

SMASH



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Unifying Inflation with the Axion, Dark Matter, Baryogenesis, and the Seesaw Mechanism

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A minimal extension of the standard model (SM) with a single new mass scale and providing a complete and consistent picture of particle physics and cosmology up to the Planck scale is presented. We add to the SM three right-handed SM-singlet neutrinos, a new vectorlike color triplet fermion, and a complex SM-singlet scalar σ that stabilizes the Higgs potential and whose vacuum expectation value at $\sim 10^{11}$ GeV breaks lepton number and a Peccei-Quinn symmetry simultaneously. Primordial inflation is produced by a combination of σ (nonminimally coupled to the scalar curvature) and the SM Higgs boson. Baryogenesis proceeds via thermal leptogenesis. At low energies, the model reduces to the SM, augmented by seesaw-generated neutrino masses, plus the axion, which solves the strong CP problem and accounts for the dark matter in the Universe. The model predicts a minimum value of the tensor-to-scalar ratio $r \simeq 0.004$, running of the scalar spectral index $\alpha \simeq -7 \times 10^{-4}$, the axion mass $m_A \sim 100 \mu\text{eV}$, and cosmic axion background radiation corresponding to an increase of the effective number of relativistic neutrinos of ~ 0.03 . It can be probed decisively by the next generation of cosmic microwave background and axion dark matter experiments.

1. What is the dark matter?

2. Horizon and flatness problems

3. Matter/anti-matter asymmetry

$$\frac{n_b - n_{\bar{b}}}{n_\gamma} \simeq 10^{-9}$$

(CMB)

$$n_p/n_{\bar{p}} \sim 10^4$$

(Galactic cosmic rays)

4. Smallness of the neutrino masses

$$\sum m_\nu \lesssim 0.2 \text{ eV}$$

5. Strong CP problem

$$\mathcal{L}_{\text{QCD}} \in -\frac{\theta_0}{32\pi^2} G\tilde{G}$$

$$\theta \equiv \theta_0 - \arg(\det M) \lesssim 10^{-10}$$

(neutron e.d.m.)

Small neutrino masses



Matter/anti-matter asymmetry



Strong CP problem



Dark matter



Inflation

Standard

Model

Axion

See-saw

Higgs portal (inflation)

$$\text{SMASH} = \text{SM} +$$

★ Three singlet neutrinos: N_i

★ A complex scalar: σ

★ Two Weyl fermions: Q and \tilde{Q} in the $\mathbf{3}$ and $\bar{\mathbf{3}}$ of $SU(3)_c$

Dias, Machado, Nishi, Ringwald and Vaudrevange 2014

★ New $U(1)$ symmetry: PQ and lepton number

q	u	d	L	N	E	Q	\tilde{Q}	σ
$1/2$	$-1/2$	$-1/2$	$1/2$	$-1/2$	$-1/2$	$-1/2$	$-1/2$	1

Strong CP problem

$$Q, \tilde{Q}$$

complex scalar, σ

modulus

phase

Axion

Dark Matter

Baryogenesis

(via thermal leptogenesis)

