

SUSY Parameters Determination with ATLAS

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Abstract. The plan for mass and spin measurement of SUSY particles with the ATLAS detector is presented. The measurements of kinematical distributions, such as edges in the invariant mass of leptons and jets, could be used to constrain the model of SUSY that may be discovered at the LHC. Examples from a few points in the mSUGRA scenario are provided with an emphasis on measurements that can be conducted within the first few years of data taking.

PACS. 11.30.Pb Supersymmetry – 12.60.Jv Supersymmetric models

1 Introduction

Discovering Supersymmetry (SUSY) is one of the motivations for building the Large Hadron Collider (LHC) which is scheduled to take data in 2008. ATLAS (A Torodial LHC ApparatuS) is one of the two general purpose experiments at the LHC. Currently the ATLAS experiment is revisiting its SUSY studies in the framework of the CSC (“Computing System Commissioning”). For this exercise event samples for a set of SUSY benchmark points were produced with the detailed simulation of the ATLAS detector. The aim of the work is to better understand the impact of realistic experimental conditions including trigger efficiencies and imperfect calibration/alignment on the SUSY potential of the experiment.

If SUSY is discovered by inclusive searches, next step will be to measure the masses and the spins of the SUSY particles through the analysis of exclusive decay channels. In this note SUSY mass and spin measurement techniques and methods for determining underlying SUSY model parameters are described within the mSUGRA framework.

2 mSUGRA Framework

The minimal SUSY extension of the SM (MSSM) brings 105 additional free parameters into the theory thus making a systematic study of the full parameter space difficult. A specific well-motivated model framework is usually assumed in which generic signatures can be studied. In the minimal Supergravity (mSUGRA) framework, SUSY is broken by the gravitational interactions and the masses and couplings are unified at the grand unified energy scale giving five free parameters; m_0 and $m_{1/2}$ (universal scalar and gaugino mass parameters), A_0 (universal trilinear coupling), $\tan(\beta)$ (the ratio of the vacuum expectation values of the scalar fields) and $\text{sgn}(\mu)$ (the sign of the higgsino mass

term). In the R-parity conserving mSUGRA models, SUSY particles are produced in pairs and the lightest SUSY particle (LSP) is stable and undetected resulting in large missing transverse energy (dominant signature). It will be possible to detect SUSY in multiple inclusive signatures over a large part of the mSUGRA parameters space with just 1 fb^{-1} of integrated luminosity, as discussed in [1]. As a part of the ongoing CSC exercise, several mSUGRA points that are favored by the WMAP data [2] have been simulated with a detailed detector simulation. The mSUGRA parameters for these points are given in Table 1.

3 Mass Measurement Techniques

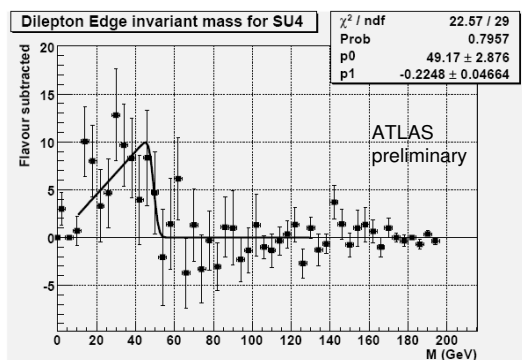
In the R-parity conserving mSUGRA models all SUSY events contain two invisible neutralinos ($\tilde{\chi}_1^0$) which, being only weakly interacting, escape the detector; therefore no mass peaks can be reconstructed directly. However kinematic endpoints and thresholds in the invariant mass distributions of the visible decay products can be measured. The values of these kinematic features can be expressed as a function of the masses of the involved sparticles [3]. The mass measurement strategy is to exploit kinematics of long decay chains originating from gluino or squark production. Specifically, from the reconstruction of the decay chain $\tilde{q}_L \rightarrow q\tilde{\chi}_2^0 \rightarrow q\tilde{l}_R^\pm l^\mp \rightarrow q\tilde{\chi}_1^0 l^\pm l^\mp$ one can measure the masses of \tilde{q}_L , \tilde{l}_R , $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^0$. Examples of mass measurements using the endpoint method are given below.

