

# Ultralight Dark Matter and the Precision Frontier of Particle Physics

Peter Graham

Stanford

# Precision Experiments

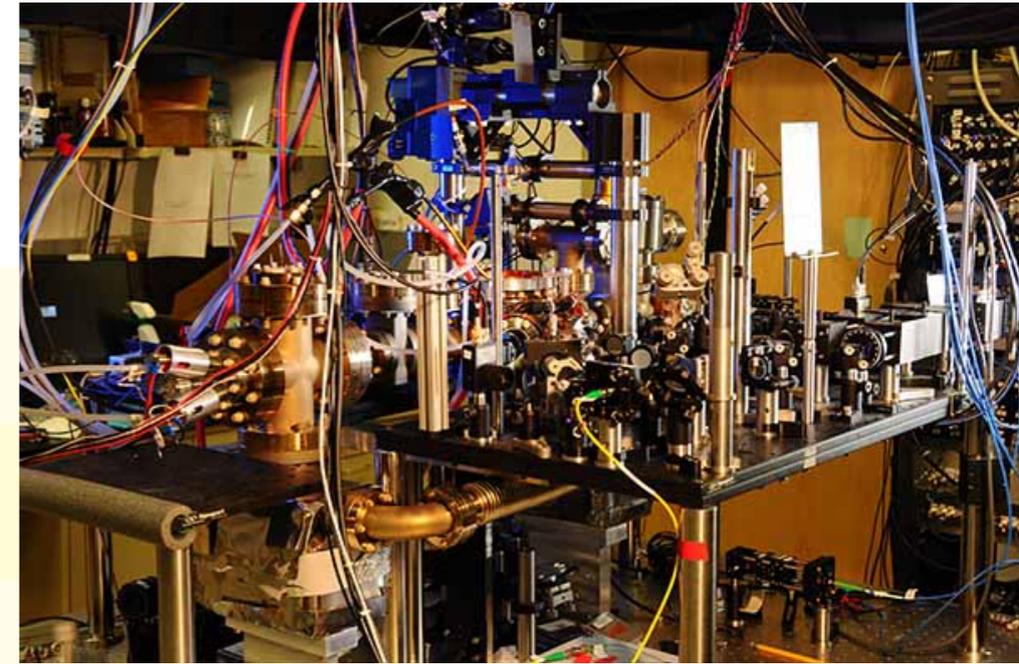
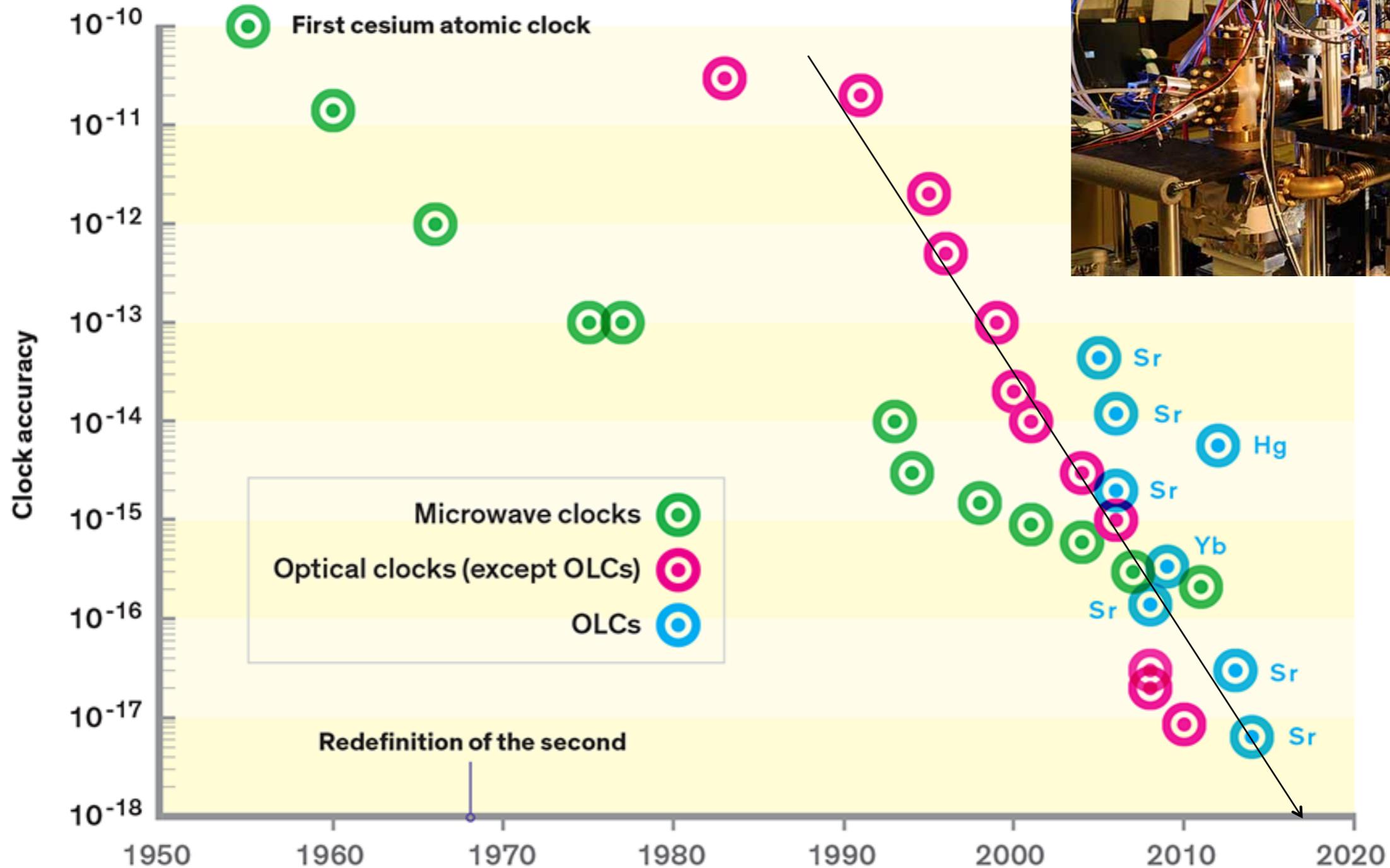
Precision measurement offers a powerful new approach for particle physics  
beyond conventional particle colliders/detectors

not completely new (e.g. EDMs, new forces, etc.),  
but small compared to traditional particle detection

- New technologies rapidly pushing precision measurement
  - e.g. atomic clocks have 18 digit precision
- Required for axions, gravitational waves...
  - critical questions such as hierarchy problem or nature of dark matter may not be answered at weak scale

Many promising, unexplored directions

# Atomic Clock Sensitivity



current technology already allows many new searches, and will improve by orders of magnitude

# Outline

## 1. Dark Matter Detection

## 2. Axion Dark Matter

1. Cosmic Axion Spin Precession Experiment (CASPER) (D. Budker + A. Sushkov)

2. DM Radio (K. Irwin)

## 3. Other techniques

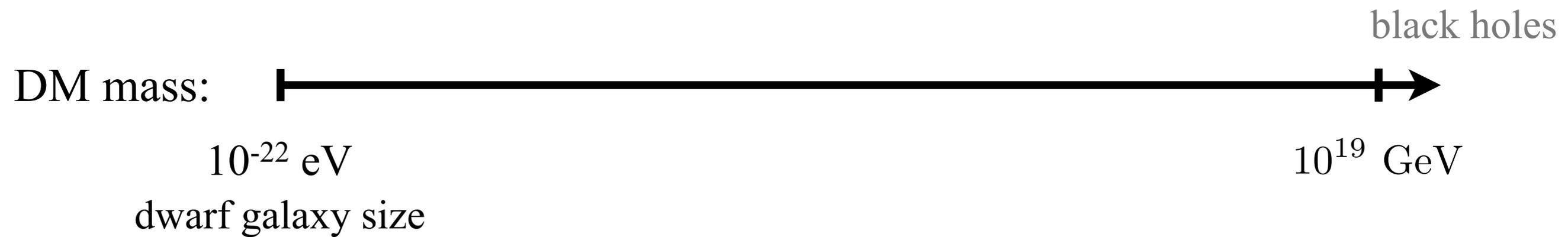
## 4. Gravitational waves and dark matter with atom interferometry

(M. Kasevich + J. Hogan)

# Dark Matter Detection

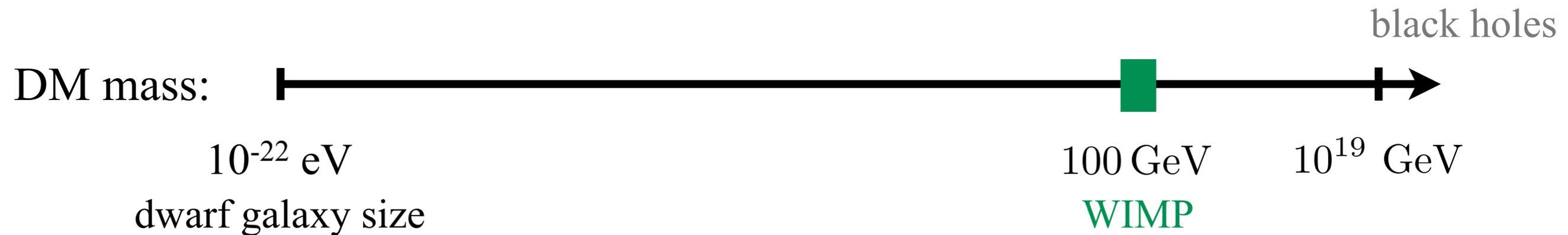
# Dark Matter Candidates

What do we know about dark matter?



# Dark Matter Candidates

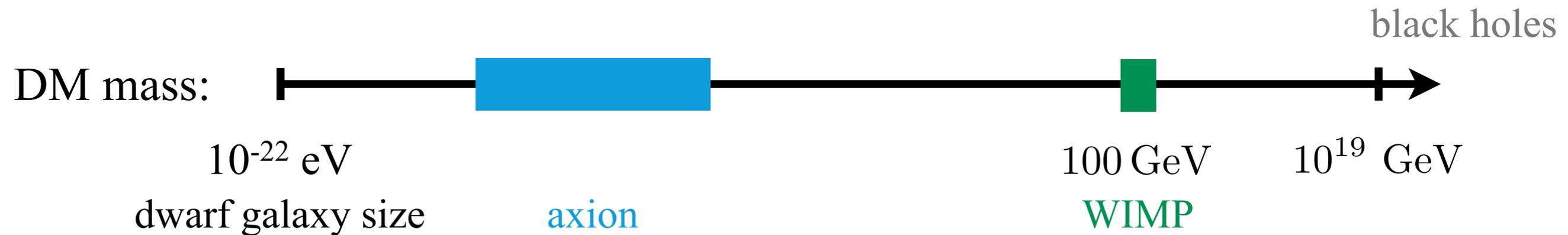
What do we know about dark matter?



WIMP is well-motivated, significant direct detection effort focused on WIMPs

# Dark Matter Candidates

What do we know about dark matter?

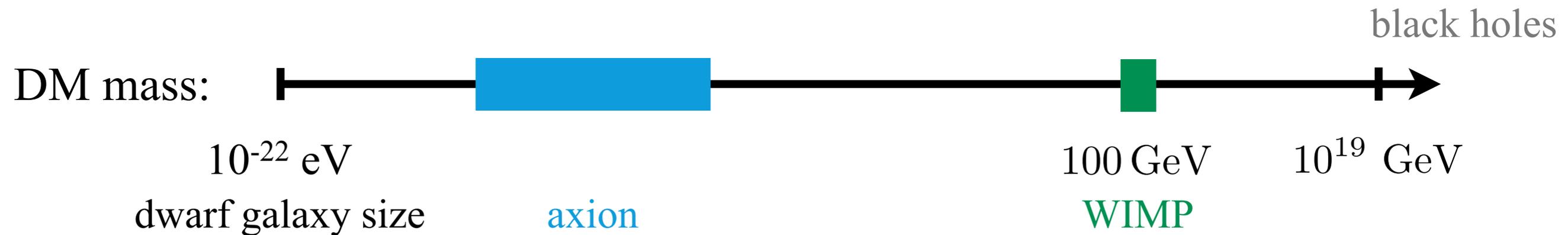


WIMP is well-motivated, significant direct detection effort focused on WIMPs

Axion is other best-motivated candidate, only a small fraction of parameter space covered

# Dark Matter Candidates

What do we know about dark matter?



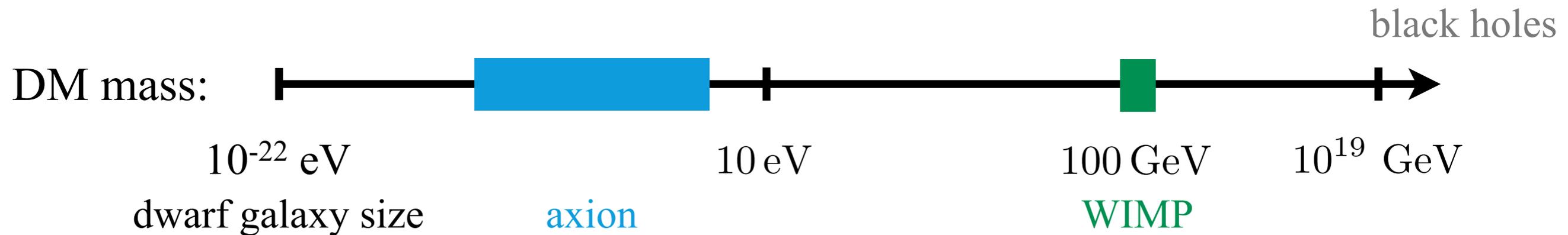
WIMP is well-motivated, significant direct detection effort focused on WIMPs

Axion is other best-motivated candidate, only a small fraction of parameter space covered

Huge DM parameter space currently unexplored!

# Direct Detection

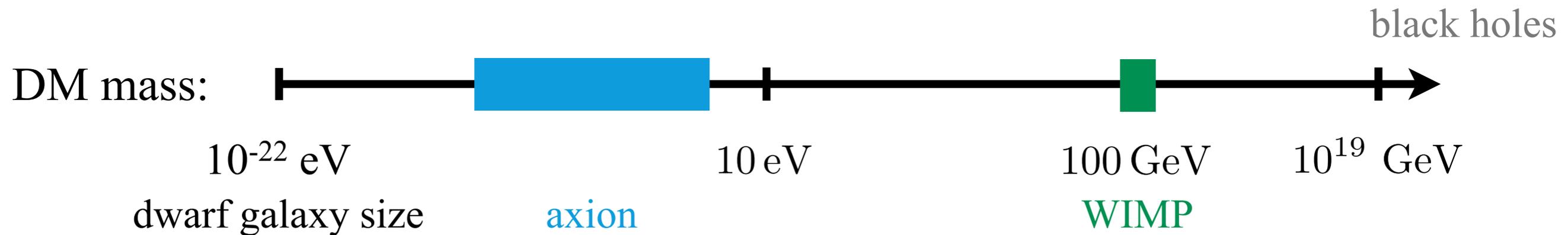
How can we detect DM?



$$\rho_{\text{DM}} \approx 0.3 \frac{\text{GeV}}{\text{cm}^3} \approx (0.04 \text{ eV})^4 \rightarrow \text{high phase space density if } m \lesssim 10 \text{ eV}$$

# Direct Detection

How can we detect DM?



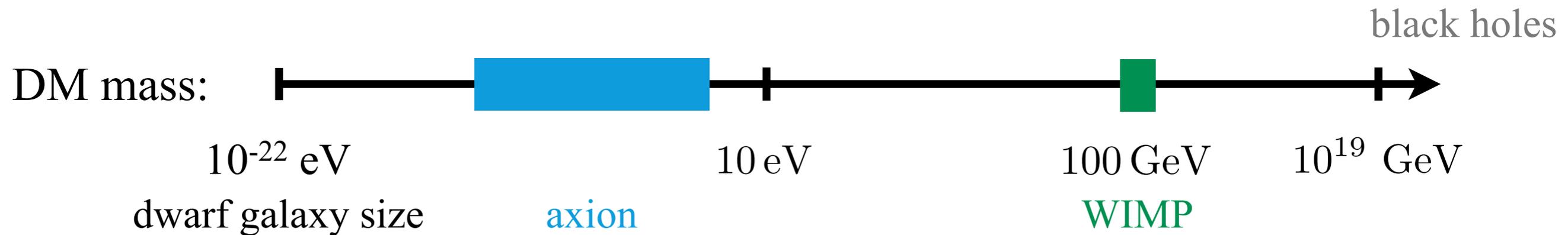
$$\rho_{\text{DM}} \approx 0.3 \frac{\text{GeV}}{\text{cm}^3} \approx (0.04 \text{ eV})^4 \rightarrow \text{high phase space density if } m \lesssim 10 \text{ eV}$$

field-like (e.g. axion)  
new detectors required

particle-like (e.g. WIMP)  
particle detectors best

# Direct Detection

How can we detect DM?



$$\rho_{\text{DM}} \approx 0.3 \frac{\text{GeV}}{\text{cm}^3} \approx (0.04 \text{ eV})^4 \rightarrow \text{high phase space density if } m \lesssim 10 \text{ eV}$$

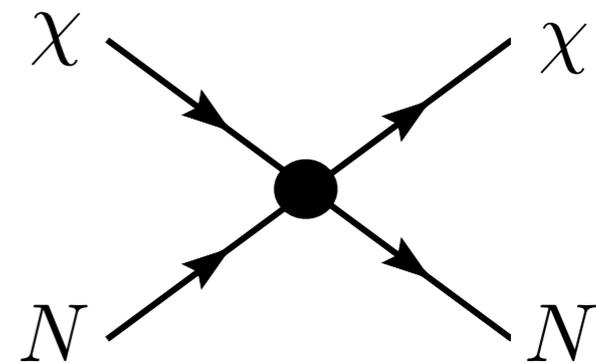
field-like (e.g. axion)  
new detectors required

Detect coherent effects of entire field  
(like gravitational wave detector)



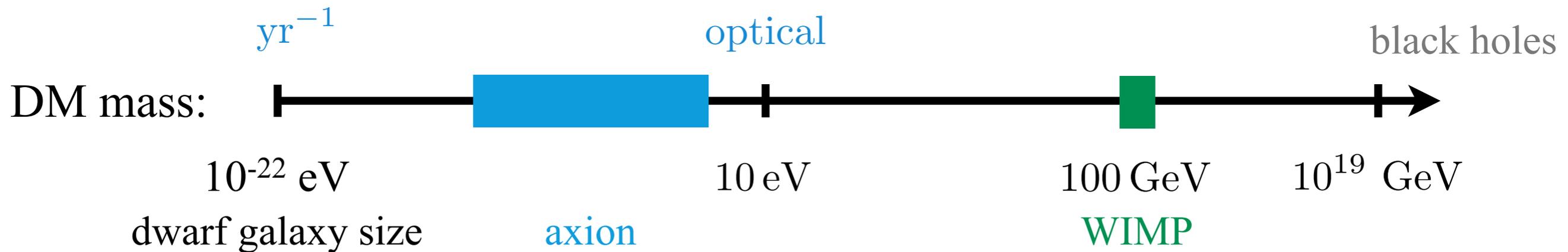
particle-like (e.g. WIMP)  
particle detectors best

Search for single, hard particle scattering



# Direct Detection

How can we detect DM?



$$\rho_{\text{DM}} \approx 0.3 \frac{\text{GeV}}{\text{cm}^3} \approx (0.04 \text{ eV})^4 \rightarrow \text{high phase space density if } m \lesssim 10 \text{ eV}$$

field-like (e.g. axion)  
new detectors required

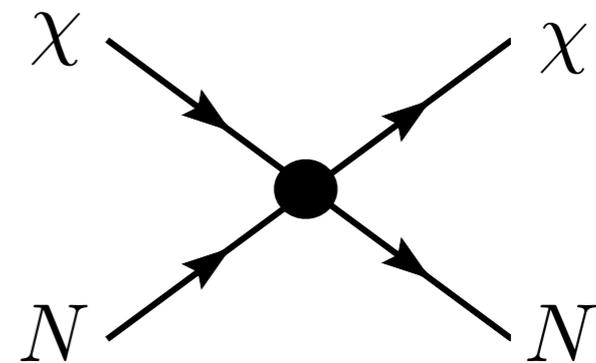
Detect coherent effects of entire field  
(like gravitational wave detector)



Frequency range accessible!

particle-like (e.g. WIMP)  
particle detectors best

Search for single, hard particle scattering

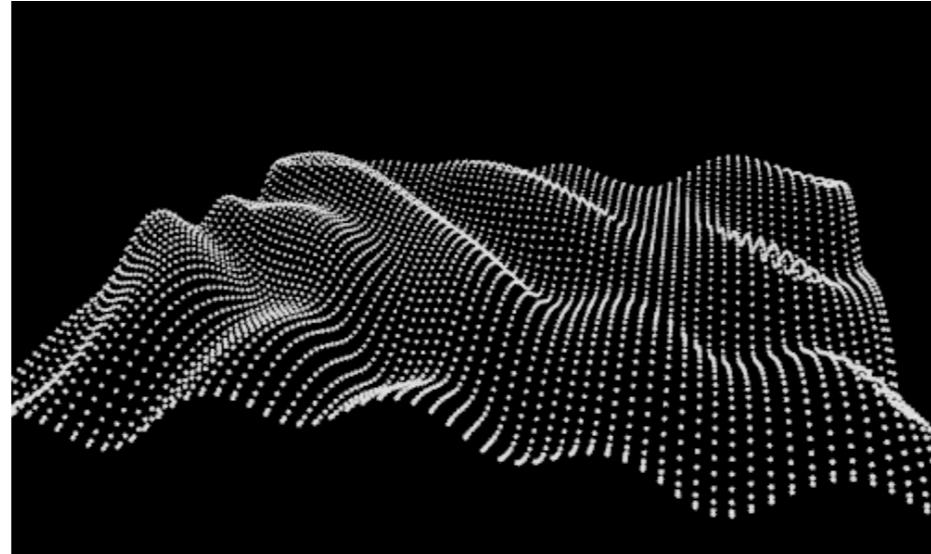


# “Field” Dark Matter

particle DM



DM at long deBroglie wavelength  
useful to picture as a “coherent” field:

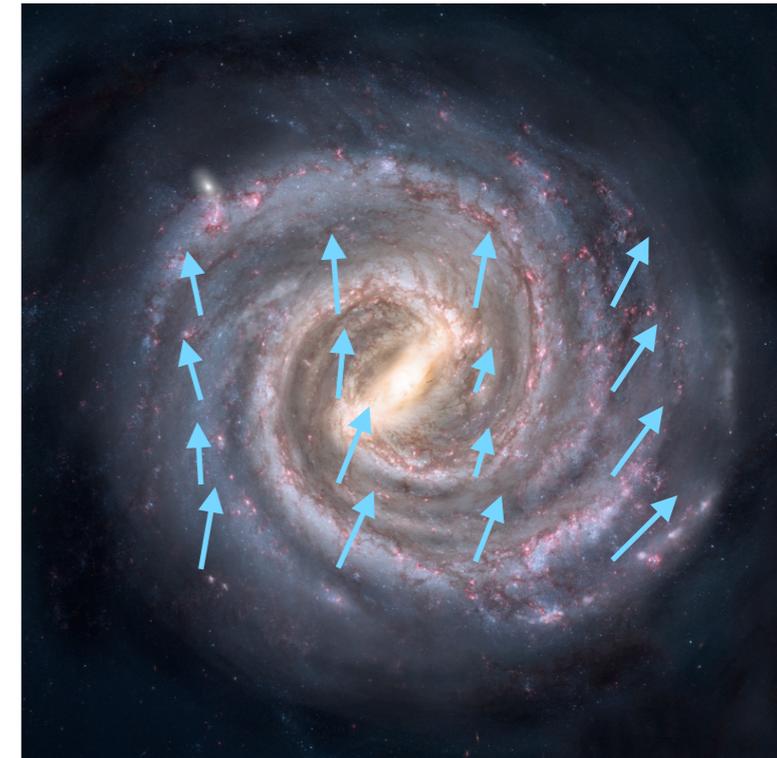
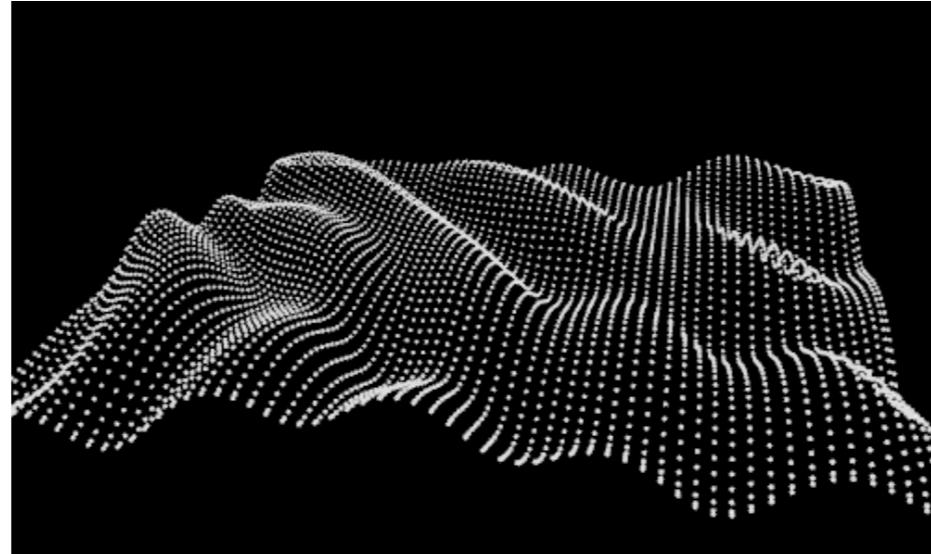


