

2 **Charmed particle production in tau-neutrino CC** 3 **interaction detected by OPERA**

4 **OPERA collaboration**

5 **ABSTRACT:**

6 A peculiar event topology with two secondary decay vertices compatible with short
7 lived particles was found in the analysis of neutrino interactions in the OPERA target.
8 Only few neutrino interactions can yield two heavy particles in the final state: ν_τ CC
9 interaction with charm production or ν NC interaction with $c\bar{c}$ pair production.

10 A dedicated analysis was set-up to identify the underlying process. A new Monte
11 Carlo was developed and several innovative procedures were introduced in the kinematic
12 reconstruction. Multivariate analysis techniques were used to achieve an optimal signal to
13 background separation.

14 Most likely this event is a ν_τ CC interaction with charm production. The significance
15 of this observation is 3.5σ .

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27 1 Introduction

28 Charmed hadron production in neutrino interactions have been studied in two ways: dilep-
29 ton searches in calorimeter detectors [1] and identification of charm decay topologies in
30 nuclear emulsions [2–4]. Emulsion based experiments allow highly detailed reconstruction
31 of the event topology such that a background rejection of the order of 10^{-4} can be achieved.
32 The background arises from pions and kaons decaying in flight or hadron interactions with-
33 out any visible nuclear break-up.

34 The OPERA experiment was designed to observe $\nu_\mu \rightarrow \nu_\tau$ neutrino oscillations in the
35 CNGS beam by the detection of tau leptons produced in ν_τ charged current (CC) inter-
36 actions. The experiment has been searching for neutrino interactions with one secondary
37 short-lived particle as a signature of the τ lepton. In 2015, OPERA reported the discovery
38 of the ν_τ appearance with a significance of 5.1σ [5].

39 An interesting muon-less event with two secondary vertices was identified in the target
40 of the OPERA experiment. Both vertices can be interpreted as heavy particle decay
41 ($c\tau \sim 80 \mu\text{m}$). According to the Standard Model such an event can originate either from
42 ν_τ CC interaction with charm production or from ν NC interaction with $c\bar{c}$ pair production
43 at the CNGS beam energy. The first one has never been directly observed, while the
44 CHORUS experiment observed three charm pairs produced in ν NC interactions [6]. In
45 OPERA the expected number of such events is smaller than one [7].

46 In this paper the analysis and interpretation of the interesting event is reported. After a
47 brief description of the OPERA apparatus (section 2), the event measurement and analysis
48 are reported in sections 3 and 4, respectively. The statistical significance of the observation
49 is discussed in section 5.

2 The OPERA experiment

The OPERA detector was located at the LNGS underground laboratory and was exposed to the CERN neutrino to Gran Sasso (CNGS) beam [8]. The experiment profited of a 730 km long baseline; the average neutrino energy was 17 GeV. The beam exposure started in 2008 and ended in 2012 (1.8×10^{20} p.o.t.). In the target fiducial volume 19 505 neutrino interactions were recorded.

2.1 The apparatus

In order to observe and fully reconstruct decay topologies of short-lived particles a spatial resolution at the micrometer scale is required. The OPERA experiment target was made of lead plates inter-spaced with nuclear emulsion films acting as high accuracy tracking devices [9], called Emulsion Cloud Chamber (ECC). The OPERA target was segmented in 150 000 units (bricks) of 57 nuclear emulsion plates alternating with 1 mm thick lead plates. The brick cross section is $10.2 \text{ cm} \times 12.7 \text{ cm}$; its thickness is 7.5 cm corresponding to about 10 radiation lengths. The brick mass is 8.3 kg. The achieved spatial resolution is of $\sim 1 \mu\text{m}$ and the angular resolution is of $\sim 2 \text{ mrad}$. Charged particle momentum is measured by Multiple Coulomb Scattering (MCS) in the lead plates [10]. Two additional emulsion sheets, called Changeable Sheet doublets (CS), are glued on the brick downstream surface [11].

The active mass of the OPERA target amounted to 1.25 kt. Bricks were hosted in a modular detector made of two identical Super Modules (SM) [9]. Each SM was composed by a target section and a muon spectrometer, as shown in Figure 1. In each SM, the bricks were arranged in 29 vertical walls orthogonal to the beam direction and alternated with electronic detectors [12] consisting of planes of plastic scintillators, called Target Tracker (TT). TT planes were made up of scintillator strips 2.6 cm wide and 1 cm thick arranged perpendicular to each other. The TT system was used to select the brick in which the neutrino interaction occurred. It also provided an estimation of the energy deposited by hadronic and electromagnetic cascades. Spectrometers were designed to measure muon charge and momentum.

2.2 Data taking and event reconstruction

OPERA emulsion analysis is performed by fast automatic scanning systems, based on microscopes equipped with a computer-controlled motorised stage, a dedicated optical system and a camera mounted on top of the optical tube [13–15].

