of tan beta, that are the values that lead to naturalness with perturbativity up to the GUT scale

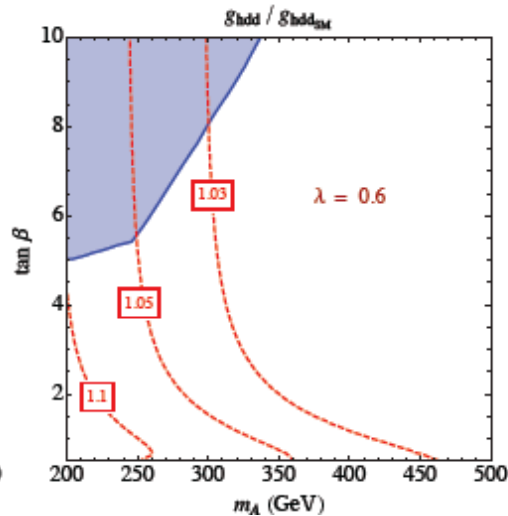
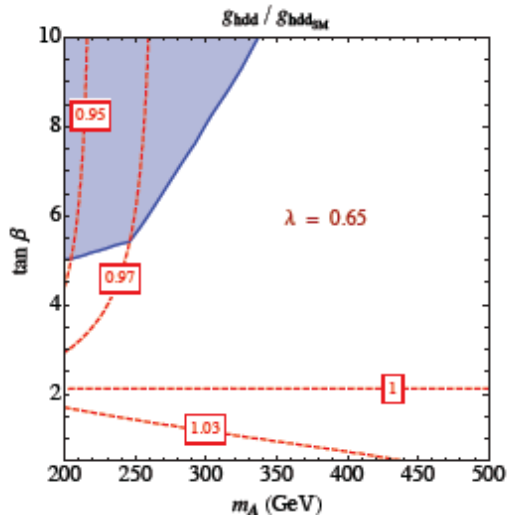
$$\lambda^2 = \frac{m_h^2 - M_Z^2 \cos 2\beta}{v^2 \sin^2 \beta}$$



(iii)



(iv)



M. Carena, I. Low, N.R.S, C. Wagner, '13

It is clear from these plots that the NMSSM does an amazing job in aligning the MSSM-like CP-even sector, provided **lambda is of about 0.65**

- In the (“MSSM m_A ”, Singlet) basis:

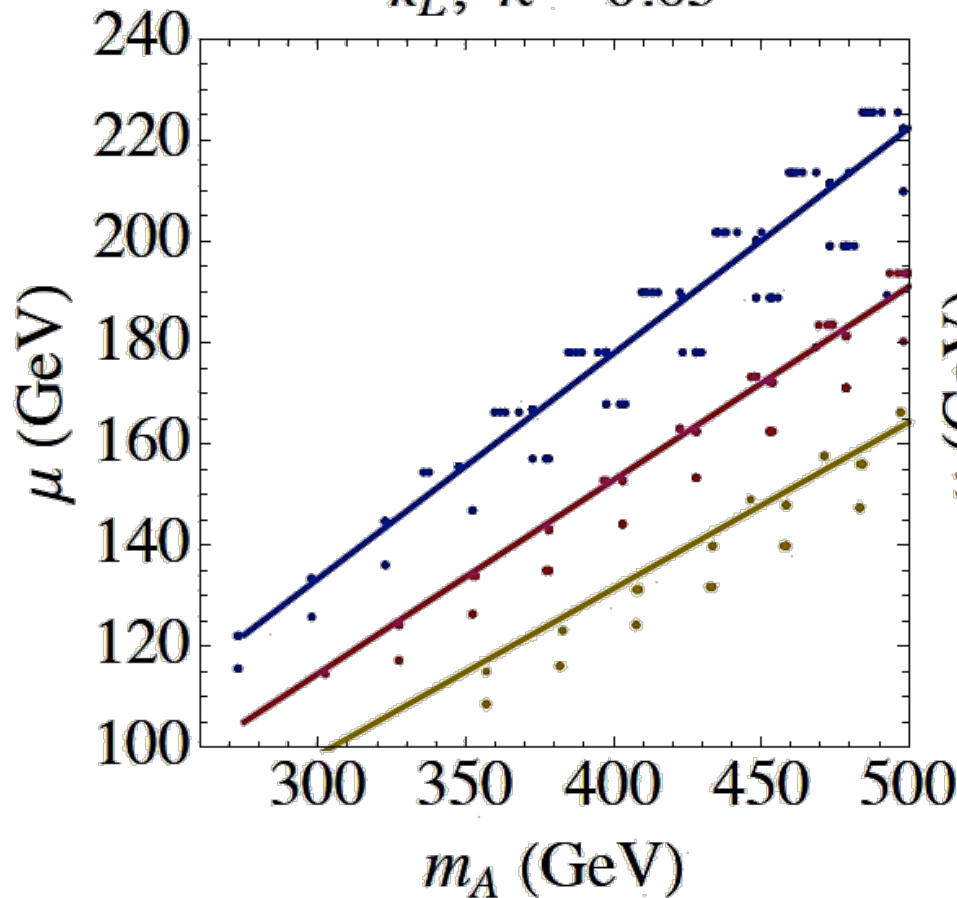
$$\mathcal{M}_P^2 = \begin{pmatrix} m_A^2 & \frac{\lambda v}{\sqrt{2}} \left(\frac{m_A^2}{2\mu} s_{2\beta} - \frac{3\kappa\mu}{\lambda} \right) \\ \frac{1}{2} \lambda^2 v^2 s_{2\beta} \left(\frac{m_A^2}{4\mu^2} s_{2\beta} + \frac{3\kappa}{2\lambda} \right) & - \frac{3\kappa A_\kappa \mu}{\lambda} \end{pmatrix}$$

$$m_A^2 = \frac{\mu}{s_\beta c_\beta} \left(A_\lambda + \frac{\kappa\mu}{\lambda} \right).$$

