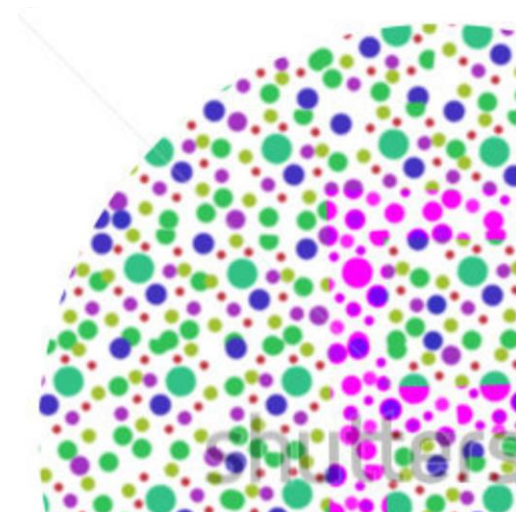
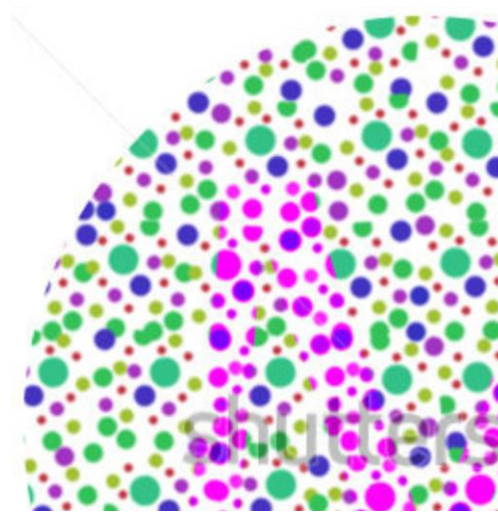
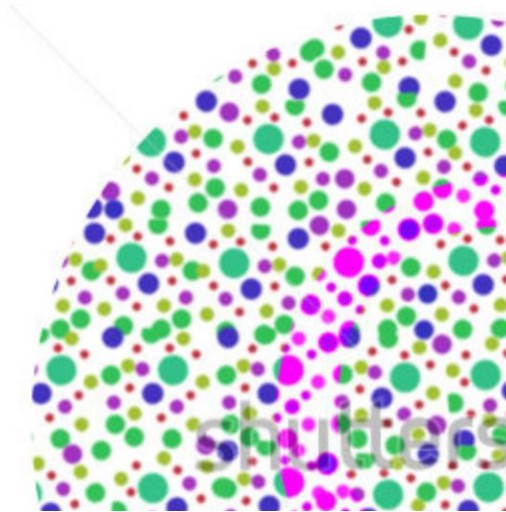


Systematics of the CME search and backgrounds

Gang Wang (UCLA)



Chiral Magnetic effect:

magnetic field + chirality = current

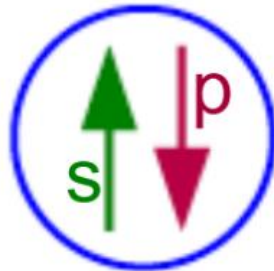
spin alignment in B-field:
opposite directions for
opposite charges

chirality

left

right

handedness:
momentum and spin,
aligned or anti-aligned



+

charge



-

negative goes up
positive goes down

positive goes up
negative goes down

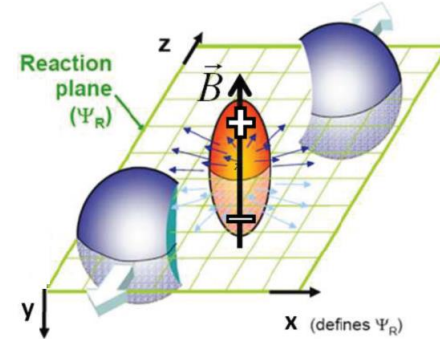
courtesy of P.Sorensen

An excess of right or left handed quarks lead to a current flow along the magnetic field.

$$\vec{J} = \frac{e^2}{2\pi^2} \mu_5 \vec{B}$$

CME observable: γ correlator

S. Voloshin, PRC 70 (2004) 057901



$$\gamma = \langle \cos(\phi_\alpha + \phi_\beta - 2\psi_{RP}) \rangle$$

$$= \left[\langle v_{1,\alpha} v_{1,\beta} \rangle + B_{in} \right] - \left[\langle a_\alpha a_\beta \rangle + B_{out} \right]$$

*background effects:
largely cancel out*

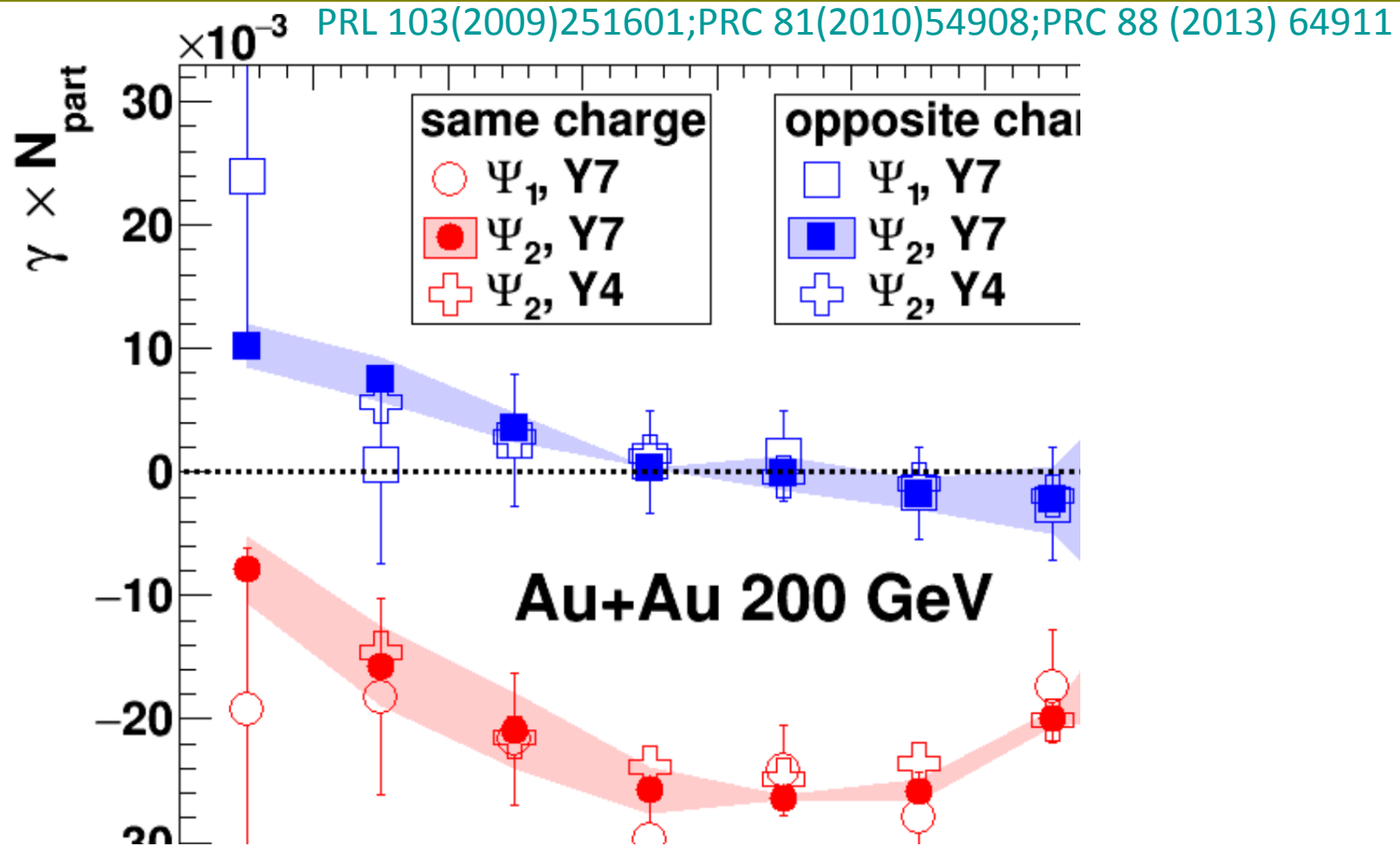
*P-even quantity:
still sensitive to
charge separation*

*Directed flow: expected to
be the same for SS and OS*

$$\frac{B_{in} - B_{out}}{B_{in} + B_{out}} = v_{2,cl} \frac{\langle \cos(\phi_\alpha + \phi_\beta - 2\phi_{cl}) \rangle}{\langle \cos(\phi_\alpha - \phi_\beta) \rangle}$$

both flow (global collectivity w.r.t the reaction plane)
and non-flow (correlations unrelated to the reaction plane:
jet, decay, HBT, momentum conservation ...)

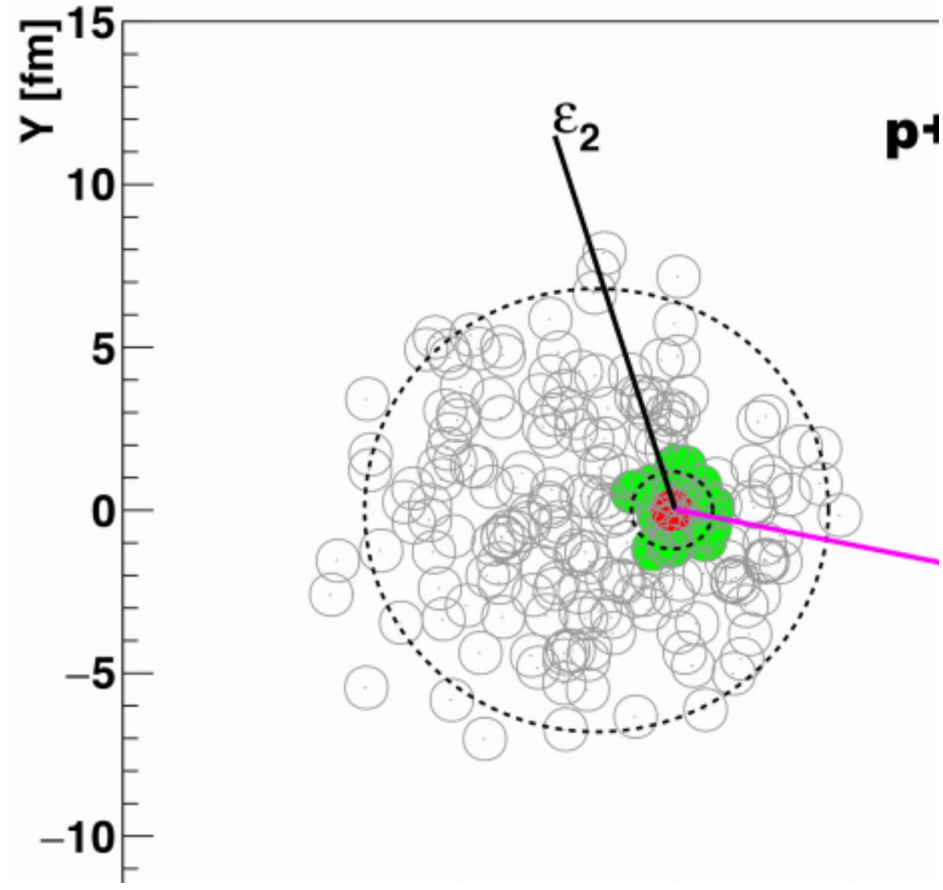
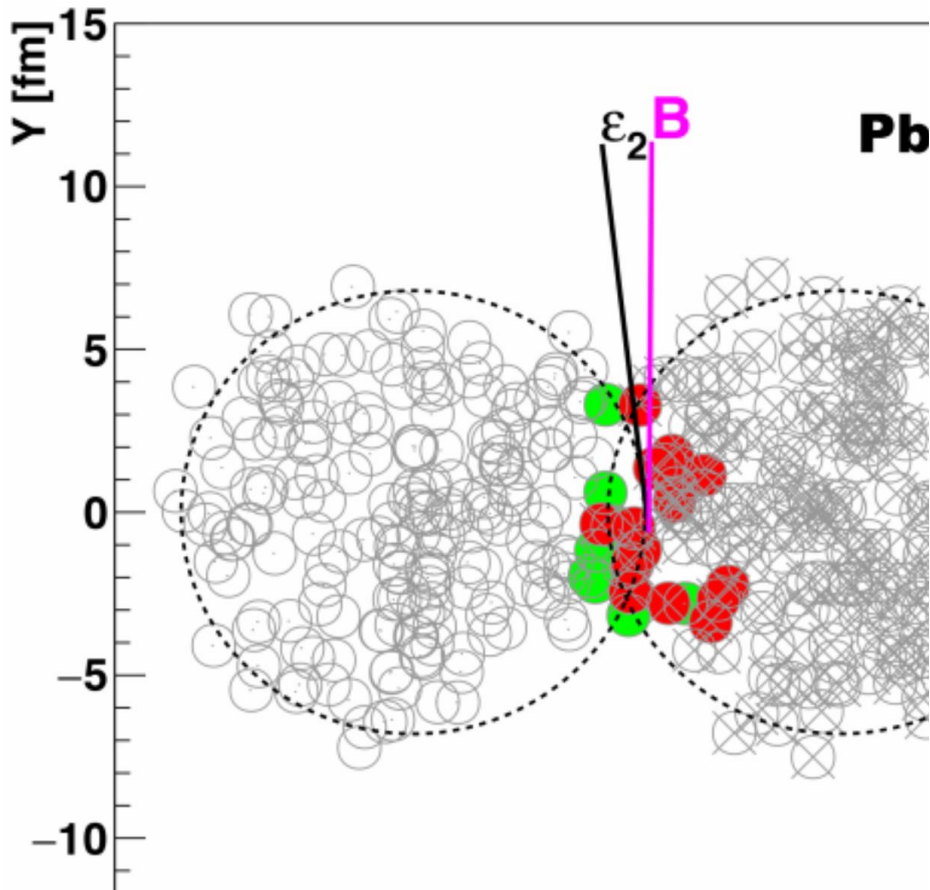
Charge separation signal



- $Y_{os} > Y_{ss}$, consistent with CME expectation
- Confirmed with 1st-order EP (from spectator neutron ν_1)
 - non-flow is not a dominant contribution in large systems
 - what about small systems?

Small systems

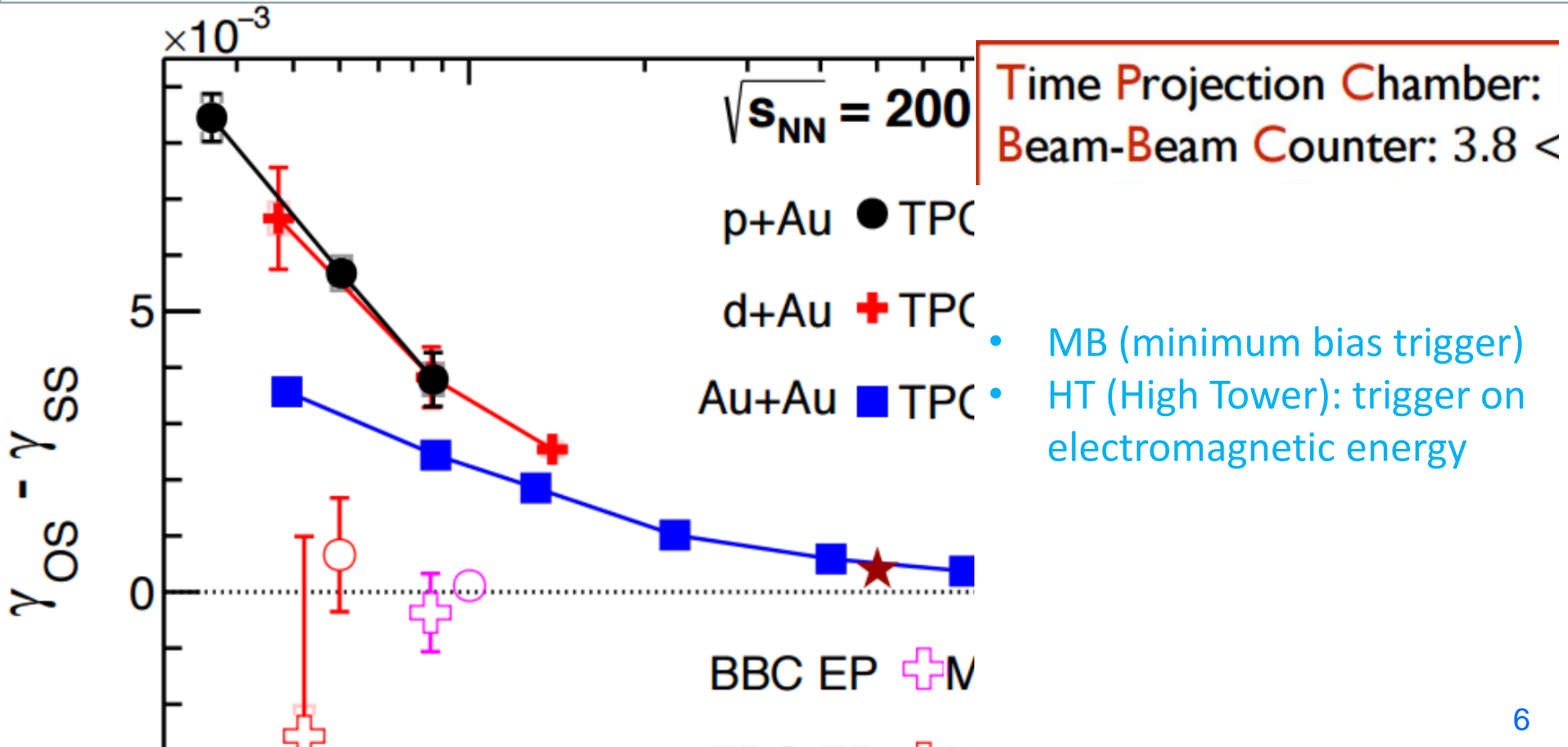
R. Belmont and J.L. Nagle, arXiv:1610.07964



