

PRESENTATION

**Does the detection of primordial
gravitational waves exclude
low energy inflation?**

Tomohiro Fujita (Stanford/Kyoto)

Based on [arXiv:1608.04216](https://arxiv.org/abs/1608.04216), [1705.01533](https://arxiv.org/abs/1705.01533)
w/ Dimastrogiovanni(CWRU) & Fasiello(Stanford);
Namba(McGill)&Tada(IAP)

In prep w/Komatsu&Agrawal(MPA);
Thone(Oxford),Hazumi(KEK),Katayama(IPMU)
Komatsu&Shiraishi(Kagawa)

22nd/June/2017@PASCOS



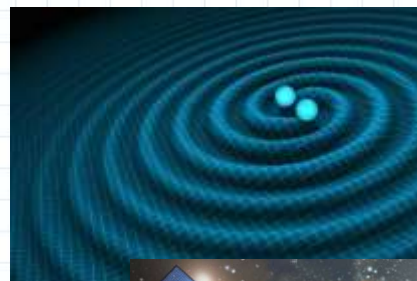
京都大学
KYOTO UNIVERSITY



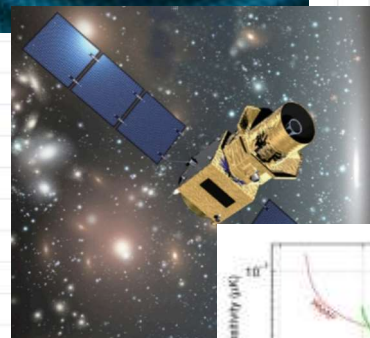
Era of PGW



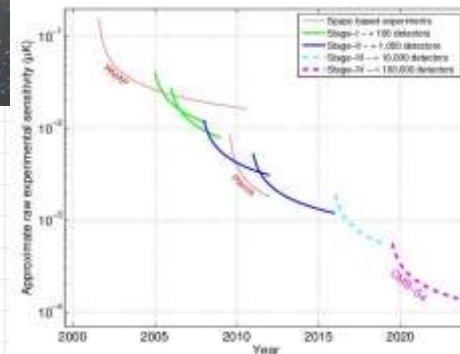
- GW exists!
aLIGO detected GW from BH binary

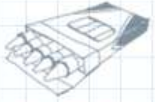


- Primordial GW soon?
PGW observed by the CMB B-mode



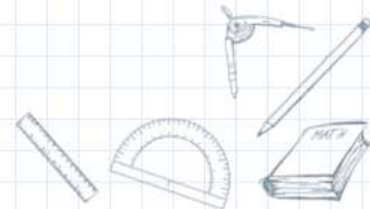
- New Obs aiming $r \geq 10^{-3}$
e.g. CMB-S4(US), LiteBIRD (Japan), etc..





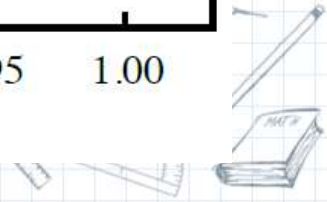
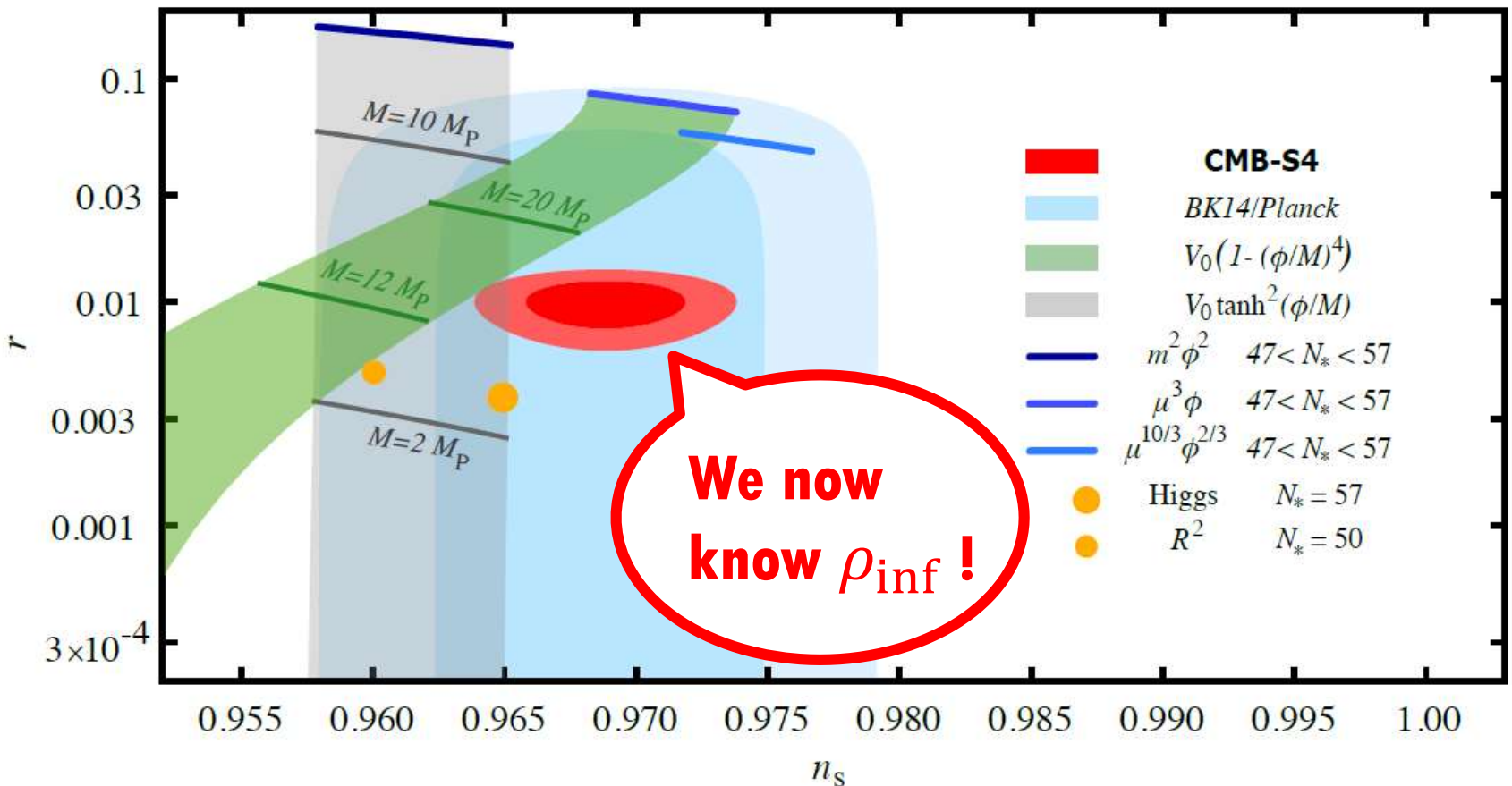
Simple relationship

$$\rho_{\text{inf}}^{1/4} \approx 6 \times 10^{15} \text{ GeV} \left(\frac{r}{0.001} \right)^{1/4}$$





What if r is detected?

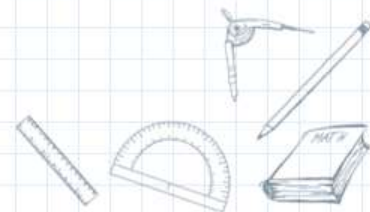




Simple relationship

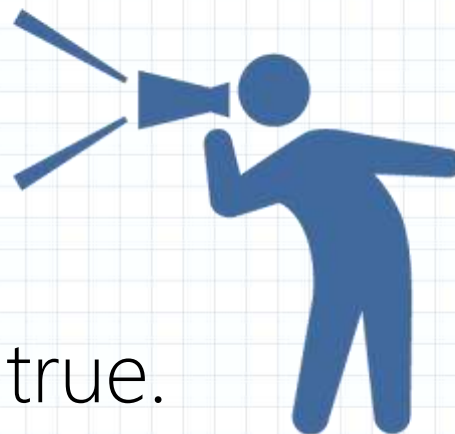
$$\rho_{\text{inf}}^{1/4} \approx 6 \times 10^{15} \text{ GeV} \left(\frac{r}{0.001} \right)^{1/4}$$

- Is this relation **robust**?
- Detection of B-mode **kills** low energy inflation models?

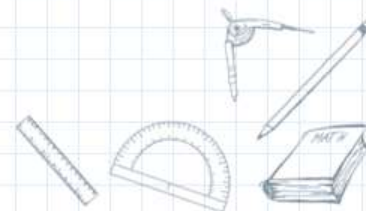


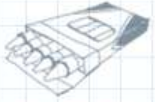


No!



That's not necessarily true.



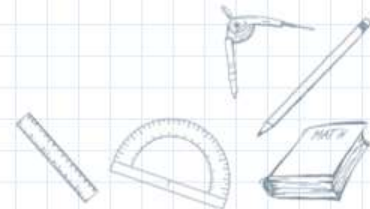


This simple relationship

$$\rho_{\text{inf}}^{1/4} \approx 6 \times 10^{15} \text{ GeV} \left(\frac{r_{\text{obs}}}{0.001} \right)^{1/4}$$

is derived under

2 Assumptions





2 assumptions

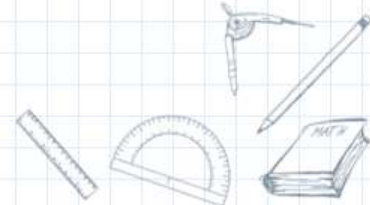
[Creminelli et al.(2014) PRL113,231301;
T.F., X.Gao & J.Yokoyama. JCAP1602.014]

① **GW is described by GR**

2nd order tensor Lagrangian is same as GR.

② **No dominant source effect**

Inhomogeneous solution of EoM is insignificant.





2 assumptions

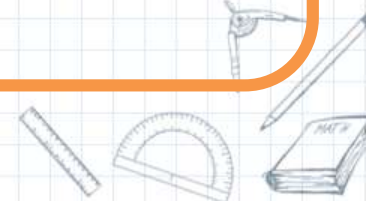
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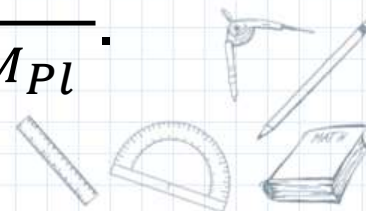
2nd assumptions

② **No dominant source effect**

Inhomogeneous solution of EoM is insignificant.

