

A new Dirac neutrino mass model with dark matter and electroweak baryogenesis

In preparation

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Tiny neutrino masses

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- It has been established that **neutrinos have tiny masses** because of the discovery of neutrino oscillations.

$$\Delta m_{21}^2 = 7.46 \times 10^{-5} \text{eV}^2 \quad \text{SNO Collaboration, PRC 88 (2013) 025501}$$

$$|\Delta m_{32}^2| = +2.51 \times 10^{-3} \text{eV}^2 \quad \text{T2K Collaboration, PRD91 (2015) 072010} \quad \Delta m_{ij}^2 = m_i^2 - m_j^2$$

- In the Standard Model (SM), Quarks and Leptons obtain masses after electroweak symmetry breaking, but neutrino mass is zero.

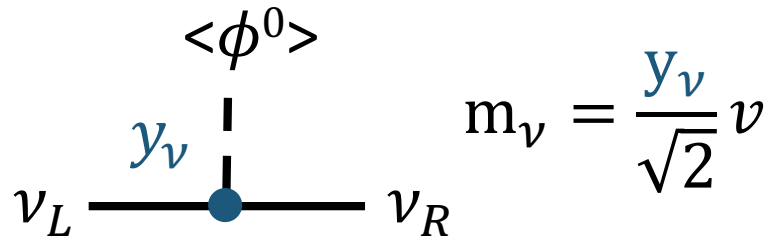
$$\frac{y_f}{\sqrt{2}} \bar{f}_L h f_R \rightarrow \frac{y_f}{\sqrt{2}} \langle h \rangle \bar{f}_L f_R$$

The SM don't include ν_R .

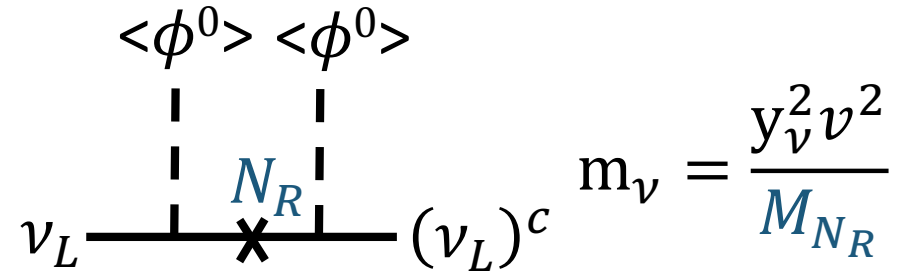
↳ What mechanism are neutrino masses m_ν generated by ?

Some mechanisms generating m_ν

① The same mechanism with quarks and leptons (Dirac mass)

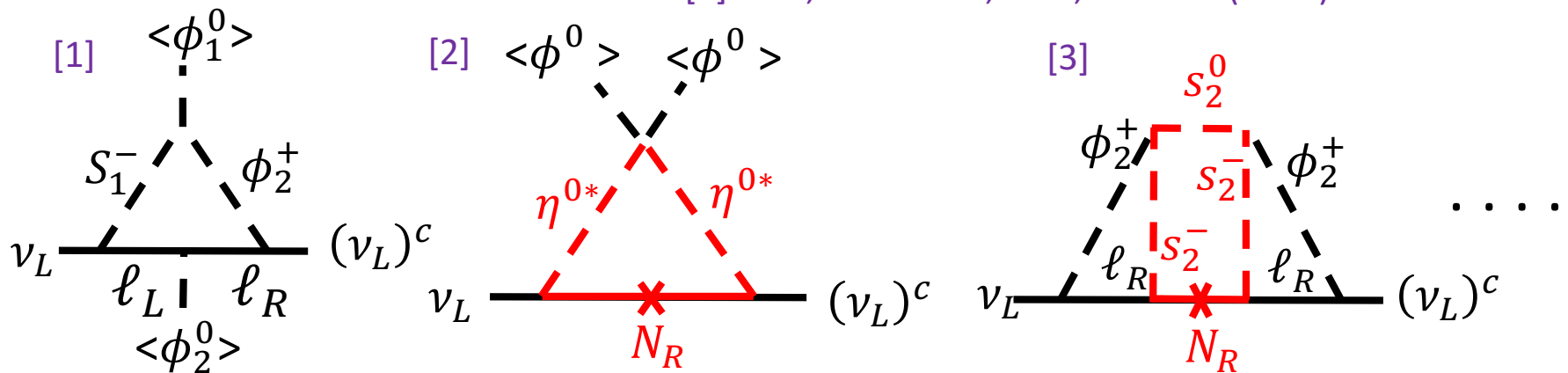


② Seesaw mechanism



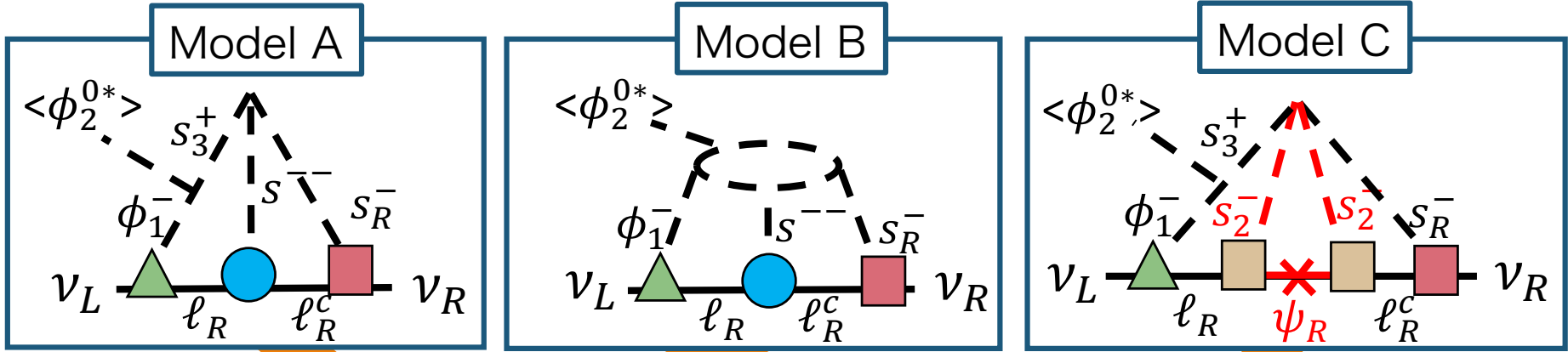
③ Radiative seesaw mechanism

[1] Zee, PLB93 (1980) 389. [2] Ma, PRD 73 (2006) 077301.
[3] Aoki, Kanemura, Seto, PRL 102 (2009) 051805.



