

# NEARBY PULSARS AND THE COSMIC RAY POSITRON EXCESS

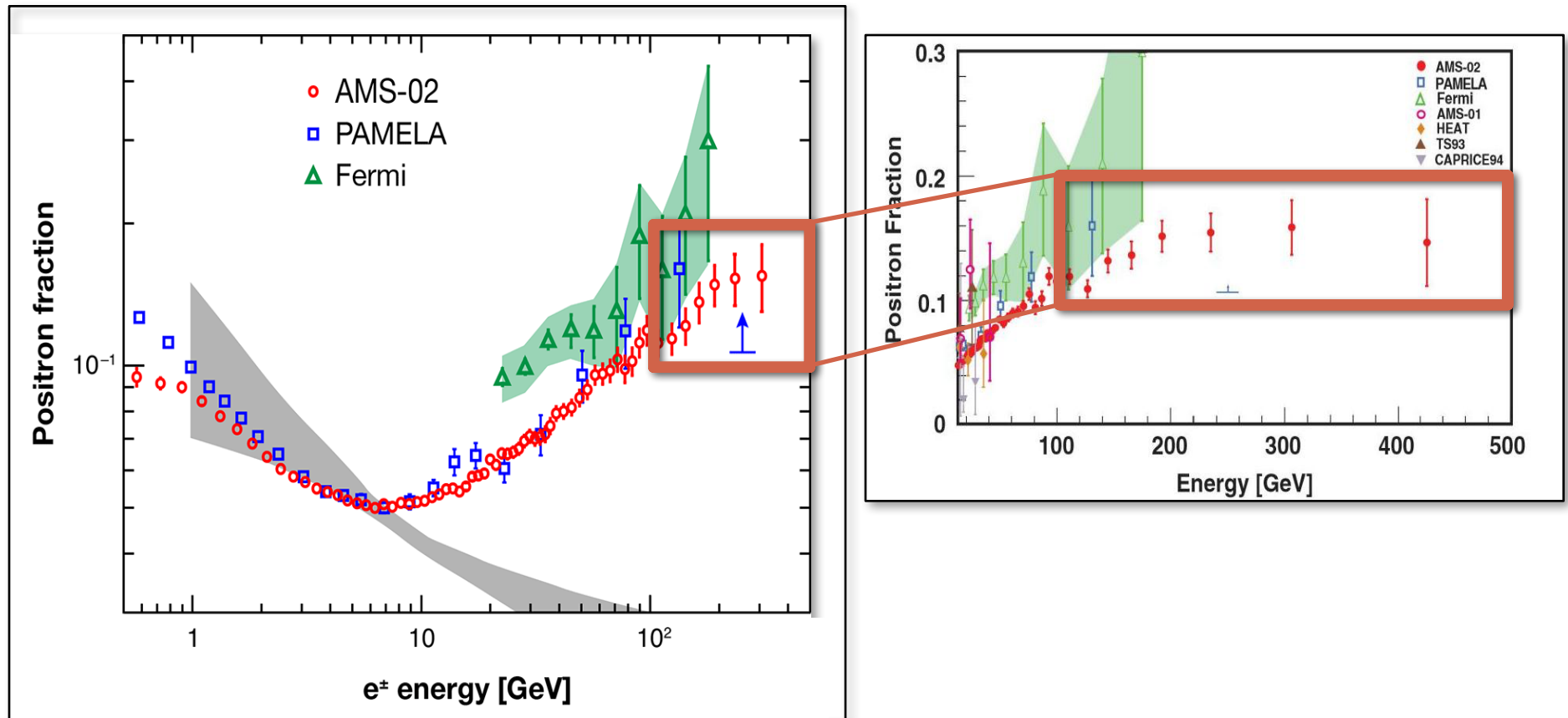
---

*Dan Hooper* – Fermilab and the University of Chicago  
TeV Particle Astrophysics Workshop, Columbus  
August 9, 2017

(Based on work with Ilias Cholis, Ke Feng and Tim Linden)

# The Cosmic Ray Positron Excess

- In 2008, PAMELA reported a surprisingly large quantity of positrons in the cosmic ray spectrum, now confirmed with much greater precision by AMS

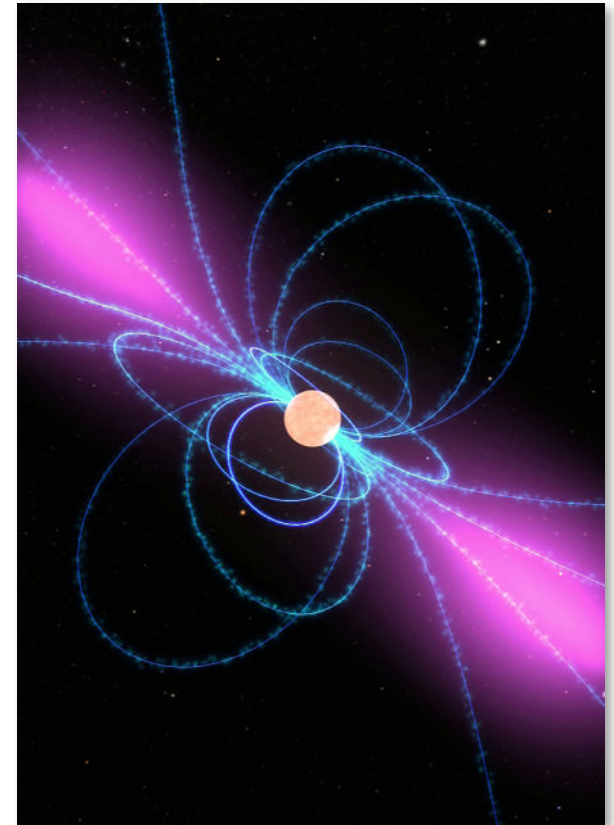


# Where Do The Excess Positrons Come From?

- The main contribution to the positron flux was anticipated to be that generated by cosmic ray interactions with gas in the ISM, yielding positrons through charged pion decay (*ie.* “*secondary*” *positrons*); this cannot account for the observed positrons
- Instead, three basic ideas have been proposed to account for the excess positrons:
  - 1) Annihilating or decaying dark matter particles
  - 2) The acceleration of secondary positrons within cosmic-ray sources (*ie.* supernova remnants)
  - 3) Nearby *primary* sources of high-energy positrons (*ie.* pulsars)