Gives mass to RH neutrinos

$$N_i$$

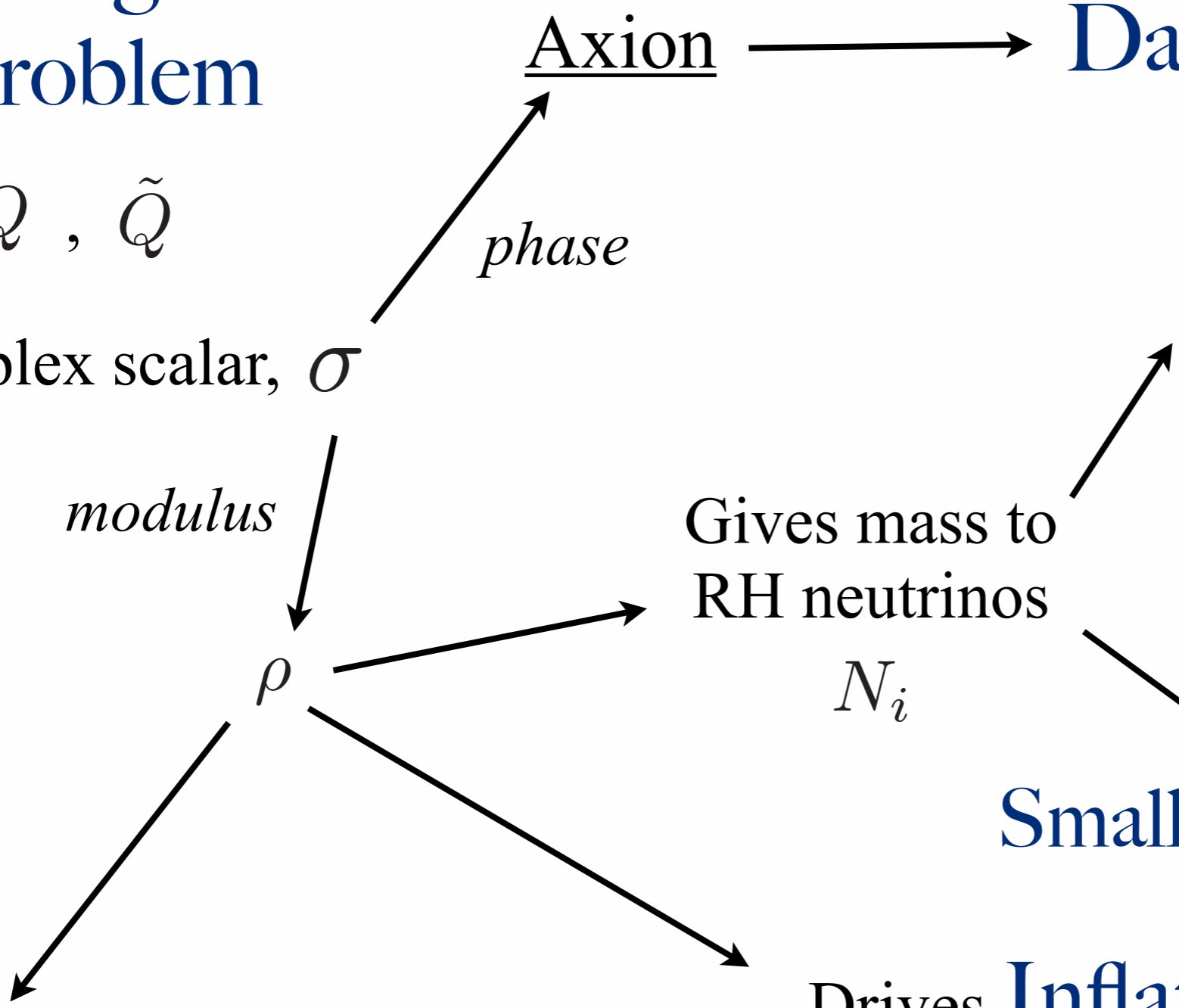
see-saw mechanism

Small neutrino masses

Drives Inflation

(together with the Higgs)
and **reheats** the Universe

Stabilizes the Higgs potential



SM

$$\mathcal{L} \supset - \left[Y_{u_{ij}} q_i \epsilon H u_j + Y_{d_{ij}} q_i H^\dagger d_j + G_{ij} L_i H^\dagger E_j + F_{ij} L_i \epsilon H N_j + \frac{1}{2} Y_{ij} \sigma N_i N_j \right. \\ \left. + y \tilde{Q} \sigma Q + y_{Q_{di}} \sigma Q d_i + h.c. \right],$$

Strong CP problem (and DM)

