



ASTRO-TS



# High-redshift post-reionisation cosmology with 21 cm intensity mapping

[arXiv:1709.07893](https://arxiv.org/abs/1709.07893)

Andrej Obuljen (SISSA)

with:

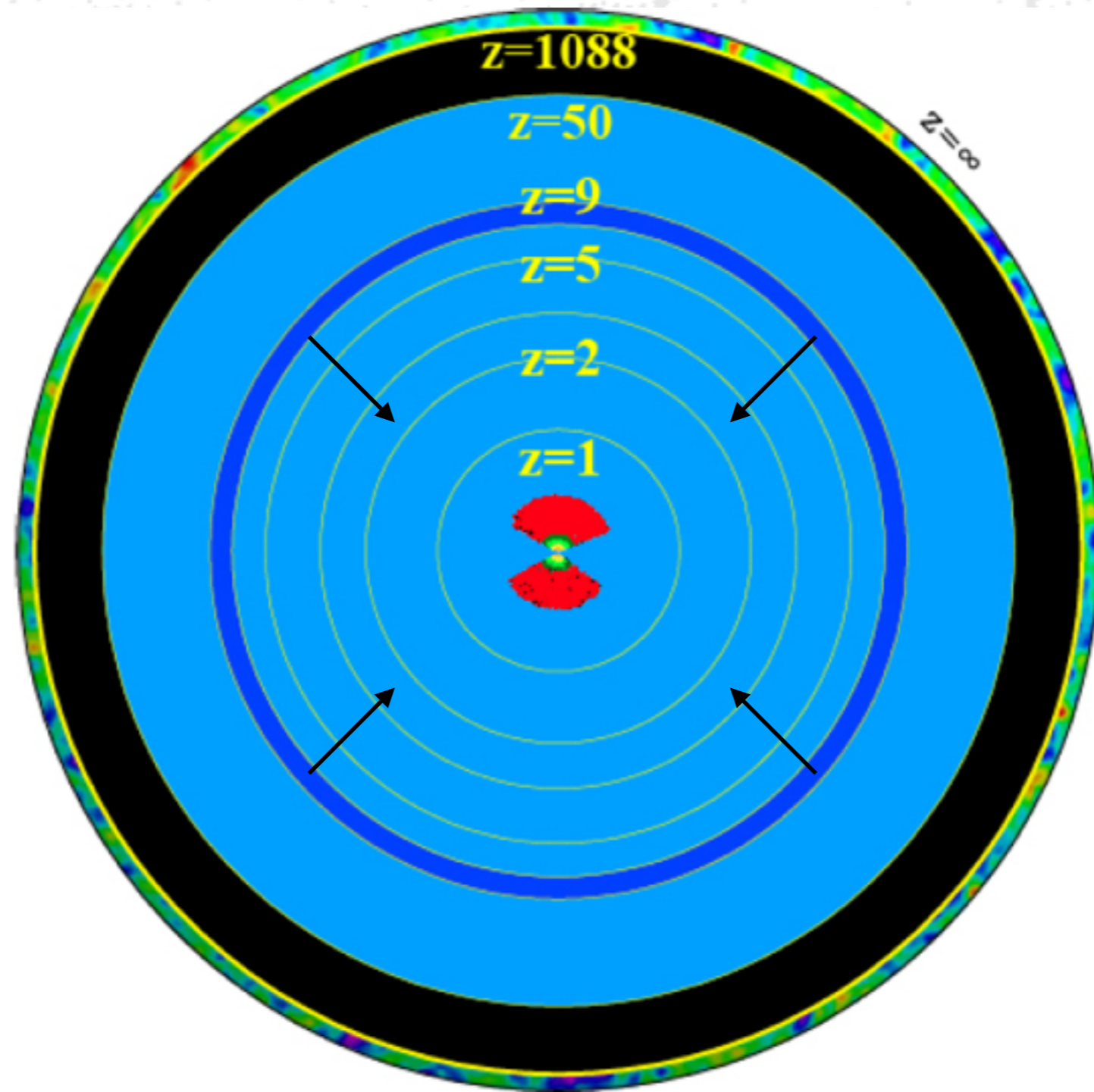
Francisco Villaescusa-Navarro (CCA)

Emanuele Castorina (BCCP)

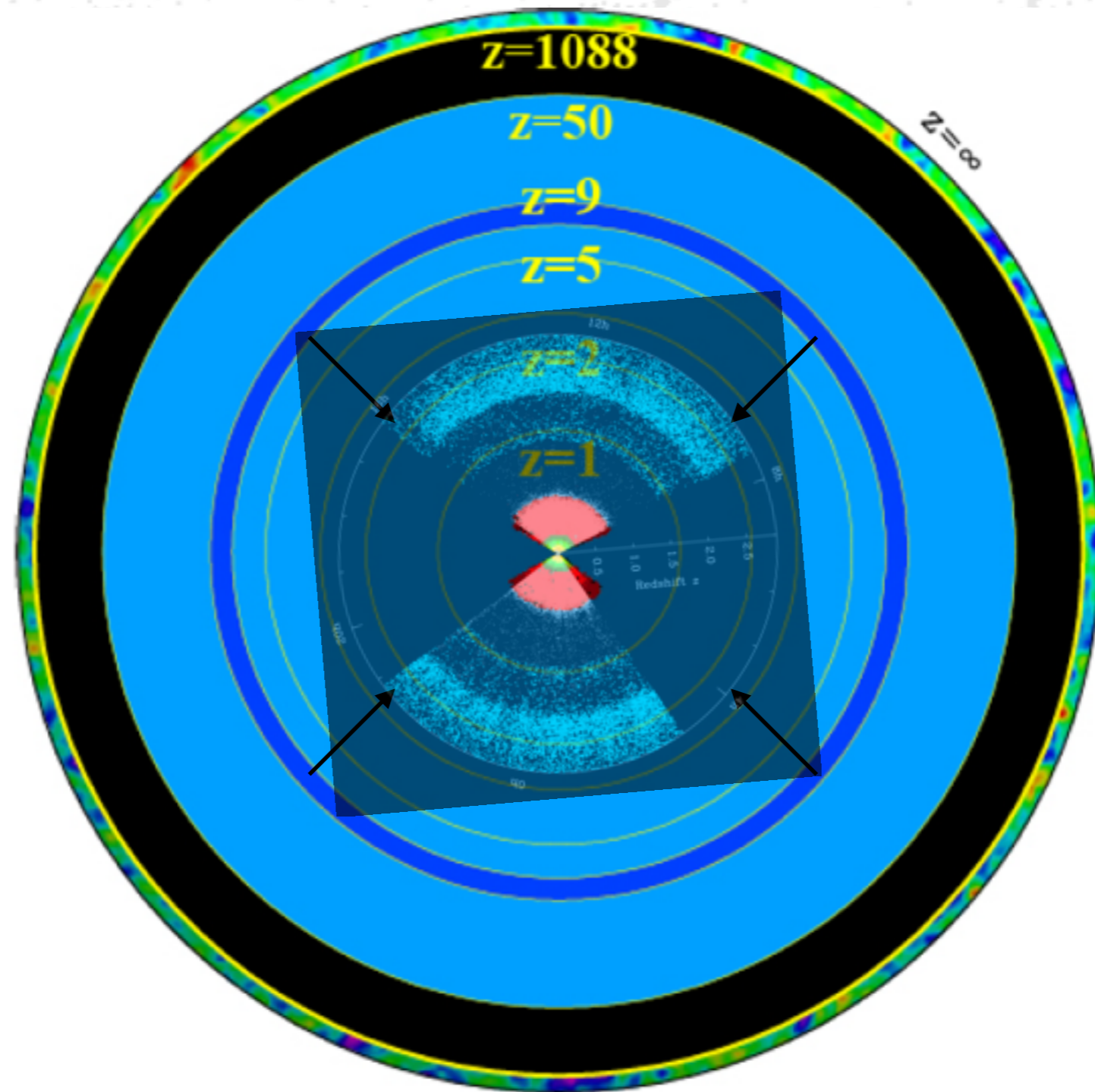
Matteo Viel (SISSA)

