

# MSSM PRECISION PHYSICS AT THE Z RESONANCE

Arne M. Weber

Max-Planck-Institute for Physics  
(Werner-Heisenberg-Institute)



MAX-PLANCK-GESellschaft

[based on a collaboration with: S.Heinemeyer, W.Hollik, G.Weiglein]

SUSY '07

# OUTLINE

- 1 WHAT, WHY, HOW?
- 2 STATUS
- 3 RESULTS
- 4 CONCLUSIONS

# What, why, how?

# Electroweak precision calculations, why bother?

## Electroweak precision calculations, why bother?

	central value	absolute error	relative error
$\sin^2 \theta_{\text{eff}}$	0.23153	$\pm 0.00016$ ILC: $\pm 0.000013$	$\pm 0.07\%$
$\Gamma_Z [\text{GeV}]$	2.4952 [GeV]	$\pm 0.0023$ ILC: $\pm 0.001$	$\pm 0.09\%$
$M_W [\text{GeV}]$	80.398	$\pm 0.025$ TEV/LHC: $\pm 0.020/0.015$ ILC: $\pm 0.007$	$\pm 0.03\%$

- Negligible errors for  $\alpha$ ,  $G_F$ ,  $M_Z$ , ...

## Electroweak precision calculations, why bother?

	central value	absolute error	relative error
$\sin^2 \theta_{\text{eff}}$	0.23153	$\pm 0.00016$ ILC: $\pm 0.000013$	$\pm 0.07\%$
$\Gamma_Z [\text{GeV}]$	2.4952 [GeV]	$\pm 0.0023$ ILC: $\pm 0.001$	$\pm 0.09\%$
$M_W [\text{GeV}]$	80.398	$\pm 0.025$ TEV/LHC: $\pm 0.020/0.015$ ILC: $\pm 0.007$	$\pm 0.03\%$

- Negligible errors for  $\alpha$ ,  $G_F$ ,  $M_Z$ , ...

⇒ Precise predictions needed to match this accuracy.

## Electroweak precision calculations, why bother?

	central value	absolute error	relative error
$\sin^2 \theta_{\text{eff}}$	0.23153	$\pm 0.00016$ ILC: $\pm 0.000013$	$\pm 0.07\%$
$\Gamma_Z [\text{GeV}]$	2.4952 [GeV]	$\pm 0.0023$ ILC: $\pm 0.001$	$\pm 0.09\%$
$M_W [\text{GeV}]$	80.398	$\pm 0.025$ TEV/LHC: $\pm 0.020/0.015$ ILC: $\pm 0.007$	$\pm 0.03\%$

- Negligible errors for  $\alpha$ ,  $G_F$ ,  $M_Z$ , ...

⇒ Precise predictions needed to match this accuracy.

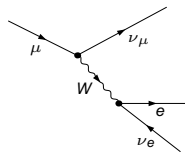
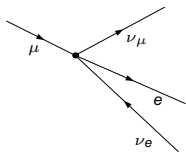
- Theory predictions sensitive to new physics via quantum corrections.

⇒ New physics can already be observed.

# $W$ boson mass from $\mu$ -decay

*Born level*

$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2}$$

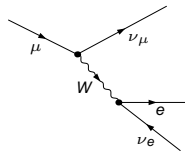
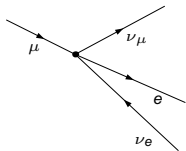




# $W$ boson mass from $\mu$ -decay

Born level

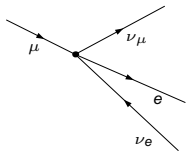
$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2}$$



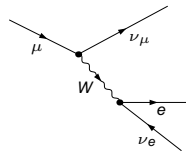
- Need to take radiative corrections into account.
- Summarise electroweak radiative corrections by  $\Delta r$ .  
[Marciano, Sirlin]

# $W$ boson mass from $\mu$ -decay

Born level



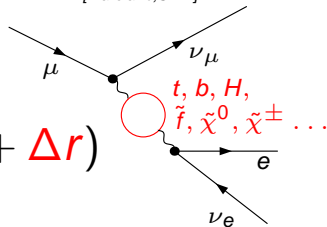
$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2}$$