We use the homogeneous solution

$$\left[\partial_{\tau}^2 + k^2 - \frac{2}{\tau^2} \right] ah_k = 0, \quad ah_k = \frac{2}{\sqrt{2k}} \frac{e^{-ik\tau}}{M_{Pl}}$$





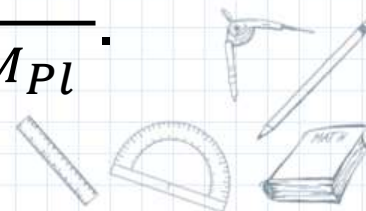
2nd assumptions

② No dominant source effect

Inhomogeneous solution of EoM is insignificant.

What if we have a significant **source term**??

$$\left[\partial_{\tau}^2 + k^2 - \frac{2}{\tau^2} \right] ah_k = S_k, \quad ah_k = \frac{2}{\sqrt{2k}} \frac{e^{-ik\tau}}{M_{Pl}}$$

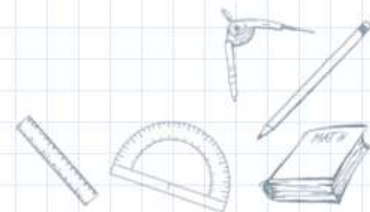




Axion-SU(2) model

- Larger GW than h^{vac} can be produced

$$r_{\text{obs}} = r_{\text{vac}} + r_{\text{add}}$$

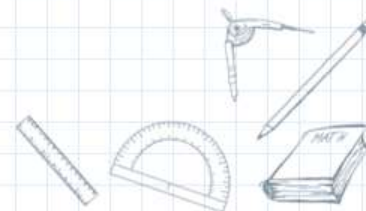


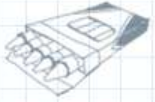


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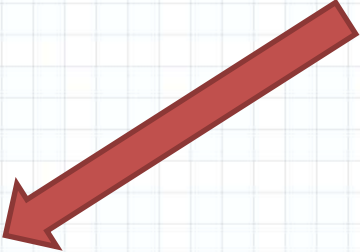


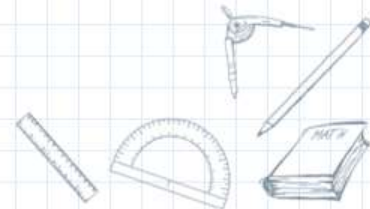


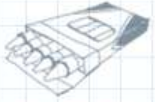
Axion-SU(2) model

- Larger GW than h^{vac} can be produced

$$r_{obs} = r_{vac} + r_{add}$$


$$~~r_{obs} = r_{vac} \approx 10^{-3} \left(\frac{H_{inf}}{10^{13} \text{GeV}} \right)^2~~$$

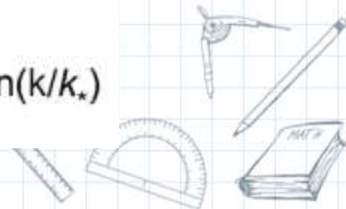
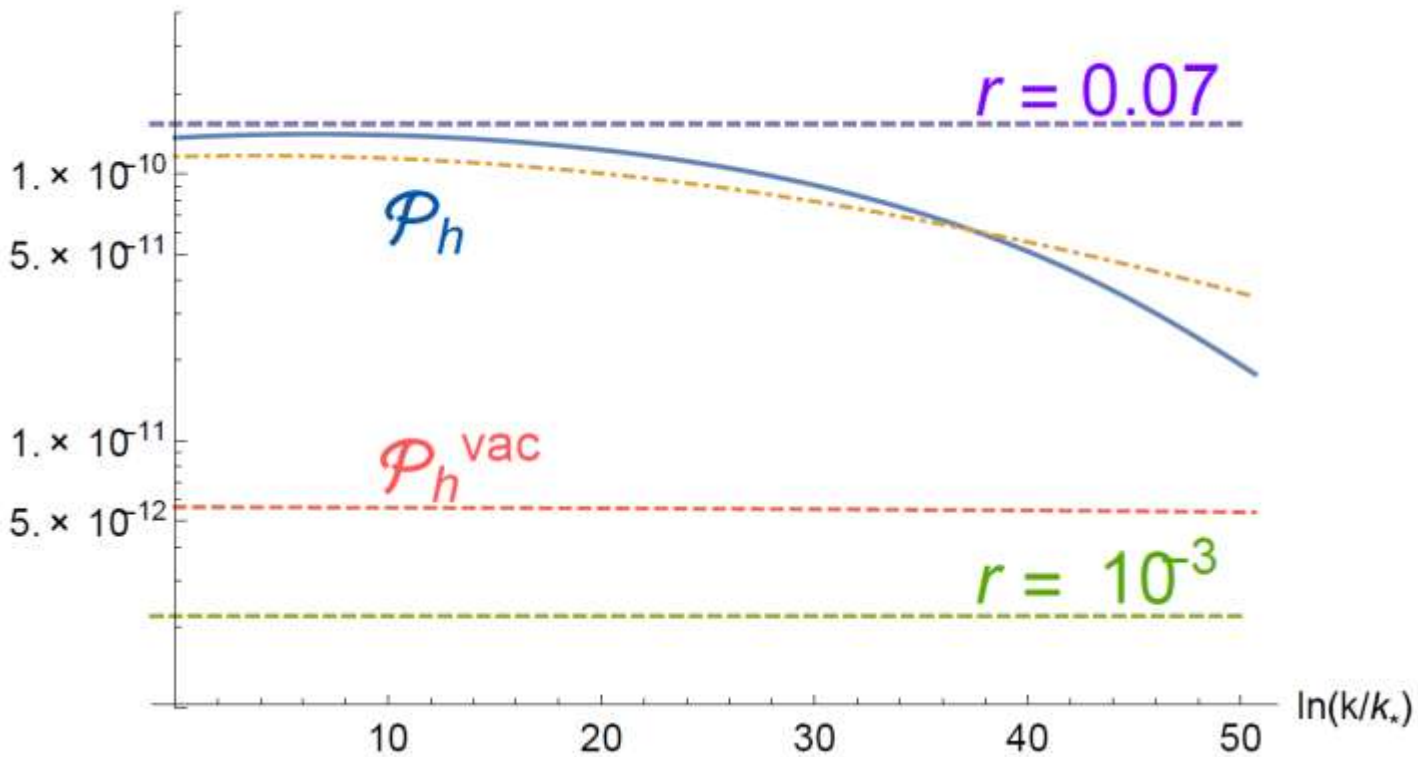




Result

○ $H_{inf} \approx 1 \times 10^{13} \text{ GeV}$

✗ $H_{inf} \approx 6 \times 10^{13} \text{ GeV}$

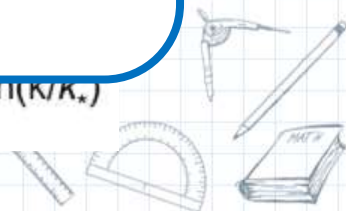
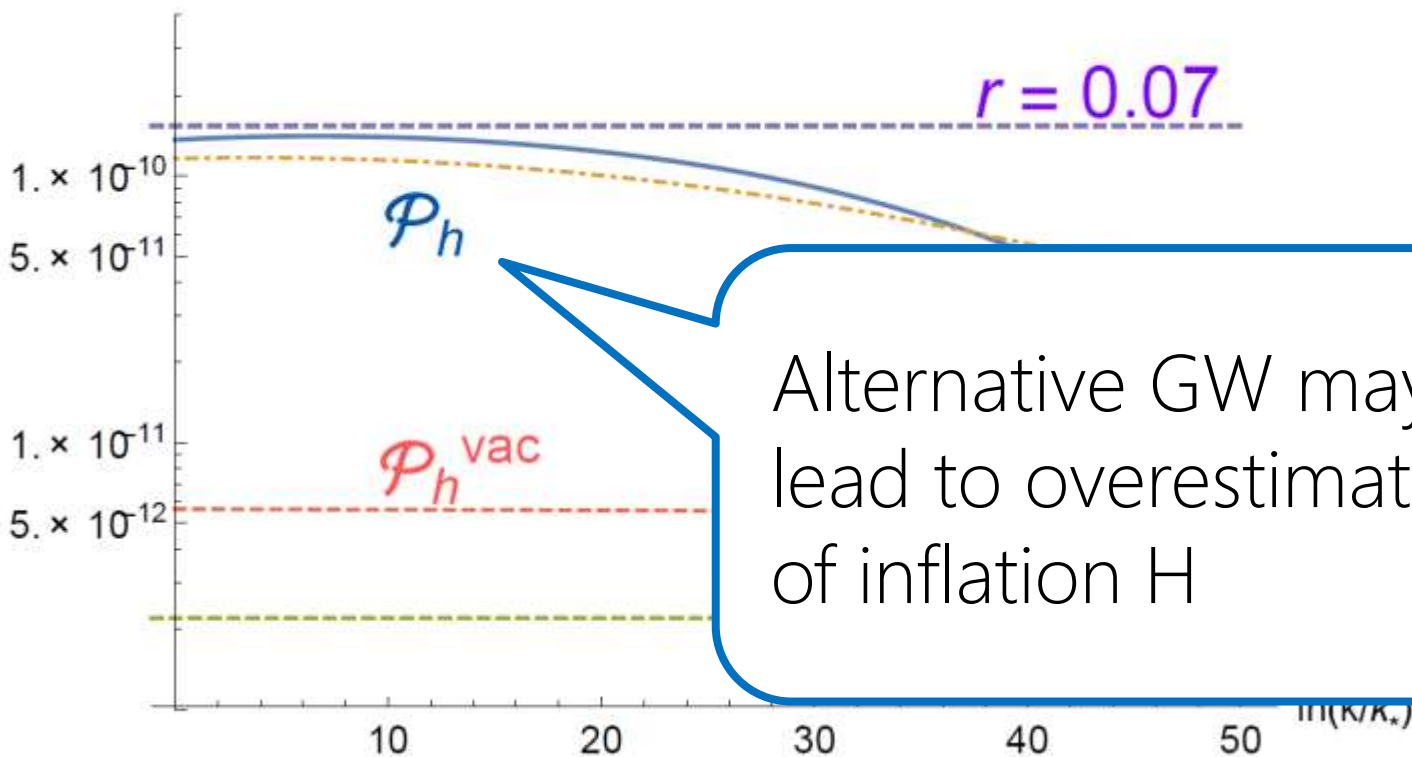




