Leiden
Oct 25 2017

Particle acceleration in shocks: insights from kinetic simulations

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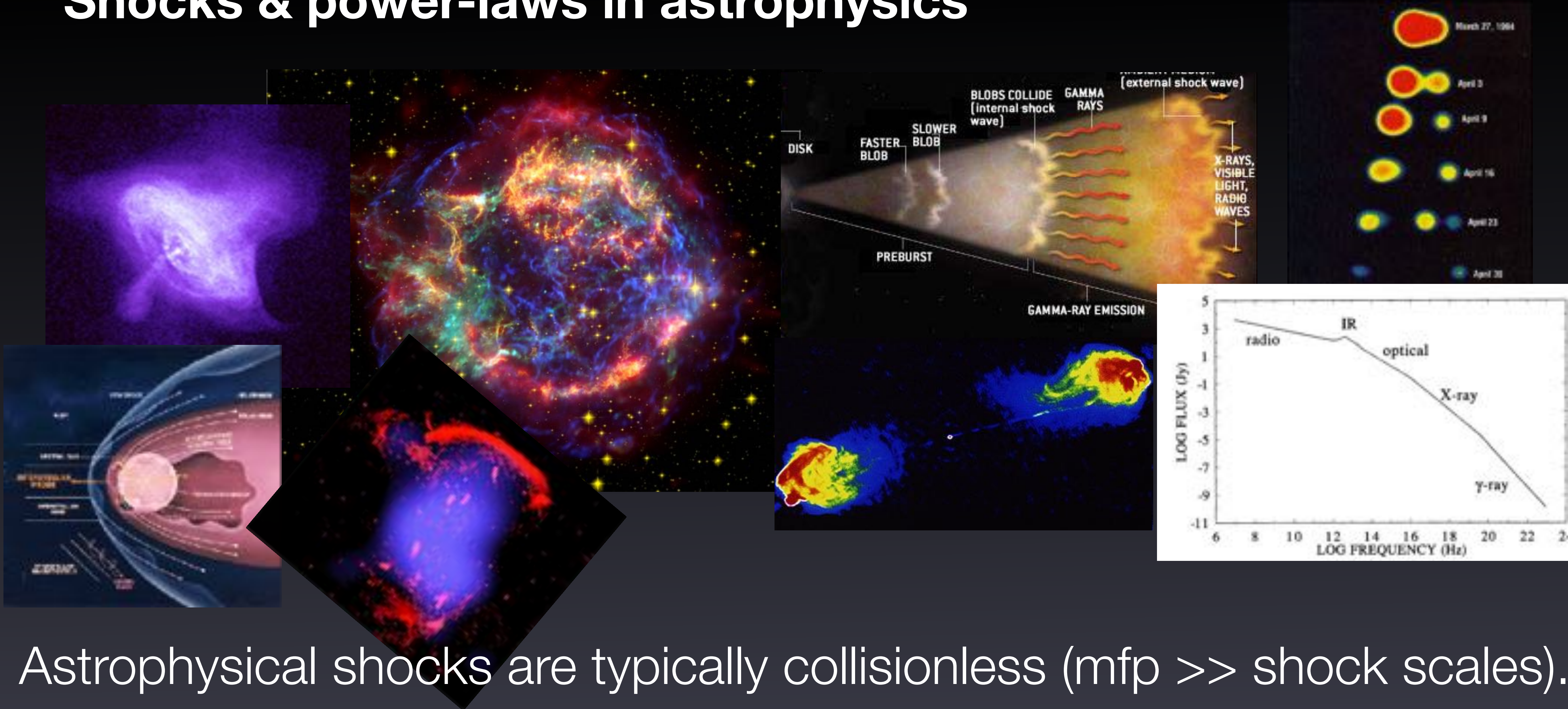
Shocks & power-laws in astrophysics

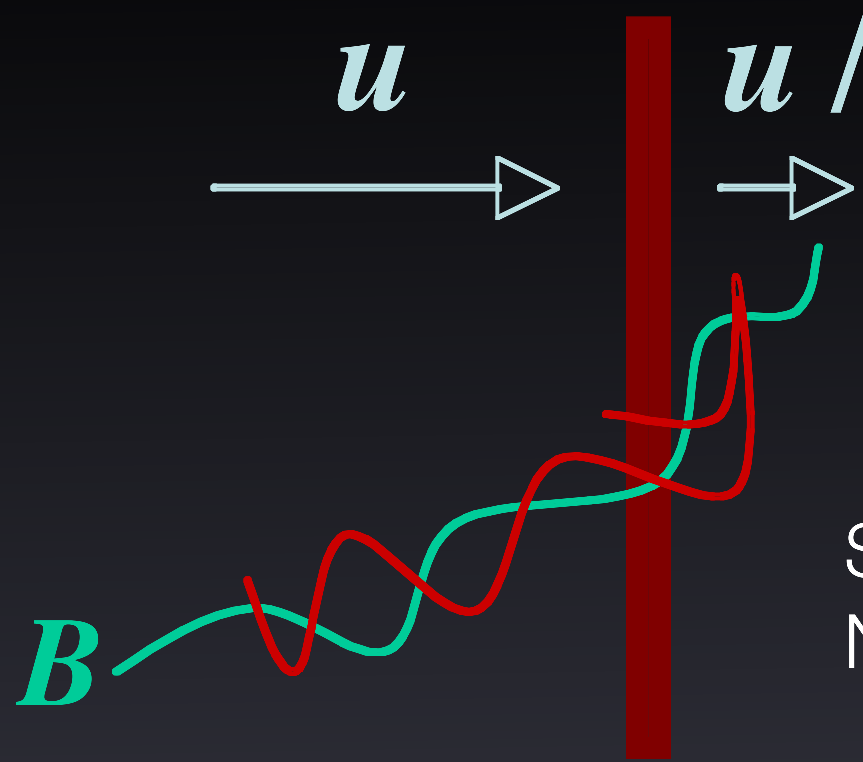


Astrophysical shocks are typically collisionless (mfp >> shock scales).
Many astrophysical shocks are inferred to:

1) accelerate particles to power-laws
2) amplify magnetic fields
3) exchange energy between electrons and ions

How do they do this? Mechanisms, efficiencies, conditions?…

Particle acceleration:



$\Delta E/E \sim v_{shock}/c$

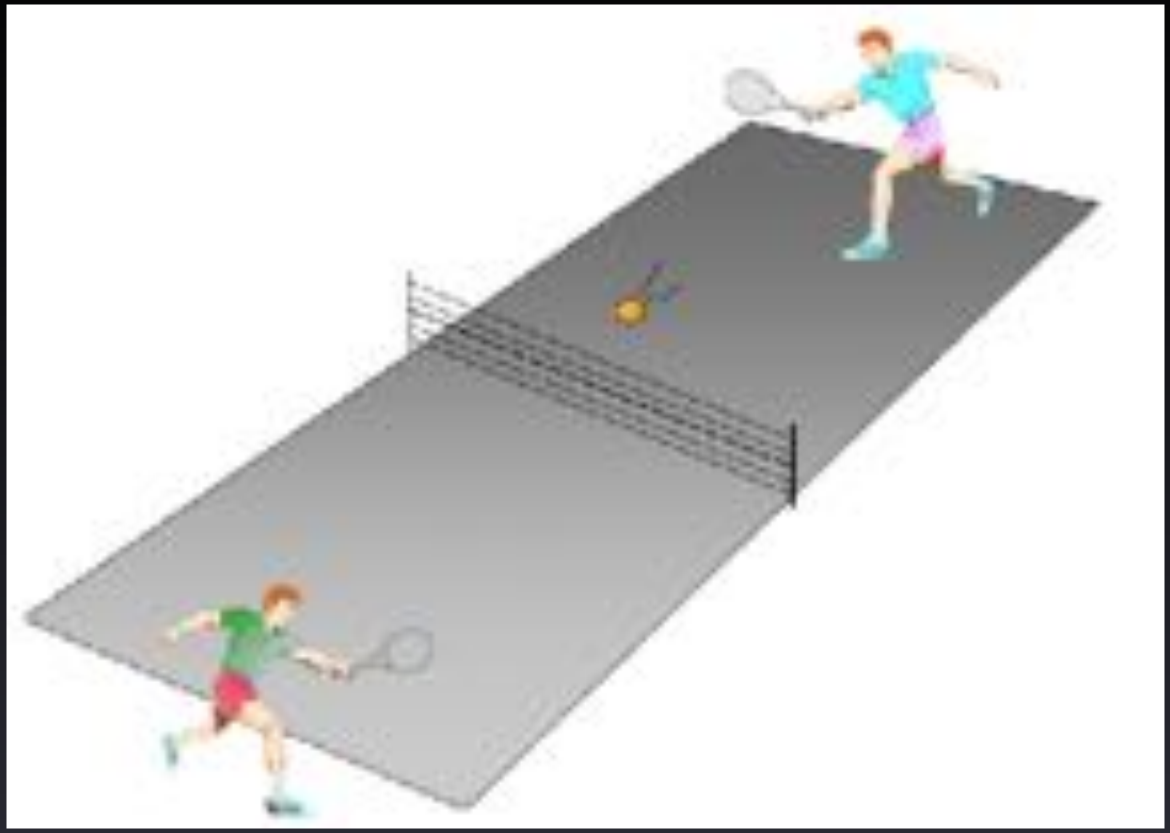
$N(E) \sim N_0\, E^{-K(r)}$

Strong shock:
$N(E) \sim N_0\, E^{-2}$

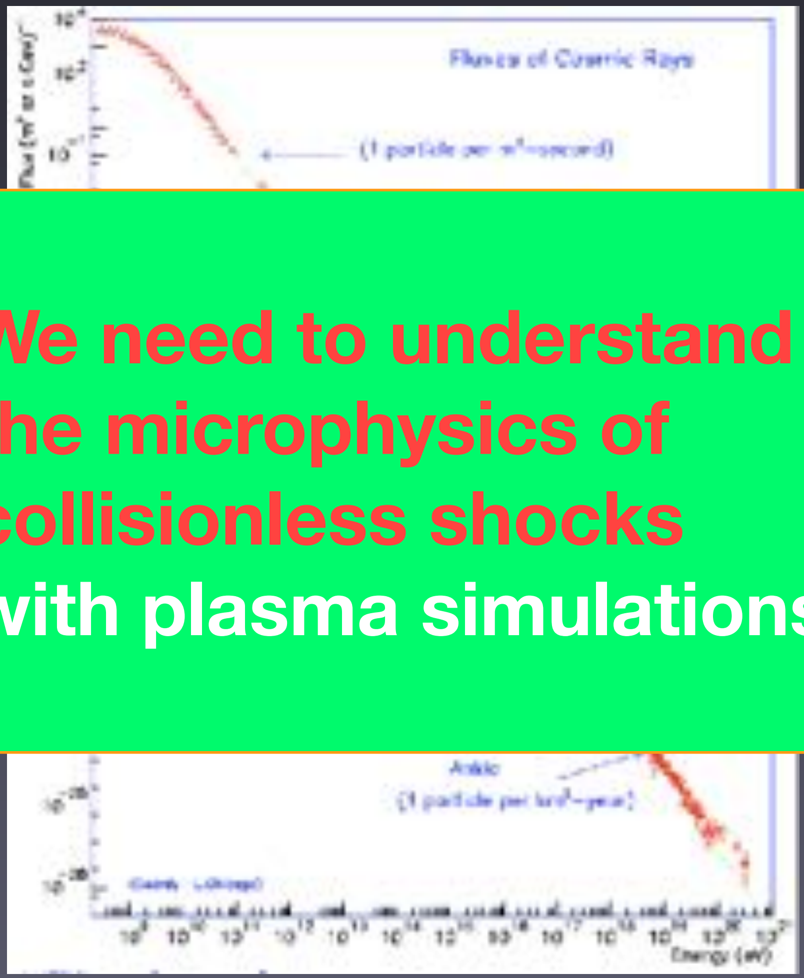
- Original idea -- Fermi (1949) -- scattering off moving clouds. Too slow (second order in v/c) to explain CR spectrum, because clouds both approach *and* recede.

- In shocks, acceleration is first order in v/c, because flows are always converging (Blandford & Ostriker 78,Bell 78, Krymsky 77)

- Efficient scattering of particles is required. Particles diffuse around the shock. Monte Carlo simulations show that this implies very high level of turbulence. Is this realistic? Are there specific conditions?



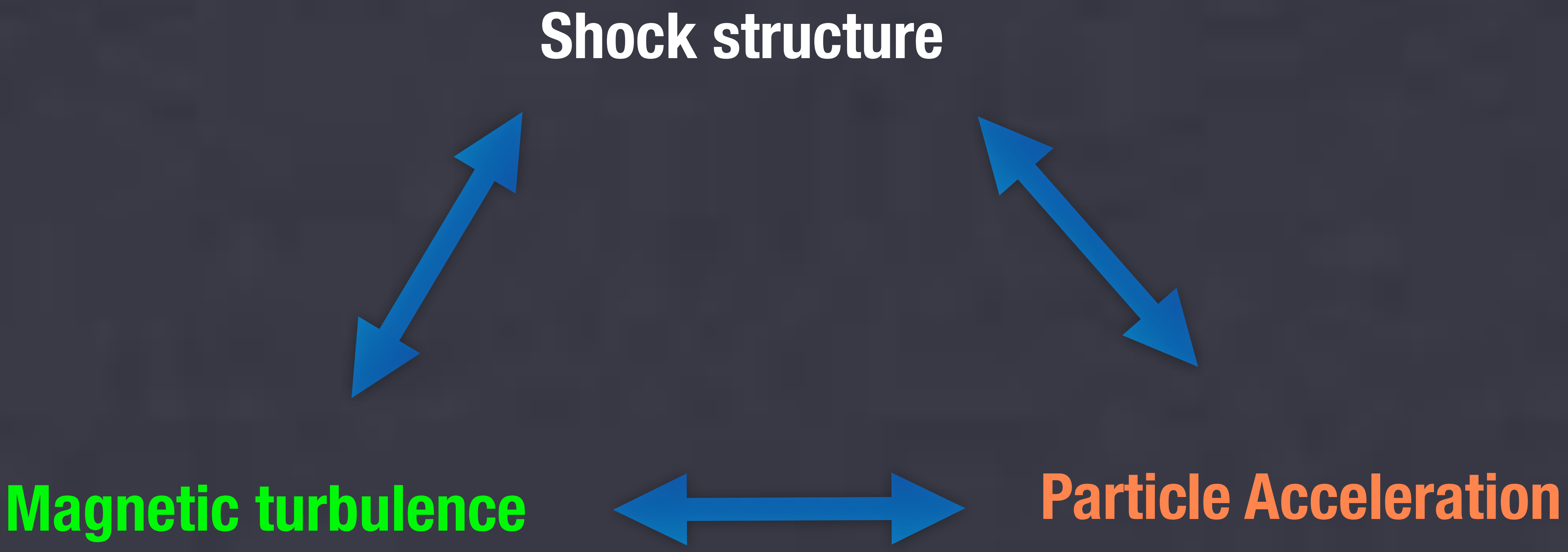
Free energy: converging flows



We need to understand the microphysics of collisionless shocks with plasma simulations

Collisionless shocks

- Complex interplay between micro and macro scales and nonlinear feedback



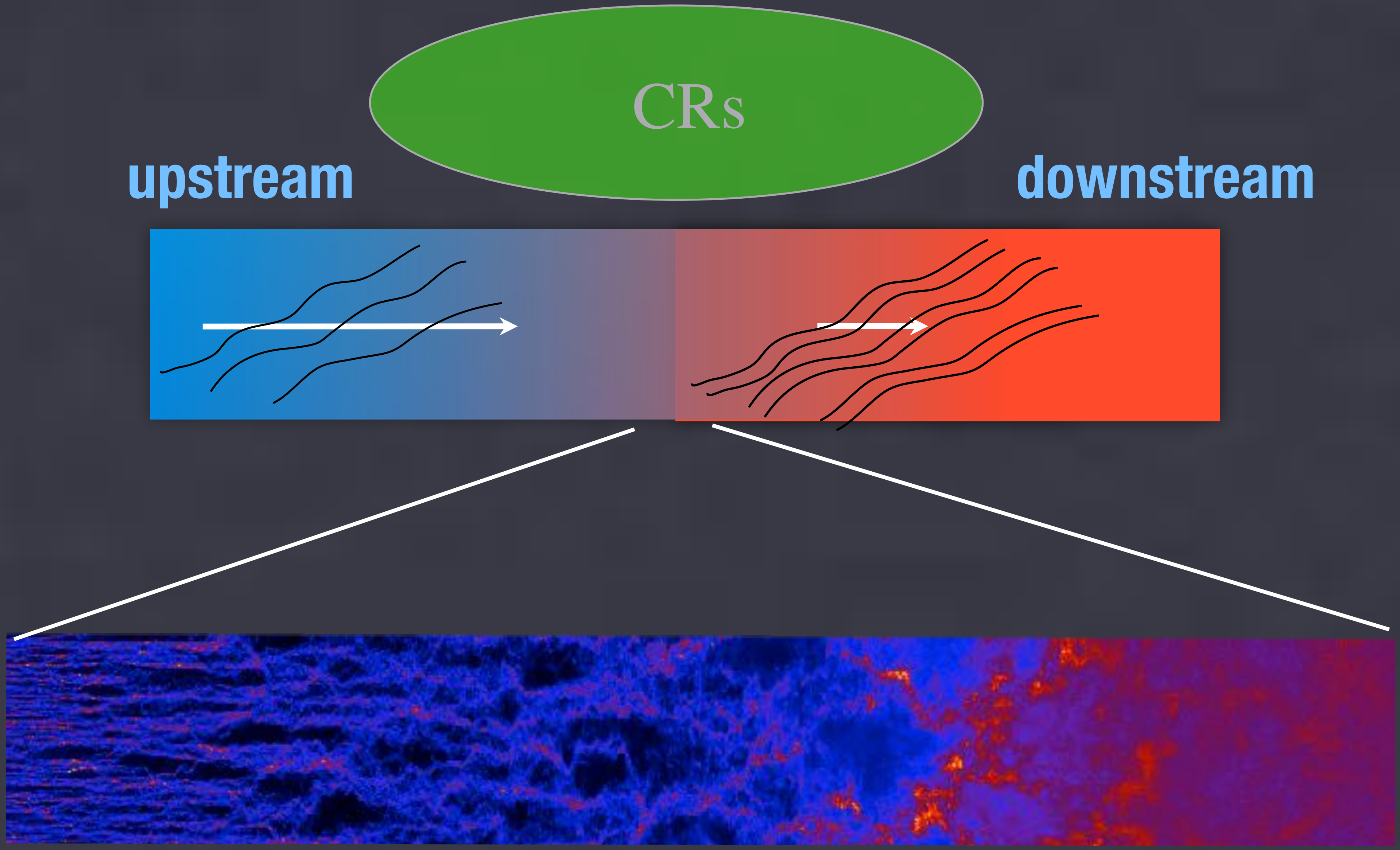
Shock structure

Magnetic turbulence

Particle Acceleration

Collisionless shocks

- Complex interplay between micro and macro scales and nonlinear feedback

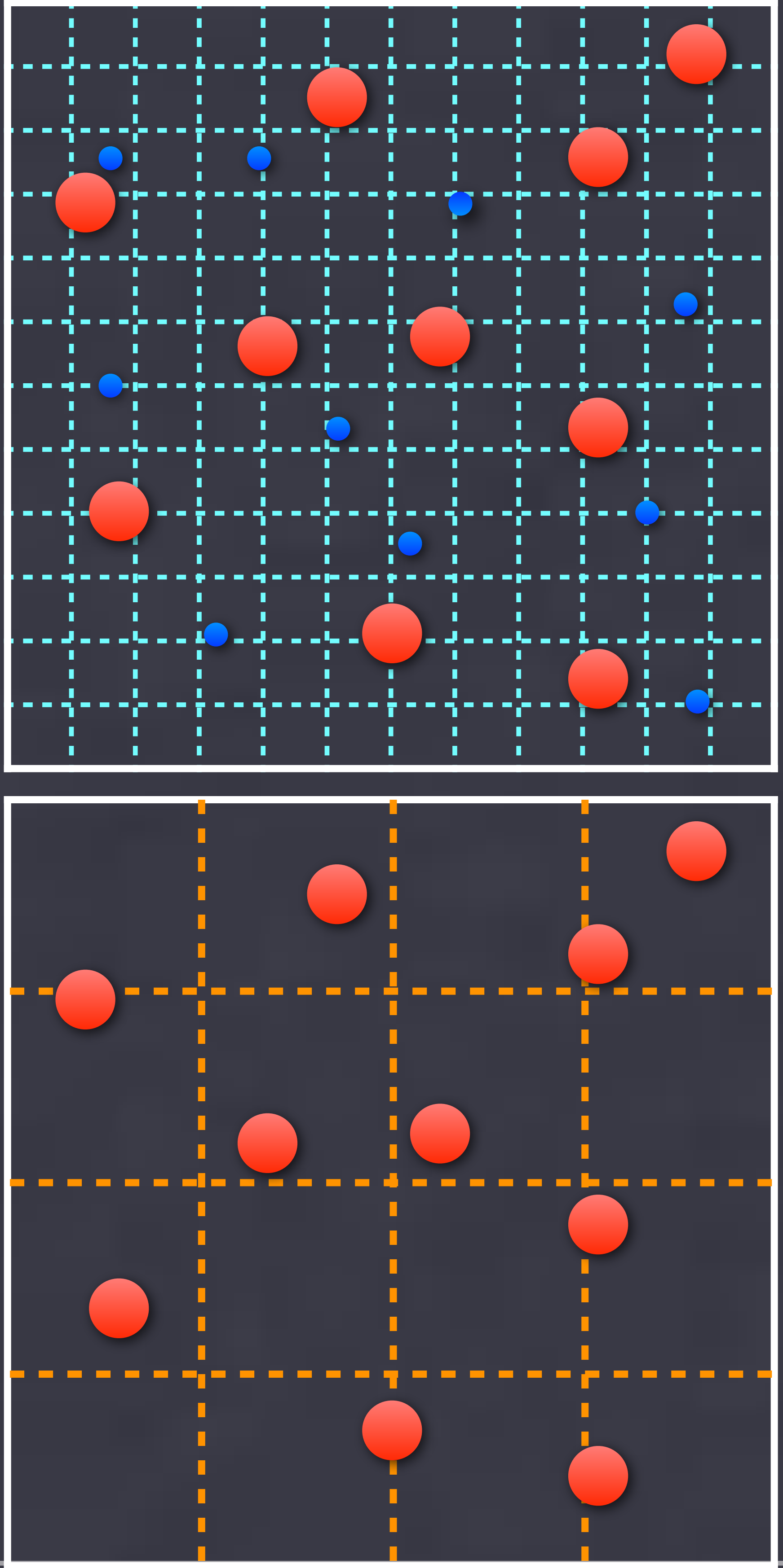


Collisionless shocks from first principles

- **Full particle in cell**: TRISTAN-MP code

  (Spitkovsky 2008, Niemiec+2008, Stroman+2009, Amano & Hoshino 2007-2010, Riquelme & Spitkovsky 2010, Sironi & Spitkovsky 2011, Park+2012, Niemiec+2012, Guo+14,...)

  - Define electromagnetic field on a grid
  - Move particles via Lorentz force
  - Evolve fields via Maxwell equations
  - Computationally expensive!



- **Hybrid approach**: dHybrid code

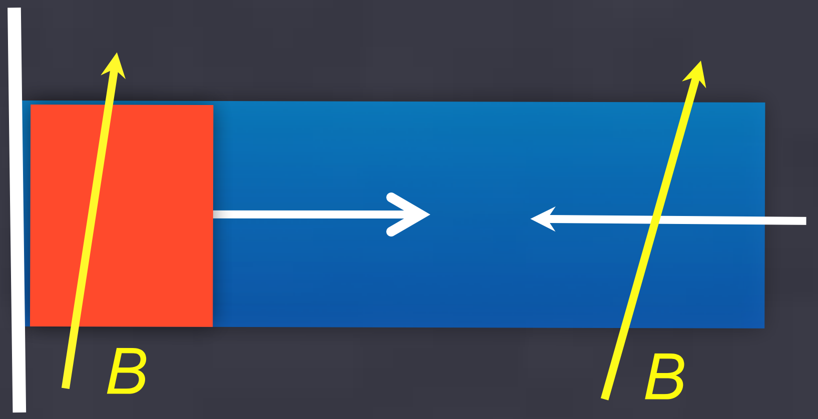
  Fluid electrons - Kinetic protons

  (Winske & Omidi; Lipatov 2002; Giacalone et al.; Gargaté & Spitkovsky 2012, DC & Spitkovsky 2013, 2014)

  - massless electrons for more macroscopic time/length scales

Survey of Collisionless Shocks

We simulated relativistic and nonrelativistic shocks for a range of upstream B fields and flow compositions, ignoring pre-existing turbulence.



Main findings:

Dependence of shock mechanism on upstream magnetization

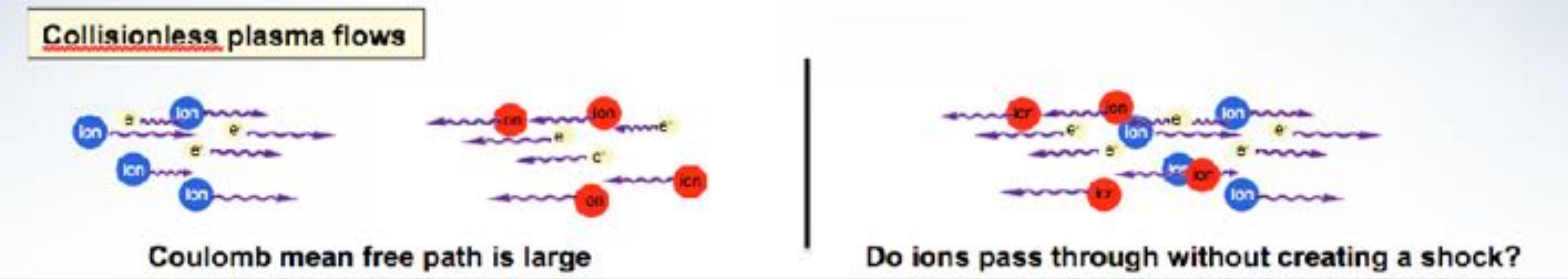
Ab-initio particle acceleration in relativistic shocks

Shock structure and acceleration in non-relativistic shocks

- Ion acceleration vs Mach # in quasipar shocks; DSA; D coeff.
- Evidence for simultaneous e-ion acceleration in parall. shks
- Electron acceleration in quasiperpendicular shocks

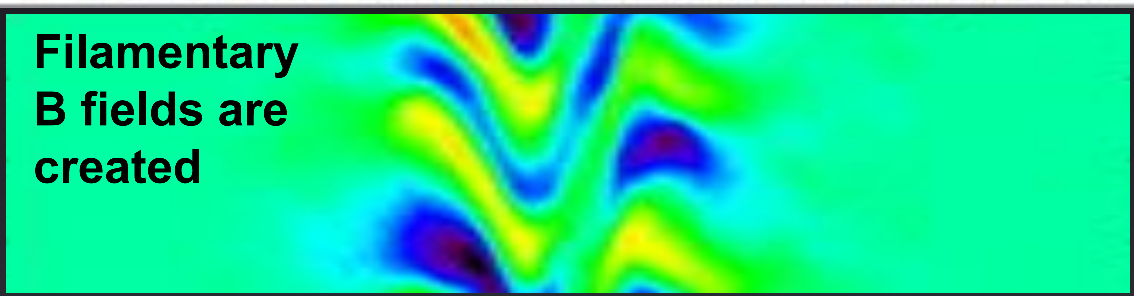
FIeld amplification and CR-induced instabilities

How collisionless shocks work



Coulomb mean free path is large

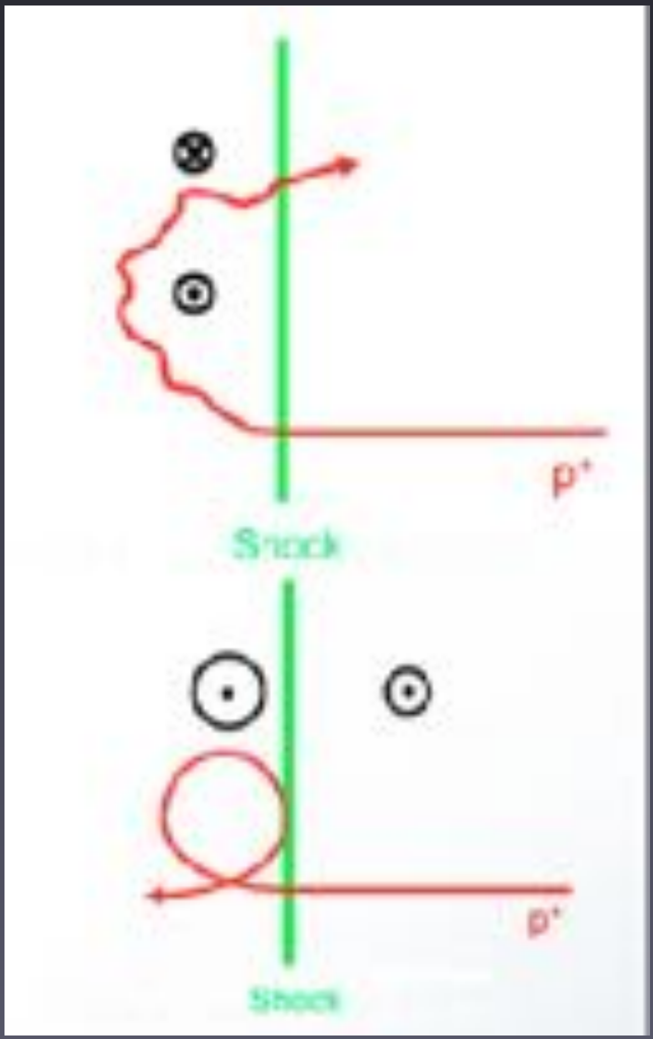
Do ions pass through without creating a shock?


Filamentary B fields are created

Two main mechanisms for creating collisionless shocks:

1) For low initial B field, particles are deflected by self-generated magnetic fields (filamentation/Weibel instability); Alvenic Mach # > 100

2) For large initial B field, particles are deflected by compressed pre-existing fields; Alfvenic Mach # < 100

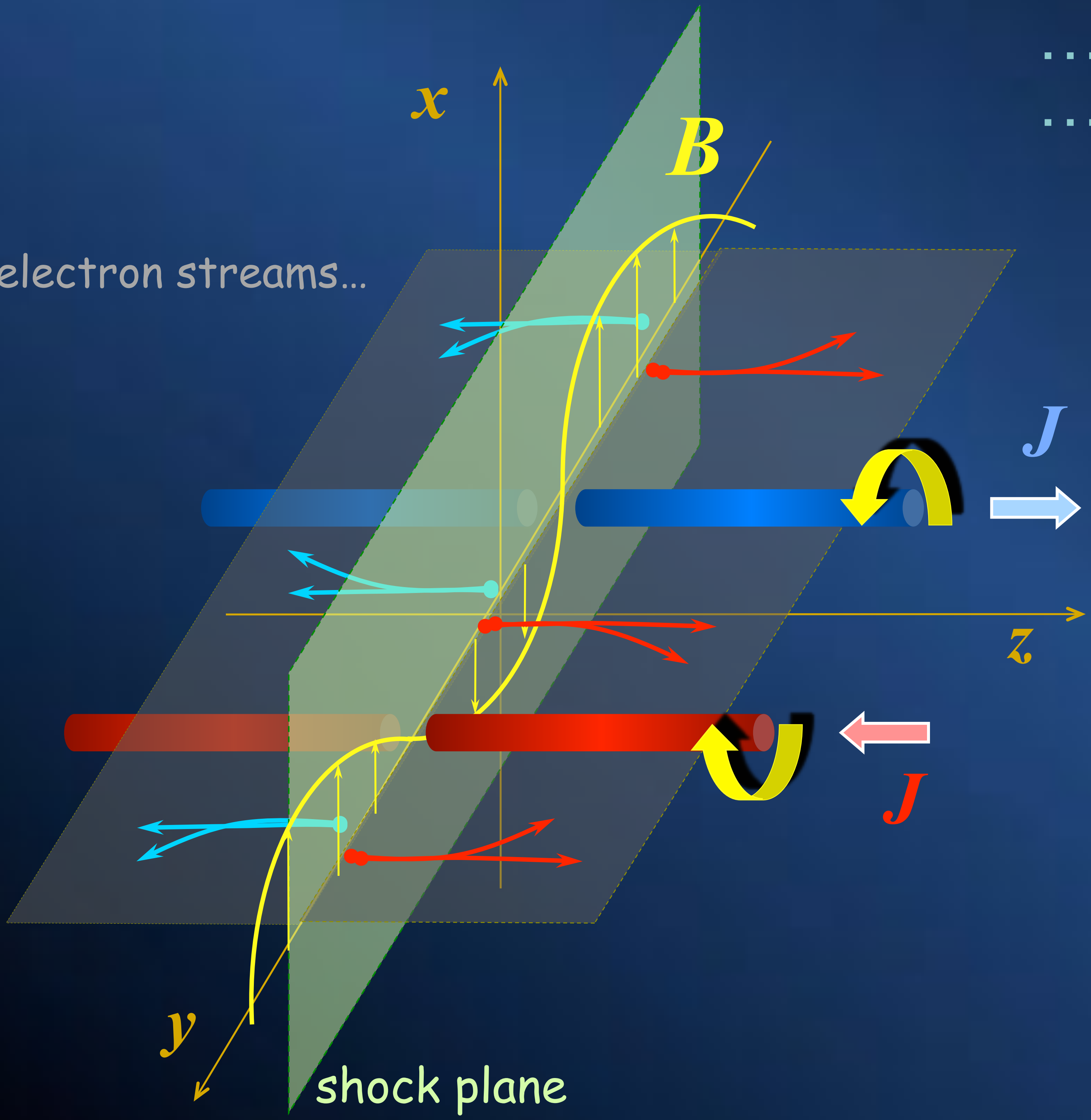


WEIBEL INSTABILITY

(Weibel 1956, Medvedev & Loeb, 1999, ApJ)

… current filamentation …
… *B* – field is generated …

For electron streams…



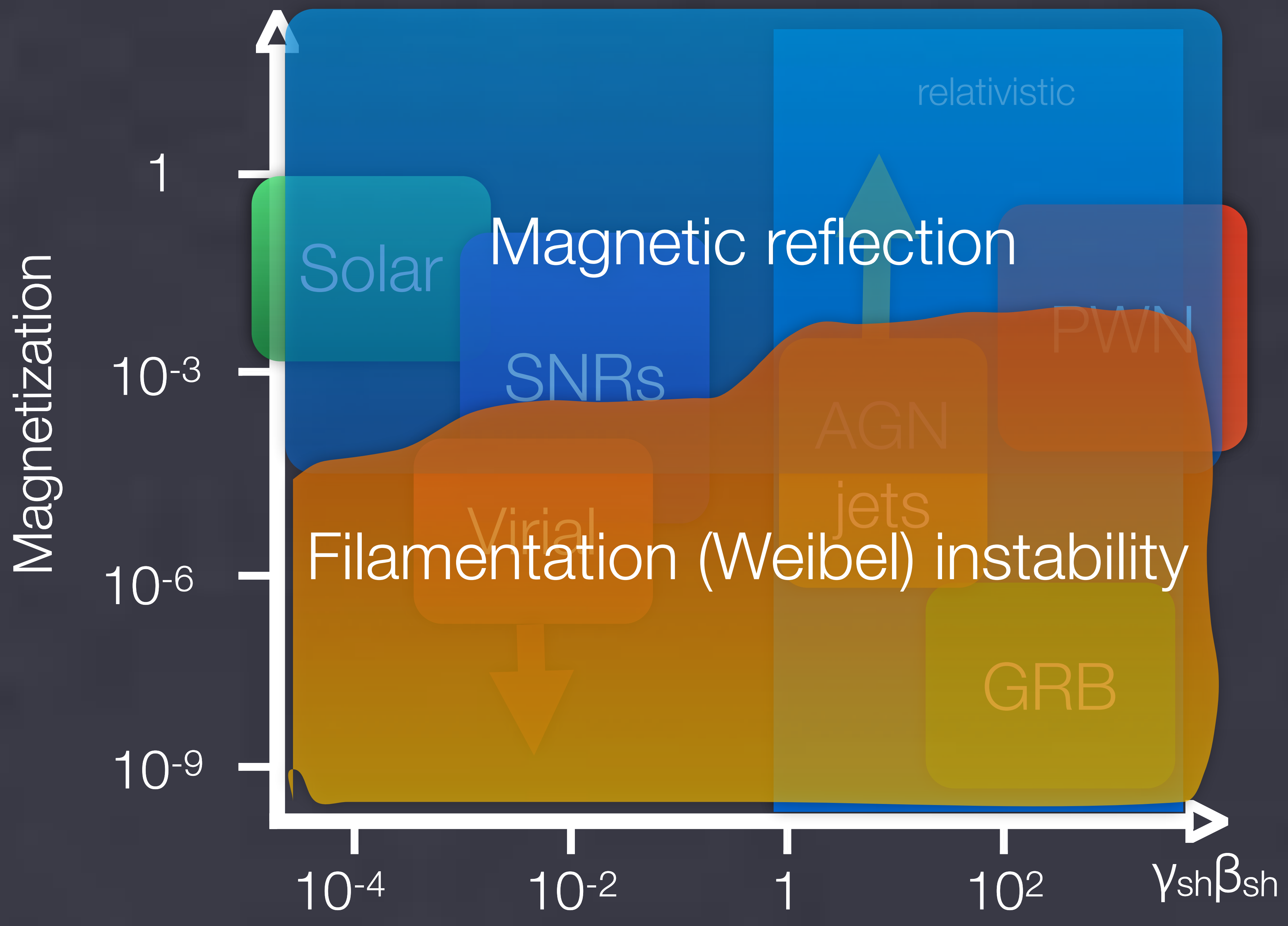
$$\Gamma^2_{\max} \simeq \frac{\omega_p^2}{\gamma}$$

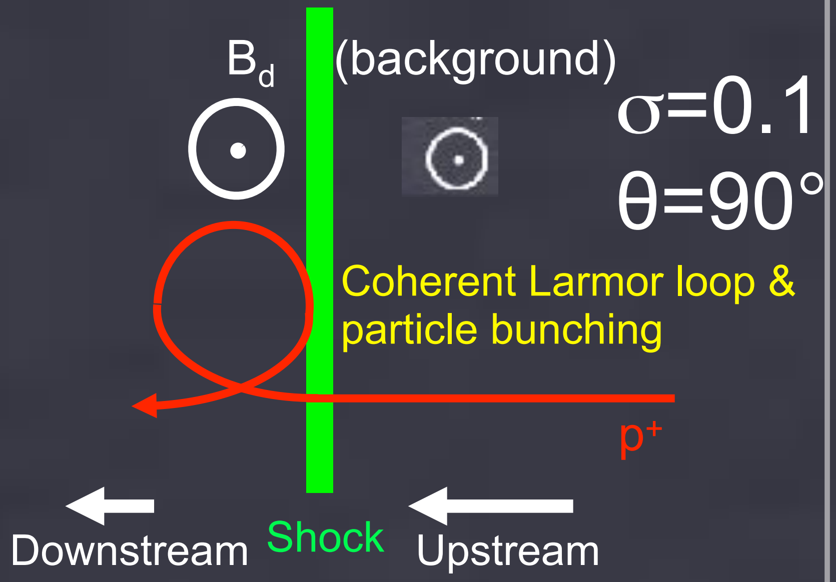
$$k^2_{\max} \simeq \frac{1}{\sqrt{2}} \frac{\omega_p^2}{\gamma_\perp c^2}$$

