

NLO Event Simulation for Chargino Production at the ILC

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- 1 Introduction and Motivation
 - Charginos and Neutralinos in the MSSM
 - Experimental accuracy and NLO results

- 2 Inclusion of NLO results in WHIZARD
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 - Photons: fixed order vs resummation
 - Results

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Chargino and Neutralino sector: Reconstruction of SUSY parameters

- Charginos $\tilde{\chi}_i^\pm$ and Neutralinos $\tilde{\chi}_i^0$:
superpositions of gauge and Higgs boson superpartners
- Chargino/ Neutralino sector:

$\tan\beta$, μ (Higgs sector), M_1 , M_2 (soft breaking terms)

can be reconstructed from

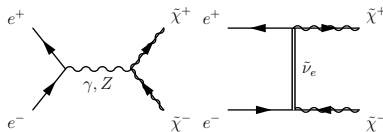
masses of $\tilde{\chi}_1^\pm$, $\tilde{\chi}_2^\pm$, $\tilde{\chi}_1^0$, 2σ in the $\tilde{\chi}^\pm$ sector

(Choi et al 98, 00, 01)

- low-scale parameters + evolution to high scales (RGEs):
 \Rightarrow hint at SUSY breaking mechanism (Blair et al, 02)
- requires high precision in ew-scale parameter determination

Chargino production at the ILC

- **ILC:** future e^+e^- collider, $\sqrt{s} = 500$ GeV (1 TeV)
“clean” environment, low backgrounds \Rightarrow high precision
- Charginos: (typically) light in the MSSM
 \Rightarrow easily accessible at colliders (ILC/ LHC) \Leftarrow
- LO production at the ILC:



- decays: typically long decay chains

$$\text{e.g. } e^+ e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow \tilde{\tau}_1^+ \tilde{\tau}_1^- \nu_\tau \bar{\nu}_\tau (\rightarrow \tau^+ \tau^- \nu_\tau \bar{\nu}_\tau \tilde{\chi}_1^0 \tilde{\chi}_1^0)$$

Experimental accuracy and theoretical next-to-leading-order (NLO) corrections

- experimental errors: obtained from simulation studies (LHC/ ILC study, Weiglein ea, 04)
- generate “experimental data” with known SUSY input parameters
- errors: combination of statistical and systematic errors

combined **LHC + ILC**: ‰

same \mathcal{O} errors from fitting routines determining SUSY parameters

- **Theory:**

Full NLO SUSY corrections for $\sigma(ee \rightarrow \tilde{\chi} \tilde{\chi})$ at ILC:
in the ‰ regime (Fritzsche ea 04, Öller ea 04, 05)

⇒ include complete NLO contributions in analyses⇐

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From σ_{tot} to Monte Carlo event generators

MC event generators: Generate event samples
(same form as experimental outcome)

- experiments: see final decay products
- need to compare with simulated event samples
- also: important irreducible background effects
(e.g. Hagiwara et al, 05, → talk by Jürgen Reuter)

⇒ include NLO results in Monte Carlo Generators ⇐

- MC Generator WHIZARD (W. Kilian, LC-TOOL-2001-039):
- so far: LO Monte Carlo Event Generator for $2 \rightarrow n$ particle processes
- includes various physical models (SM, MSSM, non-commutative geometry, little Higgs models), initial state radiation, parton shower models,...

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NLO cross section contributions

σ_{tot} contributions and dependencies:

- σ_{born}
- virtual $\mathcal{O}(\alpha)$ corrections: $\sigma_{\text{virt}}(\lambda)$
- emission of soft/ hard collinear/ hard non-collinear photons:
$$\sigma_{\text{soft}}(\Delta E_\gamma, \lambda) + \sigma_{\text{hc}}(\Delta E_\gamma, \Delta\theta_\gamma) + \sigma_{2 \rightarrow 3}(\Delta E_\gamma, \Delta\theta_\gamma)$$
- higher order initial state radiation: $\sigma_{\text{ISR}} - \sigma_{\text{ISR}}^{\mathcal{O}(\alpha)}(Q)$
 λ : photon mass, ΔE_γ : soft cut, $\Delta\theta_\gamma$: collinear angle

