

# Big Bang Nucleosynthesis-Post Planck

- BBN and the WMAP/Planck determination of  $\eta$ ,  $\Omega_B h^2$
- Observations and Comparison with Theory
  - D/H
  - $^4\text{He}$
  - $^7\text{Li}$
- The Li Problem
- Neutrinos

# Historical Perspective

## Intimate connection with CMB

Alpher  
Herman  
Gamow

## Conditions for BBN:

Require  $T > 100 \text{ keV} \Rightarrow t < 200 \text{ s}$

$$\sigma v(p + n \rightarrow D + \gamma) \approx 5 \times 10^{-20} \text{ cm}^3/\text{s}$$

$$\Rightarrow n_B \sim 1/\sigma v t \sim 10^{17} \text{ cm}^{-3}$$

Today:

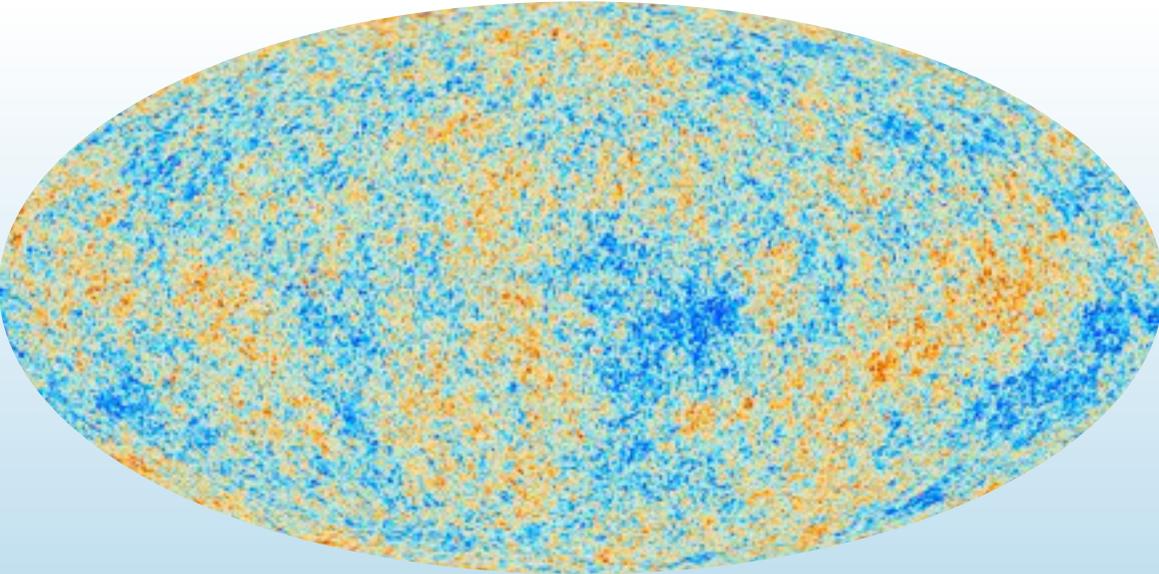
$$n_{B_0} \sim 10^{-7} \text{ cm}^{-3}$$

and

$$n_B \sim R^{-3} \sim T^3$$

Predicts the CMB temperature

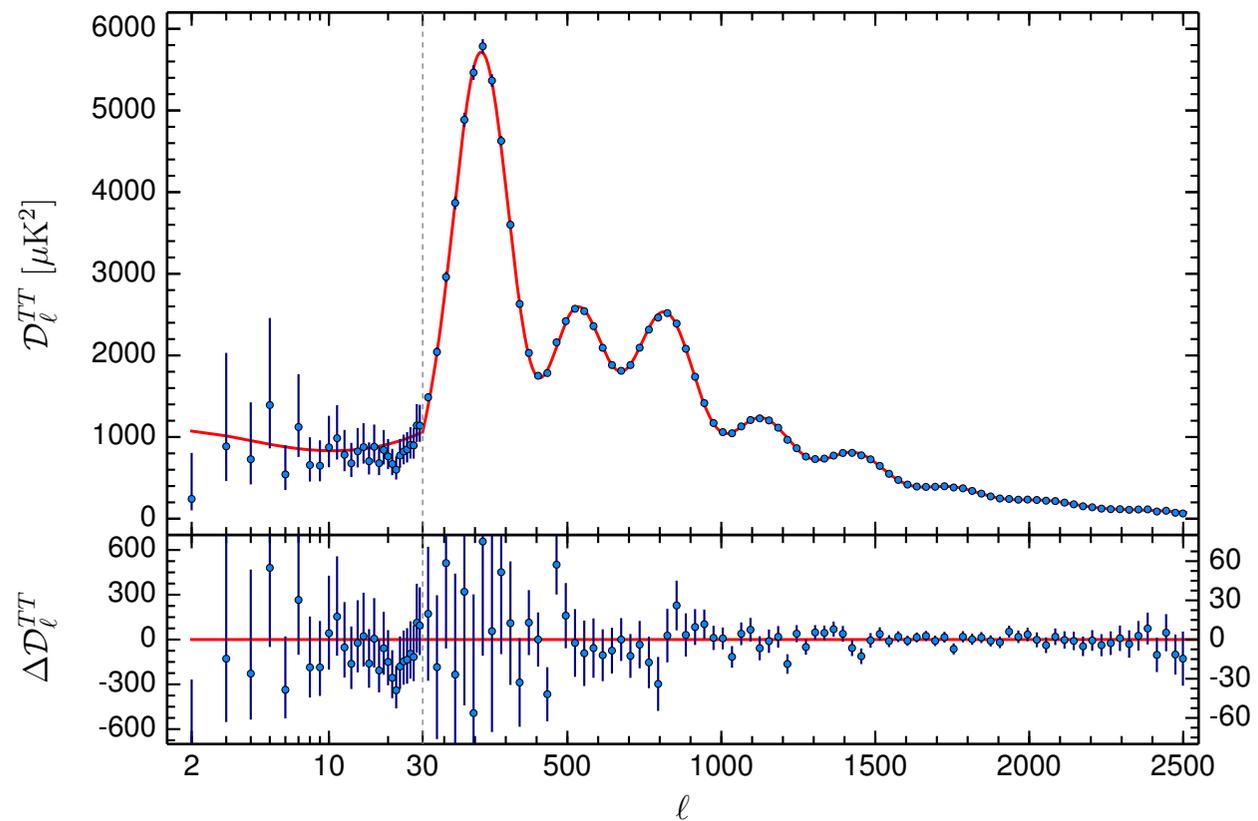
$$T_0 = (n_{B_0} / n_B)^{1/3} T_{\text{BBN}} \sim 10 \text{ K}$$

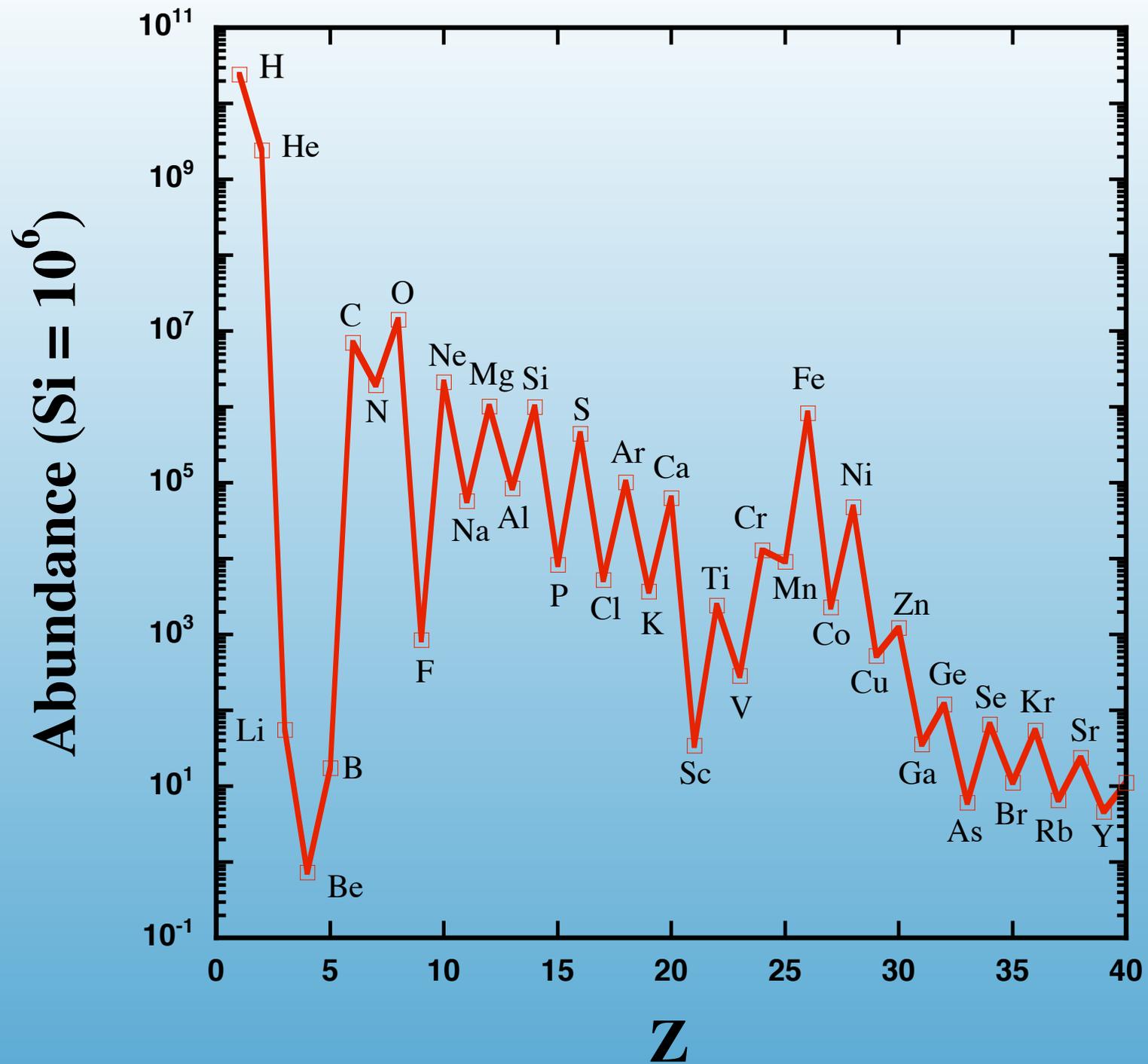


# Planck best fit

$$\Omega_B h^2 = 0.02226 \pm 0.00023$$

$$\eta_{10} = 6.09 \pm 0.06$$





BBN could not explain the abundances (or patterns) of all the elements.

⇒ growth of stellar nucleosynthesis

But,

Questions persisted:

25% (by mass) of  $^4\text{He}$  ?  
D?

Resurgence:

BBN could successfully account for the abundance of

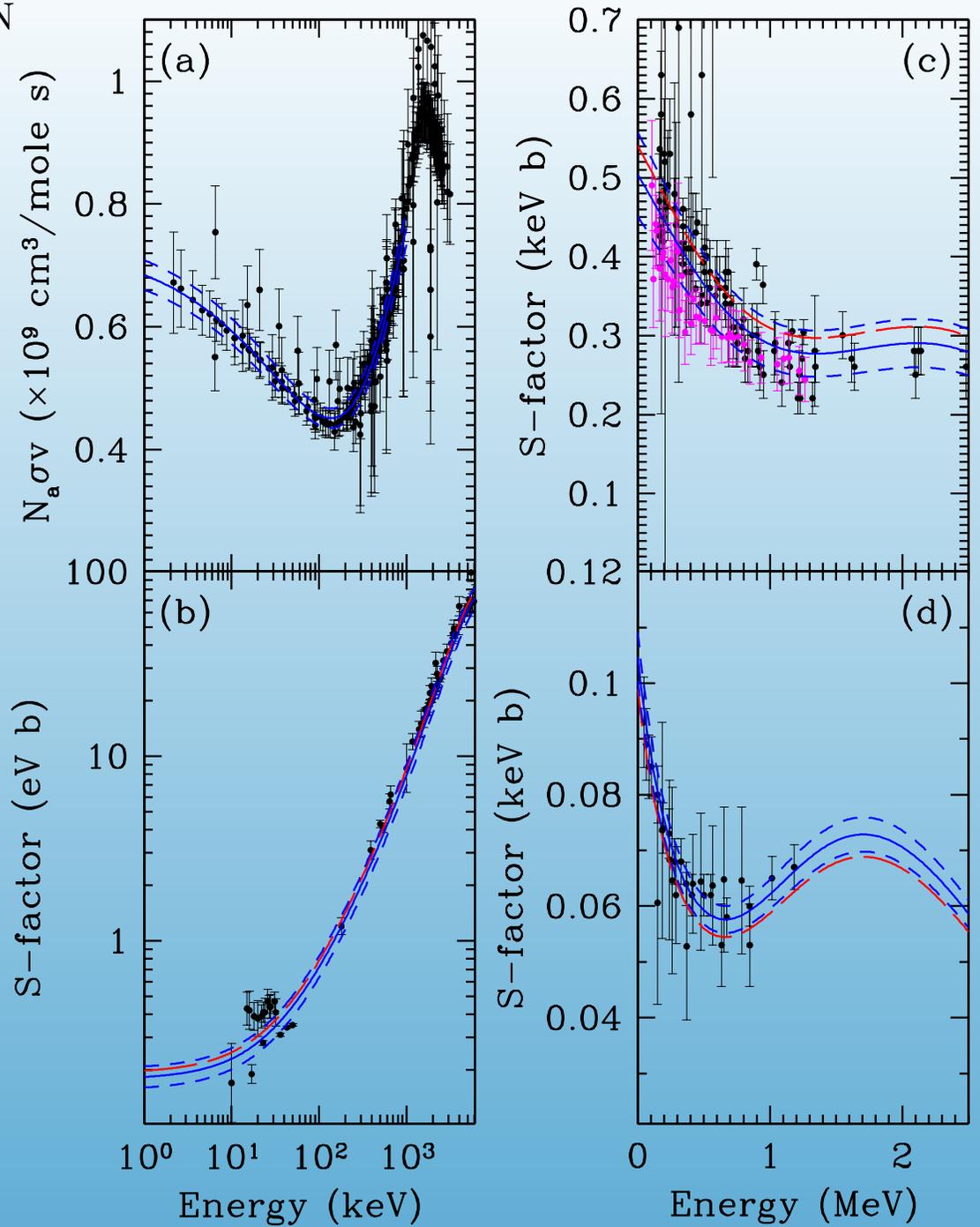
D,  $^3\text{He}$ ,  $^4\text{He}$ ,  $^7\text{Li}$ .



Table 1: Key Nuclear Reactions for BBN

Source	Reactions	
NACRE	$d(p, \gamma)^3\text{He}$	(b)
	$d(d, n)^3\text{He}$	
	$d(d, p)t$	
	$t(d, n)^4\text{He}$	
	$t(\alpha, \gamma)^7\text{Li}$	(d)
	$^3\text{He}(\alpha, \gamma)^7\text{Be}$	(c)
SKM	$^7\text{Li}(p, \alpha)^4\text{He}$	
	$p(n, \gamma)d$	
	$^3\text{He}(d, p)^4\text{He}$	
This work	$^7\text{Be}(n, p)^7\text{Li}$	
	$^3\text{He}(n, p)t$	(a)
PDG	$\tau_n$	

NACRE  
 Cyburt, Fields, KAO  
 Nollett & Burles  
 Coc et al.



# Conditions in the Early Universe:

$$T \gtrsim 1 \text{ MeV}$$

$$\rho = \frac{\pi^2}{30} \left( 2 + \frac{7}{2} + \frac{7}{4} N_\nu \right) T^4$$

$$\eta = n_B/n_\gamma \sim 10^{-10}$$

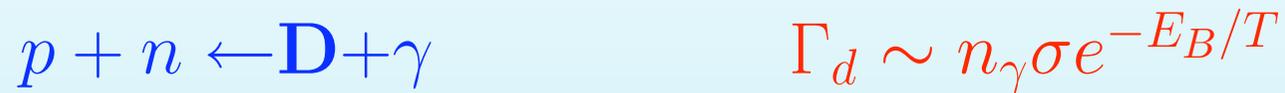
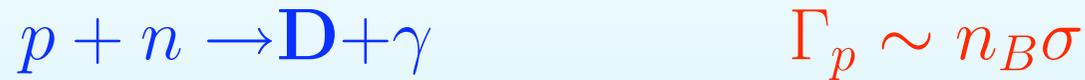
$\beta$ -Equilibrium maintained by weak interactions

Freeze-out at  $\sim 1 \text{ MeV}$  determined by the competition of expansion rate  $H \sim T^2/M_p$  and the weak interaction rate  $\Gamma \sim G_F^2 T^5$



At freezeout  $n/p$  fixed modulo free neutron decay,  $(n/p) \simeq 1/6 \rightarrow 1/7$

# Nucleosynthesis Delayed (Deuterium Bottleneck)



Nucleosynthesis begins when  $\Gamma_p \sim \Gamma_d$

$$\frac{n_\gamma}{n_B} e^{-E_B/T} \sim 1 \quad @ T \sim 0.1 \text{ MeV}$$

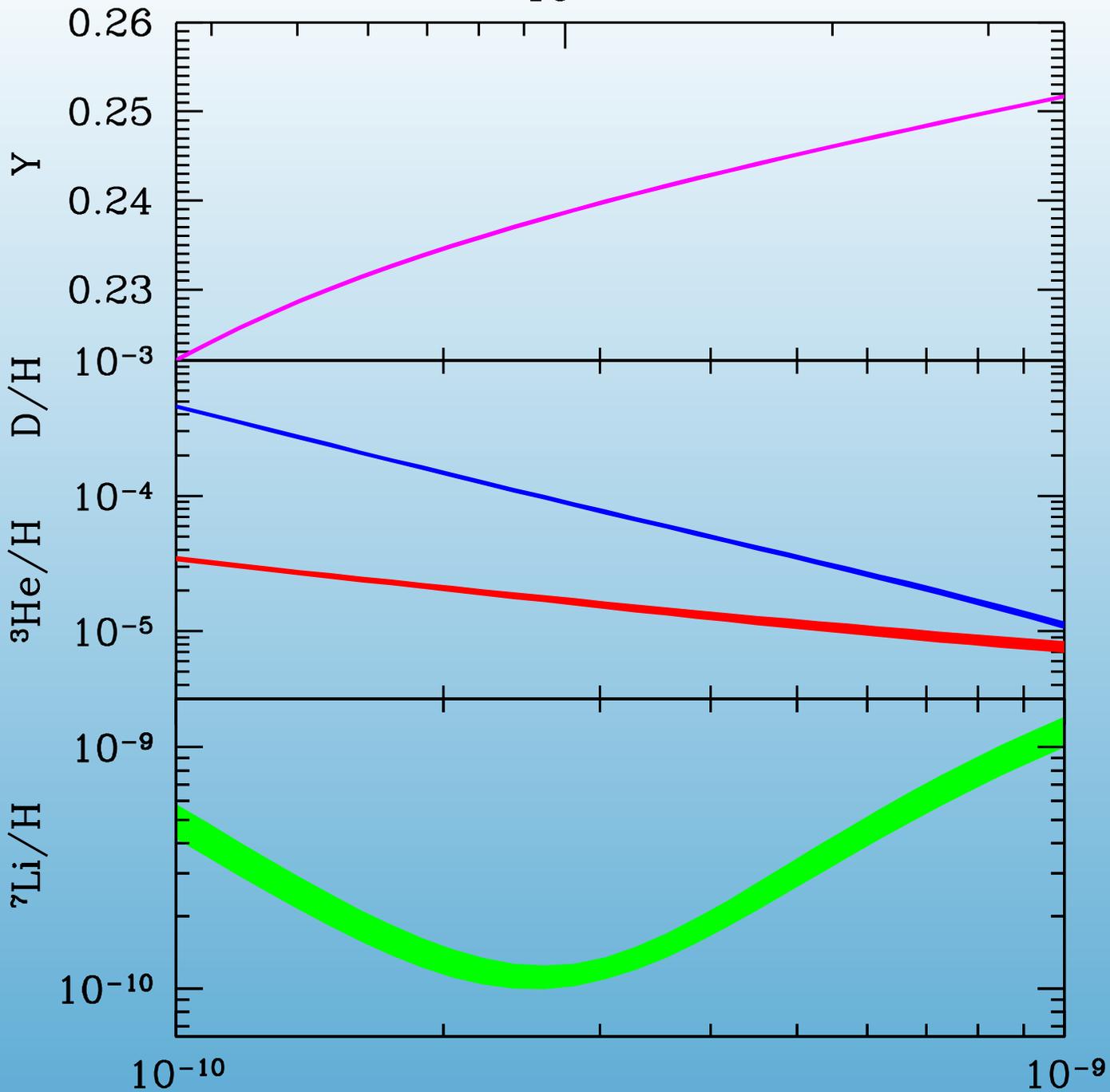
All neutrons  $\rightarrow$   ${}^4\text{He}$

$$Y_p = \frac{2(n/p)}{1 + (n/p)} \simeq 25\%$$

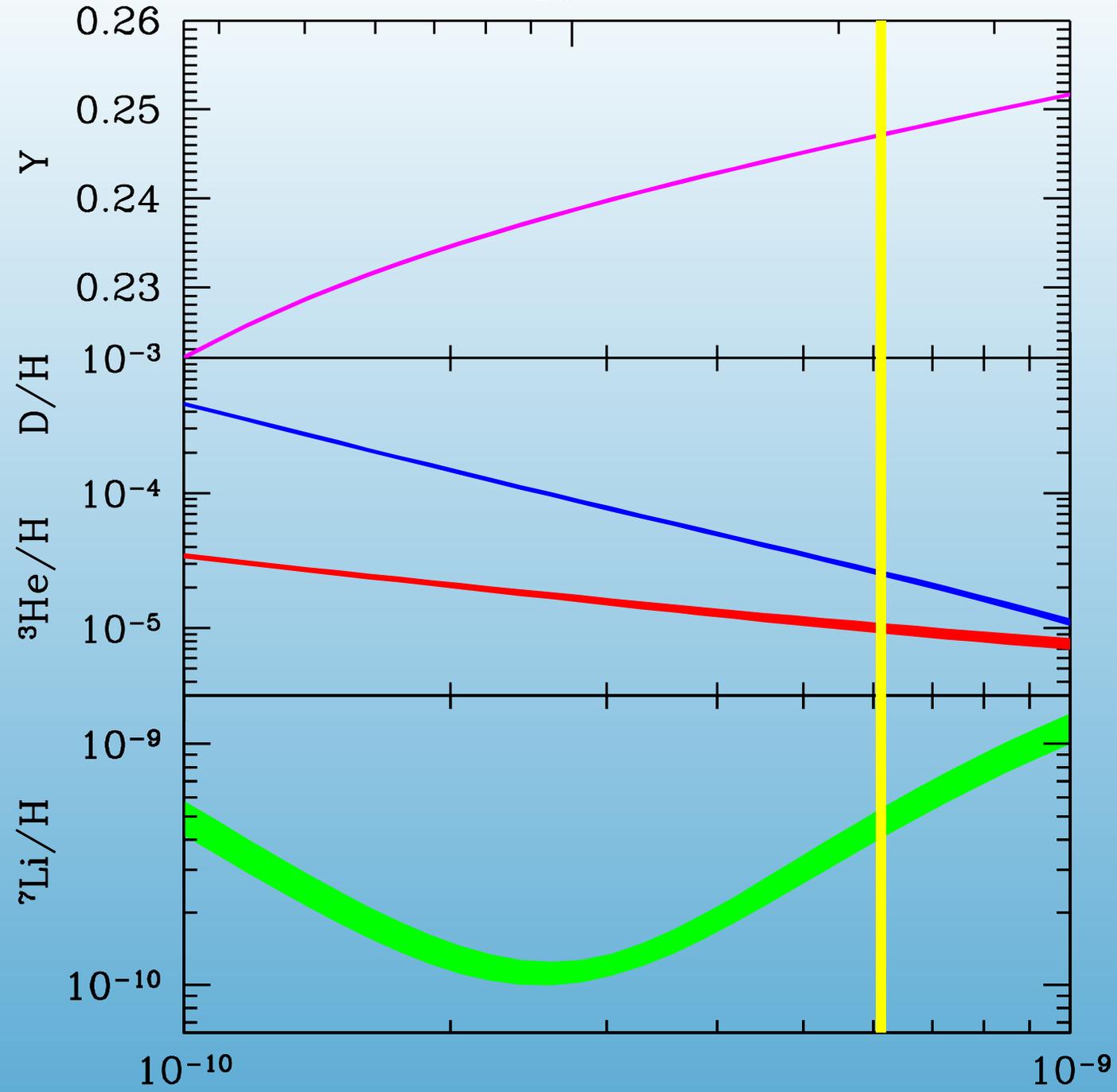
Remainder:

$\text{D}$ ,  ${}^3\text{He} \sim 10^{-5}$  and  ${}^7\text{Li} \sim 10^{-10}$  by number

baryon density  $\Omega_b h^2$   
 $10^{-2}$



baryon density  $\Omega_b h^2$   
 $10^{-2}$



# Big Bang Nucleosynthesis

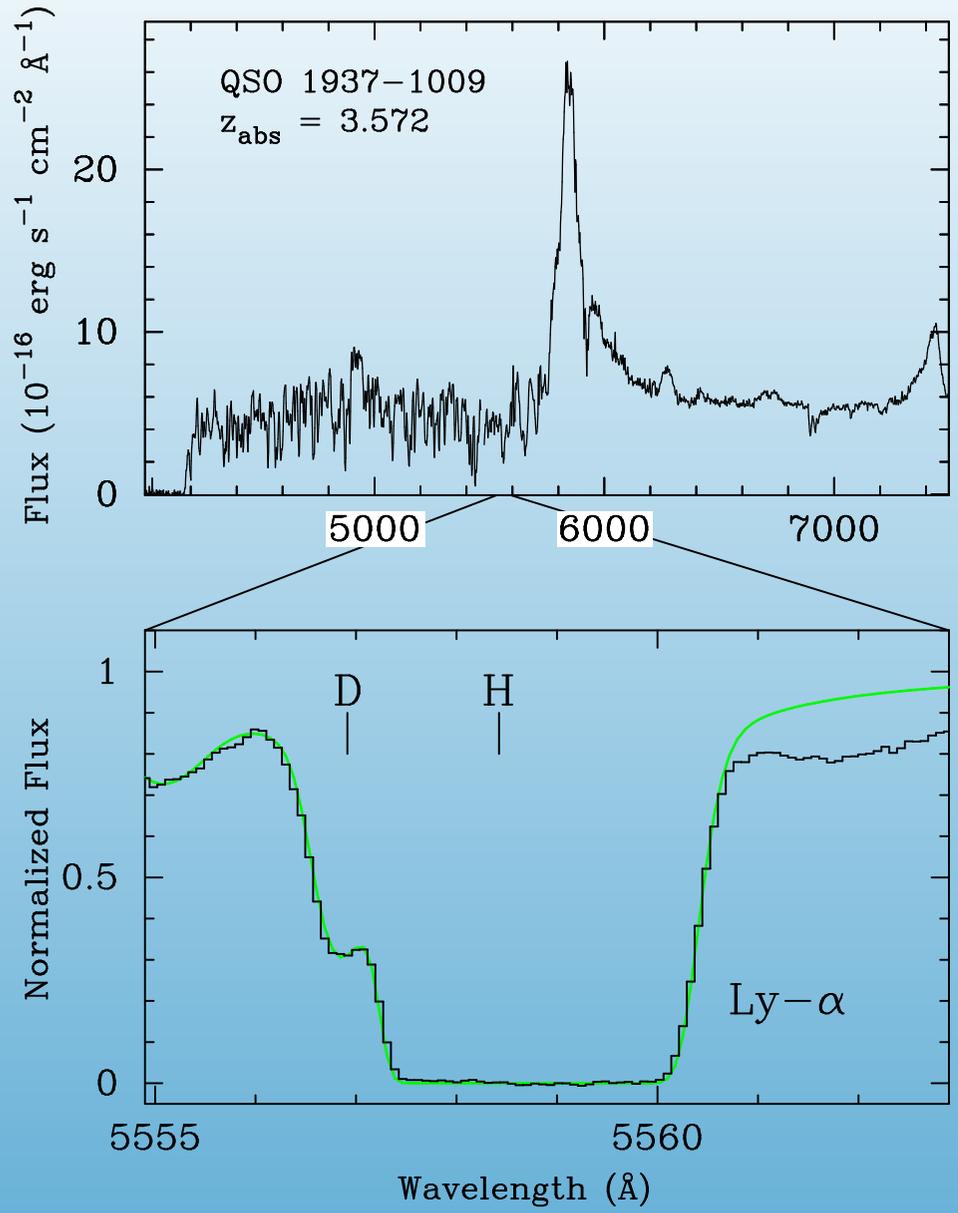
- Production of the Light Elements: D,  $^3\text{He}$ ,  $^4\text{He}$ ,  $^7\text{Li}$ 
  - $^4\text{He}$  observed in extragalactic HII regions:  
abundance by mass = 25%
  - $^7\text{Li}$  observed in the atmospheres of dwarf halo stars:  
abundance by number =  $10^{-10}$
  - D observed in quasar absorption systems (and locally):  
abundance by number =  $3 \times 10^{-5}$
  - $^3\text{He}$  in solar wind, in meteorites, and in the ISM:  
abundance by number =  $10^{-5}$

# D/H

- All Observed D is Primordial!
- Observed in the ISM and inferred from meteoritic samples (also HD in Jupiter)
- D/H observed in Quasar Absorption systems

THE PRECISION SAMPLE OF D/H MEASUREMENTS IN QSO ABSORPTION LINE SYSTEMS

QSO	$z_{\text{em}}$	$z_{\text{abs}}$	[O/H] <sup>a</sup>	Literature		This work		Ref. <sup>b</sup>
				$\log N(\text{H I})$ ( $\text{cm}^{-2}$ )	$\log(\text{D}/\text{H})$	$\log N(\text{H I})$ ( $\text{cm}^{-2}$ )	$\log(\text{D}/\text{H})$	
HS 0105+1619	2.652	2.53651	-1.77	$19.42 \pm 0.01$	$-4.60 \pm 0.04$	$19.426 \pm 0.006$	$-4.589 \pm 0.026$	1, 2
Q0913+072	2.785	2.61829	-2.40	$20.34 \pm 0.04$	$-4.56 \pm 0.04$	$20.312 \pm 0.008$	$-4.597 \pm 0.018$	1, 3, 4
SDSS J1358+6522	3.173	3.06726	-2.33	...	...	$20.495 \pm 0.008$	$-4.588 \pm 0.012$	1
SDSS J1419+0829	3.030	3.04973	-1.92	$20.391 \pm 0.008$	$-4.596 \pm 0.009$	$20.392 \pm 0.003$	$-4.601 \pm 0.009$	1, 5, 6
SDSS J1558-0031	2.823	2.70242	-1.55	$20.67 \pm 0.05$	$-4.48 \pm 0.06$	$20.75 \pm 0.03$	$-4.619 \pm 0.026$	1, 7



# D/H abundances in Quasar absorption systems

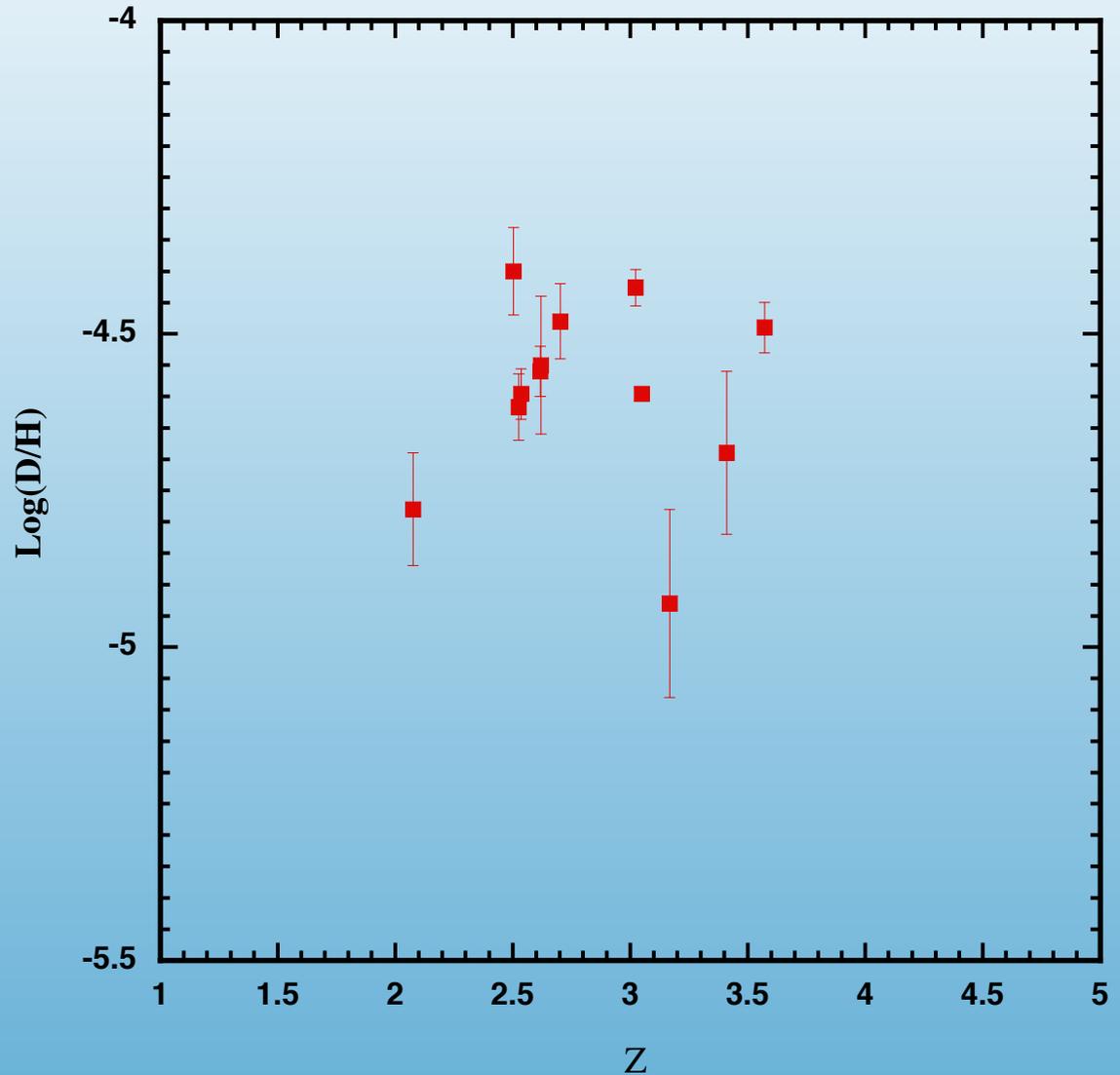
BBN Prediction:

$$10^5 D/H = 2.58 \pm 0.13$$

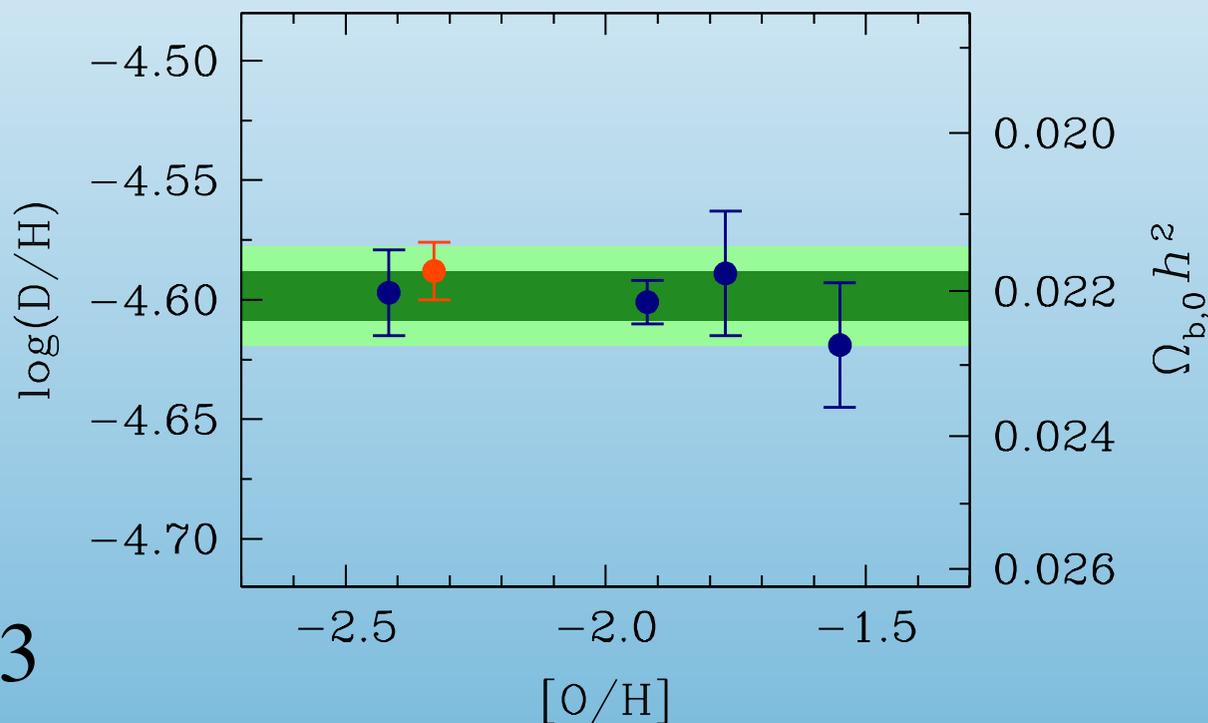
Obs Average:

$$10^5 D/H = 3.01 \pm 0.21$$

(0.68 sample variance)



# D/H abundances in Quasar absorption systems



BBN Prediction:

$$10^5 D/H = 2.58 \pm 0.13$$

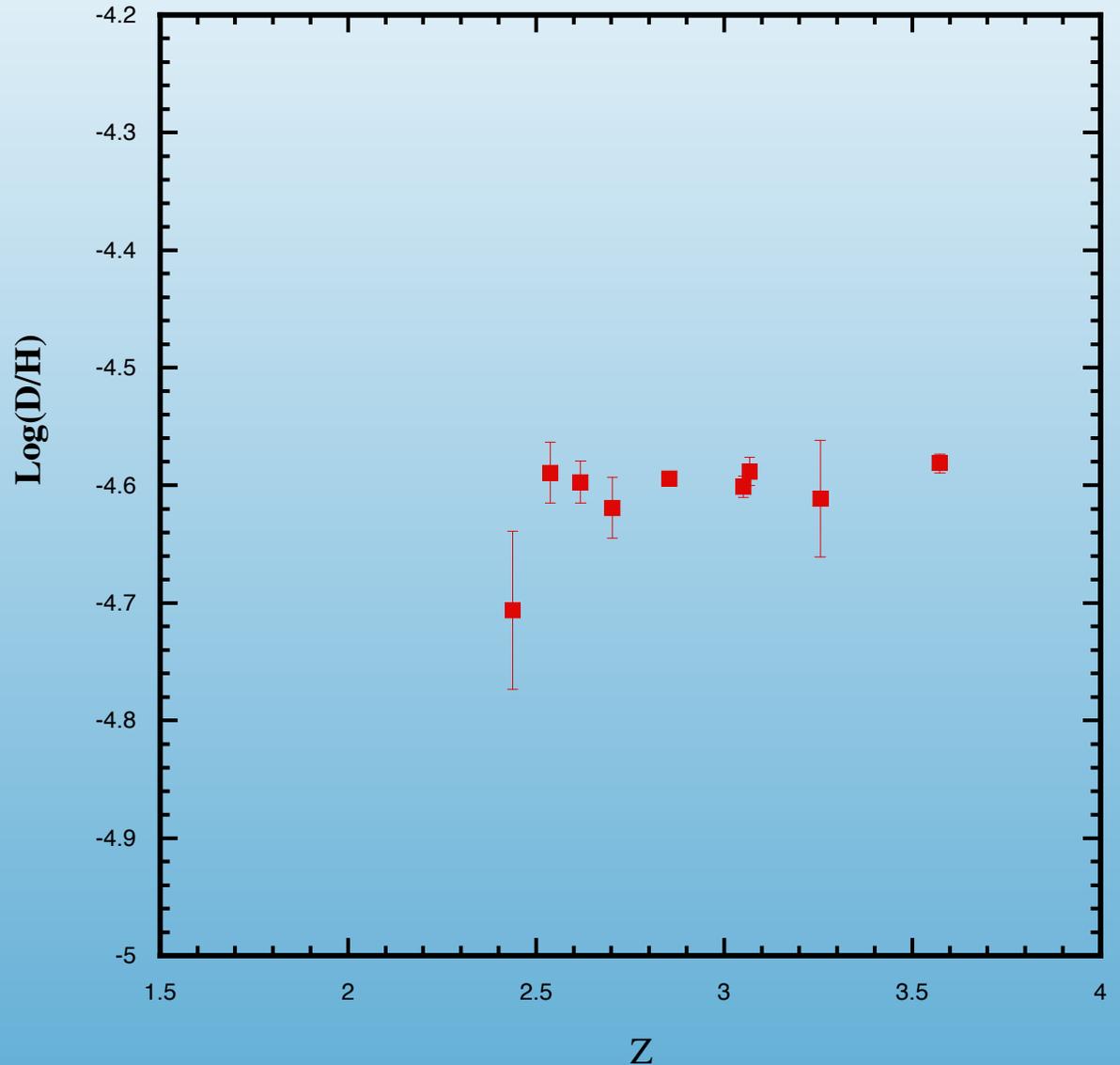
Obs Average:

$$10^5 D/H = 2.53 \pm 0.04$$

# Updated D/H abundances in Quasar absorption systems

BBN Prediction:  
 $10^5 \text{ D/H} = 2.58 \pm 0.13$

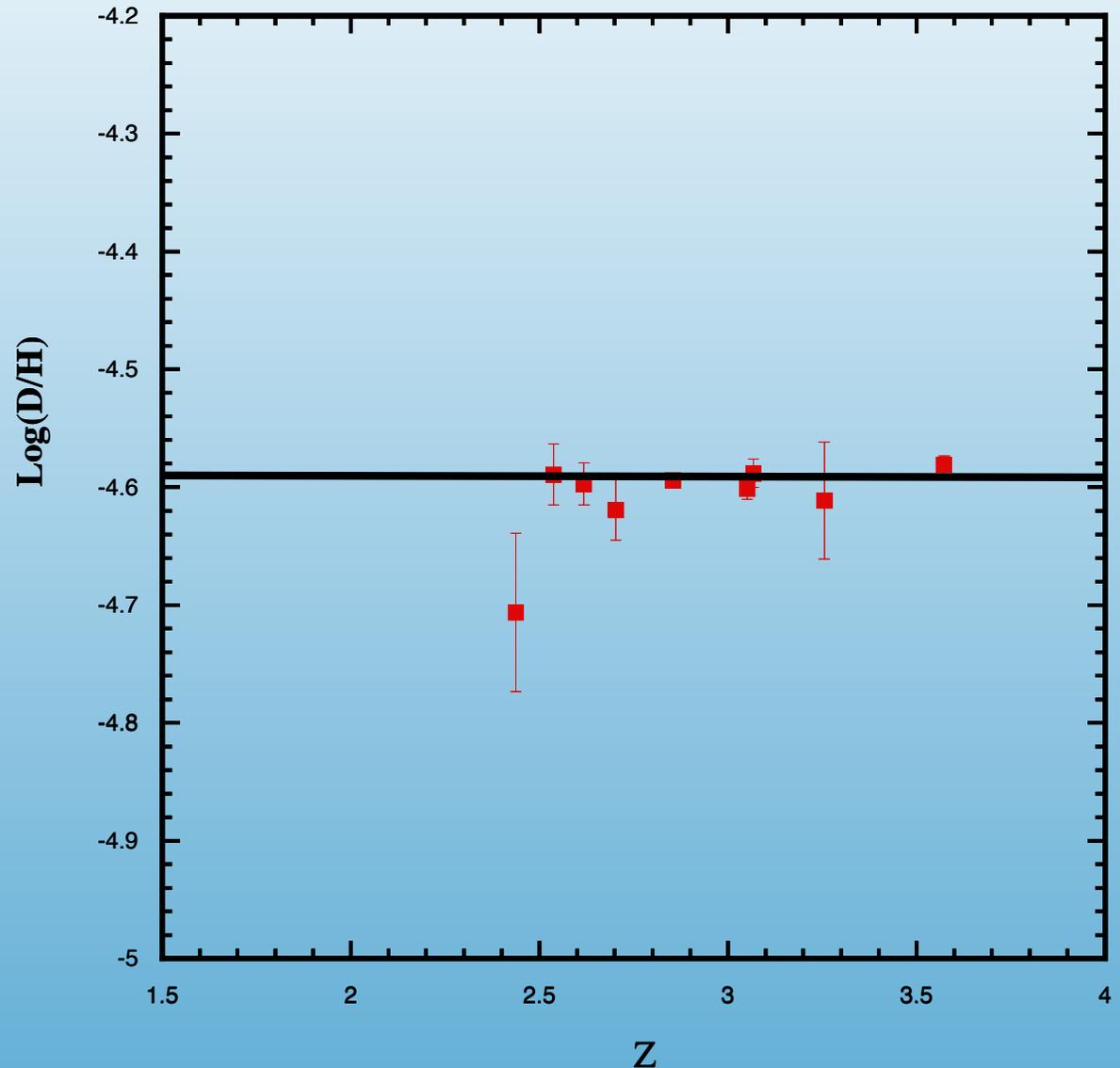
Obs Average:  
 $10^5 \text{ D/H} = 2.55 \pm 0.02$

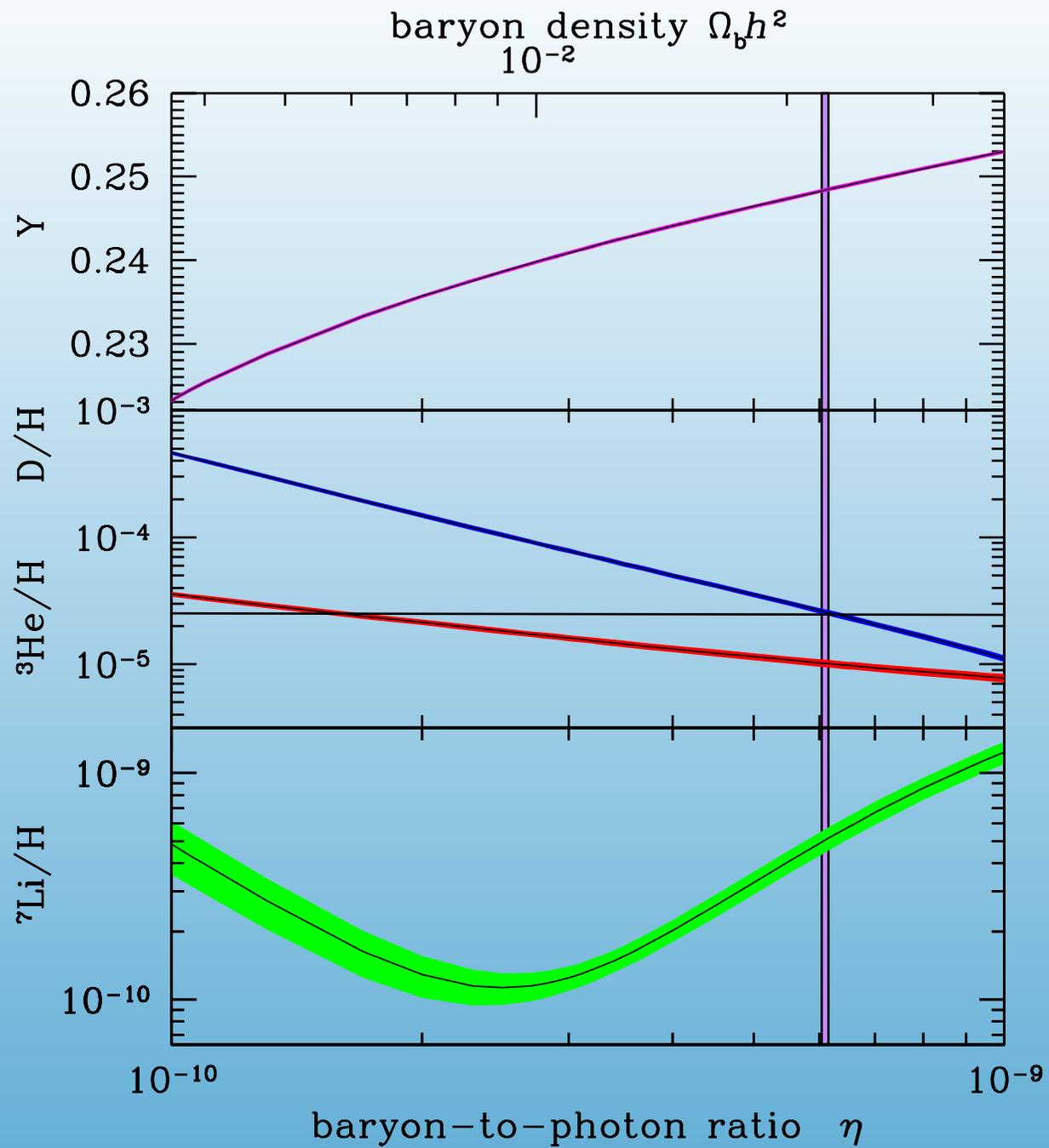


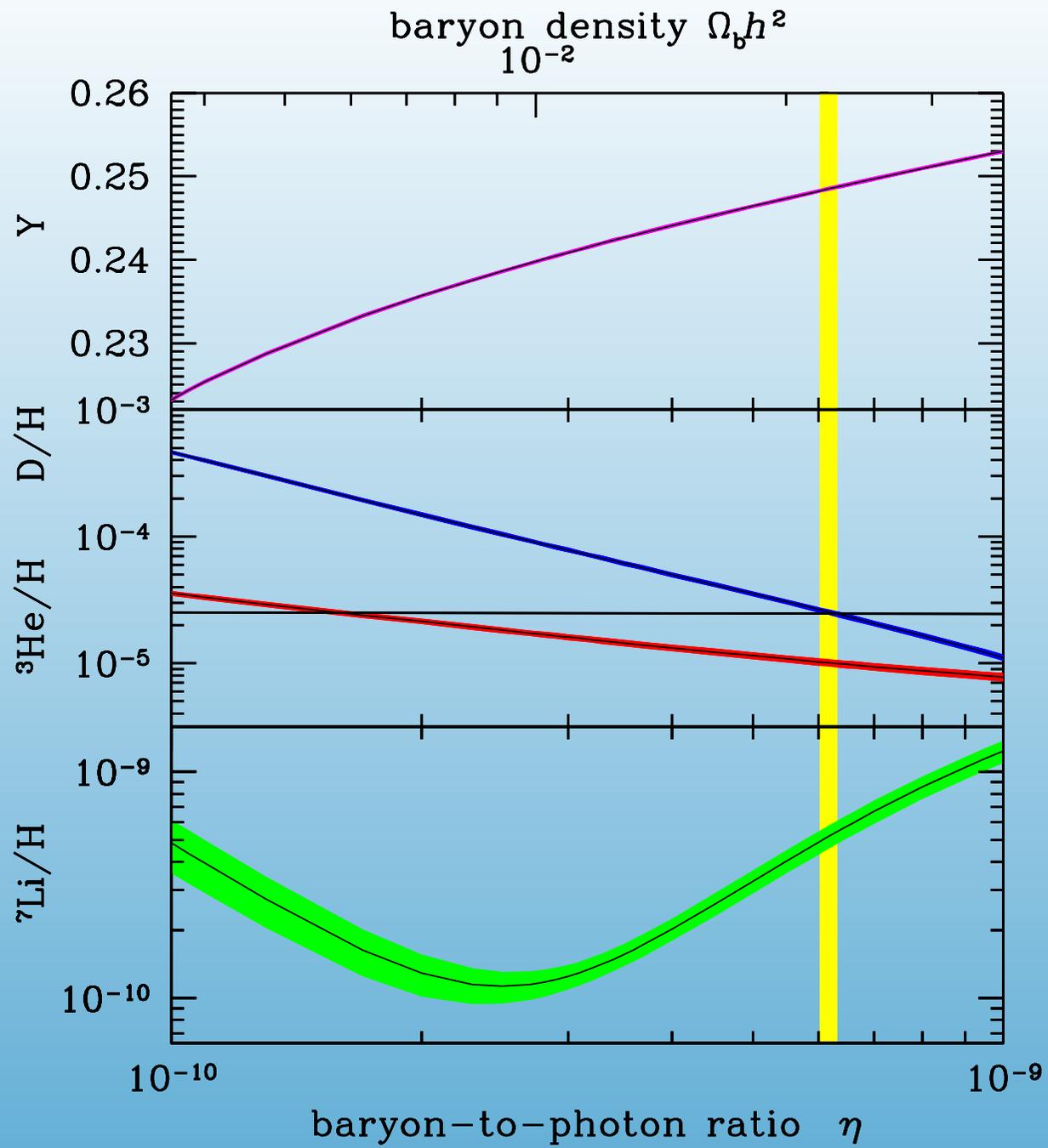
# Updated D/H abundances in Quasar absorption systems