- ψ_1 reflects B direction \rightarrow related to CME
- ψ_2 reflects ϵ_2 (eccentricity) \rightarrow related to backgrounds
- ψ_1 and ψ_2 are correlated in large systems, not in small systems
- non-flow also plays different roles in small and large systems

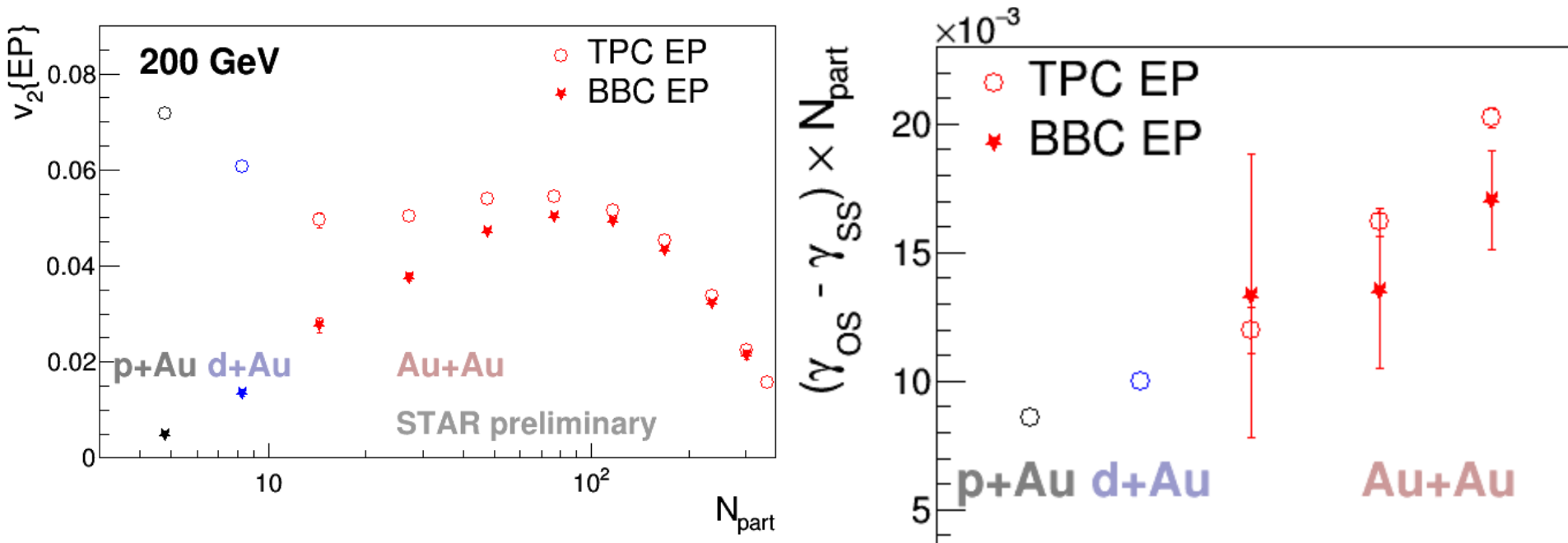
ψ_1 and ψ_2 , and non-flow

- Sizable $\Delta\gamma$ in p+Au and d+Au w.r.t 2nd-order EP fom TPC
 - similar or higher magnitudes than peripheral Au+Au
- $\Delta\gamma$ disappears in p+Au w.r.t 1st-order EP from ZDC
- $\Delta\gamma$ also disappears in p+Au w.r.t BBC EP
- short range non-flow with TPC EP



Non-flow in small systems

- Enhancement of v_2 {TPC EP} in p+Au and d+Au: non-flow
- v_2 {BBC EP} instead decreases toward p+Au/d+Au
- Important to have η gap between EP and particles of interest

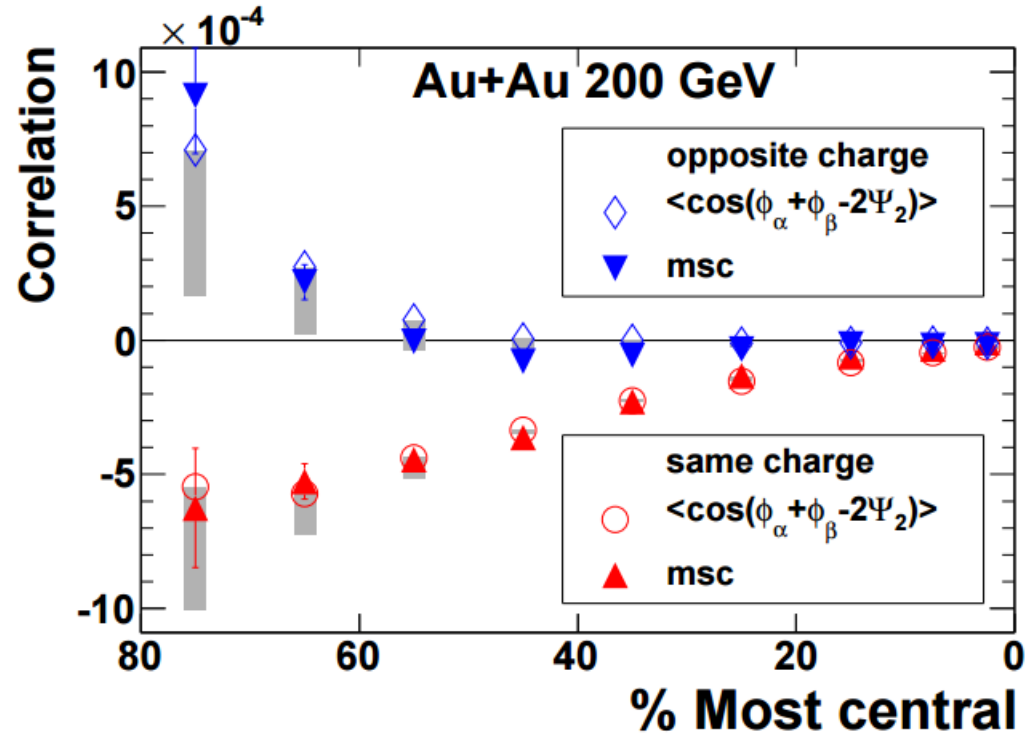
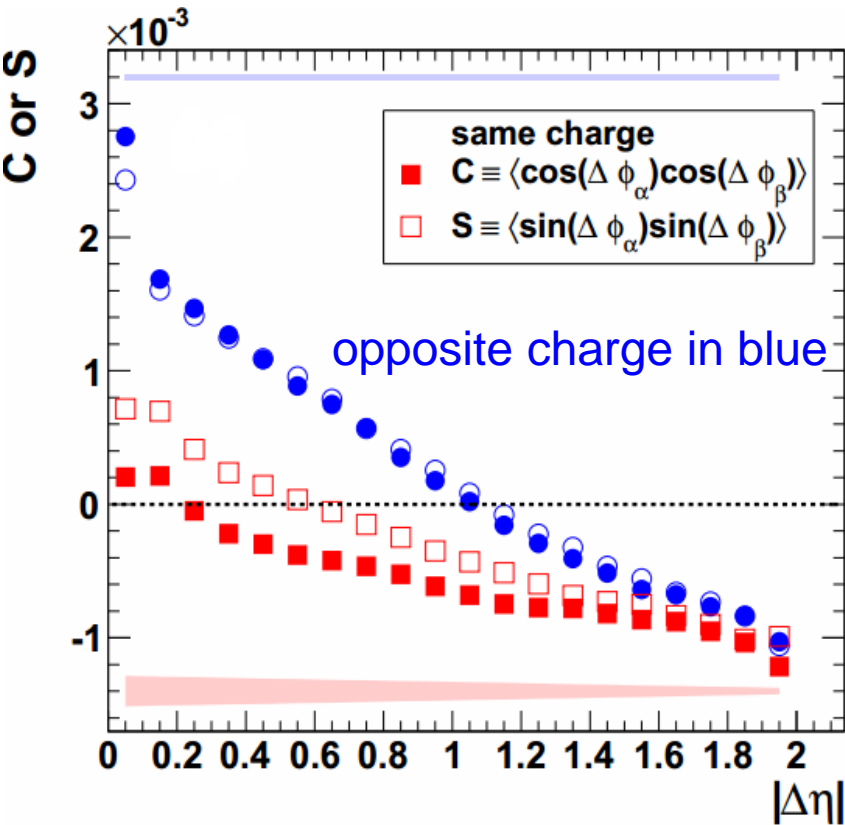


- $\Delta\gamma$ {TPC EP} mostly short range non-flow in p+Au/d+Au
- In larger systems, TPC EP and BBC EP give consistent results
 - non-flow is not a dominant contribution in large systems
 - corroborate the conclusion with ZDC EP

Short range correlations

200 GeV Au+Au: 40 - 60%

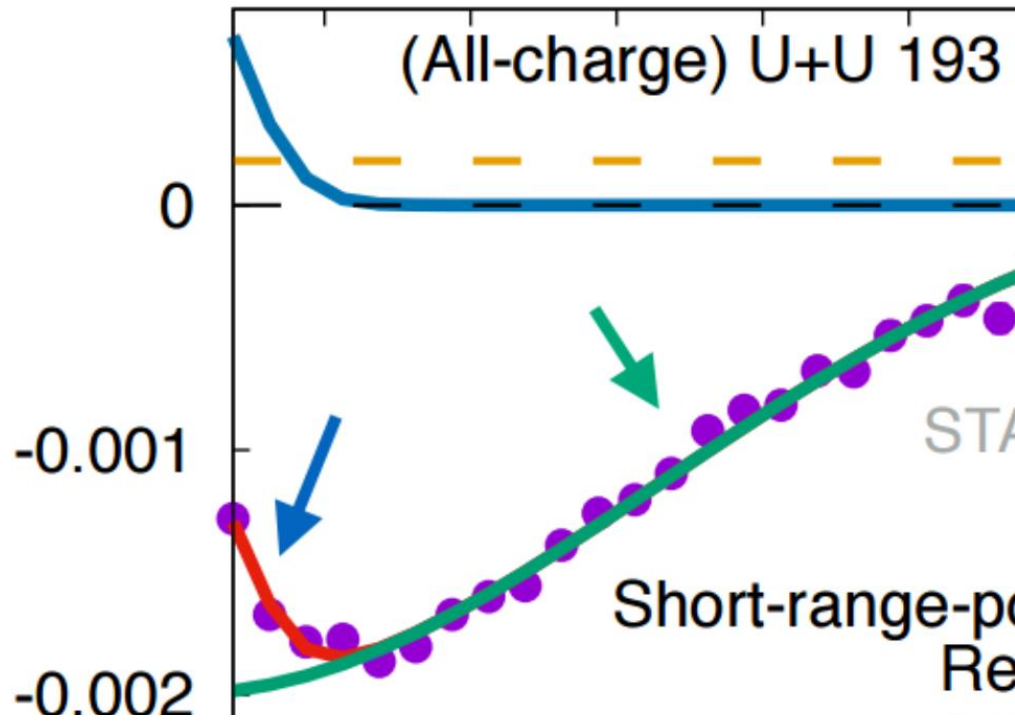
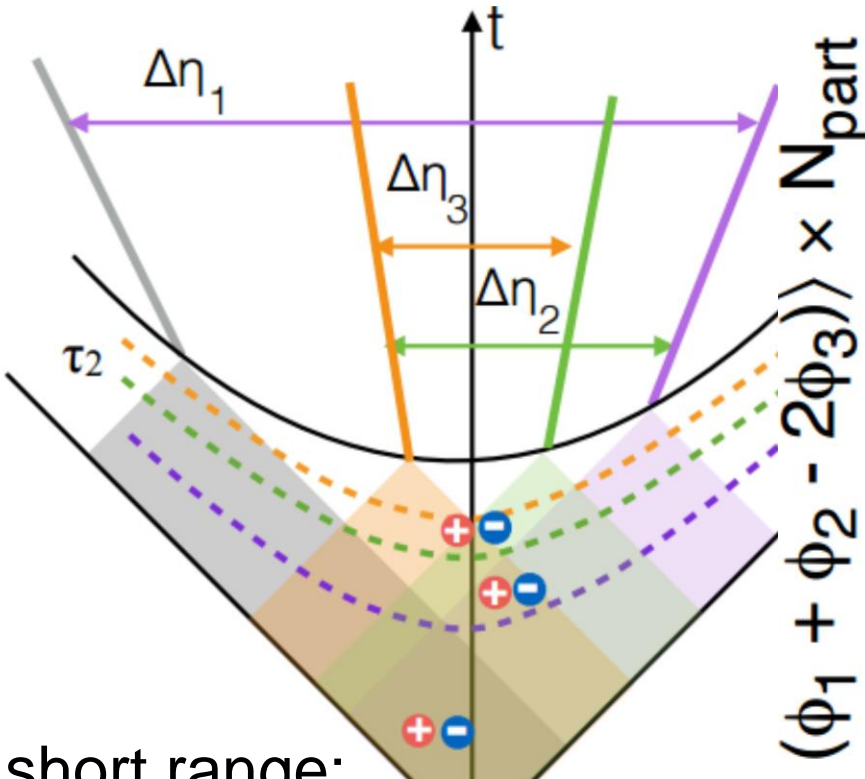
Phys. Rev. C 88 (2013) 64911



- Prominent correlations exist at small Δp_T and $\Delta \eta$
- Probably due to HBT+Coulomb
- Excluding small Δp_T and $\Delta \eta$:
 - significant effect in peripheral collisions (small systems)

Closer look at short-range correlations

early-time charge separation \rightarrow long-range correlations in $\Delta\eta$
 (between two particles)



short range:

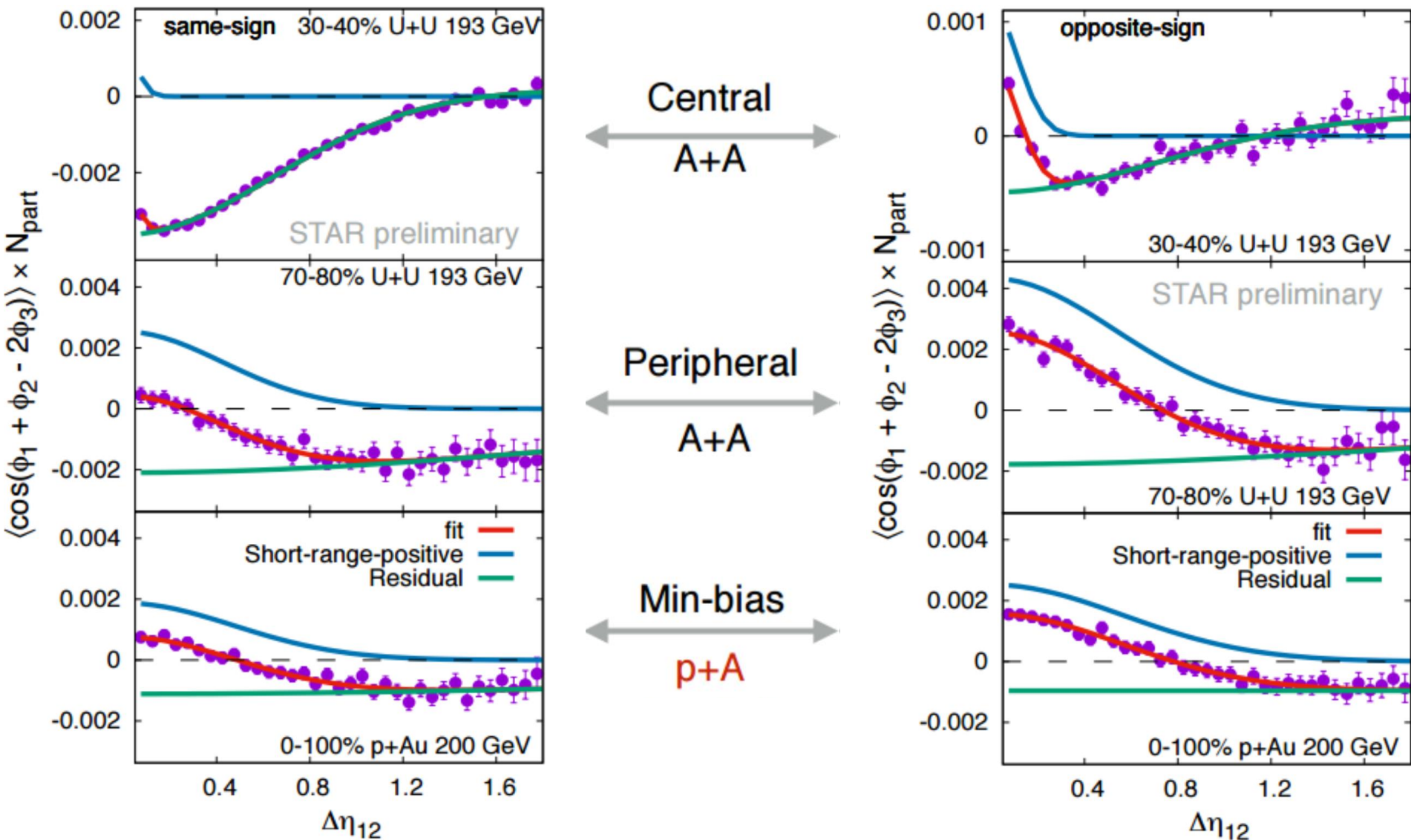
$$\Delta\phi \rightarrow 0, \Delta\eta \rightarrow 0 : C_{112} = \langle \cos(\phi_1(\eta_1) + \phi_2(\eta_2) - 2\phi_3) \rangle$$

$$C_{112}(\Delta\eta_{12}) = A_{SR}^+ e^{-(\Delta\eta)^2 / 2\sigma_{SR}^2} - A_{IR}^- e^{-(\Delta\eta)^2 / 2\sigma_{IR}^2}$$

Large and small systems

same charge

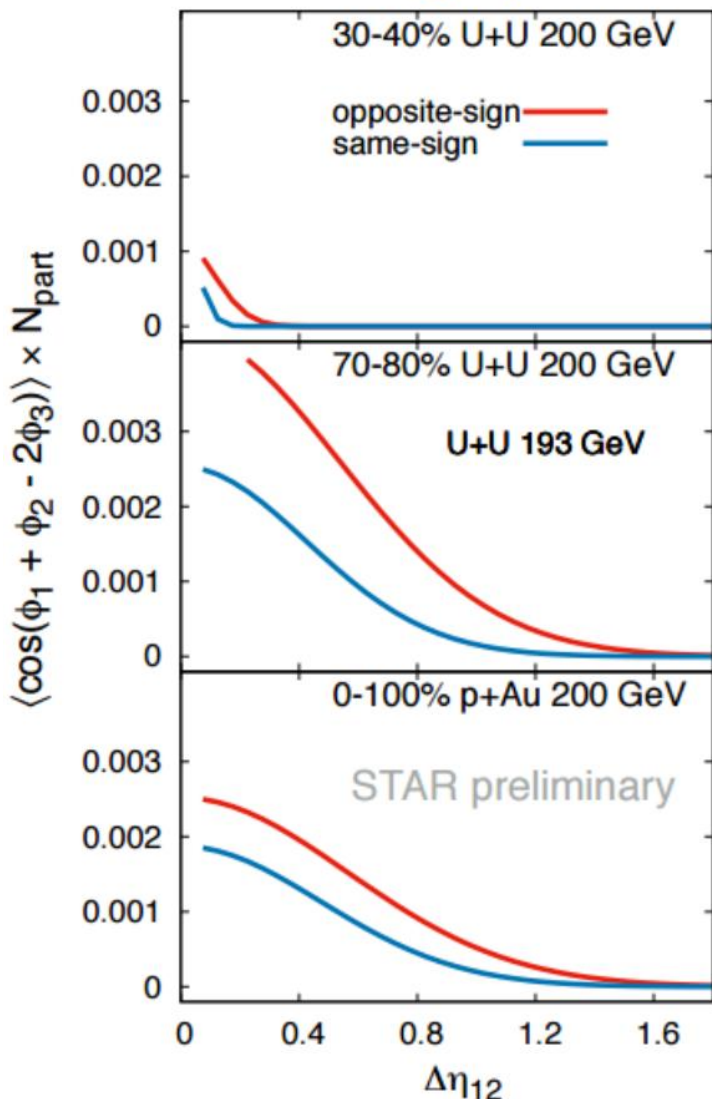
opposite charge



p+Au data are very similar to peripheral U+U.