- Need to take radiative corrections into account.
- Summarise electroweak radiative corrections by  $\Delta r$ .  
[Marciano, Sirlin]

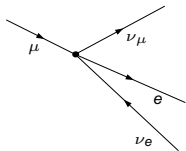
Loop order

$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2} (1 + \Delta r)$$

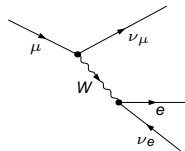


# $W$ boson mass from $\mu$ -decay

Born level



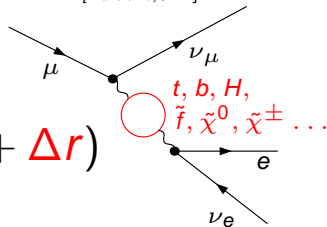
$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2}$$



- Need to take radiative corrections into account.
- Summarise electroweak radiative corrections by  $\Delta r$ .  
[Marciano, Sirlin]

Loop order

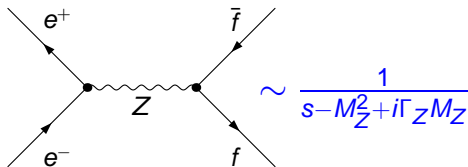
$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8\left(1 - \frac{M_W^2}{M_Z^2}\right)M_W^2} (1 + \Delta r)$$



$$\Delta r = \Delta r(M_W, M_Z, m_t, \alpha, \alpha_s, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^{0,\pm}} \dots)$$

## Z pole observables

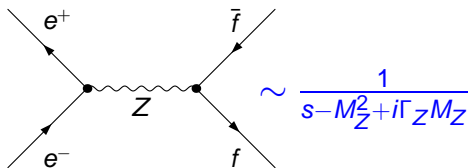
- LEP1 / SLC / GigaZ:  $e^+ e^- \rightarrow f \bar{f} @ s \sim M_Z^2$ .



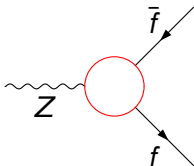
- Radiative corrections can be absorbed into effective couplings (up to small non-resonant contributions).

## Z pole observables

- LEP1 / SLC / GigaZ:  $e^+e^- \rightarrow f\bar{f}$  @  $s \sim M_Z^2$ .



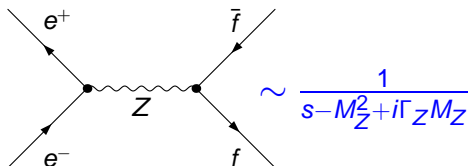
- Radiative corrections can be absorbed into effective couplings (up to small non-resonant contributions).



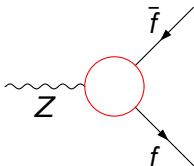
$$\Rightarrow \mathcal{M}_{\text{eff}} = \bar{u}_f \gamma_\alpha [\mathcal{G}_V^{\text{eff}} - \gamma_5 \mathcal{G}_A^{\text{eff}}] v_f \epsilon_Z^\alpha$$

## Z pole observables

- LEP1 / SLC / GigaZ:  $e^+ e^- \rightarrow f \bar{f} @ s \sim M_Z^2$ .



- Radiative corrections can be absorbed into effective couplings (up to small non-resonant contributions).



$$\Rightarrow \mathcal{M}_{\text{eff}} = \bar{u}_f \gamma_\alpha [\mathcal{G}_V^{\text{eff}} - \gamma_5 \mathcal{G}_A^{\text{eff}}] v_f \epsilon_Z^\alpha$$

$$\mathcal{G}_{V,A}^{\text{eff}} = \mathcal{G}_{V,A}^{\text{eff}}(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^{0,\pm}} \dots)$$

- Pseudo observables *defined* in terms of  $g_{\{V,A\}}^{\text{eff}}$ .