BBN Prediction:  
 $10^5 \text{ D/H} = 2.58 \pm 0.13$

Obs Average:  
 $10^5 \text{ D/H} = 2.55 \pm 0.02$



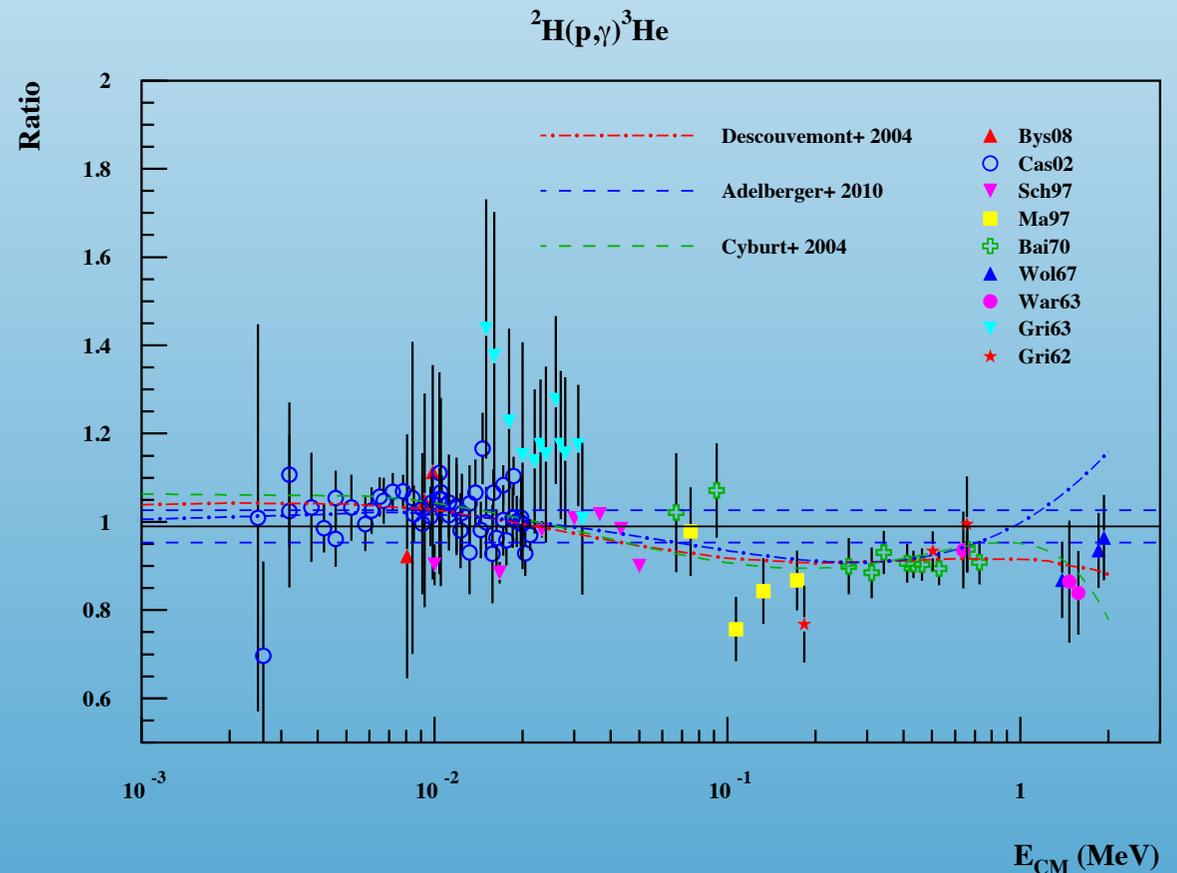
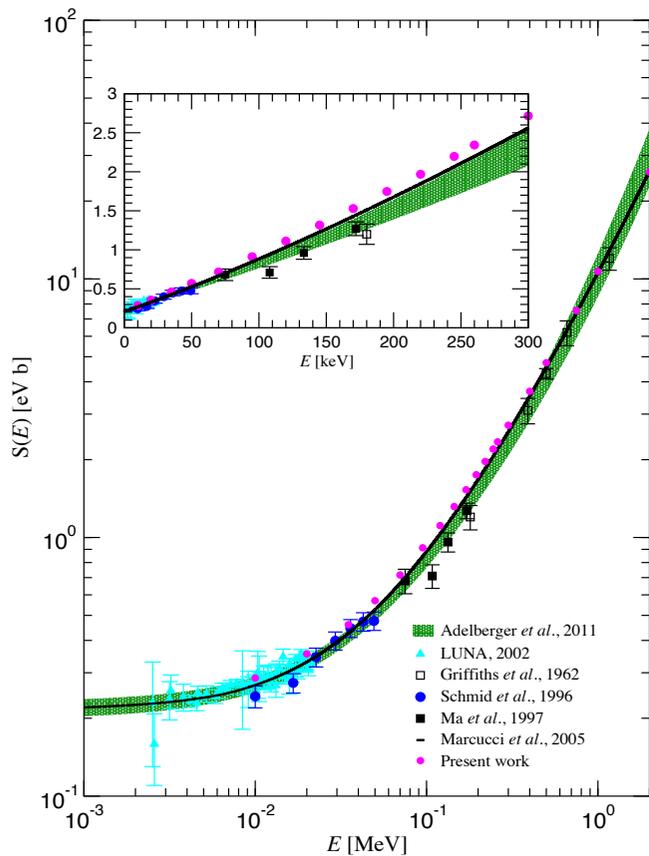




# A word about $d(p,\gamma)^3\text{He}$

Some recent claims (Coc et al.; Cooke et al.) claim a discrepancy with theory and observation in D/H.

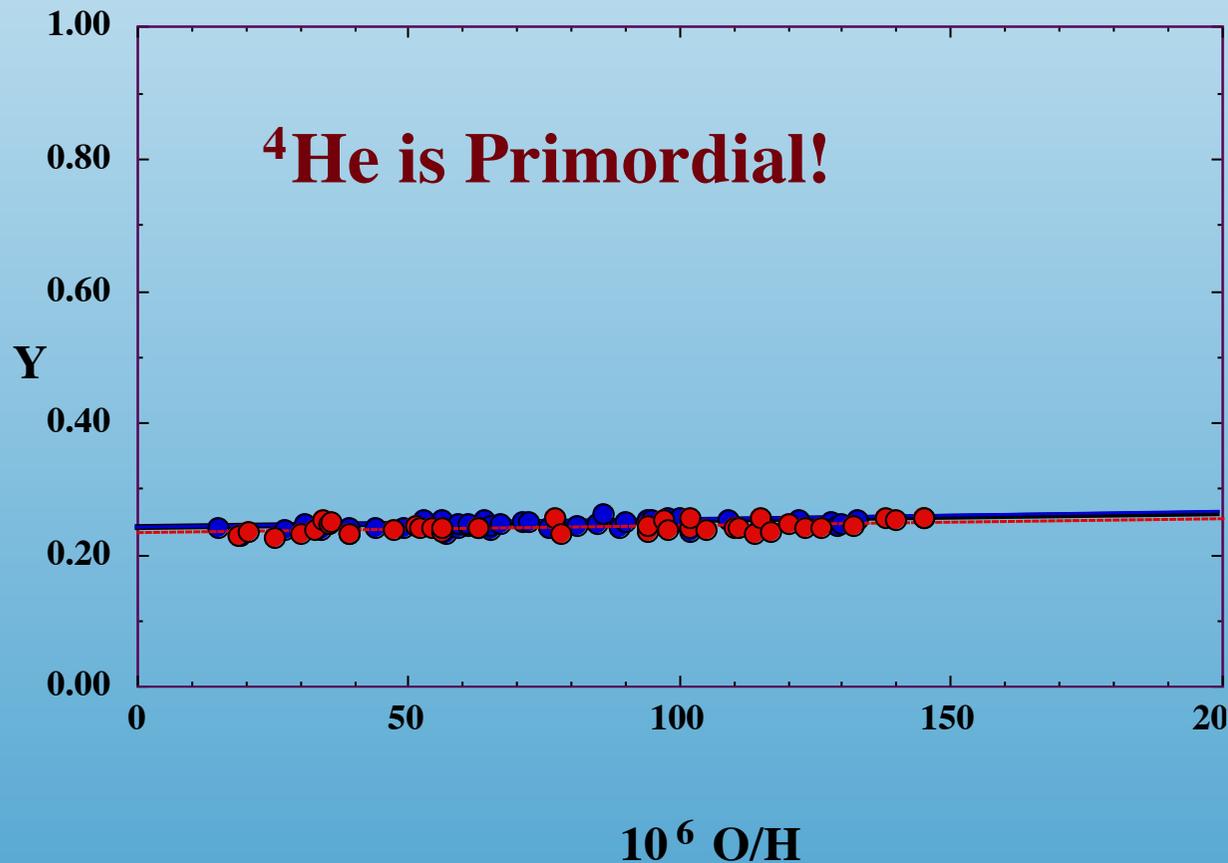
Based on fit to theoretical S-factor (Marucci et al.)

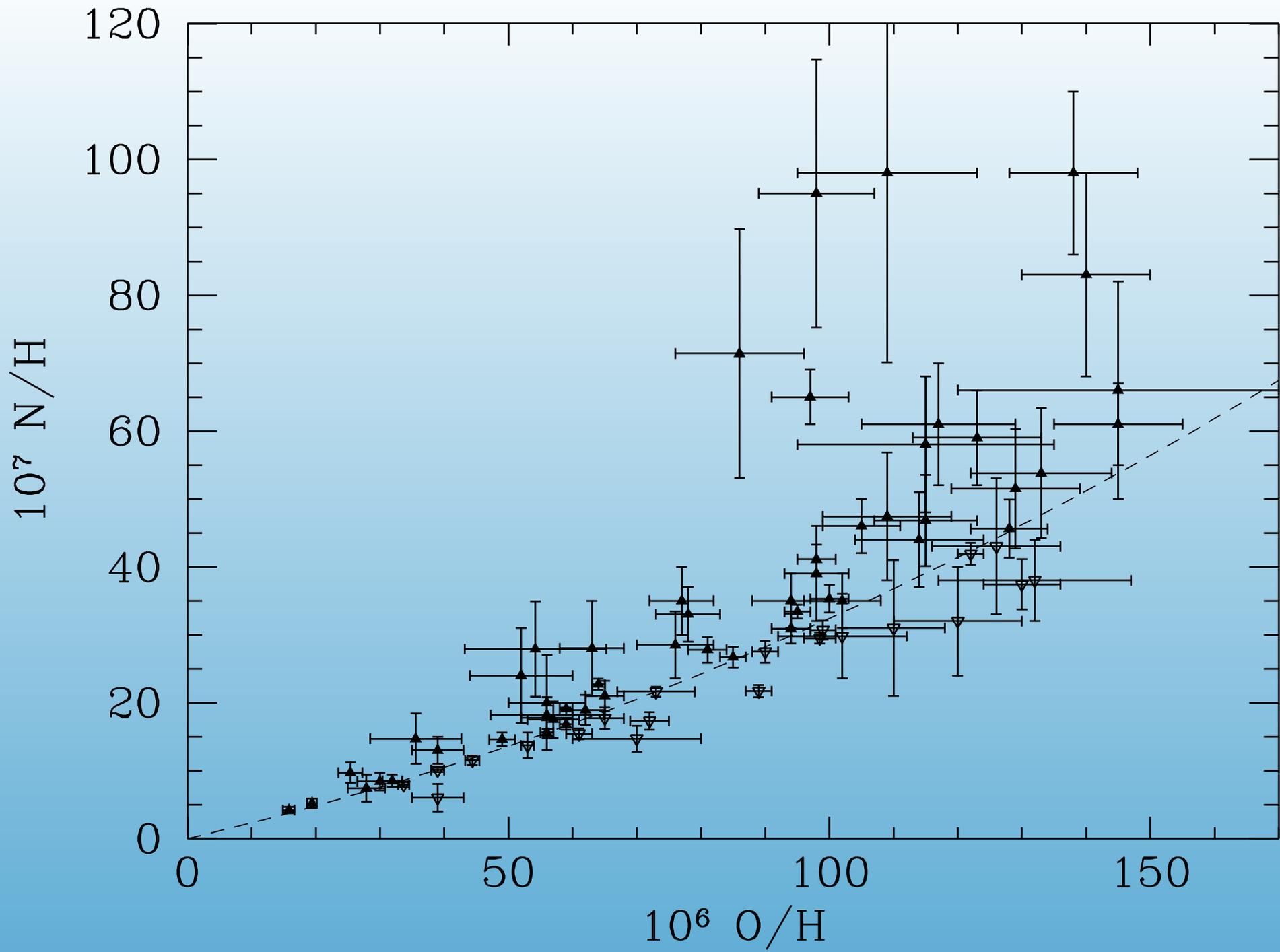


# $^4\text{He}$

Measured in low metallicity extragalactic HII regions ( $\sim 100$ ) together with O/H and N/H

$$Y_P = Y(\text{O/H} \rightarrow 0)$$





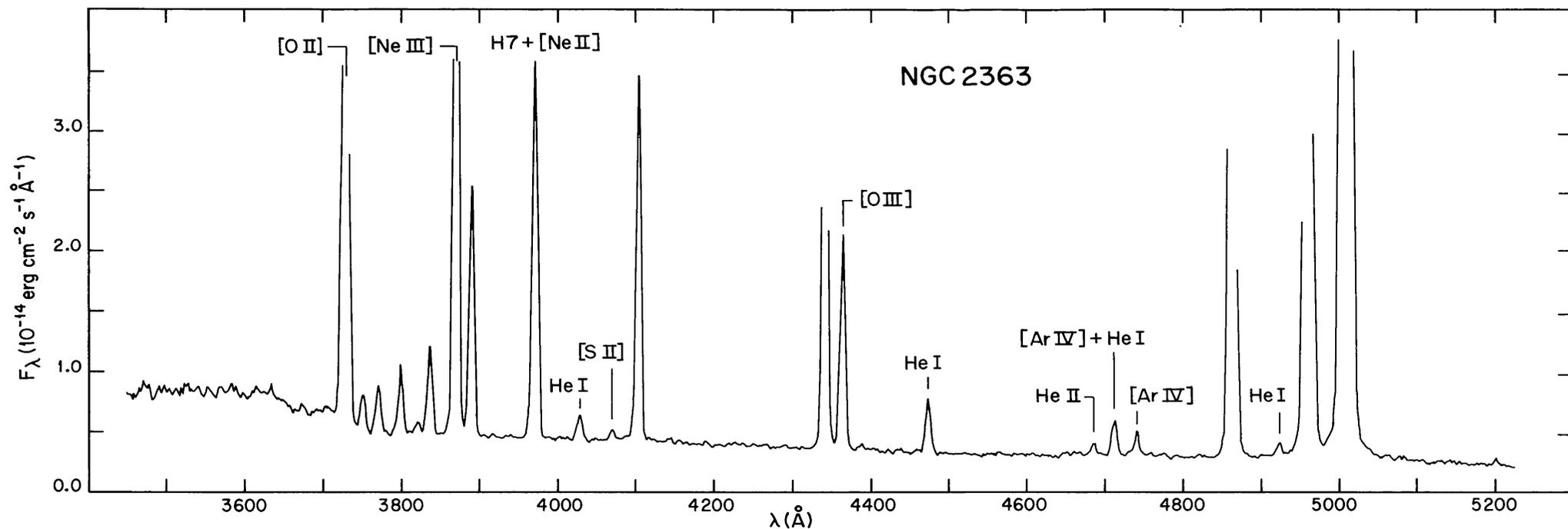


Fig. 1. Low dispersion blue spectrogram of NGC 2363, showing the faintest lines measured

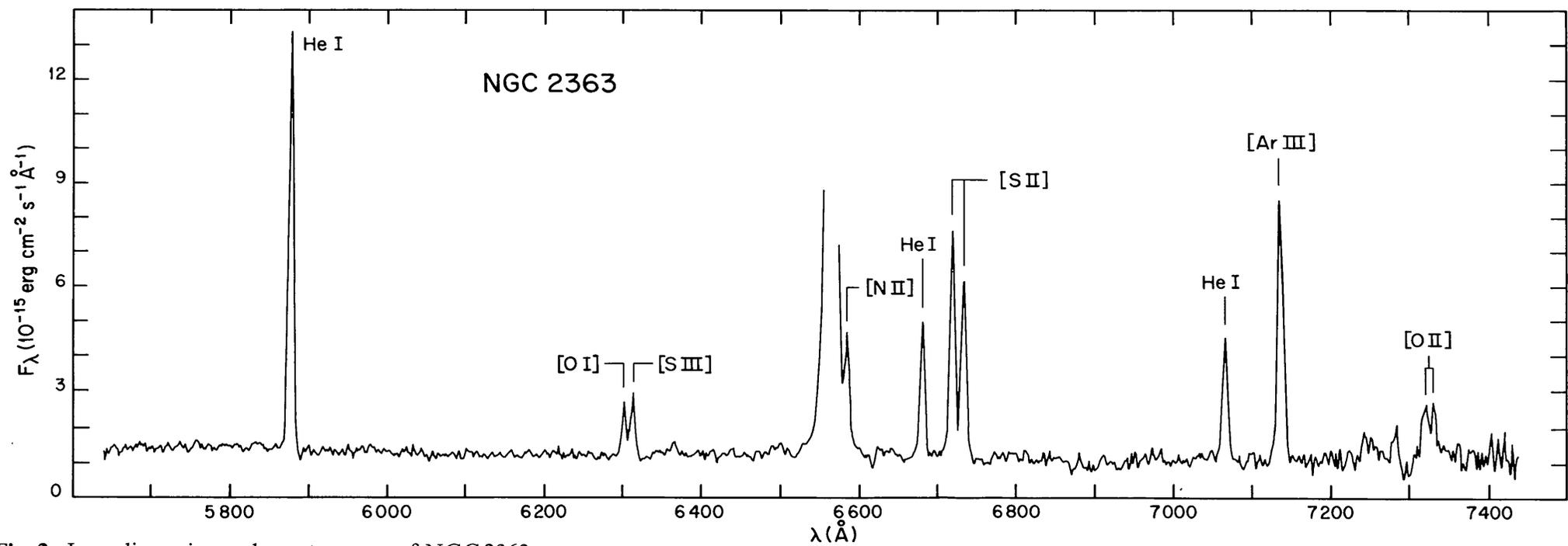
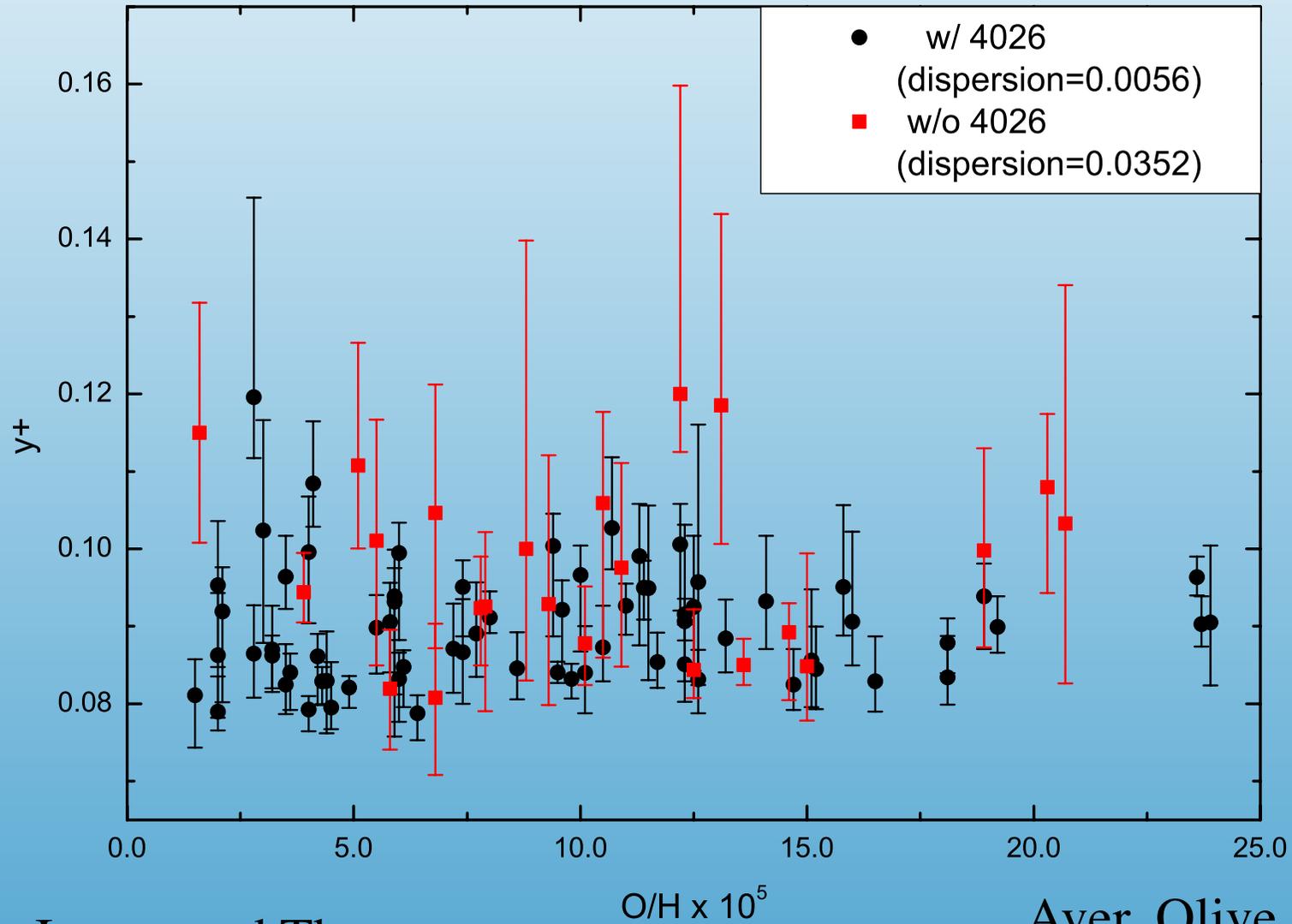


Fig. 2. Low dispersion red spectrogram of NGC 2363

# $^4\text{He}$

## Importance of a full set of line fluxes



Data from Izotov and Thuan

Aver, Olive, Skillman

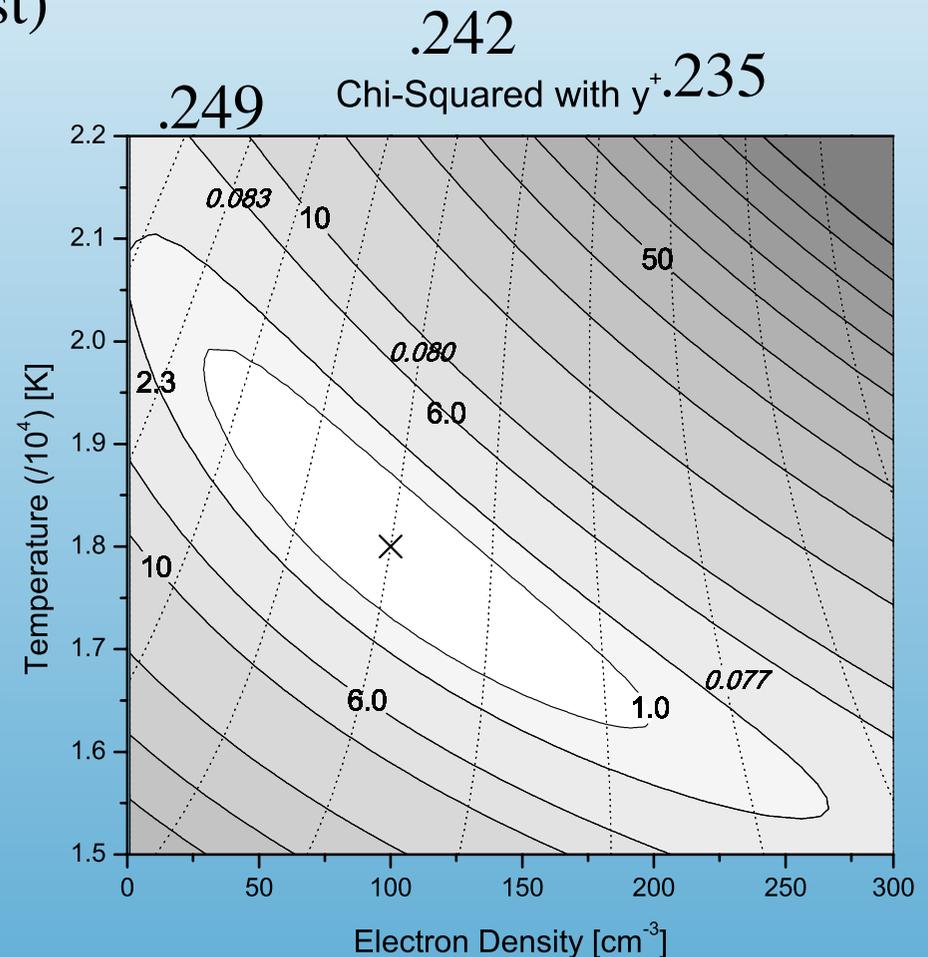
# Results for He dominated by systematic effects

- Interstellar Redding (scattered by dust)
- Underlying Stellar Absorption
- Radiative Transfer
- Collisional Corrections

MCMC statistical techniques have proven effective in parameter estimation

$$\frac{F(\lambda)}{F(H\beta)} = y^+ \frac{E(\lambda)}{E(H\beta)} \frac{\frac{W(H\beta) + a_H(H\beta)}{W(H\beta)}}{\frac{W(\lambda) + a_{He}(\lambda)}{W(\lambda)}} f_\tau(\lambda) \frac{1 + \frac{C}{R}(\lambda)}{1 + \frac{C}{R}(H\beta)} 10^{-f(\lambda)C(H\beta)}$$