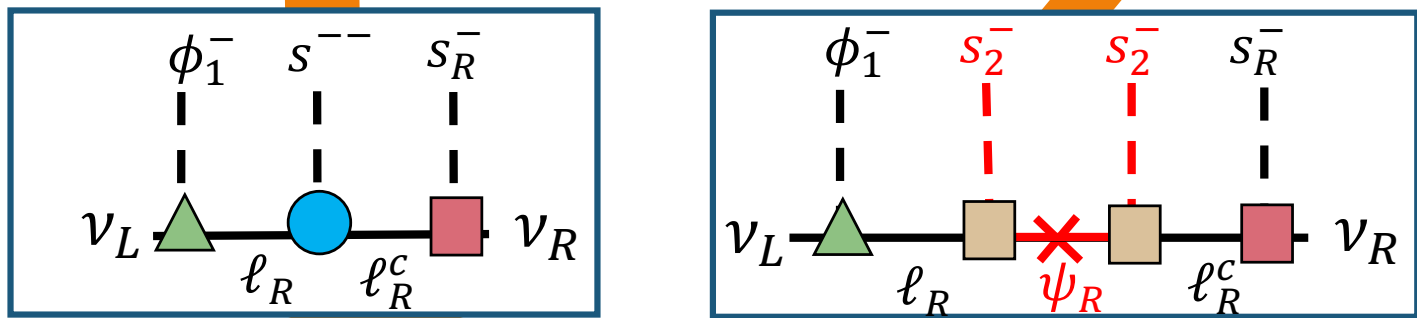
- There are many many models .
- In order to efficiently test models of m_ν , we classified the models into some groups, then, examined testability of the groups.

Classifications by flavor structure of m_ν



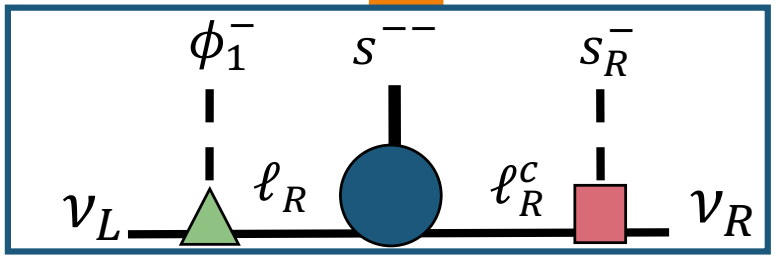
Level 1: Classification by combination of Yukawa matrices

A and B are same.

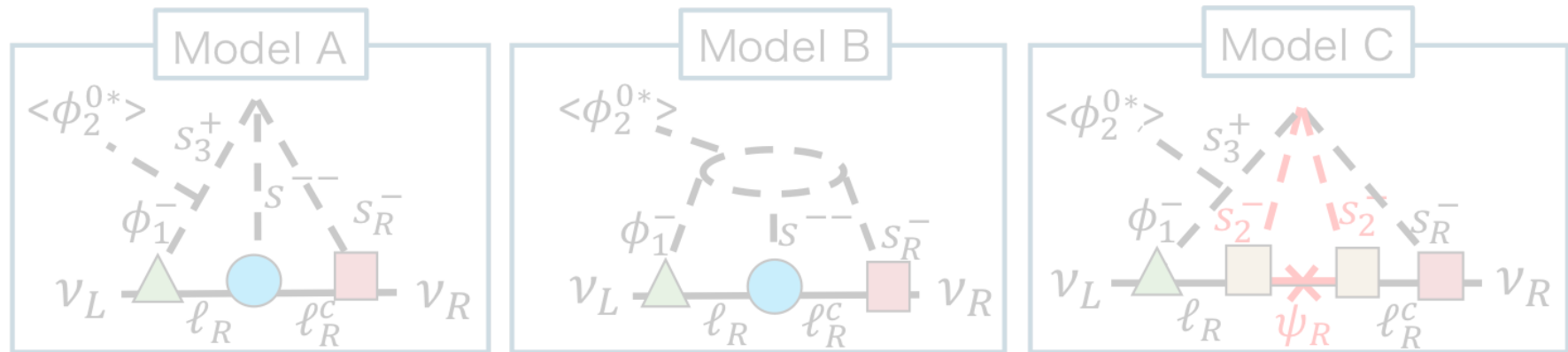


Level 2: Classification by a structure of flavor matrices

A, B and C are same.