Residual

Short-range-positive

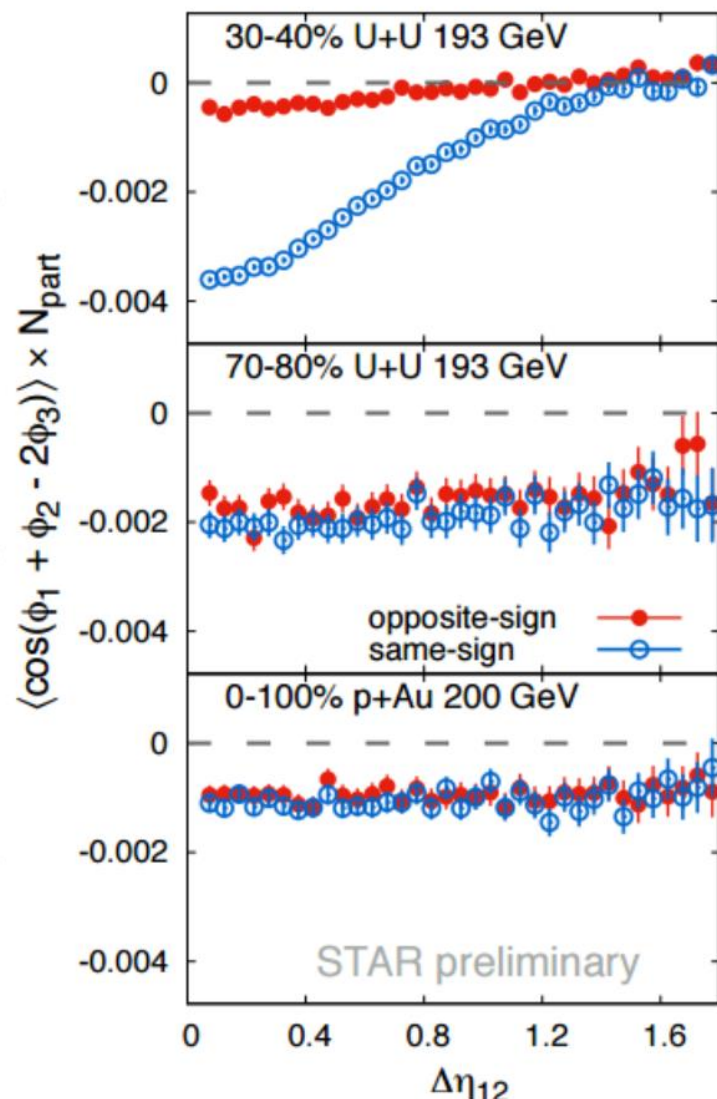


Central
A+A

Peripheral
A+A

Min-bias
p+A

Residual



- p+A & peripheral A+A → dominated by short-range correlations
- Finite residual in large systems → not necessarily true CME signal.¹¹

Flow-related background

An example with no charge separation:

v_2 + local charge conservation/decay + momentum conservation



$$\left. \begin{array}{l} \gamma_{SS} = -1 \\ \delta_{SS} = -1 \end{array} \right\} \longrightarrow H_{SS}^{\kappa=1}=0$$

$$v_2 = 1$$

$$\left. \begin{array}{l} \gamma_{OS} = 0 \\ \delta_{OS} = 0 \end{array} \right\} \longrightarrow H_{OS}^{\kappa=1}=0$$

$$\gamma \equiv \langle \cos(\phi_1 + \phi_2 - 2\Psi_{RP}) \rangle = \kappa v_2 F - H \longrightarrow H^\kappa = (\kappa v_2 \delta - \gamma) / (1 + \kappa v_2)$$

$$\delta \equiv \langle \cos(\phi_1 - \phi_2) \rangle = F + H,$$

A. Bzdak, V. Koch and J. Liao, Lect. Notes Phys. 871, 503 (2013).

With $\kappa=1$, H tells the truth.

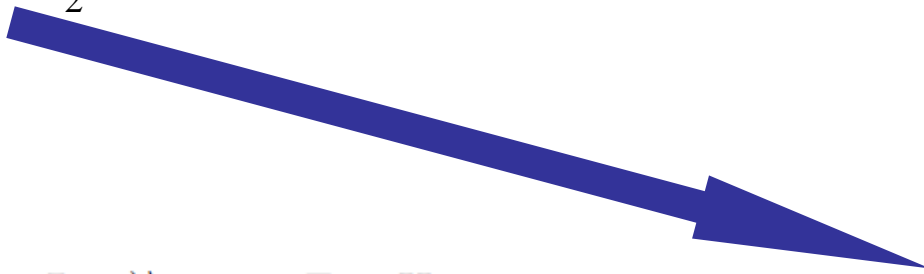
But this is an over-simplified example. In reality, κ could deviate from 1.

A cumulant way

$$\begin{aligned} & \cos(\varphi_1 + \varphi_2 - 2\psi_{\text{RP}}) \\ &= \cos(\varphi_1 - \varphi_2 + 2\varphi_2 - 2\psi_{\text{RP}}) \\ &= \cos(\varphi_1 - \varphi_2) \cos(2\varphi_2 - 2\psi_{\text{RP}}) - \sin(\varphi_1 - \varphi_2) \sin(2\varphi_2 - 2\psi_{\text{RP}}) \end{aligned}$$

If we take the "cumulant" approach, a " v_2 -free" correlator will be

$$\begin{aligned} \gamma^{\text{cumulant}} &= \langle\langle \cos(\varphi_1 + \varphi_2 - 2\psi_{\text{RP}}) \rangle\rangle \\ &= \langle \cos(\varphi_1 + \varphi_2 - 2\psi_{\text{RP}}) \rangle - \langle \cos(\varphi_1 - \varphi_2) \rangle \cdot \langle \cos(2\varphi - 2\psi_{\text{RP}}) \rangle \\ &= \gamma - \delta \cdot v_2 \end{aligned}$$



$$\begin{aligned} \gamma &\equiv \langle \cos(\phi_1 + \phi_2 - 2\Psi_{\text{RP}}) \rangle = \kappa v_2 F - H \\ \delta &\equiv \langle \cos(\phi_1 - \phi_2) \rangle = F + H, \end{aligned} \quad \longrightarrow \quad H^\kappa = (\kappa v_2 \delta - \gamma) / (1 + \kappa v_2)$$

The cumulant approach indicates $\kappa \sim 1$.

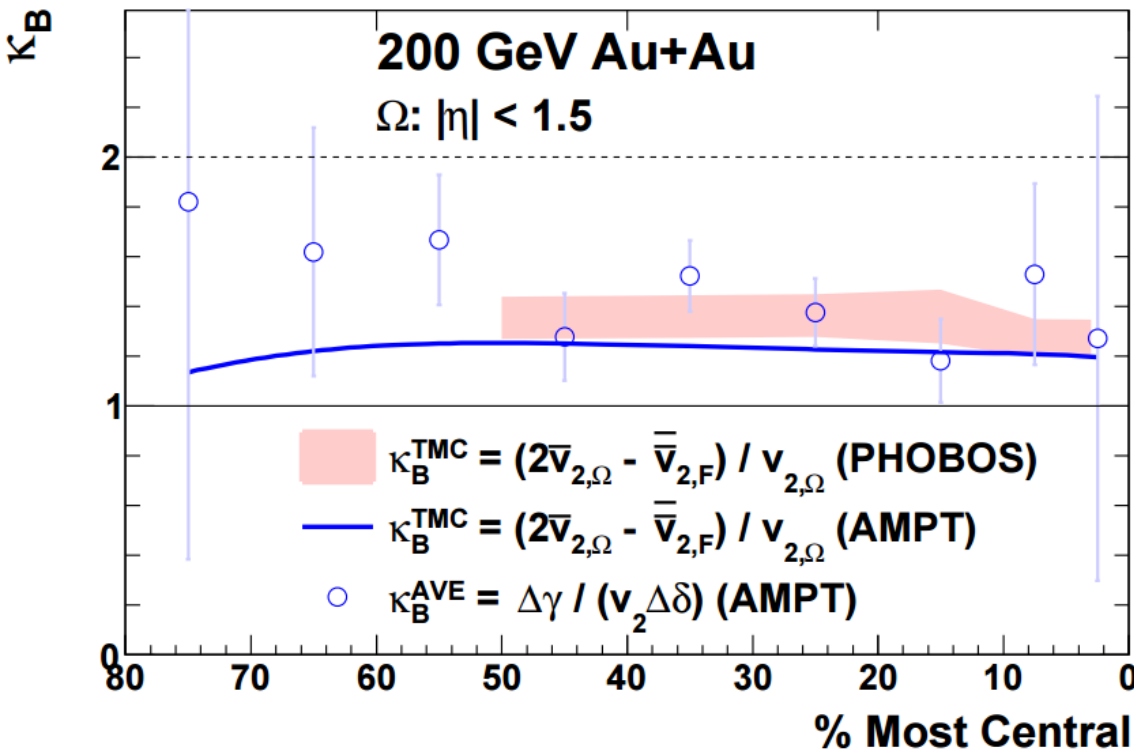
κ_B : background level

If γ measurements are dominated by v_2 + trans. momentum conservation,

$$\gamma / \delta \approx 2\bar{v}_{2,\Omega} - \bar{\bar{v}}_{2,F}$$

A. Bzdak, V. Koch and J. Liao, Lect. Notes Phys. 871, 503 (2013).

where F and Ω denote particle averages in the full phase-space and the detector acceptance, respectively. TMC: $\kappa_B^{\text{TMC}} \approx (2\bar{v}_{2,\Omega} - \bar{\bar{v}}_{2,F}) / v_{2,\Omega}$



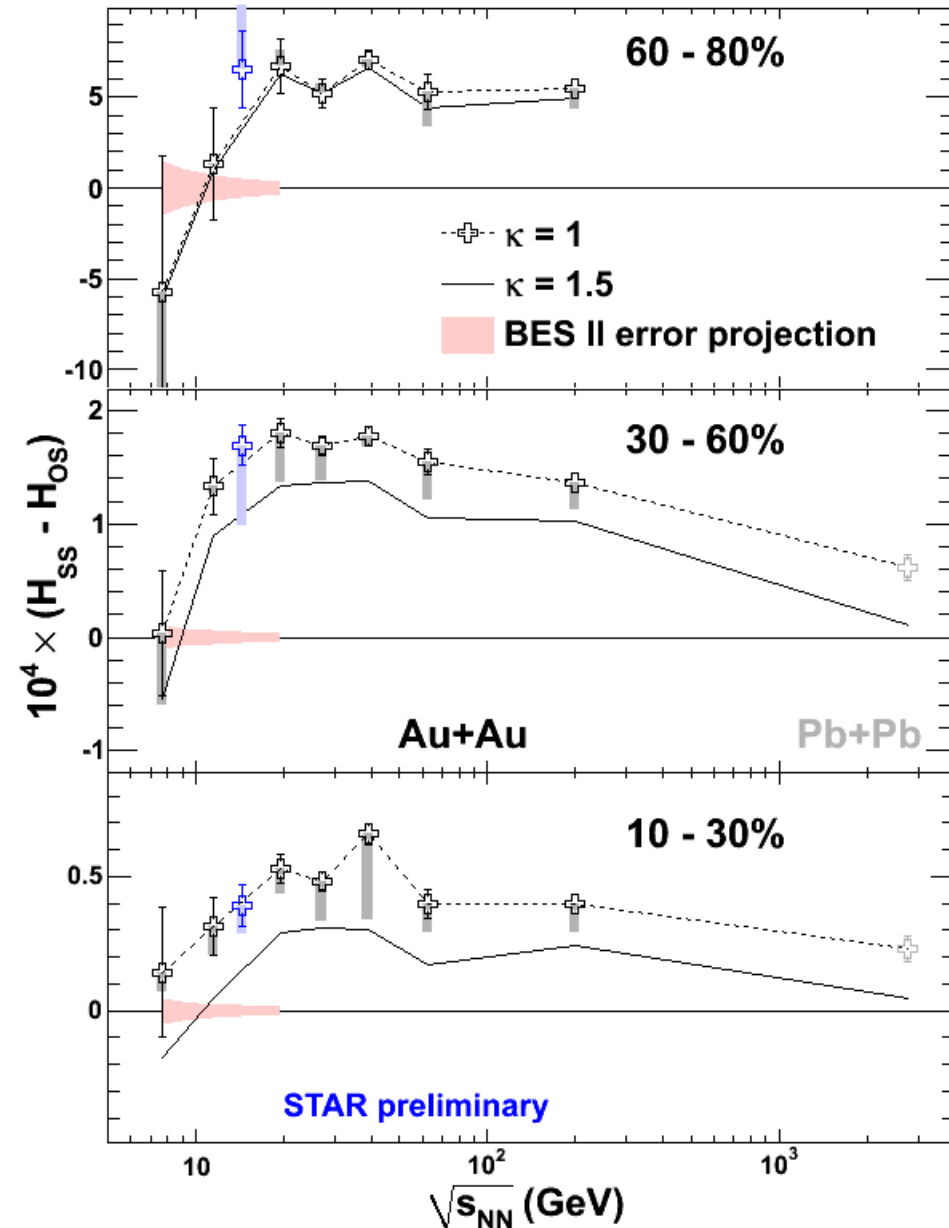
PHOBOS v_2 @ 200 GeV
Au+Au $\rightarrow \kappa_B^{\text{TMC}} \in [1.2, 1.4]$.

Other effects:
Local Charge Conservation
(LCC) and resonance decay.

AMPT shows similar κ_B
with v_2 or with $\Delta\gamma / (v_2\Delta\delta)$.

ΔH^κ at Beam Energy Scan

STAR, PRL 113 (2014) 052302



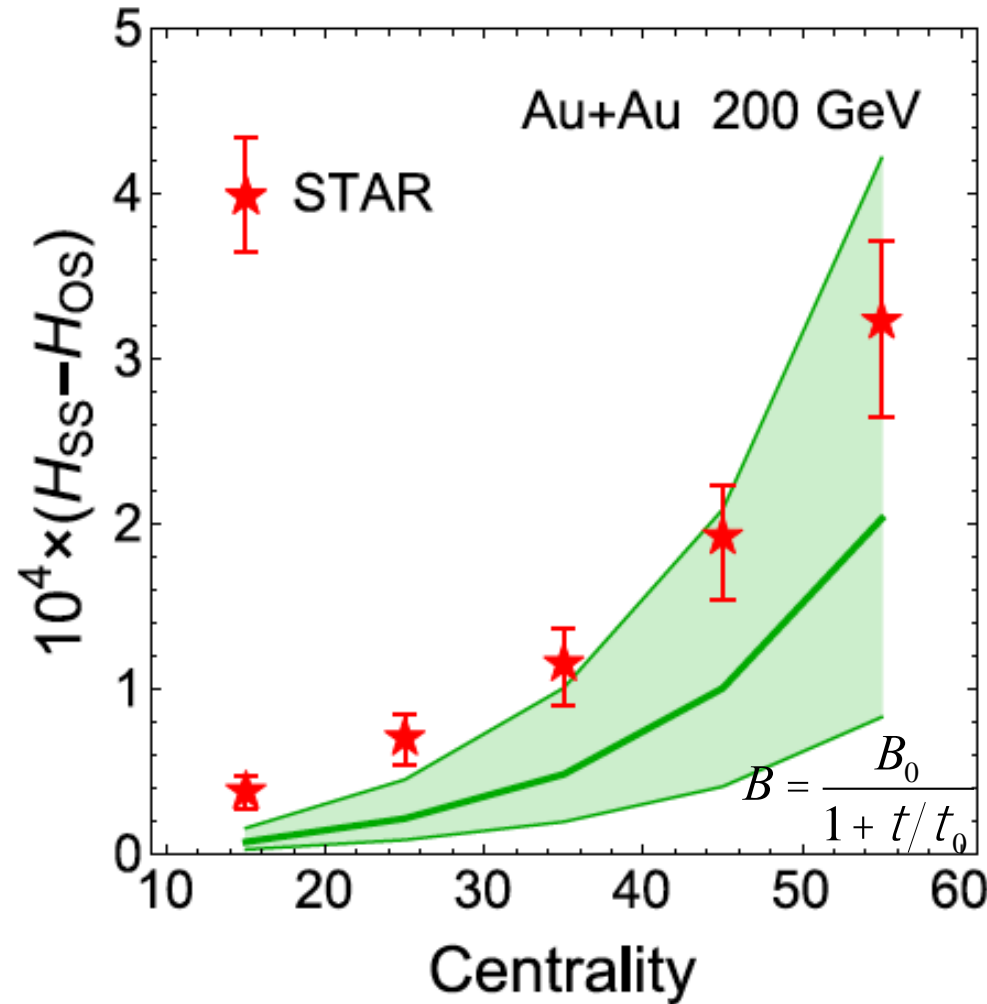
$$H^\kappa = (\kappa v_2 \delta - \gamma) / (1 + \kappa v_2)$$

A. Bzdak, V. Koch and J. Liao, Lect. Notes Phys. 871, 503 (2013).

- κ_B is roughly contained in the range of [1, 1.5].
- CME signal (ΔH) decreases to 0 from 19.6 to 7.7 GeV
- Probable domination of hadronic interactions over partonic ones
- Need more more statistics
- Another way to look at it ...