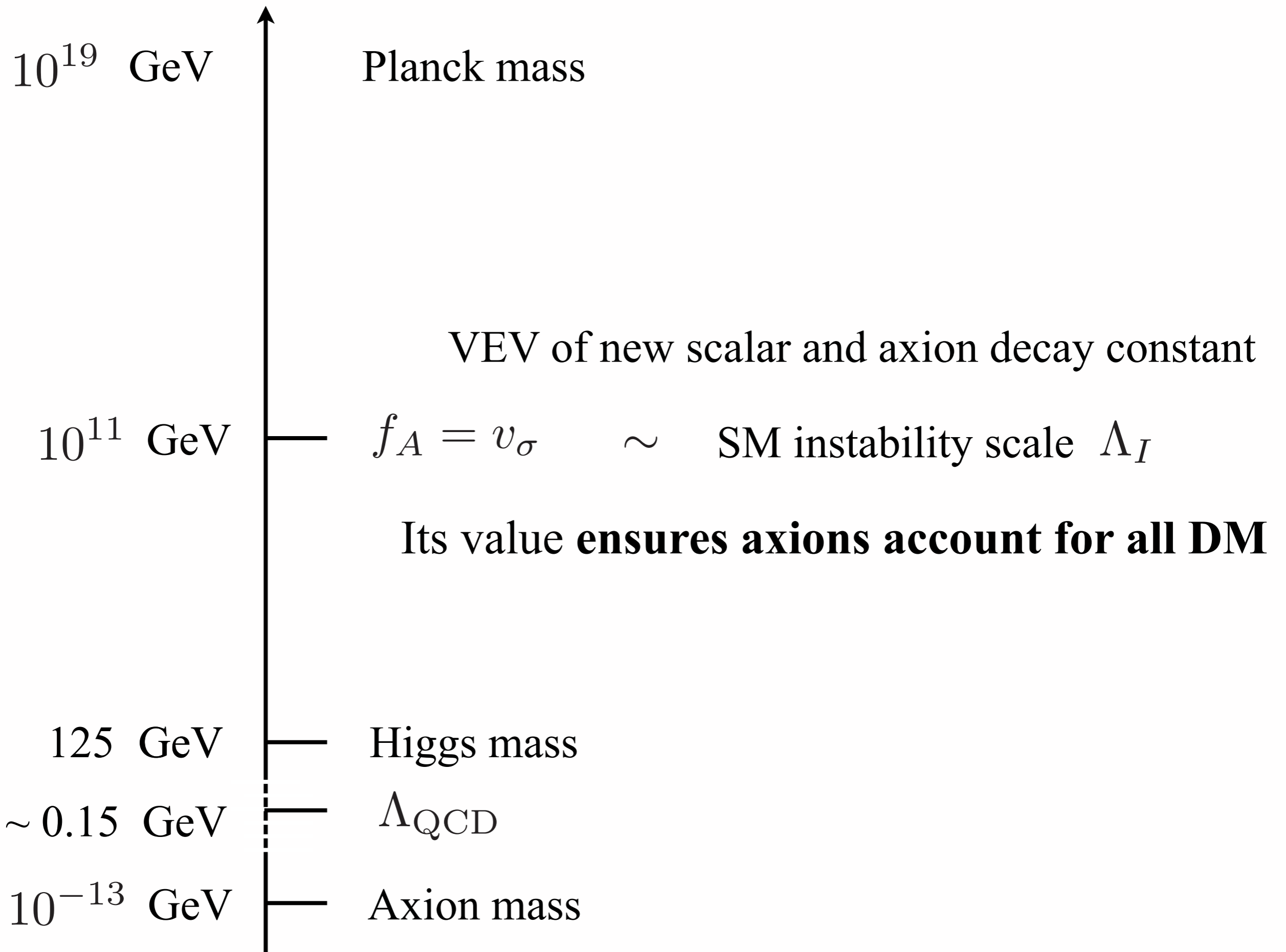


Neutrino masses

$$V(H, \sigma) = \lambda_H \left(H^\dagger H - \frac{v^2}{2} \right)^2 + \lambda_\sigma \left(|\sigma|^2 - \frac{v_\sigma^2}{2} \right)^2 + 2\lambda_{H\sigma} \left(H^\dagger H - \frac{v^2}{2} \right) \left(|\sigma|^2 - \frac{v_\sigma^2}{2} \right)$$



Stability, inflation and reheating



Axion mass: $m_A = (57.2 \pm 0.7) \left(\frac{10^{11} \text{ GeV}}{f_A} \right) \mu\text{eV}$

Borsanyi et al. 2016 from lattice QCD

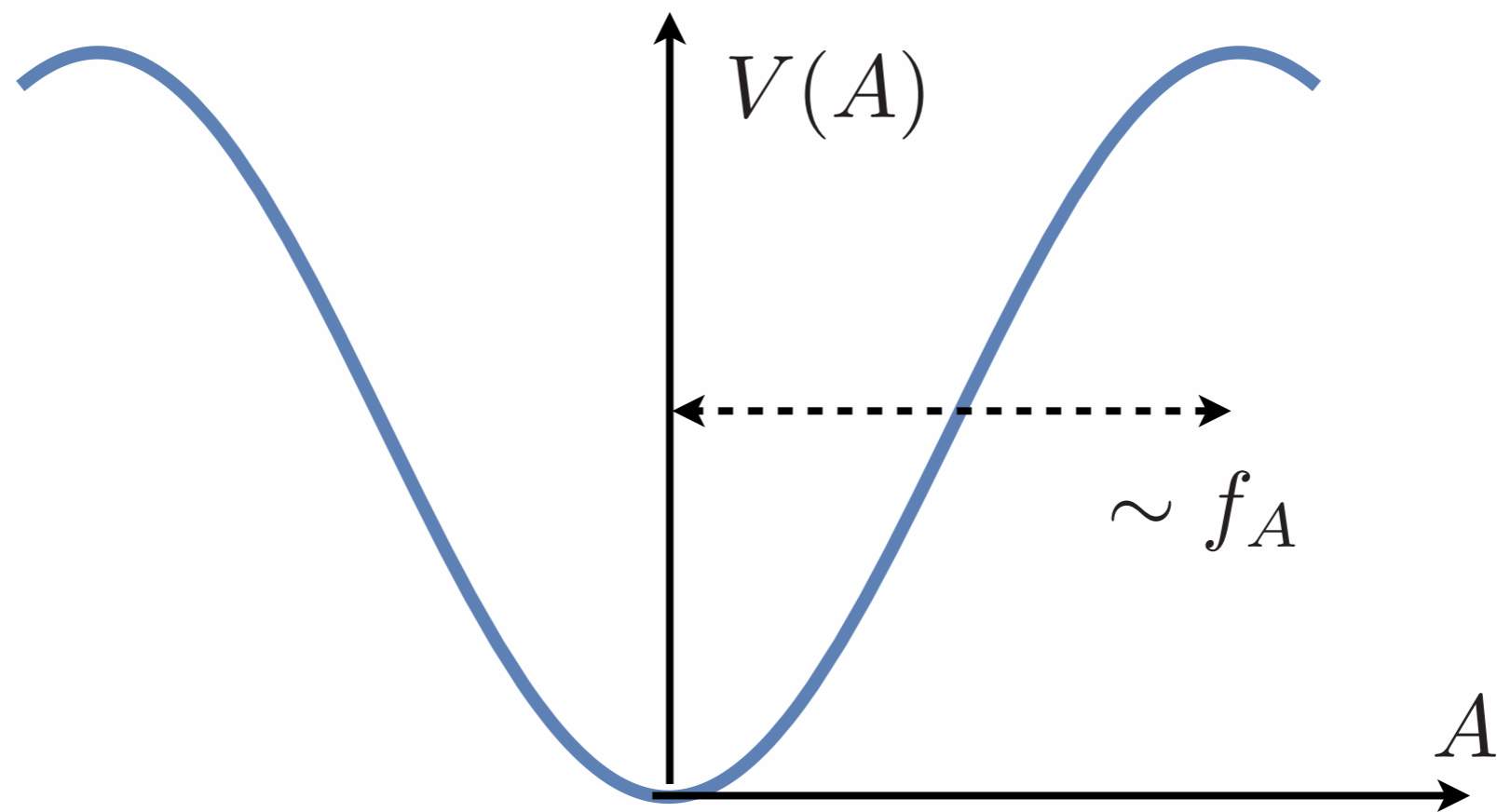
$$\frac{M_{N_i}}{Y} \sim \frac{m_Q}{y} \sim \frac{m_\rho}{\sqrt{\lambda_\sigma}} \sim v_\sigma + \mathcal{O}(v) \sim 10^{11} \text{ GeV}$$

