1 **Standard Model and Higgs Physics Results from ATLAS and CMS**

Presented at the SLAC Summer Institute 2017

SSI

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The Standard Model 24 elementary matter particles 3 forces 26 parameters

- Although the SM has many input parameters, it is very predictive
- Survived many experimental tests over a wide energy range
- With the discovery of the Higgs boson, it is fair to say that it is a good e ffective theory no matter what happens next

The Higgs boson in the Standard Model 508. G. Guralnik, C. Hagen, and T.W.B. Kibble, PRL

- **Local gauge invariance** forbids explicit mass terms in the Lagrangian – but experimentally both gauge bosons and fermions have mass
- Introduce a new field with a very specific potential that keeps the full Lagrangian invariant but makes the vacuum not invariant
- **Higgs mechanism** predicts existence of a new, neutral boson: the **Higgs boson**
	- SM parameters: mass (**µ** or **m**H₎ and vacuum expectation value, **v**

$$
\mathcal{L} = |D^{\mu}\phi|^2 - y_i q_L^i q_R^i \phi - \mu^2 \phi^2 - \lambda \phi^4 + \dots
$$

perturbations around this minimum. To do this it is more natural to introduce a field ⌘

Standard Model Lagrangian

 $\mathcal{L}_{SM}=-\tfrac{1}{2}\partial_{\nu}g_{\mu}^{a}\partial_{\nu}g_{\mu}^{a}-g_{s}f^{abc}\partial_{\mu}g_{\nu}^{a}g_{\mu}^{b}g_{\nu}^{c}-\tfrac{1}{4}g_{s}^{2}f^{abc}f^{ade}g_{\mu}^{b}g_{\nu}^{c}g_{d}^{d}g_{\nu}^{e}-\partial_{\nu}W_{\mu}^{+}\partial_{\nu}W_{\mu}^{-}$ $M^{2}W_{\mu}^{+}W_{\mu}^{-} - \frac{1}{2}\partial_{\nu}Z^{0}_{\mu}\partial_{\nu}Z^{0}_{\mu} - \frac{1}{2c^{2}}M^{2}Z^{0}_{\mu}Z^{0}_{\mu} - \frac{1}{2}\partial_{\mu}A_{\nu}\partial_{\mu}A_{\nu} - igc_{w}(\partial_{\nu}Z^{0}_{\mu}(W^{+}_{\mu}W^{-}_{\nu} W^+_\nu W^-_\mu) - Z^0_\nu (W^+_\mu \partial_\nu W^-_\mu - W^-_\mu \partial_\nu W^+_\mu) + Z^0_\mu (W^+_\nu \partial_\nu W^-_\mu - W^-_\nu \partial_\nu W^+_\mu))$ $ig s_w (\partial_\nu A_\mu (W^+_\mu W^-_\nu - W^+_\nu W^-_\mu) - A_\nu (W^+_\mu \partial_\nu W^-_\mu - W^-_\mu \partial_\nu W^+_\mu) + A_\mu (W^+_\nu \partial_\nu W^-_\mu W_{\nu}^- \partial_{\nu} W_{\mu}^+) - \frac{1}{2} g^2 W_{\mu}^+ W_{\mu}^- W_{\nu}^+ W_{\nu}^- + \frac{1}{2} g^2 W_{\mu}^+ W_{\nu}^- W_{\mu}^+ W_{\nu}^- + g^2 c_w^2 (Z_u^0 W_{\mu}^+ Z_\nu^0 W_{\nu}^- Z^0_\mu Z^0_\mu W^+_{\nu} W^-_\nu + g^2 s^2_w (A_\mu W^+_\mu A_\nu W^-_\nu - A_\mu A_\mu W^+_\nu W^-_\nu) + g^2 s_w c_w (A_\mu Z^0_\nu W^+_\mu W^-_\nu W_{\nu}^{+}W_{\mu}^{-} - 2A_{\mu}Z_{\mu}^{0}W_{\nu}^{+}W_{\nu}^{-} - \frac{1}{2}\partial_{\mu}H\partial_{\mu}H - 2M^{2}\alpha_{h}H^{2} - \partial_{\mu}\phi^{+}\partial_{\mu}\phi^{-} - \frac{1}{2}\partial_{\mu}\phi^{0}\partial_{\mu}\phi^{0} \beta_h\left(\frac{2M^2}{a^2}+\frac{2M}{a}H+\frac{1}{2}(H^2+\phi^0\phi^0+2\phi^+\phi^-)\right)+\frac{2M^4}{a^2}\alpha_h$ $g\alpha_h M (H^3 + H\phi^0\phi^0 + 