

Non-standard heavy vector bosons at the LHC

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1. Introduction
2. Z' boson phenomenology with non-standard decays at LHC
3. Doubly-charged bilepton signal at LHC
4. Conclusions

G.C. and S. Gentile, Nucl. Phys. B886 (2013) 293; G.C., EPJ C75 (2015) 264

Work in progress in collaboration with J. Araz, M. Frank and B. Fuks

G.C., C. Corianò, A. Costantini and P.H. Frampton, arXiv:1707.01381.

Searches for heavy gauge bosons Z' among the main objectives of LHC

GUT-inspired $U(1)'$, Sequential Standard Model, Kaluza–Klein models

LHC analyses focus on SM decays, e.g. high-mass dileptons or dijets

CMS Dileptons: $m(Z'_{\text{SSM}}) > 4.0 \text{ TeV}$ $m(Z'_{\text{GUT}}) > 3.5 \text{ TeV}$ Dijets: $m_{Z'} > 2.1\text{--}3.3 \text{ TeV}$

ATLAS Dileptons: $m(Z'_{\text{SSM}}) > 4.1 \text{ TeV}$ $m(Z'_{\text{GUT}}) > 3.4\text{--}3.7 \text{ TeV}$ Dijets $m_{Z'} > 2\text{--}2.6 \text{ TeV}$

In BSM analyses, one may consider BSM Z' decays, e.g. in supersymmetry

Lower SM branching ratios with BSM decays \Rightarrow lower Z' mass exclusion limits

Z' standard decays still useful for searches, BSM modes for supersymmetry

Z' constrains sparticle invariant masses, e.g. $Z' \rightarrow \tilde{\ell}^+ \tilde{\ell}^- \Rightarrow m_{Z'} = m_{\tilde{\ell}^+ \tilde{\ell}^-}$

Supersymmetric Z' decays allow study of unexplored phase space

Decays $Z' \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$: monojet events and Dark Matter candidates

Related work on supersymmetric Z' decays:

Gherghetta et al ('98), Kang & Langacker ('05), Baumgart et al ('07), Chang et al ('11)

$U(1)'$ gauge groups in GUT-inspired models:

$$E_6 \rightarrow SO(10) \times U(1)'_\psi \rightarrow SU(5) \times U(1)'_\chi \times U(1)'_\psi$$

$$Z'(\theta) = Z'_\psi \cos \theta - Z'_\chi \sin \theta$$

$$E_6 \rightarrow SM \times U(1)'_\eta \quad \theta = \arccos \sqrt{5/8} \Rightarrow Z'_\eta$$

Orthogonal combination to Z'_η : $\theta = \arccos \sqrt{5/8} - \pi/2 \Rightarrow Z'_I$

Secluded model (singlet S): $\theta = \arctan(\sqrt{15}/9) - \pi/2 \Rightarrow Z'_S$

Z'_N : Z'_χ -like with ‘unconventional’ $SO(10)$ representations (**10** vs **6**, $\theta \rightarrow \theta + \arctan 15$)

Model	θ
Z'_χ	$-\pi/2$
Z'_ψ	0
Z'_η	$\arccos \sqrt{5/8}$
Z'_I	$\arccos \sqrt{5/8} - \pi/2$
Z'_N	$\arctan \sqrt{15} - \pi/2$
Z'_S	$\arctan(\sqrt{15}/9) - \pi/2$

Analysis will be carried out for Z'_ψ and Z'_η models, which yield higher cross sections

Minimal Supersymmetric Standard Model and U(1)' (a.k.a. UMSSM)

Extra singlet S to break U(1)' and give mass to the Z'

$$H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}, \quad H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix}, \quad S = S^0$$

Higgs sector after EWSB: h, H, A, H^\pm (MSSM) and a new scalar H'

Three vacuum expectation values v_u, v_d, v_S , $\tan \beta = v_u/v_d$

Gauginos: new \tilde{Z}' and \tilde{H}' imply two new neutralinos: $\tilde{\chi}_1^0, \dots, \tilde{\chi}_6^0$ ($\tilde{\chi}_{5,6}^0$ very heavy)

Chargino sector is unchanged, as the Z' is neutral

D-term correction to sfermion masses: $\tilde{m}^2 = \tilde{m}_0^2 + \Delta \tilde{m}^2$ (\tilde{m}_0 soft mass at Z' scale)

$$\Delta \tilde{m}_a^2 = g'^2 Q'_a (Q'_{H_u} v_u^2 + Q'_{H_d} v_d^2 + Q'_S v_S^2)/2 \quad ; \quad g' = \sqrt{\frac{5}{3}} g_1 \text{ (GUT)}$$

