

A novel LHC dark matter search to test the Galactic Centre Excess [arXiv:1705.09670](https://arxiv.org/abs/1705.09670)

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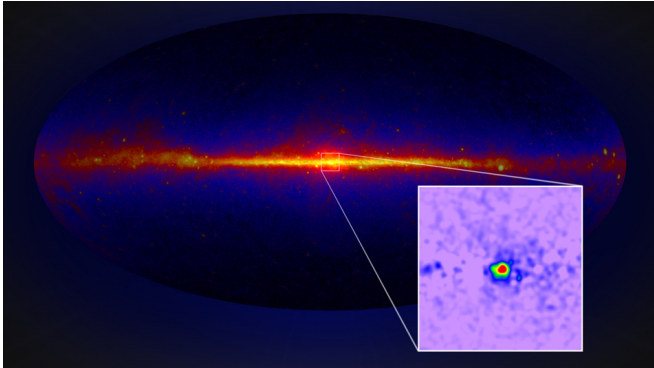
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SLAP meeting



Overview

- Simplified models of Dark Matter (DM).
- Gauge-invariant extension of pseudoscalar mediated models, featuring a Two Higgs Doublet Model (2HDM).
- A new LHC search to test parameter space that explains the Galactic Centre Excess (GCE).



Pseudoscalar Simplified Model

$$\mathcal{L} = \frac{m_{a_0}^2}{2} a_0^2 + m_\chi \bar{\chi} \chi + y_\chi a_0 \bar{\chi} i \gamma^5 \chi + \sum_q y_q a_0 \bar{q} i \gamma^5 q$$

χ is DM, a_0 the pseudoscalar and y_q are the standard model yukawa couplings.

Final term does not respect $SU(2)_L \times U(1)_Y$ gauge invariance.

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Introduce 2 Higgs Doublets $H_i = \left(\phi_i^+, (v_i + h_i + \eta_i)/\sqrt{2} \right)^T$

The vevs are related as $v_2/v_1 \equiv \tan\beta$

The spectrum contains a pseudoscalar $A_0 = \cos\beta \eta_2 - \sin\beta \eta_1$, a charged scalar $H^\pm = \cos\beta \phi_2^\pm - \sin\beta \phi_1^\pm$ and two neutral CP-even scalars $h = \cos\alpha h_2 - \sin\alpha h_1$, $H_0 = -\sin\alpha h_2 - \cos\alpha h_1$

h is the 125 GeV SM Higgs state in the limit $\beta - \alpha = \pi/2$.

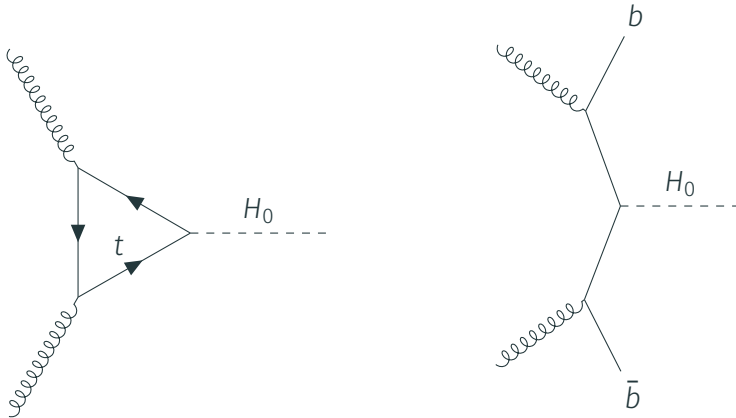
The pseudoscalar A_0 mixes with a_0 via

$$V_{\text{portal}} = i \kappa a_0 H_1^\dagger H_2 + \text{h.c.} \quad (1)$$

to give two mass eigenstates $a = c_\theta a_0 - s_\theta A_0$, $A = c_\theta A_0 + s_\theta a_0$, with $c_\theta \equiv \cos\theta$ and $s_\theta \equiv \sin\theta$.