Classifications by flavor structure of m_ν

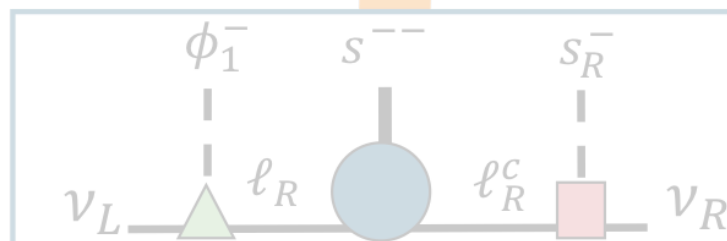


Level 1: Classification by combination of Yukawa matrices

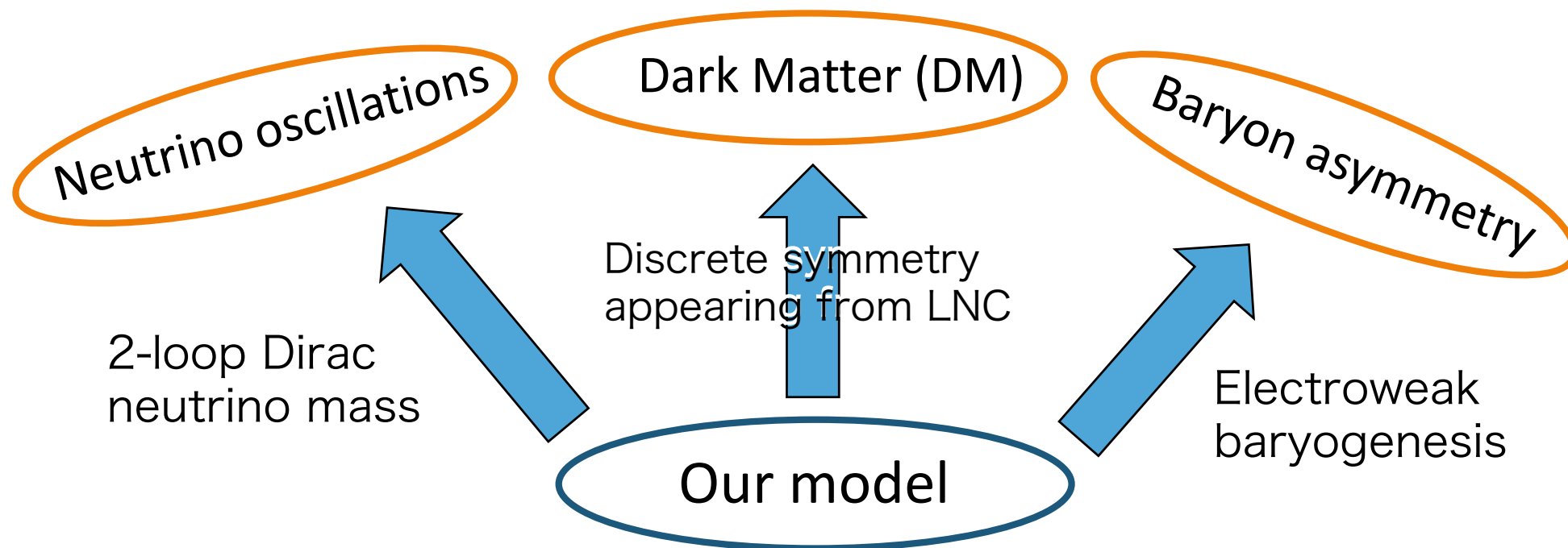
We discuss details of a new model which was found by our previous studies.

Level 2: Classification by a structure of flavor matrices

A, B and C are same.



Our model



LNC : Lepton Number Conservation

We discuss whether or not there is a benchmark scenario which can simultaneously explain three Beyond the SM (BSM) phenomena.

Setup

- New field in addition to SM field

Higgs doublet : Φ_1, Φ_2

Scalar singlet : s_1^0, s_2^+, s_3^+

Gauge singlet fermion : $\psi_{R1}, \psi_{R2}, \psi_{R3}$

Right-handed neutrino : $\nu_{R1}, \nu_{R2}, \nu_{R3}$

$$\Phi_{1,2} = \begin{pmatrix} \omega_{1,2}^+ \\ \frac{1}{\sqrt{2}}(v_{1,2} + \phi_{1,2}^0 + iz_{1,2}) \end{pmatrix}$$

Red character denote Z_2 odd particles.

$$Z_2 \text{ odd : } \psi_{Ra} \rightarrow -\psi_{Ra}, s_1^0 \rightarrow -s_1^0, s_2^+ \rightarrow -s_2^+$$

$$Z_2 \text{ even : } \Phi_{1,2} \rightarrow +\Phi_{1,2}, s_3^+ \rightarrow +s_3^+, \nu_R \rightarrow +\nu_R$$

- Symmetries: $Z'_2 \times U(1)_L \times Z_2$

Field	$SU(2)_L$	$U(1)_Y$	L#	Z'_2	Z_2
ν_R	<u>1</u>	0	1	Odd	Even
Φ_2	<u>2</u>	1/2	0	Even	Even
s_3^+	<u>1</u>	1	0	Even or Odd	Even
s_1^0	<u>1</u>	0	-1	Odd	Odd
s_2^+	<u>1</u>	1	-1	Even	Odd
ψ_R	<u>1</u>	0	0	Even	Odd

$U(1)_L$: Lepton number conservation

Z'_2 : Prohibition for $Y_\nu \bar{L} \tilde{\Phi} \nu_R$

Z_2 : Stability for DM

- Physical states: $\nu_1, \nu_2, \nu_3, H_1^\pm, H_2^\pm, H, A, s_2^\pm, s_1^0, \psi_{R1}, \psi_{R2}, \psi_{R3}$