New Z' decay modes besides the SM ones:

$$Z' \rightarrow \tilde{q}\tilde{q}^*, \tilde{\ell}^+\tilde{\ell}^-, \tilde{\nu}\tilde{\nu}^*, \tilde{\chi}_i^0\tilde{\chi}_j^0, \tilde{\chi}_{1,2}^+\tilde{\chi}_{1,2}^-, ZH, Zh, H^+H^-$$

Benchmark: $m_{Z'} = 2$ TeV, consistency with SUSY exclusion and 125 GeV Higgs

$M_1 = 400$ GeV $\simeq M_2/2$, $M' = 1$ TeV, $\tan \beta = 30$, $\mu = 200$ GeV, $A_f \simeq 4$ TeV

$U(1)'_\psi$: $m_{\tilde{\ell}}^0 = m_{\tilde{\nu}_\ell}^0 = 1.2$ TeV, $m_{\tilde{q}}^0 = 5.5$ TeV ($q = u, d, c, s$),
 $m_{\tilde{b}}^0 = m_{\tilde{t}}^0 = 2.2$ TeV ($q_{1,2} \simeq q_{L,R}$, $\ell_{1,2} \simeq \ell_{L,R}$) A. Arbey et al, arXiv:1112.3028

SARAH computes mass matrices at NLO, SPheno creates model files in the UFO format

$m_{\tilde{d}_1}$	$m_{\tilde{u}_1}$	$m_{\tilde{s}_1}$	$m_{\tilde{c}_1}$	$m_{\tilde{b}_1}$	$m_{\tilde{t}_1}$
5609.8	5609.4	5609.9	5609.5	2321.7	2397.2
$m_{\tilde{d}_2}$	$m_{\tilde{u}_2}$	$m_{\tilde{s}_2}$	$m_{\tilde{c}_2}$	$m_{\tilde{b}_2}$	$m_{\tilde{t}_2}$
5504.9	5508.7	5504.9	5508.7	2119.6	2036.3

$m_{\tilde{\ell}_1}$	$m_{\tilde{\ell}_2}$	$m_{\tilde{\tau}_1}$	$m_{\tilde{\tau}_2}$	$m_{\tilde{\nu}_{\ell,1}}$	$m_{\tilde{\nu}_{\ell,2}}$	$m_{\tilde{\nu}_{\tau,1}}$	$m_{\tilde{\nu}_{\tau,2}}$
1392.4	953.0	1398.9	971.1	1389.8	961.5	1395.9	961.5

m_h	m_H	$m_{H'}$	m_A	m_{H^\pm}
125.0	1989.7	4225.0	4225.0	4335.6

$m_{\tilde{\chi}_1^+}$	$m_{\tilde{\chi}_2^+}$	$m_{\tilde{\chi}_1^0}$	$m_{\tilde{\chi}_2^0}$	$m_{\tilde{\chi}_3^0}$	$m_{\tilde{\chi}_4^0}$	$m_{\tilde{\chi}_5^0}$	$m_{\tilde{\chi}_6^0}$
204.8	889.1	197.2	210.7	408.8	647.9	889.0	6193.5

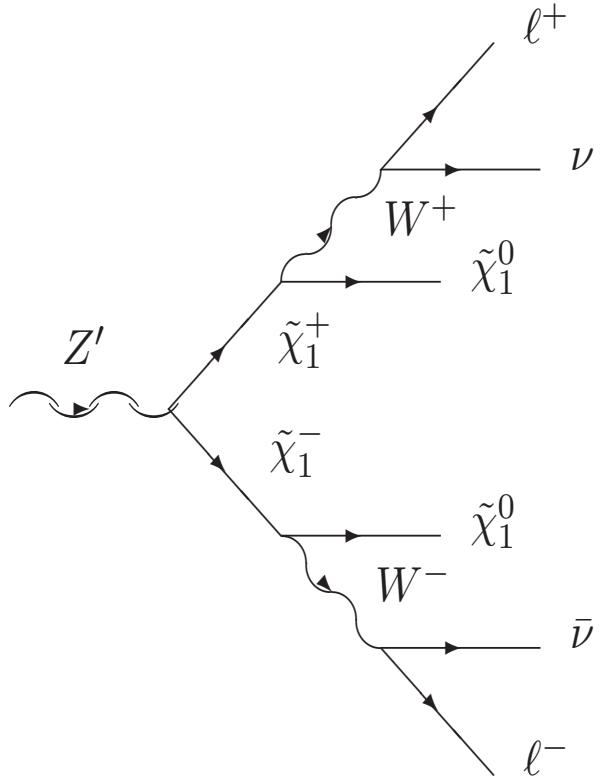
Branching ratios of Z'_ψ into SM ($\sim 70\%$) and BSM ($\sim 30\%$) final states