Including FormCalc $\mathcal{O}(\alpha)$ results in WHIZARD

- use FeynArts / FormCalc generated code for

$$\begin{aligned} \mathcal{M}_{\text{virt}}(\lambda) &: \text{ virtual corrections} \\ f_s(\Delta E_\gamma, \lambda) &: \text{ soft photon factor} \\ (\mathcal{M}_{\text{born}} &: \text{ born contribution}) \end{aligned}$$

- fixed order: integrate over effective matrix element:

$$|\mathcal{M}_{\text{eff}}|^2(\Delta E_\gamma) = (1 + f_s(\Delta E_\gamma, \lambda)) |\mathcal{M}_{\text{born}}|^2 + 2 \text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*(\lambda))$$

ΔE_γ : soft photon cut, λ : photon mass

- in practice: create library from FormCalc code, link this to WHIZARD

(1): Fixed $\mathcal{O}(\alpha)$ contributions

- integrate $|\mathcal{M}_{\text{eff}}|^2$ (born/ virtual/ soft photonic part)
- hard collinear photons: collinear approximation ($\mathcal{M}_{\text{born}}$)
- hard non-collinear photons: explicit $e e \rightarrow \tilde{\chi} \tilde{\chi} \gamma$ process ($\mathcal{M}_{\text{born}}^{2 \rightarrow 3}$)
- corresponds to analytic results in literature (Fritzsche ea/ Öller ea)

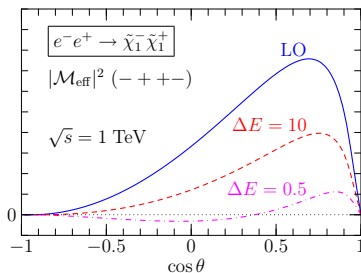
Photons: fixed order vs resummation

(1): Fixed $\mathcal{O}(\alpha)$ contributions

- integrate $|\mathcal{M}_{\text{eff}}|^2$ (born/ virtual/ soft photonic part)
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problem: too low energy cuts: $|\mathcal{M}_{\text{eff}}|^2 < 0$
 \Rightarrow use negative weights
 or set $\mathcal{M}_{\text{eff}} = 0$

event generator
specific problem
 $(\sigma_{\text{tot}} \geq 0)$



\mathcal{M}^2 behaviour, different cuts [GeV]

(2): Resumming leading logs to all orders

- idea: subtract $\mathcal{O}(\alpha)$ soft + virtual collinear contributions in \mathcal{M}_{eff} :

$$|\widetilde{\mathcal{M}}_{\text{eff}}|^2 = (1 + f_s(\Delta E_\gamma)) |\mathcal{M}_{\text{born}}|^2 + 2 \text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*) - 2 f_s^{\text{ISR}, \mathcal{O}(\alpha)}(\Delta E_\gamma) |\mathcal{M}_{\text{born}}|^2$$

- fold this with ISR structure function:

$$\int d\Gamma \int_0^1 dx_1 \int_0^1 dx_2 f^{\text{ISR}}(x_1) f^{\text{ISR}}(x_2) |\widetilde{\mathcal{M}}_{\text{eff}}|^2(s, x_i)$$

- $f^{\text{ISR}}(x)$: Initial state radiation (Jadach, Skrzypek, Z.Phys. 1991)
 \Rightarrow describes collinear (real + virtual) photons in leading log accuracy \Leftarrow
- $f_s^{\text{ISR}, \mathcal{O}(\alpha)}$: soft integrated $\mathcal{O}(\alpha)$ contribution

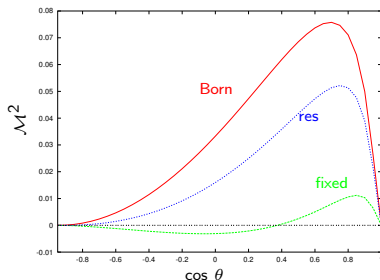
Photons: fixed order vs resummation

Resumming: What do we get ??