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .
- Effective mixing angles:

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right)$$



- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .
- Effective mixing angles:

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .
- Effective mixing angles:

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

$\updownarrow$   
 $\Delta\kappa(m_t, \alpha, M_h, M_A, m_{\tilde{t}}, m_{\tilde{\chi}^{0,\pm}} \dots)$

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .

- Effective mixing angles:

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

$$s_w^2 = 1 - \frac{M_W^2(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)}{M_Z^2}$$

$$\Delta\kappa(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)$$

$\updownarrow$  (between  $M_Z^2$  and  $\Delta\kappa$ )  
 $\updownarrow$  (between  $\Delta\kappa$  and  $\mathcal{G}_A^{\text{eff}}$ )

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .

- Effective mixing angles:

$$s_w^2 = 1 - \frac{M_W^2(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)}{M_Z^2}$$

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

$$\Delta\kappa(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)$$

- Partial widths:

$$\Gamma_f := N_c^f \frac{\alpha}{3} M_Z (|\mathcal{G}_V^{\text{eff}}|^2 R_V^{\text{QED}} + |\mathcal{G}_A^{\text{eff}}|^2 R_A^{\text{QED}})$$

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .

- Effective mixing angles:

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

$$s_w^2 = 1 - \frac{M_W^2(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)}{M_Z^2}$$

$$\Delta\kappa(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)$$

$\updownarrow$  (between  $M_Z^2$  and  $\Delta\kappa$ )  
 $\updownarrow$  (between  $\Delta\kappa$  and  $\mathcal{G}_A^{\text{eff}}$ )

- Partial widths:

$$\Gamma_f := N_c^f \frac{\alpha}{3} M_Z (|\mathcal{G}_V^{\text{eff}}|^2 R_V^{\text{QED}} + |\mathcal{G}_A^{\text{eff}}|^2 R_A^{\text{QED}})$$

$$\equiv N_c^f \bar{\Gamma}_0 |\rho_f| (4(I_3^f - 2Q_f s_w^2 |k_f|)^2 R_V^{\text{QED}} + R_A^{\text{QED}})$$

- Pseudo observables *defined* in terms of  $\mathcal{G}_{\{V,A\}}^{\text{eff}}$ .

- Effective mixing angles:  $s_w^2 = 1 - \frac{M_W^2(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)}{M_Z^2}$

$$\sin^2 \theta_{\text{eff}}^f := \frac{1}{4|Q_f|} \left( 1 - \text{Re} \frac{\mathcal{G}_V^{\text{eff}}}{\mathcal{G}_A^{\text{eff}}} \right) \equiv (1 + \Delta\kappa) s_w^2$$

$\Delta\kappa(m_t, \alpha, M_h, M_A, m_{\tilde{f}}, m_{\tilde{\chi}^0, \pm} \dots)$

- Partial widths:

$$\begin{aligned} \Gamma_f &:= N_c^f \frac{\alpha}{3} M_Z (|\mathcal{G}_V^{\text{eff}}|^2 R_V^{\text{QED}} + |\mathcal{G}_A^{\text{eff}}|^2 R_A^{\text{QED}}) \\ &\equiv N_c^f \bar{\Gamma}_0 |\rho_f| (4(I_3^f - 2Q_f s_w^2 |k_f|)^2 R_V^{\text{QED}} + R_A^{\text{QED}}) \end{aligned}$$

- Total Z width  $\Gamma_Z$ , pole asymmetries  $\mathcal{A}_{FB,LR}^0$ , peak cross sections  $\sigma_{had,l}^0$ , ratios of partial widths  $R_b, R_l \dots$  determined by  $\sin^2 \theta_{\text{eff}}^f$  and  $\Gamma_f$ .

# Status

## Status of $M_W$ & $\sin^2 \theta_{\text{eff}}$ calculation

### Standard Model:

- Full electroweak calculation up to two-loop order, full  $\mathcal{O}(\alpha\alpha_s)$ ,  $\mathcal{O}(\alpha\alpha_s^2)$  result.
- Leading universal terms in the relevant parameters at  $\mathcal{O}(\alpha_s G_F^2 m_t^4)$ ,  $\mathcal{O}(G_F^3 m_t^6)$ ,  $\mathcal{O}(\alpha\alpha_s^3)$ .

