

CKM and Tri-bimaximal MNS Matrices in a $SU(5) \times (d)T$ Model

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Introduction

- Neutrino Oscillation Parameters [Circa 2006]

$$U_{MNS} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\sin^2 \theta_{23} = 0.5 (0.38 - 0.64), \quad \sin^2 \theta_{13} = 0 (< 0.028) \quad \sin^2 \theta_{12} = 0.30 (0.25 - 0.34)$$

- Tri-bimaximal neutrino mixing:

$$U_{\text{TBM}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -\sqrt{1/6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -\sqrt{1/6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix} \quad \sin^2 \theta_{\text{atm,TBM}} = 1/2 \quad \sin \theta_{13,\text{TBM}} = 0.$$

$$\sin^2 \theta_{\odot,\text{TBM}} = 1/3$$

$$\tan^2 \theta_{\odot,\text{exp}} = 0.429$$

$$\tan^2 \theta_{\odot,\text{TBM}} = 1/2$$

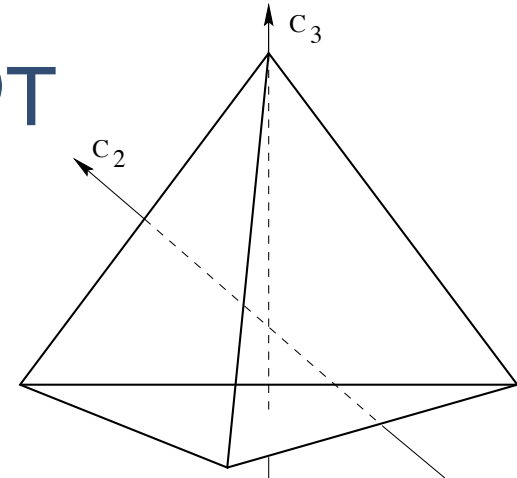
Tri-bimaximal Neutrino Mixing from A_4

- even permutations of 4 objects (invariance group of tetrahedron)
- generated by two basic permutations: (1234)

$$S=(4321) ; T=(2314)$$

- $4! / 2 = 12$ elements
- defining relations: $S^2 = T^3 = (ST)^3 = 1$
- four in-equivalent representations: $1, 1', 1'', 3$
- successfully give rise to tri-bimaximal leptonic mixing:
 - ★ lepton doublets ~ 3
 - ★ RH charged leptons: $1, 1', 1''$
- generalization to quark sector:
 - ★ no CKM mixing -- lack of doublet representations
 - ★ mass hierarchy -- need additional $U(1)$ symmetry

Group Theory of $(d)T$



- Double covering of A_4
- in-equivalent representations:

A_4 : 1, 1', 1'', 3
 other: 2, 2', 2''

→ 2 + 1 assignments for quarks

- generators:

$$S^2 = R, T^3 = 1, (ST)^3 = 1, R^2 = 1$$

R=1: 1, 1', 1'', 3
 R=-1: 2, 2', 2''

- product rules:

$$1^0 \equiv 1, 1^1 \equiv 1', 1^{-1} \equiv 1''$$

$$1^a \otimes r^b = r^b \otimes 1^a = r^{a+b} \quad \text{for } r = 1, 2 \quad a, b = 0, \pm 1$$

$$1^a \otimes 3 = 3 \otimes 1^a = 3$$

$$2^a \otimes 2^b = 3 \oplus 1^{a+b}$$

$$2^a \otimes 3 = 3 \otimes 2^a = 2 \oplus 2' \oplus 2''$$

$$3 \otimes 3 = 3 \oplus 3 \oplus 1 \oplus 1' \oplus 1''$$

The Model

- Symmetry: $SU(5) \times {}^{(d)}T$
- Particle Content $10(Q, u^c, e^c)_L \quad \bar{5}(d^c, \ell)_L$

	T_3	T_a	\bar{F}	H_5	H'_5	Δ_{45}	ϕ	ϕ'	ψ	ψ'	ζ	N	ξ	η
$SU(5)$	10	10	$\bar{5}$	5	$\bar{5}$	45	1	1	1	1	1	1	1	1
${}^{(d)}T$	1	2	3	1	1	1'	3	3	2'	2	1''	1'	3	1