- $\mathcal{O}(\alpha)$: equivalent to fixed order method

⇒ got rid of
 $|\mathcal{M}|^2 < 0$
 effects !!

**no negative
 weights**



(-++-),

$\Delta E_\gamma = 0.5 \text{ GeV}$

- higher orders:
 higher order ISR for $|\mathcal{M}_{\text{born}}|^2$ as well as $\text{Re}(\mathcal{M}_{\text{born}} \mathcal{M}_{\text{virt}}^*)$!!!
 ⇒ new higher order effects ⇐
- additional possibility: also fold 2 → 3 process with ISR
 (“res+”)

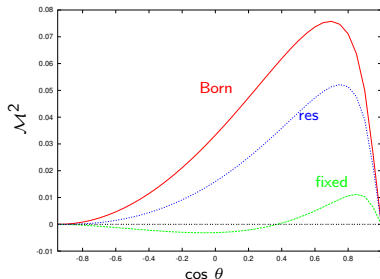
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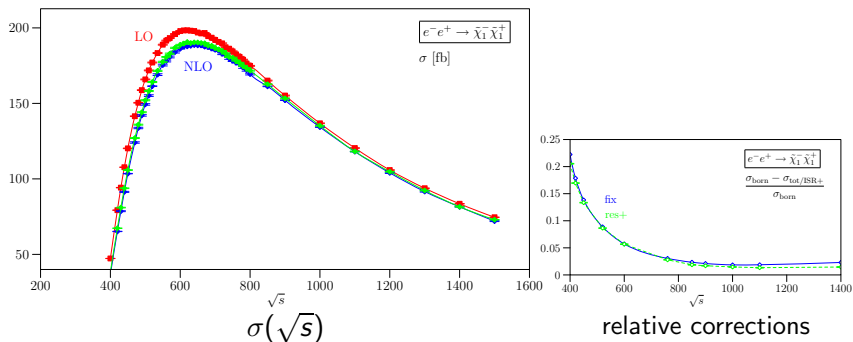


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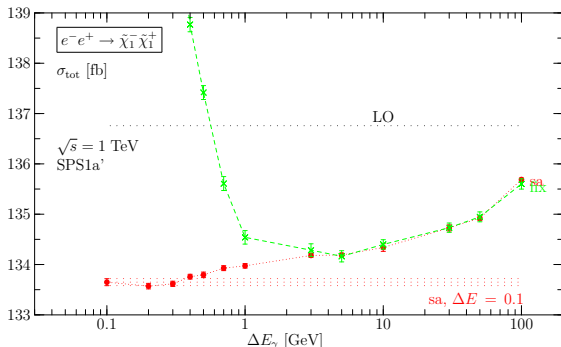
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Results: cross sections



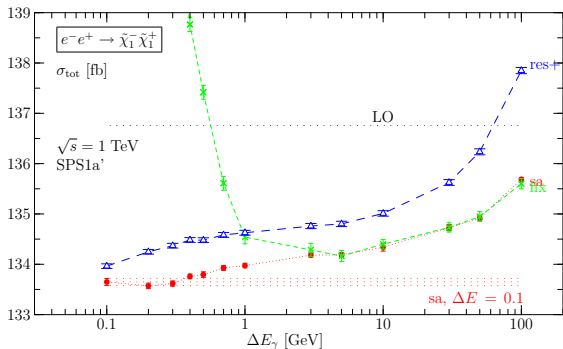
agrees with results in the literature (Fritzsche ea, Öller ea)