Upper limit on Yukawas Y, y for stability

Typically: $10^{-13} \lesssim \frac{\lambda_\sigma}{5} \lesssim 10^{-10}$ from inflation

The strong CP problem and the axion

$$\mathcal{L} \in \frac{1}{2} \partial_\mu A \partial^\mu A + i \frac{A}{32\pi^2} G\tilde{G} + V(A)$$



The axion potential is generated by non-perturbative QCD physics

KSVZ-like axion

$$\mathcal{L} \in \frac{1}{2} \partial_\mu A \partial^\mu A + i \frac{A}{32\pi^2} G\tilde{G} + V(A)$$

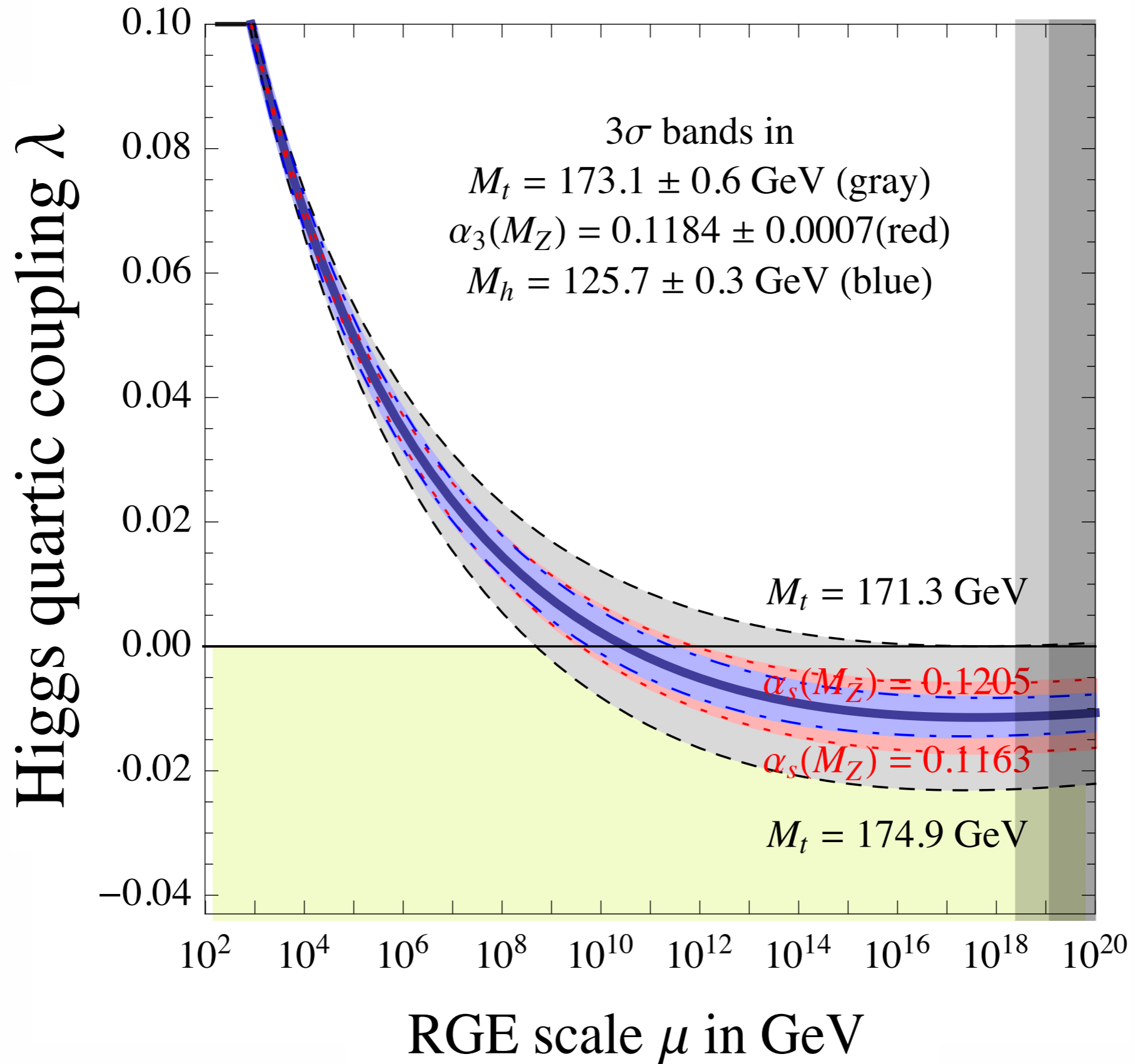
The coupling of the axion to QCD is a dim. 5 operator.

