



Bounds on R -Parity Violation from Rare Decay Data

Ben O'Leary

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Outline

1 Introduction to RPV



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2 Implications of RPV



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- 3 Bounds on RPV



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- 4 Summary



R -Parity conserving “Yukawa” superpotential



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$$\mu H_u H_d + Y_{jk}^e H_d L_j L_{kR}^c$$

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Implications of R -Parity Violation

Baryon and lepton number violation:

- Single sparticle production
- Lightest Supersymmetric Partner is unstable
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How To Bound R -Parity Violation

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- HERA (sparticle-mass-dependent bounds on $\lambda'_{ijk} \lambda'_{lmn}$, or more usually λ' -dependent bounds on sparticle masses)



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Atomic parity violation data, $\mu \rightarrow e$ conversion in atoms, ...



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- decays of τ leptons and muons into three charged leptons

Bounds from Rare Decay Data

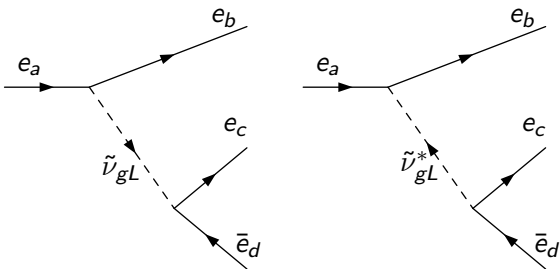
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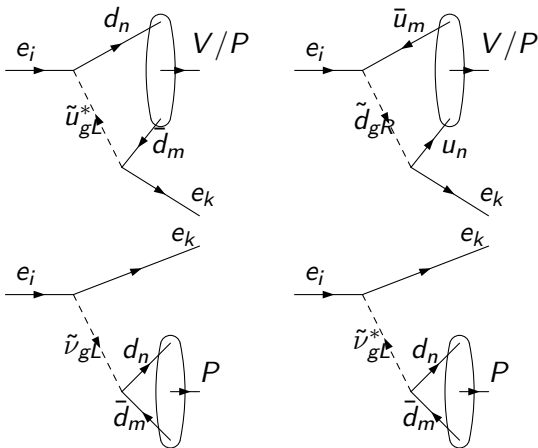
All λ''_{ijk} set to 0 (through baryon triality):

- prevents proton decay while allowing for neutrino mass
- lepton-number parity is not discrete gauge anomaly-free
- even R -Parity does not prevent proton decay through dimension-5 operators

Upper bounds on RPV



Upper bounds on RPV



Upper bounds on RPV

Coupling combination	Bound	Decay
$\lambda_{g23}\lambda'_{g12}$	$1.0 \times 10^{-3} [\tilde{\nu}_{gL}]^2$	$\tau \rightarrow \mu K_S$
$\lambda_{g32}\lambda'_{g21}$		
$\lambda'_{1g1}\lambda'_{3g2}$	$2.3 \times 10^{-3} [\tilde{u}_{gL}]^2$	$\tau \rightarrow e K_S$
$\lambda'_{1g2}\lambda'_{2g2}$	$1.5 \times 10^{+2} [\tilde{u}_{gL}]^2$	$\eta \rightarrow \mu \bar{e} + e \bar{\mu}$
$\lambda'_{1g2}\lambda'_{3g2} \dagger^{(-)}$	$1.2 \times 10^{-3} [\tilde{u}_{gL}]^2$	$\tau \rightarrow e \eta$
$\lambda'_{1g2}\lambda'_{3g2}$	$3.4 \times 10^{-3} [\tilde{u}_{gL}]^2$	$\tau \rightarrow e \phi$
$\lambda'_{2g2}\lambda'_{3g2} \dagger^{(-)}$	$1.6 \times 10^{-3} [\tilde{u}_{gL}]^2$	$\tau \rightarrow \mu \eta$

Table: Example coupling combinations which had no bounds previously.

Upper bounds on RPV

Coupling combination	Bound	Bound
$\lambda_{g13} \lambda'_{g22},$ $\lambda_{g31} \lambda'_{g22}$	$4.6 \times 10^{-4} [\tilde{\nu}_{gL}]^2$ $\tau \rightarrow e\eta$	$1.6 \times 10^{-2} [\tilde{\nu}_{gL}]^2$ $\tau \rightarrow e\eta$
$\lambda_{g13} \lambda'_{g21},$ $\lambda_{g31} \lambda'_{g12}$	$9.7 \times 10^{-4} [\tilde{\nu}_{gL}]^2$ $\tau \rightarrow eK_S$	$8.5 \times 10^{-2} [\tilde{\nu}_{gL}]^2$ $\tau \rightarrow eK^0$

Table: Example coupling combinations which have improved by a factor of 30 or more compared to those published previously.

Not all bounds were improved, e.g.

Coupling combination	Bound	Bound
$\lambda_{121} \lambda'_{111}$	$1.2 \times 10^{-2} [\tilde{\nu}_{1L}]^2$ $\pi^0 \rightarrow e \bar{\mu}$	$2.1 \times 10^{-8} [\tilde{\nu}_{1L}]^2$ $\mu \rightarrow e$ in ^{48}Ti
$\lambda_{211} \lambda'_{223}$	$2.3 \times 10^{-4} [\tilde{\nu}_{2L}]^2$ $B_s^0 \rightarrow e \bar{e}$	$1.4 \times 10^{-4} [\tilde{\nu}_{2L}]^2$ $B_d^0 \rightarrow K^0 e \bar{e}$



Summary and Outlook





R -Parity bans some awkward terms in the MSSM.

Allowing the R -Parity-violating terms can provide a mechanism for neutrino mass without adding a right-handed neutrino field.

