

Study of DiMuon Rare Beauty Decays with ATLAS and CMS

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Abstract. The LHC experiments will perform sensitive tests of physics beyond the Standard Model (BSM). Investigation of decays of beauty hadrons represents an alternative approach in addition to direct BSM searches. The ATLAS and CMS efforts concentrate on those B -decays that can be efficiently selected already at the first and second level trigger. The most favorable trigger signature will be for B -hadron decays with muons in the final state. Using this trigger, ATLAS and CMS will be able to accommodate unprecedentedly high statistics in the rare decay sector. These are purely dimuon decays, and families of semimuonic exclusive channels. Already with $1fb^{-1}$ the sensitivity in the dimuon channels will be comparable to today world statistics. The strategy is to carry on the dimuon channel program up to nominal LHC luminosity. In particular the $B_s \rightarrow \mu\mu$ signal with ~ 5 sigma significance can be measured combining low luminosity $10^{33} cm^{-2} s^{-1}$ samples with those of one year of LHC operation at a luminosity of $10^{34} cm^{-2} s^{-1}$.

PACS. 13.30.Ce Leptonic, semileptonic, and radiative decays – 13.20.He Decays of bottom mesons

1 Introduction

Rare leptonic and semileptonic B -decays, produced by FCNC transitions, are forbidden at the tree level in the Standard Model (SM). These decays occur at the lowest order only through one-loop “penguin” and “box” diagrams. The branching ratios of these decays are very small: from 4×10^{-5} for rare radiative decay $B_d^0 \rightarrow K^* \gamma$ to 10^{-15} for rare Cabibbo suppressed leptonic decay $B_d^0 \rightarrow e^+ e^-$.

The careful investigation of rare B -decays is mandatory for testing ground of the Standard Model and offer a complementary strategy in the search of new physics. The probing of loop-induced couplings provide means of testing the detailed structure of the SM at the level of radiative corrections. In particular, FCNC involving $b \rightarrow s, d$ transitions and $B \rightarrow ll$ decays, provide an excellent probe of new indirect effects by yielding informations on the masses and coupling of the virtual particles running in the loops. In SUSY models, the branching fraction for $B_{s(d)}^0 \rightarrow \mu^+ \mu^-$ has a strong dependence on $\tan\beta$. A precise measurement of such decays will allow to constrain the supersymmetric extensions of the SM.

To date the decay modes $B_{s(d)}^0 \rightarrow \mu^+ \mu^-$ have not yet been observed. The current best upper limits on branching ratio come from D0 [1] and CDF [2] collaborations and are 9.3×10^{-8} and 5.8×10^{-8} respectively at 95%CL. The searches for rare B decays at the B -factories CLEO, Belle and BaBar have not sensitivity to B_s decays.

Table 1. Branching ratios in SM for rare B -decays with $\mu^+ \mu^-$ pair in final states.

Decay channel	Br. ratio	Ref.
$B_s^0 \rightarrow \mu^+ \mu^-$	3.4×10^{-9}	[6]
$B^+ \rightarrow K^+ \mu^+ \mu^-$	3.5×10^{-7}	[7]
$B^+ \rightarrow K^{*+} \mu^+ \mu^-$	$\sim 10^{-6}$	[7]
$A_b \rightarrow \Lambda \mu^+ \mu^-$	2.0×10^{-6}	[8],[9]
$B_d^0 \rightarrow K^{0*} \mu^+ \mu^-$	1.3×10^{-6}	[10],[11]
$B_s^0 \rightarrow \phi \mu^+ \mu^-$	$\sim 10^{-6}$	[10],[11]

In the last years the B -factories BaBar and Belle presented the first results for $B \rightarrow (K^*, K) l^+ l^-$ branching ratios and forward-backward asymmetry (A_{FB}) in these rare semileptonic decays [3],[4],[5] but still affected by large statistical errors.

In this report we will pay our attention to the (semi)leptonic decays with $\mu^+ \mu^-$ pair in final state where ATLAS and CMS can give a significant contribution. We discuss the simulation results and the perspectives of measurements. The branching ratios in SM of the decays studied can be found in Table 1.