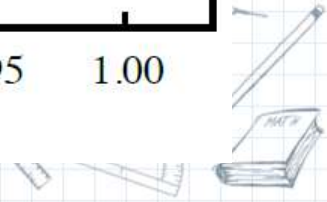
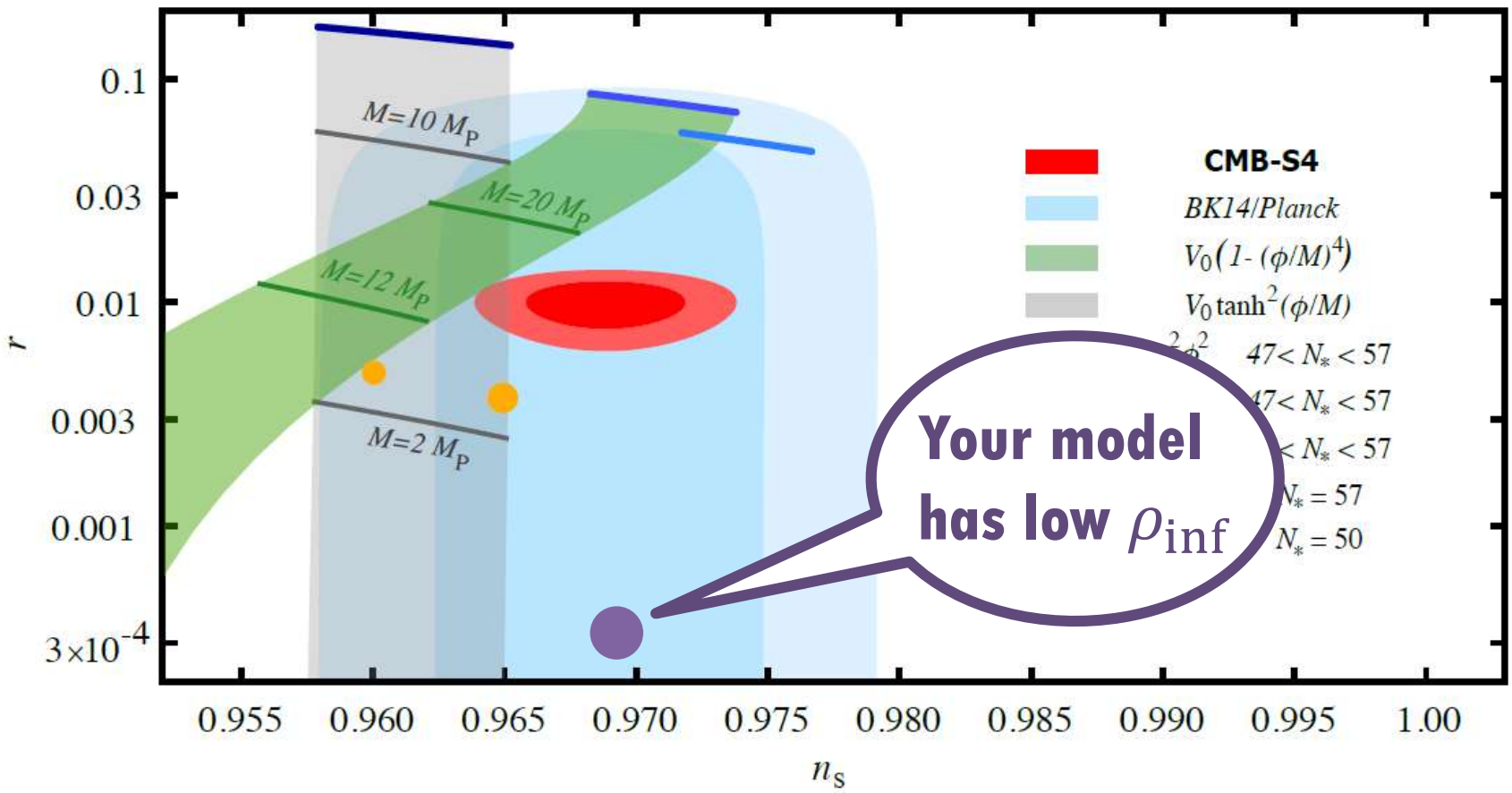
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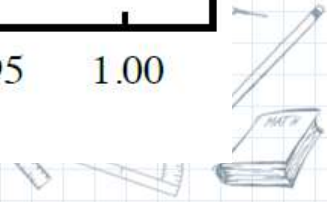
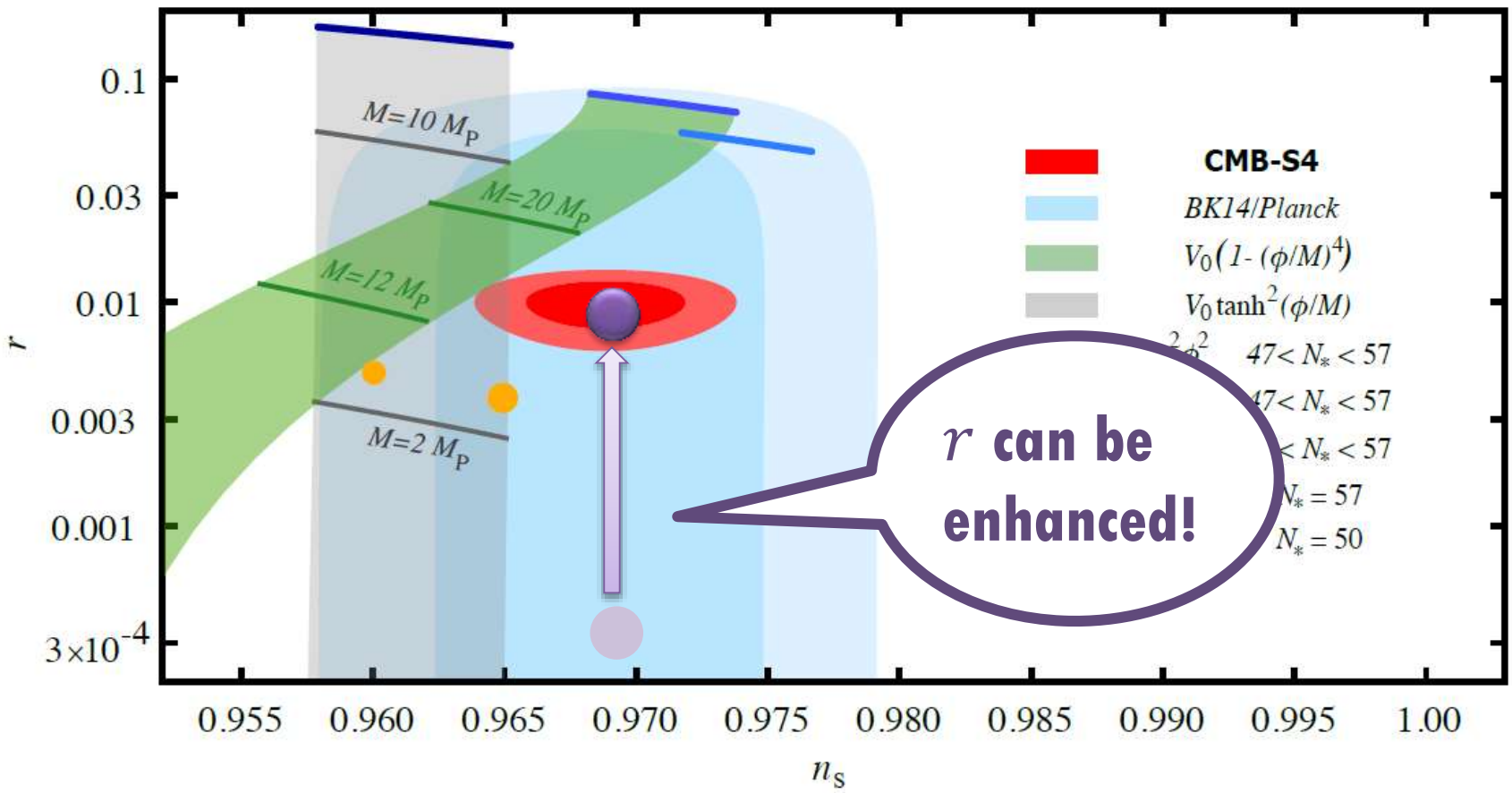
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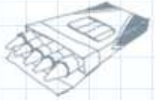


Low ρ_{inf} model can be rescued!



Low ρ_{inf} model can be rescued!





Our scenario

GW version of **Curvaton** mechanism

[Enqvist&Sloth(2002),
Lyth&Wands(2002),
Moroi&Takahashi(2002)]

$$\mathcal{L}_{\text{inflaton}} + \mathcal{L}_{\text{spectator}}$$

$$\begin{array}{ccc}
 \rho_{\text{inf}} & \gg & \rho_{\text{spec}} \\
 \downarrow H_{\text{inf}} & & \downarrow \text{Source} \\
 \mathcal{P}_{\text{GW}}^{\text{vac}} & \ll & \mathcal{P}_{\text{GW}}^{\text{spec}}
 \end{array}$$

\mathcal{L}_{inf} is arbitrary and responsible for ζ generation.

$\mathcal{L}_{\text{spec}}$ is added just to produce GW during inf.





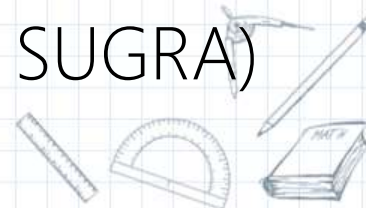
How's it work?

