F-term uplifting via consistent D-terms

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The Setup: SUGRA

Our theoretical framework is low energy effective $\mathcal{N} = 1$, $D = 4$ supergravity. The object of interest to us is the scalar potential:

$$V = e^K \left(K_{i\bar{j}} F^i F^{\bar{j}} - 3|W|^2 \right) + \frac{1}{2} \frac{8\pi^2}{Re(T)} |D|^2$$
$$= V_F + V_W + V_D$$

where $F_i = \frac{\partial W}{\partial \phi_i} + \frac{\partial K}{\partial \phi_i} W$ and $D = iK_i X^i$. X^i are Killing vectors, given by $\delta\phi_i = X^i \epsilon$.
For a $U(1)$ symmetry under which the fields $\phi_i \to e^{iq_i\epsilon}\phi_i$ we recover, for a canonical Kähler potential, the usual D-term potential:

$$V_D = g^2 \left(\sum_i q_i |\phi_i|^2 \right)$$

String inspired framework

The low energy effective SUGRA is an approximation to the KKLT class of models in which stringy effects are assumed to stabilise the shape of the extra dimensions, but not their size. The final modulus is stabilised using

$$W_{KKLT} = Ae^{-T/N} + W_0$$

at a supersymmetric point with non-zero $\langle W \rangle$.
Unfortunately this implies $\langle V \rangle < 0$.

Uplifting via D-terms

Clearly more work is required to break SUSY and lift the potential to positive values.
We consider the effects of gauging the shift of $T \to T + i\delta\Lambda$. This has two immediate effects:

- e^{-T} is no longer gauge invariant
- $U(1)$ D-terms appear in the potential

To restore gauge invariance we introduce additional fields M_1 and M_2 which transform like $M \to e^{-i\delta\Lambda}M$.

Racetrack

To allow tuning of the cosmological constant we introduce a second exponential:

$$W_{Race.} = A_1 N_1 \left(\frac{e^{-T}}{M_1} \right)^{\frac{1}{N_1}} - A_2 N_2 \left(\frac{e^{-T}}{M_2} \right)^{\frac{1}{N_2}} + W_0,$$

where W_0, $A_1 N_1$ and $A_2 N_1$ are constants in principle given by the fundamental theory, but in practice treated as input parameters. M_1 and M_2 are the determinants of the quark bi-linears, which are treated as independent fields. Finally N_1 and N_2 are the number of colours in the two groups.

Solving the Racetrack

In addition to the superpotential we need a Kähler potential given by

$$K = -3\ln\left(T + \bar{T}\right) + |M_1|^2 + |M_2|^2$$

Then the racetrack superpotential, in the absence of D-terms, stabilises T. It is sufficient to consider only

$$V = e^K\left(K_{i\bar{j}}F^i F^{\bar{j}}\right)$$

when stabilising t. The solution to $F_T = 0$ gives, in the limit of large t,

$$t \sim \ln\left(\frac{|A_1|x_2^{1/N_2}}{|A_2|x_1^{1/N_1}}\right)\frac{N_1 N_2}{N_2 - N_1}$$

where $T = t + ia$, $M_1 = x_1 e^{iN_1\phi_1}$ and $M_2 = x_2 e^{iN_2\phi_2}$.

Indirect Lifting

V_D has two natural values, M_P^4 and zero. We now show that a D-term that minimises to give $V_D = 0$ can lift $V_F + V_W$, albeit indirectly.

We add one chiral field C that appears with the opposite sign charge to M_i, $C \rightarrow e^{iq\delta\Lambda} C$, hence generating the following D-term:

$$V_D = \frac{\pi^2 \delta^2}{t} \left(\frac{3}{2t} + x_1^2 + x_2^2 - qc^2 \right)^2$$

If we allow a natural $\delta \sim 1 - 10^{-2}$ then V_D dominates the shape of the potential, $qc^2 = \frac{3}{2t} + x_1^2 + x_2^2$ is enforced and a large mass term is introduced for $\frac{3}{2t} + x_1^2 + x_2^2 - qc^2$, splitting the moduli masses.

Constraint equation

It is clear that, if c has no potential, the constraint is satisfied trivially: c aligns to set $D = 0$. However, even if $\frac{\partial W}{\partial C} = 0$, a potential is generated for c via the Kähler derivatives:

$$V = e^K c^2 |W|^2$$

This introduces an energetic preference for small c. However

$$\frac{\partial V}{\partial c} = -q\frac{4c\pi^2\delta^2}{t}\left(\frac{3}{2t} + x_1^2 + x_2^2 - qc^2\right) + 2c|W|^2 + \frac{\partial e^K}{\partial c}c^2|W|^2$$

is dominated by the derivative of V_D. Enforcing $V_D = 0$ as a constraint, is well approximated by $c^2 = \frac{1}{q}\left(\frac{3}{2t} + x_1^2 + x_2^2\right)$ and is true in the limit $\delta \to \infty$.

Minimisation

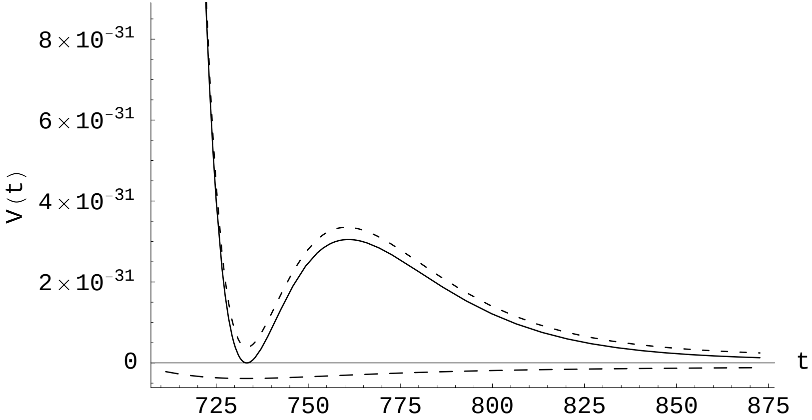
The procedure for finding a minimum is as follows:

- D-term is set to zero
- c^2 is replaced by the constraint
- Constrained potential is minimised
- Solution is input into unconstrained potential

The unconstrained minimum is, by definition, the lowest value of the potential locally. Because the constraint equation has its own potential, with $|W|^2$ acting as a mass term, a balance must be achieved by moving to a higher minimum.
As a result the potential is lifted and the cosmological constant can be tuned with arbitrary precision by varying the input parameters.

Minimum of the potential

This plot is of the constrained potential, in this region the *D*-term is parabolic.



Input parameter values

$\lvert A_1 \rvert$	2.62
$\lvert A_2 \rvert$	0.3
N_1	25
N_2	27
δ	1
W_0	2.13×10^{-12}
q	$1/16$

Minimisation results

$m_{3/2}$	478.142 GeV	m_t	1.630×10^5 GeV
$V_0^{1/4}$	0	m_a	1.646×10^5 GeV
$V_F^{1/4}$	3.41239×10^{10} GeV	m_{x_1}	3.141×10^2 GeV
$V_D^{1/4}$	2292.40 GeV	m_{x_2}	2.427×10^{17} GeV
t	733.325	m_c	7.853×10^2 GeV
x_1	0.173619	m_{ϕ_1}	6.000×10^2 GeV
x_2	0.237810	m_{ϕ_2}	0
ϕ_1	π	m_V	8.56×10^{17} GeV
ϕ_2	π		

Conclusions

- All fields are stabilised with a realistic gravitino mass and vanishing or slightly positive cosmological constant
- Cancellable D-terms introduce a constraint indirectly lifting the potential and allowing for a natural δ