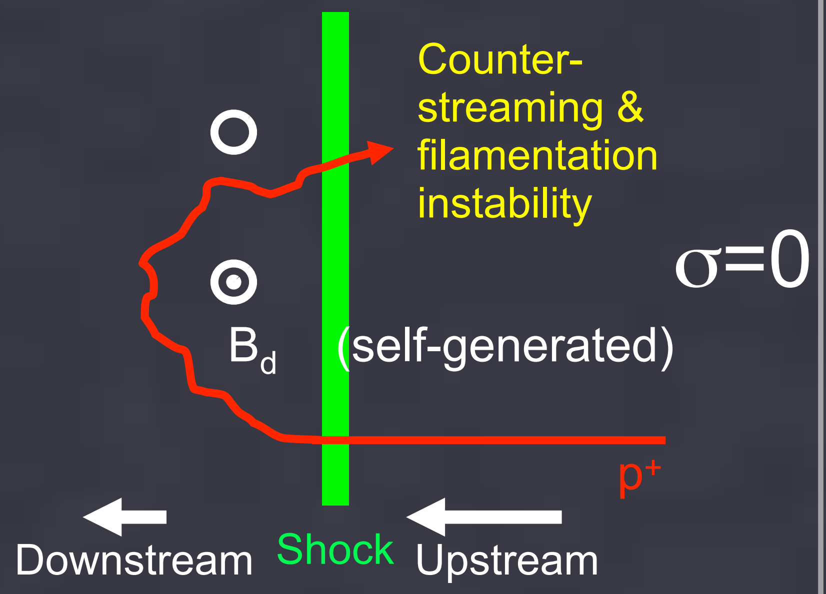
Parameter Space of shocks

$$\sigma \equiv \frac{B^2/4\pi}{(\gamma-1)nmc^2} = \frac{1}{M_A^2} = \left(\frac{\omega_c}{\omega_p}\right)^2 \left(\frac{c}{v}\right)^2 = \left[\frac{c/\omega_p}{R_L}\right]^2$$

$$M_A = \frac{v}{v_A} \qquad M_s = \frac{v}{v_{th}} \qquad \beta = \frac{M_A^2}{M_s^2} \qquad \frac{m_i}{m_e}$$



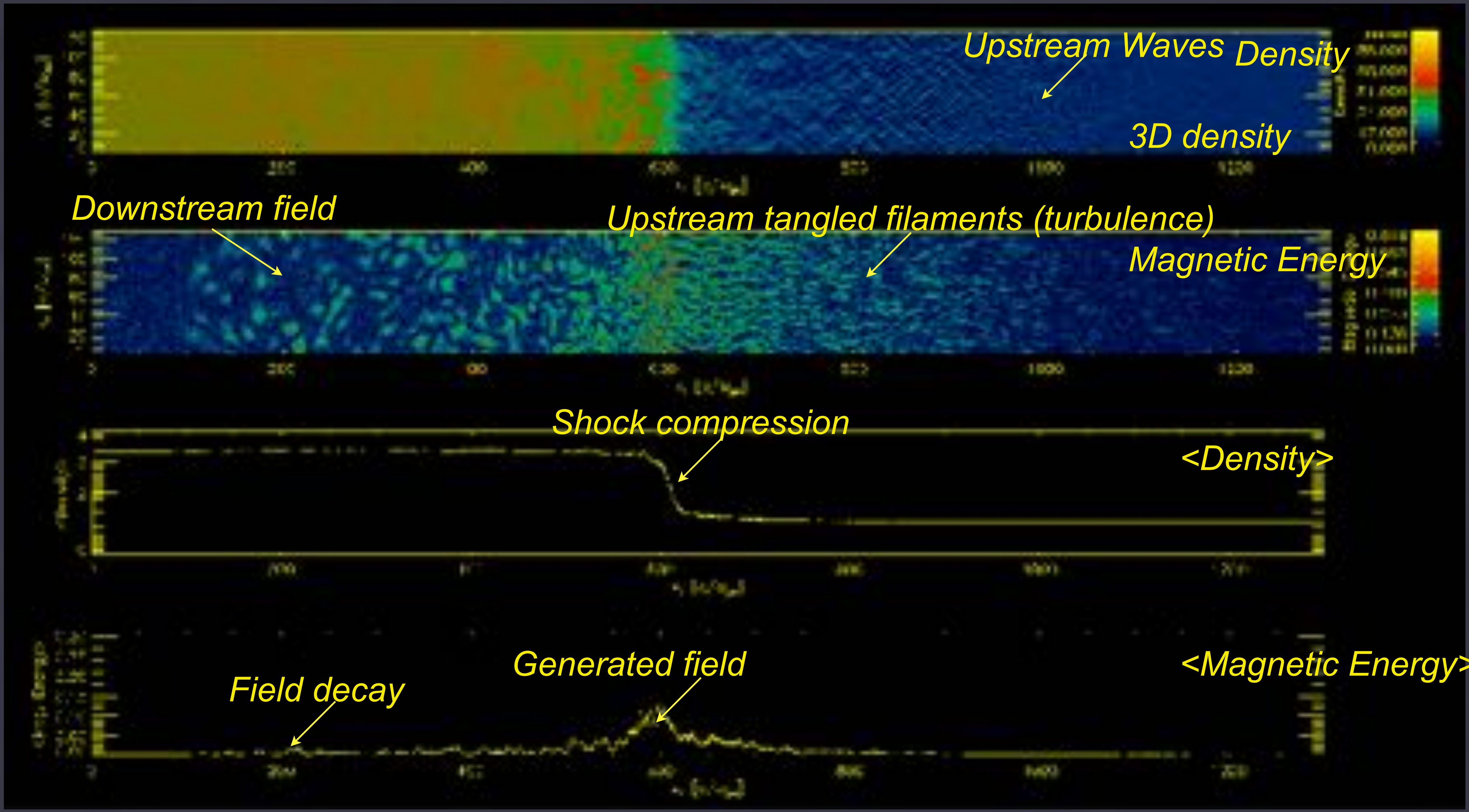




Collisionless shocks

Structure of an unmagnetized relativistic pair shock

min max



Upstream Waves

Density

3D density

Downstream field

Upstream tangled filaments (turbulence)

Magnetic Energy

Shock compression

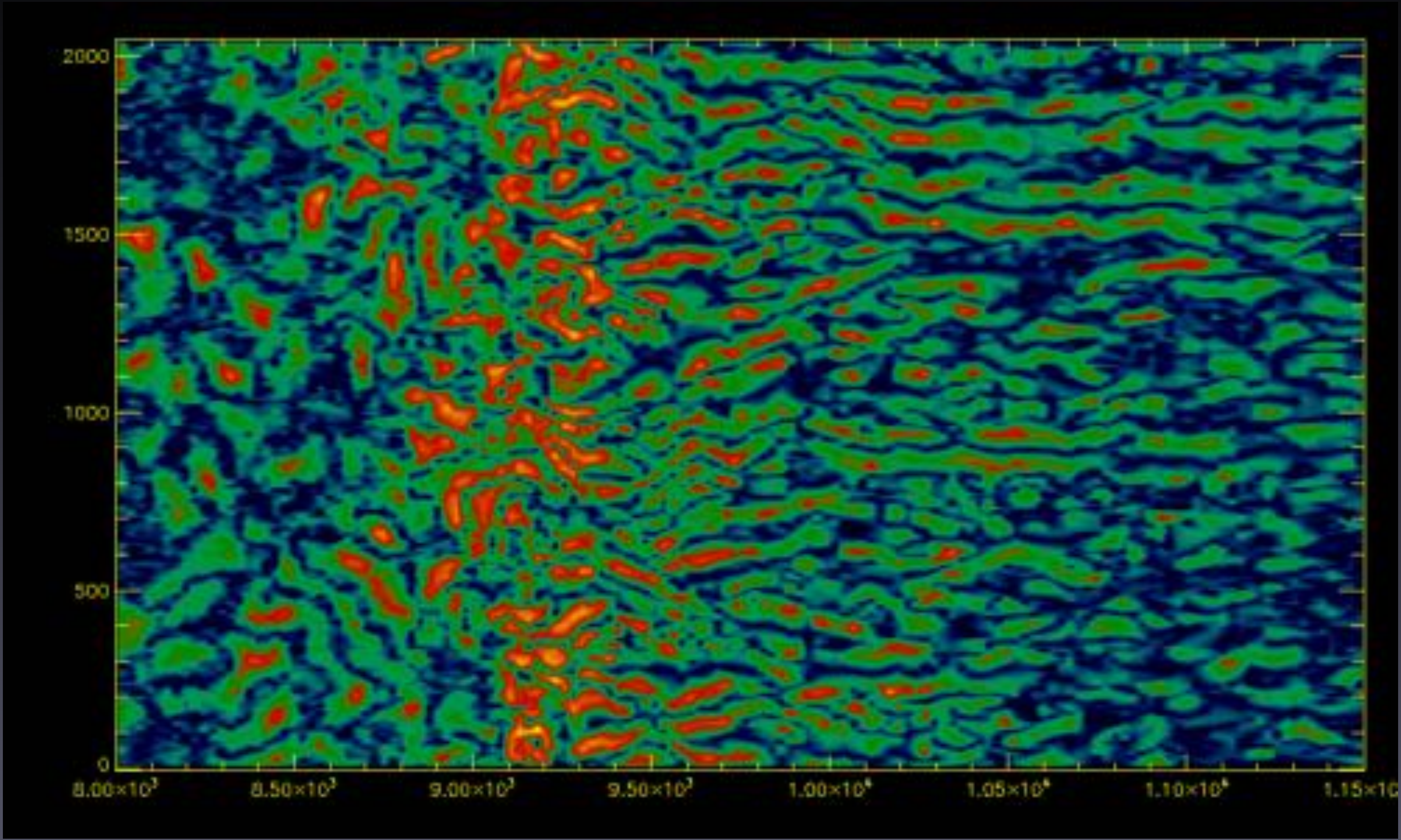
<Density>

Field decay

Generated field

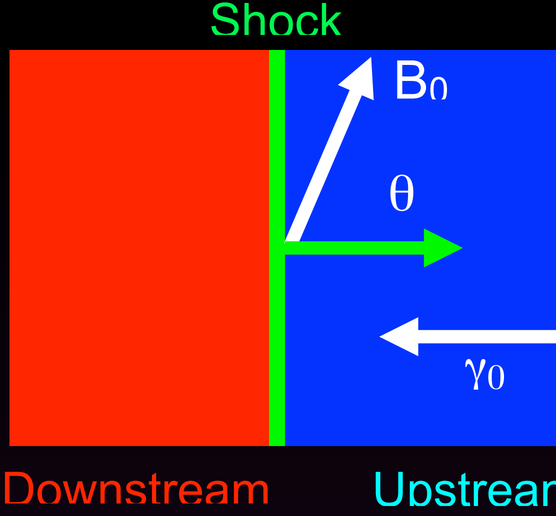
<Magnetic Energy>

Unmagnetized pair shock: particle trajectories



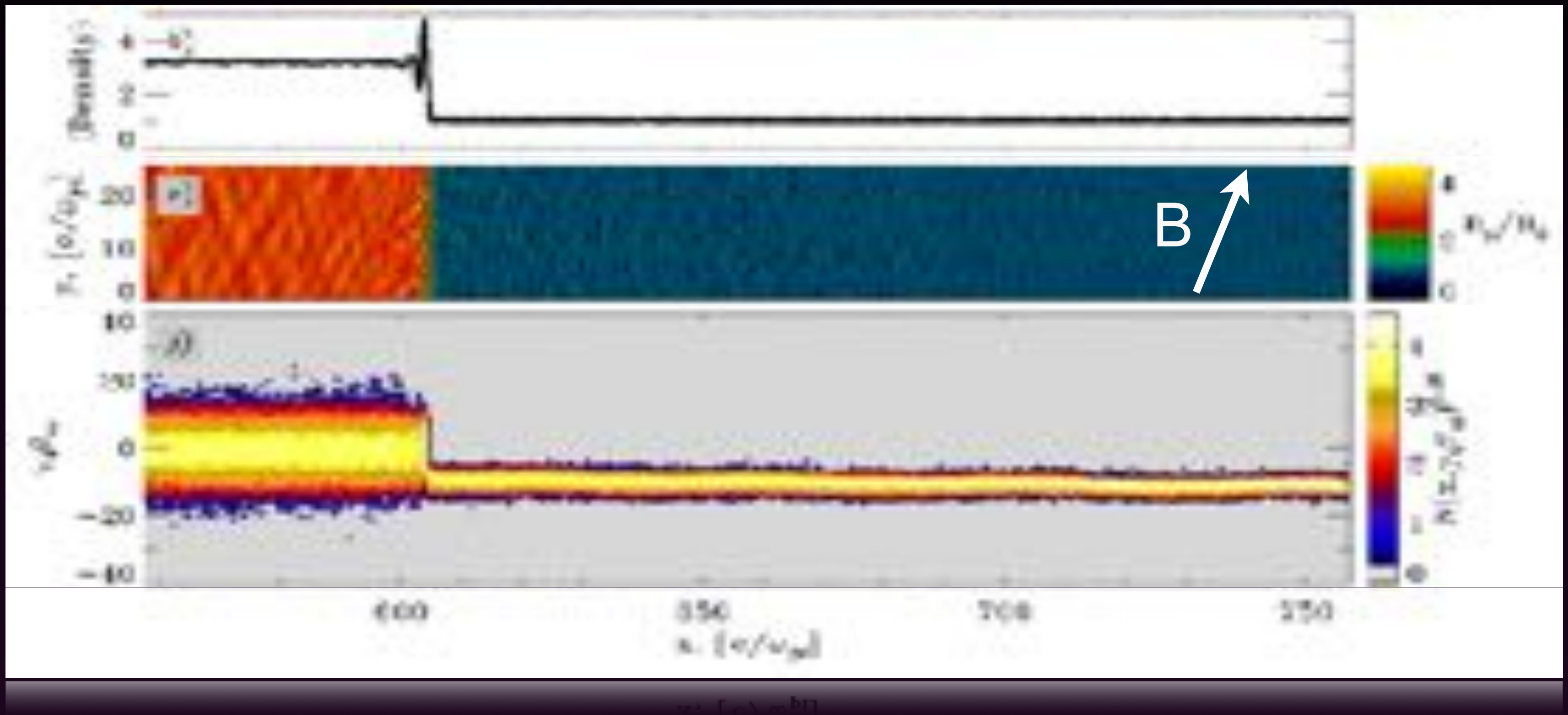
color: magnetic energy density;

Perpendicular vs parallel shocks



- Quasi-perpendicular shocks: mediated by magnetic reflection

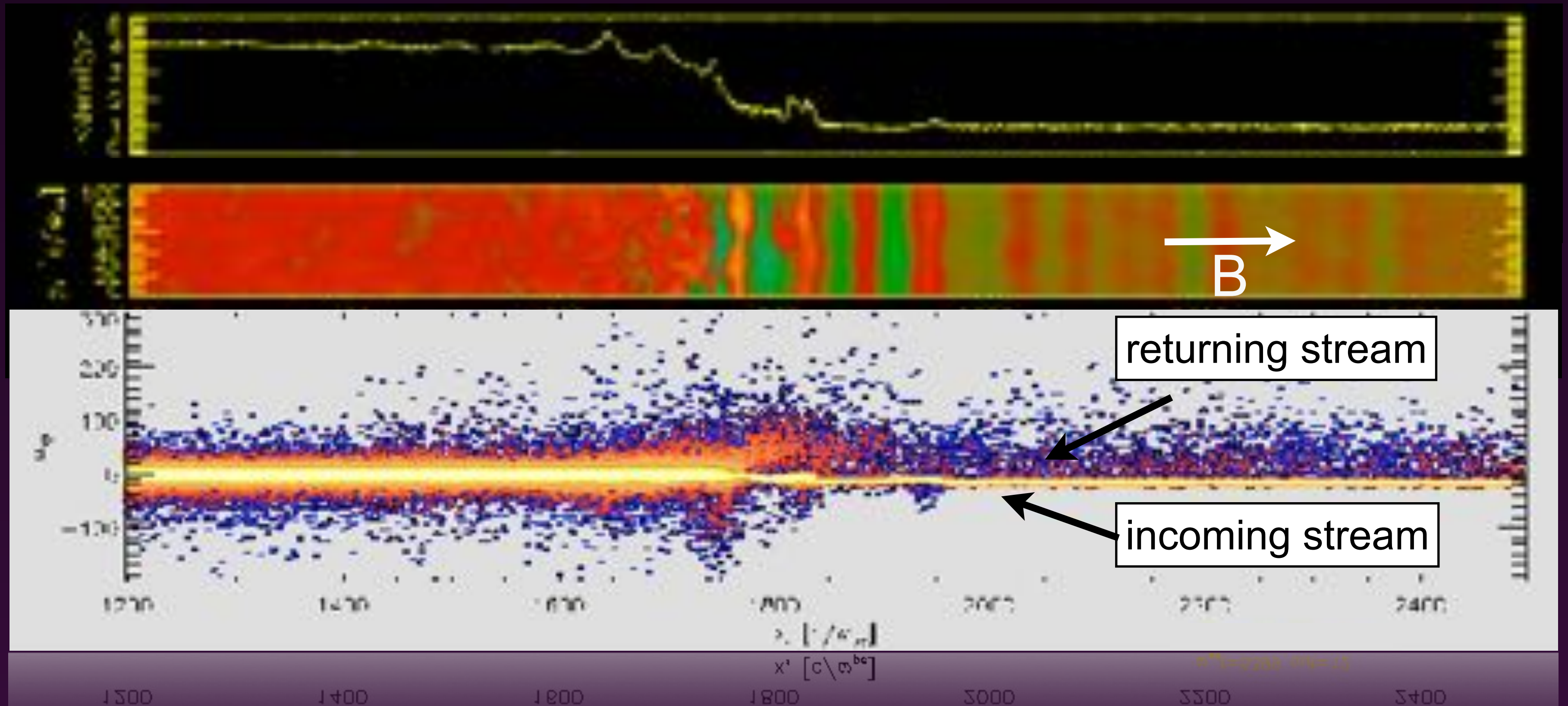
$\sigma=0.1$
$\theta=75°$
$\gamma_0=15$
e^--p^+



(Sironi and AS 11)

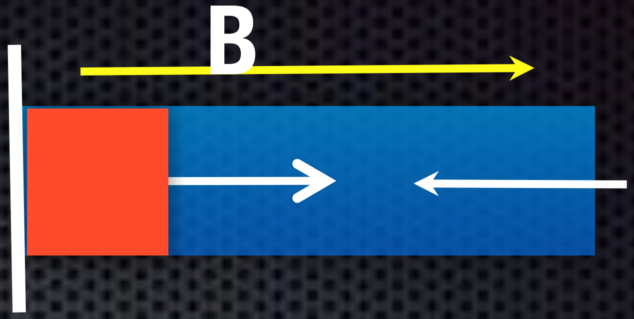
- Quasi-parallel shocks: instabilities amplify transverse field component

$\sigma=0.1$
$\theta=15°$
$\gamma_0=15$
e^--p^+



(Sironi & AS 11)

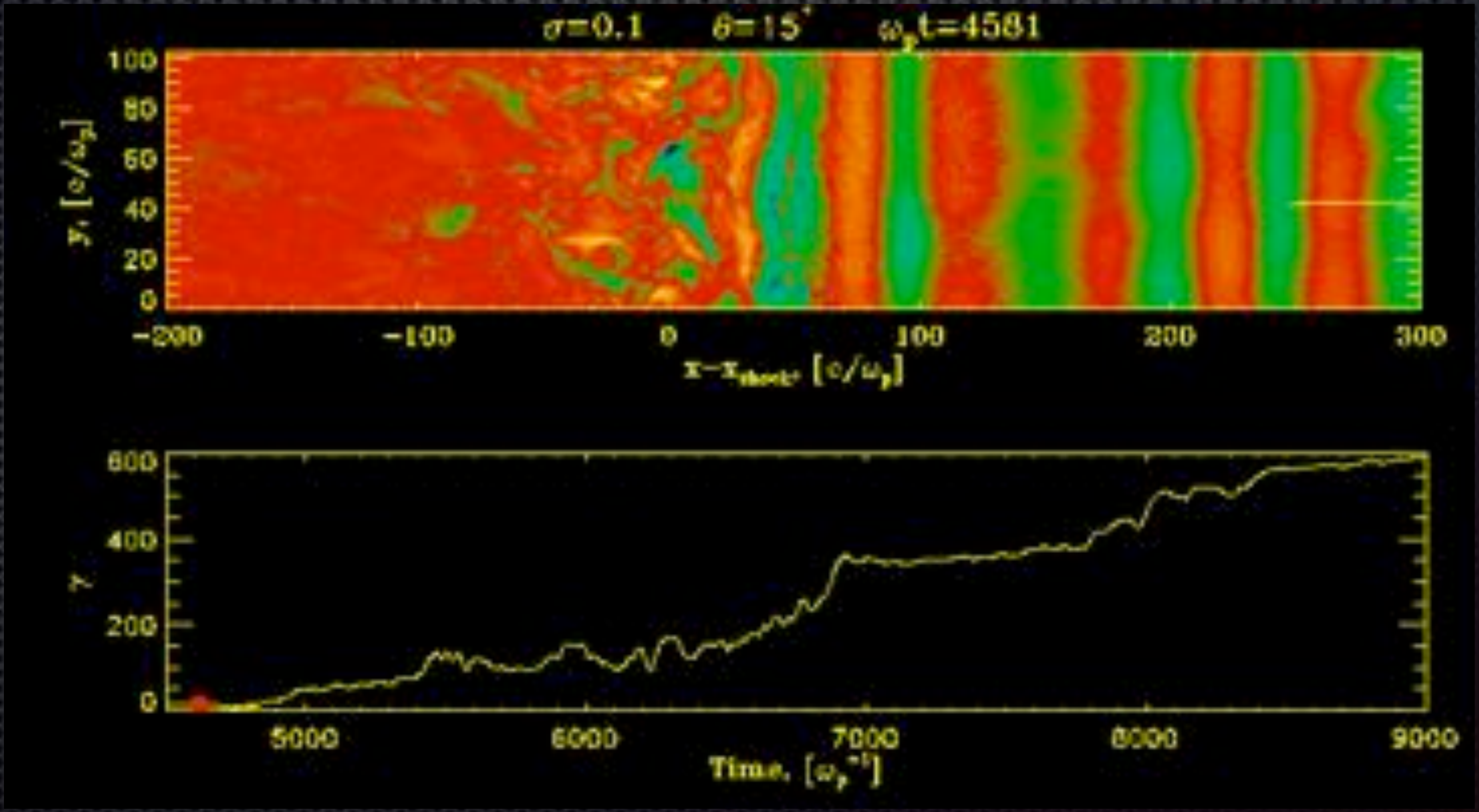
Particle acceleration



Magnetized shock (parallel, e-p): scattering on self-generated upstream waves

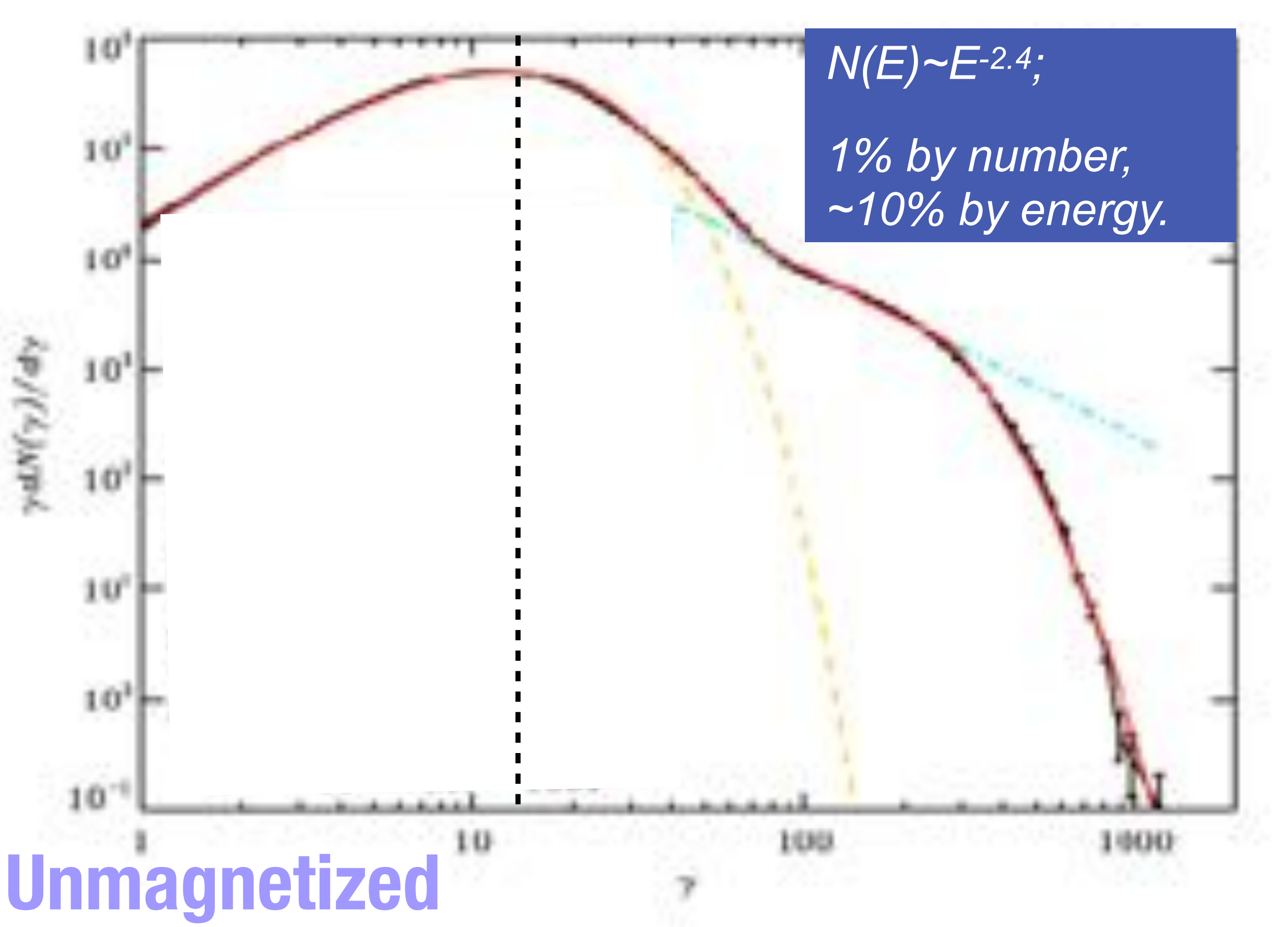
Transverse Magnetic Field

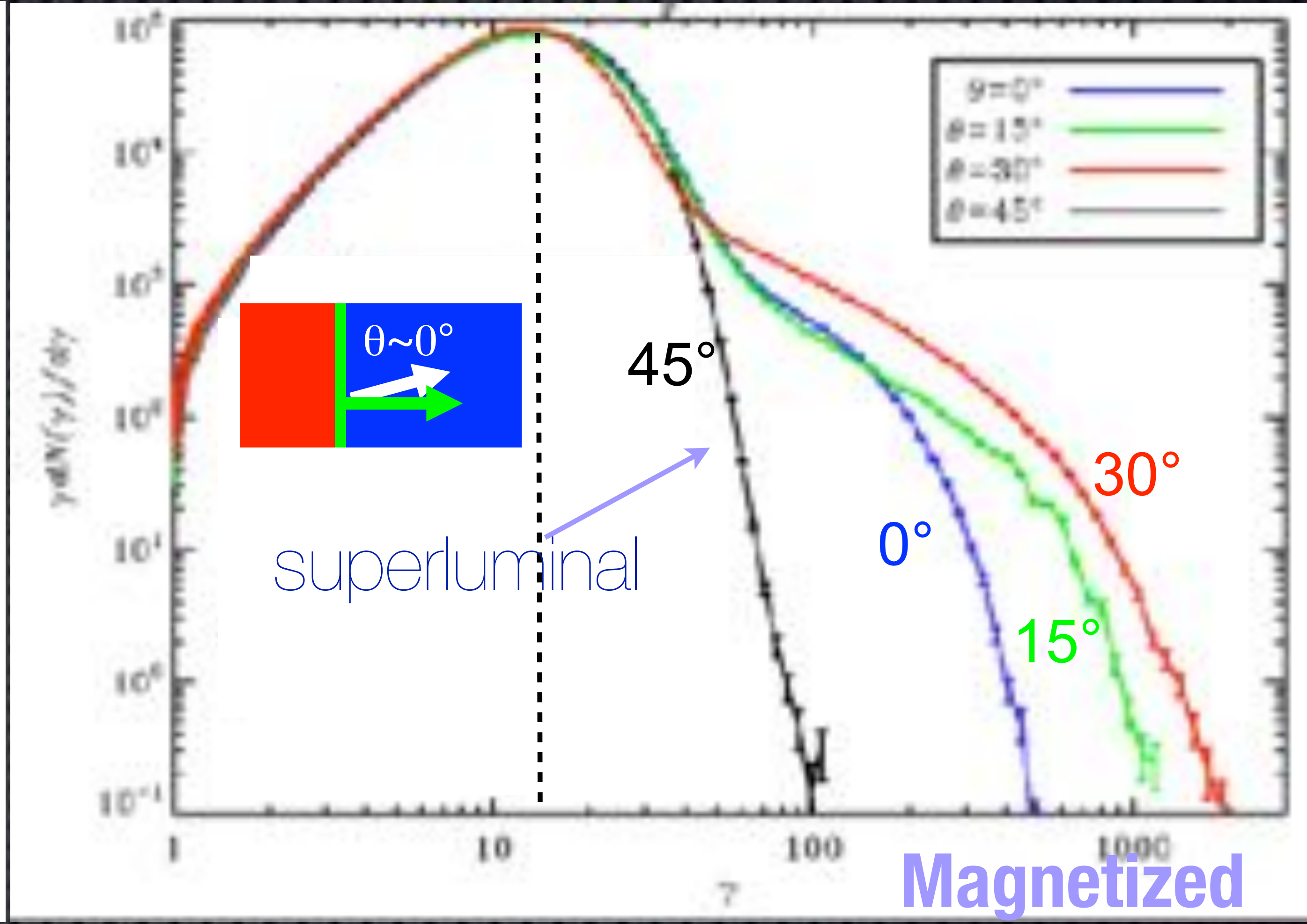
Particle energy

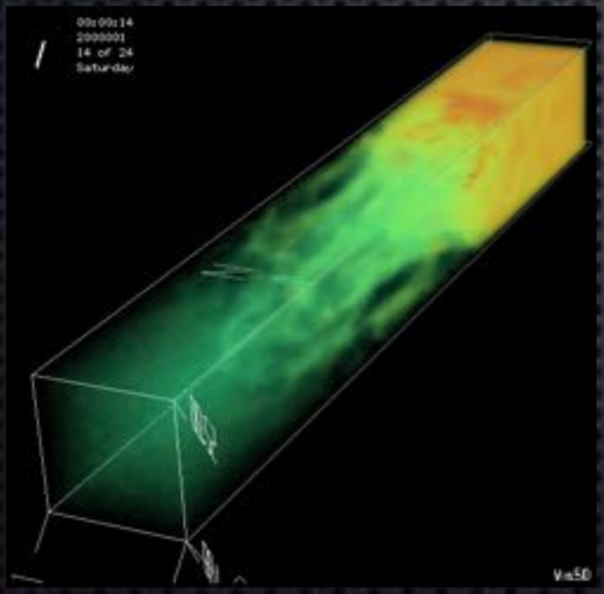


Particle acceleration

Sironi & AS 09





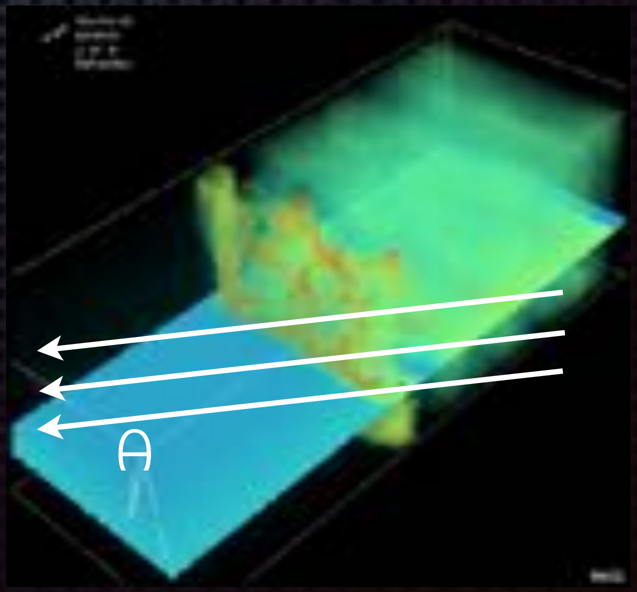


Conditions for acceleration in relativistic shocks:

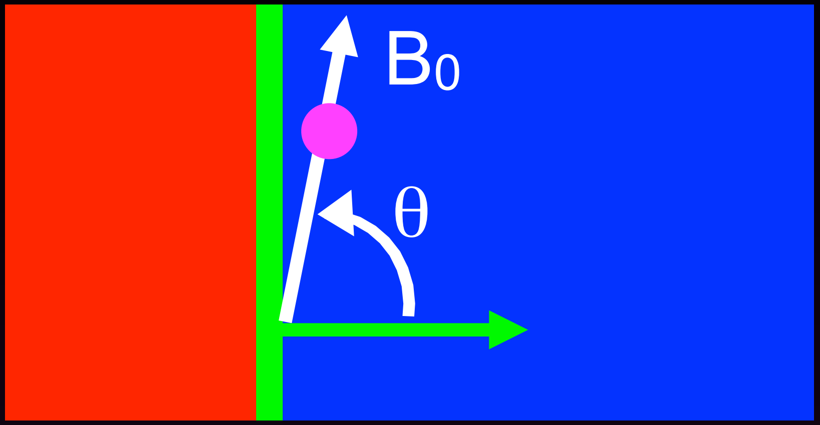
low magnetization of the flow

or quasi-parallel B field (θ<34°/Γ);

electrons & ions behave similarly



Superluminal vs subluminal shocks



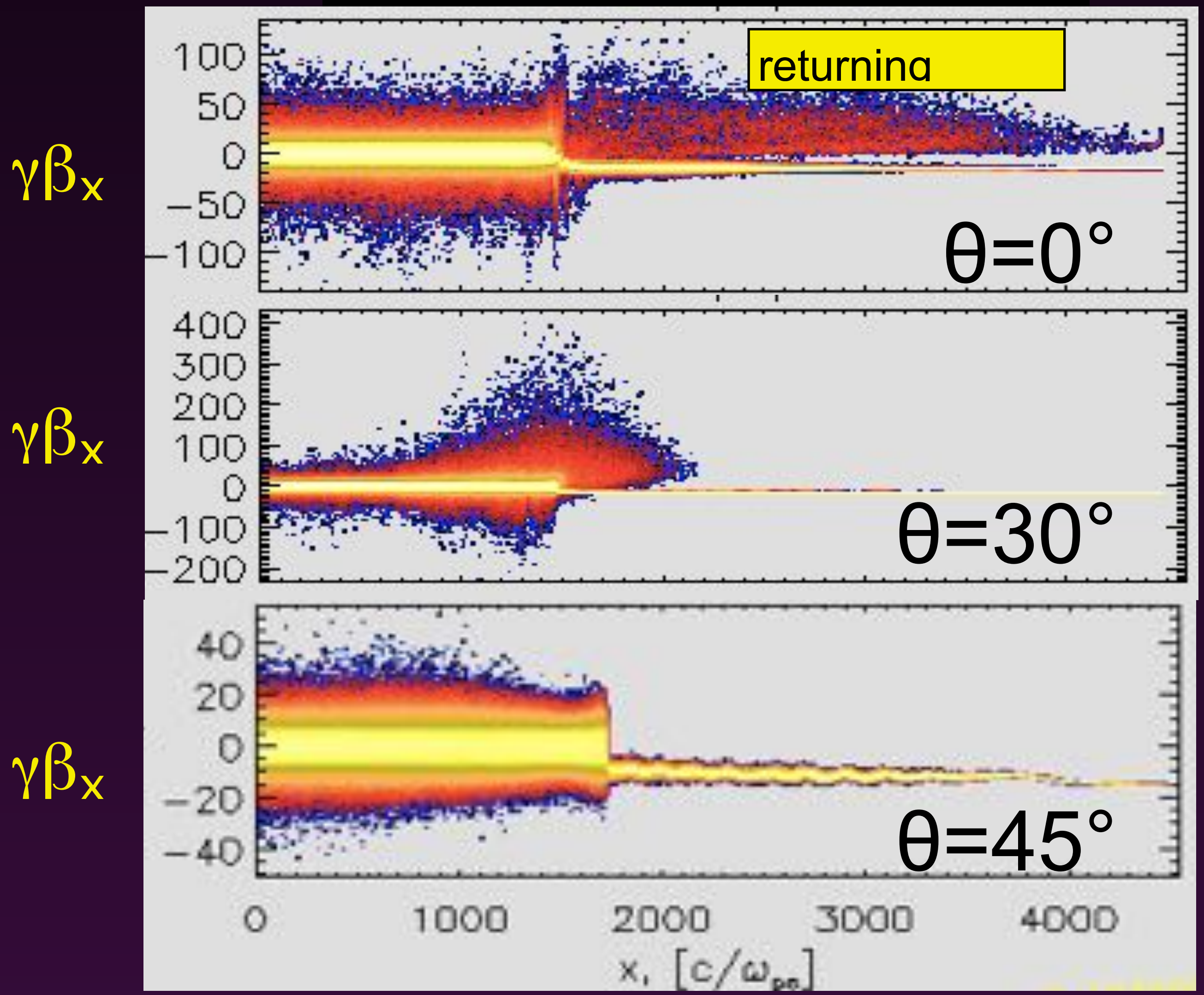
σ is large → particles slide along field lines

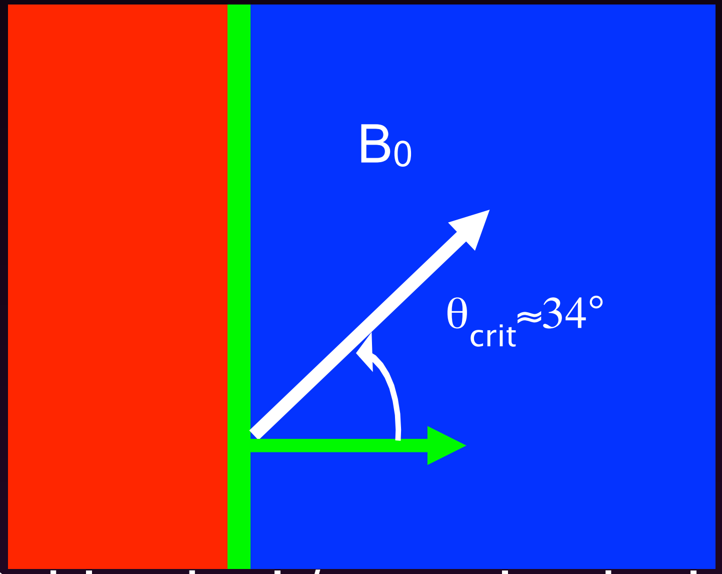
θ is large → particles cannot outrun the shock

unless v>c ("superluminal" shock)

⇒ no returning particles in superluminal shocks

$\sigma=0.1$ $\gamma_0=15$ e^--p^+ shock





Subluminal / superluminal boundary at $\theta\sim 34°$

→ Fermi acceleration should be suppressed

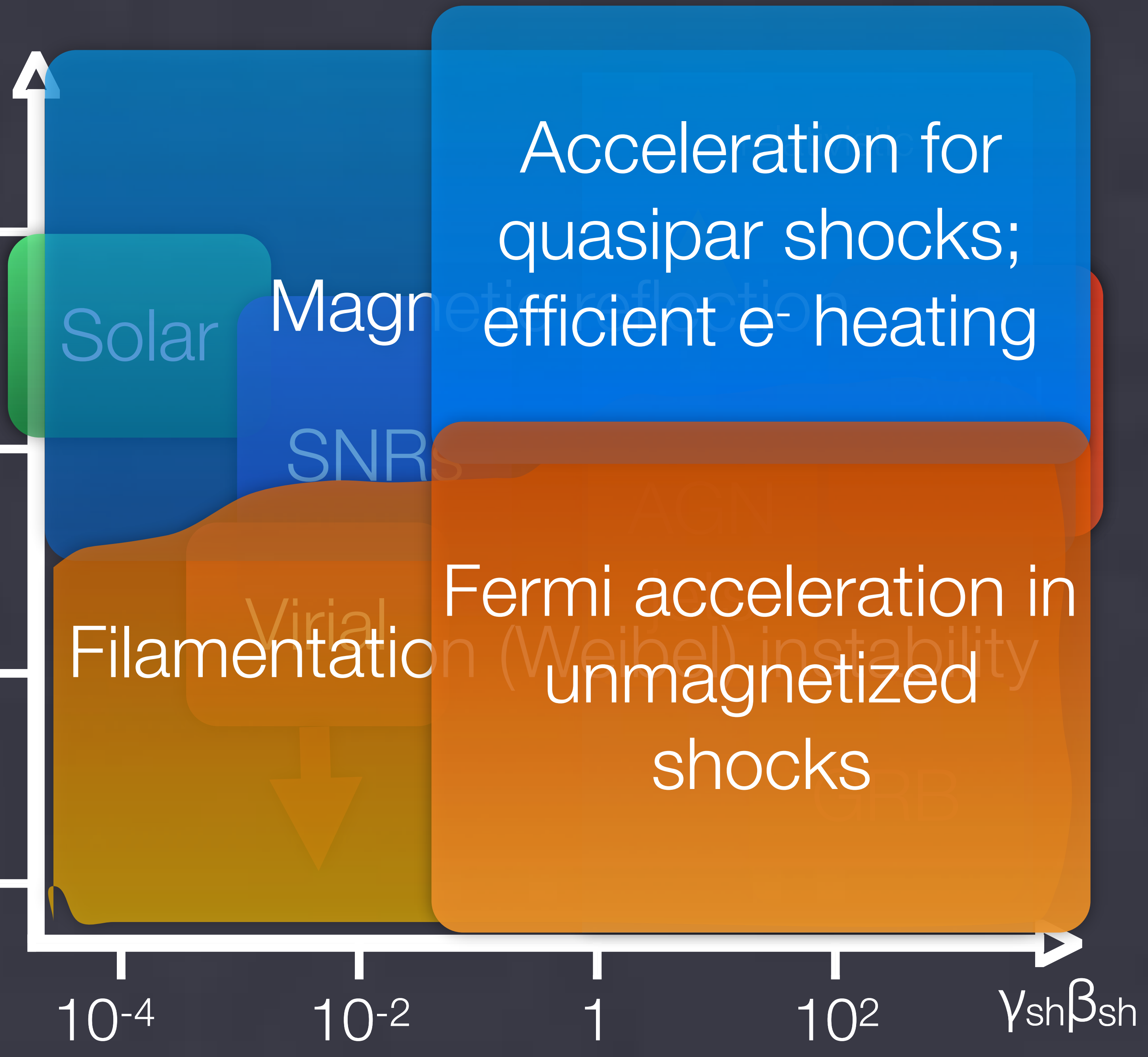
If $\sigma>10^{-3}$, particle acceleration only for:

$\theta<\theta_{crit}\approx 34°$ (downstream frame)

$\theta'<34°/\gamma_0<<1$ (upstream frame)

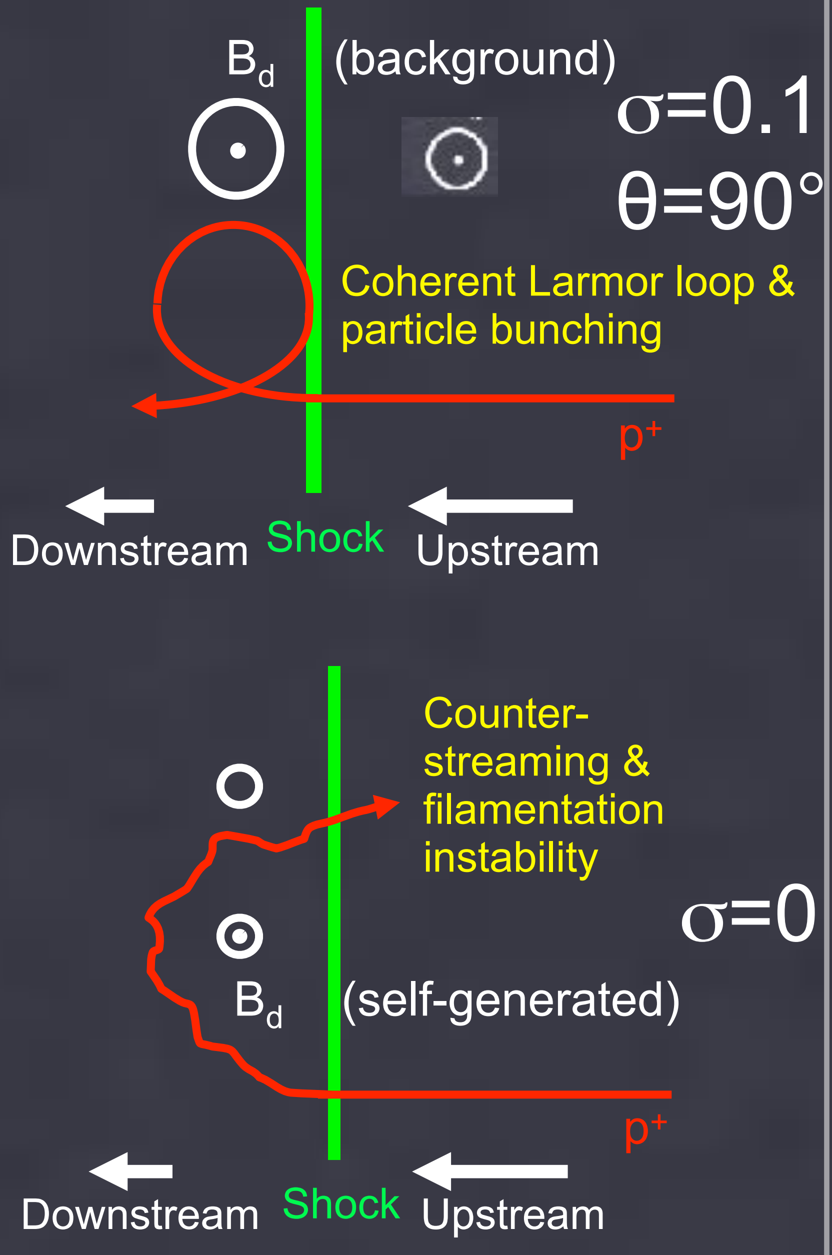
Parameter Space of shocks

$$\sigma \equiv \frac{B^2/4\pi}{(\gamma-1)nmc^2} = \frac{1}{M_A^2} = \left(\frac{\omega_c}{\omega_p}\right)^2\left(\frac{c}{v}\right)^2 = \left[\frac{c/\omega_p}{R_L}\right]^2$$



Acceleration for quasipar shocks; efficient e- heating

Fermi acceleration in unmagnetized shocks



Parameter Space of shocks

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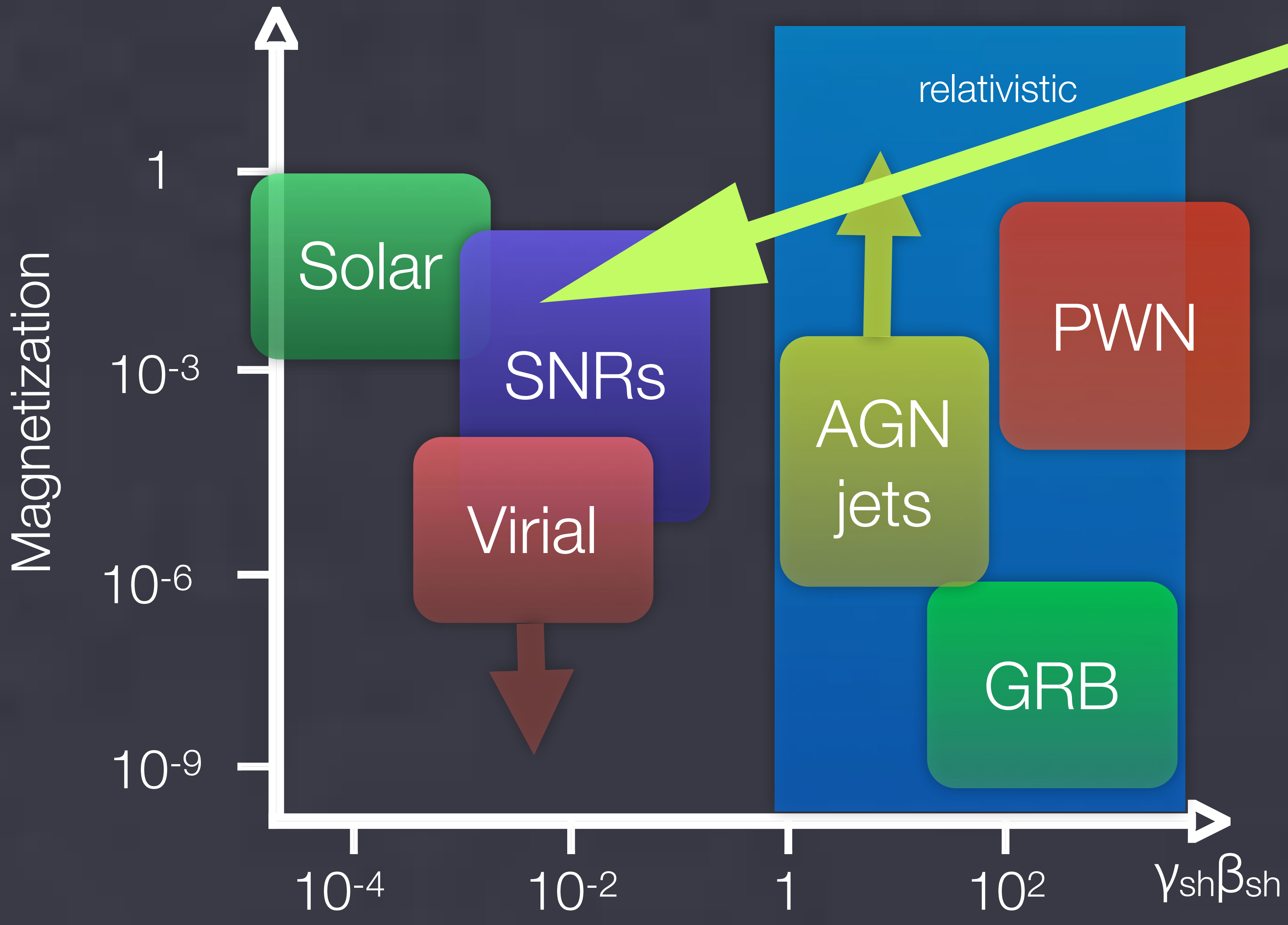


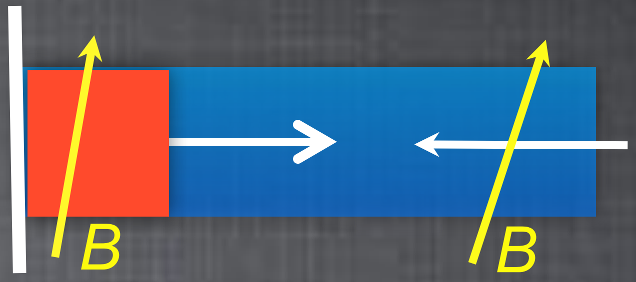
nonrelativistic shocks

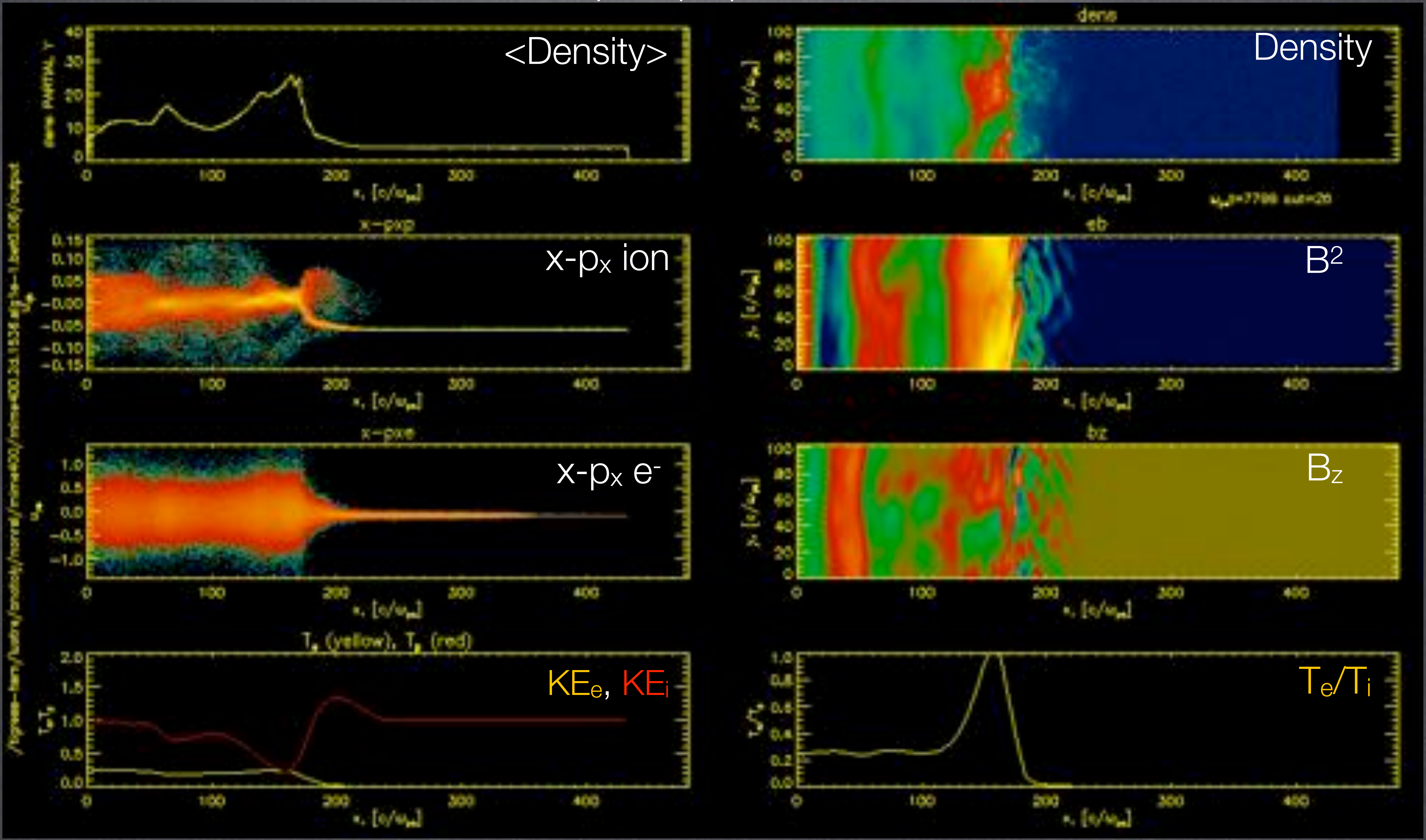
e and ions are different in non-relativistic case

most of our PIC runs are still mildly relativistic (v/c~0.03-0.1c)

Nonrelativistic shocks: shock structure

mi/me=400, v=18,000km/s, Ma=5, quasi-perp 75° inclination

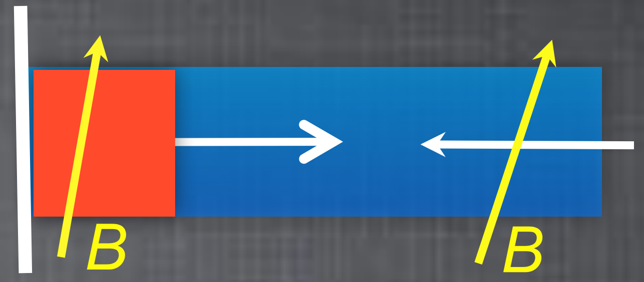


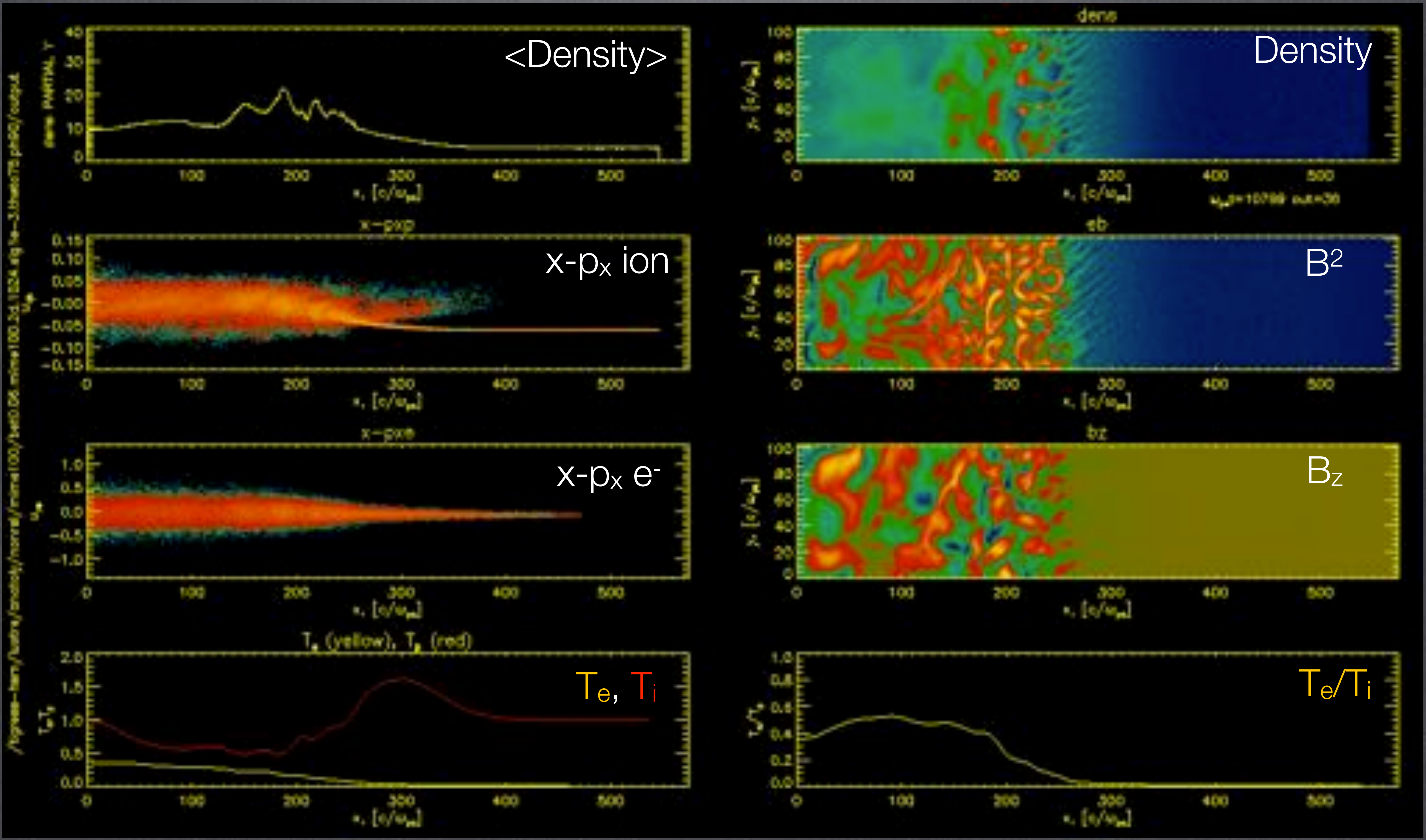


PIC simulation: Shock foot, ramp, overshoot, returning ions, electron heating, whistlers

Nonrelativistic shocks: shock structure

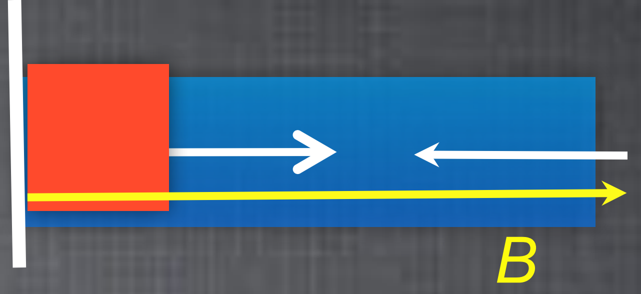
mi/me=100, v=18,000km/s, Ma=45 quasi-perp 75° inclination

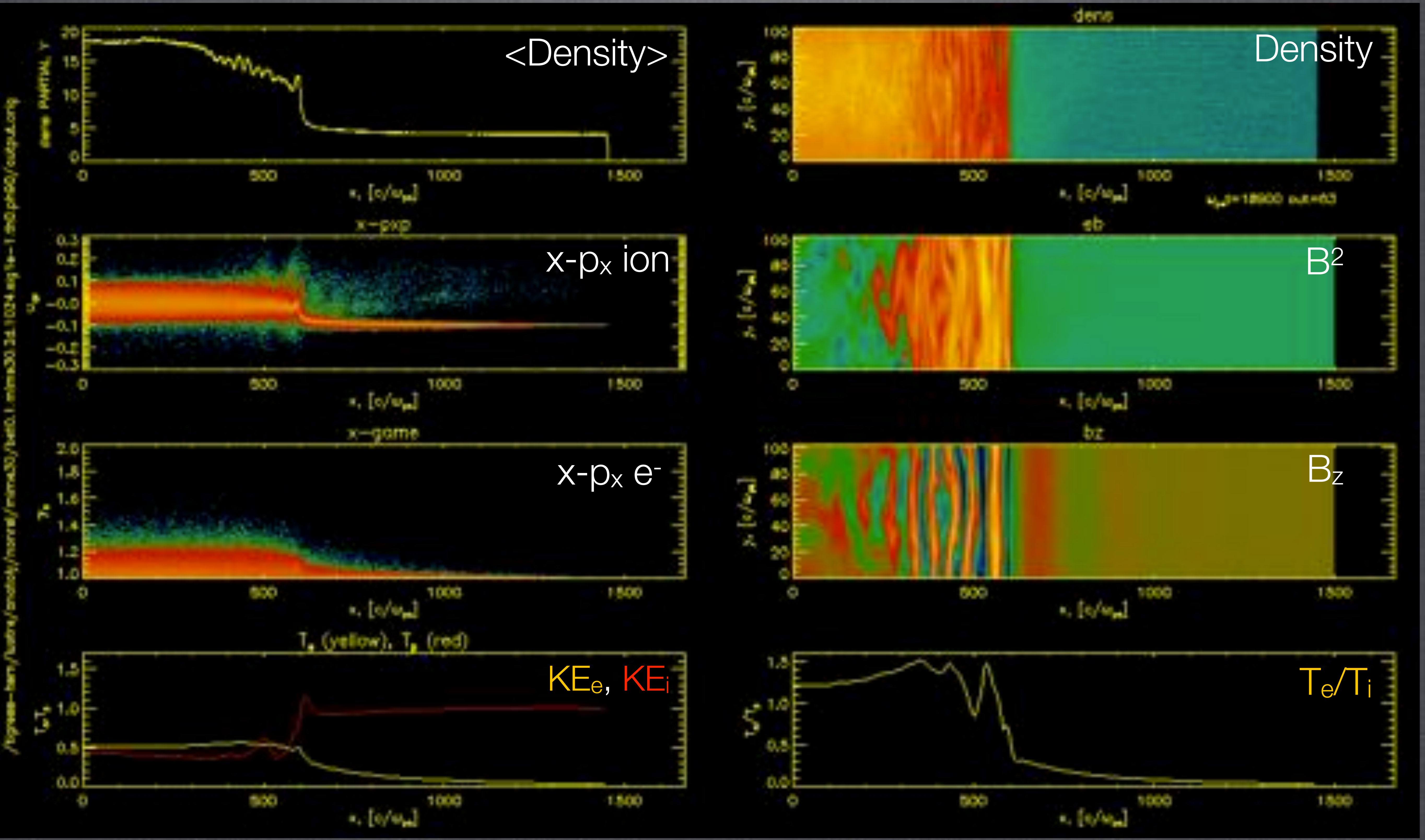




Nonrelativistic shocks: quasiparallel shock

mi/me=30, v=30,000km/s, Ma=5 parallel 0° inclination





PIC simulation: returning ions, reorientation of B field, shock reformations

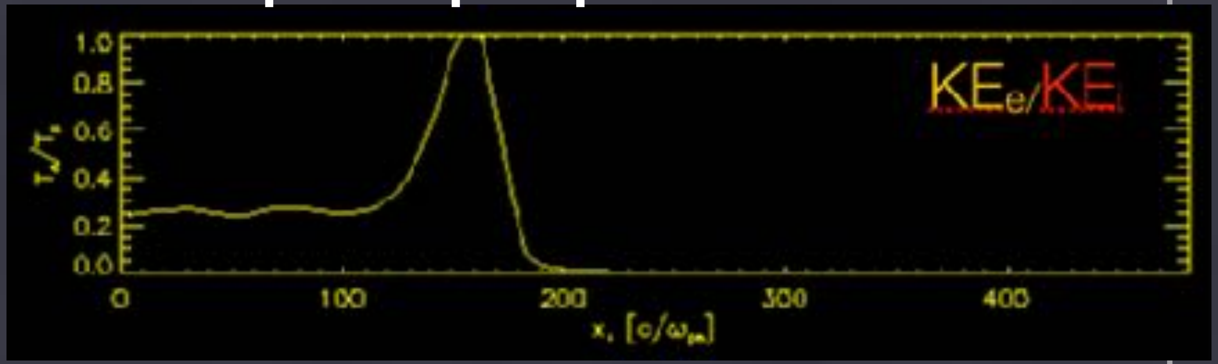
Temperature equilibration?

In full PIC simulations we see very efficient energy exchange between ions and electrons:

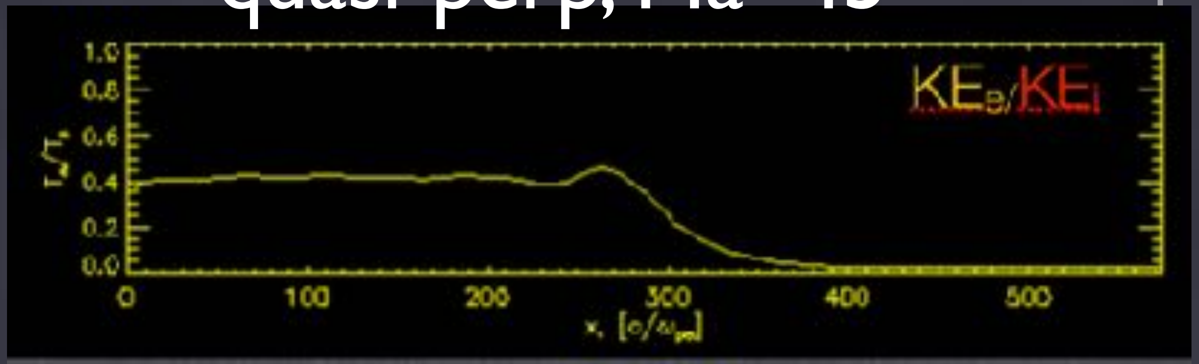
Te/Ti~0.1-0.3 for quasi-perp shocks
Te/Ti ~0.5-1 for quasi-parallel shocks

Physics: shock transition instabilities and upstream electron pre-heating in ion-driven turbulence

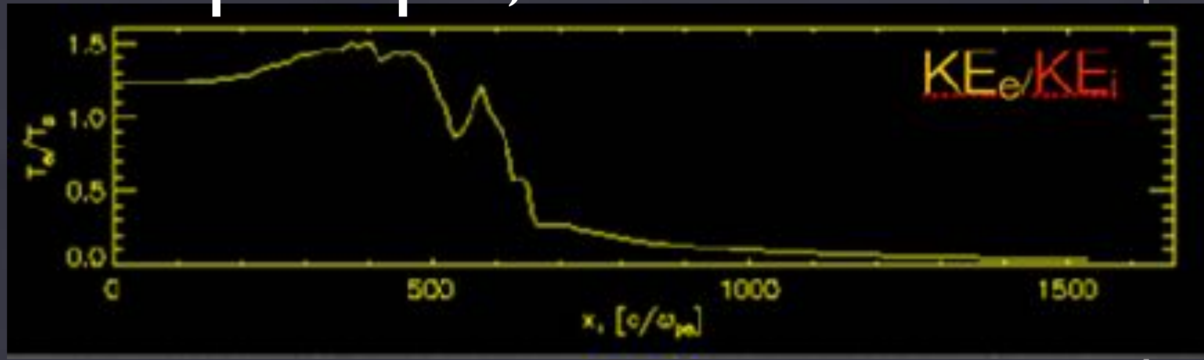
quasi-perp, Ma=10



quasi-perp, Ma=45



quasi-par, Ma=10



Shock acceleration

Two crucial ingredients:

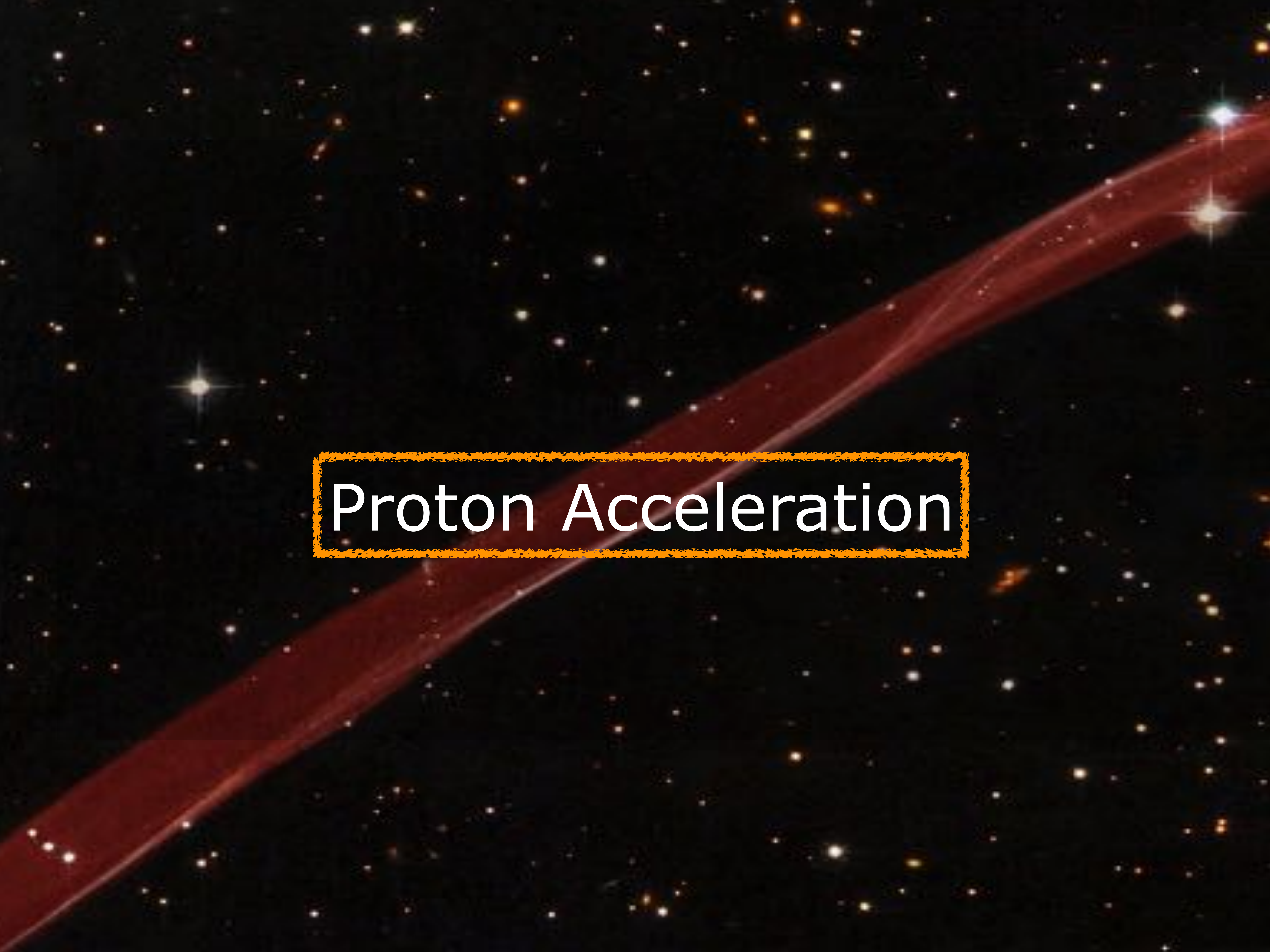
1) ability of a shock to reflect particles back into the upstream (injection)

2) ability of these particles to scatter and return to the shock (pre-existing or generated turbulence)

Generically, parallel shocks are good for ion and electron acceleration, while perpendicular shocks mainly accelerate electrons. *There are many sub-regimes, not fully mapped yet.*

Outline

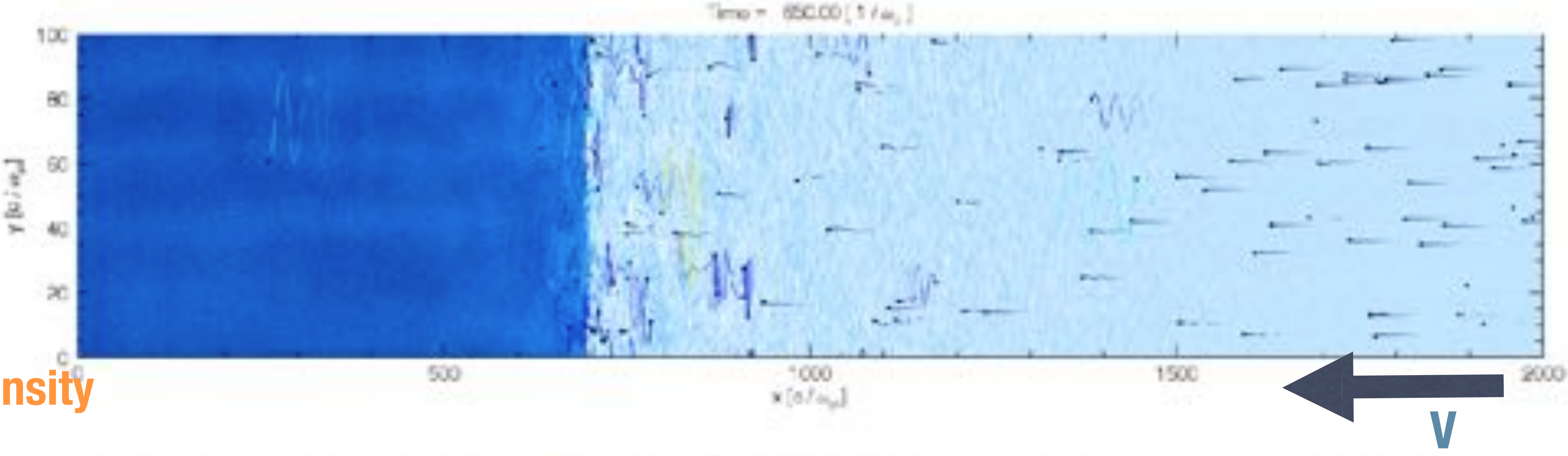
1) Proton injection physics
2) Electron injection physics and proton/electron ratio in CRs
3) Injection of heavy ions
4) Reacceleration of CRs



Proton Acceleration

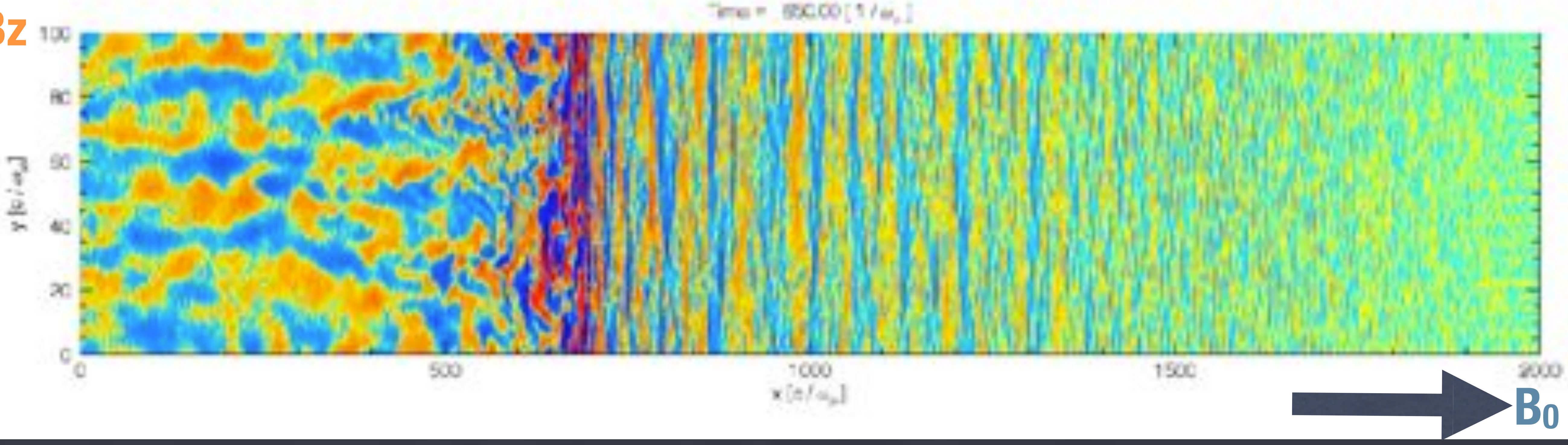
Proton acceleration

M_A=5, parallel shock; hybrid simulation. Quasi-parallel shocks accelerate ions and produce self-generated waves in the upstream.



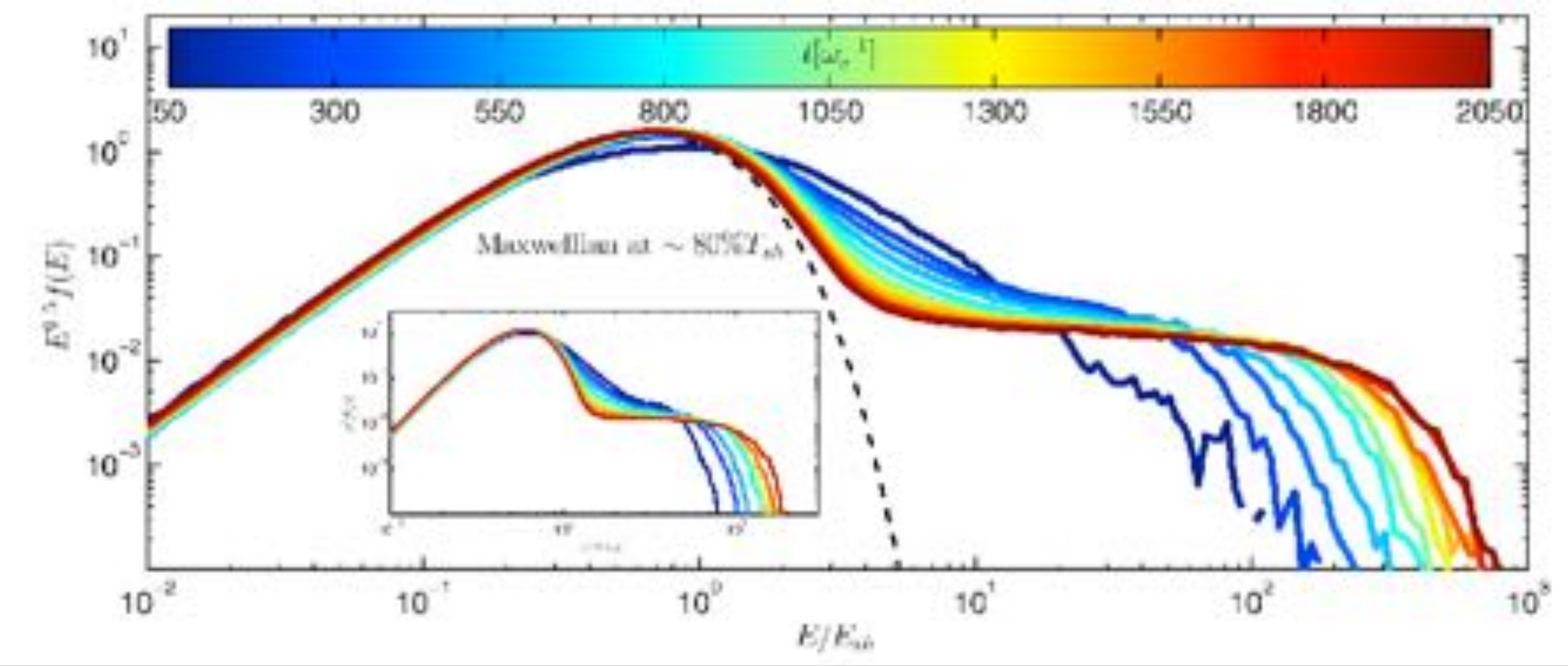
Density

Bz



Proton spectrum

Long term evolution: Diffusive Shock Acceleration spectrum recovered



First-order Fermi acceleration: $f(p) \propto p^{-4}$ $4\pi p^2 f(p) dp = f(E) dE$

$f(E) \propto E^{-2}$ (relativistic) $f(E) \propto E^{-1.5}$ (non-relativistic)

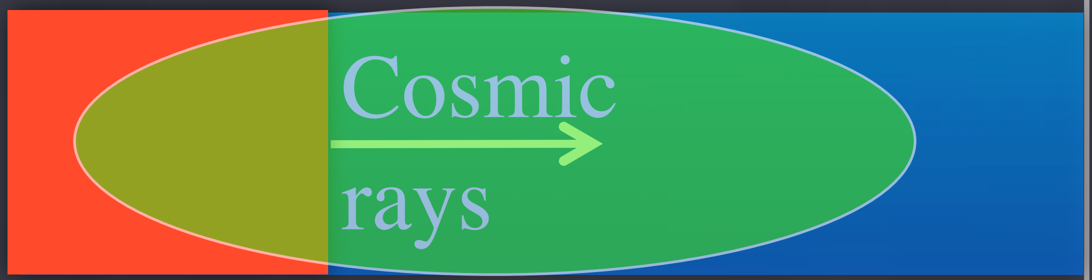
CR backreaction is affecting downstream temperature

Caprioli & Spitkovsky 2014

Field amplification

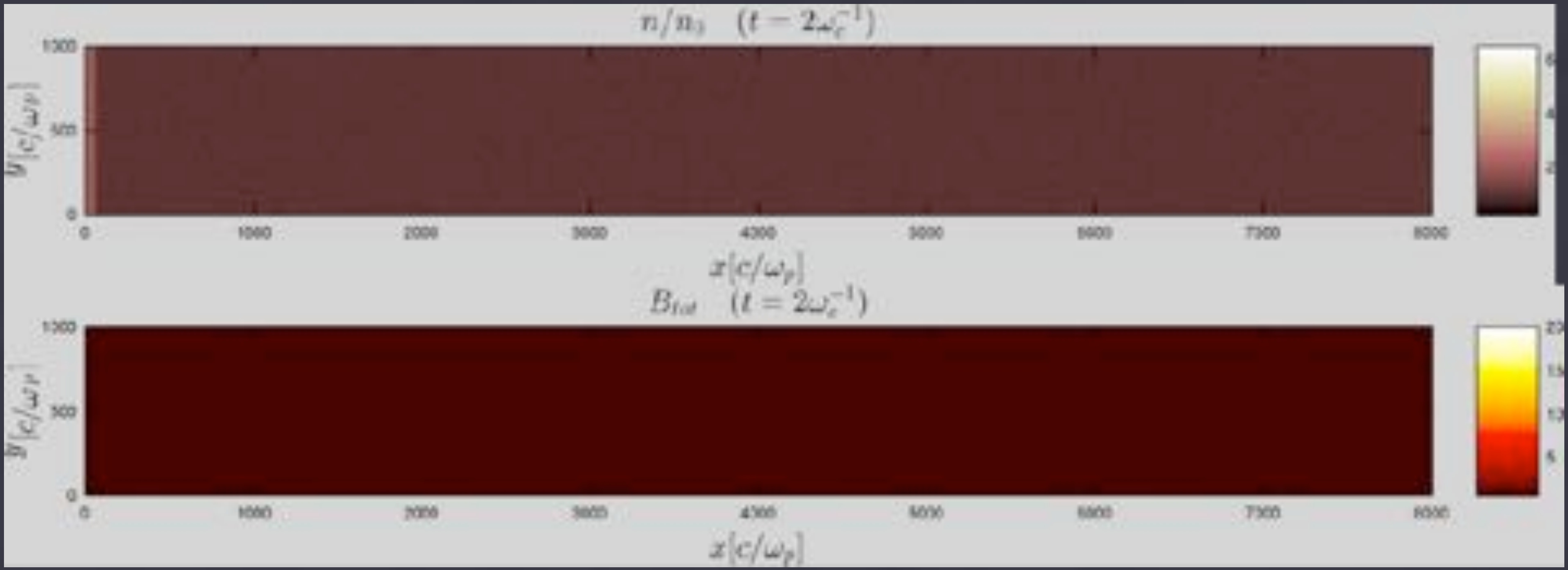
We see evidence of CR effect on upstream.

