

ADD gravitons & black holes at the LHC

Sezen Sekmen
METU, CMS



SUSY07, Karlsruhe,
25/07 – 01/08/2007
Exotics WG

CONTENTS

2. ADD gravitons

1.1. ADD graviton emission

1.2. Virtual ADD graviton exchange

2. Black holes

2.1. Production and decay

2.2. LHC signatures and significance

2.3. BH reconstruction

2.4. Learning about spacetime



ADD gravitons - introduction

- Arkani-Hamed, Dimopoulos, Dvali model introduces n compactified LEDs with radius R to solve the Hierarchy problem where gravitation unifies with the other forces at a lowered, \sim TeV order scale

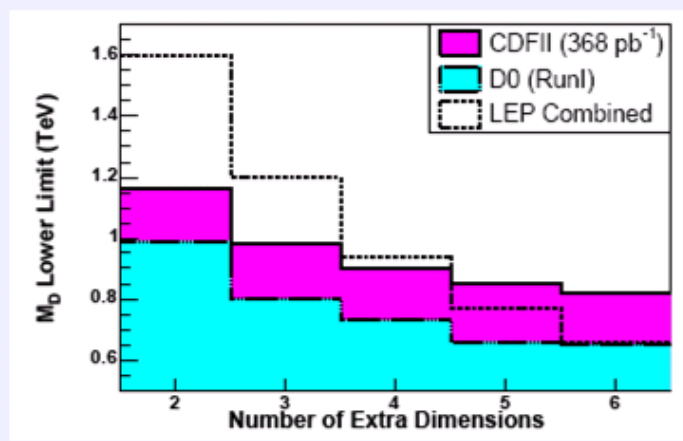
$$M_4^2 = M_{4+n}^{2+n} R^n$$

- Only gravitons probe the bulk space while others are confined to the 3-brane. Hence, a KK tower of graviton modes: 0-mode is massless; higher modes are massive spin-2 particles that can couple to SM matter. Following consequences are possible:

- Direct graviton emission

$$q\bar{q}/gg \rightarrow g/\gamma G_{KK}^{ADD}, \quad gq \rightarrow q G_{KK}^{ADD}$$

LEP & Tevatron constraints on $M_{(4+n)}$ from jet+G channel



- Virtual graviton exchange:

$$q\bar{q} \rightarrow \gamma/Z^0 / G_{KK}^{ADD} \rightarrow ll/\gamma\gamma/jj$$

Find the deviation from SM caused by virtual graviton exchange

%95CL lower limits on $M_{(4+n)}$ for LED with 200pb⁻¹ D0 Run II diEM sample and the combined limit for Run I and II

	GRW	HLZ					Hewett	
		$n=2$	$n=3$	$n=4$	$n=5$	$n=6$	$n=7$	$\lambda=+1$
Run II	1.36	1.56	1.61	1.36	1.23	1.14	1.08	1.22
Run I+II, combined	1.43	1.67	1.70	1.43	1.29	1.20	1.14	1.28



ADD graviton emission in γ +MET channel: I



J.Weng, G.Quast, C.Saout, A.de Roeck, M.Spiropulu, CMS-NOTE 2006/129

$$q\bar{q}/gg \rightarrow \gamma G_{KK}^{ADD}$$

Signatures:

A single high- p_T photon in the central η region.

High missing p_T back-to-back with the photon in the azimuthal plane

Not strongly dependent on ADD model parameters.

Trigger: Single photon (L1+HLT), $E > 80\text{ GeV}$, %100

Event selection:

$E_{T\text{miss}} > 400\text{ GeV}$

$p_T > 400\text{ GeV}$

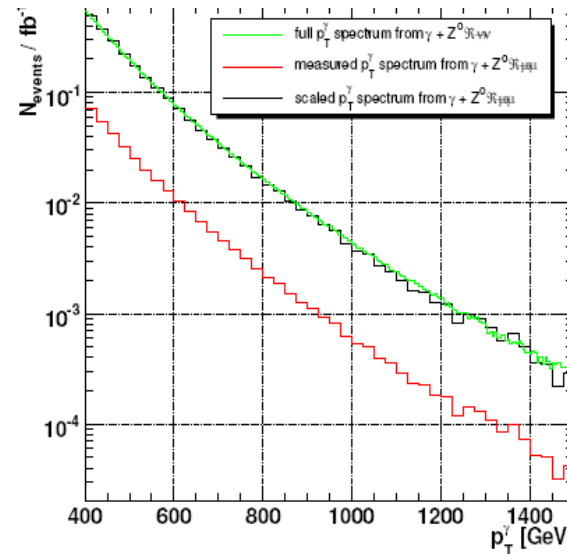
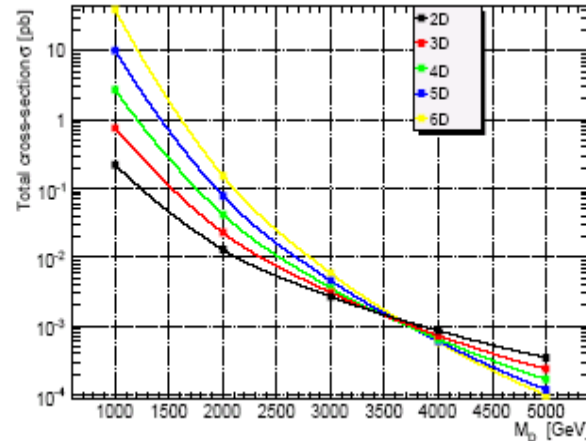
$\Delta\Phi(E_{T\text{miss}}, \gamma) > 2.5$

$|\eta| < 2.4$

Track veto $> 40\text{ GeV}$

γ likelihood > 0.2

Signal cross section: PYTHIA



Background	σ for $p_T > 400\text{ GeV}$
$Z^0\gamma \rightarrow \nu\bar{\nu} + \gamma$	2.16 fb
$W^\pm \rightarrow e\nu$	18.2 fb
$W^\pm \rightarrow \mu\nu$	18.2 fb
$W^\pm \rightarrow \tau\nu$	18.2 fb
$W^\pm\gamma \rightarrow e\nu + \gamma$	0.83 fb
γ +Jets	2.50 pb
QCD	2.15 nb
di- γ born	5.20 fb
di- γ box	0.14 fb
Z^0 + jets	0.69 pb

Estimation of γ +Z($\rightarrow \nu\nu$) background from γ +Z($\rightarrow \mu\mu$) processes.