- only top mass allowed at renormalizable level

- need to break ${}^{(d)}T$ to generate all other fermion masses

The Model

$H_5T_3T_a$	ψ', ψ $\psi\phi, \psi\phi', \psi'\phi, \psi'\phi', \psi'\zeta, \psi'N, \psi N$ $\psi^3, \psi\psi'^2, \psi\phi^2, \psi\phi'^2, \psi\phi\zeta, \psi\phi'\zeta, \psi'^3, \psi'\psi^2, \psi'\phi^2, \psi'\phi'^2, \psi'\phi\zeta, \psi'\phi'\zeta,$ $\psi\phi N, \psi\phi'N, \psi'\phi N, \psi'\phi'N$
	$\psi\xi, \psi'\xi, \psi\xi^2, \psi\xi\phi, \psi\xi\phi', \psi\xi\zeta, \psi'\xi^2, \psi'\xi\phi, \psi'\xi\phi', \psi'\xi\zeta, \psi\xi N, \psi'\xi N, \psi'\eta, \psi\phi\eta, \psi\phi'\eta, \psi\xi\eta,$ $\psi'\phi\eta, \psi'\phi'\eta, \psi'\xi\eta, \psi\eta, \psi\phi\eta, \psi\phi'\eta, \psi'\phi\eta, \psi'\phi'\eta, \psi\phi\eta, \psi\phi'\eta, \psi'\phi\eta, \psi'\phi'\eta$
$H_5T_aT_a$	ϕ, ϕ' $\phi'^2, \psi^2, \psi'^2, \phi\phi', \psi\psi'$ $\phi^3, \phi^2\zeta, \phi\zeta^2, \phi'^2\zeta, \phi'\zeta^2, \phi\phi'\zeta, \phi\phi'^2, \phi'\phi^2, \phi N^2, \phi'N^2, \phi'^2N, \phi\phi'N, \phi N\zeta, \phi'N\zeta$
	$\xi, \xi^2, \xi\zeta, \xi N, \xi\eta, \xi^2, \xi\phi, \xi\phi', \xi^3, \xi^2\zeta, \xi^2\eta, \xi^2\zeta, \xi N\zeta, \xi N\eta, \xi\zeta\eta, \xi\phi^2, \xi\phi'^2, \xi\phi\phi',$ $\xi^2\phi, \xi^2\phi', \xi\phi N, \xi\phi\eta, \xi\phi\zeta, \xi\phi'N, \xi\phi'\eta, \xi\phi'\zeta, \phi^2\eta, \phi\eta^2, \phi\eta N, \phi\eta\zeta, \phi'\eta^2, \phi'\eta N,$ $\phi'\eta\zeta, \phi\eta, \phi'\eta, \xi N^2, \xi\eta^2, \xi\zeta^2$
$H_5^L\bar{F}T_3$	ϕ, ϕ' $\psi^2, \phi^2, \phi'^2, \phi'\phi, \psi'^2, \psi\psi', \phi'\zeta, \phi'N, \phi N$ $\phi^3, \phi'^3, \phi^2\phi', \phi\phi'^2, \phi\zeta^2, \phi'\zeta^2, \phi\psi^2, \phi'\psi'^2, \zeta\psi^2, \zeta\psi'^2, \phi'\psi^2, \phi\psi^2,$ $\phi N^2, \phi'N^2, \phi N\zeta, \phi'N\zeta, N\psi^2, \zeta\psi^2, \zeta\psi\psi', N\psi\psi'$
	$\xi, \xi^2, \xi N, \xi\zeta, \xi\eta, \xi\phi, \xi\phi', \xi^3, \xi^2N, \xi^2\zeta, \xi^2\eta, \xi^2\phi, \xi^2\phi', \xi\phi^2,$ $\xi\phi'^2, \xi\phi\phi', \xi\phi N, \xi\phi\zeta, \xi\phi\eta, \xi\phi'N, \xi\phi'\zeta, \xi\phi'\eta, \phi\eta^2, \phi\eta N, \phi\eta\zeta, \phi'\eta^2, \phi'\eta N, \phi'\eta\zeta, \eta\psi^2,$ $\eta\psi'^2, \phi\eta, \phi\eta N, \phi\eta\zeta, \phi'\eta^2, \phi'\eta N, \eta\psi\psi'$
$H_5^L\bar{F}T_a$	ψ, ψ' $\psi\phi', \psi'\phi, \psi'\phi', \phi\psi$ $\psi\phi^2, \psi\phi\zeta, \psi'\phi\zeta, \psi\phi'^2, \psi'\phi'^2, \psi\phi\phi', \psi'\phi\phi', \psi\phi'\zeta, \psi'\phi'\zeta, \psi\phi N, \psi'\phi N, \psi\phi'N, \psi'\phi'N$
	$\psi\xi, \psi'\xi, \psi\xi^2, \psi'\xi^2, \psi\xi\phi, \psi\xi\phi', \psi'\xi\phi, \psi'\xi\phi',$ $\psi\xi N, \psi\xi\eta, \psi\xi\zeta, \psi'\xi\zeta, \psi'\xi\eta, \psi'\xi N, \psi\phi\eta, \psi'\phi\eta, \psi'\phi'\eta, \psi\phi'\eta, \psi'\phi'\eta, \psi\phi\eta, \psi'\phi\eta$