Model comparison

$$H^\kappa = (\kappa v_2 \delta - \gamma) / (1 + \kappa v_2)$$



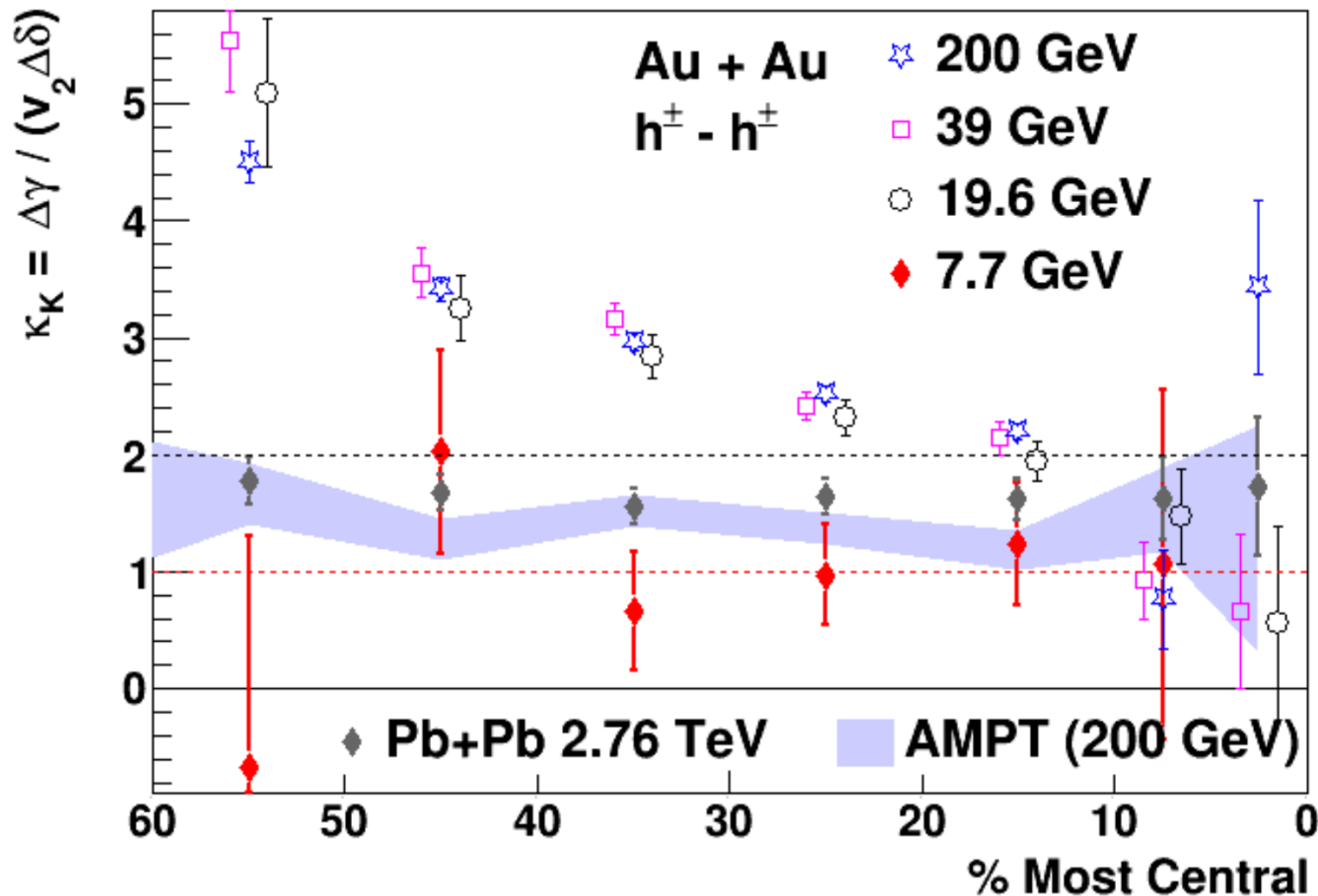
- CME requires a magnetic field, massless quarks, and axial charge **at the same time**
- Anomalous hydrodynamic calculations with $\kappa_B = 1.2$, with the range of [1, 1.5].
- Model and STAR data are compatible
- Another way to look at it ...

κ_K : normalized (signal + background)

$$\kappa_B \equiv \frac{\Delta\gamma + \Delta H}{v_2(\Delta\delta - \Delta H)}, \quad \kappa_K \equiv \kappa_B (\Delta H = 0) = \frac{\Delta\gamma}{v_2\Delta\delta}.$$

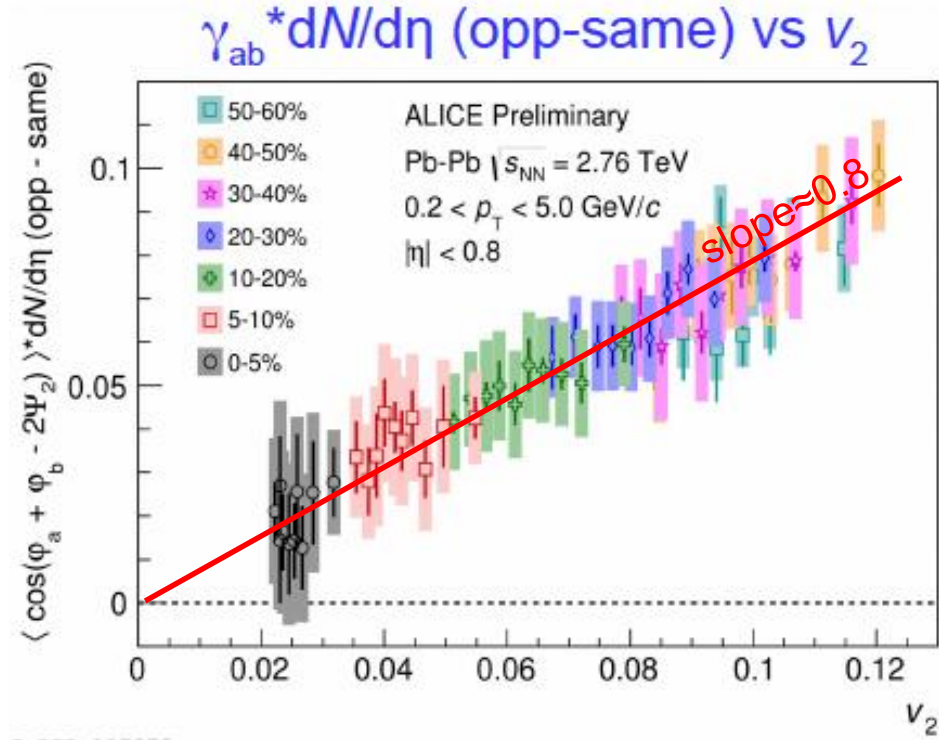
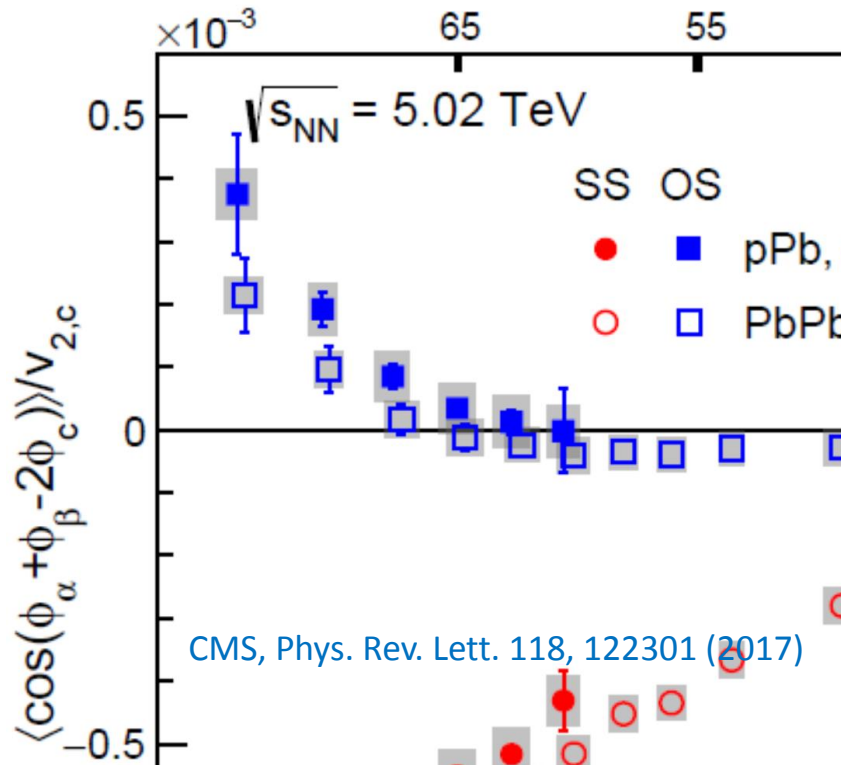
If $\kappa_K > \kappa_B$ for real data, there could be extra physics like the CME.

STAR, PRL113 (2014) 052302; ALICE, PRL110 (2013) 021301.



No CME at high energies?

ALICE (2.76 TeV Pb+Pb) event-shape engineering results show the signal is consistent with pure v_2 background.



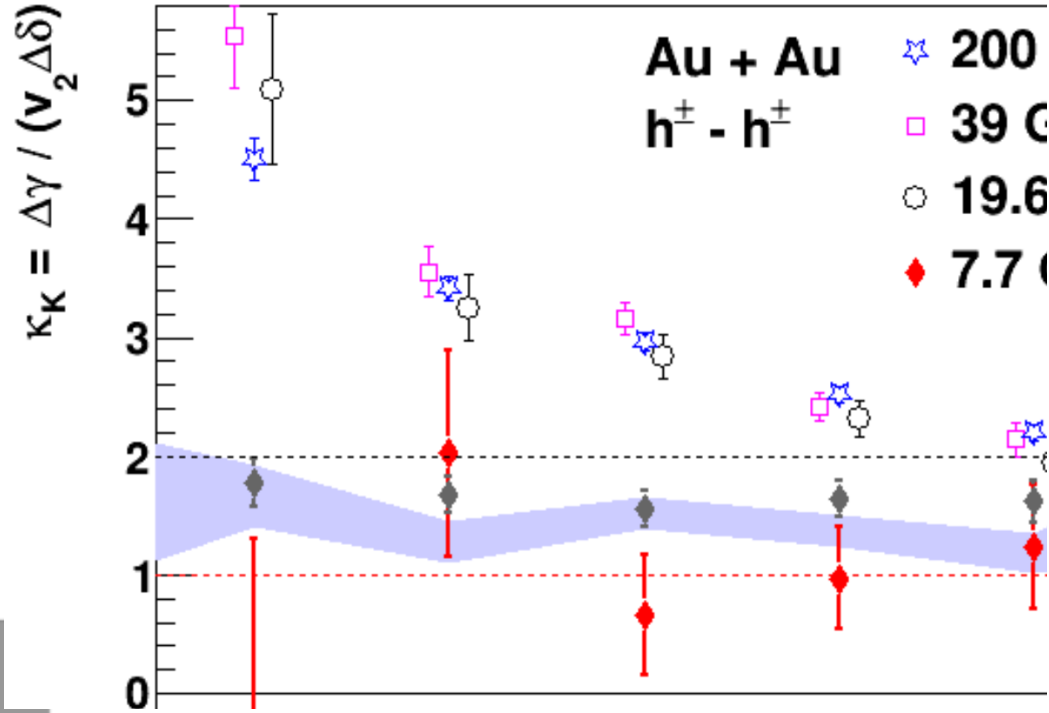
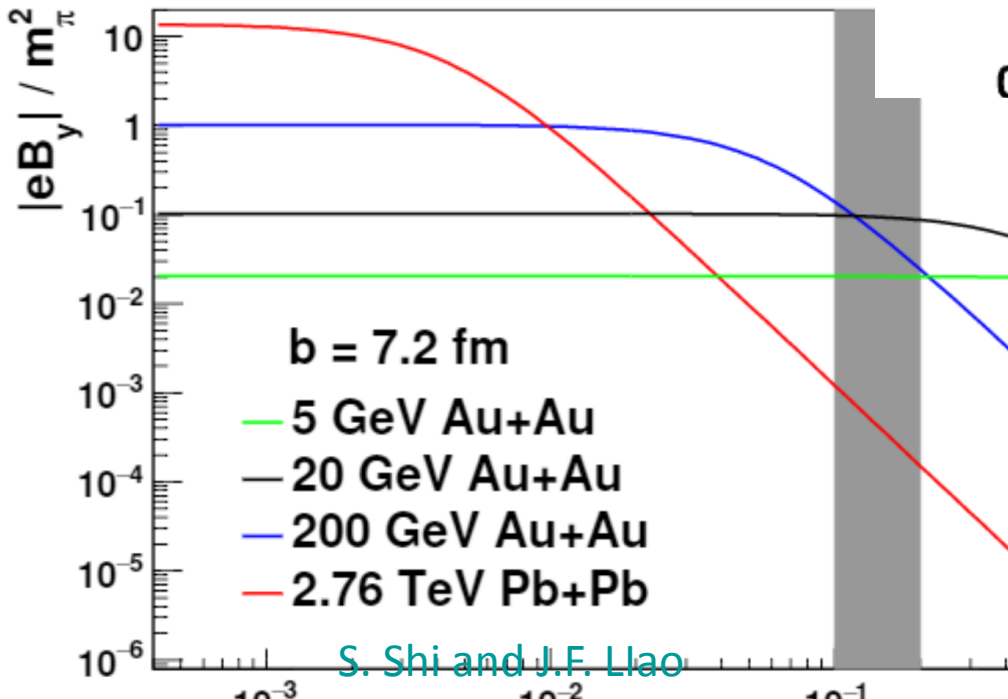
CMS (5 TeV) results show γ consistent between p+Pb and Pb+Pb: both are pure flow-backgrounds?

No CME at high energies?

$$\kappa_K \equiv \frac{\Delta\gamma}{v_2\Delta\delta}$$

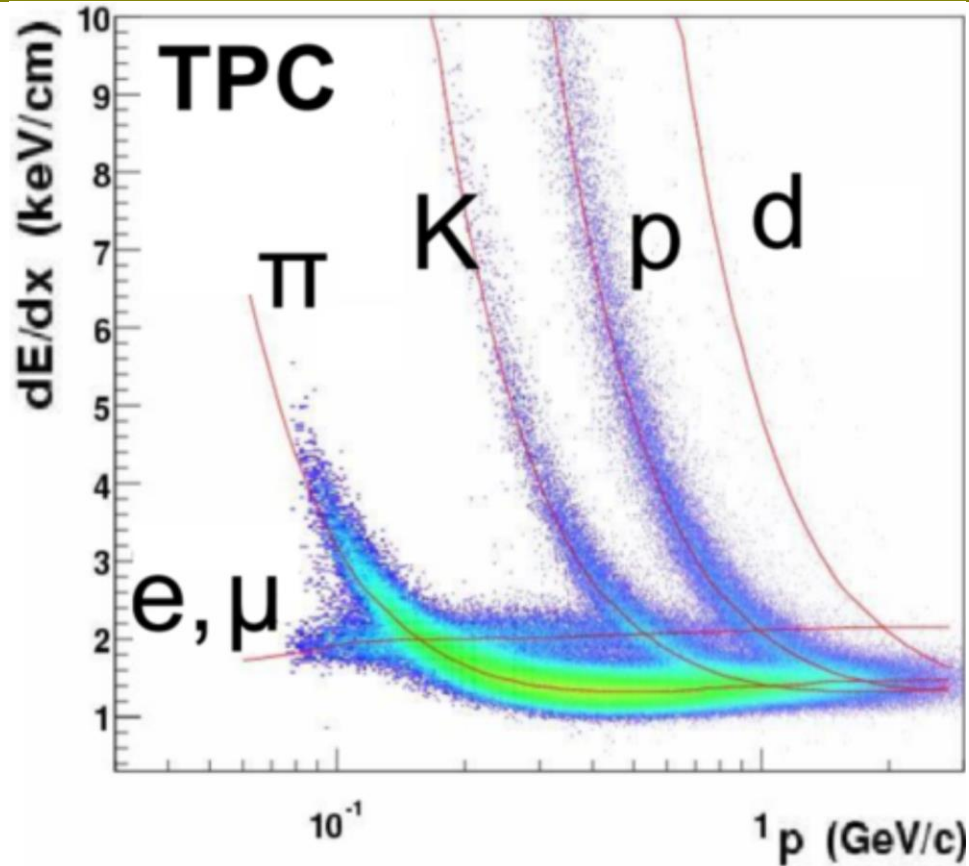
$\kappa_K \sim \kappa_B$ for both very low and very high collision energies:

- low energies, no QGP
- high energies, no B field?