Final State	Z'_ψ Branching ratio (%)
$\tilde{\chi}_1^+ \tilde{\chi}_1^-$	10.2
$\tilde{\chi}_1^0 \tilde{\chi}_1^0$	4.9
$\tilde{\chi}_2^0 \tilde{\chi}_2^0$	5.1
$\tilde{\chi}_4^0 \tilde{\chi}_4^0$	8.0
hZ	1.4
W^+W^-	2.9
$\sum_i q\bar{q}$	50.1
$\sum_i \nu_i \bar{\nu}_i$	8.3
$\sum_i \ell_i^+ \ell_i^-$	8.3

$Z'_\psi \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-$ exhibits the highest branching ratio: need to consider $\tilde{\chi}_1^\pm$ rates

Final State	$\tilde{\chi}_1^+$ branching ratio (%)
$\tilde{\chi}_1^0 u\bar{d}$	34.3
$\tilde{\chi}_1^0 u\bar{c}$	1.8
$\tilde{\chi}_1^0 c\bar{d}$	1.6
$\tilde{\chi}_1^0 c\bar{s}$	29.3
$\tilde{\chi}_1^0 \ell^+ \nu_\ell$	32.9

Final states with leptons ($\ell = e, \mu$) and missing transverse energy



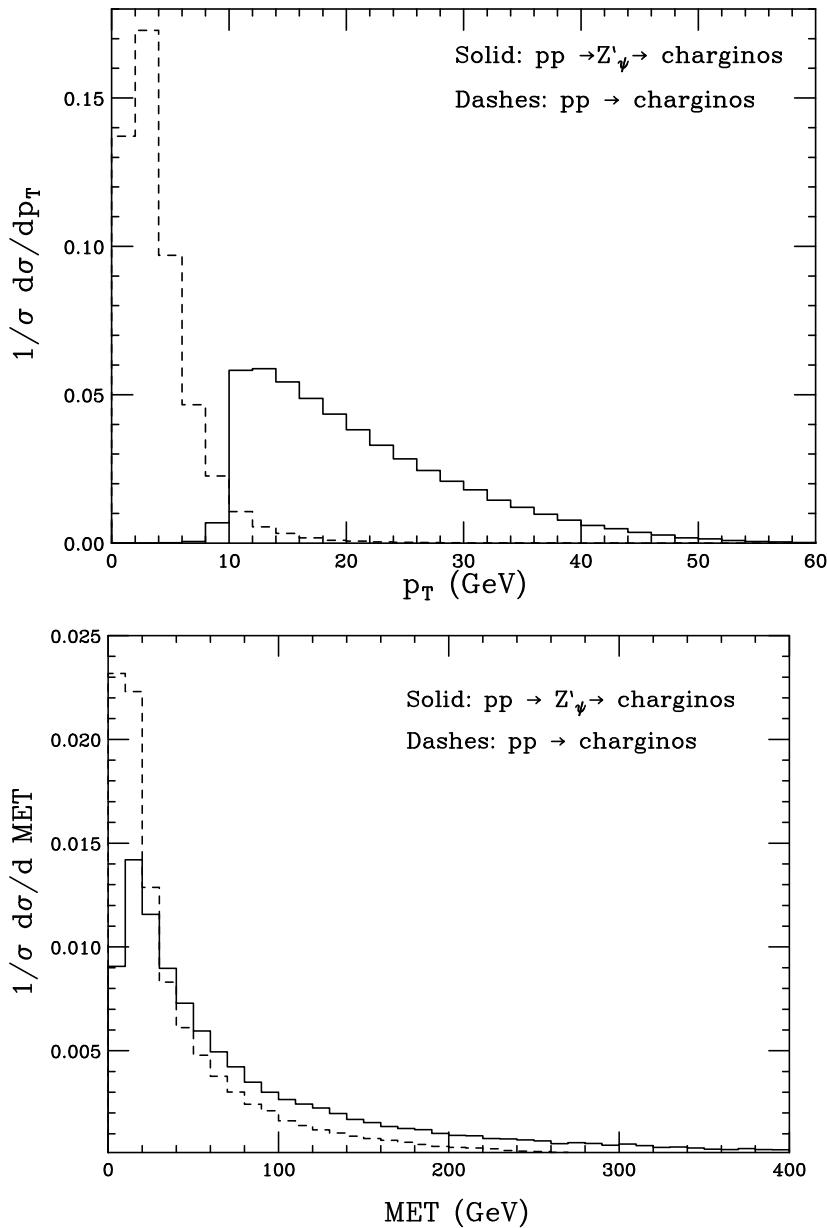
In the reference point, at $\sqrt{s} = 14$ TeV, using MadGraph and LO CTEQL1:

$$\sigma(pp \rightarrow Z'_\psi) \simeq 0.13 \text{ pb} ; \text{ BR}(Z'_\psi \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-) \simeq 10.2\% ; \text{ BR}(\tilde{\chi}_1^+ \rightarrow \tilde{\chi}_1^0 \ell^+ \nu_\ell) \simeq 24\%$$

$$\sigma(pp \rightarrow Z'_\psi \rightarrow \ell^+ \ell^- + \text{MET}) \simeq 8 \times 10^{-4} \text{ pb} \Rightarrow N \simeq 80 \text{ (100 fb}^{-1}\text{)} , N \simeq 240 \text{ (300 fb}^{-1}\text{)}$$

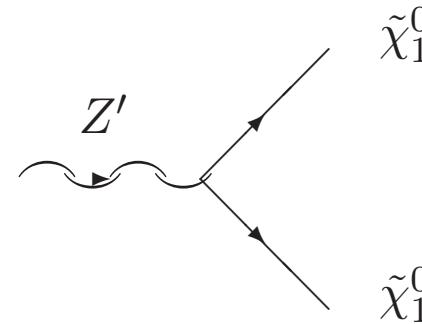
Competitive process: $pp \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow (\tilde{\chi}_1^0 \ell^+ \nu_\ell)(\tilde{\chi}_1^0 \ell^- \bar{\nu}_\ell)$ ($\sigma \simeq 1.15 \times 10^{-2} \text{ pb}$)

Phenomenology - $Z'_\psi \rightarrow \text{charginos}$ (MadGraph+HERWIG - $\sqrt{s} = 14 \text{ TeV}$)



$$\text{MET} = \sqrt{\left(\sum_i p_{x,i}\right)^2 + \left(\sum_i p_{y,i}\right)^2} \quad (i = \nu, \tilde{\chi}_1^0); \quad m_T = \sqrt{\left(\sum_j E_{T,j}\right)^2 - \left(\sum_j \vec{p}_{T,j}\right)^2} \quad (j = \ell, \nu, \tilde{\chi}_1^0)$$

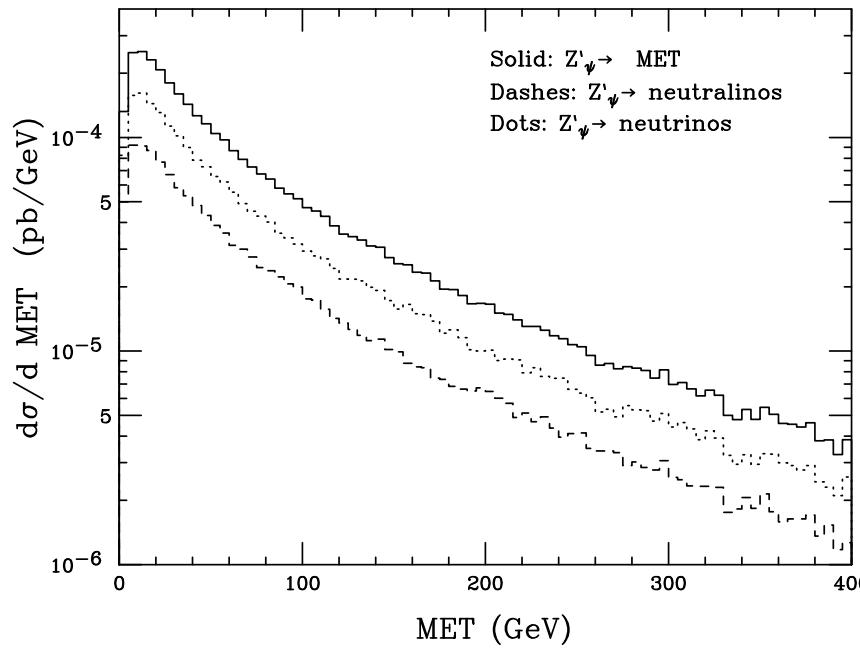
DM signals in Z' decays: $Z'_\psi \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$
 (MadGraph+HERWIG – $\tilde{\chi}_1^0$ mostly higgsino)



$$\text{BR}(Z'_\psi \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0) \simeq 10\% \Rightarrow \sigma(pp \rightarrow Z'_\psi \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0) \simeq 6.4 \times 10^{-3} \text{ pb at 14 TeV}$$