2 Theoretical description

From the theoretical point of view, the $b \rightarrow q$ ($q = s, d$) transitions are described using the effective Hamiltonian

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$$\mathcal{H}_{eff} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{tq}^* \sum_{i=1}^{10} C_i(\mu) O_i(\mu) \quad (1)$$

in the form of Wilson expansion (see e.g. [12]). The set of Wilson coefficients $C_i(\mu)$ depends on the current model and contains the lowest order model contributions and perturbative QCD corrections. The scale parameter μ is approximately equal to the mass of b -quark ($\sim 5\text{GeV}$). This parameter separates the perturbative and nonperturbative contributions of the strong interactions. The nonperturbative contribution is contained in the matrix elements of basis operators $O_i(\mu)$ between the initial and final hadronic states. For the calculation of these matrix elements it is necessary to use different decay-specific nonperturbative methods: QCD Sum Rules, Heavy Quarks Effective Theory, Quark Models and Lattice calculations. The accuracy in nonperturbative calculations depends on the method, but it is not less than 15%. The accuracy of the Wilson coefficient for with NLO and NNLO QCD corrections [13] is not greater than 15% if μ parameter ranges in $[m_b/2, 2m_b]$.

In the SM the decay width of the rare muonic decays is:

$$\Gamma(B_q^0 \rightarrow \mu^+ \mu^-) = \frac{G_F^2 \alpha_{em}^2}{16\pi^3} |V_{tq}^* V_{tb}|^2 \cdot \left(f_{B_q} m_\mu C_{10A} \right)^2 \sqrt{M_{B_q}^2 - 4m_\mu^2}. \quad (2)$$

This expression contains only one nonperturbative constant f_{B_q} . The value of this constant is known with high accuracy (about 5-10%). Furthermore the Wilson Coefficient C_{10A} in the NLO approach is not dependent on the scale parameter μ , and does not add any uncertainty in the theoretical predictions.

3 Trigger strategies for rare decays

Details of the ATLAS and CMS experiments can be found in Refs. [14,15].

3.1 The ATLAS trigger

ATLAS has a three level trigger system [16] which reduces the 40MHz bunch crossing rate to about 100Hz of events to be recorded. The first level trigger (LVL1) is hardware-based and makes a fast decision (in $2.5\mu\text{s}$) about which events are interesting for further processing. Coarse granularity informations from calorimeter and muon spectrometer are used to identify region of interest (RoI) of the detector which contain interesting signals (high energy electrons, muons and taus and jets with large transverse or missing energy). The RoIs are used to guide the later stages of the trigger. After LVL1 the trigger rate will be reduced to less than 75kHz.

The high level trigger (HLT) is software-based and is split into two levels. At the level 2 (LVL2) the full

granularity of the detector is used to confirm the LVL1 decisions and then to combine informations from different sub-detectors within the LVL1 RoIs. Fast algorithms are used for the reconstruction at this stage and the rate is reduced to $\sim 2\text{kHz}$ with an average time of execution of $\sim 10\text{ms}$. At the level 3, the Event Filter (EF), the whole event is available and offline-like algorithms are used with better alignment and calibration informations to form the final decision. The rate is reduced to 100Hz with an execution time of $\sim 1\text{s}$.

3.1.1 ATLAS trigger for rare decays

The B -trigger is expected to account for 5-10% of the total trigger resources. The core of the B -trigger is the LVL1 muon trigger which is based on the measurement of the muon transverse momentum (p_T). The efficiency of the muon trigger is expected at about 85% at the plateau. The dimuon LVL1 trigger (two muons with p_T above 6GeV), used for rare decay selection, is expected to have a rate of about 500Hz.

The LVL1 dimuons will be confirmed at the LVL2 firstly in the muon system by means of the precision tracking chamber and then by combining muon and inner detector tracks. Finally the two muons can be combined and mass cuts are applied. At the EF the tracks are refitted in the RoIs and vertex reconstruction is performed. Cuts are applied on decay length and invariant B -hadron mass. For $B_s^0 \rightarrow \mu^+ \mu^-$ events containing two muons with $p_T > 6\text{GeV}$, efficiencies of 60-70% are expected.

3.2 The CMS trigger

CMS has a two level trigger [15] which reduces the bunch crossing rate down to about 150Hz for recording. The Level-1 (L1) trigger uses muon detector and calorimeter informations and is hardware-based with an output rate of about 100kHz and a latency of $3.5\mu\text{s}$. The HLT is software-based with the required 150Hz output rate. It uses reconstruction algorithms similar to the offline with a mean execution time per event of about 40ms. To speed up reconstruction in the HLT, a partial track reconstruction is performed: the track resolution becomes asymptotic after 5-6 hits are used in the track fit.