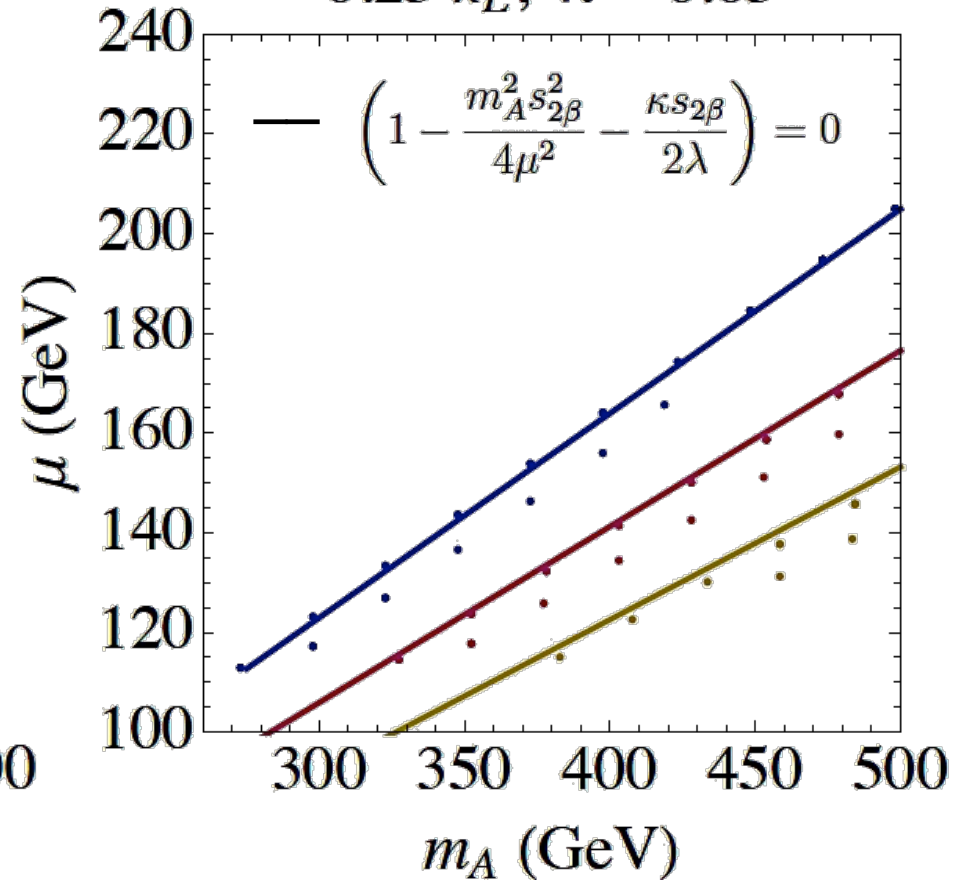
CP-Odd Mass Matrix

. $\tan \beta = 2$. $\tan \beta = 2.5$. $\tan \beta = 3$

$\kappa_L, \lambda = 0.65$



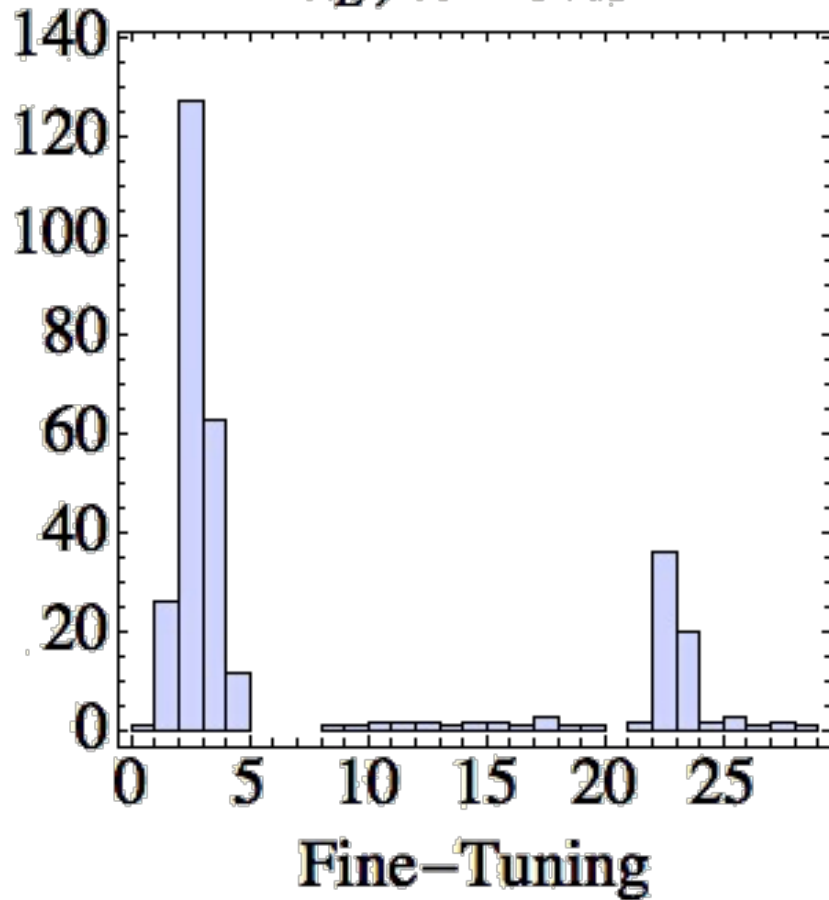
$0.25 \kappa_L, \lambda = 0.65$



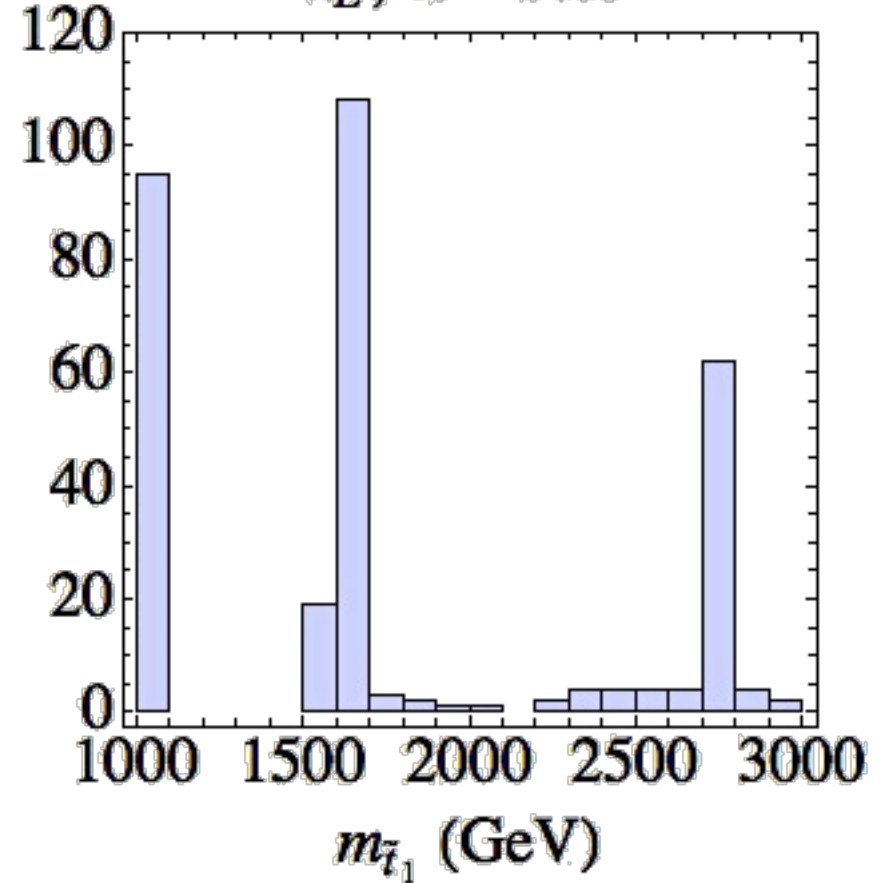