$$\frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma^* + \lambda_\sigma \left(|\sigma|^2 - \frac{v_\sigma^2}{2} \right)^2 + y \tilde{Q} \sigma Q + h.c.$$

$$\sigma \rightarrow e^{i\alpha} \sigma, \quad Q \rightarrow e^{-i\frac{\alpha}{2} \gamma_5} Q, \quad \alpha = A/v_\sigma$$

and integrate out Q and $|\sigma|$ below $v_\sigma = f_A$

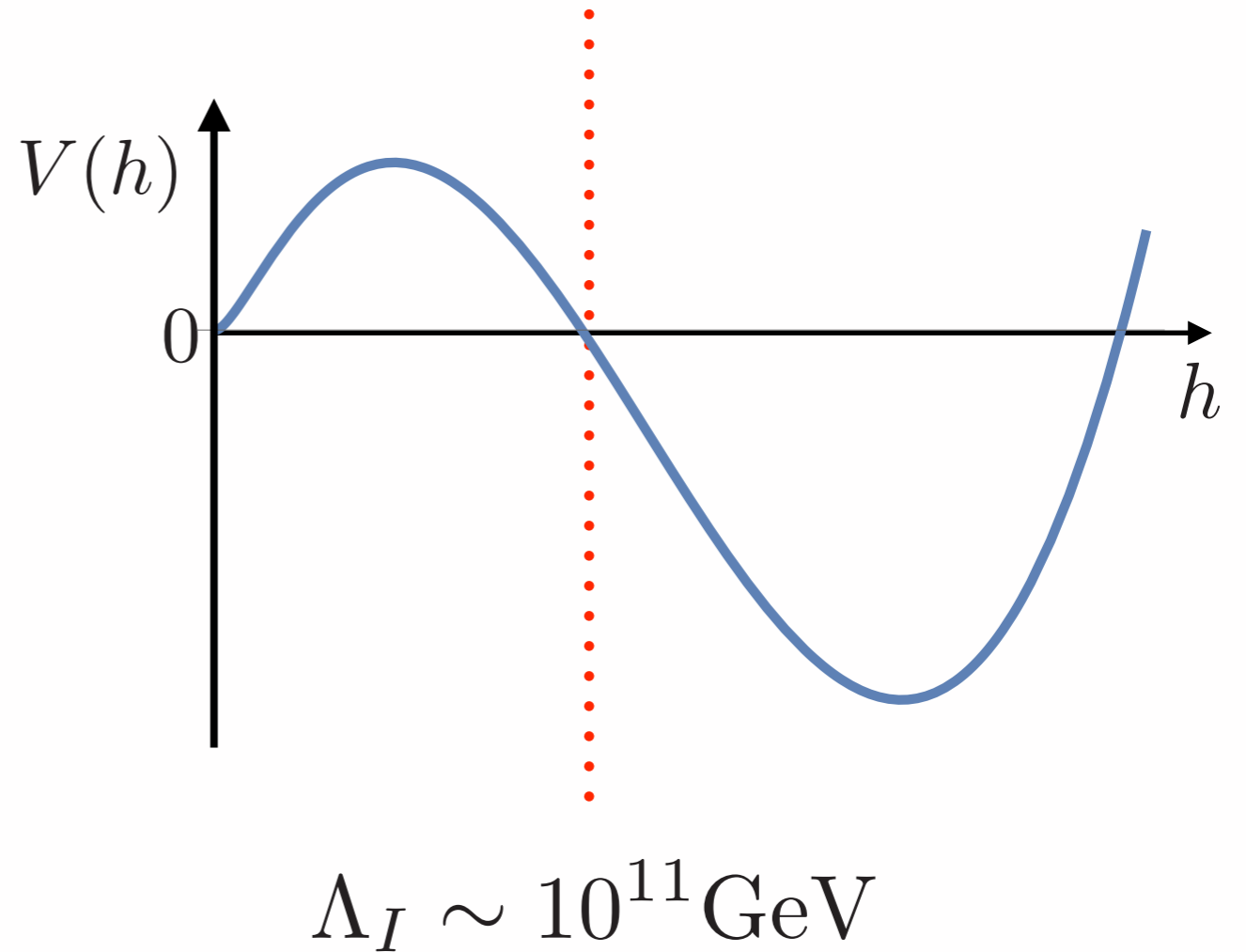
Inflation and the SM instability



Degrassi et al. arXiv:1205.6497

Inflation and the SM instability

$$V(h) \simeq \frac{\lambda_H(h)}{4} h^4$$



Fluctuations during inflaton:

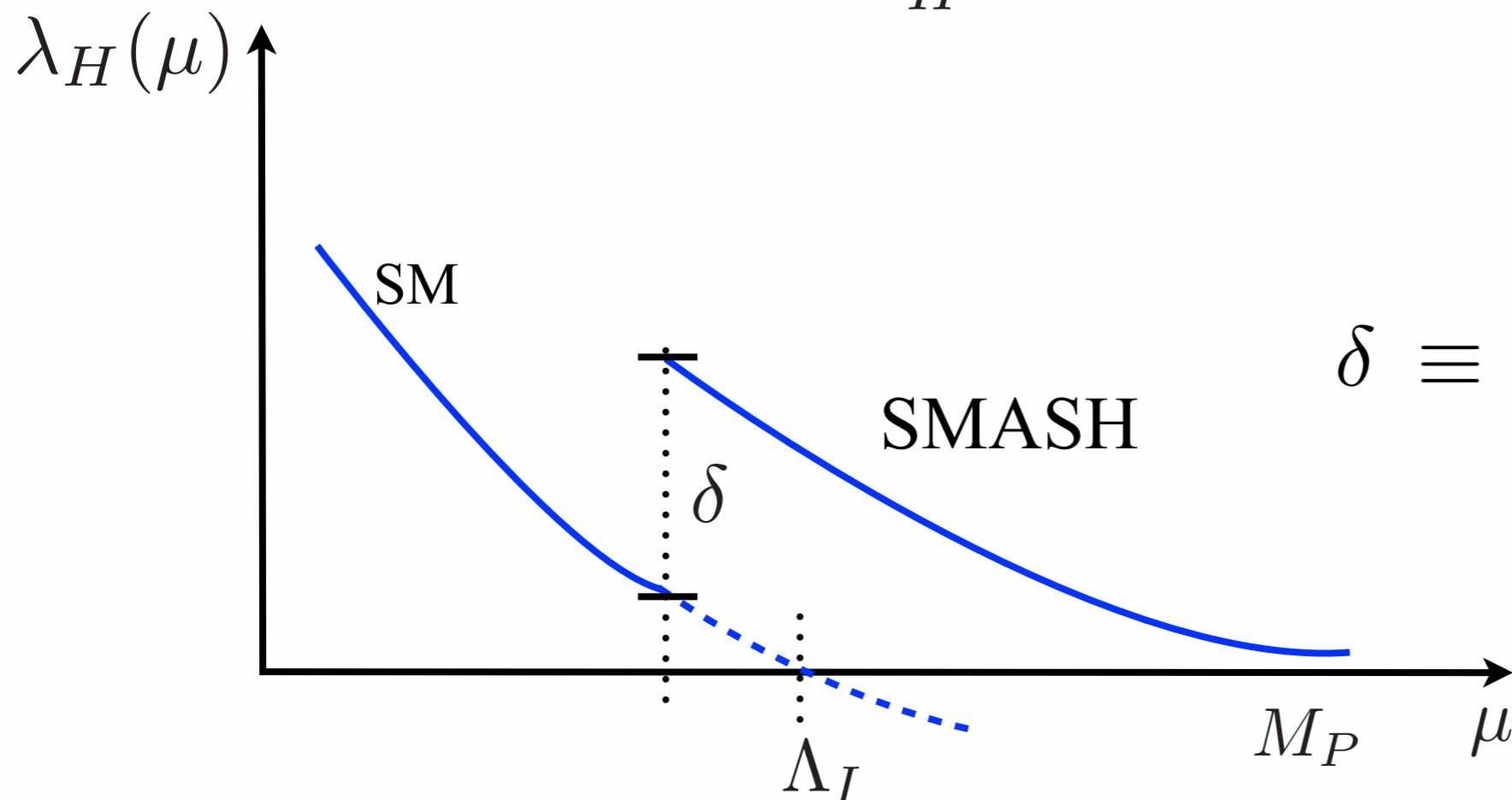
$$\sqrt{\langle h^2 \rangle} \sim \mathcal{H} \sim 10^{-5} M_P \sim 10^{14} \text{ GeV} \gg \Lambda_I$$

Threshold stabilization

$$V(H, \sigma) = \lambda_H \left(H^\dagger H - \frac{v^2}{2} \right)^2 + \lambda_\sigma \left(|\sigma|^2 - \frac{v_\sigma^2}{2} \right)^2 + 2\lambda_{H\sigma} \left(H^\dagger H - \frac{v^2}{2} \right) \left(|\sigma|^2 - \frac{v_\sigma^2}{2} \right)$$

At low energies, below the mass of $|\sigma|$

$$\lambda_H^{(SM)} = \lambda_H - \delta$$



$$\delta \equiv \lambda_{H\sigma}^2 / \lambda_\sigma \sim 10^{-2}$$

Inflation from the Higgs?

$$S \supset - \int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_H H^\dagger H + \dots \right] R$$

$$\tilde{V} \sim \frac{\lambda_H}{\xi_H^2} M_P^4$$

CMB temperature fluctuations $\longrightarrow \xi_H \sim 10^5 \sqrt{\lambda_H} \sim 10^4$

Breaking of perturbative unitarity:

$$\Lambda_U = \frac{M_P}{\xi_H} \sim 10^{14} \text{ GeV} \ll \frac{M_P}{\sqrt{\xi_H}} \sim 10^{16} \text{ GeV}$$