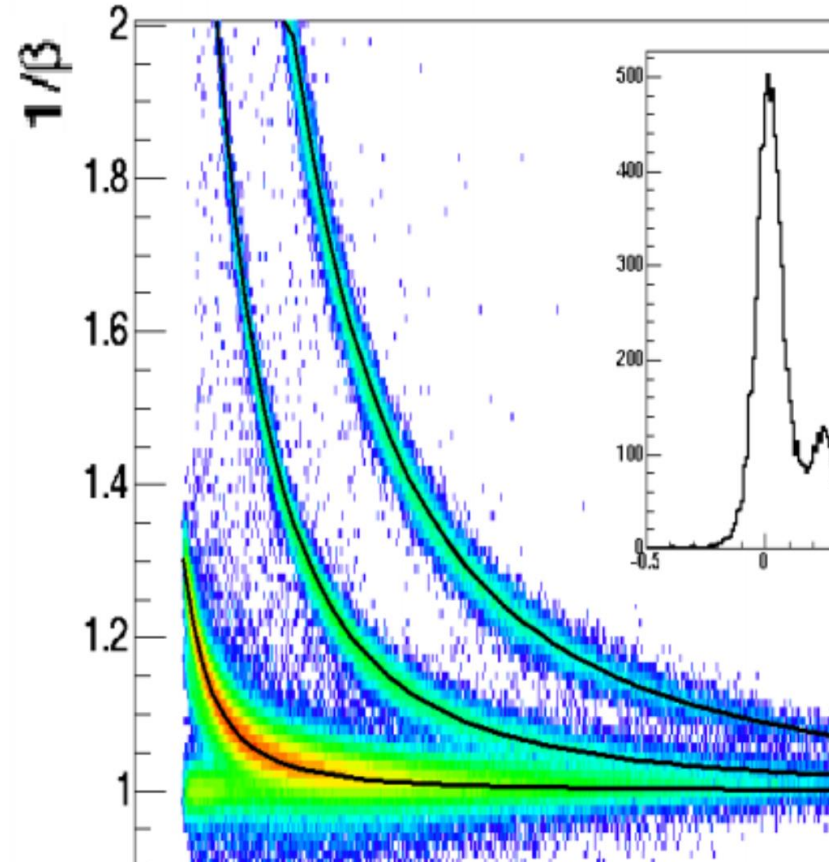


- $t=0$: B@2.76 TeV is 13 times stronger than B@200 GeV
- $t=0.01 \text{ fm/c}$, they are the same
- $t=0.1 \sim 0.2 \text{ fm/c}$, B@2.76 TeV is lower than B@200 GeV by a few orders of magnitude

Particle identification

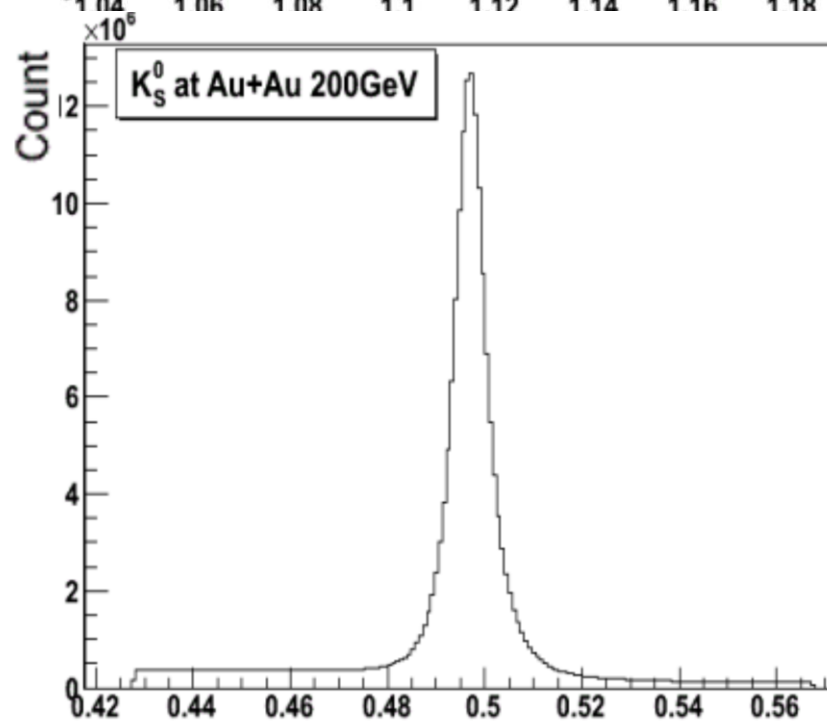
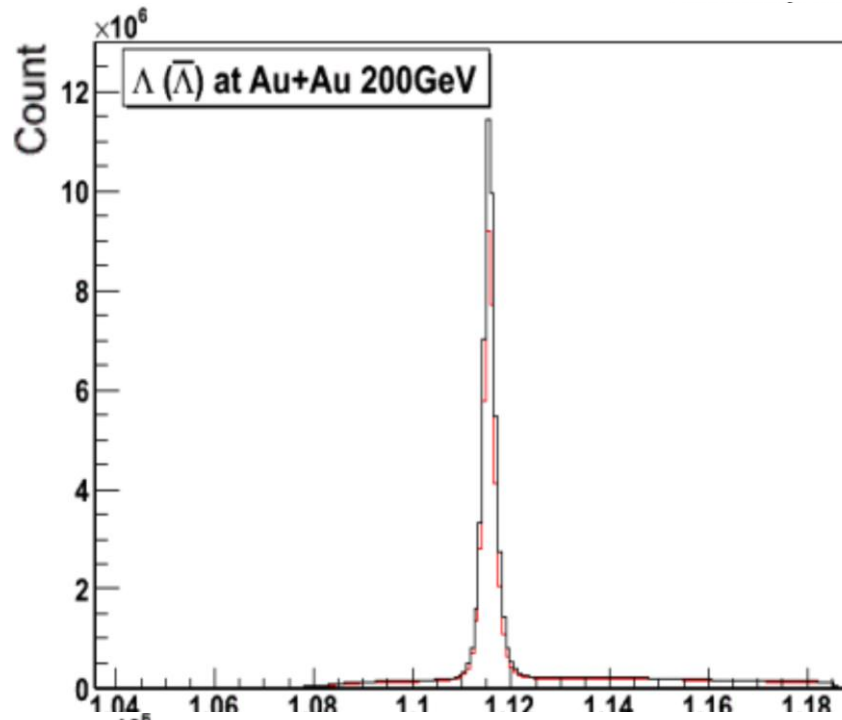
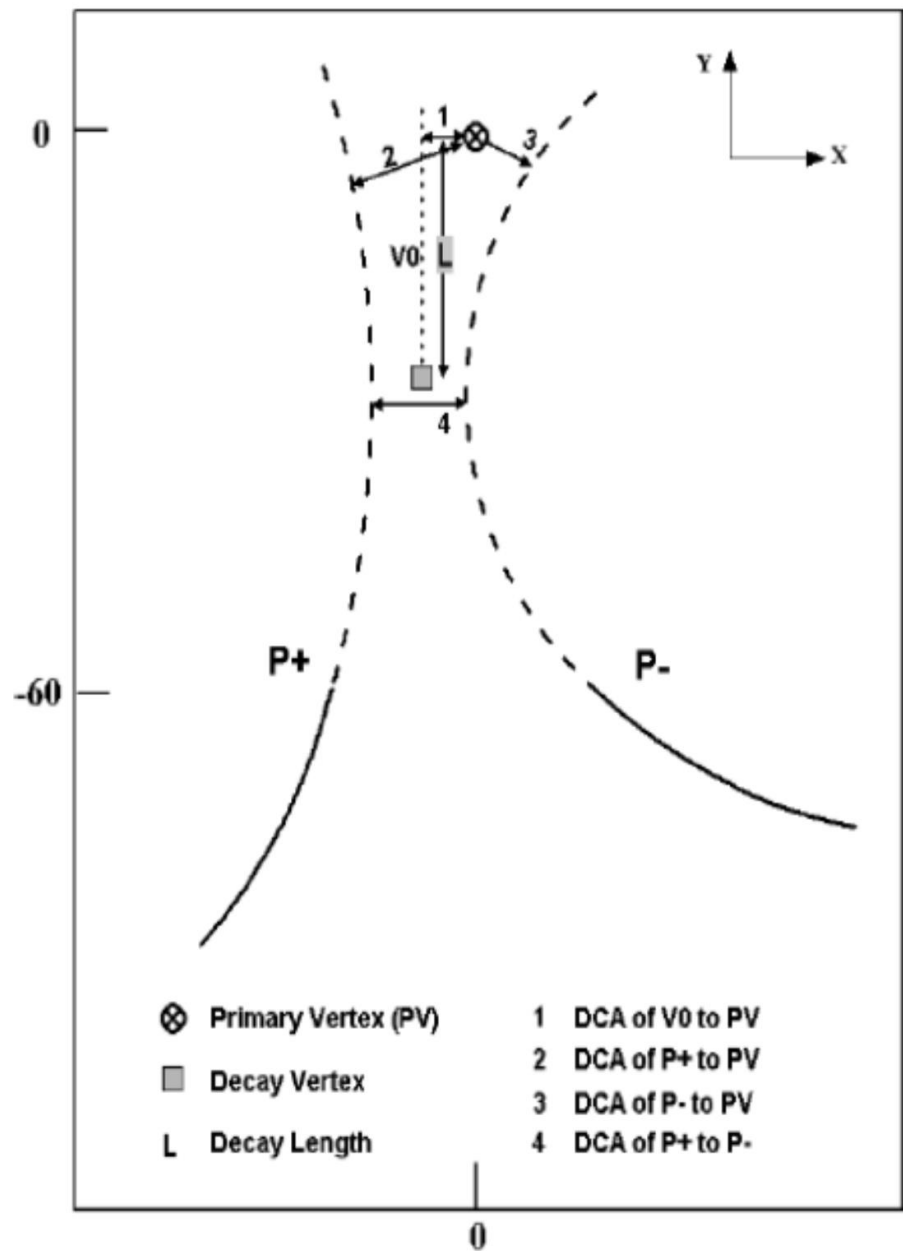


Time Projection Chamber (TPC) measures the mean energy loss per distance travelled of swift charged particles.

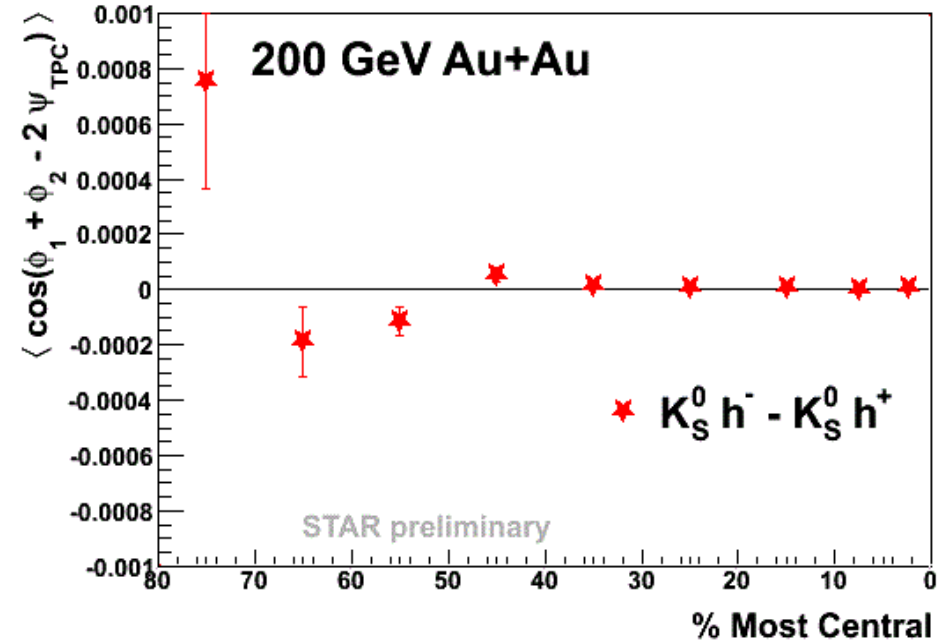
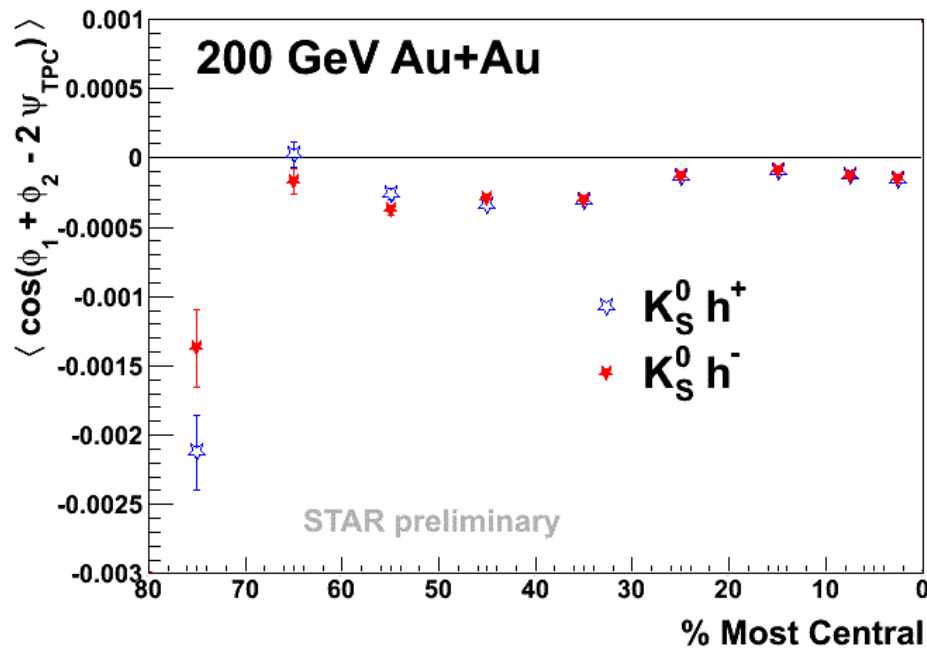


Time-of-Flight (TOF) measures the flight duration, hence velocity ($\beta = v/c$) of swift charged particles.

$\Lambda (K^0_S)$

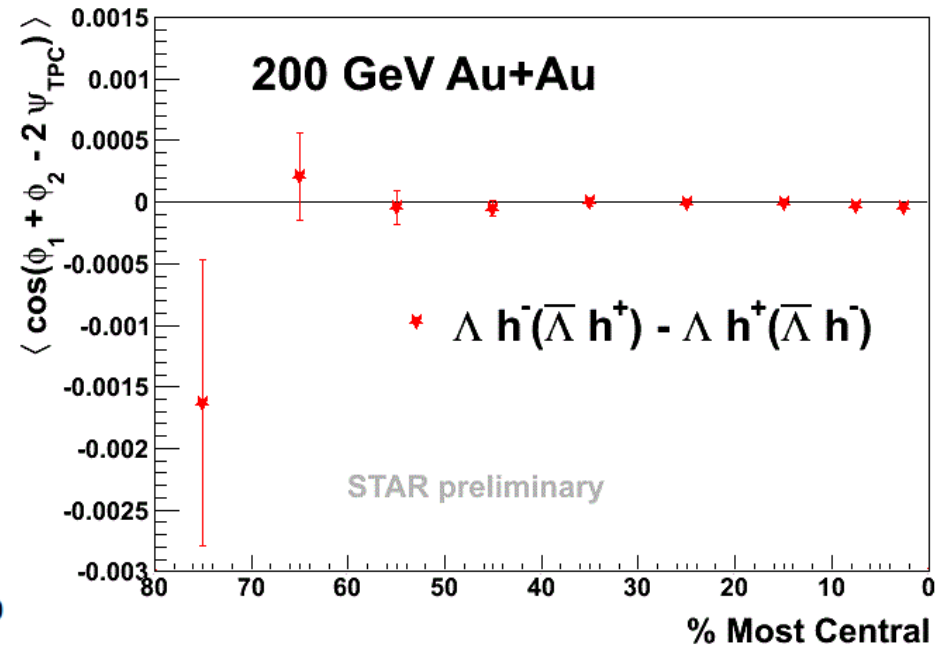
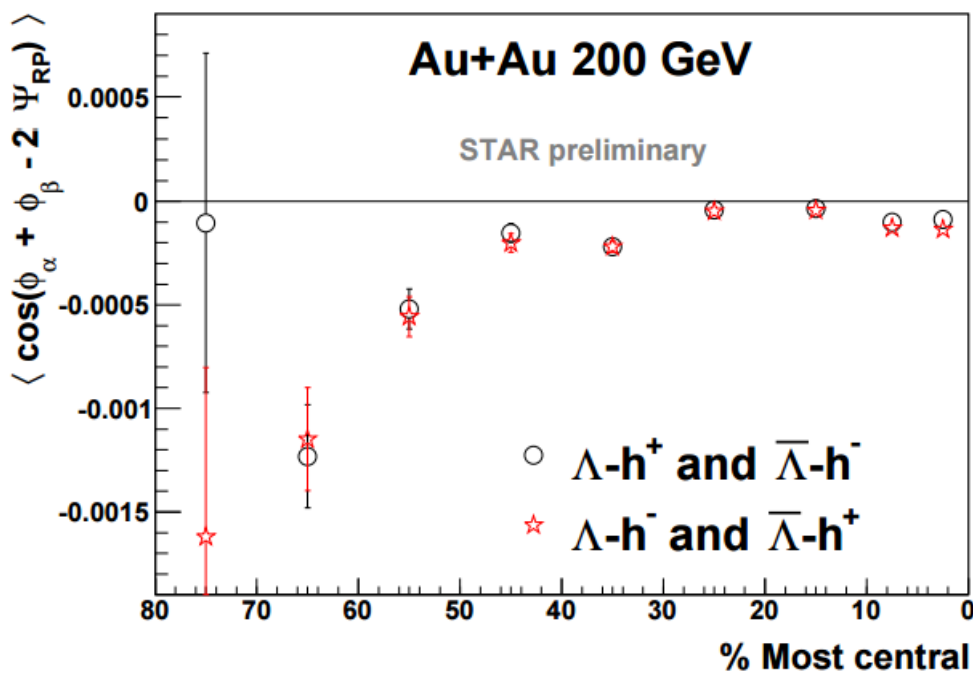


K_S^0 -hadron correlation



- $K_S^0 h^-$ consistent with $K_S^0 h^+$: no charge-dependence
- The separation observed in $h^\pm h^\pm$ is due to electric charge
- Strange quarks participate in the chiral dynamics in the same way as u and d.