3.1 Dilepton Endpoint

A clear endpoint is exhibited by the dilepton invariant mass distribution for the decay channel $\tilde{\chi}_2^0 \rightarrow \tilde{l}_R^\pm l^\mp \rightarrow \tilde{\chi}_1^0 l^\pm l^\mp$. Fully simulated samples for the SU4 point (see Table 1) and $t\bar{t}$ background are considered. Events are selected by requiring at least two leptons

Table 1. mSUGRA points chosen for full detector simulation in the CSC exercise. ISAJET 7.71 is used and the top quark mass is set to 175 GeV.

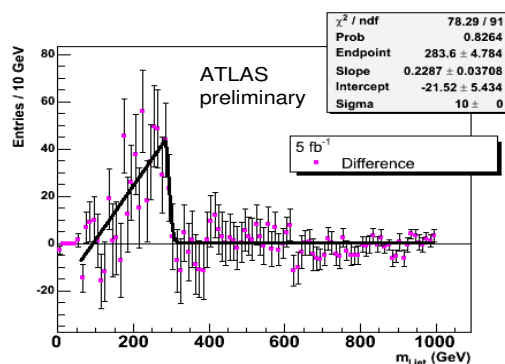
Point Name	mSUGRA Region	m_0 (GeV)	$m_{1/2}$ (GeV)	A_0 (GeV)	$\tan(\beta)$	$\text{sgn}(\mu)$	σ (pb)
SU1	Coannihilation	70	350	0	10	+	7.43
SU2	Focus	3550	300	0	10	+	4.86
SU3	Bulk	100	300	-300	6	+	18.59
SU4	Low mass	200	160	-400	10	+	262
SU6	Funnel	320	375	0	50	+	4.48
SU8.1	Coannihilation	210	360	0	40	+	6.44
SU8.2	Coannihilation	215	360	0	40	+	6.40
SU8.3	Coannihilation	225	360	0	40	+	6.32

**Fig. 1.** The dilepton invariant mass distribution for a full simulation sample of the SU4 point with an integrated luminosity of 0.35 fb^{-1} . A triangular function convoluted with a Gaussian is fitted to estimate the edge position.

(e or μ) with $p_T > 10$ GeV and $|\eta| < 2.5$, a calorimetric energy deposit $E_T < 5$ GeV in a $\eta - \phi$ cone of size 0.3 around the lepton direction, and at least four jets with $E_T > 100, 50, 50, 50$ GeV and the missing transverse energy $E_T^{\text{miss}} > 120$ GeV. The effective mass variable M_{eff} , defined as $M_{\text{eff}} = E_T^{\text{miss}} + \sum_{\text{jet}} E_{T,\text{jet}}$ taking into account the first four leading jets, is required to be $M_{\text{eff}} > 550$ GeV. The combinatorial background from $t\bar{t}$ and $\tilde{\chi}^\pm$ decays cancel in the combination $e^+e^- + \mu^+\mu^- - e^\pm\mu^\mp$ (so called flavor subtraction) which is plotted in Figure 1. The dilepton invariant mass distribution is shown for an integrated luminosity of 0.35 fb^{-1} . The mass distribution is fitted to a triangular function convoluted with a Gaussian to estimate the edge position. The endpoint fit gives a value of (49.2 ± 2.9) GeV which is within 1.6σ of the expected value of 53.7 GeV. The signal significance is calculated as 16.5 for 100 pb^{-1} of data.