Inflation with the new singlet

$$S \supset - \int d^4x \sqrt{-g} \left[\frac{M^2}{2} + \xi_H H^\dagger H + \xi_\sigma \sigma^* \sigma \right] R,$$

$$\tilde{V} \sim \frac{\lambda}{\xi_\sigma} M_P^4, \quad \xi_\sigma \lesssim 1 \quad \text{and also} \quad \xi_H \lesssim 1$$

$$\lambda_{H\sigma} > 0 \longrightarrow \text{inflaton} = |\sigma|, \quad \lambda = \lambda_\sigma$$

$$\lambda_{H\sigma} < 0 \longrightarrow \text{inflaton} = |\sigma| + \text{small Higgs component},$$
$$\lambda = \lambda_\sigma - \lambda_{H\sigma}^2 / \lambda_H$$

Reheating after inflation

A small Higgs component in the inflaton of SMASH guarantees successful reheating

$$N_{\nu}^{\text{eff}} = 3.04 \pm 0.18$$

from CMB and BAO data

$$\lambda_{H\sigma} > 0, \quad T_R \sim 10^7 \text{ GeV}$$

Axions remain decoupled
from thermal bath

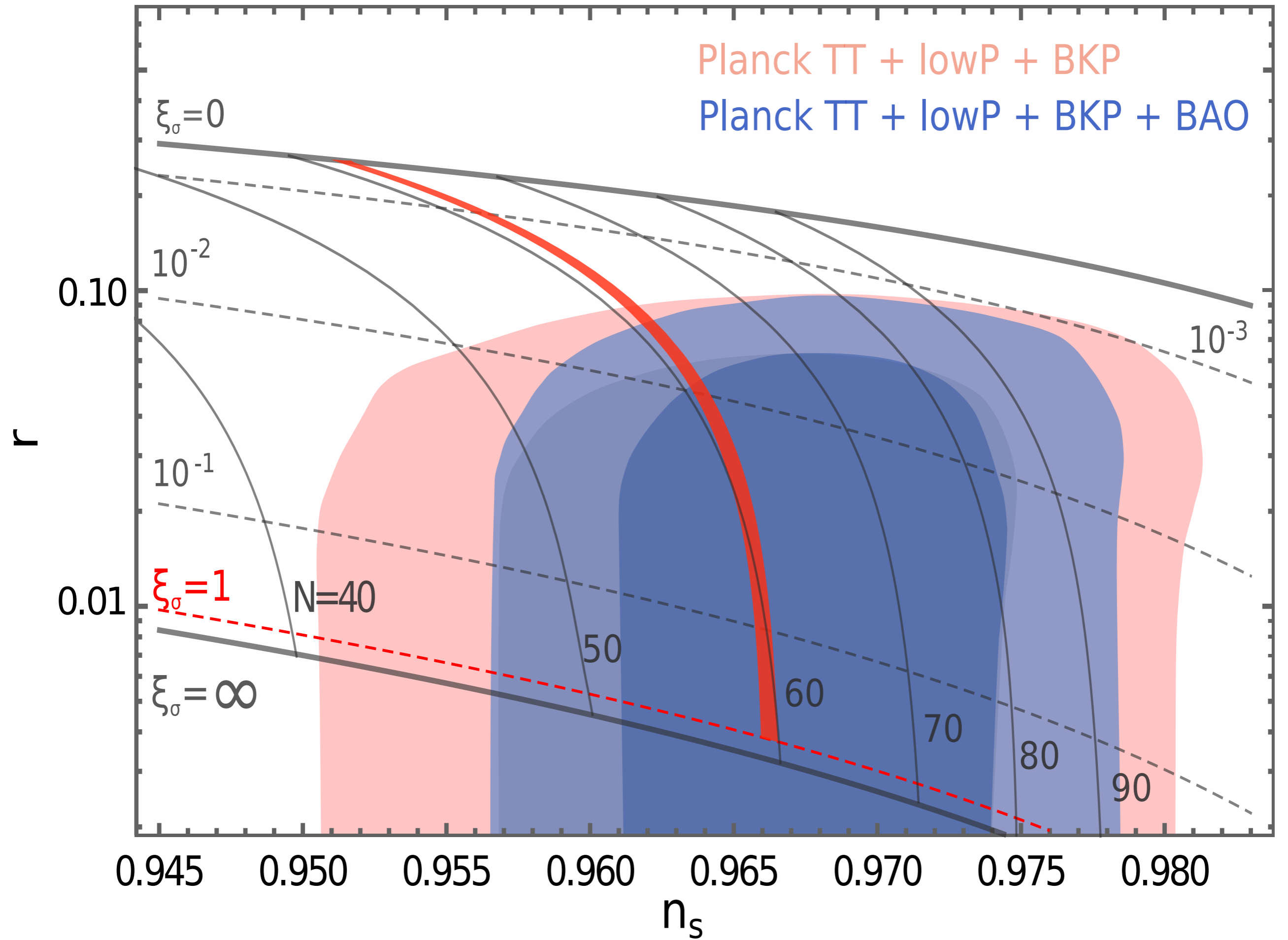
$$\Delta N_{\text{eff}} \sim 1$$

Too much axion radiation

$$\lambda_{H\sigma} < 0, \quad T_R \sim 10^{10} \text{ GeV}$$

$$\Delta N_{\text{eff}} \sim 0.03$$





Predictions for cosmology

$$\text{CMB + unitarity: } 0.04 \lesssim r \lesssim 0.07$$

(CORE, LiteBird, Pixie, CMB S4)

$$5 \times 10^{-13} \lesssim \lambda \lesssim 5 \times 10^{-10}$$

Small non-Gaussianities and isocurvature

$$0.962 \lesssim n_s \lesssim 0.966$$

$$\text{Spectral index running: } \alpha \simeq -7 \times 10^{-4}$$

(21 cm line of neutral Hydrogen)

