Implication of Higgs Factory Precision Measurements on New Physics Models

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Outline

- **Higgs precision measurements**
- **Global fit framework**
- **Perturbative models**
	- **– SM with a real singlet extension**
	- **– 2HDM (tree + loop, Higgs + Zpole)**
	- **– MSSM**
- **Strong dynamics models (skip in this talk, see Jiayin's)**
- **Complementarity with direct search @ 100 pp** ¥
- **Conclusion**

Higgs Precision Measurements

LHC: 14 TeV, 300 fb-1, 3000 fb-1

Higgs Precision Measurements Figgs Frecision wiedsureniems

linear colliders), such a function process becomes significantly more important and can provide α

- **CEPC / FCC / ILC** crucial complementary information. For p*s >* 500 GeV, *tth* production can also be used as

S. Su CEPC-preCDR, TLEP Design Study Working Group, ILC Operating Scenarios. T p. Su CEPC-preCDR, TLEP Design Study working Group, ILC Operating Scenarios.

Higgs Precision Measurements Figgs Frecision wiedsureniems

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Kappa framework and EFT Framework ¹⁹⁸³ **2.4 Coupling Extractions and Combinations** by data.

S. Su 6 ²⁰⁰³ of the exotic decay or treating it as an independent parameter (essentially assuming it can not be $\overline{\text{S}}$: $\overline{\text{S}}$ u

New Physics Implication

Kappa Framework and EFT Framework

-

limitations of model-independent approaches

- **•large level of degeneracy parameter space for specific model much smaller**
- **•correlation matrix often not provided over conservative estimation when not include correlation**
- **•assumptions and simplifications may not be valid for a particular model**

New Physics Implication

New Physics Implication

New Physics Implication precision corrections (*e.g.* by imposing custodial symmetries) such that these constraints are ing Tmplingtion studies. The detailed in

Perturbative Models

- **•SM with a real singlet extension**
- **•2HDM (Type I, II, L, F)**
- **•MSSM**

SM + Real Scalar Singlet

-

\bullet SM + real scalar singlet The general Lagrangian for this model is, of SM + real scalar singlet model parameters, we do not single out spontaneous *Z*2 breaking as a separate out spontaneous *D*

the *Z*² limit, it is still possible to have spontaneous *Z*² breaking once *S* acquires a vacuum

$$
\mathcal{L} = \mathcal{L}_{\rm SM} + \frac{1}{2} (\partial_{\mu} S)^2 - \frac{1}{2} m_S^2 S^2 - \Lambda_{SH} S(H^{\dagger} H) - \frac{1}{2} \lambda_{SH} S^2(H^{\dagger} H) - \frac{1}{3!} \Lambda_S S^3 - \frac{1}{4!} \lambda_S S^4
$$

This simplest extension to the SM could already induce many test scenarios. Such model ⊙ after EWSB, 2 physical Higgse: CP-even Higgses: hsm, singlet S

where $\overline{}$ is the SM Higgs doublet and $\overline{}$ is the new real singlet field. The new real singlet field $\overline{}$

Example 22 breaking: mixing between hsm and S

$$
h_{125} = \cos\theta \ h_{\rm SM} + \sin\theta \ S \qquad \qquad \kappa_i = g_i^{\rm SM + singlet} / g_i^{SM} = \cos\theta
$$

purpose of this study is to focus on the Higgs physics implications, instead of the extraction

SM + Real Scalar Singlet

-

๏ **fit to sin θ**

SM + Real Scalar Singlet

-

๏ **fit to sin θ**

SM + Real Scalar Singlet the Effect of the possible e the changes of the SM Higgs property in models with an extra real scalar singlet, we adopted α the EFT language to EFT language to each \mathcal{S} M + Real Scalar \mathcal{S} \sim changes property in models with an extra real singlet, we are also real singlet, we ar $SM + Real Scalar$ Singlet the changes of the SM Higgs property in models with an extra real scalar singlet, we adopt

After integrating out the singlet field, the singlet field, the general EFT with dimension-six operators can be