the lightest Z_2 odd particle is dark matter candidate

Lagrangian

- Scalar potential : we consider the CP conserving potential

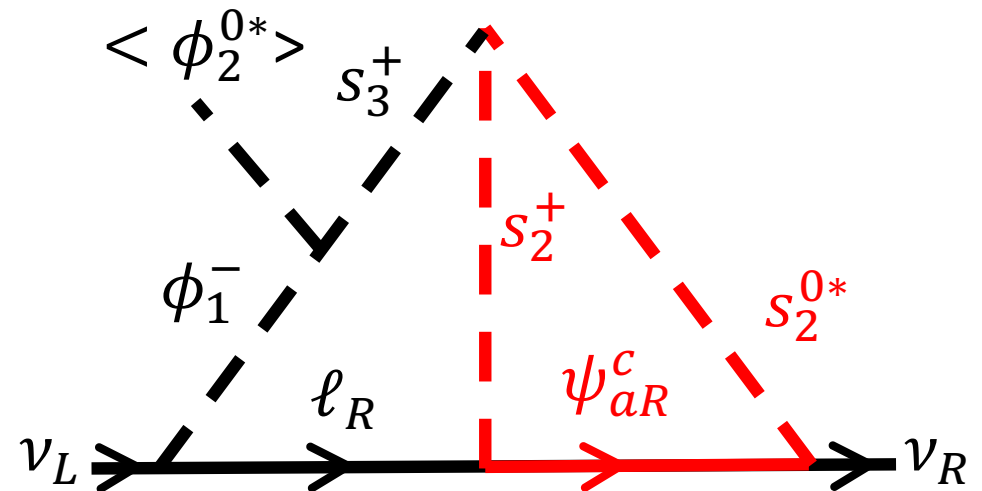
$$\begin{aligned}
 V = & \sum_{i=1}^2 \left(m_i^2 |\Phi_i|^2 + \frac{\lambda_i}{2} |\Phi_i|^4 \right) - m_3^2 (\Phi_1^\dagger \Phi_2 + \text{h.c.}) \\
 & + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{1}{2} \lambda_5 [(\Phi_1^\dagger \Phi_2)^2 + \text{h.c.}] \quad \left. \vphantom{\sum_{i=1}^2} \right\} \text{Two Higgs doublet model} \\
 & + \sum_{i=1}^3 \left(m_{s_i}^2 |s_i|^2 + \frac{\lambda_{s_i}}{2} |s_i|^4 \right) + \sum_{i < j}^3 \lambda_{s_i s_j} |s_i|^2 |s_j|^2 \quad \left. \vphantom{\sum_{i=1}^3} \right\} \text{scalar singlet} \\
 & + \left(\mu_3 \Phi_2^\dagger \epsilon \Phi_1^* s_3^+ + \text{h.c.} \right) + \left(\mu'_3 s_3^- s_2^+ s_1^0 + \text{h.c.} \right) \\
 & + \sum_{i=1}^3 \left(\lambda_{\phi 1 s_i} |\Phi_1|^2 |s_i|^2 + \lambda_{\phi 2 s_i} |\Phi_2|^2 |s_i|^2 \right) \quad \left. \vphantom{\sum_{i=1}^3} \right\} \text{Interaction terms}
 \end{aligned}$$

- Yukawa interactions: Additional Z_2 symmetry is imposed to prohibit FCNC.

$$\begin{aligned}
 \mathcal{L}_Y = & Y_\ell \bar{L} \Phi_1 \ell_R + Y_u \bar{Q} \Phi_2 u_R + Y_d \bar{Q} \Phi_2 d_R \quad \left. \vphantom{Y_\ell} \right\} \text{Type-X} \\
 & + Y_\psi^0 \left[\overline{(\nu_{iR})^c} \psi_{jR}^0 s_1^0 \right] + Y_\psi^+ \left[\overline{(\ell_R)^c} \psi_{iR}^0 s_2^+ \right] + \text{h.c.} \quad \left. \vphantom{Y_\psi^0} \right\} \text{New Yukawa}
 \end{aligned}$$

Neutrino mass

$$m_{\nu il} = (Y_{\psi}^{0*})_{ib} (Y_{\psi}^{+T})_{bl} \mu'_3 \tan \beta \frac{\sqrt{2} m_{\ell} \sin 2\theta_+}{v} \frac{1}{2} \\ \times \left(\frac{1}{16\pi^2} \right)^2 F(m_{H_1^{\pm}}, m_{H_2^{\pm}}, m_{\psi_b}, m_{s_1^0})$$



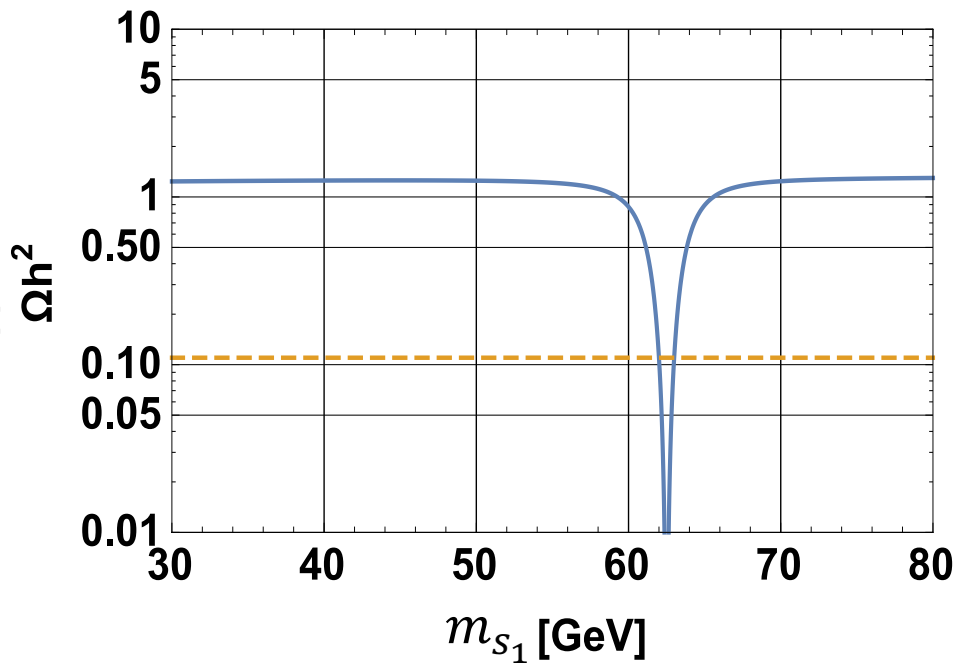
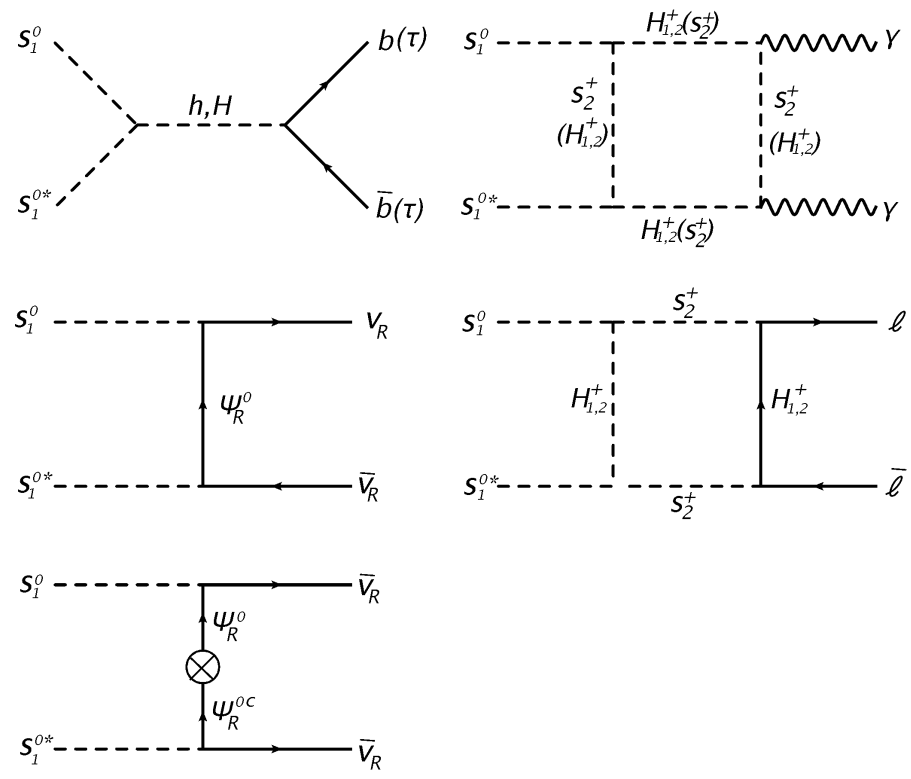
- m_{ν} is suppressed by 5 coupling constants + 2-loop suppression factor + a loop function F
- In order to satisfy the data of neutrino oscillations and constraint of LFV (Lepton Flavor Violation), the following Eq. is required