$N \simeq 640$ (100 fb^{-1}) or 2×10^3 (300 fb^{-1}) with possible Dark Matter candidates

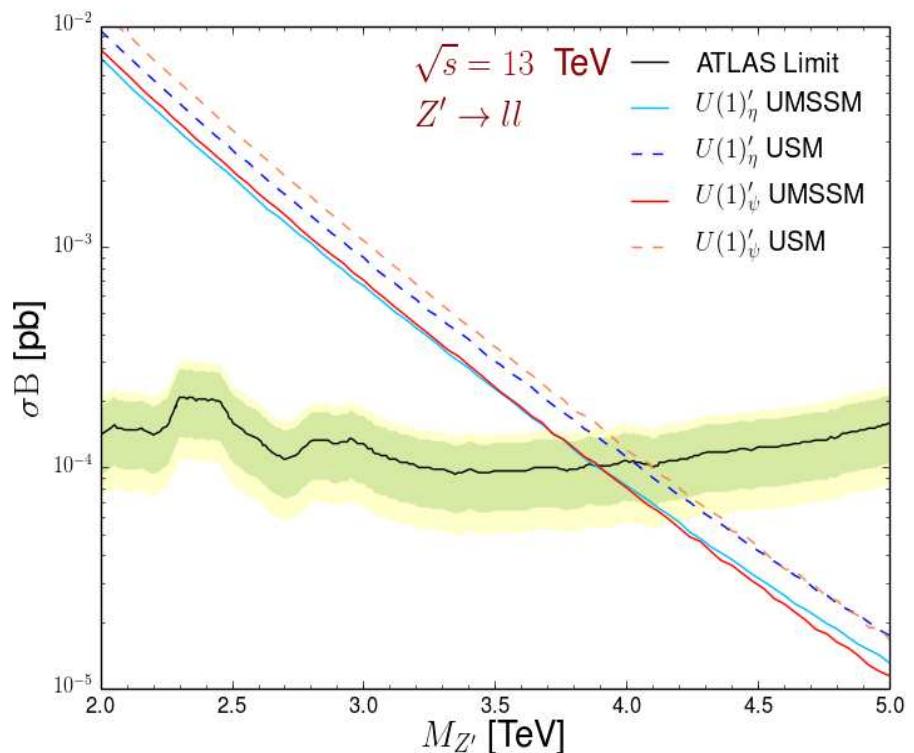
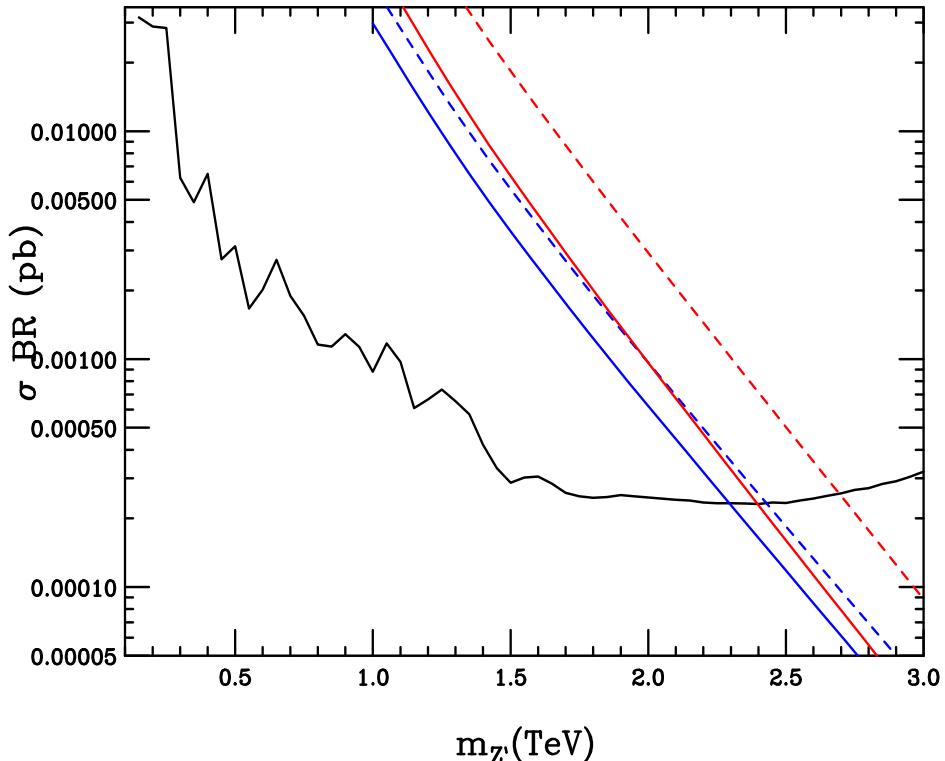
Competitive process: $Z'_\psi \rightarrow \nu \bar{\nu}$: $\sigma \simeq 1.1 \times 10^{-2}$; $N \simeq \mathcal{O}(10^3)$