[Freitas, Hollik, Walter, Weiglein], [Awramik, Czakon], [Onishchenko, Veretin],

[Awramik, Czakon, Freitas, Weiglein],[Hollik, Meier, Uccirati],[Chetyrkin, Kühn, Steinhauser],

[Faisst, Kühn, Seidensticker, Veretin], [Schröder, Steinhauser],[Chetyrkin, Faisst, Kühn, Maierhöfer, Sturm], *and many more!*



## Status of $M_W$ & $\sin^2 \theta_{\text{eff}}$ calculation

### Standard Model:

- Full electroweak calculation up to two-loop order, full  $\mathcal{O}(\alpha\alpha_s)$ ,  $\mathcal{O}(\alpha\alpha_s^2)$  result.
- Leading universal terms in the relevant parameters at  $\mathcal{O}(\alpha_s G_F^2 m_t^4)$ ,  $\mathcal{O}(G_F^3 m_t^6)$ ,  $\mathcal{O}(\alpha\alpha_s^3)$ .

[Freitas, Hollik, Walter, Weiglein], [Awramik, Czakon], [Onishchenko, Veretin],

[Awramik, Czakon, Freitas, Weiglein],[Hollik, Meier, Uccirati],[Chetyrkin, Kühn, Steinhäuser],

[Faisst, Kühn, Seidensticker, Veretin], [Schröder, Steinhäuser],[Chetyrkin, Faisst, Kühn, Maierhöfer, Sturm], *and many more!*

### MSSM in previous analyses:

- One-loop with assumptions (real parameters, . . . ).  
[Dabelstein, Hollik, Mööle]
- Leading SM and universal SUSY  $\mathcal{O}(\alpha\alpha_s)$  terms.  
[Dabelstein, Hollik] & [Heinemeyer, Weiglein]

***New:***

*[Heinemeyer, Hollik, Stöckinger, Weiglein, AMW] – JHEP08(2006)052 & forthcoming pub.*

- Independent MSSM one-loop calculation for  $M_W$  & observables at the  $Z$  resonance.
  - No restrictions on MSSM parameters (complex parameters, no assumptions about SUSY breaking...).

***New:***

*[Heinemeyer, Hollik, Stöckinger, Weiglein, AMW] – JHEP08(2006)052 & forthcoming pub.*

- Independent MSSM one-loop calculation for  $M_W$  & observables at the  $Z$  resonance.
  - No restrictions on MSSM parameters (complex parameters, no assumptions about SUSY breaking...).
  - Higgs sector: Loop corrected masses and CP mixing.

*[Frank, Hahn, Heinemeyer, Hollik, Rzehak, Weiglein]*

## *New:*

*[Heinemeyer, Hollik, Stöckinger, Weiglein, AMW] – JHEP08(2006)052 & forthcoming pub.*

- Independent MSSM one-loop calculation for  $M_W$  & observables at the  $Z$  resonance.
  - No restrictions on MSSM parameters (complex parameters, no assumptions about SUSY breaking. . .).
  - Higgs sector: Loop corrected masses and CP mixing.  
*[Frank, Hahn, Heinemeyer, Hollik, Rzehak, Weiglein]*
  - Resummation of  $\tan \beta$  enhanced bottom Yukawa couplings ( $Z \rightarrow b\bar{b}$ ).  
*[Carena, Garcia, Nierste, Wagner]*

## New:

[Heinemeyer, Hollik, Stöckinger, Weiglein, AMW] – JHEP08(2006)052 & forthcoming pub.

- Independent MSSM one-loop calculation for  $M_W$  & observables at the  $Z$  resonance.
  - No restrictions on MSSM parameters (complex parameters, no assumptions about SUSY breaking...).
  - Higgs sector: Loop corrected masses and CP mixing.  
*[Frank, Hahn, Heinemeyer, Hollik, Rzehak, Weiglein]*
  - Resummation of  $\tan \beta$  enhanced bottom Yukawa couplings ( $Z \rightarrow b\bar{b}$ ).  
*[Carena, Garcia, Nierste, Wagner]*
  - Full one-loop calculation for  $Z \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$ .