2H\phi^+\phi^-)$ $\frac{1}{8}q^2\alpha_h(H^4 + (\phi^0)^4 + 4(\phi^+\phi^-)^2 + 4(\phi^0)^2\phi^+\phi^- + 4H^2\phi^+\phi^- + 2(\phi^0)^2H^2)$ $gM W^+_\mu W^-_\mu H - \frac{1}{2} g \frac{M}{c^2} Z^0_\mu Z^0_\mu H \frac{1}{2}ig\left(W_\mu^+(\phi^0\partial_\mu\phi^- - \phi^-\partial_\mu\phi^0) - W_\mu^-(\phi^0\partial_\mu\phi^+ - \phi^+\partial_\mu\phi^0)\right) +$ $\frac{1}{2}g\left(W^+_\mu(H\partial_\mu\ddot{\phi^-}-\dot{\phi^-}\partial_\mu H)+W^-_\mu(H\partial_\mu\phi^+-\phi^+\partial_\mu H)\right)+\frac{1}{2}g\frac{1}{c_w}(Z^0_\mu(H\partial_\mu\phi^0-\phi^0\partial_\mu H)+$ $M\left(\frac{1}{c_w}Z^0_{\mu}\partial_{\mu}\phi^0+W^+_{\mu}\partial_{\mu}\phi^-+W^-_{\mu}\partial_{\mu}\phi^+\right)-ig\frac{s_w^2}{c_w}MZ^0_{\mu}(W^+_{\mu}\phi^- -W^-_{\mu}\phi^+)+ig s_w MA_{\mu}(W^+_{\mu}\phi^- \begin{array}{c} W^-_\mu \phi^+) - ig \frac{1-2c_w^2}{2c_w} Z^0_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) + ig s_w A_\mu (\phi^+ \partial_\mu \phi^- - \phi^- \partial_\mu \phi^+) \ - \frac{1}{4} g^2 W^+_\mu W^-_\mu \left(H^2 + (\phi^0)^2 + 2 \phi^+ \phi^- \right) - \frac{1}{8} g^2 \frac{1}{c_w^2} Z^0_\mu Z^0_\mu \left(H^2 + (\phi^0)^2 + 2 (2 s_w^2 - 1)^2 \phi^+ \phi^- \right) - \end{array}$ $\frac{1}{2}g^2\frac{s_w^2}{c_w}Z^0_\mu\phi^0(W^+_\mu\phi^- + W^-_\mu\phi^+) - \frac{1}{2}ig^2\frac{s_w^2}{c_w}Z^0_\mu H(W^+_\mu\phi^- - W^-_\mu\phi^+) + \frac{1}{2}g^2s_wA_\mu\phi^0(W^+_\mu\phi^- + W^-_\mu\phi^+)$ $W^{-}_{\mu}\phi^{+}\big) + \frac{1}{2}ig^{2}s_{w}A_{\mu}H(W^{+}_{\mu}\phi^{-} - W^{-}_{\mu}\phi^{+}) - g^{2}\frac{s_{w}}{c_{w}}(2c_{w}^{2} - 1)Z^{0}_{\mu}A_{\mu}\phi^{+}\phi^{-}$ $g^2s_w^2A_\mu A_\mu\phi^+\phi^- + \frac{1}{2}ig_s\lambda_{ii}^a(\bar{q}_i^\sigma\gamma^\mu q_i^\sigma)g_\mu^a - \bar{e}^{\lambda}(\gamma\partial+m_e^{\lambda})\bar{e}^{\lambda} - \bar{\nu}^{\lambda}(\gamma\partial+m_\nu^{\lambda})\nu^\lambda - \bar{u}_i^{\lambda}(\gamma\partial+m_\nu^{\lambda})\nu^\lambda$ $m_u^{\lambda} u_i^{\lambda} - d_i^{\lambda} (\gamma \partial + m_d^{\lambda}) d_i^{\lambda} + ig s_w A_\mu \left(- (\bar{e}^{\lambda} \gamma^\mu e^{\lambda}) + \frac{2}{3} (\bar{u}_i^{\lambda} \gamma^\mu u_i^{\lambda}) - \frac{1}{3} (\bar{d}_i^{\lambda} \gamma^\mu d_i^{\lambda}) \right) +$ $\frac{ig}{4c}Z_u^0\{(\bar{\nu}^{\lambda}\gamma^{\mu}(1+\gamma^5)\nu^{\lambda})+(\bar{e}^{\lambda}\gamma^{\mu}(4s_w^2-1-\gamma^5)e^{\lambda})+(\bar{d}_i^{\lambda}\gamma^{\mu}(\frac{4}{3}s_w^2-1-\gamma^5)d_i^{\lambda})+$ $(\bar{u}_j^{\lambda}\gamma^{\mu}\tilde{(1-\frac{8}{3}s_w^2+\gamma^5)u_j^{\lambda})}\}+\frac{ig}{2\sqrt{2}}W^+_{\mu}\left((\bar{\nu}^{\lambda}\gamma^{\mu}(1+\gamma^5)U^{lep}{}_{\lambda\kappa}e^{\kappa})+(\bar{u}_j^{\lambda}\gamma^{\mu}(1+\gamma^5)C_{\lambda\kappa}d_j^{\kappa})\right)+$ $\frac{ig}{2\sqrt{2}}W^{-}_{\mu}\left((\bar{e}^{\kappa}U^{lep}^{\dagger}_{\kappa\lambda}\gamma^{\mu}(1+\gamma^5)\nu^{\lambda}) + (\bar{d}_{j}^{\kappa}C^{\dagger}_{\kappa\lambda}\gamma^{\mu}(1+\gamma^5)u_{j}^{\lambda}) \right) +$ $\frac{ig}{2M\sqrt{2}}\phi^+\left(-m_e^{\kappa}(\bar{\nu}^{\lambda}U^{lep}{}_{\lambda\kappa}(1-\gamma^5)e^{\kappa})+m_{\nu}^{\lambda}(\bar{\nu}^{\lambda}U^{lep}{}_{\lambda\kappa}(1+\gamma^5)e^{\kappa}\right)+$ $\frac{ig}{2M\sqrt{2}}\phi^{-}\left(m_e^{\lambda}(\bar{e}^{\lambda}U^{lep}^{\dagger}_{\lambda\kappa}(1+\gamma^5)\nu^{\kappa})-m_{\nu}^{\kappa}(\bar{e}^{\lambda}U^{lep}^{\dagger}_{\lambda\kappa}(1-\gamma^5)\nu^{\kappa}\right)-\frac{g}{2}\frac{m_{\nu}^{\lambda}}{M}H(\bar{\nu}^{\lambda}\nu^{\lambda}) \frac{g}{2}\frac{m_e^{\lambda}}{M}H\left(\bar{e}^{\lambda}e^{\lambda}\right)+\frac{ig}{2}\frac{m_{\nu}^{\lambda}}{M}\phi^0(\bar{\nu}^{\lambda}\gamma^5\nu^{\lambda})-\frac{ig}{2}\frac{m_{e}^{\lambda}}{M}\phi^0(\bar{e}^{\lambda}\gamma^5e^{\lambda})-\frac{1}{4}\bar{\nu}_{\lambda}M_{\lambda\kappa}^R(1-\gamma_5)\hat{\nu}_{\kappa} -\frac{1}{4}\overline{\nu}\lambda M_{\lambda\kappa}^R(1-\gamma_5)\hat{\nu}\kappa + \frac{ig}{2M_{\lambda}\beta}\phi^+\left(-m_d^{\kappa}(\bar{u}_i^{\lambda}C_{\lambda\kappa}(1-\gamma^5)d_i^{\kappa}) + m_u^{\lambda}(\bar{u}_i^{\lambda}C_{\lambda\kappa}(1+\gamma^5)d_i^{\kappa}) + \right.$ $-\frac{ig}{2M\sqrt{2}}\phi^{-}\left(m_d^{\lambda}(\bar{d}_j^{\lambda}C^{\dagger}_{\lambda\kappa}(1+\gamma^5)u_j^{\kappa})-m_u^{\kappa}(\bar{d}_j^{\lambda}C^{\dagger}_{\lambda\kappa}(1-\gamma^5)u_j^{\kappa}\right)-\frac{g}{2}\frac{m_u^{\lambda}}{M}H(\bar{u}_j^{\lambda}u_j^{\lambda}) \frac{g}{2}\frac{m_d^{\lambda}}{M}H(\bar{d}_i^{\lambda}d_i^{\lambda}) + \frac{ig}{2}\frac{m_u^{\lambda}}{M}\phi^0(\bar{u}_i^{\lambda}\gamma^5u_i^{\lambda}) - \frac{ig}{2}\frac{m_d^{\lambda}}{M}\phi^0(\bar{d}_i^{\lambda}\gamma^5d_i^{\lambda}) + \bar{G}^a\partial^2G^a + g_sf^{abc}\partial_{\mu}\bar{G}^aG^bg_u^c +$ $\bar{X}^+ (\partial^2 - M^2) X^+ + \bar{X}^- (\partial^2 - M^2) X^- + \bar{X}^0 (\partial^2 - \frac{M^2}{c^2}) X^0 + \bar{Y} \partial^2 Y + i g c_w W^+_u (\partial_\mu \bar{X}^0 X^- \partial_u \bar{X}^+ X^0) + ig s_w W_u^+ (\partial_u \bar{Y} X^- - \partial_u \bar{X}^+ Y) + ig c_w W_u^- (\partial_u \bar{X}^- X^0 \partial_{\mu}\bar{X}^{0}X^{+}\big)+ig s_{w}\dot{W}_{\mu}^{-}\left(\partial_{\mu}\bar{X}^{-}Y-\partial_{\mu}\bar{Y}X^{+}\right)+ig c_{w}Z_{\mu}^{0}\left(\partial_{\mu}\bar{X}^{+}X^{+}-\right)$ $\partial_{\mu}\bar{X}^{-}X^{-})+ig s_{w}A_{\mu}(\partial_{\mu}\bar{X}^{+}X^{+} \partial_{\mu}\bar{X}^{-}X^{-})-\frac{1}{2}gM\left(\bar{X}^{+}X^{+}H+\bar{X}^{-}X^{-}H+\frac{1}{c_{-}^{2}}\bar{X}^{0}X^{0}H\right)+\frac{1-2c_{w}^{2}}{2c_{w}}igM\left(\bar{X}^{+}X^{0}\phi^{+}-\bar{X}^{-}X^{0}\phi^{-}\right)+$ $\frac{1}{2c} i g M \left(\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^- \right) + i g M s_w \left(\bar{X}^0 X^- \phi^+ - \bar{X}^0 X^+ \phi^- \right) +$ $\frac{1}{2}igM\left(\bar{X}^+X^+\phi^0-\bar{X}^-X^-\phi^0\right)$.