Misalignment

- Larger $\tan \beta \Rightarrow$ Heavier Stops \Rightarrow Larger Fine Tuning

$\kappa_L, \lambda = 0.65$



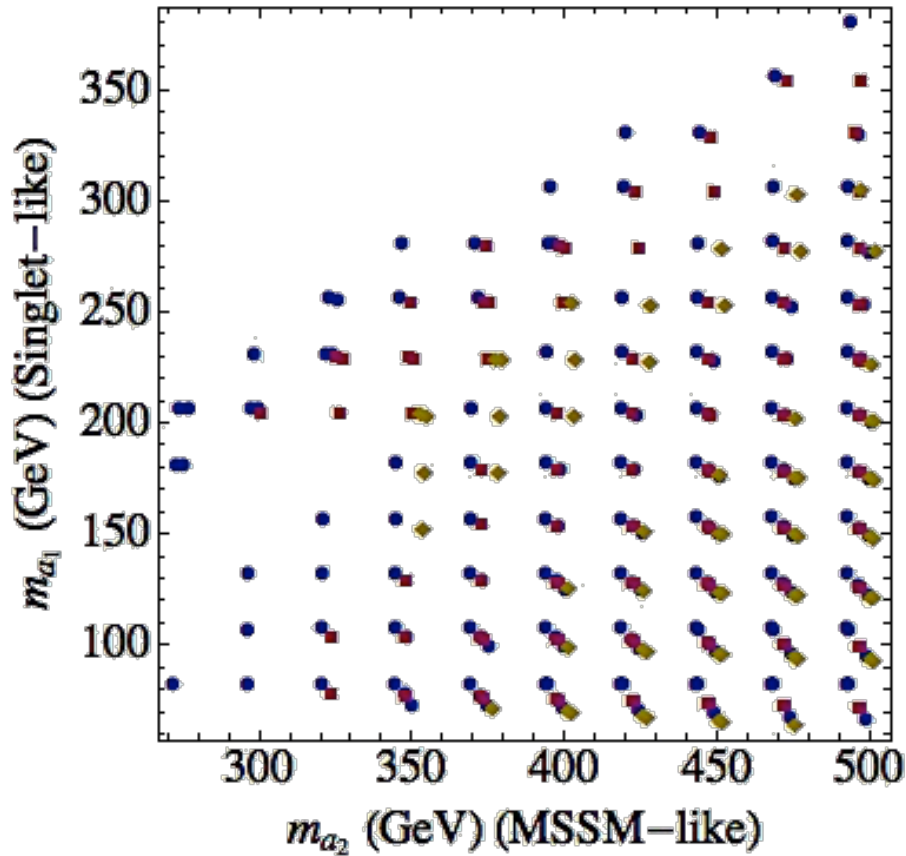
$\kappa_L, \lambda = 0.65$



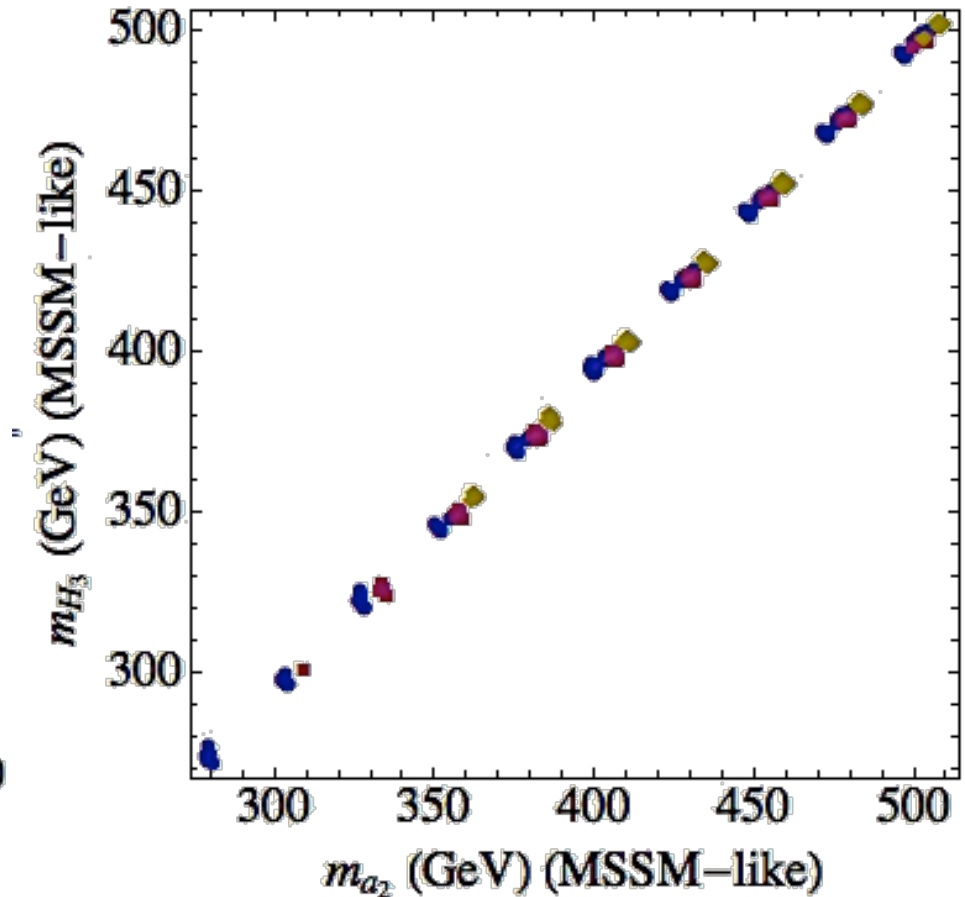
Fine-Tuning (No-mixing in stops)