**New:**

[Heinemeyer, Hollik, Stöckinger, Weiglein, AMW] – JHEP08(2006)052 & forthcoming pub.

- Independent MSSM one-loop calculation for  $M_W$  & observables at the  $Z$  resonance.
  - No restrictions on MSSM parameters (complex parameters, no assumptions about SUSY breaking...).
  - Higgs sector: Loop corrected masses and CP mixing.
 

[Frank, Hahn, Heinemeyer, Hollik, Rzehak, Weiglein]
  - Resummation of  $\tan \beta$  enhanced bottom Yukawa couplings ( $Z \rightarrow b\bar{b}$ ).
 

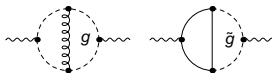
[Carena, Garcia, Nierste, Wagner]
  - Full one-loop calculation for  $Z \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0$ .



- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

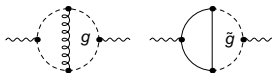
*[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]*



- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]



- Yukawa enhanced universal  $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$  terms.

[Haestier, Heinemeyer, Stöckinger, Weiglein]

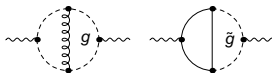




- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]



- Yukawa enhanced universal  $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$  terms.

[Haestier, Heinemeyer, Stöckinger, Weiglein]



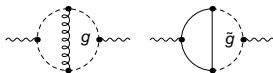
- Resummation of reducible one-loop contributions.

[Consoli, Hollik, Jegerlehner]

- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]



- Yukawa enhanced universal  $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$  terms.

[Haestier, Heinemeyer, Stöckinger, Weiglein]



- Resummation of reducible one-loop contributions.

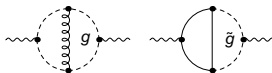
[Consoli, Hollik, Jegerlehner]

- Inclusion of full SM result.

- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]



- Yukawa enhanced universal  $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$  terms.

[Haestier, Heinemeyer, Stöckinger, Weiglein]



- Resummation of reducible one-loop contributions.

[Consoli, Hollik, Jegerlehner]

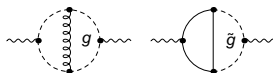
- Inclusion of full SM result.

⇒ Most precise & most general MSSM prediction for  $W$  boson mass and  $Z$  observables!

- All available beyond one-loop contributions from SM and MSSM incorporated.

- Leading universal  $\mathcal{O}(\alpha\alpha_s)$  terms.

[Djouadi, Gambino, Heinemeyer, Hollik, Jünger, Weiglein]



- Yukawa enhanced universal  $\mathcal{O}(\alpha_t^2, \alpha_t\alpha_b, \alpha_b^2)$  terms.

[Haestier, Heinemeyer, Stöckinger, Weiglein]



- Resummation of reducible one-loop contributions.

[Consoli, Hollik, Jegerlehner]

- Inclusion of full SM result.

⇒ Most precise & most general MSSM prediction for  
 $W$  boson mass and  $Z$  observables!

$$\delta M_W^{\text{th}} \approx 4 \dots 10 \text{ MeV}$$

$$\delta \sin^2 \theta_{\text{eff}}^{\text{th}} \approx (4.9 \dots 7.1) \times 10^{-5}$$

