

# Dynamically determining the weak scale from inflation [JCAP, 1701.09167]

**Tevong You**

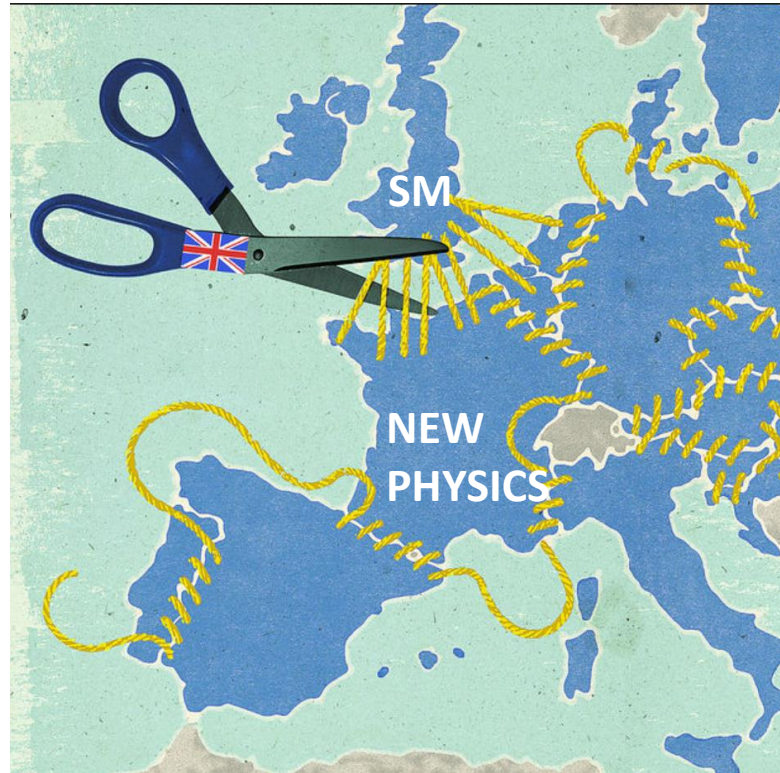


# Naturalness

- Usual arguments are old and well known but **still as valid** now as back when they were made
  - After all, our understanding of QFT has not changed so why should the conclusions drawn from it be any different now?
- But what about null experimental data? The issue is with the **solutions**, not the **hierarchy problem**, which is as motivated as ever
  - Philosophical waffling about unknown priors for fine-tuning and all that matters are observables etc. are irrelevant to the fact that there is something strange going on
  - c.f. Newton:  $m_{\text{inertial}} = m_{\text{gravity}}$
- Interesting to consider alternative solutions assuming a desert above the weak scale

# Naturalness

- Assume new physics is decoupled (SMEXIT)



- Can we still have a naturally light Higgs mass?

# Cosmological Relaxation

P. W. Graham, D. E. Kaplan and S. Rajendran,  
[arXiv:1504.07551]

L. F. Abbott, Phys. Lett. B 150  
(1985) 427

- Higgs mass is naturally at large cut-off  $M$

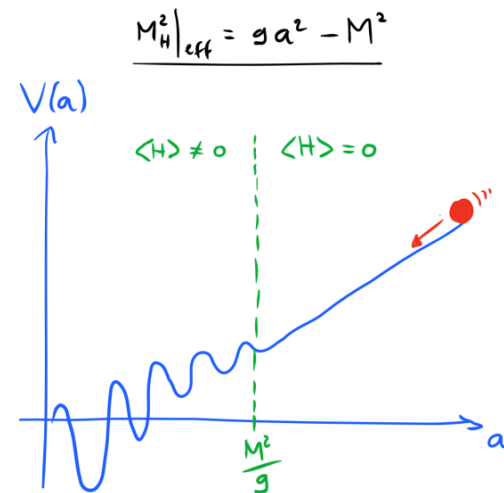
$$V_{\text{soft}}(a) \simeq (ga - M^2)|h|^2 + gM^2a + \dots$$

- Axion-like particle  $a$  protected by shift symmetry, explicitly broken through technically-small parameter  $g$

- Scans an effective Higgs mass

- Barriers switch on after EWSB

$$V_{\text{cos}}(a) = \Lambda_G^4 \cos(a/f) \quad \Lambda_G^4 \equiv \Lambda_G^{4-n} v^n$$