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-

 \bullet fit to c₆ and c_H After integrating out the singlet field, the singlet field, the general EFT with dimension-six operators can be f_{i+1} to be end on

the changes of the SM Higgs property in models with an extra real scalar singlet, we adopt

Perturbative Models

-

•SM with a real singlet extension

- **•2HDM (Type I, II, L, F)**
- **•MSSM**

2HDM in one slide ⎛ ⎞ 2 Scenarios with large *H* ! *AZ* or *A* ! *HZ* discrete *Z*2 symmetry imposed on the Lagrangian, we are left with six free parameters, we are left with six free

H^u =

H⁰

⎛

Higgses, and the ratio of the two vacuum expectation values, tan = *v*2*/v*1. In the case in which

1

Two types of couplings that are of particular interest are the couplings of a Higgs to two gauge to two

u

⎞

 \equiv d

 $\frac{1}{\sqrt{2}}$, He $\frac{1}{\sqrt{2}}$, He $\frac{1}{\sqrt{2}}$

- $|$ (CP. **CC** H⁺ $\overline{}$ \circ Two Higgs Doublet Model (CP-conserving) $\sum_{i=1}^{\infty}$ iwo inggs boublet model (or which can be chosen as four Higgs masses (*mh*, *mH*, *mA*, *mH[±]*), the mixing angle ↵ **e** Two Higgs Doublet Model (CP-conserving)

which can be chosen as four Higgs masses (*mh*, *mH*, *mA*, *mH[±]*), the mixing angle ↵

$$
\Phi_{i} = \begin{pmatrix} \phi_{i}^{+} \\ (v_{i} + \phi_{i}^{0} + iG_{i})/\sqrt{2} \end{pmatrix} \qquad \begin{aligned} v_{u}^{2} + v_{d}^{2} &= v^{2} = (246 \text{GeV})^{2} \\ \tan \beta &= v_{u}/v_{d} \end{aligned}
$$
\n
$$
\begin{pmatrix} H^{0} \\ h^{0} \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} \phi_{1}^{0} \\ \phi_{2}^{0} \end{pmatrix}, \qquad A = -G_{1} \sin \beta + G_{2} \cos \beta \\ H^{\pm} = -\phi_{1}^{\pm} \sin \beta + \phi_{2}^{\pm} \cos \beta \end{aligned}
$$

 \mathbf{S} %&&'¹ [−] ^m² **ש**
 d Higgses:
= $\frac{m_V^2}{m_V^2} \sin \theta$ CP -even Higgses: hº, Hº, CP -odd Higgs: Aº, Charged Higgses: H $^{\pm}$ 1 er after EWSB, 5 physical Higgses and the set of the set Two types of couplings that are of particular interest are *ZAH*0*/h*⁰ couplings and

χ[−] 2

CP-even Higgses: h⁰, H⁰, CP-odd Higgs: A⁰, Charged Higgses: H[±]
\n• h⁰/H⁰ VV coupling
$$
g_{H^0VV} = \frac{m_V^2}{v} \cos(\beta - \alpha), \quad g_{h^0VV} = \frac{m_V^2}{v} \sin(\beta - \alpha).
$$

S. Su and the set of the ⎜⎜⎜⎜⎜⎜⎜⎝ χ0 2 χ0 ⎟⎟⎟⎟⎟⎟⎟⎠ = $\overline{}$ ^O(m^Z ^µ)(m^Z ^M²)(s2^β ⁺ ^M² ^µ) 1 ^O(m^Z $\frac{1}{15}$ $\overline{}$ alignment limit: $cos(\beta-\alpha)=0$, h^o is the SM Higgs with SM couplings. ⎟⎟⎟⎟⎟⎟⎟⎠ **A**
S. Su *A* ² cos *.* (2.2) $T_{\text{S}(\text{C}(\text{C}))}$ for a set $\text{C}(\text{C}(\text{C}(\text{C}))$ is the SM Higgs with SM coupling $\overline{}$ 15 **alignment limit: cos(β-α)=0, h0 is the SM Higgs with SM couplings.**

2HDM parameters

-

@ parameters (CP-conserving, flavor limit, Z2 symmetry)

Tree-level 2HDM fit

-

2HDM, LHC/FCC fit

Tree-level 2HDM fit

2HDM Model Distinction *µ*-fit for all three types of 2HDM. For Type II and Type-F, the CEPC -fit results are similar

2HDM: Loop in the Alignment Limit factory seep in the thighment similar

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® theoretical constraints and Perturbativity and Perturbativity and Perturbativity and Perturbativity and Per

masses for heavy Higgses deviate from the simplified relation.