# Results

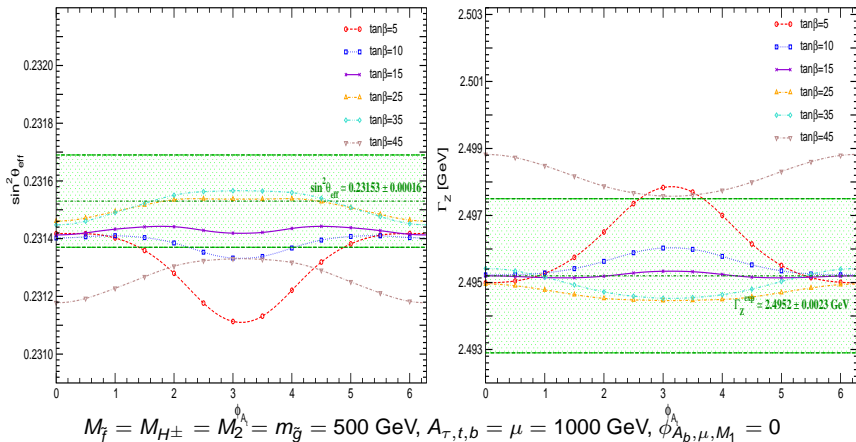
- $\mathcal{CP}$ -phases, in particular  $\phi_{\mu}$ , strongly constrained by EDMs.
  - However,  $\phi_{A_{t,b}}$  almost unconstrained. [Barger, Falk, Han, Jiang, Plehn]
- ⇒ Phases often assumed to be negligible/unobservable in EW physics.

■  $\mathcal{CP}$ -phases, in particular  $\phi_{\mu}$ , strongly constrained by EDMs.

■ However,  $\phi_{A_{t,b}}$  almost unconstrained. [Barger,Falk,Han,Jiang,Plehn]

⇒ Phases often assumed to be negligible/unobservable in EW physics.

### Effect of $\phi_{A_t}$ on $\sin^2 \theta_{\text{eff}}$ and $\Gamma_Z$



⇒ Large shifts can be induced by complex parameters.

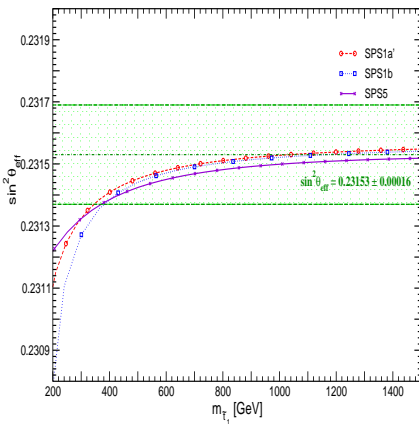
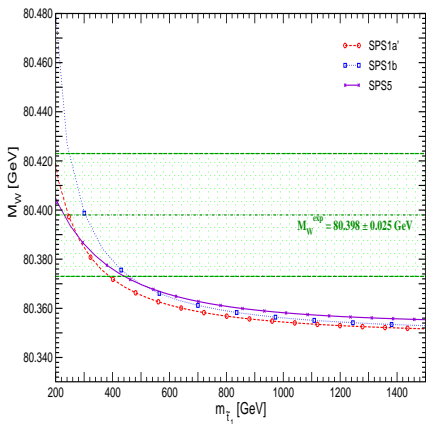
## SPS benchmark scenarios

- Benchmark points within “typical” constrained MSSM scenarios.
- SPS scenarios fix low-energy MSSM parameters.
- *here:*

$$\begin{aligned}
 M_{A^0} &= \text{scalefactor} \cdot M_{A^0}^{\text{SPS}}, & M_{\tilde{F}, \tilde{F}'} &= \text{scalefactor} \cdot M_{\tilde{F}, \tilde{F}'}^{\text{SPS}}, \\
 A_{t,b} &= \text{scalefactor} \cdot A_{t,b}^{\text{SPS}}, & \mu &= \text{scalefactor} \cdot \mu^{\text{SPS}}, \\
 M_{1,2,3} &= \text{scalefactor} \cdot M_{1,2,3}^{\text{SPS}}.
 \end{aligned}$$