# “Field” Dark Matter

particle DM



DM at long deBroglie wavelength  
useful to picture as a “coherent” field:

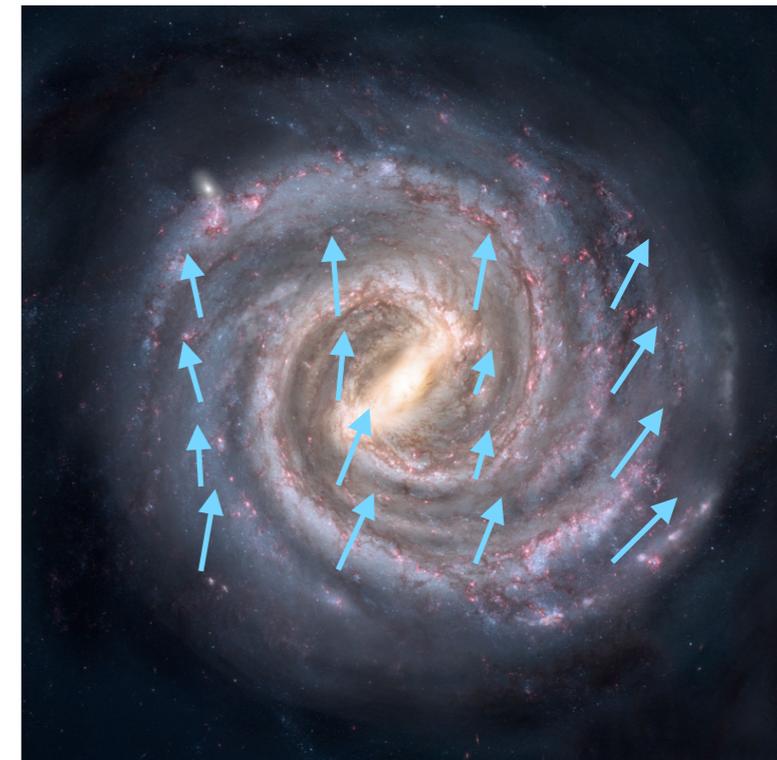
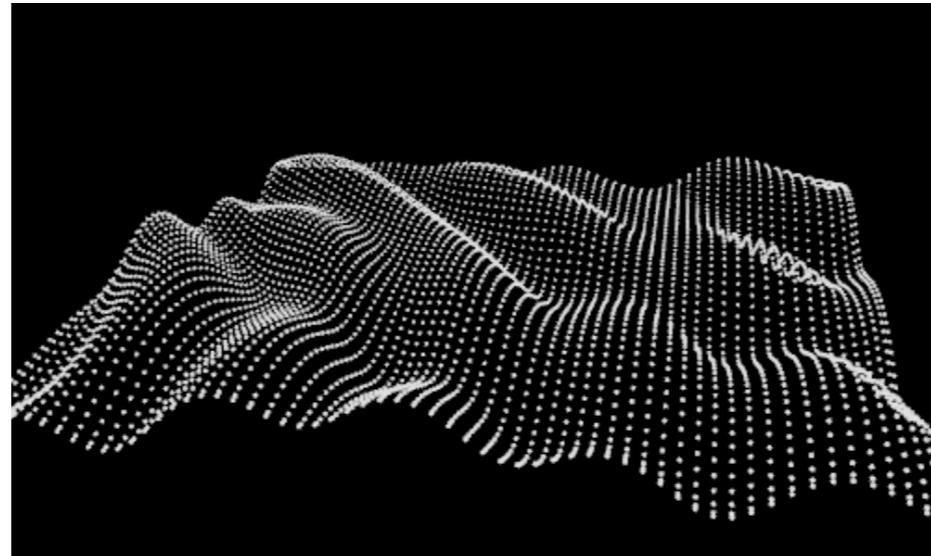


# “Field” Dark Matter

particle DM



DM at long deBroglie wavelength  
useful to picture as a “coherent” field:



signal frequency = DM mass =  $m$

spread by DM kinetic energy  $\sim mv^2$

galactic virial velocity  $v \sim 10^{-3}$   $\rightarrow$  line width  $\sim 10^{-6}m$

$\rightarrow$  coherence time,  $Q \sim 10^6$  periods

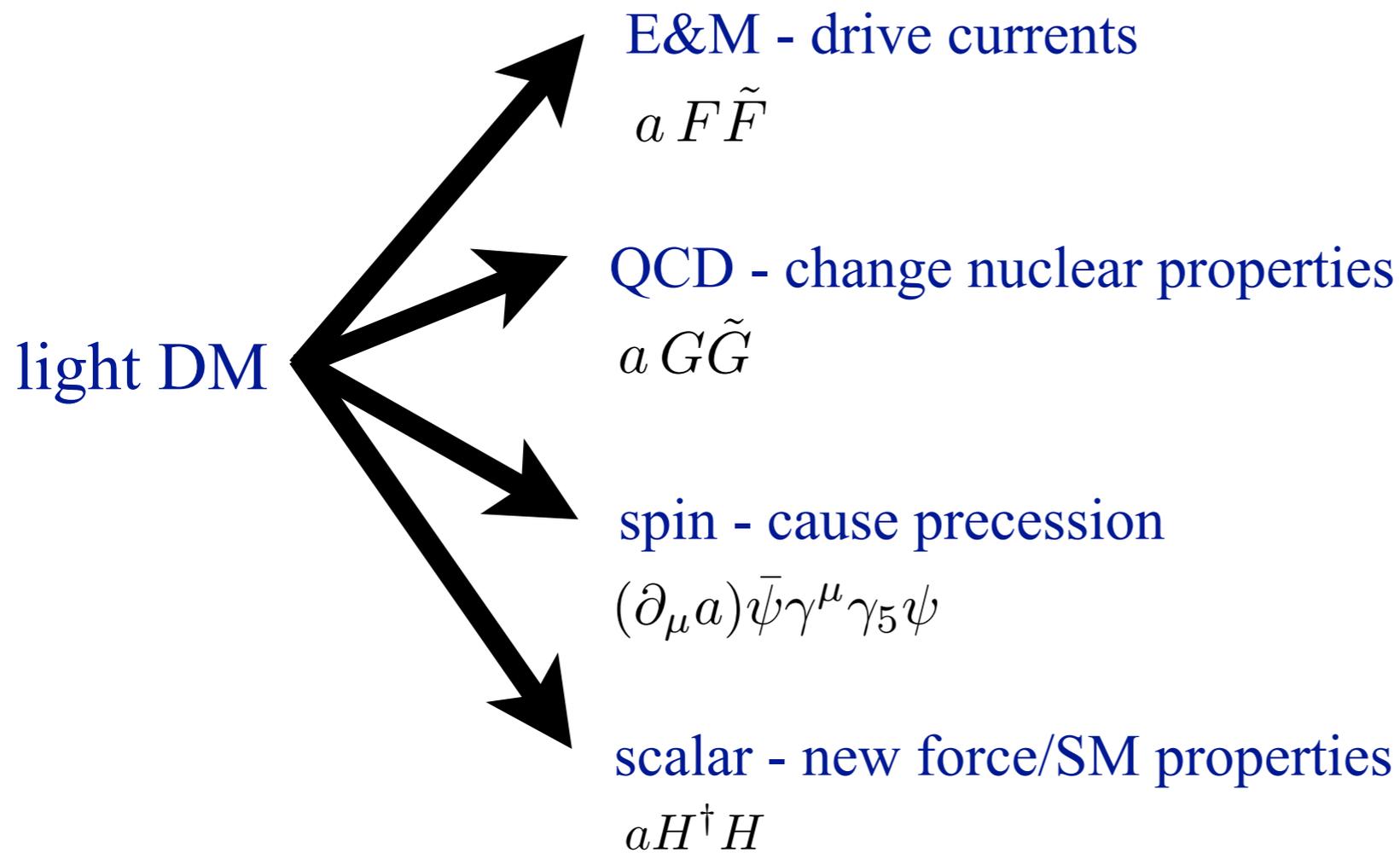
# Possibilities for Light Dark Matter

Effective field theory → only a few possible couplings to us  
four types of experiments:

light DM

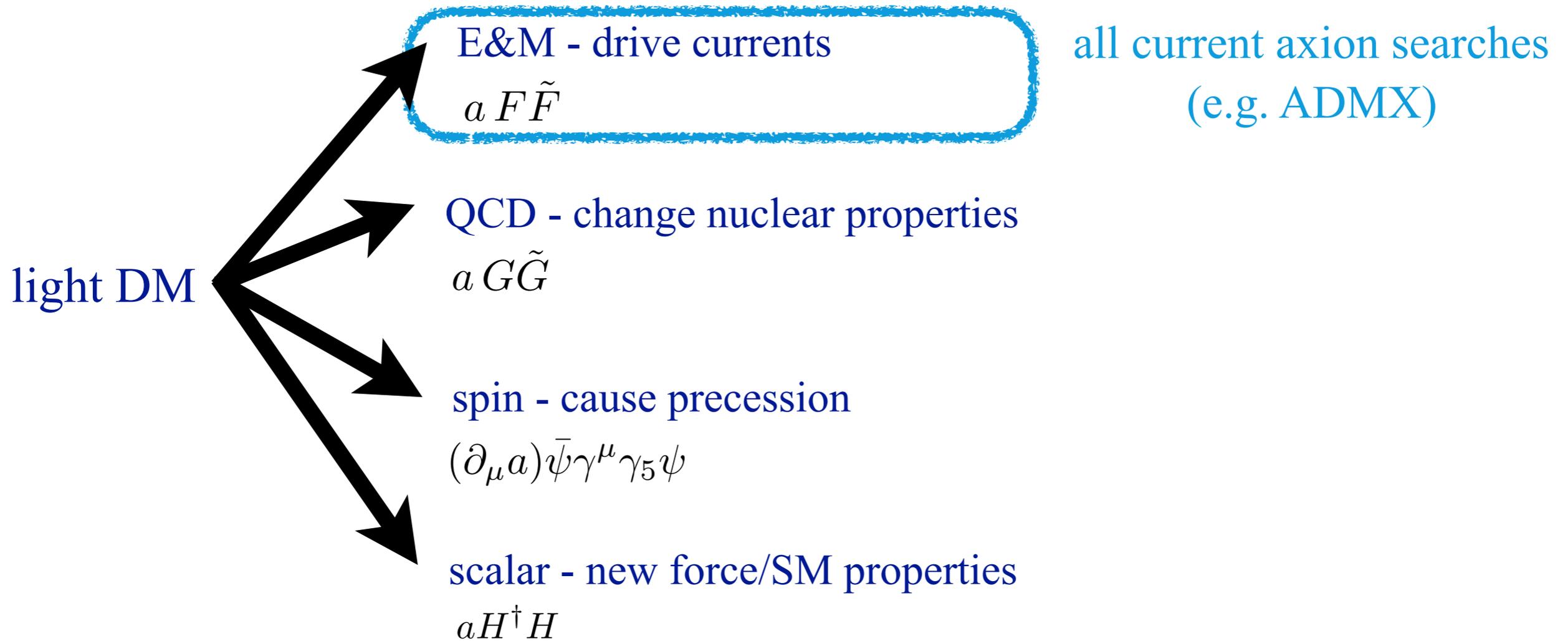
# Possibilities for Light Dark Matter

Effective field theory → only a few possible couplings to us  
four types of experiments:



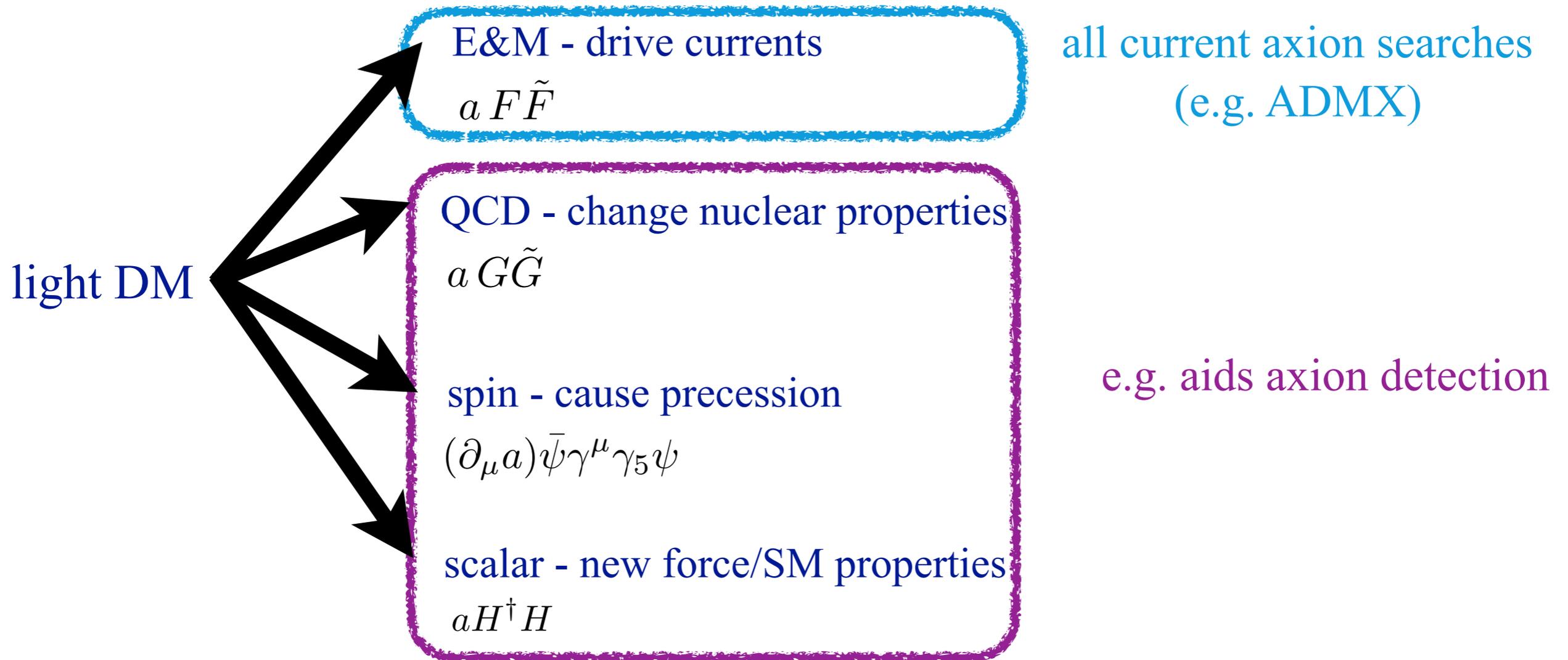
# Possibilities for Light Dark Matter

Effective field theory → only a few possible couplings to us  
four types of experiments:



# Possibilities for Light Dark Matter

Effective field theory → only a few possible couplings to us  
four types of experiments:



Can cover all these possibilities

# Axion Detection

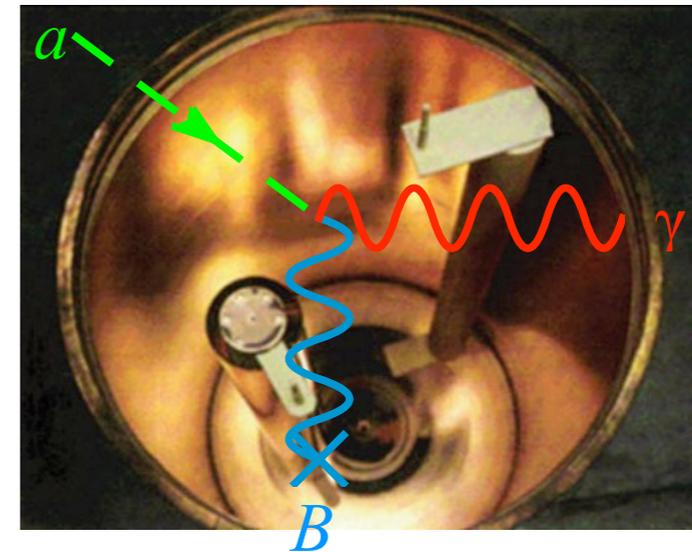
# Existing Axion Searches

all existing experiments rely on axion coupling to E&M (photons):  $\mathcal{L} \supset a F \tilde{F} = a \vec{E} \cdot \vec{B}$

drives cavity at frequency  $m_a$

ADMX focuses on axions  $\sim 0.5 - 10$  GHz

axion Compton wavelength  $\sim$  size of cavity



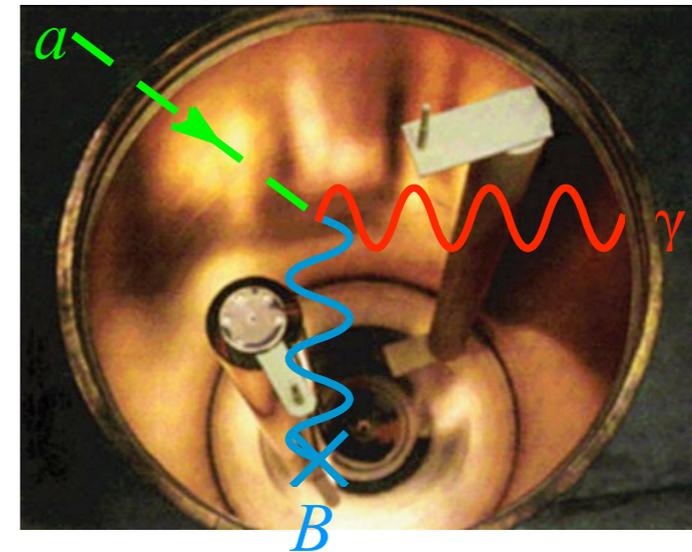
# Existing Axion Searches

all existing experiments rely on axion coupling to E&M (photons):  $\mathcal{L} \supset a F \tilde{F} = a \vec{E} \cdot \vec{B}$

drives cavity at frequency  $m_a$

ADMX focuses on axions  $\sim 0.5 - 10$  GHz

axion Compton wavelength  $\sim$  size of cavity



$a F \tilde{F} \sim a \partial (A \partial A)$  is a derivative operator

integrate by parts  $\rightarrow$  all effects depend on derivative of axion field

all effects suppressed by  $\sim \frac{\text{experiment size}}{\text{axion wavelength}}$

at lower masses, axion wavelength  $\rightarrow 300$  km

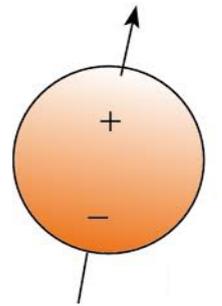
(axions from fundamental scales near Planck scale)

how cover the full axion mass range? a different operator

# The Axion

Strong CP problem:

$\mathcal{L} \supset \theta G\tilde{G}$  creates nucleon EDM  $d \sim 3 \times 10^{-16} \theta e \text{ cm}$  measurements  $\rightarrow \theta \lesssim 10^{-9}$



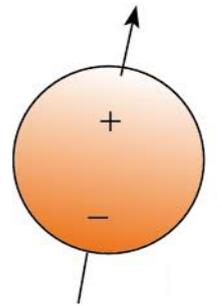
# The Axion

Strong CP problem:

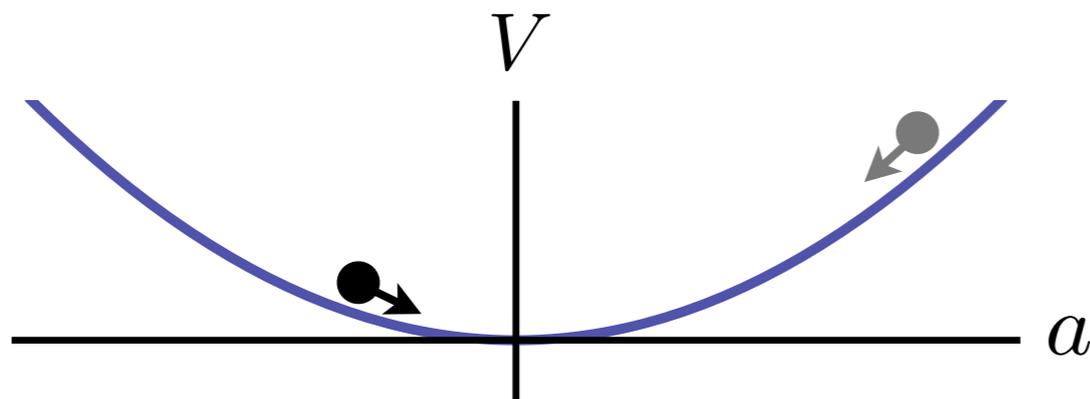
$\mathcal{L} \supset \theta G\tilde{G}$  creates nucleon EDM  $d \sim 3 \times 10^{-16} \theta e \text{ cm}$  measurements  $\rightarrow \theta \lesssim 10^{-9}$

Axion solution:

make it dynamical  $\mathcal{L} \supset \frac{a}{f_a} G\tilde{G}$  so damps down towards zero



$$a(t) \sim a_0 \cos(m_a t)$$



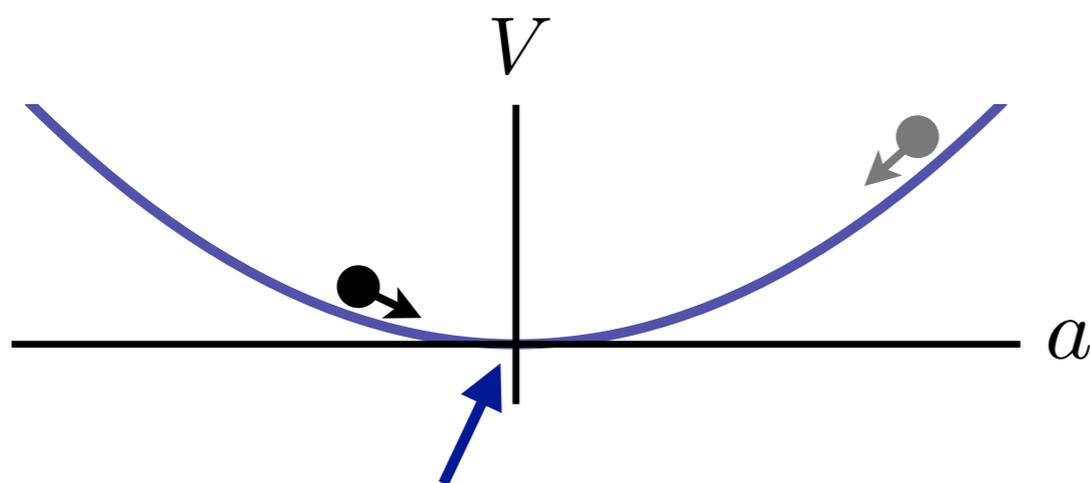
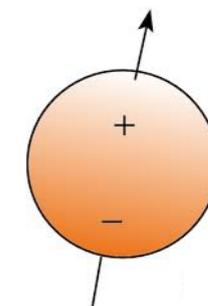
# The Axion

Strong CP problem:

$\mathcal{L} \supset \theta G\tilde{G}$  creates nucleon EDM  $d \sim 3 \times 10^{-16} \theta \text{ e cm}$  measurements  $\rightarrow \theta \lesssim 10^{-9}$

Axion solution:

make it dynamical  $\mathcal{L} \supset \frac{a}{f_a} G\tilde{G}$  so damps down towards zero



$$a(t) \sim a_0 \cos(m_a t)$$

calculate  $a_0$ :

$$m_a^2 a_0^2 \sim \rho_{\text{DM}} \sim 0.3 \frac{\text{GeV}}{\text{cm}^3}$$

still has small residual oscillations today  $\rightarrow$  Axion is a natural dark matter candidate

Preskill, Wise & Wilczek; Abott & Sikivie; Dine & Fischler (1983)

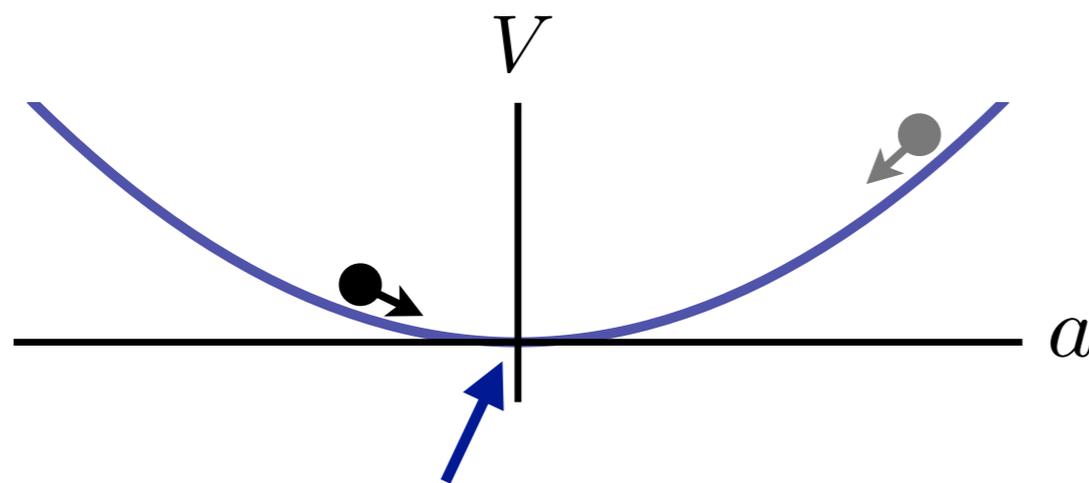
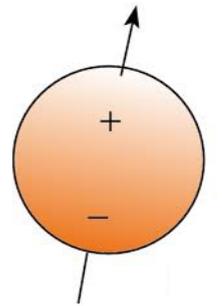
# The Axion

