# Einstein-Yang-Mills-Dirac Systems from the discretized Kaluza-Klein Theory

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#### Abstract

A unified theory of the non-Abelian gauge interactions with gravity in the framework of a discretized Kaluza-Klein theory is constructed with modified Dirac Operator and wedge product. All the couplings of the chiral spinors to non-Abelian gauge fields emerge naturally as components of the couplings of the chiral spinors to the generalized gravity together with some new interactions. In particular, the currently prevalent gravity-QCD-quark and gravity –electroweak –quark –lepton models are shown to follow as special cases of the general framework.

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## OUTLINE

- Introductory Remarks
- Fundamental Basics of NCG; Spectral Triple
- Algebraic Formulation of NCG Geometry
- Equivalence Principle and Vielbeins
- A set of constraints; Cartan Structure Equations
- Determination of connection coefficients
- Ricci curvature leading to Lagrangians
- Physical consequences and Conclusions

#### Introductory Remarks

- Our current picture of space-time is inadequate for the description of phenomena at short distances. The twin pillars of modern physics, GR and QFT are both structured on a continuum picture of space-time.
- Algebraically geometrized pseudo-Riemannian manifold with a metric based on a continuum picture underlies General Relativity.
- Likewise, Quantum Fields and their interactions are local operators based on continues functions of space-time coordinates.

Noncommutative Geometry (NCG) a la Connes has provided a new suitable framework for a Geometrical description of all elementary particle interactions.

### General Framework of NCG

Consists of Three Basic Points: Spectral Triple

(A, H, D)

A- An involutive algebra, commutative or noncommutative.

- H- A Hilbert space as the carrier space for A
- D- A self-adjoint operator acting on A

The Algebra A generalizes the smooth functions C(M). The Operator D allows one to build a Differential structure associated with any associative algebra.

# Extended space-time $\mathcal{M}^4 \times Z_2$ and Dirac operator

•  $Z_2$  algebra has only two elements represented as follows

$$\mathbf{e} \hspace{.1in} = \hspace{.1in} \begin{bmatrix} 1 \hspace{.1in} 0 \\ 0 \hspace{.1in} 1 \end{bmatrix}, \hspace{.1in} \mathbf{r} = \begin{bmatrix} 1 \hspace{.1in} 0 \\ 0 \hspace{.1in} -1 \end{bmatrix}$$

• Generalized spinors and function operators

$$\Psi = \begin{bmatrix} \psi_L \\ \psi_R \end{bmatrix}, \quad F = \begin{bmatrix} f_L & 0 \\ 0 & f_R \end{bmatrix}$$
(2)

• Generalized Dirac operator and Dirac matrices  $M, N = \mu, 5$ 

$$D = \Gamma^{M}\partial_{M} = \Gamma^{\mu}\partial_{\mu} + \Gamma^{5}\partial_{5} = d + Q = \begin{bmatrix} \gamma^{\mu}\partial_{\mu} & im_{2}\gamma^{5} \\ -im_{2}\gamma^{5} & \gamma^{\mu}\partial_{\mu} \end{bmatrix}, \quad D_{5} = \begin{bmatrix} 0 & m \\ -m & 0 \end{bmatrix},$$
  
$$\partial_{5}F = \sigma^{\dagger}[D_{5}, F] = m(f_{L} - f_{R})\mathbf{r}, \quad \Gamma^{\mu} = \begin{bmatrix} \gamma^{\mu} & 0 \\ 0 & \gamma^{\mu} \end{bmatrix}, \quad \Gamma^{\mu} = \begin{bmatrix} 0 & i\gamma^{5} \\ -i\gamma^{5} & 0 \end{bmatrix}$$
(3)

(1)

## Differential forms and wedge product

• Exterior derivative of 0-form

$$DF = [D, F] = \begin{bmatrix} df_L^u & im_2\gamma^5(f_L^u - f_R^u) \\ im_2\gamma^5(f_L^u - f_R^u) & df_R^u \end{bmatrix}, D^2 = 0.$$
(4)

- Generalized 1-form  $U = \Gamma^{M}U_{M} = \Gamma^{\mu}U_{\mu} + \Gamma^{5}U_{5} = \begin{bmatrix} \gamma^{\mu}u_{L\mu} & im_{2}\gamma^{5}u_{5L} \\ -im_{2}u_{5R} & u_{R\mu} \end{bmatrix}$ . For gravity it enough to impose the hermitian 1-form  $u_{5L} = u_{5R}$ . For nonabelian gauge fields  $u_{L,R\mu}$  can have values in some nonabelian Lie-algebra, while  $u_{L5} = u_{R5}^{\dagger}$ .
- Wedge product

$$\Gamma^{\mu} \wedge \Gamma^{M} = -\Gamma^{M} \wedge \Gamma^{\mu}, \quad \Gamma^{5} \wedge \Gamma^{5} \neq 0, \quad U \wedge V = \Gamma^{M} \wedge \Gamma^{N} (U \wedge V)_{MN}$$
(5)

- It is straight forward to write down all the components of  $DU = [D, U] = \Gamma^M \wedge \Gamma^N (DU)_{MN}$  and  $U \wedge V = \Gamma^M \wedge \Gamma^N (U \wedge V)_{MN}$ .
- Generalized 2-form  $S = \Gamma^M \wedge \Gamma^N S_{MN}$
- NCG is determined by a spectral quartet not triplet!