$$Y_{\psi}^{+} \sim O(0.01), Y_{\psi}^{0} \sim O(0.001), m_{\psi_{Ra}} \sim O(1) \text{TeV}$$

A mass of heaviest particles

Relic abundance for DM

- s_1^0 is DM candidate.
- We focus on a mass region $m_{DM} < 80$ GeV in which WW channel don't open.
- There are 5 annihilation processes :



$$\Omega h^2 = 0.1199 \pm 0.0020$$

Plank, arXiv:1502.01589

↳ $m_{s_1^0} = 62$ or 63 GeV are consistent with Plank exp. data.

Electroweakbaryogenesis^{12/16}

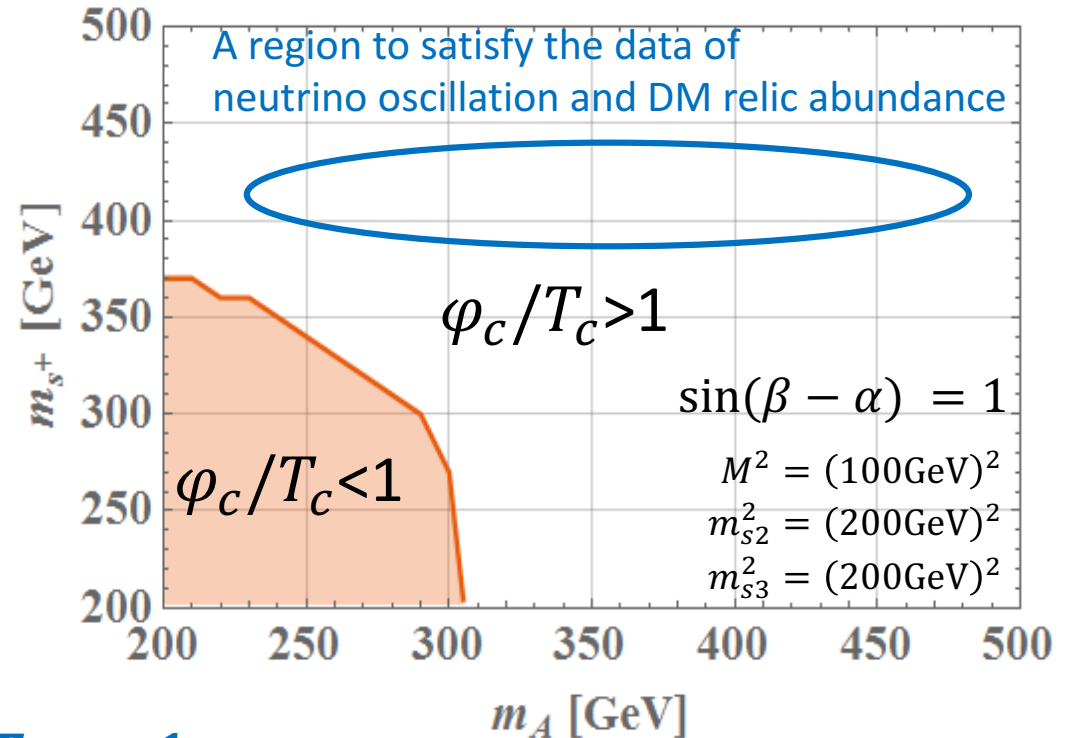
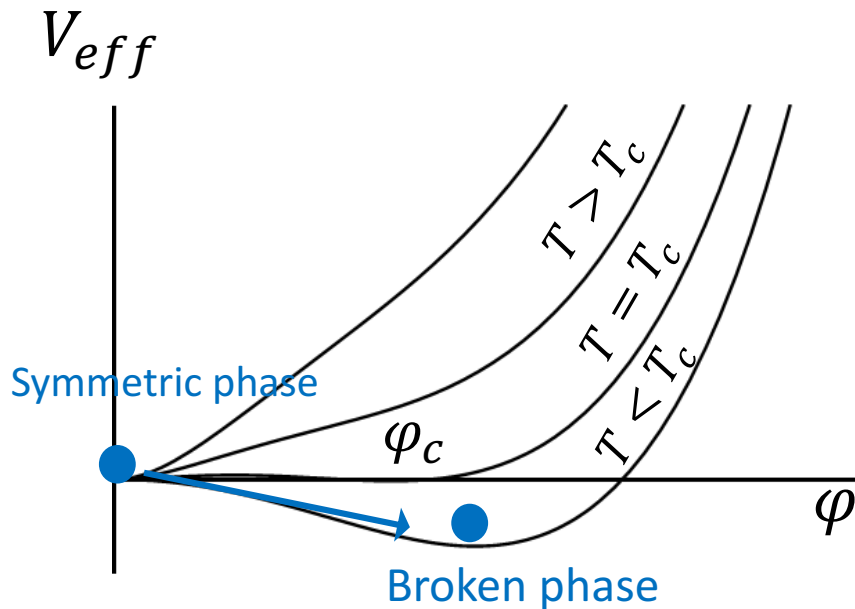
Sakharov's conditions