Strong CP problem:

$\mathcal{L} \supset \theta G\tilde{G}$  creates nucleon EDM  $d \sim 3 \times 10^{-16} \theta \text{ e cm}$     measurements  $\rightarrow \theta \lesssim 10^{-9}$

Axion solution:

make it dynamical  $\mathcal{L} \supset \frac{a}{f_a} G\tilde{G}$  so damps down towards zero



$$a(t) \sim a_0 \cos(m_a t)$$

calculate  $a_0$ :

$$m_a^2 a_0^2 \sim \rho_{\text{DM}} \sim 0.3 \frac{\text{GeV}}{\text{cm}^3}$$

still has small residual oscillations today  $\rightarrow$  Axion is a natural dark matter candidate

Preskill, Wise & Wilczek; Abbott & Sikivie; Dine & Fischler (1983)

adiabatic approximation  $\rightarrow d \sim 3 \times 10^{-16} \frac{a}{f_a} \text{ e cm}$

Axion DM causes oscillating nucleon EDM today, not a derivative effect!

completely changes axion detection

generally light bosonic DM causes oscillating fundamental “constants”

# Cosmic Axion Spin Precession Experiment (CASPEr)

with

Dmitry Budker  
Micah Ledbetter  
Surjeet Rajendran  
Alex Sushkov



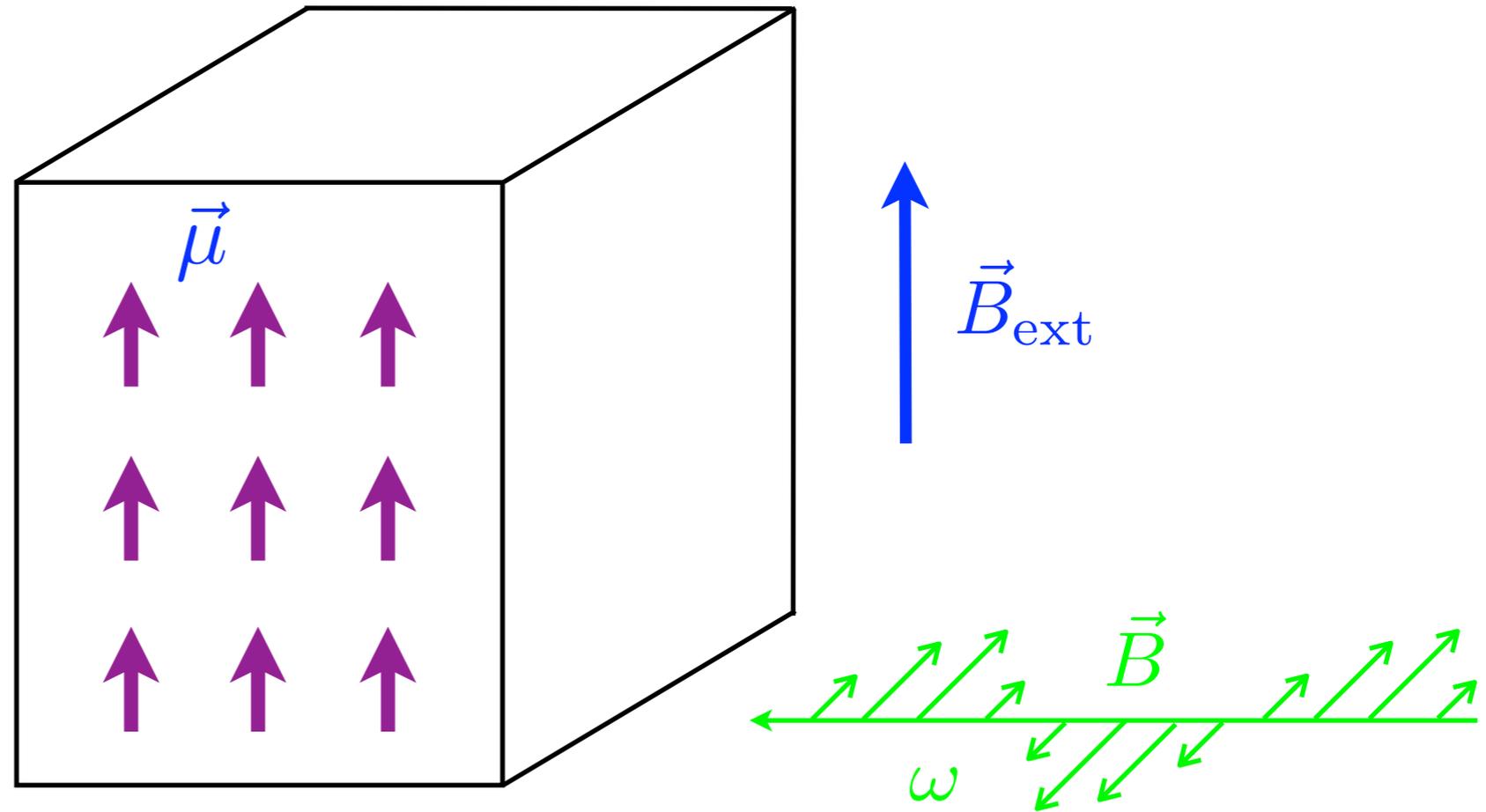
HEISING - SIMONS  
FOUNDATION

SIMONS FOUNDATION

**DFG** Deutsche  
Forschungsgemeinschaft

PRX **4** (2014) arXiv:1306.6089  
PRD **88** (2013) arXiv:1306.6088  
PRD **84** (2011) arXiv:1101.2691

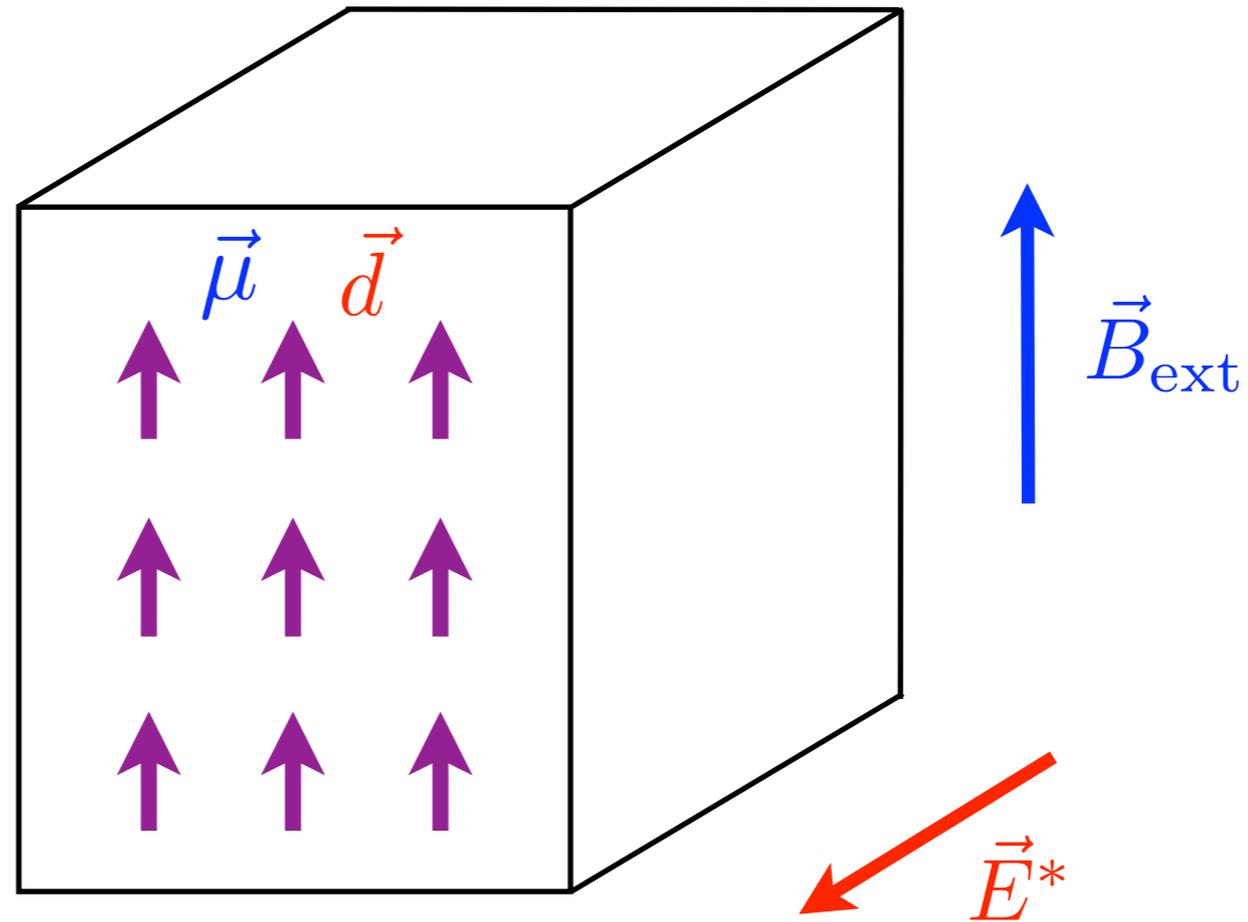
# Axions with NMR



NMR resonant spin flip when Larmor frequency  $2\mu B_{\text{ext}} = \omega$

# Cosmic Axion Spin Precession Experiment (CASPER)

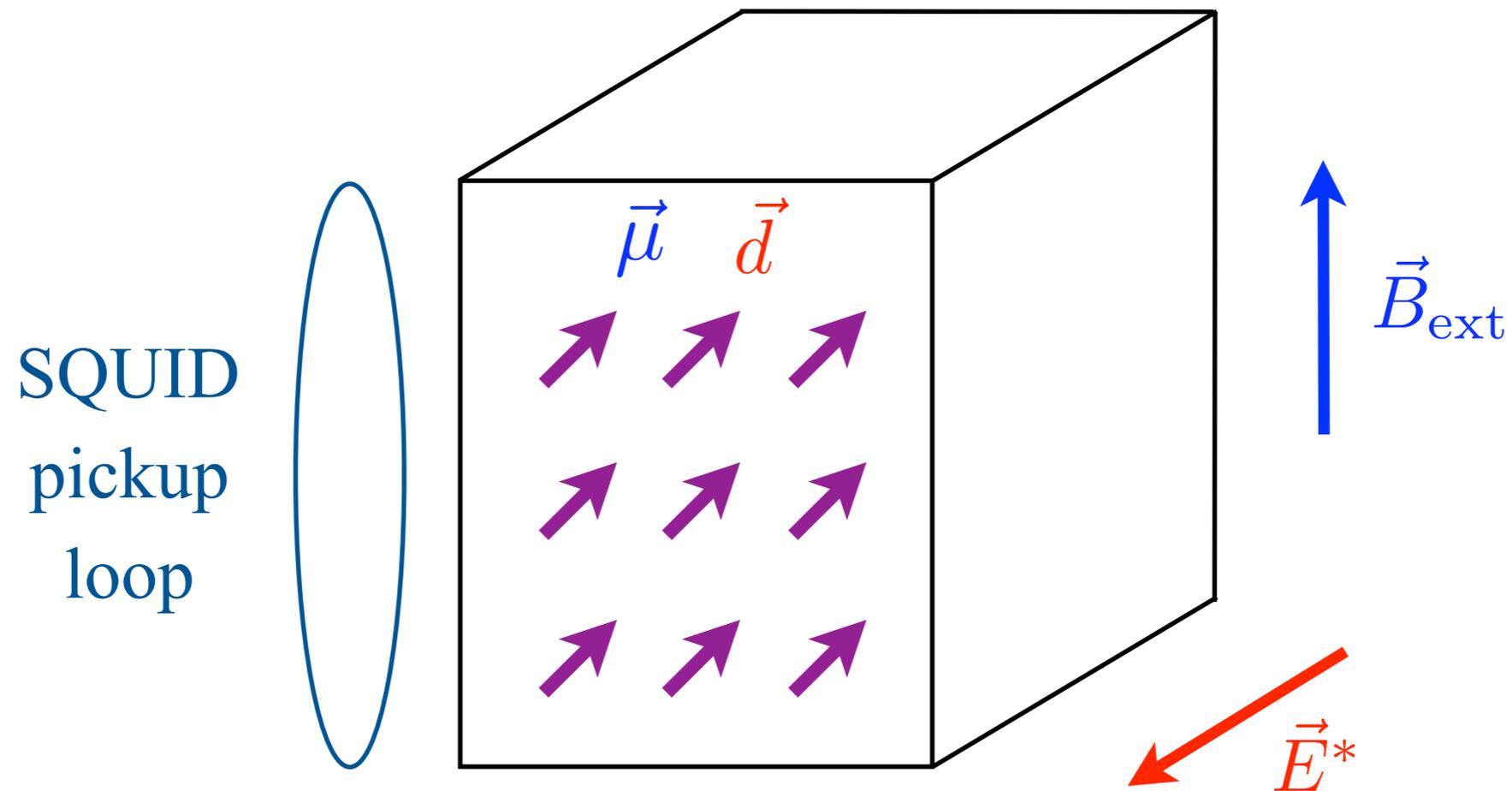
NMR techniques + high precision magnetometry



Larmor frequency = axion mass  $\rightarrow$  resonant enhancement

# Cosmic Axion Spin Precession Experiment (CASPEr)

NMR techniques + high precision magnetometry

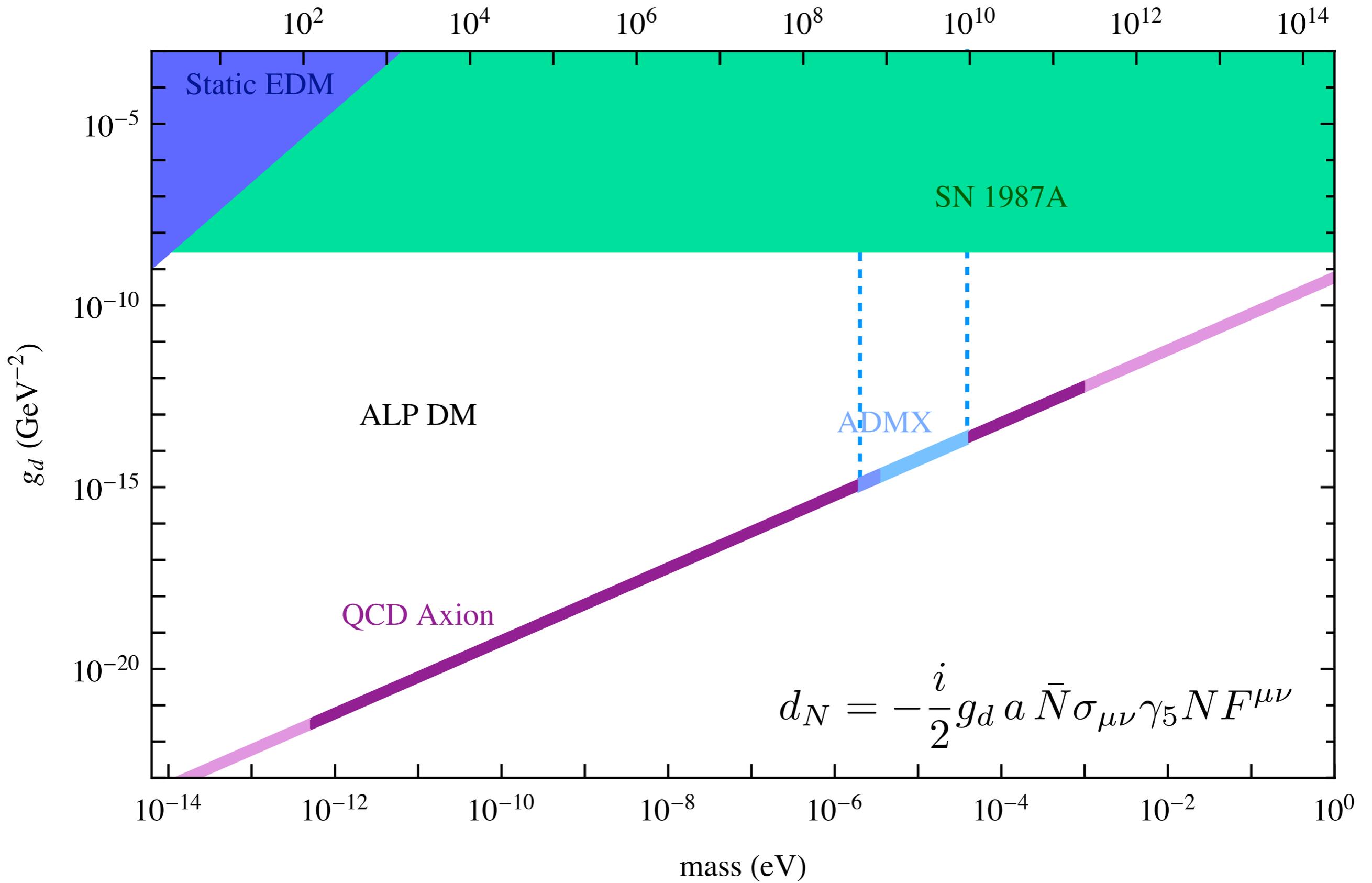


Larmor frequency = axion mass  $\rightarrow$  resonant enhancement

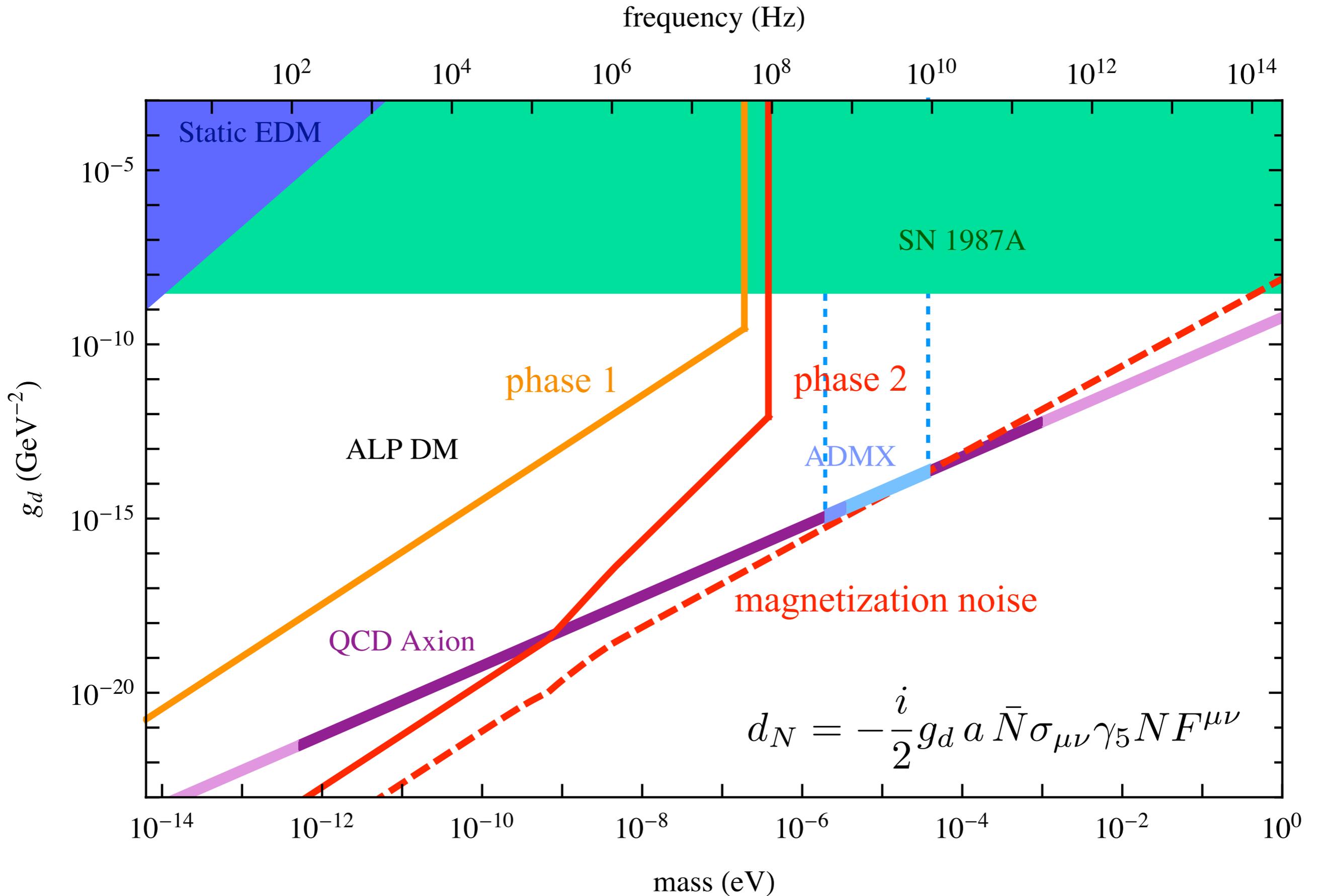
SQUID measures resulting transverse magnetization

ferroelectric (e.g.  $\text{PbTiO}_3$ ), NMR pulse sequences (spin-echo,...),...  
quantum spin projection (magnetization) noise small enough

# Axion Limits on $\frac{a}{f_a} G\tilde{G}$



# CASPEr Sensitivity



# Cosmic Axion Spin Precession Experiment (CASPEr)

New field of axion direct detection, similar to early stages of WIMP direct detection

No other way to search for light axions

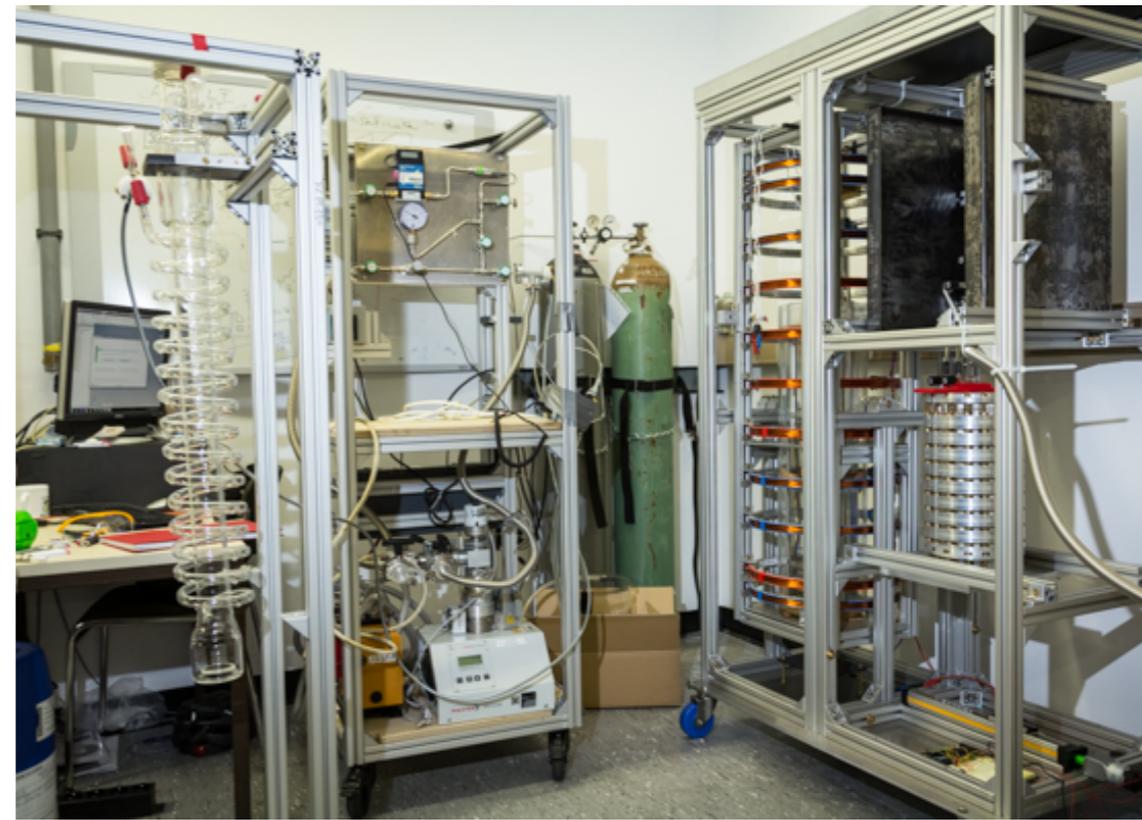
Would be the discovery of dark matter and glimpse into physics at high energies  $\sim 10^{16} - 10^{19}$  GeV



Dmitry Budker  
Alexander Sushkov  
Peter W. Graham  
Surjeet Rajendran  
Derek J. Kimball  
Arne Wickenbrock  
John Blanchard  
Marina Gil Sendra  
Gary Centers  
Nataniel Figueroa  
Deniz Aybas  
Adam Pearson  
Hannah Mekbib  
Tao Wang



under construction at Mainz and BU



SIMONS FOUNDATION



HEISING - SIMONS  
FOUNDATION



Alfred P. Sloan  
FOUNDATION

DFG Deutsche  
Forschungsgemeinschaft

# DM Radio

with

Kent Irwin

Saptarshi Chaudhuri

Jeremy Mardon

Surjeet Rajendran

Yue Zhao

+ collaborating with Tony Tyson + Mani Tripathi's groups (Davis)

The logo for SLAC (Stanford Linear Accelerator Center) in a bold, dark red font.The logo for KIPAC (Kavli IPAC) in a blue font, featuring a stylized orange and red wave above the letter 'I'.