$(y^+, n_e, a_{He}, \tau, T, C(H\beta), a_H, \xi)$



Aver, Olive, Skillman

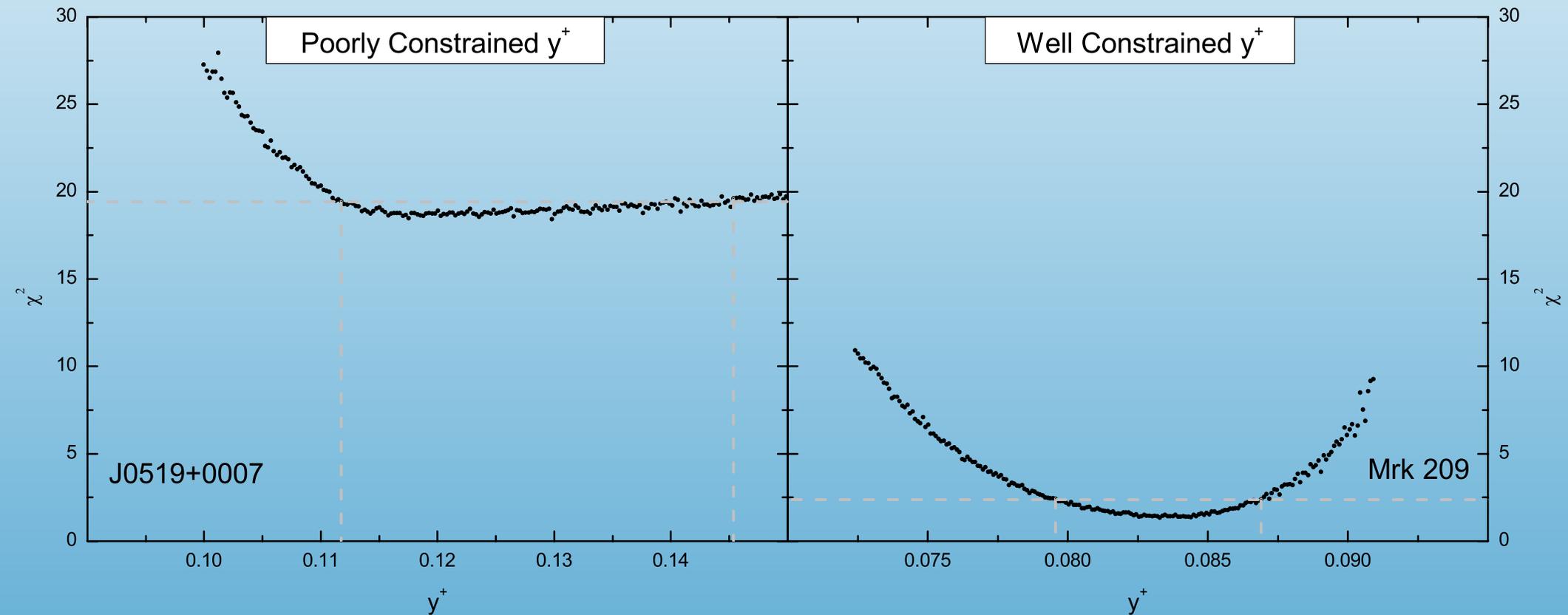
# Results for He dominated by systematic effects

$$\chi^2 = \sum_{\lambda} \frac{\left( \frac{F(\lambda)}{F(H\beta)} - \frac{F(\lambda)}{F(H\beta)}_{\text{meas}} \right)^2}{\sigma(\lambda)^2}$$

9 observables

$(y^+, n_e, a_{He}, \tau, T, C(H\beta), a_H, \xi)$

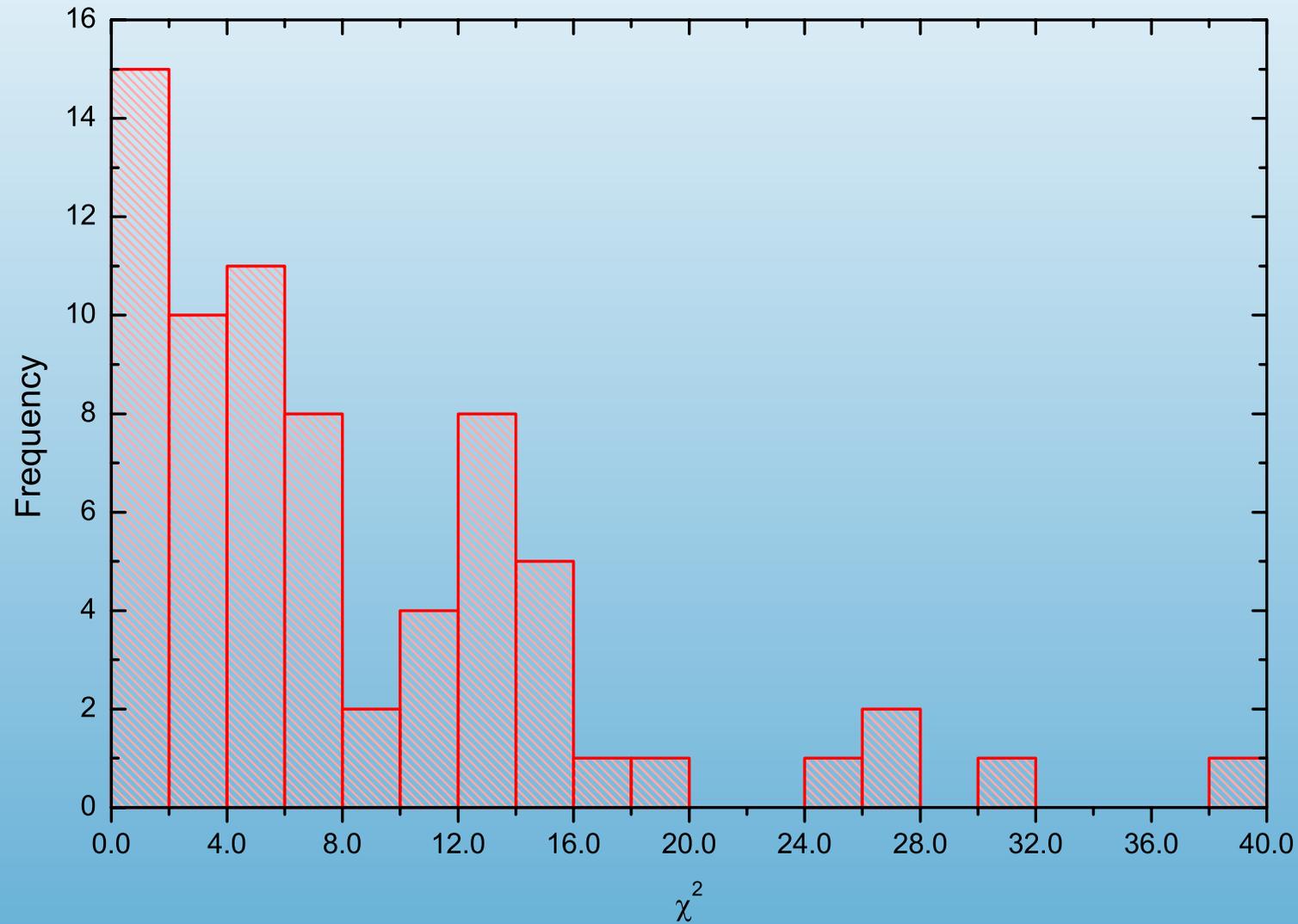
8 parameters



$$\frac{F(\lambda)}{F(H\beta)} = y^+ \frac{E(\lambda)}{E(H\beta)} \frac{\frac{W(H\beta) + a_H(H\beta)}{W(H\beta)}}{\frac{W(\lambda) + a_{He}(\lambda)}{W(\lambda)}} f_{\tau}(\lambda) \frac{1 + \frac{C}{R}(\lambda)}{1 + \frac{C}{R}(H\beta)} 10^{-f(\lambda)C(H\beta)}$$

Aver, Olive, Skillman

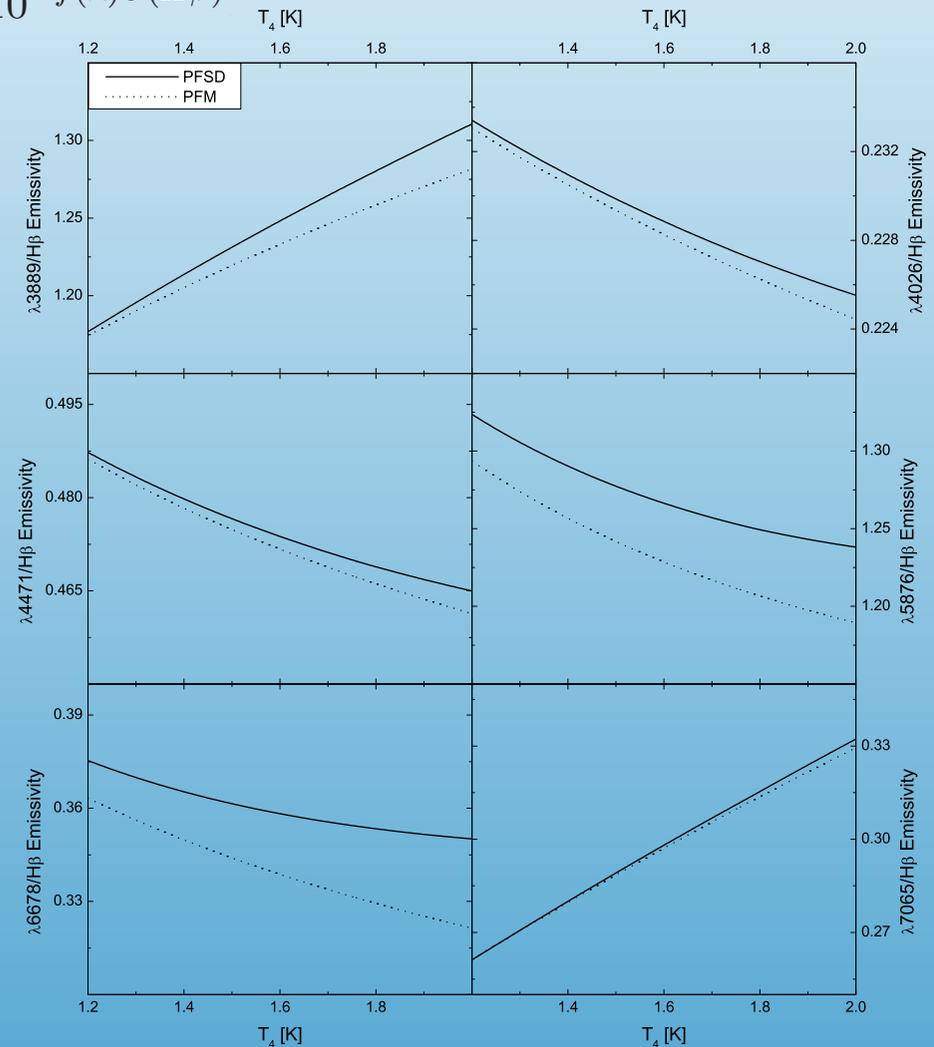
# Using $\chi^2$ as a discriminator



# New Emissivities

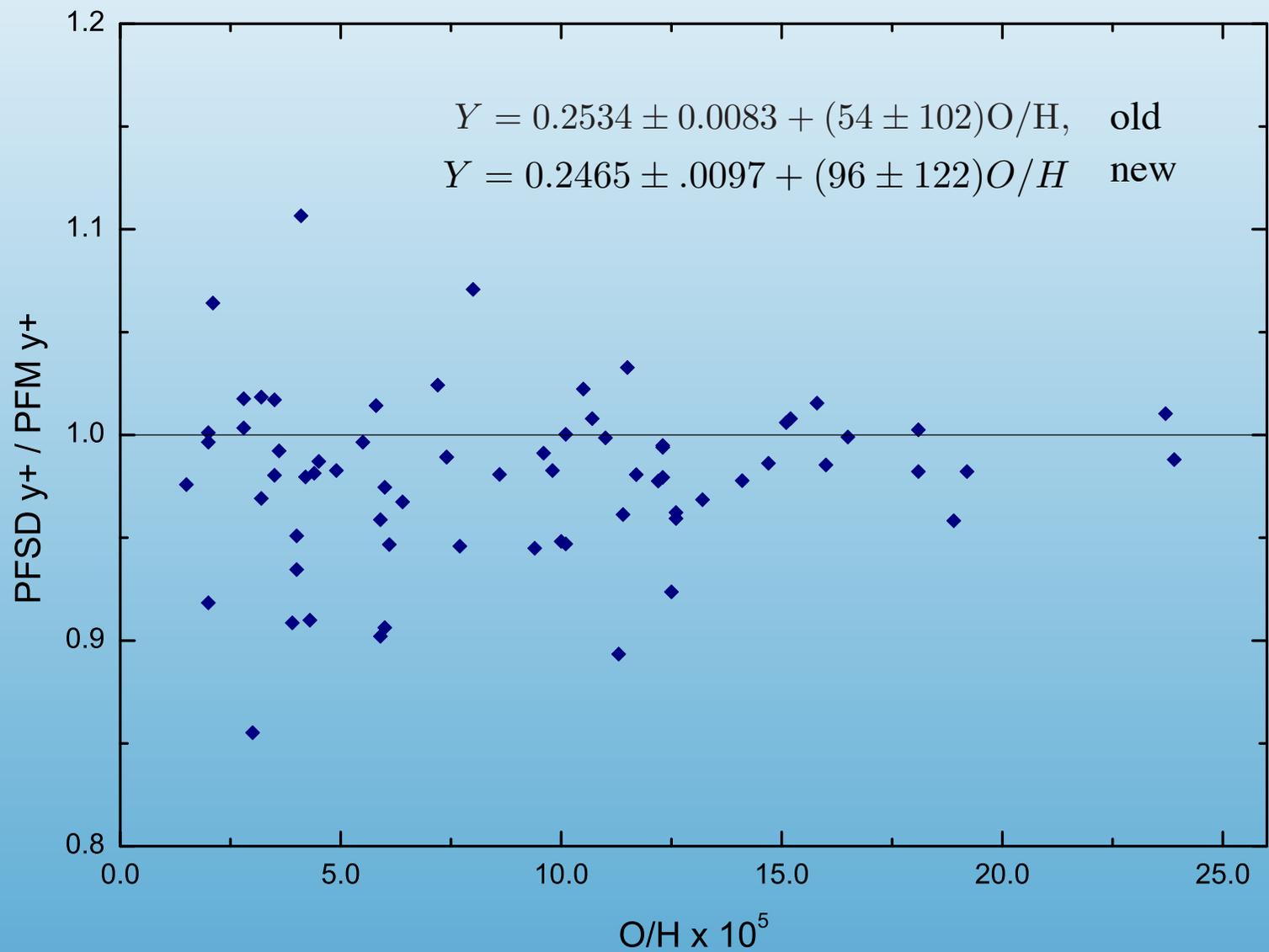
$$\frac{F(\lambda)}{F(H\beta)} = y^+ \frac{E(\lambda)}{E(H\beta)} \frac{\frac{W(H\beta)+a_H(H\beta)}{W(H\beta)}}{\frac{W(\lambda)+a_{He}(\lambda)}{W(\lambda)}} f_{\tau}(\lambda) \frac{1 + \frac{C}{R}(\lambda)}{1 + \frac{C}{R}(H\beta)} 10^{-f(\lambda)C(H\beta)}$$

New emissivities from  
Porter et al. (2012,2013) using  
improved cross sections  
(and corrected errors!)



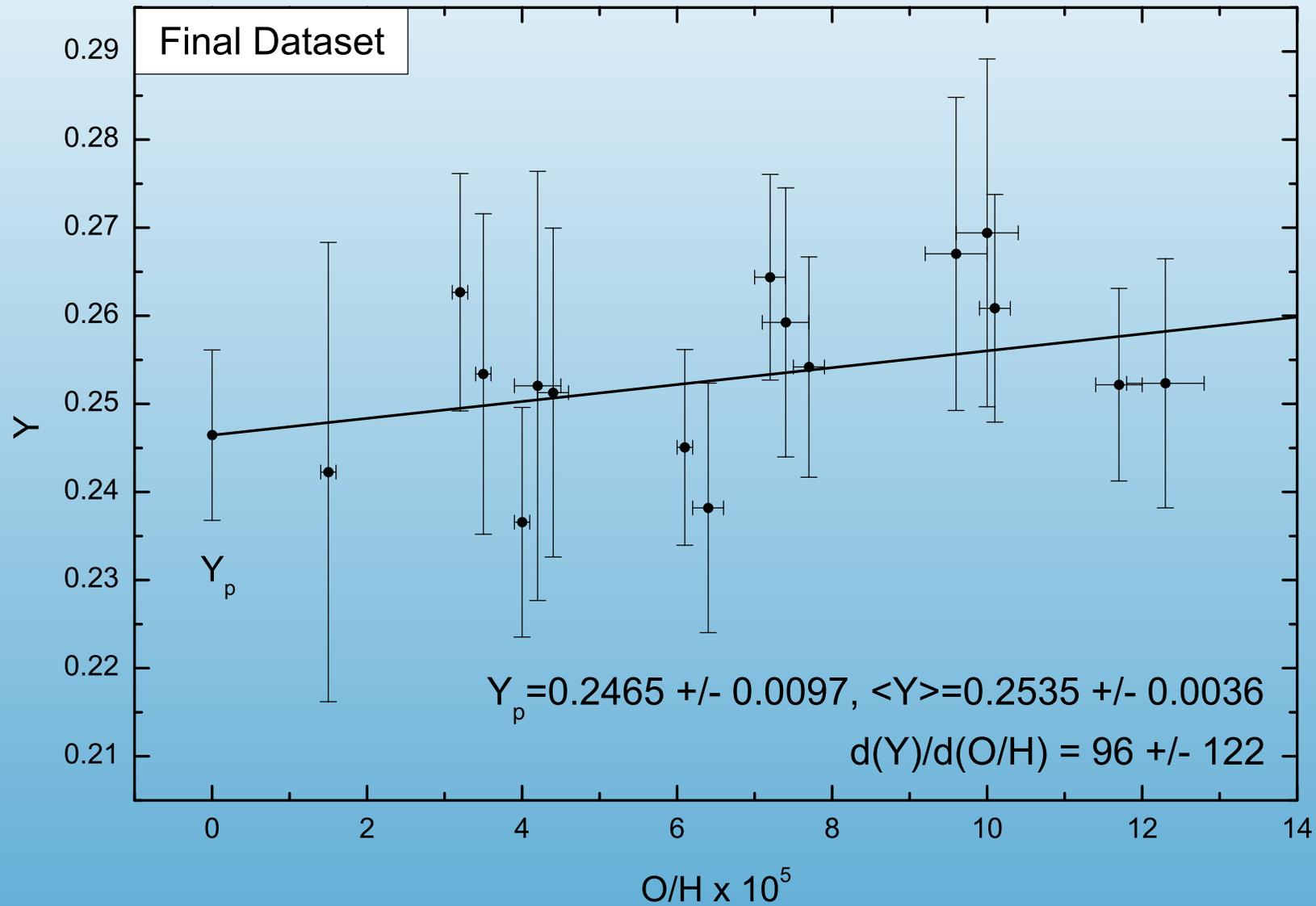
Aver, Olive, Porter, Skillman

Higher emissivities  $\Rightarrow$  lower  
He abundances



Aver, Olive, Porter, Skillman

# $^4\text{He}$



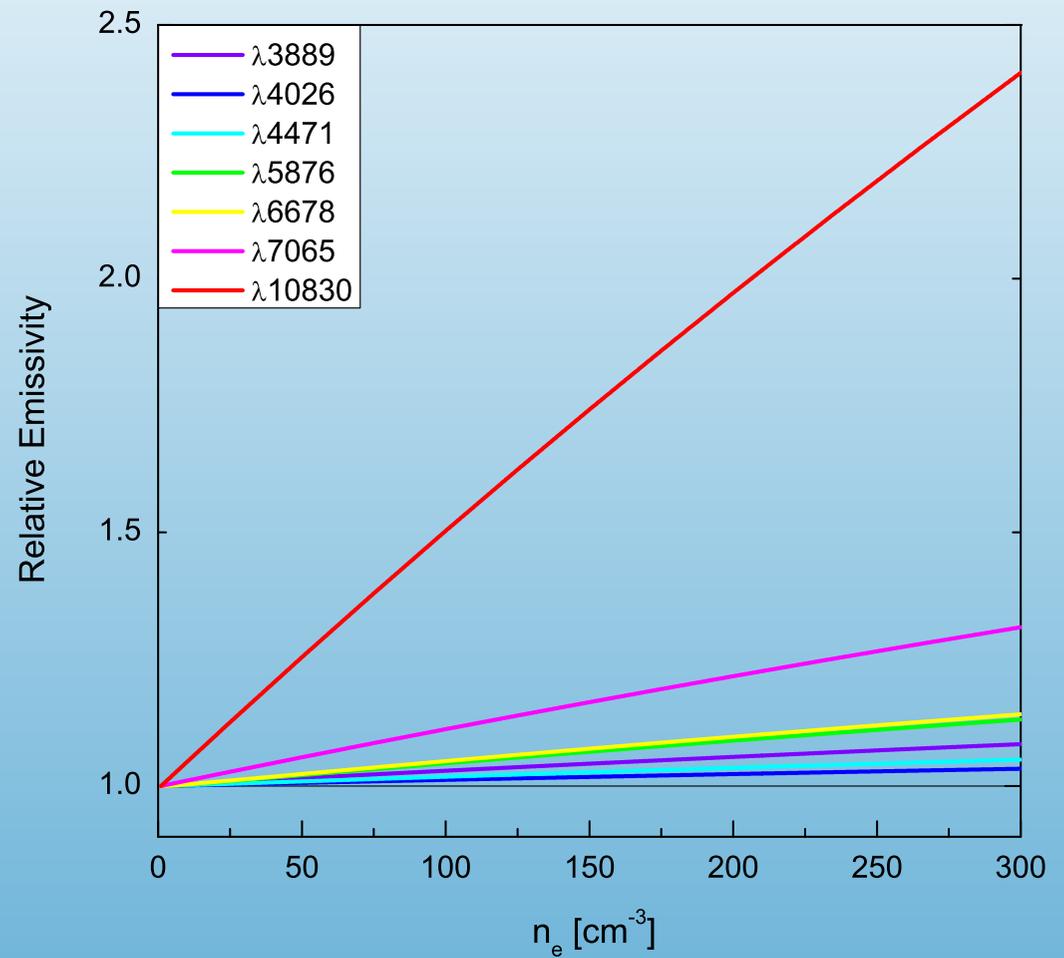
old emissivities:  $Y_p = .2534 \pm 0.0083$