# Cosmic Ray Positrons From Pulsars

- Shortly after the PAMELA excess was reported, it was suggested that the positrons might originate from pulsars
- Pulsars are rapidly spinning neutron stars, which gradually convert their rotational kinetic energy into radio, X-ray, and gamma-ray emission, and into  $e^+e^-$  pairs
- Newly formed pulsars typically exhibit periods on the order of  $\sim 0.01$ - $0.1$  second, although most observed pulsars have higher periods (between  $\sim 0.1$  and a few seconds)
- The rate of a pulsar's spin-down evolution (and power) depends on the strength of its magnetic field (which transfers rotational kinetic energy into radiation via magnetic dipole braking)



DH, Blasi, Serpico, PRD, arXiv:0810.1527;  
Yuksel, Kistler, PRL, arXiv:0810.2784

(see also Zhang, Cheng, A&A, 2001; Grimani, A&A, 2007)

# Pulsars Emission Models

- Considerable research activity has been directed toward understanding exactly how pulsars generate their observed emission
- There are a number of basic elements that are found across a wide range of proposed models:
  - Electrons are accelerated by the strong magnetic fields, somewhere in the magnetosphere (the location is model dependent)
  - These electrons then induce electromagnetic cascades through the emission of curvature radiation
  - This results in the production of photons with energies above the threshold for pair production in the strong magnetic field
  - These electrons and positrons then escape the magnetosphere through open field lines, or after reaching the pulsar wind
- There is no consensus on what fraction of a pulsar's power is likely to go into the production of energetic  $e^+e^-$  pairs
- As high as ~20-30% of the energy budget? Or perhaps ~0.01%?

# Which Pulsars Contribute to the Positron Flux?

Consider the standard cosmic-ray transport equation:

$$\frac{\partial}{\partial t} \frac{dn_e}{dE_e}(E_e, r, t) = \vec{\nabla} \cdot \left[ D(E_e) \vec{\nabla} \frac{dn_e}{dE_e}(E_e, r, t) \right] + \frac{\partial}{\partial E_e} \left[ \frac{dE_e}{dt}(r) \frac{dn_e}{dE_e}(E_e, r, t) \right] + \delta(r) Q(E_e, t)$$

# Which Pulsars Contribute to the Positron Flux?

Consider the standard cosmic-ray transport equation:

$$\frac{\partial}{\partial t} \frac{dn_e}{dE_e}(E_e, r, t) = \vec{\nabla} \cdot \left[ D(E_e) \vec{\nabla} \frac{dn_e}{dE_e}(E_e, r, t) \right] + \frac{\partial}{\partial E_e} \left[ \frac{dE_e}{dt}(r) \frac{dn_e}{dE_e}(E_e, r, t) \right] + \delta(r) Q(E_e, t)$$

Diffusion:  $D(E_e) = D_0 E_e^\delta$

Energy Losses: (ICS, Synchrotron)

$$-\frac{dE_e}{dt}(r) = \sum_i \frac{4}{3} \sigma_T \rho_i(r) S_i(E_e) \left( \frac{E_e}{m_e} \right)^2 + \frac{4}{3} \sigma_T \rho_{\text{mag}}(r) \left( \frac{E_e}{m_e} \right)^2$$

$$\equiv b(E_e, r) \left( \frac{E_e}{\text{GeV}} \right)^2$$

Injection Spectrum: (burst-like approximation)

$$Q(E_e, t) = \delta(t) Q_0 E_e^{-\alpha} \exp(-E_e/E_c)$$

# Which Pulsars Contribute to the Positron Flux?

The solution to this equation is as follows:

$$\frac{dn_e}{dE_e}(E_e, r, t) = \frac{Q_0 E_0^{2-\alpha}}{8\pi^{3/2} E_e^2 L_{\text{dif}}^3(E_e, t)} \exp\left[\frac{-E_0}{E_c}\right] \exp\left[\frac{-r^2}{4L_{\text{dif}}^2(E_e, t)}\right]$$

where

$$L_{\text{dif}}(E_e, t) \equiv \left[ \frac{D_0}{b(E_e/\text{GeV})^{1-\delta}(1-\delta)} \left( 1 - (1 - E_e b t)^{1-\delta} \right) \right]^{1/2}$$

The sources that make the maximum contribution to the local positron flux are those located at a distance of  $r \sim 2.4 L_{\text{dif}}$ , which for ISM-like diffusion parameters yields:

$$r \sim 2.4 L_{\text{dif}} \sim 100 \text{ pc} \left( \frac{t}{10^5 \text{ yr}} \right) \left( \frac{E_e}{100 \text{ GeV}} \right)^{0.7}$$

**Conclusion: The pulsars which contribute most to the local positron flux are  $\sim 10^5$  years old and are located at a distance of  $\sim 100$  pc**