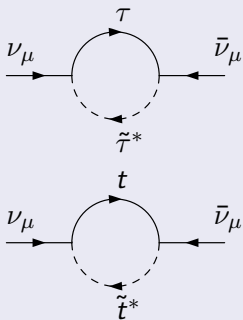
Allowing these terms also provides a variety of interesting new flavor physics, though severely constrained by experiment.

No strict lower bound exists for these couplings, but upper bounds could rule out neutrino mass from RPV; most upper bounds are still about an order of magnitude above the value consistent with neutrino masses.

Lower bounds

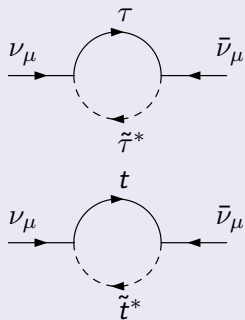
-  P. Fayet, Nucl. Phys. B **90** (1975) 104.
-  R. E. Marshak and R. N. Mohapatra, Phys. Lett. B **91** (1980) 222.
-  R. Barbier *et al.*, Phys. Rept. **420** (2005) 1 [arXiv:hep-ph/0406039].
-  H. K. Dreiner, M. Kramer and B. O'Leary, Phys. Rev. D **75** (2007) 114016 [arXiv:hep-ph/0612278].

Are there lower bounds on these couplings?



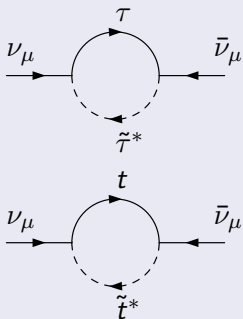
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 $|\lambda_{233}|^2 / m_{\tilde{\tau}_L}^2 \approx (4\pi)^2 m_{\nu_\mu} / (\mu \tan(\beta) m_\tau^2)$

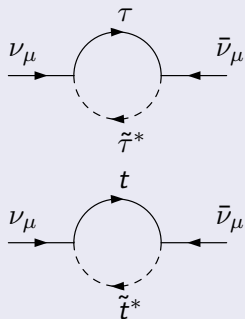
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Assuming $\mu \approx m_{\tilde{\tau}_L} \approx 100$ GeV, $\tan\beta \approx 30$,
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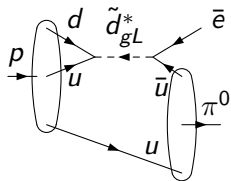
$$|\lambda_{233}|^2 \approx 10^{-5} (m_{\tilde{\tau}_L} / 100 \text{ GeV})^2$$

Proton decay

All λ 's and λ ''s present and of order 1, with sparticle masses of order 1 TeV, leads to proton decay with a disappointingly short lifetime ($\sim 10^{50}$ times too short (R. Barbier *et al.*, Phys. Rept. **420** (2005) 1 [arXiv:hep-ph/0406039])).

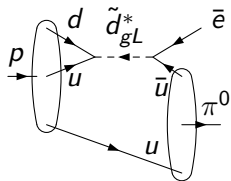
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However, it is sufficient to enforce either lepton number conservation (though this is not anomaly-free) by setting $\lambda_{ijk} = \lambda'_{ijk} = 0$, or baryon number conservation by setting $\lambda''_{ijk} = 0$.



How Collider Bounds Compare to Decay Bounds

The OPAL collaboration published upper bounds on lepton flavor–violating two–to–two scatterings from the LEP2 run at $\sqrt{s} \sim 200$ GeV (G. Abbiendi *et al.* [OPAL Collaboration], Phys. Lett. B **519** (2001) 23 [arXiv:hep-ex/0109011]).

These can be translated into upper bounds on λ – λ couplings as a function of the sneutrino mass.

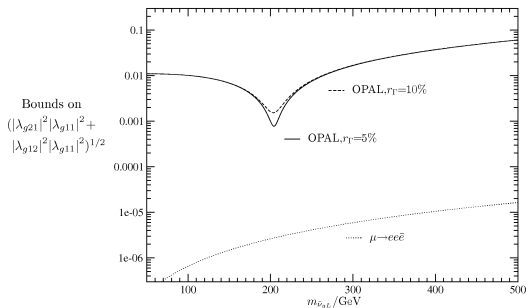


Figure: The bounds on $(|\lambda_{g21}|^2|\lambda_{g11}|^2 + |\lambda_{g12}|^2|\lambda_{g11}|^2)^{1/2}$ as a function of $m_{\tilde{\nu}_{gL}}$ from the search for $e\bar{e} \rightarrow \mu\bar{e}/e\bar{\mu}$ by OPAL in the range $200 \text{ GeV} \leq \sqrt{s} \leq 209 \text{ GeV}$.

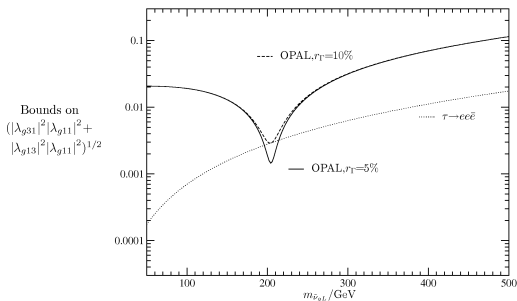


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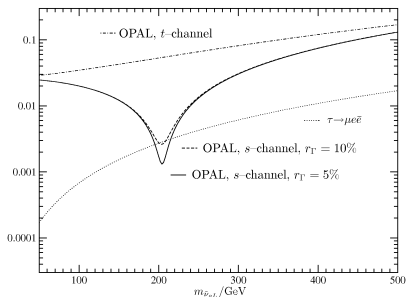


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“Right-handed” sneutrino vacuum expectation value:

$$Y_{jk}^\nu H_u L_{jL} \nu_{kR}^c \rightarrow m'_{jk} e_j^T \tilde{h}_u^+$$