3.2.1 CMS trigger for rare decays

As in ATLAS, CMS trigger for B events uses single and dimuon triggers. The L1 dimuon trigger has a low p_T threshold of 3GeV which ensures a high efficiency for events with two muons in the final state with a rate of 0.9kHz at $2 \times 10^{33} \text{cm}^{-2} \text{s}^{-1}$. At the HLT the L1 decision will be confirmed using the full muon system and an improved momentum measurement with the tracker. Primary vertex (PV) reconstruction is available from the pixel detector and also track reconstruction is performed in cones around the L1 muons using the partial

Table 2. Mass and proper time resolution obtained on the $B_s^0 \rightarrow \mu^+ \mu^-$ signal MonteCarlo event sample for CMS and ATLAS.

Experiment	Mass res.(GeV)	Proper time res.(fs)
ATLAS	0.084	91
CMS	0.036	95

track reconstruction. The exclusive rare decay is then reconstructed and cuts on invariant mass, vertex fit quality and decay length are applied.

4 Muonic decays in ATLAS and CMS

As purely leptonic B -decays are theoretically very clean, they provide an ideal channel for seeking indirect hints of new physics effects. However they are very difficult to observe because of their small branching ratio (see Table 1). Most probably they will not be visible before the LHC data taking. ATLAS and CMS will start sensitive measurements at $10^{33} \text{cm}^{-2} \text{s}^{-1}$. Even at high design luminosity ($10^{34} \text{cm}^{-2} \text{s}^{-1}$) the trigger for $B \rightarrow \mu^+ \mu^-$ decay is not problematic.

A good background rejection is necessary for the signal selection. The main contributions to the background [17] come from the processes $b\bar{b}(b\bar{b}b\bar{b}, b\bar{b}c\bar{c}) \rightarrow X\mu^+\mu^-$ with the muons originating mainly from semi-leptonic b and c quark decays. This background can be estimated extrapolating the Tevatron data on heavy quark production to the LHC energies. It would be remarked that both ATLAS and CMS analysis are limited by the statistics of the MonteCarlo background sample.

The event selection relies on topological variables related to the PV, the muon candidates and the B_s secondary vertex and is very similar in both the experiments [18,19]. Simple cuts can be applied to distinguish combinatorial background from the signal:

- B -hadron invariant mass;
- secondary vertex length and quality;
- pointing of B -hadron momentum to PV;
- track isolation.

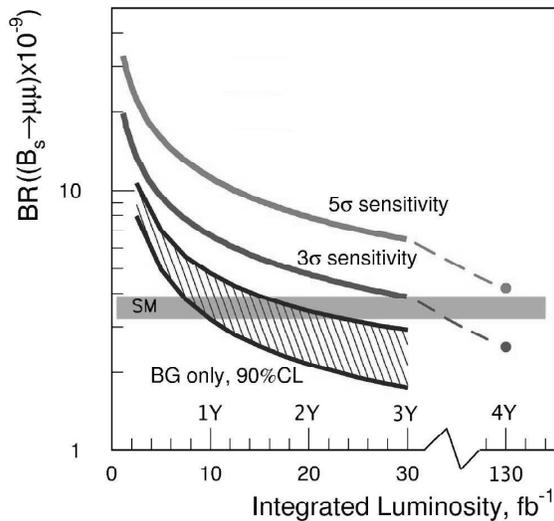
Table 2 summarizes the mass resolution and the proper time resolution obtained on the signal MonteCarlo event sample for CMS and ATLAS.

The main result of the analysis is the expected upper limit on branching ratio that can be achieved in the first 10fb^{-1} as summarized in Table 3. Even if these upper limits are above the SM expectation, they allow to set stringent constraint on new physics models with the first 10fb^{-1} of integrated luminosity.

Figure 1 shows the ATLAS expectation for $B_s^0 \rightarrow \mu^+ \mu^-$ branching ratio as a function of the integrated luminosity (or equivalently as a function of the time). The SM expectation can be reached with ~ 5 sigma significance combining low luminosity $10^{33} \text{cm}^{-2} \text{s}^{-1}$ samples with those of one year of LHC operation at the nominal luminosity of $10^{34} \text{cm}^{-2} \text{s}^{-1}$.

Table 3. Expected upper limit on $B_s^0 \rightarrow \mu^+ \mu^-$ branching ratio after 10fb^{-1} .