A closer look: ΔE_γ dependence of σ_{tot}



- **semianalytic (FormCalc)**: tests soft approximation, shifts : 2 - 5 ‰ ($\Delta E_\gamma \leq 10 \text{ GeV}$)
- **fixed order result (WHIZARD)**: same as 'sa' for $\Delta E_\gamma \geq 3 \text{ GeV}$, smaller values: $|\mathcal{M}_{\text{eff}}|^2 \leq 0$ effects

ΔE_γ dependence: resummation



$\sigma_{\text{tot}}(\Delta E_\gamma)$:
resummation includes
 higher order effects
 5‰ difference to 'sa'
 for $\Delta E_\gamma \leq 10 \text{ GeV}$

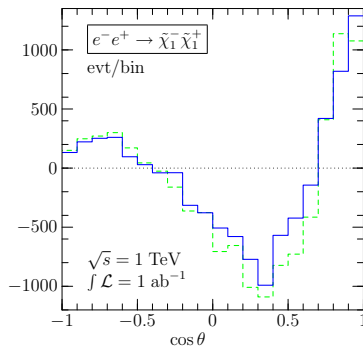
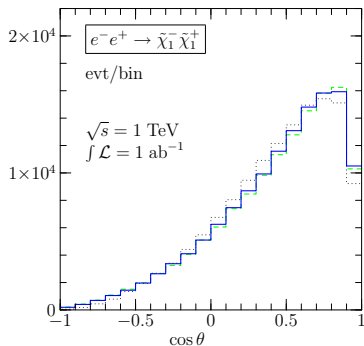
In summary:

shift in ΔE_γ leads to ‰ effects, match ILC accuracy
 \Rightarrow careful choice of ΔE_γ , method important

“best” choice: fully resummed version with low energy cut

Results: simulated events

simulation results: angular distributions

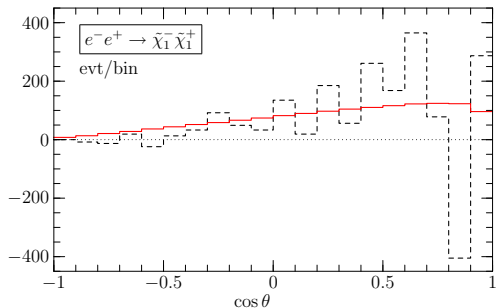


Born, fixed order, resummation

!! more than 1σ deviation !! $\sqrt{n_{\max}} \approx \mathcal{O}(10^2)$; nbins = 20

Results: simulated events

Angular distributions: higher orders



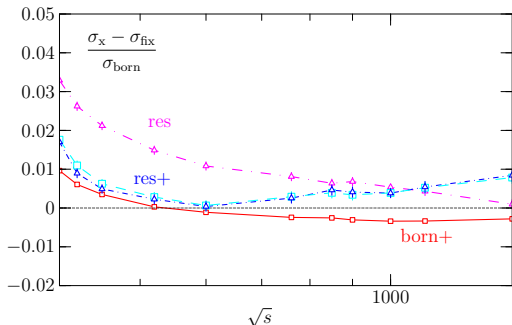
$N_{\text{res},+} - N_{\text{ex}}$
red: 1 standard dev
from Born result

N_{res}^+ : resummation, additionally 2 \rightarrow 3 folded w ISR; most complete

also higher order contributions statistically significant

Results: higher order effects

\sqrt{s} dependence of different higher order contributions



relative difference:

$$\frac{\sigma_x - \sigma_{\text{fix}}}{\sigma_{\text{Born}}}$$

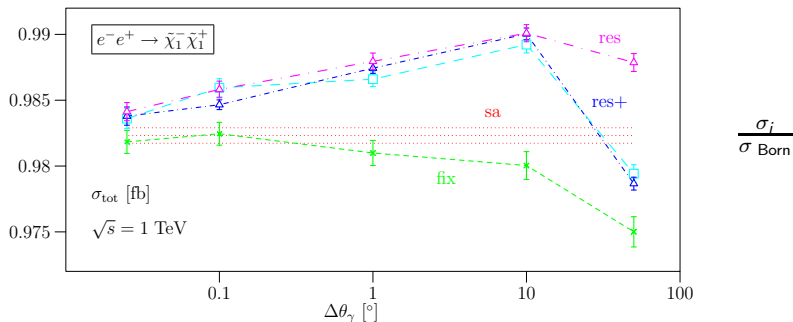
Born+: only Born folded w ISR (standard way in the literature),
 fully resummed result: subtraction, also fold 2 → 3 part with ISR
 difference between **Born+** and fully resummed result: multiple
 photon emission from interaction term