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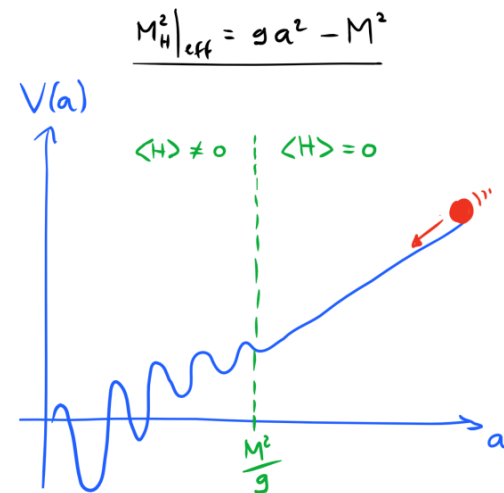
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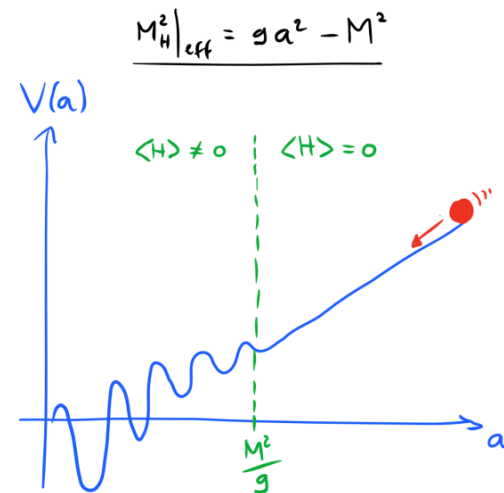
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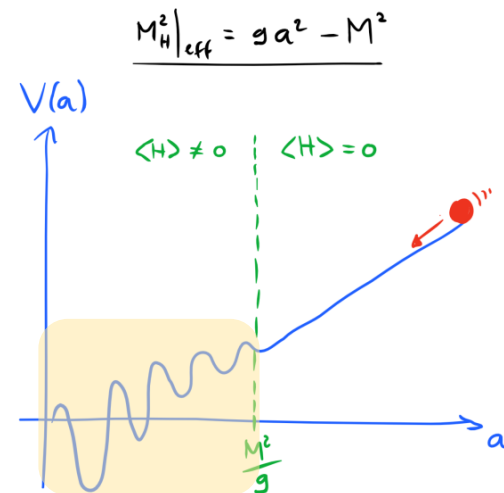
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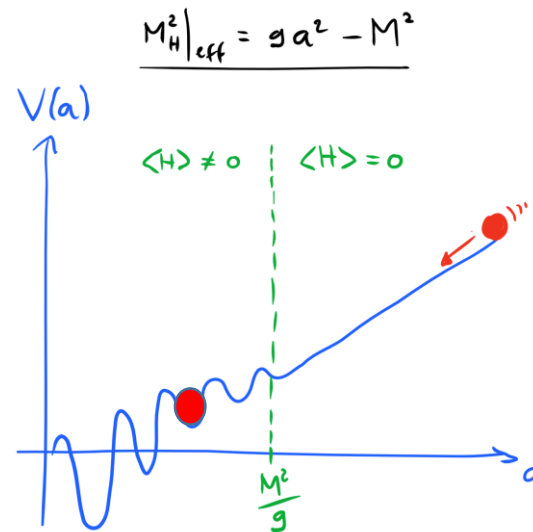
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$$gM^2 \sim \frac{\Lambda_G^{4-n} v^n}{f_\phi}$$