Experiment	Br. ratio
ATLAS	$7.0 \times 10^{-9} @ 90\% CL$
CMS	$1.4 \times 10^{-8} @ 90\% CL$

**Fig. 1.** ATLAS perspective of measurements for $B_s^0 \rightarrow \mu^+ \mu^-$ branching ratio as a function of the integrated luminosity (or equivalently as a function of the time).

5 SemiMuonic decays in ATLAS

Thanks to the dimuon final state (see Table 1), the semimuonic decays, as the purely dimuonic ones, are easy to select at the trigger level. The observation of semileptonic decays give access to a number of observables. The precise measurements of such observables could give very interesting informations for new physics reach. A_{FB} is one of the most promising parameters.

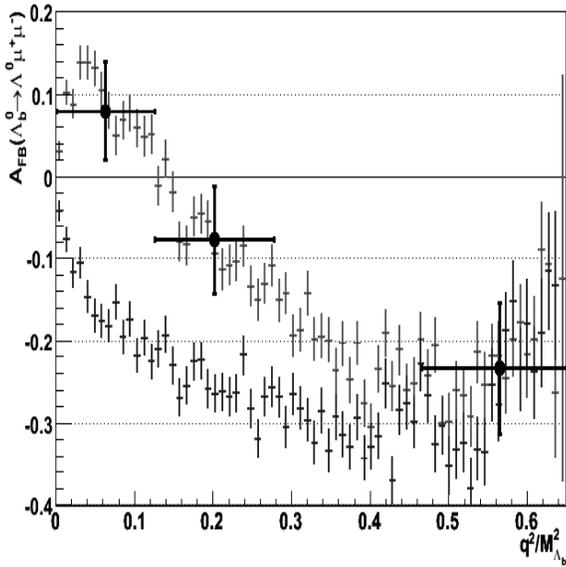
The small branching ratios of semimuonic decays require a powerful background rejection. Semileptonic decays with $c\bar{c}$ resonances decaying into two muons represent an irreducible background source. A cut on dimuon invariant mass around nominal values for resonances removes this background. Combinatorial background arises from muons originating mainly from semi-leptonic decays of b and c quarks. Specific decay channels can represent background sources due mainly to hadron misidentification as muons, but their contribution is expected to be less important than the previous two sources.

The event selection [20,21] is related to topological variables:

- vertex quality and invariant mass of the dimuon system;
- displacement and quality of vertices and mass of the secondary hadrons;

Table 4. Expected statistics for semimuonic decays events and background in ATLAS after 30fb^{-1} .

Decay	Signal events	Background events
$B^+ \rightarrow K^+ \mu^+ \mu^-$	4000	<10000
$B^+ \rightarrow K^{*+} \mu^+ \mu^-$	2300	<10000
$\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$	800	<4000
$B_d^0 \rightarrow K^{0*} \mu^+ \mu^-$	2500	<10000
$B_s^0 \rightarrow \phi \mu^+ \mu^-$	900	<10000

**Fig. 2.** Forward-backward asymmetry for $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ after 30fb^{-1} (see Section 5).

– pointing of B hadron momentum to PV.

Statistics expected after three years of data taking at low luminosity (30fb^{-1}) are summarized in Table 4. The background level estimation is only limited by the low MonteCarlo statistics available at the moment. It should be pointed out that, thanks to the muon pair in the final state, semimuonic decays will be also studied at high luminosity, so that a larger statistics than those of Table 4 will be collected.

The expected precision on A_{FB} after 30fb^{-1} is presented on Figure 2 for $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$ decay. The three dots with error bars correspond to simulated data after offline analysis. Upper point set corresponds to theoretical SM predictions, and lower set corresponds to MSSM predictions. The statistical error in the low dimuon invariant mass region is at level of 6%. The statistical errors expected on branching ratio measurements are at the level of 3.5% and 6.5% respectively for $B \rightarrow K \mu^+ \mu^-$ and $B \rightarrow K^* \mu^+ \mu^-$ decays. These errors on the branching ratio measurements are much smaller than the current experimental and theoretical ones.

6 Conclusions

The results obtained for $B_s^0 \rightarrow \mu^+ \mu^-$ by ATLAS and CMS are comparable and promise an interesting startup analysis with the possibility of setting tight constraints on new physics models beyond the SM. The simulation studies show that ATLAS detector will be capable to extract signals of semimuonic B -decays and reach a good sensitivity to new physics beyond the SM.

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