

NLO+PS matching for $t\bar{t}b\bar{b}$ production with massive b -quarks

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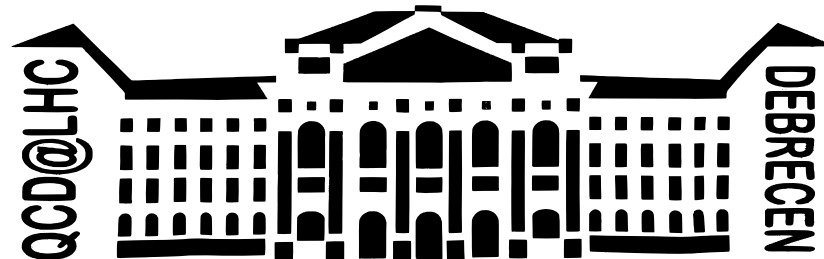
In collaboration with:

J. Lindert, N. Moretti and S. Pozzorini

QCD@LHC, August 29 2017



**Universität
Zürich**^{UZH}



NLO+PS matching for $t\bar{t}b\bar{b}$ production with massive b -quarks

- Summarize results of $t\bar{t}H$ search at the LHC
- Discuss various ways of simulating $t\bar{t}b\bar{b}$
- Present a LO study on the relative importance of IS and FS $g \rightarrow b\bar{b}$ splittings
- Review existing NLO+PS tools for $t\bar{t}b\bar{b}$ and compare them
- Present our results

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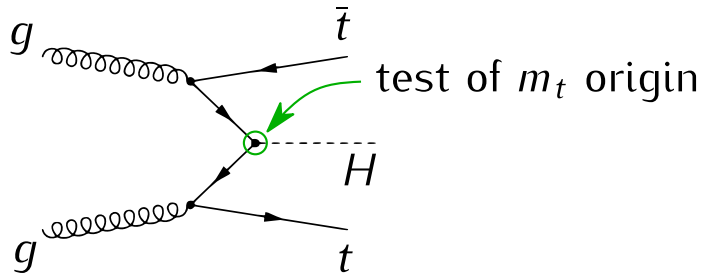
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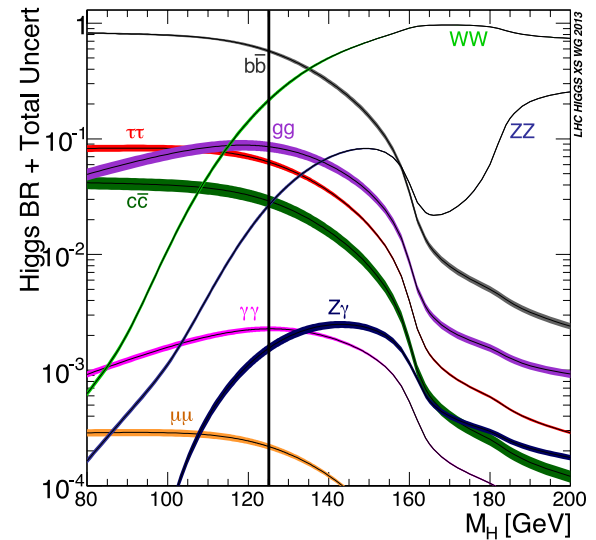
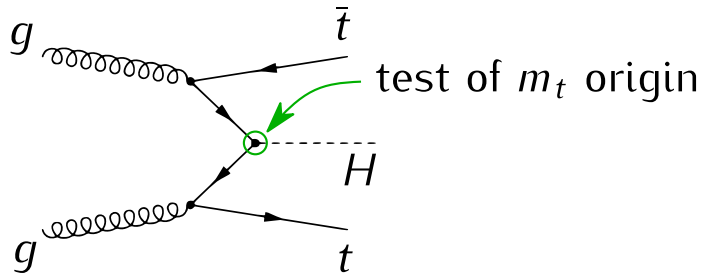
$t\bar{t}H$ @ LHC run 2

- Direct probe of top-quark Yukawa coupling



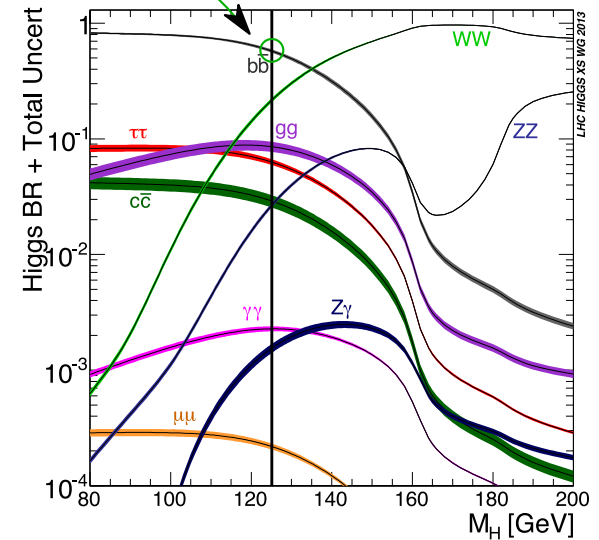
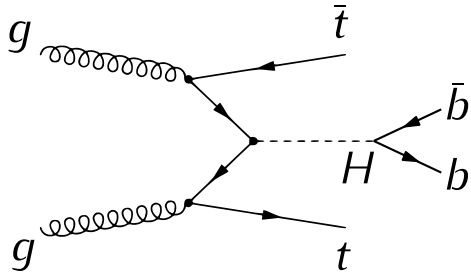
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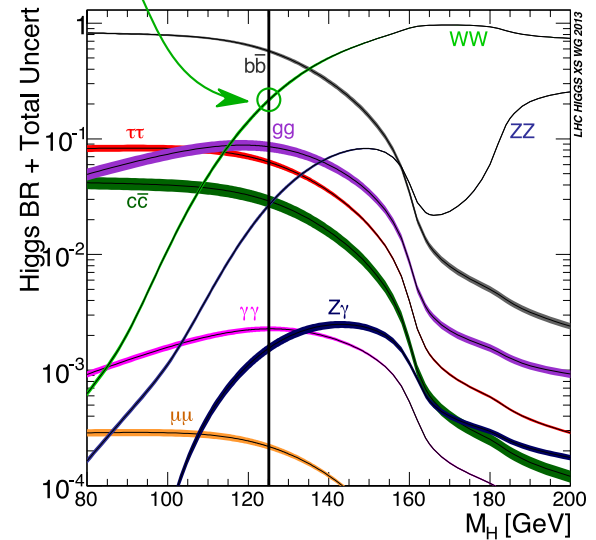


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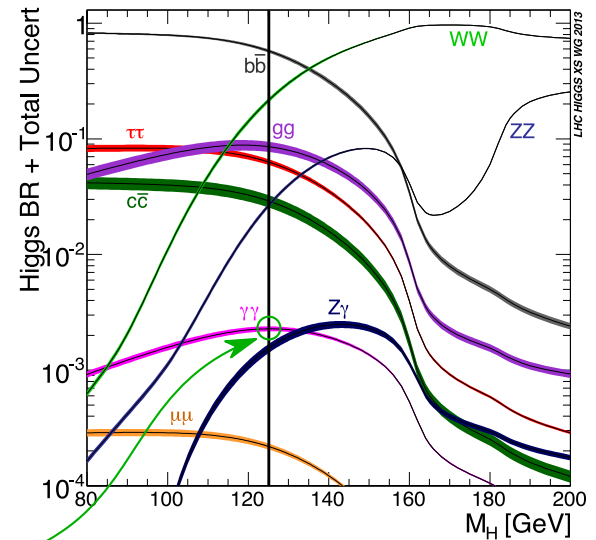
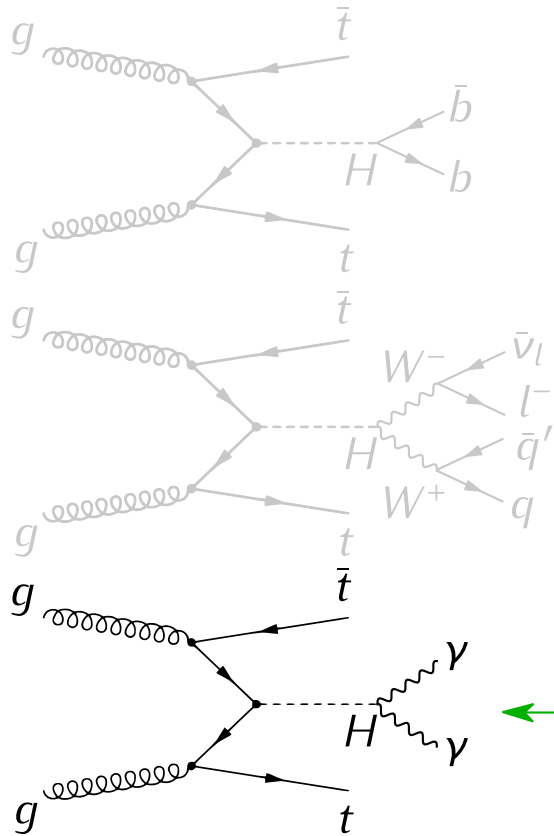


-
- The figure contains two Feynman diagrams. The top diagram illustrates the production of a Higgs boson (H) through gluon fusion (gg) and its subsequent decay into two photons ($\gamma\gamma$) via a top quark loop (t). The bottom diagram illustrates the production of a Higgs boson (H) through gluon fusion (gg) and its subsequent decay into two photons ($\gamma\gamma$) via a top quark loop (t) and a W boson loop (W).