26. September 2017.

# 21 cm cosmology

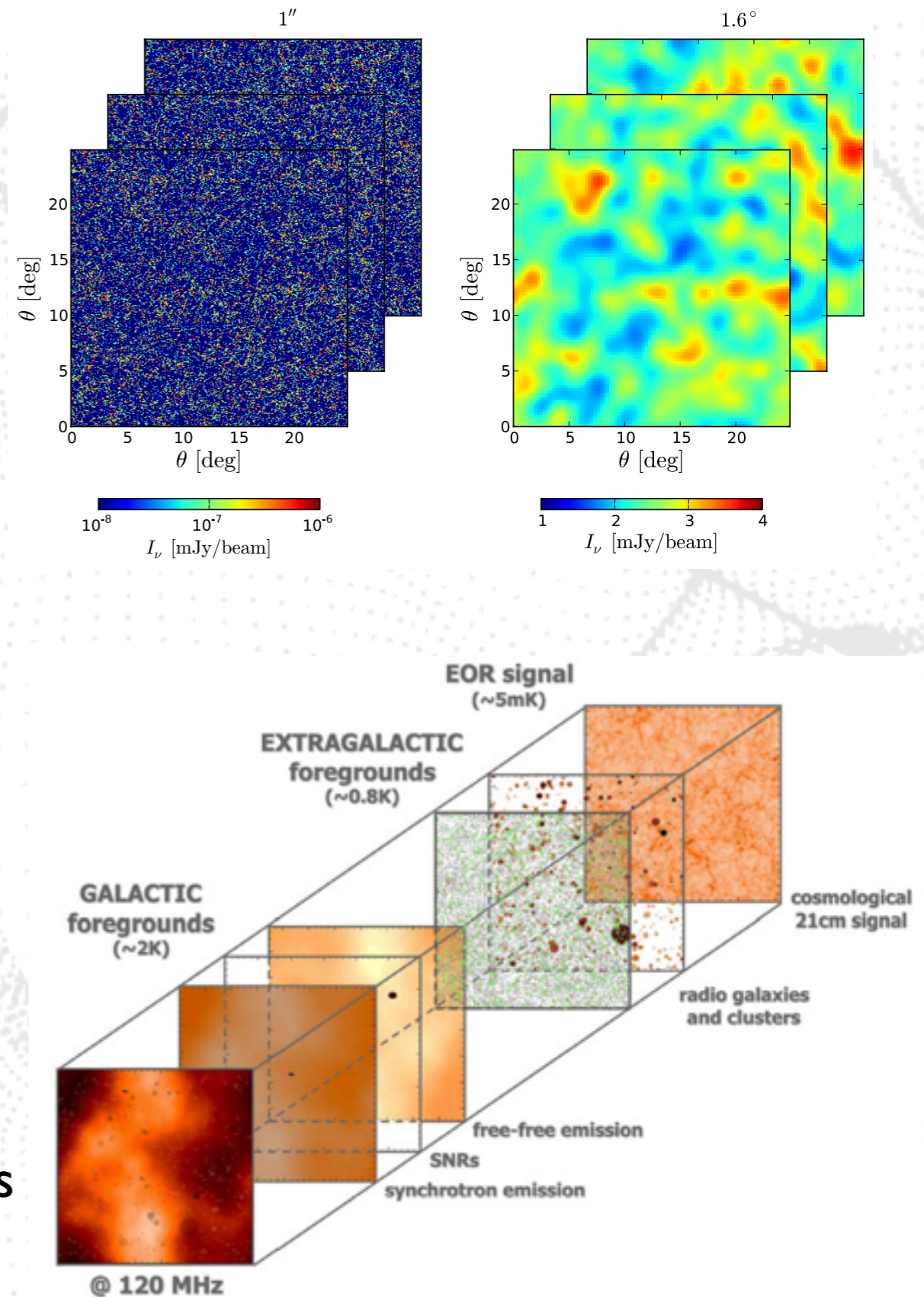


# 21 cm cosmology



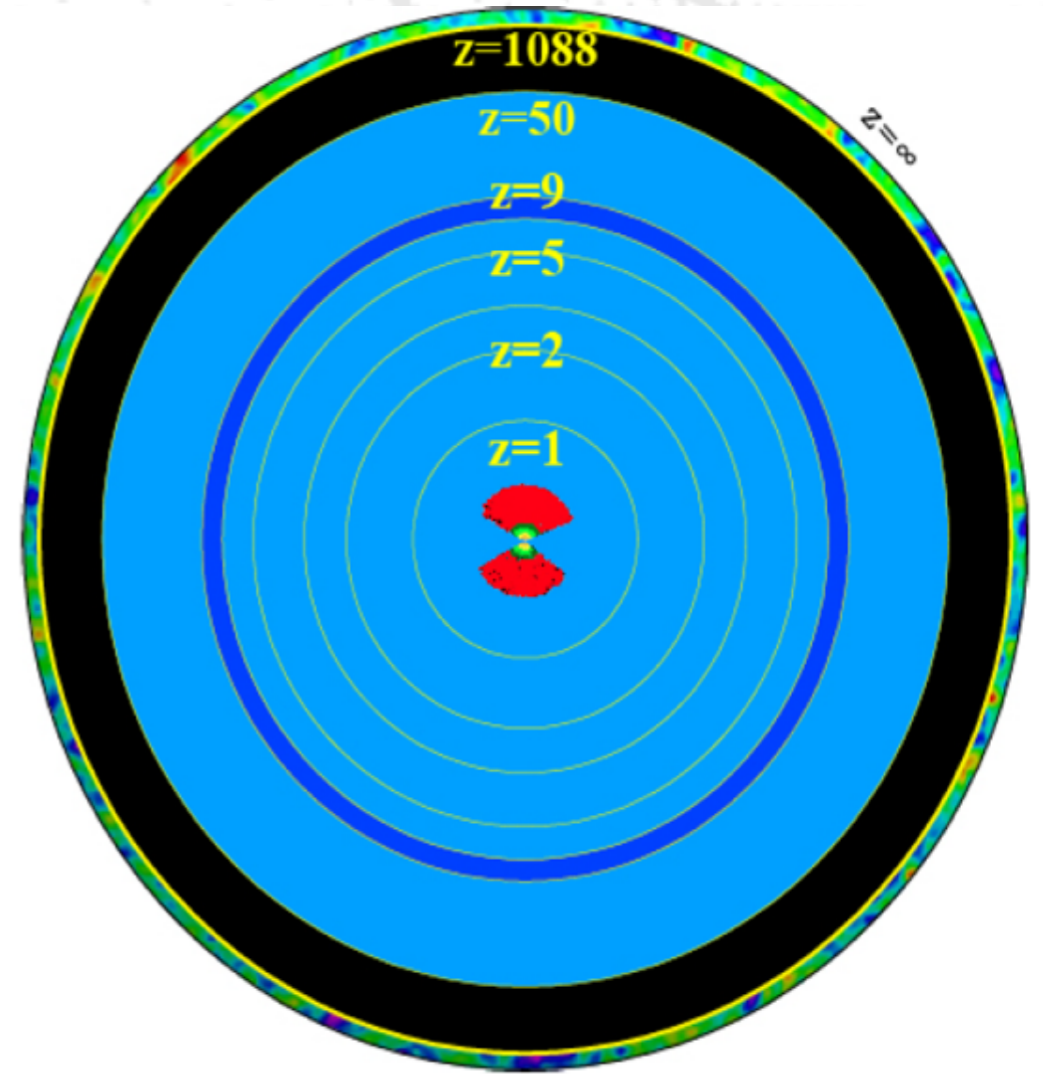
# 21-cm Cosmology

- Promising probe of LSS
- Intensity mapping (IM) technique
- Foregrounds problem!
- Reionization probes: PAPER, HERA, LOFAR, MWA etc.
- We will focus on low- $z$  where HI resides in galaxies
- Many probes are targeting the BAO peak at  $z < 2.5$ : SKA, CHIME, HIRAX, Tianlai, FAST, BINGO, BAOBAB etc.
- Two approaches: single dish & interferometers