- CP-even vs. Odd

$\kappa_L, \lambda = 0.65$

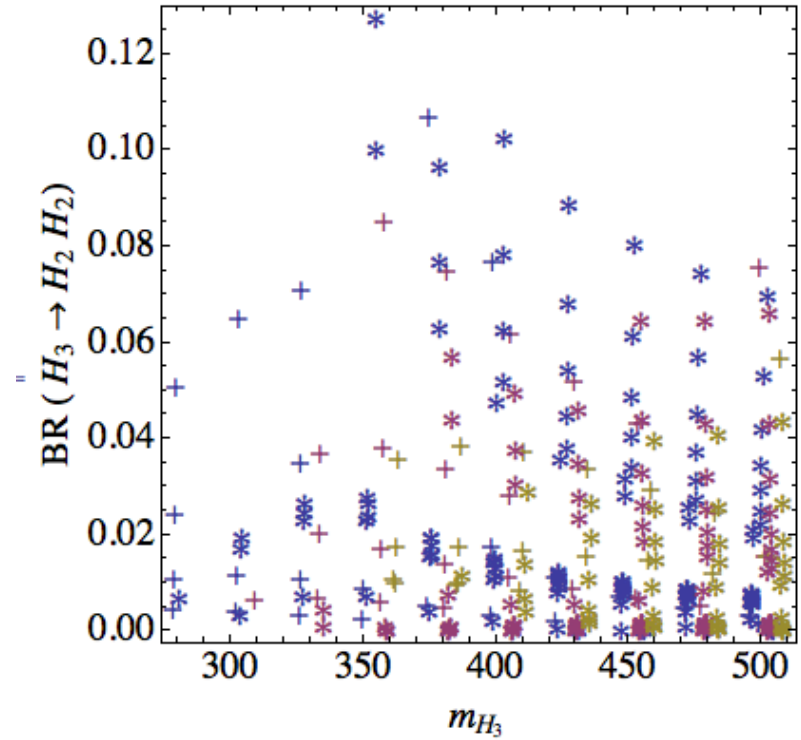
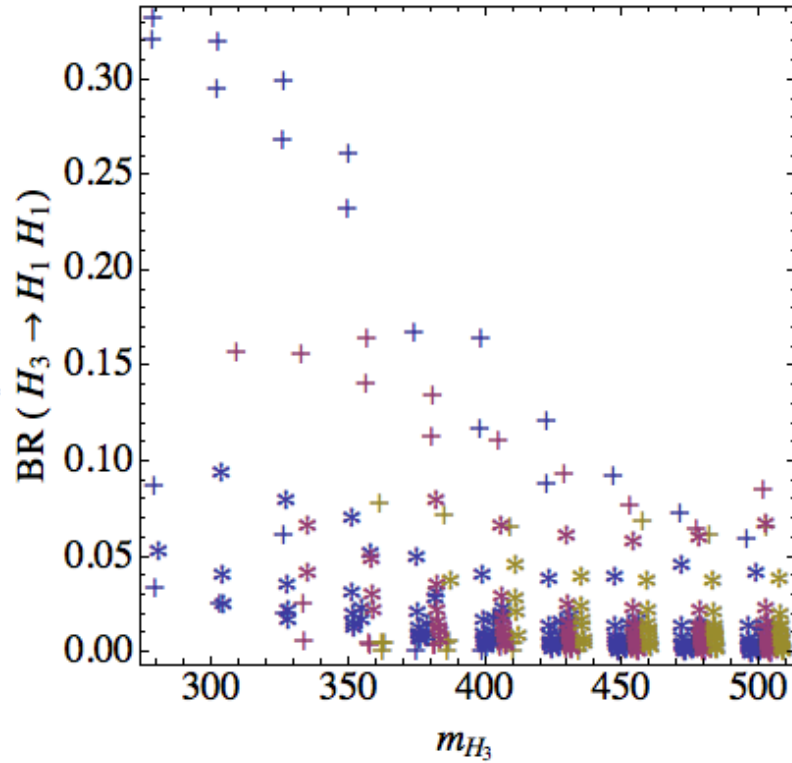