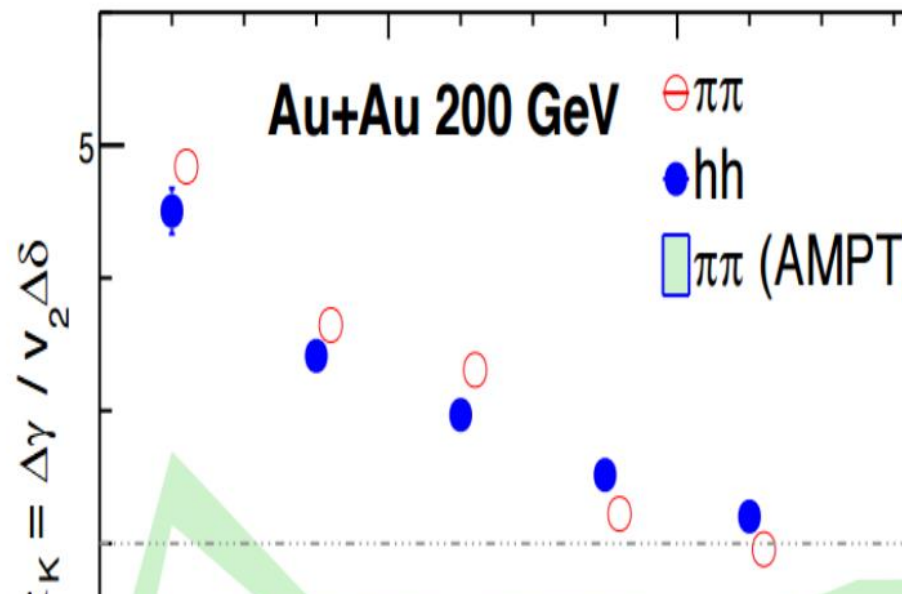
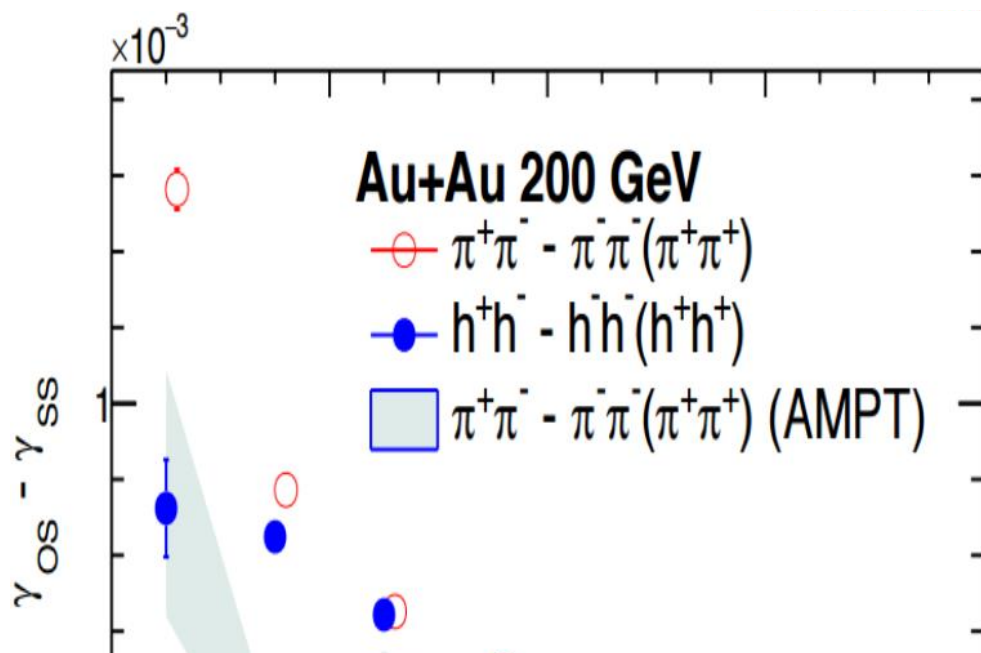
Λ -hadron correlation



- $\Lambda-h^\pm$ also show no charge-dependent separation
 - (protons and antiprotons have been excluded from h^\pm)
- s quarks participate in the chiral dynamics in a similar way as u/d
- $\Lambda-h^\pm$ also provides a baseline for $\Lambda-p$ correlations

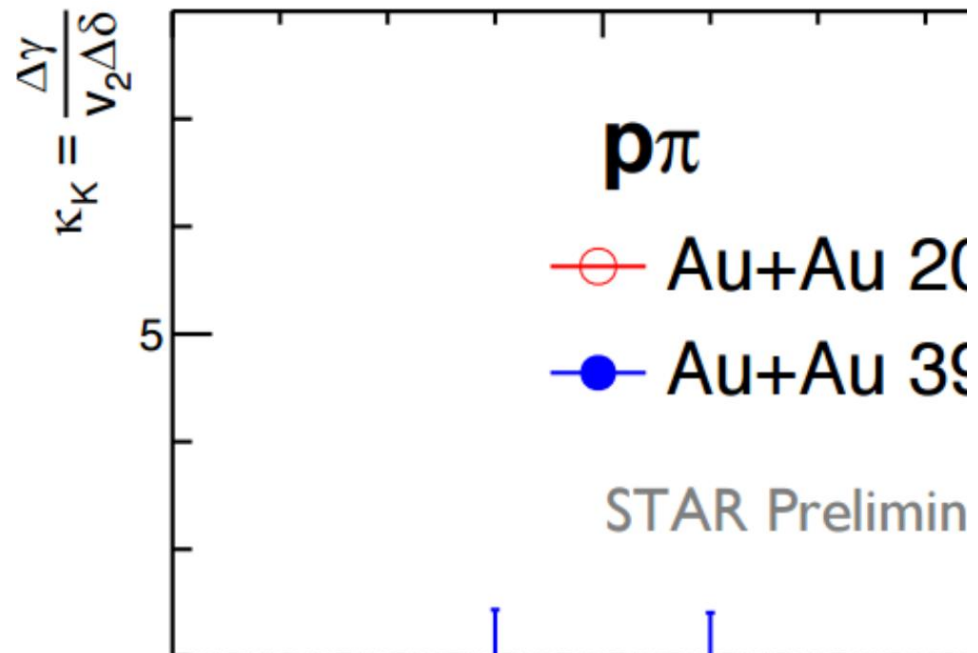
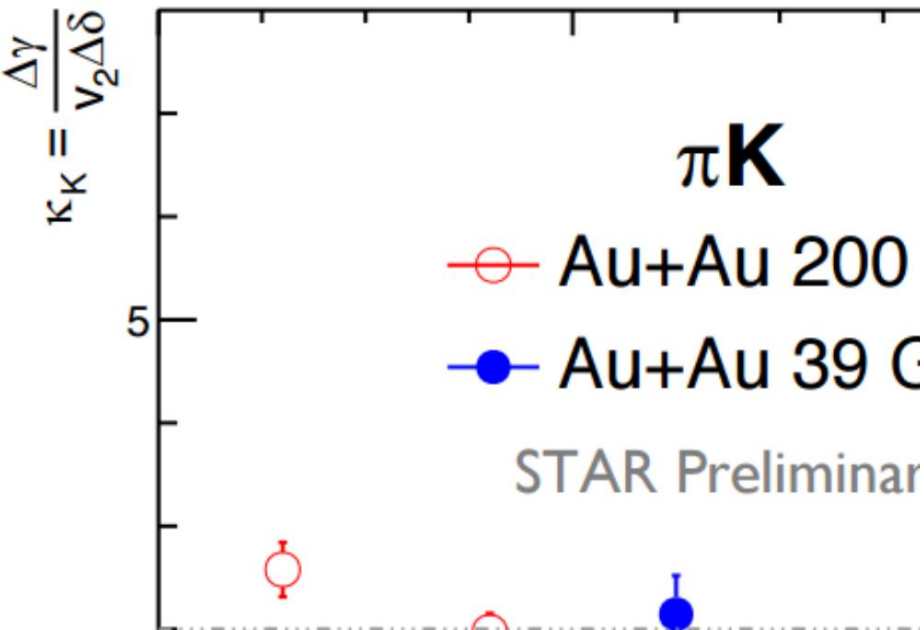
PID $\kappa_K: \pi-\pi$

- $\Delta\gamma$ for $\pi-\pi$ is similar to $h-h$ in Au+Au at 200 GeV.
 - also 39 GeV (not shown here).
- $\kappa_K > \kappa_B$ for mid-central and mid-peripheral collisions.
 - κ_B estimated from AMPT



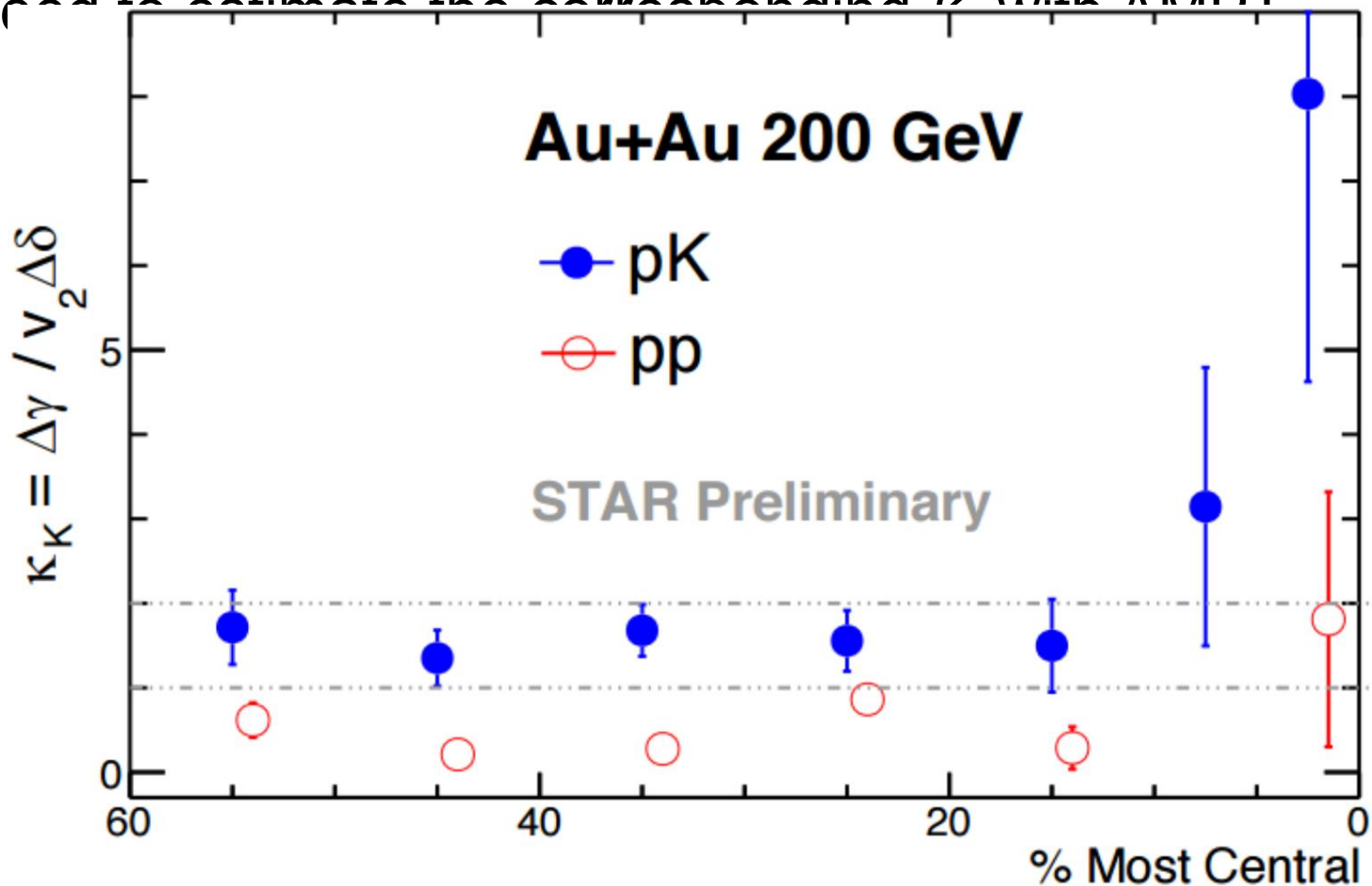
PID κ_K : π -K and p - π

- κ_K for π -K and p - π are mostly between 1 and 2 in Au+Au at 200.
 - also true at 39 GeV
- hard to distinguish the observable from the flow background



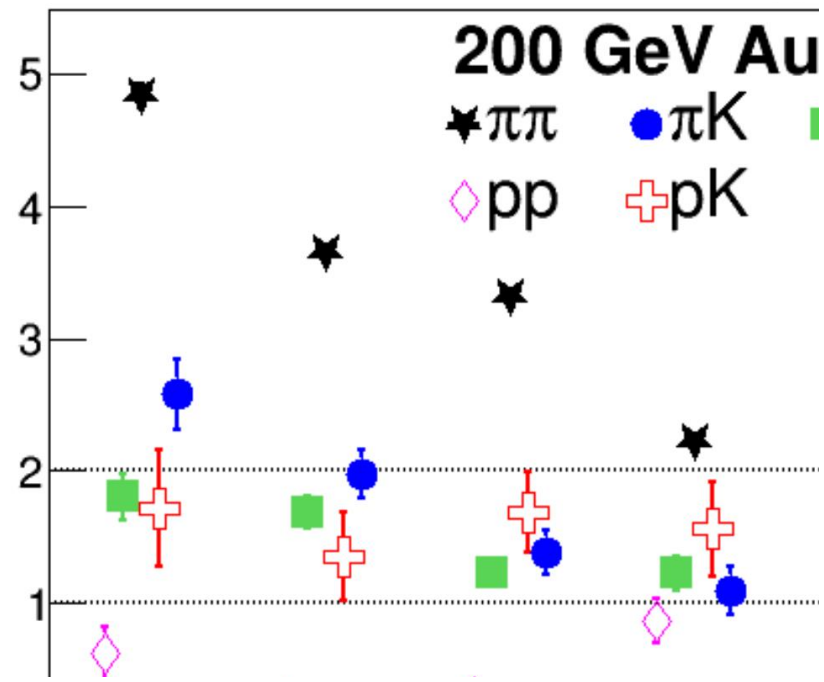
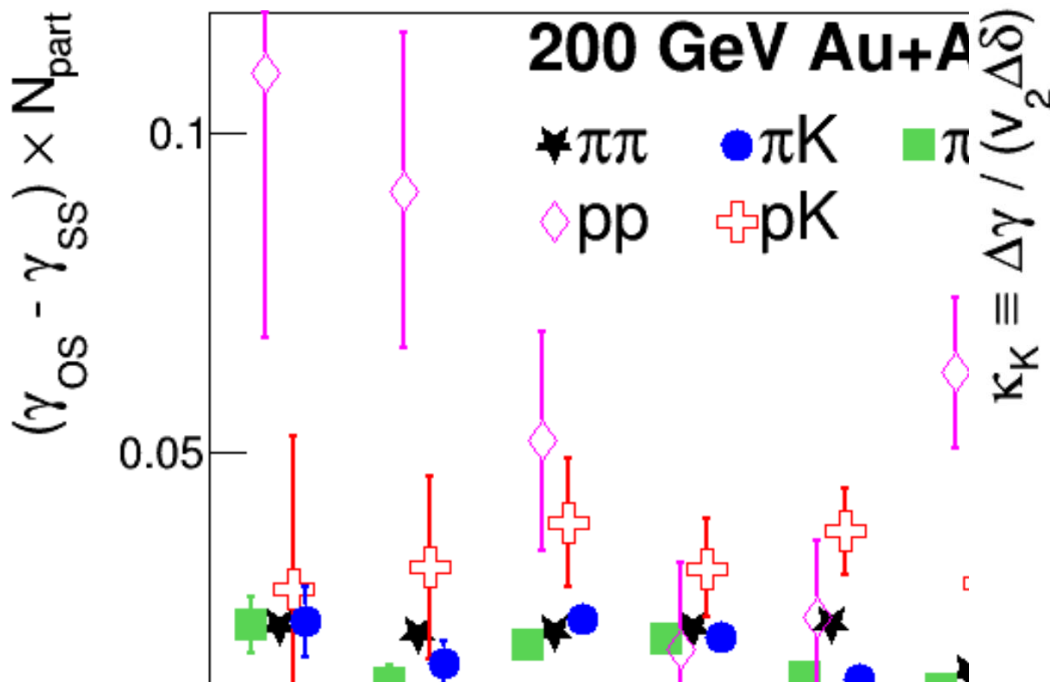
PID κ_K : p - K and p - p

- κ_K for p - K is similar to the cases for π - K and p - π in Au+Au at 200.
 - hard to distinguish the observable from the flow background
- κ_K for p - p is even below 1.
 - need to estimate the corresponding v_2 with AMBT