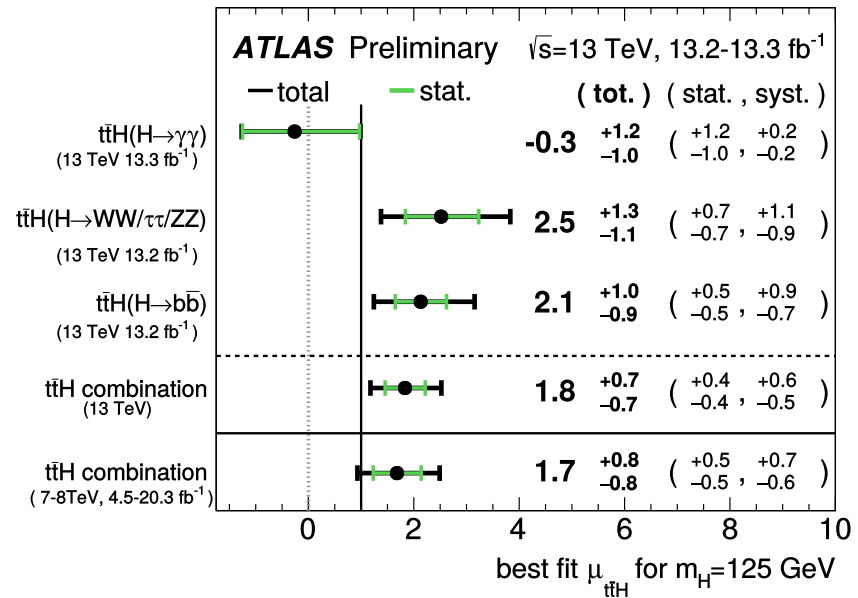
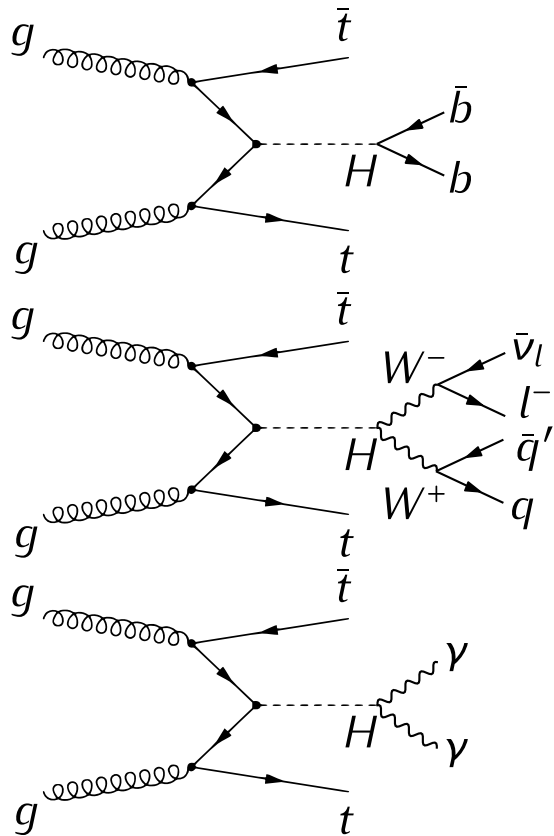
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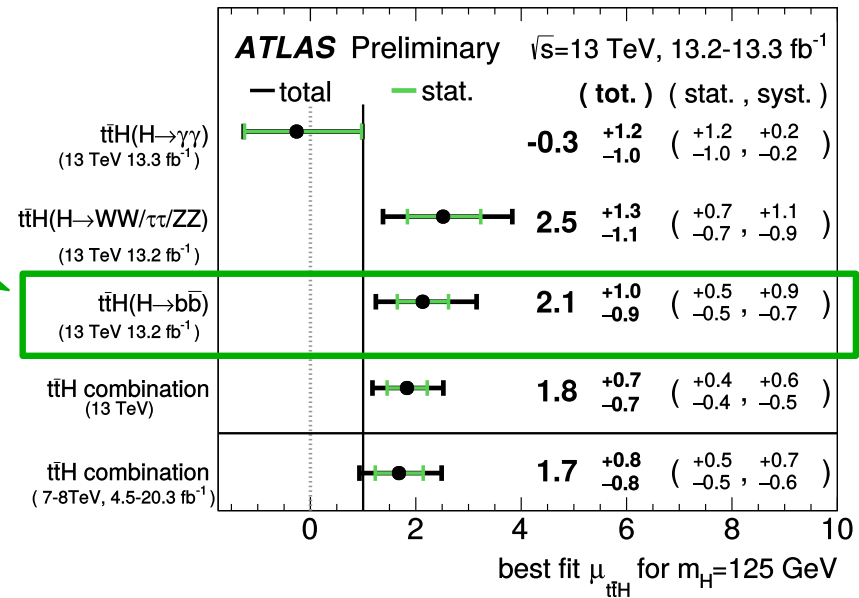
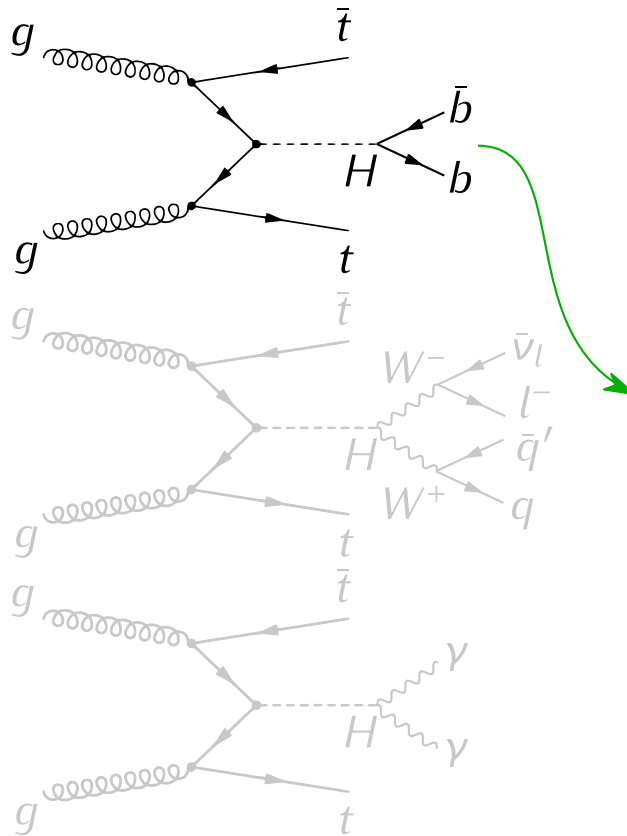
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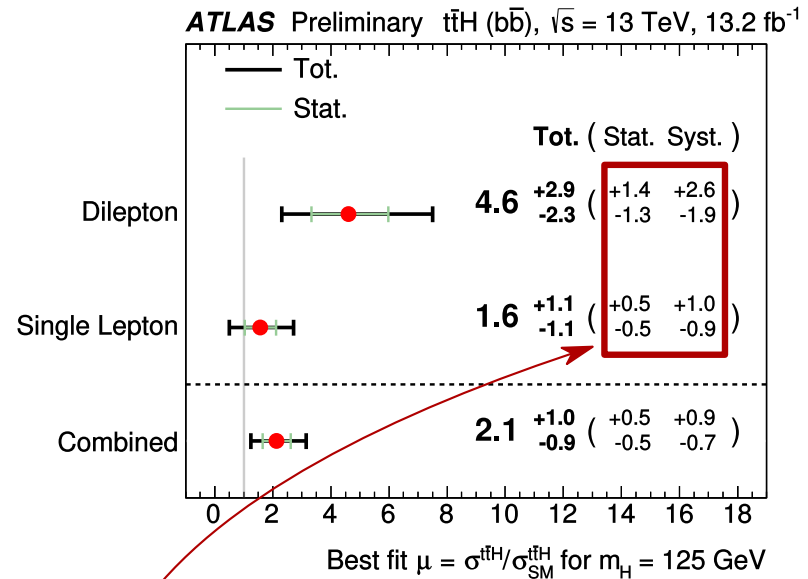
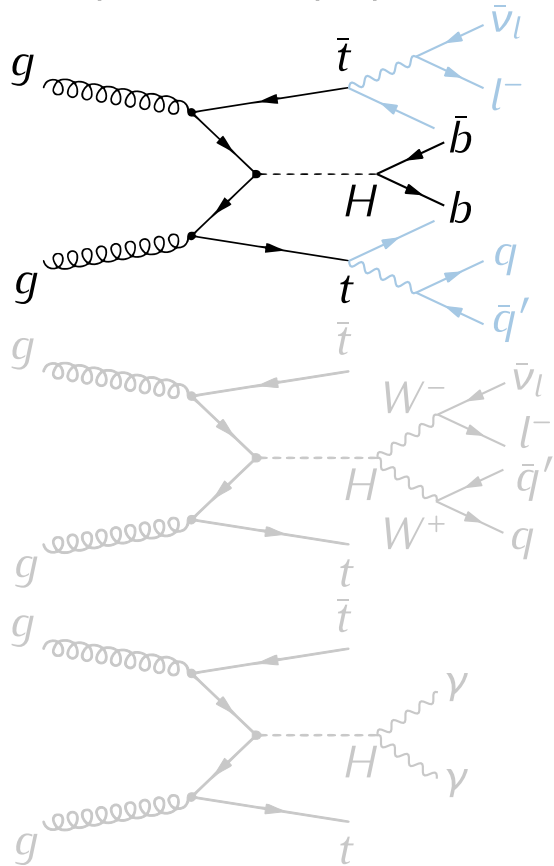
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$t\bar{t}H$ @ LHC run 2

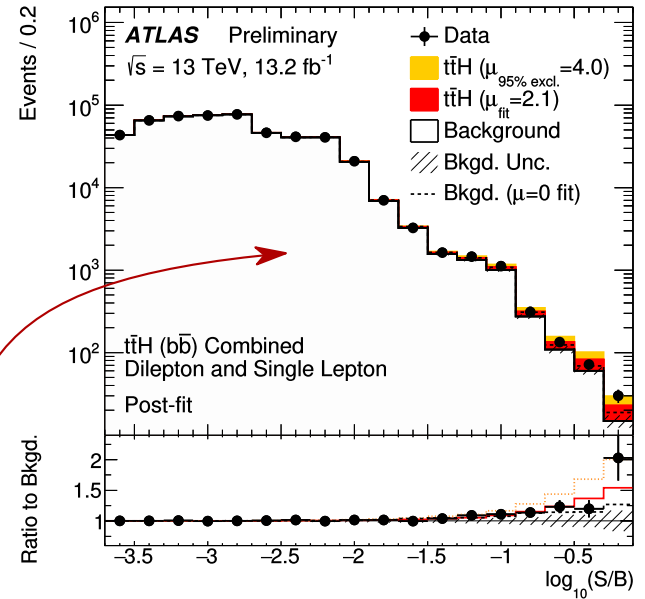
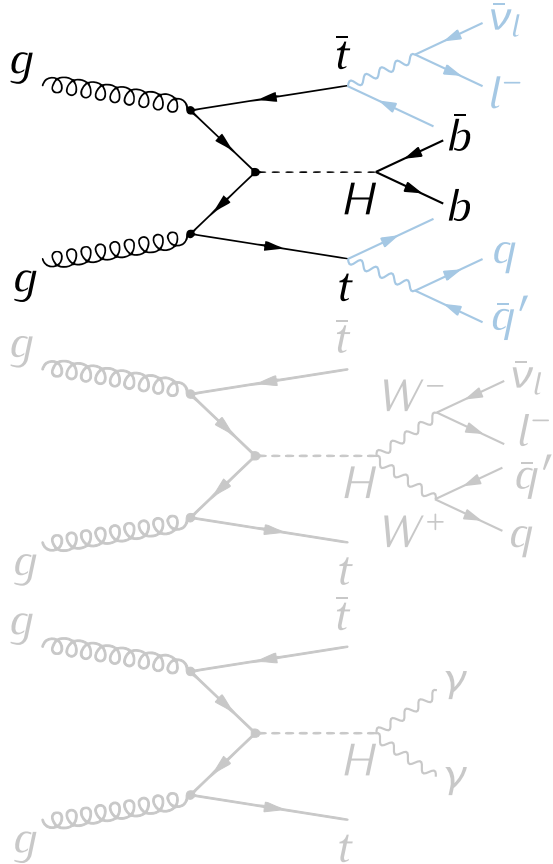
- Direct probe of top-quark Yukawa coupling



Dominated by systematics!

$t\bar{t}H$ @ LHC run 2

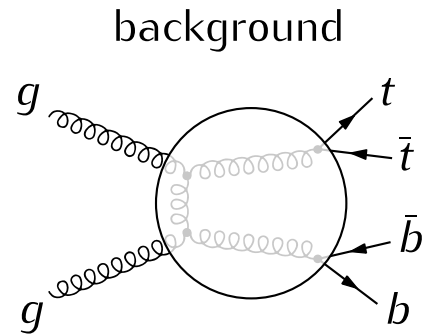
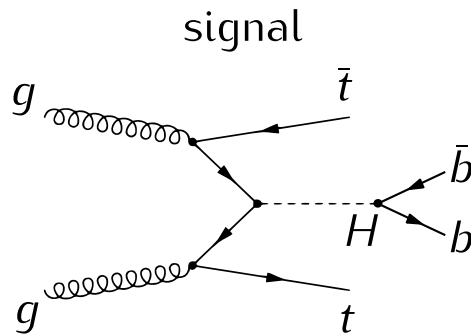
- Direct probe of top-quark Yukawa coupling



Background dominated!

$t\bar{t}H$ @ LHC run 2

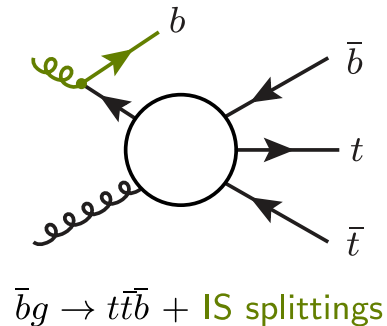
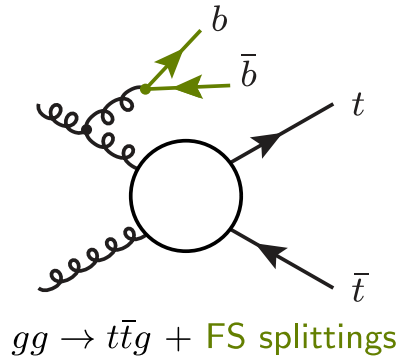
- Large $t\bar{t} + b$ -jets background and its theory uncertainties are bottleneck of $t\bar{t}H(b\bar{b})$ searches



- Modern tools support automated $t\bar{t}b\bar{b}$ simulations, but it remains highly nontrivial multi-particle multi-scale process
- Realistic estimates of theory uncertainties necessitate understanding of dynamics governing $pp \rightarrow t\bar{t}b\bar{b}$ as well as technical aspects related to:
 - ▶ 5F/4F scheme choice
 - ▶ NLO+PS matching
 - ▶ PS effects

How to simulate $t\bar{t} + b$ -jets?