Similar shapes ($m_{\tilde{\chi}_1^0} \ll m_{Z'}$), but $\sigma(pp \rightarrow \text{MET})$ increases by 60% adding neutralinos

In progress: implementation of jet/photon clustering algorithms

Mass exclusion limits in the SUSY reference point



Solid: SM+BSM decays ; Dashes: only SM decays;

Left: ATLAS at 7 TeV Red: Z'_SSM ; Blue: Z'_ψ

Right: ATLAS at 13 TeV Red: Z'_ψ ; Blue: Z'_η (Plot by J.Araz, preliminary)

Excluded-mass shift: Z'_SSM : $\Delta m \simeq 300 \text{ GeV}$; $Z'_{\psi,\eta}$: $\Delta m \simeq 200 \text{ GeV}$

Bilepton Model (Frampton'92, Pisano-Pleitez'92) : $SU(3)_c \times SU(3)_L \times U(1)_X$

$$\text{Fermions} : Q_1 = \begin{pmatrix} u_L \\ d_L \\ D_L \end{pmatrix}, \quad Q_2 = \begin{pmatrix} c_L \\ s_L \\ S_L \end{pmatrix}, \quad Q_{1,2} \in (3, 3, -1/3), \quad q(D, S) = -\frac{4}{3}$$

$$Q_3 = \begin{pmatrix} t_L \\ b_L \\ T_L \end{pmatrix}, \quad Q_3 \in (3, \bar{3}, 2/3), \quad q(T) = +\frac{5}{3}, \quad \ell = \begin{pmatrix} \ell_L \\ \nu_\ell \\ \bar{\ell}_R \end{pmatrix}, \quad \ell \in (1, \bar{3}, 0)$$

$$(d, s, b)_R \in (\bar{3}, 1, 1/3), \quad (u, c, t)_R \in (\bar{3}, 1, -2/3), \quad (D, S)_R \in (\bar{3}, 1, 4/3), \quad T_R \in (\bar{3}, 1, -5/3)$$

$$\text{Scalars} : \rho = \begin{pmatrix} \rho^{++} \\ \rho^+ \\ \rho^0 \end{pmatrix} \in (1, 3, 1), \quad \eta = \begin{pmatrix} \eta^+ \\ \eta^0 \\ \eta^- \end{pmatrix} \in (1, 3, 0), \quad \chi = \begin{pmatrix} \chi^0 \\ \chi^- \\ \chi^{--} \end{pmatrix} \in (1, 3, -1)$$

Anomaly cancellation between families as $N_C = N_{\text{families}}$; asymmetric 3rd family

Electroweak symmetry breaking : $SU(3)_L \times U(1)_X \rightarrow SU(2)_L \times U(1)_Y \rightarrow U(1)_{\text{em}}$

ρ^0 , η^0 and χ^0 get vevs and give mass to new vector bosons Z' , $Y^{\pm, \pm}$ and Y^\pm