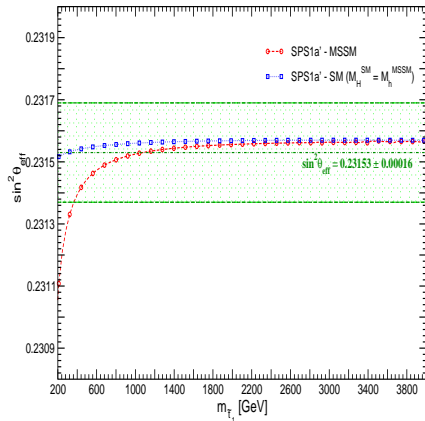
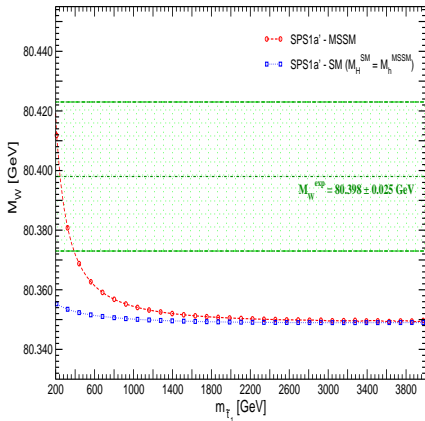


# $M_W$ and $\sin^2 \theta_{\text{eff}}$ for varied SUSY-scale



- $\Rightarrow$  Slight preference for light SUSY from  $M_W$ .
- $\Rightarrow$  No clear preference for light SUSY from  $\sin^2 \theta_{\text{eff}}$ .

# $M_W$ and $\sin^2 \theta_{\text{eff}}$ for varied SUSY-scale



- $\Rightarrow$  MSSM slightly favoured over Standard Model from  $M_W$ .
- $\Rightarrow$  No preference for MSSM from  $\sin^2 \theta_{\text{eff}}$ .

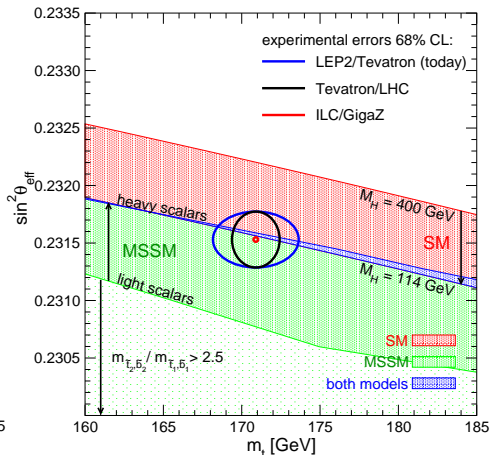
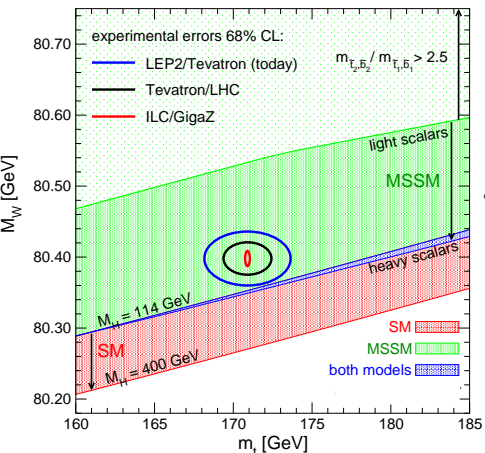
## Scatter plots for $M_W$ & $\sin^2 \theta_{\text{eff}}$