This will lead to “turbulent” shock with effectively lower Alfvenic Mach number with locally 45 degree inclined fields.

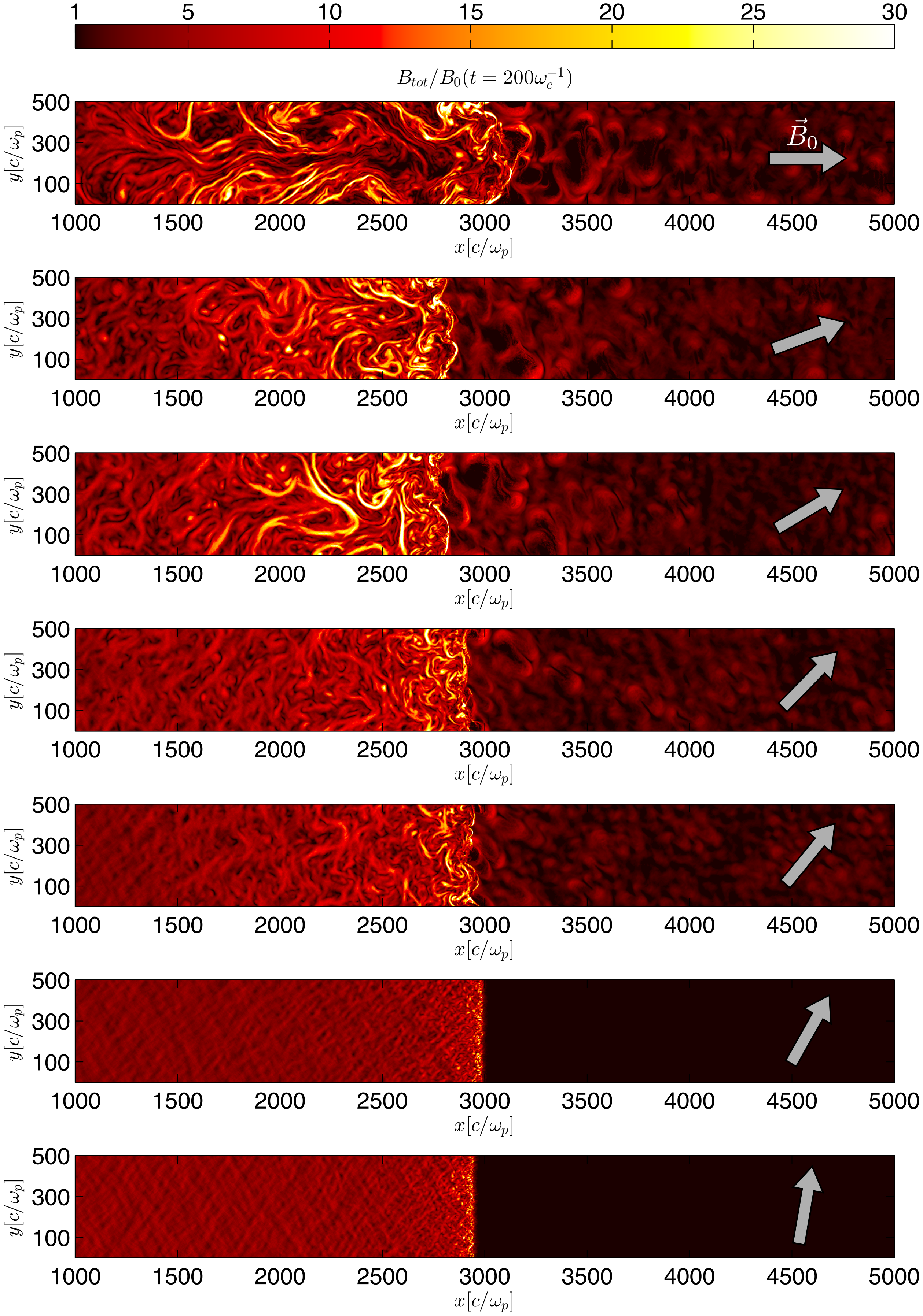


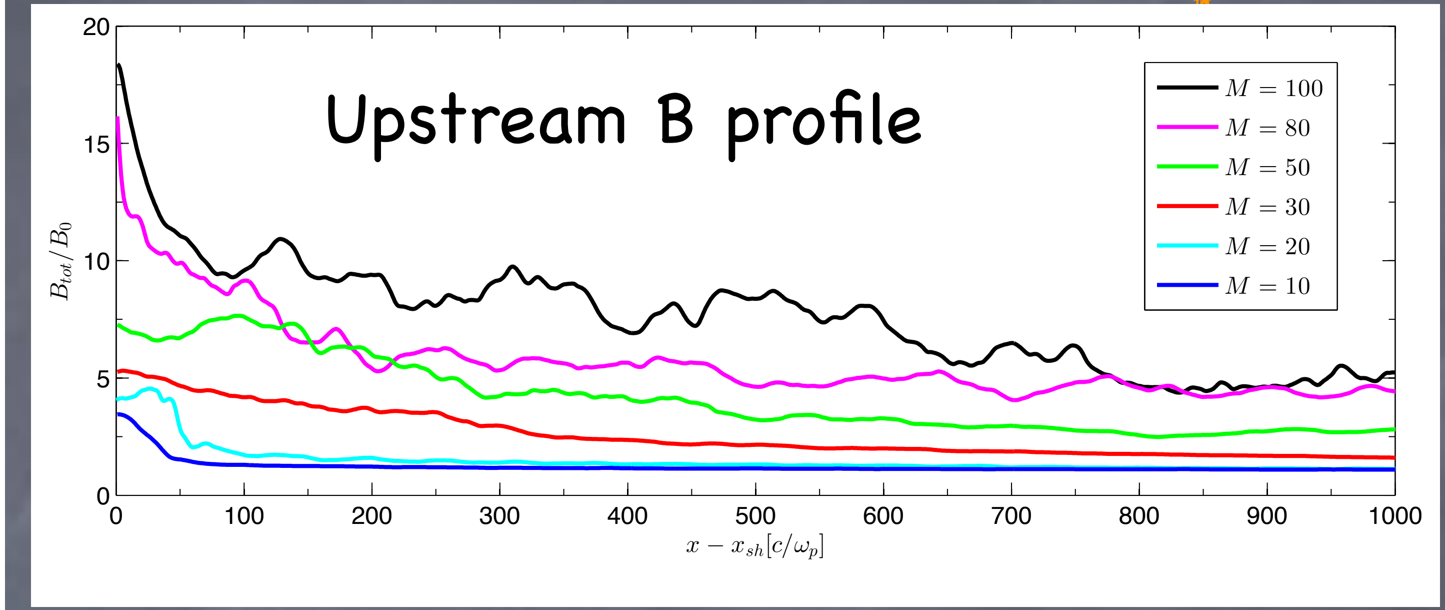
Cosmic ray current $J_{cr}=en_{cr}v_{sh}$

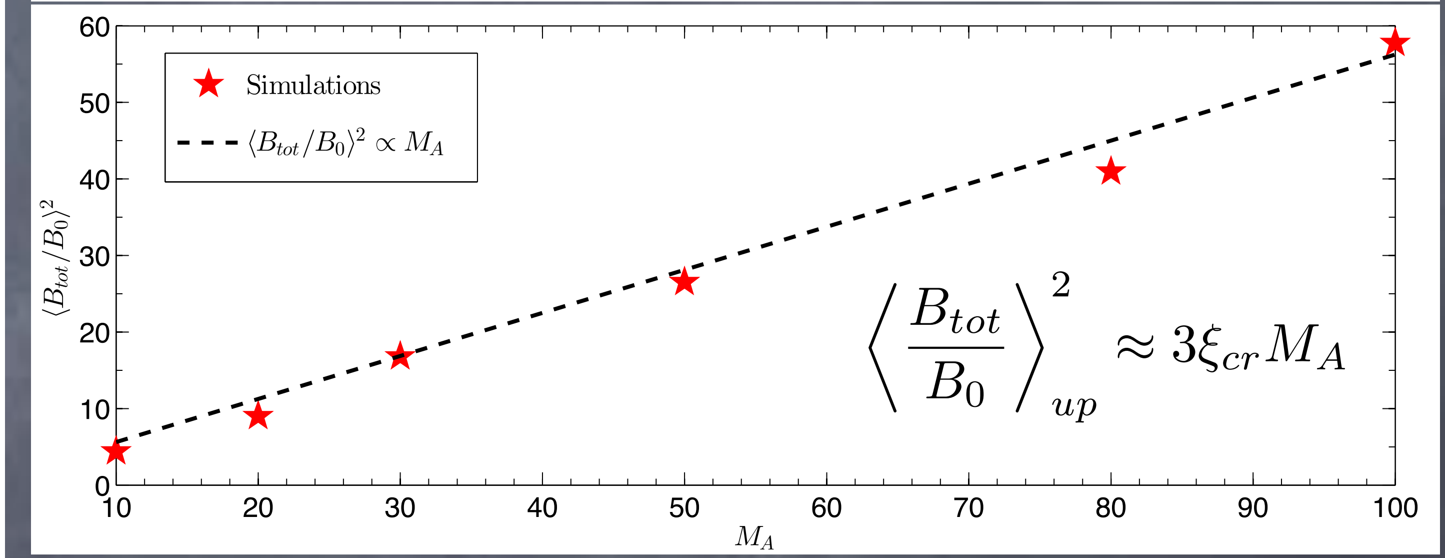
Combination of nonresonant (Bell), resonant, and firehose instabilities + CR filamentation



Dependence of field amplif. on inclination and M



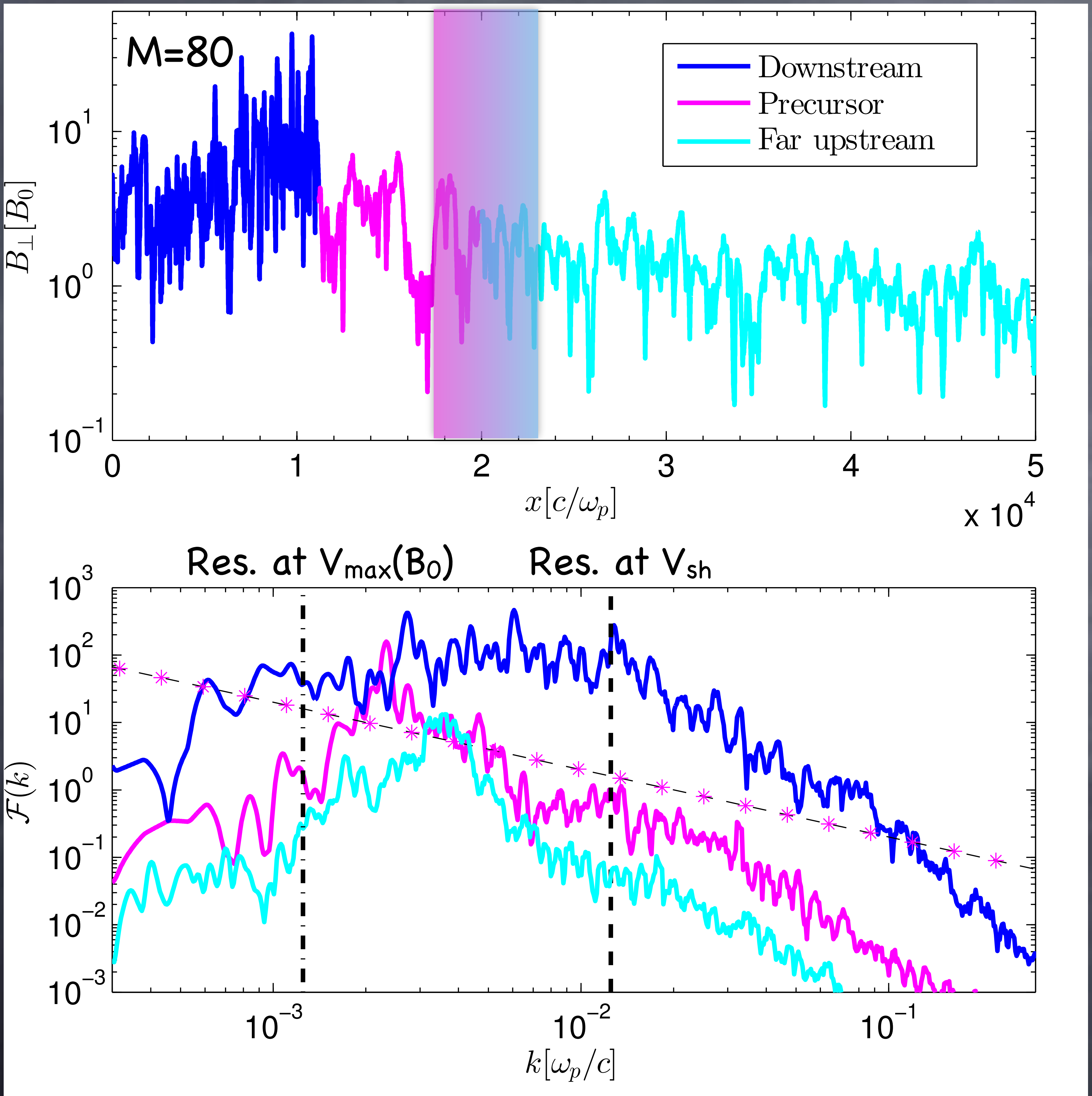




- In agreement with the prediction of resonant streaming instability

More B-field amplification for stronger shocks!

Magnetic field spectrum, high M_A



- Bell modes (short-wavelength, right-handed) grow faster than resonant
- Far upstream: escaping CRs at $\sim p_{max}$ (Bell)
- For large $b=\delta B/B_0$ $k_{max}(b) \sim k_{max,0}/b^2$
- There exist a b^* such that $k_{max}(b^*) r_L(p_{esc}) \sim 1$

Free escape boundary

- Precursor: diffusion + resonant

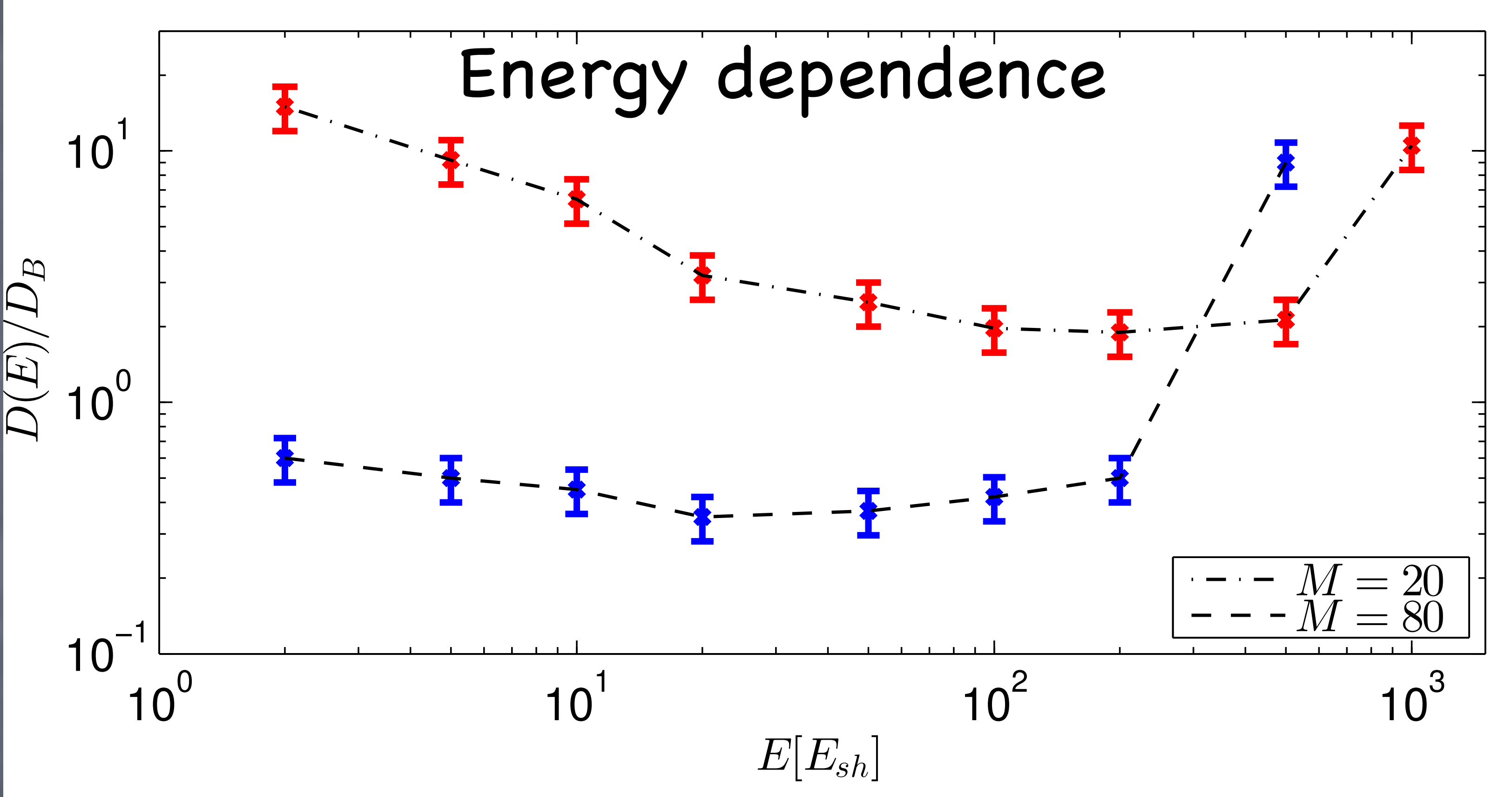
Caprioli & AS, 2014b

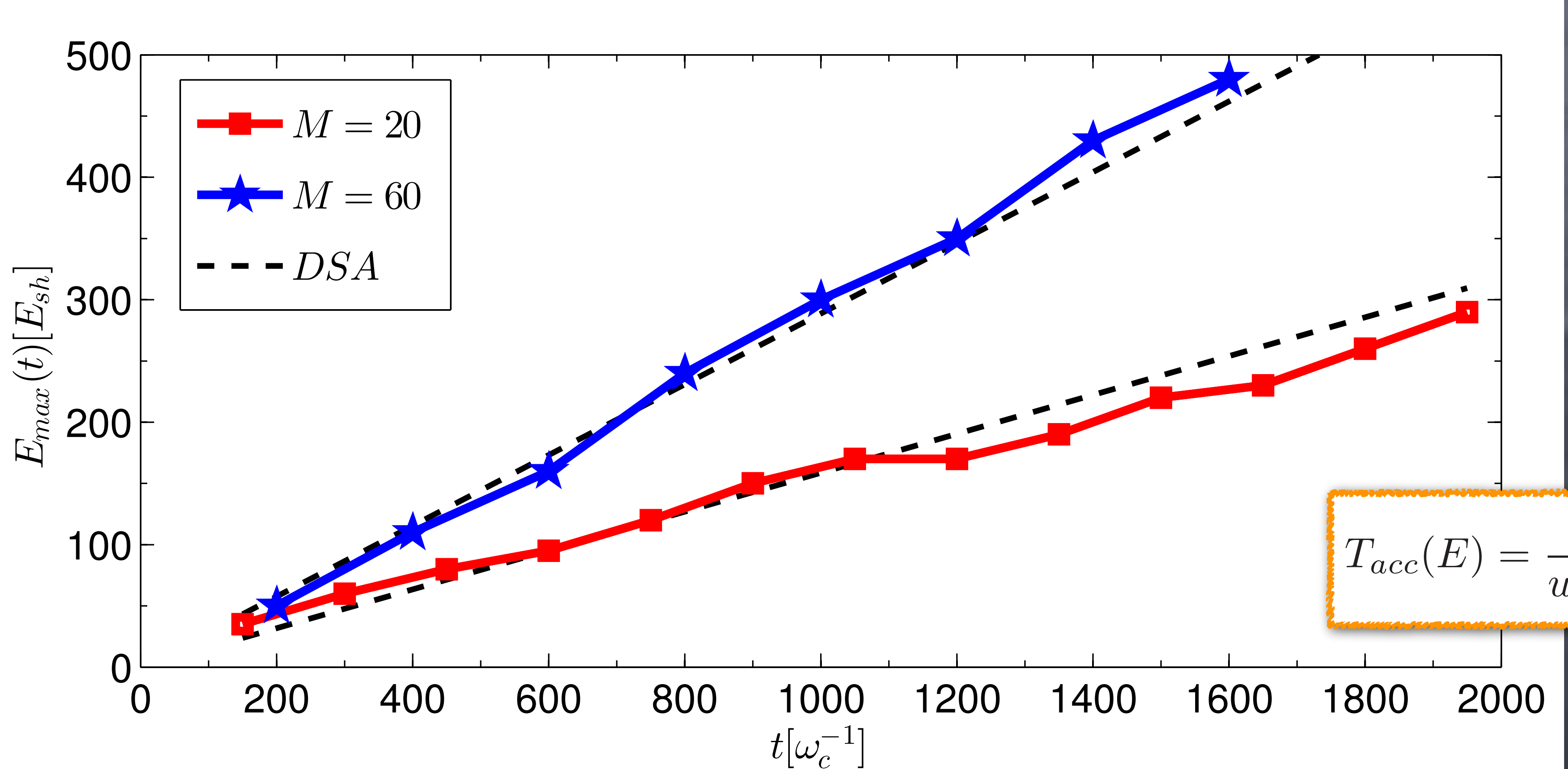
Diffusion coefficient

- Directly measurable in simulations:

$$D(E) \equiv \lim_{t\to\infty} D(E,t) = \lim_{t\to\infty}\sum_{n=1}^{N}\frac{|x_n(t)-x_n(0)|^2}{2tN}.$$

Bohm diffusion in the amplified B



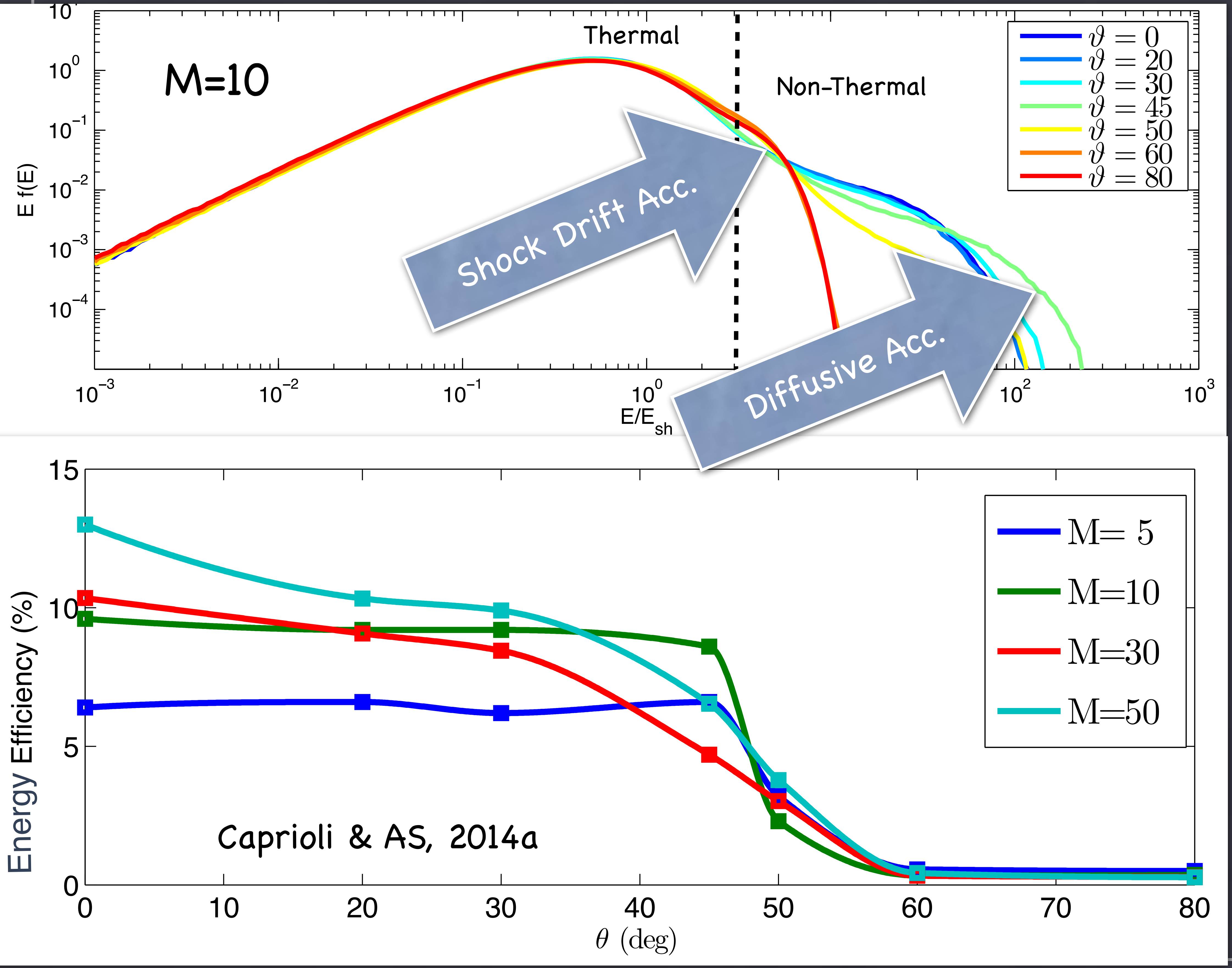


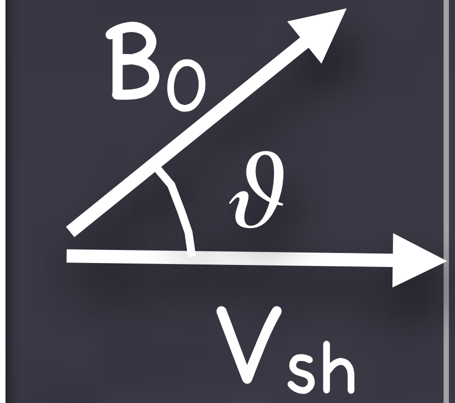
- Evolution of $E_{max}(t)$ according to DSA (e.g., Drury 1983)

$$T_{acc}(E) = \frac{3}{u_1-u_2}\left[\frac{D_1(E)}{u_1}+\frac{D_2(E)}{u_2}\right] \simeq \frac{3r^3}{r^2-1}\frac{D(E)}{v_{sh}^2}.$$

Caprioli & AS, 2014c

Acceleration in parallel vs oblique shocks

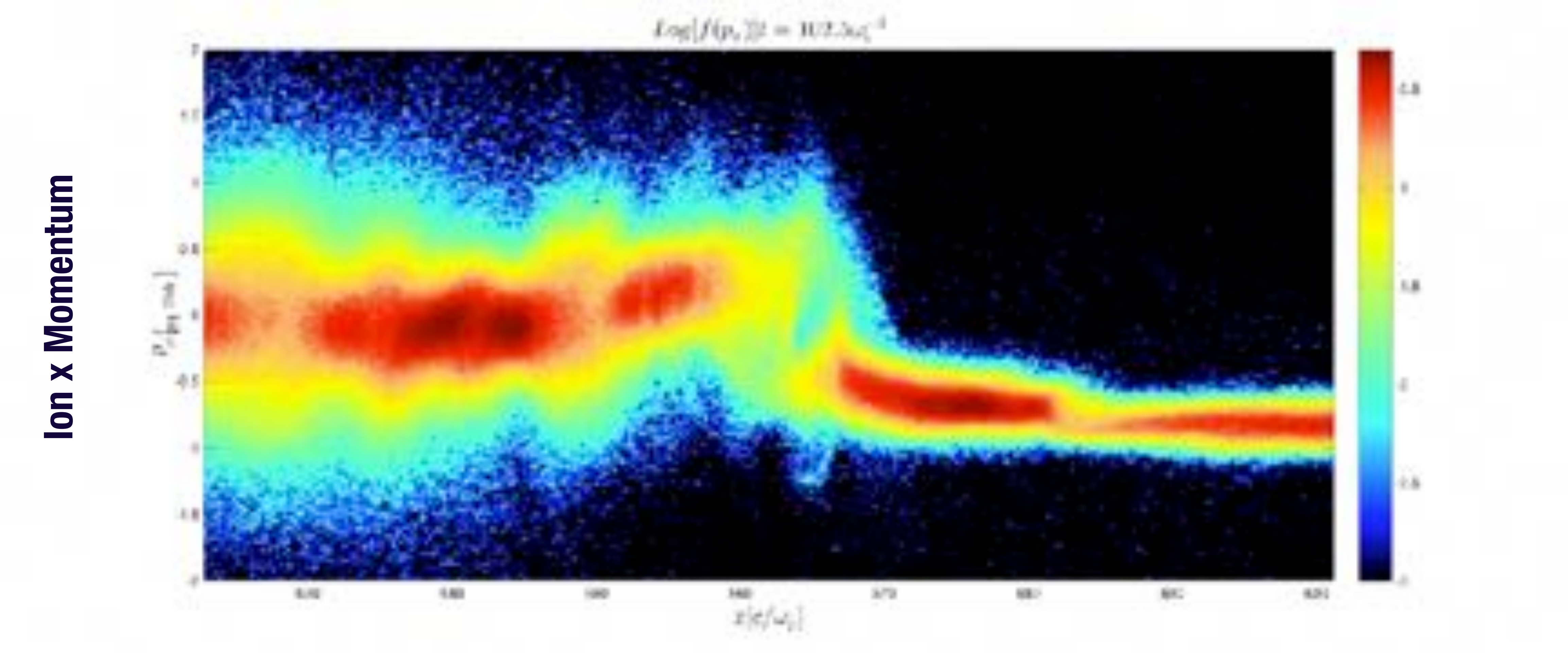




About 1% accelerated protons by number, what is causing that?

Shock structure & injection

Quasiparallel shocks look like intermittent quasiperp shocks

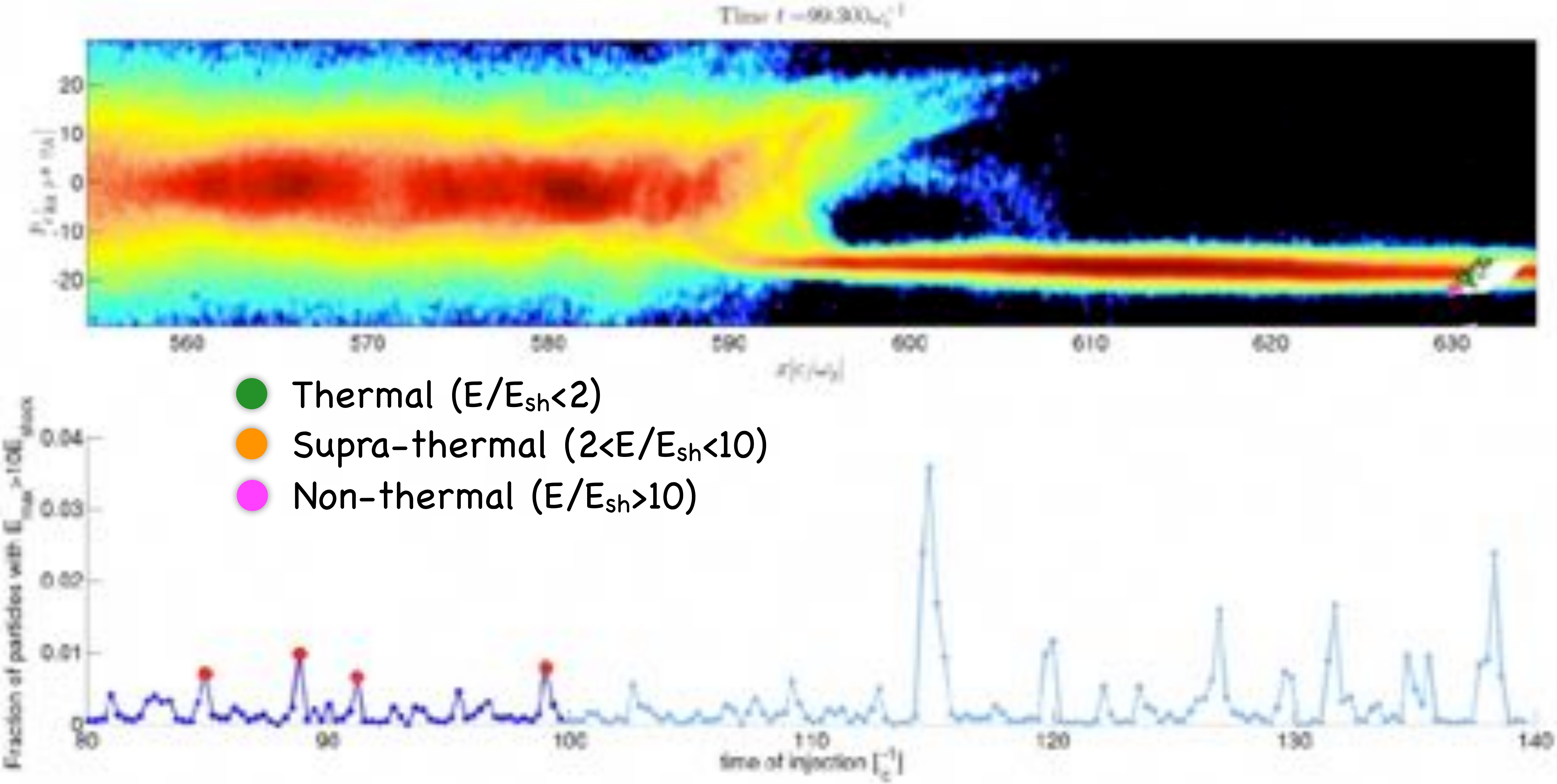


Injection of ions happens on first crossing due to specular reflection from reforming magnetic and electric barrier and shock-drift acceleration.

Multiple cycles in a time-dependent shock structure result in injection into DSA; no “thermal leakage” from downstream.