3.2 Lepton-Jet Endpoint

For the SU1 point, the decay $\tilde{q}_L \rightarrow q\tilde{\chi}_1^\pm \rightarrow ql^\pm\tilde{\nu}_l \rightarrow ql^\pm\nu_l\tilde{\chi}_1^0$ gives rise to an endpoint in the lepton-jet invariant mass distribution. A mixed event technique is used to subtract the combinatorial jet background. The technique makes use of randomly pairing the jets from a different event (satisfying same event selection)

**Fig. 2.** The invariant mass distribution of lepton-jet for the SU1 full simulation sample with an integrated luminosity of 5 fb^{-1} . A triangular function convoluted with a Gaussian is fitted to estimate the edge position.

with the lepton and then subtracting the mixed-event-jet distribution from the same-event-jets distribution (with a normalization correction applied) to obtain an inferred “correct jet” distribution. Fully simulated samples for the SU1 point and $t\bar{t}$ background are considered. Events are selected by requiring one lepton (e or μ) with $p_T > 20$ GeV and $|\eta| < 2.5$, and a lepton isolation cut of $E_T < 10$ GeV in a cone of 0.45. Additional cuts are applied to reduce the $t\bar{t}$ background; the leading and second leading jets are selected with $E_T > 200$ GeV, the transverse mass $M_T < 60$ GeV or $M_T > 100$ GeV, $E_T^{\text{miss}} > 250$ GeV. The invariant mass of the lepton-jet is shown in Figure 2 for an integrated luminosity of 5 fb^{-1} . The mass distribution is fitted to a triangular function convoluted with a Gaussian to estimate the edge position. The endpoint fit gives a value of (283.6 ± 4.8) GeV compared to the expected value of 284 GeV. The endpoint can be determined with 5σ statistical significance for 5 fb^{-1} of data.

3.3 Ditau Endpoint

The mass difference between $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^0$ can also be measured from the endpoint of the ditau invariant mass distribution from the decay channel $\tilde{\chi}_2^0 \rightarrow \tilde{\tau}_1^\pm\tau^\mp \rightarrow \tilde{\chi}_1^0\tau^\pm\tau^\mp$. The SU3 point has a factor of four larger

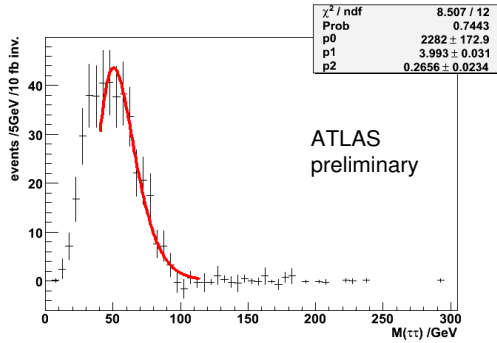


Fig. 3. The ditau invariant mass distribution for the SU3 fast simulation sample with an integrated luminosity of 10 fb^{-1} . It is fitted to a function as described in the text to estimate the endpoint.

branching ratio to $\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 \tau^\pm \tau^\mp$ than to e/μ . A fast simulation sample for the SU3 point together with Z+jets, W+jets, $t\bar{t}$, $b\bar{b}$ +jets, dijets and multijets backgrounds is considered. Events are selected by requiring at least one jet with $E_T > 220 \text{ GeV}$, at least three jets with $E_T > 50 \text{ GeV}$ and at least four jets with $E_T > 40 \text{ GeV}$ and $E_T^{\text{miss}} > 230 \text{ GeV}$. The angular separation between the taus is chosen as $\Delta R(\tau, \tau) < 2$. Only the hadronic tau decays are considered. The ditau invariant mass distribution calculated from $\tau^+ \tau^- - \tau^\pm \tau^\pm$ (to cancel the background from $\tilde{\chi}_1^\pm$ decays) is shown in Figure 3 for an integrated luminosity of 10 fb^{-1} . Since the measurement of the endpoint by a linear fit has a strong dependence from the fitting range or binning of the distribution, a new approach is considered to estimate the endpoint. A fit function in the form of

$$y = \frac{p_0}{x} \cdot \exp \left[\frac{-1}{2p_2^2} (\ln(x) - p_1)^2 \right] \quad (1)$$

is used (modified from [4]). The inflection point is calculated as:

$$x_{IP} = \exp \left[\frac{1}{2} p_2^2 \left(-3 + \sqrt{1 + \frac{4}{p_2^2}} \right) + p_1 \right] \quad (2)$$

The measurement of the inflection point is more stable under variations of the fitting range or binning. A calibration line is driven between the inflection point and the endpoint by varying the involved masses in the decay chain. This calibration line is found to be $y = (0.47 \pm 0.02)x + (15 \pm 2)$. The endpoint value is then calculated as $(105 \pm 4) \text{ GeV}$ compared to the expected value of 98.3 GeV .