# DM Radio

open axion frequency range below ADMX and above CASPEr

use LC circuit S. Thomas, B. Cabrera, P. Sikivie

want to cover wide DM frequency (mass) range, and all possible candidates (scalar and vector)

# DM Radio

open axion frequency range below ADMX and above CASPER

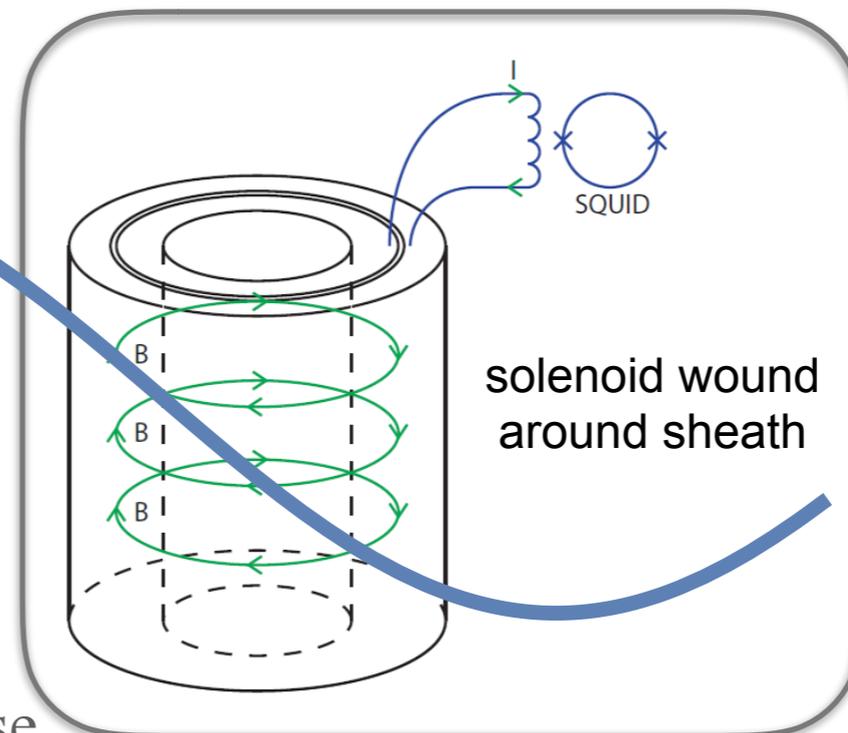
use LC circuit S. Thomas, B. Cabrera, P. Sikivie

want to cover wide DM frequency (mass) range, and all possible candidates (scalar and vector)

DM Radio:

oscillating  
DM field

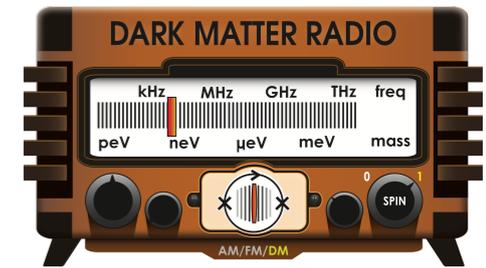
superconducting  
shield for EM noise



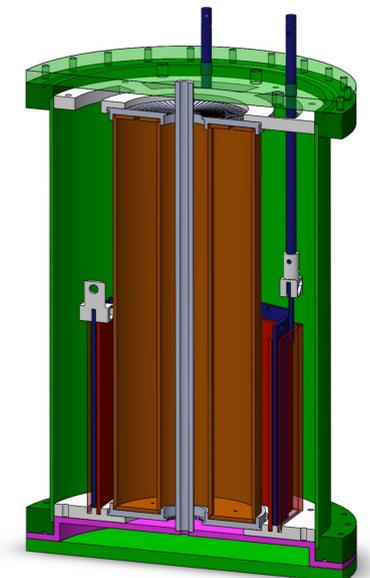
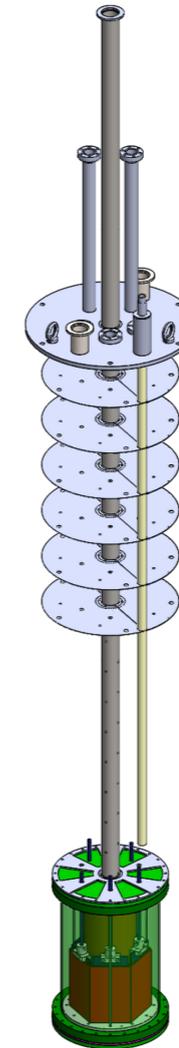
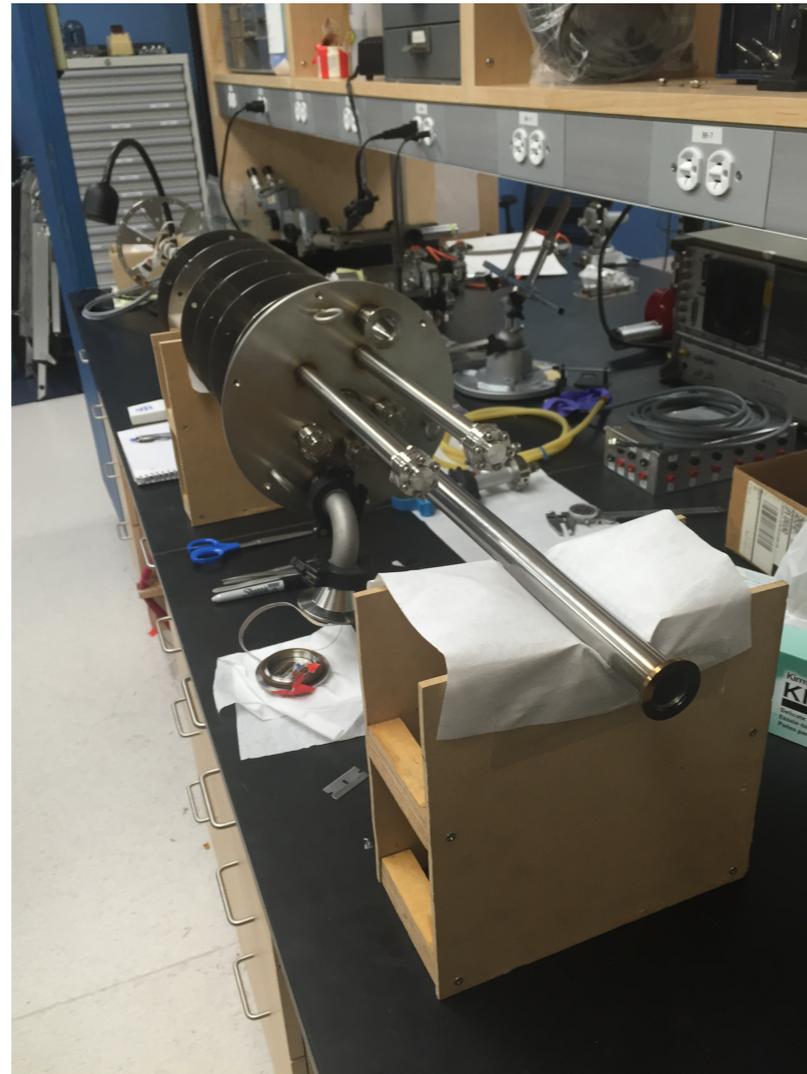
this geometry necessary for DM signal

# DM Radio Experiment

Widely tunable, lumped element resonator  $Q \sim 10^6$



Kent Irwin  
Peter W. Graham  
Surjeet Rajendran  
Jeremy Mardon  
Saptarshi Chaudhuri  
Arran Phipps  
Dale Li  
Sherry Cho  
Betty Young  
Stephen Kuenstner  
Harvey Mosley  
Richard Mule  
Max Silva-Feaver  
Zach Steffen  
Sarah Stokes Kernasovskiy

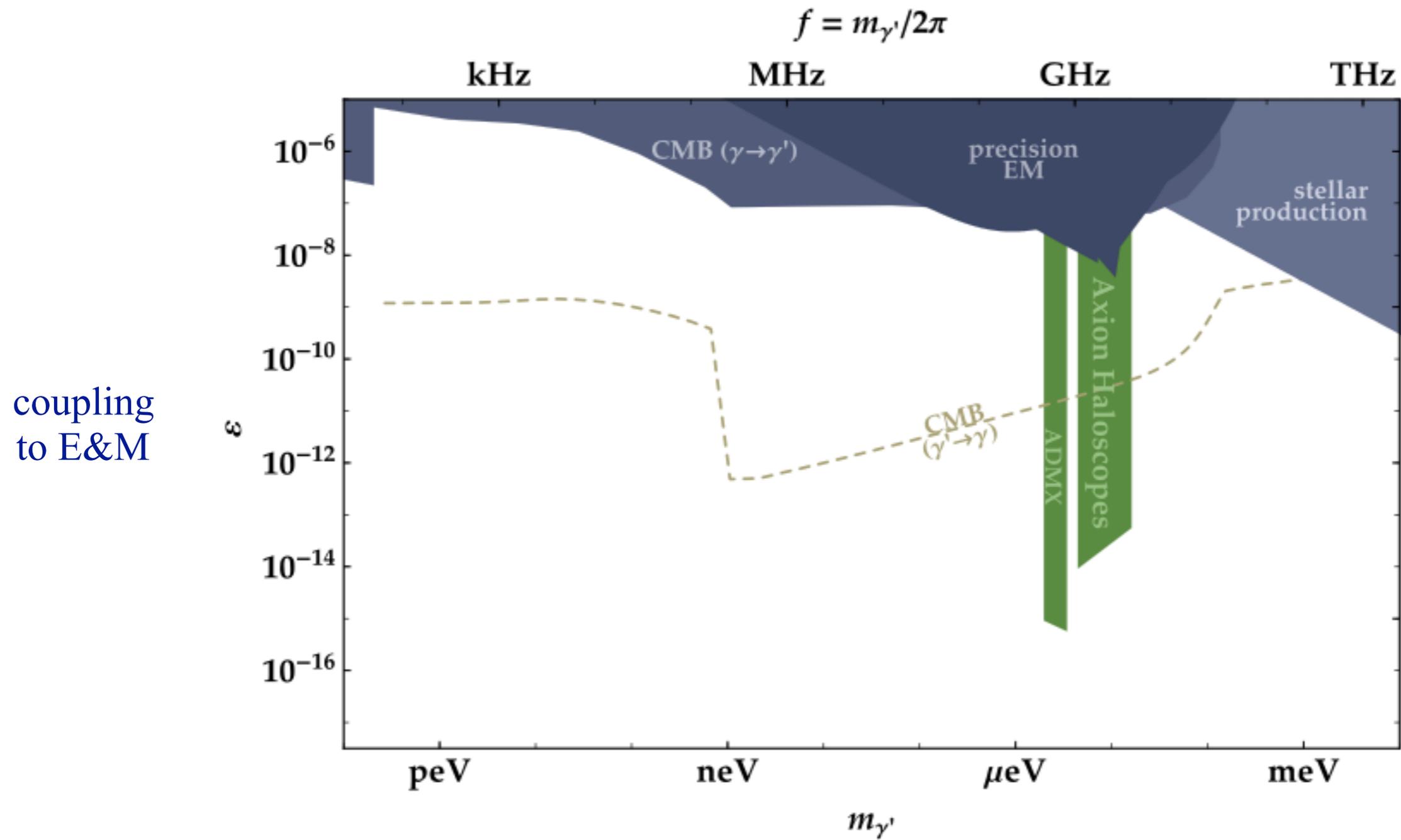


Pathfinder: 4 K 300 cm<sup>3</sup> under construction, initial results ~ 2017

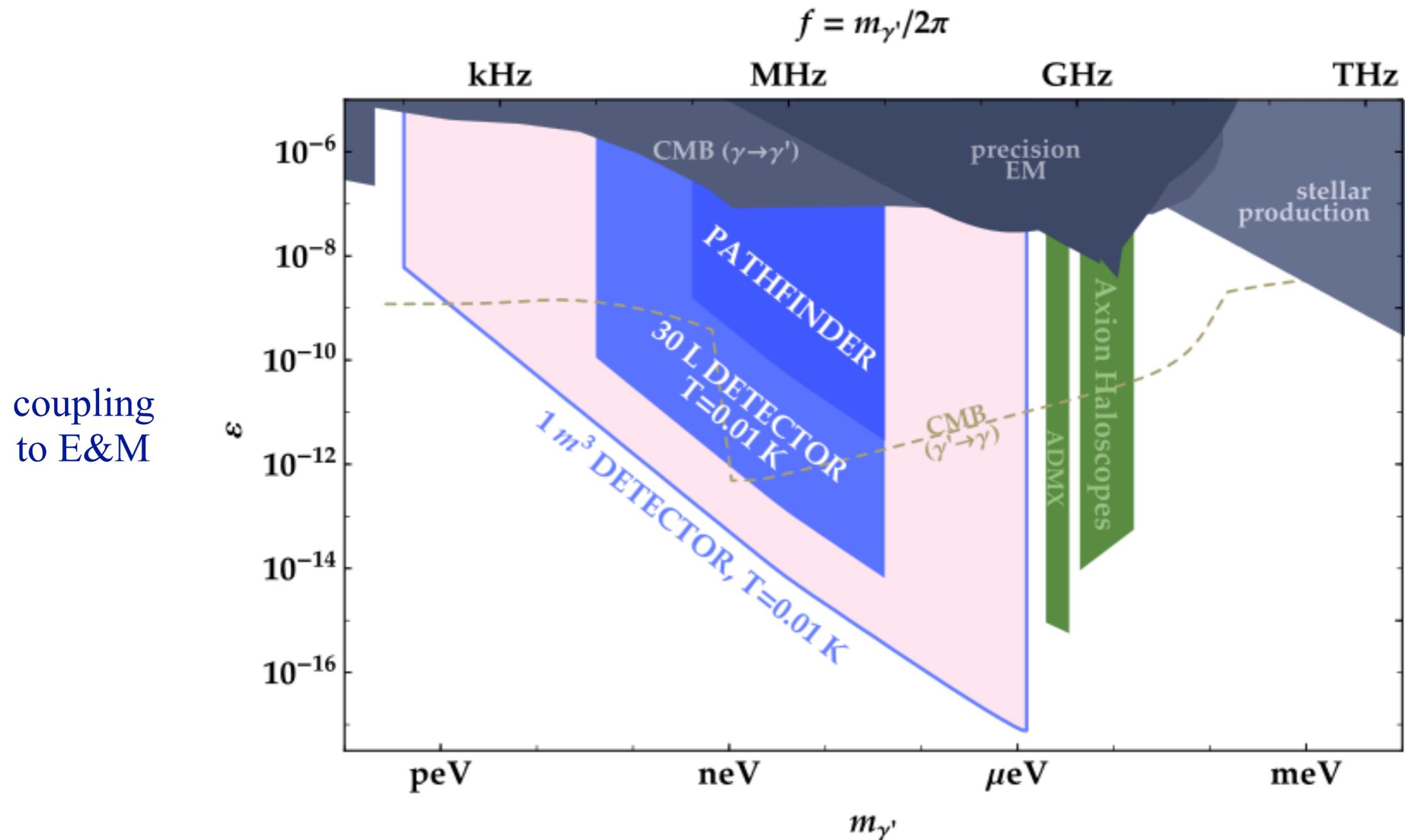
start with hidden photon (vector) detection, later add B field for axion detection

complementary to accelerator searches for heavier hidden photons (e.g. HPS)

# DM Radio Sensitivity to Hidden Photons



# DM Radio Sensitivity to Hidden Photons



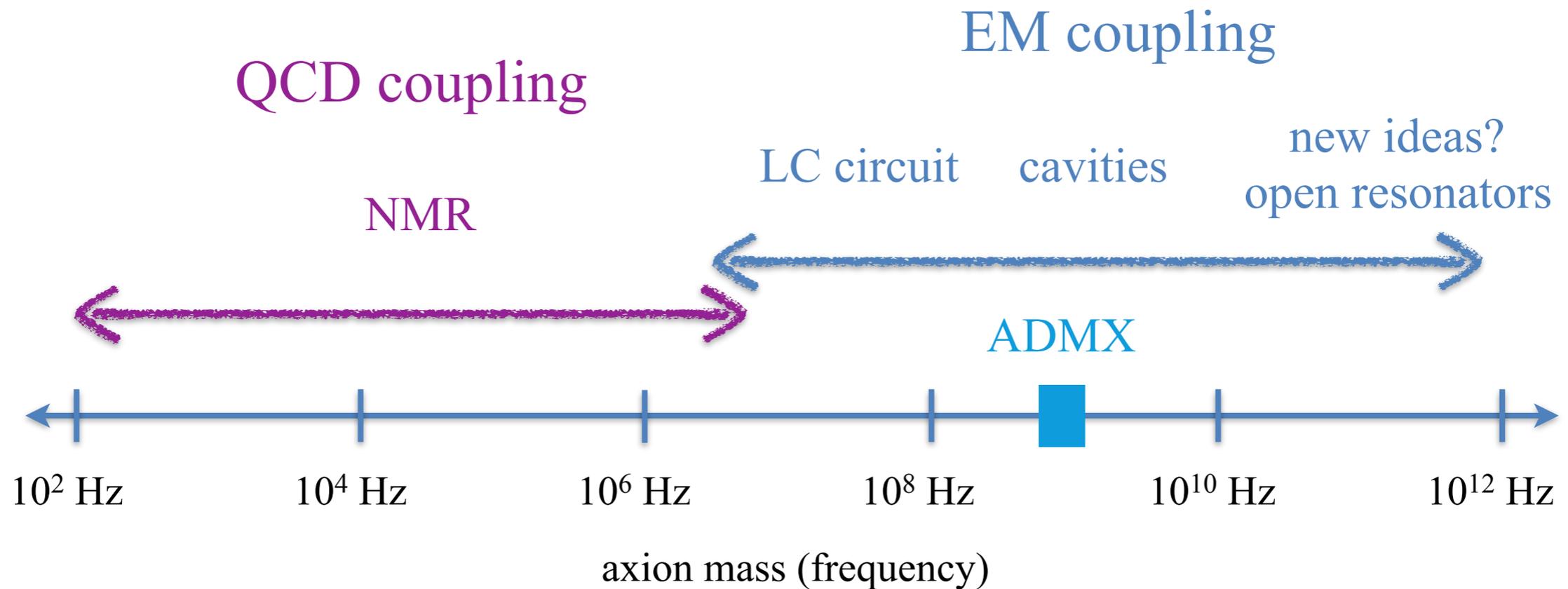
we found hidden photon DM is produced by inflation, and in this frequency range

PWG, Mardon, Rajendran PRD **93** (2016)

a discovery allows measurement of DM power spectrum:  
 verify quantum fluctuation production  
 and measure scale of inflation

# QCD Axion Dark Matter

May be able to cover all of QCD axion dark matter:

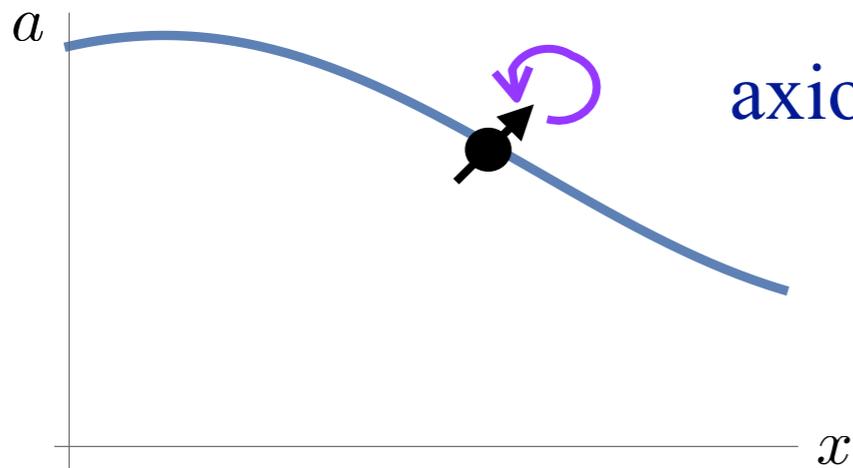


many more new ideas beyond these for axion detection in general!

# Other Couplings & Techniques

# Axion DM Effects

spin coupling:  $(\partial_\mu a)\bar{\psi}\gamma^\mu\gamma_5\psi \rightarrow H \ni \nabla a \cdot \vec{\sigma}_N$



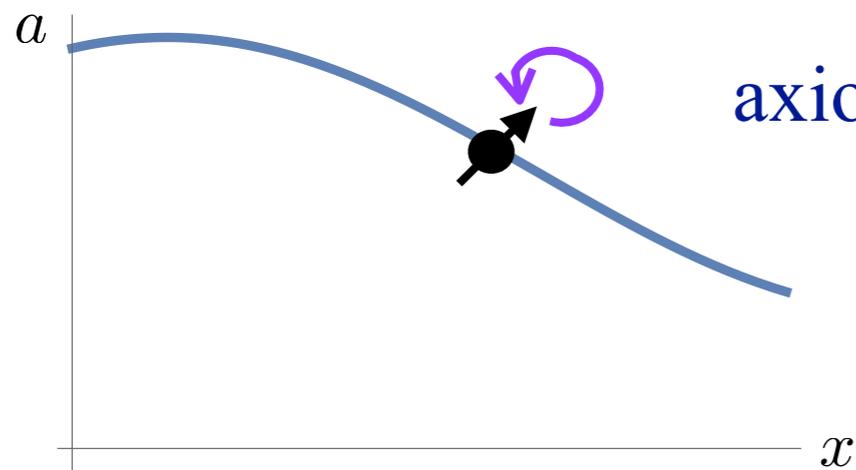
axion DM field gradient torques electron and nucleon spins

oscillates with axion frequency

proportional to axion momentum (“wind”)

# Axion DM Effects

spin coupling:  $(\partial_\mu a)\bar{\psi}\gamma^\mu\gamma_5\psi \rightarrow H \ni \nabla a \cdot \vec{\sigma}_N$

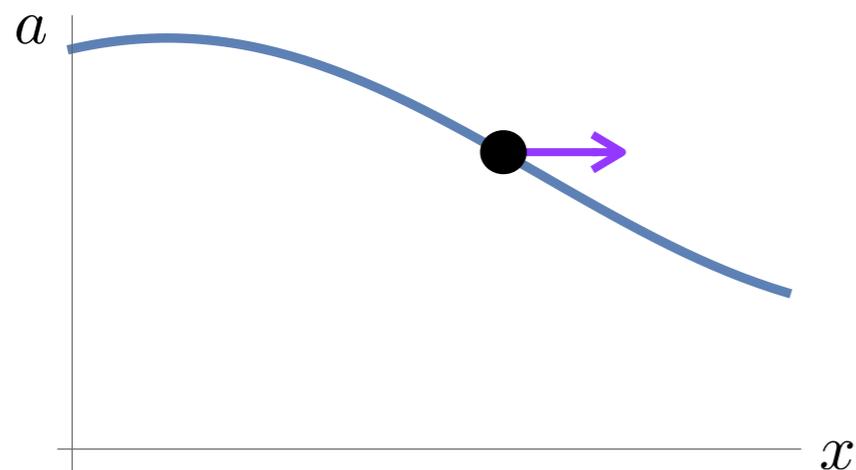


axion DM field gradient torques electron and nucleon spins

oscillates with axion frequency

proportional to axion momentum (“wind”)

scalar coupling:  $\alpha H^\dagger H$  e.g. change electron mass



axion DM field gradient can exert a force

oscillatory and violates equivalence principle

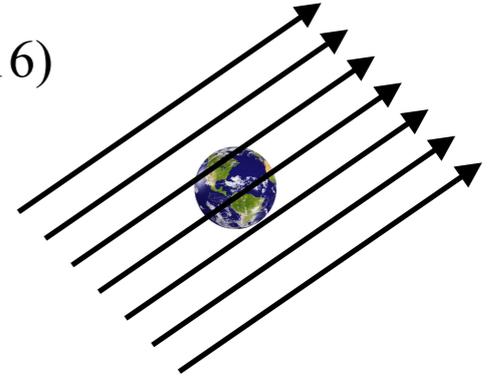
same effects allow searches for hidden photons

# Force/Torque from Dark Matter

with D.E.Kaplan, J.Mardon, S.Rajendran, & W.A.Terrano PRD **93** (2016)

New oscillatory force/torque from dark matter

New Direct Detection Experiments:

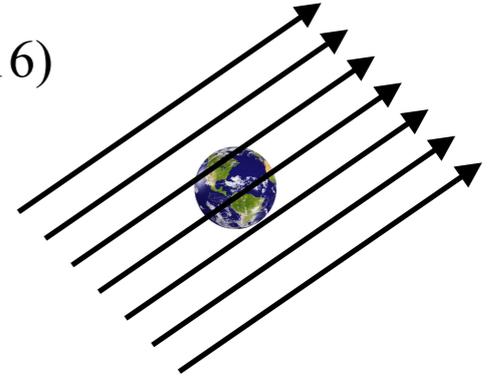


# Force/Torque from Dark Matter

with D.E.Kaplan, J.Mardon, S.Rajendran, & W.A.Terrano PRD **93** (2016)