The Model

- Symmetry: $SU(5) \times {}^{(d)}T$
- Particle Content $10(Q, u^c, e^c)_L \quad \bar{5}(d^c, \ell)_L$

	T_3	T_a	\bar{F}	H_5	H'_5	Δ_{45}	ϕ	ϕ'	ψ	ψ'	ζ	N	ξ	η
$SU(5)$	10	10	$\bar{5}$	5	$\bar{5}$	45	1	1	1	1	1	1	1	1
${}^{(d)}T$	1	2	3	1	1	1'	3	3	2'	2	1''	1'	3	1
Z_{12}	ω^5	ω^2	ω^5	ω^2	ω^2	ω^5	ω^3	ω^2	ω^6	ω^9	ω^9	ω^3	ω^{10}	ω^{10}
Z'_{12}	ω	ω^4	ω^8	ω^{10}	ω^{10}	ω^3	ω^3	ω^6	ω^7	ω^8	ω^2	ω^{11}	1	1

$$\omega = e^{i\pi/6}.$$

- additional $Z_{12} \times Z'_{12}$ symmetry:
 - ★ predictive model: only 9 operators allowed up to at least dim-7
 - ★ vacuum misalignment: neutrino sector vs charged fermion sector
 - ★ mass hierarchy: lighter generation masses allowed only at higher dim

The Model

- Abelian subgroups of ${}^{(d)}T$:

$$\begin{aligned} Z_3 : & G_T \\ Z_4 : & G_{TST^2} \end{aligned}$$

$$T = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix}$$

$$\omega = e^{2\pi i/3}$$

$$TST^2 = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}$$

- ${}^{(d)}T$ breaking:

$${}^{(d)}T \longrightarrow G_{TST^2} : \quad \langle \xi \rangle = \xi_0 \Lambda \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}, \quad \langle \phi' \rangle = \phi'_0 \Lambda \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}, \quad \longrightarrow$$

tri-bimaximal mixing
in neutrino sector

$${}^{(d)}T \longrightarrow G_T : \quad \langle \phi \rangle = \phi_0 \Lambda \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad \langle \psi \rangle = \psi_0 \Lambda \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$${}^{(d)}T \longrightarrow \text{nothing} : \quad \langle \psi' \rangle = \psi'_0 \Lambda \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$${}^{(d)}T \longrightarrow G_S : \quad \langle \zeta \rangle = \zeta_0, \quad \langle N \rangle = N_0$$

$${}^{(d)}T - \text{invariant} : \quad \langle \eta \rangle = u$$

charged fermion
sector

The Model

- Lagrangian: only 9 operators allowed!!

$$\mathcal{L}_{\text{Yuk}} = \mathcal{L}_{\text{TT}} + \mathcal{L}_{\text{TF}} + \mathcal{L}_{\text{FF}}$$

$$\mathcal{L}_{\text{TT}} = y_t H_5 T_3 T_3 + \frac{1}{\Lambda^2} y_{ts} H_5 T_3 T_a \psi \zeta + \frac{1}{\Lambda^2} y_c H_5 T_a T_a \phi^2 + \frac{1}{\Lambda^3} y_u H_5 T_a T_a \phi'^3$$

$$\mathcal{L}_{\text{TF}} = \frac{1}{\Lambda^2} y_b H'_5 \bar{F} T_3 \phi \zeta + \frac{1}{\Lambda^3} \left[y_s \Delta_{45} \bar{F} T_a \phi \psi N + y_d H'_5 \bar{F} T_a \phi^2 \psi' \right]$$

$$\mathcal{L}_{\text{FF}} = \frac{1}{M_x \Lambda} \left[\lambda_1 H_5 H_5 \bar{F} \bar{F} \xi + \lambda_2 H_5 H_5 \bar{F} \bar{F} \eta \right],$$

Neutrino Sector

- **Operators:** $\mathcal{L}_{\text{FF}} = \frac{1}{M_x \Lambda} \left[\lambda_1 H_5 H_5 \bar{F} \bar{F} \xi + \lambda_2 H_5 H_5 \bar{F} \bar{F} \eta \right]$