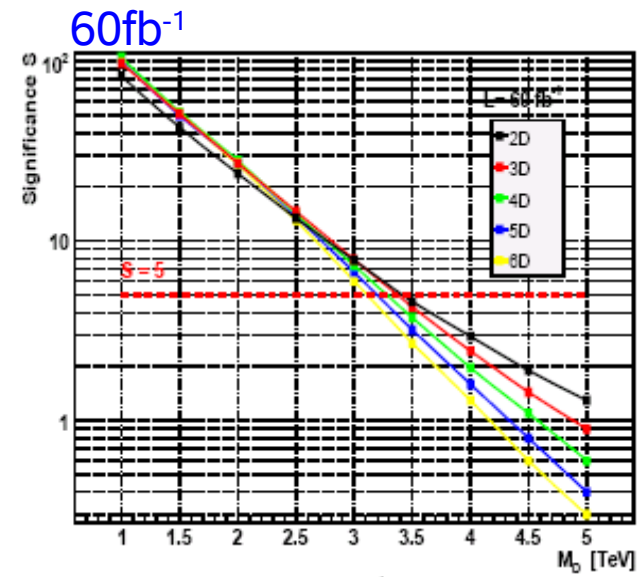
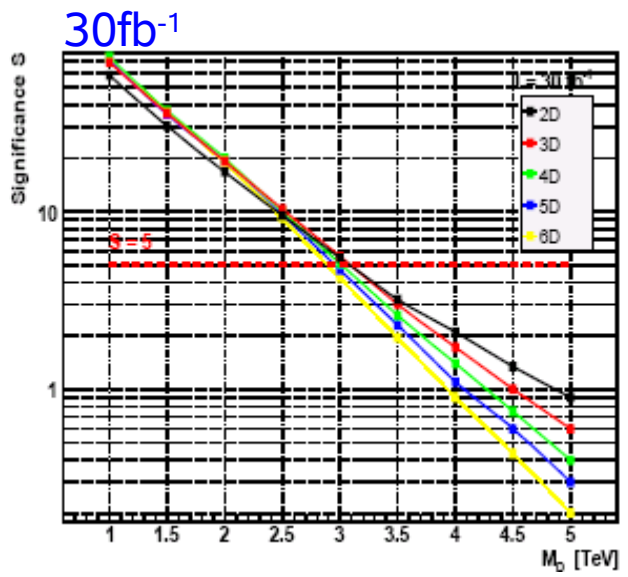
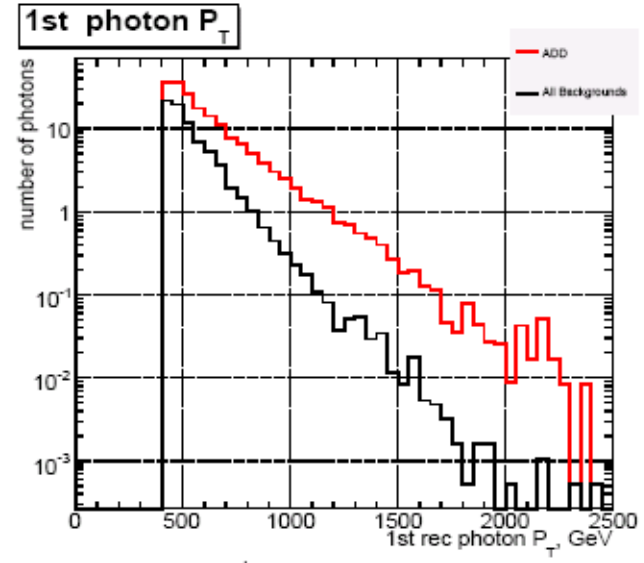
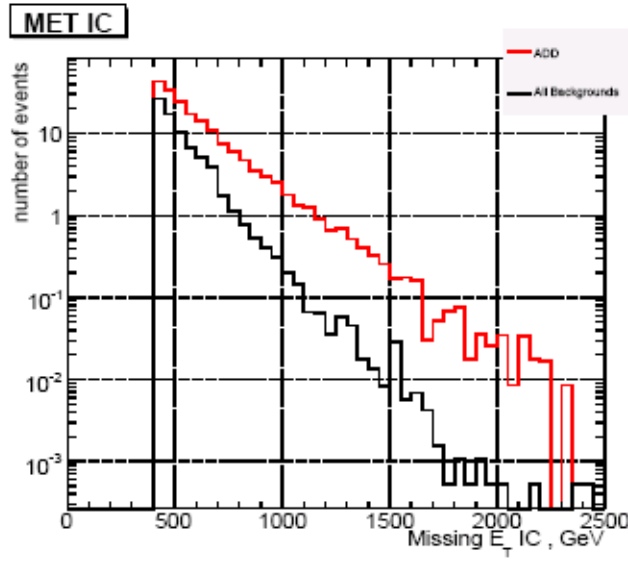


ADD graviton emission in γ +MET channel: II



$M_{4+n} = 2.5\text{TeV}$,
 $n = 2$
PYTHIA+CMKIN
+FAMOS

Significances:



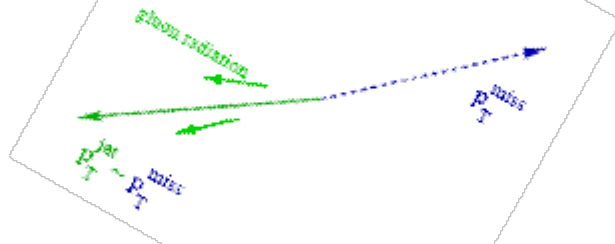
ADD graviton emission in jet+MET channel: I



L.Vacavant, I.Hinchliffe, ATLAS-PHYS-2000-016

$$qg \rightarrow qG_{KK}^{ADD}, \quad gg/q\bar{q} \rightarrow gG_{KK}^{ADD}$$

Signature: Monojet back-to-back in azimuth with missing p_T . Additional jets come from ISR/FSR



Trigger:

Single jet + MET

jet: $|\eta| < 3.2, p_T > 50$ GeV

$p_{T\text{miss}} > 50$ GeV (low L)

$p_{T\text{miss}} \rightarrow 100$ GeV for high L

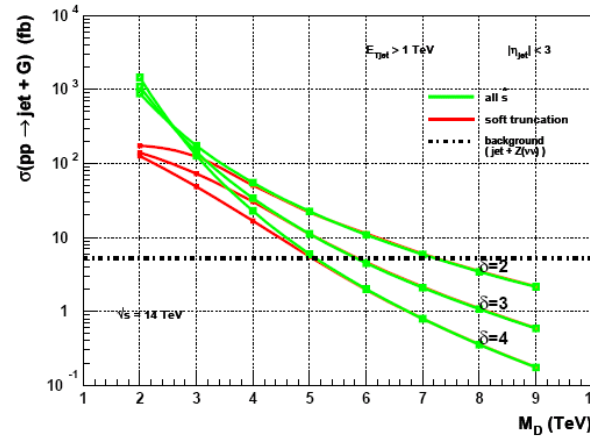
Event selection:

Lepton veto:

Electrons: $p_T > 5$ GeV, $|\eta| < 2.5$

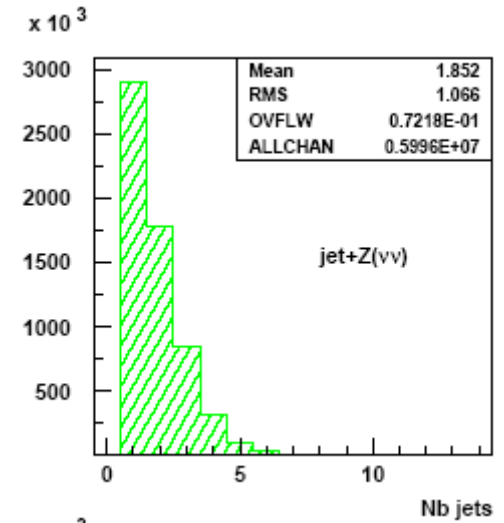
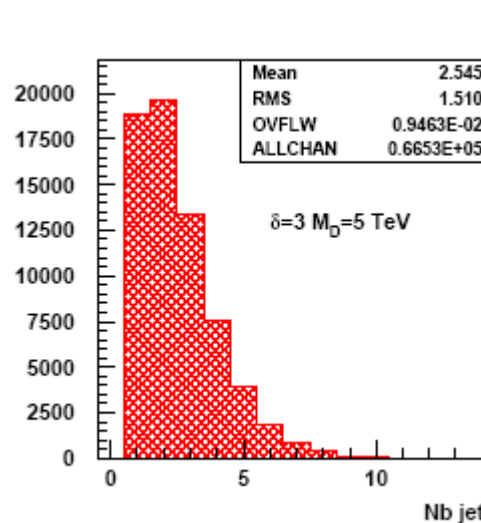
Muons: $p_T > 6$ GeV, $|\eta| < 2.5$

S&B cross sections for jets with $E_T > 1$ TeV -: ISAJET



Backgrounds:

- Jet + Z ($\rightarrow \nu\nu$)
- jet + W ($\rightarrow \tau\nu$)
- jet + W ($\rightarrow \mu\nu$)
- jet + W ($\rightarrow e\nu$)

