# Hunting for WIMPs: how low should we go?

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## WIMP with a capital W

- Cosmology and direct detection are really controlled by interactions with gauge/ SM Higgs boson
- e.g. Singlet-doublet model, split SUSY...

https://arxiv.org/pdf/1109.2604.pdf













### Spin Independent Scattering



Dirac DM coupling to Z ~ 10  $\sigma \approx \text{few} \times 10^{-40}$ 



## Z-mediated Dark Matter

$$\mathcal{L} \supset \frac{c}{2\Lambda^2 c} (iH^{\dagger}D_{\mu}H + \text{h.c.})\bar{\chi}\gamma^{\mu}\gamma^5\chi$$
$$\mathcal{L} \supset \frac{2\Lambda^2 c}{2\Lambda^2} (iH^{\dagger}D_{\mu}H + \text{h.c.})\bar{\chi}\gamma^{\mu}\gamma^5\chi$$

$$\mathcal{L} \supset -\frac{g_2}{4c_W} \frac{cv^2}{\Lambda^2} Z_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi$$
$$-\frac{4c_W}{4c_W} \Lambda^2 \gamma^{-\mu} \chi^{-\mu} \chi^{$$

### Also $\chi$ - $\chi$ -Z-h and $\chi$ - $\chi$ -Z-h-h contact interactions(!)

Also:

de Simone et al, arXiv:1402.6287;

Arcadi, Mambrini and Richard, arXiv:1411.2985,

Berlin, Escudero, Hooper and Lin, arXiv 1609.09079;







J. Kearney and N. Orlofsky, AP 1611.05048





### Contribution to T parameter

 $\mathcal{L} \supset \frac{c}{2\Lambda^2} (iH^{\dagger}D_{\mu}H + \text{h.c.})\bar{\chi}\gamma^{\mu}\gamma^5\chi$ 











#### Precision Electroweak Constraints





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## Motivates inclusion of new EW states



#### Scissors credit: J. Kearney



## Singlet-Doublet Dark Matter

$$\mathcal{L} \supset -yDHN - y^c D^c \tilde{H}N - M_D DD^c - \frac{M_N}{2}N^2 + \text{h.c.}$$

- Dirac doublet, D/Dc and Majorana N.
- Similar to Higgsino/Bino sector of the MSSM, but without all the pesky symmetry.
- Gives couplings to h and Z
- Ensures approximate unification (cf. split SUSY)

Arkani-Hamed, Dimopoulos, and Kachru hep-th/0501082;

Mahbubani, Senatore [hep-ph/0510064] D'Eramo [arXiv:0705.4493]

Enberg et al. [arXiv:0706.0918] Cohen, Kearney, AP, Tucker-Smith [arXiv: 1109.2604]



## Spin-Independent Coupling?

• There is a direct detection "blind spot" Cohen, Kearney, AP, Tucker-Smith [arXiv:1109.2604] Cheung, Hall, Pinner, Ruderman [arXiv:1211.4873]

$$y_{\rm BS}^c = -y \frac{M_N}{M_D} \left( 1 \pm \sqrt{1 - \left(\frac{M_N}{M_D}\right)^2} \right)^{-1}$$

Can be found by "low energy theorem"

$$\begin{aligned} \mathcal{L}_{h\chi\chi} &= \frac{1}{2} m_{\chi_i} (v+h) \chi_i \chi_i \\ &= \frac{1}{2} m_{\chi_i} (v) \chi_i \chi_i + \frac{1}{2} \frac{\partial m_{\chi_i} (v)}{\partial v} h \chi_i \chi_i + \mathcal{O}(h^2), \\ &\det \left( M_{\chi} - \mathbb{1} m_{\chi_i} (v) \right) = 0. \end{aligned}$$















 $E \longrightarrow W \rightarrow \nu_1$ 

 $\rightarrow$   $\,$  No tree-level direct detection analog  $\,$ 

No tree-level direct detection analog



## Question:

 Suppose Higgs coupling is small (near the blindspot), can we expect to see the Dark Matter through its spin-dependent scattering?



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### Singlet Doublet $m_D$

### In blind spot (fixes y<sup>c</sup>)

Relic density thermal (fixes  $M_D$ )



# Breaking the Crossing Symmetry: Co-annihilation

If  $\chi$  and Y simultaneously inhabit the thermal bath at freeze-

out

(Boltzman suppression not too large)

$$e^{-\Delta m/T_{FO}} \approx e^{-20\frac{\Delta m}{m}}$$

 $\sigma(\chi\chi \to SM) \ \sigma(\chi Y \to SM) \ \sigma(YY \to SM)$ 

Griest and Seckel, Phys.Rev. D43 (1991) 3191-3203







## Loop Induced Dark Matter Couplings







WIMP Mass  $[\text{GeV}/c^2]$ 

Snowmass 1310.8327



### **Mass Splitting**





#### Mass Splittings









## But...

 It is always possible that there could be "some Higgsino" in the dark matter, in which case, direct detection may have nothing to do with the cosmology.

$$\sigma_{SI} \approx 3 \times 10^{-47} \text{cm}^2 \left(\frac{1 \text{ TeV}}{\mu}\right)^4 \left(\frac{m_{\chi}}{500 \text{ GeV}}\right)^2 \left(1 + \frac{\mu s_{2\beta}}{m_{\chi}}\right)^2 \left(1 - \frac{m_{\chi}^2}{\mu^2}\right)^{-2}$$



## Conclusion WIMPs: a Status Report

- Higgs-centric cosmology getting squeezed
- Z-centric cosmology is **exciting now** 
  - Symmetry reason for blind spot?
- Co-annihilation-centric cosmology (stop or otherwise) will be very hard for the foreseeable future, but we could get lucky.
- Why co-annihilation? (AP, Kearney, Phys.Rev. D88 (2013) no.9, 095009)



## Extra Slides





Cohen, Kearney, AP, Tucker-Smith 1109.2604





Cohen, Kearney, AP, Tucker-Smith 1109.2604



### Effective Field Theory Validity





## Both gluino mass signs







- Long list of 2-loop (and more!) computations:
  - Carena, Degrassi, Ellis, Espinoza, Haber, Harlander, Heinemeyer, Hempfling, Hoang, Hollik, Hahn,
    Martin, Pilaftsis, Quiros, Ridolfi, Rzehak, Slavich, Wagner, Weiglein, Zhang, Zwirner.



Carena and Haber, hep-ph/0208209



### Charge Color-Breaking Vacuua



Morrissey and Blinov 1310.4174





## Double Blind spot

- There is a spot where both the SI and SD vanish.
- Requiring these two couplings to vanishes, along with reproducing the thermal abundance, sets MN =MD=850 GeV. (model building?)







### Another caveat

 If willing to tune to the Higgs blindspot, there is one more out (not requiring SD): CP Violation

$$\mathcal{C} \to \chi \gamma_5 \chi h \to (\bar{q}q)(\chi \gamma_5 \chi)$$



r= ratio of yukawas