- **Symmetry breaking:**

$${}^{(d)}T \longrightarrow G_{\text{TST}^2} : \quad \langle \xi \rangle = \xi_0 \Lambda \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} \quad {}^{(d)}T - \text{invariant} : \quad \langle \eta \rangle = u$$

- **Resulting mass matrix:**

$$M_\nu = \frac{\lambda v^2}{M_x} \begin{pmatrix} 2\xi_0 + u & -\xi_0 & -\xi_0 \\ -\xi_0 & 2\xi_0 & u - \xi_0 \\ -\xi_0 & u - \xi_0 & 2\xi_0 \end{pmatrix}$$

$$V_\nu^T M_\nu V_\nu = \text{diag}(u + 3\xi_0, u, -u + 3\xi_0) \frac{v_u^2}{M_x} \quad U_{\text{TBM}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -\sqrt{1/6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -\sqrt{1/6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

Up Quark Sector

- **Operators:** $\mathcal{L}_{\text{TT}} = y_t H_5 T_3 T_3 + \frac{1}{\Lambda^2} y_{ts} H_5 T_3 T_a \psi \zeta + \frac{1}{\Lambda^2} y_c H_5 T_a T_a \phi^2 + \frac{1}{\Lambda^3} y_u H_5 T_a T_a \phi'^3$
- top mass: allowed by $^{(d)}T$
- lighter family acquire masses thru operators with higher dimensionality

➡ dynamical origin of mass hierarchy

- **symmetry breaking:**

$$^{(d)}T \longrightarrow G_T : \quad \langle \phi \rangle = \phi_0 \Lambda \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad \langle \psi \rangle = \psi_0 \Lambda \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$^{(d)}T \longrightarrow G_{\text{TST}^2} : \quad \langle \phi' \rangle = \phi'_0 \Lambda \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$

- **Mass matrix:**

$$M_u = \begin{pmatrix} i\phi_0'^3 & \frac{1-i}{2}\phi_0'^3 & 0 \\ \frac{1-i}{2}\phi_0'^3 & \phi_0'^3 + (1 - \frac{i}{2})\phi_0^2 & y'\psi_0\zeta_0 \\ 0 & y'\psi_0\zeta_0 & 1 \end{pmatrix} y_t v_u$$

Down Quark Sector

- operators: $\mathcal{L}_{\text{TF}} = \frac{1}{\Lambda^2} y_b H'_5 \bar{F} T_3 \phi \zeta + \frac{1}{\Lambda^3} \left[y_s \Delta_{45} \bar{F} T_a \phi \psi N + y_d H'_5 \bar{F} T_a \phi^2 \psi' \right]$
- generation of b-quark mass: breaking of $^{(d)}T$: dynamical origin for hierarchy between m_b and m_t
- lighter family acquire masses thru operators with higher dimensionality

➡ dynamical origin of mass hierarchy

- symmetry breaking:

$$^{(d)}T \longrightarrow G_T : \quad \langle \phi \rangle = \phi_0 \Lambda \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad \langle \psi \rangle = \psi_0 \Lambda \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad ^{(d)}T \longrightarrow \text{nothing} : \quad \langle \psi' \rangle = \psi'_0 \Lambda \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

- mass matrix:

$$M_d = \begin{pmatrix} 0 & (1+i)\phi_0\psi'_0 & 0 \\ -(1-i)\phi_0\psi'_0 & \psi_0 N_0 & 0 \\ \phi_0\psi'_0 & \phi_0\psi'_0 & \zeta_0 \end{pmatrix} y_b v_d \phi_0, \quad M_e = \begin{pmatrix} 0 & -(1-i)\phi_0\psi'_0 & \phi_0\psi'_0 \\ (1+i)\phi_0\psi'_0 & -3\psi_0 N_0 & \phi_0\psi'_0 \\ 0 & 0 & \zeta_0 \end{pmatrix} y_b v_e \phi_0$$

- consider 2nd, 3rd families only: TBM exact

- Georgi-Jarlskog relations: $m_d \simeq 3m_e \quad m_\mu \simeq 3m_s \rightarrow$ corrections to TBM

Quark and Lepton Mixing Matrices

- CKM mixing matrix:

$$M_u = \begin{pmatrix} i\phi_0^3 & \frac{1-i}{2}\phi_0^3 & 0 \\ \frac{1-i}{2}\phi_0^3 & \phi_0^3 + (1-i)\phi_0^2 & y'\psi_0\zeta_0 \\ 0 & y'\psi_0\zeta_0 & 1 \end{pmatrix} y_t v_u \quad M_d = \begin{pmatrix} 0 & (1+i)\phi_0\psi'_0 & 0 \\ -(1-i)\phi_0\psi'_0 & \psi_0 N_0 & 0 \\ \phi_0\psi'_0 & \phi_0\psi'_0 & \zeta_0 \end{pmatrix} y_b v_d \phi_0,$$