$\kappa_L, \lambda = 0.65$



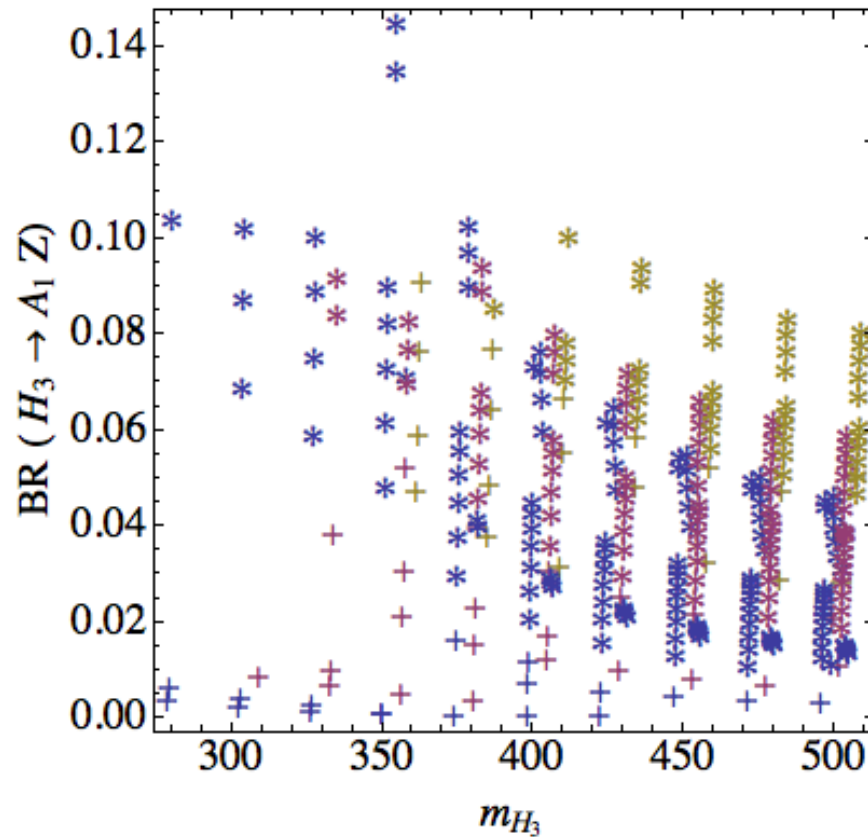
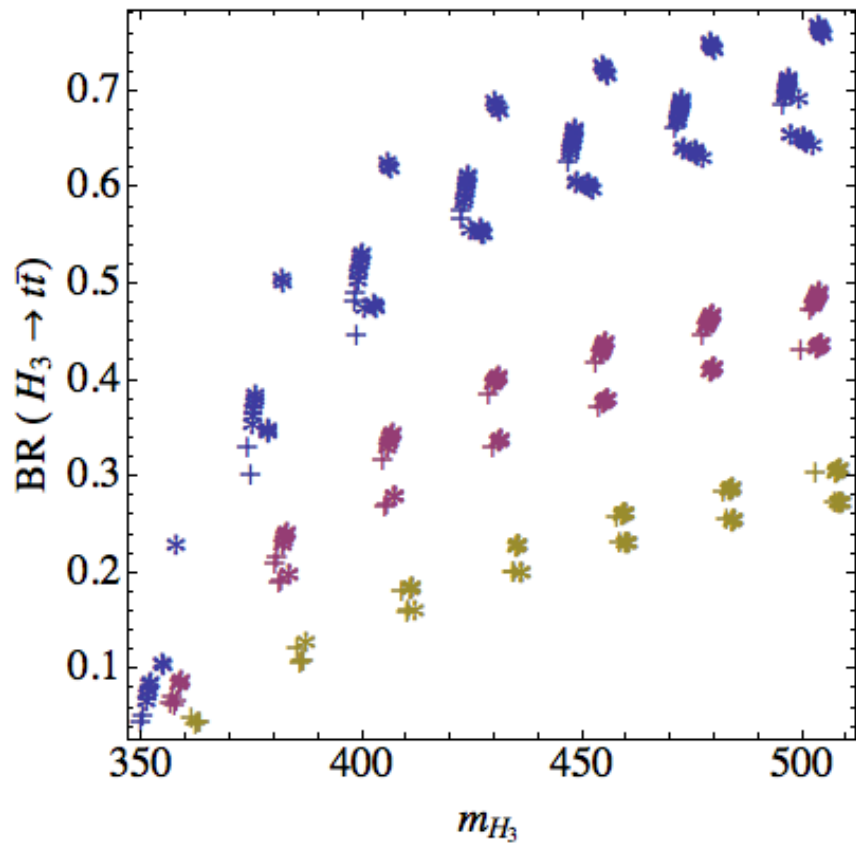
Higgs Spectra