[T. Gutierrez](http://nuclear.ucdavis.edu/~tgutierr/files/stmL1.html)

Standard Model Lagrangian 5

$$
\mathcal{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{8} tr(\mathbf{W}_{\mu\nu} \mathbf{W}^{\mu\nu}) - \frac{1}{2} tr(\mathbf{G}_{\mu\nu} \mathbf{G}^{\mu\nu})
$$
\n
$$
+ (\bar{v}_L, \bar{e}_L) \tilde{\sigma}^{\mu} i D_{\mu} \begin{pmatrix} v_L \\ e_L \end{pmatrix} + \bar{e}_R \sigma^{\mu} i D_{\mu} e_R + \bar{\nu}_R \sigma^{\mu} i D_{\mu} v_R + (\text{h.c.})
$$
\n(lepton dynamical term)\n
$$
-\frac{\sqrt{2}}{v} \begin{bmatrix} (\bar{v}_L, \bar{e}_L) \phi M^e e_R + \bar{e}_R \bar{M}^e \bar{\phi} \begin{pmatrix} v_L \\ e_L \end{pmatrix} \end{bmatrix}
$$
\n
$$
+ (\bar{u}_L, \bar{d}_L) \tilde{\sigma}^{\mu} i D_{\mu} \begin{pmatrix} u_L \\ d_L \end{pmatrix} + \bar{u}_R \sigma^{\mu} i D_{\mu} u_R + \bar{d}_R \sigma^{\mu} i D_{\mu} d_R + (\text{h.c.})
$$
\n(neutrino mass term)\n
$$
+ (\bar{u}_L, \bar{d}_L) \tilde{\sigma}^{\mu} i D_{\mu} \begin{pmatrix} u_L \\ d_L \end{pmatrix} + \bar{u}_R \sigma^{\mu} i D_{\mu} u_R + \bar{d}_R \sigma^{\mu} i D_{\mu} d_R + (\text{h.c.})
$$
\n(quark dynamical term)\n
$$
-\frac{\sqrt{2}}{v} \begin{bmatrix} (\bar{u}_L, \bar{d}_L) \phi M^d d_R + \bar{d}_R \bar{M}^d \bar{\phi} \begin{pmatrix} u_L \\ d_L \end{pmatrix} \end{bmatrix}
$$
\n(down, strange, bottom mass term)\n
$$
-\frac{\sqrt{2}}{v} \begin{bmatrix} (-\bar{d}_L, \bar{u}_L) \phi M^d u_R + \bar{u}_R \bar{M}^u \phi^T \begin{pmatrix} -d_L \\ u_L \end{pmatrix} \end{bmatrix}
$$
\n(Higgs dynamical and mass term)\n
$$
+ (\overline{D}_{\mu} \phi) D^{\mu} \phi - m_h^
$$