#### Generalized equivalence principle and vielbeins

- In curved space-time the curvi-linear basis is  $\Gamma^M(x)$
- The new equivalence principle: there exist a locally flat basis Γ<sup>A</sup>, A = a, 5 which is related to Γ<sup>M</sup> by the local transformation

$$\Gamma^{A} = \Gamma^{M} E^{A}_{M}(x), \quad \Gamma^{M} = \Gamma^{A} E^{M}_{A}(x)$$
(6)

•  $E_A^M(x)$  vielbeins defined the metric as follows

$$G^{MN} = Tr(\Gamma^{M}\Gamma^{N}) = E^{M}_{A}(x)Tr(\Gamma^{A}\Gamma^{B})E^{N}_{B}(x) = E^{M}_{A}(x)\eta^{AB}E^{N}_{B}(x)$$
(7)

• We will use the hermitian vielbein

$$E^{a}_{\mu} = \begin{bmatrix} e^{a}_{L,\mu} & 0\\ 0 & e^{a}_{R,\mu} \end{bmatrix} , \ E^{a}_{\mu} = 0 , \ E^{\dot{5}}_{\mu} = \begin{bmatrix} a_{L\mu} & 0\\ 0 & a_{R\mu} \end{bmatrix} , \ E^{\dot{5}}_{5} = \begin{bmatrix} \phi & 0\\ 0 & \phi \end{bmatrix},$$
(8)

#### Levi-Civita connection and Cartan structure equation

- Metric compatible and hermitian Levi-Civita connection  $\Omega_{AB} = -\Omega_{BA} = \Omega_{AB}^{\dagger}$
- Generalized structure equation

$$T^{A} = DE^{A} + E^{B} \wedge \Omega^{A}_{B} , \ R^{AB} = D\Omega^{AB} + \Omega^{AC} \wedge \Omega^{B}_{C}$$
(9)

• Viet-Wali's constraint (1995): The torsion free  $T^E = 0$  leads to restricted metric. Therefore, we have

$$\mathcal{T}^{a} = 0 , \quad \mathcal{T}^{\dot{5}}_{AB} = t^{\dot{5}}_{AB} \mathbf{r}$$
 (10)

allows us to compute torsion and connection and then Ricci tensor from vielbeins.

• Generalized Hilbert-Einstein action

$$S_{HE} = M_{5,PI}^2 \int \sqrt{-\det G} \, Tr(R_5), \quad R_5 = \eta^{AC} R_{ABCD} \eta^{BD}$$
(11)

• Finite spectrum of bigravity, one Brans-Dicke and the gauge sector from  $a_{L,R\mu}$  possibly non-abelian

### Bigravity coupled to Brans-Dicke

• In this case we choose

$$a_{L,R} = 0, \quad \phi = \phi_0 \exp\left(\frac{\sigma(x)}{\sqrt{2}M_{Pl}}\right), \quad (12)$$
  
$$e_a^{\mu}(x) = \frac{1}{2}(e_{La}^{\mu}(x) + e_{Ra}^{\mu}(x)), \quad h_{\mu\nu} = \frac{1}{2m_h}(e_{L\mu}^a - e_{R\mu}^a)e_{a\nu}, \quad (13)$$

• The Hilbert-Einstein action reduces to

$$S_{HE}(5) = \int dx^{4} (\sqrt{-detg} \left( \frac{M_{Pl}^{2}}{2} (r_{L4} + r_{R4}) + \frac{1}{2} g^{\mu\nu} \partial_{\mu} \sigma(x) \partial_{\nu} \sigma(x) \right) \\ + \frac{2m^{2} M_{Pl}^{2}}{\phi_{0}^{2} m_{h}^{2}} (h_{\nu}^{\mu} h_{\mu}^{\nu} - (h_{\mu\nu} g^{\mu\nu})^{2}) + \mathcal{L}_{int}(\sigma(x), h_{a\mu}(x), e_{a}^{\mu}(x)), \quad (14)$$

• The Pauli-Fierz mass term for the massive gravity  $h_{\mu\nu}$  with mass  $\frac{\sqrt{2}mM_{Pl}}{\phi_0 m_h}$  and  $M_{Pl} = M_{5,Pl}\sqrt{\phi_0}$ 