- H1 H1 /H2 H2



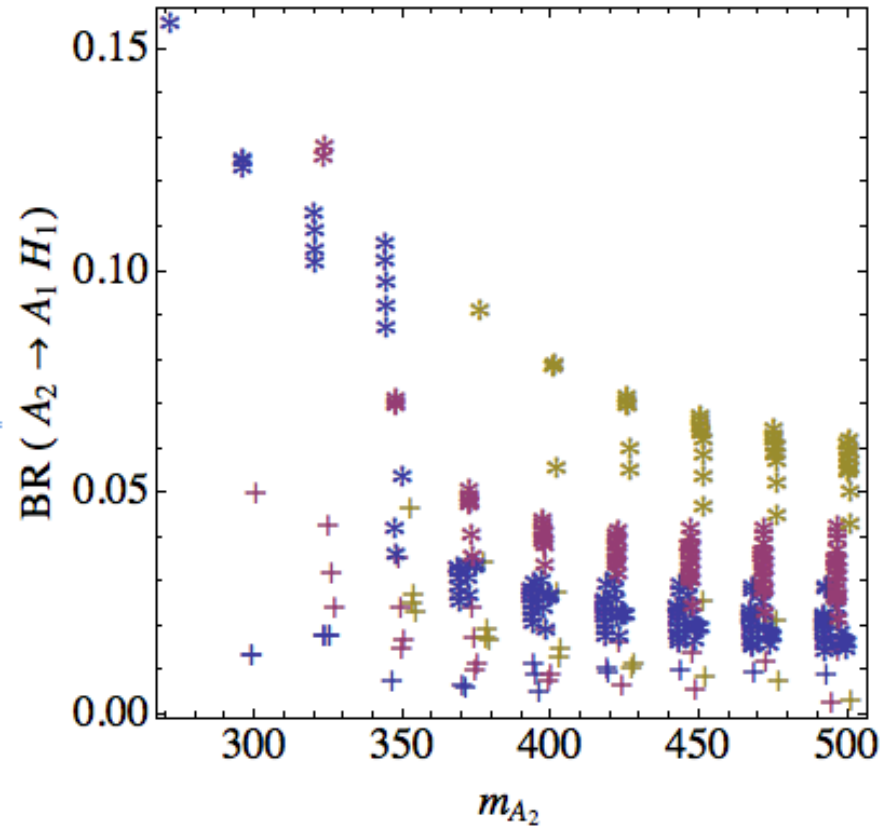
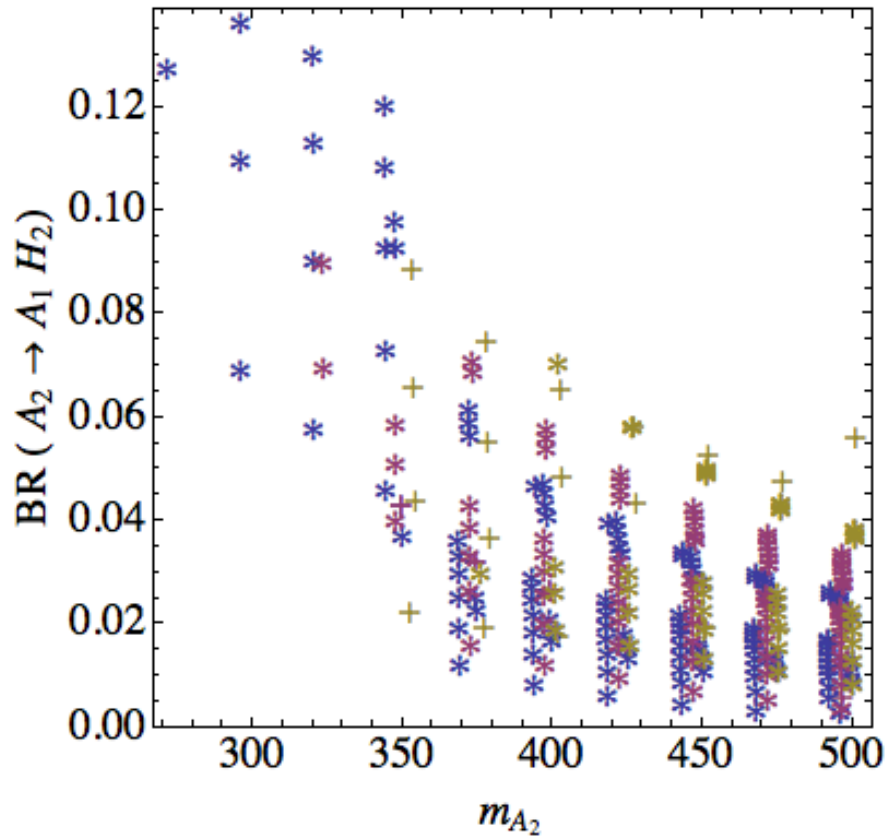
H3 BR