where (h.c.) means Hermitian conjugate of preceeding terms, $\bar{\psi} = (\text{h.c.})\psi = \psi^{\dagger} = \psi^{*T}$, and the derivative operators are $D_\mu\left(\begin{array}{c} \nu_L \\ e_L \end{array}\right) = \left[\partial_\mu - \frac{ig_1}{2}B_\mu + \frac{ig_2}{2}\mathbf{W}_\mu\right]\left(\begin{array}{c} \nu_L \\ e_L \end{array}\right), \quad D_\mu\left(\begin{array}{c} u_L \\ d_L \end{array}\right) = \left[\partial_\mu + \frac{ig_1}{6}B_\mu + \frac{ig_2}{2}\mathbf{W}_\mu + ig\mathbf{G}_\mu\right]\left(\begin{array}{c} u_L \\ d_L \end{array}\right),$ (2) $D_\mu \nu_R = \partial_\mu \nu_R, \quad D_\mu e_R = \left[\partial_\mu - ig_1 B_\mu \right] e_R, \quad D_\mu u_R = \left[\partial_\mu + \frac{i2g_1}{3} B_\mu + ig \mathbf{G}_\mu \right] u_R, \quad D_\mu d_R = \left[\partial_\mu - \frac{ig_1}{3} B_\mu + ig \mathbf{G}_\mu \right] d_R,$ (3) $D_\mu \phi = \left[\partial_\mu {+} \frac{ig_1}{2} B_\mu {+} \frac{ig_2}{2} {\bf W}_\mu \right] \phi.$ (4)

This talk

- Briefly introduce the LHC and the ATLAS and CMS detectors
- Select a few key Standard Model particles (W, Z, top and Higgs)
- Highlight key ATLAS and CMS results
	- See how this has improved our understand of the Standard Model

The Large Hadron Collider and the ATLAS and CMS Detectors

The Large Hadron Collider (LHC) ⁹ points contain equipment used for beam contained for beam contained and 7 and 7 and 7 and 7 contains radio-fre

cavities; and P is the location of the location of the beam dump. Beam dump. Beam dump. In the beam dump.

Side note: Cross-sections and luminosity ¹⁰

Already ~12 fb-1 of 2017 data!

Cross-sections are measured in <u>barns</u>: 1 barn = 10^{-28} m² (100 $fm²$ Range at the LHC: mb to fb

Measuring Particles 11

The General Purpose Detectors 12

ATLAS CMS

The Detectors 13

$ATLAS$

A Large Toroidal ApparatuS (ATLAS)

http://atlas.cern/

The Compact Muon Solenoid (CMS)¹⁵

https://cms.cern/detector

Something new: CMS Pixel Upgrade CMS Pixel TD

[CMS Pixel TDR](https://cds.cern.ch/record/1481838?ln=en)

- CMS pixel detector was completely and simulations in the simulation of the s replaced during the 2016-2017 shut-
that the pixel hit position resolution resolution resolution study to be seen in this simulation resolution resolution resolution resolution resolution resolution resolution resolution r down for the current detector than what is currently achievable with the 2011/2012 data. Details for 2011/2012 data the configuration of the track reconstruction used is given in Section 2.1.2.
- **Additional pixel layer**
	- 3→4 barrel layers (smaller radius) a conceptual layout for the Phase 1 upgrade pixel detector. The current 3-layer
	-
- e.g. 50% improvement in do resolution

The Vector Bosons 17

The Z boson

- Discovered in 1983 at the SPS at CERN
- Carrier of the weak force
- Reconstruct from a pair of leptons of the same flavour but with opposite charge typically with $p_T > 20$ GeV
- One of the easiest processes to identify
	- Almost background free
- Widely used for lepton calibration
- Highly accurate tests of the Standard Model

[SMP-12-011](http://cms-results.web.cern.ch/cms-results/public-results/publications/SMP-12-011/index.html)

Z Boson Properties 19
19

[SMP-14-012](http://cms-results.web.cern.ch/cms-results/public-results/publications/SMP-14-012/index.html) [STDM-2014-10](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/STDM-2014-10/)

Measure Z boson angular distributions to probe QCD dynamics

[STDM-2014-18](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/STDM-2014-18/)

The W boson

- Also discovered in 1983 at the CERN SPS
- Other weak force mediator
- Most precise measurements reconstruct W boson from decay to a lepton and a neutrino

- Cannot reconstruct the full mass because the neutrino is only detected indirectly (missing energy)
- Backgrounds from multijet (fake lepton) and top quarks