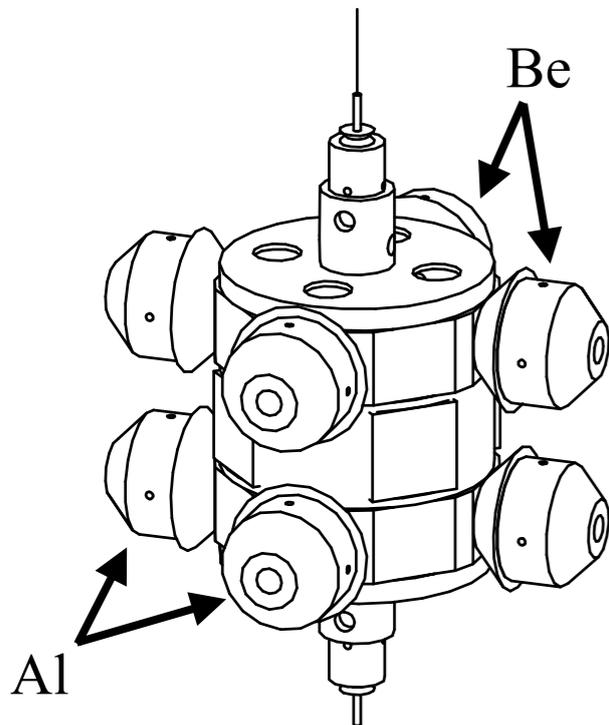
New oscillatory force/torque from dark matter

New Direct Detection Experiments:



## Torsion Balances

scalar balance for force  
spin-polarized for torque



Eot-Wash analysis underway

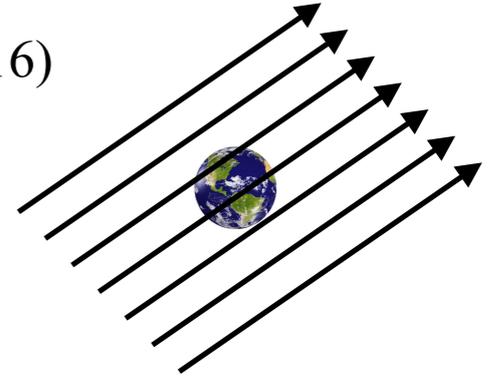
with Will Terrano

# Force/Torque from Dark Matter

with D.E.Kaplan, J.Mardon, S.Rajendran, & W.A.Terrano PRD **93** (2016)

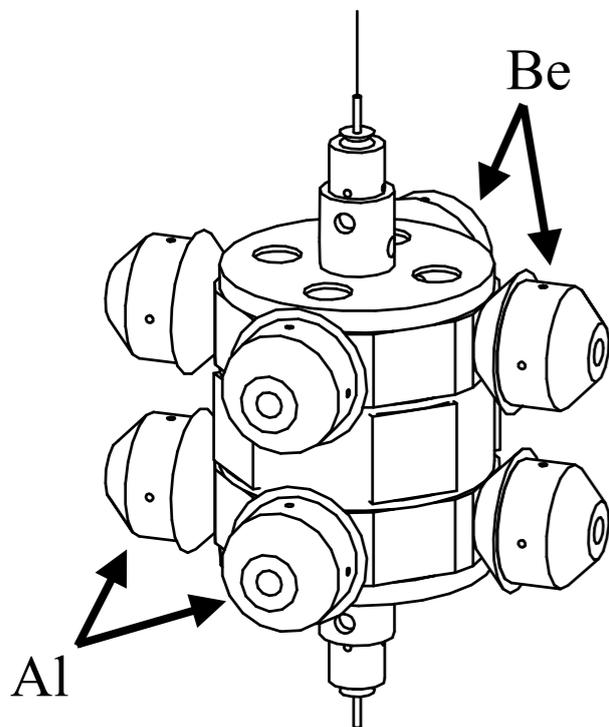
New oscillatory force/torque from dark matter

New Direct Detection Experiments:



## Torsion Balances

scalar balance for force  
spin-polarized for torque

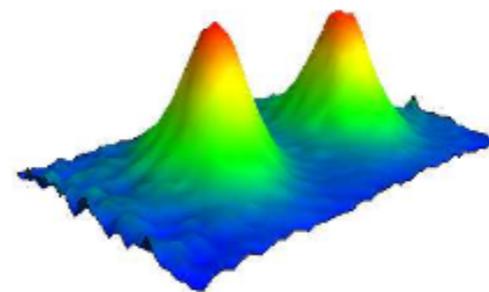
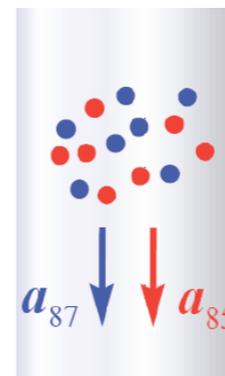


## Atom Interferometers

split + recombine atom wavefunction  
measure atom spin and acceleration



$^{85}\text{Rb}$ - $^{87}\text{Rb}$



Eot-Wash analysis underway

with Will Terrano

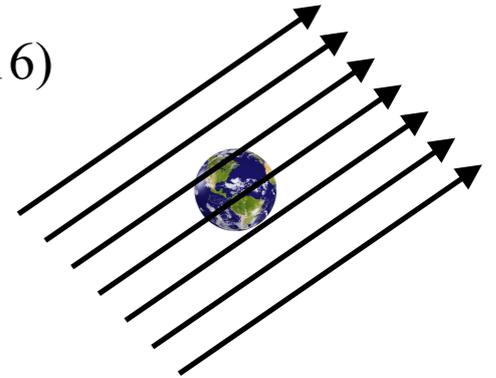
In construction Kasevich/Hogan groups

# Force/Torque from Dark Matter

with D.E.Kaplan, J.Mardon, S.Rajendran, & W.A.Terrano PRD 93 (2016)

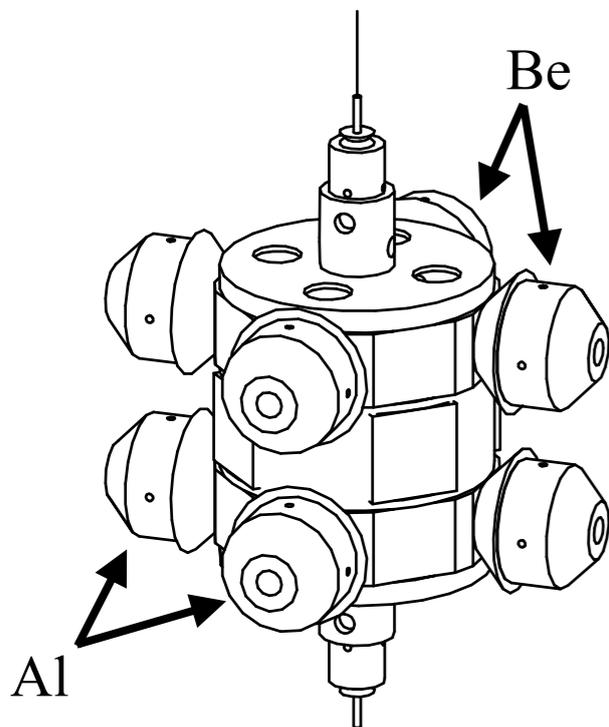
New oscillatory force/torque from dark matter

New Direct Detection Experiments:



## Torsion Balances

scalar balance for force  
spin-polarized for torque



Eot-Wash analysis underway

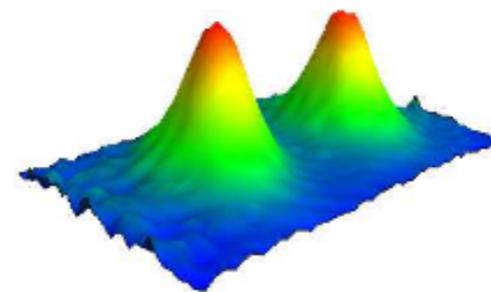
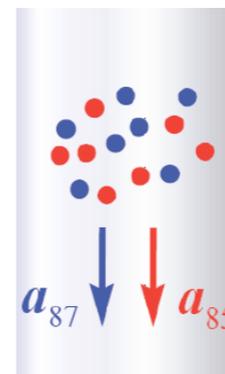
with Will Terrano

## Atom Interferometers

split + recombine atom wavefunction  
measure atom spin and acceleration



$^{85}\text{Rb}$ - $^{87}\text{Rb}$

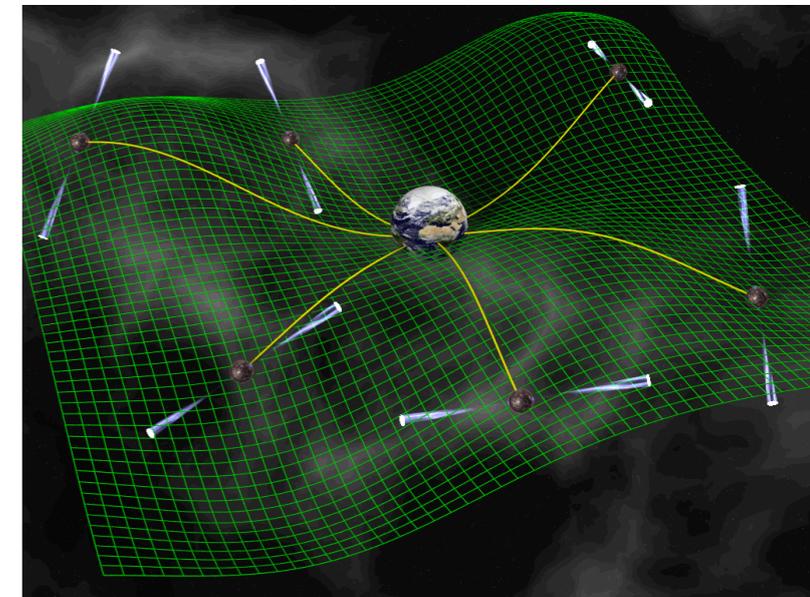


In construction Kasevich/Hogan groups

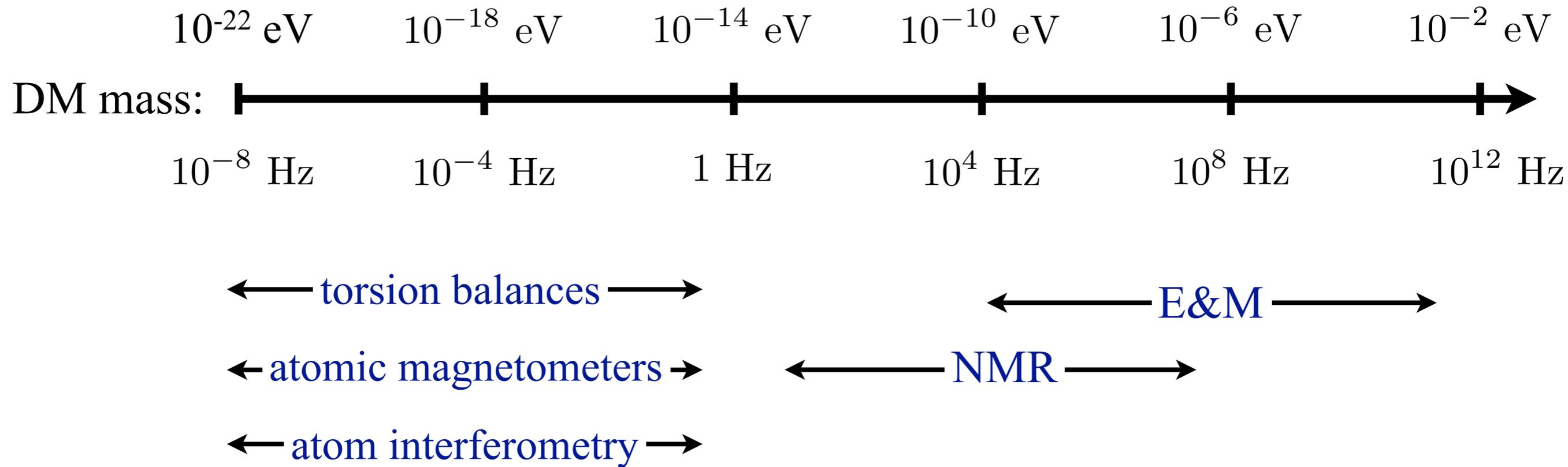
covers frequency range  $\sim 10$  Hz down to  $\text{yr}^{-1}$

## Pulsar Timing Arrays

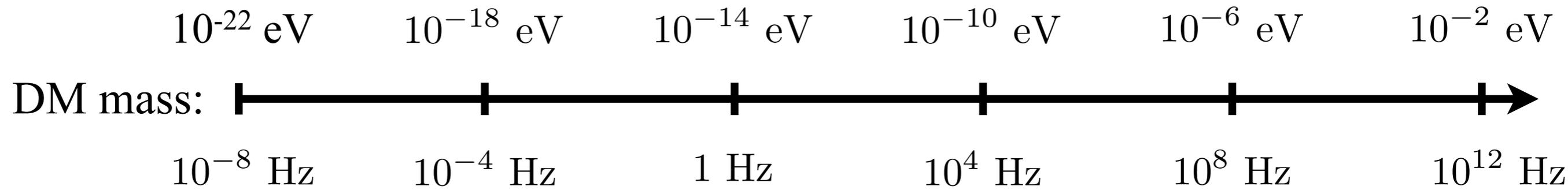
DM and gravitational wave  
detection similar



# DM Direct Detection



# DM Direct Detection



← torsion balances →

← E&M →

← atomic magnetometers →

← NMR →

← atom interferometry →

Coupling:

E&M

Eot-Wash (spin)

CASPER-Electric

ADMX

QCD

Eot-Wash (scalar)

CASPER-Wind

HAYSTAC

Spin

LC Circuit

Scalar

Atom Interferometry (spin)

DM Radio

Atom Interferometry (scalar)

ABRACADABRA

# Gravitational Wave Detection with Atom Interferometry

with

Savas Dimopoulos

Jason Hogan

Mark Kasevich

Surjeet Rajendran

PRD **94** (2016) arXiv:1606.01860

PRL **110** (2013) arXiv:1206.0818

GRG **43** (2011) arXiv:1009.2702

PLB **678** (2009) arXiv:0712.1250

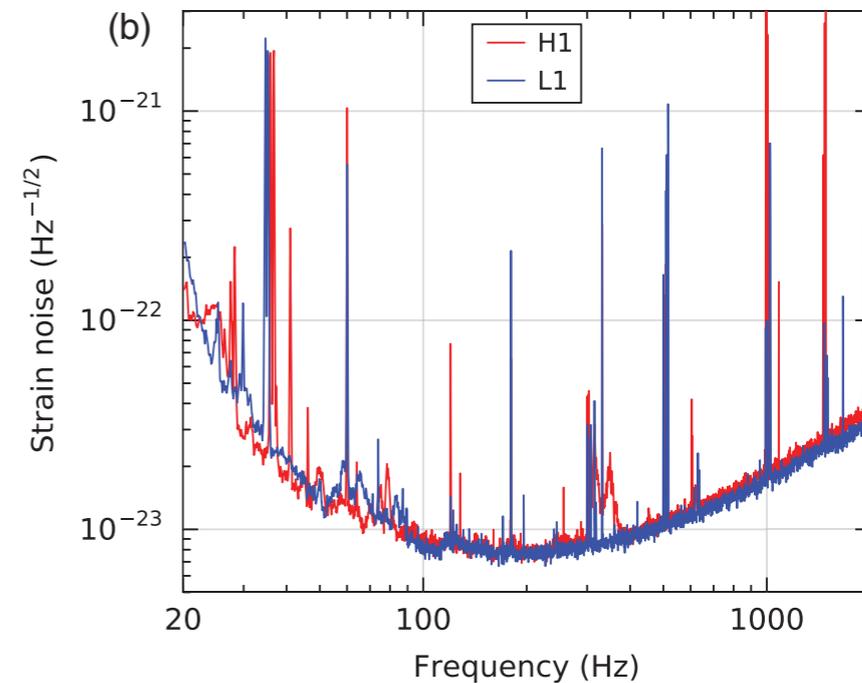
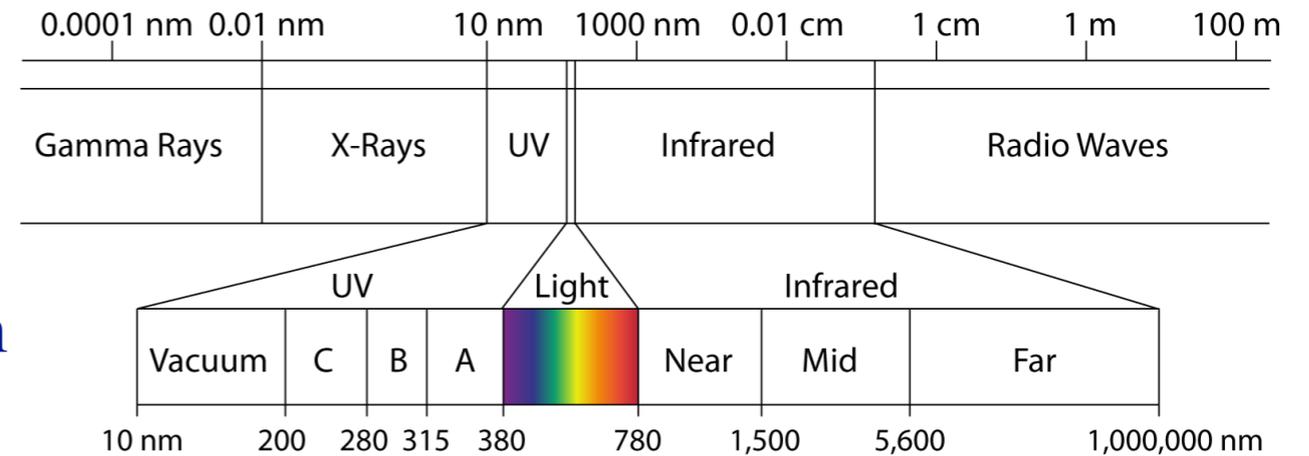
PRD **78** (2008) arXiv:0806.2125



# Gravitational Spectrum

Gravitational waves open a new window to the universe

Every new EM band opened has revealed unexpected discoveries, gravitational waves give a new spectrum



Advanced LIGO can only detect GW's  $> 10$  Hz  $\rightarrow$  How look at lower spectrum?

New detectors?

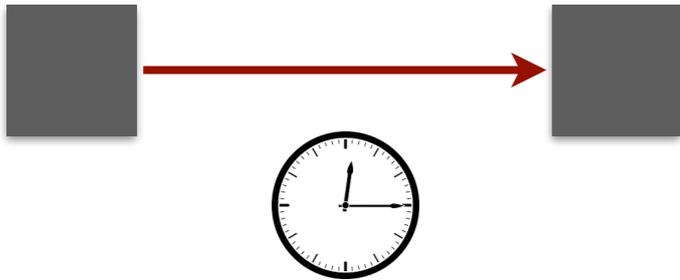
# Gravitational Wave Detection

## Gravitation Wave Detector

inertial test masses

baseline

good clock



# Gravitational Wave Detection

Gravitation Wave Detector

LIGO

inertial test masses

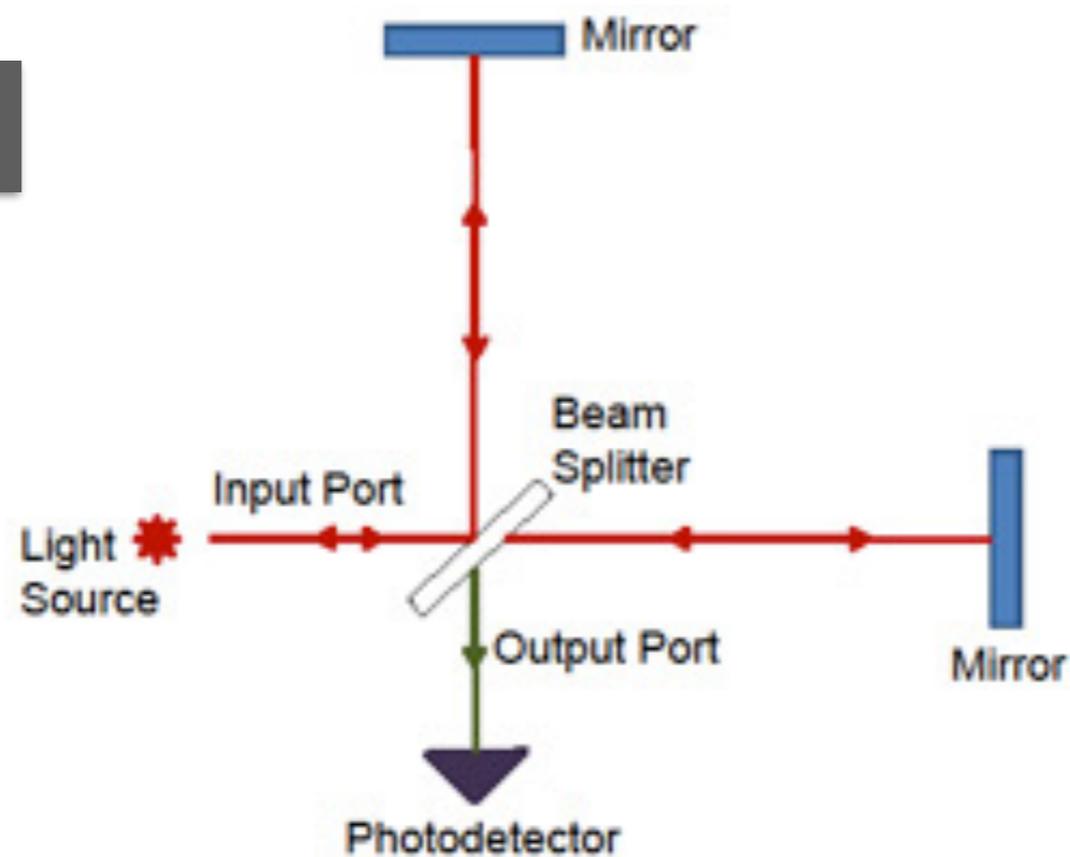
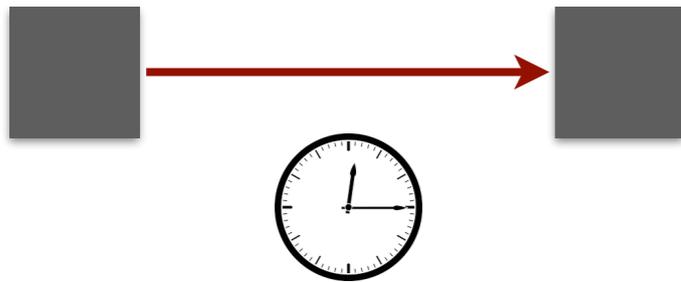
mirrors