Aver, Olive, Porter, Skillman

# Adding $\lambda 10380$

Izotov, Thuan, Guseva

Increased sensitivity to  $n_e$

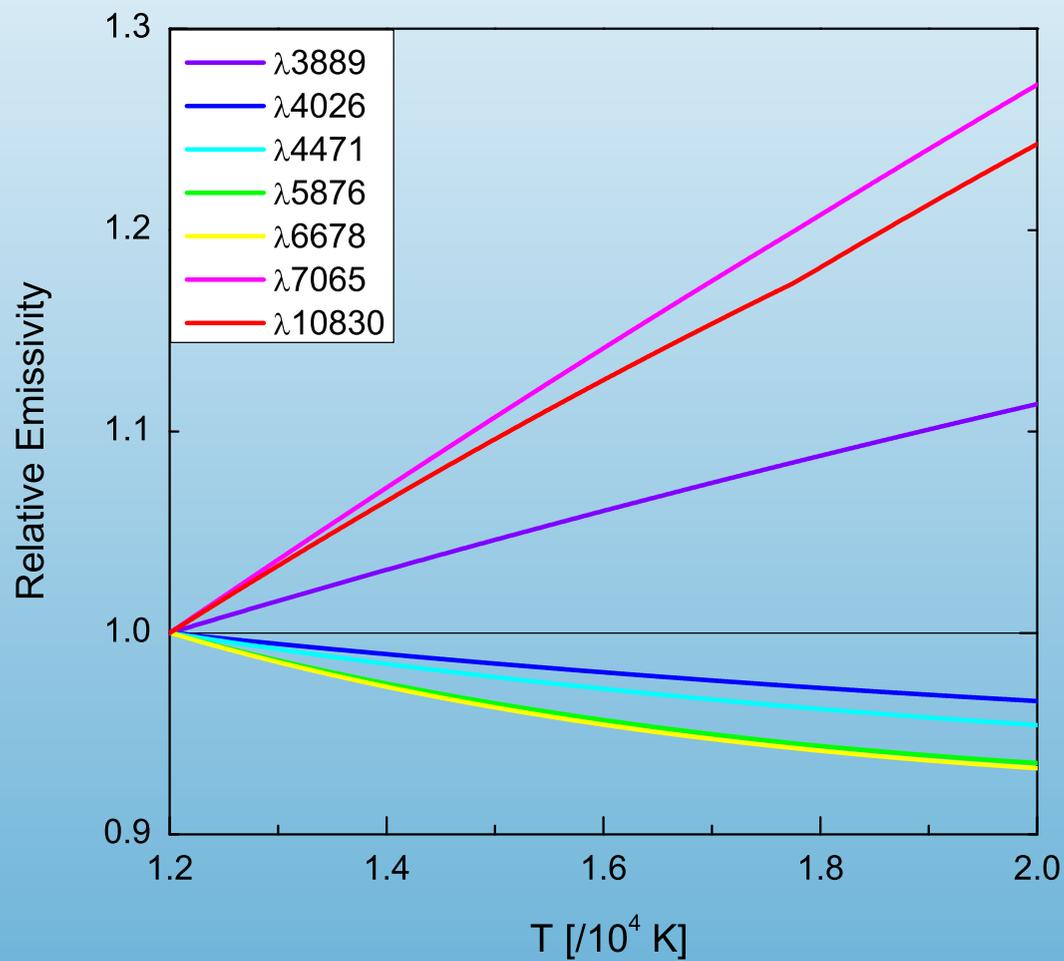


Aver, Olive, Skillman

# Adding $\lambda 10380$

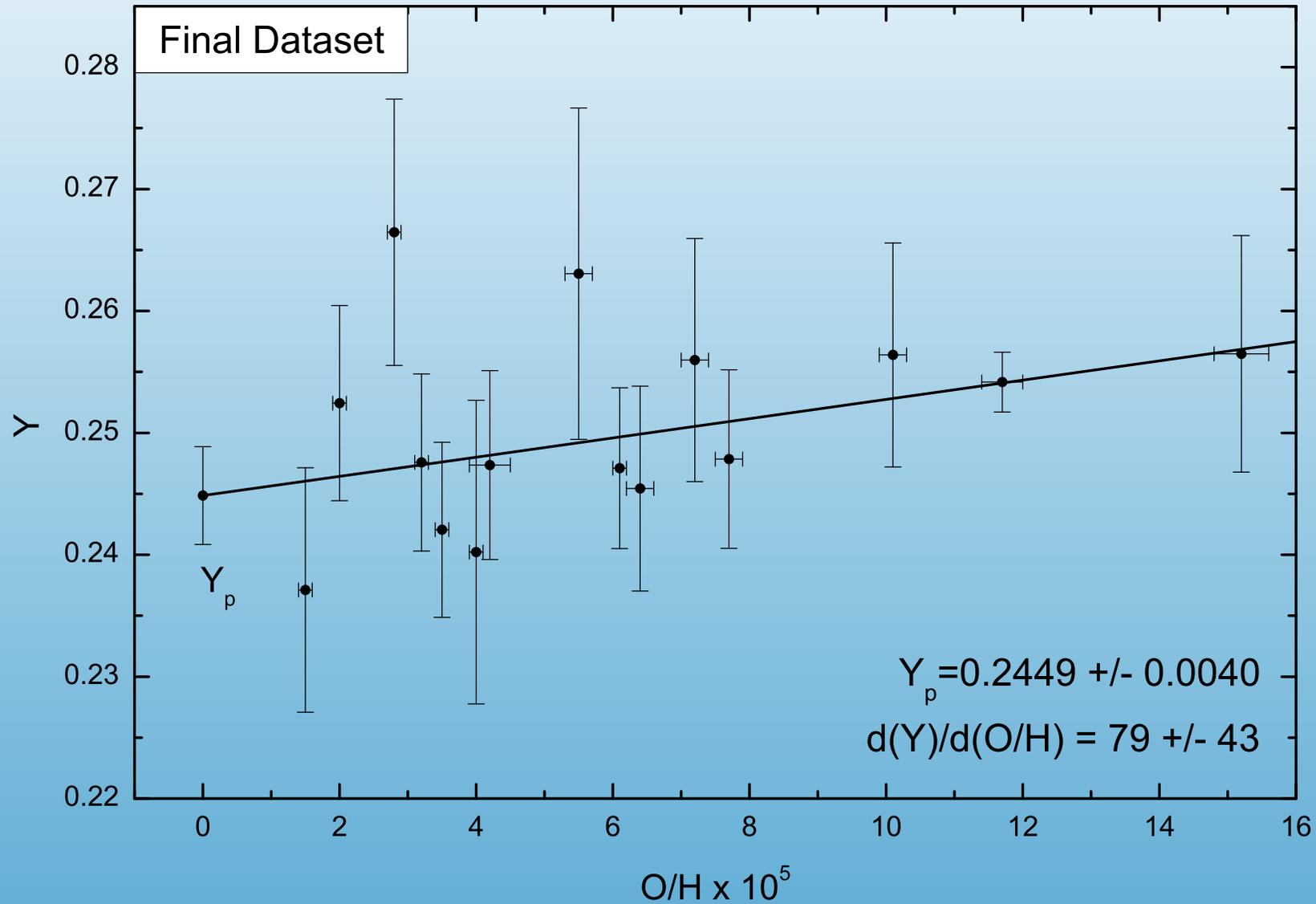
Izotov, Thuan, Guseva

Strong sensitivity to T



Aver, Olive, Skillman

# $^4\text{He}$



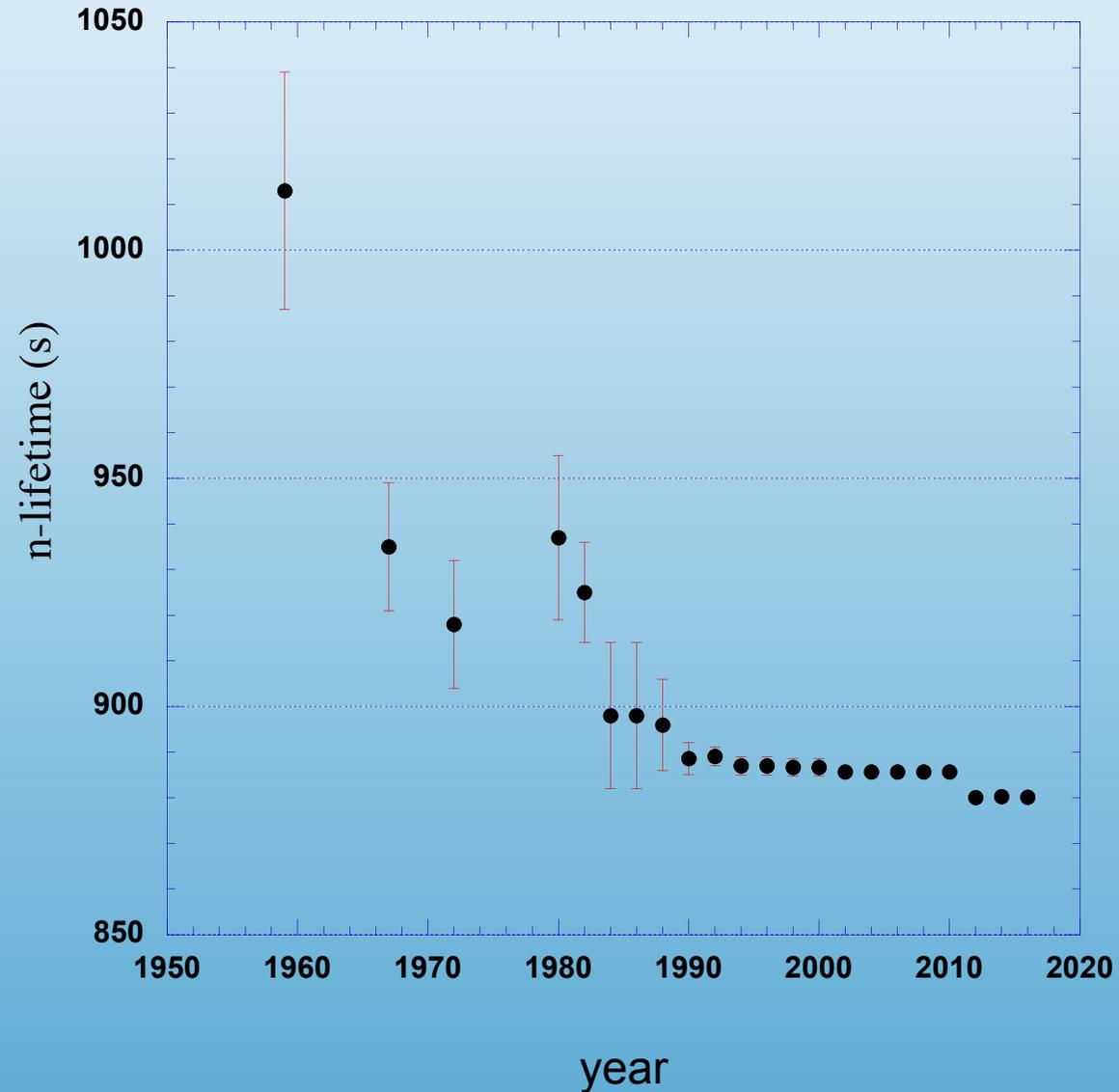
without 10830:  $Y_p = .2465 \pm 0.0097$

Aver, Olive, Skillman

# Neutron Lifetime

$$\tau = 885.7 \rightarrow Y = .2481$$

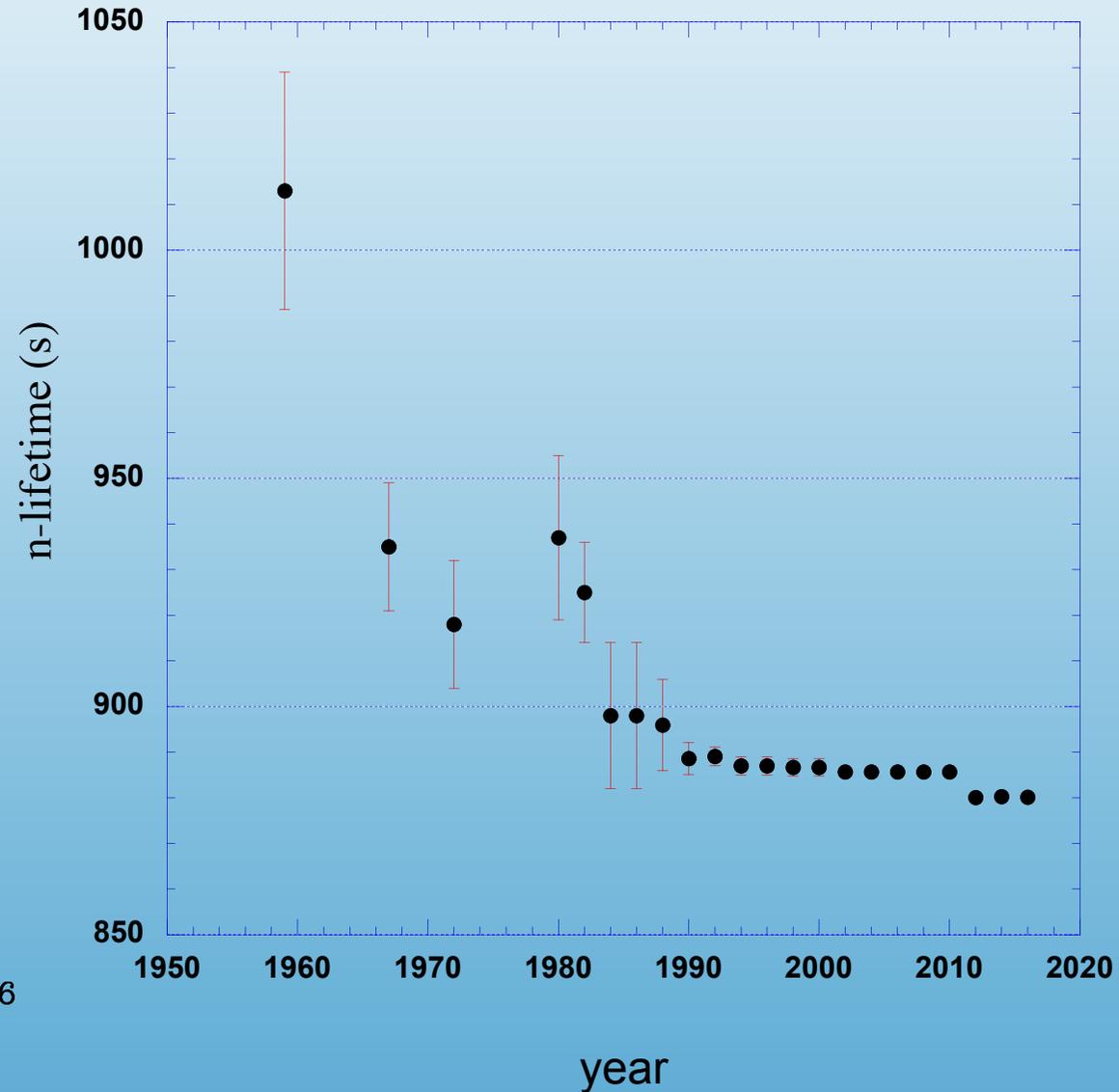
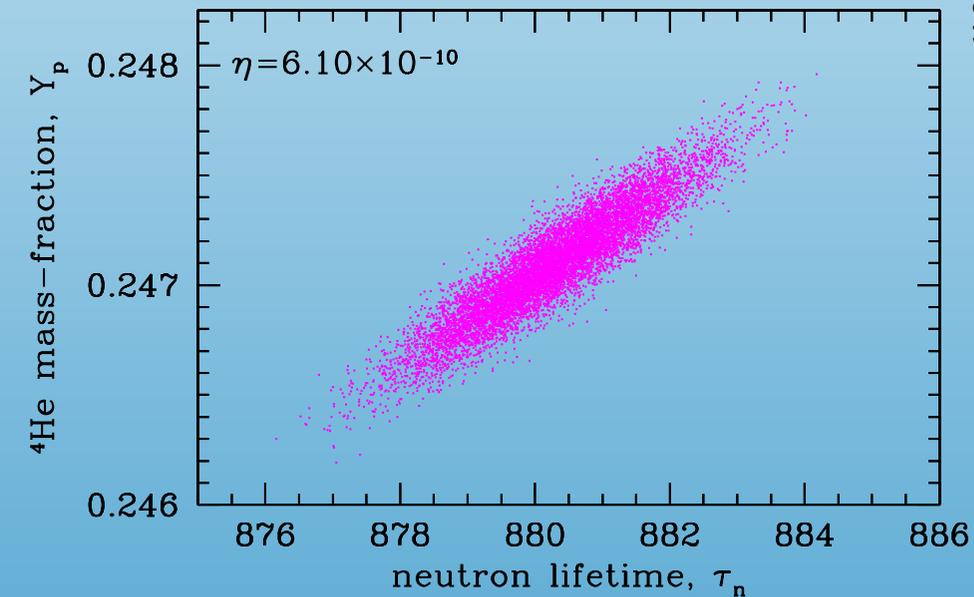
$$\tau = 880.2 \rightarrow Y = .2470$$



# Neutron Lifetime

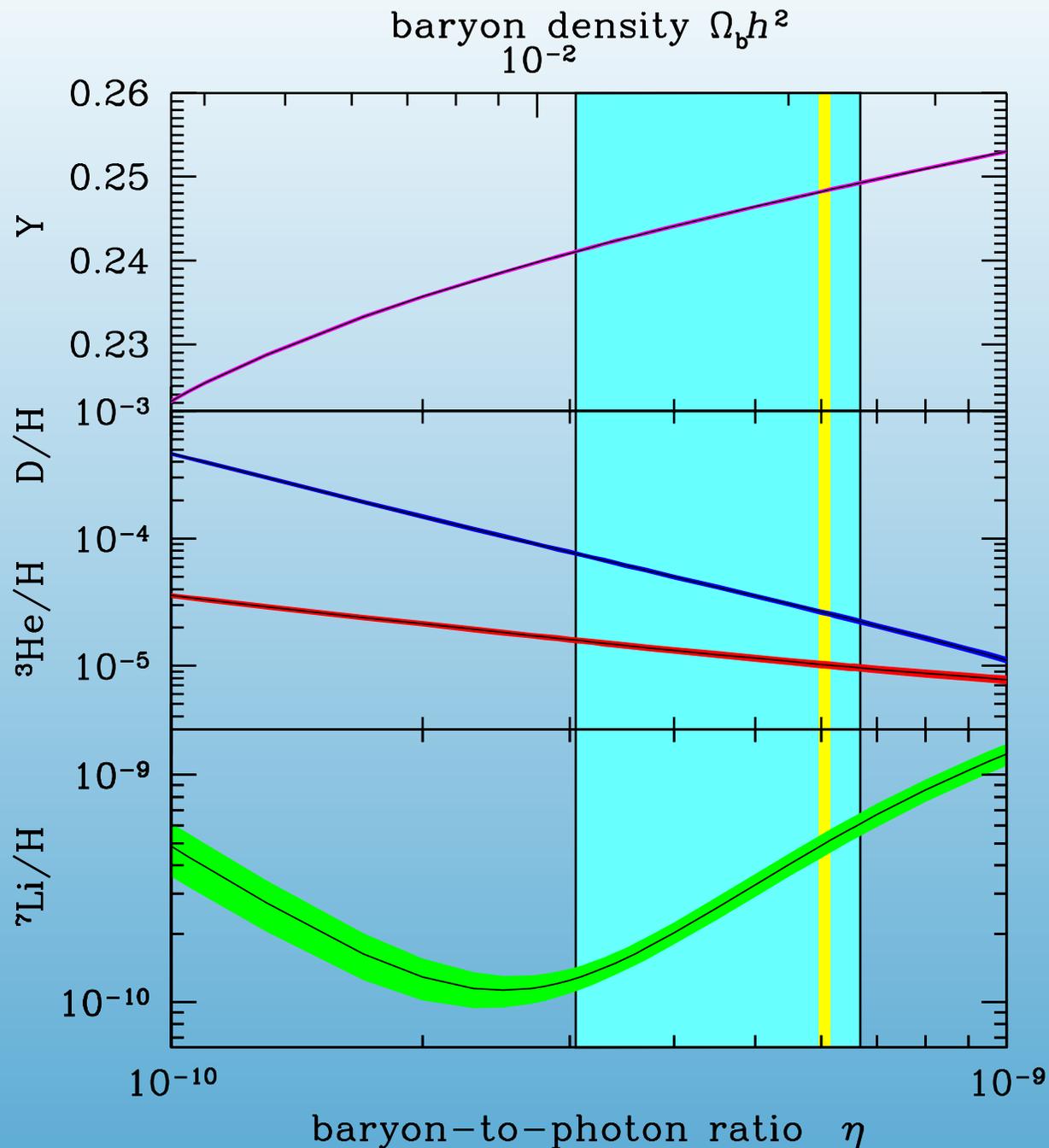
$$\tau = 885.7 \rightarrow Y = .2481$$

$$\tau = 880.2 \rightarrow Y = .2470$$



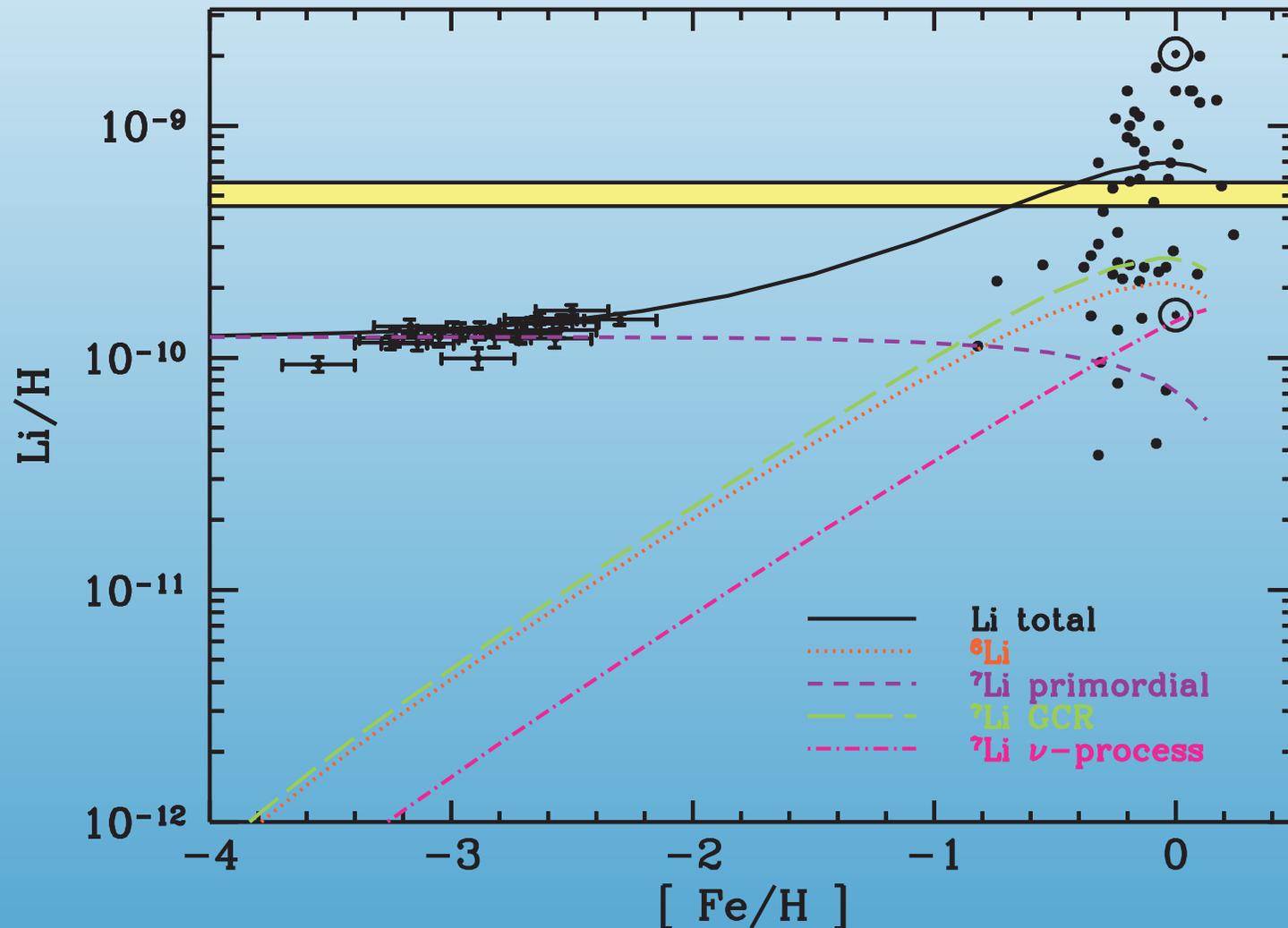
$^4\text{He}$  Prediction:  
 $0.2471 \pm 0.0002$

Data: Regression:  
 $0.2449 \pm 0.0040$

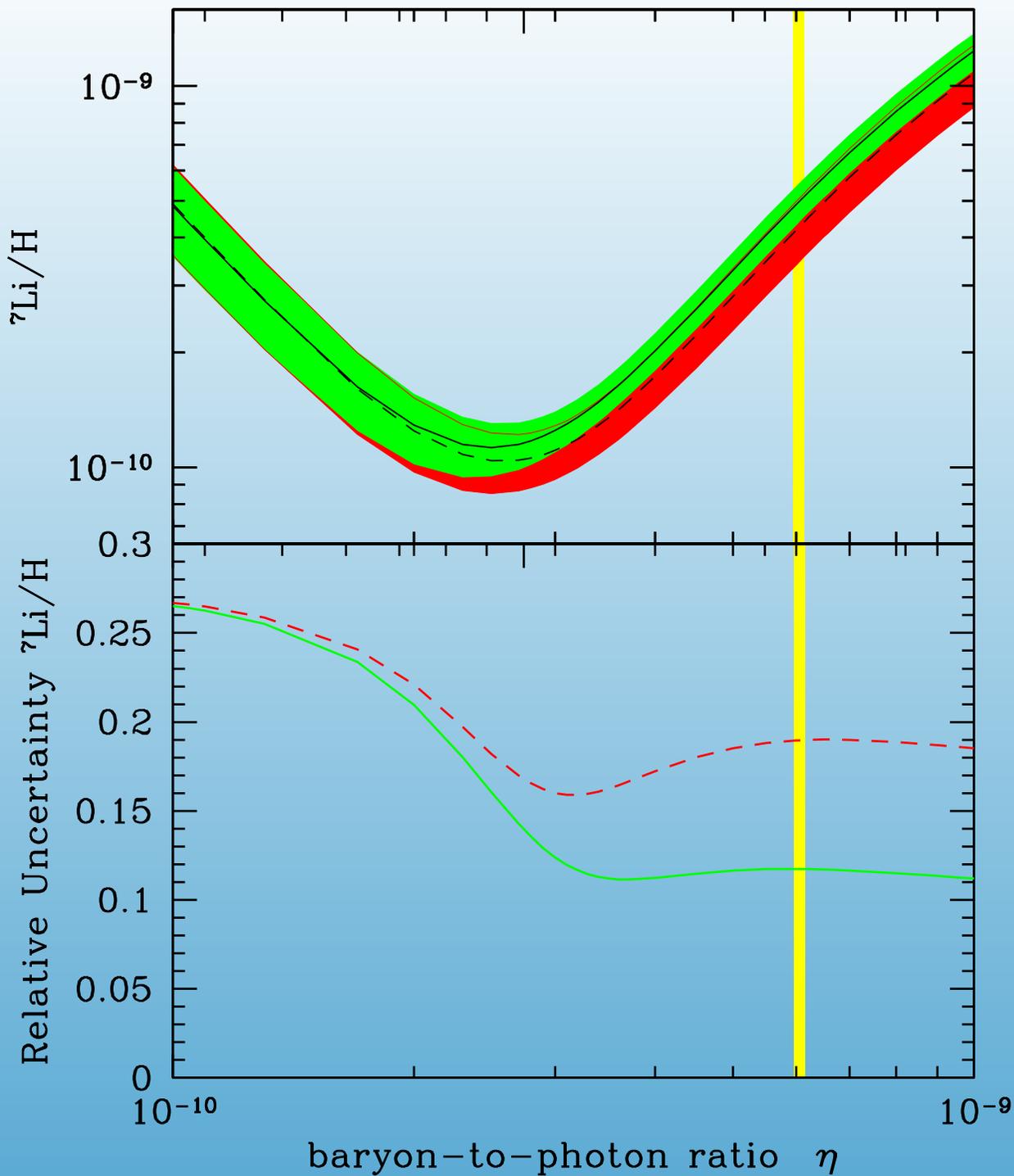


# Li/H

Measured in low metallicity dwarf halo stars  
(over 100 observed)



baryon density  $\Omega_b h^2$   
0.01



At the Planck  
value for  $\eta$ :

$\text{Li}/\text{H} =$

$$(4.88^{+0.71}_{-0.62}) \times 10^{-10}$$

cf. data at

$$\left(\frac{\text{Li}}{\text{H}}\right)_{\text{halo}^*} = (1.23^{+0.34}_{-0.16}) \times 10^{-10},$$

$$\left(\frac{\text{Li}}{\text{H}}\right)_{\text{Gl.Cl.}} = (2.34 \pm 0.05) \times 10^{-10},$$

improved  ${}^3\text{He}(\alpha, \gamma){}^7\text{Be}$

Cyburt, Fields, KAO

# Possible sources for the discrepancy

- Nuclear Rates
  - Restricted by solar neutrino flux

Coc et al.  
Cyburt, Fields, KAO  
Boyd, et al.

# Resonant Reactions

Is there a missing excited state providing a resonant reaction?