The W Mass

- Precision measurement of W mass tests consistency of Standard Model
- Extremely challenging measurement at a hadron collider
- Template fit to distributions sensitive to the W mass
- Requires careful calibration and detailed understanding of reconstructed objects
- $m_W = 80370 \pm 19$ MeV

[STDM-2014-18](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/STDM-2014-18/)

Observation of WW scattering

- In the SM without the Higgs, the cross-section for WW scattering was predicted to diverge at high energies
	- One component of the "no lose theorem" which argued that the LHC had to find something
- Exactly two leptons of the same charge and two jets with a rapidity gap
- Counting experiment in 6 categories by lepton flavour
- Observed (expected) significance is 5.5 (5.7)σ

More than one: Diboson Results 23

Massive Triboson ?

24

W W W W

- Events / 40 GeV Events / 40 GeV 20 \rightarrow Data **WWW** $18\Box$ $\sqrt{s} = 8$ TeV, 20.3 fb⁻¹ WZ Fake L. 16 wh w Vγ Charge Flip L. 14 *0+1+2 SFOS SR* Other Bkg. $f_{S,0}/\Lambda^4 = 2000 \text{ TeV}^4$ 12 $\left\vert \right\vert \wedge_{\mathsf{FF}}$ = ∞ $\rm{f}_{\rm{S},4}^{\rm{}}/\Lambda^{\rm{4}}$ = 2000 TeV $^{\rm{-4}}$ 10 $f_{S,0}/\Lambda^4 = 2000 \text{ TeV}^4$ 8 $\rm{f}_{\rm{S},1}^{\rm{}}/\Lambda^{\rm{4}}$ = -6000 TeV⁻⁴ 6 4 2 Ω B
Data/S+B
0.5
0.5 l.5 \natural [GeV] lll |||
|||| 1 0 100 200 300 400 500 600 700 800 900 1000 m_T^{3l} [GeV]
- Currently only a limit from ATLAS on WWWW coupling

The Top Quark 25

The Top Quark

- The heaviest particle in the Standard Model with a Yukawa (Higgs) coupling of \sim 1
- Discovered at the Tevatron, but large production rate at the LHC allows its properties to be studied in detail
- Typically study ttbar production
- Each top decays to a W-boson and a b-quark
	- Either leptonic or hadronic W decay
- Production cross-section measured to 4%
	- Consistent with theoretical predictions

The Top Mass

- Another key parameter of the Standard Model
- Measured using analogous techniques to the W mass measurement, but more challenging as it requires both leptons and jets
- Theory interpretation is challenging
	- Measured top mass \neq theoretical pole mass

Single top production 28

q^q *Measurements of the top produced by itself*

Top and other particles

- Probe the production of the top together with other particles
- Measurement of ttW and ttZ crosssections:

 $\sigma(\text{t\bar{t}W}) = 0.80^{+0.12}_{-0.11}$ (stat.) $^{+0.13}_{-0.12}$ (sys.) pb $\sigma(\text{t\bar{t}Z}) = 1.00^{+0.09}_{-0.08}$ (stat.) $^{+0.12}_{-0.10}$ (sys.) pb

- Measurement of tty cross-section
	- 139 \pm 7(stat.) \pm 17(syst.) fb

[TOPQ-2015-21](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/TOPQ-2015-21/)

[TOP-17-005](http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/TOP-17-005/index.html)

The Higgs Boson 30

The Higgs Boson

- Predictions date from the 1960s
- Discovered at CERN by ATLAS and CMS in 2012
- Only known elementary scalar
- Particle associated with the Higgs mechanism which provides elementary particles with their mass

http://www.elsevier.com/locate/physlet b

Producing the Higgs 32

$$
\mathcal{L} = |D^{\mu}\phi|^2 - y_i q_L^i q_R^i \phi - \mu^2 \phi^2 - \lambda \phi^4 + \dots
$$

Massive gauge boson? …then it couples to the Higgs Massive fermion? …then it couples to the Higgs

Gauge bosons Fermions

Massive gauge boson?