Injection mechanism: importance of timing

Caprioli, Pop & AS 2015

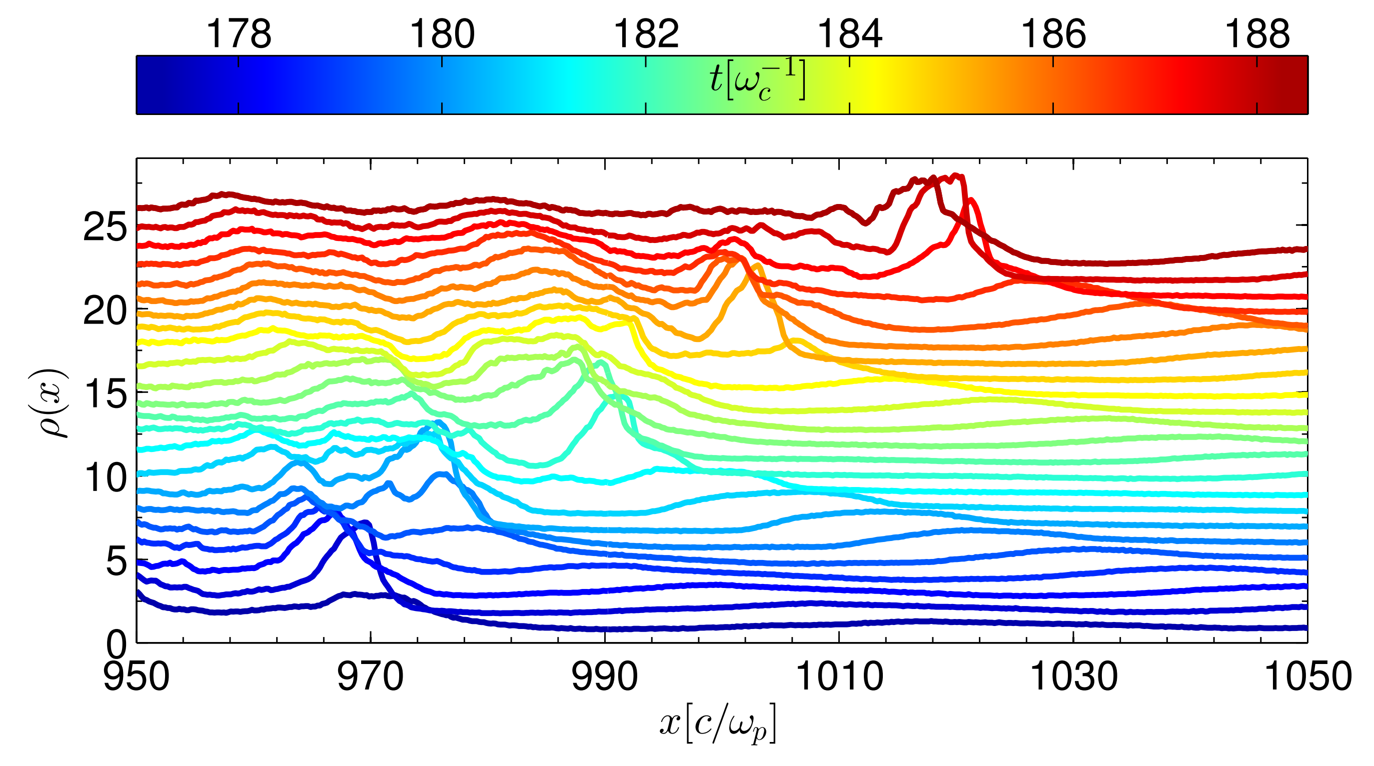


- ● Thermal ($E/E_{sh}<2$)
- ● Supra-thermal ($2<E/E_{sh}<10$)
- ● Non-thermal ($E/E_{sh}>10$)

Ion injection: theory

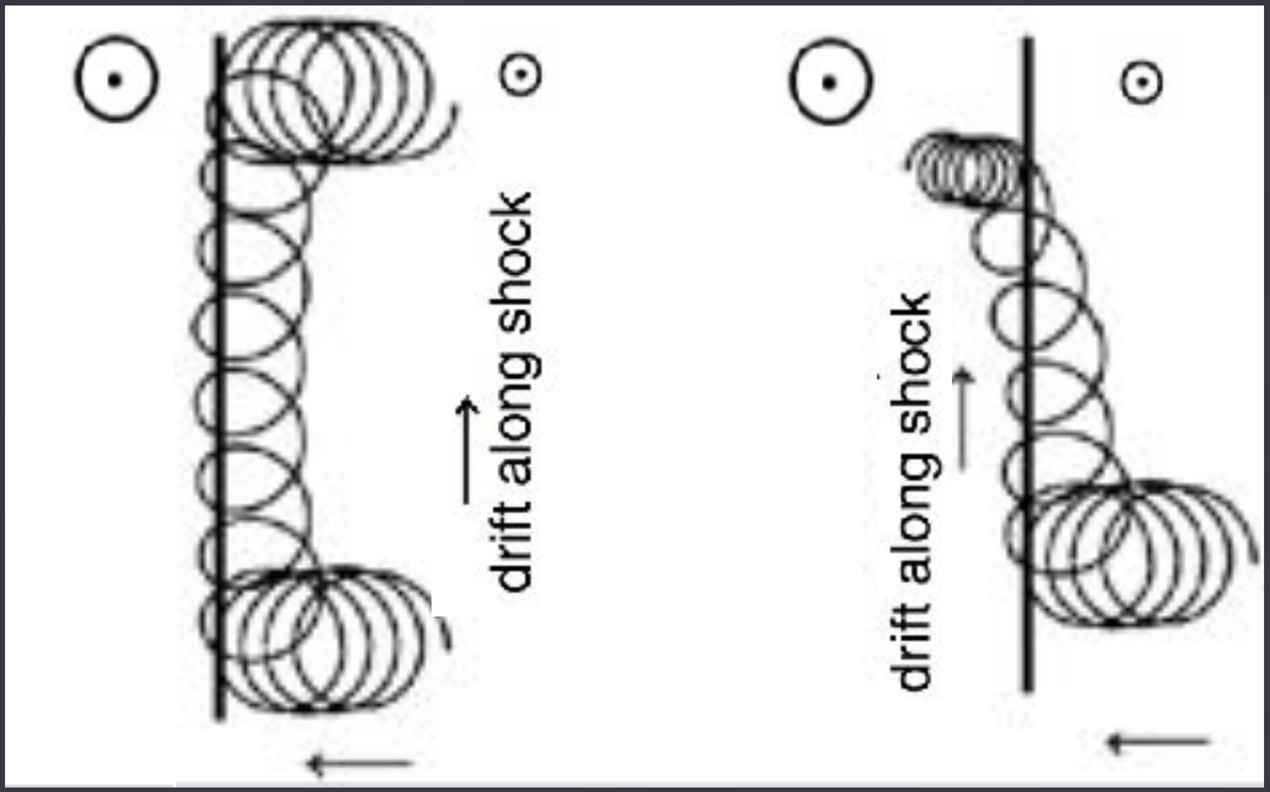
Caprioli, Pop & AS 2015

- Reflection off the shock potential barrier (stationary in the downstream frame)
- For reflection into upstream, particle needs certain minimal energy for given shock inclination;
- Particles first gain energy via shock-drift acceleration (SDA)
- Several cycles are required for higher shock obliquities
- Each cycle is "leaky", not everyone comes back for more
- Higher obliquities less likely to get injected



Shock-drift acceleration:

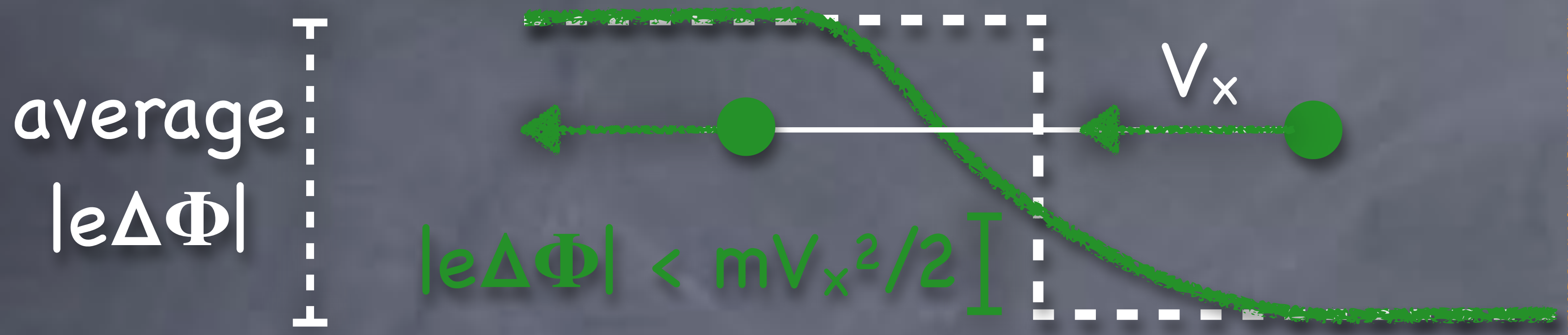
downstream upstream Larger B Smaller B



Path of incoming particle

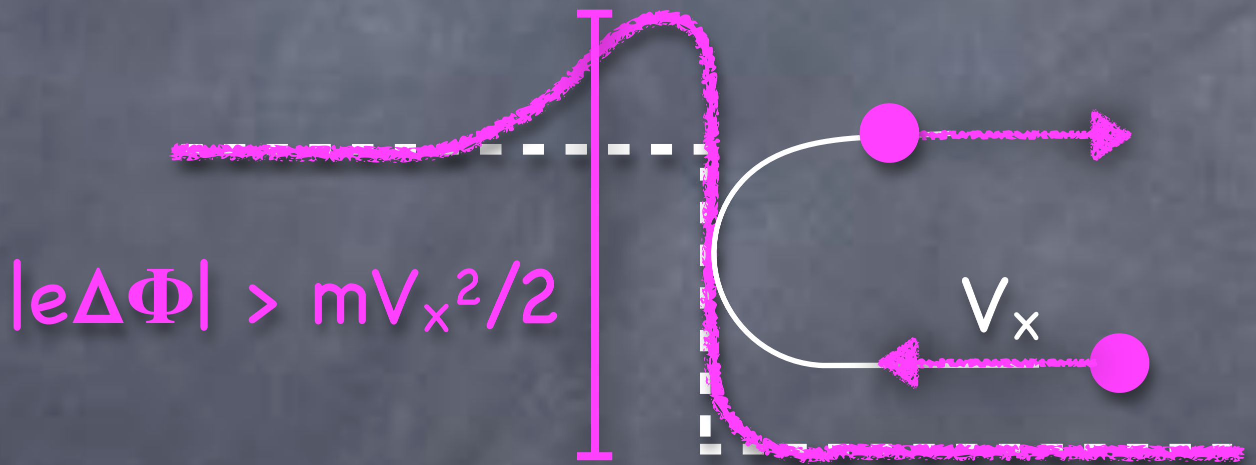
Encounter with the shock barrier

- Low barrier (shock reforming)



Particles are advected downstream, and thermalized

- High barrier (overshoot)



Particles are reflected upstream, and energized via Shock Drift Acc.

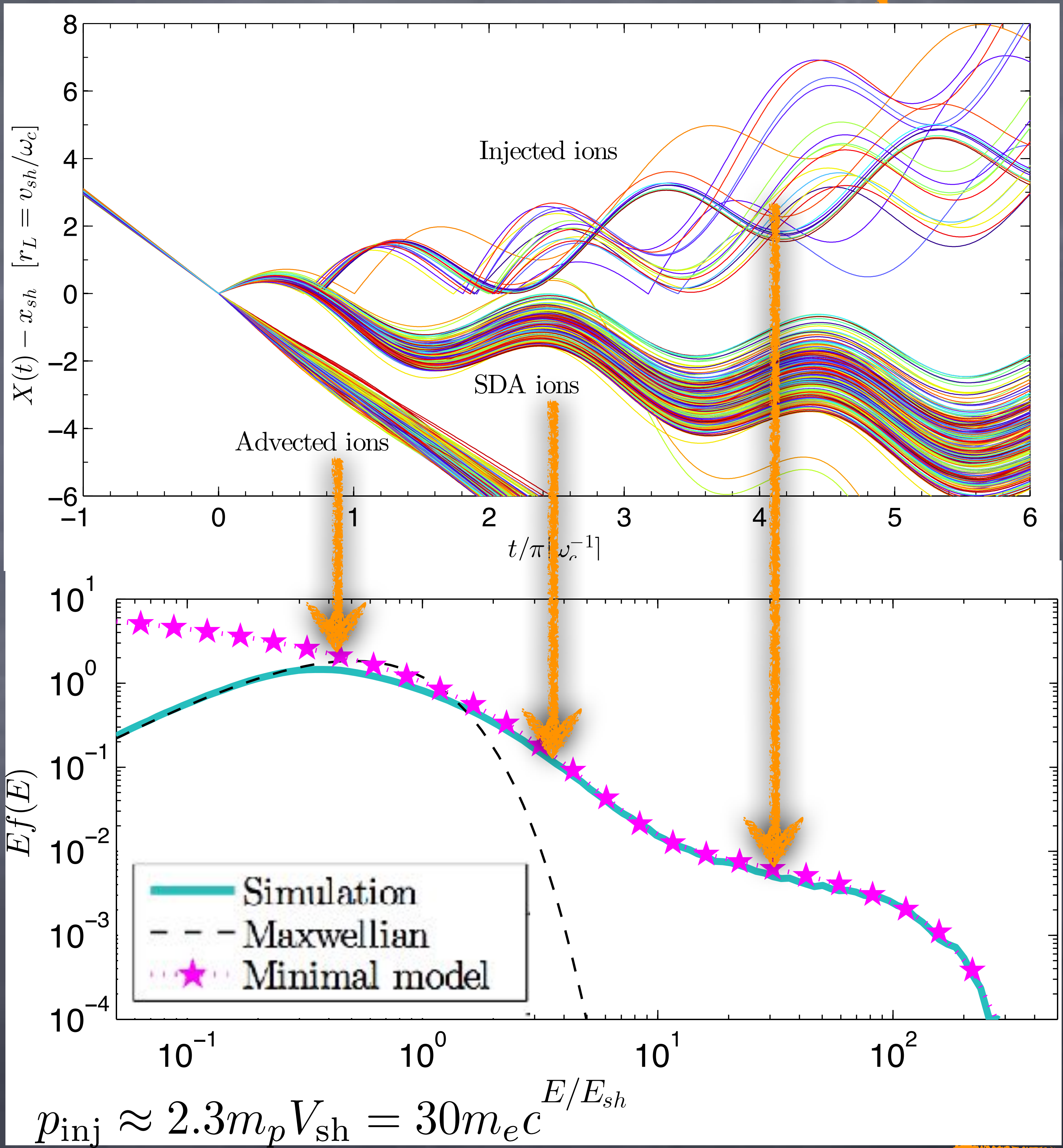
- To overrun the shock, proton need a minimum $E_{inj,}$ increasing with ϑ
- Particle fate determined by barrier duty cycle (~25%) and shock inclination
- After N SDA cycles, only a fraction $\eta \sim 0.25^N$ has not been advected
  - For $\vartheta=45^\circ$, $E_{inj} \sim 10E_0$, which requires N~3 -> $\eta \sim 1\%$

Minimal Model for Ion Injection

- Time-varying potential barrier
  - **High state** (duty cycle 25%)
    -> Reflection
    -> Shock Drift Acceleration
  - **Low-state** -> Thermalization
- Spectrum à la Bell (1978)

$$f(E) \propto E^{-1-\gamma}; \quad \gamma \equiv -\frac{\ln(1-\mathcal{P})}{\ln(1+\mathcal{E})}$$

- P=probability of being advected
- ε=fractional energy gain/cycle



$p_{\rm inj} \approx 2.3 m_p V_{\rm sh} = 30 m_e c$

Caprioli, Pop & Spitkovsky, 2015

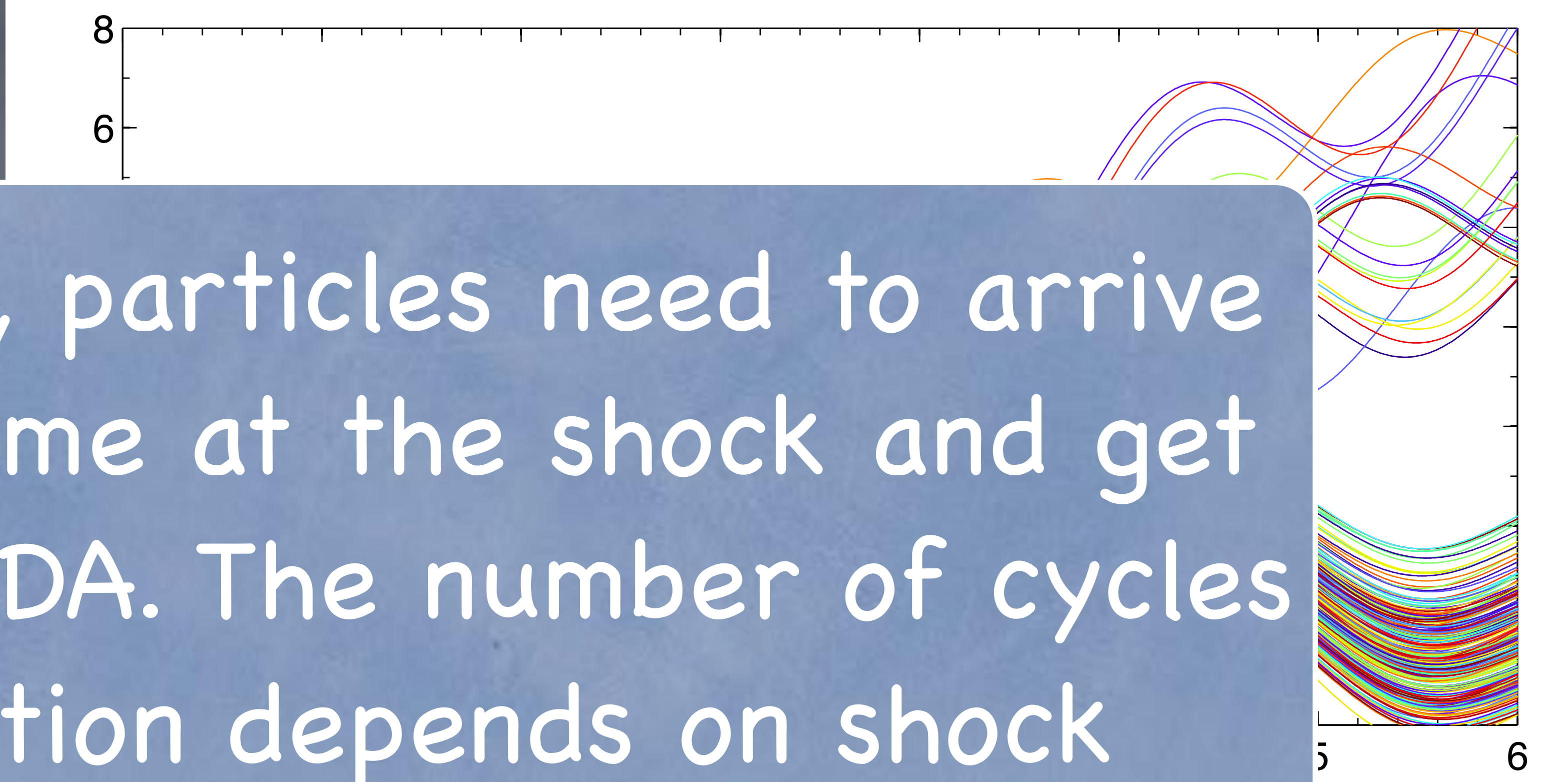
Minimal Model for Ion Injection

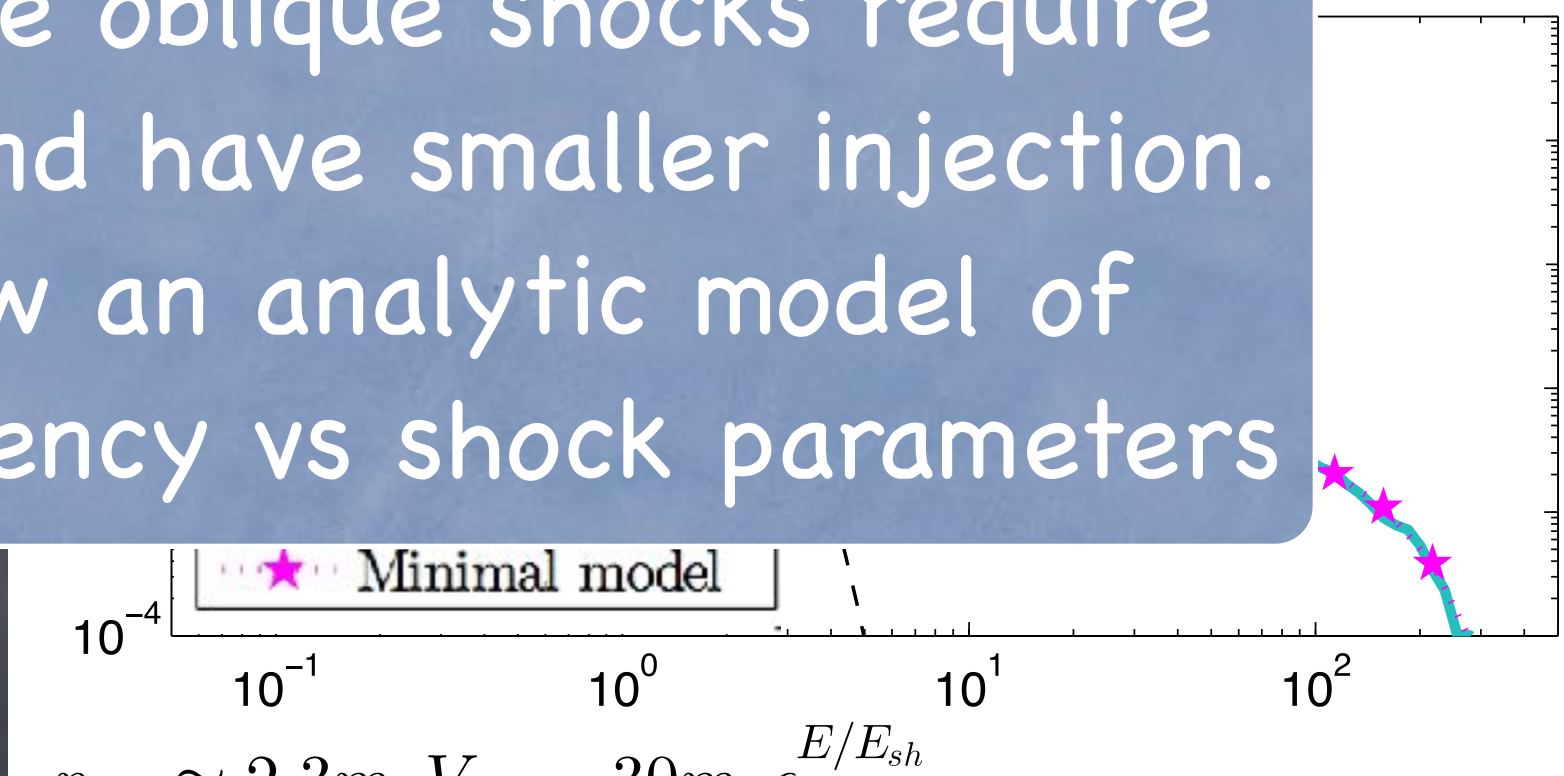
- Time-varying potential barrier
  - **High**
    -> R
    -> S
  - **Low-**
- Spectru

$f(E) \propto$

- P=proba
- ε=fractional energy gain/cycle

To be injected, particles need to arrive at the right time at the shock and get energized by SDA. The number of cycles of energization depends on shock obliquity. More oblique shocks require more cycles, and have smaller injection. There is now an analytic model of injection efficiency vs shock parameters





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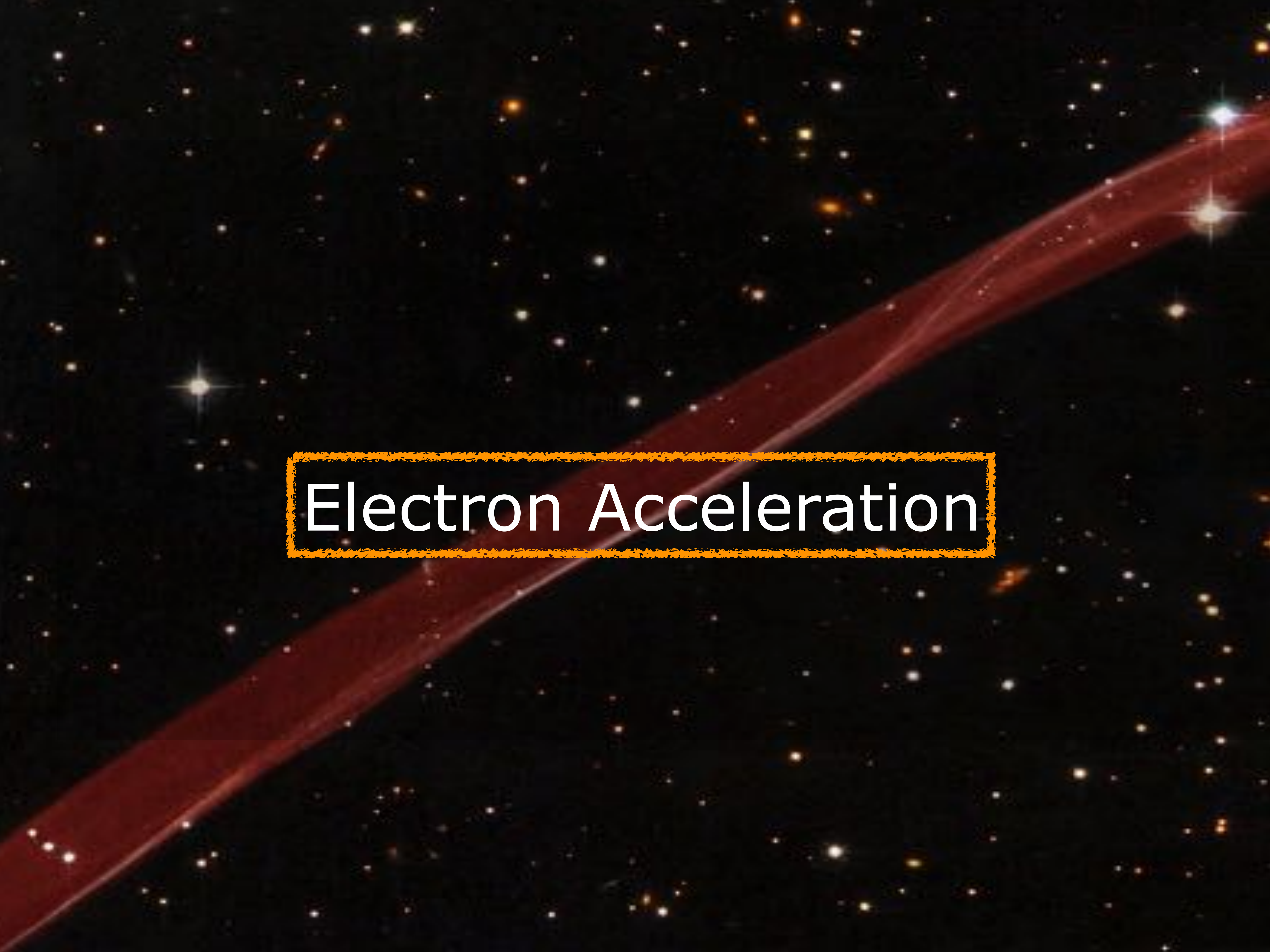
Caprioli, Pop & Spitkovsky, 2015

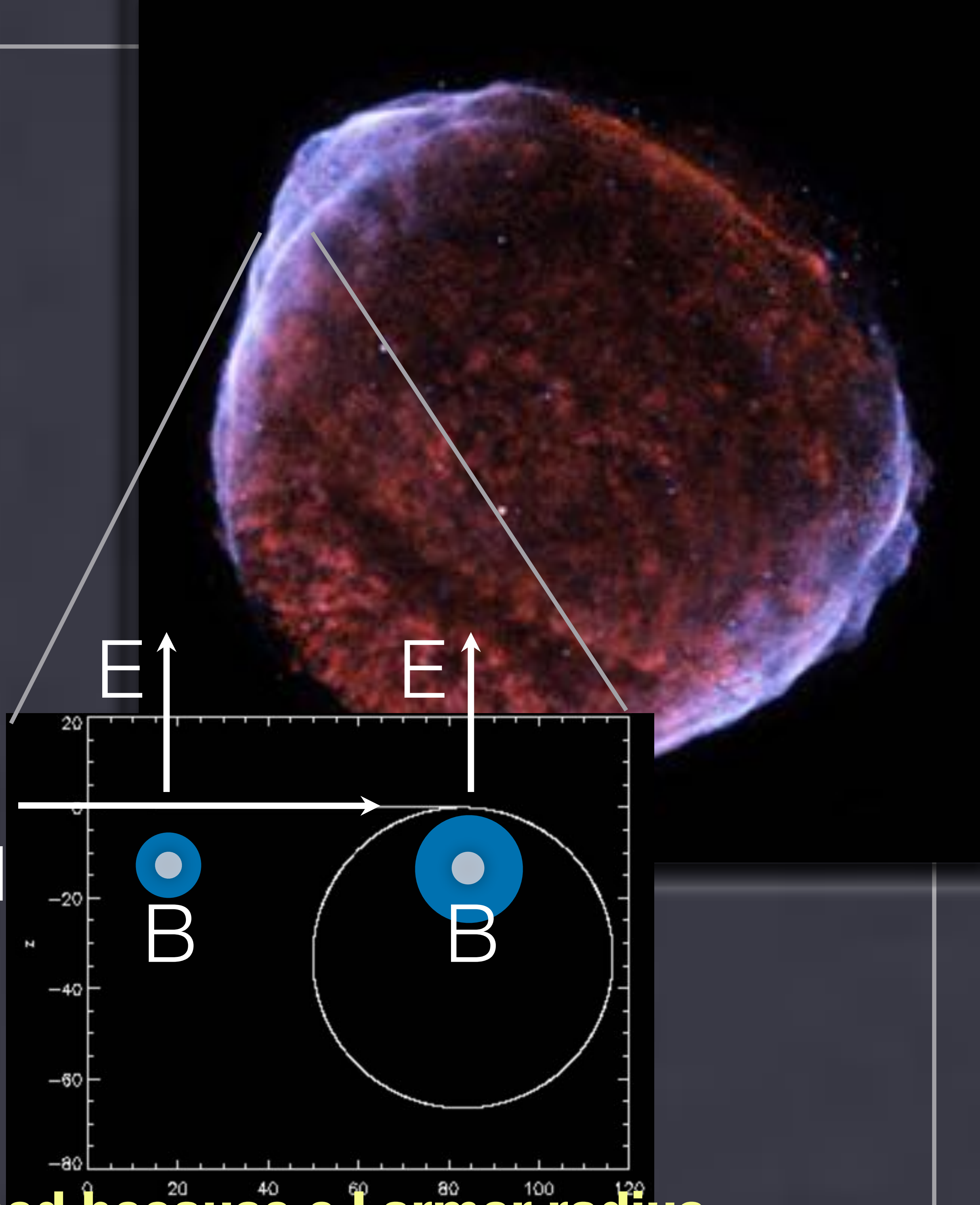


Electron Acceleration

WHAT ACCELERATES ELECTRONS?

Electrons are notorious for being difficult to inject because of the disparity in the Larmor scales with ions.

Shock is driven on ion scales, electrons need to be pre-accelerated to be injected. But how?



Typically electron acceleration is suppressed because e Larmor radius is << ion Larmor radius. Need pre-acceleration of electrons.

This means trapping at the shock, and turbulence upstream. Is it self-generated?

More dimensionless numbers

$$\sigma \equiv \frac{B^2/4\pi}{(\gamma-1)nmc^2} = \frac{2}{M_A^2} = \left(\frac{\omega_c}{\omega_p}\right)^2 \left(\frac{c}{v}\right)^2 = \left[\frac{c/\omega_p}{R_L}\right]^2$$

$$M_A = \frac{v_{sh}}{v_A}$$

$$\beta = \frac{P_{th}}{P_B} = \left(\frac{M_A}{M_{si}}\right)^2 = \left(\frac{v_{th,i}}{v_A}\right)^2$$

$$M_{si} = \frac{v_{sh}}{v_{th,i}}$$

$$M_{se} = M_{si}\left(\frac{m_e}{m_i}\right)^{1/2} = \frac{M_A}{\sqrt{\beta}}\left(\frac{m_e}{m_i}\right)^{1/2}$$

$$M_{se} = \frac{v_{sh}}{v_{th,e}}$$

$$v_A = \frac{B}{\sqrt{4\pi n m_i}}$$

The ability of electrons to be reflected at the shock depends on $v_{th,e}/v_{sh} = 1/M_{se}$. So, lower sonic Mach number shocks are better! Some pre-heating will help.

What pre-heats? Depends on parameters.

Electron acceleration:

All shocks are quasiperp close-up. So, first reflection will be guided by mirroring at the quasi-perp shock. Reflection may lead to couple of cycles of shock-drift, and then transmission into downstream, or streaming back upstream.

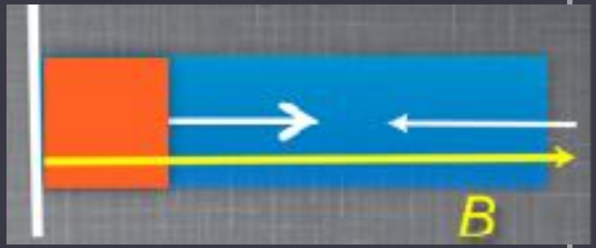
If electrons can be confined at the shock either by ion-produced or self-produced turbulence, electrons can get extra acceleration, and even enter DSA. Recent progress:

Kato 2014; Park, Caproli, AS 2015 -- initially quasipar shock gives electron trapping due to ion-driven waves. DSA transition observed.

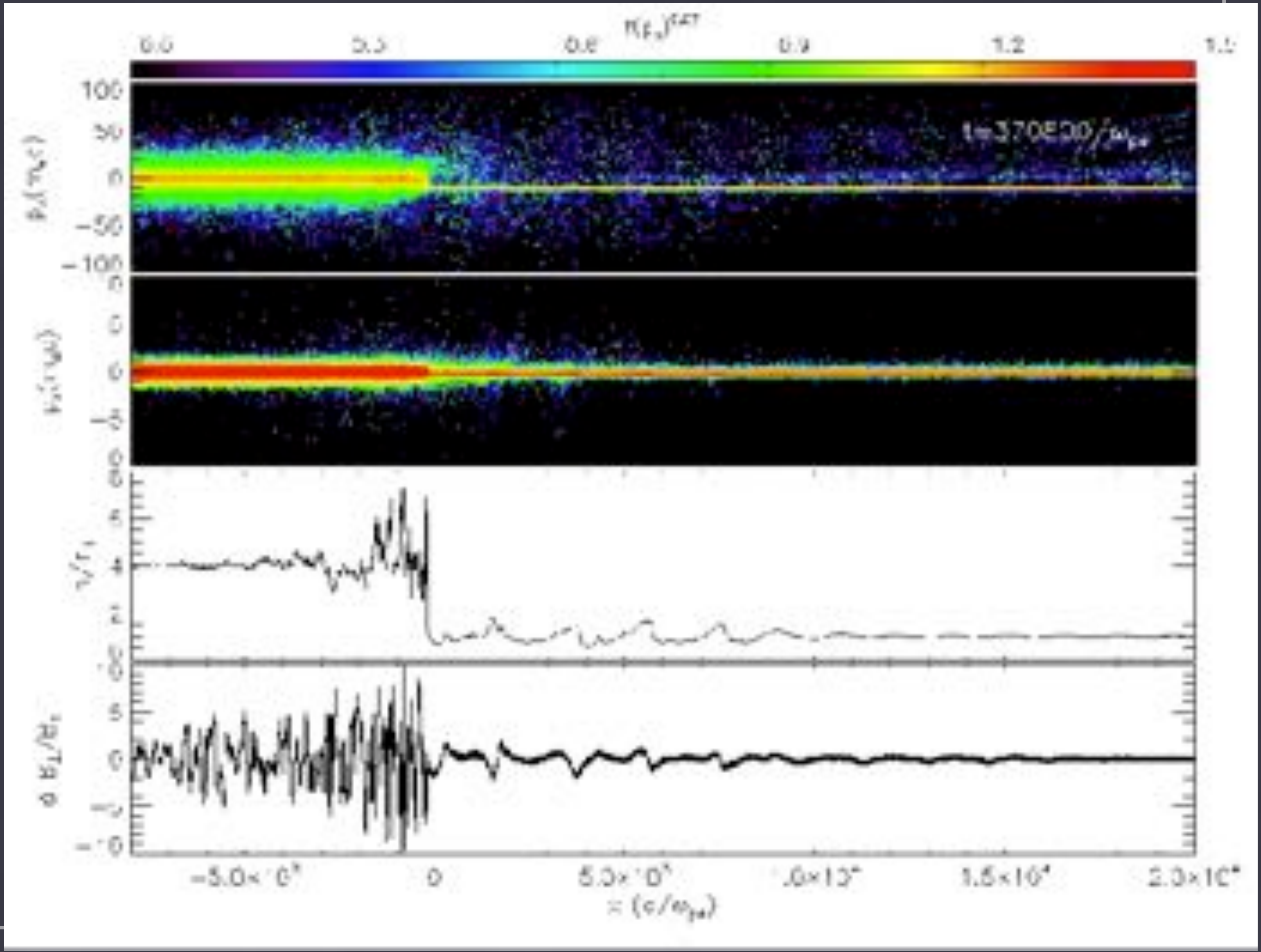
Riquelme & AS 2011: ion-driven whistler waves in quasiperp shock can trap electrons, even at high sonic Machs. Waves exist mainly for Ma < sqrt(mi/me) ~ 40; Amano & Hoshino 2011 -- different ion-driven waves.

Guo & Sironi 2014: low sonic Mach -- quasiperp; electron reflection into upstream, electrons drive their own waves! Firehose?