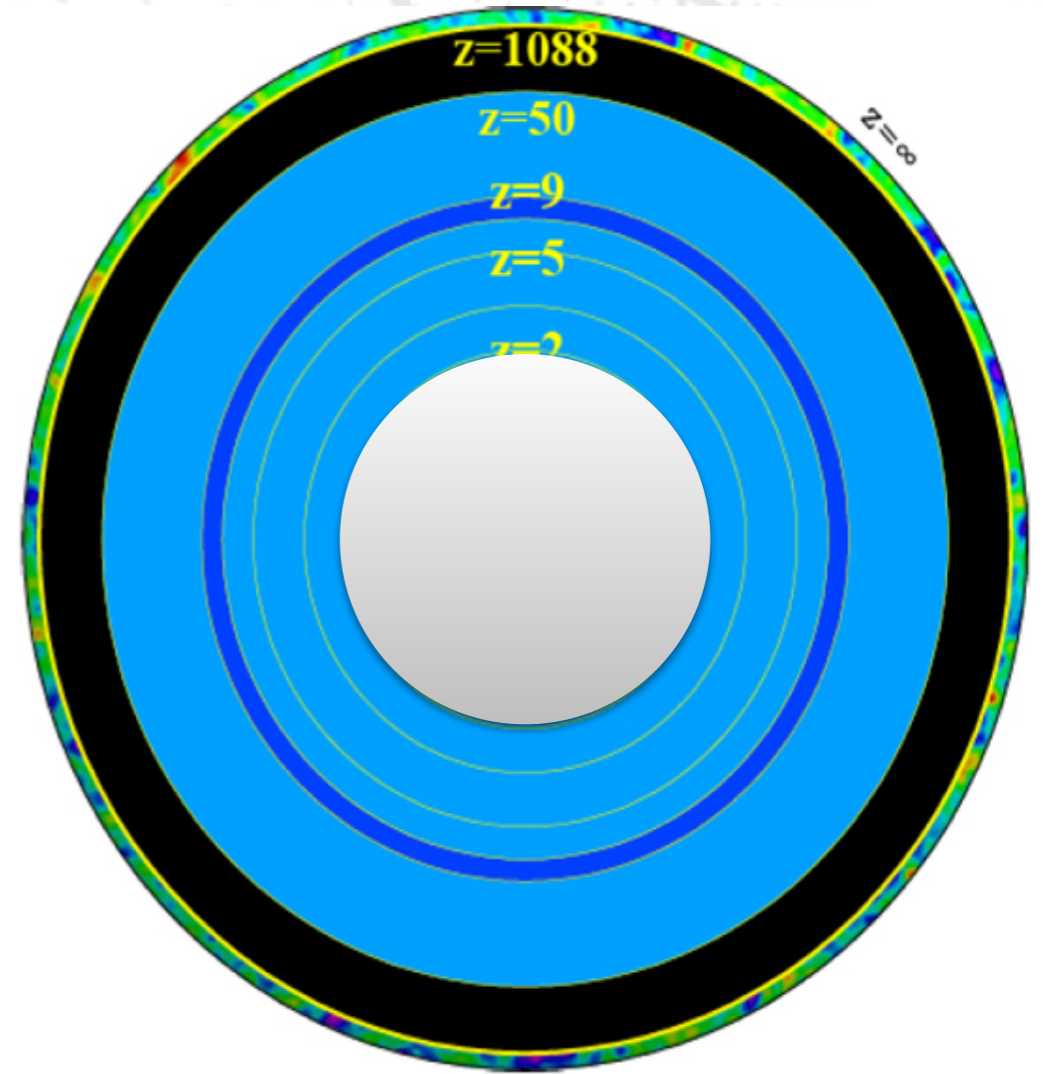
# High-redshift cosmology with 21 cm IM

$2.5 < z < 5.0$



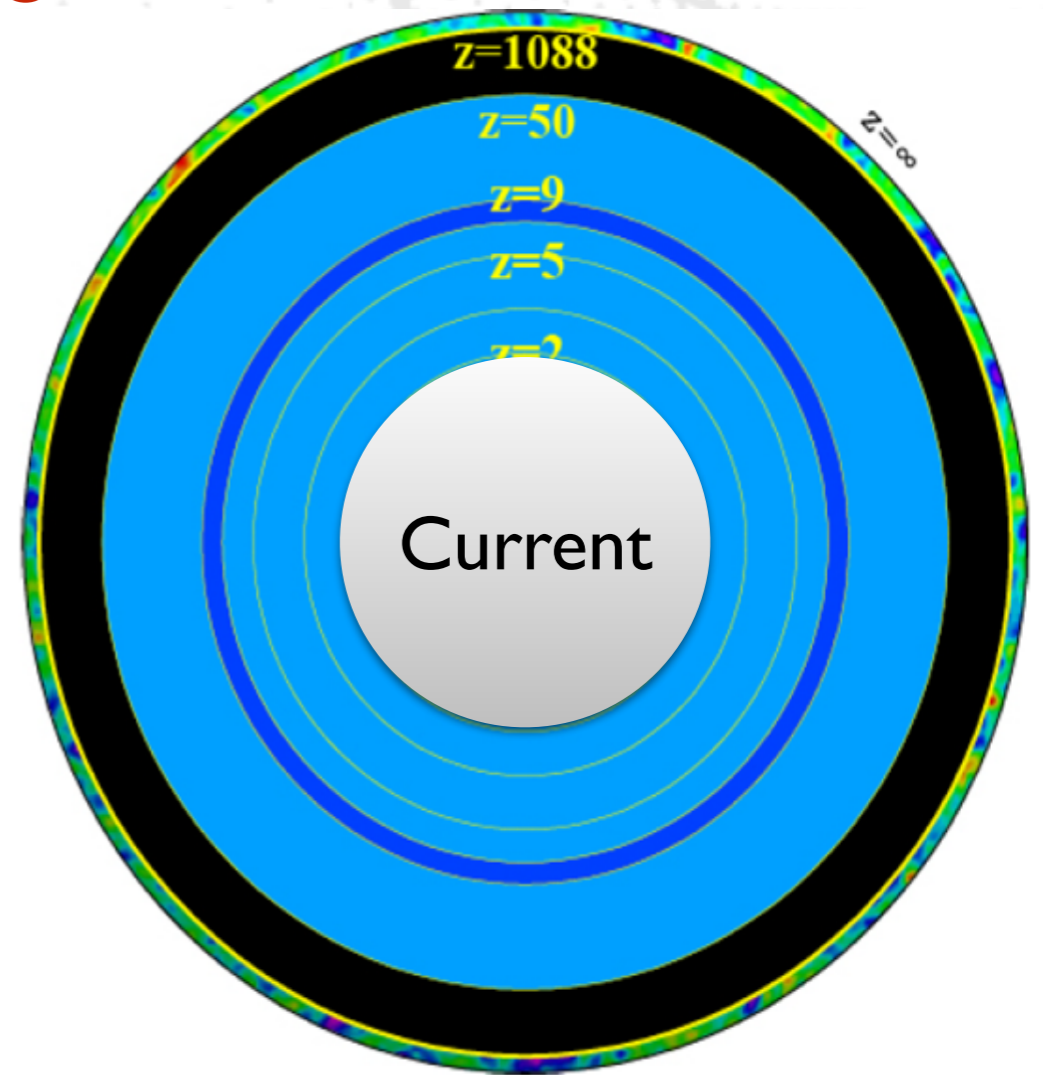
# High-redshift cosmology with 21 cm IM

$2.5 < z < 5.0$



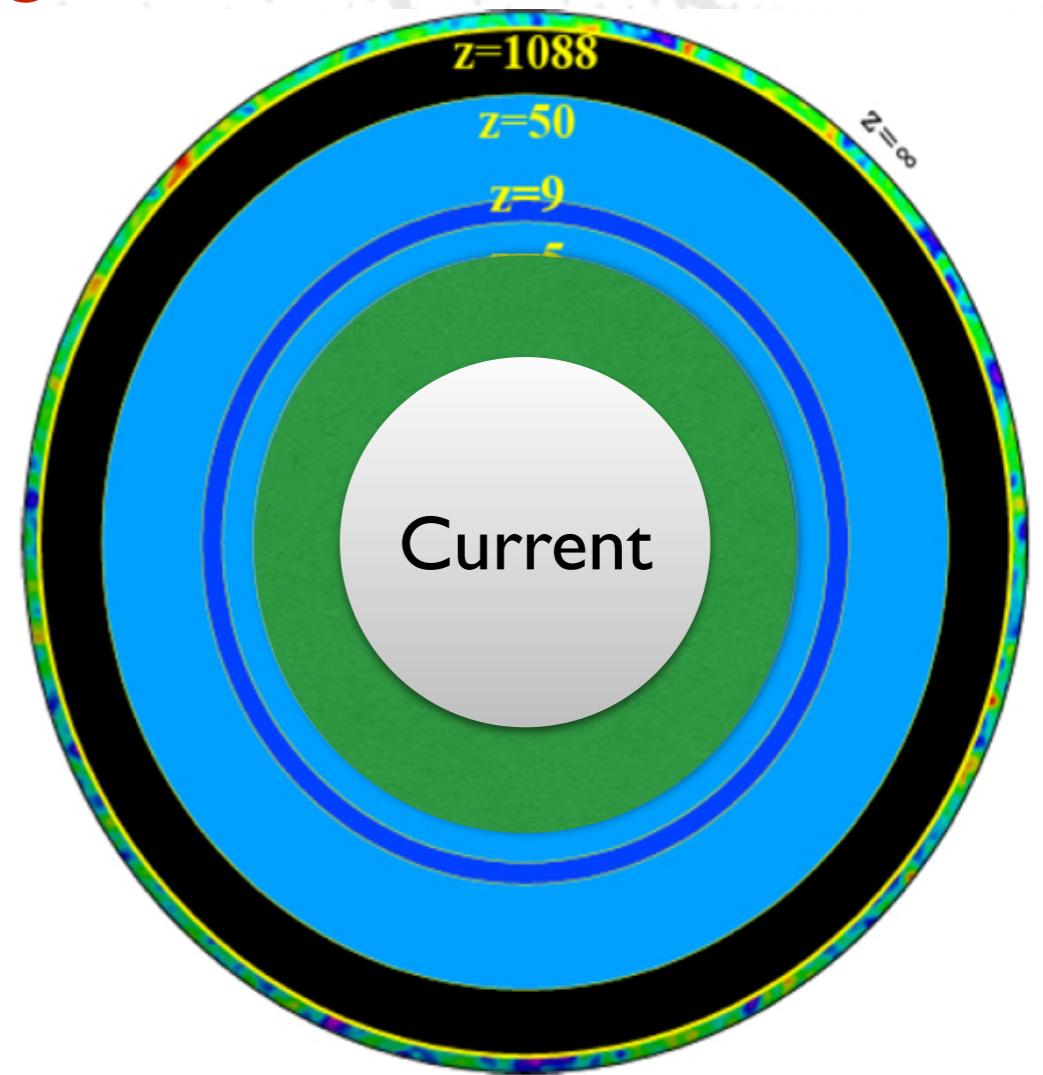
# High-redshift cosmology with 21 cm IM

$2.5 < z < 5.0$



# High-redshift cosmology with 21 cm IM

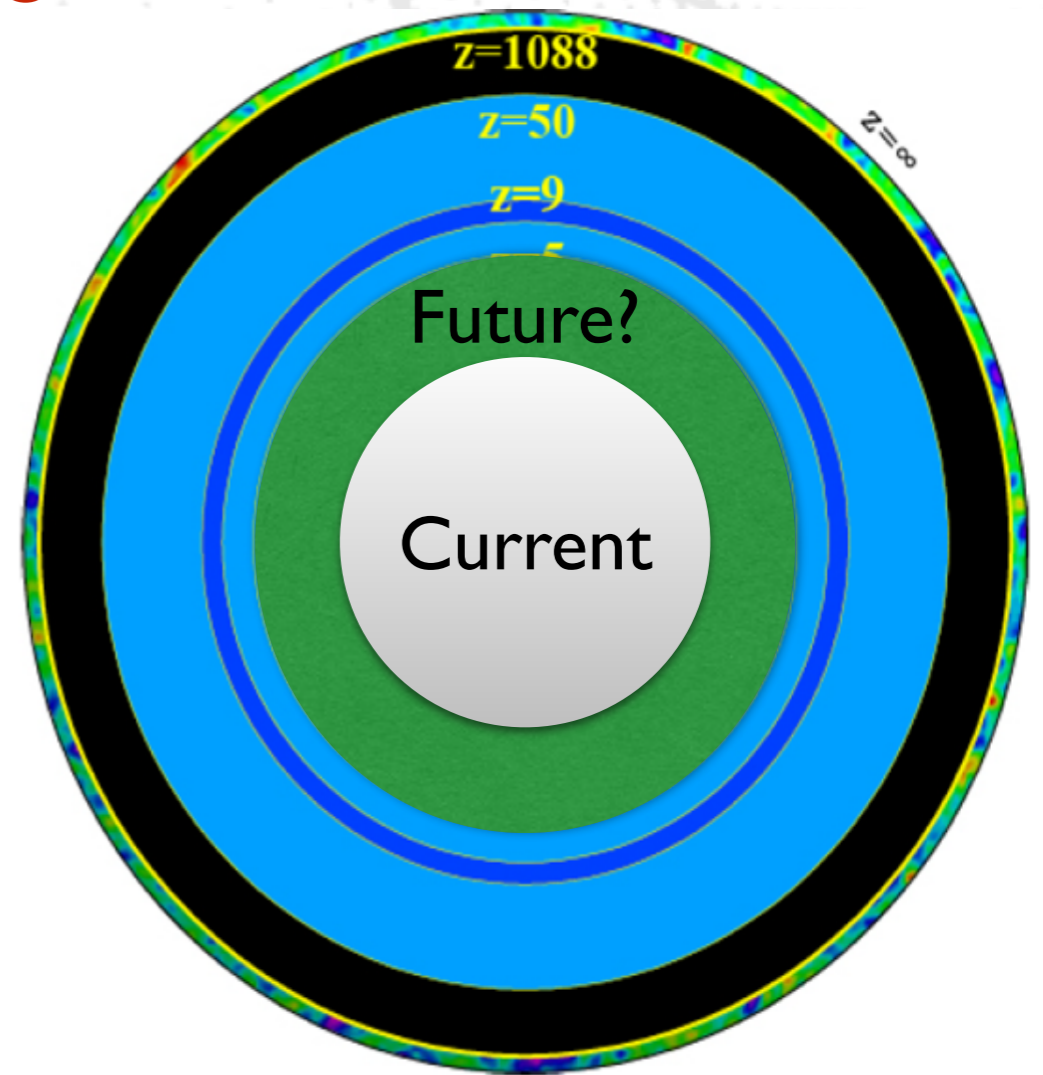
$2.5 < z < 5.0$





# High-redshift cosmology with 21 cm IM

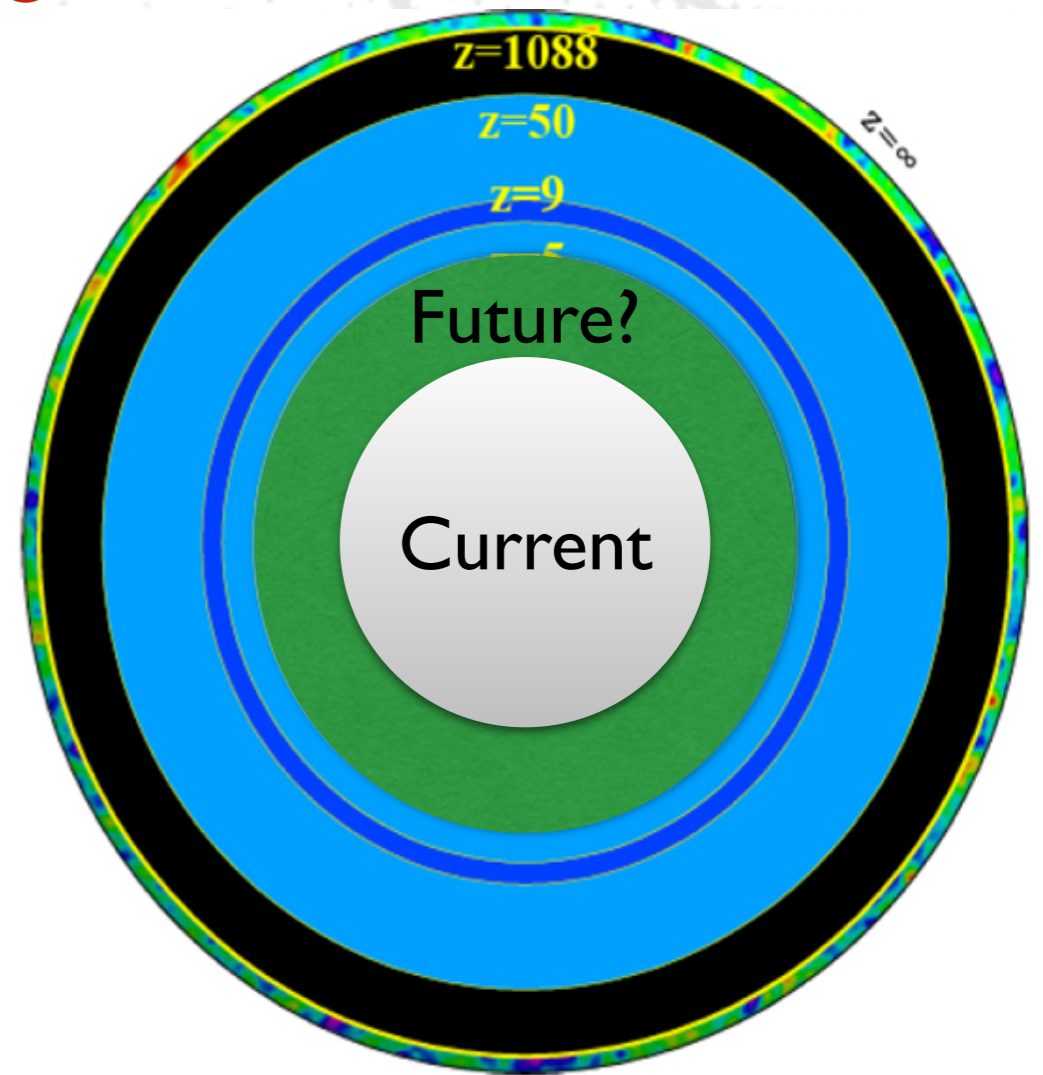
$2.5 < z < 5.0$



# High-redshift cosmology with 21 cm IM

## $2.5 < z < 5.0$

- Advantages
- More comoving volume:
- More linear
- You are running out of galaxies anyway

