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# New Insights on Particle Acceleration at Non-relativistic Shocks

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# Non-Relativistic Collisionless Shocks

Mediated by collective electromagnetic interactions Show prominent non-thermal activity

### Heliospheric

### Extra-Galactic

### Propagate in environments likely rich in energetic particles (seeds)

#### Galactic





# Astroplasmas from first principles

Full-PIC approach
Define electromagnetic fields on a grid
Move particles via Lorentz force
Evolve fields via Maxwell equations
Computationally very challenging!

Hybrid approach: Fluid electrons - Kinetic protons (Winske & Omidi; Burgess et al., Lipatov; Giacalone et al.; DC & Spitkovsky,....)

massless electrons for more macroscopical time/length scales









# Hybrid simulations of collisionless shocks



dHybrid code (Gargaté et al, 2007; DC & Spitkovsky 2014)

Time =  $880.00 [1/\omega_p]$ 





# **CR-induced Magnetic Field Amplification**



 $x[c/\omega_p]$ 

### Initial B field M<sub>s</sub>=M<sub>A</sub>=30

DC & Spitkovsky, 2013





# Spectrum evolution

• Diffusive Shock Acceleration: non-thermal tail with universal spectrum  $f(p) \propto p^{-4}$ 

#### Acceleration efficiency: ~15% of the shock bulk energy!



DC & Spitkovsky, 2014a









### Ion Injection at Shocks

# • The fraction of injected particles depends on $v_{in}$ and $\vartheta$ (DC, Pop & Spitkovsky, 2015)



Ions injected into DSA undergo specular reflection at the reforming shock barrier and shock drift acceleration



# Injection via Shock-Drift Acceleration (SDA)



#### lons advected downstream, and thermalized



High barrier (overshoot)

#### $|e \Delta \Phi| > m V_x^2/2$



Reflection probability ~ barrier duty cycle (~25%)

To overrun the shock, ions need a minimum energy E<sub>ini</sub>, increasing with  $\vartheta$ , which they may achieve via multiple SDA cycles

• After N cycles, only a fraction  $\eta \sim 0.25^{N}$  has not been advected

Sor  $\vartheta ≤ 45^\circ$ ,  $E_{ini} ≤ 10E_{sh}$ , which requires N ≤ 3 → η~1%

 $\odot$  For  $\vartheta > 45^{\circ}$ ,  $E_{ini} > 10E_{sh}$ , hence N > 3 and  $\eta <<1\%$ 



# Hybrid Simulations: Summary

DSA efficient at q-parallel, strong shocks
 CRs amplify B via streaming instability
 Injection of thermal ions at q-parallel shocks via specular reflection and SDA

What if there are already energetic seeds?
How does injection depend on mass/charge?







# What if there are already energetic particles (seeds)?

The Martin The Carrie Star

### **Diffusive Shock Re-Acceleration**

### $= \frac{9}{60^{\circ}}$ shock with isotropic seeds $E_{CR} = 10E_{sh}$ ; $n_{CR} = 0.01$ (DC, Zhang, Spitkovsky, JPP submitted)

#### Seeds are effectively reflected at the shock, amplify the upstream B, and undergo DSA: DSRA!







Efficiency

80

 $\circ$  Seed DSRA independent of  $\vartheta$ , about 4x the initial energy density Also electrons are reaccelerated!

 $\oslash$  A ( $\vartheta$ <45°): As without seeds  $\oslash$  B (45°< $\vartheta$ <70°): Boosted to few %  $OC(\vartheta > 70^\circ)$ : No proton DSA







# Quasi-Perpendicular SEEDED Shocks

Image: steeper start with seeds E<sub>CR</sub>=10E<sub>sh</sub>
 Seeds diffuse but their spectrum is steeper than DSA
 <u>Non-thermal protons only downstream</u>









# Quasi-Perpendicular SEEDED Shocks

 $\oslash \ \vartheta = 80^{\circ}$  quasi-perp shock with seeds  $E_{CR} = 10E_{sh}$ Seeds diffuse but their spectrum is steeper than DSA Solution Non-thermal protons only downstream







# The Current in Reflected CRs

### $\circ$ It depends on the fraction of reflected seeds, n, and their speed, v<sub>r</sub>









# A Universal Current in Reflected CRs



 $\circ \eta$  and  $v_r$  balance their dependence on  $\vartheta$  and M exactly:  $J_{CR} = n_{CR}V_{sh}$ Easy explanation: CR anisotropy conserved at the shock crossing, in the shock frame For SNRs and Galactic CRs: T<sub>Bell</sub>~10yr Minimum level of B-amplification for shocks in the ISM





SN1006 1517.500 MHZ

#### Radio (GeV electrons)

TeV acceleration only where quasi-par, but seed DSRA can produce GeV electrons where oblique/quasi-perp

### SN 1006

X-rays. Red: thermal White: synchrotron (TeV electrons)







#### DC & Spitkovsky, 2014a



# How does DSA depend on the ion mass/charge ratio?

# Chemical Composition of Galactic CRs

"Urban legend": similar to solar (Simpson 1983)
Depends on volatility, on atomic mass A, on first ionization potential..
Above 1 TeV, fluxes of H, He, CNO, and Fe are comparable!



Nuclei heavier than H must be injected more efficiently





### Hybrid Simulations: Acceleration of Heavy lons





# The Onset of Ion Acceleration

#### Early times

 $\log_{10}[f_1(p_x)/\chi_1]~~(t=200\omega_c^{-1})$ 





#### Late times

 $\log_{10}[f_1(p_x)/\chi_1]$   $(t = 700\omega_c^{-1})$ 

2700

2750

2850

2800

2600

2650

with singly-ionized nuclei Ions injected by being isotropized just downstream: no shock reflection! Heavy ion injection after the onset of self-generated B turbulence DC, Yi, Spitkovsky 2017

-3

-3.5

2950



### No Injection at Quasi-perpendicular Shocks

 $\oslash$  M=20, oblique ( $\vartheta$ =60°) shock: no injection into DSA! Having a large gyroradius (large A/Z) is not sufficient for injection
 Seed ions can still enter DSA (e.g., solar energetic particles/solar flares Tylka+05)





Seed Diffusive Shock Re-Acceleration effective Streaming instability with universal current  $J_{CR} = n_{CR}V_{sh}$ Can trigger proton DSA for oblique shocks New phenomena at quasi-perpendicular shocks: steep seed spectra & proton acceleration downstream! CR reacceleration must happen in SNRs (e.g., W44,
 IC443, see Uchiyama+10, Cardillo+16) When proton DSA and B amplification are effective, heavy ions are preferentially injected Nuclei enhancement depends on A/Z and on the shock Mach number

Summary



magnetic turbulence







# Ion Injection at Shocks: a Minimal Model









# Earth's bow shock (AMPTE/IRM): Monte Carlo (Ellison, Möbius, Paschmann 1990)