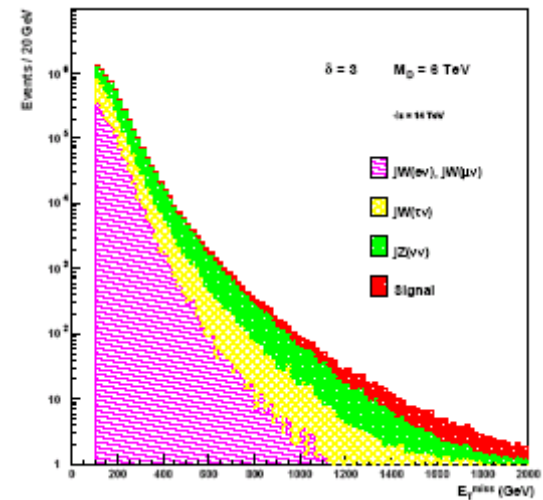
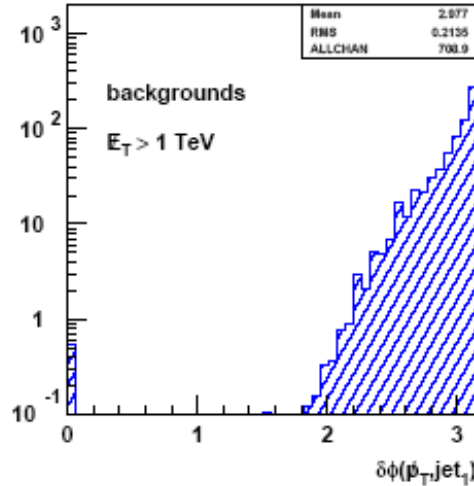
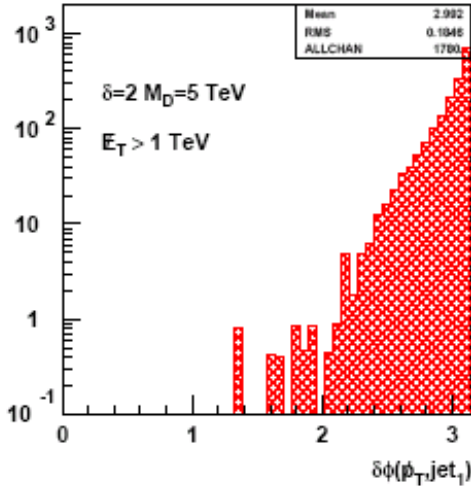


of jets: ISAJET + ATLFast

ADD graviton emission in jet+MET channel: II



ISAJET + ATLFAST



δ	M_D	Low luminosity, $30 fb^{-1}$			High luminosity, $100 fb^{-1}$		
		S	S/\sqrt{B}	$S/\sqrt{7B}$	S	S/\sqrt{B}	$S/\sqrt{7B}$
2	4	1036.4	81.6	30.8	3542.2	150.2	56.8
	5	417.0	32.9	12.4	1426.9	60.4	22.8
	6	205.9	16.3	6.2	700.6	29.6	11.2
	7	111.3	8.8	3.3	379.4	16.1	6.1
	8	65.3	5.2	2.0	222.5	9.4	3.5
3	4	641.8	50.6	19.1	2168.4	92.0	34.8
	5	211.5	16.6	6.3	706.0	30.0	11.3
	6	85.1	6.8	2.6	287.5	12.1	4.6
	7	39.3	3.1	1.2	134.0	5.7	2.2
4	4	436.2	34.3	13.0	1473.4	62.5	23.6
	5	113.0	8.8	3.3	383.4	16.3	6.2
	6	37.8	2.9	1.1	128.5	5.4	2.0

Remaining signal events after the selection cuts + $E_{T_{miss}} > 1$ TeV, and statistical significance

B estimation via normalization is 7 times smaller than expected due to calibration effects, hence $S/(7B)^{1/2}$ is also used as a worst case scenario.

ADD graviton contribution in dimuon channel: I

I.Belotelov, I.Golotvin, A.Lanyov, E.Rogalev, M.Savina, S.Shmatov, D.Bourikov, CMS-NOTE 2006/076

$$q\bar{q} \rightarrow \gamma/Z^0/G_{KK}^{ADD} \rightarrow \mu\mu$$

Signatures:
Two high- p_T muons.

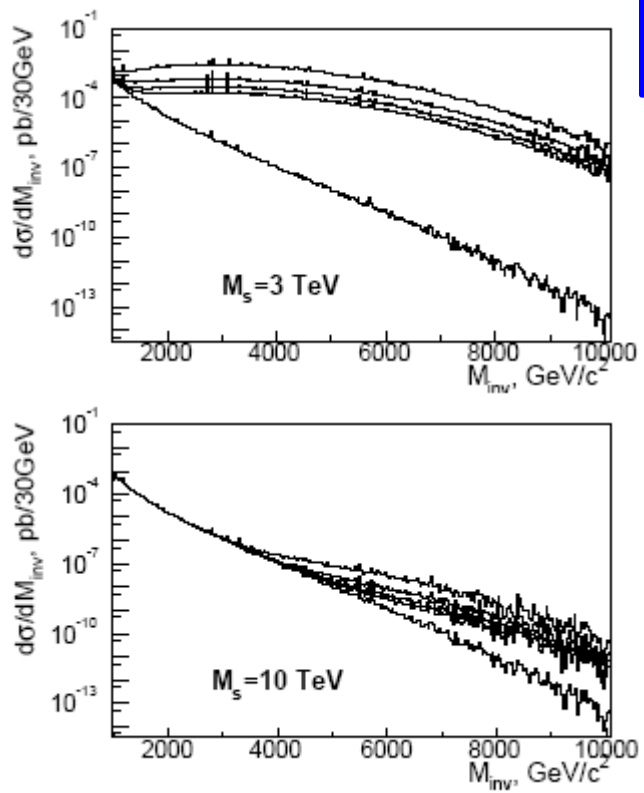
Trigger:
Single muon & dimuon (L1+HLT)
 $p_T > 7\text{ GeV}(\mu\mu), 19\text{ GeV}(\text{single } \mu)$
Efficiency > %98

Event selection:
 $M(\mu\mu)_{\text{inv}} > M_{\text{min}}$
Different M_{min} for different M_{4+n} :
 $M_{\text{min}} = 1\text{ TeV}$ for $M_{4+n} = 3\text{ TeV}$
 $M_{\text{min}} = 1.5\text{ TeV}$ for $M_{4+n} = 5\text{ TeV}$

LO DY cross section

$M_S, \text{TeV}/c^2$	3.0	4.0	5.0	7.0
$n = 3$	1.5×10^3	160	32.1	8.1
$n = 6$	103	11.4	10.1	6.4

Backgrounds:
For $M(\mu\mu)_{\text{inv}} > 1\text{ TeV}$
ZZ/WZ/WW:
 $\sigma = 2.59 \times 10^{-4} \text{ fb}^{-1}$
tt:
 $\sigma = 2.88 \times 10^{-4} \text{ fb}^{-1}$