Adding **axion-SU(2)** gauge spectator sector

$$\begin{aligned} \mathcal{L} = & \mathcal{L}_{\text{inflaton}} \\ & + \frac{1}{2} (\partial\chi)^2 - \mu^4 \left(\cos \frac{\chi}{f} + 1 \right) \\ & - \frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} - \frac{\lambda}{4f} \chi F_{\mu\nu}^a \tilde{F}^{a\mu\nu} \end{aligned}$$

Very **well motivated terms** in HEP (e.g. String, SUGRA)

[cf. Chromo-natural inflation: Adshead&Wyman(2012)]





How's it work?

Adding axion-SU(2) gauge spectator sector

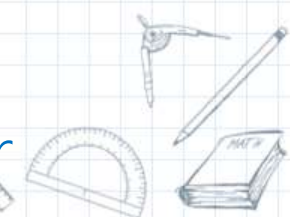
$\mathcal{L} = \mathcal{L}_{\text{inflaton}}$ Inflaton sector → \mathcal{P}_ζ

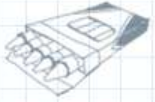
Decoupled

$$\begin{aligned}
 & + \frac{1}{2} (\partial\chi)^2 - U(\chi) \\
 & - \frac{1}{4} FF - \frac{\lambda}{4f} \chi F \tilde{F}
 \end{aligned}$$

→ $\mathcal{P}_{\text{GW}}^{\text{spec}}$

Axion- SU(2) gauge spectator sector



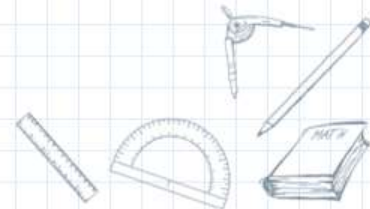


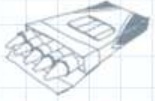
Why $SU(2)$?

SVT Decomposition Theorem:

At the 1st order cosmological perturbation, scalar, vector and tensor are decoupled.

$$\delta S, \delta V_i \xrightarrow{\text{Source}} \delta T_{ij}$$

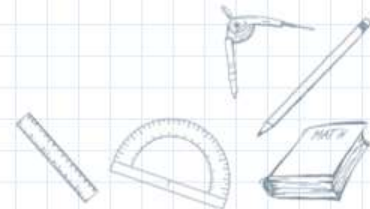
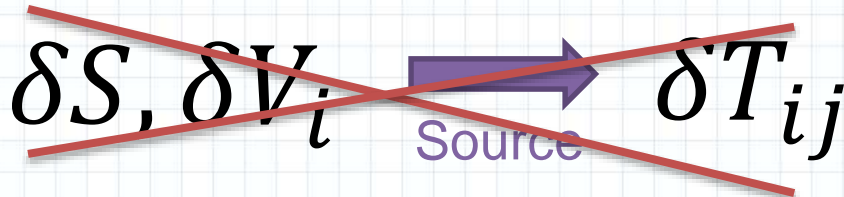




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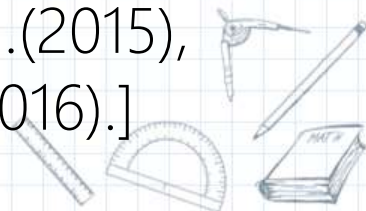


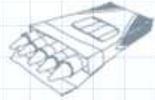
Why SU(2)?

Using the 2nd order pert., GW can be sourced.
But it's hard to generate $\mathcal{P}_{\text{GW}}^{\text{spec}} \gg \mathcal{P}_{\text{GW}}^{\text{vac}}$.

$$\partial_i \delta S \partial_j \delta S, \delta V_i \delta V_j \xrightarrow{\text{Source}} \delta T_{ij}$$

[Biagetti et al.(2013), Mukohyama et al.(2014), TF et al.(2015),
Ferreira et al.(2015), Choi et al.(2015), Namba et al.(2016).]





Way Out

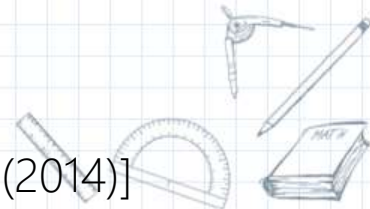
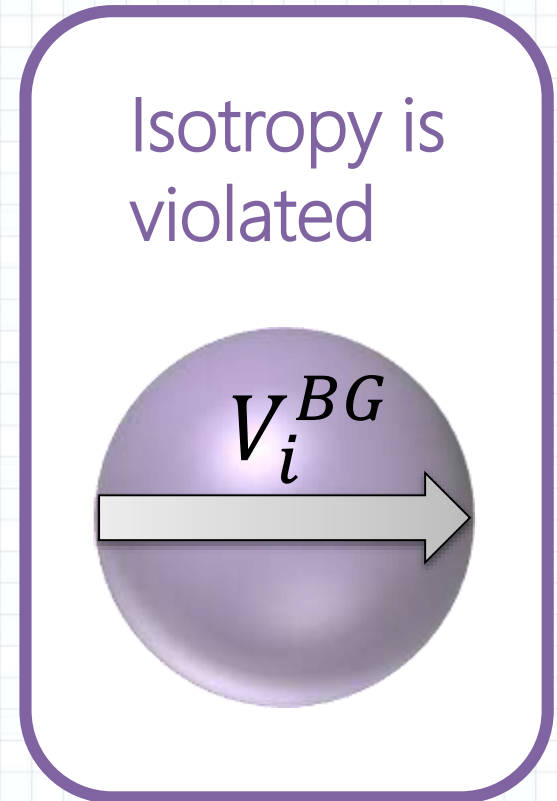
Background vector field V_i^{BG} helps.

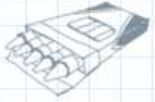
$$V_i^{BG} \delta V_j \xrightarrow{\text{Source}} \delta T_{ij}$$

CMB says the universe is isotropic.

U(1) gauge \rightarrow Anisotropic BG

SU(2) gauge \rightarrow Isotropic BG is **Attractor**.





Way Out

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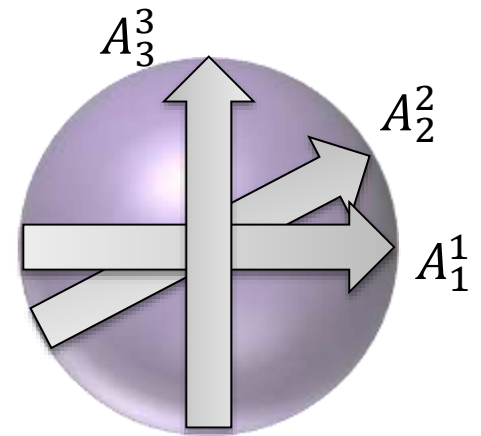
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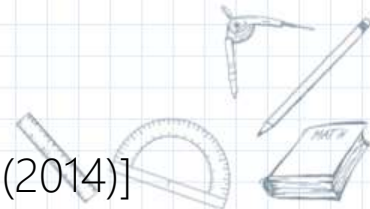
U(1) gauge \rightarrow Anisotropic BG

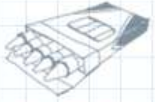
SU(2) gauge \rightarrow Isotropic BG is **Attractor**.

Isotropy is conserved



$$A_i^a = a A_{BG}(t) \delta_i^a$$

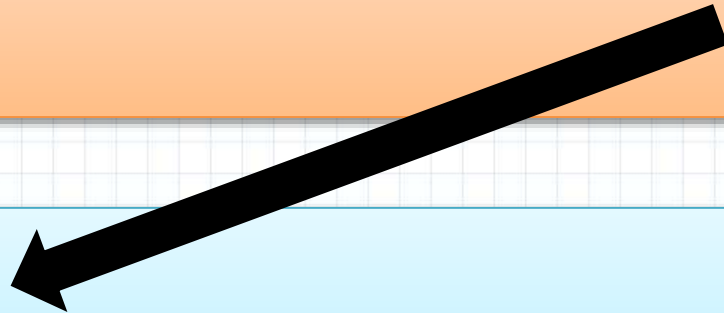


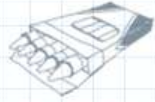


Flow of dynamics

BG: Axion χ_{BG} \longrightarrow Vector A_i^{BG}

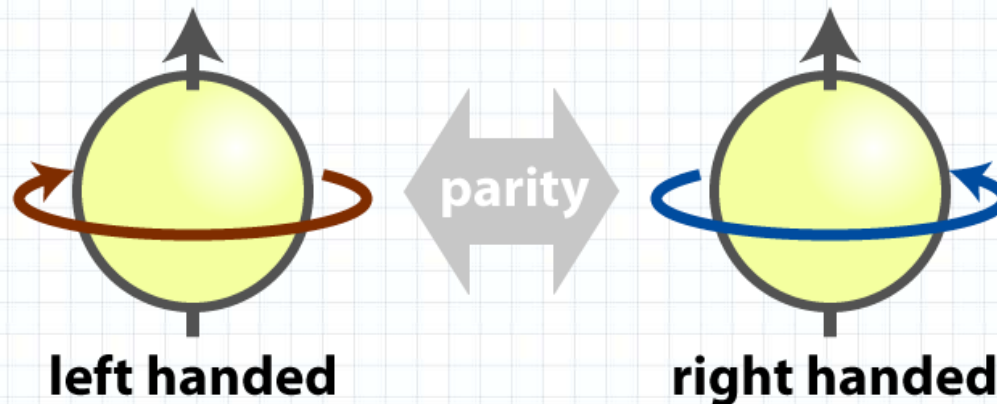
Pert: $t_{ij} = A_i^{BG} \delta A_j \longrightarrow$ GW $h_{ij}^{(s)}$





Parity-breaking

Parity symmetry is spontaneously broken while χ is rolling



Either one of two circular polarization of GW is amplified



Instability of Chiral Tensor

The EoMs for tensor perturbations are

$$h''_{R,L} + \left(1 - \frac{2}{x^2}\right) h_{R,L} = \mathcal{O}\left(\Omega_A^{1/2}\right) t_{R,L}$$

$$t''_{R,L} + \left(1 + \frac{2m_Q \xi}{x^2} \mp \frac{2}{x} (m_Q + \xi)\right) t_{R,L} = \mathcal{O}\left(\Omega_A^{1/2}\right) h_{R,L}$$

$$FF \supset g \epsilon^{ijk} A^i A^j \partial A^k, \quad \chi F \tilde{F} \supset \dot{\chi} \epsilon_{ijk} A_i \partial_j A_k$$

where we have used $\epsilon_{ijk} k_i e_{jl}^{R,L}(\hat{\mathbf{k}}) = \mp k e_{kl}^{R,L}(\hat{\mathbf{k}})$.