Axion dark matter

$\lambda_{H\sigma} < 0$: SSB of PQ symmetry after inflation

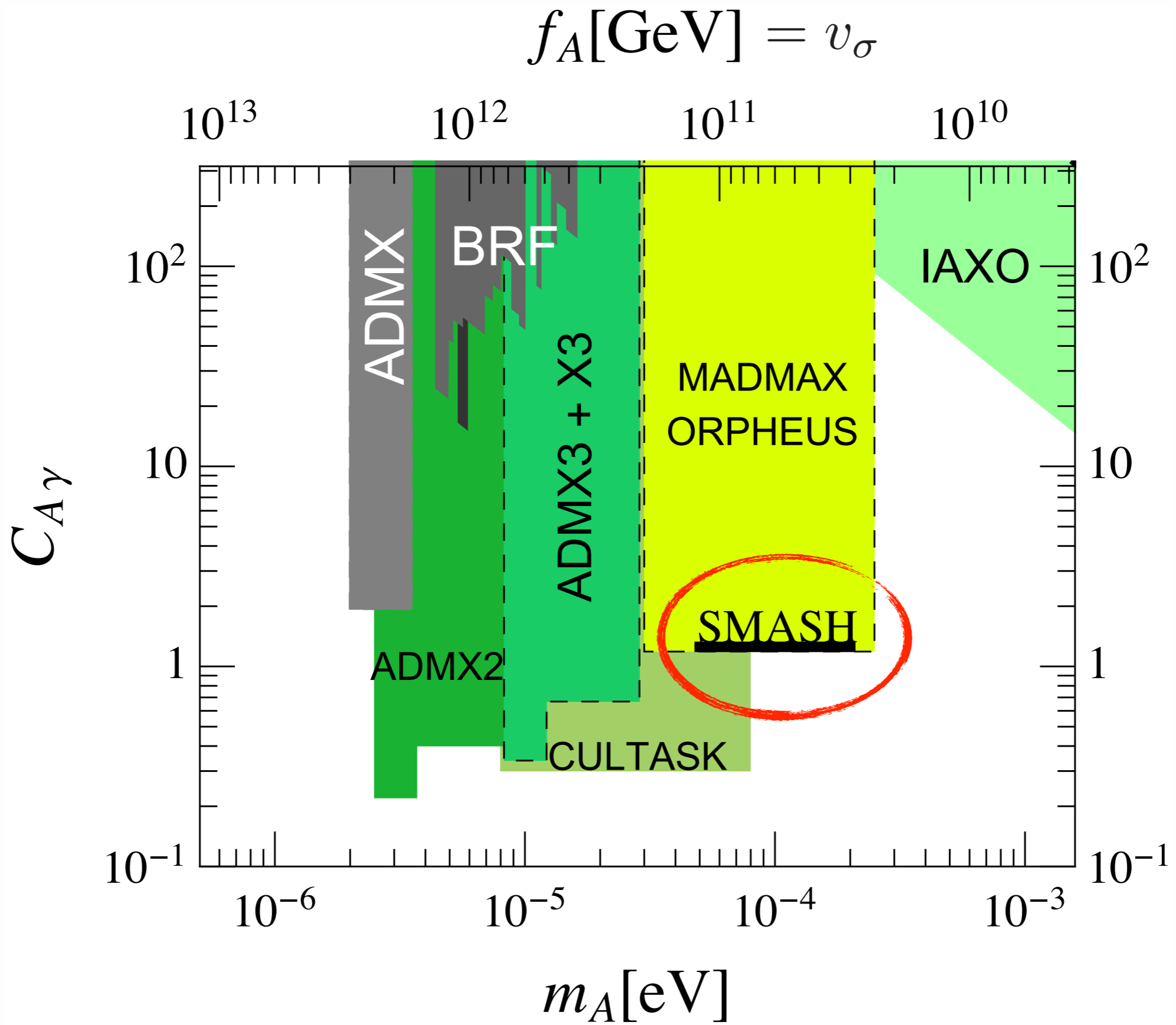
vacuum misalignment: $\ddot{A} + 3\mathcal{H}\dot{A} + m_A^2 A = 0$

and decay of Peccei-Quinn strings

$$3 \times 10^{10} \text{ GeV} \lesssim v_\sigma \lesssim 1.2 \times 10^{11} \text{ GeV},$$



$$50 \mu\text{eV} \lesssim m_A \lesssim 200 \mu\text{eV}$$



SMASH

Solves *the strong CP problem* with a *KSVZ-like axion*,

explains:

the nature of dark matter (with *the axion*),

the smallness of neutrino masses (through *the see-saw*),

baryogenesis (via *leptogenesis*)

and

gives a candidate for *primordial inflation*.

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