If energy released in producing  $C^*$  is

$$Q_C = \Delta({}^7\text{Be}) + \Delta(A) - \Delta(C^{\text{g.s.}})$$

$$\Delta = m - Am_u$$

mass defect

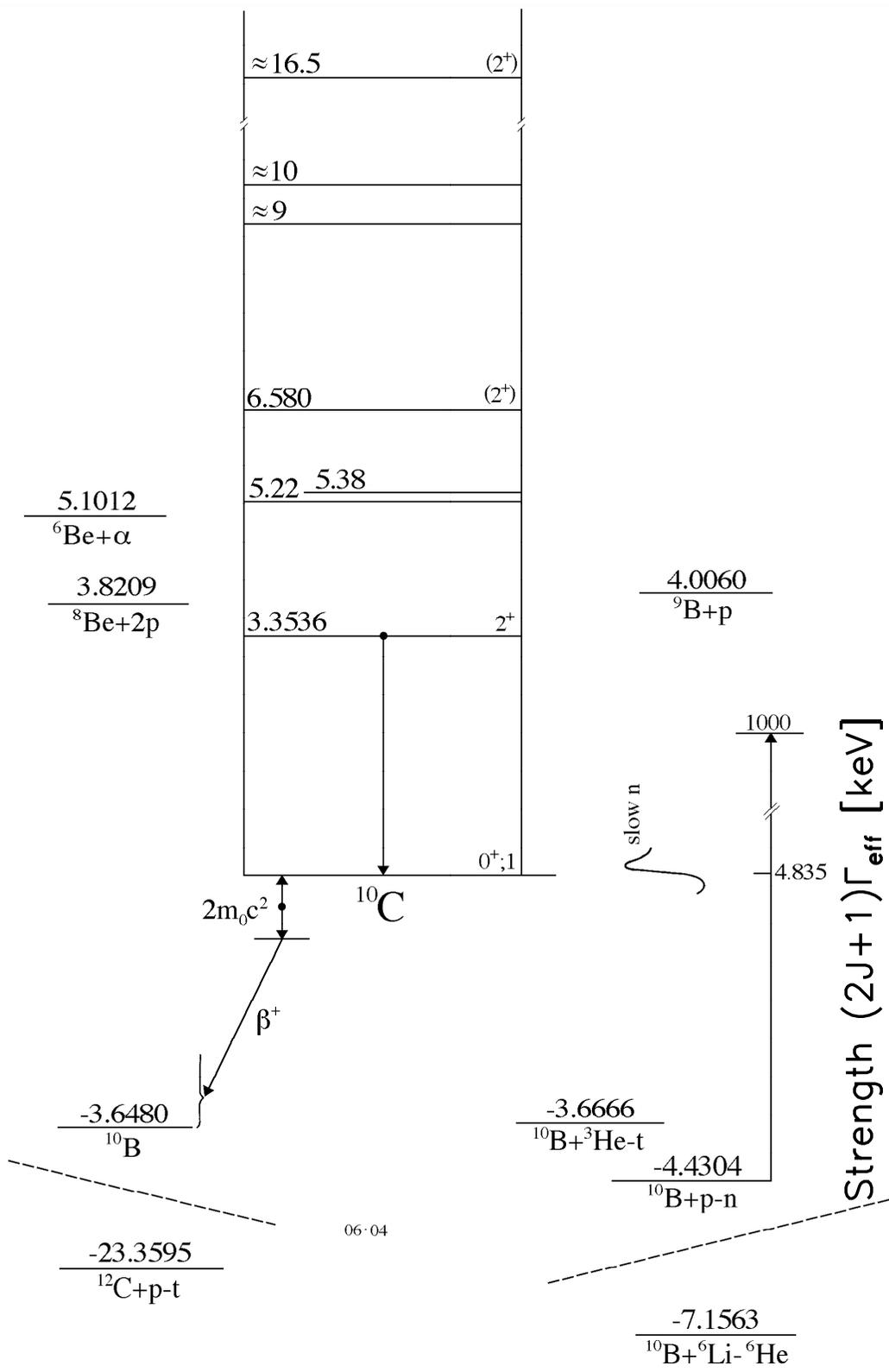
Then the resonant energy is

$$E_{\text{res}} \equiv E_{\text{ex}} - Q_C$$

In principle, long list of possible resonance candidates

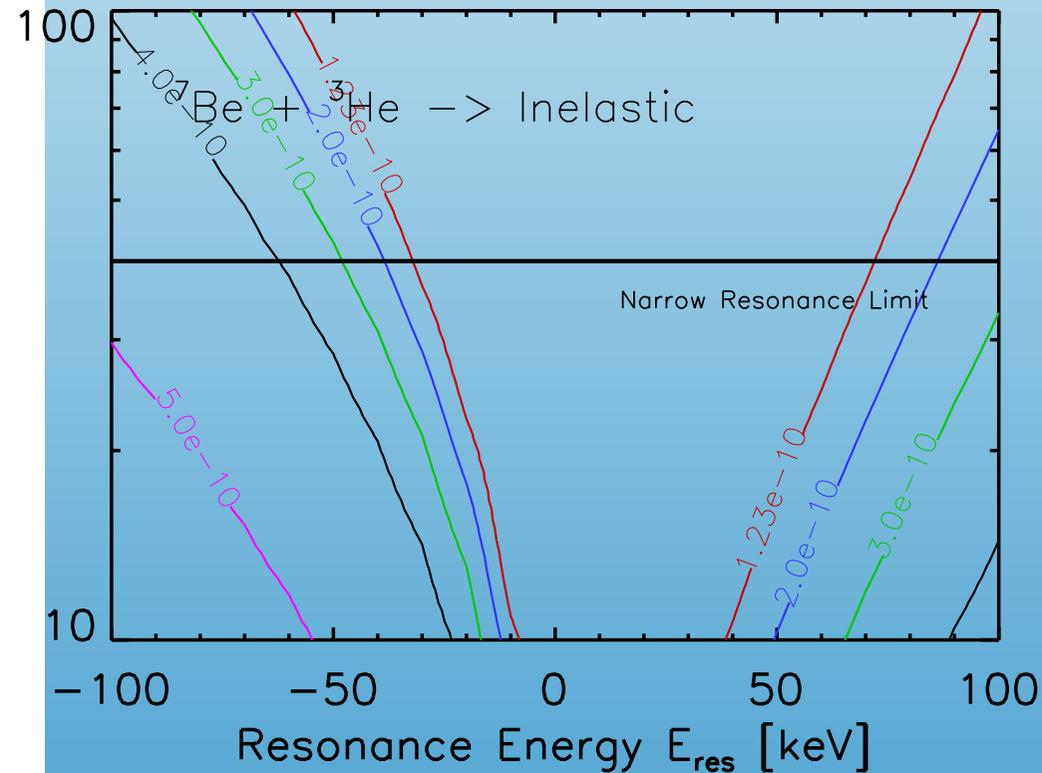
- Excited states of  ${}^8\text{Li}$  (included)
- ${}^8\text{Be}$  (some included) - large  $E_{\text{res}}$
- ${}^8\text{B}$  (included)
- ${}^9\text{B}$  - interesting state at 16.71 MeV
- ${}^{10}\text{B}$  - interesting state at 18.80 MeV
- ${}^{10}\text{C}$  - potentially interesting state at 15 MeV
- ${}^{11}\text{C}$  - negligible effect





$$\frac{15.0}{^7\text{Be} + ^3\text{He}}$$

eg. if a 1- or 2- excited state of  $^{10}\text{C}$  were near 15.0 MeV .....



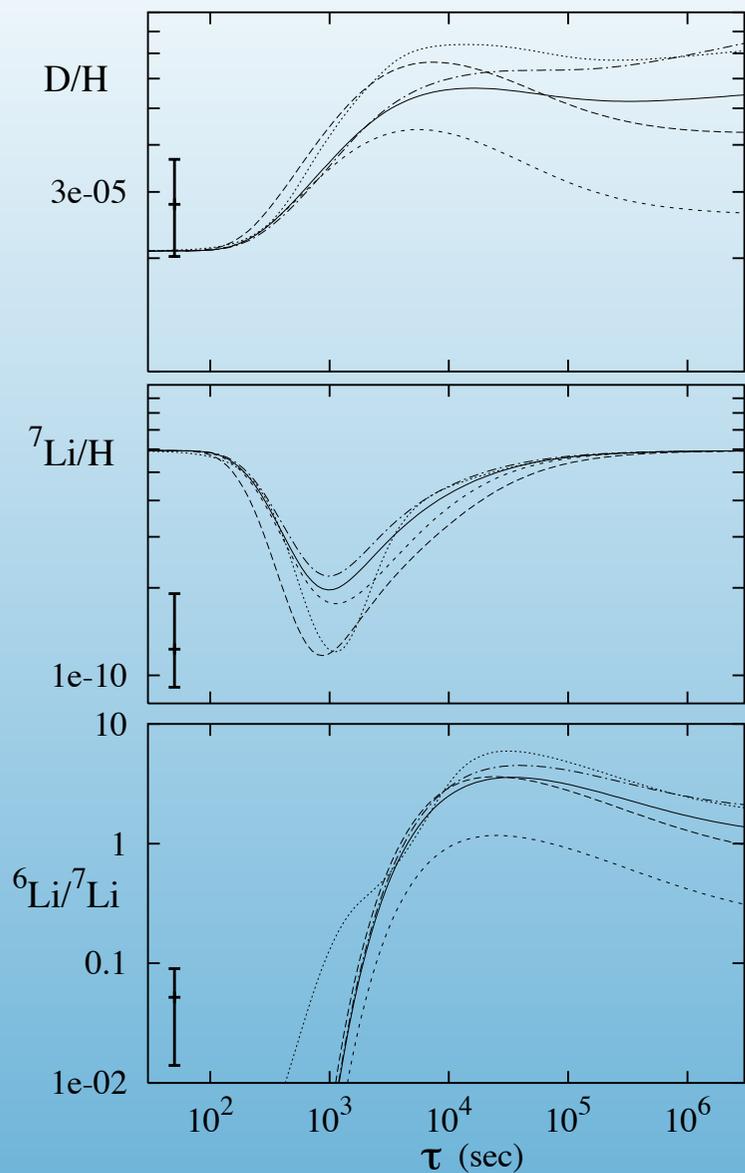


# Limits on Unstable particles due to Electromagnetic/Hadronic Production and Destruction of Nuclei

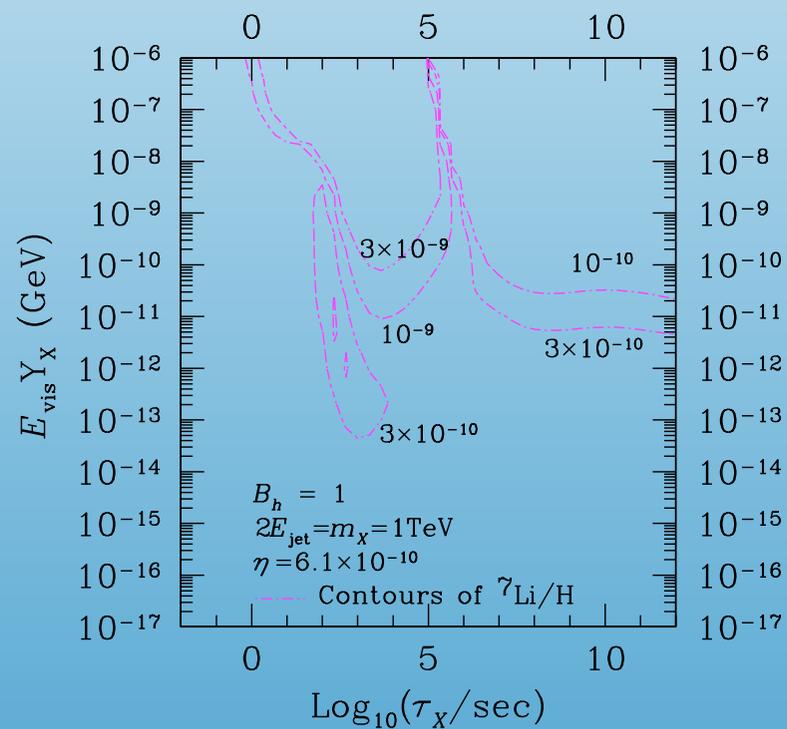
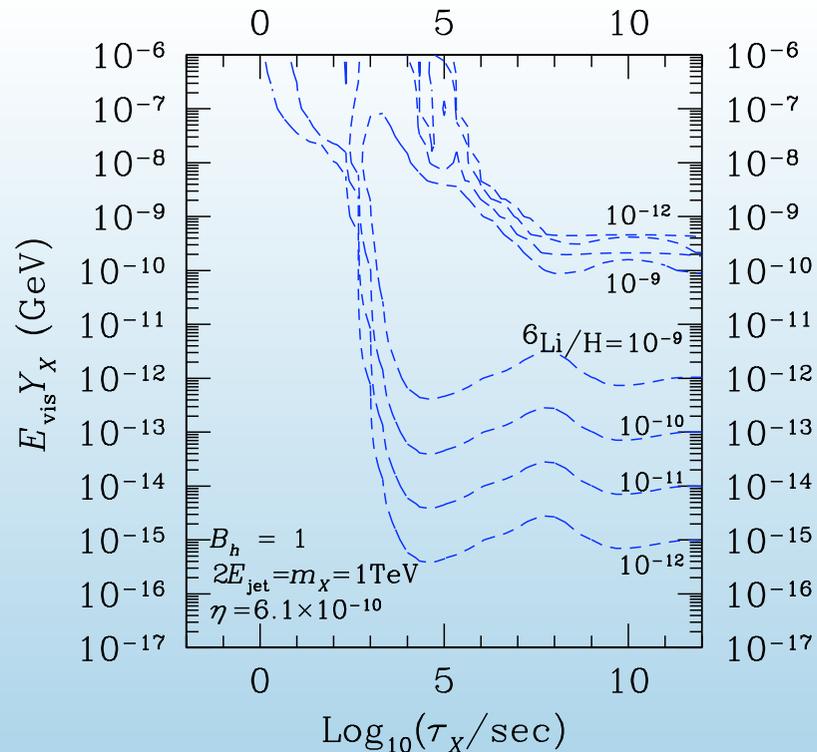
3 free parameters

$$\zeta_X = n_X m_X / n_\gamma = m_X Y_X \eta, \quad m_X, \\ \text{and } \tau_X$$

- Start with non-thermal injection spectrum (Pythia)
- Evolve element abundances including thermal (BBN) and non-thermal processes.



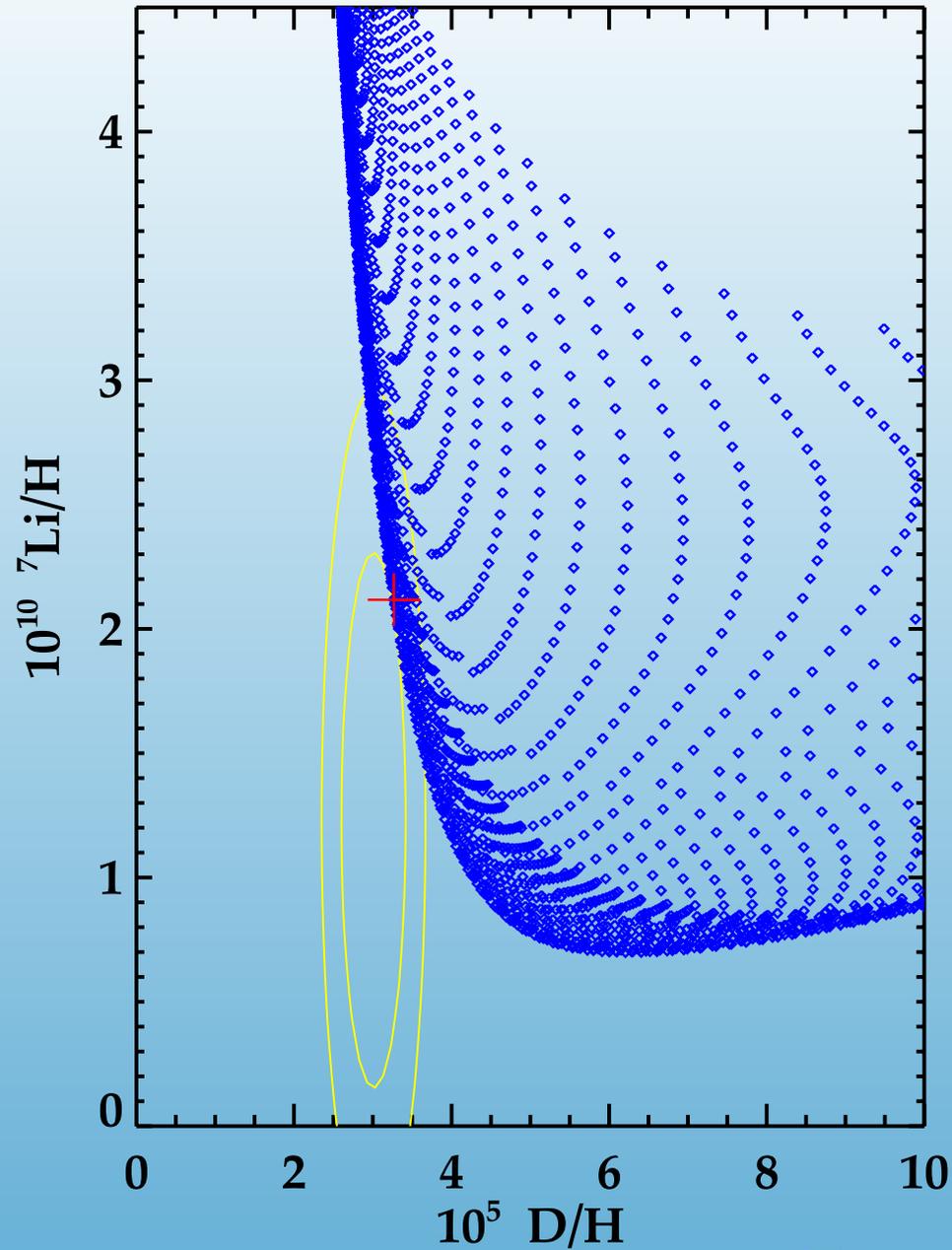
Jedamzik



Kawasaki, Kohri, Moroi

CMSSM

$\tan\beta=16$



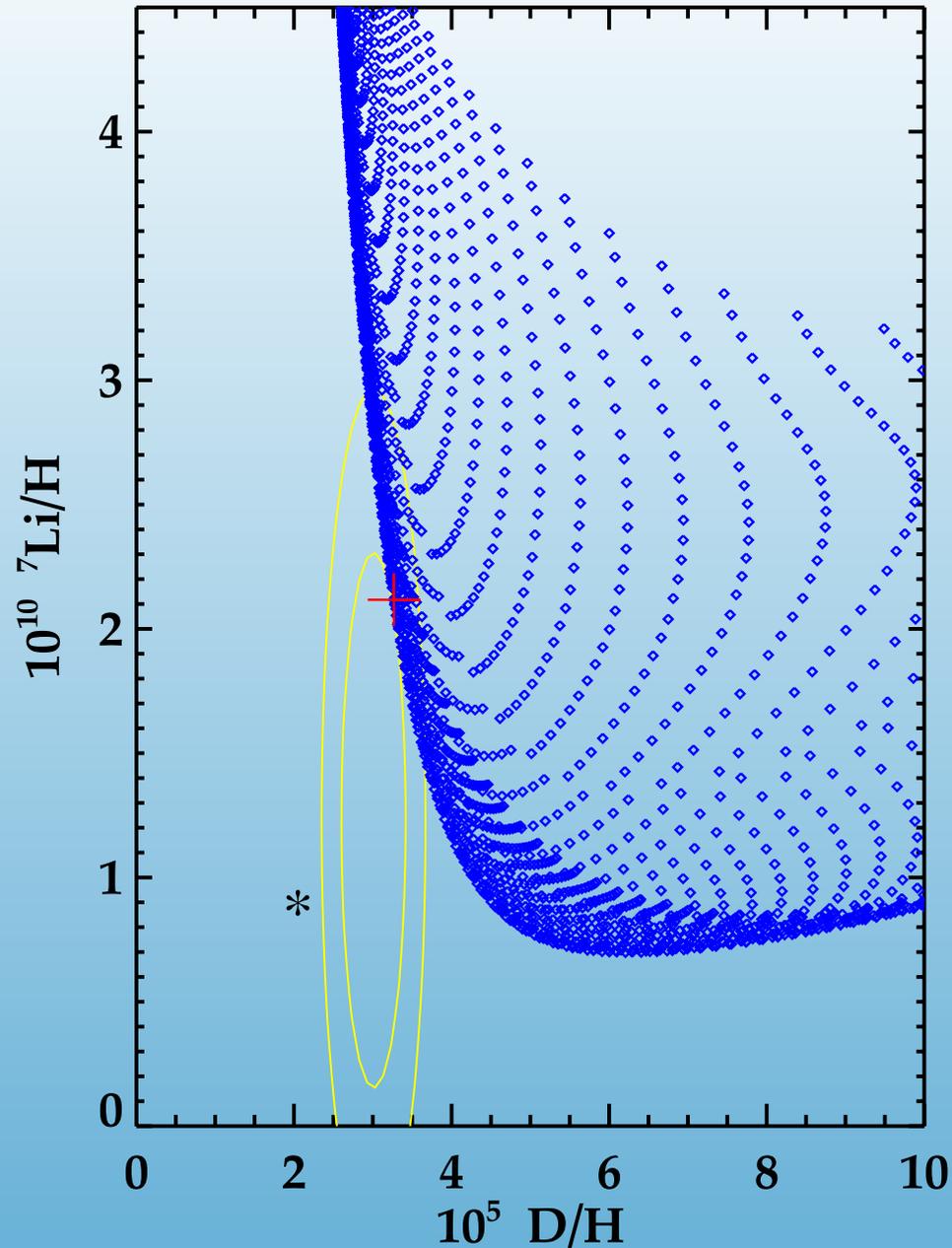
Injection of p,n with  
timescale of  $\sim 1000$  s

${}^7\text{Be}(n,p){}^7\text{Li}$   
followed by  
 ${}^7\text{Li}(p,\alpha){}^4\text{He}$

Olive, Petitjean, Vangioni, Silk  
Cyburt, Ellis, Fields, Luo, Olive, Spanos

CMSSM

$\tan\beta=16$



Injection of p,n with  
timescale of  $\sim 1000$  s

${}^7\text{Be}(n,p){}^7\text{Li}$   
followed by  
 ${}^7\text{Li}(p,\alpha){}^4\text{He}$

\*Uses  $\sigma_{D/H} = \pm 0.27$  should now be  $\pm 0.13$

Olive, Petitjean, Vangioni, Silk  
Cyburt, Ellis, Fields, Luo, Olive, Spanos

# Possible sources for the discrepancy

- Stellar Depletion

- lack of dispersion in the data,  ${}^6\text{Li}$  abundance
- standard models ( $< .05$  dex), models (0.2 - 0.4 dex)

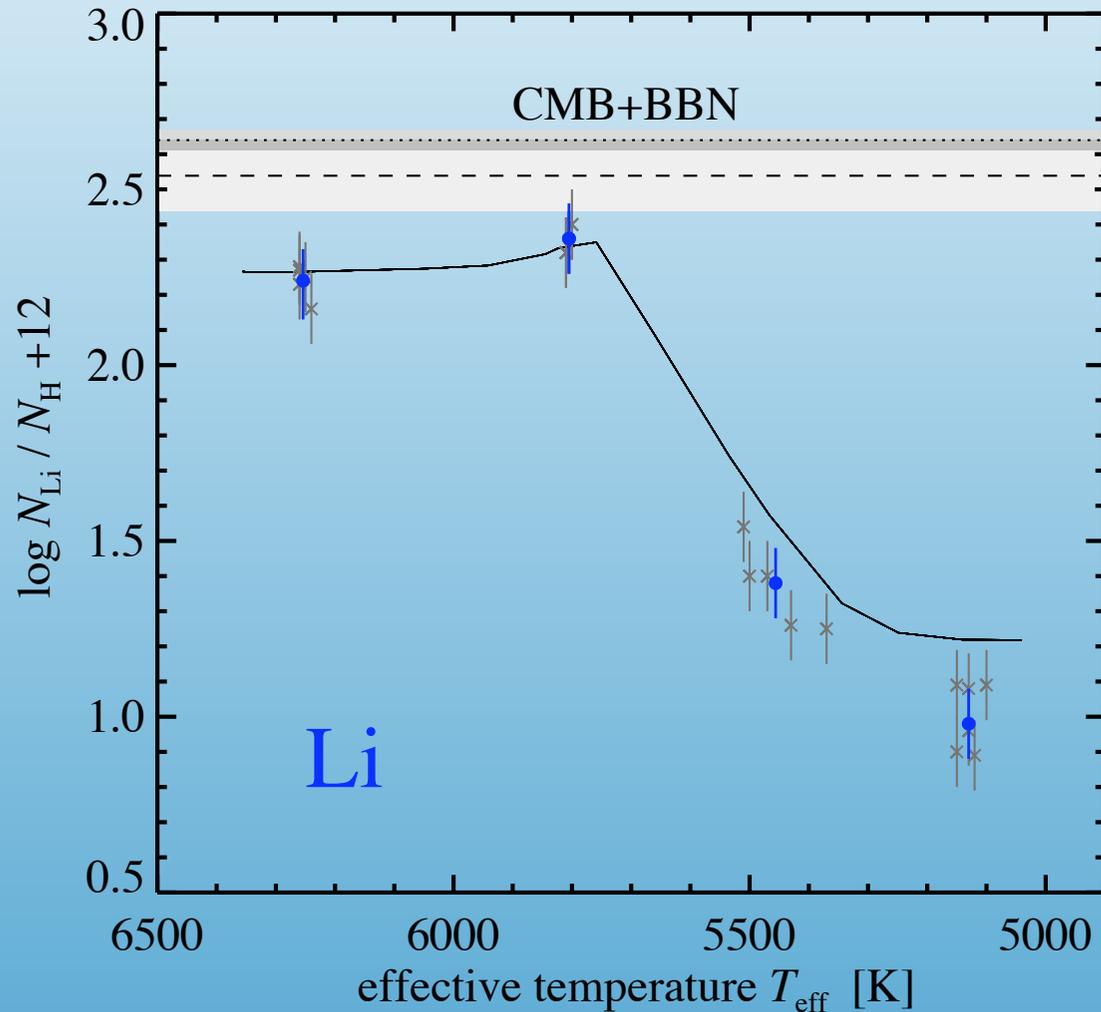
Vauclaire & Charbonnel

Pinsonneault et al.