baseline

laser

good clock

second arm



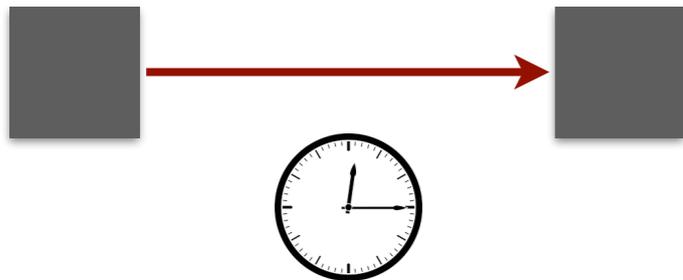
# Gravitational Wave Detection

Gravitation Wave Detector

inertial test masses

baseline

good clock

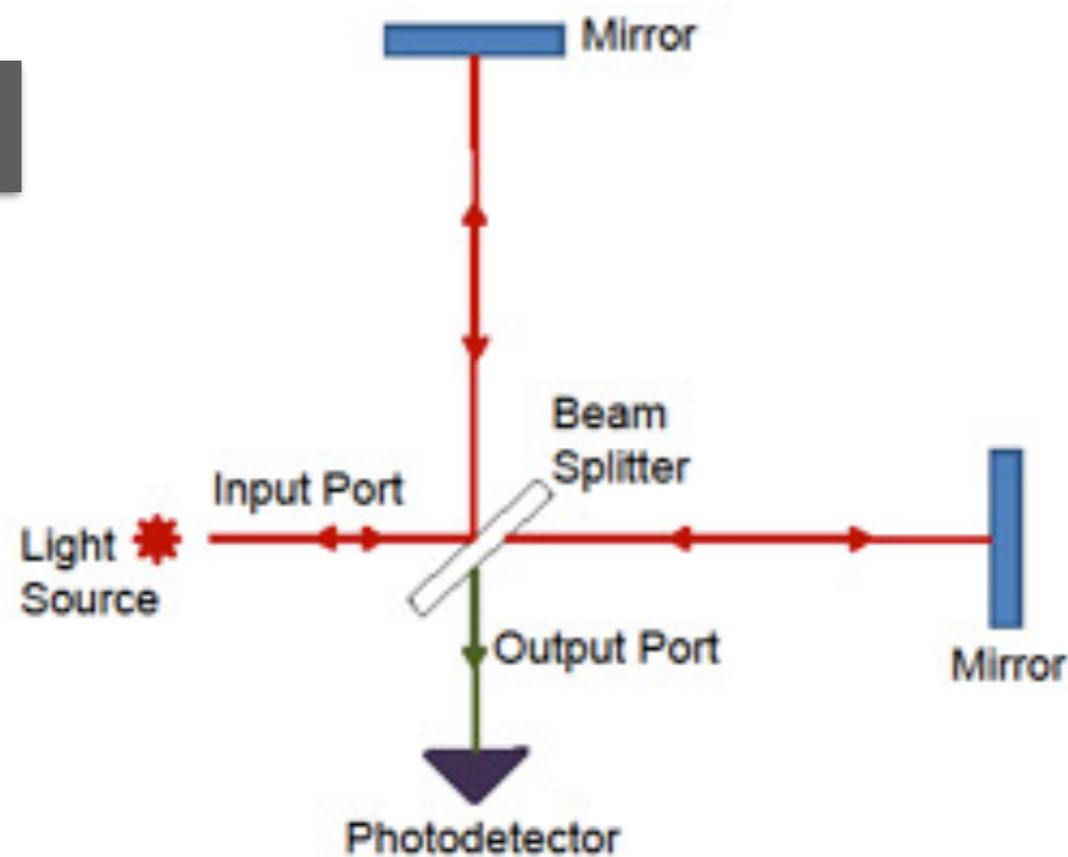


LIGO

mirrors

laser

second arm



Atom Interferometry

atoms

laser

atoms



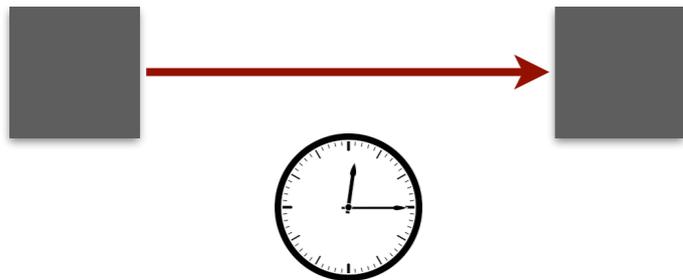
# Gravitational Wave Detection

Gravitation Wave Detector

inertial test masses

baseline

good clock

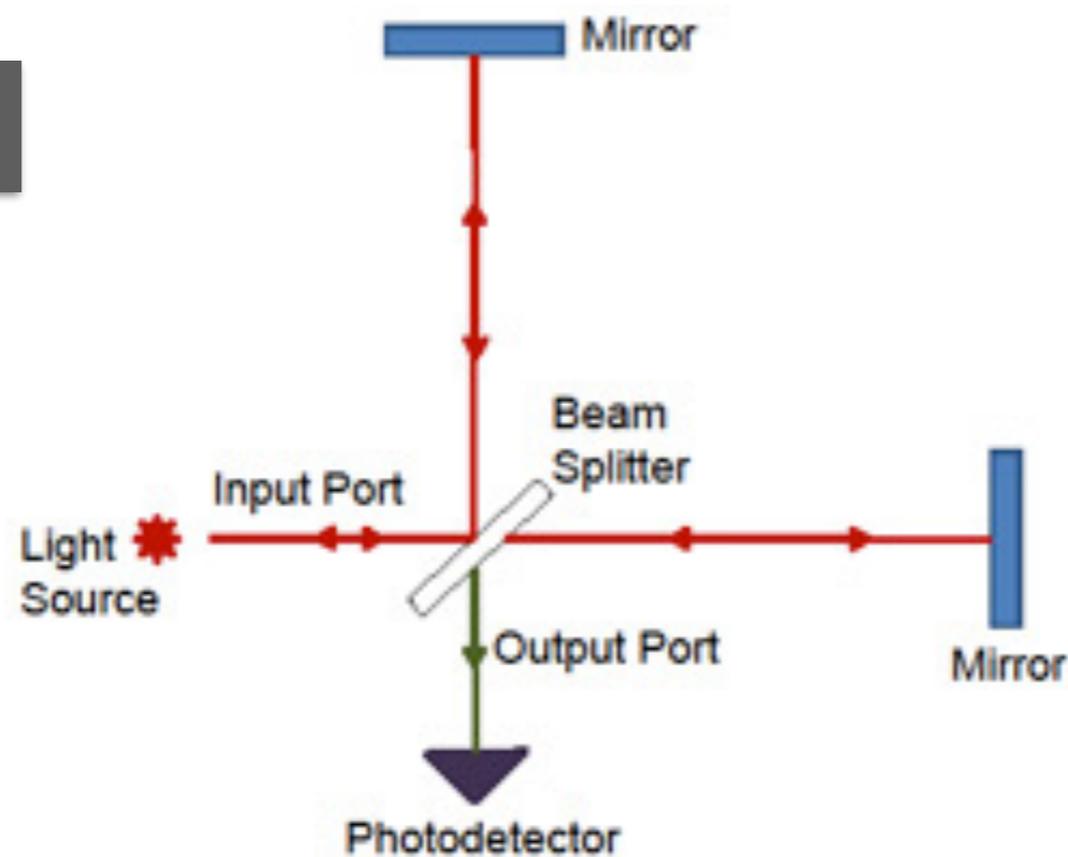


LIGO

mirrors

laser

second arm

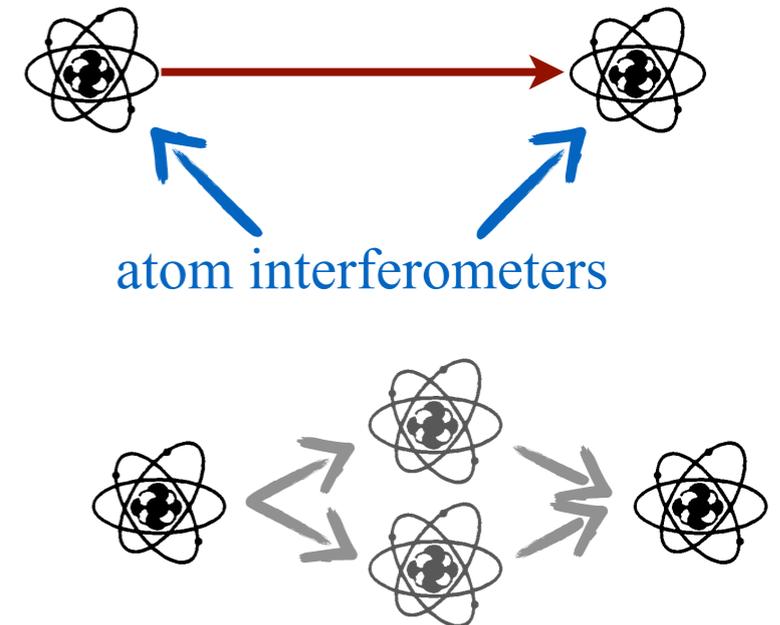


Atom Interferometry

atoms

laser

atoms



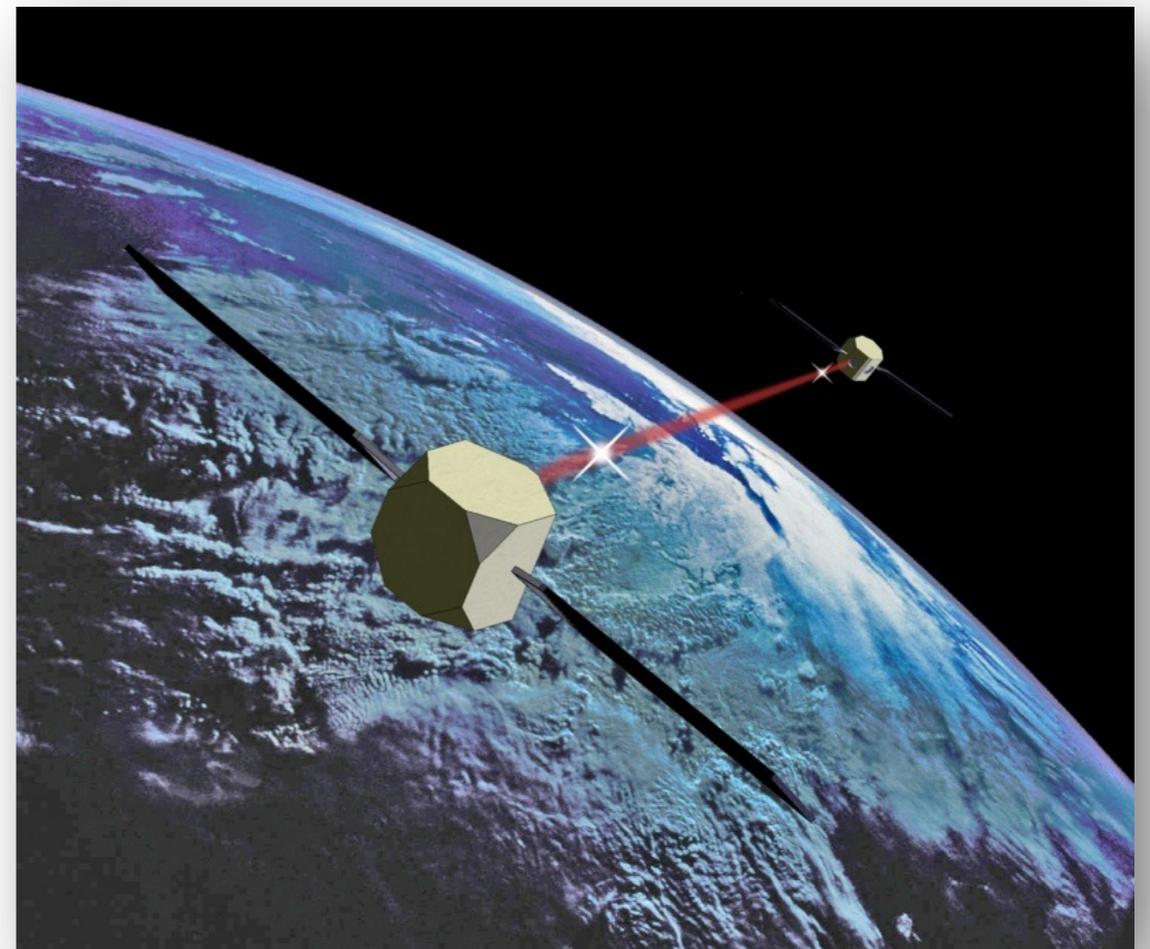
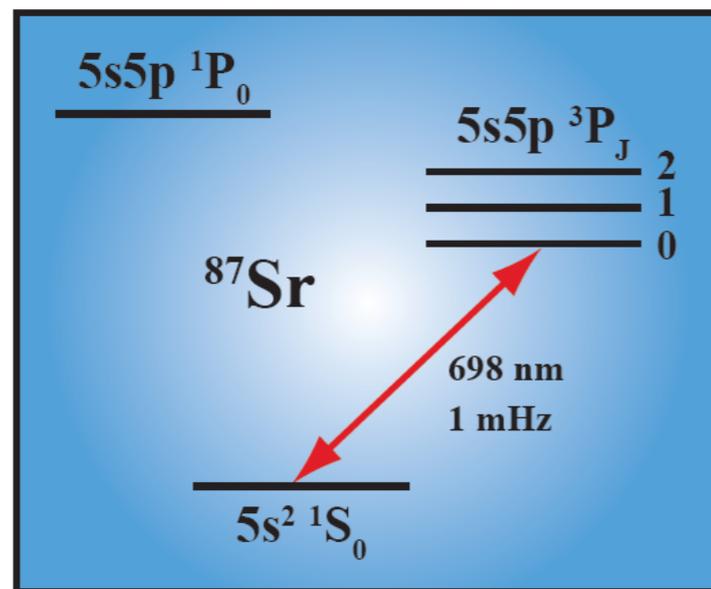
# Atom Interferometry for Gravitational Waves

run atom interferometer as hybrid clock/accelerometer

PWG, Hogan, Kasevich, Rajendran PRL **110** (2013)

atoms act as clocks, measure light travel time

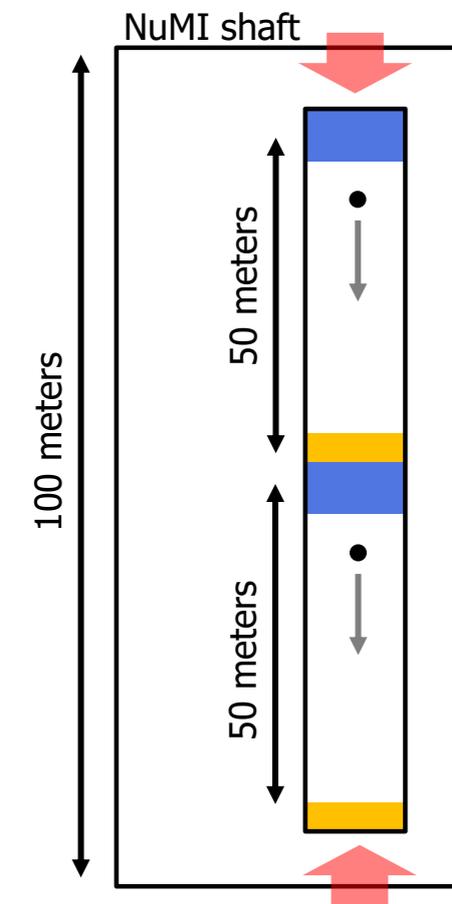
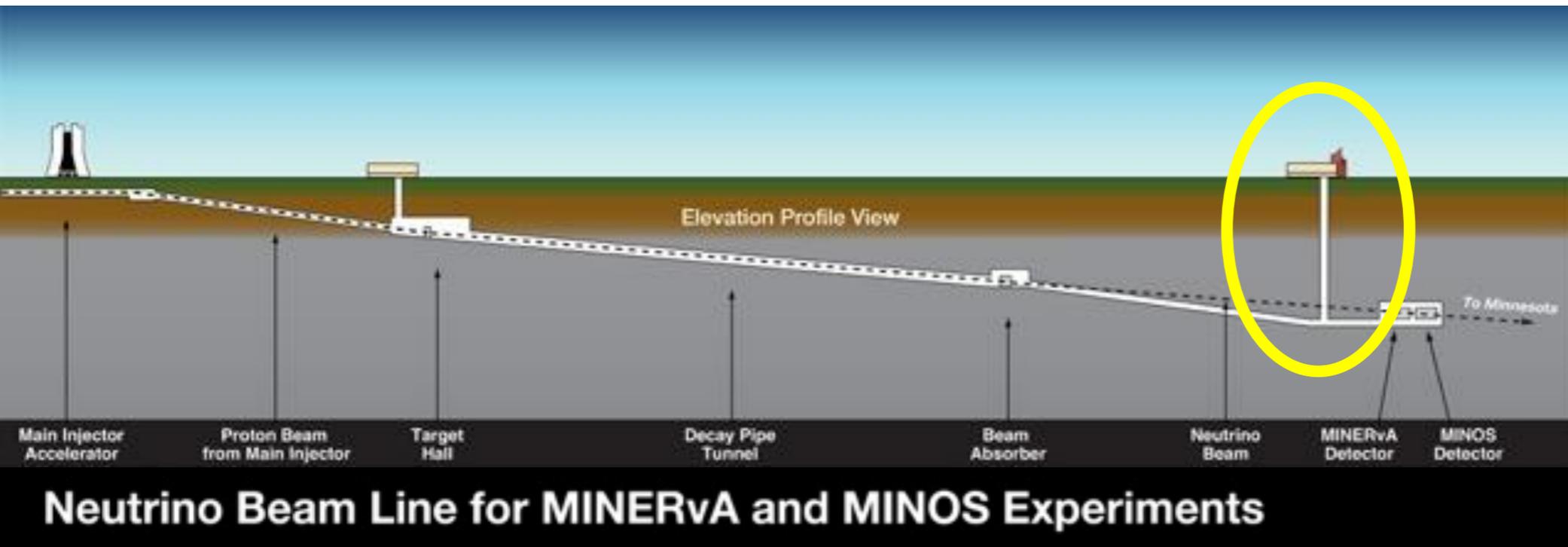
remove laser noise with single baseline



accelerometer → atoms are good inertial test masses

terrestrial? seismic noise automatically removed

# 100 m Detector Proposal at Fermilab



**Source 1**

- 100 m atom interferometer (accelerometer) drop tower
- $>3$  s drop time to split and recombine atomic wavefunctions

**Detector 1**  
**Source 2**

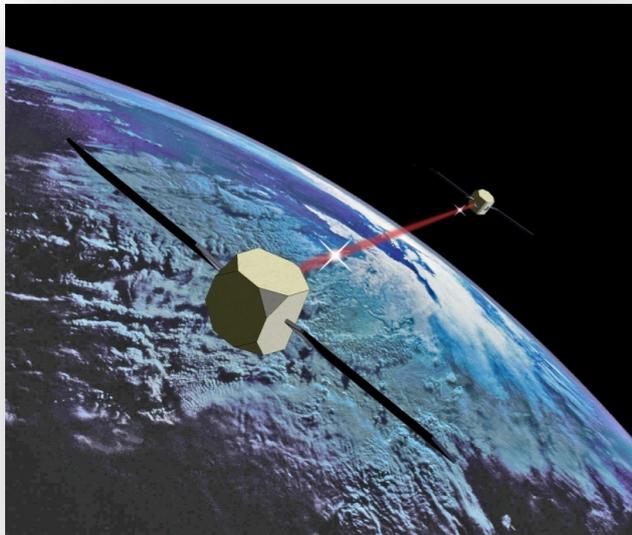
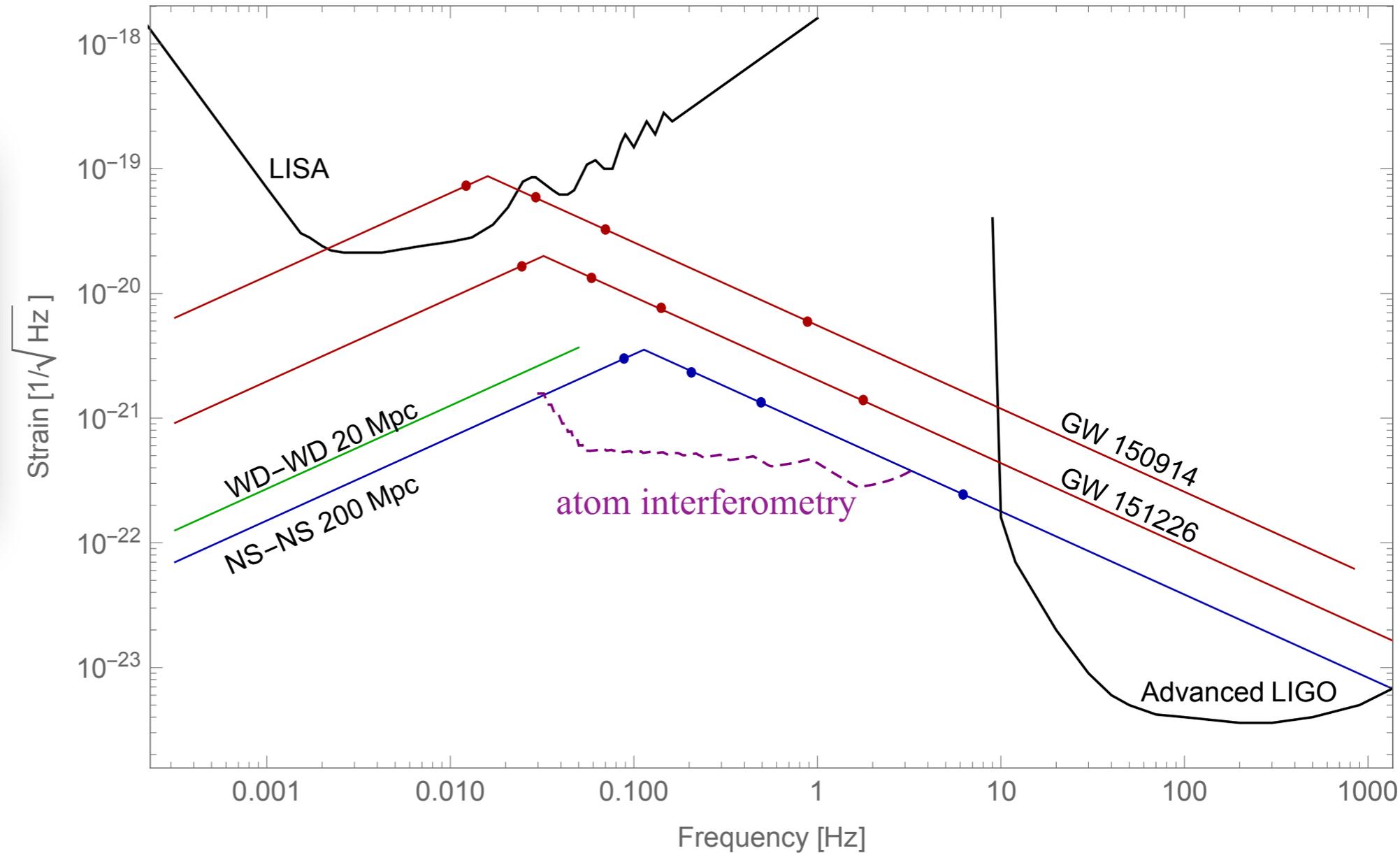
- Detect dark matter through oscillatory force
- Also gravitational waves from unknown sources

**Detector 2**

- Lead to  $\sim$ km scale detector for GW's (e.g. BH mergers) and DM, opens band below LIGO and above LISA ( $\sim 0.1 - 10$  Hz)

# Atom Interferometry for Gravitational Waves

Atoms could access mid-frequency band



earth orbit allows  
polarization measurement  
with single detector

for example this band allows:

localize sources on the sky (e.g. sub-degree accuracy) and predict  
BH and NS binary mergers for other telescopes to observe

with Sunghoon Jung

may measure initial BH spins and orbital eccentricity

# Recent Experimental Results

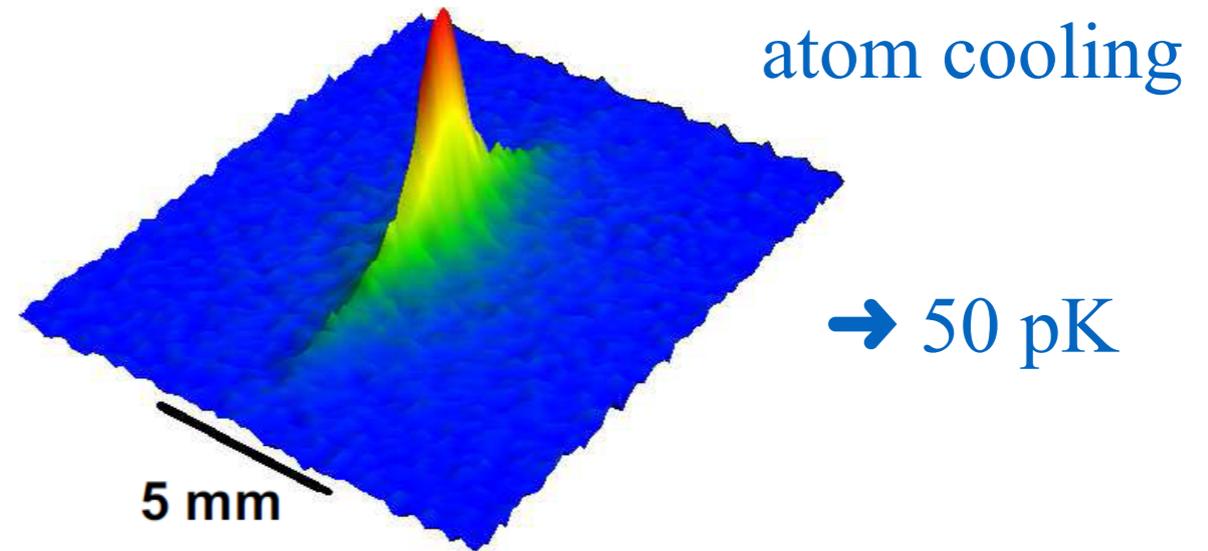
(Kasevich and Hogan groups)



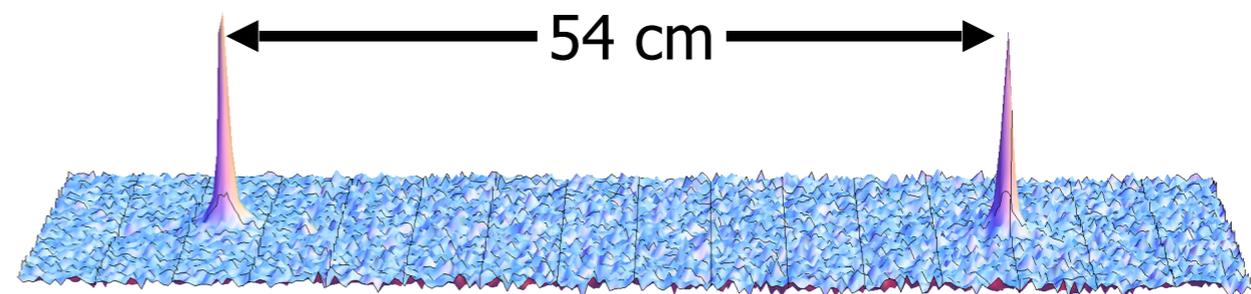
## Stanford Test Facility



demonstrate necessary technologies:



Macroscopic splitting of atomic wavefunction:



Kovachy et. al, *Nature* (2015)

# Summary

Precision measurement is a powerful tool for particle physics and cosmology

new technologies beyond traditional particle detectors

e.g. combination of several experiments will cover QCD axion dark matter fully

Light dark matter (axions) and gravitational wave detection similar:

detect coherent effects of entire field, not single particles

- laser interferometry
- atom interferometry (clocks)
- NMR
- high-precision magnetometry (SQUIDs, atomic systems)
- torsion pendulums
- optically-levitated dielectric spheres
- ...

Many more possibilities we haven't thought of yet...