Electron acceleration at parallel shocks



Recent evidence of electron acceleration in quasi parallel shocks.
PIC simulation of quasiparallel shock. Very long simulation in 1D.
Alfven Mach = Sonic Mach = 20; mi/me=100-400;
Ion-driven Bell waves drive electron acceleration: correct polarization



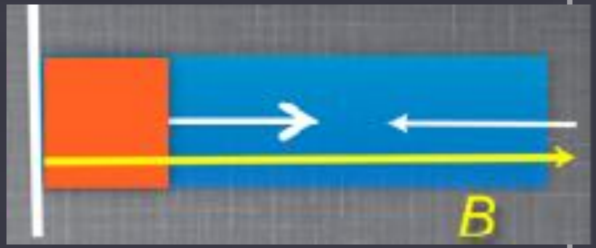
Ion phase space

Electron phase space

Density

Transverse Magnetic field

Electron acceleration at parallel shocks



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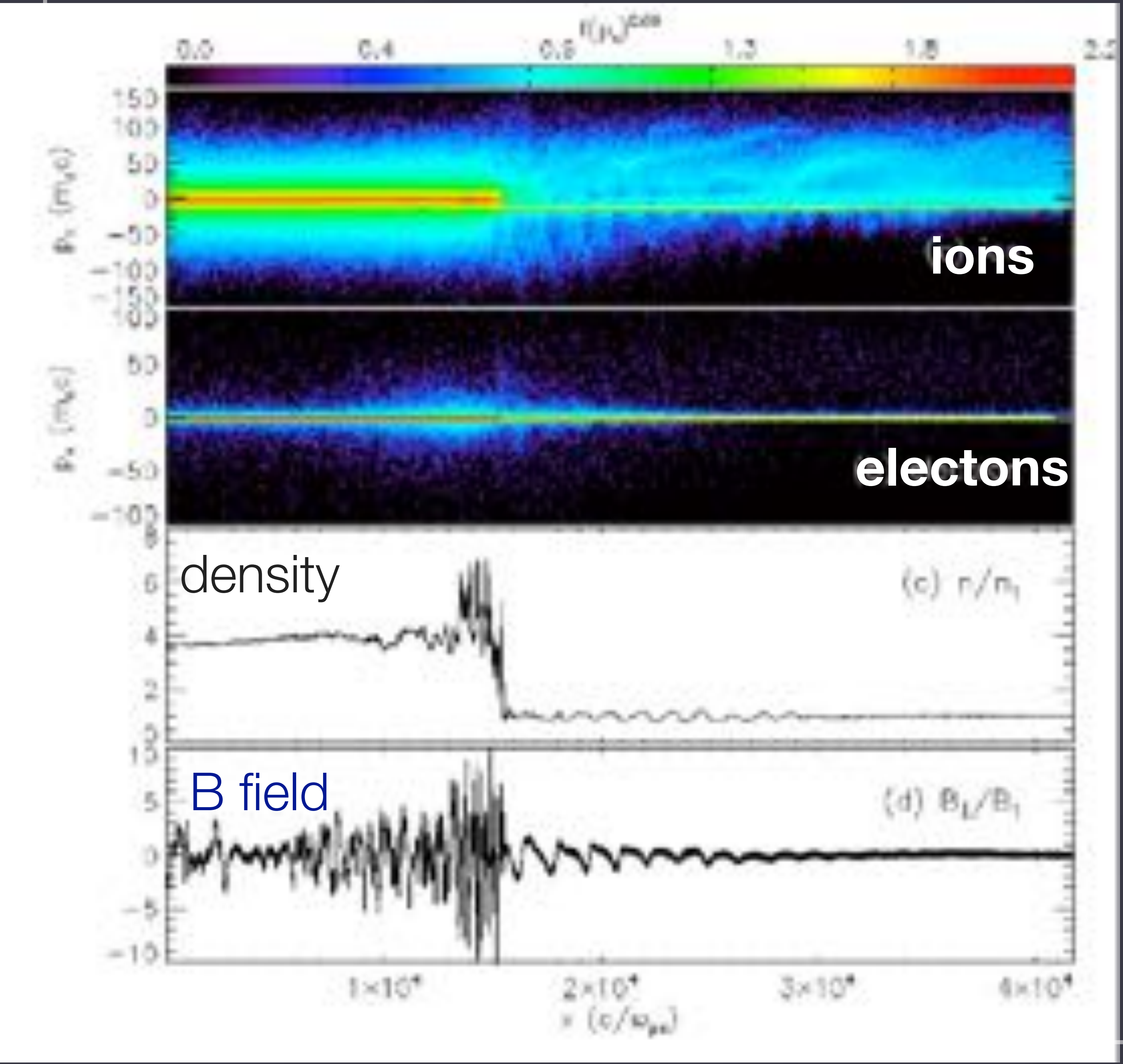
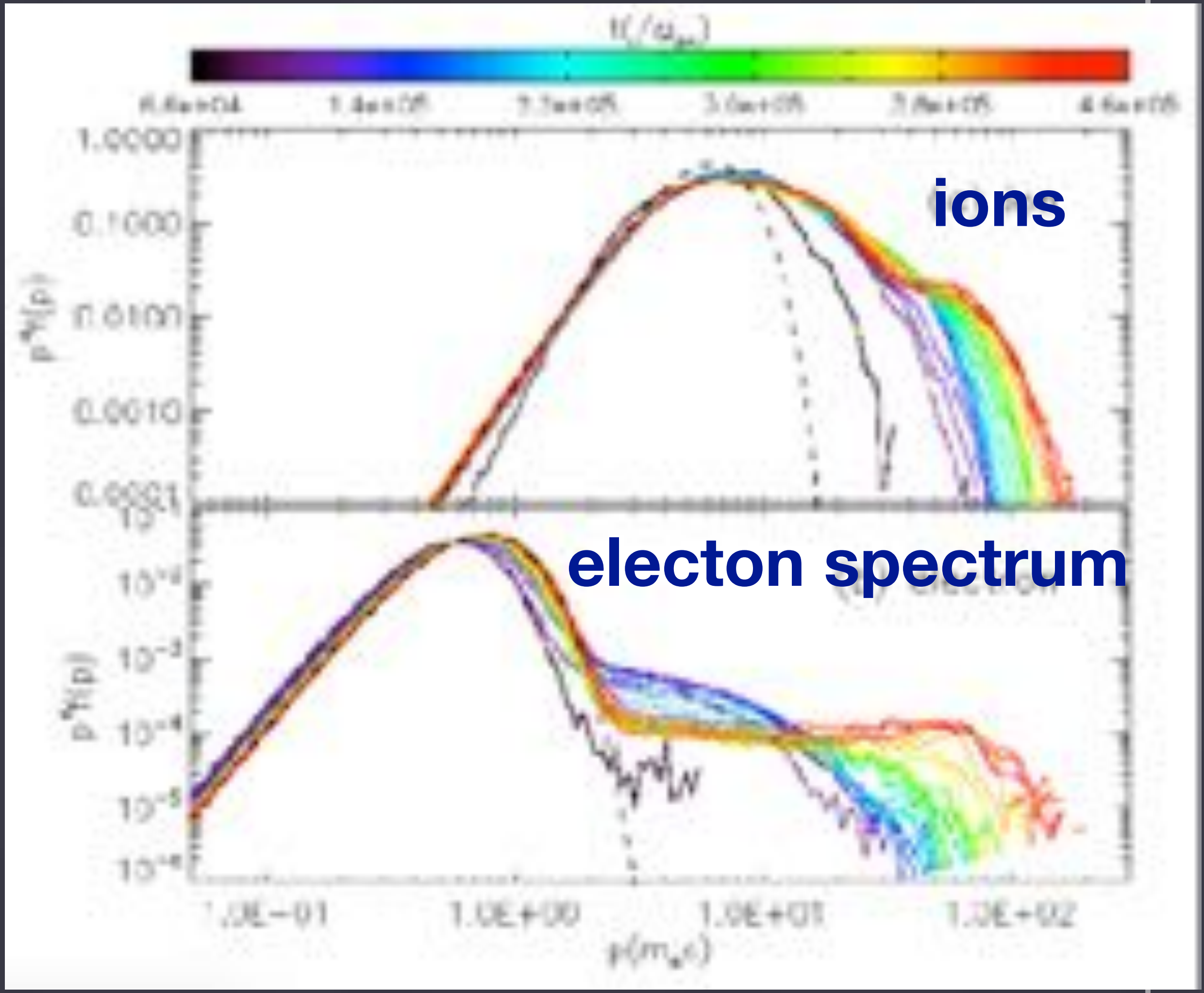


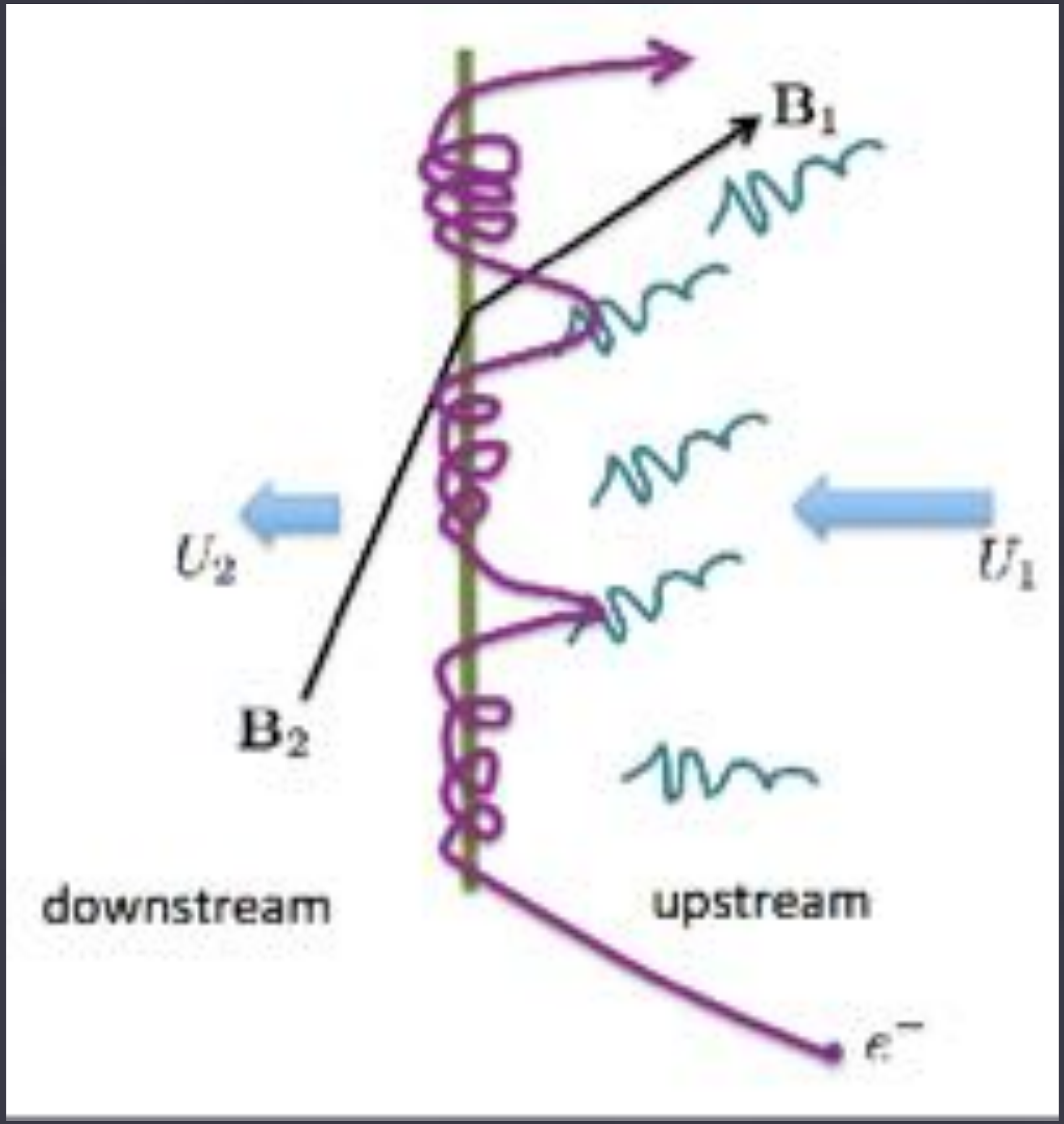


DSA spectrum recovered in _both_ electrons and ions. Electron-proton ratio obtained: $\mathbf{K_{ep}=10^{-3} -10^{-2}}$

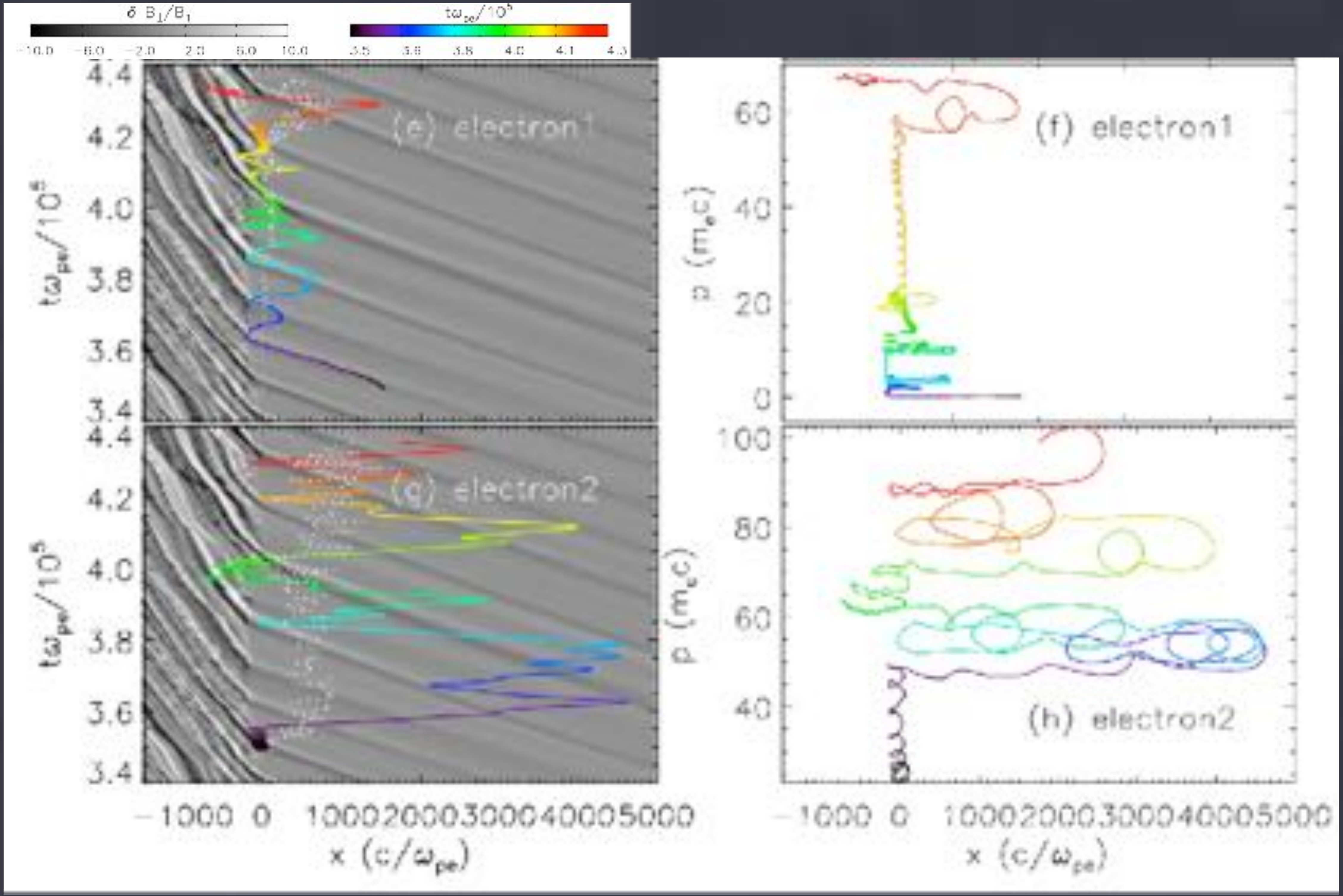
Park, Caprioli, AS (2015)

Electron acceleration at parallel shocks

Multi-cycle shock-drift acceleration, with electrons returning back due to upstream ion-generated waves.



Electron acceleration mechanism: shock drift cycles+ diffusion in upstream



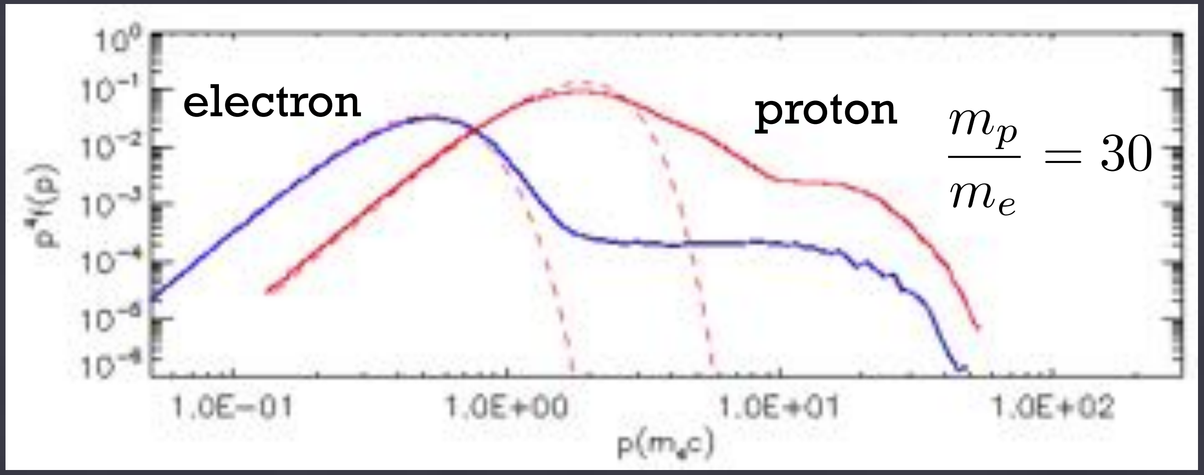
Shock-drift

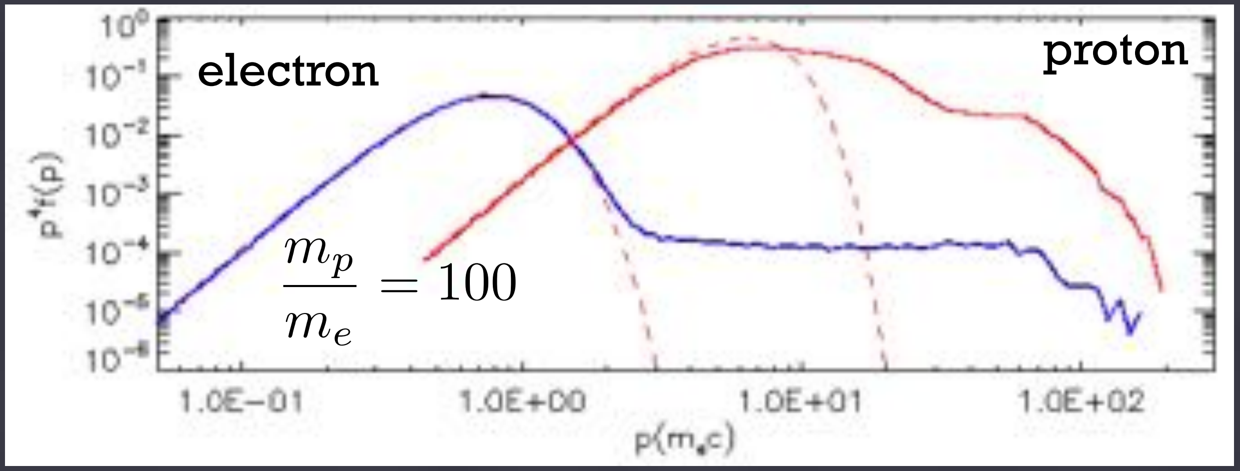
Diffusive

Electron track from PIC simulation.

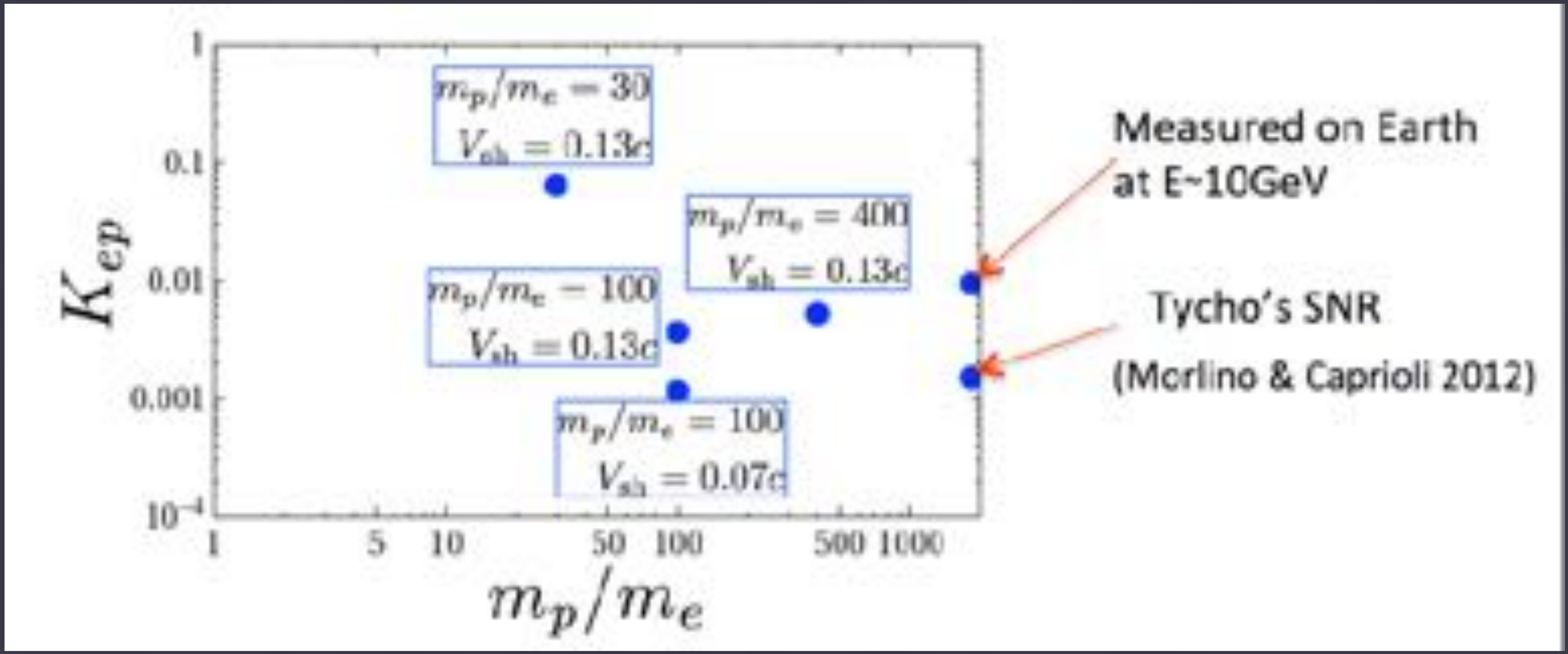
Electron-proton ratio K_{ep}:

Park, Caprioli, AS (2015)





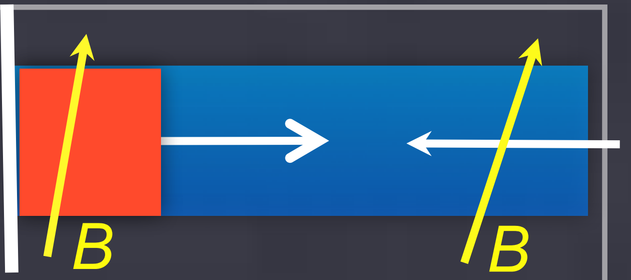
$$K_{\rm ep} \equiv \frac{f_e(p)}{f_p(p)} = \text{const for } p > p_{\rm inj} \quad K_{\rm ep} \approx 3.8 \times 10^{-3} \text{ for } \frac{m_p}{m_e} = 100$$



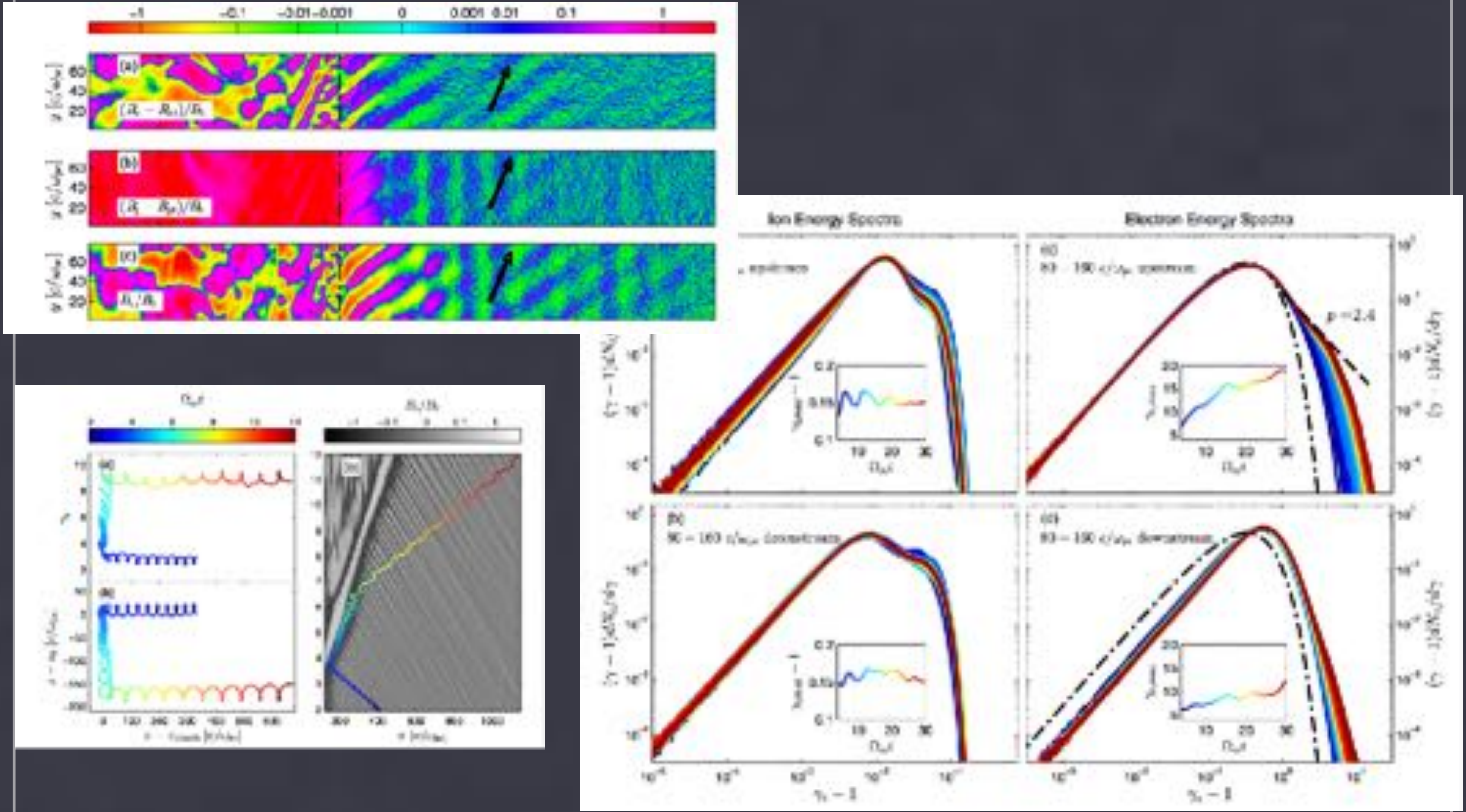
Measured on Earth at E~10GeV

Tycho's SNR (Morlino & Caprioli 2012)

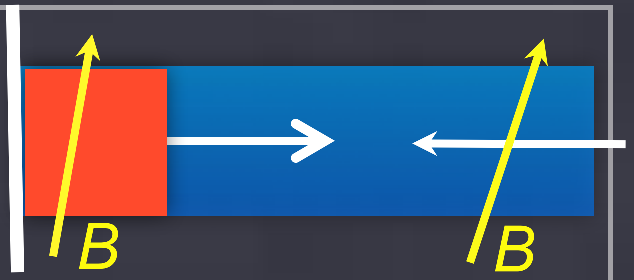
Electron acceleration at ⊥-shocks



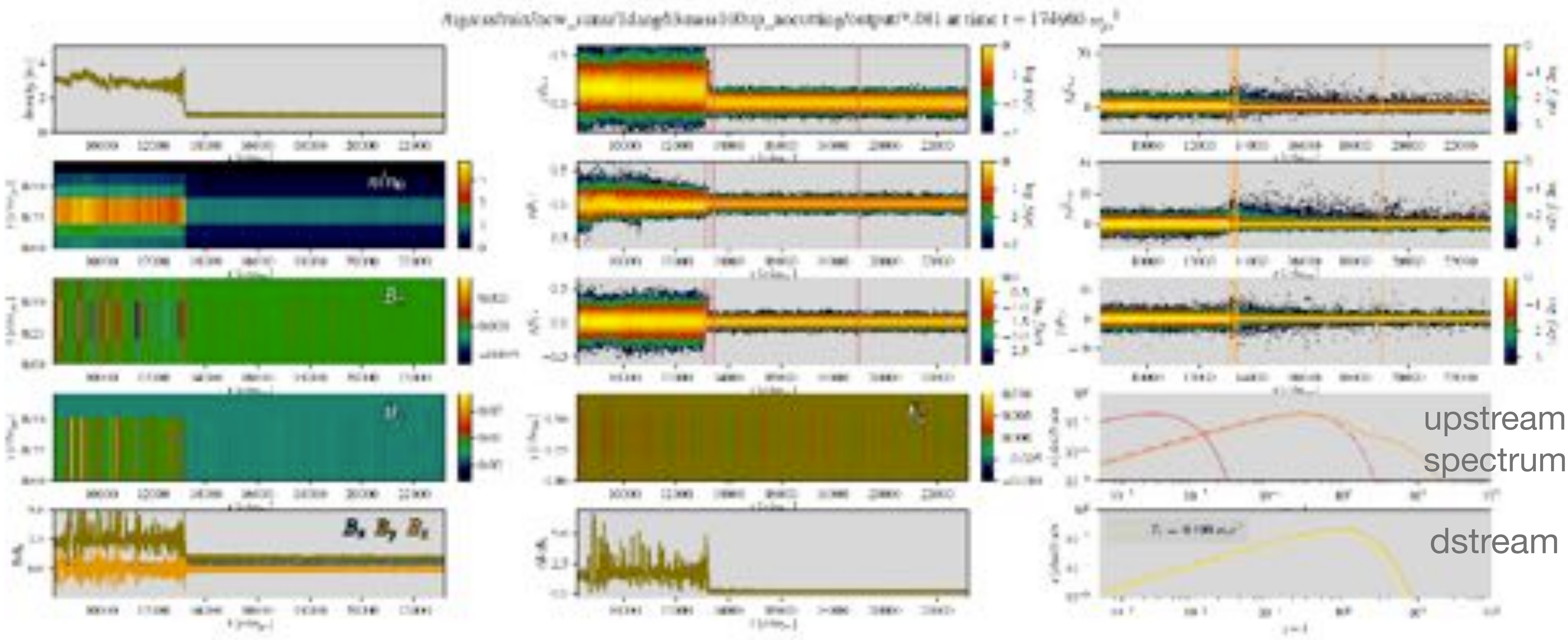
Guo, Sironi, Narayan (2014):
Low sonic Mach # = 2; 63 degrees shock inclination,
mi/me=100, M_A=20; electron-driven waves upstream



Electron acceleration at ⊥-shocks



Xu, Caprioli, AS (in prep):
Low sonic Mach # = 2; 63 degrees shock inclination,
mi/me=100, M_A=20; electron-driven waves upstream



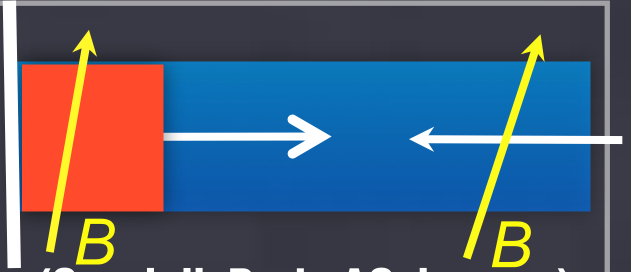
upstream spectrum

dstream

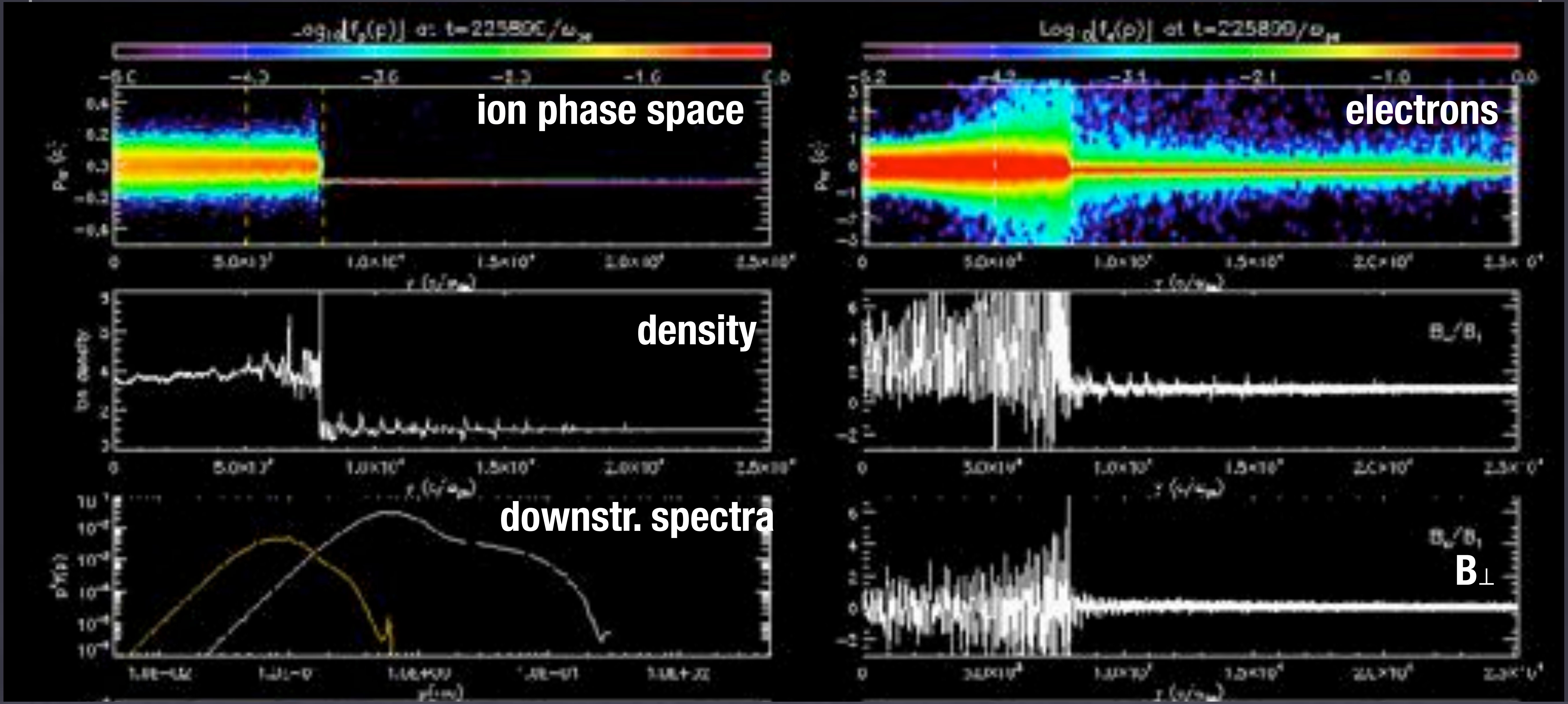
saturation?

Electron acceleration at ⊥-shocks

Higher sonic Mach: 60 degrees shock inclination, mi/me=100, $M_A=M_S=20$; electron-driven waves upstream



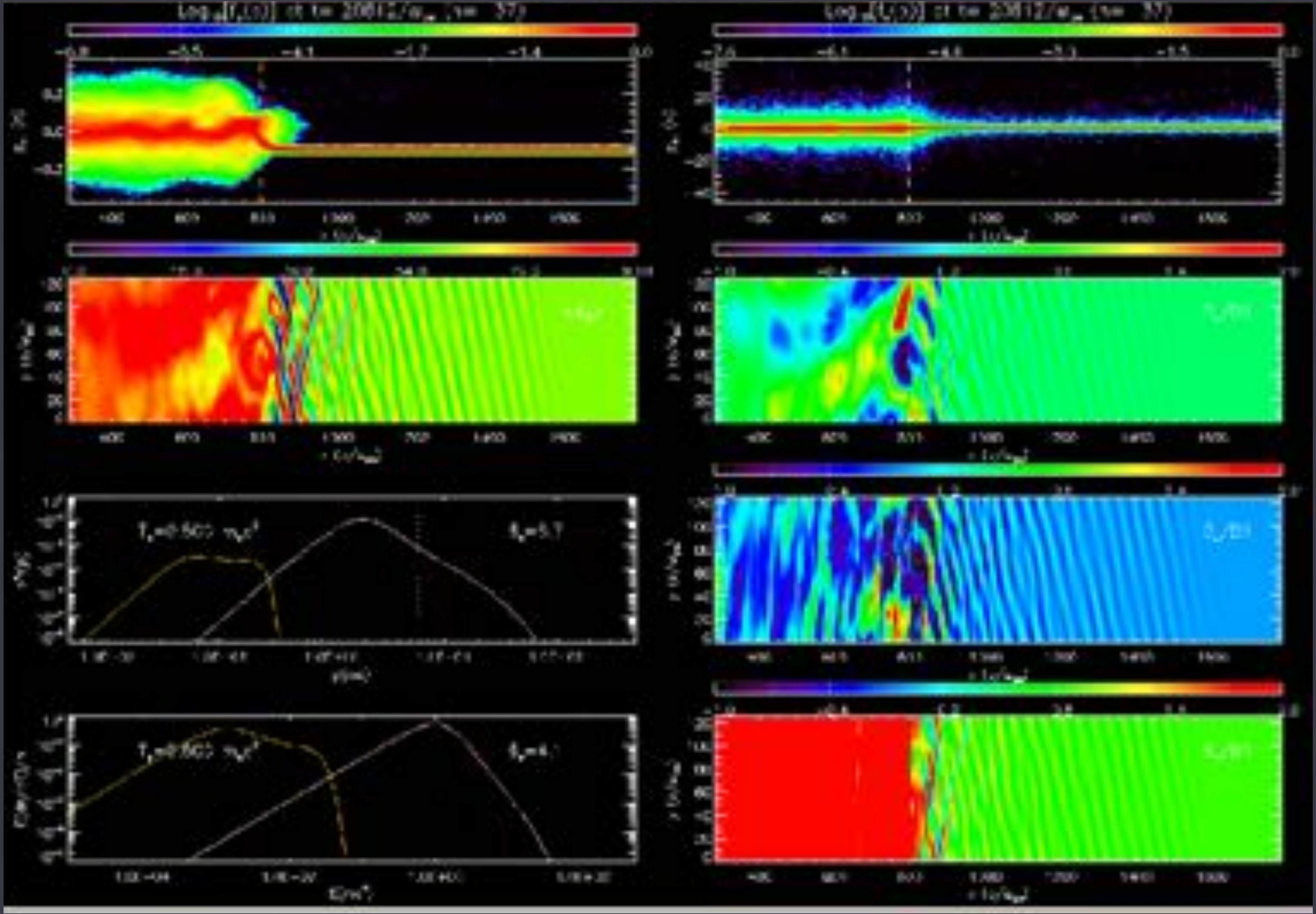
(Caprioli, Park, AS, in prep)



Ions are not injected or accelerated into DSA, while electrons drive their own Bell-type waves. Electrons are reflected from shock due to magnetic mirroring.

Recover DSA electron spectrum, 0.1-4% in energy, <1% by number.

Electron acceleration at ⊥-shocks: 2D



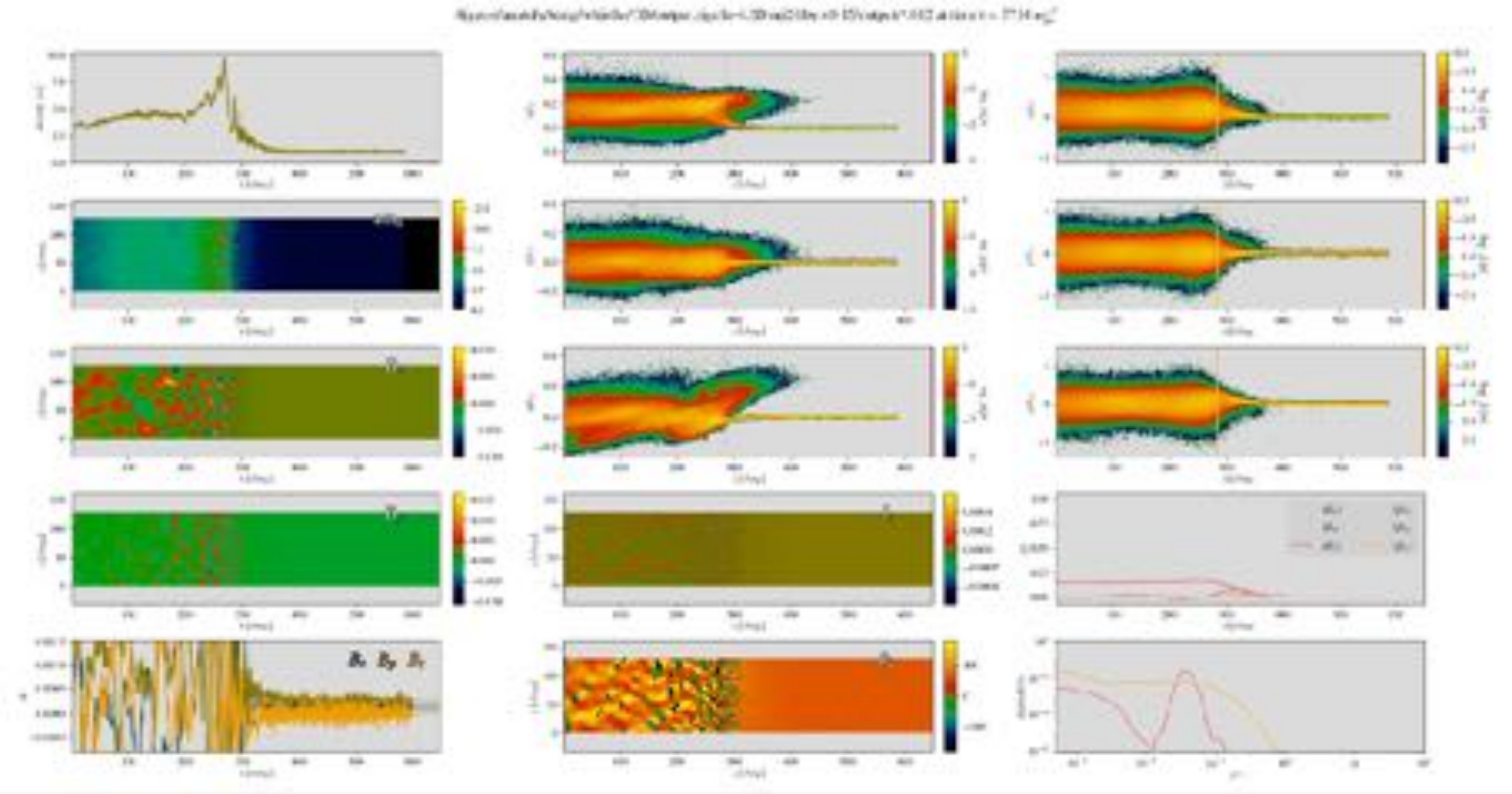
Low-M_A shocks; Whistler waves in the shock foot for $M_A<m_i/m_e$;

Electron DSA! Large-amplitude Electron-driven modes! Oblique firehose? (Guo+ 2014). Or whistlers?