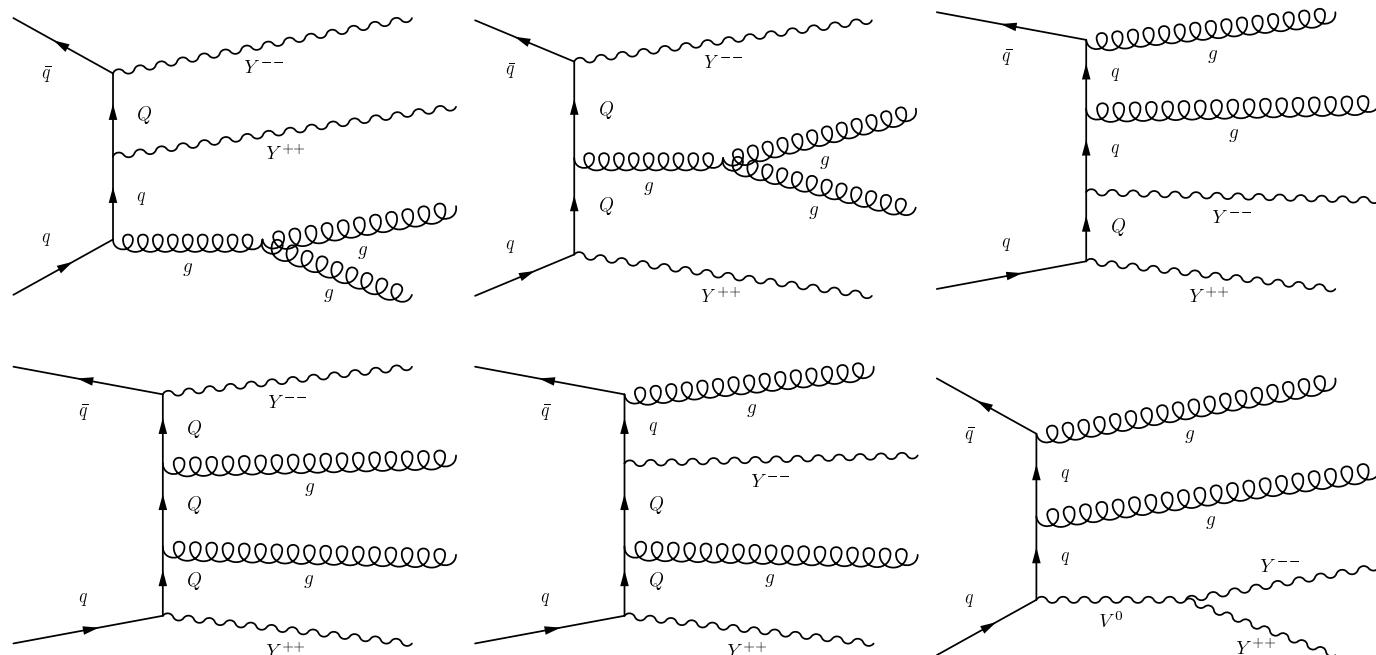
(Y^{++}, Y^+) (Y^{--}, Y^-) bileptons $L_Y = \pm 2$ $m_Y > 250\text{-}550$ GeV from $H^{\pm, \pm}$ Nepomuceno'16

D , S and T : exotic quarks with $L(D, S) = +2$ and $L(T) = -2$

Higgs sector: h_1 , h_2 , h_3 (scalar); a_1 (pseudoscalar); h_1^\pm , h_2^\pm and $h_1^{\pm\pm}$

Benchmark Point (SARAH)		
$m_{h_1} = 125.1 \text{ GeV}$	$m_{h_2} = 3172 \text{ GeV}$	$m_{h_3} = 3610 \text{ GeV}$
$m_{a_1} = 3595 \text{ GeV}$		
$m_{h_1^\pm} = 1857 \text{ GeV}$	$m_{h_2^\pm} = 3590 \text{ GeV}$	
$m_{h_1^{\pm\pm}} = 3734 \text{ GeV}$		
$m_{Y^\pm\pm} = 873.3 \text{ GeV}$	$m_{Y^\pm} = 875.7 \text{ GeV}$	
$m_{Z'} = 3229 \text{ GeV}$		
$m_D = 1650 \text{ GeV}$	$m_S = 1660 \text{ GeV}$	$m_T = 1700 \text{ GeV}$