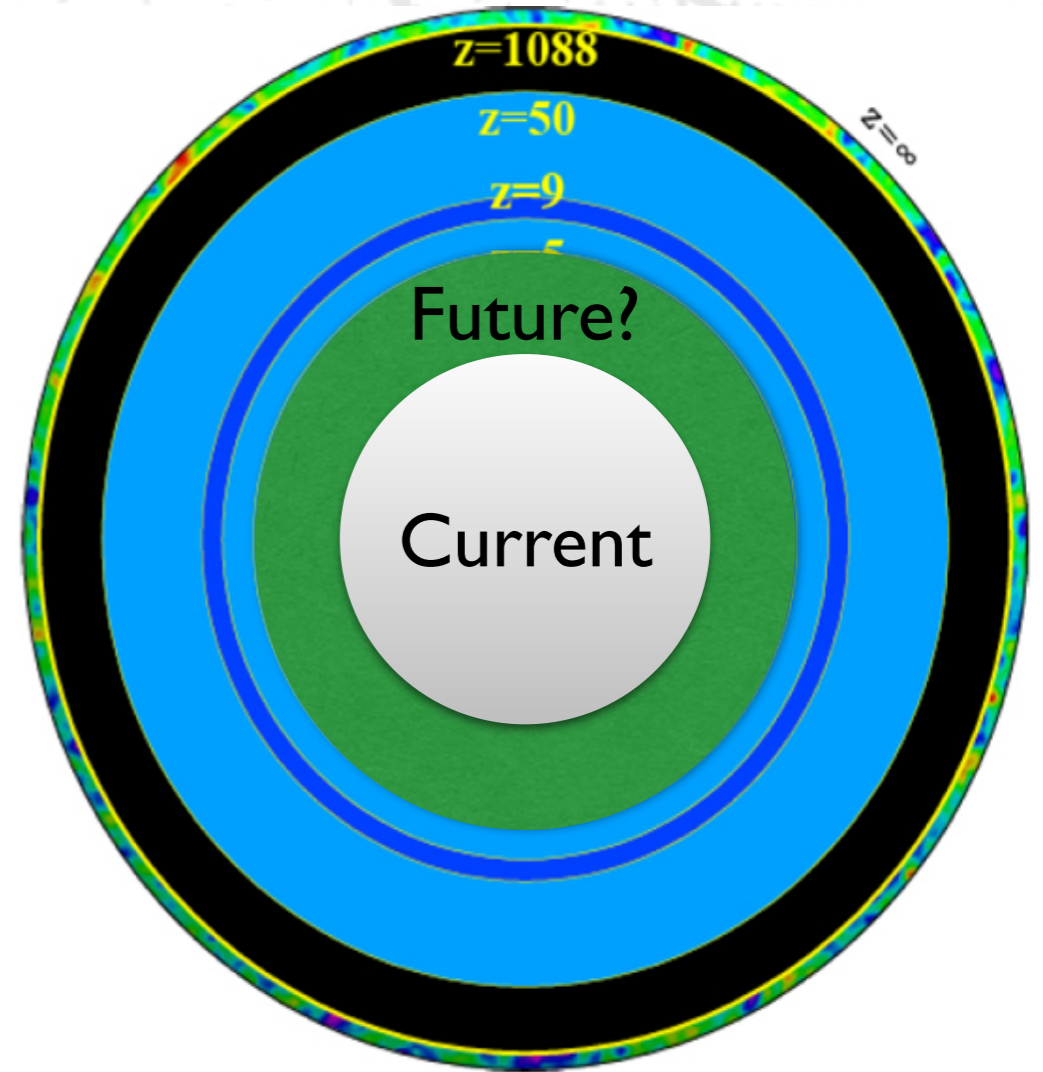


# High-redshift cosmology with 21 cm IM

## $2.5 < z < 5.0$

- **Advantages**

- More comoving volume:  $\frac{V(2.5 < z < 5)}{V(z < 2.5)} \approx 1.4$
- More linear
- You are running out of galaxies anyway



# High-redshift cosmology with 21 cm IM

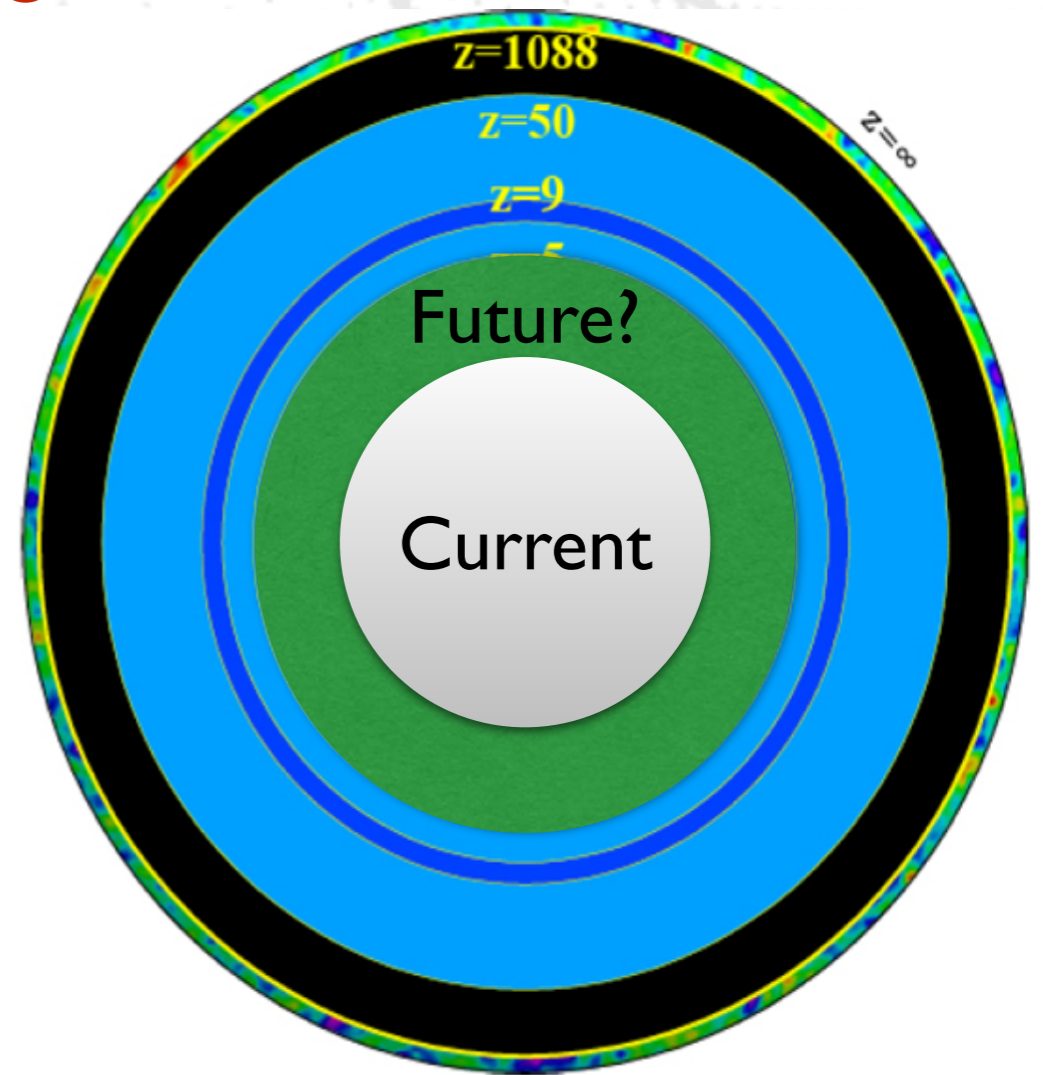
## $2.5 < z < 5.0$

- **Advantages**

- More comoving volume:  $\frac{V(2.5 < z < 5)}{V(z < 2.5)} \approx 1.4$
- More linear
- You are running out of galaxies anyway

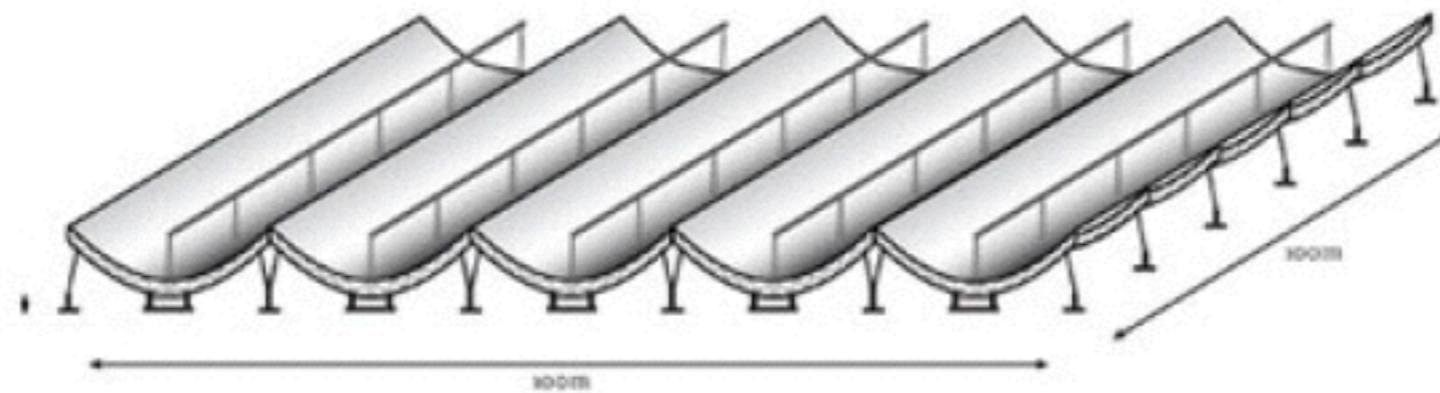
- **Disadvantages**

- Foreground problems
- Sky temperature blows up
- Wedge for interferometers
- Angular resolution for single-dish



# Expand current experiments?

CHIME



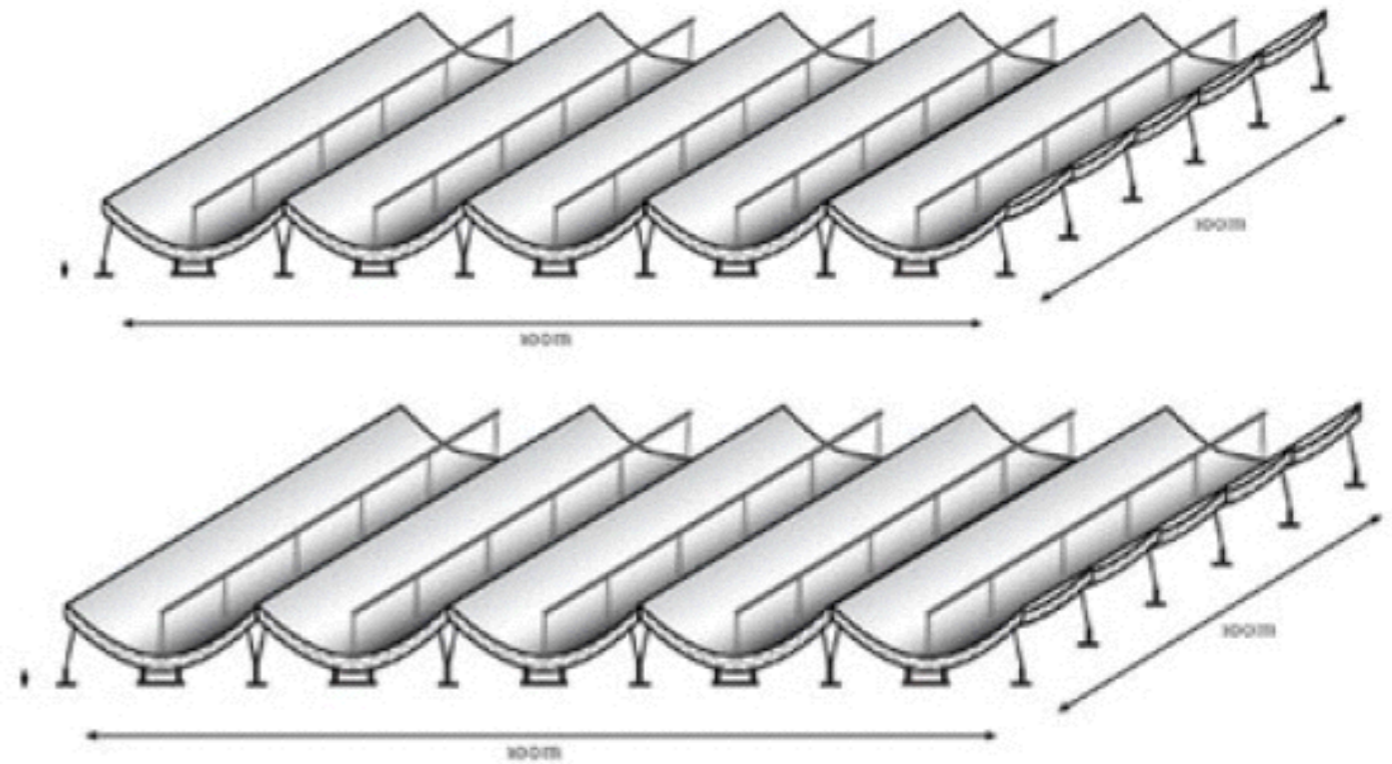
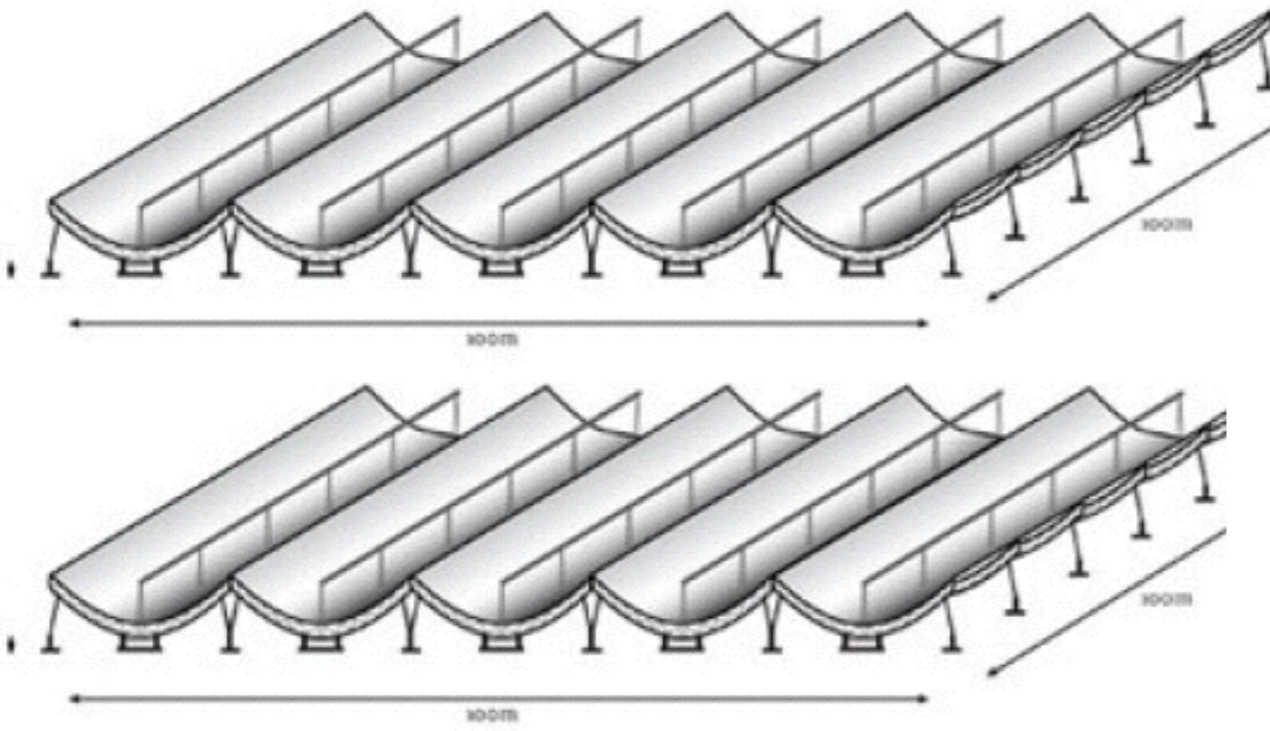
FAST