- $tt + a Z$



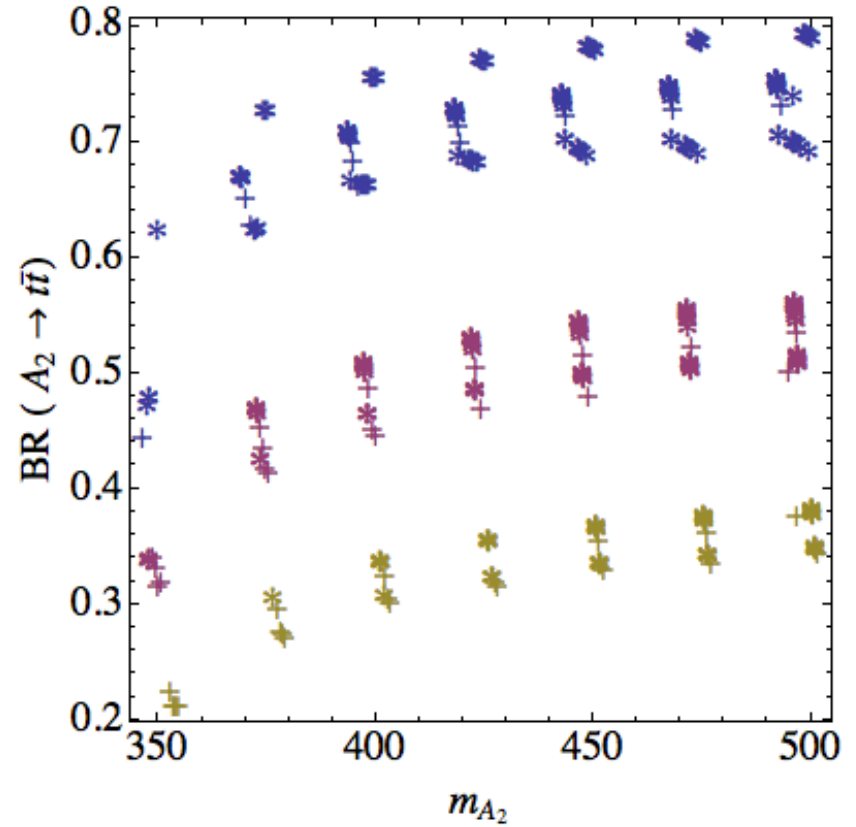
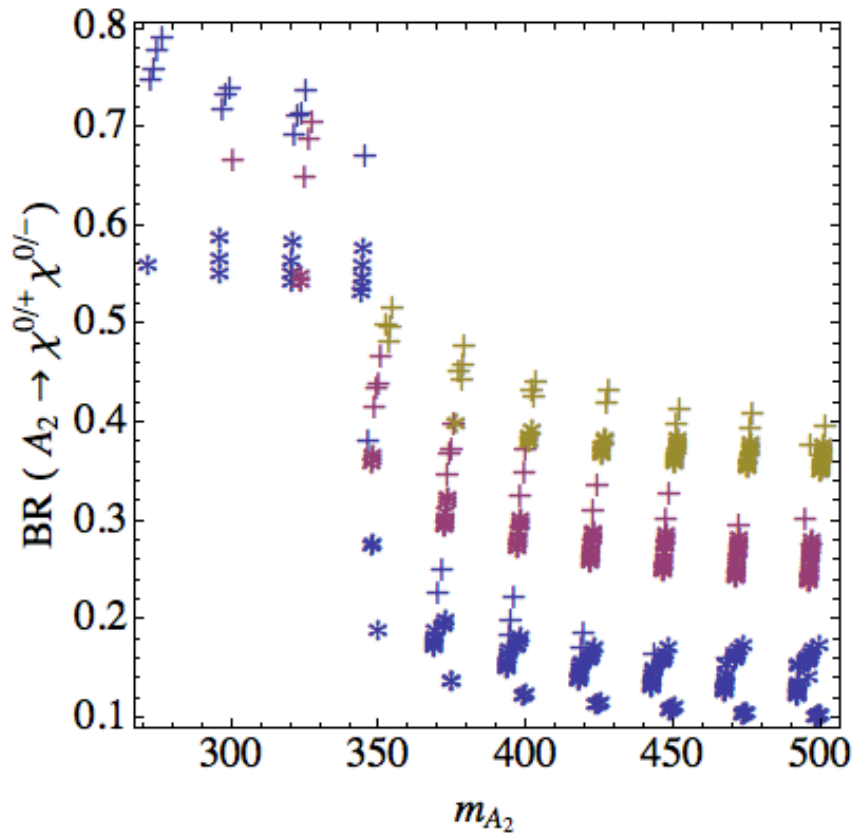
BR H3