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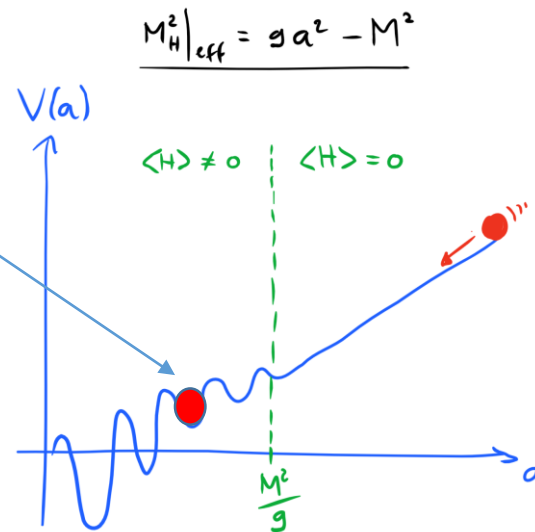
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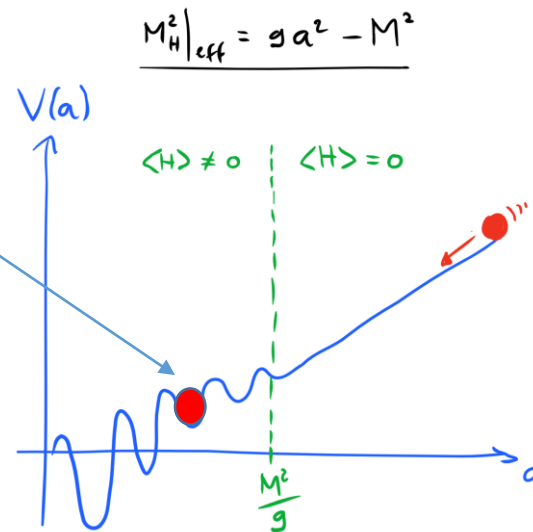
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Constraints:  $H < v$ , classical rolling vs quantum, inflaton energy density dominates relaxation, etc.

Very small  $g$  and natural scanning range lead to super-planckian field excursions, exponential e-foldings...

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# Relaxation Models

(apologies for lack of references)

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- **n=1 models** Graham et al [arXiv:1504.07551]
  - G=QCD: Theta problem
  - G=QCD with inflaton coupling: not technically natural
  - New gauge group G: new physics at weak scale + coincidence problem
- **n=2 models** Espinosa et al [arXiv:1506.09217]
  - G can be at higher scales, raises M cut-off too
  - Requires second scalar to relax relaxation barriers: double-scanning mechanism
- **n=0 models** Hook and Marques-Tavares [arXiv:1607.01786], TY [arXiv:1701.09167]
  - More promising, make use of axial gauge coupling

$$\mathcal{L} = \frac{1}{32\pi^2} \frac{a}{f} \epsilon^{\mu\nu\rho\sigma} \text{Tr} G_{\mu\nu} G_{\rho\sigma}$$

- (see also Toyokazu Sekiguchi talk for use of dissipation in n=1,2 models)

# Relaxation backreaction on inflation

TY [arXiv:1701.09167]

- Minimal relaxion setup, **no v-dependence in relaxion sector**

$$\mathcal{L} \supset (M^2 - g\phi) |h|^2 + gM^2\phi + \dots + \Lambda_G^4 \cos\left(\frac{\phi}{f_\phi}\right) - \frac{\alpha_D}{f_D} \phi F_{\mu\nu} \tilde{F}^{\mu\nu},$$

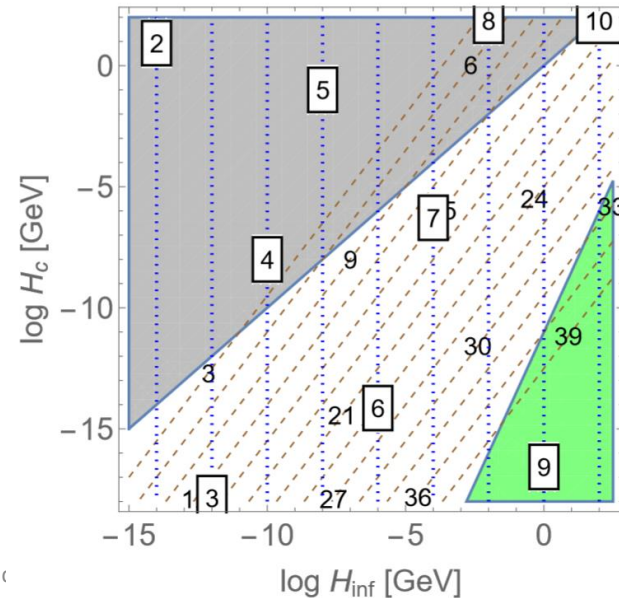
- Backreaction instead ends inflation

- e.g. Inflation supported by electroweak dissipation  $\mathcal{L} \supset -\frac{\alpha}{f} \sigma F_{\mu\nu} \tilde{F}^{\mu\nu}$

$$\ddot{\sigma} + 3H\dot{\sigma} + V'_\sigma(\sigma) = -I \frac{\alpha}{f} \left(\frac{H}{\xi}\right)^4 e^{2\pi\xi}, \quad \xi \equiv \frac{\alpha}{2f} \frac{\dot{\sigma}}{H}$$