The first step of the event reconstruction is the *location* of the primary neutrino interaction inside the brick [16]. The vertex location procedure in one brick starts from a set of predictions provided by the electronic detectors; then, tracks of secondary particles produced in a neutrino interaction are followed back in the brick, film by film, from the most downstream one to the interaction point where they originate. Whenever a track disappearance hint is detected (the track is not found in three consecutive films), a volume of 1 cm^2 for 5 films upstream and 10 films downstream of the last observed track seg-

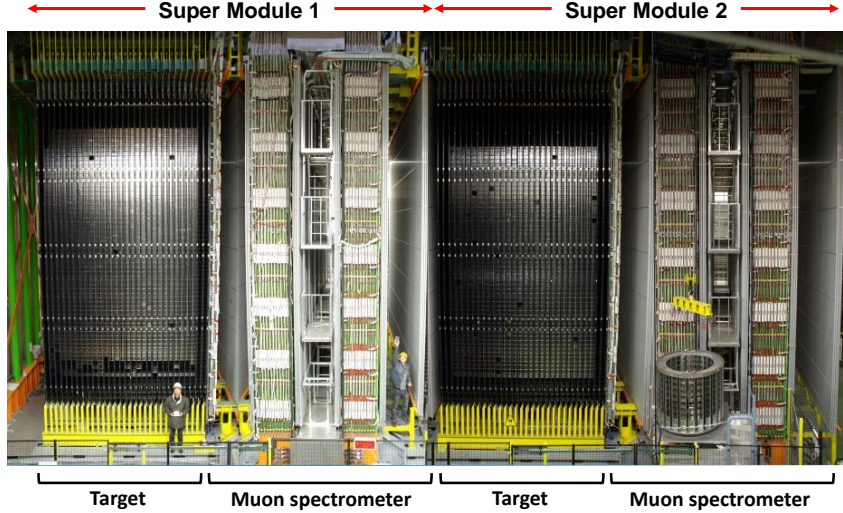


Figure 1: Side view of the OPERA detector; the neutrino beam entered from the left. The upper horizontal lines indicate the position of the two identical supermodules (SM1 and SM2). The target was made of walls filled with bricks interleaved with planes of plastic scintillators (TT).

89 ment origin is scanned in order to fully reconstruct the event and find any decay candidate
 90 through a dedicated decay search procedure.

91 In this *decay search* procedure [4], secondary vertices are searched for by fine manual
 92 measurements of tracks. The hint of a decay topology is the observation of an impact
 93 parameter larger than $10 \mu\text{m}$, defined as the minimum distance between a track and the
 94 reconstructed vertex, excluding low momentum tracks. Decay topologies are classified
 95 either as *short* or *long* decays. Short decays are those with a particle decaying in the same
 96 lead plate where the primary interaction took place; the remaining ones are long decays.
 97 In long decays the signature is the change in slope of a charged particle attached to the
 98 primary vertex, without nuclear recoils or knock-on electrons attached to the secondary
 99 vertex point.

100 The precision obtained in the vertex position is affected by particle scattering through
 101 lead plates evaluated by Monte Carlo (MC) simulations. Figure 2 shows the impact pa-
 102 rameters of reconstructed tracks as a function of the primary vertex depth in lead. If
 103 secondary vertices or decays are found, a full kinematical analysis is performed combining
 104 the measurements in the nuclear emulsion with data from the electronic detectors.

105 The appearance of the τ lepton is identified by the detection of its characteristic decay
 106 topologies, either in one prong (electron, muon or hadron) or in three prongs. Kinematical
 107 selection criteria are applied according to the decay channel [17], shown in Table 1.

108 The detection and reconstruction efficiencies are evaluated by detailed MC simula-
 109 tions [4].

Table 1: Selection criteria for tau candidates. The value denoted by * is used when a reconstructed EM shower is connected to the kink.

variable	$\tau \rightarrow h$	$\tau \rightarrow 3h$	$\tau \rightarrow \mu$	$\tau \rightarrow e$
lepton-tag	No μ or e at the primary vertex			
z_{dec} (μm)	[44; 2600]	< 2600	[44; 2600]	< 2600
p_T^{miss} (GeV/c)	< 1	< 1	–	–
ϕ_{lH} (rad)	> $\pi/2$	> $\pi/2$	–	–
p_T^{2ry} (GeV/c)	> 0.6 (0.3)*	–	> 0.25	> 0.1
p^{2ry} (GeV/c)	> 2	> 3	> 1	> 1
θ_{kink} (mrad)	> 20	< 500	> 20	> 20
m, m_{min} (GeV/c ²)	–	> 0.5 and < 2	–	–

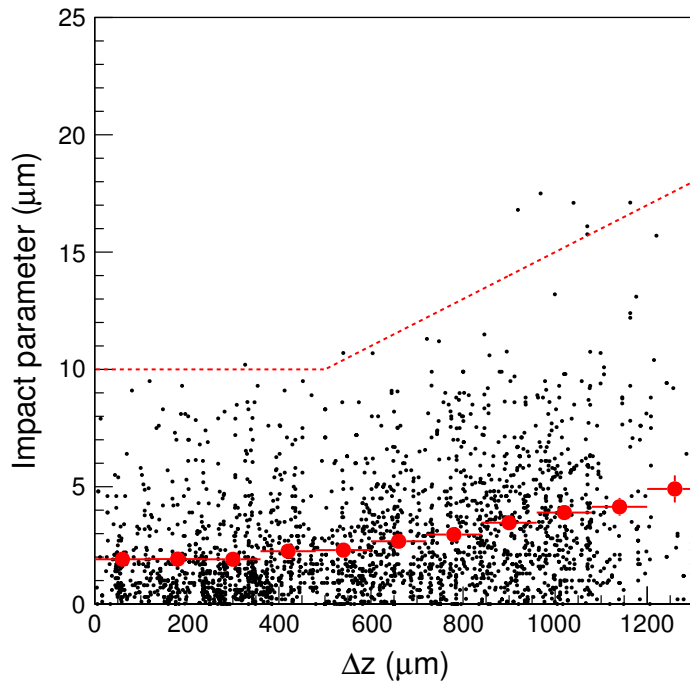


Figure 2: Track impact parameter as a function of the longitudinal distance from the neutrino interaction vertex (2008–2009 data). The red bullets show the average value for each bin. The dotted red line represents the cut applied to select possible short-decay daughter tracks.