### ■ SUSY parameters:

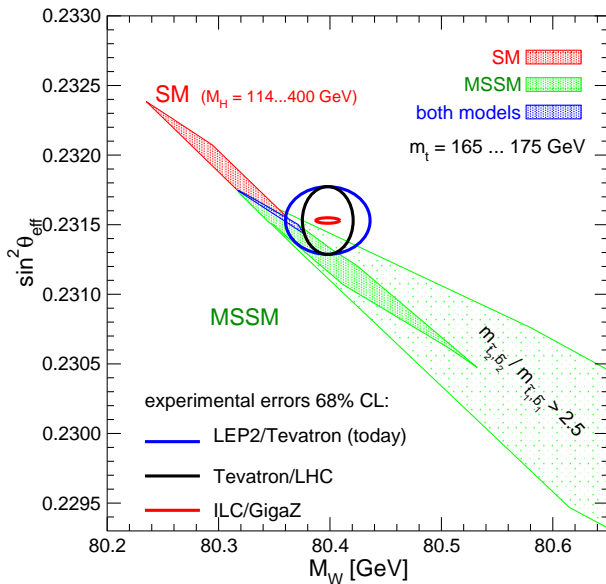
- sleptons :  $M_{\tilde{F}, \tilde{F}'} = 100 \dots 2000 \text{ GeV}$
- light squarks :  $M_{\tilde{F}, \tilde{F}'_{\text{up/down}}} = 100 \dots 2000 \text{ GeV}$
- $\tilde{t}/\tilde{b}$  doublet :  $M_{\tilde{F}, \tilde{F}'_{\text{up/down}}} = 100 \dots 2000 \text{ GeV}$
- $A_{t,b} = -2000 \dots 2000 \text{ GeV}$
- gauginos :  $M_{1,2} = 100 \dots 2000 \text{ GeV}$
- $m_{\tilde{g}} = 195 \dots 1500 \text{ GeV}$
- $\mu = -2000 \dots 2000 \text{ GeV}$
- Higgs :  $M_A = 90 - 1000 \text{ GeV}$
- $\tan \beta = 1.1 \dots 60$

- Unconstrained scan, only Higgs mass required to be in agreement with LEP data.

# $M_W(m_t)$ and $\sin^2 \theta_{\text{eff}}(m_t)$ in the MSSM



- $\Rightarrow$  Preference of MSSM over SM from  $M_W$ .
- $\Rightarrow$  MSSM and SM equally good for  $\sin^2 \theta_{\text{eff}}$ .



⇒ Combination of  $M_W$  and  $\sin^2 \theta_{\text{eff}}$  slightly favours MSSM.

# Conclusions

- Why precision physics?
  - Match experimental accuracy.
  - Test Standard Model and its extensions.

- Why precision physics?

- Match experimental accuracy.
- Test Standard Model and its extensions.

⇒ Most precise & most general MSSM predictions for  $W$  boson mass and  $Z$  observables.



- Why precision physics?
  - Match experimental accuracy.
  - Test Standard Model and its extensions.
- ⇒ Most precise & most general MSSM predictions for  $W$  boson mass and  $Z$  observables.
- $Z$  Observables and  $M_W$  alone:
  - Slight preference for light SUSY.
  - Slight preference for MSSM over Standard Model.

- Why precision physics?
    - Match experimental accuracy.
    - Test Standard Model and its extensions.
  - ⇒ Most precise & most general MSSM predictions for  $W$  boson mass and  $Z$  observables.
  - $Z$  Observables and  $M_W$  alone:
    - Slight preference for light SUSY.
    - Slight preference for MSSM over Standard Model.
  - Electroweak observables and dark matter constraints:
    - Best fit for to current data for low mass scales.  
*[Ellis, Heinemeyer, Olive, AMW, Weiglein ]*
- ⇒ *G. Weiglein, SUSY 07*

## ■ Why precision physics?

- Match experimental accuracy.
- Test Standard Model and its extensions.

⇒ Most precise & most general MSSM predictions for  $W$  boson mass and  $Z$  observables.

## ■ $Z$ Observables and $M_W$ alone:

- Slight preference for light SUSY.
- Slight preference for MSSM over Standard Model.

## ■ Electroweak observables and dark matter constraints:

- Best fit for to current data for low mass scales.

[Ellis, Heinemeyer, Olive, AMW, Weiglein ]

⇒ *G. Weiglein, SUSY 07*

## ■ Global mSUGRA fits:

- Preference for light SUSY for "natural" priors.

[Allanach, Cranmer, Lester, AMW]

⇒ *B.C. Allanach, SUSY 07*

## Outlook & current projects

- Extend computation and analysis to non-minimal models.
- Preparation of public computer code “SUSY-POPE”.
- LHC precision analysis projects:
  - “KISMET”  
*[Allanach, Cranmer, Lester, AMW], see B.C. Allanach SUSY '07*
  - “Mastercode”  
*[Buchmüller, Cavanaugh, Heinemeyer, Isidori, Paradisi, Ronga, AMW, Weiglein], see O. Buchmüller/F. Ronga SUSY '07*