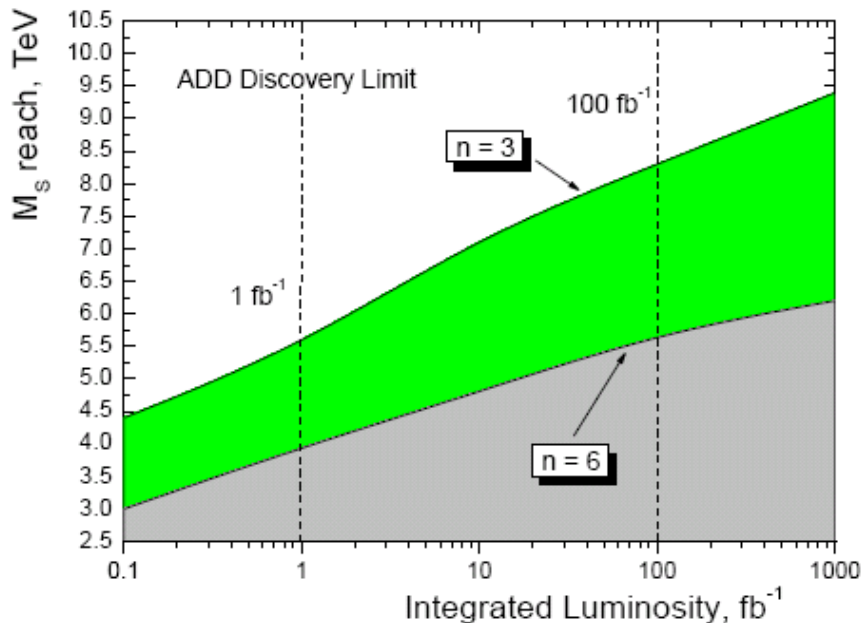
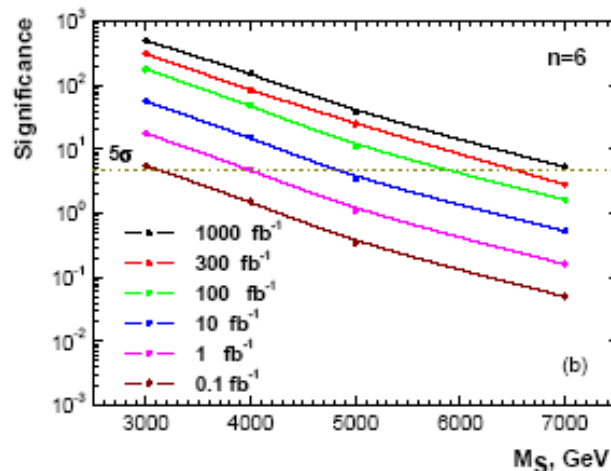
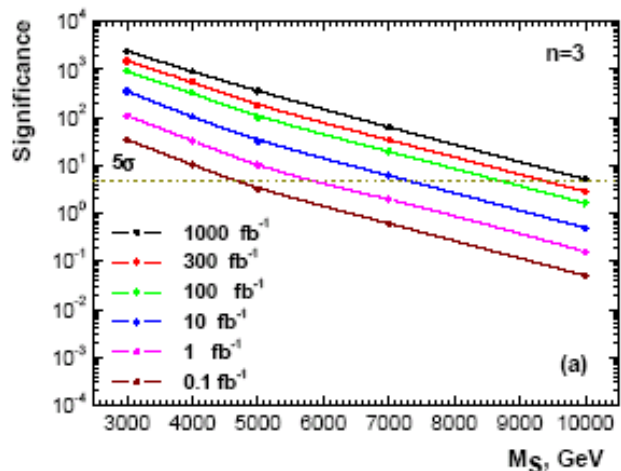


Landsberg code + STAGEN + PHYTIA + ORCA + OSCAR
Bottom to top:
SM, n = 6,5,4,3

ADD graviton contribution in dimuon channel: II



Significance:



5σ limit on M_{4+n} for the number of extra dimensions $n = 3, 4, 5, 6$

ADD graviton contribution in $\gamma\gamma/\ell\ell$ channels: I



V.Kabachenko, A. Miagkov, A. Zenin ATLAS-PHYS-2001-012

$$q\bar{q}/gg \rightarrow \gamma/Z^0 / G_{KK}^{ADD} \rightarrow \gamma\gamma/\ell\ell$$

Signature: Two photons or leptons

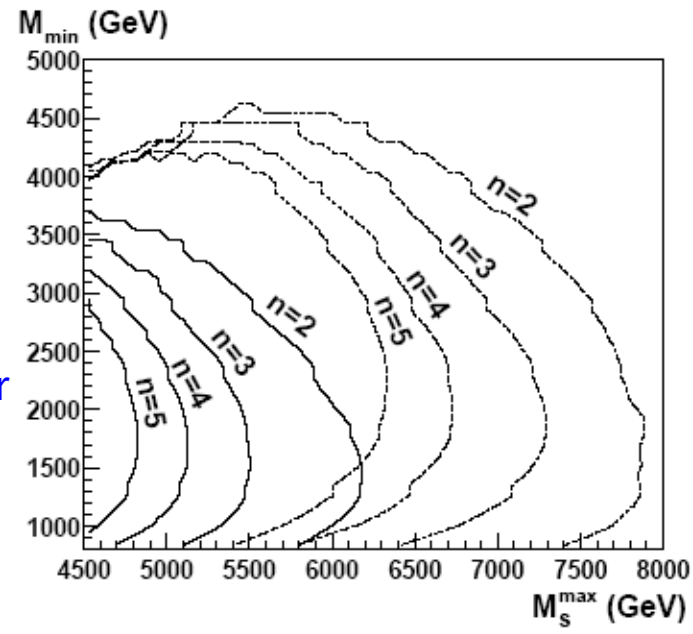
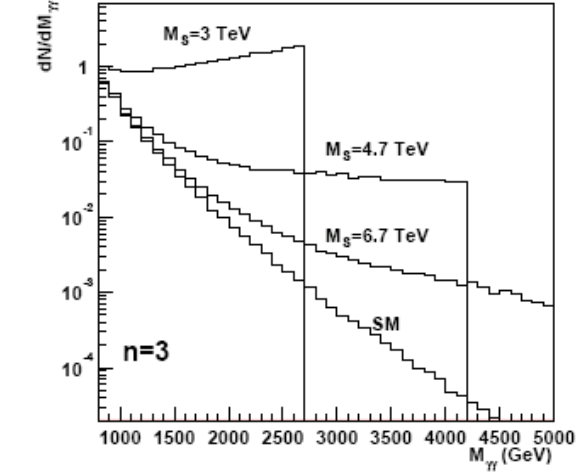
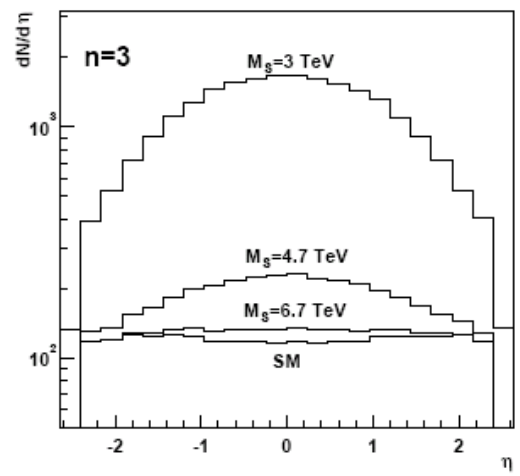
Trigger:

-

Event selection:

- $M(\gamma\gamma/\ell\ell)_{inv} < 0.9M_{4+n}$
- $M(\gamma\gamma/\ell\ell)_{inv} > M_{min} > 0.8TeV$
- $p_T(\gamma/\ell) > 50GeV$
- $|\eta| < 2.5$

ATLFAST



Diphoton production, maximal reach at 5σ level in the scale M_{4+n} as a function of the lower cut M_{min} . Solid lines: $10fb^{-1}$, dashed lines $100fb^{-1}$.

Higher-dimensional black holes - introduction

Schwarzschild radius

FORMATION

$4+n$ d

$$r_{4+n} = \frac{1}{\sqrt{\pi} M_{4+n}} \left(\frac{M_{BH}}{M_{4+n}} \left(\frac{8\Gamma((n+3)/2)}{n+2} \right) \right)^{\frac{1}{n+1}}$$

4d

$$r_4 = \frac{1}{M_4} \left(\frac{M_{BH}}{M_4} \right) = \frac{1}{M_{4+n}} \left(\frac{M_{BH}}{M_{4+n}} \right) \frac{1}{(M_{4+n} R)^n}$$

$$r_4 < r_{4+n} < R$$

4d

RADIATION

$4+n$ d

$$T_4 = \frac{dE}{dS}$$

$$T_4 = \frac{dE}{dS}$$

$$\tau_4 \sim \frac{1}{M_4} \left(\frac{M_{BH}}{M_4} \right)^3$$

Hawking temperature

$$T_4 > T_{4+n}$$

Lifetime

$$\tau_4 < \tau_{4+n}$$

$$T_{4+n} = \frac{dM}{dA}$$

$$T_{4+n} \sim M_{4+n} \left(\frac{M_{4+n}}{M_{BH}} \right)^{\frac{1}{n+1}}$$

$$\tau_{4+n} \sim \frac{1}{M_{4+n}} \left(\frac{M_{BH}}{M_{4+n}} \right)^{\frac{n+3}{n+1}}$$

BH production – cross sections

For $M_{BH} \gg M_{4+n}$, parton level BH formation cross section is given by semi-classical arguments through a geometrical approach:

$$\sigma(s = M_{BH}^2) = F_n \pi r_{4+n}^2 = \frac{F_n}{M_{4+n}^2} \left(\frac{M_{BH}}{M_{4+n}} \frac{8\Gamma((n+3)/2)}{n+2} \right)^{\frac{2}{n+1}}$$

where F_n is the formation factor. However, the definition of F_n is yet uncertain. σ_{BH} **dominates in high energies.**

For $M_{BH} \sim M_{4+n}$ quantum gravity effects come into play.