Summary on PID

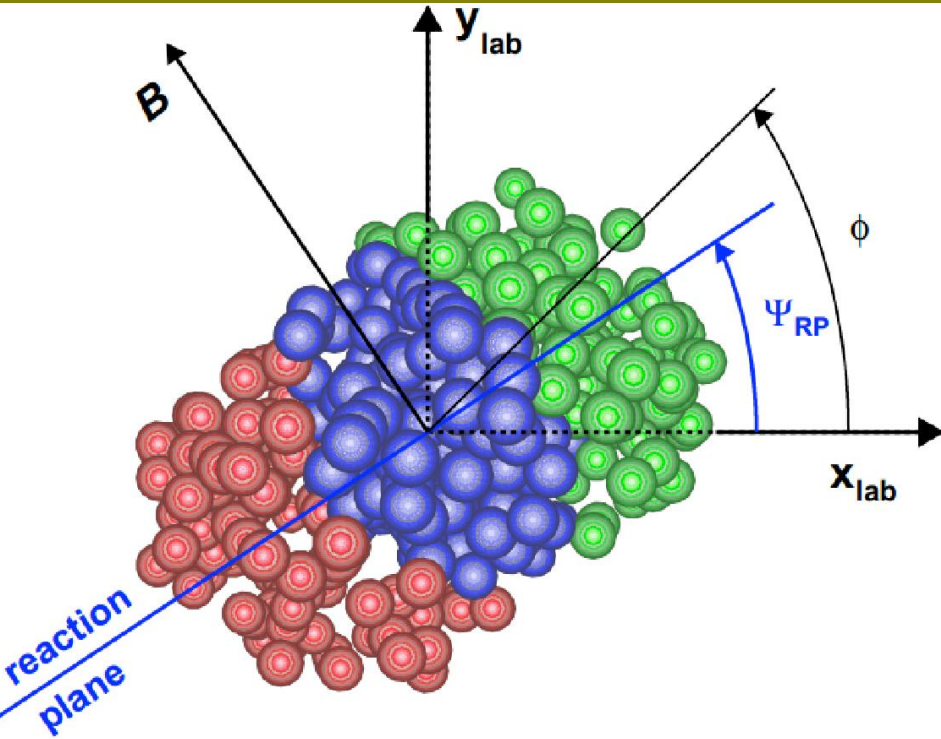
- $\Delta\gamma$ for π - π / π - K / p - π / p - K / p - p all show sizable charge separation signals for mid-central and mid-peripheral collisions at 200 GeV.
- κ_K values, however, fall into 3 groups.
 - π - π : higher than 2 for 20-60% most central collisions
 - π - K / p - π / p - K : between 1 and 2
 - p - p : below 1





Back-up slides

Event plane

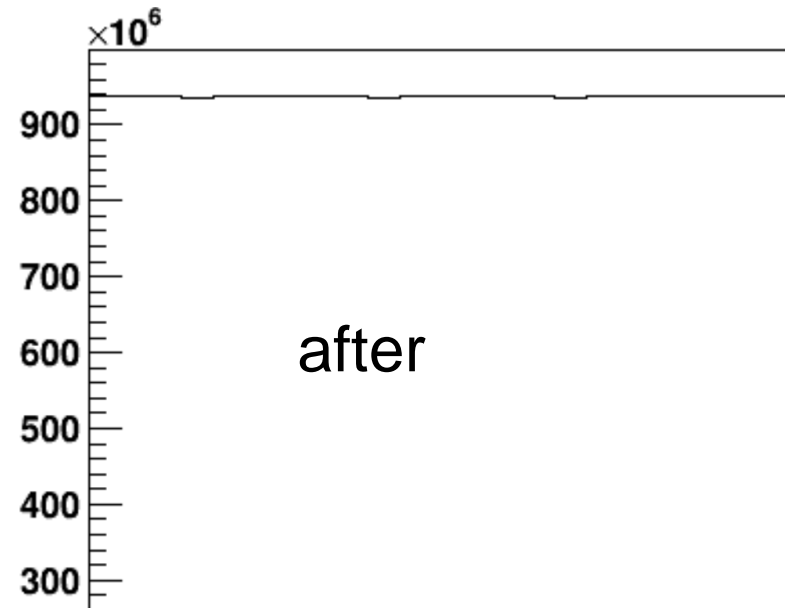
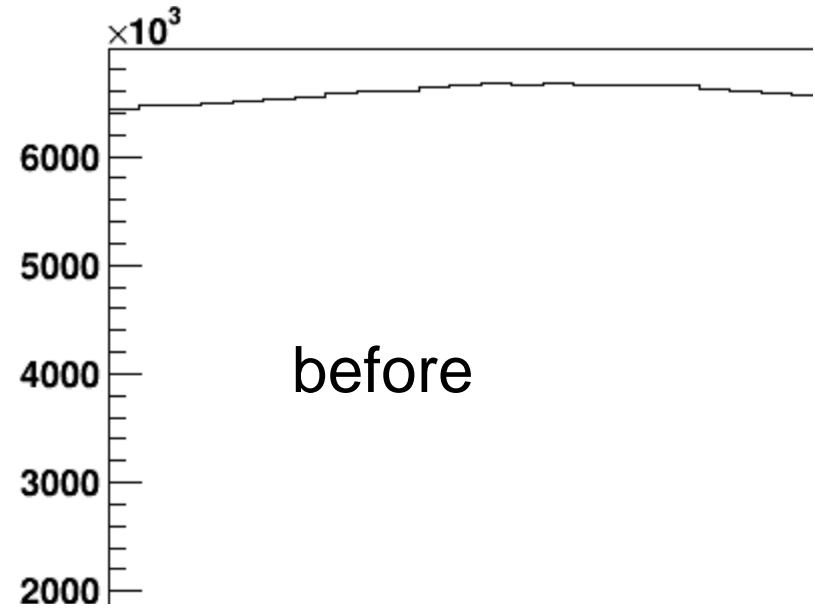


The estimated reaction plane is called the event plane.

$$Q_n \cos(n\Psi_n) = Q_x = \sum_i w_i \cos(n\phi_i)$$

$$Q_n \sin(n\Psi_n) = Q_y = \sum_i w_i \sin(n\phi_i)$$

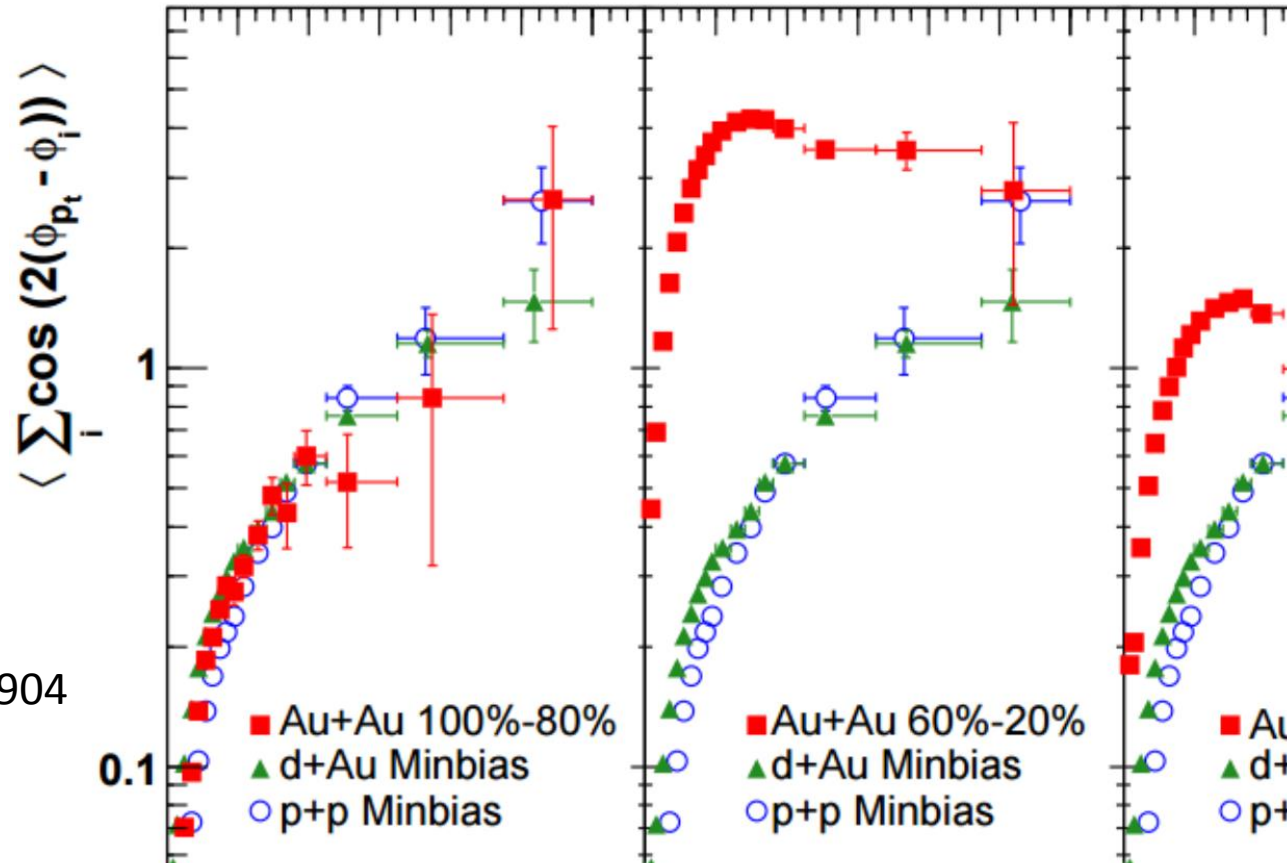
$$\Psi_n = \left(\tan^{-1} \frac{Q_y}{Q_x} \right) / n$$



Collectivity vs non-flow

What is collectivity?
A working definition:
multiple particles
correlated across
rapidity due to a
common source

STAR, Phys. Rev. C 72 (2005) 14904

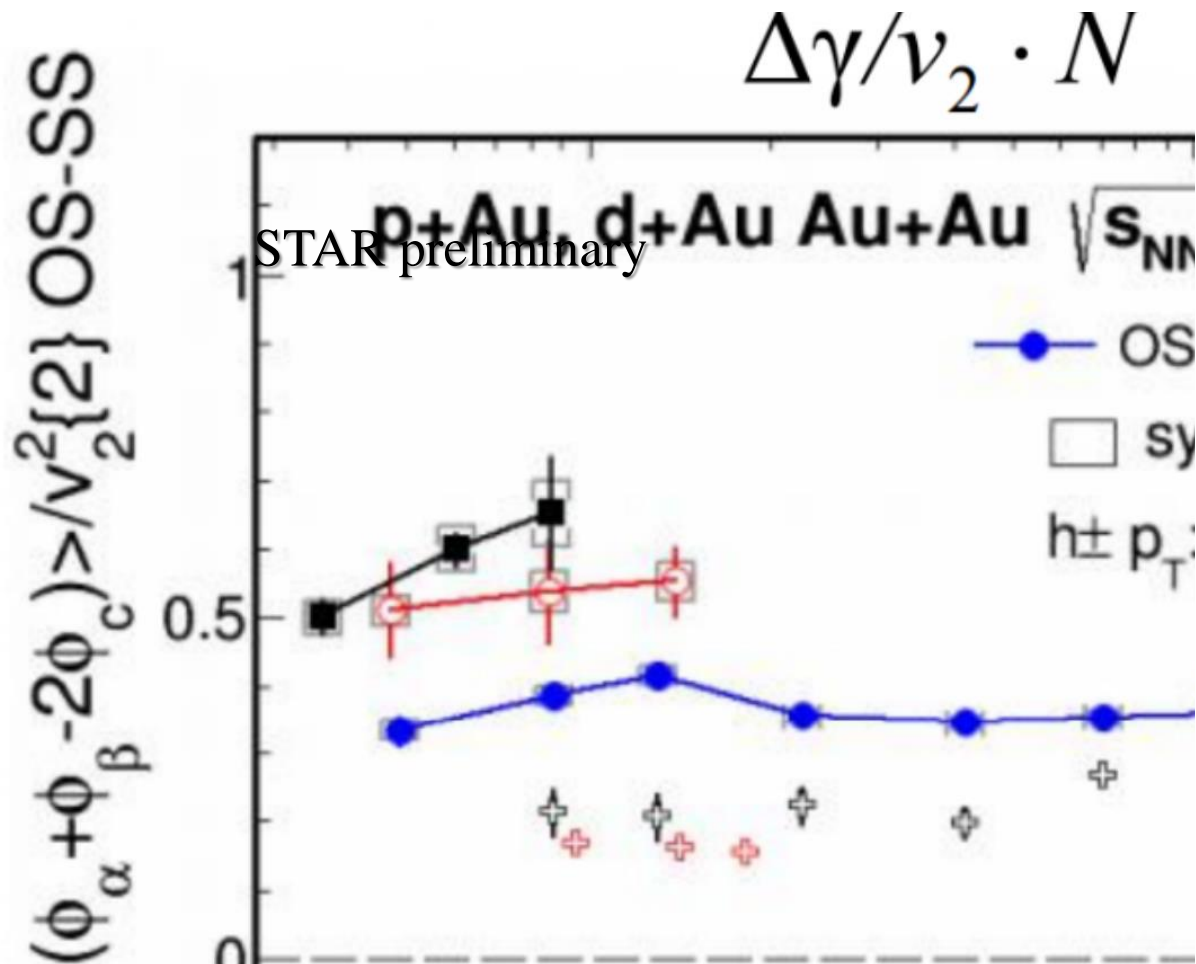


➤ Note 1: collectivity does not imply a specific physical interpretation (i.e. collectivity \neq hydro)

➤ Note 2: correlations between particles which do not have a “collective” origin (jets, resonance decays, momentum conservation) are commonly called “non-flow”...

p+Au/d+Au

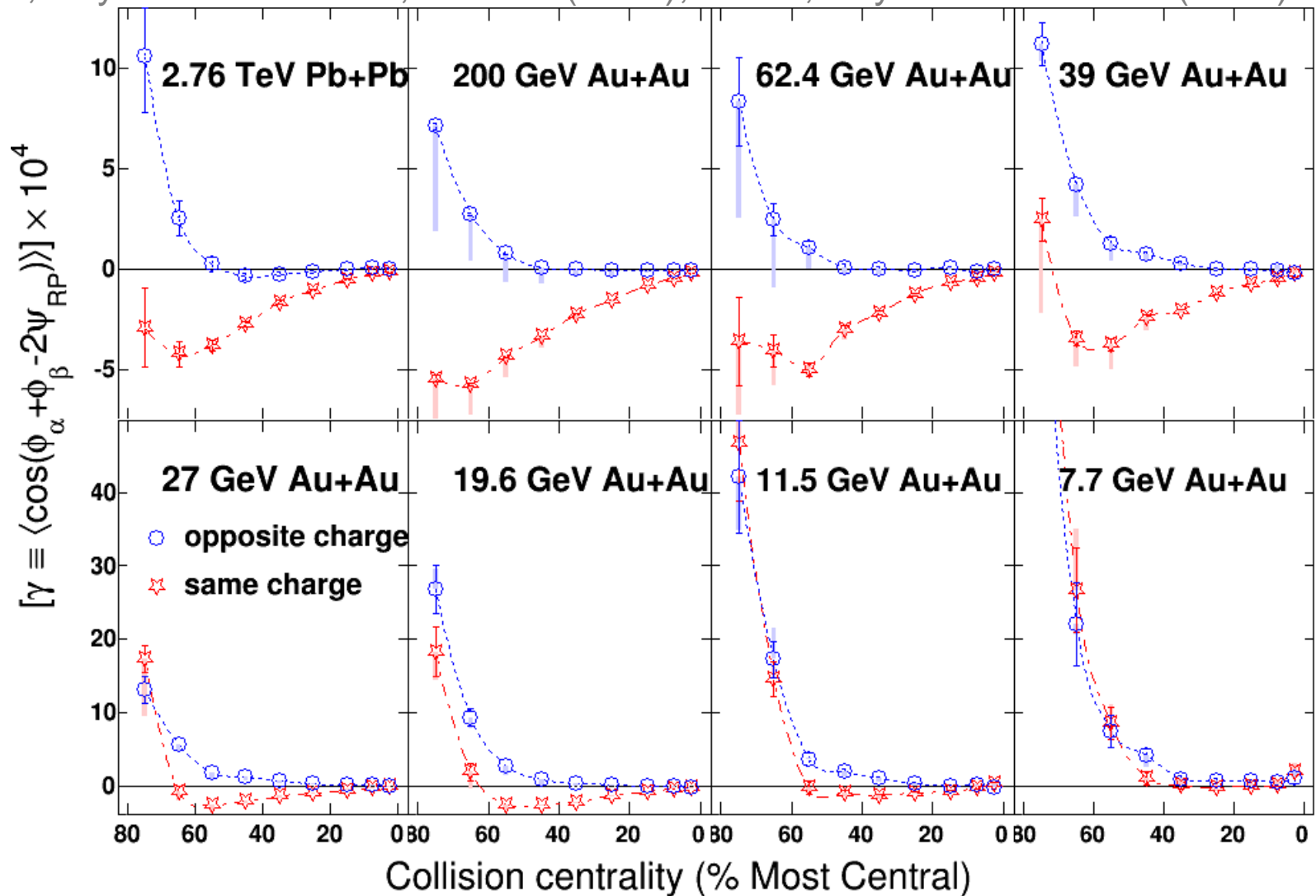
- $\Delta\gamma/v_2 \cdot N$ is another measure for the flow background.
- AMPT accounts for $\sim 2/3$ (1/3) of the observed data in Au+Au (d+Au)
- Central Au+Au events are not well described by AMPT.



This version of AMPT turned off hadronic scattering.