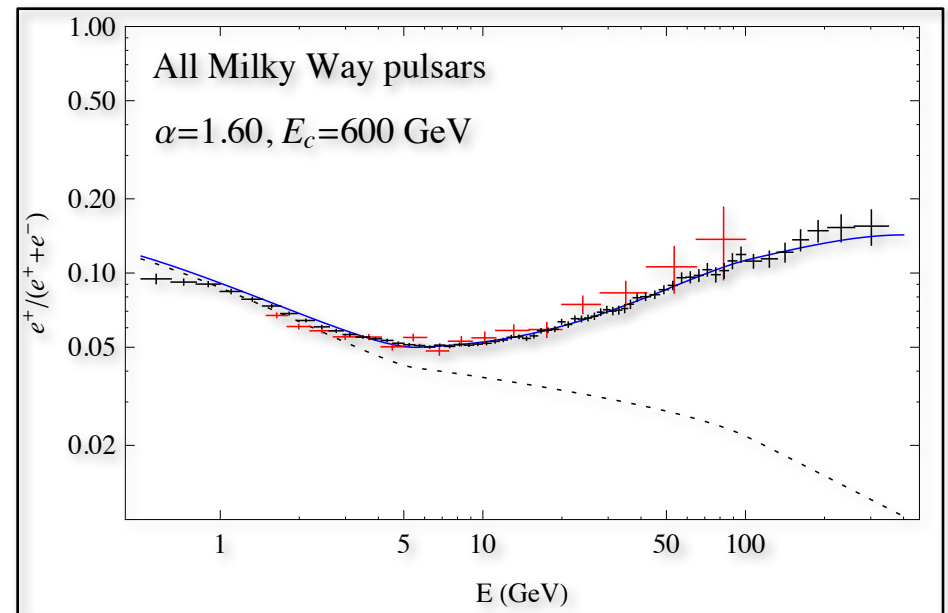
# Cosmic Ray Positrons From Pulsars

- From these considerations, there are two known pulsars which stand out as the strongest potential sources of  $\sim 100$  GeV cosmic-ray positrons:

**Geminga**, age  $\sim 370,000$  yrs, distance  $\sim 250$  pc

**B0656+14** (*ie.* monogem), age  $\sim 110,000$  yrs, distance  $\sim 280$  pc

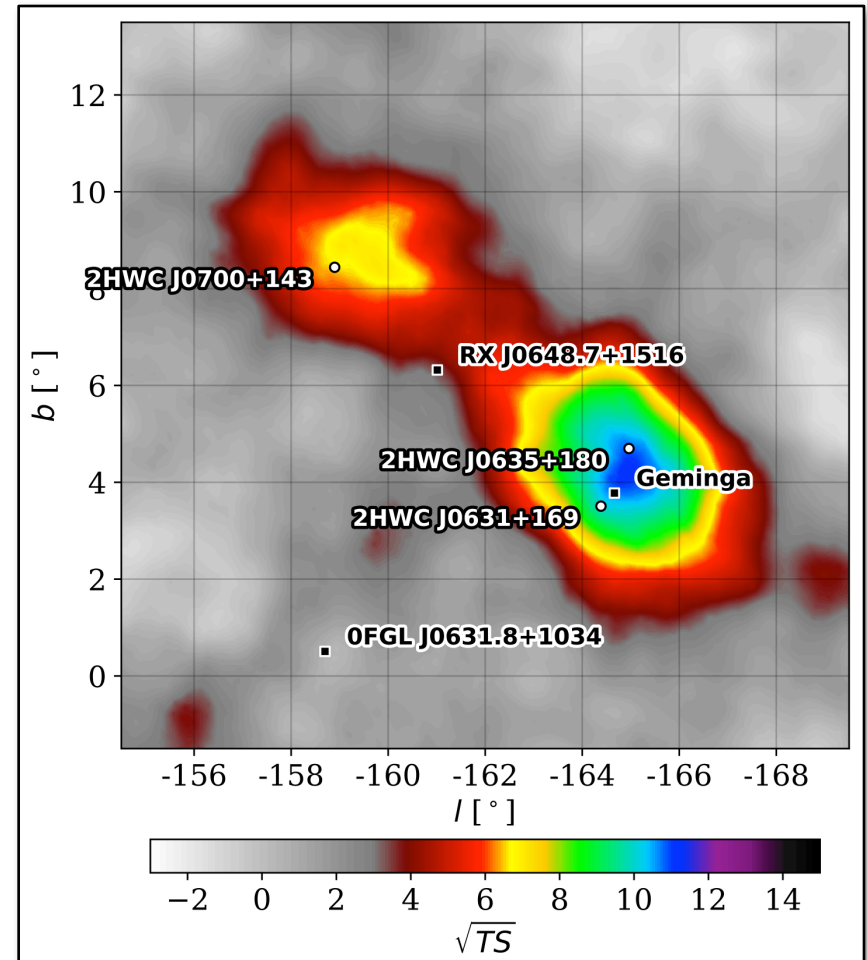
- If  $\sim 10$ - $20\%$  of the spin-down power of these pulsars is transferred into pairs, they could plausibly dominate the observed positron spectrum



DH, Blasi, Serpico, PRD, arXiv:0810.1527;  
 Yuksel, Kistler, PRL, arXiv:0810.2784;  
 Cholis, DH, PRD, arXiv:1304.1840