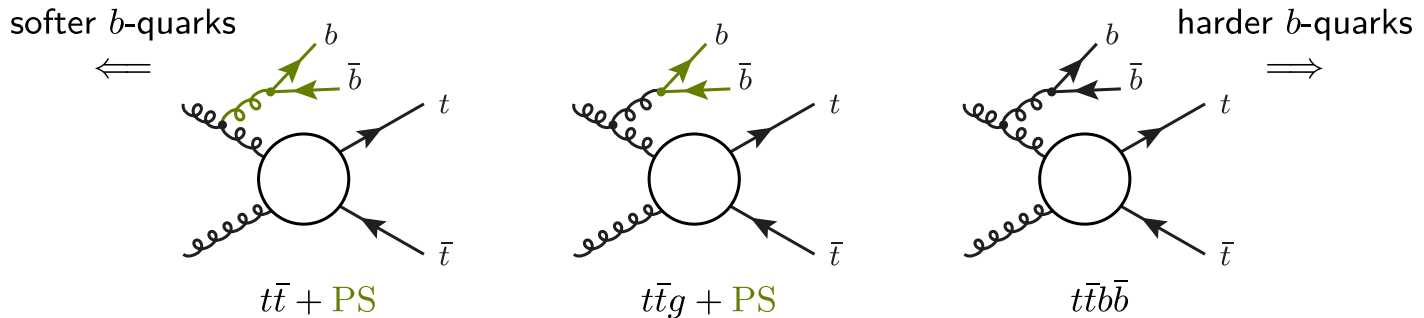
- Option 1: NLO+PS $t\bar{t}$ 5F
 - ▶ $t\bar{t}j$ tree MEs + $g \rightarrow b\bar{b}$ shower splittings



- ▶ Not even LO precision (although PS allows for accurate tuning to data)
- ▶ Description based on $t\bar{t}b\bar{b}$ MEs crucial for realistic theory uncertainty estimates

How to simulate $t\bar{t} + b$ -jets?

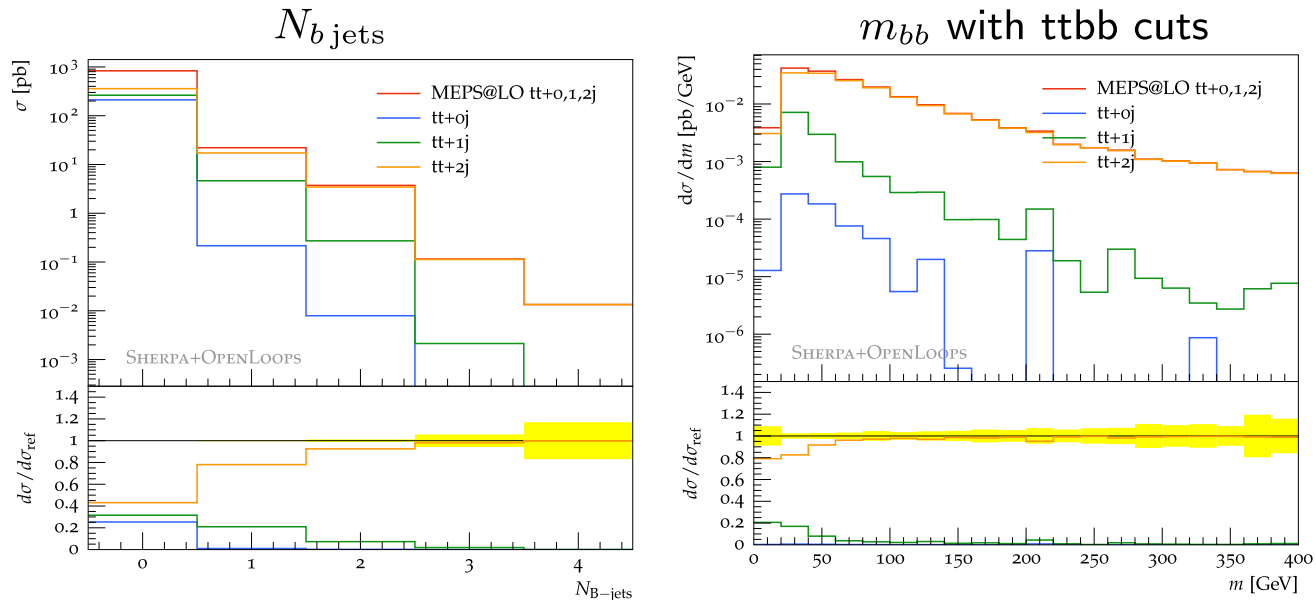
- Option 1: NLO+PS $t\bar{t}$ 5F ... insufficient precision
- Option 2: (N)LO merging $t\bar{t} + 0, 1, 2$ jets 5F
 - ▶ $t\bar{t} + 0, 1, 2$ jet MEs and $g \rightarrow b\bar{b}$ splittings



- ▶ Precision and CPU cost strongly dependent on the merging cut Q_{cut}
- ▶ Does this describe $t\bar{t} + b$ -jets mostly through $t\bar{t}b\bar{b}$ MEs though?

Amount of $t\bar{t}$ +jets ME information

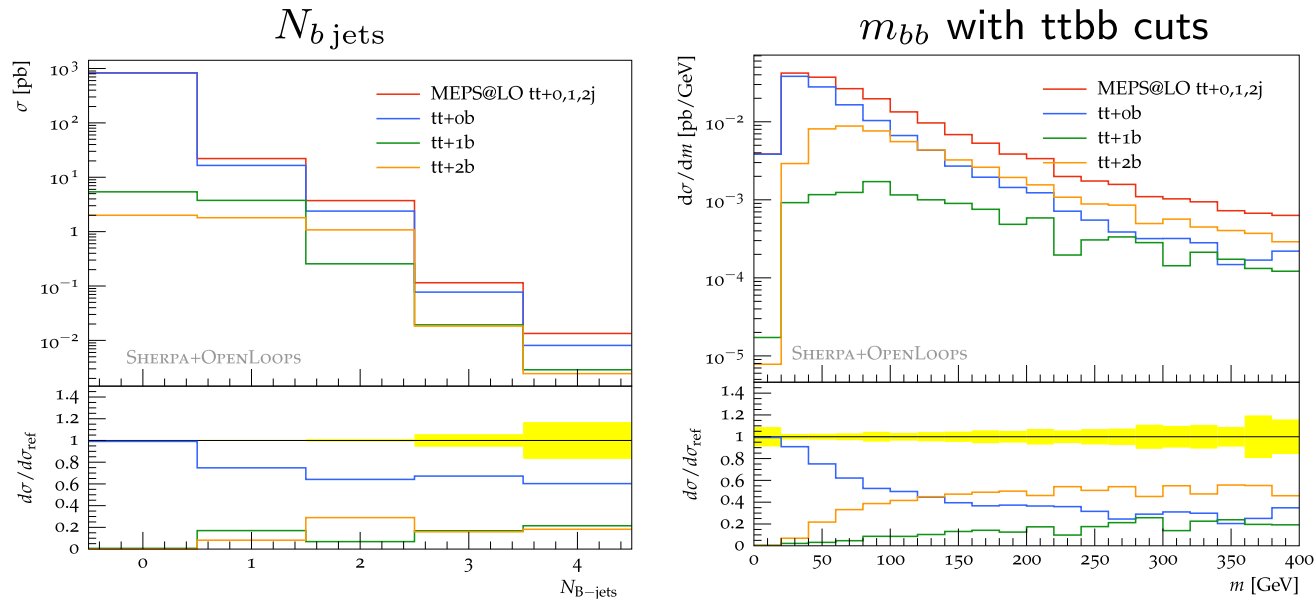
- $t\bar{t} + 0, 1, 2$ jet LO merging with $Q_{\text{cut}} = 20$ GeV



- Observables with ≥ 1 additional b -jets
 - dominated by $t\bar{t} + 2\text{jet}$ MEs (suggesting ME precision)

Amount of $t\bar{t}+b$ -jets ME information

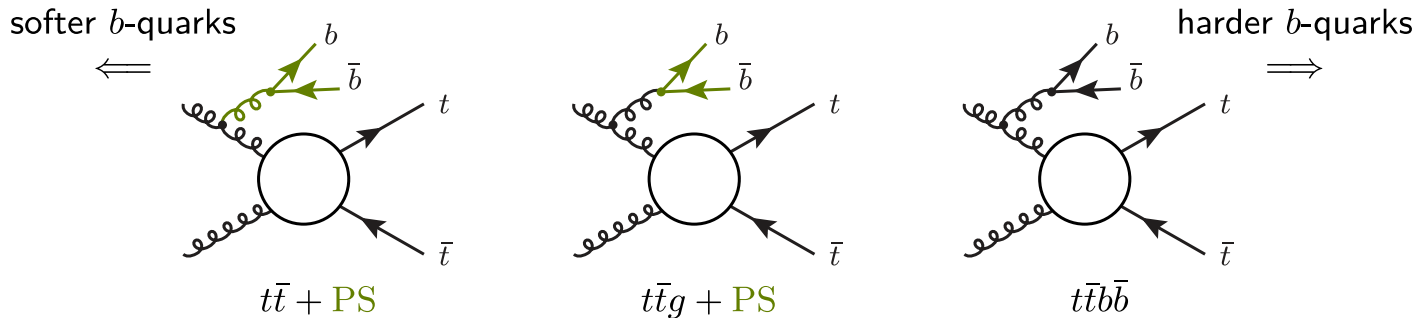
- $t\bar{t} + 0, 1, 2$ jet LO merging with $Q_{\text{cut}} = 20$ GeV



- Observables with ≥ 1 additional b -jets
 - ▶ actually dominated by MEs with 2 light jets and no b -jets (up to $Q \sim 100$ GeV)!

How to simulate $t\bar{t} + b$ -jets?

- Option 1: NLO+PS $t\bar{t}$ 5F ... **insufficient precision**
- Option 2: (N)LO merging $t\bar{t} + 0, 1, 2$ jets 5F
 - ▶ $t\bar{t} + 0, 1, 2$ jet MEs and $g \rightarrow b\bar{b}$ splittings



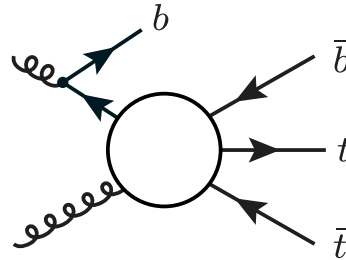
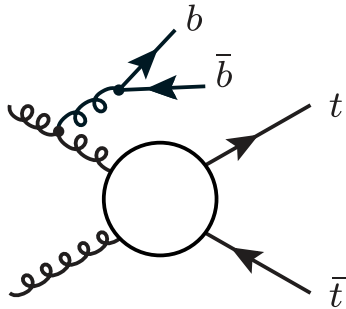
- ▶ Precision and CPU cost strongly dependent on the merging cut Q_{cut}
- ▶ Does this describe $t\bar{t} + b$ -jets mostly through $t\bar{t}b\bar{b}$ MEs though?

No!

Direct description in terms of $t\bar{t}b\bar{b}$ MEs preferable.

How to simulate $t\bar{t} + b$ -jets?

- Option 1: NLO+PS $t\bar{t}$ 5F ... insufficient precision
- Option 2: (N)LO merging $t\bar{t} + 0, 1, 2$ jets 5F ... observables with ≥ 1 additional b -jets dominated by MEs with 2 light jets
- Option 3: $t\bar{t}b\bar{b}$ at NLO+PS

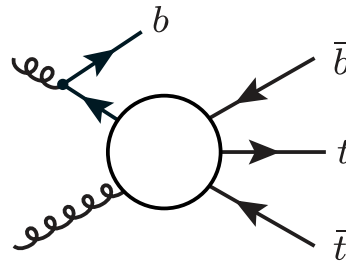
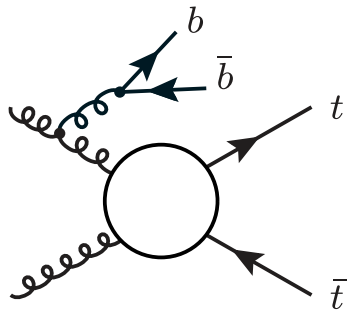


- NLO+PS precision for $t\bar{t} + 2b$ -jet and $t\bar{t} + 1b$ -jet observables

How to simulate $t\bar{t} + b$ -jets?