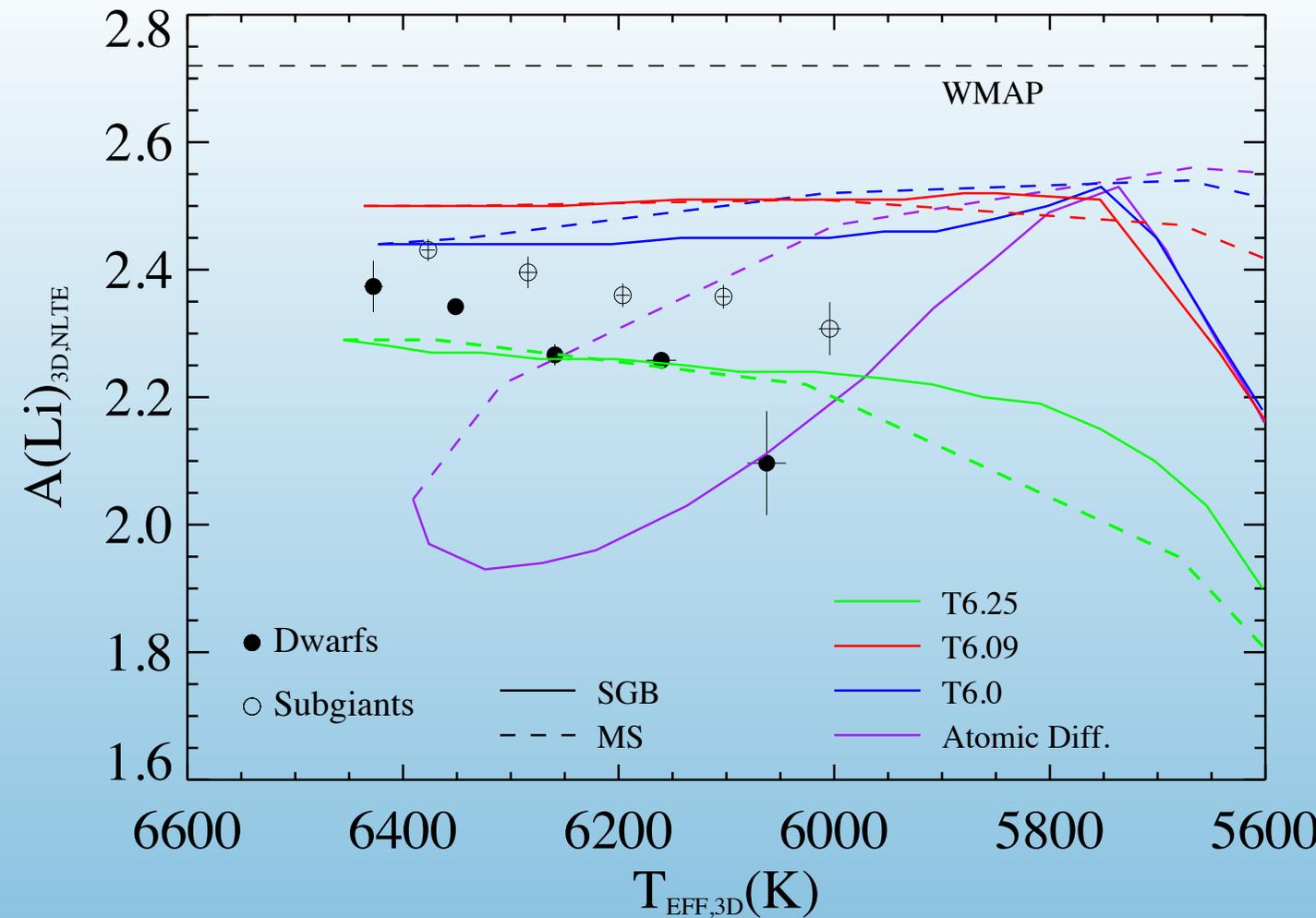
Richard, Michaud, Richer

Korn et al.

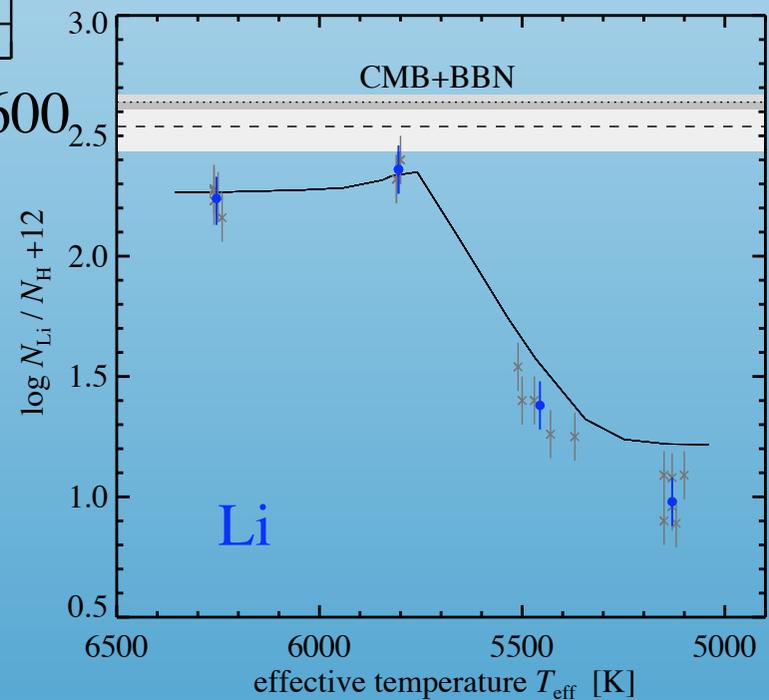
# Stellar Depletion in the Turbulence Model of Korn et al.



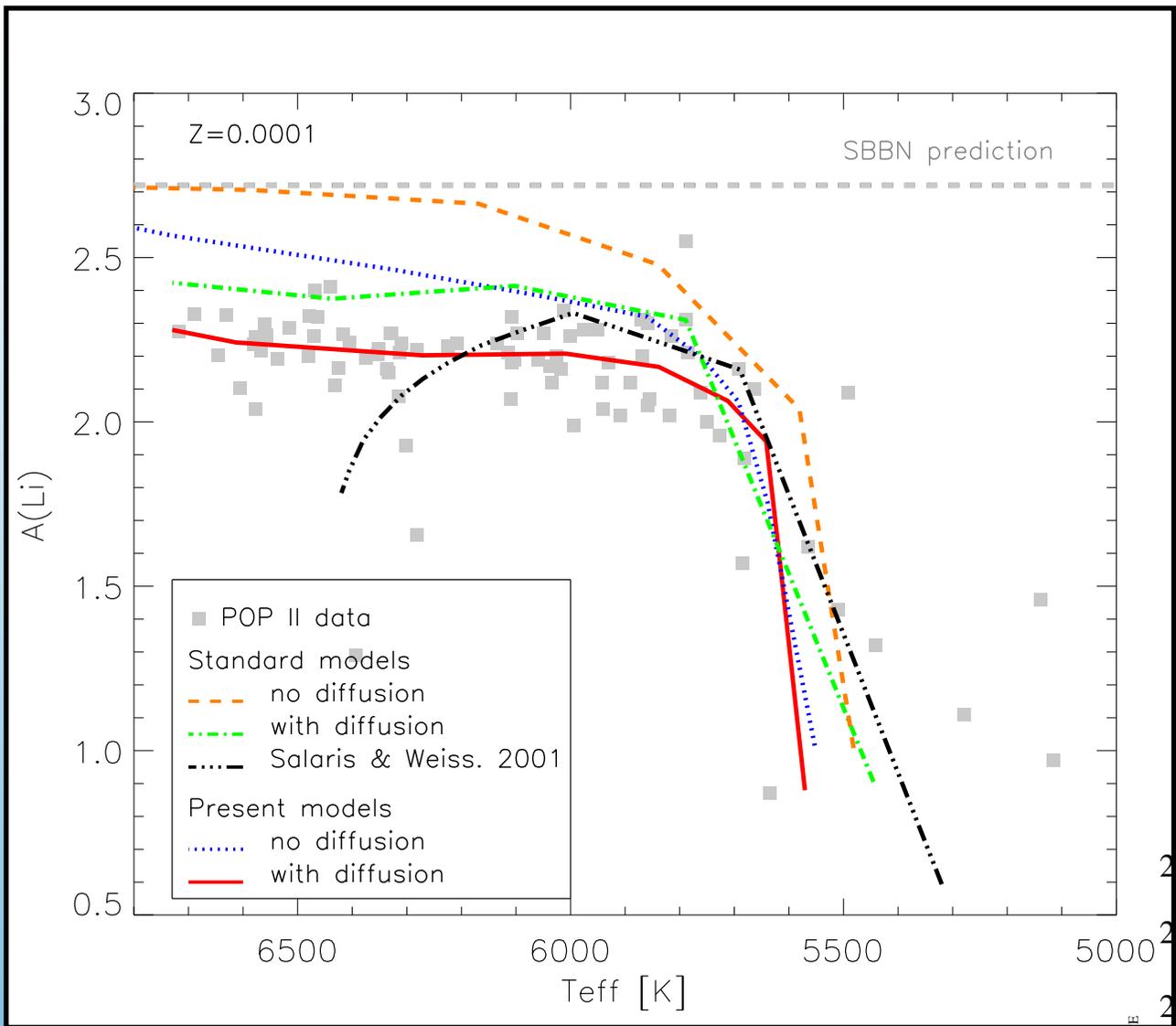
Note new BBN Li result  
pushes primordial value up from  
2.63 to 2.72



But,

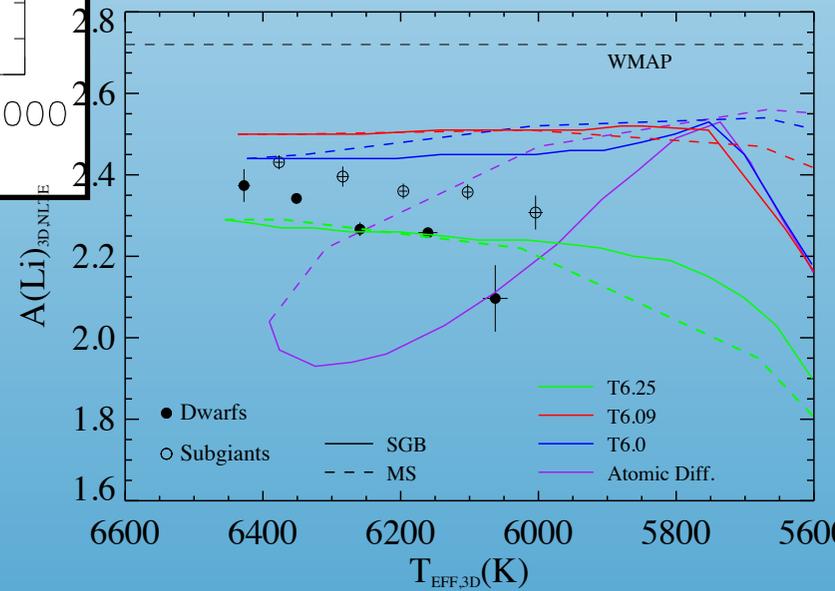


from Gonzáles Hernández et al. (2009)



from Fu et al. (2015)

But,

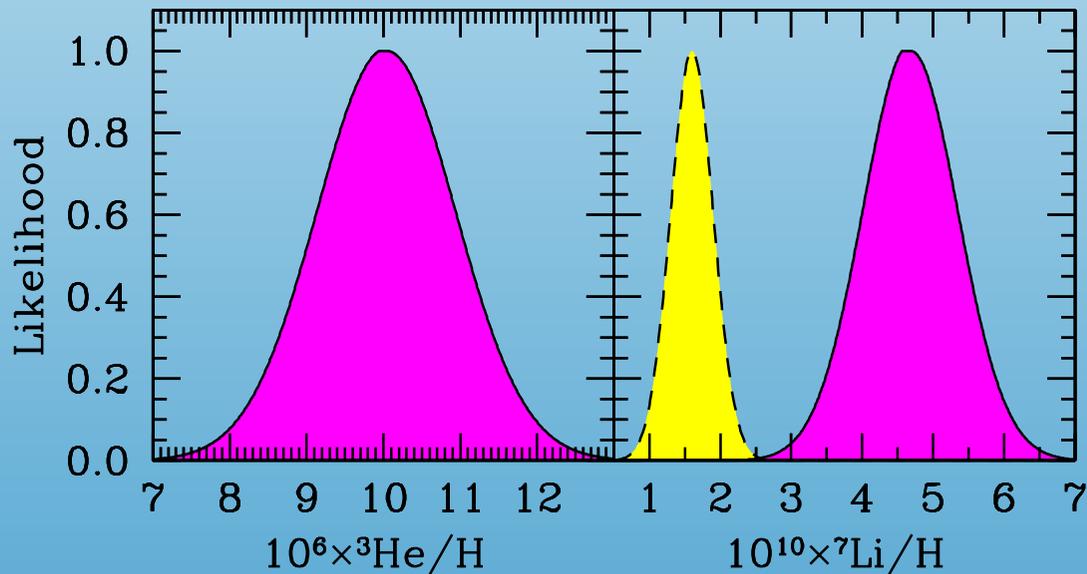
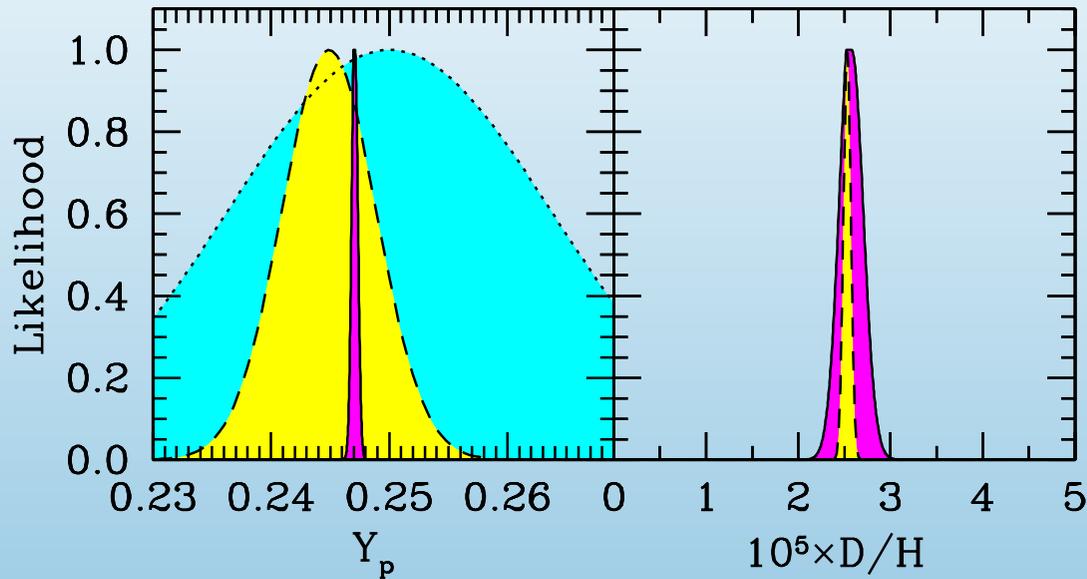


# Other possible sources for the discrepancy

- Stellar parameters
- Axion Cooling
- Variable Constants

# BBN and the CMB

Monte-Carlo approach combining BBN rates, observations and CMB



# BBN and the CMB

## Sensitivities

$$Y_p = 0.24703 \left( \frac{10^{10}\eta}{6.10} \right)^{0.039} \left( \frac{N_\nu}{3.0} \right)^{0.163} \left( \frac{G_N}{G_{N,0}} \right)^{0.35} \left( \frac{\tau_n}{880.3s} \right)^{0.73}$$
$$\times [p(n, \gamma)d]^{0.005} [d(d, n)^3\text{He}]^{0.006} [d(d, p)t]^{0.005}$$

$$\frac{D}{H} = 2.579 \times 10^{-5} \left( \frac{10^{10}\eta}{6.10} \right)^{-1.60} \left( \frac{N_\nu}{3.0} \right)^{0.395} \left( \frac{G_N}{G_{N,0}} \right)^{0.95} \left( \frac{\tau_n}{880.3s} \right)^{0.41}$$
$$\times [p(n, \gamma)d]^{-0.19} [d(d, n)^3\text{He}]^{-0.53} [d(d, p)t]^{-0.47}$$
$$\times [d(p, \gamma)^3\text{He}]^{-0.31} [^3\text{He}(n, p)t]^{0.023} [^3\text{He}(d, p)^4\text{He}]^{-0.012}$$

$$\frac{^3\text{He}}{H} = 9.996 \times 10^{-6} \left( \frac{10^{10}\eta}{6.10} \right)^{-0.59} \left( \frac{N_\nu}{3.0} \right)^{0.14} \left( \frac{G_N}{G_{N,0}} \right)^{0.34} \left( \frac{\tau_n}{880.3s} \right)^{0.15}$$
$$\times [p(n, \gamma)d]^{0.088} [d(d, n)^3\text{He}]^{0.21} [d(d, p)t]^{-0.27}$$
$$\times [d(p, \gamma)^3\text{He}]^{0.38} [^3\text{He}(n, p)t]^{-0.17} [^3\text{He}(d, p)^4\text{He}]^{-0.76} [t(d, n)^4\text{He}]^{-0.009}$$

$$\frac{^7\text{Li}}{H} = 4.648 \times 10^{-10} \left( \frac{10^{10}\eta}{6.10} \right)^{2.11} \left( \frac{N_\nu}{3.0} \right)^{-0.284} \left( \frac{G_N}{G_{N,0}} \right)^{-0.73} \left( \frac{\tau_n}{880.3s} \right)^{0.43}$$
$$\times [p(n, \gamma)d]^{1.34} [d(d, n)^3\text{He}]^{0.70} [d(d, p)t]^{0.065}$$
$$\times [d(p, \gamma)^3\text{He}]^{0.59} [^3\text{He}(n, p)t]^{-0.27} [^3\text{He}(d, p)^4\text{He}]^{-0.75} [t(d, n)^4\text{He}]^{-0.023}$$
$$\times [^3\text{He}(\alpha, \gamma)^7\text{Be}]^{0.96} [^7\text{Be}(n, p)^7\text{Li}]^{-0.71} [^7\text{Li}(p, \alpha)^4\text{He}]^{-0.056} [t(\alpha, \gamma)^7\text{Li}]^{0.030}$$

# BBN and the CMB

## Convolved Likelihoods

From Planck:

$\mathcal{L}_{\text{PLA-base-yhe}}(\omega_b, Y_p)$  and  $\mathcal{L}_{\text{PLA-base-nnu-yhe}}(\omega_b, Y_p, N_\nu)$

$$\omega_b = 0.022305 \pm 0.000225$$

$$Y_p = 0.25003 \pm 0.01367$$

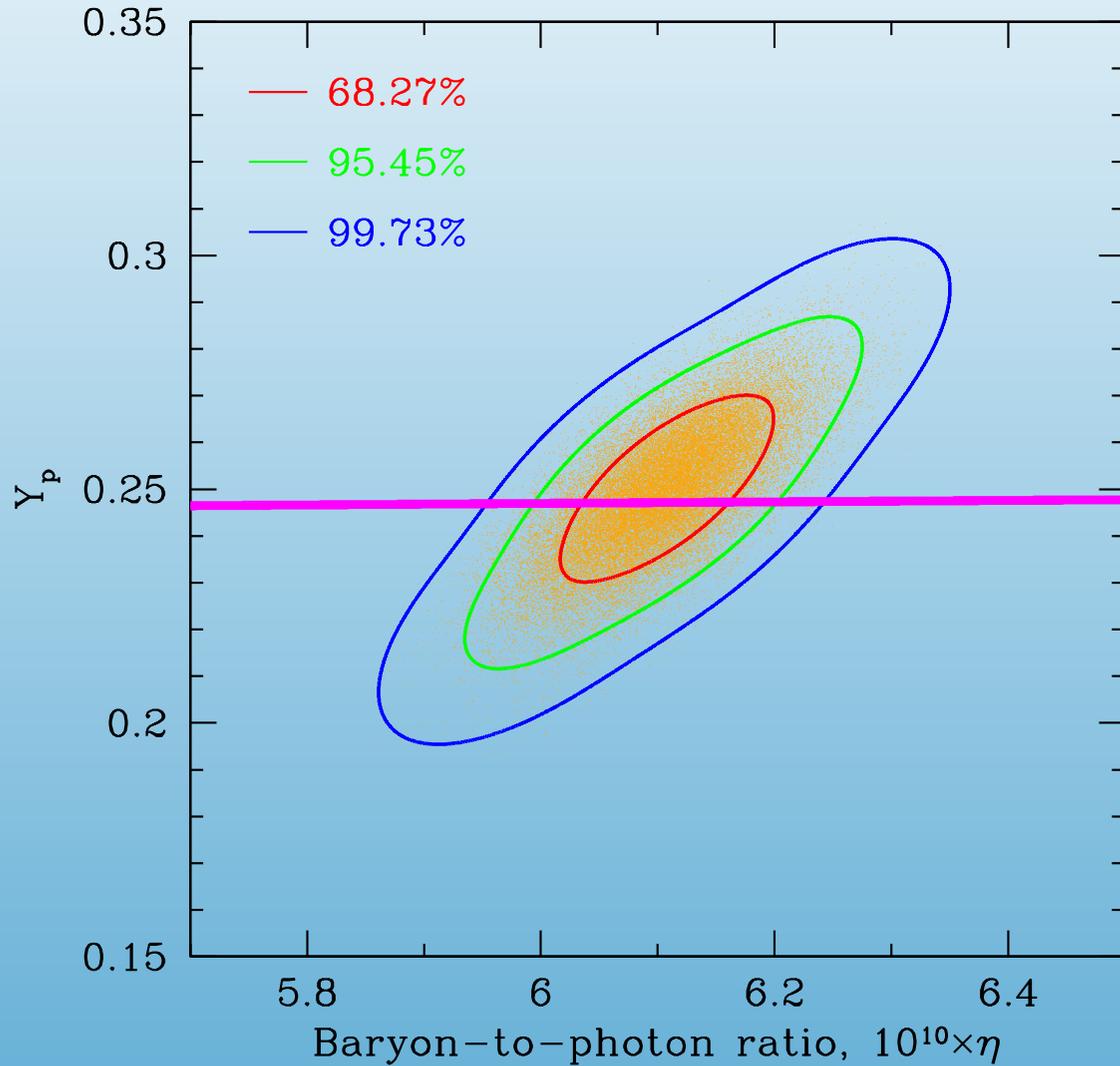
$$\omega_b = 0.022212 \pm 0.000242$$

$$N_{\text{eff}} = 2.7542 \pm 0.3064$$

$$Y_p = 0.26116 \pm 0.01812$$

# BBN and the CMB

$N_v = 3$



CMB only determination  
of  $\eta$  and  $Y_P$

$3\sigma$  BBN Prediction

# BBN and the CMB

## Convolved Likelihoods

### Results for $\eta$ ( $N_\nu$ )

Constraints Used	$\eta \times 10^{10}$
CMB-only	$6.108 \pm 0.060$
BBN+ $Y_p$	$4.87^{+2.46}_{-1.54}$
BBN+D	$6.180 \pm 0.195$
BBN+ $Y_p$ +D	$6.172 \pm 0.195$
CMB+BBN	$6.098 \pm 0.042$
CMB+BBN+ $Y_p$	$6.098 \pm 0.042$
CMB+BBN+D	$6.102 \pm 0.041$
CMB+BBN+ $Y_p$ +D	$6.101 \pm 0.041$

Constraints Used	$\eta_{10}$	$N_\nu$
CMB-only	$6.08 \pm 0.07$	$2.67^{+0.30}_{-0.27}$
BBN+ $Y_p$ +D	$6.10 \pm 0.23$	$2.85 \pm 0.28$
CMB+BBN	$6.08 \pm 0.07$	$2.91 \pm 0.20$
CMB+BBN+ $Y_p$	$6.07 \pm 0.06$	$2.89 \pm 0.16$
CMB+BBN+D	$6.07 \pm 0.07$	$2.90 \pm 0.19$
CMB+BBN+ $Y_p$ +D	$6.07 \pm 0.06$	$2.88 \pm 0.16$

# BBN and the CMB

## Convolved Likelihoods

### Results for $\eta$ ( $N_\nu$ )

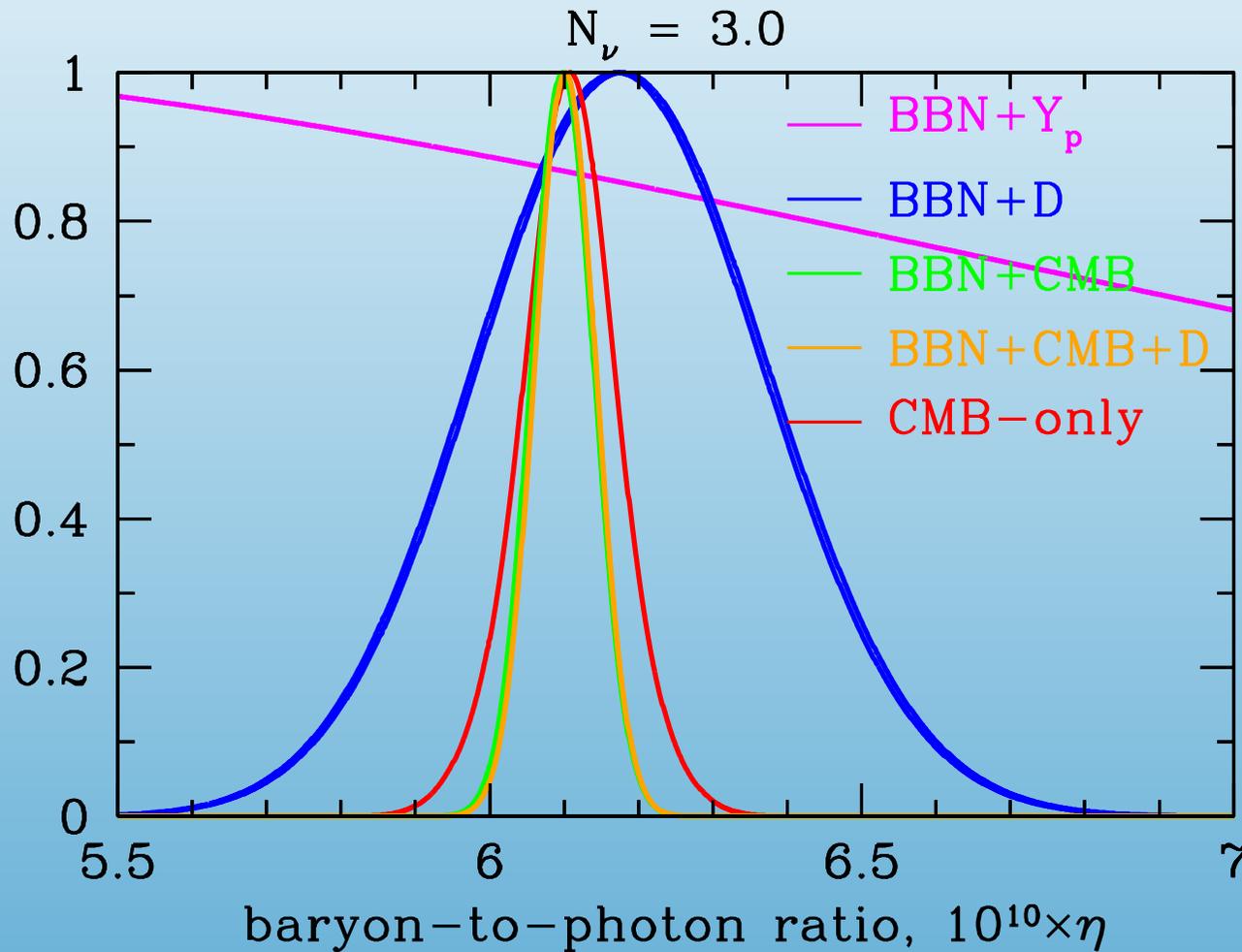
Constraints Used	$\eta \times 10^{10}$
CMB-only	$6.108 \pm 0.060$
BBN+ $Y_p$	$4.87^{+2.46}_{-1.54}$
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$N_\nu < 3.2$  (95% CL)

