

Aspects of String Compactifications and Moduli Stabilisation

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Introduction

The Standard Model of Particle Physics \rightarrow gauge group $SU(3) \times SU(2) \times U(1)$
– works well at energies of order 100 GeV (includes Higgs boson)

It is just an effective theory – need “Beyond SM” at higher energies.

Possibilities: Supersymmetry \rightarrow **Minimal Supersymmetric Standard Model**
(MSSM)

GUT/susy GUT

Supersymmetry = fermionic symmetry: boson \leftrightarrow fermion

(Super)Multiplets \rightarrow combinations of fields with different spin

Matter \rightarrow chiral supermultiplets $\Phi = (\phi, \psi)$

Lagrangian for these fields is given by three functions

- Kähler potential $K(\Phi, \bar{\Phi})$
- superpotential $W(\Phi)$
- gauge coupling function $f_{ab}(\Phi)$

$$\mathcal{L} \sim -g_{i\bar{j}}\partial_\mu\phi^i\partial^\mu\bar{\phi}^{\bar{j}} - \frac{1}{4}\text{Im}f_{ab}F_{\mu\nu}^aF^{b\mu\nu} + \frac{i}{4}\text{Re}f_{ab}F_{\mu\nu}^a\tilde{F}^{b\mu\nu} - V ,$$

$$g_{i\bar{j}} = \partial_{\Phi^i}\partial_{\bar{\Phi}^{\bar{j}}}K(\Phi, \bar{\Phi}) ,$$

$$V = e^K (D_iW\overline{D_jW}g^{i\bar{j}} - 3|W|^2) + \frac{1}{2}\text{Im}f_{ab}^{-1}D^aD^b$$

$$D_iW = \partial_{\Phi^i}W + (\partial_{\Phi^i}K)W .$$

Supersymmetric solutions: $D_iW = 0$.

String theory

Replaces point particles by 1d extended objects – strings.

Supposed to be valid at energies of order $M_{Pl} = 10^{19} GeV$.

Contains gravity – good candidate for a theory of quantum gravity.

In the low energy limit \rightarrow supergravity in 10 space-time dimensions

Compactifications on 6-dimensional manifolds \rightarrow supergravity in 4d: K , W and f can be computed in string theory

There exist 5 consistent superstring theories: type IIA/B, type I, heterotic $SO(32)/E_8 \times E_8$.

4d requirements

N=1 supersymmetry

Standard Model/GUT

- gauge grup $G \supset SU(3) \times SU(2) \times U(1)$
- chiral matter

Type II \rightarrow need additional constructions: intersecting branes, singularities etc.

$SO(32)$ – does not have the right representations for matter fields in 4d

E_8 \rightarrow works pretty well

Compactification

Simple example 5d \rightarrow 4d on S^1

5d scalar field $\phi \rightarrow$ 4d scalar field

5d vector field $\rightarrow \begin{cases} A_\mu & - & 4\text{d vector field} \\ A_5 & - & 4\text{d scalar} \end{cases}$

5d metric $\rightarrow \begin{cases} g_{\mu\nu} & - & 4\text{d metric} \\ g_{\mu 5} & - & 4\text{d vector field } (U(1) \text{ gauge field}) \\ g_{55} & - & 4\text{d scalar} \end{cases}$

Structure of the 4d theory depends on the internal (6d) manifold.

$N = 1$ susy in 4d requires a globally defined (covariantly constant) spinor in 6d

Need $SU(3)$ holonomy \Rightarrow Calabi–Yau manifolds (complex, Kähler, Ricci-flat manifolds)

$$\mathbf{4} = \mathbf{1} \oplus \mathbf{3}$$

Solve Einstein equations for 4d Minkowski space

Heterotic E_8 string \rightarrow realistic 4d models e.g. gauge group E_6 and chiral matter

Problems: moduli

Deformations of the internal geometry which preserve the background

δg_{mn} such that $\delta R_{mn}(\delta g) = 0$.

Kähler class deformations $\delta g_{a\bar{b}}$

Complex structure deformations δg_{ab} .

All couplings and masses in 4d depend on the vev's of the moduli.

Large vacuum degeneracy

Need to fix moduli!

Turning on fluxes

10d tensor fields of various degree – usually set to zero in KK compactifications

Flux = vev for the field strengths of these fields

$$\int_{\gamma_p} F_p = n$$

Generate superpotentials for the moduli

$$W = \int_{CY} \sum_{p\text{-even}} F_p \wedge e^J + F_3 \wedge \Omega$$

J - two form, Ω - three form specific to CY manifolds – depend on moduli.

Moduli stabilisation

Not all moduli enter the flux superpotential

Need generalisations to manifolds with $SU(3)$ structure – geometric fluxes

Very few cases where all moduli appear in W

Almost always need to rely on non-perturbative effects (gaugino condensation, world-sheet or brane instantons)

$$W = W_{\text{flux}} + e^{-S}$$

Fluxes are quantised $\rightarrow W_{\text{flux}} \sim \mathcal{O}(1)$

Usually $S \gg 1 \rightarrow e^{-S} \ll 1$