HIRAX 32x32

# Expand current experiments?

**2x2x**CHIME = Ext-CHIME



**highz**FAST



**HIRAX 64x64** = Ext-HIRAX

# High redshift cosmology with 21 cm IM

## $2.5 < z < 5.0$

We are interested in the following quantities:

- Growth of structures —  $f\sigma_8$
- BAO Alcock-Paczynski parameters —  $H(z) \& D_A(z)$
- Sum of the neutrino masses —  $\Sigma m_\nu$
- Effective number of the neutrino species —  $N_{\text{eff}}$

# Fisher forecasts

Fisher matrix:

$$F_{ij} = \frac{1}{8\pi^2} \int_{-1}^1 d\mu \int k^2 dk \frac{\partial \ln P_{21}(k, \mu)}{\partial p_i} \frac{\partial \ln P_{21}(k, \mu)}{\partial p_j} V_{\text{eff}}(k, \mu)$$

$$V_{\text{eff}}(k, \mu) = V_{\text{sur}} \left( \frac{P_{21}(k, \mu) W(k, \mu)}{P_{21}(k, \mu) W(k, \mu) + P_{\text{N}}^{\text{tot}}(k, \mu)} \right)^2$$

It works under the assumption of Gaussian likelihood and no theoretical uncertainties!

One should be careful about the range of wavenumber considered ( $k_{\text{min}}, k_{\text{max}}$ )



# HI power spectrum model

$$P_{21}(k, z, \mu) = \bar{T}_b^2(z) (b_{\text{HI}}(z) + f(z)\mu^2)^2 P(k, z)$$

$$\bar{T}_b(z) \propto \Omega_{\text{HI}}(z)$$

# HI power spectrum model

$$P_{21}(k, z, \mu) = \bar{T}_b^2(z) (b_{\text{HI}}(z) + f(z)\mu^2)^2 P(k, z)$$

$$\bar{T}_b(z) \propto \Omega_{\text{HI}}(z)$$

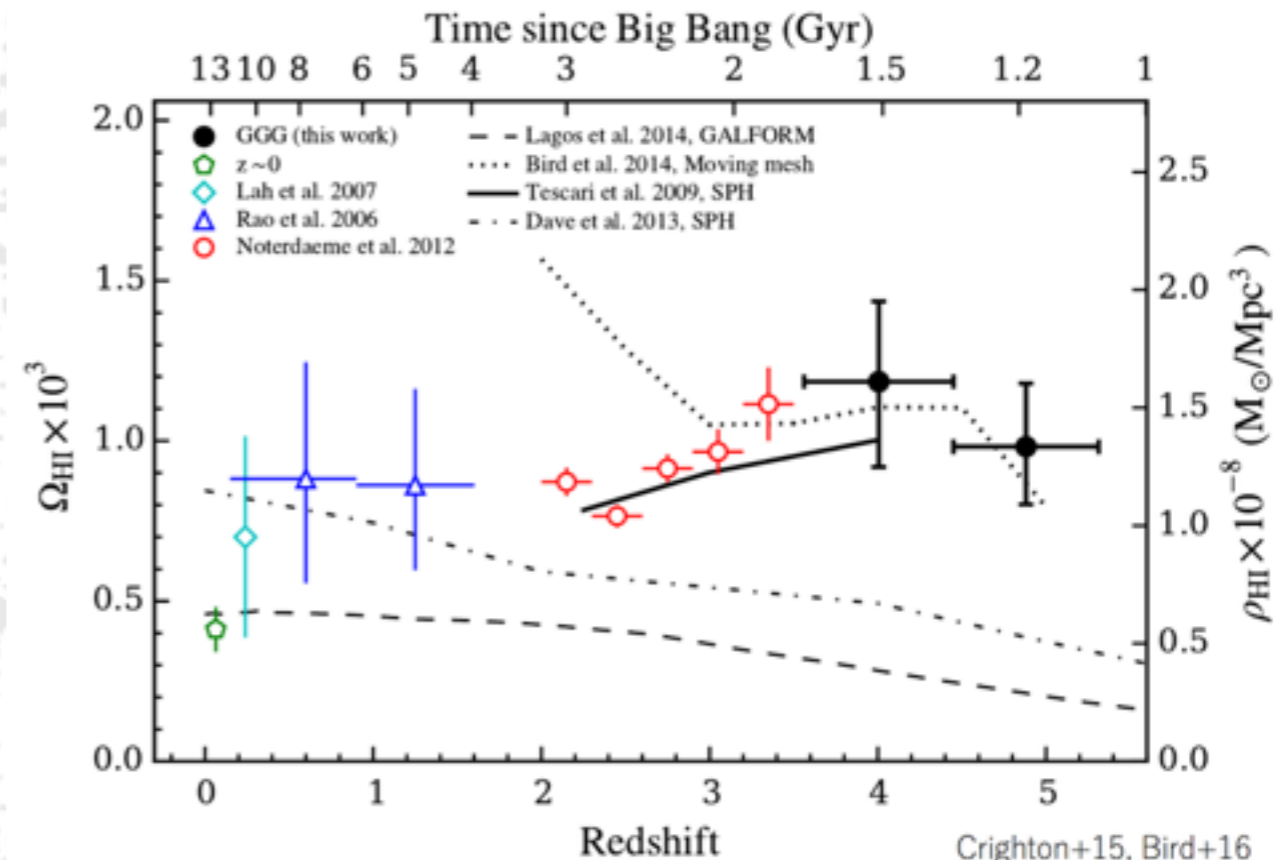
There is a complete degeneracy on the amplitude of the power spectrum!  
In principle, no RSD measurements are possible!

# HI power spectrum model

$$P_{21}(k, z, \mu) = \bar{T}_b^2(z) (b_{\text{HI}}(z) + f(z)\mu^2)^2 P(k, z)$$

$$\bar{T}_b(z) \propto \Omega_{\text{HI}}(z)$$

There is a complete degeneracy on the amplitude of the power spectrum!  
In principle, no RSD measurements are possible!



# HI power spectrum model

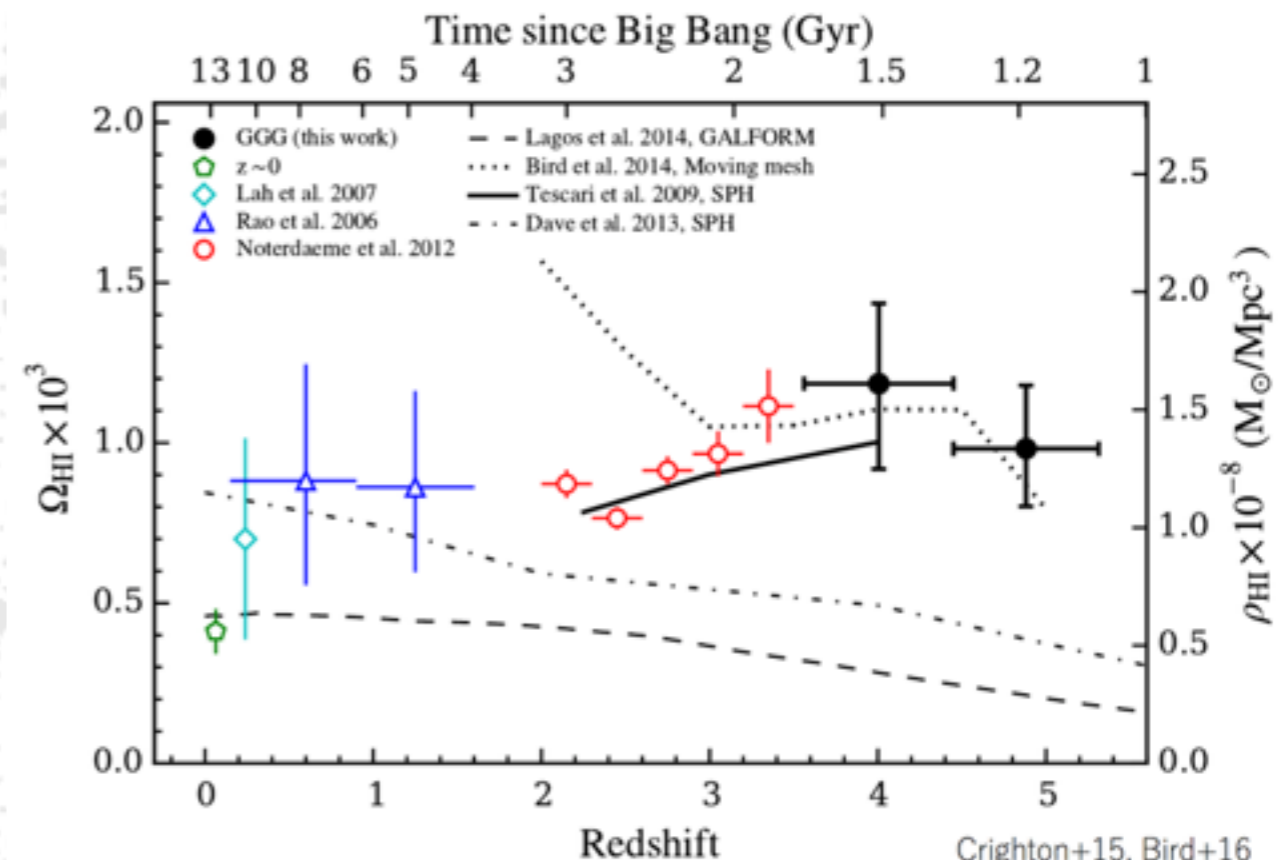
$$P_{21}(k, z, \mu) = \bar{T}_b^2(z) (b_{\text{HI}}(z) + f(z)\mu^2)^2 P(k, z)$$

$$\bar{T}_b(z) \propto \Omega_{\text{HI}}(z)$$

There is a complete degeneracy on the amplitude of the power spectrum!  
In principle, no RSD measurements are possible!