# BBN and the CMB

Convolved Likelihoods

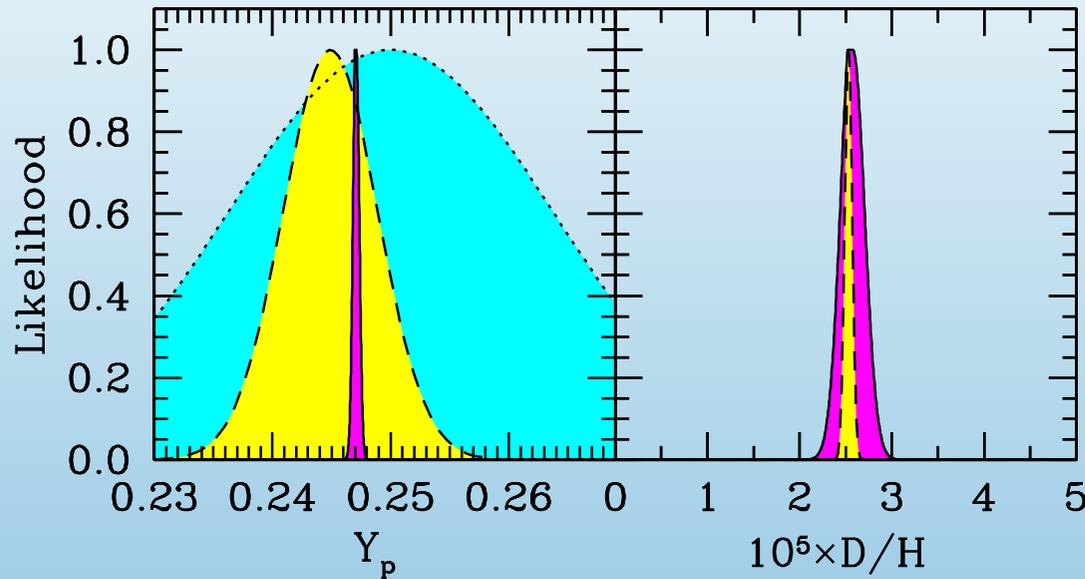


Determination of  $\eta$

# BBN and the CMB

$$\mathcal{L}(X_i) \propto \int \mathcal{L}_{\text{PLA-base-yhe}}(\omega_b, Y_p) \mathcal{L}_{\text{BBN}}(\eta; \{X_i\}) d\eta$$

recall  $\omega_b$  and  $\eta$  are related



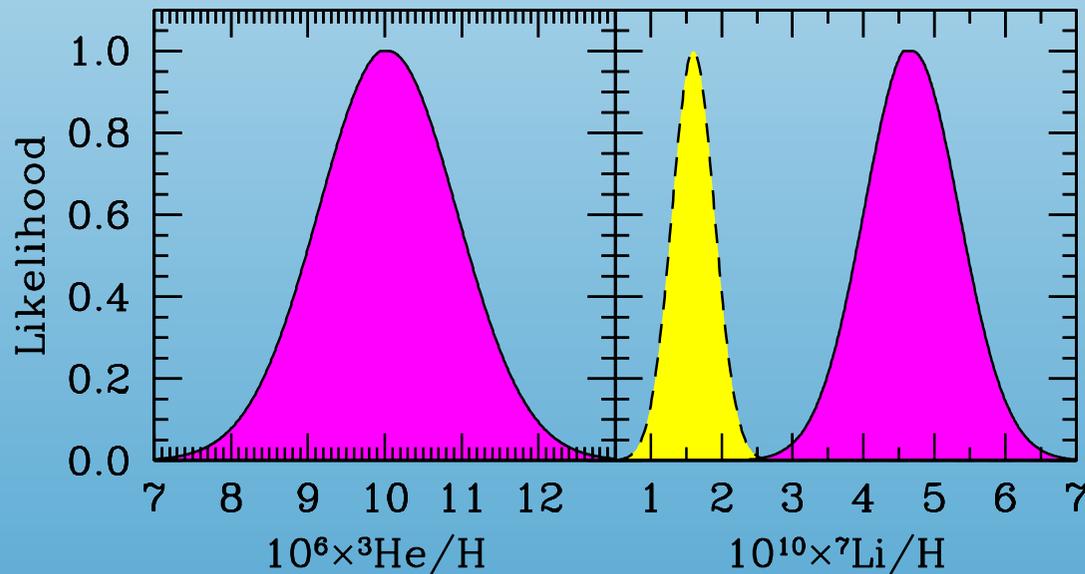
$$Y_p = 0.24709 \pm 0.00025$$

$$D/H = (2.58 \pm 0.13) \times 10^{-5}$$

$${}^3\text{He}/\text{H} = (10.039 \pm 0.090) \times 10^{-5}$$

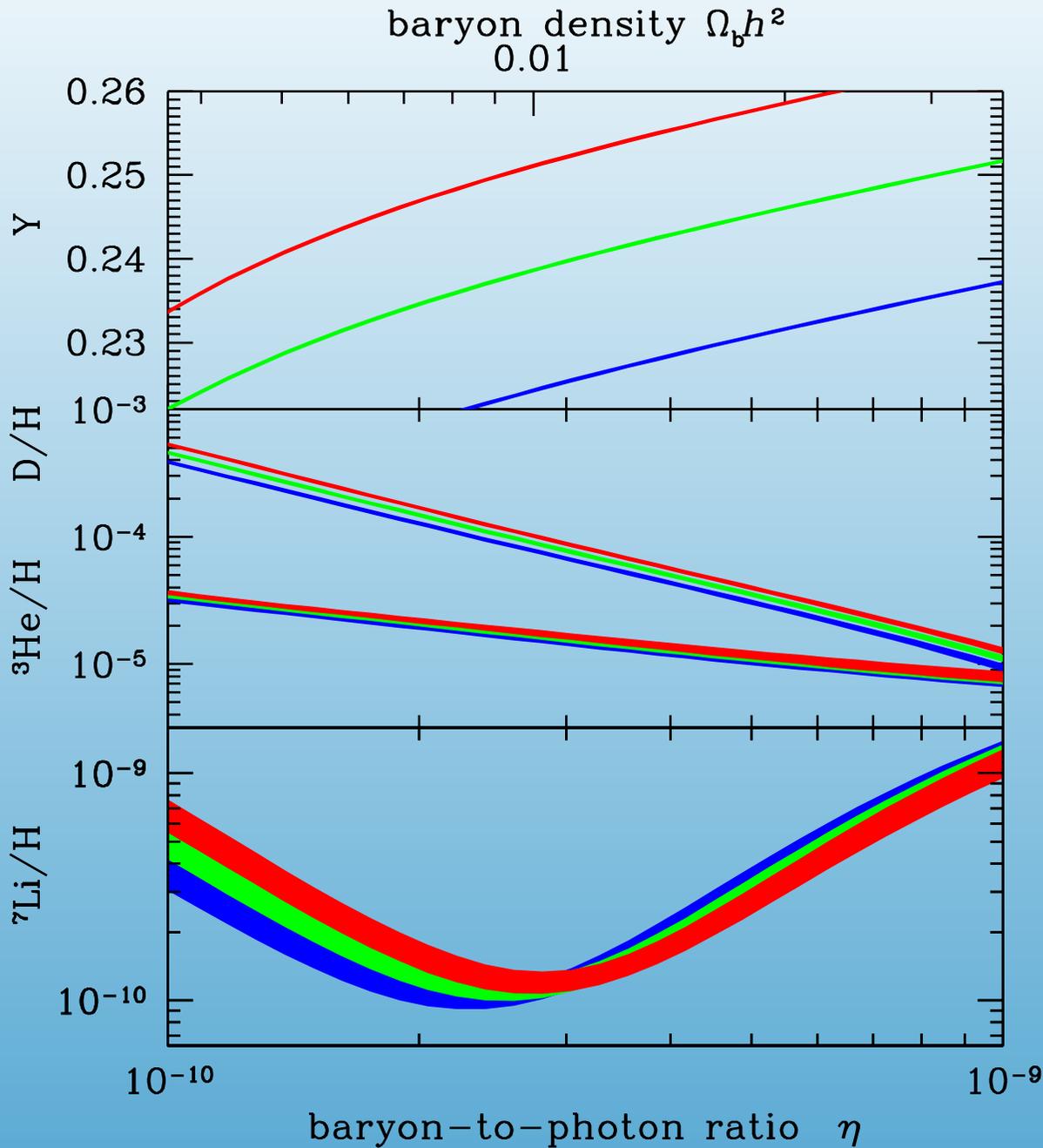
$${}^7\text{Li}/\text{H} = (4.68 \pm 0.67) \times 10^{-10}$$

$$\log_{10} ({}^6\text{Li}/\text{H}) = -13.89 \pm 0.20$$



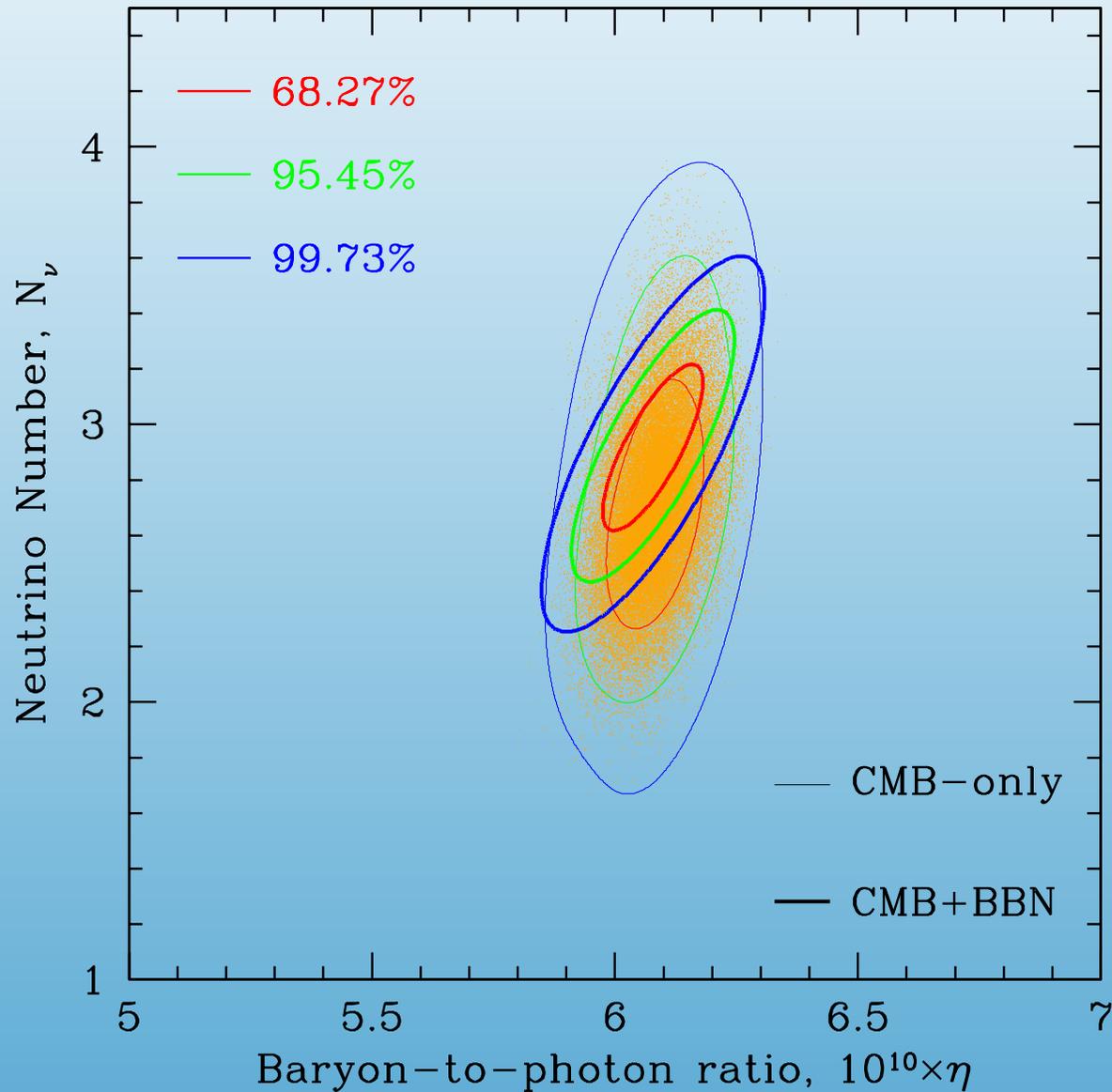
Cyburt, Fields, Olive, Yeh

# BBN and the CMB



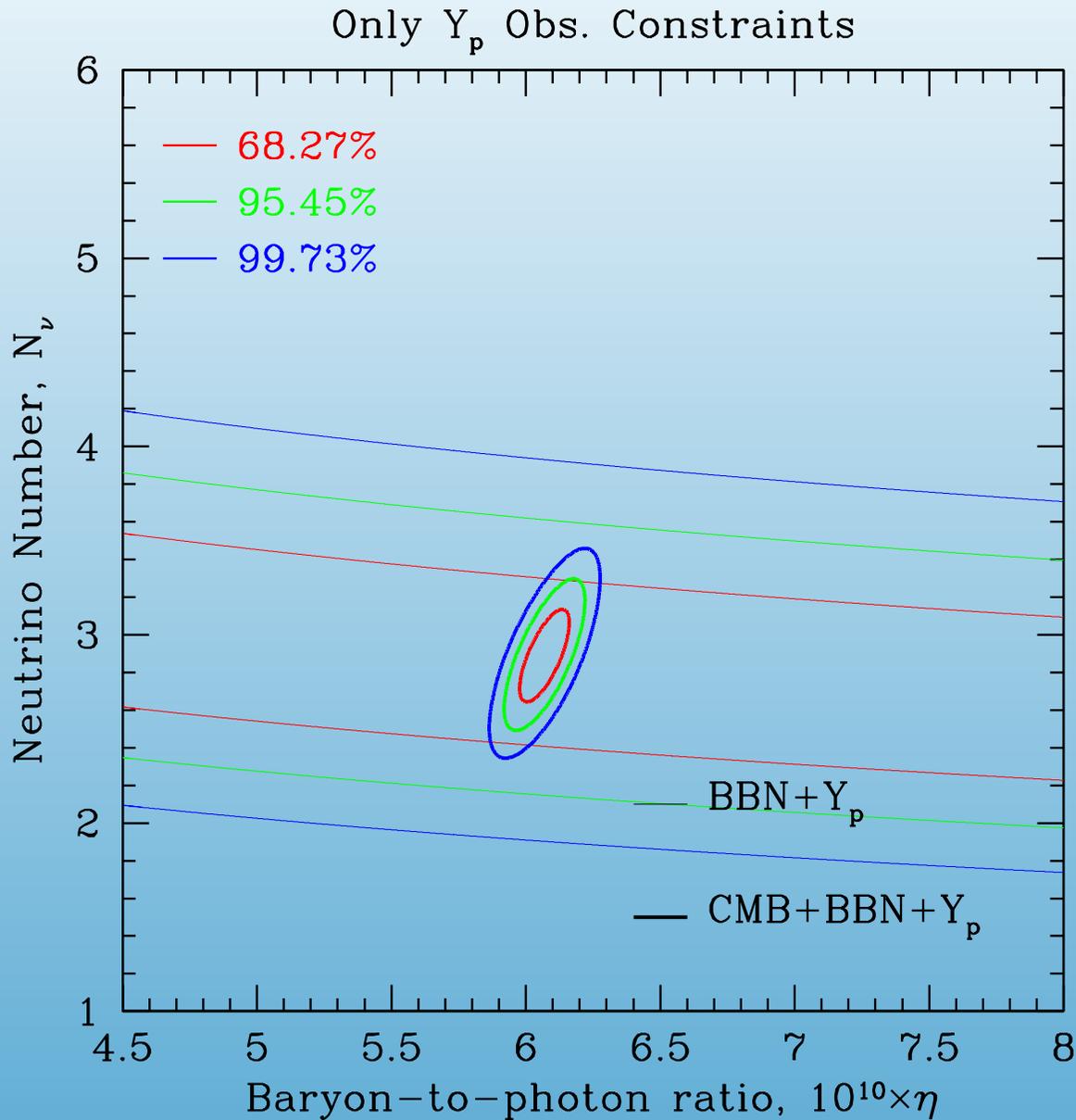
Sensitivity to  $N_\nu$

# BBN and the CMB



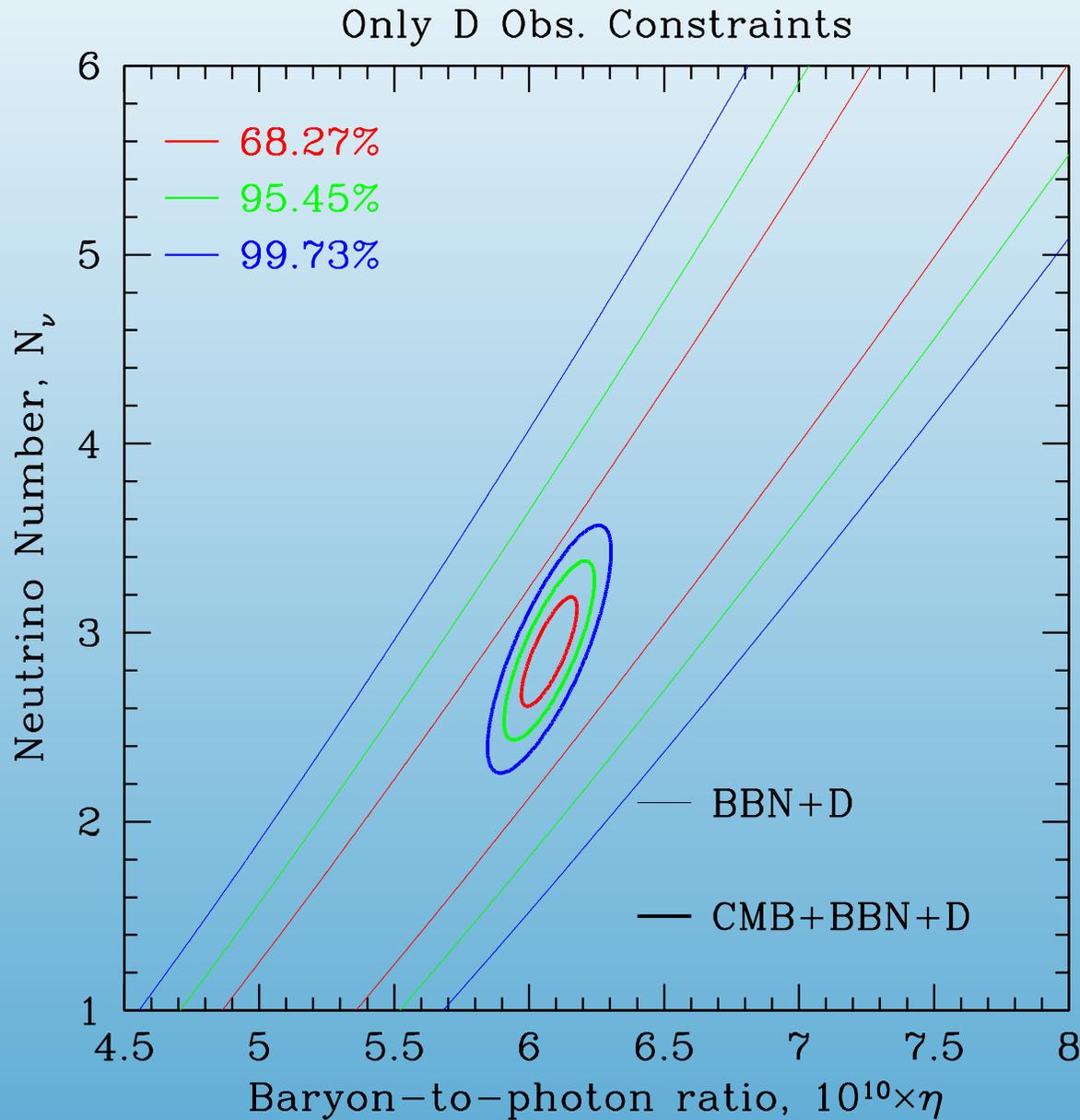
CMB only determination  
of  $\eta$  and  $N_\nu$

# BBN and the CMB



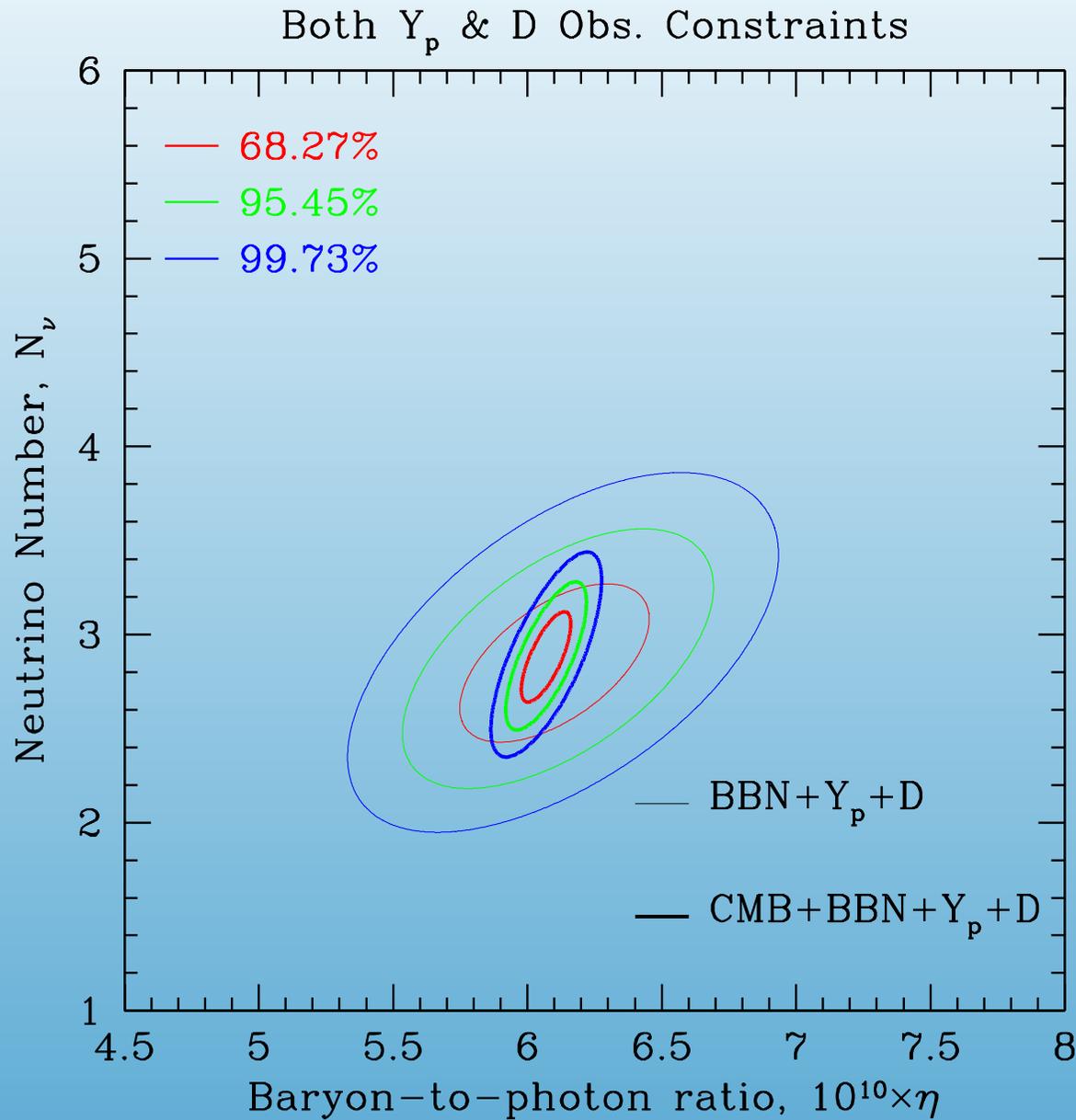
CMB and BBN determination  
of  $\eta$  and  $N_\nu$

# BBN and the CMB



CMB and BBN determination  
of  $\eta$  and  $N_\nu$

# BBN and the CMB



CMB and BBN determination  
of  $\eta$  and  $N_\nu$

# Summary

- BBN and CMB are in excellent agreement wrt D and He
- Li: Problematic
  - BBN  ${}^7\text{Li}$  high compared to observations
- Important to consider:
  - Nuclear considerations
  - Depletion (tuned)
  - BSM solutions?
- Standard Model ( $N_\nu = 3$ ) is looking good!