- Option 1: NLO+PS $t\bar{t}$ 5F ... insufficient precision
- Option 2: (N)LO merging $t\bar{t} + 0, 1, 2$ jets 5F ... observables with ≥ 1 additional b -jets dominated by MEs with 2 light jets

- Option 3: $t\bar{t}b\bar{b}$ at NLO+PS

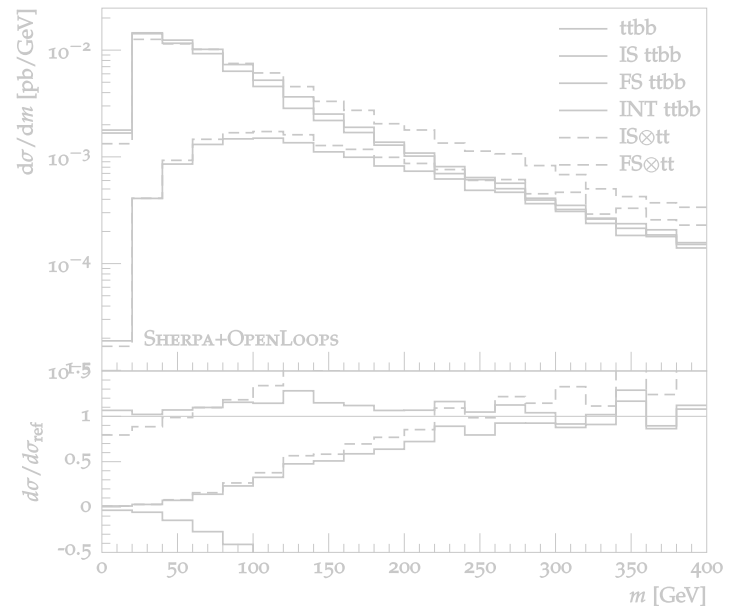
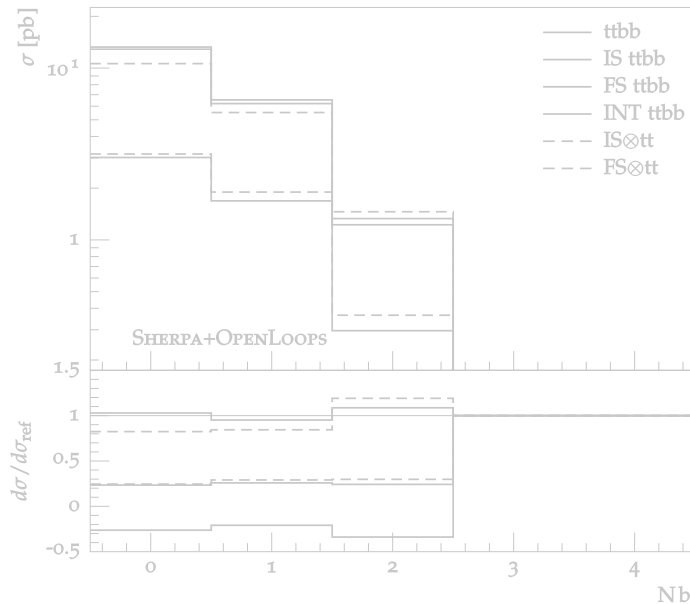


- NLO+PS precision for $t\bar{t} + 2b$ -jet and $t\bar{t} + 1b$ -jet observables

 This talk

QCD production of $t\bar{t}b\bar{b}$

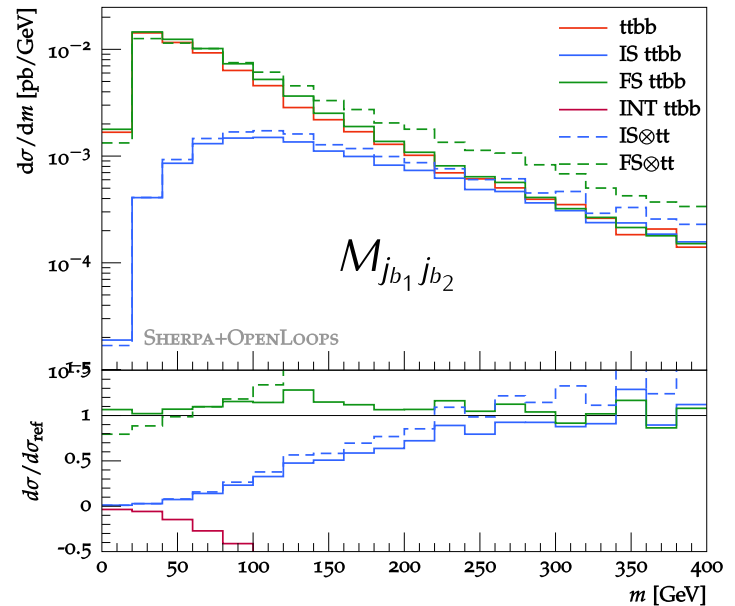
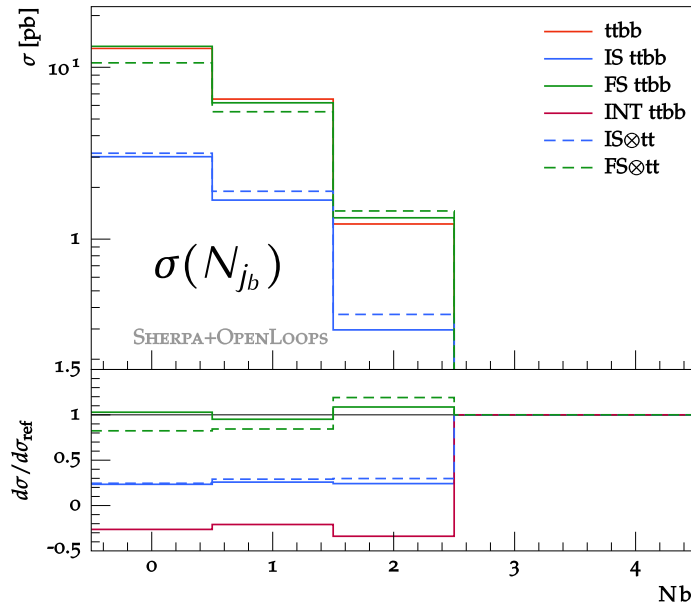
- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:
 - ▶ 6 external coloured partons, $\sigma_{t\bar{t}b\bar{b}} \propto \alpha_S^4(\mu_R)$
 - ▶ 34 LO diagrams, multiple scales from 5 to 500 GeV
 - ▶ Dominated by topologies with FS $g \rightarrow b\bar{b}$ splittings



QCD production of $t\bar{t}b\bar{b}$

- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:

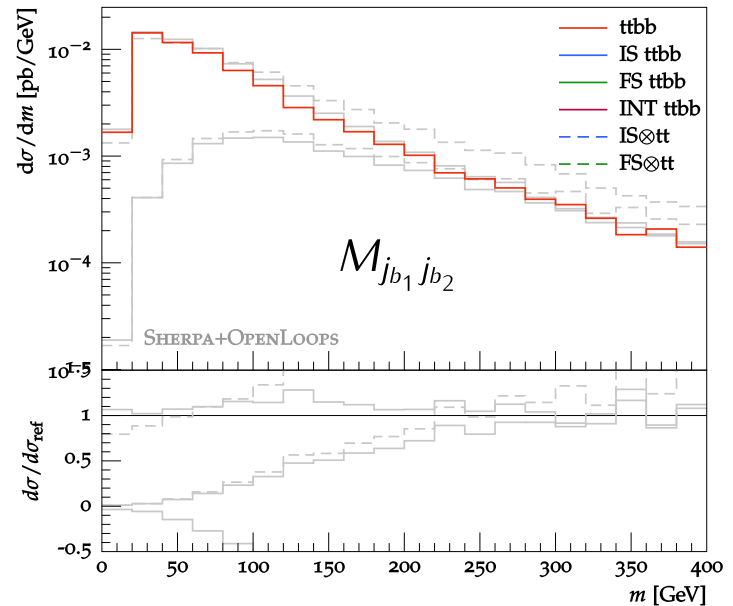
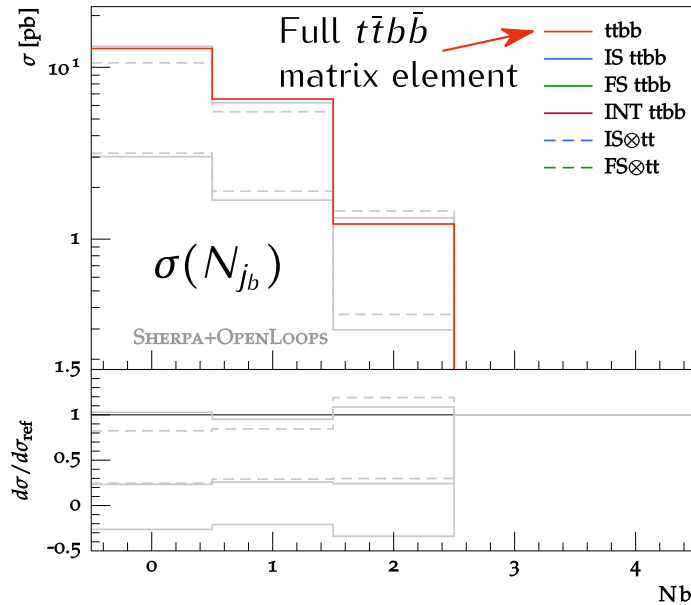
► Dominated by topologies with FS $g \rightarrow b\bar{b}$ splittings



QCD production of $t\bar{t}b\bar{b}$

- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:

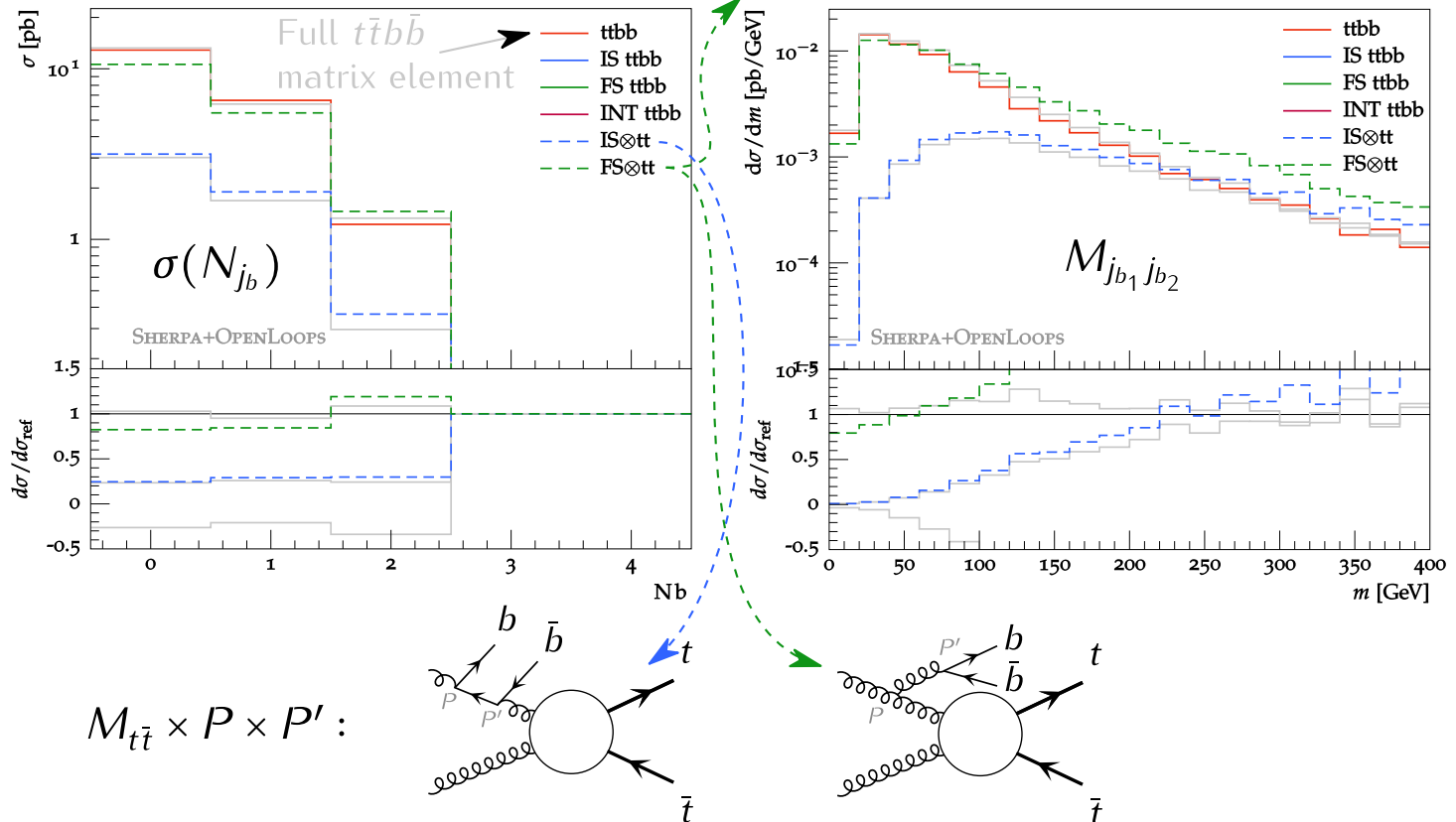
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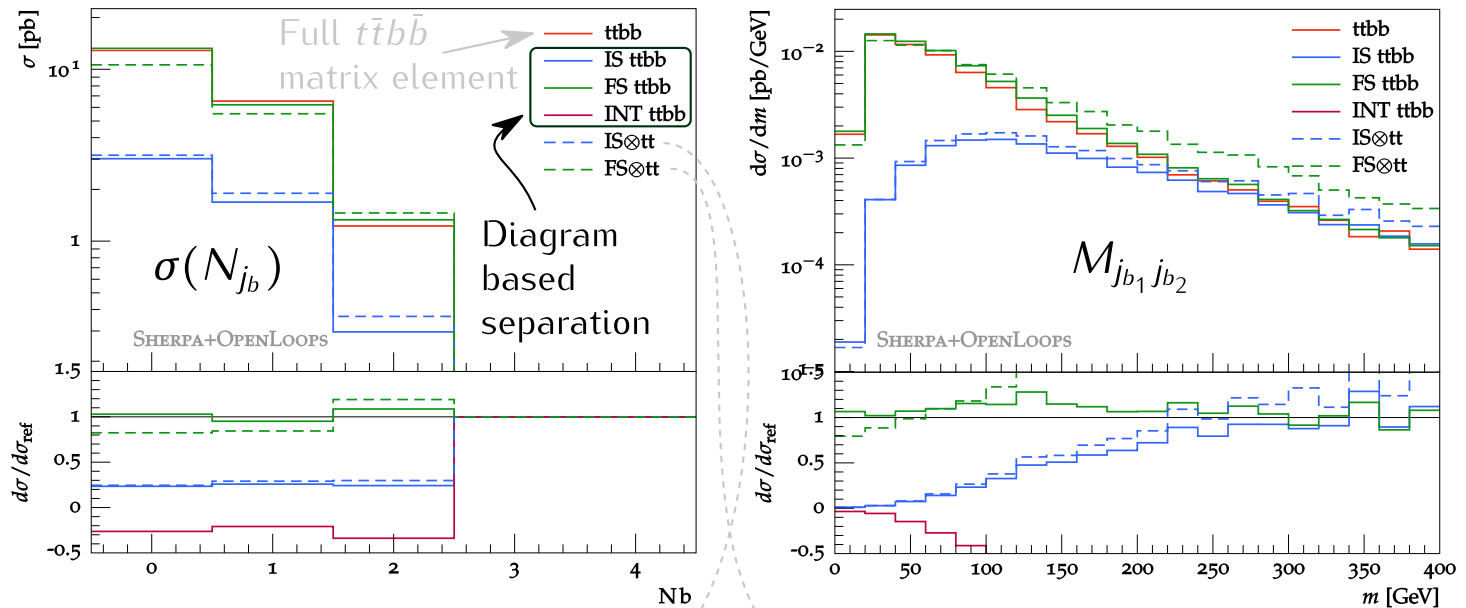
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QCD production of $t\bar{t}b\bar{b}$

- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:

- Dominated by topologies with FS $g \rightarrow b\bar{b}$ splittings



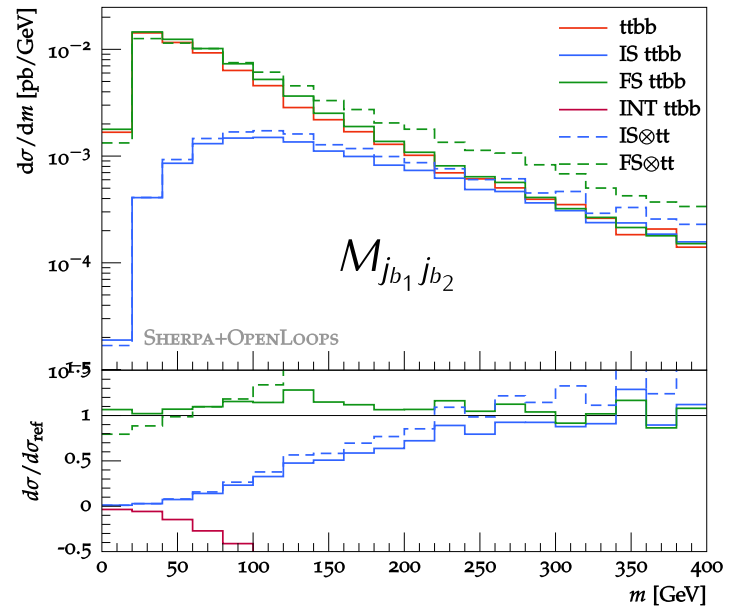
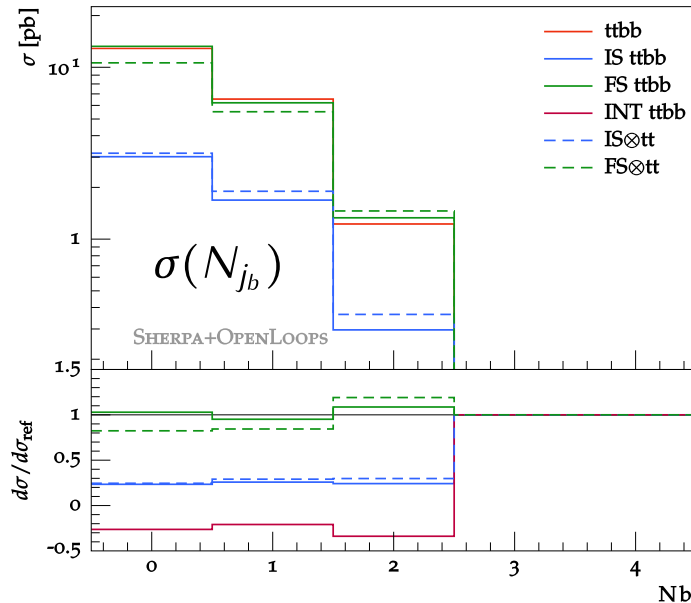
$$M_{t\bar{t}} \times P \times P' :$$



QCD production of $t\bar{t}b\bar{b}$

- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:

► Dominated by topologies with FS $g \rightarrow b\bar{b}$ splittings

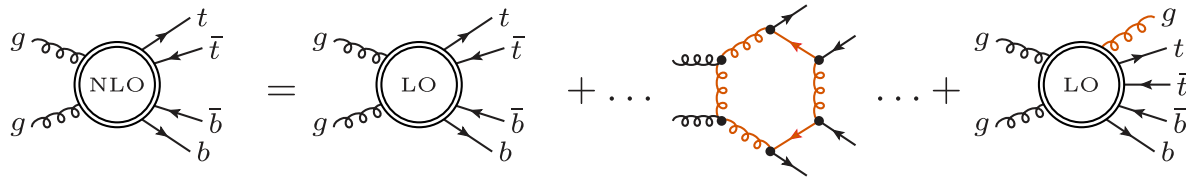


- FS $g \rightarrow b\bar{b}$ dominant, also away from collinear regime
- IS $g \rightarrow b\bar{b}$ subdominant (no need for 5F resummation)

supports choice of 4F scheme with $m_b > 0$ and no b -quark PDF

QCD production of $t\bar{t}b\bar{b}$ @NLO

- $t\bar{t}b\bar{b}$ @ NLO QCD:

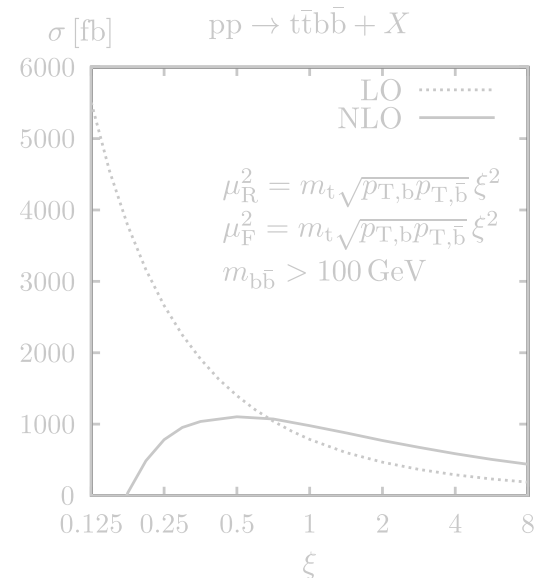


- ▶ 5FNS ($m_b = 0$): [Bredenstein et al. '09-'10; Bevilacqua et al. '10]
- ▶ 4FNS ($m_b > 0$): [Cascioli et al. '13]

- $\sigma_{t\bar{t}b\bar{b}} \propto \alpha_S^4(\mu_R) \Rightarrow$ scale uncertainty:

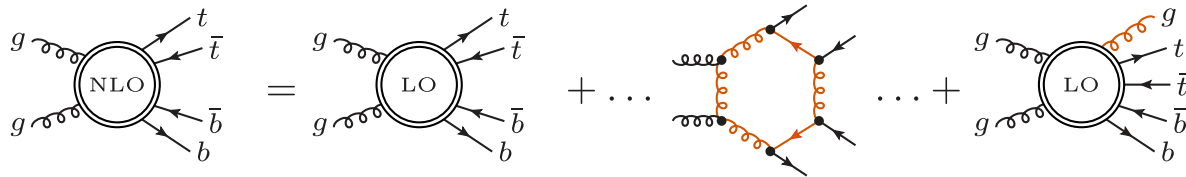
- ▶ $\sim 80\%$ @ LO
- ▶ 20 – 30% @ NLO

- NLO+PS predictions mandatory for realistic analysis

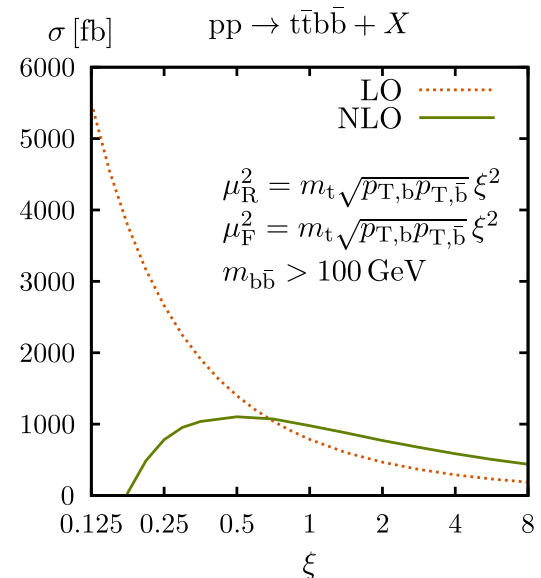


QCD production of $t\bar{t}b\bar{b}$ @NLO

- $t\bar{t}b\bar{b}$ @ NLO QCD:



- ▶ 5FNS ($m_b = 0$): [Bredenstein et al. '09-'10; Bevilacqua et al. '10]
- ▶ 4FNS ($m_b > 0$): [Cascioli et al. '13]
- $\sigma_{t\bar{t}b\bar{b}} \propto \alpha_S^4(\mu_R) \Rightarrow$ scale uncertainty:
 - ▶ $\sim 80\%$ @ LO
 - ▶ $20 - 30\%$ @ NLO
- NLO+PS predictions mandatory for realistic analysis



QCD production of $t\bar{t}b\bar{b}$ @NLO+PS



- Available $t\bar{t}b\bar{b}$ calculations @NLO+PS:
 - ▶ Powhe1 [Garzelli et al. '13/'14]
 - ▷ POWHEG matching
 - ▷ 5F scheme, $m_b = 0$
 - ▷ requires a generation cut
 - ▶ Sherpa+OpenLoops [Cascioli et al. '13]
 - ▷ S-MC@NLO matching
 - ▷ 4F scheme, $m_b > 0$
 - ▶ POWHEG-BOX+OpenLoops [upcoming]
 - ▷ POWHEG matching
 - ▷ 4F scheme, $m_b > 0$

QCD production of $t\bar{t}b\bar{b}$ @NLO+PS

- Available $t\bar{t}b\bar{b}$ calculations @NLO+PS:

- ▶ Powhe1 [Garzelli et al. '13/'14]

- POWHEG matching
- 5F scheme, $m_b = 0$
- requires a generation cut

MEs cannot describe quasi-collinear $g \rightarrow b\bar{b}$ splittings

- ▶ Sherpa+OpenLoops [Cascioli et al. '13]

- S-MC@NLO matching
- 4F scheme, $m_b > 0$

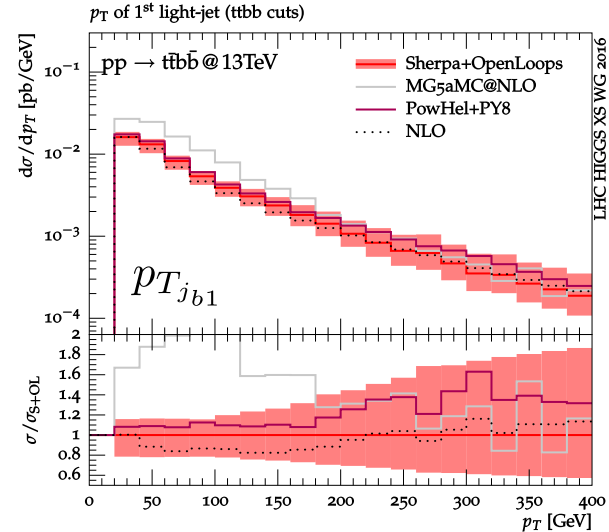
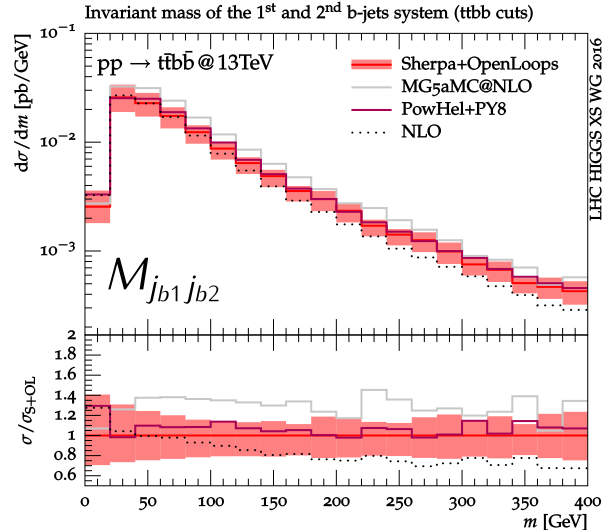
- ▶ POWHEG-BOX+OpenLoops [upcoming]

- POWHEG matching
- 4F scheme, $m_b > 0$

MEs cover full b -quark phase space

Why another $t\bar{t}b\bar{b}$ @NLO+PS?

- YR4 [[arXiv:1610.07922](https://arxiv.org/abs/1610.07922)]:



- Sherpa+OpenLoops vs. PowHel+PY8

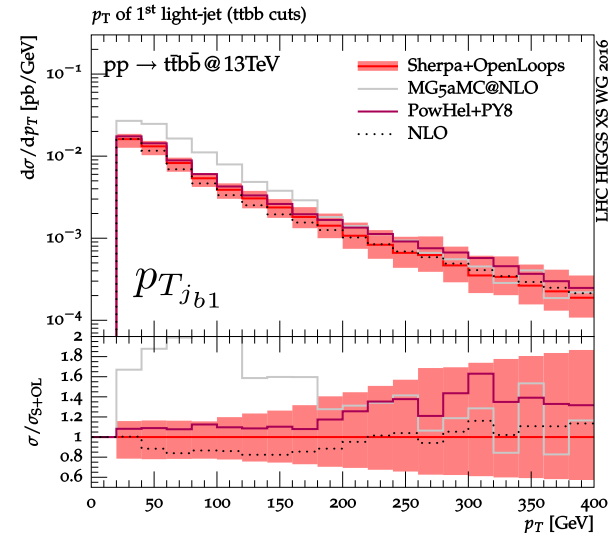
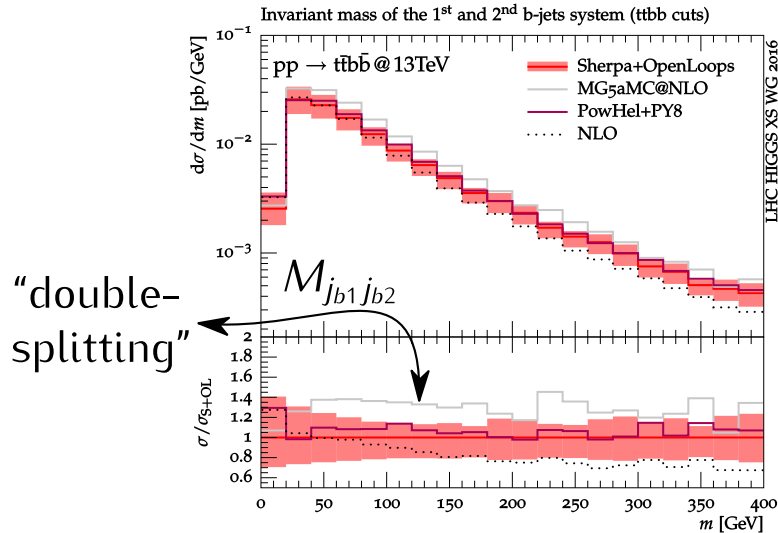
- ▶ Good agreement also in observables with large NLO+PS corrections

- Sherpa+OpenLoops vs. MG5_aMC@NLO+PY8 [[arXiv:1405.0301](https://arxiv.org/abs/1405.0301)]

- ▶ Sizable differences in NLO radiation pattern
- ▶ Strong resummation-scale sensitivity of $t\bar{t}b\bar{b}$ +jet in MG5_aMC@NLO+PY8

Why another $t\bar{t}b\bar{b}$ @NLO+PS?

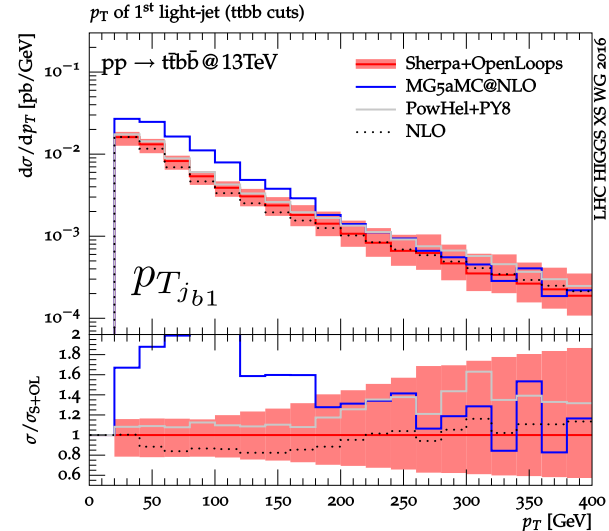
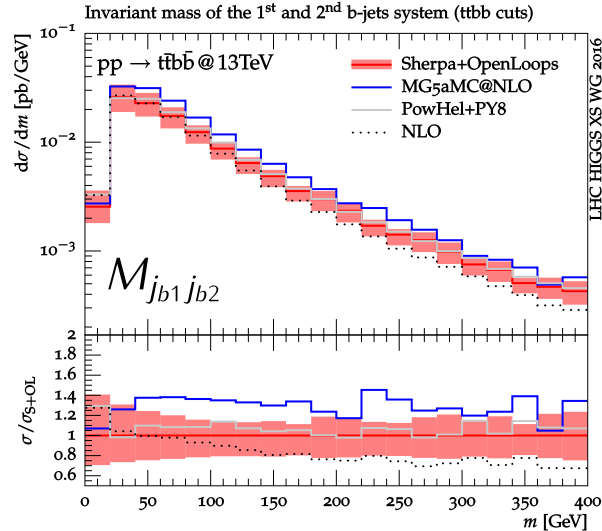
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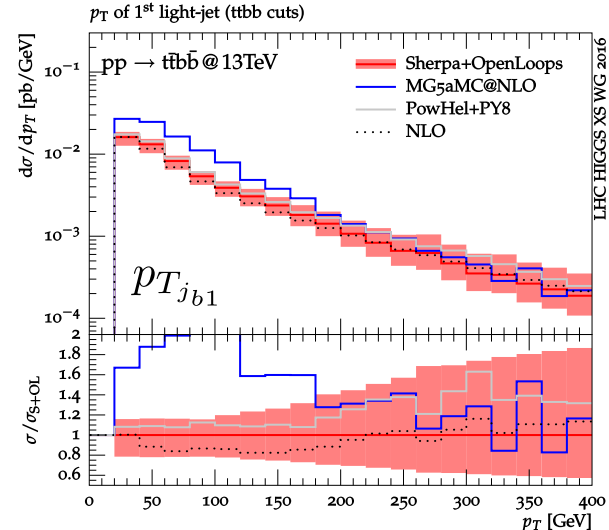
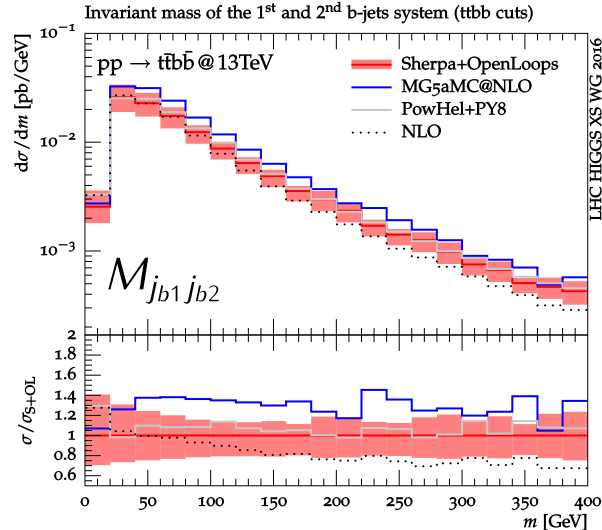
- ▶ Good agreement also in observables with large NLO+PS corrections

- Sherpa+OpenLoops vs. MG5_aMC@NLO+PY8 [[arXiv:1405.0301](https://arxiv.org/abs/1405.0301)]

- ▶ Sizable differences in NLO radiation pattern
- ▶ Strong resummation-scale sensitivity of $t\bar{t}b\bar{b}$ +jet in MG5_aMC@NLO+PY8

Why another $t\bar{t}b\bar{b}$ @NLO+PS?

- YR4 [[arXiv:1610.07922](https://arxiv.org/abs/1610.07922)]:



- Sherpa+OpenLoops vs. PowHel+PY8

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- Sherpa+OpenLoops vs. MG5_aMC@NLO+PY8 [[arXiv:1405.0301](https://arxiv.org/abs/1405.0301)]

- ▶ Sizable differences in NLO radiation pattern
- ▶ Strong resummation-scale sensitivity of $t\bar{t}b\bar{b}$ +jet in MG5_aMC@NLO+PY8
- ▶ New: MG5_aMC@NLO+HW++ in good agreement with Sherpa+OpenLoops

Why another $t\bar{t}b\bar{b}$ @NLO+PS?