# VHE Gamma-Ray Observations of Geminga

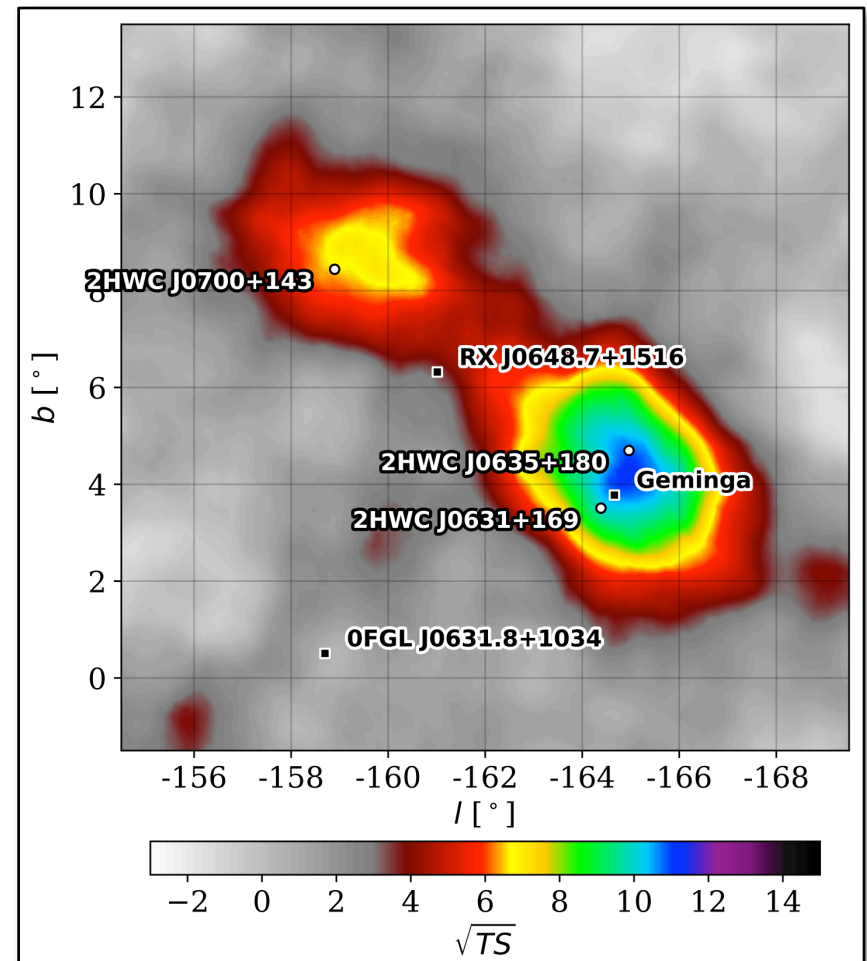
- Milagro detected VHE emission from Geminga and reported the “definitive detection of extended emission” from this source, with a full-width-half-max of  $2.6^{+0.7}_{-0.9}$  degrees
- Very recently, the HAWC Collaboration confirmed Milagro’s detection of Geminga, and its spatial extension, finding a radius of  $\sim 2^\circ$
- Furthermore, HAWC also reports  $\sim 2^\circ$  extended emission from the pulsar B0656+14 (2HWC J0700+143), not detected by Milagro (or by Fermi)



HAWC Collaboration, arXiv:1702.02992

# What Produces These Gamma Rays?

- The spatial extension of this emission indicates that the observed gamma rays **do not originate from the pulsar itself**, but from a surrounding region several parsecs in extent
- The only diffuse emission mechanisms that can produce such high-energy photons are inverse Compton scattering and pion production
- A pion production origin would require an implausibly large quantity of  $\sim 10^2$  TeV protons ( $>10^{46}$  erg), which would have to somehow be confined to the region for  $>10^5$  years
- In light of these considerations, inverse Compton scattering is almost certainly responsible for this emission



HAWC Collaboration, arXiv:1702.02992

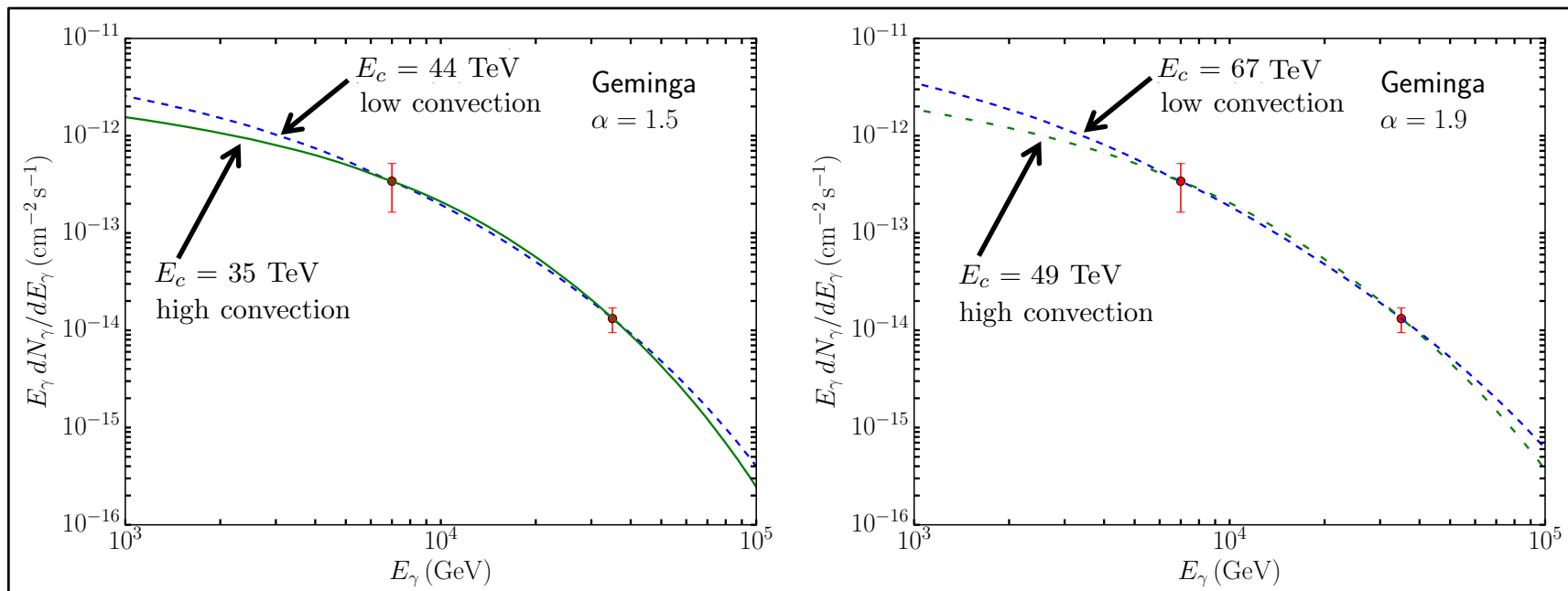
# HAWC Measurements Are Essential To Solving The Mystery Of The Positron Excess