Differential production cross section in pp collisions is calculated using parton luminosity approach:

$$\frac{d\sigma(pp \rightarrow BH + X)}{dM_{BH}} = \frac{dL}{dM_{BH}} \sigma(ab \rightarrow BH)|_{\hat{s}=M_{BH}^2}$$
$$\frac{dL}{dM_{BH}} = \frac{2M_{BH}}{s} \sum_{a,b} \int_{M_{BH}^2/s}^1 \frac{dx_a}{x_a} f_a(x_a) f_b\left(\frac{M_{BH}^2}{sx_a}\right)$$

where $f_i(x_i)$ are the PDFs.

Hawking radiation and multiplicity

f: flux, N: multiplicity, $x = E/T_H$

$$\frac{df}{dx} \sim \frac{x^3}{e^x + c}$$

$$\frac{dN}{dE} \sim \frac{1}{E} \frac{df}{dE} \sim \frac{x^2}{e^x + c}$$

$$\left\langle \frac{1}{E} \right\rangle = \frac{1}{T_H} \frac{\int_0^\infty dx \frac{x^2}{e^x + c}}{\int_0^\infty dx \frac{1}{x} \frac{x^2}{e^x + c}} = \frac{a}{T_H}$$

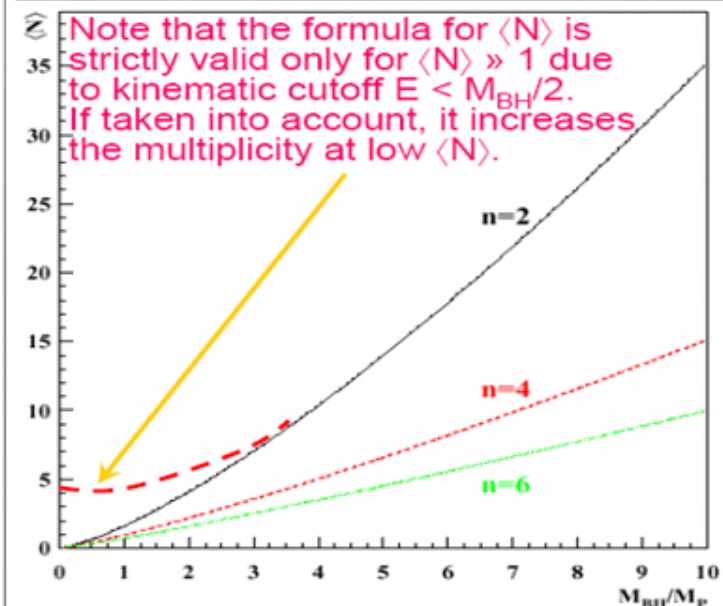
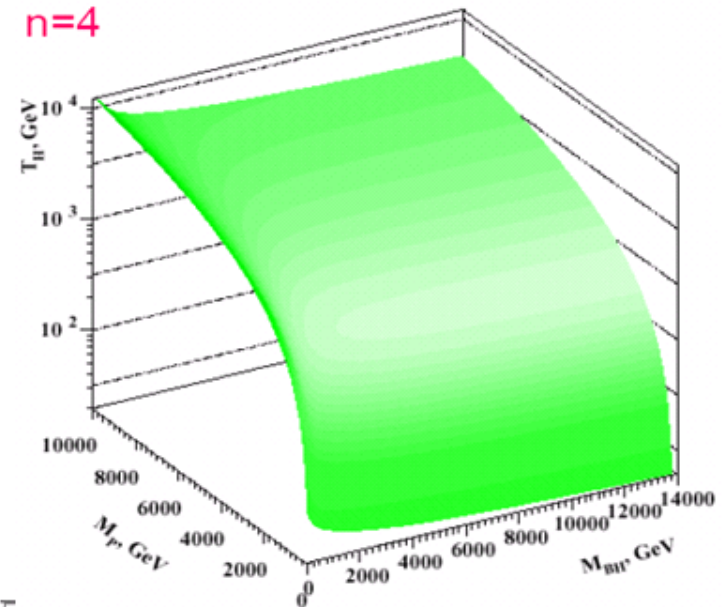
Use $a = 0.5$ for Boltzman statistics.

$$\langle N \rangle = \left\langle \frac{M_{BH}}{E} \right\rangle = \frac{M_{BH}}{2T_H}$$

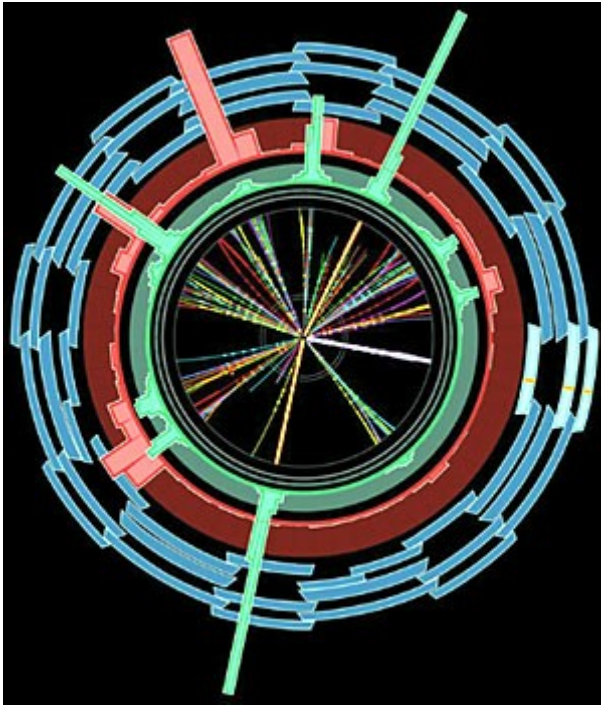
BHs decay democratically to all SM particles

Landsberg, Dimopoulos,
PRL 87, 161602 (2001)

$n=4$



Reconstructing BHs at the LHC



ATLAS

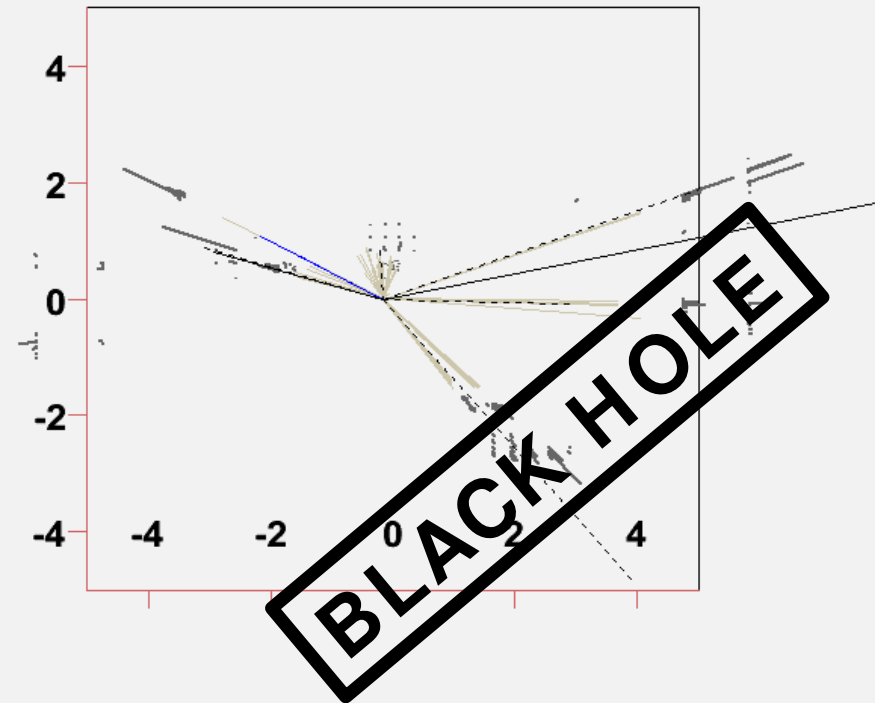
BH event generators:

- TRUNOIR (Landsberg, Dimopoulos)
- Unnamed MC by Japan group
- Charybdis (Harris, Richardson, Webber)
- Catfish (Cavaglia, Cremaldi, Godang, Summers)

Event 017:-1:2982.6%

H. Gamsızkan, CMS, Charybdis

$M_{4+n} = 2, n = 3$





BH signatures

H.Gamsızkan, A.de Roeck, S.Sekmen, M.T.Zeyrek, CMS-AN 2006/088

Signature: Due to thermal decay, BHs have high sphericity, high Σp_T and high multiplicity compared to SM events – and due to democratic decay, jets/leptons ratio $\sim 5 : 1$

Backgrounds: tt+jets, W/Z+jets, QCD, WW+jets

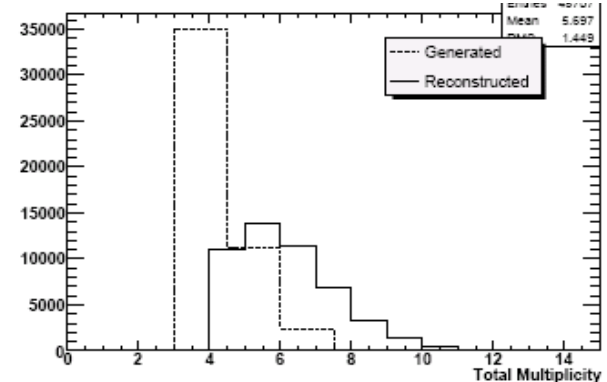
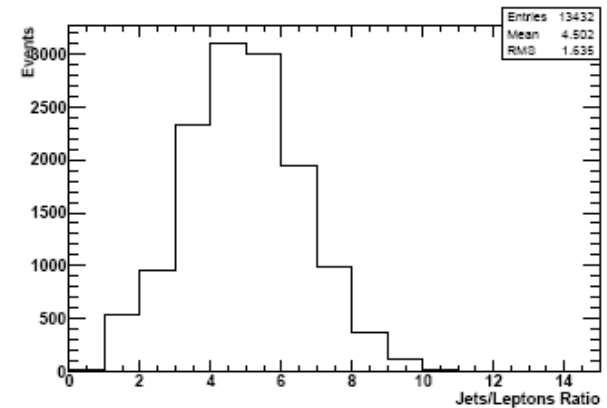
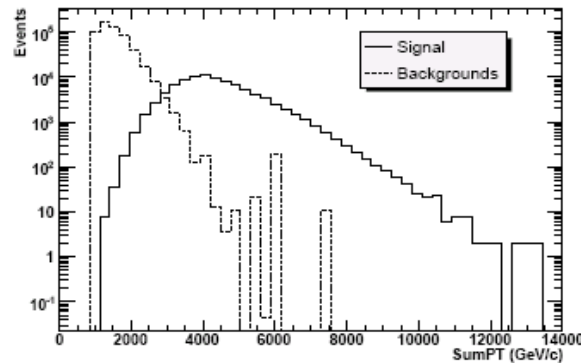
Trigger: >4 jets, softest having $E_T > 120$ GeV (efficiency $> \%93$). BHs can be tagged efficiently by requiring a prompt photon or lepton.

Event selection:
 $M_{BH}(\text{reco}) > 2$ TeV
Multiplicity > 4
 $E(x_i) < M_{BH}/2$
Sphericity < 0.28

$M_{4+n} = 2, n = 3, M_{BH} = 4-14,$
 $\sigma = 18.85$ pb
Charybdis+ORCA+OSCAR

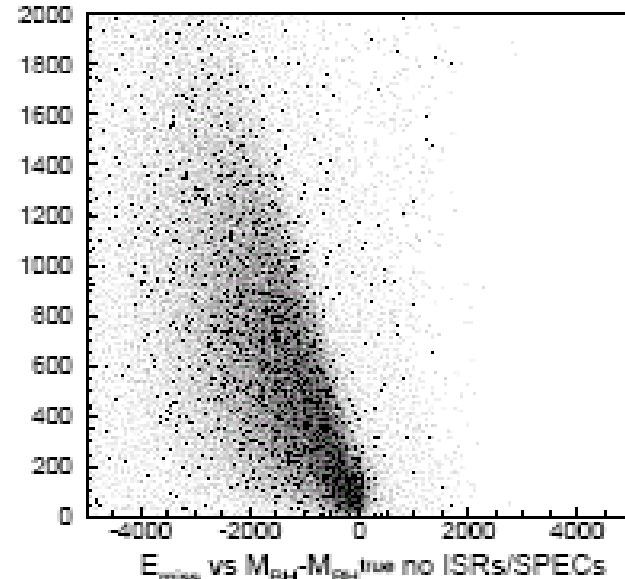
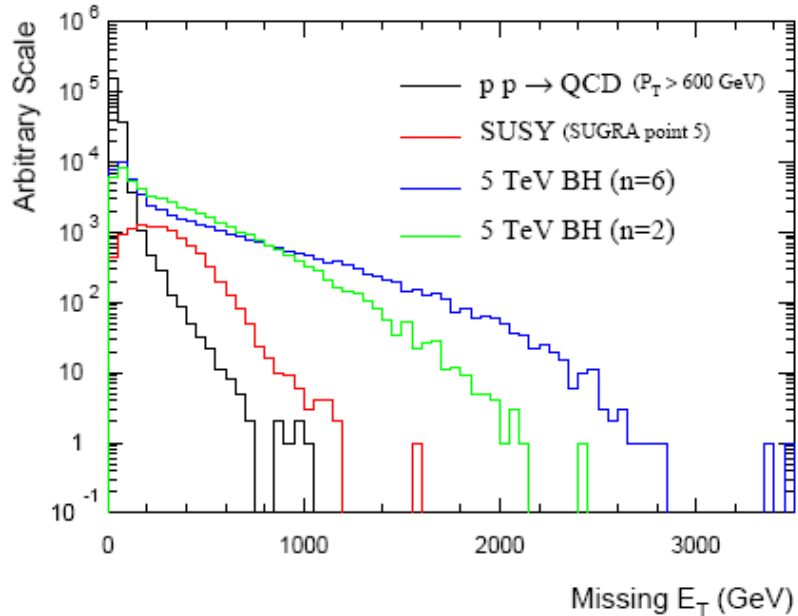
$$\sum p_T = \sum_{i < 4} |p_T(\text{jet}_i)|$$

(4 hardest jets)



BH signatures: High missing energy

BHs would have high missing E_T due to high energies of democratically emitted neutrinos.



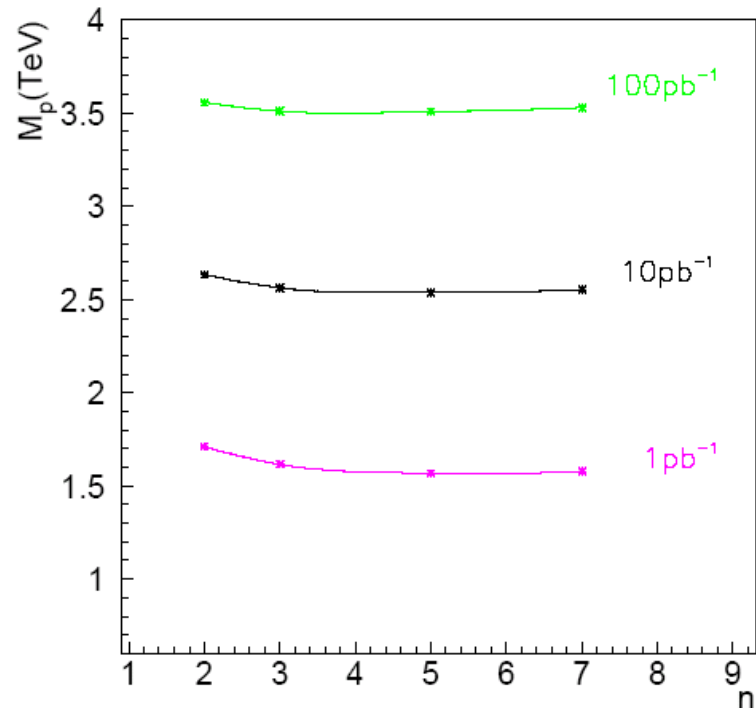
Harris, Palmer, Parker, Richardson,
Sabetfakhri, Webber, ATLAS, JHEP 0505
(2005) 053
Charybdis + ATLFAST

Tanaka, Yamamura, Asai, Kanzaki,
ATLAS, Eur. Phys. J. C41 (2005) 19-23
Unnamed MC
 $M_{4+n} = 1$, $n = 3$, $M_{BH} = 1-14$

Most analyses use $E_{T_{miss}} > 100$ GeV cut.