ϵ_{ijk} terms make t_R instable (but not t_L)

 $h_R \gg h_L$: Chiral GW is generated!

$$x \equiv -k\eta$$

$$m_Q \equiv g A^{BG} / H$$

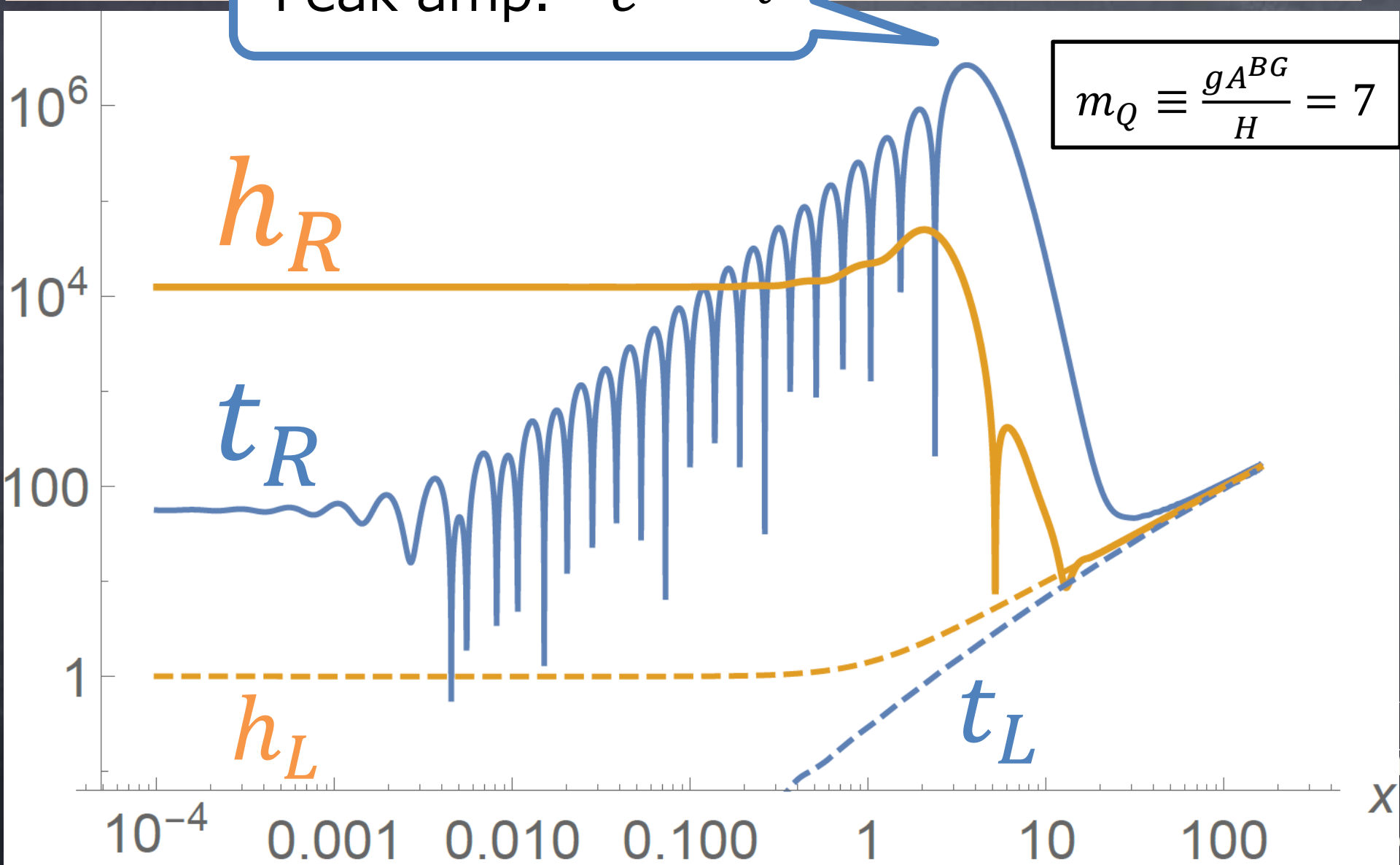
$$\xi \equiv \lambda \dot{\chi} / 2fH$$

$$A_i^a \equiv a \delta_i^a A^{BG}$$

Insta

Peak amp. $\sim e^{1.8m_Q}$

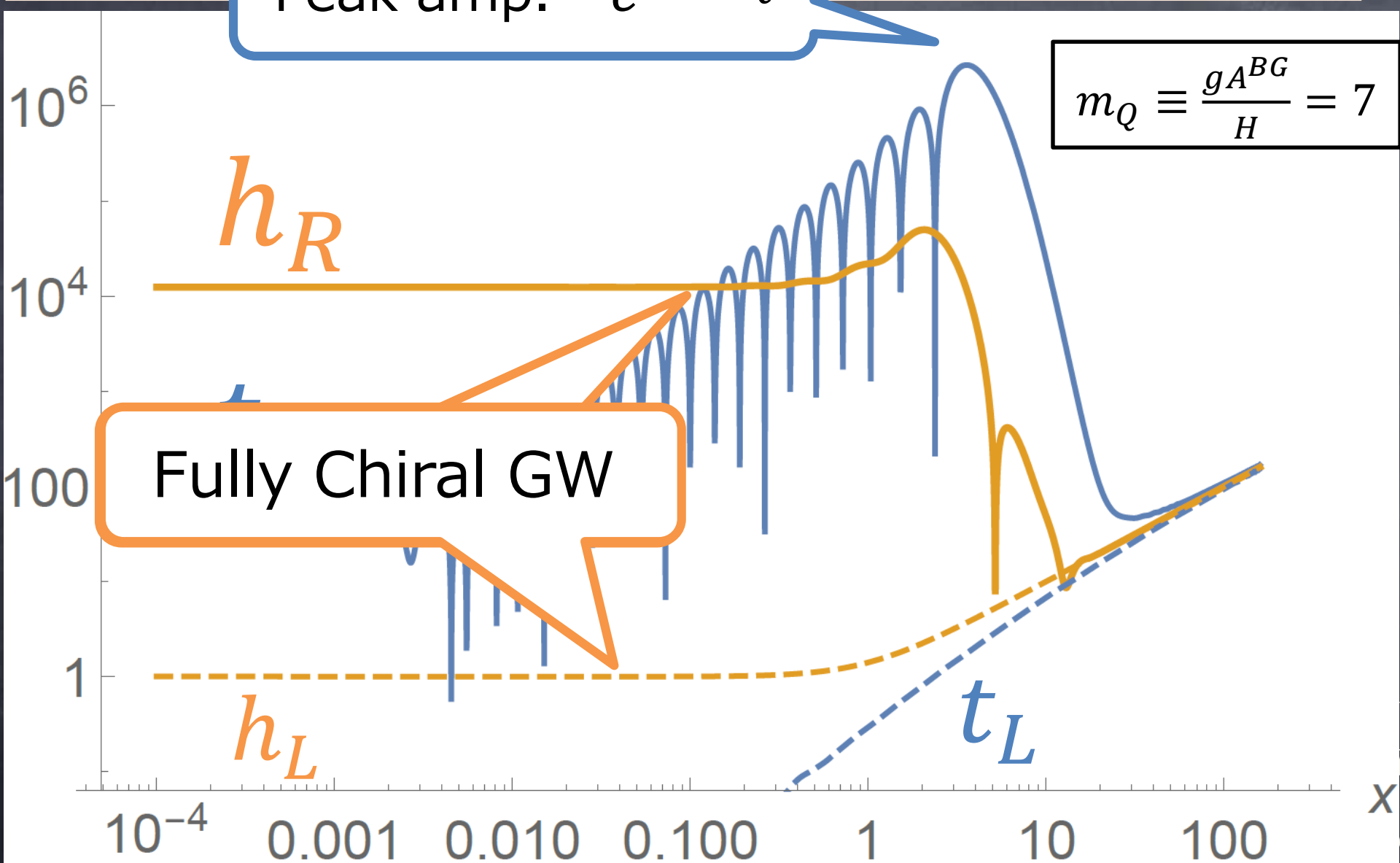
nsor



Insta

Peak amp. $\sim e^{1.8m_Q}$

nsor



$$m_Q \equiv \frac{g^{ABG}}{H} = 7$$

h_R

Fully Chiral GW

h_L

t_L

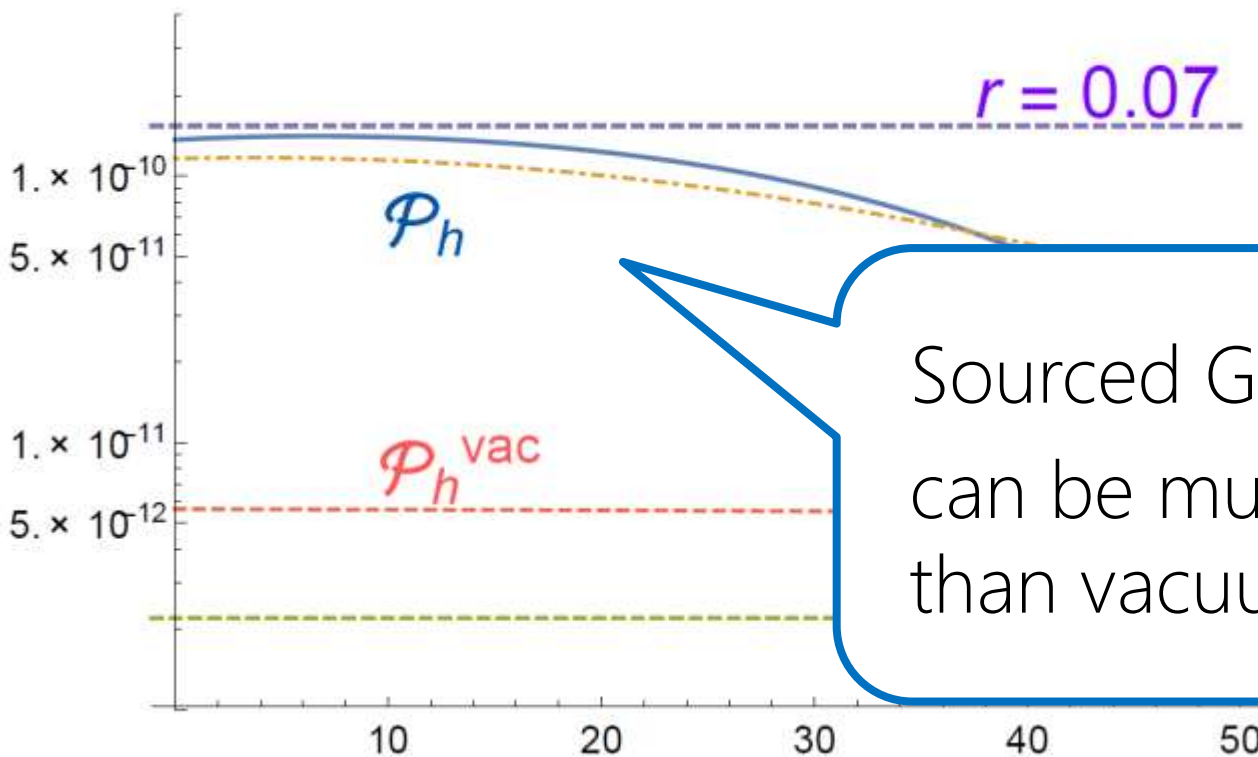
x



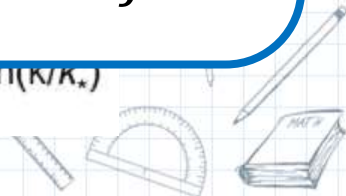
Result

○ $H_{\text{inf}} \approx 1 \times 10^{13} \text{ GeV}$

✗ $H_{\text{inf}} \approx 6 \times 10^{13} \text{ GeV}$



Sourced GW $h_{ij}^{(s)}$
 can be much larger
 than vacuum h_{ij}^{vac}

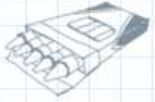




Short summary of the model

- ① It naturally realizes $r_{\text{add}} > r_{\text{vac}}$
- ② GW can be sourced at linear order with **SU(2)** gauge field
- ③ SSB of the parity symmetry leads to the Chiral GW





Axion-SU(2) model

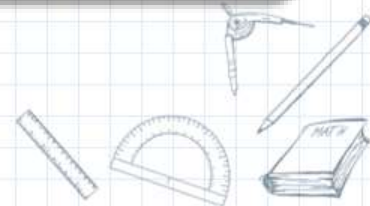
- Larger GW than h^{vac} can be produced

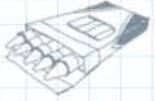
$$r_{\text{obs}} = r_{\text{vac}} + r_{\text{add}}$$

r_{obs} doesn't fix H_{inf}

- Distinguishable w/
 - Polarization $h_R \neq h_L$
 - Non-Gaussianity $\langle hhh \rangle$
 - Tensor tilt n_t

Let's observe the signature



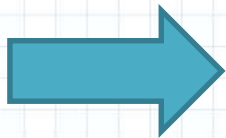