See e.g. Anber and Sorbo 0908.4089

- Hubble falls
- Dark dissipation increases
- Relaxion loses KE and is trapped



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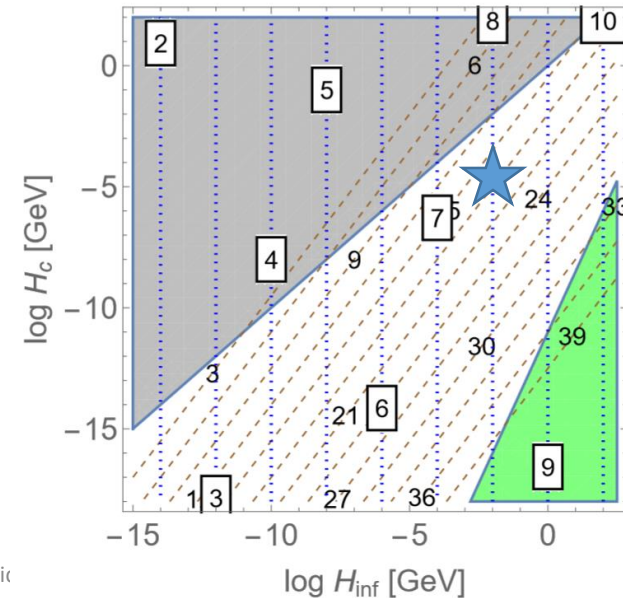
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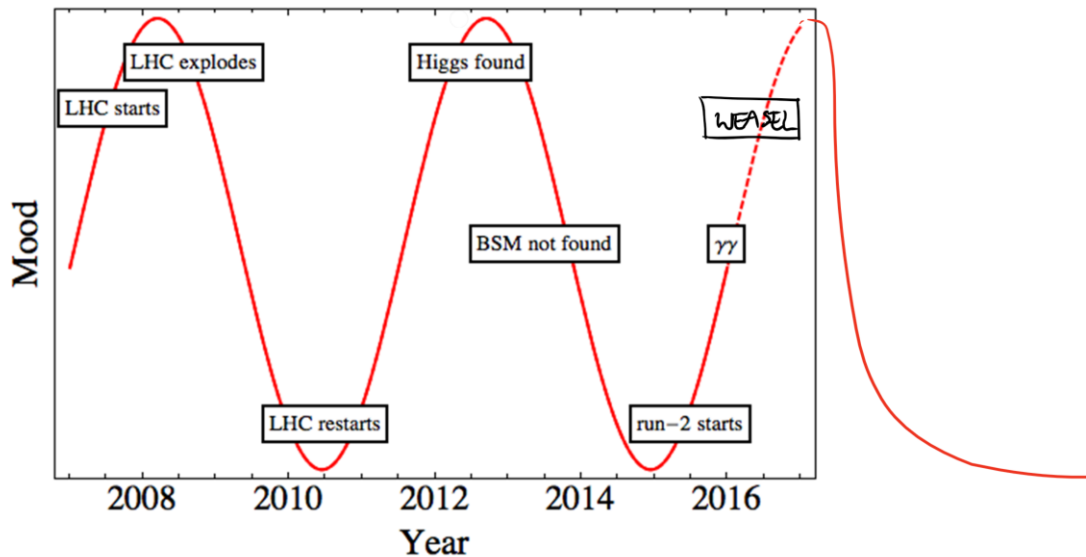
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	$M$	$g$	$H_I$	$H_c$	$N_e$	$\Lambda_G$	$f_\phi$	$f_D/\alpha_D$
$\sim$ [GeV]	$10^8$	$10^{-11}$	$10^{-2}$	$10^{-5}$	$10^{18}$	$10^{3.5}$	$10^9$	$10^{15}$



# Conclusion

<http://resonaances.blogspot.com.es/2016/01/do-or-die-year.html>



- Perhaps we are missing something in the way we think about the hierarchy problem?
- Models of natural weak scale with high cut-off are interesting to explore, realistic or not
- Motivates searching for alternative ways Higgs vev backreaction could dynamically select weak scale