2HDM: Loop in the Alignment Limit **POM: LOOP IN THE ANGINIER LIMIT OF COSTAL EXPLORER**

 W define the normalized Higgs coupling including loop e \mathcal{A}

2HDM: Loop in the Alignment Limit

-

๏ **Type II, varying luminosity**

2HDM: Tree + Loop

2HDM: Tree + Loop

- **Varying λv2**

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2HDM: Tree + Loop

work in progress

Direct Search of Heavy Higgses @ 100 pp

2HDM: non-degenerate

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Perturbative Models

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•SM with a real singlet extension

- **•2HDM (Type I, II, L, F)**
- **•MSSM**

MSSM

-

MSSM Higgs sector **• Higgs mass**

$$
M_h^2 = m_h^{2,\text{tree}} + \frac{3}{2} \frac{G_F \sqrt{2}}{\pi^2} m_t^4 \left\{ -\ln\left(\frac{\overline{m}_t^2}{M_S^2}\right) + \frac{X_t^2}{M_S^2} \left(1 - \frac{1}{12} \frac{X_t^2}{M_S^2}\right) \right\}
$$

 \sim 3 GeV uncertainties (higher loops, m_t,...)

Package: FeynHiggs gaage and , andwa coupings **• gauge and Yukawa couplings • hgg and hγγ**

• hgg and hyy

$$
\binom{H}{h} = \begin{pmatrix} \cos \alpha_{eff} & \sin \alpha_{eff} \\ -\sin \alpha_{eff} & \cos \alpha_{eff} \end{pmatrix} \begin{pmatrix} H^d \\ H^u \end{pmatrix}
$$

k anβ, M_S, X_t, **µ=500 GeV, other irrelevant** $A \subset A$ **MSSM parameters:**

 m_A vs. X_t

-

tanβ=30, µ=500 GeV, mA=2000 GeV

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m_A vs. tan $β$

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m_A vs. tan β

Complementary to LHC direct search

m_A vs. M_s

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Conclusion

- **Higgs factory reach impressive precision** ¥
- ₽ **Kappa-scheme/EFT scheme/model specific fit**
- ¥ **indirect constraints on new physics models**
- ₽ **complementary to Zpole precision program**
- ₽ **complementary to direct search @ 100 TeV pp**

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Conclusion

Conclusion

An exciting journey ahead of us!

Backup Slides

- **•Minimum composite Higgs Model (MCHM)**
- **•General EFT patterns of strong interacting models with a light Higgs**

Composite Higgs in one slide The McHMs representation of the McHMs representation of the Higgs as a pseudo-Nambu-Goldstone as a pseudo-Namb oon posito in ggo in ond direct

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- © Higgs is the PNGB of the spontaneous breaking of G⇒H boson of the global symmetry *SO*(5)*/SO*(4) that respects custodial symmetry. We investigate ϵ ringers is the rived of the spontaneous breaking of ϵ -Tr 6.1 Minimal Composite Higgs Models
- ๏ **EWSB is induced by vacuum misalignment, parametrized by ξ=v2/f2** following the notation and calculations detailed in Ref. [98]. In the minimal coset *SO*(5)*/SO*(4), gauge invariance fixes the rescaled gauge coupling \bullet EWSB is induced by vacuum misalignment, parametrized by ξ =v
- **O** mass of SM fermion generated by mixing with composite states and find a way to the community of the means of the manual representation representations in the MCHMS, in the
- \bullet light top partners can be searched at the LHC *g*CH *hV V g*SM \bullet light top partners can be searched at the LHC
- ๏ **minimal composite Higgs Model (MCHM): SO(5)/SO(4)** *hV V* W_{2} minimal composite ringgs model (morning). $\sigma(\sigma_{\text{1}})$ In the minimal coset *SO*(5)*/SO*(4), gauge invariance fixes the rescaled gauge coupling
	- **- hVV**

underlying assumptions in the following sections.