- Sherpa+OpenLoops vs. MG5_aMC@NLO+PY8
 - ▶ Sizable differences in NLO radiation pattern
 - ▶ Considerable resummation-scale sensitivity in MG5_aMC@NLO+PY8
 - ▶ Strong shower dependence in MG5_aMC@NLO (PY8 vs HW++)

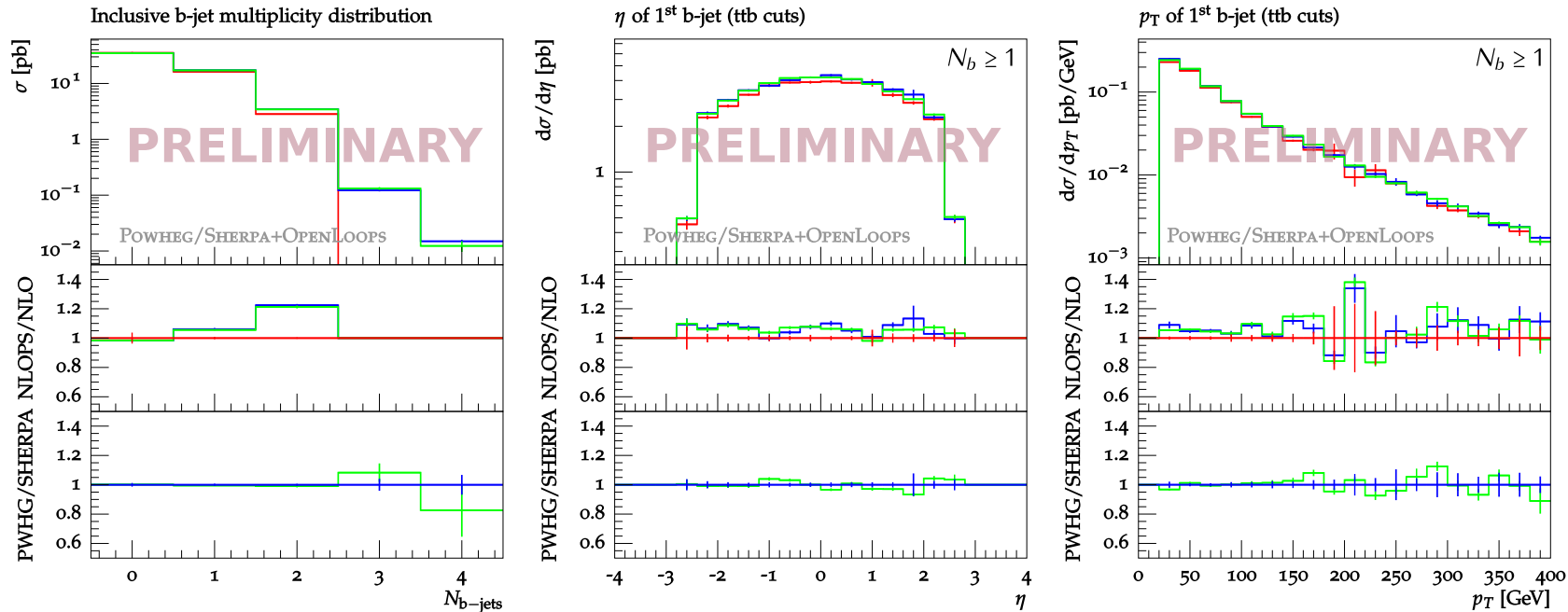


Surprisingly large matching/shower uncertainty? Issue in either Sherpa or MG5_aMC@NLO?

- How about we try a different matching method?
 - ▶ NLO+PS matching with POWHEG BOX (no resum.-scale dependence)
 - ▶ Matrix elements from OpenLoops
 - ▶ 4F scheme: $m_b > 0$ and no b PDF
 - ▶ Compare against Sherpa and study hdamp and shower dependence

POWHEG BOX vs SHERPA

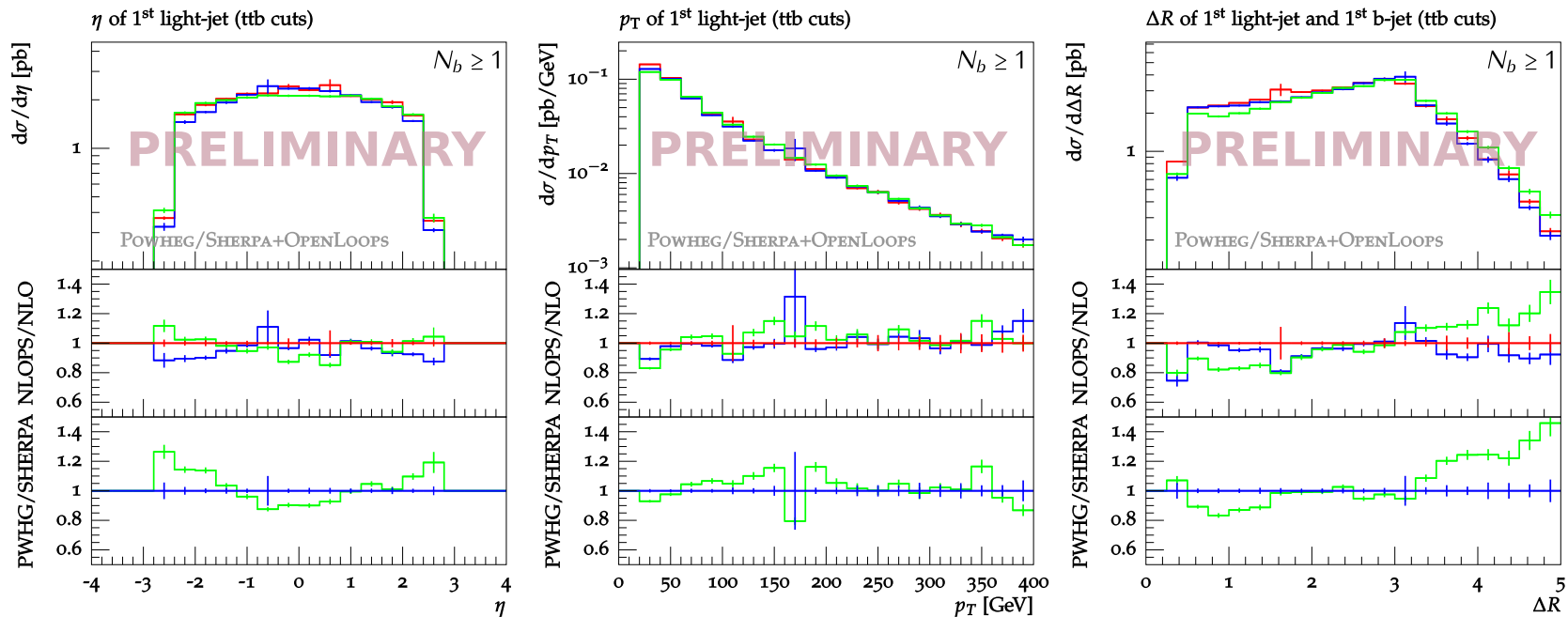
—+— POWHG NLO —+— POWHG NLOPS —+— SHERPA NLOPS



- Remarkable agreement for NLO accurate ttb observables
 - ▶ Agreement well under 5%; expected scale uncertainty $\sim 20\%$
- Good agreement also in LOPS accurate bins 3 and 4 of $\sigma(N_{b\text{-jets}})$

POWHEG BOX vs SHERPA

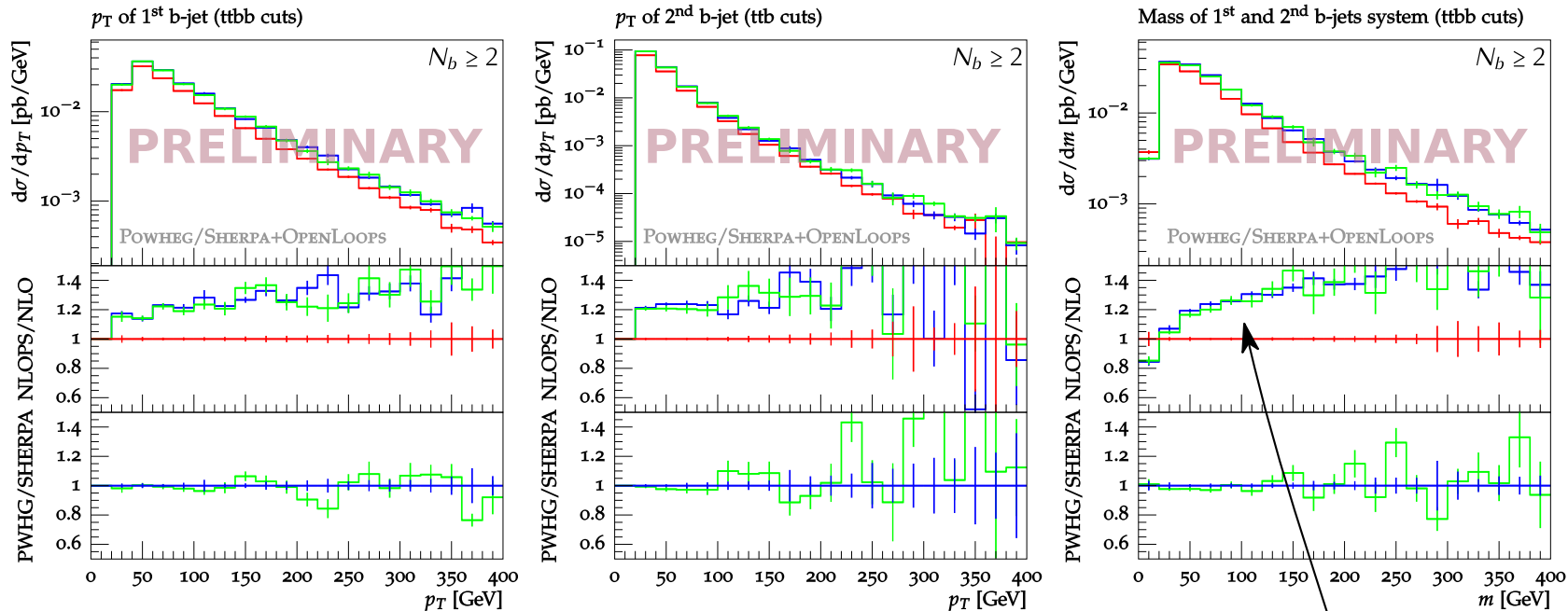
+ POWHEG NLO
 + POWHEG NLOPS
 + SHERPA NLOPS



- Good agreement for LOPS accurate $ttbj$ observables
 - Agreement to $\sim 20\%$; expected scale uncertainty $\sim 50\%$

POWHEG BOX vs SHERPA

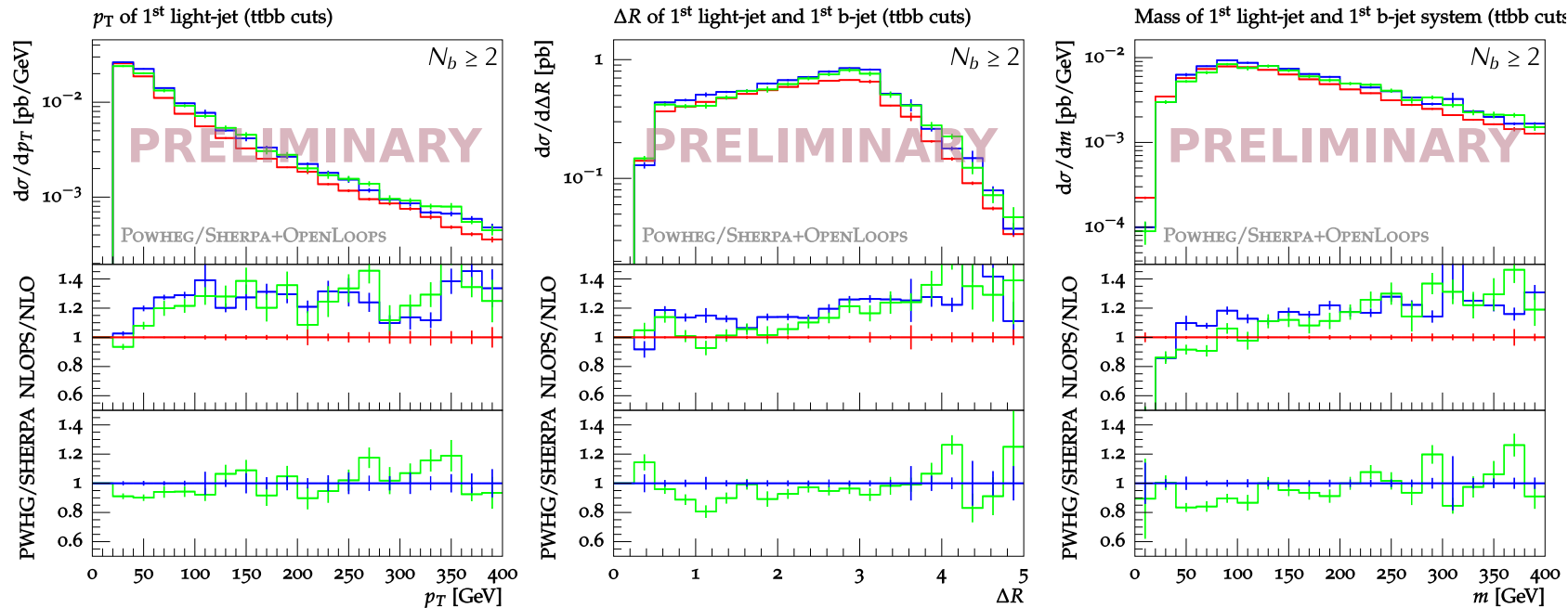
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- Remarkable agreement for NLO accurate ttbb observables
 - ▶ Agreement well under 5%; expected scale uncertainty $\sim 20\%$
- POWHEG BOX RES confirms the “double splitting” enhancement