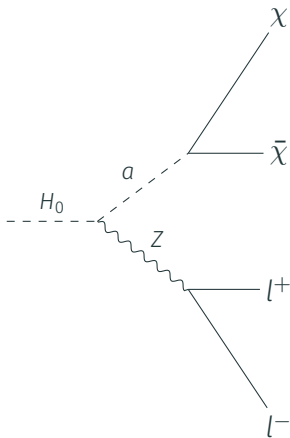
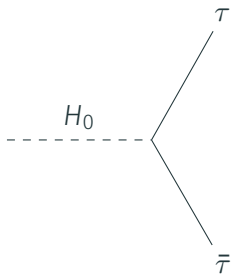
- Take $\cos(\beta - \alpha) = c_{\beta-\alpha} = 0$ to make h SM like
- Type-II 2HDM gives couplings to b quarks and τ leptons $\propto \tan \beta = t_\beta$. Coupling to t quark $\propto 1/t_\beta$
- Flavour constraints from $\bar{B} \rightarrow X_s \gamma$ and electroweak precision observables require $m_A, m_{H_0}, m_{H^\pm} \gtrsim 500$ GeV.
- Unitarity requires $m_A \leq 1.4$ TeV, $m_{H_0} \leq 1$ TeV.
- $\langle \sigma v \rangle \simeq 3 \times 10^{-26} \text{cm}^3/\text{s}$ gives the right relic density and fits the Galactic Centre Excess.
- Fix $m_\chi = 30$ GeV to be compatible with the GCE

Production mechanisms



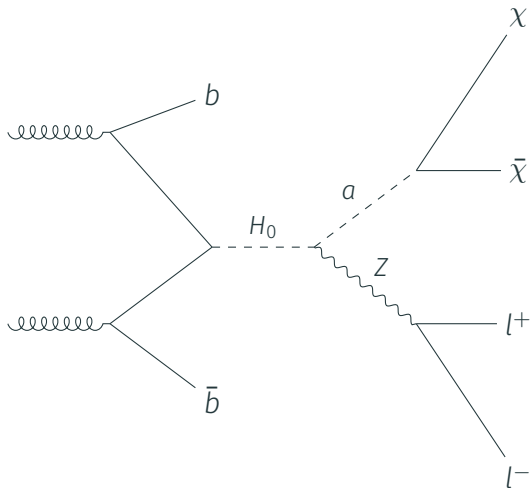
Production via gluon fusion (suppressed at high t_β) or $b\bar{b}$ associated (enhanced at high t_β)

Decay mechanisms



Decay into τ pairs or into mono- Z + missing energy

Our new search



Our new search

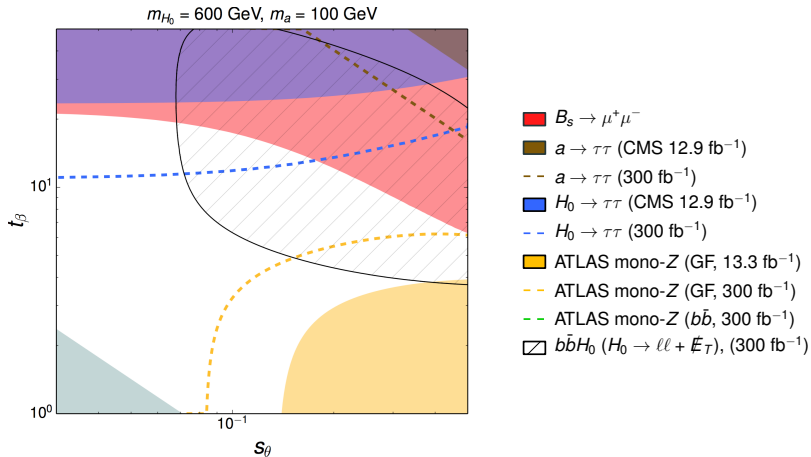
With a $b \bar{b} l^+ l^- + E_T$ final state the dominant background is from $t \bar{t}$ decaying leptonically.

We can take advantage of the fact our leptons come from a Z and impose $m_{\ell\ell} \in [76, 106]$.