However:

- 1) future measurements of Damped Ly Alpha systems will constrain  $\Omega_{\text{HI}}$
- 2) cross-correlating weighted DLA and Ly forest (or 21cm) can give  $b_{\text{HI}}$



# HI power spectrum model

$$P_{21}(k, z, \mu) = \bar{T}_b^2(z) (b_{\text{HI}}(z) + f(z)\mu^2)^2 P(k, z)$$

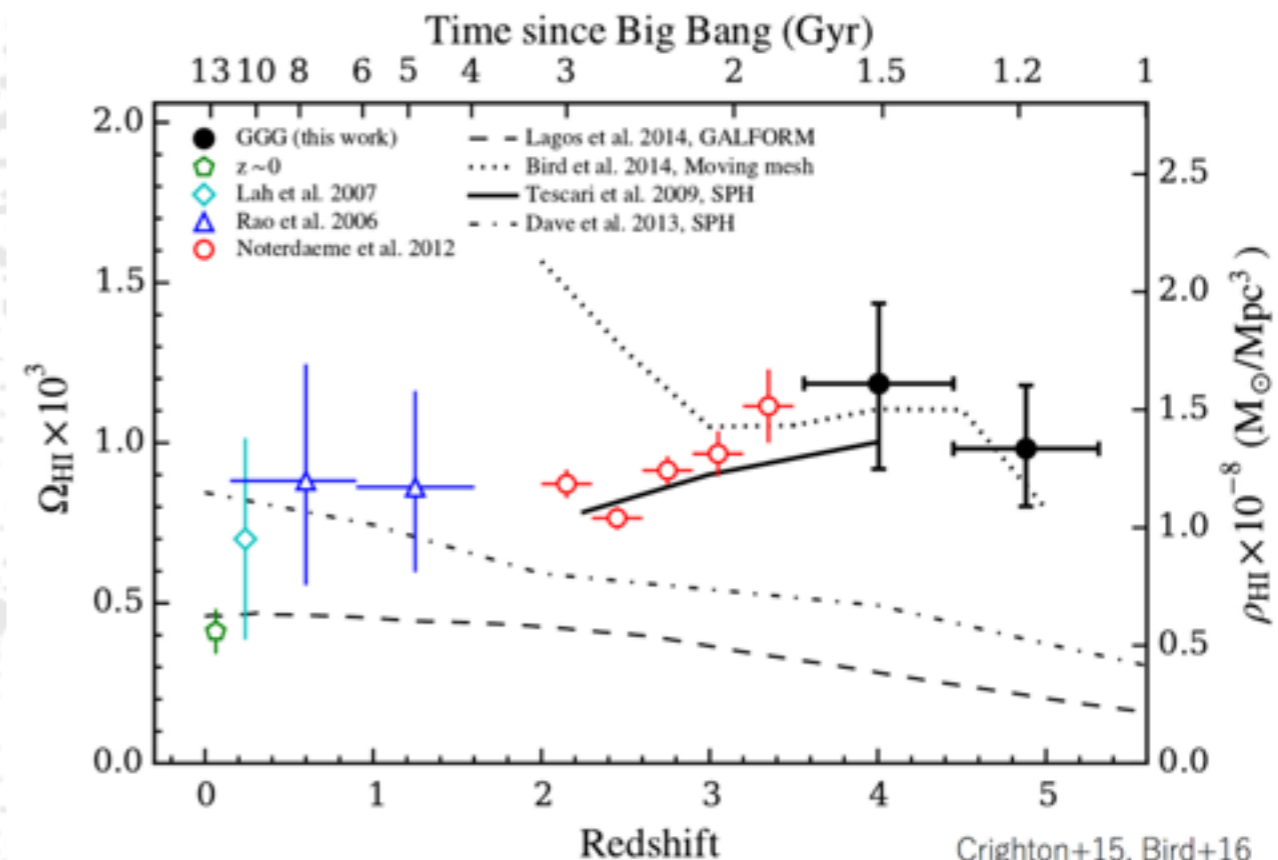
$$\bar{T}_b(z) \propto \Omega_{\text{HI}}(z)$$

There is a complete degeneracy on the amplitude of the power spectrum!  
In principle, no RSD measurements are possible!

However:

- 1) future measurements of Damped Ly Alpha systems will constrain  $\Omega_{\text{HI}}$
- 2) cross-correlating weighted DLA and Ly forest (or 21 cm) can give  $b_{\text{HI}}$

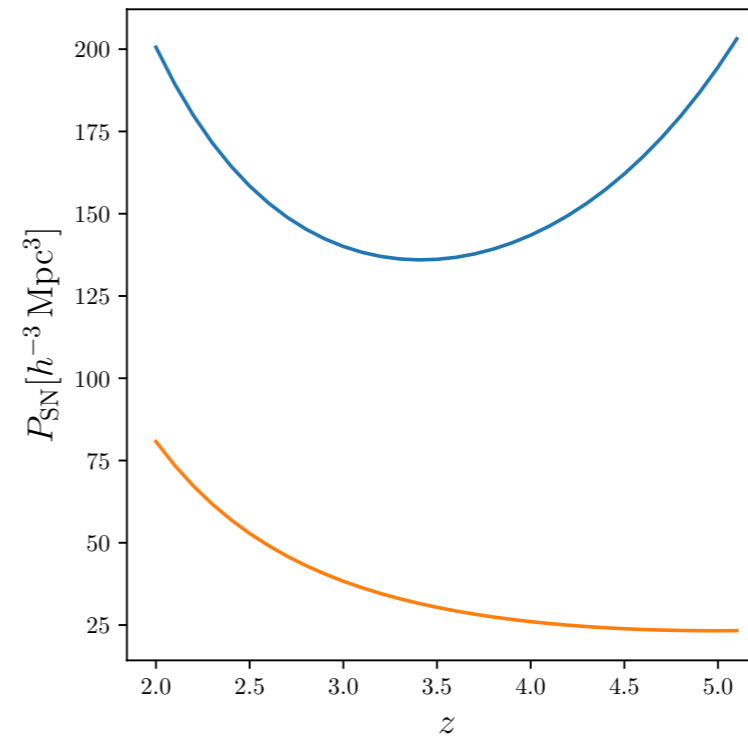
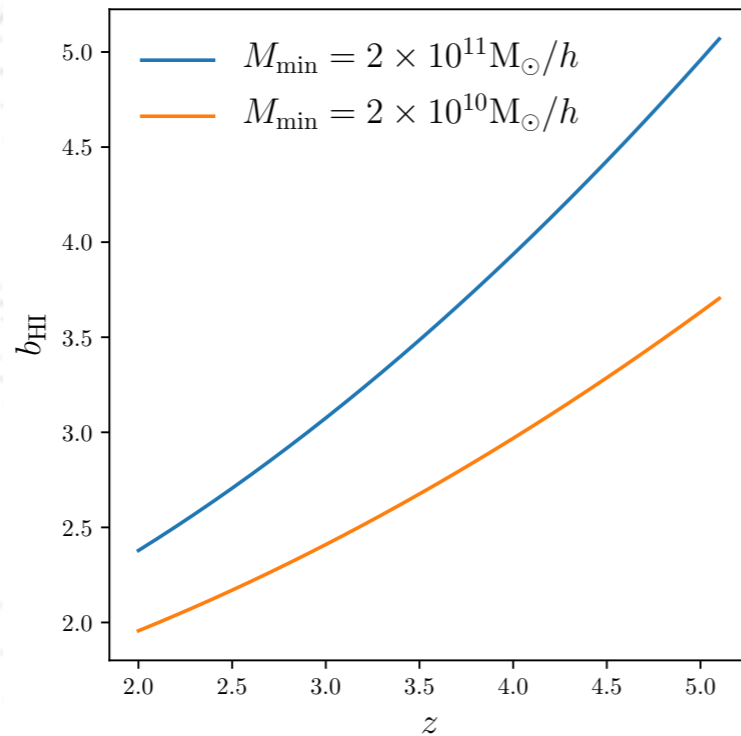
We will assume external  
2%, 5% & 10% priors  
on both  $\Omega_{\text{HI}}$  &  $b_{\text{HI}}$



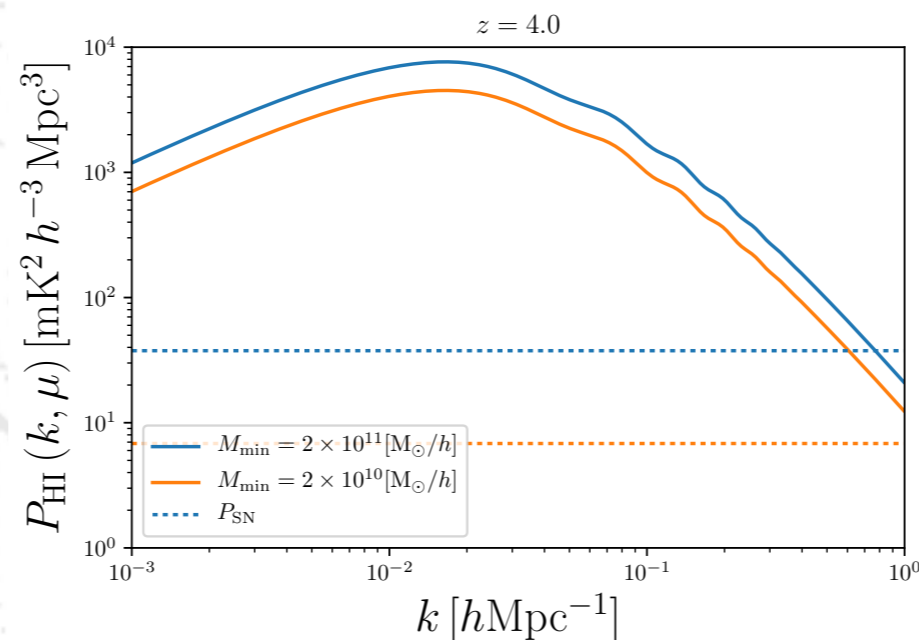
# HI bias model & the shot-noise

We follow Castorina et al. (2016)

$$M_{\text{HI}}(M) \propto e^{-M_{\text{min}}/M} M^{\alpha}$$
$$\alpha = 1$$



HI shot-noise is  
subdominant



# Thermal noise

## Interferometers

$$P_N^{\text{th}}(z) = \frac{T_{\text{sys}}^2(z) X^2(z) Y(z) \lambda^4(z) S_{21}}{A_{\text{eff}}^2 \text{FOV}(z) t_0 n_{\text{pol}} n(\mathbf{u}, z)}$$

$$T_{\text{sky}}(z) = 60 \text{ K} \times (\nu_{21}(z)/300\text{MHz})^{-2.55}$$

## Single-dish

$$P_N^{\text{th}}(k, \mu) = \frac{T_{\text{sys}}^2 V_{\text{pix}} W^{-2}(k_{\perp})}{\Delta\nu t_{\text{obs}} \Omega_{\text{pix}} / S_{21} N_{\text{dish}} N_{\text{beam}}}$$