110 **3 Event 1114301850**

111 The event 1114301850 was recorded on May 23rd, 2011 in the first SM. The event display
 112 is shown in Figure 3: the number of fired TT planes is 9. No muon track is reconstructed
 113 therefore the event is tagged as *0-muon*. The energy reconstructed by the TT system is
 114 equivalent to 20 ± 6 GeV of hadronic energy released in the lead/emulsion target [12].

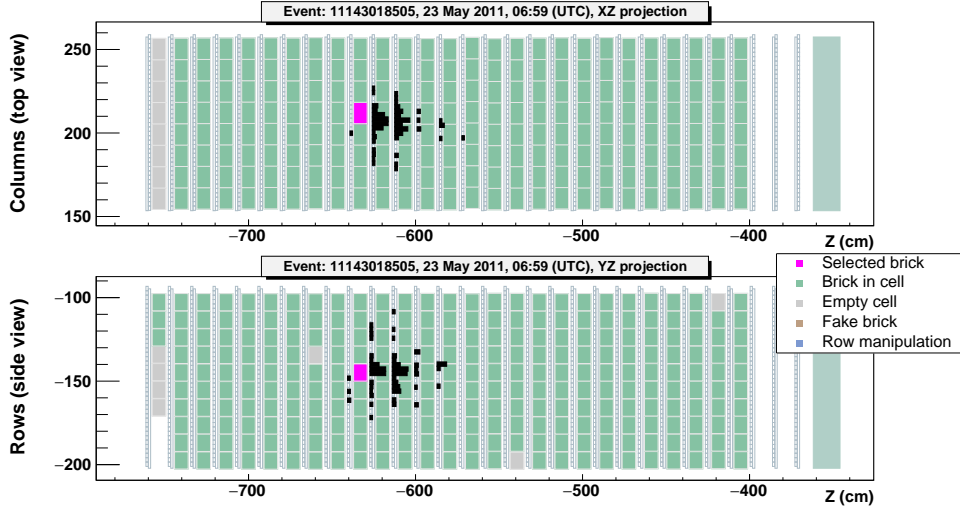


Figure 3: Electronic detector display of the event 1114301850. The neutrino beam direction is along the Z axis entering from the left. Black dots indicates TT hits over threshold and the pink region shows the selected brick (1077152) located in wall 12 of the first SM.

115 3.1 Data acquisition, event building and reconstruction

116 The neutrino interaction occurred in brick 1077152, which was extracted from the apparatus
 117 for the emulsion film development and measurement. The analysis of the CS doublet
 118 reveals a converging pattern of 27 tracks. Out of them, 11 are found also in the brick most
 119 downstream plate (plate 57). These tracks are clustered in a region few hundreds micron
 120 large, hinting to an electromagnetic activity possibly related to the neutrino interaction.
 121 All tracks are followed upstream along the brick; the majority of them are just few emulsion
 122 plates long. By visual inspection they are confirmed as electron-positron pairs candidates.

123 In the location and decay search procedures (see section 2.2), a primary stopping point
 124 of 5 tracks is found in plate 32. The reconstructed angular coordinates are reported for each
 125 track in Table 2. Following these measurements the neutrino interaction point is located
 126 in the lead plate between emulsion plates 31 and 32. Since track 4 impact parameter
 127 (IP) is over the $10 \mu\text{m}$ threshold, a 5-prongs primary vertex topology is discarded. The
 128 most probable topology, based on the particle momenta is a double vertex event with the
 129 primary neutrino vertex (I) formed by tracks 4, 5, 2 and a secondary vertex (II) formed
 130 by tracks 1 and 3. Other configurations have smaller probabilities. Figure 4 shows the
 131 superimposition of several emulsion images taken in plate 32; only grains belonging to the
 132 event are selected.

133 Two additional measurements are performed, both allowing a better resolution than
 134 the standard one: i) manual measurement with a higher magnification objective and ii)
 135 improved automatic image acquisition and analysis. In the first case, the tracks are mea-
 136 sured in plate 32 and 33 under a Zeiss 100x objective mounted on the OPERA microscope;
 137 thus achieving a $0.3 \mu\text{m}$ resolution on the position coordinates (X, Y) [18]. In the second
 138 case an improved scanning procedure, detailed in [19] is applied. The achieved accuracy in

Table 2: Slopes of tracks at plate 32 and their impacts parameters (IPs), evaluated assuming a single vertex topology and a two vertices one. The single vertex, V_0 has coordinates $x = 15083.4 \mu\text{m}$, $y = 59151.5 \mu\text{m}$ and $z = -32999.0 \mu\text{m}$. Vertices V_I and V_{II} are defined in Table 3.