# Backup Slides

# Possibilities for Light Dark Matter

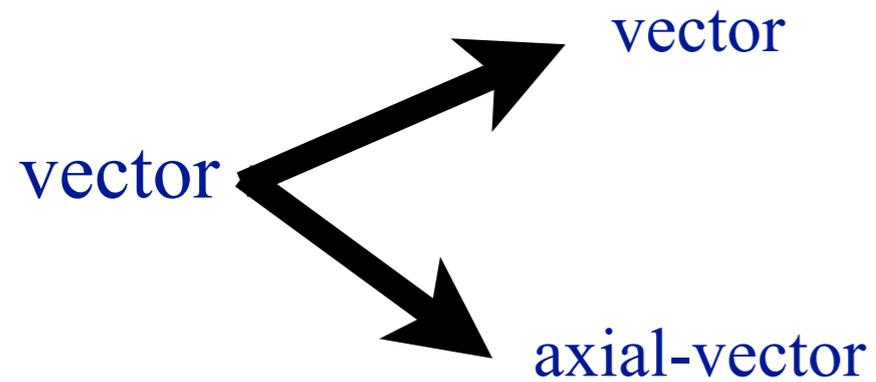
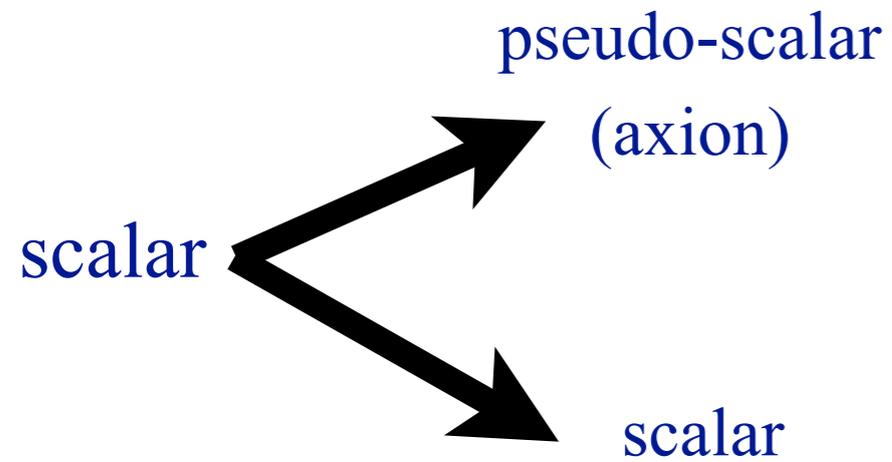
All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):

scalar

vector

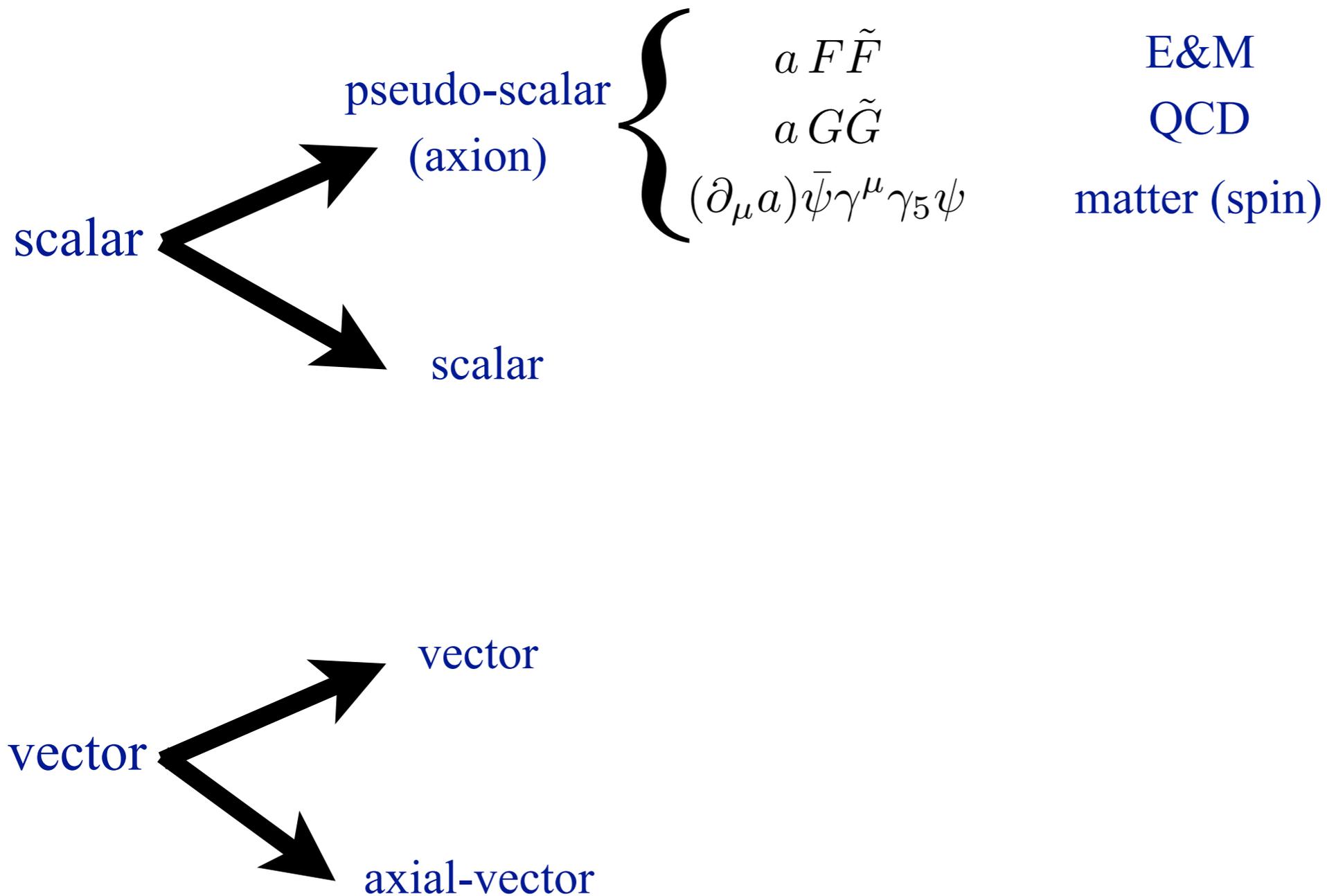
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



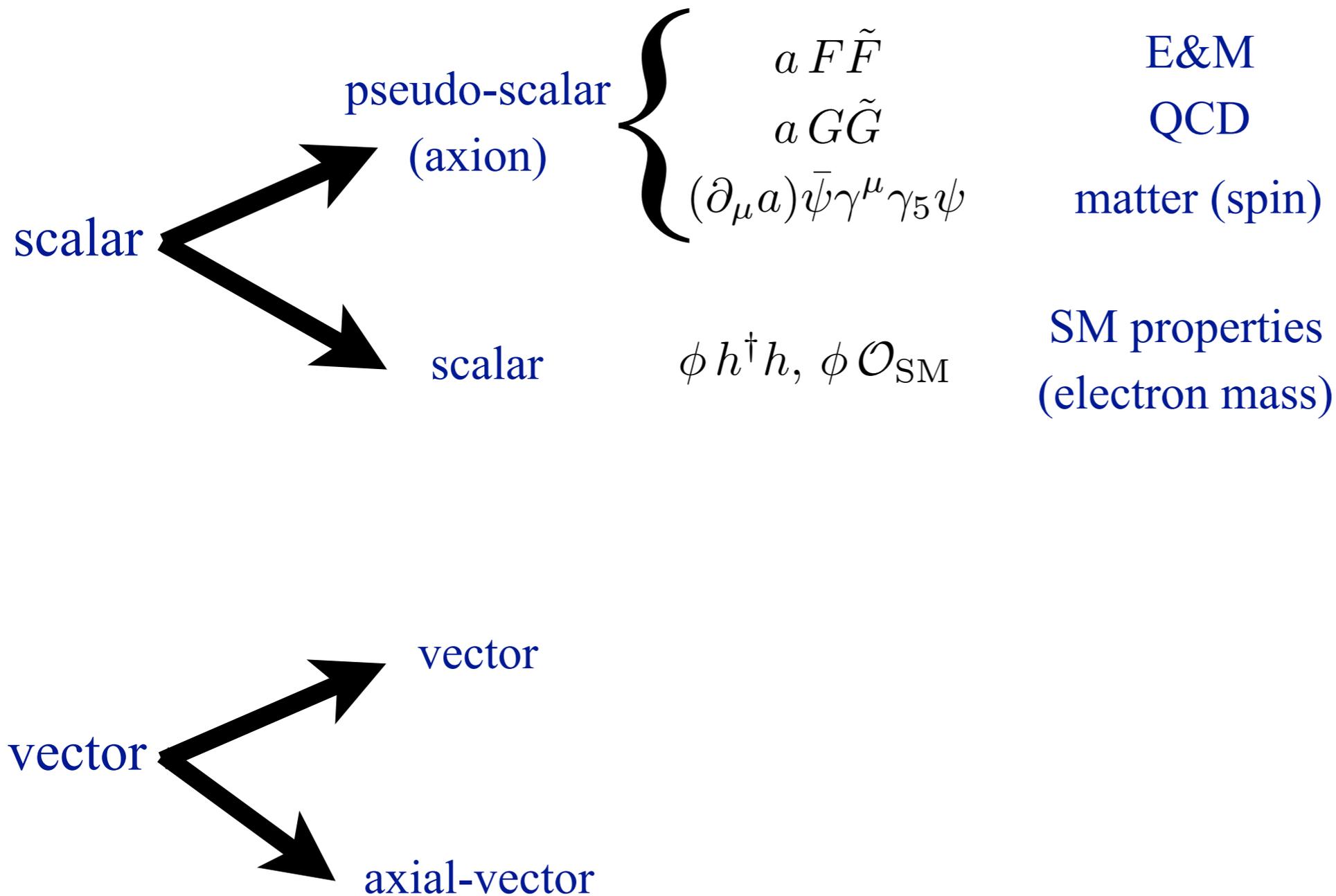
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



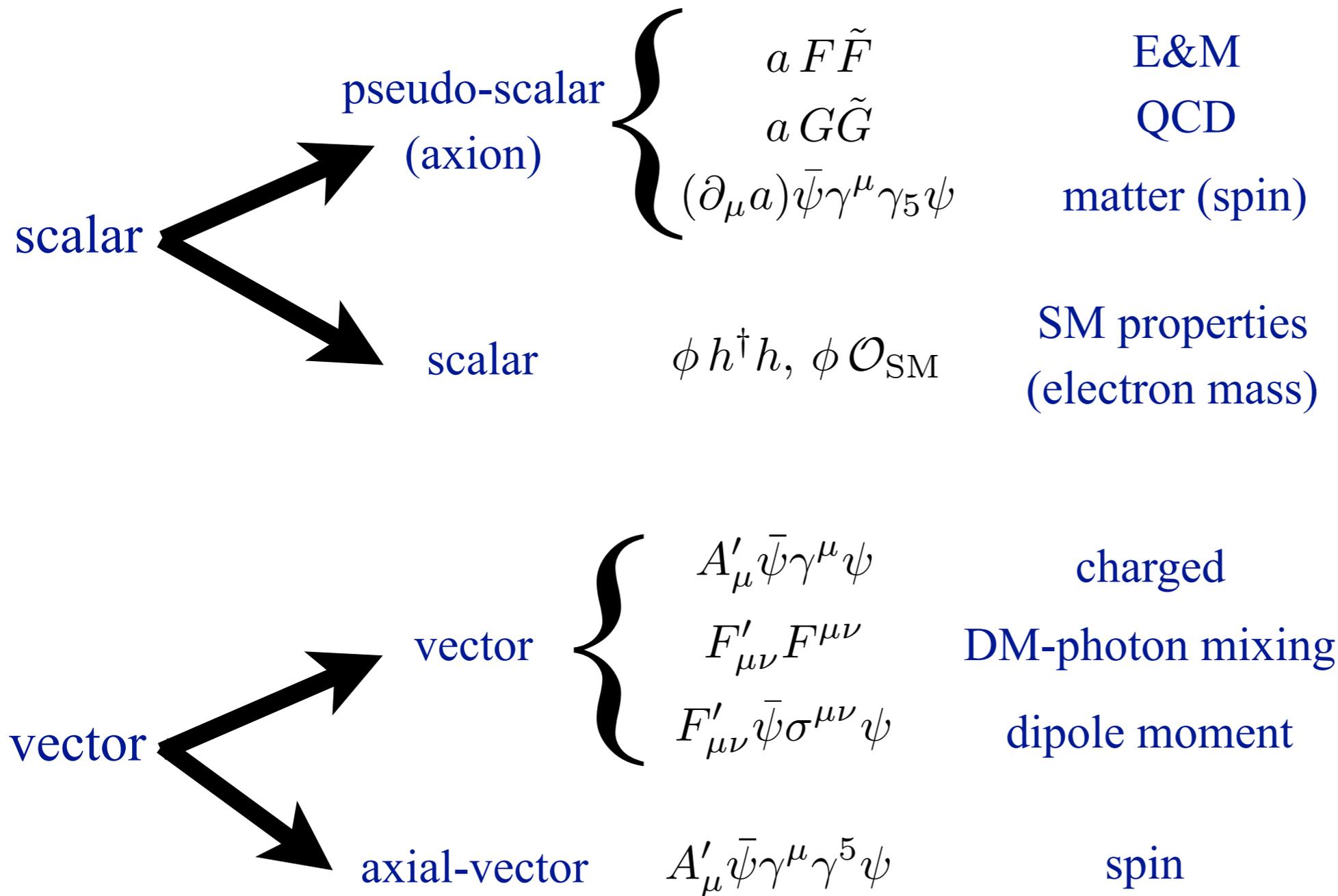
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



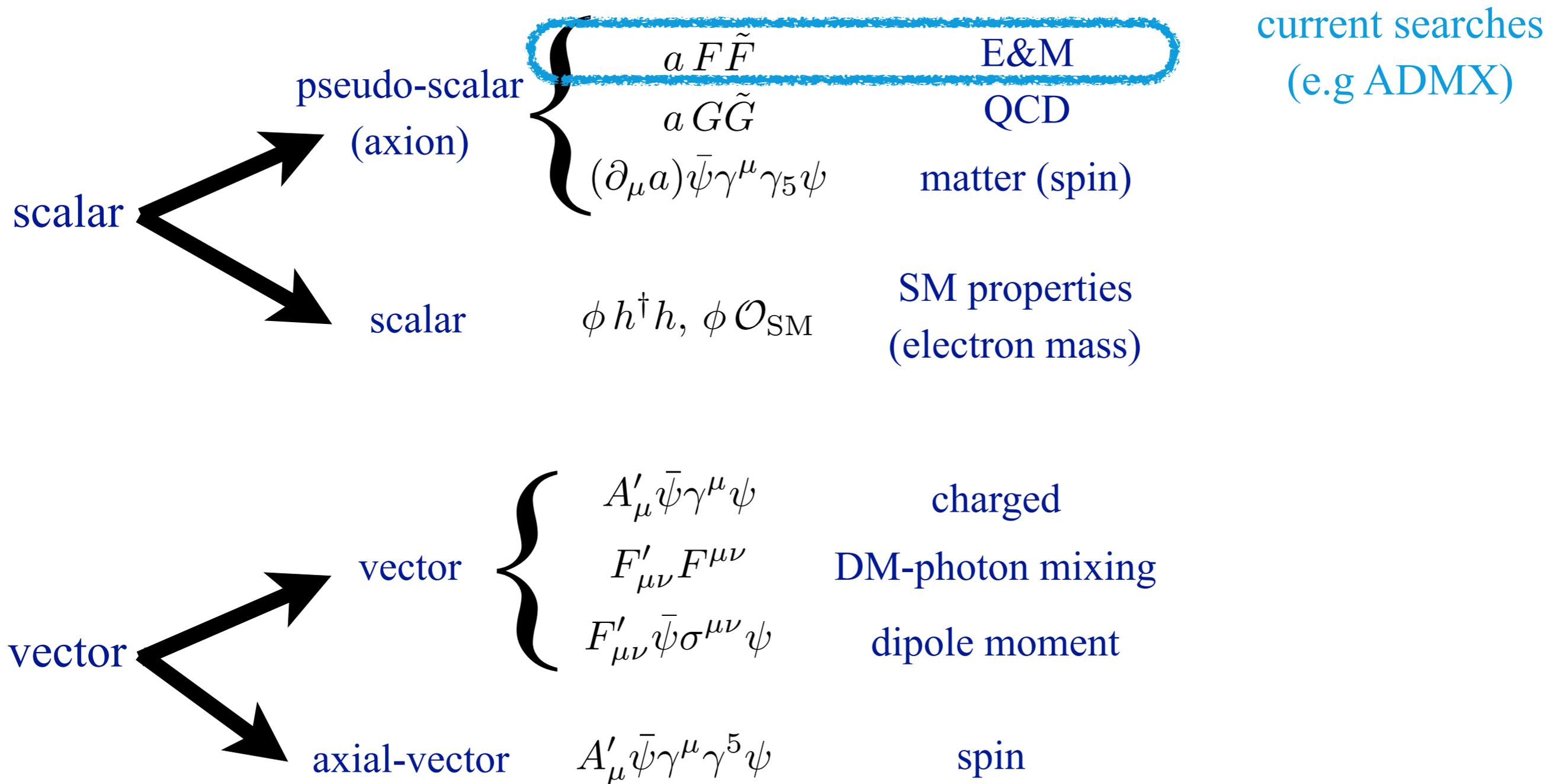
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



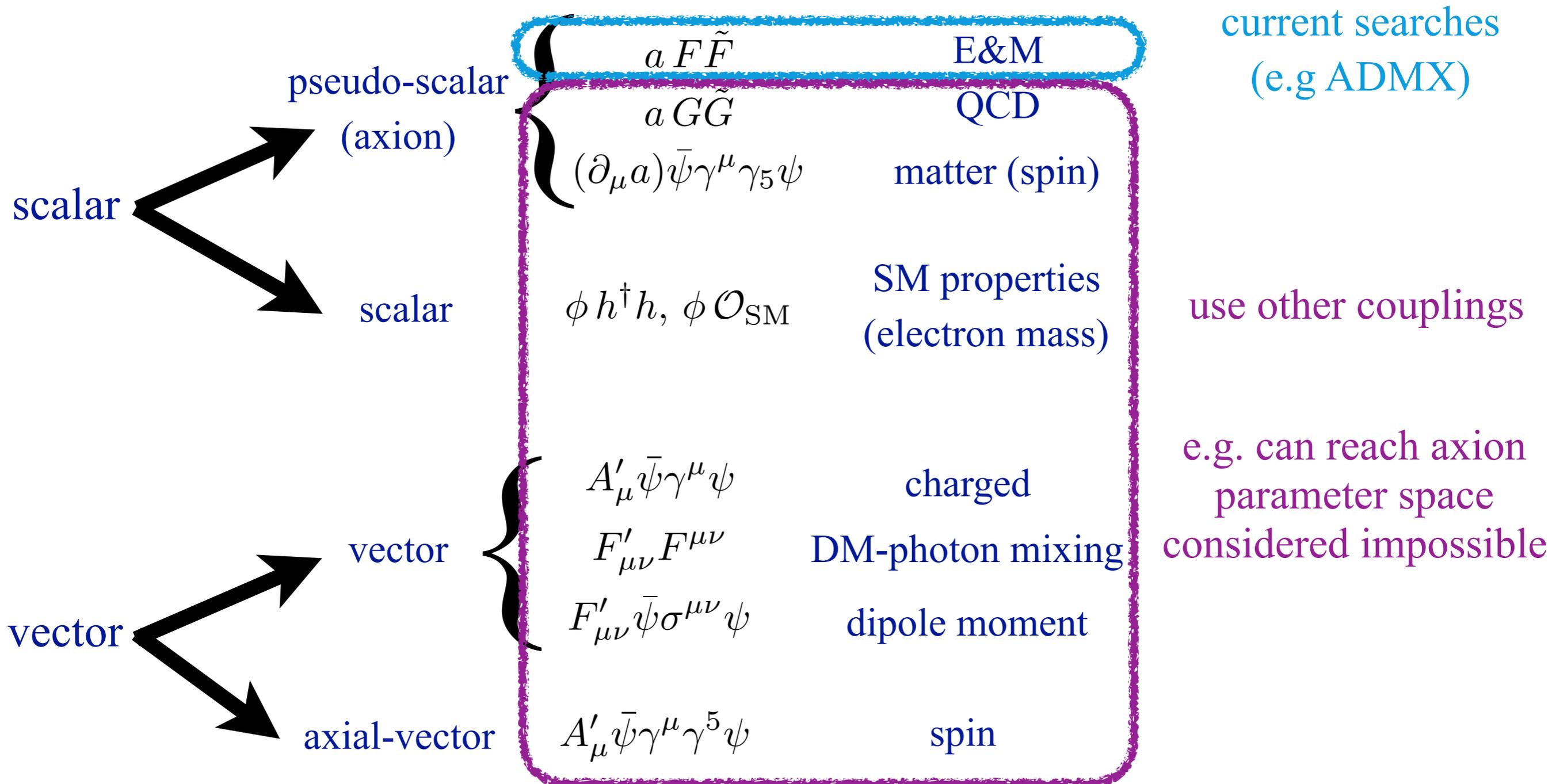
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



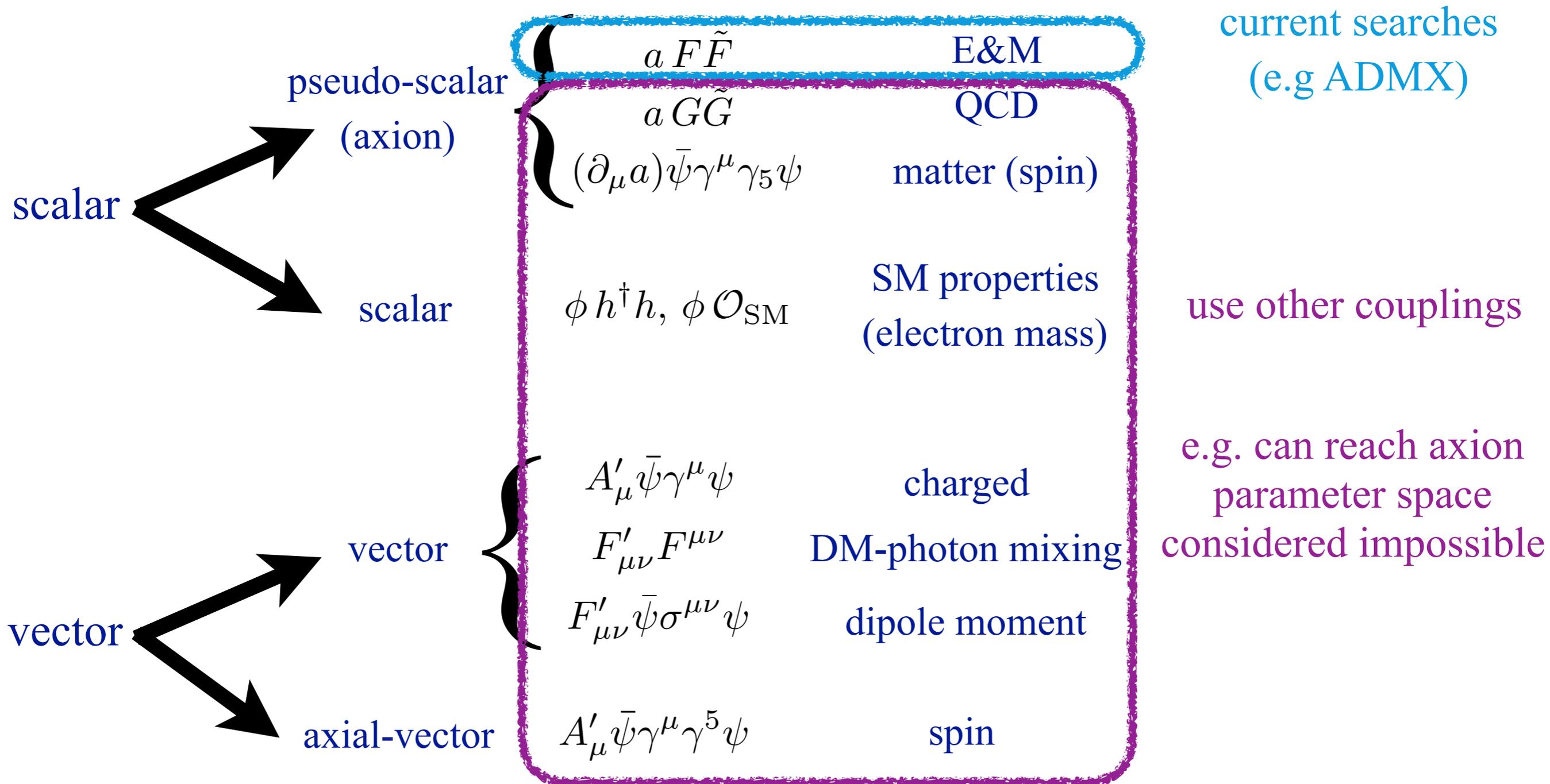
# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



# Possibilities for Light Dark Matter

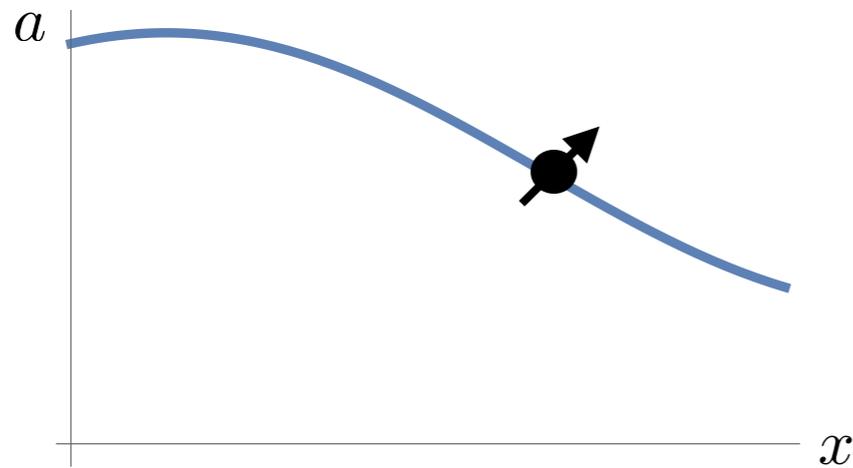
All UV theories summarized by only a few possibilities  
(symmetry, effective field theory):



Can cover all these possibilities!