TB, EB correlation

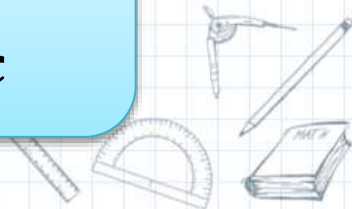
Chiral GW induces TB & EB cross correlations

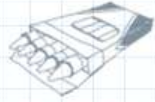
$$\langle TT \rangle, \langle TE \rangle, \langle EE \rangle, \langle BB \rangle \propto \langle h_R h_R \rangle + \langle h_L h_L \rangle$$

$$\langle TB \rangle, \langle EB \rangle \propto \langle h_R h_R \rangle - \langle h_L h_L \rangle$$



By detecting TB & EB correlations,
we can distinguish $h^{(s)}$ from h_{vac}





PRESENTATION

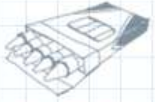
LiteBIRD

- CMB satellite mission
- Will be launched in 2020s
- Aims to detect $r \geq 10^{-3}$



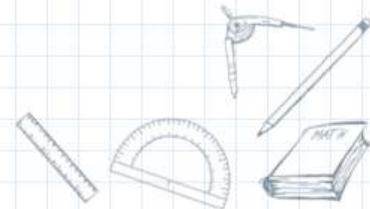
Channel (GHz)	θ_{FWHM} (amin)	$\sigma_{\text{T}}(\nu)$ [μKamin]	$\sigma_{\text{P}}(\nu)$ [μKamin]
40.0	69.0	0.0	36.8
50.0	56.0	0.0	23.6
60.0	48.0	0.0	19.5
68.0	43.0	0.0	15.9
78.0	39.0	0.0	13.3
89.0	35.0	0.0	11.5
100.0	29.0	0.0	9.0
119.0	25.0	0.0	7.5
140.0	23.0	0.0	5.8
166.0	21.0	0.0	6.3
195.0	20.0	0.0	5.7
235.0	19.0	0.0	7.5
280.0	24.0	0.0	13.0
337.0	20.0	0.0	19.1
402.0	17.0	0.0	36.9

Table 3: Summary of the LiteBIRD specifications. And $f_{\text{sky}} = 0.5$



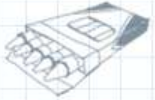
S/N for TB+EB

- w/ lensing effect (no delensing)
- LiteBIRD instrumental noise
- 2% foreground contamination



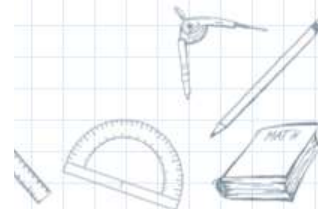
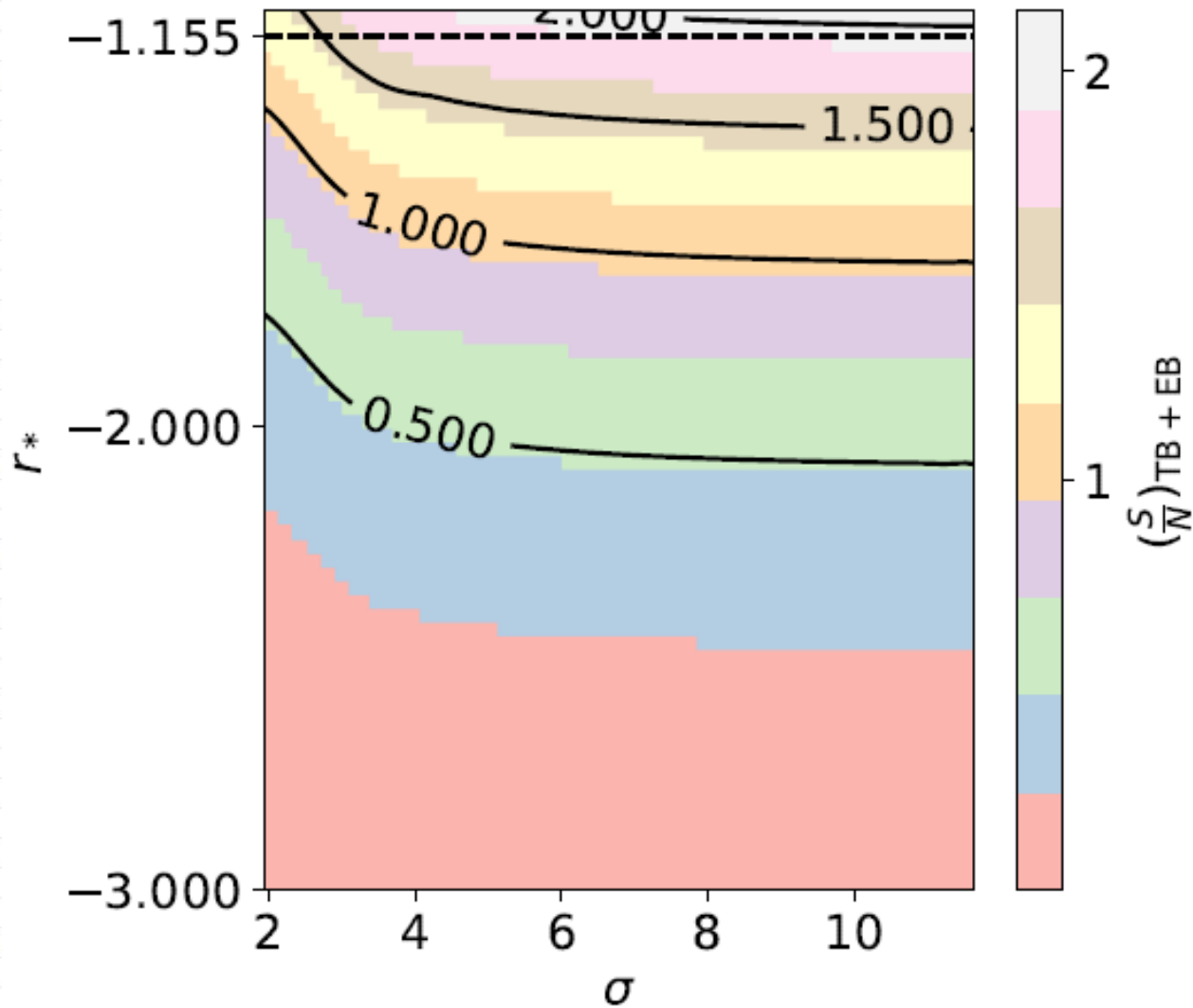