Z and a are roughly back to back, so impose $|p_T^{\ell\ell} - \cancel{E}_T|/p_T^{\ell\ell} < 0.5$

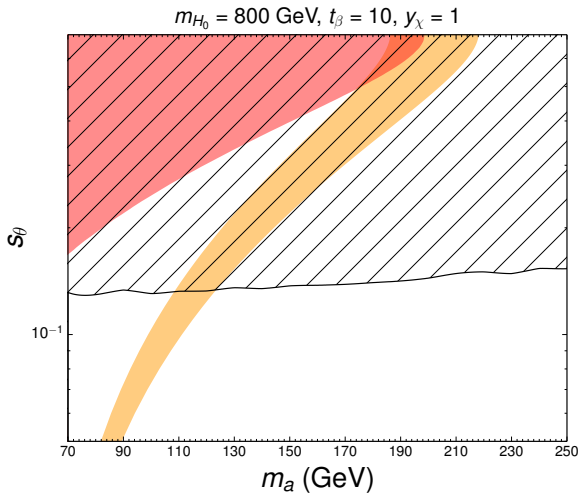
Harsh cuts of $\cancel{E}_T > 110$ GeV and $m_{T2} > 110$ GeV get the $t \bar{t}$ background to a controllable level.

Results



y_χ is fixed by the relic density. If $y_\chi > 4\pi$ the point is shaded grey.
The DM mass is $m_\chi = 30 \text{ GeV}$.

Results



Hatched region is our new search. Red: excluded by $B_s \rightarrow \mu^+ \mu^-$, orange band: relic density. The DM mass is $m_\chi = 30 \text{ GeV}$.

Conclusions

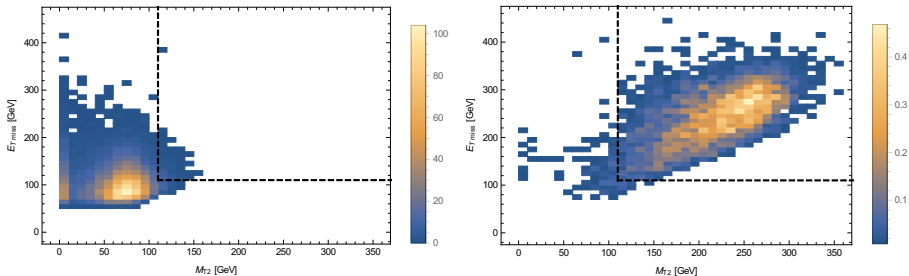
- Whether the GCE is dark matter or not, it is important to test its interpretation
- Consistent completions of simplified models are not so simple, and always offer richer phenomenology
- The implications of gauge invariance on the pseudoscalar model has opened up a new channel at the LHC

Dominated by s-channel a-exchange when $m_a \ll m_A$

$$\begin{aligned} \langle \sigma v \rangle = & \frac{y_\chi^2 m_\chi^2}{2\pi m_a^4} s_\theta^2 c_\theta^2 t_\beta^2 \left[\left(1 - \frac{4m_\chi^2}{m_a^2} \right)^2 + \frac{\Gamma_a^2}{m_a^2} \right]^{-1} \\ & \times \sum_f N_c \frac{m_f^2}{v^2} \sqrt{1 - \frac{m_f^2}{m_a^2}}. \end{aligned} \quad (2)$$

$\langle \sigma v \rangle \simeq 3 \times 10^{-26} \text{cm}^3/\text{s}$ gives the right relic density and fits the Galactic Centre Excess.

Fix $m_\chi = 30 \text{ GeV}$ to be compatible with the GCE



Requiring $\cancel{E}_T > 110$ GeV and $m_{T2} > 110$ GeV removes most of the background (left) while not affecting most of the signal (right)

We calculate m_{T2} as

$$m_{T2}^2 \equiv \min_{\vec{k}_T + \vec{q}_T = \vec{p}_T} \left\{ \max \left[m_T^2(\vec{p}_T^{\ell^+}, \vec{k}_T), m_T^2(\vec{p}_T^{\ell^-}, \vec{q}_T) \right] \right\} \quad (3)$$

where the minimisation is over all possible vectors \vec{k}_T and \vec{q}_T that satisfy $\vec{k}_T + \vec{q}_T = \vec{p}_T$