POWHEG BOX vs SHERPA

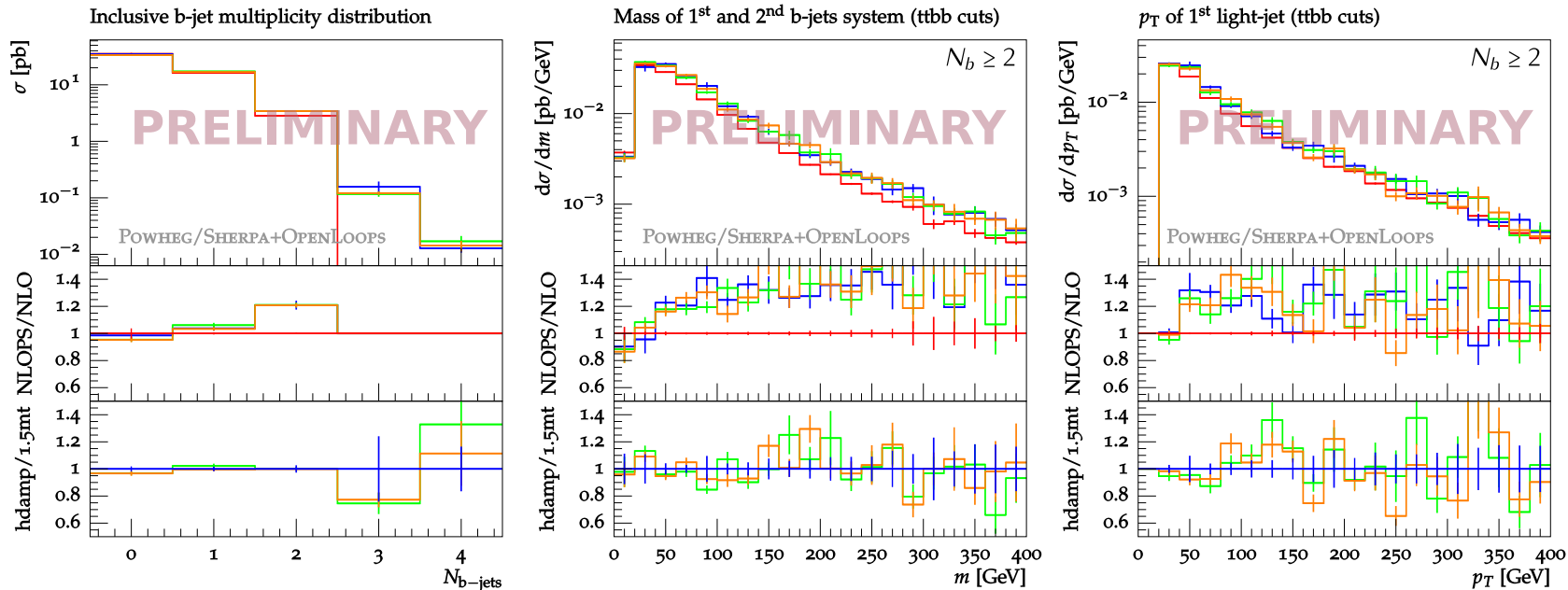
—+— POWHG NLO —+— POWHG NLOPS —+— SHERPA NLOPS



- Good agreement for LOPS accurate ttbbj observables
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hdamp dependence

+ NLO
 + NLOPS $\text{hdamp} = 1.5m_t$
+ NLOPS $\text{hdamp} = 5m_t$
+ NLOPS $\text{hdamp} = 20m_t$

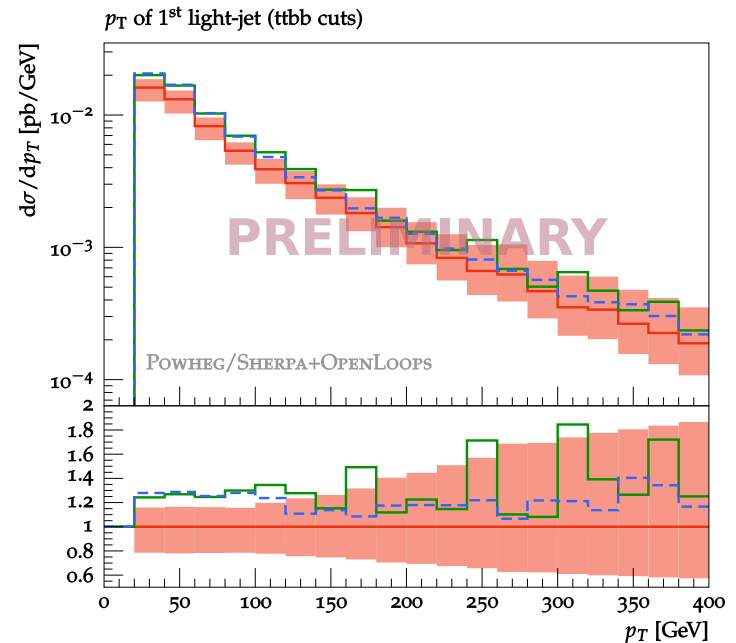
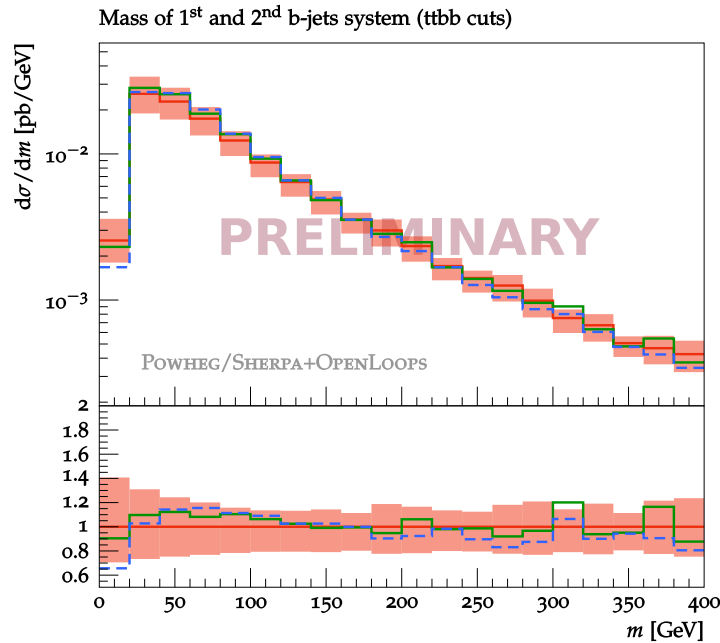


- Both NLO and LOPS accurate observables very stable with respect to hdamp^\dagger
 - Variations of at most 10% are observed

$^\dagger \text{hdamp}$ applied also to final state massive emitters

Shower dependence

— SHERPA — POWHEG BOX+Pythia - - - POWHEG BOX+Herwig



- Remarkable stability with respect to the choice of shower (Pythia 8.210 vs. Herwig 7.1.0)

Conclusions

- $t\bar{t}b\bar{b}$ @NLO+PS preferable option for $t\bar{t} + b$ -jet simulations
 - ▶ 4F scheme provides access to full $g \rightarrow b\bar{b}$ splitting phase space
 - ▶ 5F scheme resummation not too important because IS $g \rightarrow b\bar{b}$ subdominant
- Theory uncertainty of $t\bar{t}b\bar{b}$ bottleneck for $t\bar{t}H(b\bar{b})$ searches
 - HXSWG YR4 MC comparisons reveal significant matching/shower dependence
 - We now have three independent $t\bar{t}b\bar{b}$ 4F generators: Sherpa, MG5_aMC@NLO, POWHEG BOX (new!)
 - First results: POWHEG BOX shows good agreement with Sherpa and very mild hdamp and shower dependence

- POWHEG radiation formula:

$$d\sigma = \overline{B}(\Phi_B) d\Phi_B \left[\Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}^s(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d\Phi_{\text{rad}} \right] + (R_{\alpha}^r \text{ contr.})$$

- ▶ where $R_{\alpha} = R_{\alpha}^s + R_{\alpha}^r$
- Separation of the real contribution introduced to deal with “Born zeroes”
 - ▶ if (`r0.gt.5*abs(rc+rs-rs)`) then ... R_{α}^r
 - ▶ else ... R_{α}^s
- More sophisticated separation introduced in the present form:

$$R_{\alpha}^s = R_{\alpha} F(k_T^2) , \quad R_{\alpha}^r = R_{\alpha} [1 - F(k_T^2)] , \quad F(k_T^2) = \frac{h^2}{k_T^2 + h^2}$$

- In top-pair production choosing $\text{hdamp} \sim m_t$ improves the description of the data
 - ▶ ATLAS tunes $\text{hdamp} = 1.5m_t$, CMS sets to the same value

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← singular

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maybe be thought of as an analogue to μ_Q in MC@NLO

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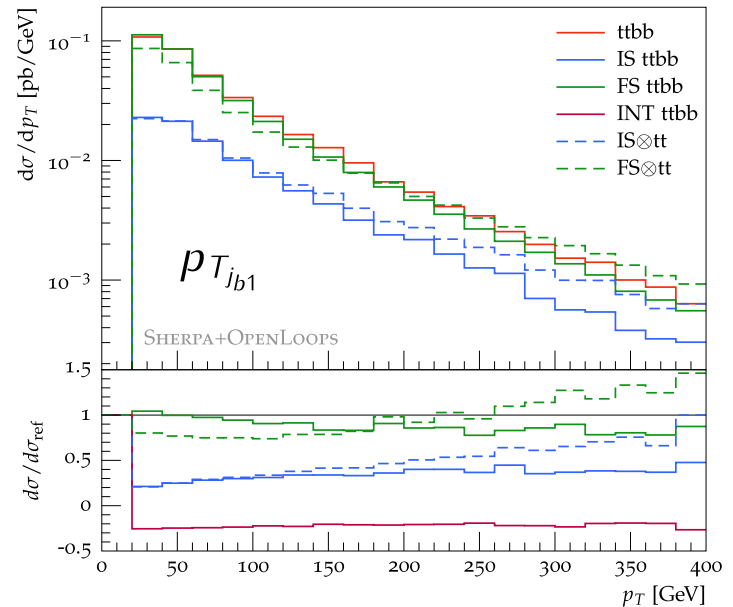
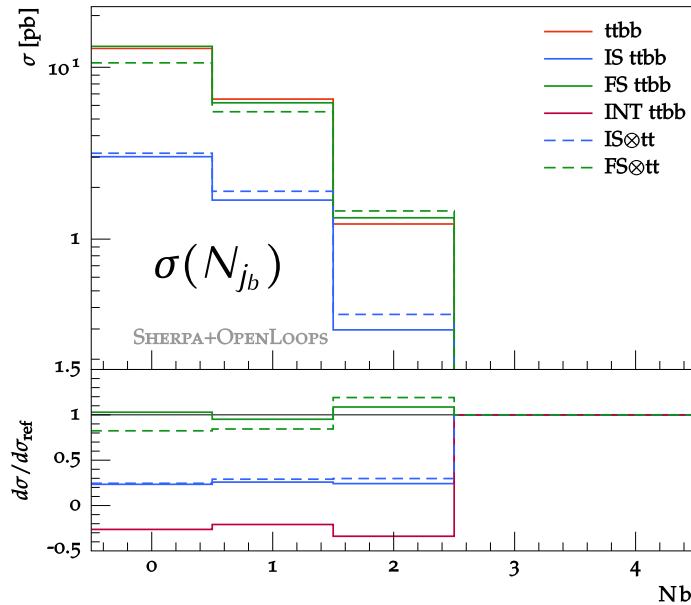
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- In $t\bar{t}b\bar{b}$:
 - Default behaviour of `hdamp` needs modifying:
 - Default “`hdamp` applied only to IS” manifests convergence issues
 - We apply `hdamp` also to massive FS, with `hdampIS` and `hdampFS` independent
 - Further investigation underway
 - New POWHEG BOX RES features could be exploited for better understanding of the `hdamp` dependence

QCD production of $t\bar{t}b\bar{b}$

- Key features of 4F $pp \rightarrow t\bar{t}b\bar{b}$:

► Dominated by topologies with FS $g \rightarrow b\bar{b}$ splittings



- FS $g \rightarrow b\bar{b}$ dominant, also away from collinear regime
- IS $g \rightarrow b\bar{b}$ subdominant (no need for 5F resummation)

supports choice of 4F scheme with $m_b > 0$ and no b -quark PDF