S/N ratio for TB+EB with noises



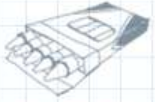
PRESENTATI

$$k_p = 7e - 05 \text{ [Mpc]}^{-1}$$



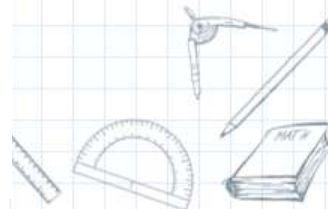
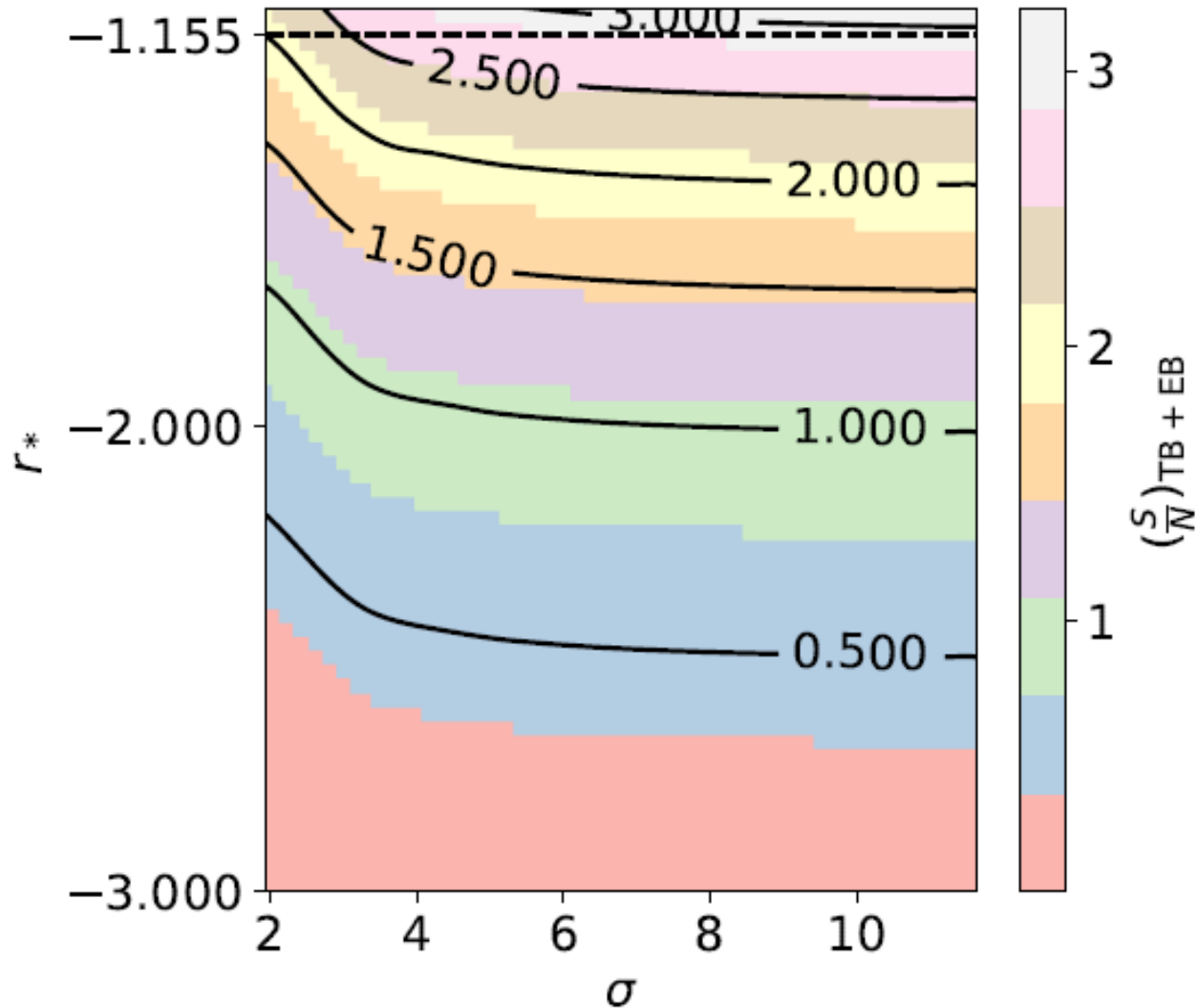


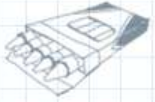
Cosmic Variance limited case



PRESENTAT

$$k_p = 7e - 05 \text{ [Mpc]}^{-1}$$





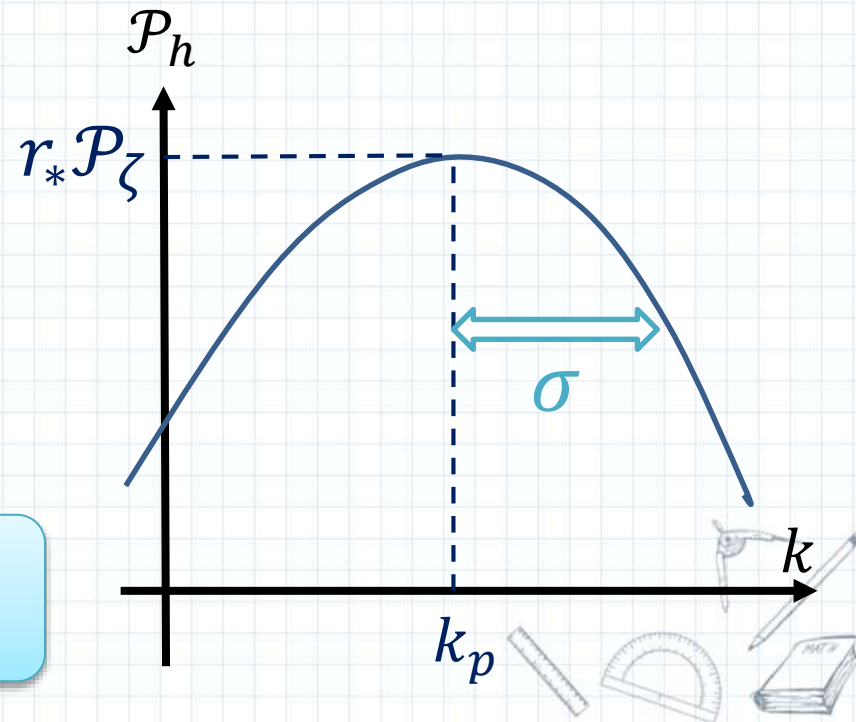
PRESENTATION

Excellent Template

$$\mathcal{P}_h^{L, \text{ Sourced}}(k) = r_* \mathcal{P}_\zeta \exp \left[-\frac{1}{2\sigma^2} \ln^2 \left(\frac{k}{k_p} \right) \right]$$

$$\sigma^2 = \frac{\Delta N^2}{2\mathcal{G}(m_*)} \cdot \Delta N \equiv \frac{\lambda}{2\xi_*}$$

$k_p = \text{arbitrary}$ (depend on χ_i)



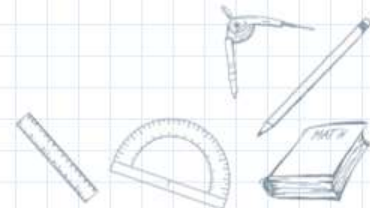
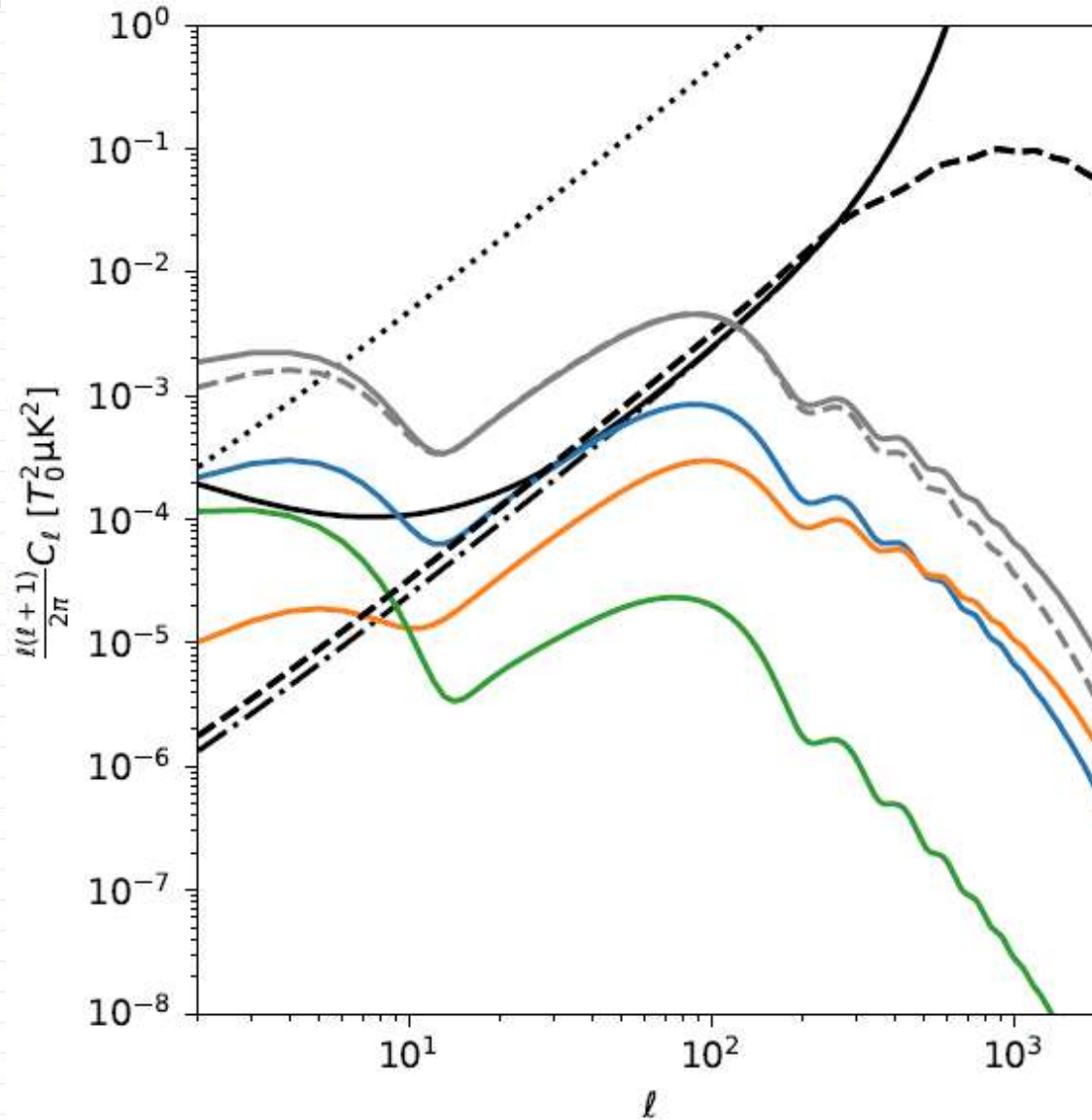
$\mathcal{P}_h^{(s)}$ takes Gaussian shape.