Track	θ_x (rad)	θ_y (rad)	One Vertex IP (μm)		Two vertices IP (μm)	
			w.r.t. V_0		w.r.t. V_I	w.r.t. V_{II}
1	-0.230	-0.275	8.3		36.2	0.1
2	0.121	-0.144	8.8		1.0	6.5
3	0.349	-0.036	4.8		25.9	0.1
4	-0.003	0.088	13.0		1.5	20.4
5	-0.003	-0.025	5.1		2.2	9.6

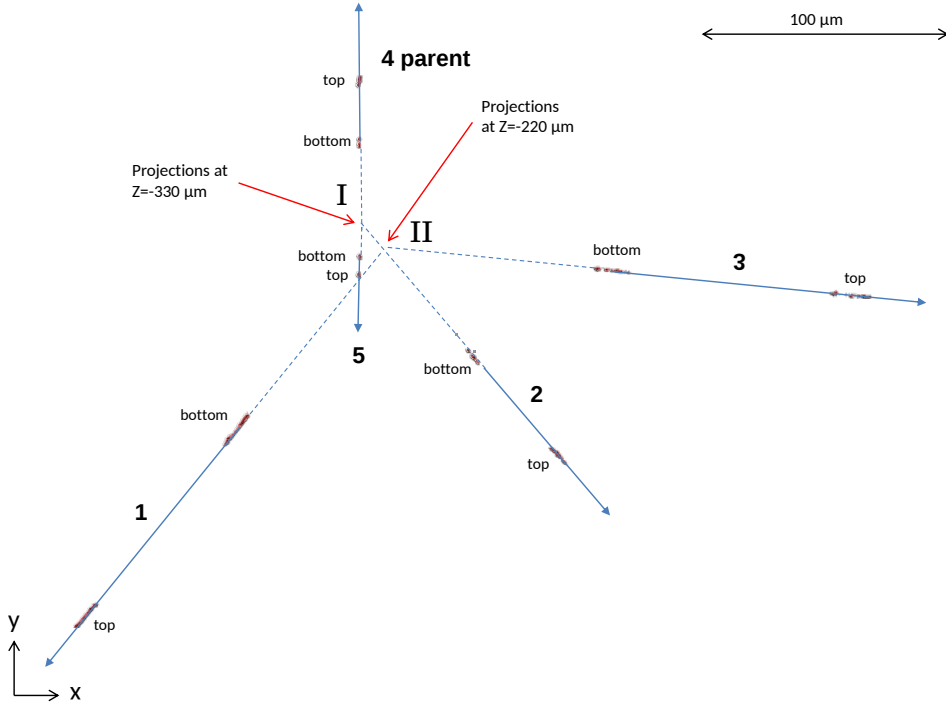


Figure 4: Superimposition of several tomographic emulsion images taken at different depth in plate 32; images are processed to remove fog and Compton electron grains and to show only grains related to the events. Each track is composed by two *groups* of clustered grains because OPERA emulsions have two sensitive layers (*top* and *bottom*).

139 3D cluster finding is better than one micron and a 3D tracking algorithm is applied.

140 The primary vertex is located $581.8 \mu\text{m}$ upstream with respect to the top emulsion layer
 141 of plate 32. The flight length of a neutral particle originated at the primary vertex with
 142 angles $(0.0862; 0.0774)$ mrad and decaying at the secondary vertex, would be $103.2 \mu\text{m}$.

143 Track 4 (*parent*) shows a kink topology (vertex III) between plates 32 and 33. The

144 minimum distance between track 4 and track 6 (*daughter*) emerging from the kink, is
 145 $(0.9 \pm 0.4) \mu\text{m}$. The kink angle is (95 ± 2) mrad; parent flight length is $(1174 \pm 5) \mu\text{m}$.

146 All tracks reconstructed at plates 32 and 33 are followed down in the brick in order
 147 to estimate their momenta and to asses that they belong to the event by confirming their
 148 presence in the CS doublet. A scheme of the full event is shown in Figure 5.

149 Track 2 stops at plate 34; track 3 undergoes a re-interaction at plate 53, while track 1, 5
 150 and 6 reach the CS plates. The coordinates of the three vertices are listed in Table 3.

151 By the decay search procedure two e^+e^- pairs are identified in plates 35 (γ_1) and 41
 152 (γ_2). An innovative procedure developed to identify and reconstruct the electromagnetic
 153 showers [19] is applied. An image data taking with 1 micron Z pitch is performed in a cone
 154 of 400 mrad aperture around the slope of the primary photon, starting from plate 31 down
 155 to plate 57 in the brick. All tracks in the volume are reconstructed using the 3D tracking
 156 algorithm. The main features of the reconstructed showers are listed in Table 4 and shown
 157 in figure Figure 6.

158 The most downstream shower tracks are reconstructed in the CS plates, confirming
 159 that the primary photons are related to the primary neutrino interaction. Their parents
 160 are not relevant for the analysis described in the section 4 which takes into account only the
 161 total visible electromagnetic energy. The accuracy of the reconstruction procedure allows
 162 to establish that γ_1 points to the kink, while γ_2 may emerge from any vertex (*I*, *II* or *III*).

163 A dedicated scanning system was used [21] to search for nuclear fragments at each
 164 vertex, within $|\tan \theta| < 3$ acceptance window. No nuclear fragment was detected.

165 3.2 Kinematics

166 Momenta of tracks 1, 3, 5, 6 are estimated by the MCS method. The alignment accuracy
 167 is evaluated from the angular and position resolution of a sample of penetrating tracks
 168 reconstructed in the scanned volume. All measurements are performed on high resolution
 169 image files. The achieved angular and position resolution is $3.4/\sqrt{2}$ mrad and $0.8/\sqrt{2} \mu\text{m}$
 170 respectively. Results are shown in Table 5.