Caveat:

Electron acceleration is sensitive to simulation dimensionality and field orientation: 2D in-plane B field reflects fewer electrons than out-of-plane B field

field in-plane



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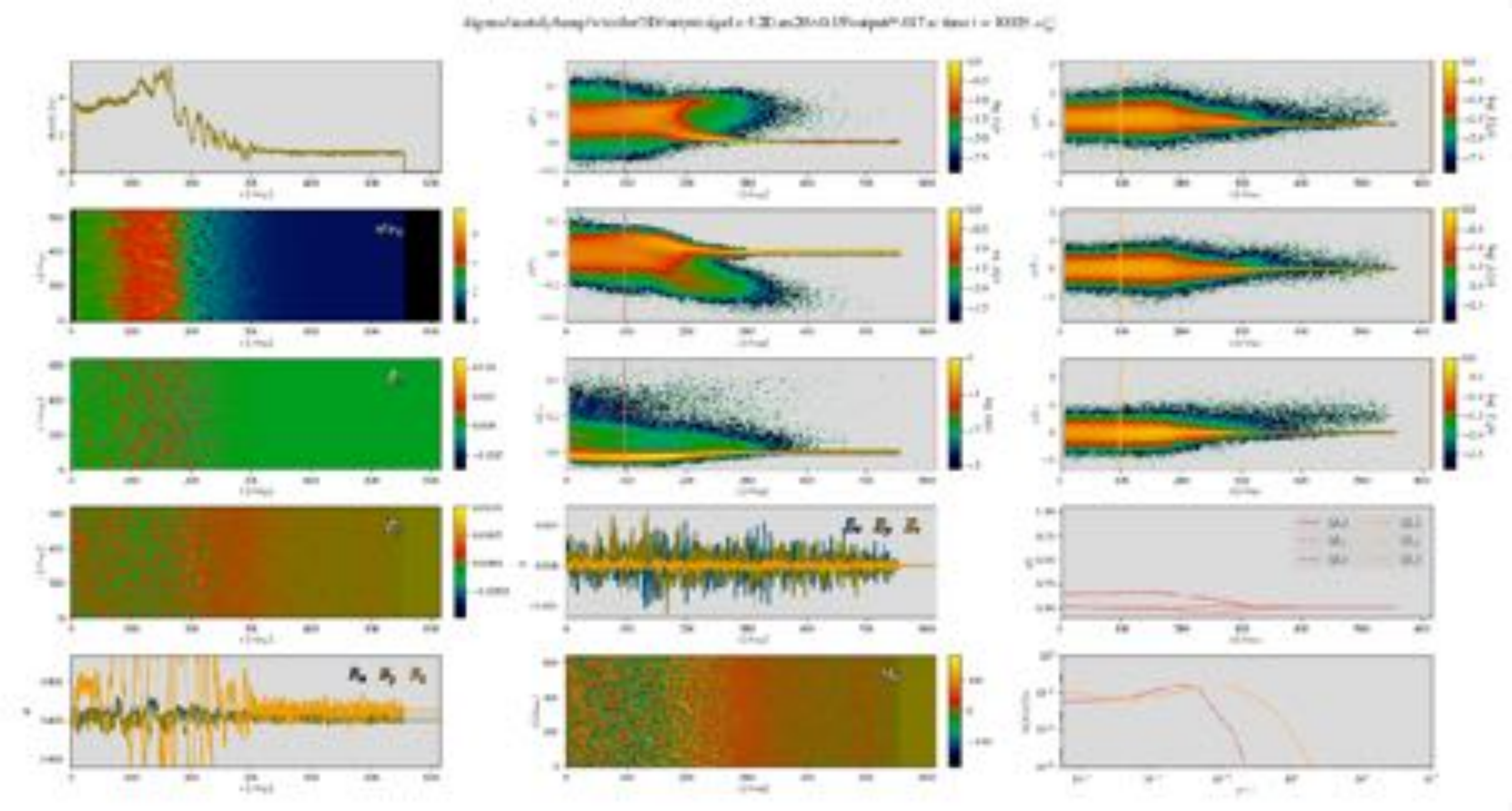
field out-of-plane



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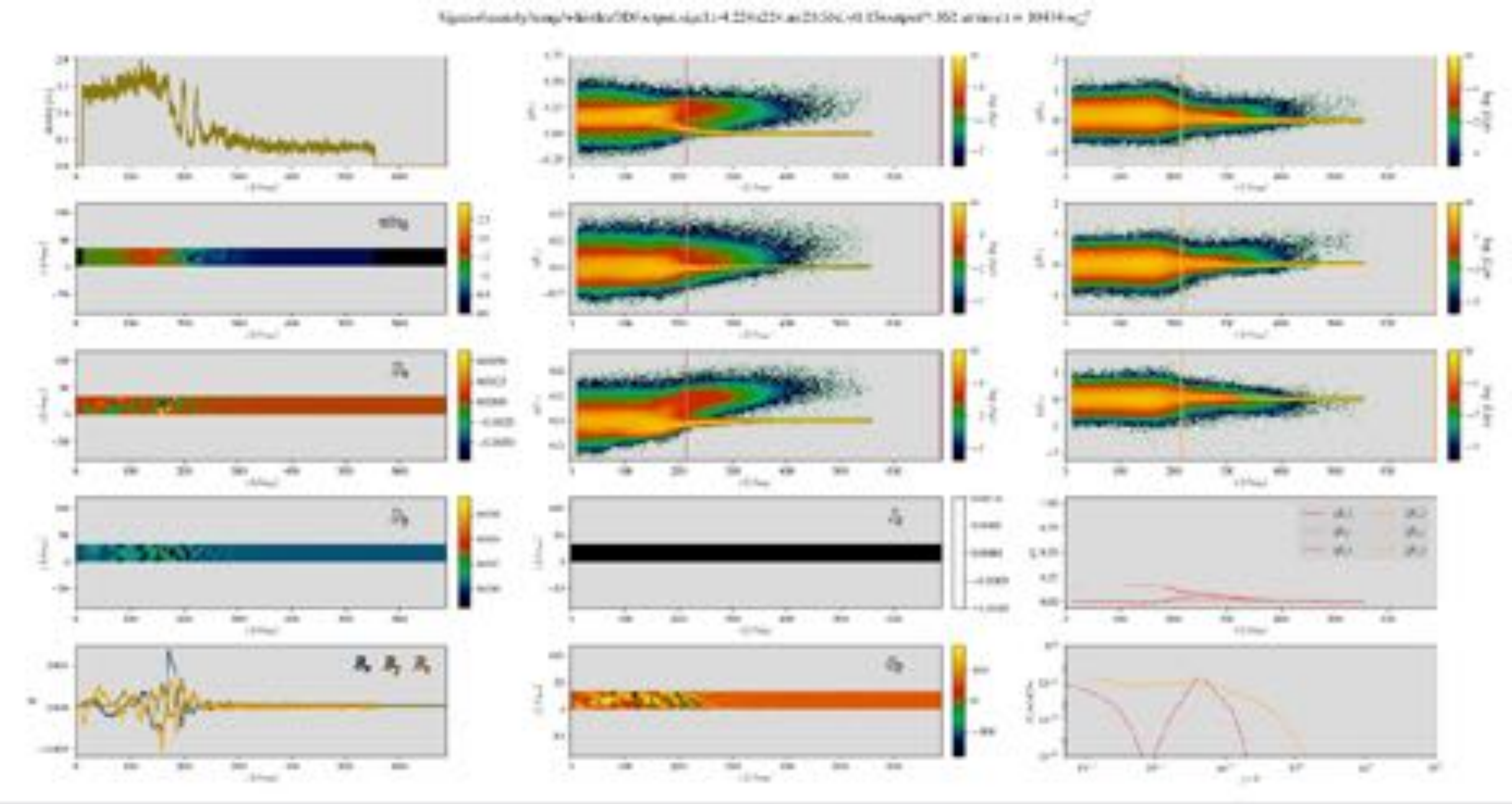
3D field



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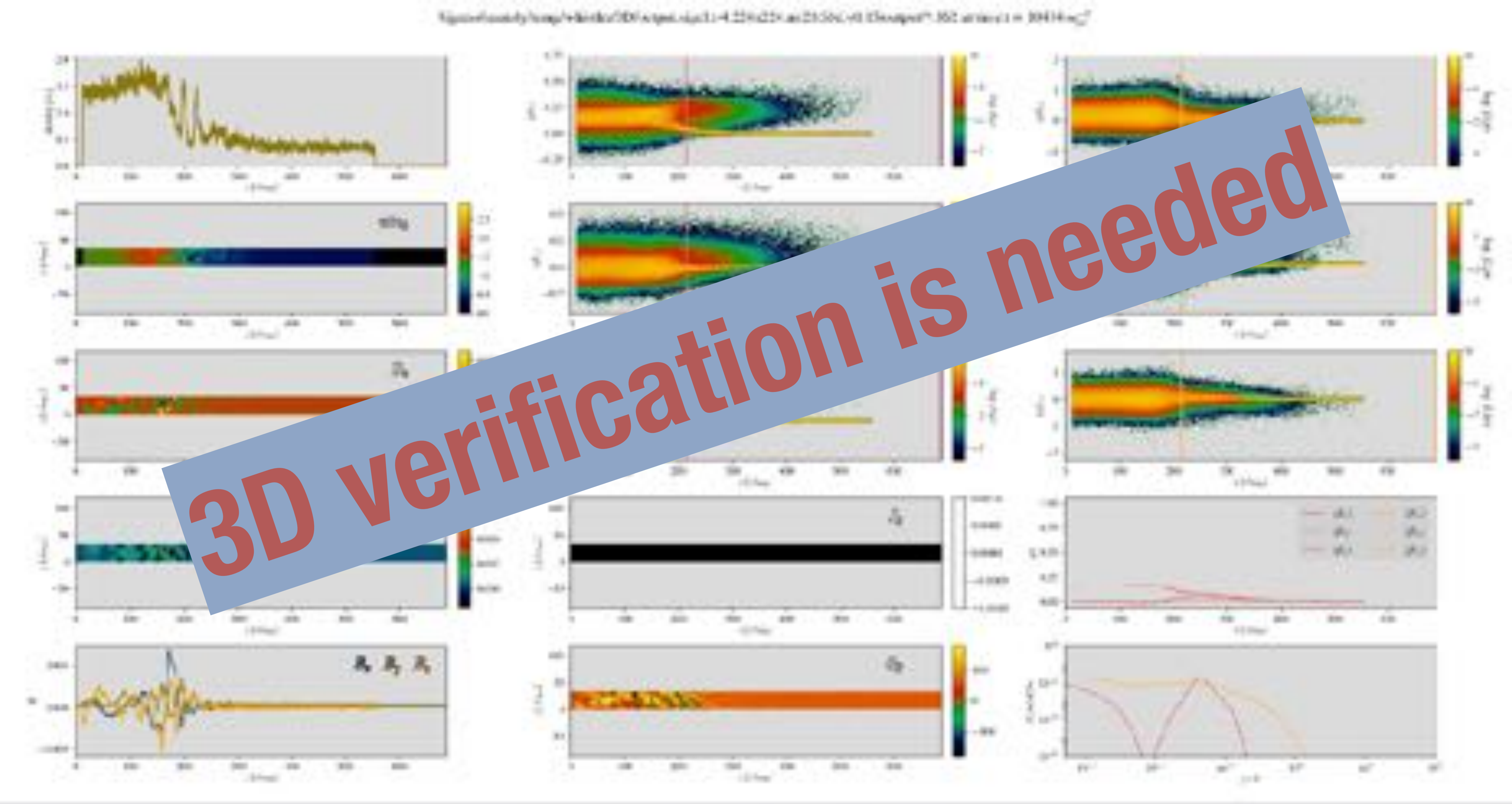
3D field



3D verification is needed

Shock acceleration: emerging picture

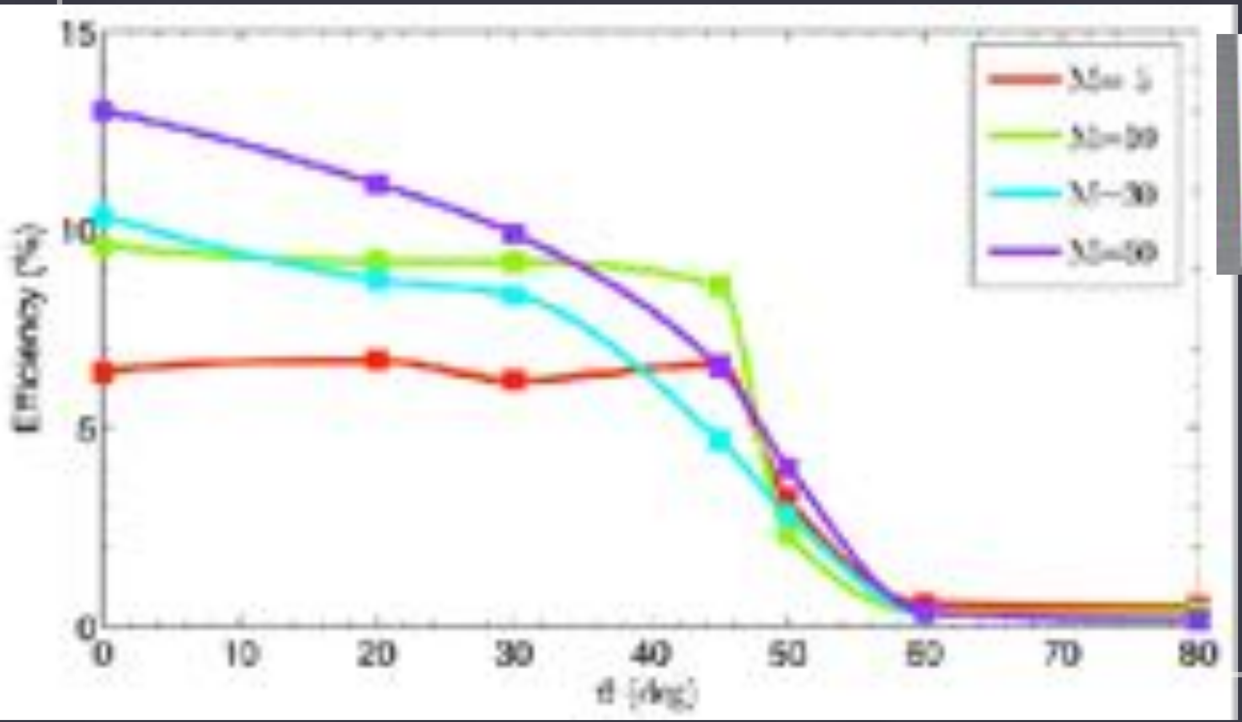
Acceleration in laminar field:

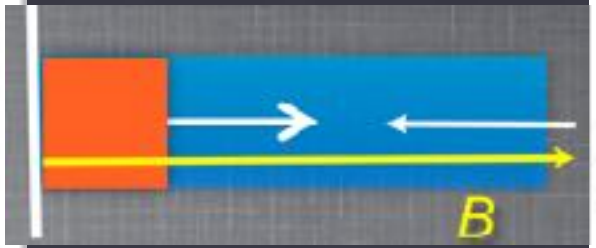
quasi-parallel -- accelerate both ions and electrons

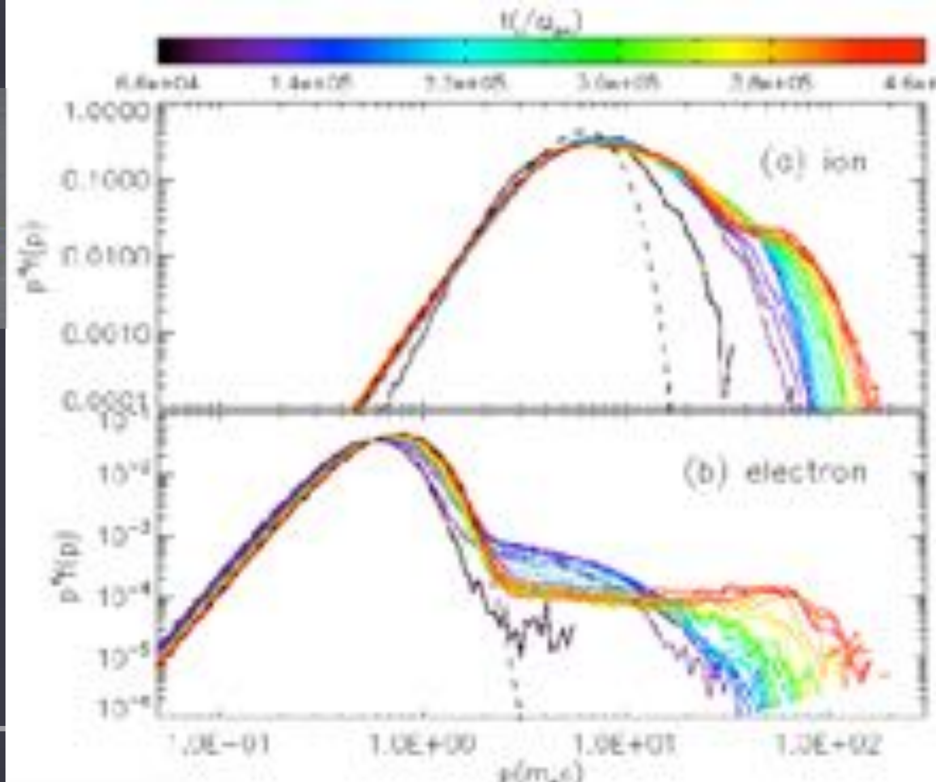
(Caprioli & AS, 2014abc; Park, Caprioli, AS 2015)

quasi-perpendicular -- accelerate mostly electrons

(Guo, Sironi & Narayan 2014; Caprioli, Park, AS in prep)







Shock acceleration: emerging picture

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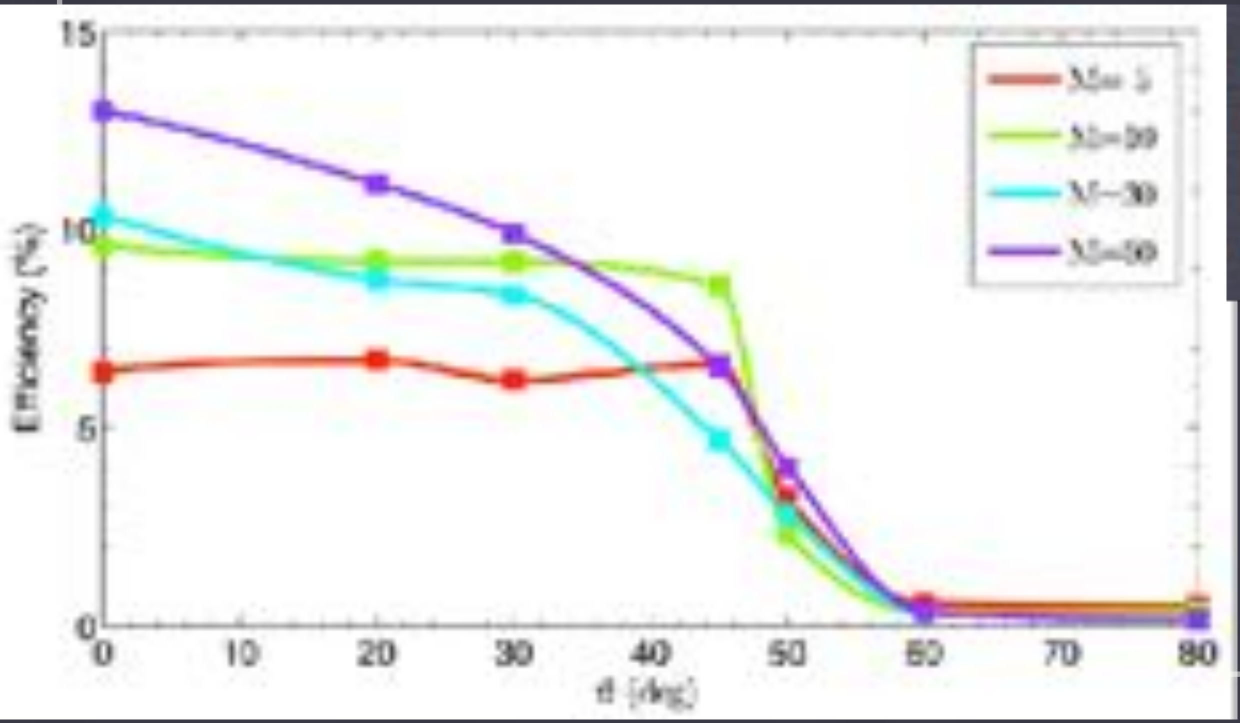
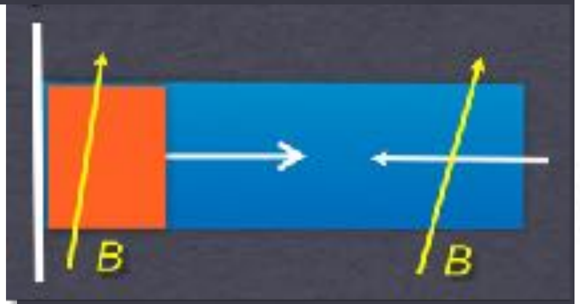
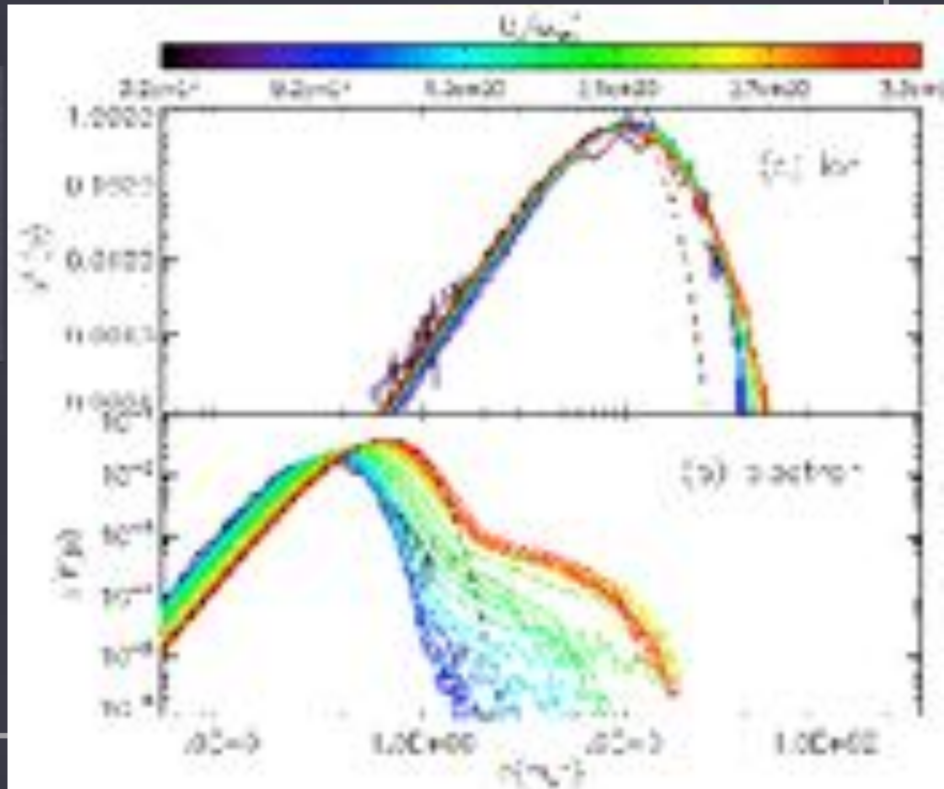


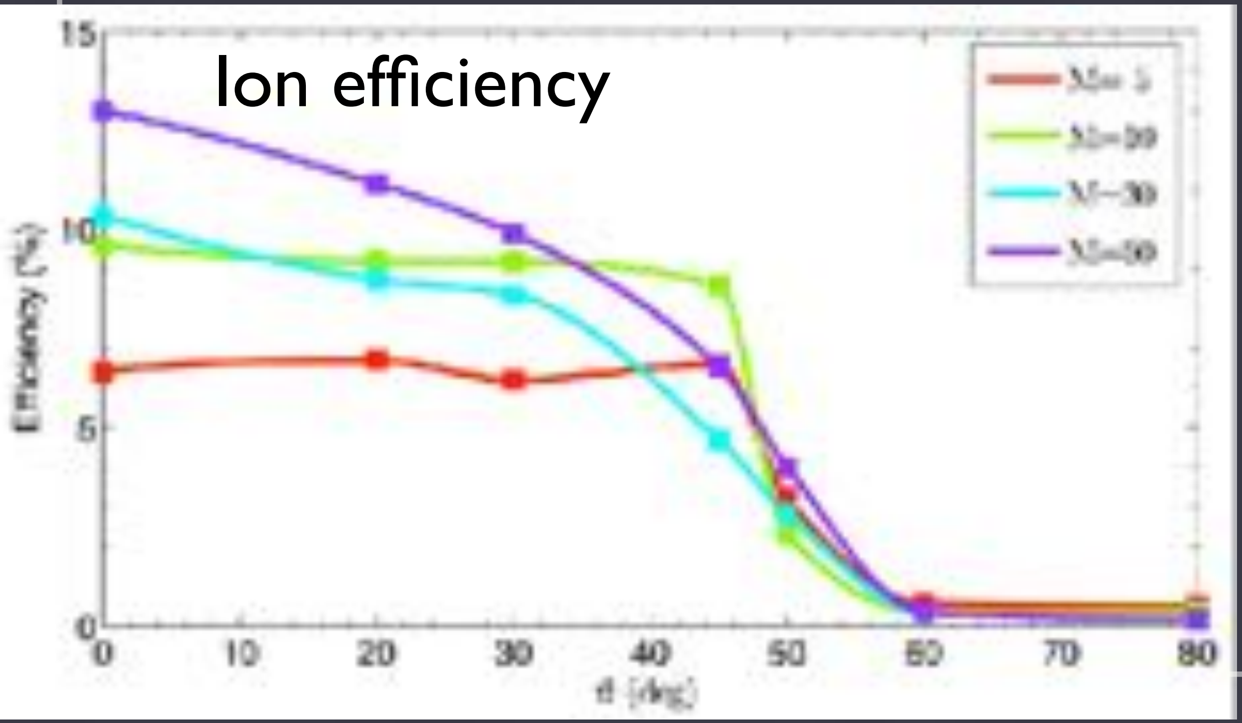


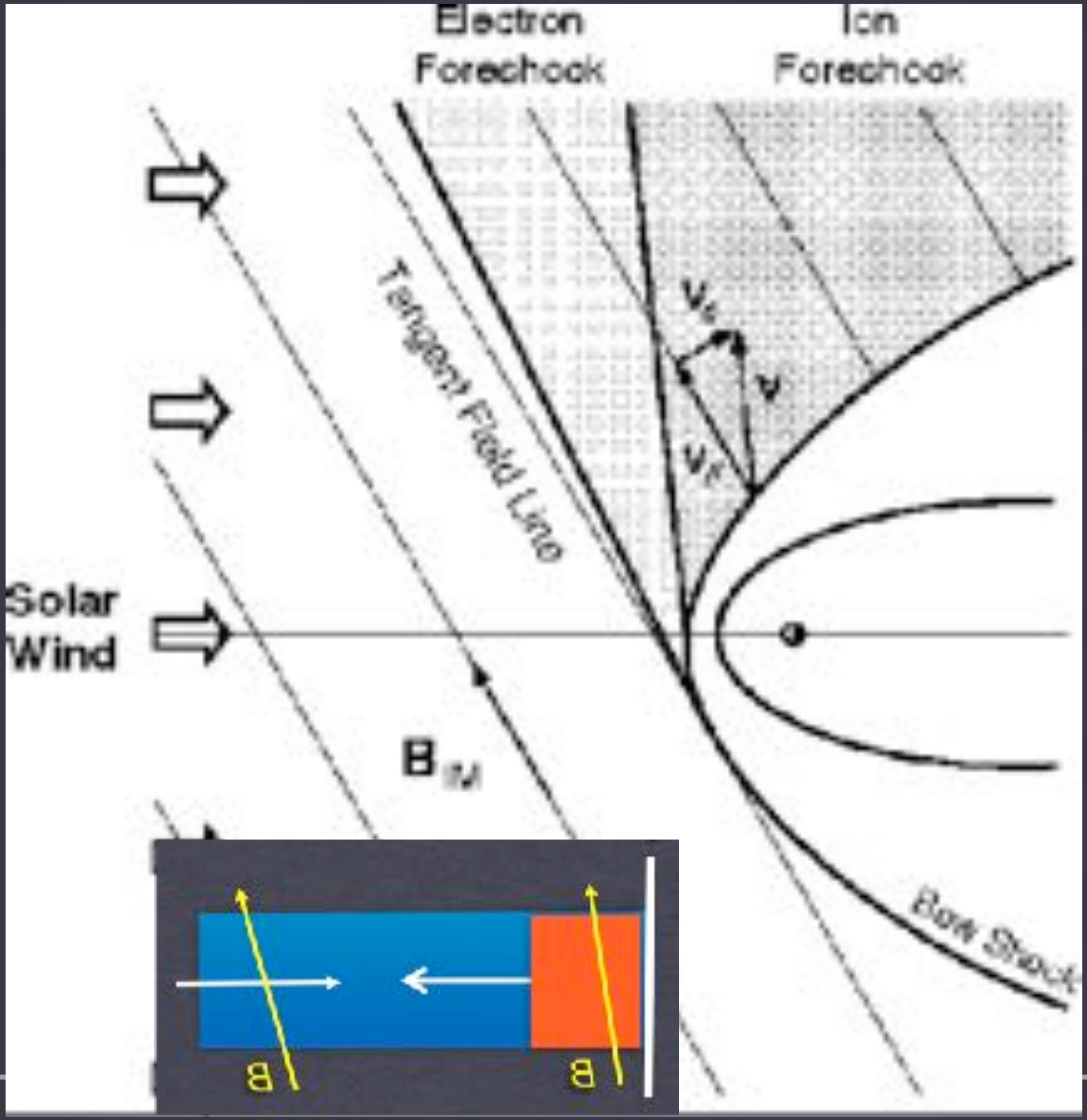


Shock acceleration: emerging picture

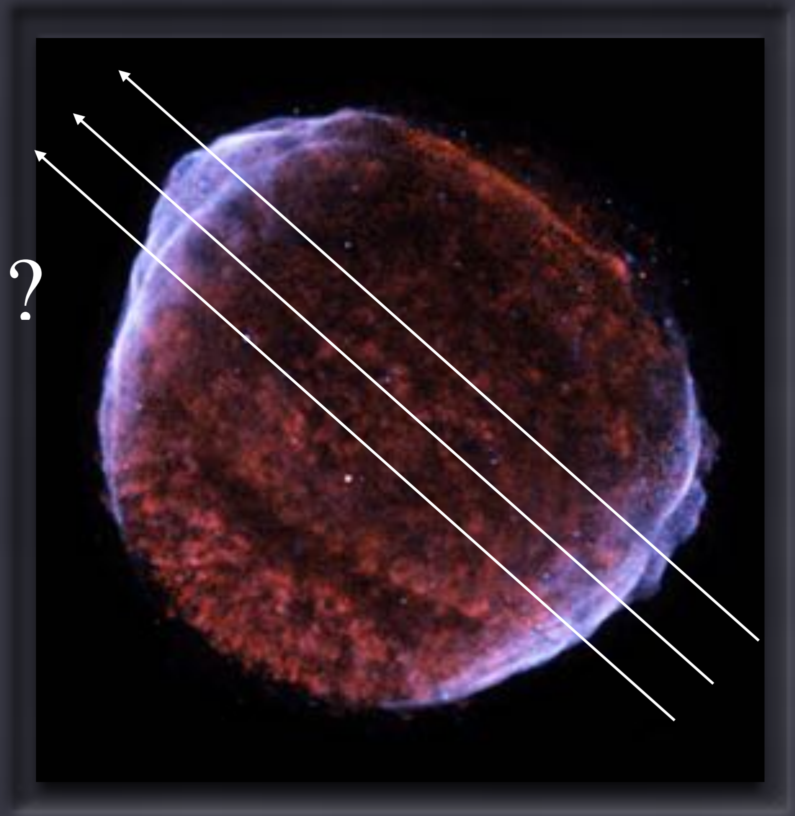
Magnetosphere
does this!

Efficiency vs ang w
explained theoret ally
(Caprioli et al 2015)





SNR story



Nonthermally-emitting SNRs likely have large scale parallel magnetic field (radial). This leads to CR acceleration and field amplification.

Locally-transverse field enters the shock, and causes electron injection and DSA.

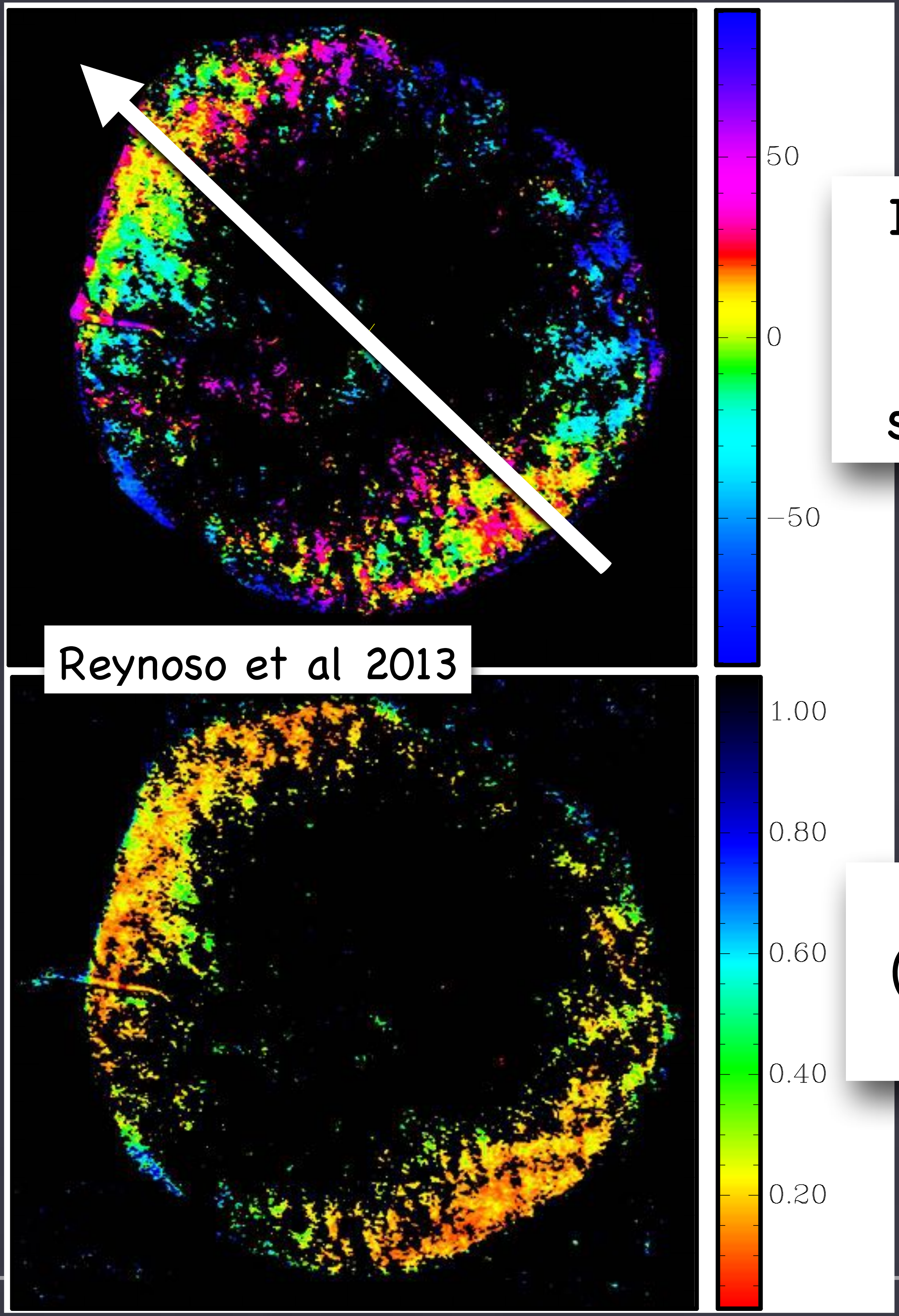
This favors large-scale radial B fields in young SNRs. Polarization in “polar caps” should be small -- field is random

Ab-initio plasma results allow to put constraints on the large-scale picture!

SN1006: a parallel accelerator

X-ray emission
(red=thermal
white=synchrotron)

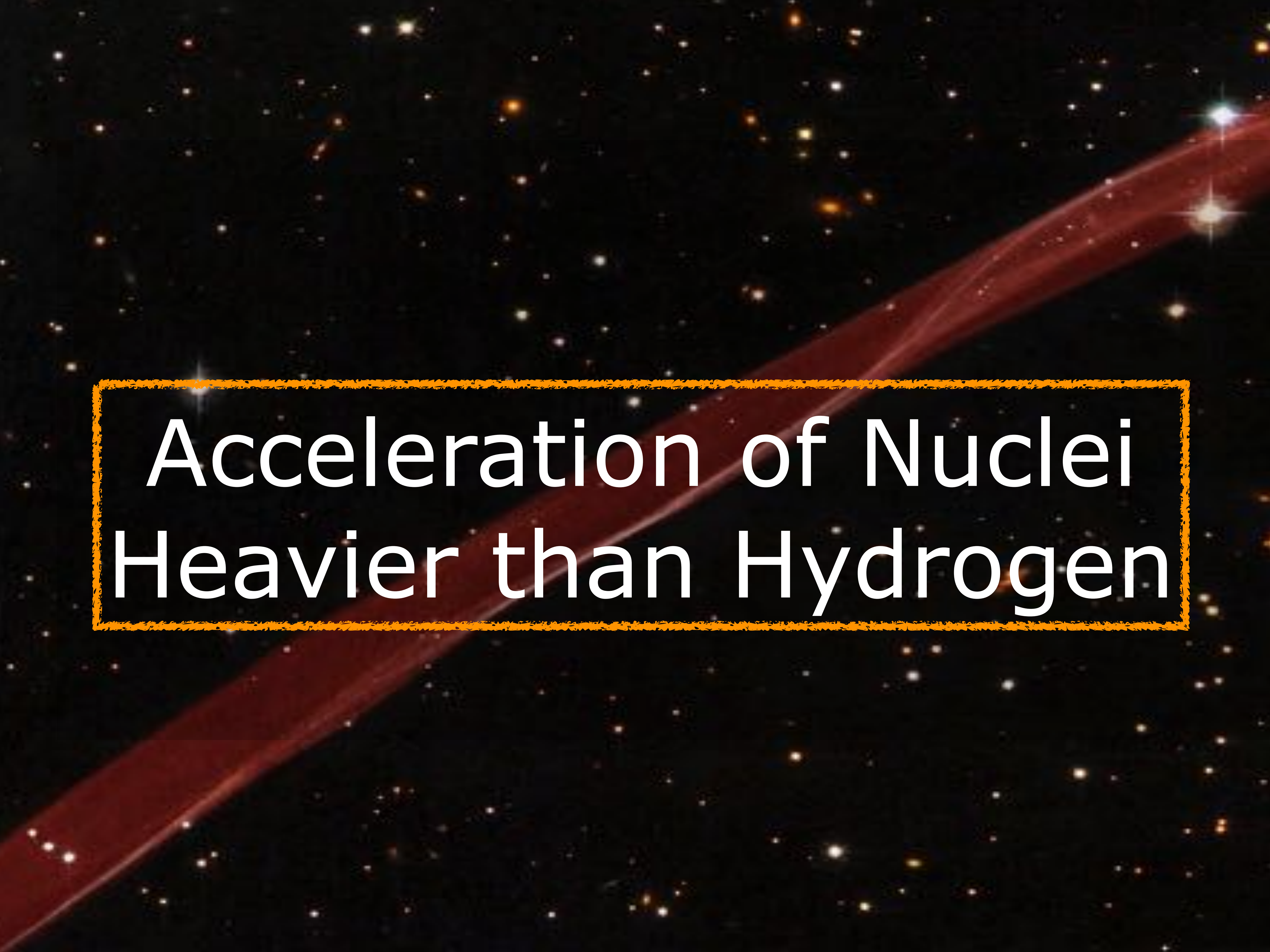
Magnetic field amplification and particle acceleration where the shock is parallel



Reynoso et al 2013

Inclination of
the B field
wrt to the
shock normal

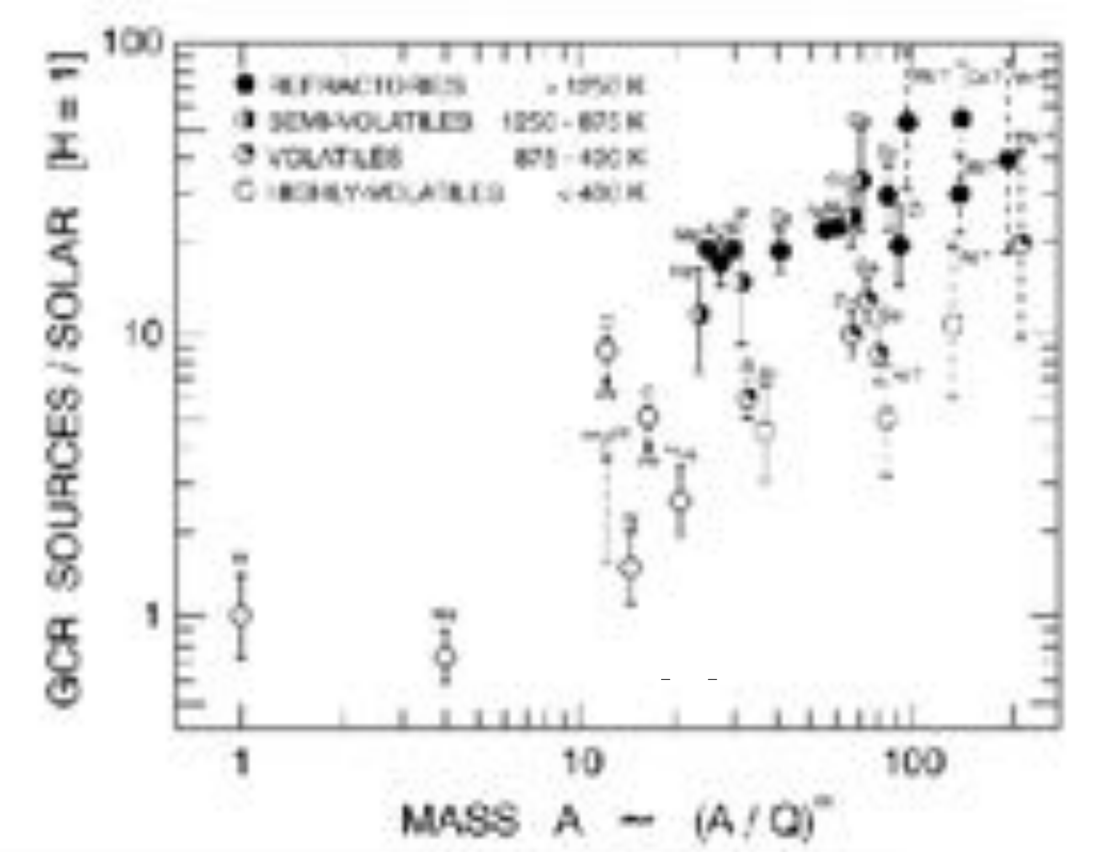
Polarization
(low=turbulent
high=ordered)



Acceleration of Nuclei Heavier than Hydrogen

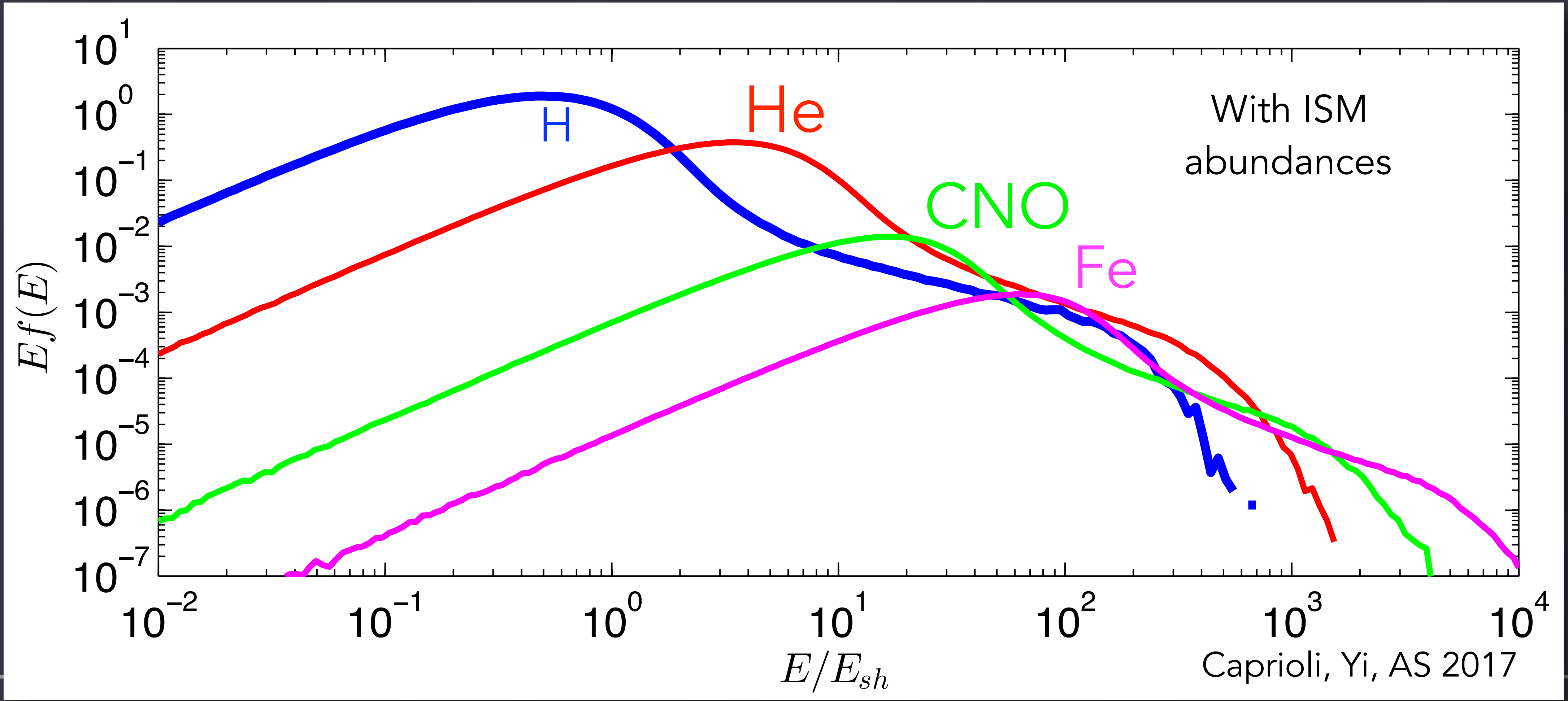
Acceleration of heavy nuclei

Nuclei heavier than H must be injected more efficiently (Meyer et al 97)



Multi-species hybrid simulations.
Max energy is proportional to charge Z;

Most nuclei have A/Z ~ 2. Investigate also A/Z>2 for partially ionized nuclei.