$\theta_c \simeq |\sqrt{m_d/m_s} - e^{i\alpha}\sqrt{m_u/m_c}| \sim \sqrt{m_d/m_s},$

- MNS matrix:

$$M_e = \begin{pmatrix} 0 & -(1-i)\phi_0\psi'_0 & \phi_0\psi'_0 \\ (1+i)\phi_0\psi'_0 & -3\psi_0 N_0 & \phi_0\psi'_0 \\ 0 & 0 & \zeta_0 \end{pmatrix} y_b v_d \phi_0 \longrightarrow \theta_{12}^e \simeq \sqrt{\frac{m_e}{m_\mu}} \simeq \frac{1}{3}\sqrt{\frac{m_d}{m_s}} \sim \frac{1}{3}\theta_c$$

$$U_{\text{MNS}} = V_{e,L}^\dagger U_{\text{TBM}} = \begin{pmatrix} 1 & -\theta_c/3 & * \\ \theta_c/3 & 1 & * \\ * & * & 1 \end{pmatrix} \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -\sqrt{1/6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -\sqrt{1/6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$

$$\tan^2 \theta_\odot \simeq \tan^2 \theta_{\odot, \text{TBM}} - \frac{1}{2}\theta_c \cos \beta$$

$$\theta_{13} \simeq \theta_c/3\sqrt{2}$$

leptonic CPV

Numerical Results

- **Experimentally:** $m_u : m_c : m_t = \epsilon_u^2 : \epsilon_u : 1$, $m_d : m_s : m_b = \epsilon_d^2 : \epsilon_d : 1$
 $\epsilon_u \simeq (1/200) = 0.005$ $\epsilon_d \simeq (1/20) = 0.05$.

- **Model Parameters:**

$$M_u = \begin{pmatrix} ig & \frac{1-i}{2}g & 0 \\ \frac{1-i}{2}g & g+h & k \\ 0 & k & 1 \end{pmatrix} y_t v_u$$

$$\frac{M_d}{y_b v_d \phi_0 \zeta_0} = \begin{pmatrix} 0 & (1+i)b & 0 \\ -(1-i)b & c & 0 \\ b & b & 1 \end{pmatrix}$$

$$k \equiv y' \psi_0 \zeta_0 = -0.032$$

$$h \equiv \psi_0^2 = 0.0053$$

$$g \equiv \phi_0^3 = -2.25 \times 10^{-5}$$

7 parameters in
charged fermion
sector

$$y_b \phi_0 \zeta_0 \simeq m_b / m_t \simeq (0.011)$$

$$c \equiv \psi_0 N_0 / \zeta_0 = 0.0474$$

$$b \equiv \phi_0 \psi'_0 / \zeta_0 = 0.00789$$

- **Mixing Matrices:**

$$|V_{\text{CKM}}| = \begin{pmatrix} 0.976 & 0.217 & 0.00778 \\ 0.216 & 0.975 & 0.040 \\ 0.015 & 0.0378 & 0.999 \end{pmatrix}$$

$$|U_{\text{MNS}}| = |V_{e,L}^\dagger U_{\text{TBM}}| = \begin{pmatrix} 0.838 & 0.545 & 0.0550 \\ 0.364 & 0.608 & 0.706 \\ 0.409 & 0.578 & 0.706 \end{pmatrix}$$

- **neutrino masses:**

cos (beta) = 2/3 : best fit values

$$u = -1.87 \times 10^{-2}, \quad \xi_0 = 1.15 \times 10^{-2}, \quad M_x \sim 10^{14} \text{ GeV}$$

2 parameters in
neutrino sector

Conclusions

- $SU(5) \times {}^{(d)}T$ symmetry: tri-bimaximal lepton mixing & realistic CKM matrix
 - $Z_{12} \times Z_{12}'$ symmetry: only 9 operators present (only 9 parameters in Yukawa sector)
 - ★ forbid proton decay
 - ★ likely linked to orbifold compactification
 - dynamical origin of mass hierarchy (including m_b vs m_t)
 - interesting sum rules:

$$\tan^2 \theta_{\odot} \simeq \tan^2 \theta_{\odot, \text{TBM}} - \frac{1}{2} \theta_c \cos \beta$$

$\theta_{13} \simeq \theta_c / 3\sqrt{2} \sim 0.05$
- could give right amount to account for discrepancy bt exp best fit value and TBM prediction