…then it couples to the Higgs

Massive fermion? …then it couples to the Higgs

Higgs Production and Decays 34

Higgs Production

Higgs Decay

Discovery Channels 35

 $for m_{11} - 125$ GeV $for m_U = 125$ GaV $\overline{101}$ $\overline{111}$ $\overline{120}$ $\overline{0}$ An excellent channel for $m_H = 125$ GeV

The two \sim two \sim two \sim two \sim $Mida$ mass range range in The two contracts were contracted as a second contract of the $\frac{1}{2}$ range inde Golden channel over a wide mass range

*n*_{po} 500

<u>a</u> 20 - 30
 n De Channels with excellent mass resolution Simple channels with excellent mass resolution

4 July 2012

What do we know about the Higgs? 75 (2015) 212, [Phys. Lett. B 726 \(2013\), pp. 120-144](http://www.sciencedirect.com/science/article/pii/S0370269313006527)

- Measure basic properties
	- **Mass** and **width**
	- **Production** rate
	- **Spin** and **parity** (only elementary scalar):
	- Measure **decays**

$J = 0$

In the following H^0 refers to the signal that has been discovered in the Higgs searches. Whereas the observed signal is labeled as a spin 0 particle and is called a Higgs Boson, the detailed properties of H^0 and its role in the context of electroweak symmetry breaking need to be further clarified. These issues are addressed by the measurements listed below.

Concerning mass limits and cross section limits that have been obtained in the searches for neutral and charged Higgs bosons, see the sections "Searches for Neutral Higgs Bosons" and "Searches for Charged Higgs Bosons (H^{\pm}) and $H^{\pm \pm}$)", respectively.

Higgs Mass Measurement

[HIGG-2014-14](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HIGG-2014-14/), [ATLAS-CONF-2017-046,](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2017-046) [HIG-16-041](http://arxiv.org/abs/1706.09936)

- Final Run-1: 125.09 ± 0.24 GeV
- CMS Run-2: 125.26 ± 0.21 GeV
- ATLAS Run-2: 124.98 ± 0.28 GeV

Higgs Mass Implications

- Measured final SM parameter !
- Good consistency with m_W and m_{top}
- m_H = 125 GeV
	- A bit too heavy for supersymmetry, but not so heavy as to exclude supersymmetry
- Perhaps a bit lighter than the mass

Readed for the Standard Medel validity: needed for the Standard Model validity
	- to Planck scale (modulo theory assumptions) SSLIN *Mt* \$ 173.1 & 0.7 GeV $\overline{}$ band in $\overline{}$ Α*s*!*MZ*" \$ 0.1184 & 0.0007
	- \cdot m_H = 125 GeV \rightarrow oğr universe may lie on the boundary between instability and stability 1014 Instability $\epsilon \omega$ le in GeV
	- No need to panic: metastability n eans that the universe is unlikely to end tor lorrow Α*s*!*MZ*" \$ 0.1205 \overline{a} 1⊖

RGE scButterntriguing, nonetheles Siggs mass M_h in GeV 10^{20} 115 120 125 130 135

Spin/Parity

- Only elementary particle with **spin-0**
- Spin and parity determine angular distributions of decay products
	- Use γγ, ZZ and WW
- Don't forget, though, that the γγ observation implies
	- Does not originate from spin 1 : Landau-Yang theorem
	- Charge conjugation is +1 (assuming C and P separately conserved)
	- WW/ZZ channels disfavour CP odd hypothesis (can occur through loops)

[HIG-13-002](http://cms-results.web.cern.ch/cms-results/public-results/publications/HIG-13-002/index.html)

[HIGG-2013-17](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HIGG-2013-17/)

⁴¹
Spin/Parity Results 41
HIGG-2013-17, HIG-14-018

Strong evidence that the Higgs is 0+ as predicted by the Standard **Model**

Both ATLAS and CMS find that the observed Higgs boson is compatible with a standard CP-even

Higgs Width

- Direct measurements of the Higgs width are limited by the detector resolution to a few GeV (SM: a few MeV)
- Can do much better with indirect measurements using the ratio of the off-shell to on-shell cross-section
	- Currently constraint width to a few tens of GeV
	- But: brings in model assumptions

Higgs Production and Decays 43

Higgs Production

Higgs Decay

-100

20

44
Higgs Results ATLAS-CONF-2017-045, ATLAS-CONF-2017-041, HIG-16-021, ATLAS Results [ATLAS-CONF-2017-045](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2017-045/), [ATLAS-CONF-2017-041](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2017-041/), [HIG-16-021](http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/HIG-16-021), [HIG-16-043.](http://cms-results.web.cern.ch/cms-results/public-results/publications/HIG-16-043/index.html) [HIG-16-041,](http://cms-results.web.cern.ch/cms-results/public-results/publications/HIG-16-041/index.html) [HIG-16-044,](http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/HIG-16-044/index.html) [ATLAS-CONF-2017-043](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/CONFNOTES/ATLAS-CONF-2017-043/)

100

 m_{\parallel} [GeV]