Summary and Outlook

- Chargino/ neutralino sector of MSSM: high precision in SUSY parameter analysis at EW scale ($\%$ at ILC)
 - same size/ larger NLO corrections
- ⇒ include NLO results in Monte Carlo Event generators
- resummation method for photons allows lower soft cuts/ inclusion of higher order contributions
 - NLO as well as higher order contributions significant !!
 - next steps: include NLO corrections to $\tilde{\chi}$ decays (→ talk by K.Rolbiecki), non-factorizing contributions (start with photonic corrections in the double-pole approximation)
 - general interface to FormCalc generated matrix elements: extendable to other processes...

cut dependencies: $\Delta\theta_\gamma$

tests: collinear photon approximation



σ_{tot} again larger for resummation method
 for higher angles: second order ISR effects between 0.05° and 0.1°
 ($\mathcal{O}(\%)$)

η , f_s , hard collinear approximation, $ISR^{\mathcal{O}(\alpha)}$

- $\eta = \frac{2\alpha}{\pi} \left(\log \left(\frac{Q^2}{m_e^2} \right) - 1 \right)$ (Q = scale of process)

-

$$f_s = -\frac{\alpha}{2\pi} \sum_{i,j=e^\pm} \int_{|\mathbf{k}| \leq \Delta E} \frac{d^3k}{2\omega_k} \frac{(\pm) p_i p_j Q_i Q_j}{p_i k p_j k},$$

(Denner 1992)

$\omega_k = \sqrt{\mathbf{k}^2 + \lambda^2}$, p_i initial/ final state momenta, k : γ momentum

- hard collinear factor (\pm helicity conserving/ flipping):

$$f^+(x) = \frac{\alpha}{2\pi} \frac{1+x^2}{(1-x)} \left(\ln \left(\frac{s(\Delta\theta)^2}{4m^2} \right) - 1 \right), \quad f^-(x) = \frac{\alpha}{2\pi} x.$$

(Dittmaier 1993)

-

$$f_s^{ISR, \mathcal{O}(\alpha)} = \left[\int_{x_0}^1 f_{ISR}(x) dx \right]_{\mathcal{O}(\alpha)} = \frac{\eta}{4} \left(2 \ln(1-x_0) + x_0 + \frac{1}{2} x_0^2 \right)$$

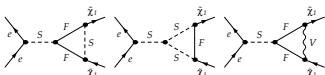
ISR in its full beauty (Skrzypek ea, 91)

$$\begin{aligned}
\Gamma_{ee}^{LL}(x, Q^2) = & \frac{\exp(-\frac{1}{2}\eta\gamma_E + \frac{3}{8}\eta)}{\Gamma(1 + \frac{\eta}{2})} \frac{\eta}{2} (1-x)^{(\frac{\eta}{2}-1)} \\
& - \frac{\eta}{4} (1+x) + \frac{\eta^2}{16} \left(-2(1-x) \log(1-x) - \frac{2 \log x}{1-x} + \frac{3}{2} (1+x) \log x - \frac{x}{2} \right. \\
& - \left. \frac{5}{2} \right) + \left(\frac{\eta}{2} \right)^3 \left[-\frac{1}{2} (1+x) \left(\frac{9}{32} - \frac{\pi^2}{12} + \frac{3}{4} \log(1-x) + \frac{1}{2} \log^2(1-x) \right. \right. \\
& - \left. \left. \frac{1}{4} \log x \log(1-x) + \frac{1}{16} \log^2 x - \frac{1}{4} \text{Li}_2(1-x) \right) \right. \\
& + \left. \frac{1}{2} \frac{1+x^2}{1-x} \left(-\frac{3}{8} \log x + \frac{1}{12} \log^2 x - \frac{1}{2} \log x \log(1-x) \right) \right. \\
& - \left. \frac{1}{4} (1-x) \left(\log(1-x) + \frac{1}{4} \right) + \frac{1}{32} (5-3x) \log x \right] ; \eta = \frac{2\alpha}{\pi} \left(\log \left(\frac{Q^2}{m_e^2} \right) - 1 \right)
\end{aligned}$$