- When a very high energy electron is injected into this environment, it emits the majority of its energy as Inverse Compton emission (along with a similar, but likely smaller quantity as synchrotron)
- The results of HAWC (and Milagro) thus provide us with a direct measurement of the energy that Geminga and B0656+14 are currently injecting into very high-energy  $e^+e^-$  pairs (as well as information pertaining to the spectral shape of these pairs)

Main Idea: ***The spatial extension of Geminga and B0656+14 allow us to measure the critically important (and until now highly uncertain) fraction of these pulsars' spindown power that goes into the production of energetic  $e^+e^-$  pairs***

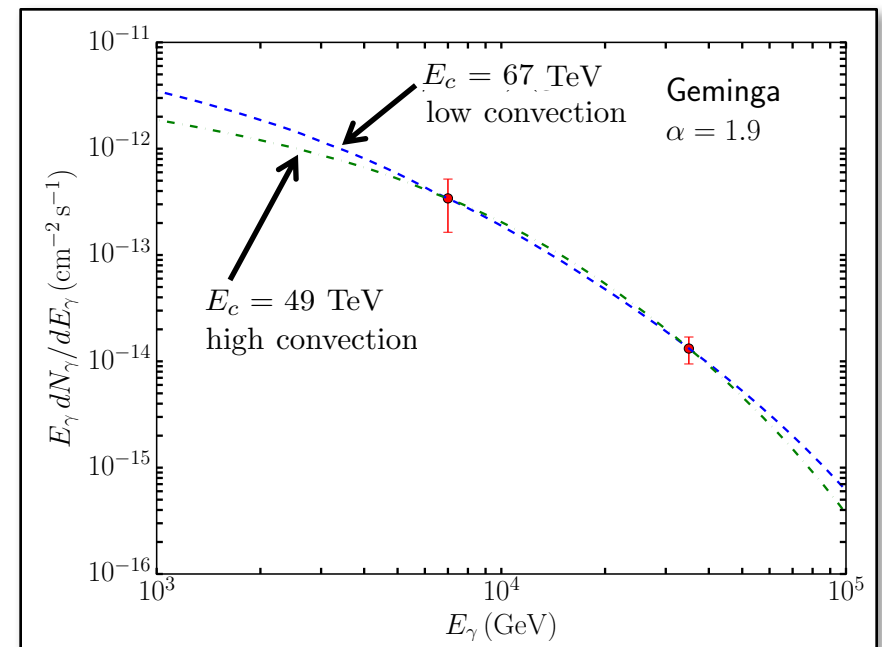
# Implications of HAWC and Milagro for the Positron Excess

- For a given spectrum of injected pairs, we calculate the resulting ICS spectrum (including all Klein-Nishina corrections), and use this to constrain the normalization, spectral index ( $\alpha$ ), and energy cutoff ( $E_c$ ) of the injected spectrum of  $e^+e^-$  pairs
- The VHE gamma-ray fluxes are best fit by  $\alpha \sim 1.5-2.0$  and  $E_c \sim 35-70$  TeV
- In these best-fit models, between 7-29% of Geminga's current spindown power goes into  $e^+e^-$  pairs – **similar to that required to generate the positron excess!**