#### Viet-Du's theorem

• The case of  $e_L = e_R$  Einstein's gravity, Brans-Dicke coupled to the gauge fields

$$R_{5} = r_{4} - \frac{1}{16} \phi^{2} g^{\mu\rho} g^{\nu\tau} \mathcal{F}_{\mu\nu} \mathcal{F}_{\rho\tau} + g^{\mu\nu} \frac{\partial_{\mu} \phi}{\phi} \frac{\partial_{\nu} \phi}{\phi}, \qquad (15)$$

$$\mathcal{F}_{\mu\nu} = f_{L\mu\nu} + f_{R\mu\nu} - m([a_{R\nu}, a_{L\mu}] + [a_{L\nu}, a_{R\mu}]), \qquad (16)$$

$$f_{L,R\mu\nu} = \partial_{\mu} a_{L,R\nu} - \partial_{\nu} a_{L,R\mu} - m[a_{L,R\mu}, a_{L,R\nu}] \qquad (17)$$

R<sub>5</sub> and S<sub>HE</sub> is gauge invariant with nonabelian gauge fields only in two cases:
 Gauge field on one sheet is abelian a<sub>Rµ</sub>. Electroweak case.

2  $a_{R\mu} \sim a_{L\mu}$ . Two gauge fields are the same on both sheets. Strong interaction case.

• Viet-Du's results are recovered with the new Dirac operator and wedge product.

#### Coupling to the chiral spinors

• The curved Dirac operator with spinor connection

$$\mathcal{D} = D + \Omega = D - \frac{1}{8} \Gamma^C \Omega_{ABC} [\Gamma^A, \Gamma^B], \qquad (18)$$

$$D = \Gamma^{\mu}\partial_{\mu} + \Gamma^{5}\sigma^{\dagger}D_{5} = \Gamma^{a}E^{\mu}_{a}(x)\partial_{\mu} - \Gamma^{a}E^{\mu}_{a}(x)A_{\mu}(x)\sigma^{\dagger}D_{5} + \Gamma^{5}\phi^{-1}(x)\sigma^{\dagger}D_{5}, \qquad (19)$$

• The Einstein-Dirac parts can be derived from the following generalized Lagrangian

$$\mathcal{L}_{f} = \bar{\Psi} \mathcal{D} \Psi = \mathcal{L}_{d+m} + \mathcal{L}_{f-g} + \mathcal{L}_{\omega} + \mathcal{L}_{\Omega}(2) + \mathcal{L}_{\Omega}(3)$$
(20)

where the Brans-Dicke modifies the quark-lepton mass by a small amount since  $\phi_0$  is large

$$\mathcal{L}_{d+m} = i\bar{\psi}(\gamma^{a}e^{\mu}_{a}(x)\partial_{\mu} - m\phi^{-1}(x))\psi, \quad \mathcal{L}_{\omega} = -\frac{1}{8}\bar{\psi}\gamma^{c}\omega_{abc}[\gamma^{a},\gamma^{b}]\psi \qquad (21)$$

$$\mathcal{L}_{\Omega}(2) = \frac{1}{16} \bar{\psi} e^{\mu}_{a} e^{\nu}_{b} (\hat{f}_{+\mu\nu} + 2m[a_{-\nu}, a_{-\mu}]) [\gamma^{a}, \gamma^{b}] \psi, \qquad (22)$$

$$\mathcal{L}_{\Omega}(3) = \frac{1}{\sqrt{2}M_{Pl}}\bar{\psi}\gamma^{a}e^{\mu}_{a}\partial_{\mu}\sigma(x)\psi$$
(23)

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### Coupling to the gauge fields

- All the gauge-chiral spinor coupling can be derived from generalized Einstein-Dirac Lagrangian.
- Matching with the SM terms, based on the currently known quark-leptons we can derived the relations

$$g = \frac{2\sqrt{2}m}{\phi_0 M_{Pl}}, \quad g' = \sqrt{1.2} g, \quad g_s = \sqrt{2}g$$
 (24)

• Parity violation by QCD can be transfered into P-violation by gravity

$$i\bar{\psi}_L\gamma^a e^{\mu}_{La}(\partial_{\mu} + ia_{L\mu})\psi_L + i\bar{\psi}_R\gamma^a e^{\mu}_{Ra}(\partial_{\mu} + ia_{R\mu})\psi_R$$
(25)

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# Summary

- With new Dirac operator and wedge product, one can derive nonabelian gauge fields from the generalized gravity under two different conditions, which allows to build the Einstein-Yang-Mills-Dirac systems for QCD and Electroweak interactions.
- The coupling of generalized gravity to chiral spinors can be reduced to couplings of gravity and gauge interactions to chiral spinors, together with the spinor connection terms and new terms.
- There are some relations between the coupling constants.
- Bigravity can be derived fron the framework.
- The theory is based on a new equivalence principle for spacetime extended by discrete dimension.
- The theory has a finite spectrum without truncation
- The hierachy problem can be solved by Brans-Dicke scalar.

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#### Questions for further studies

- The energy scale, where the theory becomes valid? (Frome the relations between the coupling constant)
- Since two conditions of Viet-Du theorem cannot be satisfied at the same time, one can must go with the extended space-time  $M^4 \times Z_2 \times Z_2$  (Viet 2015)
- w-type and v-type chiral matter (SM and dark matter?)