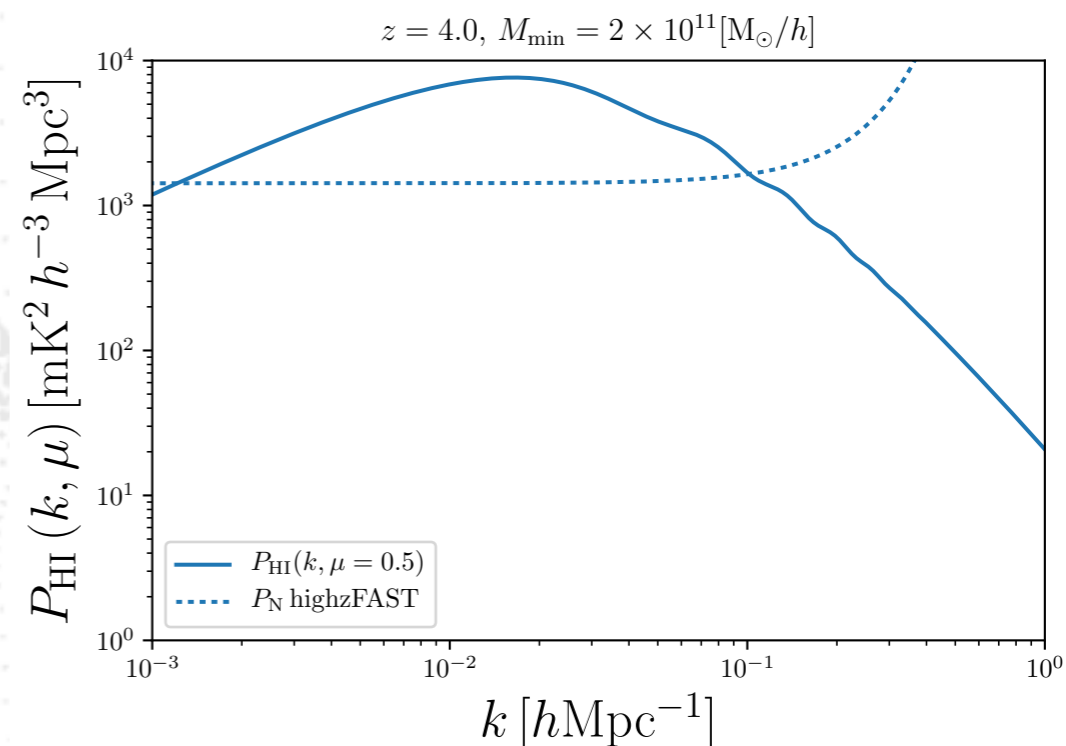
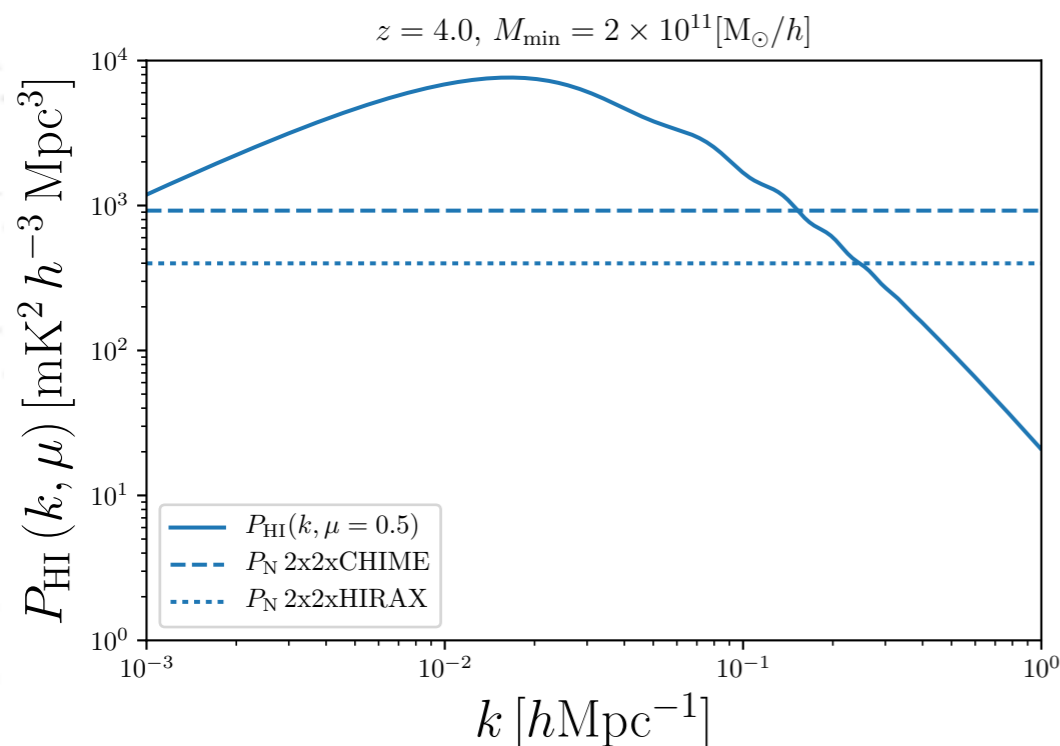
## Observable modes

$$k_{\perp}^{\text{min}}(z) = \frac{2\pi D_{\text{min}}}{D(z)\lambda(z)}$$

$$k_{\perp}^{\text{max}}(z) = \frac{2\pi D_{\text{max}}}{D(z)\lambda(z)}$$

$$k_{\perp}^{\text{min}}(z) = \frac{2\pi}{\sqrt{D(z)^2 S_{21}}}$$

$$k_{\perp}^{\text{max}}(z) = \frac{2\pi D_{\text{dish}}}{D(z)\lambda(z)}$$



# Thermal noise

## Interferometers

$$P_N^{\text{th}}(z) = \frac{T_{\text{sys}}^2(z) X^2(z) Y(z) \lambda^4(z) S_{21}}{A_{\text{eff}}^2 \text{FOV}(z) t_0 n_{\text{pol}} n(\mathbf{u}, z)}$$

$$T_{\text{sky}}(z) = 60 \text{ K} \times (\nu_{21}(z)/300\text{MHz})^{-2.55}$$

## Single-dish

$$P_N^{\text{th}}(k, \mu) = \frac{T_{\text{sys}}^2 V_{\text{pix}} W^{-2}(k_{\perp})}{\Delta\nu t_{\text{obs}} \Omega_{\text{pix}} / S_{21} N_{\text{dish}} N_{\text{beam}}}$$

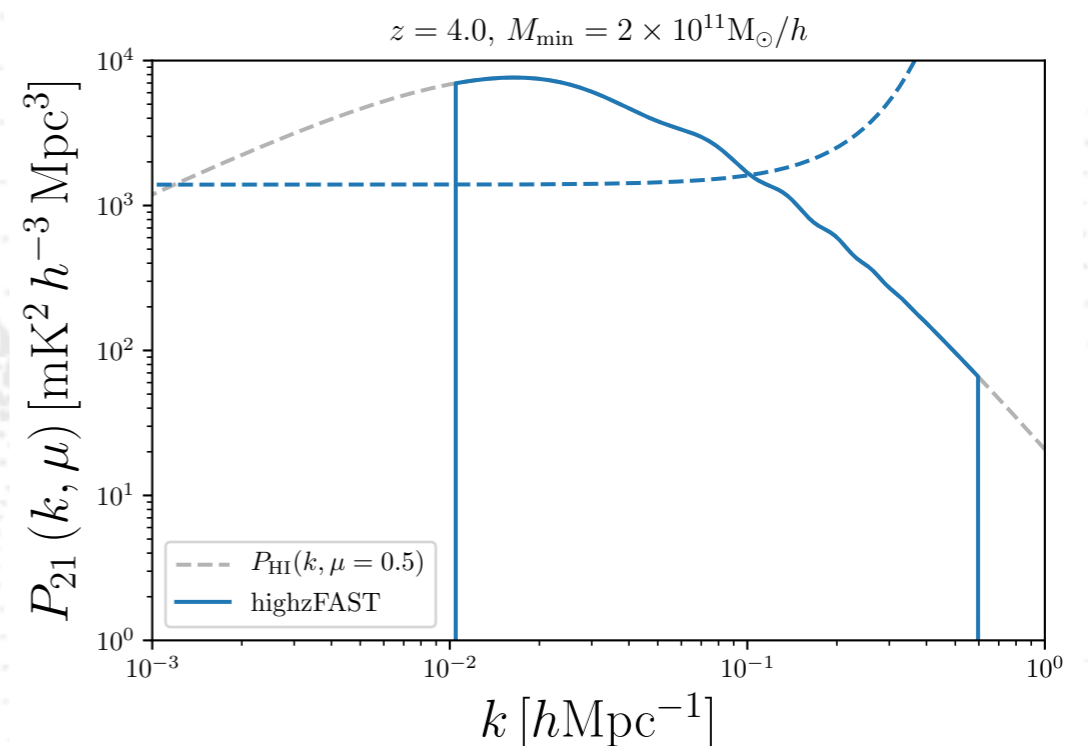
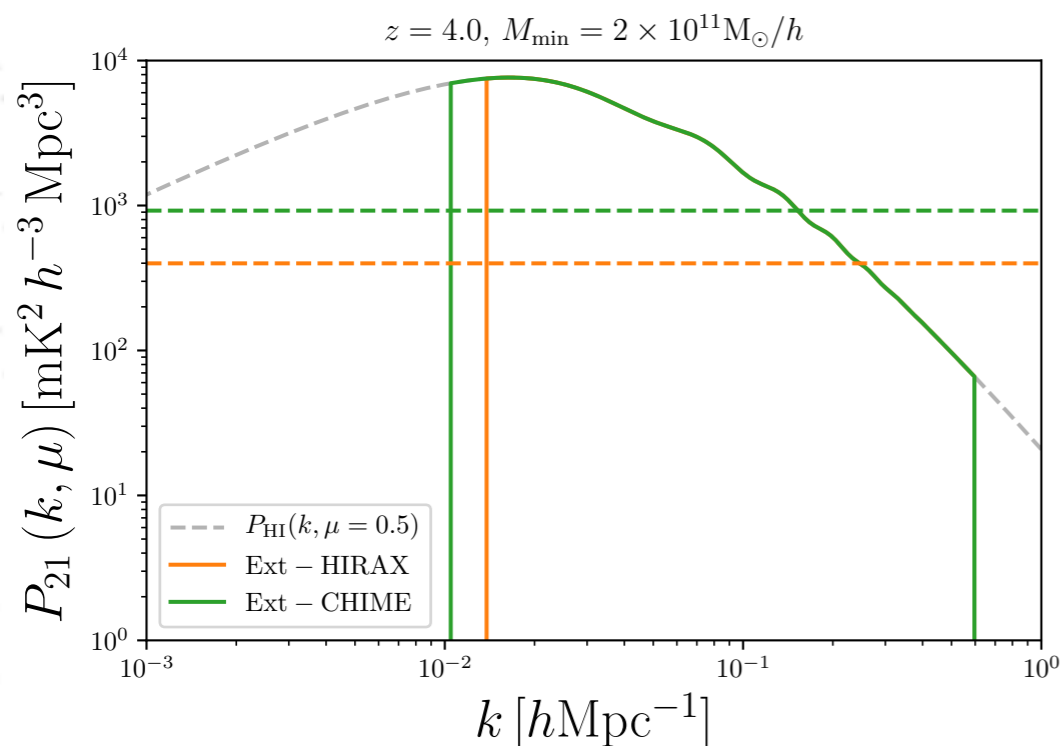
## Observable modes