- A1 H1 / A1 H2



A2 BR

- $tt + \text{chargino/neutralinos}$

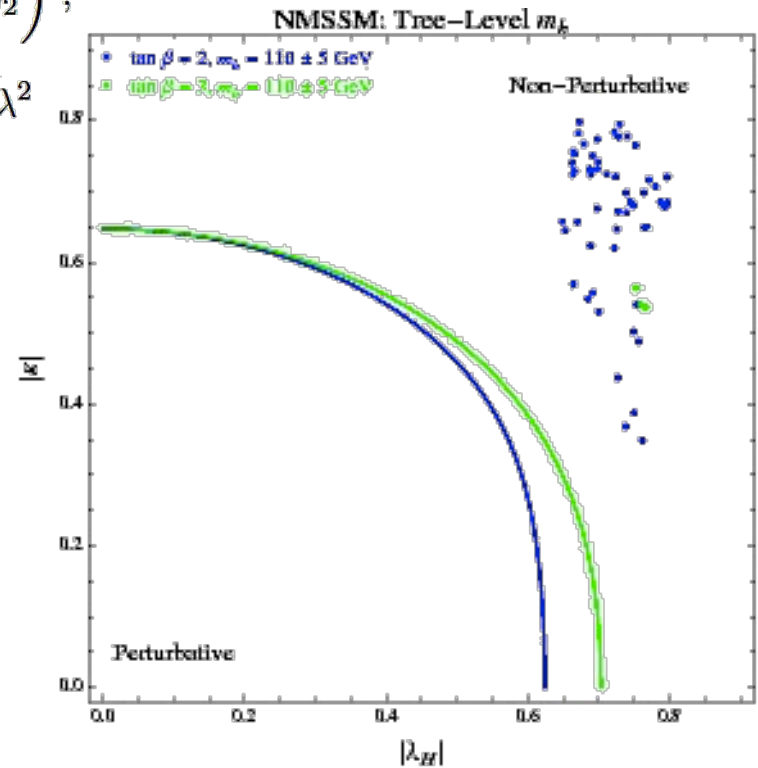
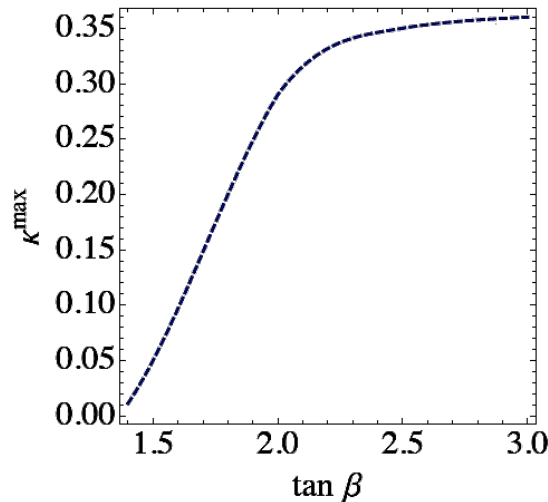


A2 BR

$$16\pi^2 \frac{d\lambda^2}{dt} = \lambda^2 \left(3h_t^2 + 3h_b^2 + h_\tau^2 + 4\lambda^2 + 2\kappa^2 - g_1^2 - 3g_2^2 \right) + \frac{\lambda^2}{16\pi^2} \left(-10\lambda^4 - 9h_t^4 - 9h_b^4 - 3h_\tau^4 - 8\kappa^4 - 9\lambda^2 h_t^2 - 9\lambda^2 h_b^2 - 3\lambda^2 h_\tau^2 - 12\lambda^2 \kappa^2 - 6h_t^2 h_b^2 + 2g_1^2 \lambda^2 + \frac{4}{3}g_1^2 h_t^2 - \frac{2}{3}g_1^2 h_b^2 + 2g_1^2 h_\tau^2 + 6g_2^2 \lambda^2 + 16g_3^2 h_t^2 + 16g_3^2 h_b^2 + \frac{23}{2}g_1^4 + \frac{15}{2}g_2^4 + 3g_1^2 g_2^2 \right),$$

$$16\pi^2 \frac{d\kappa^2}{dt} = \kappa^2 \left(6\lambda^2 + 6\kappa^2 \right) + \frac{\kappa^2}{16\pi^2} \left(-24\kappa^4 - 12\lambda^4 - 24\kappa^2 \lambda^2 - 18h_t^2 \lambda^2 - 18h_b^2 \lambda^2 - 6h_\tau^2 \lambda^2 + 6g_1^2 \lambda^2 + 18g_2^2 \lambda^2 \right).$$

$\lambda = 0.65$



RGE

1. The gluon fusion production cross section of the would be heavy MSSM states (**A** and **H**) is enhanced due to the top Yukawa contributions at low $\tan \beta$, and can be of the order of a few pb .
2. The non-standard Higgs bosons can have relevant decays into the lighter singlet like Higgs bosons as well as into the light electroweakinos.
3. The decay of non-standard Higgs bosons into tops, taus and bottoms will be suppressed due to the small values of $\tan \beta$ and the presence of additional decays.

Phenomenological Consequences