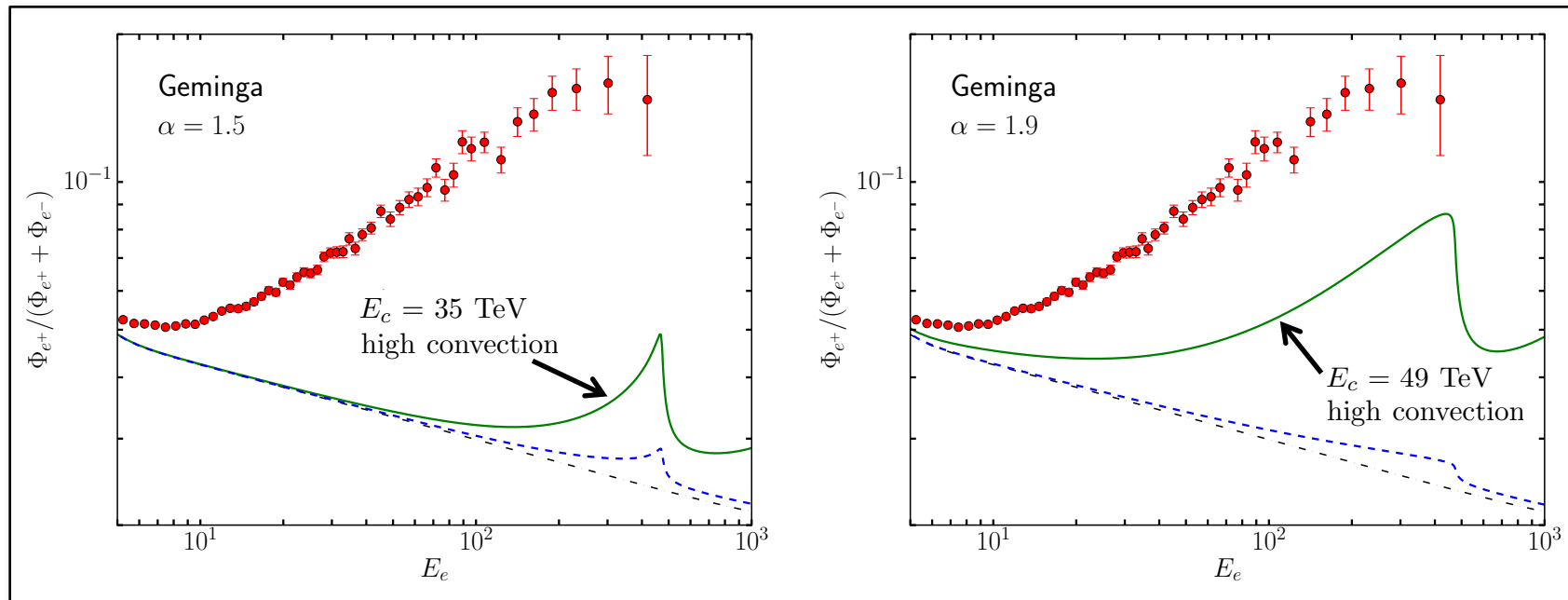
# The Role Of Convection

- The angular extension of the emission observed by HAWC indicates that diffusion is inefficient for VHE electrons/positrons in the parsecs surrounding these pulsars
- Lower energy leptons cool more slowly, and are likely to escape this region, perhaps via convection or through energy-independent diffusion
- We remain agnostic about how exactly this occurs, and simply parameterize the combination of these effects by a convection velocity
- This quantity impacts the shape of the gamma-ray spectrum, and when we take into account the spectral slope reported by HAWC ( $-2.23 \pm 0.08$ ), we find that a sizable convection velocity is required,  $v_c \sim 100\text{-}500 \text{ km/s}$
- In these plots, “*high convection*” refers to  $v_c \sim 230 \text{ km/s} \times (r_{\text{region}}/5 \text{ pc})$  – *focus your attention on these curves*



# Implications of HAWC and Milagro for the Positron Excess

- We can now use this information to calculate the contribution from Geminga to the local positron flux
- Across the range of models that provide a good fit to the HAWC and Milagro data, Geminga contributes non-negligibly to the observed excess



# Main Uncertainties

## ICS vs Synchrotron

- The fraction of energy in  $e^+e^-$  pairs goes into synchrotron rather than ICS is an uncertainty in our calculations (we adopted what we think are reasonable parameters:  $B=3 \mu\text{G}$ ,  $\rho_{\text{star}}=0.60 \text{ eV/cm}^3$ ,  $\rho_{\text{IR}}=0.60 \text{ eV/cm}^3$ , and  $\rho_{\text{UV}}=0.10 \text{ eV/cm}^3$ )
- Over a reasonable range of these parameters, we could plausibly change the net result by up to a factor of roughly  $\sim 2$  (either way)

## The Time Profile of Geminga's Emission

- HAWC and Milagro measure the energy in ICS today, and thus are sensitive to the pairs that were injected over the past  $\sim 10^4$  years
- In contrast, the positrons reaching the Solar System today were injected much longer ago, when the pulsar was young ( $\sim 10^5$  years ago)
- In our calculation, we adopt the standard magnetic dipole braking model with a spindown timescale of  $10^4$  years
- By varying our choice of this parameter, we could plausibly change the net result by an order one factor

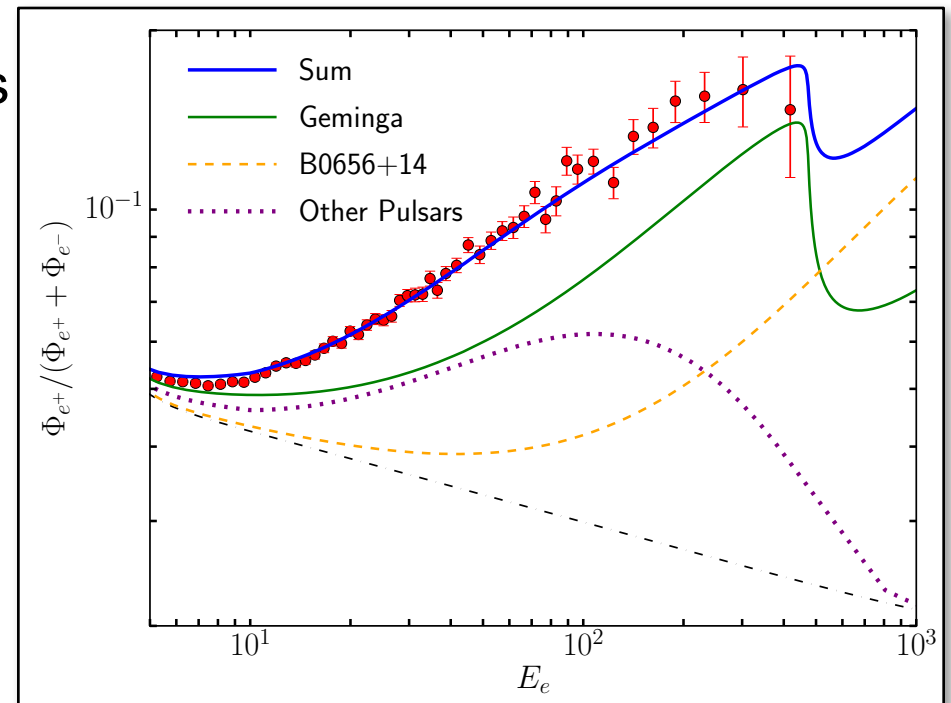


# Positrons From Geminga, B0656+14, and Other, More Distant Pulsars

- We have the most information about Geminga, and there is still an order one uncertainty as to its contribution to the local positron flux
- Larger uncertainties apply to B0656+14 and other pulsars
- That being said, can make a reasonable estimate for the total contribution