171 Being track 2 measured only in three emulsion plates, its momentum can not be es-
 172 timated by MCS. An estimation of the momentum is obtained by the NIST code [22]
 173 and yields to $\beta < 0.5$, whatever the mass of particle 2. However, this evaluation is not
 174 compatible with the visible energy loss in emulsion for this track, which is identified as a
 175 m.i.p. by counting the developed grains along the track. Therefore, a different estimation
 176 is performed considering absorption processes which have a resonance at a kinetic energy

Table 3: Position inside the brick of the reconstructed vertices. The z coordinate is evaluated with respect to the downstream side of the brick (plate 57).

Vertex ID	Parent	Daughters	x (μm)	y (μm)	z (μm)
I (primary)	-	2, 4, 5, neutral	15077.0	59157.9	-33081.8
II (secondary)	neutral	1, 3	15085.9	59149.9	-32979.2
III (kink)	4	6	15073.9	59262.4	-31926.4

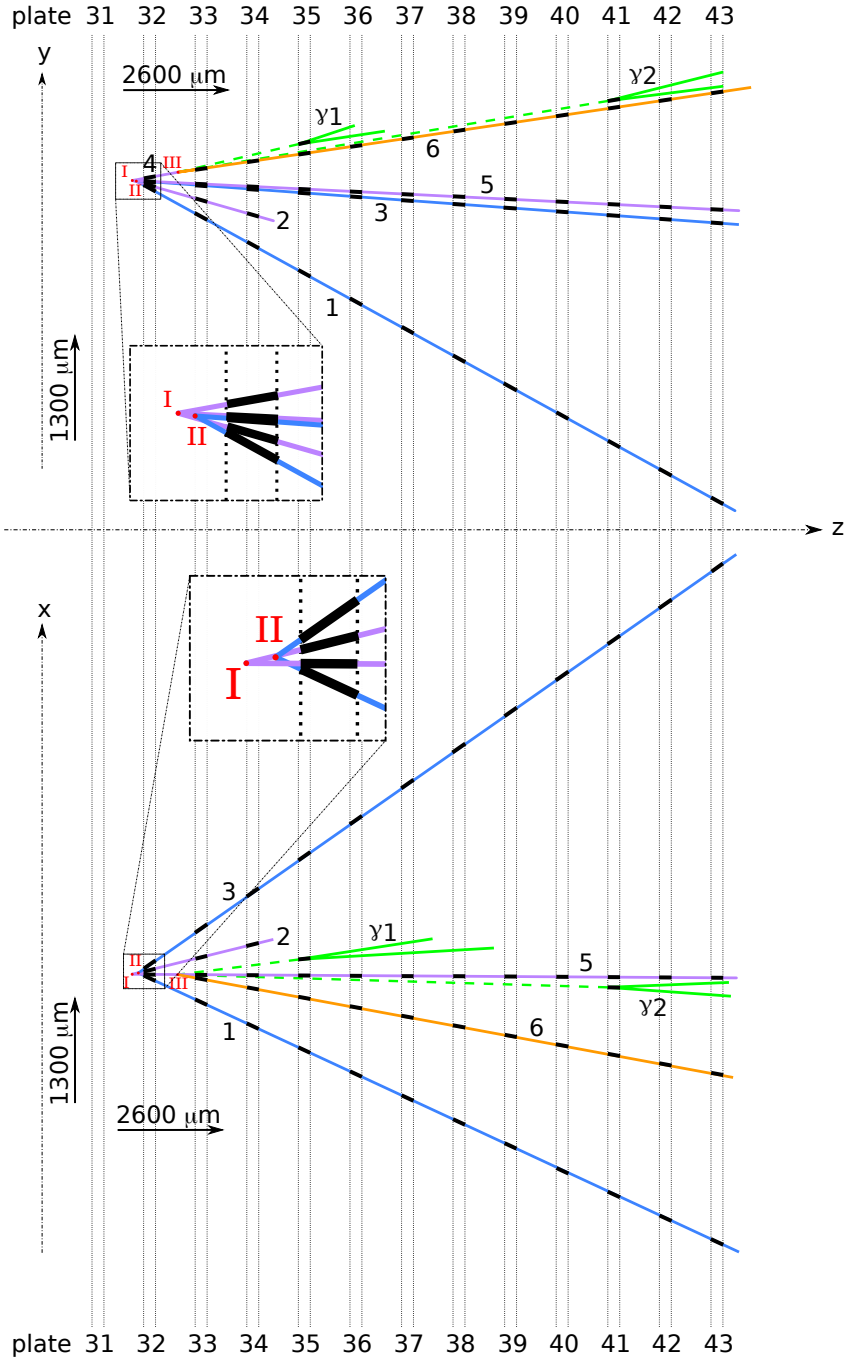


Figure 5: Projection of the event in the YZ (upper plot) and XZ (down plot) planes. Energy deposits measured in a single emulsion films are represented in black, while global reconstructed tracks are represented using colored lines: purple for tracks coming from the primary vertex (I), blue for tracks coming from vertex II and orange for the daughter of vertex III. For photons, only the first electron-positron pair is shown (black).

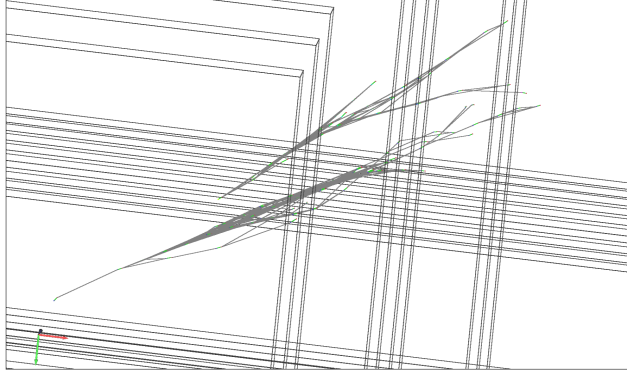


Figure 6: Reconstruction of the electromagnetic showers associated to the event.