# CASPEr-Wind

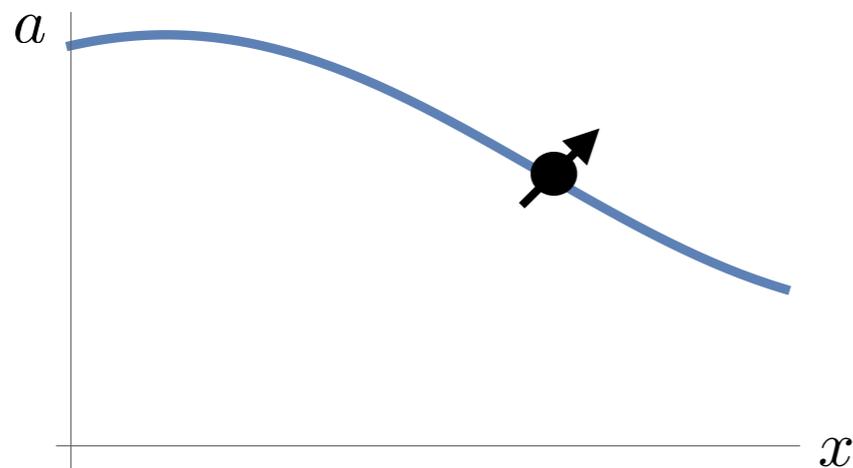
can also use direct coupling of axion to nucleons:  $(\partial_\mu a)\bar{\psi}\gamma^\mu\gamma_5\psi \rightarrow H \ni \nabla a \cdot \vec{\sigma}_N$



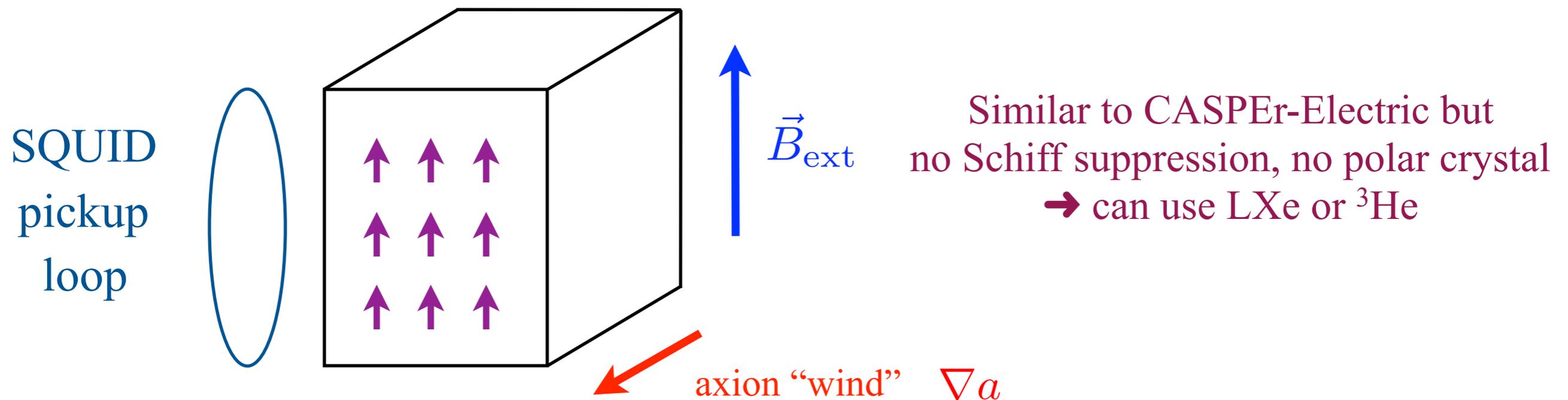
nuclear spins precess in DM axion field  
proportional to axion momentum (“wind”)

# CASPEr-Wind

can also use direct coupling of axion to nucleons:  $(\partial_\mu a)\bar{\psi}\gamma^\mu\gamma_5\psi \rightarrow H \ni \nabla a \cdot \vec{\sigma}_N$

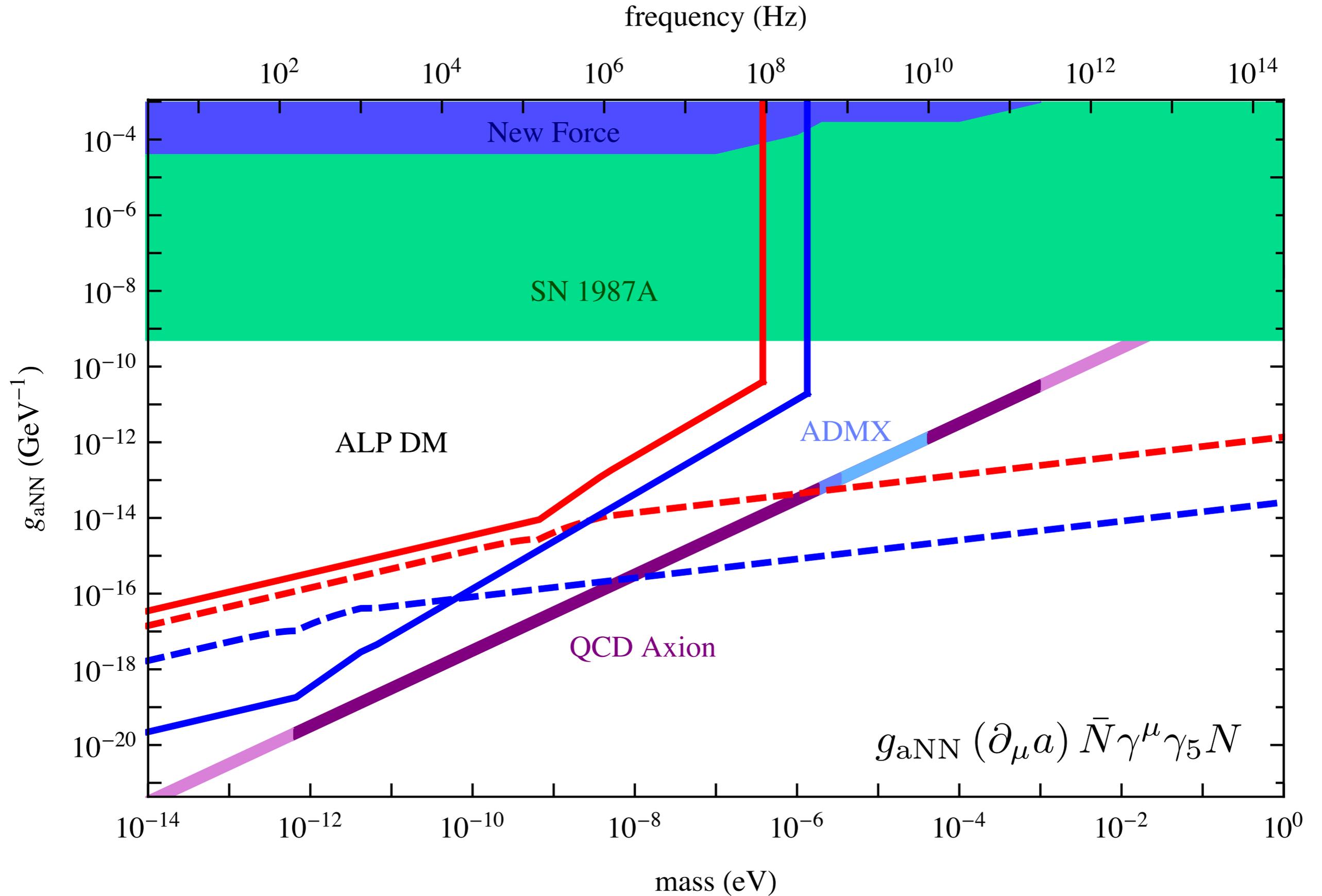


nuclear spins precess in DM axion field  
proportional to axion momentum (“wind”)



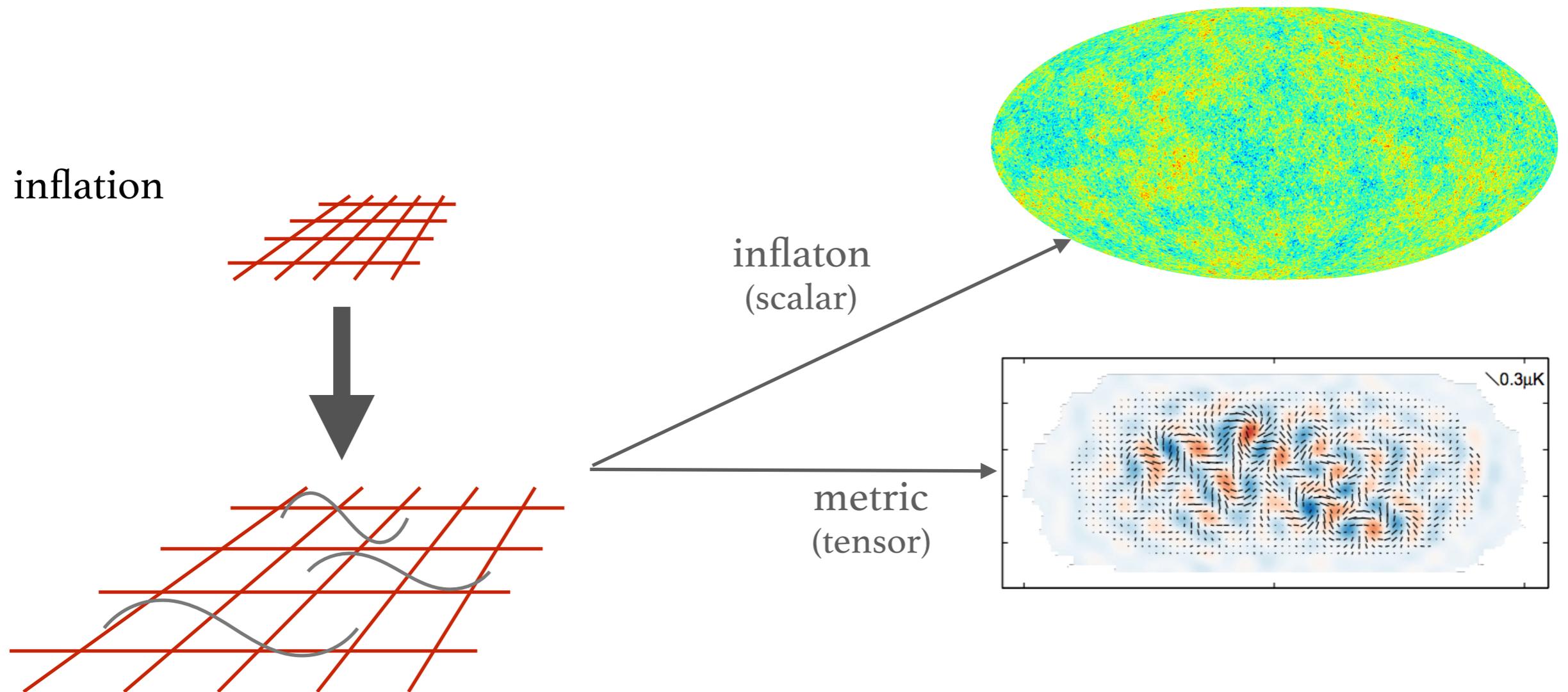
makes a directional detector for axions (and gives annual modulation)

# Limits on Axion-Nucleon Coupling

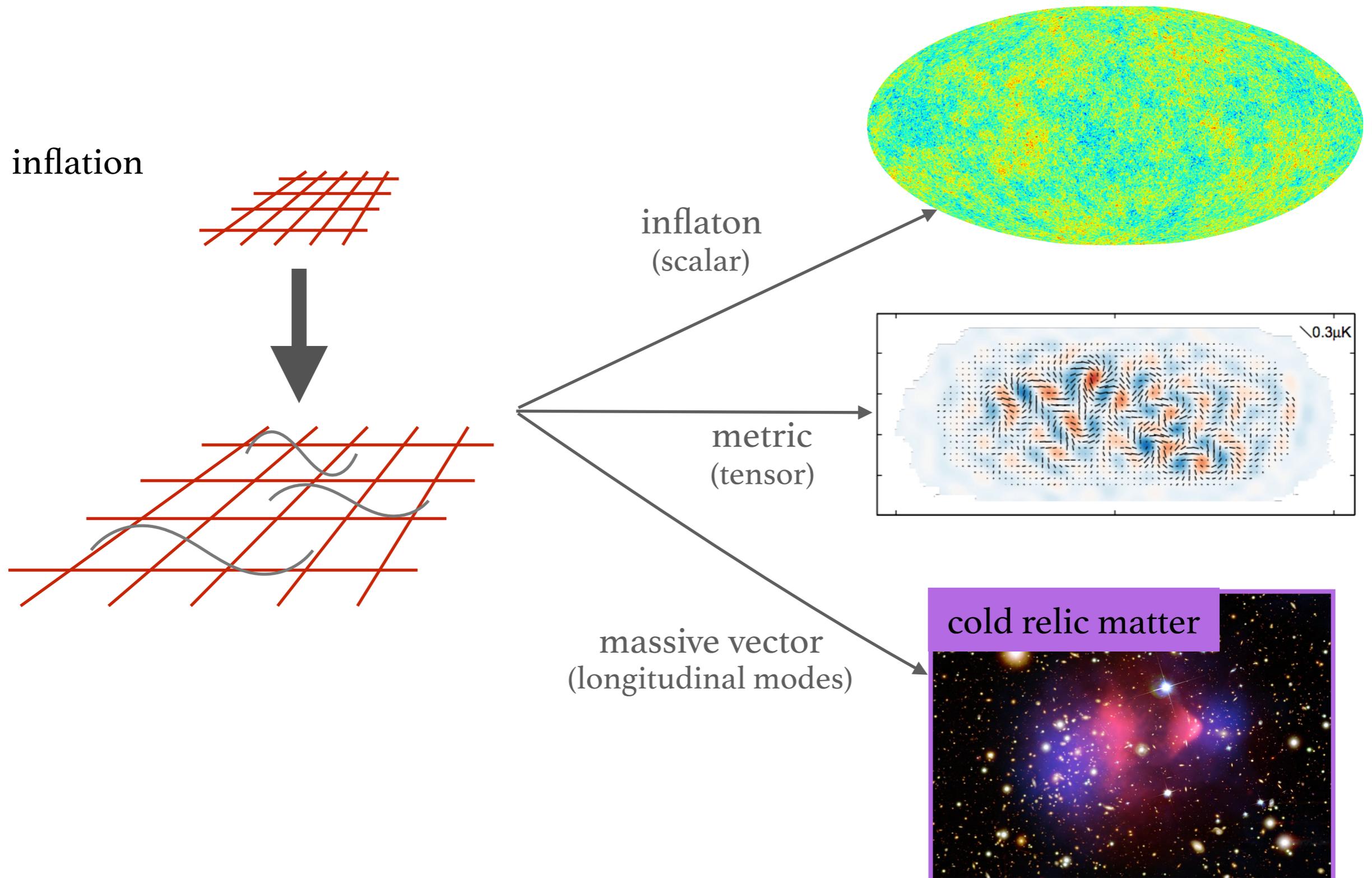


# Vectors

# Inflationary Production of Massive Vector

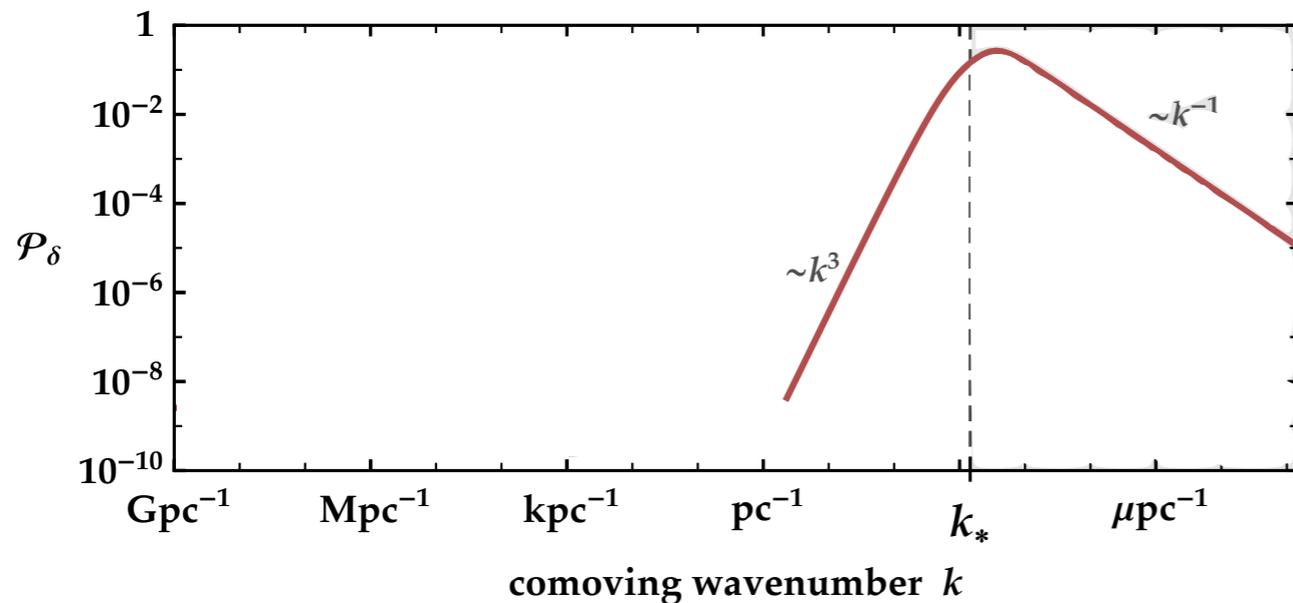


# Inflationary Production of Massive Vector

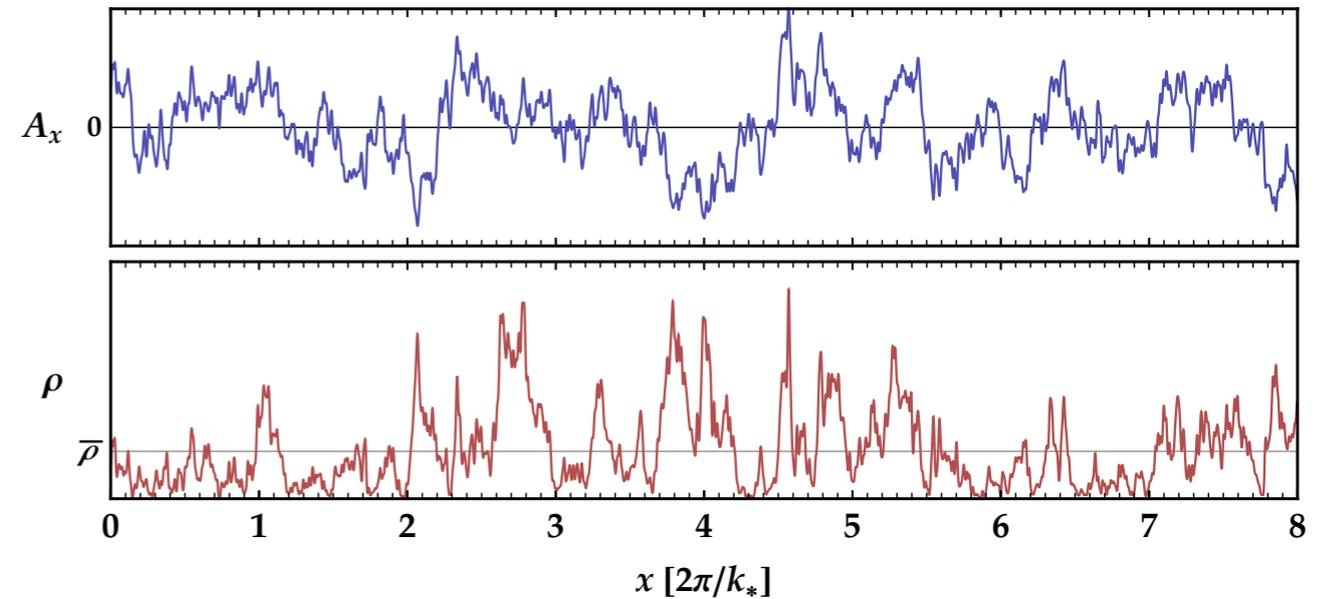


# Relic vector matter

## Power spectrum



## Large short-scale fluctuations



scalar: scale-invariant on long scales  $\rightarrow$  isocurvature problem

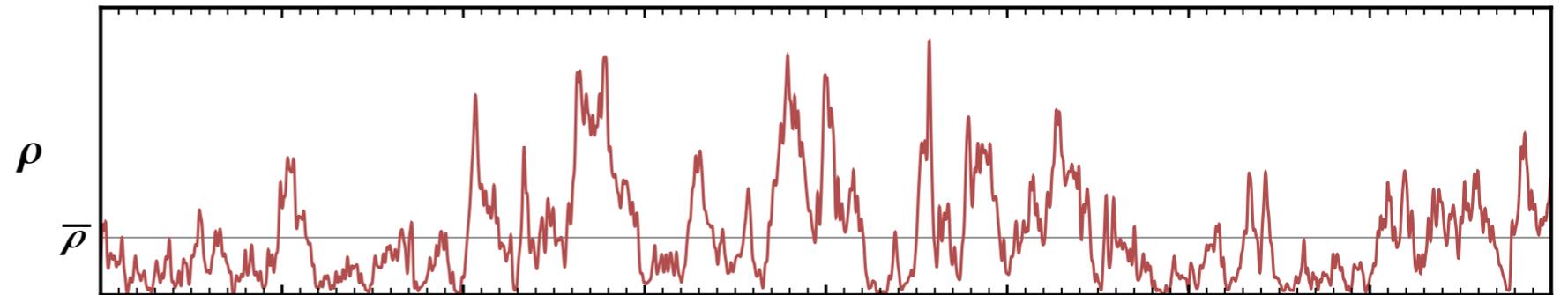
vector: suppressed power on long scales  $\rightarrow$  NO isocurvature problem

Abundance set by  $H_I$  and  $m$ :

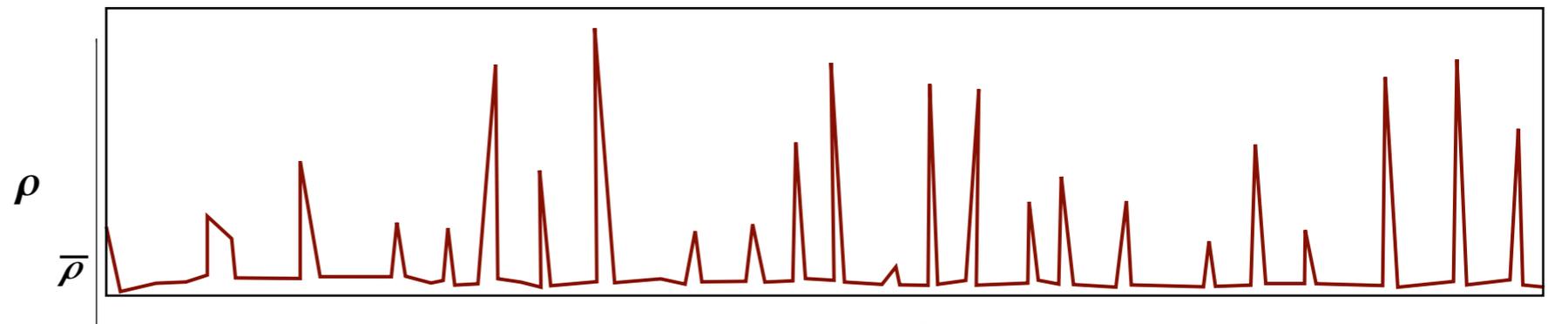
$$\frac{\Omega_A}{\Omega_{\text{cdm}}} \approx \sqrt{\frac{m}{6 \times 10^{-6} \text{ eV}}} \left( \frac{H_I}{10^{14} \text{ GeV}} \right)^2$$

# Substructure (best guess)

Primordial  
fluctuations  
(pre m.r.e.)



Self-bound  
clumps  
(post m.r.e.)



time

**small clumps**

density  $\sim \rho_{\text{mre}}$  ( $10^6$  overdense)

size  $\sim L_{\text{mre}} \ll 1 \text{A.U.}$

spacing  $\sim (a_0/a_{\text{mre}})L_{\text{mre}} \approx 1 \text{A.U.}$

# Implication for Searches?

- Large small-scale clumping is a general possibility
- $\rho$  in lab varies over time
- $\rho \rightarrow \sim 10^6 \rho$  when in clump (or more?)
- $\rightarrow$  signal is HUGE and easy to see...  
...IF at right frequency
- Best strategy: scan range fast and repeat (maximizes chance of hitting frequency when in clump)