BH significance

$M_{(4+n)}$	min M_{BH}	4+n	Cross section (pb)	Luminosity to reach 5σ (fb^{-1})
2	4	7	18.85	3.71×10^{-3}
2.5	5	7	2.661	1.20
1	1	6	27980	6.50×10^{-6}
1	1	8	23510	8.50×10^{-6}
1	1	10	24910	8.50×10^{-6}
1	2	6	3086	1.95×10^{-5}
1	2	8	2282	2.95×10^{-5}
1	2	10	2284	3.15×10^{-5}
2	2	6	486.1	1.52×10^{-4}
2	2	8	432.4	1.82×10^{-4}
2	2	10	468.3	1.74×10^{-4}
2	3	6	93.06	5.79×10^{-4}
2	3	8	75.54	8.74×10^{-4}
2	3	10	78.7	8.88×10^{-4}
2	4	6	23.3	2.30×10^{-3}
2	4	10	17.73	5.29×10^{-3}
3	3	6	31.56	2.47×10^{-3}
3	3	8	28.55	2.96×10^{-3}
3	3	10	31.15	2.74×10^{-3}
3	4	6	7.904	1.45×10^{-2}
3	4	8	6.638	2.38×10^{-2}
3	4	10	7.017	2.32×10^{-2}
4	4	6	3.67	1.21
4	4	8	3.328	1.81
4	4	10	3.636	2.10



Tanaka, et.al., ATLAS,
 Eur. Phys. J. C41 (2005) 19-23
 Unnamed MC + ATLFAST
 $M_{\text{BH}} = M_{4+n} + 1$ to 14

Gamsızkan et.al., CMS-AN 2006/088,
 Charybdis + FAMOS

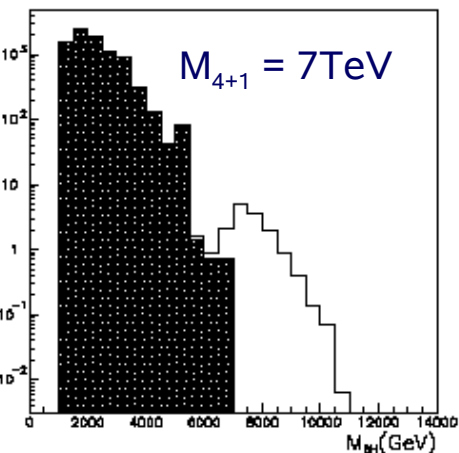
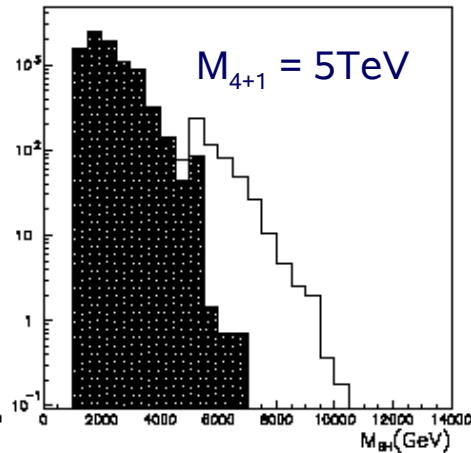
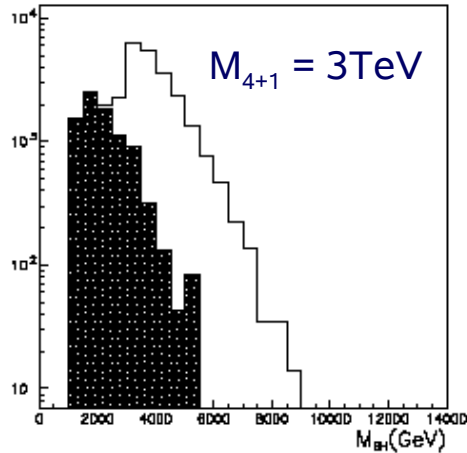
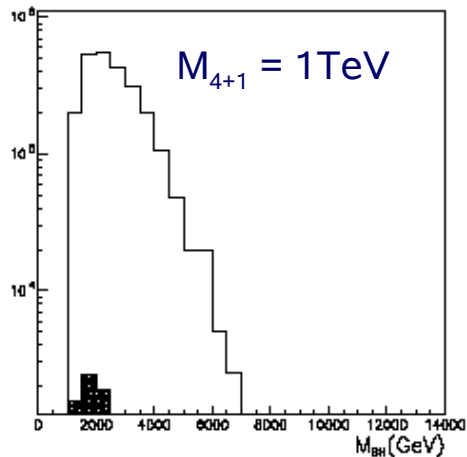
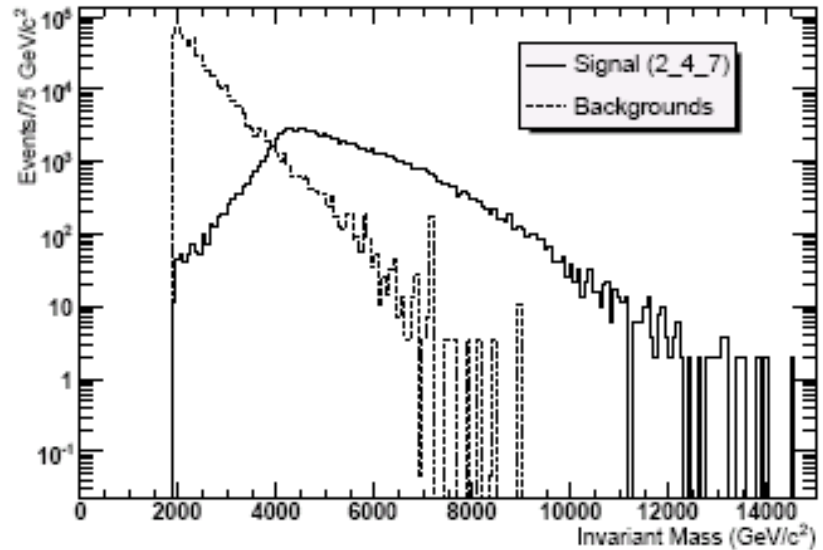


BH invariant mass

$$M_{BH} = \sqrt{(\sum_i p_i)^2}$$

Gamsızkan et.al., CMS-AN 2006/088
Charybdis + ORCA + OSCAR
 $M_{4+n} = 2, n = 3, M_{BH} = 4-14$

Tanaka et.al., ATLAS,
Unnamed MC,
 $n = 3, M_{BH} = 1-14,$
solid line: S+B, cross-hatched: B only