177 of about 200 MeV, for pions in every material [23, 24]. In this region, especially for high
 178 A nuclei, the absorption cross section is up to $\sim 40\%$ of the total cross section. From
 179 these arguments, the momentum estimation for track 2 is (0.31 ± 0.08) GeV/ c . This is the
 180 initial momentum of a pion which is absorbed after crossing 2 mm of lead and that has a
 181 kinetic energy of about 200 MeV when absorbed. The uncertainty is evaluated assuming a
 182 uniform kinetic energy distribution: the minimum is the kinetic energy such that $\beta > 0.7$;
 183 while the maximum is 300 MeV, which is the endpoint of the absorption peak.

184 The energies of the electromagnetic showers, γ_1 and γ_2 , are estimated by counting the
 185 electron tracks belonging to each shower. The procedure is calibrated with MC simulations,
 186 taking into account also background tracks in emulsions. The results are $E_{\gamma_1} = (7.2 \pm$
 187 $1.7)$ GeV and $E_{\gamma_2} = (5.3 \pm 2.2)$ GeV.

188 In conclusion, event 1114301850 is identified as a neutrino interaction with two sec-
 189 ondary vertices: vertex *II* has a 2-prong topology while Vertex *III* is a kink originated by
 190 a primary charged particle. The invariant mass at each vertex are reported in Table 6.

Table 4: Electromagnetic showers features. Reconstructed energies were estimated counting the tracks multiplicity [20].

Shower ID	γ_1	γ_2
Starting plate	35	41
θ_x (rad)	0.050	0.011
θ_y (rad)	0.122	0.085
IP _I (μm)	30 ± 22	40 ± 23
IP _{II} (μm)	28 ± 22	40 ± 23
IP _{III} (μm)	8 ± 8	40 ± 11
Opening angle (rad)	0.027	0.029
Energy (GeV)	7.1 ± 1.7	5.3 ± 2.2

Table 5: Particles momenta reconstructed by the multiple Coulomb scattering method.

Track ID	p best fit (GeV/ c)	68 % p range (GeV/ c)
1	2.1	[1.6 ; 3.1]
3	4.3	[3.1 ; 7.1]
5	0.54	[0.45 ; 0.68]
6 (daughter)	2.7	[2.1 ; 3.7]

Table 6: Secondary vertices invariant masses and minimum invariant masses. The differences are evaluated too, $\Delta \equiv M_{min} - M$. Due to the correlation between the two distributions, the errors are relatively small.

Vertex ID	Invariant Mass	Minimum Invariant Mass	Difference
	M (GeV/ c^2)	M_{min} (GeV/ c^2)	Δ (GeV/ c^2)
II	1.8 ± 0.5	2.5 ± 0.8	0.7 ± 0.4
III (kink)	0.9 ± 0.1	1.2 ± 0.2	0.3 ± 0.1

191 4 Event analysis

192 The standard OPERA analysis does not include events with double-decay topology. As-
 193 sociated charm production is taken into account only as a background in the tau search,
 194 assuming the identification of just one decay vertex.

195 According to the standard OPERA analysis neither of the two decays can be classified
 196 as a tau candidate [7, 17]. In particular, the vertex III (kink) matches all selection criteria
 197 except for the daughter transverse momentum (p_T) which should be > 0.300 GeV/ c , while
 198 the event kink daughter has $p_T = (0.24 \pm 0.07)$ GeV/ c .

199 Therefore a non-standard new analysis is necessary for a more accurate classification
 200 of event 1114301850. In CNGS kinematic conditions, two short-lived particle decays can
 201 be produced by the following processes(Figure 7):

- 202 • ν_τ CC interaction with charm production;
- 203 • ν NC interaction with $c\bar{c}$ pair production;

204 Other processes faking this topology are:

- 205 • ν_μ CC interaction with a mis-identified muon and two secondary interactions.
- 206 • ν_μ CC interaction with single charm production, a mis-identified muon and one sec-
 207 ondary interaction;
- 208 • ν NC interaction with two secondary interactions;
- 209 • ν_τ CC interaction with one secondary interaction;

210 A secondary interaction can be either i) hadronic interaction of a final state particle,
 211 ii) short decay of pions or kaons, or iii) large angle Coulomb scattering by hadrons or
 212 mis-identified muons.

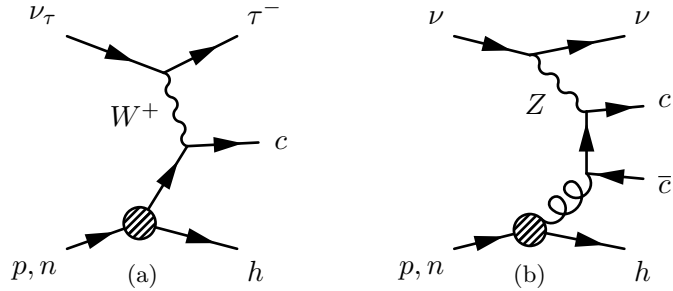


Figure 7: Leading Feynman diagrams for the production of two prompt short decaying particles: (a) tau charm production in CC interaction and (b) charm pair production in NC interaction .

213 The new analysis is based on the distributions of kinematical variables obtained through
 214 a dedicated MC production. Neutrino interactions are generated using GENIE [25] besides
 215 charm pair production, which is simulated using HERWIG [26]. Due to the high multi-
 216 plicity of event 1114301850, only DIS interactions are taken into account. In total, about
 217 300 million events are generated.