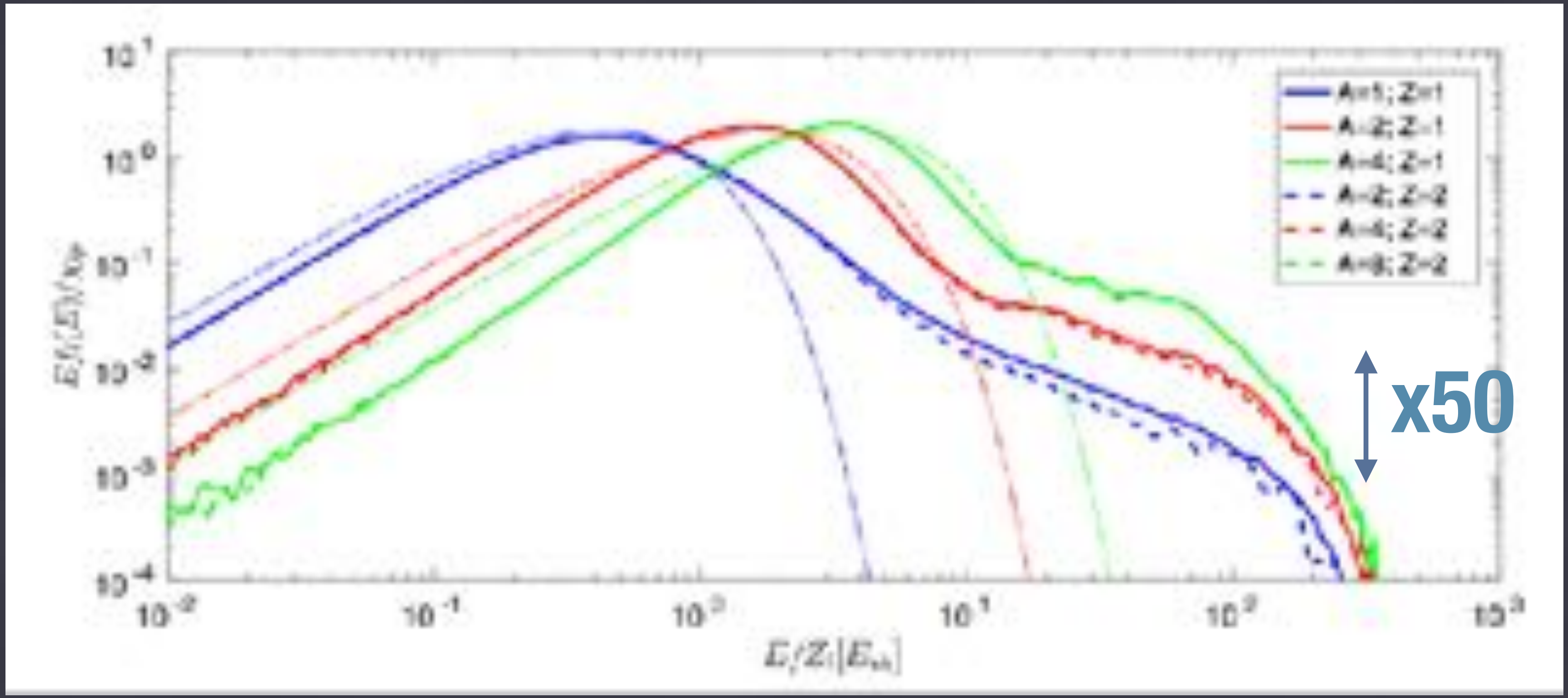
Caprioli, Yi, AS 2017

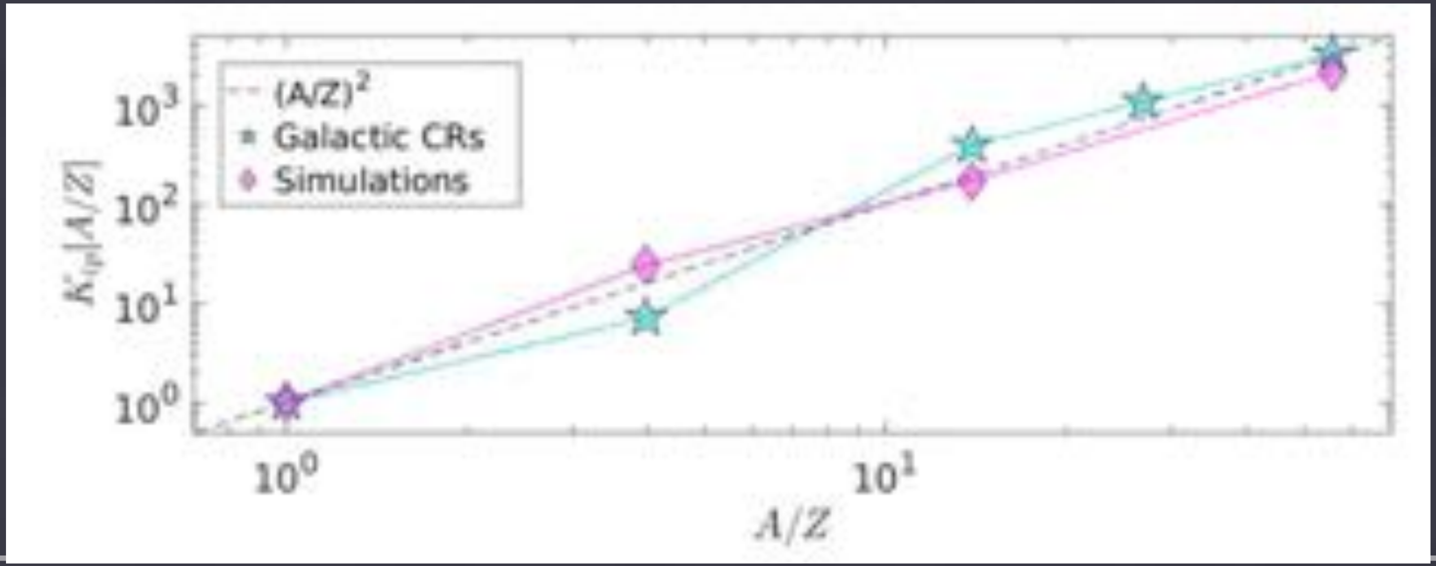
Injection of singly-ionized nuclei

In the absence of H-driven turbulence, heavies are thermalized far downstream

With B amplification from H, heavies are thermalized to kT=A $mv_{sh}^2/2$, and can recross the shock due to their large larmor radii. More chances to scatter on H fluctuations leads to higher “duty fraction” of the shock for larger A/Z.

Nuclei enhancement depends on A/Z and Mach number. Caprioli, Yi, AS arXiv: 1704.08252

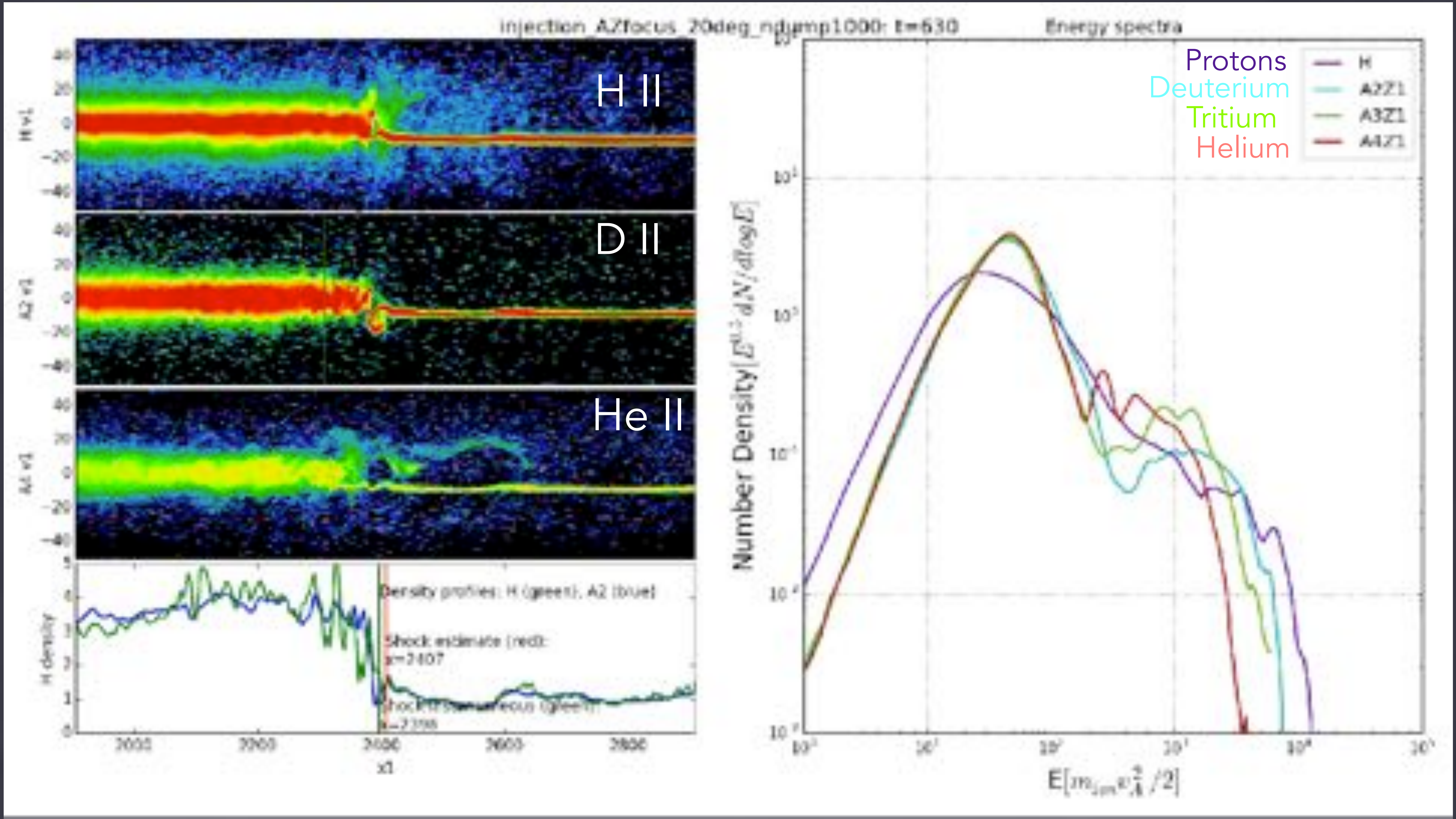




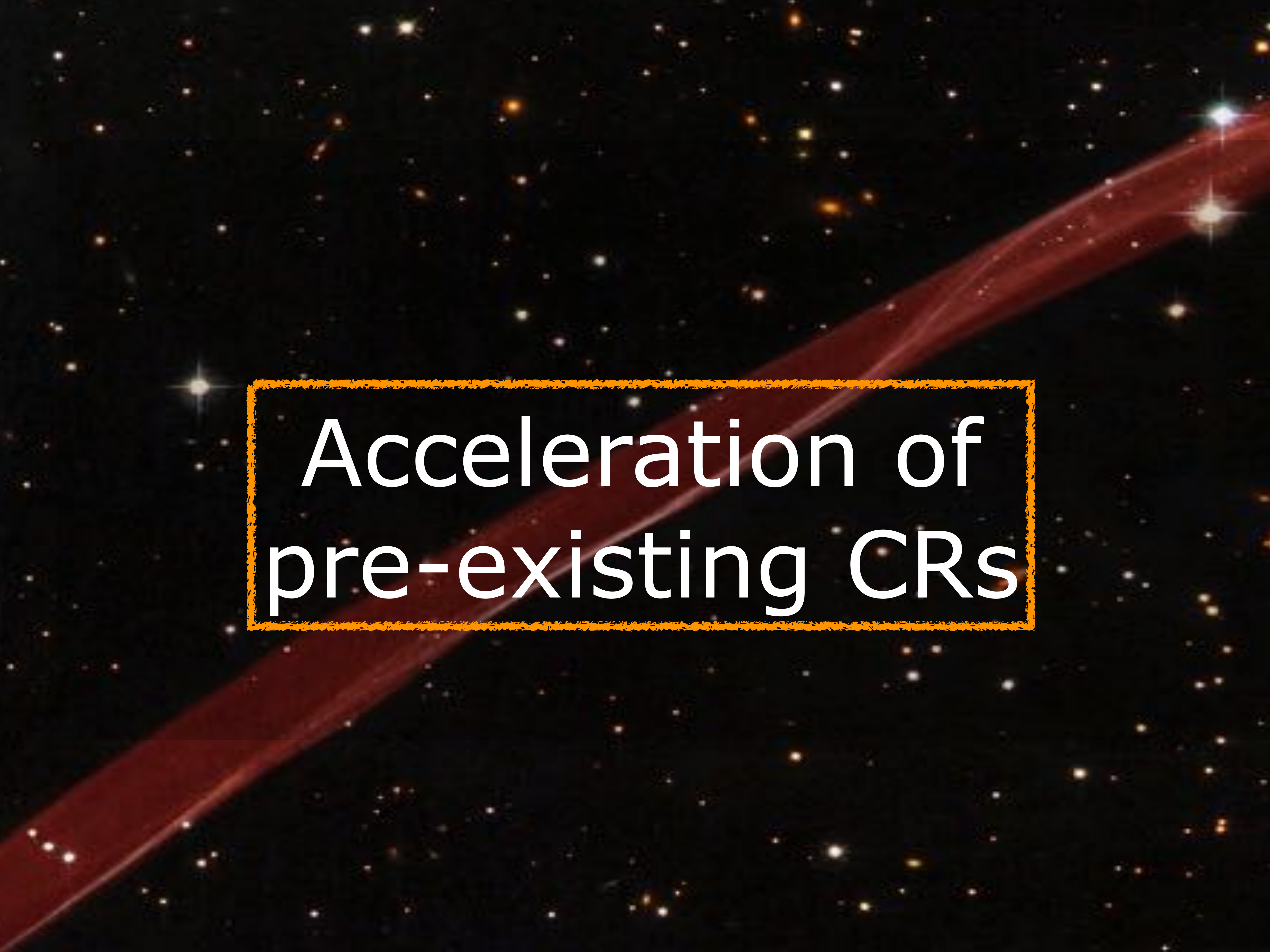
Injection fraction is larger for nuclei with larger A/Z! Pickup of incompletely ionized heavy ions may be responsible for GCR abundance with A

Injection of singly-ionized nuclei

M=10, parallel shock (Caprioli, Yi, AS 2017)



Injection fraction is larger for nuclei with larger A/Z!

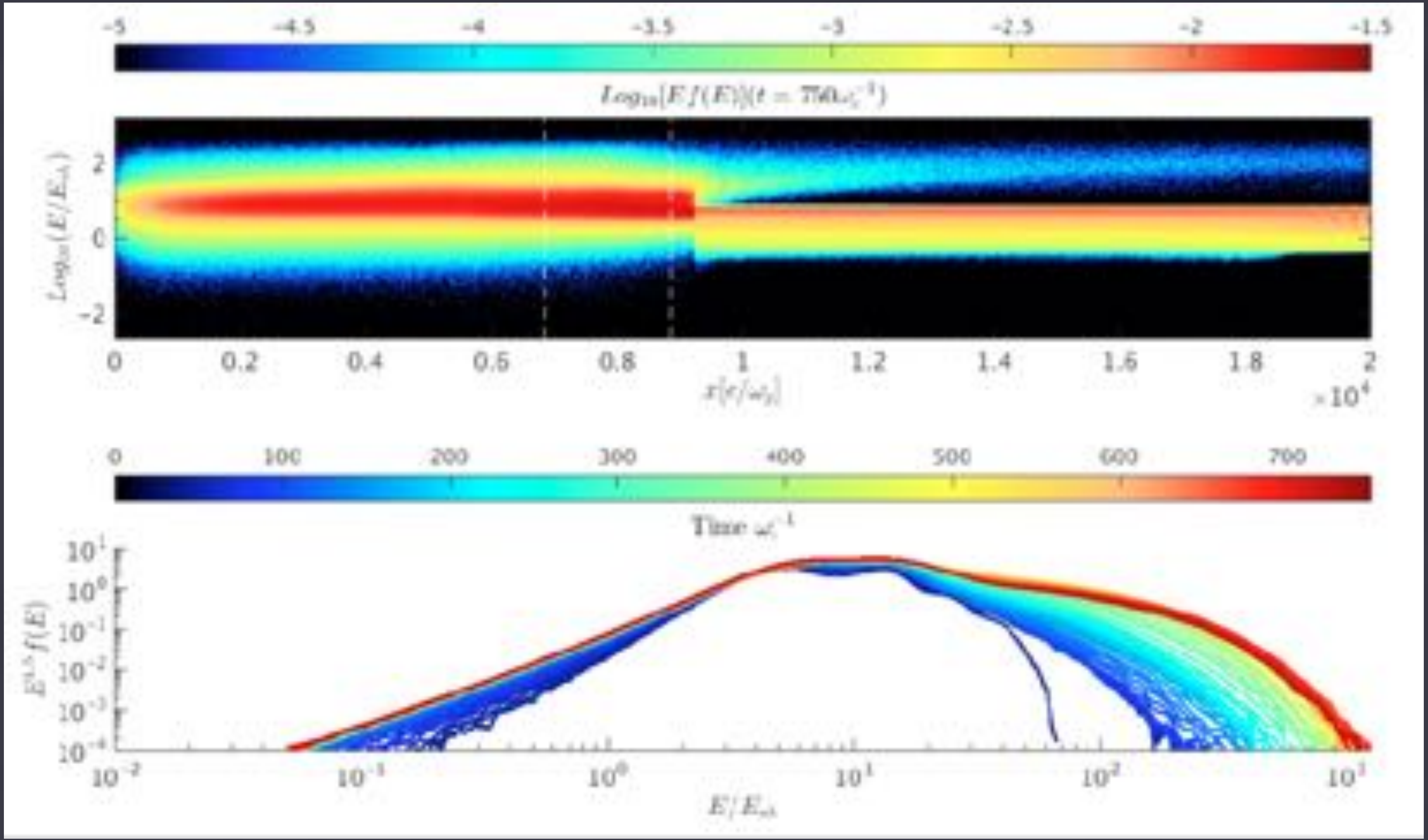


Acceleration of pre-existing CRs

Re-acceleration of pre-existing CRs

Add hot “CR” particles to upstream flow (Caprioli, Zhang, AS, in prep).

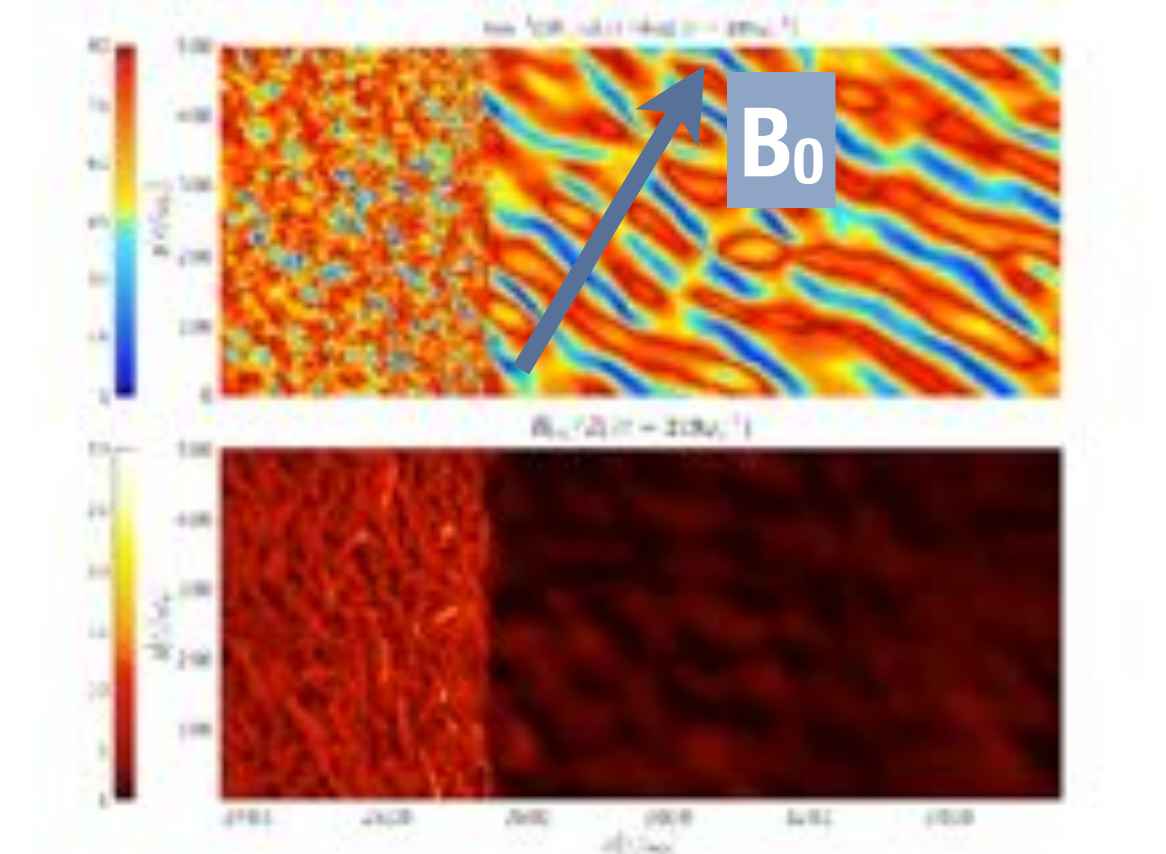
Quasi-perp shock: CRs have large Larmor radii and can recross the shock, accelerate, and be injected into diffusive acceleration process



Turbulence driven by reaccelerated CRs

Escaping CRs drive turbulence

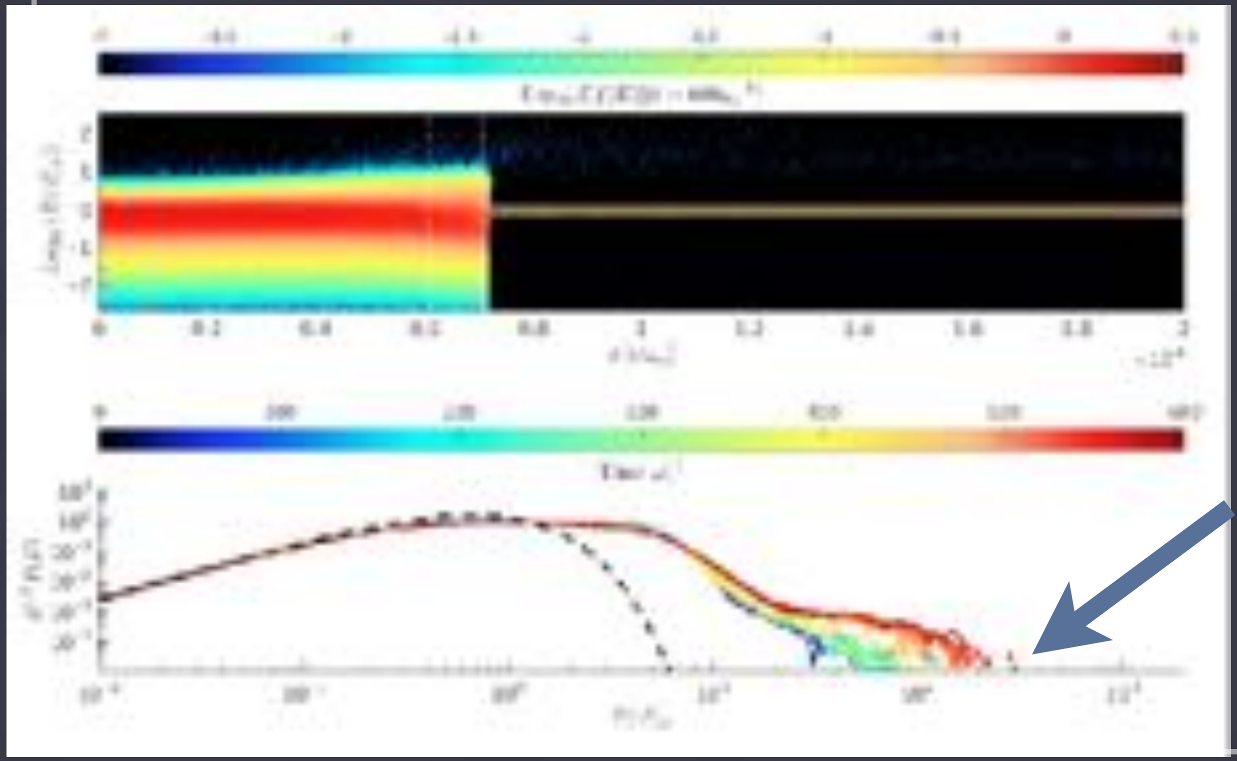
field inclination



Orientation of the field at the shock changes to regions of quasi-parallel, and efficiency of H acceleration increases.

Pre-existing CRs improve local efficiency of the shock!

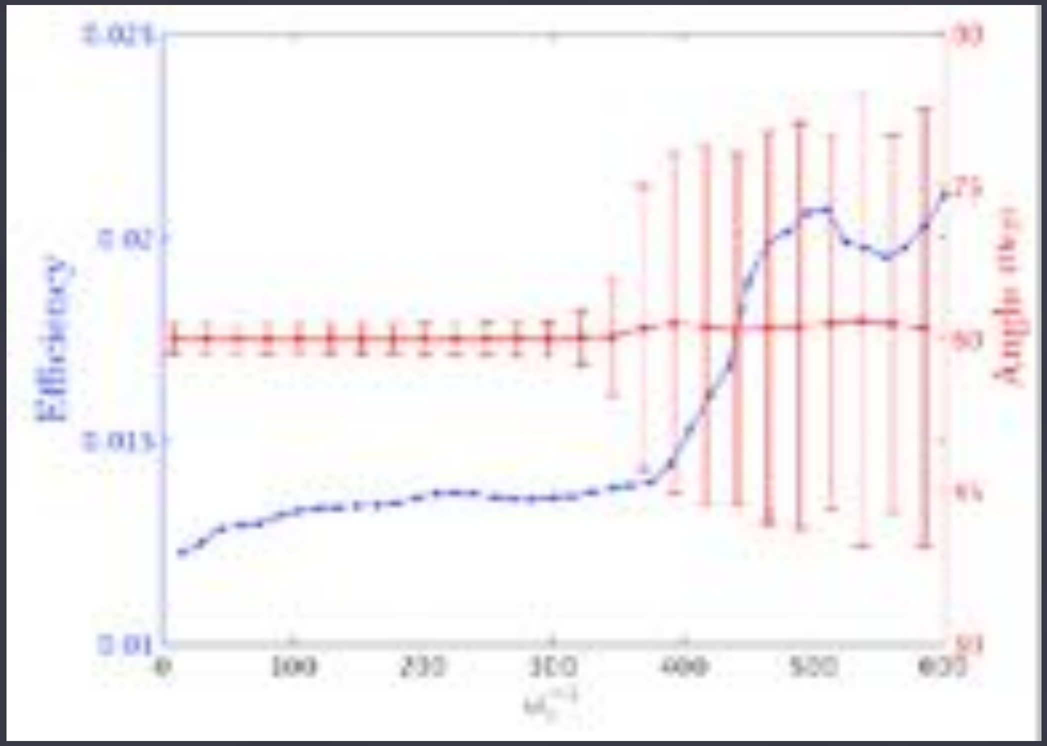
Growth time in SNR ~10yrs << age.



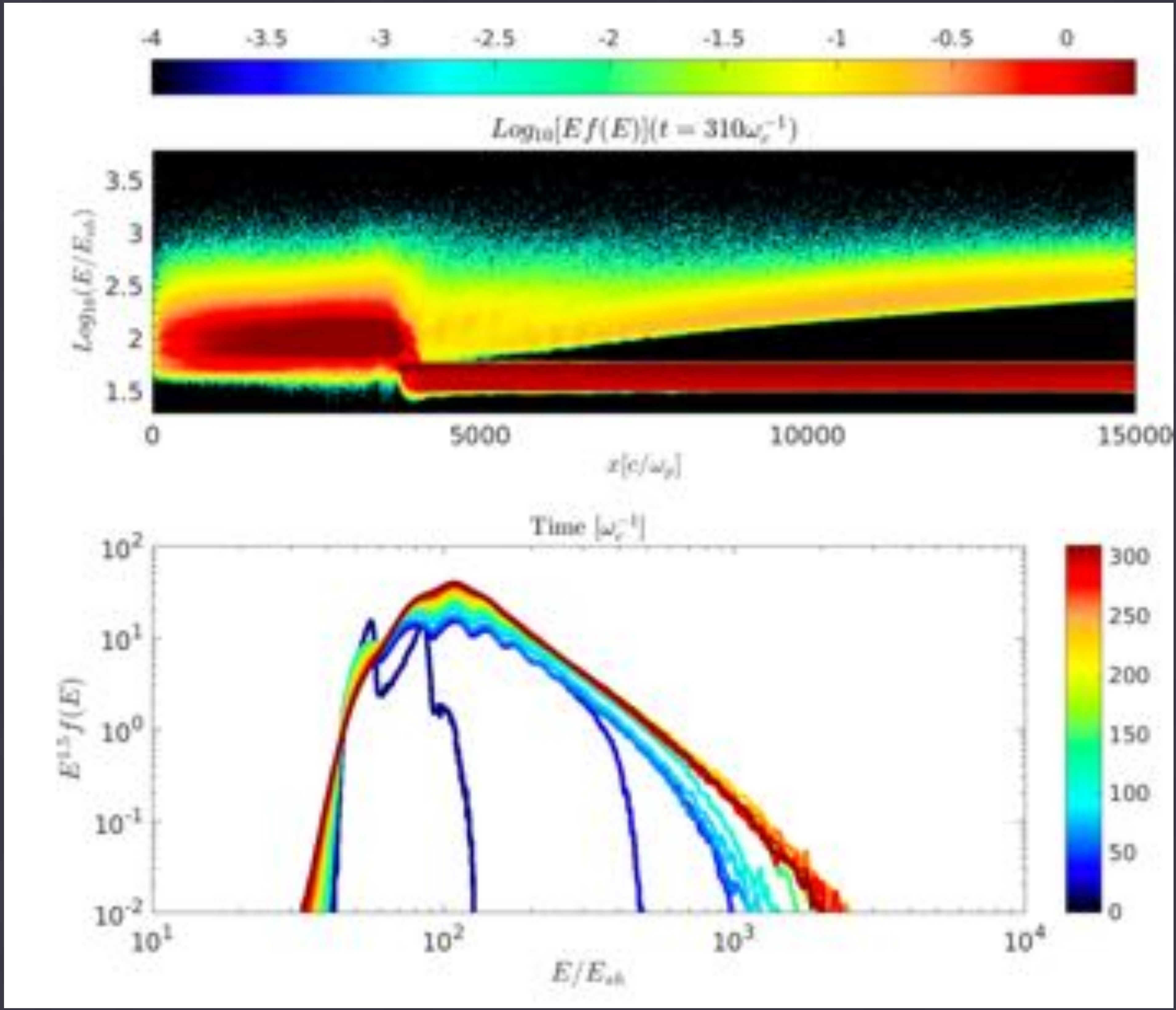
Proton spectrum

60° shock

n_{cr}=2e-3



Reacceleration of CRs: spectrum is steeper E^{-4}



Conclusions

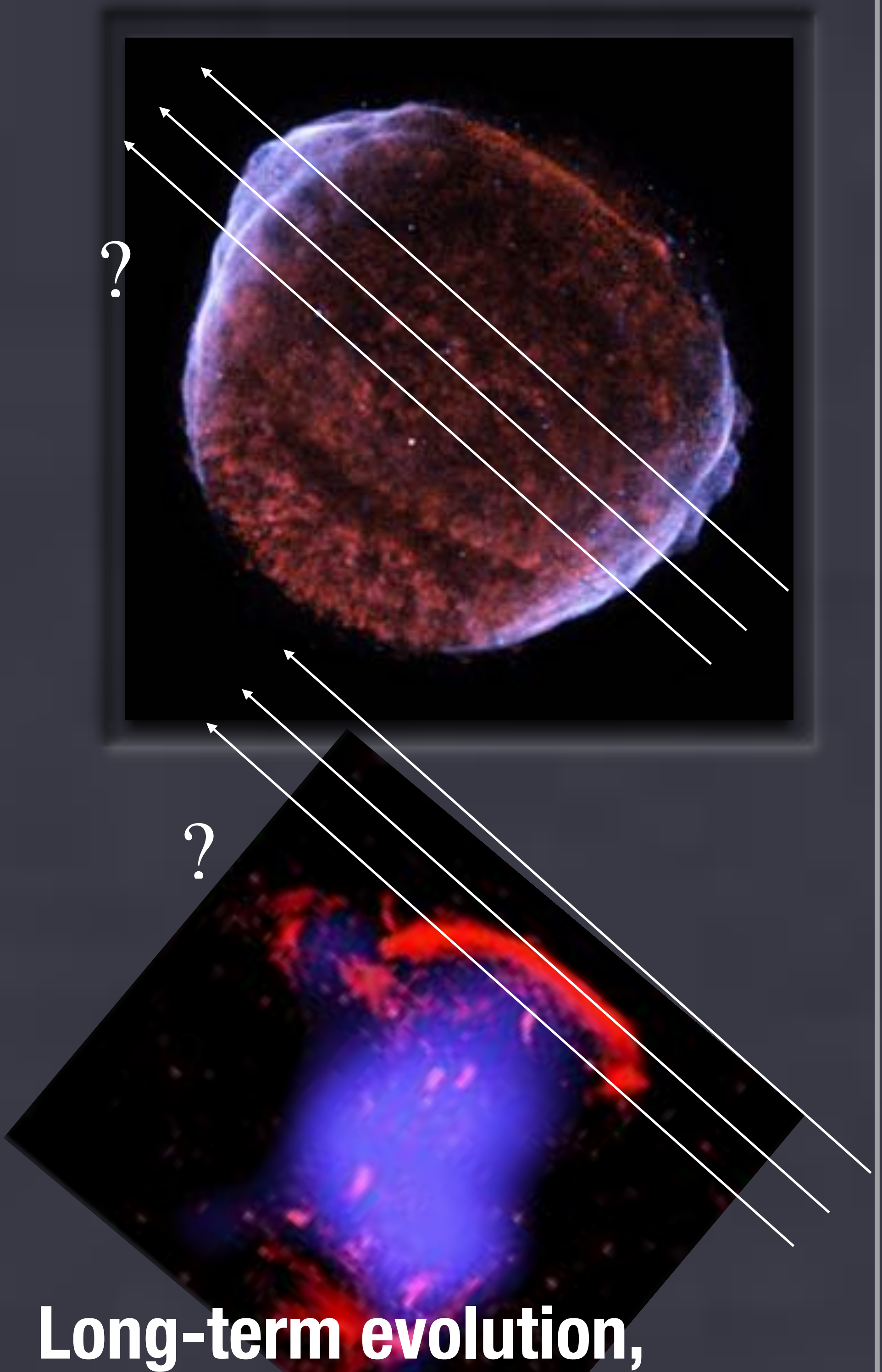
Kinetic simulations allow to calculate particle injection and acceleration from first principles, constraining injection fraction

Magnetization (Mach #) of the shock and B inclination controls the shock structure

Nonrelativistic shocks accelerate ions and electrons in quasi-par if B fields are amplified by CRs. Energy efficiency of ions 10-20%, number ~few percent; K_{ep}~10^{-3}; p^{-4} spectrum

Electrons are accelerated in quasi-perp shocks, energy several percent, number <1%. Fewer ions are accelerated at oblique shocks.

A/Z>2 species are injected more efficiently; CR re-acceleration may be important



Long-term evolution, turbulence & 3D effects need to be explored more: more advanced simulation methods are coming