- B violation ← Sphaleron process
- C and CP violation ← CP phase in the Higgs sector
- Out of equilibrium ← Strong 1st Order phase Transition (1st OPT)

Electroweakbaryogenesis

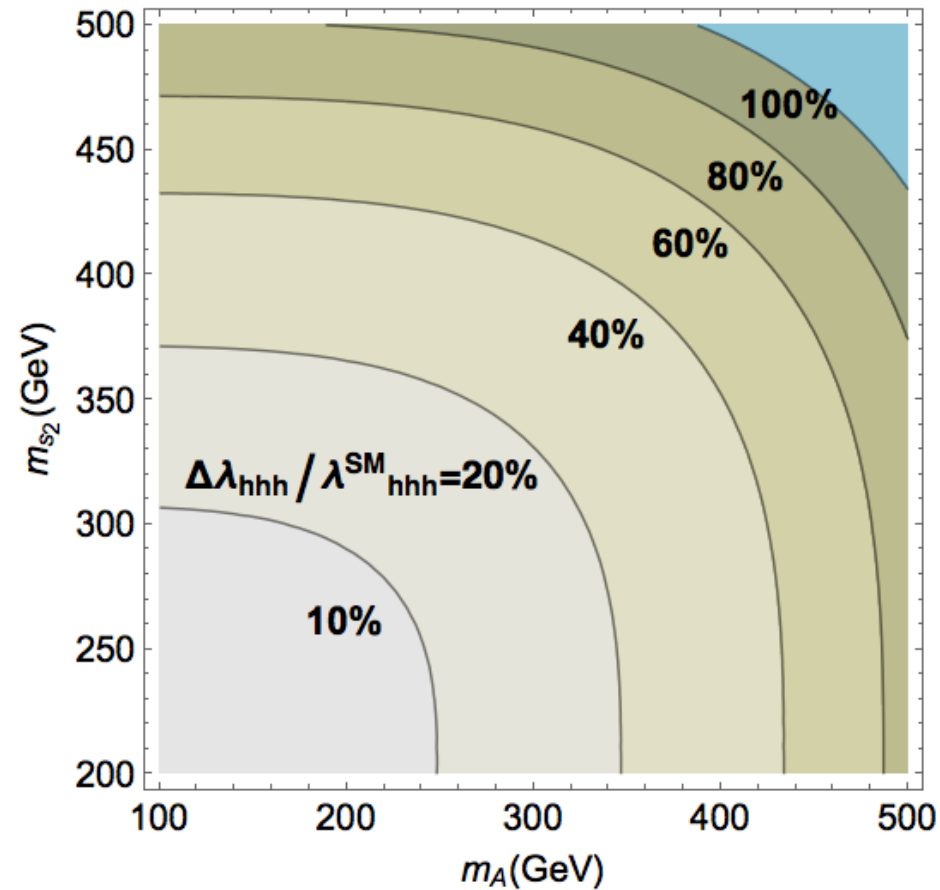
Sakharov's conditions

- B violation ← Sphaleron process
- C and CP violation ← CP phase in the Higgs sector
- Out of equilibrium ← **Strong 1st Order phase Transition (1st OPT)**



A condition to occur 1st OPT: $\varphi_c/T_c > 1$

Triple Higgs boson couplings ^{14/16}



$$\Delta\lambda_{hhh} = \lambda_{hhh} - \lambda_{hhh}^{SM}$$

- We also evaluated triple Higgs boson coupling λ_{hhh} .
- Large deviation from the SM in λ_{hhh} is occurred by non-decoupling effect of the additional scalar.

A successful scenario

$$\mu'_3 = 50\text{GeV}, \quad \sin\theta_+ = 0.1, \quad m_{\psi_{R1}} = m_{\psi_{R2}} = m_{\psi_{R3}} = 5\text{TeV},$$

$$m_{H_1^\pm} = 200\text{GeV}, \quad m_{H_2^\pm} = 300\text{GeV}, \quad m_{S_2^\pm} = 400\text{GeV},$$

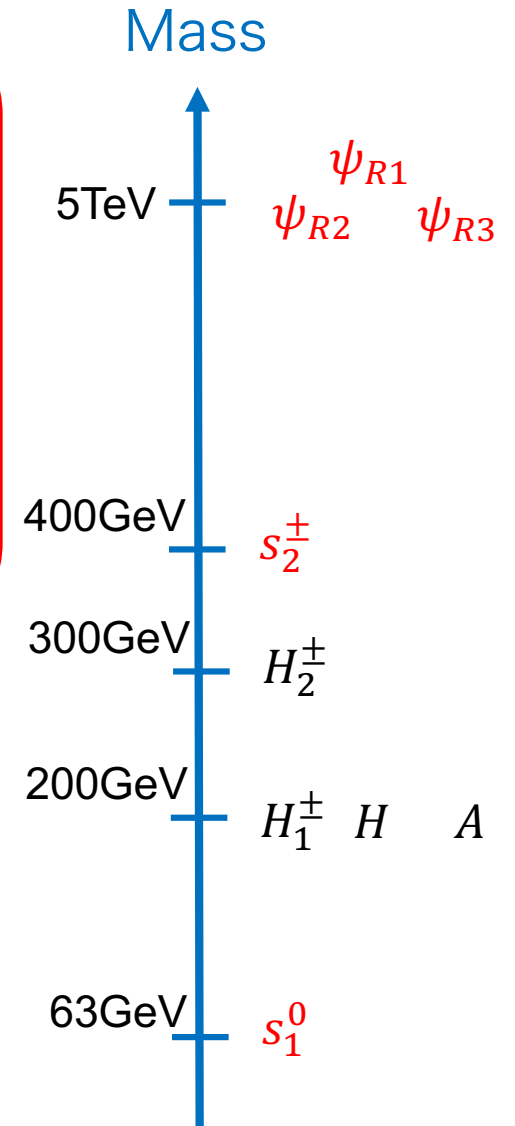
$$m_{H,A} = 200\text{GeV}, \quad m_{S_1^0} = 63\text{GeV},$$