218 Particles from neutrino interactions are propagated in few cubic centimeter volumes of
 219 the OPERA brick using the Geant4 framework [27, 28], assuming that the primary vertex
 220 has the same depth in lead as the one estimated for event 1114301850. The MCS is taken
 221 into account using a parametrization based on the standard OPERA MC. The hadronic
 222 interaction simulations are validated with dedicated test beam data [29].

223 For each process, the number of expected events is normalized to the 12 352 observed
 224 CC events with a primary vertex in the target section of the OPERA experiment con-
 225 sidering the shape of the CNGS neutrino flux [8], the oscillation probabilities and the
 226 cross section. The vertex location efficiency is determined according to a data driven pa-
 227 rameterization. The efficiencies related to the electronic detectors (brick selection, muon
 228 identification, muon momentum estimation) are evaluated using the standard OPERA MC
 229 with a parameterization based on the hadronic energy and muon momentum.

230 Simulated events are selected regardless of the multiplicity at the primary vertex by
 231 requiring:

- 232 • no muon nor electron reconstructed at the primary vertex;
- 233 • a one prong-like secondary vertex (1pr-like);
- 234 • a two prong-like secondary vertex (2pr-like);
- 235 • no fragments at any vertex.

236 Moreover, the 1pr-like daughter has to be a charged track but not a muon, neither an
 237 electron nor a positron. The 1pr-like parent is required to be charged and crossing at least
 238 one emulsion plate. No kinematical cuts are applied.

239 The total number of expected events matching the topology of event 1114301850 is
 240 ~ 0.1 ; the details for each simulated process are given in Table 7.

Table 7: Expected events with a two secondary-vertices topology as selected by the analysis.

Sample	Expected events (10^{-3})
ν_τ CC + charm	44.5 ± 0.1
ν NC + $c\bar{c}$ pair	12.59 ± 0.02
ν_μ CC + two 2ry	4.0 ± 0.5
ν_μ CC + charm + 2ry	20.5 ± 0.5
ν NC + two 2ry	3.8 ± 0.3
ν_τ CC + 2ry	9.0 ± 0.1
Total	94.4

241 A multivariate analysis is applied on selected events and the signal to back-
 242 ground discrimination is based on 12 kinematic variables:

- 243 • the *total EM energy*, that is the sum of the visible photon energy, regardless of the
 244 photon origin vertex;
- 245 • the *transverse angle* φ between the parents of the 1pr-like and 2pr-like vertices;
- 246 • the *missing transverse momentum* at primary vertex with respect to the beam direc-
 247 tion;
- 248 • the *hadronic momentum*, i.e. the sum of the primary track momenta excluding the
 249 two parents;

250 for the 1pr-like vertex:

- 251 • the *daughter's momentum*;
- 252 • the *daughter's transverse momentum* with respect to the parent direction;
- 253 • the *flight length*.
- 254 • the *kink* angle between parent and daughter;

255 for the 2pr-like vertex:

- 256 • the total *daughters' momentum*;
- 257 • the total *daughters' transverse momentum* with respect to the parent direction;
- 258 • the *flight length*.
- 259 • the *invariant mass* of the charged daughters.

260 In order to find the best method for the discrimination, several algorithms are tested:
 261 an Artificial Neural Networks (ANN) method [30], two kinds of Boost Decision Trees [31]
 262 and the Fisher Discriminant [32]. The optimal one turns out to be the ANN, whose output
 263 variable distribution is shown in Figure 8.

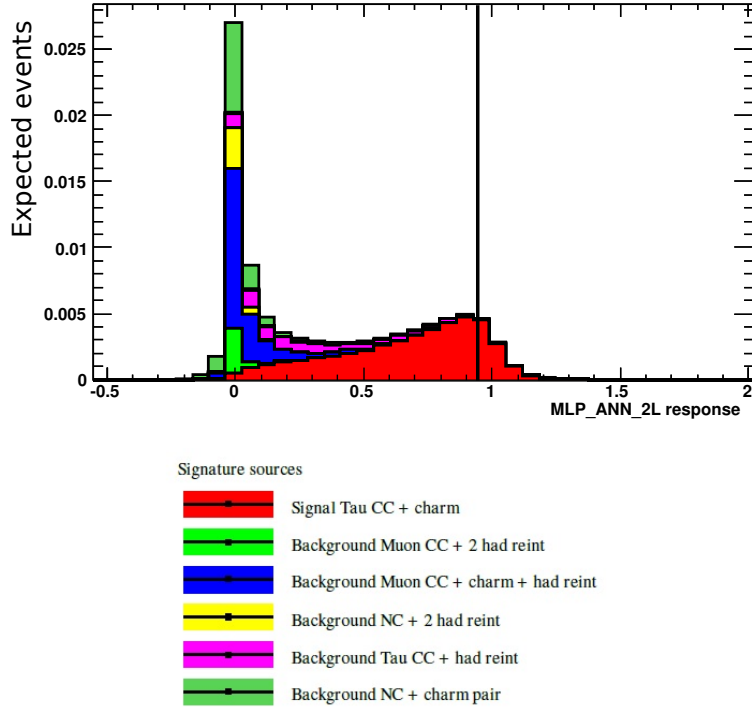


Figure 8: Distribution of the ANN output variable. The weighted contribution of each source of Table 7 is shown with a different color. The vertical black line represents the ANN output for the event 1114301850.