60

40

80

 $m_{\tau\tau}$ (GeV)

0 50 100 150 200 250 300

Production and Decay Strengths

[JHEP08\(2016\)045](http://link.springer.com/article/10.1007/JHEP08(2016)045)

Production Decay

Figure 12: Best fit results for the production signal strengths for the combination of ATLAS and CMS data. Also shown are the results from each experiment. The error bars indicate the 1 (thick lines) and 2 (thin lines) intervals. The measurements of the global strength \sim F fiter accume SM for the other F Two independent fits: assume SM for the other

A6
Interpretation as Couplings

$$
-\frac{1}{2} \sum_{\nu}^{K_{\nu\nu}} \sum_{\nu}^{\nu} \frac{1}{22\%}
$$

Assumptions: No contributions to width from BSM particles (no decay to BSM particles) No contributions to loops from BSM particles

Generally good agreement with SM

Results for **fermions** are much weaker than for bosons

Observation of H→ττ

25 [HIG-16-043](http://cms-results.web.cern.ch/cms-results/public-results/publications/HIG-16-043/index.html)

47

- CMS paper last week as first single experiment observation of the H→ττ (ATLAS+CMS observation in Run-1 combination) \boldsymbol{s} first single experiment opservation of the \boldsymbol{v} \overline{C} distributions, and does not allow merging of the two figures. The normalization of the two figures of the normalization of the two figures. The normalization of the two figures of the two figures of the two figur servation in Run-1 combination) corresponds to the result of the signal field \sim is normalized to its best fit signal strength. The mass distributions for a constant \mathbf{r}
- Only channel so far to directly observed the coupling to fermions second dimension of the signal distributions are weighted according to *S*/(*S* + *B*), where *S* irectiv observed the coupling to fermions \blacksquare bution excluding the first and last bins. The "Others" background contribution includes events
- Two channels (by τ decay): lepton-hadron, hadron-hadron W boson-nadron, nadron-nadron accounts for all sources of background uncertainty, systematic as well as \mathcal{L}
- · Exploit both gluon-gluon fusion and VBF production percepted background is to get the signal expectation. The signal expectation is not signal \mathcal{L}
- Key elements: reconstructing the Higgs mass despite the presence of neutrinos and accurately estimating the Z→ττ background NSTDA (Thailand); TUBITAK and TAEK (Turkey); NASU and SFFR (Ukraine); STFC (United / estimating the

Evidence for H→bb

- H→bb is the most common decay (58%) but, due to the large backgrounds, it is very challenging
- Recent result from ATLAS provides the evidence for H→bb with an observed (expected) significance of 3.5σ (3.0σ)
- Use associated production with a W and Z boson
	- Leptonic decays provide trigger
	- Strongly reduce backgrounds
- Cross-check by measuring VZ production with Z→bb with an observed (expected) significance of 5.8σ (5.3σ)
- CMS combination of Run1+Run2 of 4.8σ

Novel High pT H→bb Search

- Select events with a large radius jet with $p_T > 450$ GeV
	- Typically accompanied by a jet radiated off the Higgs
- Validate with Z→bb observation (5.1σ)
- Early days for the Higgs: observed (expected) significance of 1.5σ (0.7σ)

[HIG-17-010](http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/HIG-17-010/index.html)

[HIGG-2016-10](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/HIGG-2016-10/)

H→μμ

- So far, we've focussed on coupling of the Higgs to third generation fermions
- Structure of the fermionic sector is far from trivial !
- $H \rightarrow \mu\mu$ will soon provide us with a means to probe the coupling of the Higgs to the second generation
- Higgs is easily identifiable via the two muons, but there is a background many orders of magnitude larger than the Higgs from $Z \rightarrow \mu\mu$ decays
- Current limit is 2.8 (2.9) x the SM
	- Will become very interesting with more data!

Does the Higgs Decay to new particles? 51

Constraints on rate of decays to particles that we cannot see

Summary: Coupling vs Mass

Good agreement but check the yaxis carefully

Overall conclusion: Generally very good agreement with the **SM**

Thins we don't know about the Higgs

- Direct evidence for the Higgs-top Yukawa coupling
- Other rare Higgs decays, e.g. [lepton flavour violation](http://cms-results.web.cern.ch/cms-results/public-results/preliminary-results/HIG-17-001/index.html)
- Confirm the Higgs self-interaction (HH production)
- Study the Higgs potential
	- Evolution from the early universe
	- Phase transition ? Connection to electroweak baryogengesis

Beginning of

Higgs

physics

Conclusion 54

Standard Model Total Production Cross Section Measurements Status: July 2017