$$Y_\psi^+ = \begin{pmatrix} 1 & 0.01 & 0.01 \\ 0.01 & 1 & 0.01 \\ 0.01 & 0.01 & 1 \end{pmatrix}, \quad Y_\psi^0 \cong \begin{pmatrix} 0.33 & -0.0041 & -0.0032 \\ 0.28 & -0.0014 & -0.0029 \\ 0.30 & 0.0040 & -0.0026 \end{pmatrix}$$

$$\sin(\beta - \alpha) = 1, \quad \tan\beta = 3, \quad \lambda_{\phi_{1S0}} = 0.02, \quad \lambda_{\phi_{2S0}} = 0.005,$$

$$M^2 = (100\text{GeV})^2, \quad m_{S_2^0}^2 = m_{S_3^0}^2 = (200\text{GeV})^2$$

- ✓ Neutrino oscillation
- ✓ DM relic density
- ✓ Strong 1st OPT
- ✓ LFV, LU
- ✓ DM direct detection
- ✓ S, T, U parameters



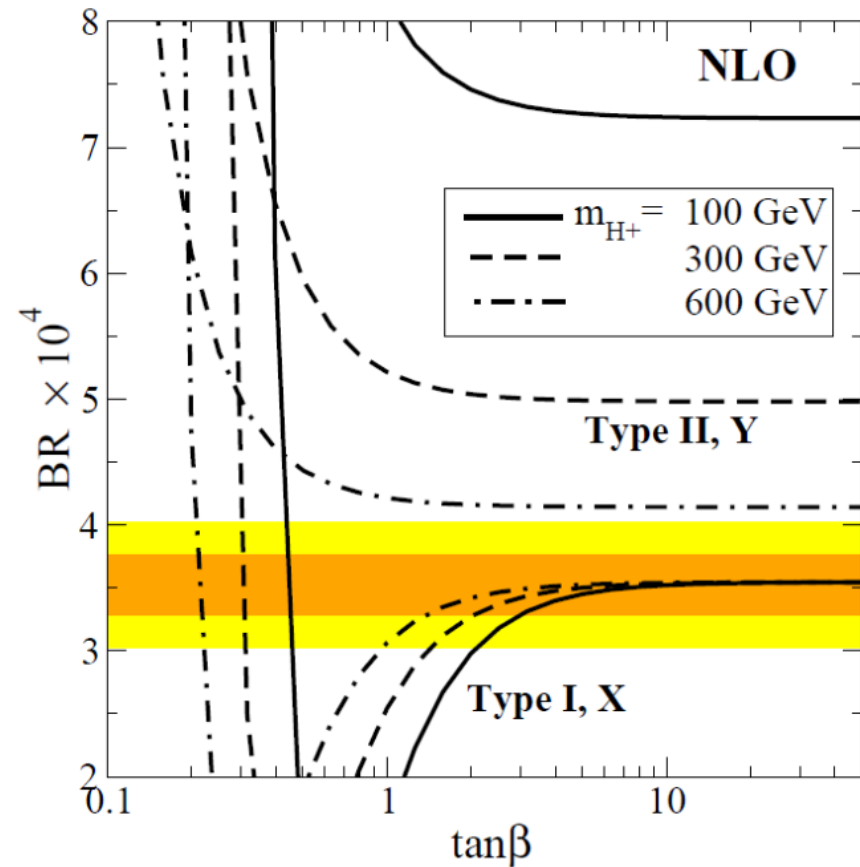
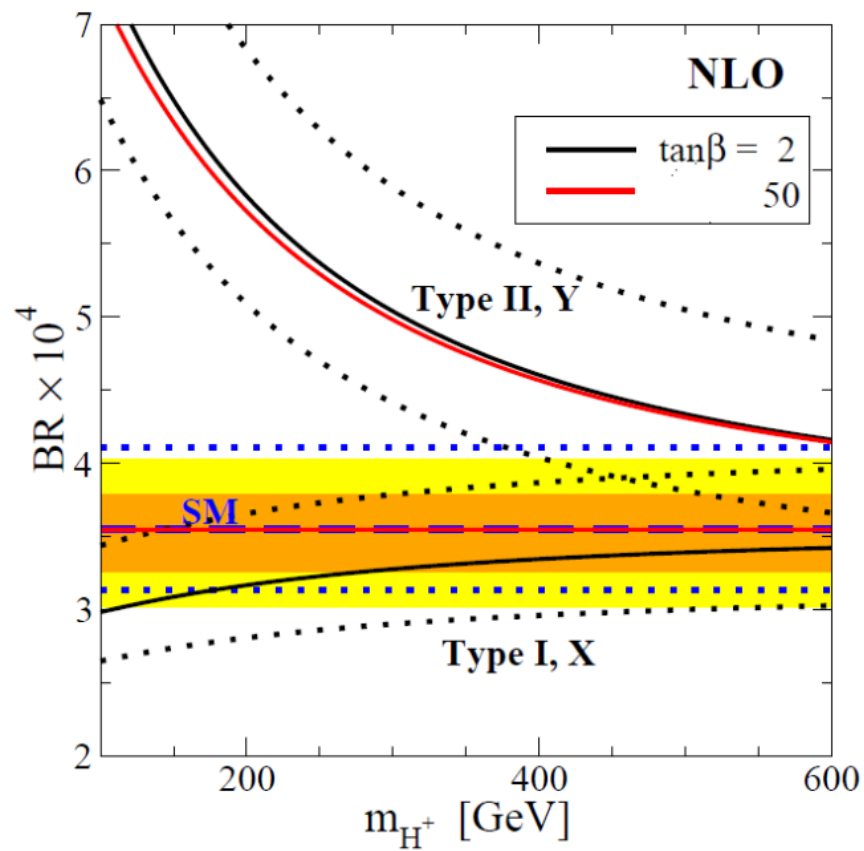
Summary