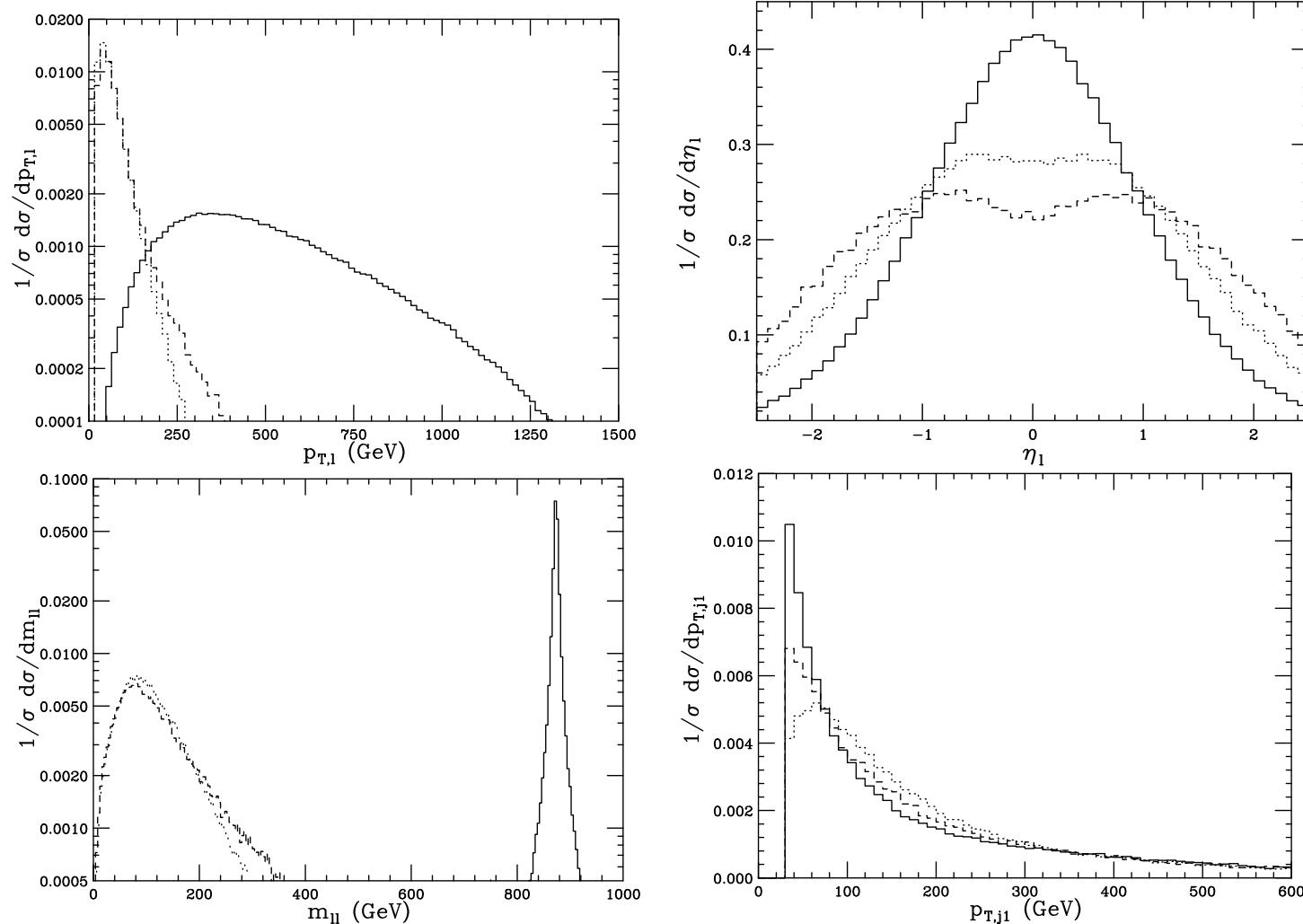
Final states with 2 jets and 2 same-sign leptons: $pp \rightarrow Y^{++}Y^{--}jj \rightarrow (\ell^+\ell^+)(\ell^-\ell^-)jj$



Simulation with MadGraph+HERWIG, $\sqrt{s} = 13$ TeV, k_T algorithm, $R = 1$

$p_{T,j} > 30$ GeV, $p_{T,\ell} > 20$ GeV, $|\eta_j| < 4.5$, $|\eta_\ell| < 2.5$, $\Delta R_{jj} > 0.4$, $\Delta R_{\ell\ell} > 0.1$, $\Delta R_{j\ell} > 0.4$

$\sigma(YYjj) \simeq 3.7$ fb; $\sigma(t\bar{t}Z) \simeq 8.6$ fb; $\sigma(ZZjj) \simeq 6.4$ fb (MET < 100 GeV for $t\bar{t}Z$)



Solid: signal; dashes: $ZZjj$; dots: $t\bar{t}Z$ (Acknowledgement to Lecce ATLAS group)

Dutta-Nandi '94, Dion et al '99, Barreto et al '09-'13, Alves et al '12, Nepomuceno'16: related work on 331 phenomenology

Conclusions and outlook

Investigation on non-standard Z' and bileptons $Y^{\pm\pm}$ at LHC

Supersymmetric Z' modes decrease SM rates; Z' constrains sparticle invariant masses

Discrimination from Z' dilepton decays and other supersymmetric modes is feasible

Z' decays into the lightest neutralinos channel: Dark Matter candidates

$(\Delta m_{Z'})_{\min} \approx 200\text{-}300 \text{ GeV}$ for a reference point in the parameter space

Bilepton signal in final states with 2 jets and 2 same-sign charged-lepton pairs

SM background can be easily suppressed

In progress:

Implementation of Z' leptophobic model to enhance SUSY rates

Investigation of DM signals in mono-X events

QCD effects in production and decays of both Z' and $Y^{\pm,\pm}$ (NLO and resummations)

Investigation of further properties of the bilepton model

Discussion with exotics/SUSY WGs to undertake an actual experimental search