264 5 Significance

265 As shown in Figure 8, according to the multivariate analysis the event can be classified as
 266 a ν_τ CC interaction with charm production with rather high probability. The significance
 267 of the result is evaluated using RooFit/RooStats libraries [33] provided by the ROOT
 268 framework [34].

269 The observable of the model is the ANN output variable x whose distribution, nor-
 270 malized to the expected number of events, is shown in Figure 8. Its shape is obtained by
 271 the sum of the contributions of each process reported in Table 7. In order to evaluate the
 272 significance, a parameter μ , the signal strength, is introduced such that the background
 273 only hypothesis corresponds to $\mu = 0$, while $\mu = 1$ represents the background plus the
 274 expected signal model. The likelihood can be written as:

$$275 \quad \mathcal{L}(\mu|x) = \sum_i n_i b_i(x) + \mu \cdot n_s s(x) \quad (5.1)$$

276 where $s(x)$ and $b_i(x)$ are the signal and backgrounds PDFs, respectively; n_i and n_s the
 277 expected number of events.

278 Systematic effects are introduced as scale factors, f_j , different for each process. Scale
 279 factors f_j depend on nuisance parameters σ_k distributed according to their PDFs, g_k .
 280 Some nuisance parameters are specifically defined for a particular process, while others

281 are common to several contributions. The first nuisance parameter is a normalization fac-
 282 tor, σ_N . The normalization is dominated by the CNGS flux uncertainty, which is about
 283 20 % [8]. Another source of uncertainty are cross sections of two processes not observed
 284 in the OPERA experiment: NC interactions with charm pair production and ν_τ CC inter-
 285 actions with single charm. These uncertainties are assumed to be 20 %. The systematic
 286 cross section error of ν_τ CC interactions without charm estimated from [35] is 6 % for ν_τ
 287 in the few tens of GeV range. The uncertainty associated to the hadron re-interaction is
 288 30 %, according to data-MC comparisons based on test beam results [29].

289 Each nuisance parameter distribution g_k depends on some constant parameters as the
 290 range boundaries or other PDF parameters like Gaussian variances. These are different for
 291 each g_k and they are labeled $\hat{\sigma}_k$. Hence, nuisance parameters distributions are identified
 292 as $g_k(\sigma_k|\hat{\sigma}_k)$.

293 Including systematics, the likelihood can be expressed as:

$$294 \quad \mathcal{L}(\mu, \boldsymbol{\sigma}|x) = \left[\sum_{i \in B} f_i(\boldsymbol{\sigma}) n_i b_i(x) + \mu f_s(\boldsymbol{\sigma}) n_s s(x) \right] \prod_k g_k(\sigma_k|\hat{\sigma}_k) \quad (5.2)$$

295 where $\boldsymbol{\sigma}$ indicates the whole set of nuisance parameters.

296 The test statistic used for the significance evaluation is the profile likelihood ratio [36,
 297 37]. A sample of MC pseudo-experiments is generated according to the background only
 298 PDF, in order to get the test statistic distribution (Figure 9). In each pseudo-experiment,
 299 nuisance parameters are varied according to their PDFs. Under the background only
 300 hypothesis, the probability of data being less likely or equal to the observed event is $(2.6 \pm$
 301 $0.2) \times 10^{-4}$. Therefore, the absence of ν_τ CC interaction with charm production can be
 302 excluded with a significance of 3.5 standard deviations.

303 The most likely interpretation is that vertex II is originated by a charm decay and
 304 vertex III by tau decay into an hadron.

305 6 Conclusions

306 A neutrino interaction event was observed in the OPERA target having a rare topology:
 307 two secondary vertices within about 1 mm from the primary one were observed. Such
 308 topology could arise from ν_τ interaction with charm production and ν NC interaction with
 309 double charm production. Dedicated scanning and analysis procedures were thus performed
 310 for this event that was not considered in the original OPERA proposal.

311 The event turned out to be very likely a ν_τ CC interaction with charm production
 312 with a significance of about 3.5σ .

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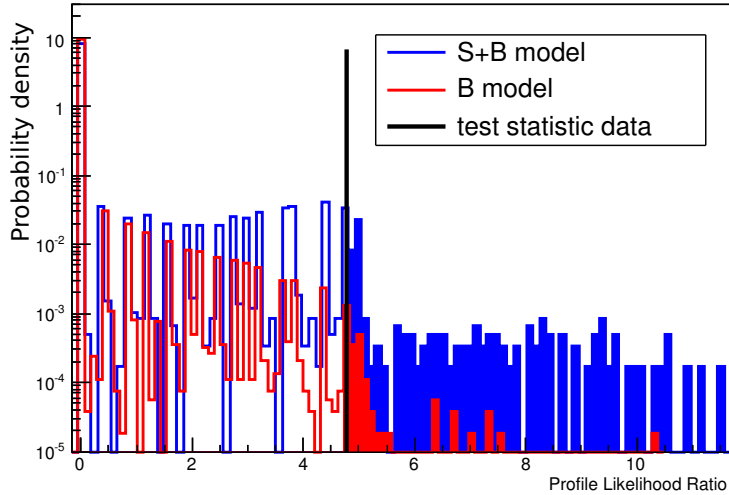


Figure 9: Expected distributions of the test statistic (profile likelihood ratio) obtained by pseudo-experiments under the background only hypothesis (B model, $\mu = 0$) and signal plus background hypothesis ($S+B$ model, $\mu = 1$). The black line indicates the value obtained by applying the test statistic on the observed data.

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