4 mSUGRA Parameters Determination

More endpoint measurements can be done for instance from the reconstruction of gluino decays; $\tilde{g} \rightarrow \tilde{\chi}_2^0 t\bar{t}$ and $\tilde{g} \rightarrow \tilde{t}_1 t \rightarrow t b \tilde{\chi}_1^\pm$ for which the details can be found in [5] and [6] respectively. Also, the mass of \tilde{q}_R can be measured from the decay $\tilde{q}_R \rightarrow \tilde{\chi}_1^0 q$ as studied

Table 2. Determination of mSUGRA parameters for the SPS1a point using the SFitter program. $\text{sgn}(\mu)$ is fixed.

	SPS1a	Δ_{masses}	Δ_{edges}
m_0 (GeV)	100	3.9	1.2
$m_{1/2}$ (GeV)	250	1.7	1.0
$\tan(\beta)$	10	1.1	0.9
A_0 (GeV)	-100	33	20

in [6]. The potential for extracting the masses from the endpoint measurements is described in [9] for the SPS1a point [10] with 300 fb^{-1} . The ultimate goal is to perform a fit of the mSUGRA parameters from a given set of mass measurements using the tools such as Fittino [7] and SFitter [8]. As an example, the measured masses and kinematical edges of the SPS1a point from [9] are fed into the SFitter program to determine the mSUGRA parameters for 300 fb^{-1} [11]. As seen in Table 2 the parameters can be determined with a precision at the percent level by using the masses (Δ_{masses} column), moreover the precision can be improved significantly by using the measured edges, thresholds and mass differences (Δ_{edges} column) in the fit instead of the masses.

5 Spin Measurements

If SUSY signals are observed at the LHC, it will be vital to measure the spins of the new particles to demonstrate that they are indeed the predicted super-partners. Two methods of measuring the spin are described below.

5.1 Neutralino Spin Measurement

The decay chain $\tilde{q}_L \rightarrow q \tilde{\chi}_2^0 \rightarrow q \tilde{l}_R^{\pm} l_{near}^{\mp} \rightarrow q \tilde{\chi}_1^0 l_{far}^{\pm} l_{near}^{\mp}$ provides a good opportunity to measure the spin of $\tilde{\chi}_2^0$ using the lepton charge asymmetry [12]. The squarks and sleptons are spin-0 particles and their decays are spherically symmetric, $\tilde{\chi}_2^0$ however has spin-1/2 and the angular distribution of its decay products is not spherically symmetric. This leads to a charge asymmetry of the invariant mass of the quark and the near lepton which is defined as:

$$A^{+-} = \frac{s^+ - s^-}{s^+ + s^-} \quad s^\pm = \frac{d\sigma}{dm(ql^\pm)} \quad (3)$$

The asymmetry is suppressed by the fact that quark jets cannot be experimentally distinguished from anti-quark jets at the LHC, however at the LHC much more squark than anti-squark will be produced, and therefore a residual asymmetry is still observable. This asymmetry is calculated for the SU3 point considering a fast simulation sample of 30 fb^{-1} together with the most relevant SM backgrounds, $t\bar{t}$, W+jets, Z+jets. Events are selected by requiring two opposite sign leptons (e or μ) with $p_T > 10 \text{ GeV}$ and $|\eta| < 2.5$, a lepton isolation cut of $E_T < 10 \text{ GeV}$ in a cone of 0.2,

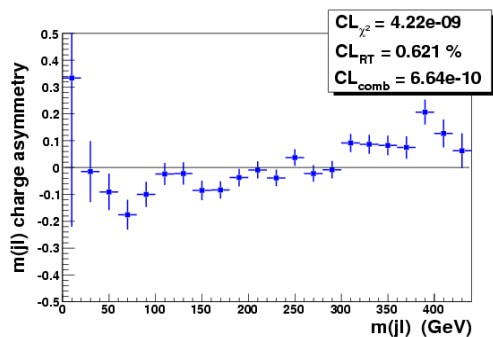


Fig. 4. The lepton charge asymmetry in the invariant mass of the lepton-jet for the SU3 point with an integrated luminosity of 30 fb^{-1} .