Some NLO matrix elements

Some NLO matrix elements

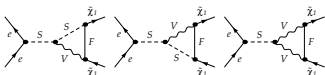
$$e e \rightarrow \tilde{\chi}_1 \tilde{\chi}_1$$



T1 G1 N1

T1 G2 N2

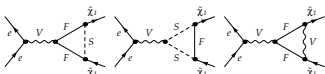
T1 G3 N3



T1 G4 N4

T1 G5 N5

T1 G6 N6

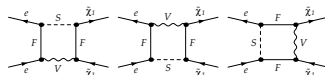


T1 G7 N7

T1 G8 N8

T1 G9 N9

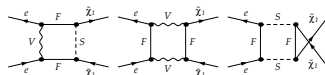
$$e e \rightarrow \tilde{\chi}_1 \tilde{\chi}_1$$



T5 G3 N37

T5 G4 N38

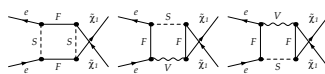
T5 G5 N39



T5 G6 N40

T5 G7 N41

T6 G1 N42



T6 G2 N43

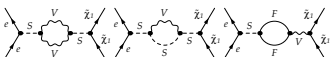
T6 G3 N44

T6 G4 N45

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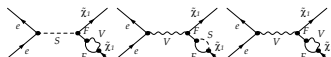


T10 G4 N64

T10 G5 N65

T10 G6 N66

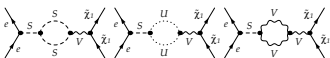
$$e e \rightarrow \tilde{\chi}_1 \tilde{\chi}_1$$



T11 G2 N82

T11 G3 N83

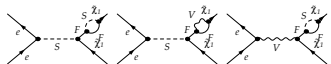
T11 G4 N84



T10 G7 N67

T10 G8 N68

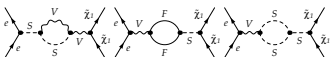
T10 G9 N69



T12 G1 N85

T12 G2 N86

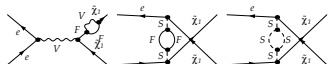
T12 G3 N87



T10 G10 N70

T10 G11 N71

T10 G12 N72



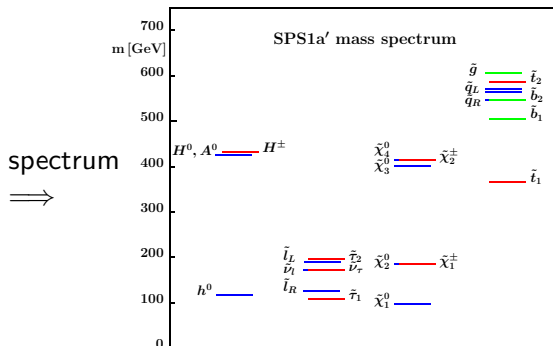
T12 G4 N88

T13 G1 N89

T13 G2 N90

Point SPS1a'

- mSUGRA scenario
- according to Snowmass Points (Allanach et al, 02), in agreement with cosmology data/ WMAP ($\tilde{\chi}_1^0$ as DM candidate)



light sleptons
heavy squarks
some light $\tilde{\chi}$ s
all masses < 1 TeV