- We proposed a new Dirac neutrino mass model which can explain dark matter and baryon asymmetry of the Universe
- This new model was found by classifications for neutrino mass model in our previous studies.
- We confirmed that there was a benchmark scenario explaining these three BSM phenomena.

back up slides

Constraint of $bs \rightarrow \gamma$ (2HDM)

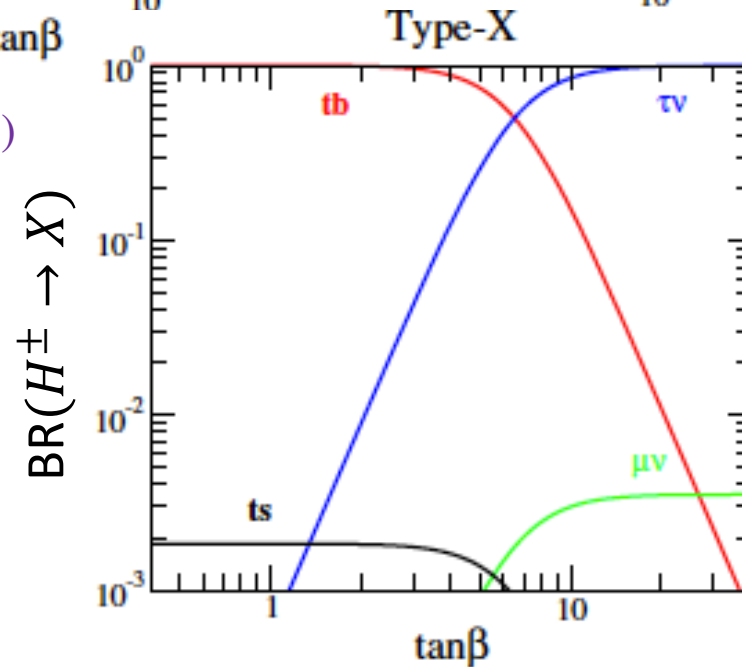
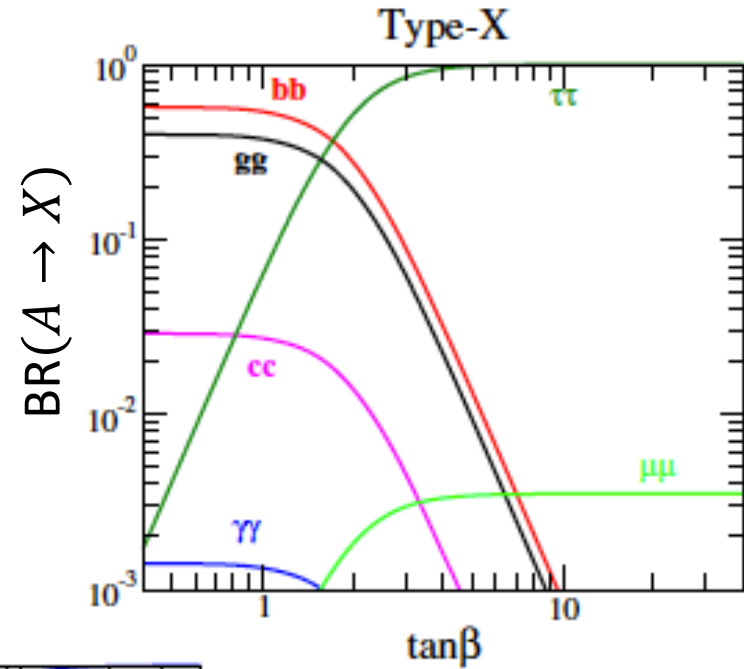
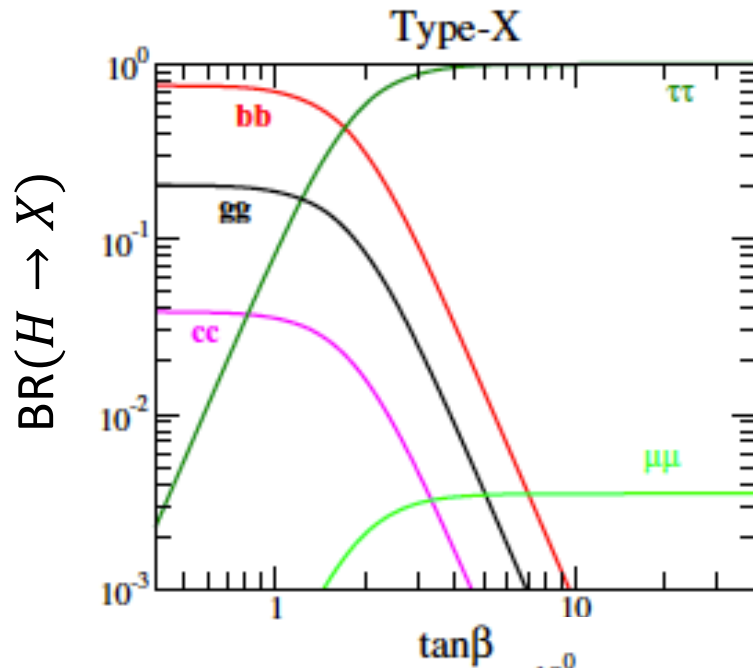
Aoki, Kanemura, Tsumura and Yagyu, Phys.Rev. D80 (2009) 015017



- : allowed regions (1σ)
- : allowed regions (2σ)

This is the reason why we adapt Type-X

Branching ratio H, A, H^\pm (Type-x THDM)



Kanemura, Tsumura, Yagyu,
Yokoya PRD 90, 075001 (2014)

$$\sin(\beta - \alpha) = 1$$

$$m_\Phi = M = 200 \text{ GeV}$$

Prospects of Direct search of A by the LHC(Type-x THDM)

Kanemura, Tsumura, Yagyū,
Yokoya PRD 90, 075001 (2014)

$$gg \rightarrow \phi^0 \rightarrow \tau^+\tau^-, gg \rightarrow b^-b\phi^0 \rightarrow b^-b\tau^+\tau^-$$

($\phi^0 = H, A$)