at least four jets with $E_T > 100, 50, 50, 50 \text{ GeV}$ and $E_T^{\text{miss}} > 100 \text{ GeV}$. The charge asymmetry is plotted in Figure 4 by considering both the near and far leptons since these are not distinguishable for the SU3 point. The confidence level for the asymmetry distribution to be flat (spin-0) is calculated using two independent statistical methods as described in [13] and found to be $\sim 10^{-9}$. It is observed that 10 fb^{-1} of data would be sufficient to detect a non-zero charge asymmetry at the 99% confidence level for the SU3 point.

5.2 Slepton Spin Measurement

The spin of the slepton can be measured using an angular variable which is sensitive to the polar angle in direct slepton pair production $q\bar{q} \rightarrow Z^0/\gamma \rightarrow \tilde{l}^+\tilde{l}^- \rightarrow \tilde{\chi}_1^0 l^+ \tilde{\chi}_1^0 l^-$ as discussed in [14]. The angular variable is defined as $\cos\theta_{ll}^* \equiv \cos(2 \tan^{-1} e^{(\Delta\eta_{l+l^-}/2)}) = \tanh(\Delta\eta_{l+l^-}/2)$. It is interpreted as the cosine of the polar angle between each lepton and the beam axis in the longitudinally boosted frame in which the pseudorapidities of the leptons are equal and opposite. This variable is on average smaller for SUSY than for the Universal Extra Dimensions (UED) so it can be employed as a spin-discriminant in slepton/Kaluza-Klein-lepton pair production in hadron colliders. The study was performed on Point 5 studied for the ATLAS TDR [15] which has a phenomenology very similar to the one of point SU3. A fast simulation sample of 200 fb^{-1} is considered together with the major SM backgrounds. Events are selected by requiring two opposite sign and same flavor electrons or muons with $p_T(l_1) > 40 \text{ GeV}$, $p_T(l_2) > 30 \text{ GeV}$, $M_{ll} < 150 \text{ GeV}$, no jet with $E_T > 100 \text{ GeV}$, no b-jets, $E_T^{\text{miss}} > 100 \text{ GeV}$, $M_{T2} < 100 \text{ GeV}$. The angular variable is plotted in Figure 5 together with the predictions for SUSY (black line), phase space (dotted blue line) and the UED (dashed red line). The yellow shaded band represent the SUSY expectation when all the SUSY particle masses are simultaneously changed by $\pm 20 \text{ GeV}$. As seen in the figure, the data points are much better matched to the slepton angular distribution than to either the phase-space one or the UED-like one. This shows that $\cos\theta_{ll}^*$ does indeed

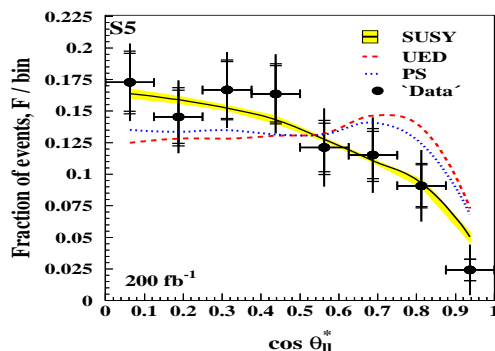


Fig. 5. $\cos\theta_{ll}^*$ distribution for the Point 5 with an integrated luminosity of 200 fb^{-1} .

measure the spin of the sleptons for this point. Further studies with the Snowmass points SPS1a, SPS1b, SPS3, SPS5 [10] allow slepton spin determination with $100\text{-}300 \text{ fb}^{-1}$ of data.

6 Conclusions

LHC brings experimental physics into a new territory. ATLAS can discover the SUSY particles up to 2-3 GeV mass scale if they exist. Understanding the detector response and the SM background will be a challenge at the beginning. Many techniques have been developed to measure the masses (edges, thresholds, mass differences) and spin of SUSY particles and to determine the underlying model parameters.

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