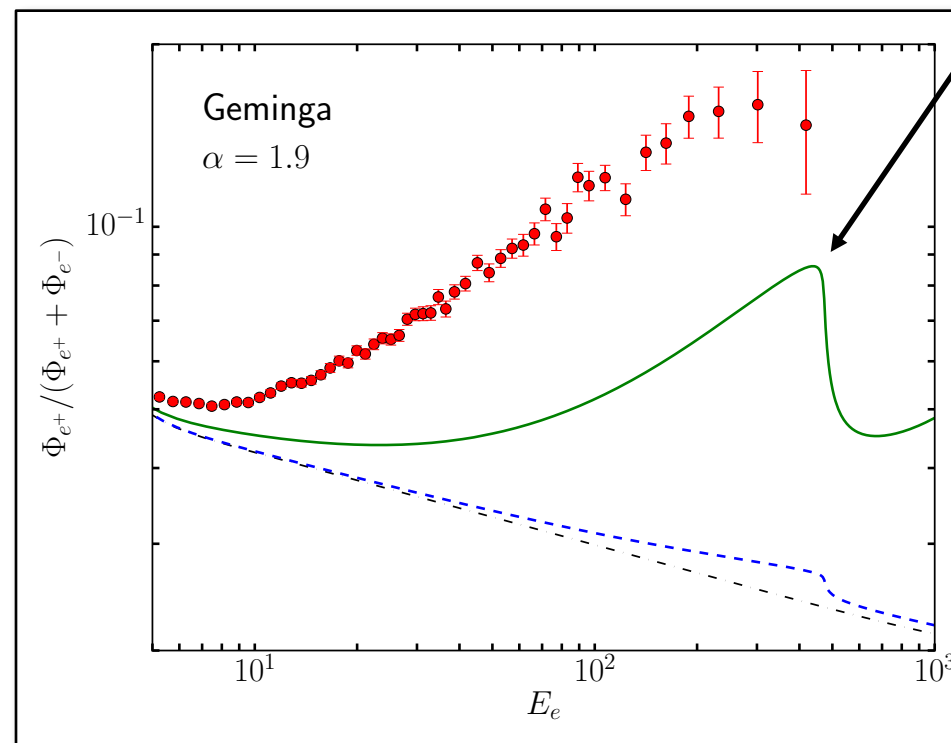
# Positrons From Geminga, B0656+14, and Other, More Distant Pulsars

- We have the most information about Geminga, and there is still an order one uncertainty as to its contribution to the local positron flux
- Larger uncertainties apply to B0656+14 and other pulsars
- That being said, can make a reasonable estimate for the total contribution
- In this figure, we have assumed that all pulsars inject  $e^+e^-$  pairs with the same efficiency and spectrum as Geminga, and adopted  $\tau \sim 4.3 \times 10^3$  years and a birth rate of 2 new pulsars per century throughout the Milky Way (adopting the Lorimer *et al.* spatial distribution)
- These assumptions might not be precisely correct, but this shows that **pulsars very plausibly generate the entire excess, and very likely provide the dominant contribution**



# A Note On Positron Spectral Features

- A great deal is often made about “edges” and other spectral features that might appear in the positron spectrum; such features are sometimes called a “smoking gun” for dark matter annihilation or decay
- In fact, a nearby pulsar could very plausibly generate an edge-like feature, at an energy of  $E \sim I/bt_{age}$  (which for Geminga is at  $\sim 350\text{-}700$  GeV)



# Summary

- Recent observations of Geminga and B0656+14 by HAWC provide a determination of the flux of very high-energy  $e^+e^-$  pairs that is currently being injected by these pulsars; equivalent to  $\sim 7-29\%$  of the total current spindown power
- This is a critical quantity, and was previously almost entirely unknown
- This new information indicates that pulsars generate an order one fraction of the positron excess, and are very likely to be responsible for the majority (or entirety) of this signal
- There is still room in the uncertainties for other contributions to the positron flux (*i.e.* dark matter or secondary acceleration), but it is now reasonably clear that pulsars are the main source of the observed positron excess

Personally, I think this is a very exciting result ... regardless of what Science Magazine has to say about it;)



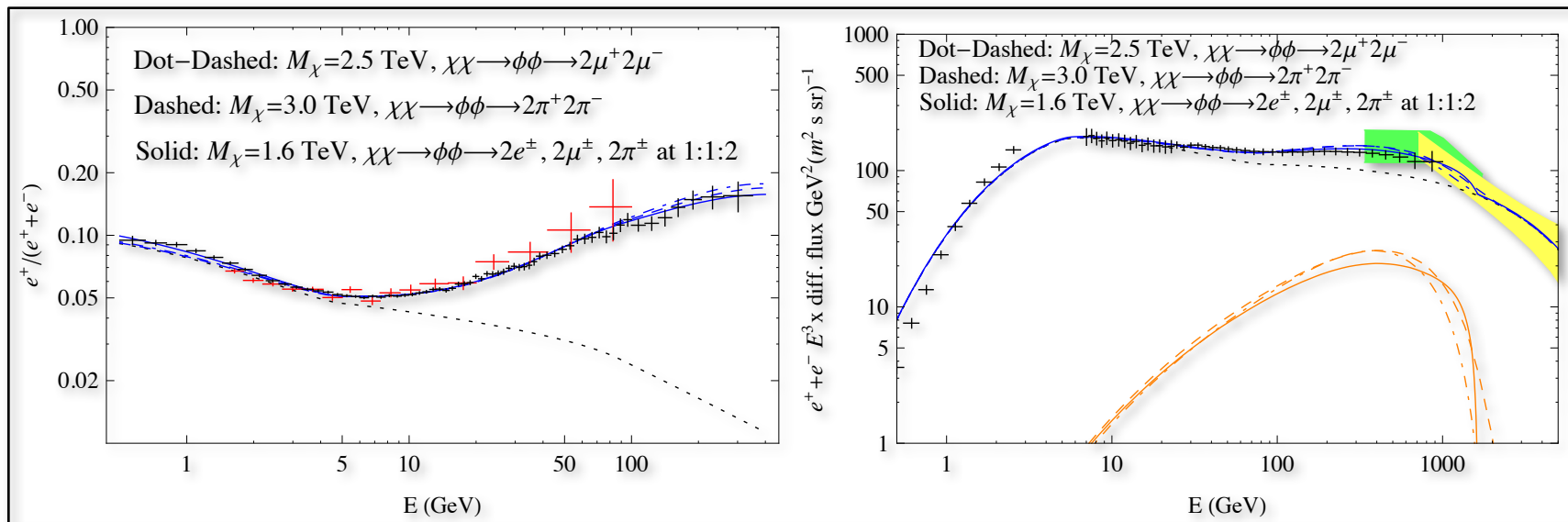
## Case weakens for antimatter sign of dark matter