- hVV
$$
\kappa_V \equiv \frac{g_{hVV}^{\text{CH}}}{g_{hVV}^{\text{SM}}} = \sqrt{1 - \xi}
$$

- hff: depends on the fermion representation

$$
F_1 \equiv \frac{1-2\xi}{\sqrt{1-\xi}}\,, \qquad F_2 \equiv \sqrt{1-\xi}
$$

$MCHM$ $\frac{1}{2}$ reality, the summation is truncated after a few tower fermions, generating after a few tower f

-

a more complicated and scattered relation between model parameters. For simplicity and as

the modifications to the Higgs couplings are dominated by the first few tower few

๏ **Fermion representation**

e Fermion representation MCHM: $\xi = v^2/f^2 < 10^{-3}$, f > 4 TeV

MCHM to be bigger than 4 TeV. For the embeddings with *^t* related to function *F*3*,*4*,*5, the 95% C.L. to be bigger than 4 TeV. For the embeddings with *^t* related to function *F*3*,*4*,*5, the 95% C.L.

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bounds of ⇠ and *f* vary in a certain range, given the extra parameter dependence.

- **•Minimum composite Higgs Model (MCHM)**
- **•General EFT patterns of strong interacting models with a light Higgs**

Strong Dynamics in EFT Language EFT framework with dimension-6 (D6) operators, parameterized by

-

๏ **EFT operators**

$$
\mathcal{L}_6 = \frac{1}{m_*^2} \sum_i c_i \mathcal{O}_i
$$

$$
\begin{array}{c|c}\n\hline\n\hline\n\mathcal{O}_H = \frac{1}{2} (\partial_\mu |H^2|)^2 \\
\mathcal{O}_W = \frac{ig}{2} (H^\dagger \sigma^a \overleftrightarrow{D}^\mu H) D^\nu W^a_{\mu\nu} \\
\mathcal{O}_B = \frac{ig}{2} (H^\dagger \overleftrightarrow{D}^\mu H) \partial^\nu B_{\mu\nu} \\
\mathcal{O}_{HM} = ig (D^\mu H)^\dagger \sigma^a (D^\nu H) W^a_{\mu\nu} \\
\mathcal{O}_{HM} = ig (D^\mu H)^\dagger \sigma^a (D^\nu H) W^a_{\mu\nu} \\
\mathcal{O}_{V_d} = Y_d |H|^2 \overline{Q}_L H d_R \\
\mathcal{O}_{HH} = ig (D^\mu H)^\dagger (\sigma^a (D^\nu H) W^a_{\mu\nu} \\
\mathcal{O}_{Y_e} = Y_e |H|^2 \overline{L}_L H e_R \\
\mathcal{O}_{HB} = ig' (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu} \\
\mathcal{O}_{3W} = \frac{1}{3!} g \epsilon_{abc} W^a_\mu W^b_\nu W^{c \rho \mu} \\
\hline\n\mathcal{O}_{BB} = g'^2 |H|^2 B_{\mu\nu} B^{\mu\nu}\n\end{array}
$$

Table 6. Operators in the SILH basis considered in our study. Assuming no corrections to electroweak

Strong Dynamics in EFT Language

Strong Dynamics in EFT Language

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Figure 12. The 95% C.L. constraints on the overall coecient of *Oⁱ* from Ref. [38], which are

Conclusion

- $-$ **MCHM:** $\xi = v^2/f^2 < 10^{-3}$, $f > 4$ TeV
- **– ALH/GSILH/SILH** translated into the SILH basis and presented in the form *m*⇤*/* precision from CEPC, ILC, and FCC-ee are used with the inclusion of HL-LHC Higgs precision.