Beam energy scan

ALICE, Phys. Rev. Lett. 110, 012301 (2013); STAR, Phys. Rev. Lett 113 (2014) 052302

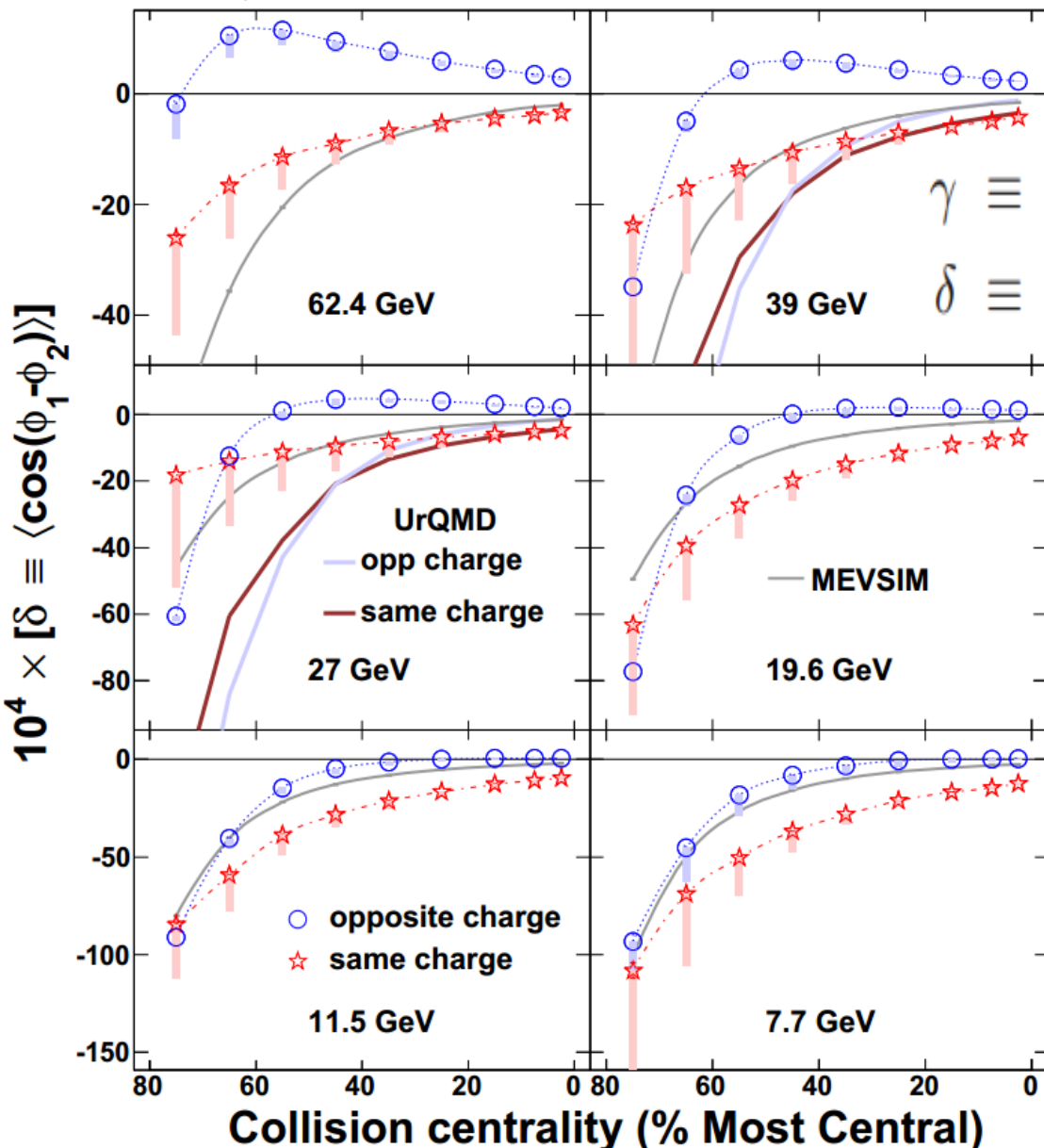


At lower beam energies, charge separation starts to diminish.

Flow-related background

STAR, Phys. Rev. Lett 113 (2014) 052302

A. Bzdak, V. Koch and J. Liao,
Lect. Notes Phys. 871, 503 (2013).



$$\gamma \equiv \langle \cos(\phi_1 + \phi_2 - 2\Psi_{\text{RP}}) \rangle = \kappa v_2 F - H$$

$$\delta \equiv \langle \cos(\phi_1 - \phi_2) \rangle = F + H,$$

- Against CME expectation, δ_{OS} is above δ_{SS}
- indicate overwhelming background larger than any possible CME effect.
- try to combine information from γ and δ to retrieve the CME contribution, H

Transverse momentum conservation

$$\gamma = -\frac{1}{N_{\text{tot}}} \frac{\langle p_t \rangle_{\Omega}^2}{\langle p_t^2 \rangle_F} \frac{2\bar{v}_{2,\Omega} - \bar{\bar{v}}_{2,F} - \bar{\bar{v}}_{2,F} (\bar{v}_{2,\Omega})^2}{1 - (\bar{\bar{v}}_{2,F})^2},$$

A. Bzdak, V. Koch and J. Liao, Lect. Notes Phys. 871, 503 (2013).

$$\delta = -\frac{1}{N_{\text{tot}}} \frac{\langle p_t \rangle_{\Omega}^2}{\langle p_t^2 \rangle_F} \frac{1 + (\bar{v}_{2,\Omega})^2 - 2\bar{\bar{v}}_{2,F} \bar{v}_{2,\Omega}}{1 - (\bar{\bar{v}}_{2,F})^2},$$

we have introduced certain weighted moments of v_2 :

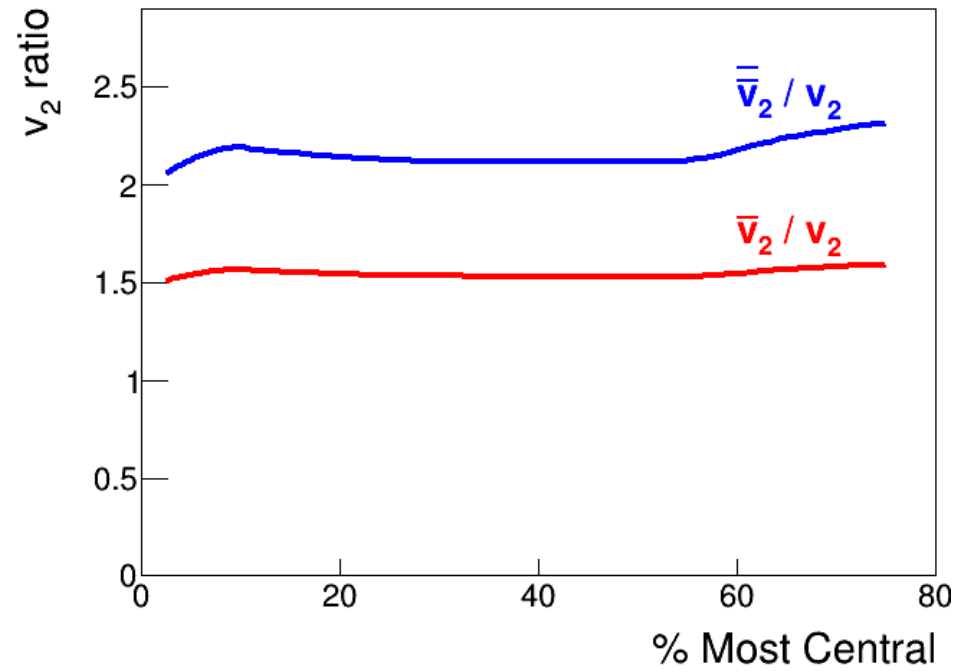
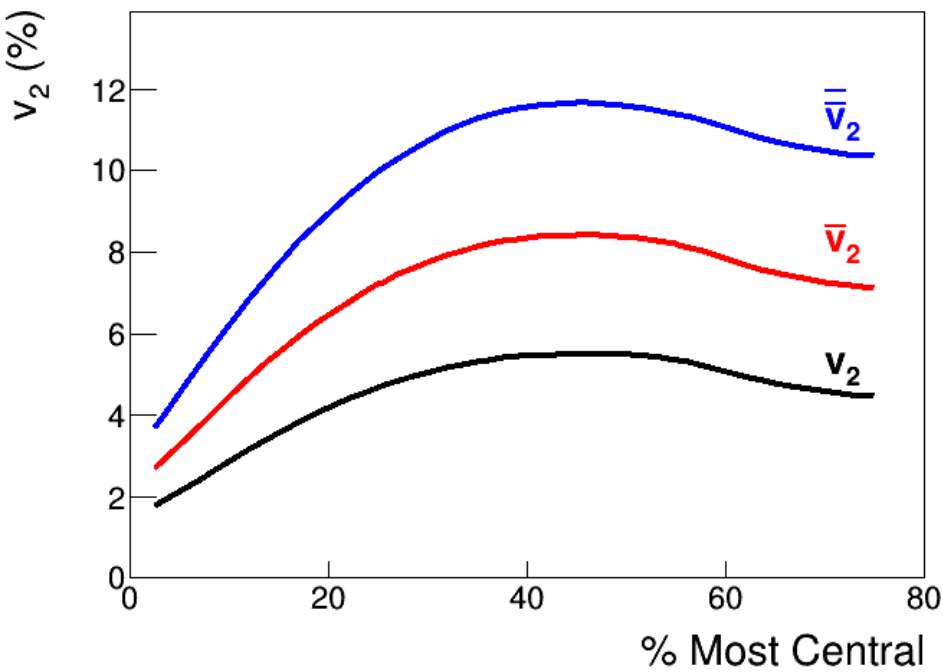
$$\bar{v}_2 = \frac{\langle v_2(p_t, \eta) p_t \rangle}{\langle p_t \rangle}, \quad \bar{\bar{v}}_2 = \frac{\langle v_2(p_t, \eta) p_t^2 \rangle}{\langle p_t^2 \rangle}.$$

If our measurements are dominated by this type of background,

$$\gamma / \delta \approx 2\bar{v}_{2,\Omega} - \bar{\bar{v}}_{2,F}$$

where F and Ω denote particle averages in the full phase-space and the detector acceptance, respectively.

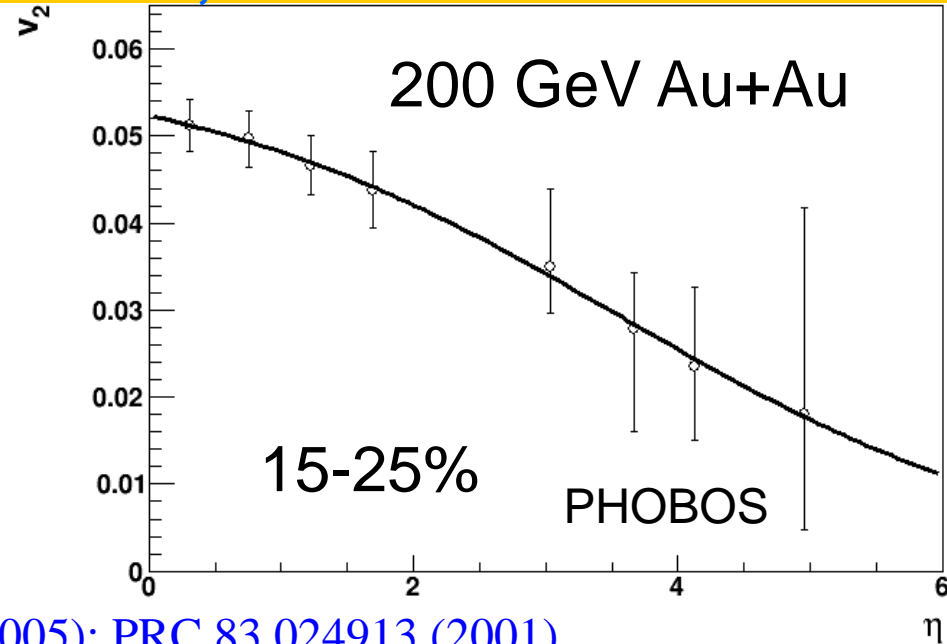
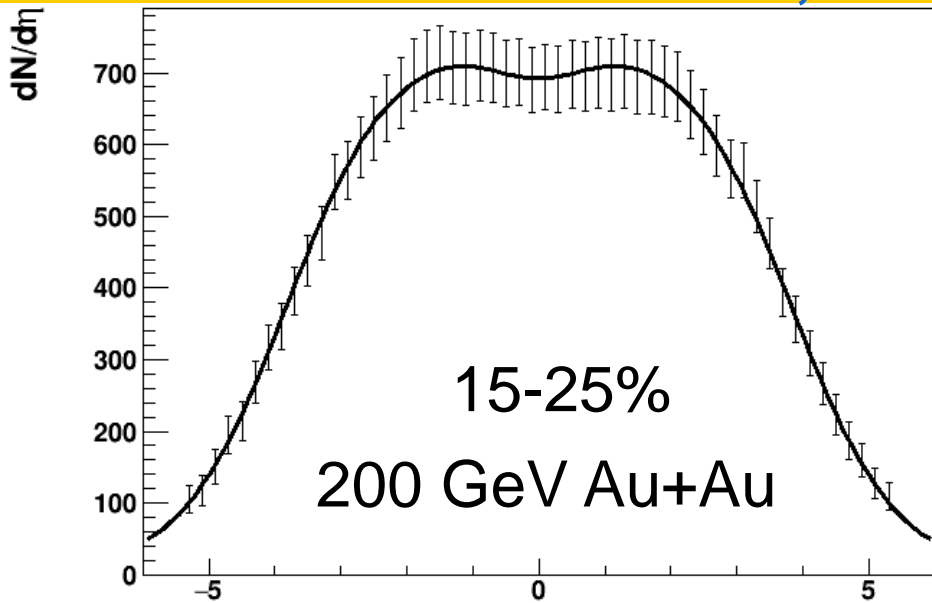
v_2 , \bar{v}_2 and $\bar{\bar{v}}_2$



$$\bar{v}_2 = \frac{\langle v_2(p_t, \eta) p_t \rangle}{\langle p_t \rangle}, \quad \bar{\bar{v}}_2 = \frac{\langle v_2(p_t, \eta) p_t^2 \rangle}{\langle p_t^2 \rangle}$$

The ratios of the p_T -weighted v_2 over conventional v_2 are almost constant over centrality.

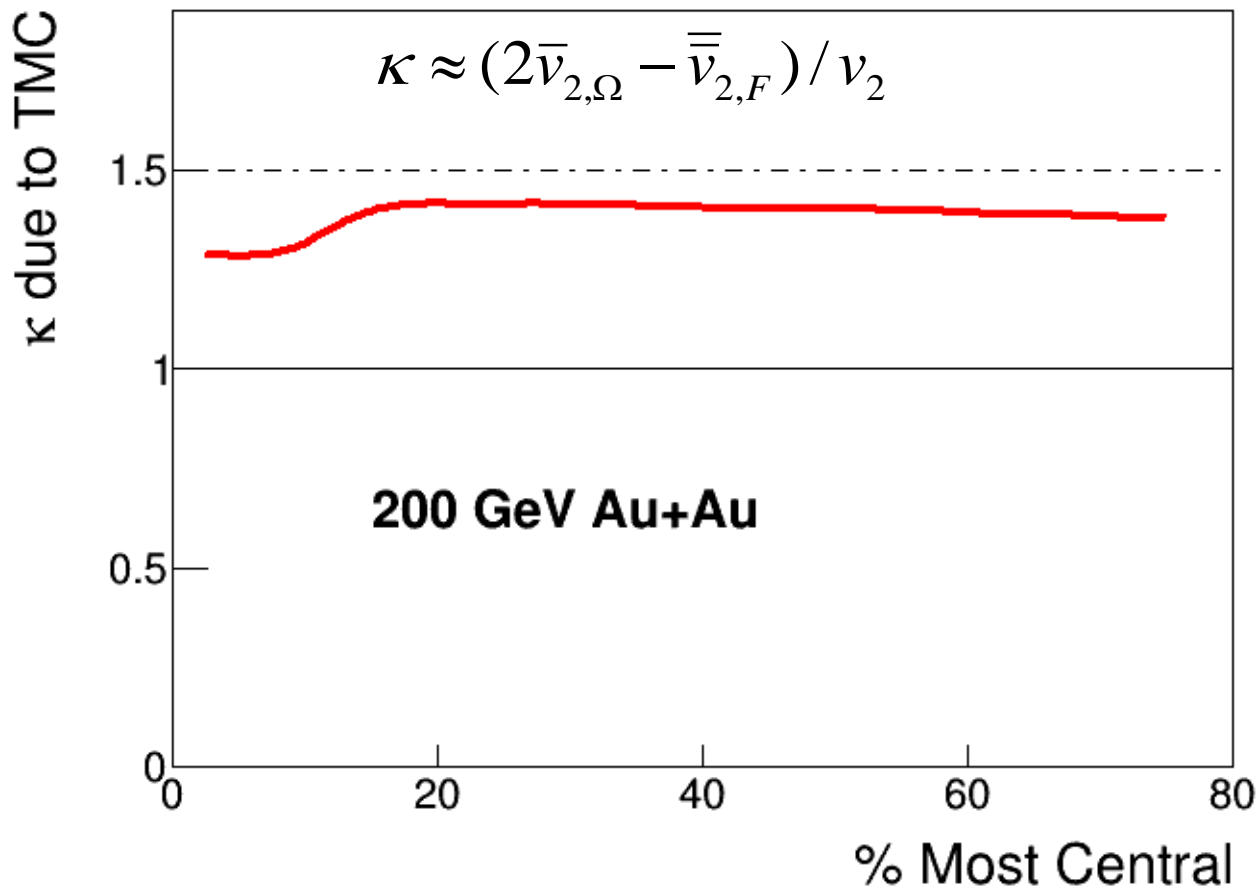
$v_{2,\Omega}$ and $v_{2,F}$



PHOBOS, PRC 72 014904 (2005); PRC 83 024913 (2001)

centrality	$v_{2,\Omega}$ (%)	$v_{2,F}$ (%)	$v_{2,F}/v_{2,\Omega}$
3-15%	3.17	2.66	0.84
15-25%	5.04	3.97	0.79
25-50%	6.21	4.87	0.78

κ due to TMC



$$\gamma \equiv \langle \cos(\phi_1 + \phi_2 - 2\Psi_{\text{RP}}) \rangle = \kappa v_2 F - H$$
$$\delta \equiv \langle \cos(\phi_1 - \phi_2) \rangle = F + H,$$

Other effects: Local Charge Conservation (LCC) and resonance decay.
A rough estimate of κ in the next slide.