By **Edwin Cartlidge** | Mar. 6, 2017 , 4:00 PM

A long debate over a mysterious surplus of antimatter—and whether it’s a sign of dark matter—may be coming to **an anticlimactic end**. For more than a decade, multiple experiments have found an unexpected excess in the number of high-energy antielectrons, or positrons, in space, and some physicists suggested it could be due to particles of dark matter annihilating one another. Others countered with **a more mundane explanation**: The positrons come from rapidly rotating neutron stars, or pulsars. Now, a team of theorists has bolstered that more prosaic explanation, showing in detail that pulsars can indeed produce most or all of the excess.

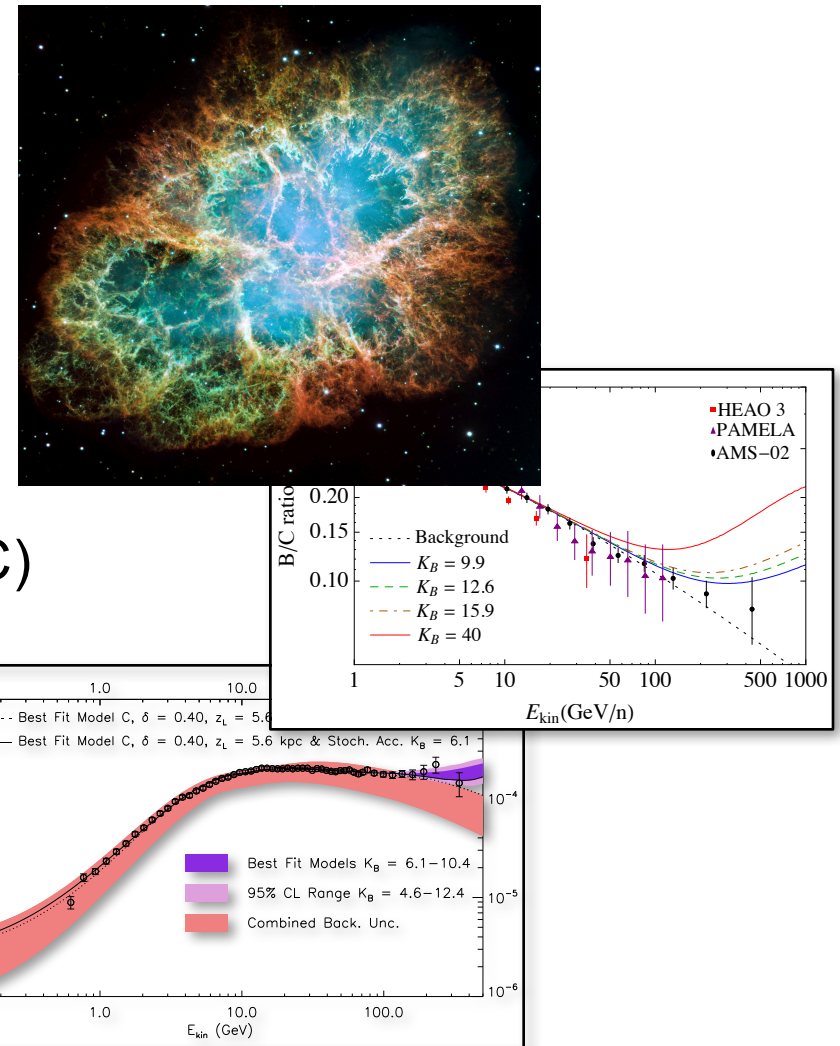
# Annihilating Dark Matter and the Positron Excess

- In light of the detailed measurements of the positron fraction from AMS (and of the electron+positron spectrum from Fermi and HESS), few dark matter models can accommodate the data
- Dark matter models that *can* accommodate the data generally consist of a  $\sim 1\text{-}3$  TeV particle that annihilates to unstable intermediate states, which then decay to electrons, muons and/or charged pions
- Large annihilation cross sections are also required ( $\sim 10^{-24}$  to  $3 \times 10^{-23}$   $\text{cm}^3/\text{s}$ ), making constraints from Fermi difficult to evade



# The Acceleration of Secondary Positrons in Supernova Remnants

- Supernova remnants could generate secondary positrons and then accelerate them before they escape into the ISM
- If secondary positrons are accelerated in supernova remnants, then secondary antiprotons and boron nuclei should be accelerated as well
- Measurements of the boron-to-carbon (B/C) and antiproton-to-proton ratios from AMS indicate that secondary acceleration cannot account for the entirety of the positron excess, but may contribute non-negligibly



P. Blasi, PRL, arXiv:0903.2794; Mertsch, Sarkar, PRL, arXiv:0905.3152;  
 Cholis, DH, PRD, arXiv:1312.2952; Cholis, DH, Linden, arXiv:1701.04406