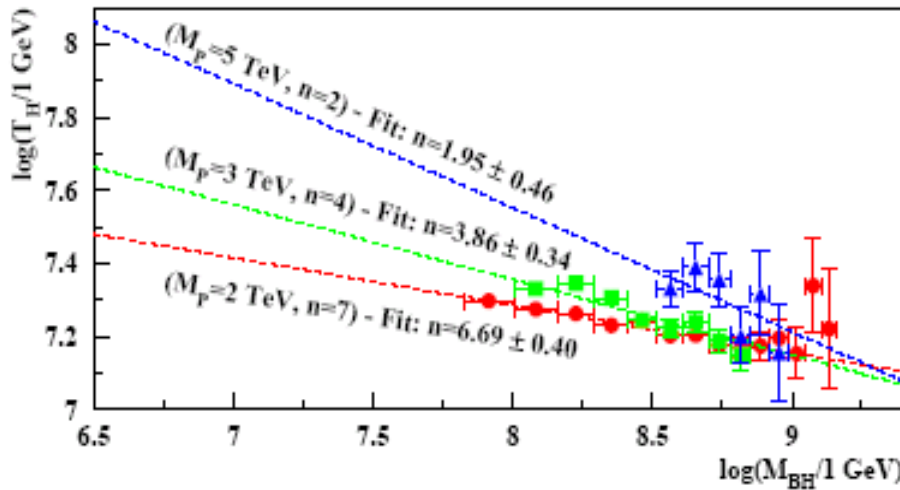


Learning about spacetime

Take the logarithm of both sides of T_H :

$$\log(T_H) = \frac{-1}{n+1} \log(M_{BH}) + const$$

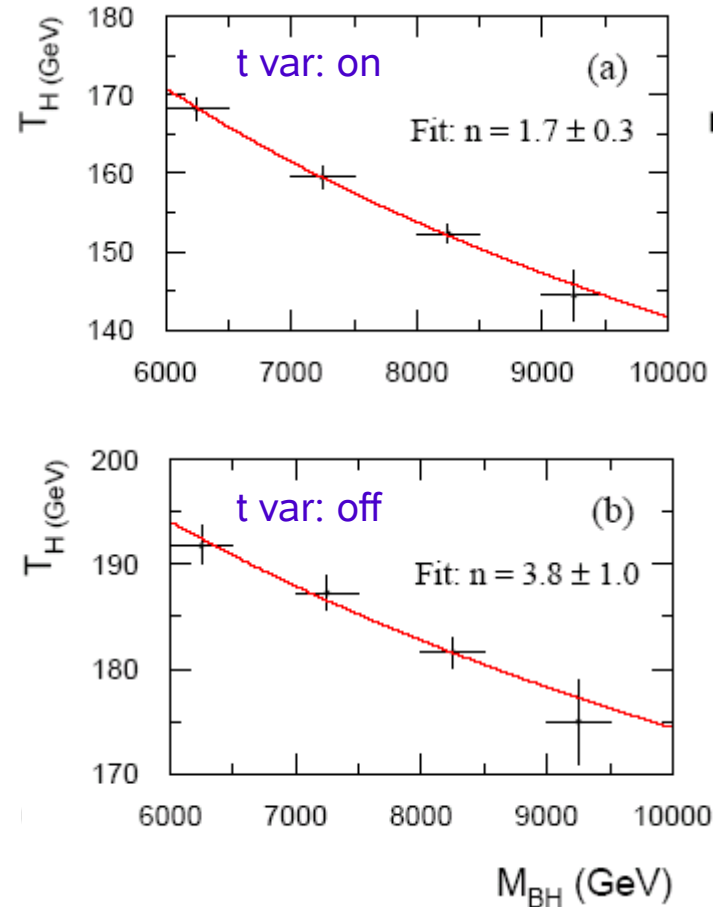
Plot M_{BH} vs. T_H and fit to a line. Find n from the slope.



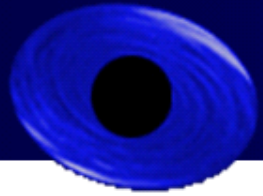
Landsberg, Dimopoulos, PRL 87, 161602 (2001)

Harris et al, ATLAS, Charybdis,
 $M_{4+n} = 1, n = 2$

But, including time variation leads to problems!



Harris et al tried a method in which they determined the order of particles emitted, reconstructed M_{BH} and T_H after each emission.

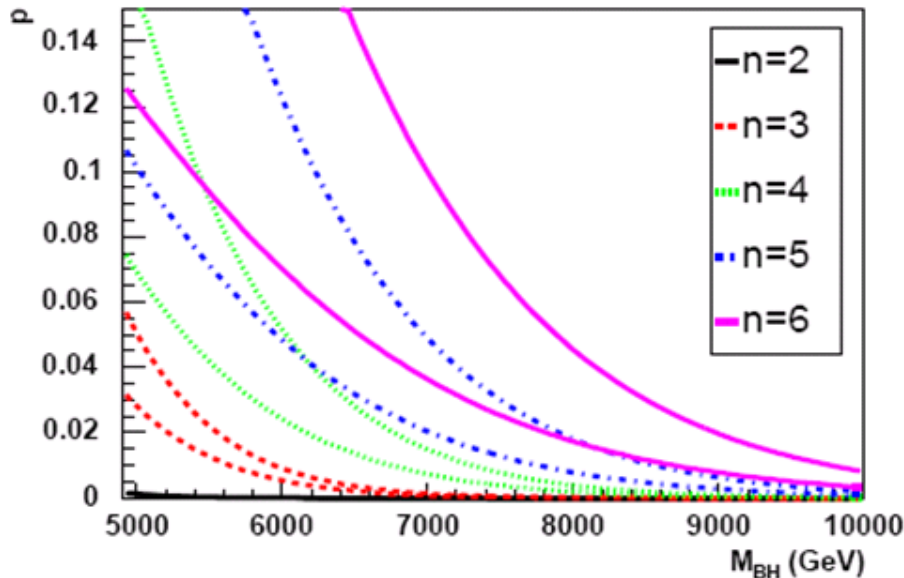


Learning about spacetime

Harris et al, ATLAS, Charybdis, $M_{4+n} = 1$, $n = 2$

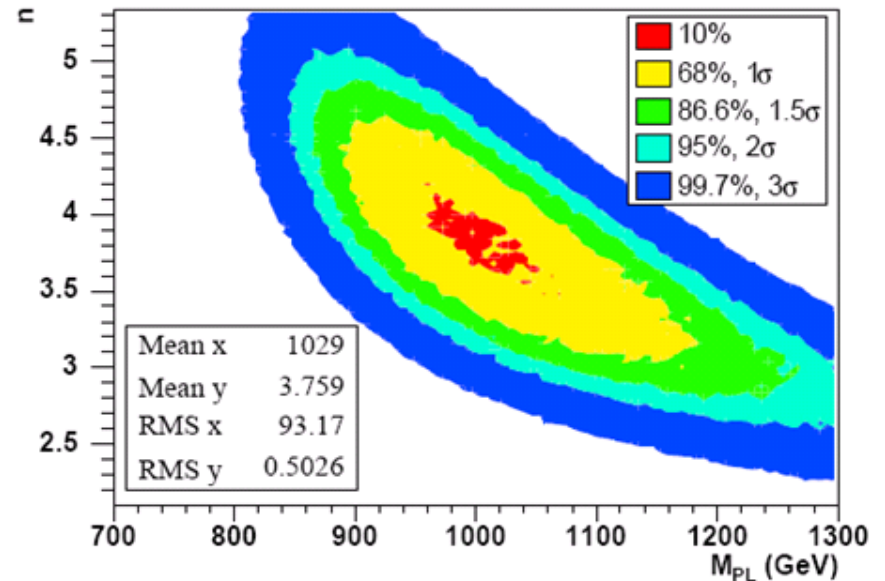
To find n : Plot the fraction p of events which emit particles with energy $E_{\text{cut}} = M_{\text{BH}}/2 - d$. This fraction is strongly dependent on n .

To find $M_{(4+n)}$: Make use of cross sections since they strongly depend on $M_{(4+n)}$ but their n -dependence is negligible.



Theoretical upper and lower limits on p :

$$p_{\text{lower}} = k \int_{E_{\text{cut}}}^{M_{\text{BH}}/2} P dE \quad p_{\text{upper}} = k \int_{E_{\text{cut}}}^{\infty} P dE$$



The final fit!

Combines “ n ” determination from p parameter and “ $M_{(4+n)}$ ” determination from cross sections

CONCLUSIONS

Both direct ADD graviton production and ADD graviton contribution to DY processes can be observed with sufficient significance for moderate values of M_{4+n}

Copious BH production might be possible at the LHC, and if this is realized, BHs will be identified by their significant signature. By looking at BH decay products, BH mass, Hawking temperature and spacetime properties might be reconstructed.