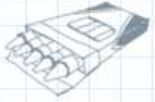


C_ℓ for BB and sensitivity curves



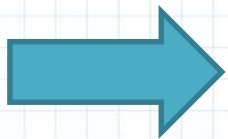
PRESENTATION



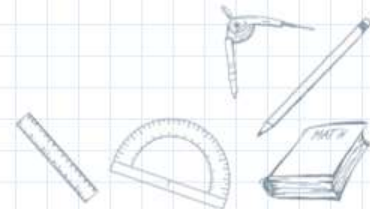


S/N for TB+EB

- w/ lensing effect (no delensing)
- LiteBIRD instrumental noise
- 2% foreground contamination



TB + EB can be detected
by LiteBIRD for $r > 0.03$.



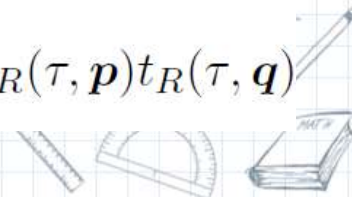


Non-Gaussianity

Large $\langle h_R h_R h_R \rangle$ ← Large $\langle t_R t_R t_R \rangle$

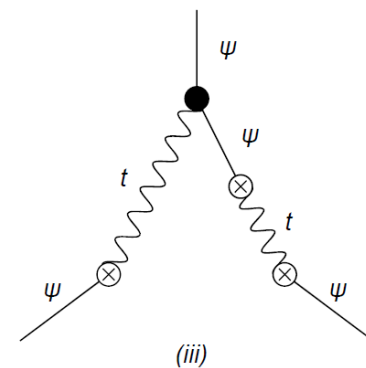
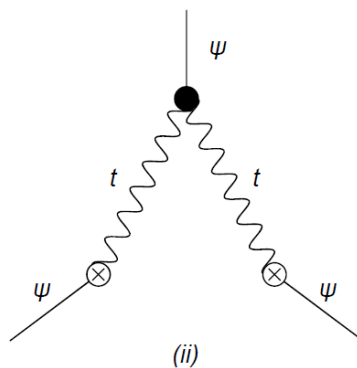
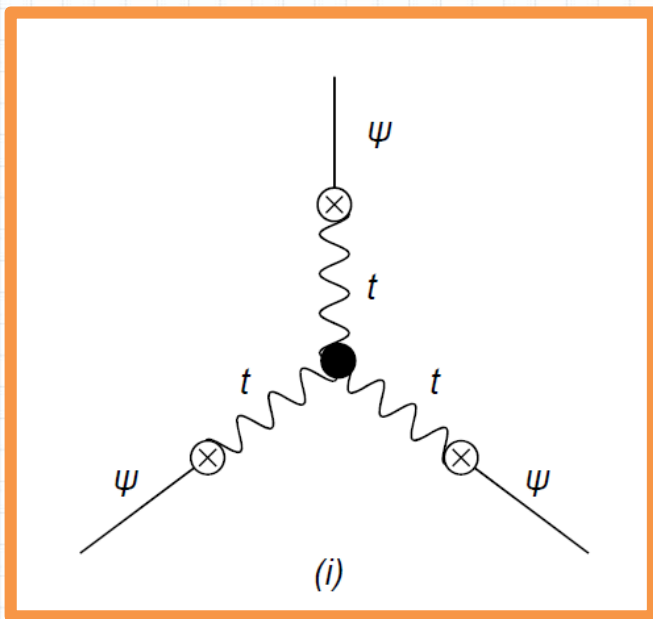
The EoM for SU(2) gauge field is non-linear

$$\begin{aligned} \partial_\tau^2 t_R(\tau, \mathbf{k}) + \left(k^2 + \frac{2m_Q \xi}{\tau^2} + 2 \frac{m_Q + \xi}{\tau} k \right) t_R(\tau, \mathbf{k}) \\ = g \left[e_{ij}^R(\hat{\mathbf{k}}) \right]^{-1} \iint \frac{d^3 p d^3 q}{(2\pi)^3} \delta(\mathbf{p} + \mathbf{q} - \mathbf{k}) Q_{ij}(\mathbf{p}, \mathbf{q}, \tau) t_R(\tau, \mathbf{p}) t_R(\tau, \mathbf{q}) \end{aligned}$$

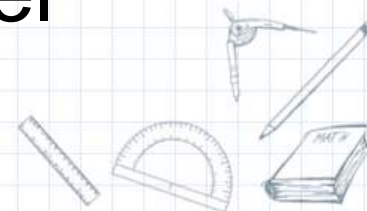


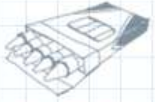


Non-Gaussianity

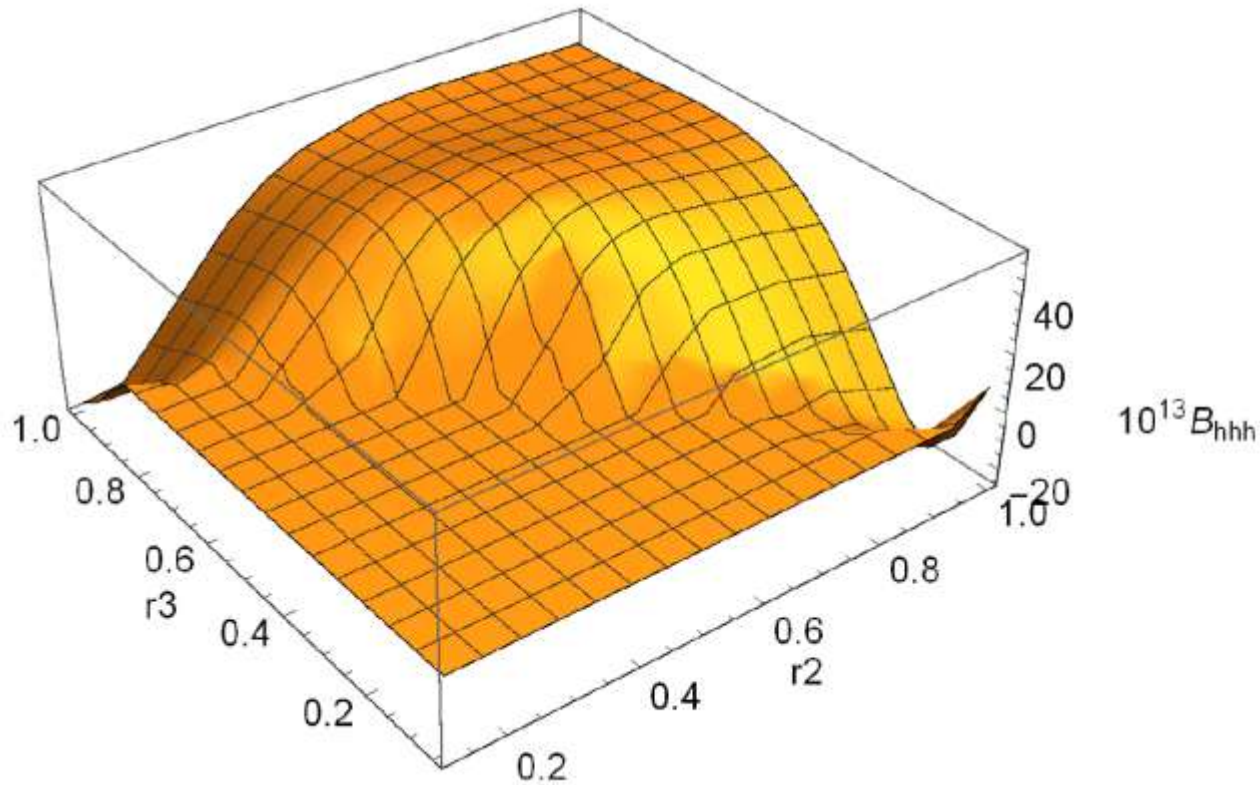


$\langle h_R h_R h_R \rangle$ is produced at Tree level

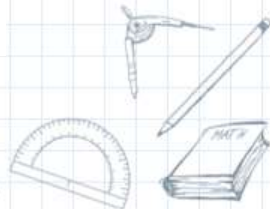




Non-Gaussianity



$\langle h_R h_R h_R \rangle$ has an interesting shape





Non-Gaussianity

Model prediction (preliminary):

$$(k_1^2 k_2^2 k_3^2) B_h^{\text{equil}} \sim -10^{-12}$$

Planck 2015 XVII constraint on $\langle h_R h_R h_R \rangle$

$$(k_1^2 k_2^2 k_3^2) B_h^{\text{equil}} = (3.5 \pm 13.2) \times 10^{-12}$$

Note: Planck analysis assumes different shape of NG





Non-Gaussianity

Model prediction (preliminary):

$$(k_1^2 k_2^2 k_3^2) B_h^{\text{equil}} \sim -10^{-12}$$

GW NG may be detected soon?

Planck 2015 XVII constraint:

$$(k_1^2 k_2^2 k_3^2) B_h^{\text{equil}} = (3.5 \pm 13.2) \times 10^{-12}$$

Note: Planck analysis assumes different shape of NG





Summary

Axion-SU(2) model



PRESENTATION

- Larger GW than h^{vac} can be produced

$$r_{\text{obs}} = r_{\text{vac}} + r_{\text{add}}$$

r_{obs} doesn't fix H_{inf}

- Distinguishable w/

Let's observe it!

- Polarization $h_R \neq h_L \Rightarrow$ TB&EB correlation

- Non-Gaussianity $\langle hhh \rangle \Rightarrow$ large B_h^{equil}

- Tensor tilt $n_t \Rightarrow n_T \neq -r/8$



- Lowest ρ_{inf} for $r = 10^{-3}$

- $r \propto \Omega_{\text{GW}}(k) \propto \Omega_{\text{source}}(k)$

$$\rho_{\text{inf}}^{1/4} \geq 0.1 \text{ GeV}$$





Fin

THE THEME
OF CHAPTER IS...

Thank you !