$$k_{\perp}^{\text{min}}(z) = \frac{2\pi D_{\text{min}}}{D(z)\lambda(z)}$$

$$k_{\perp}^{\text{max}}(z) = \frac{2\pi D_{\text{max}}}{D(z)\lambda(z)}$$

$$k_{\perp}^{\text{min}}(z) = \frac{2\pi}{\sqrt{D(z)^2 S_{21}}}$$

$$k_{\perp}^{\text{max}}(z) = \frac{2\pi D_{\text{dish}}}{D(z)\lambda(z)}$$





# Thermal noise

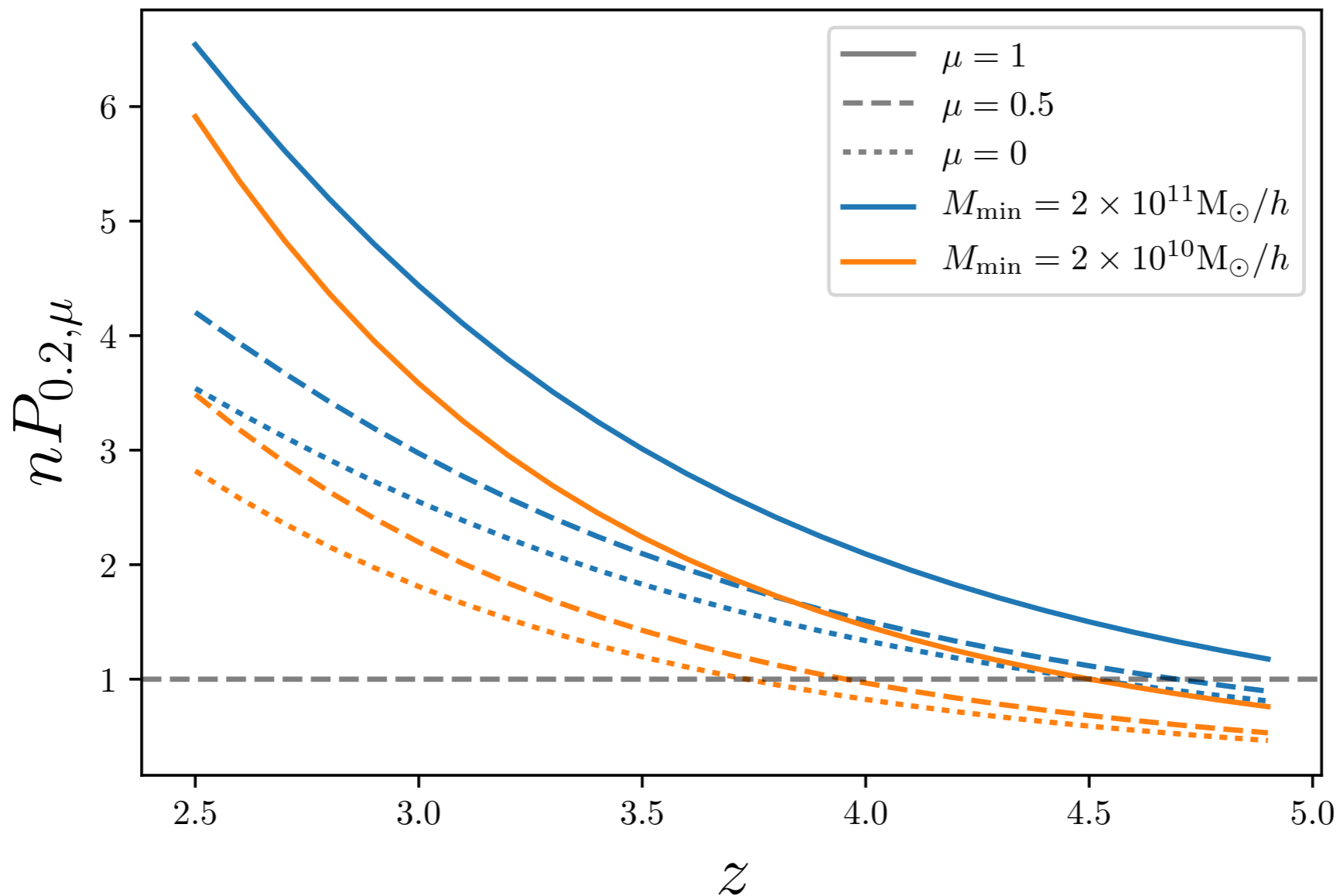
## Interferometers

$$P_N^{\text{th}}(z) = \frac{T_{\text{sys}}^2(z) X^2(z) Y(z) \lambda^4(z) S_{21}}{A_{\text{eff}}^2 \text{FOV}(z) t_0 n_{\text{pol}} n(\mathbf{u}, z)}$$

## Single-dish

$$P_N^{\text{th}}(k, \mu) = \frac{T_{\text{sys}}^2 V_{\text{pix}} W^{-2}(k_{\perp})}{\Delta\nu t_{\text{obs}} \Omega_{\text{pix}} / S_{21} N_{\text{dish}} N_{\text{beam}}}$$

Ext – HIRAX

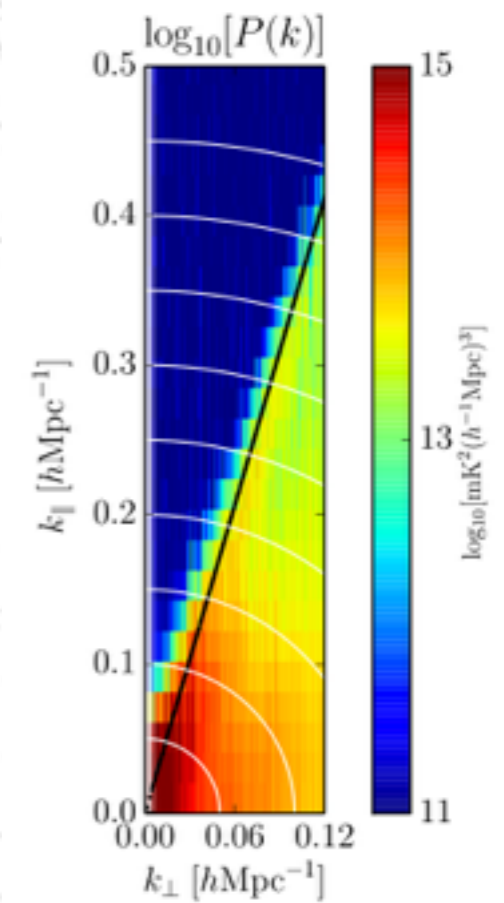


# The foreground wedge

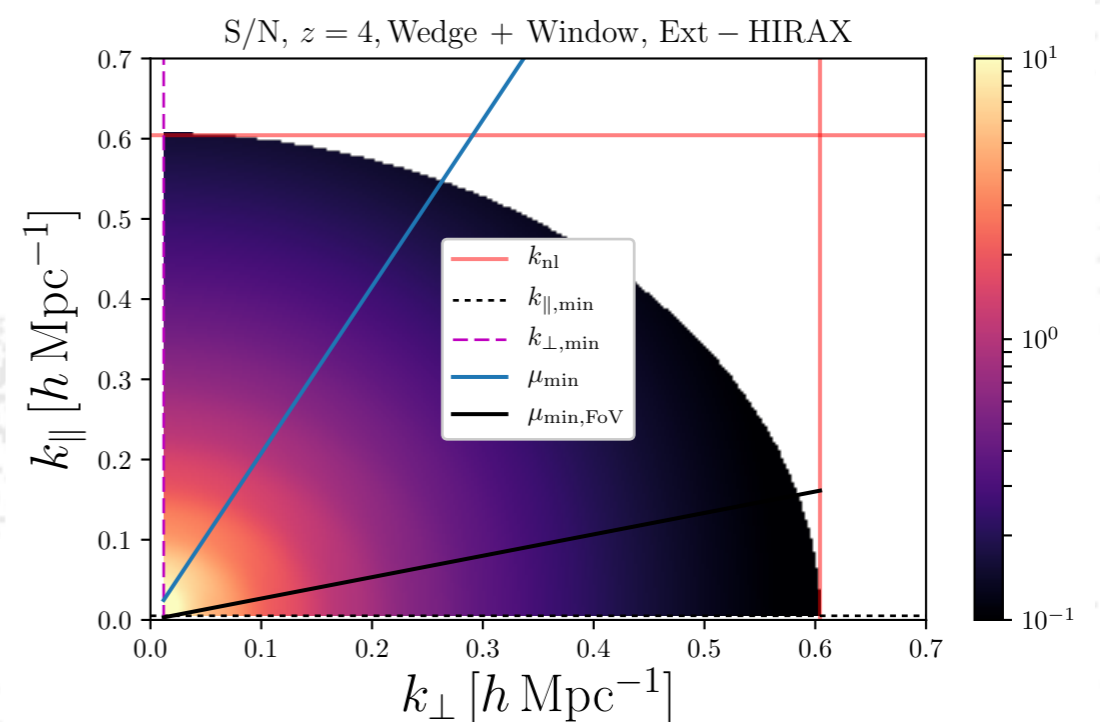
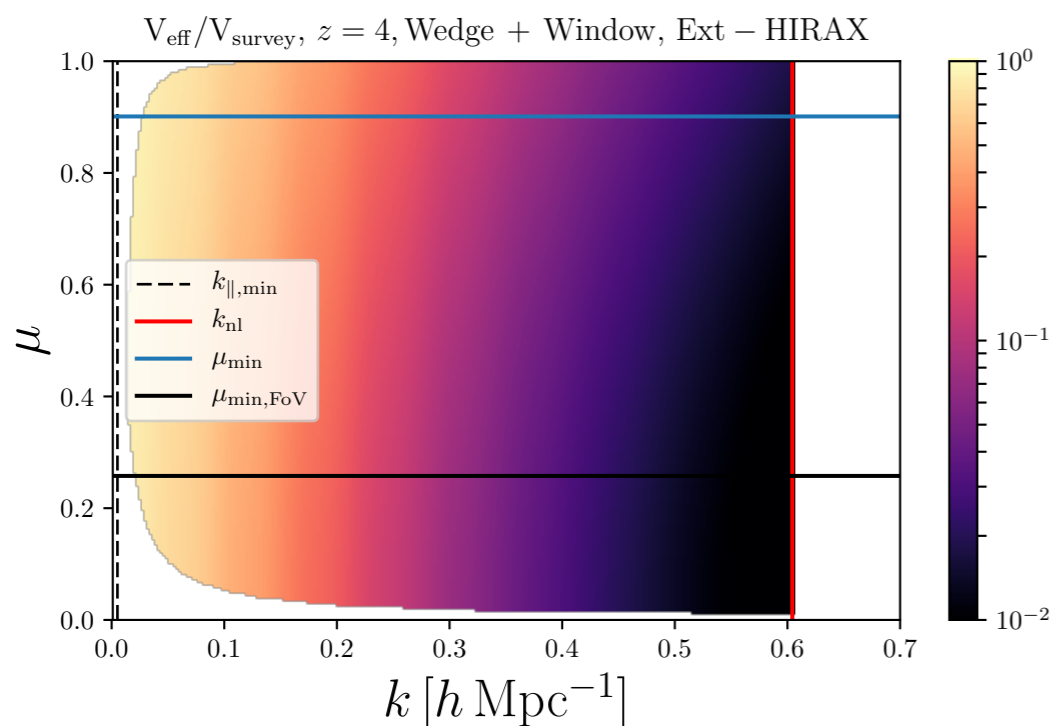
Foregrounds are expected to be smooth in frequency, should affect only low parallel modes, but the interferometer chromaticity makes them less smooth as we go to larger baselines, i.e. larger transverse modes.

The wedge: 
$$k_{\parallel} < \sin(\theta_{\text{FoV}}) \frac{D_c(z)H(z)}{c(1+z)} k_{\perp}$$

$$\mu_{\text{min}}(z) = \frac{k_{\parallel}}{\sqrt{k_{\parallel}^2 + k_{\perp}^2}}$$



Pober et al, 2015



# Fisher forecasts

Full power spectrum:

$$P_{21}(k_f, \mu_f, z) = \bar{T}_b^2(z) \frac{D_A(z)_f^2 H(z)}{D_A(z)^2 H(z)_f} (b_{\text{HI}} \sigma_8(z) + f \sigma_8(z) \mu^2)^2 \frac{P(k, z)}{\sigma_{8,f}^2}$$

Fisher matrix:

$$F_{ij} = \frac{1}{8\pi^2} \int_{-1}^1 d\mu \int k^2 dk \frac{\partial \ln P_{21}(k, \mu)}{\partial p_i} \frac{\partial \ln P_{21}(k, \mu)}{\partial p_j} V_{\text{eff}}(k, \mu)$$

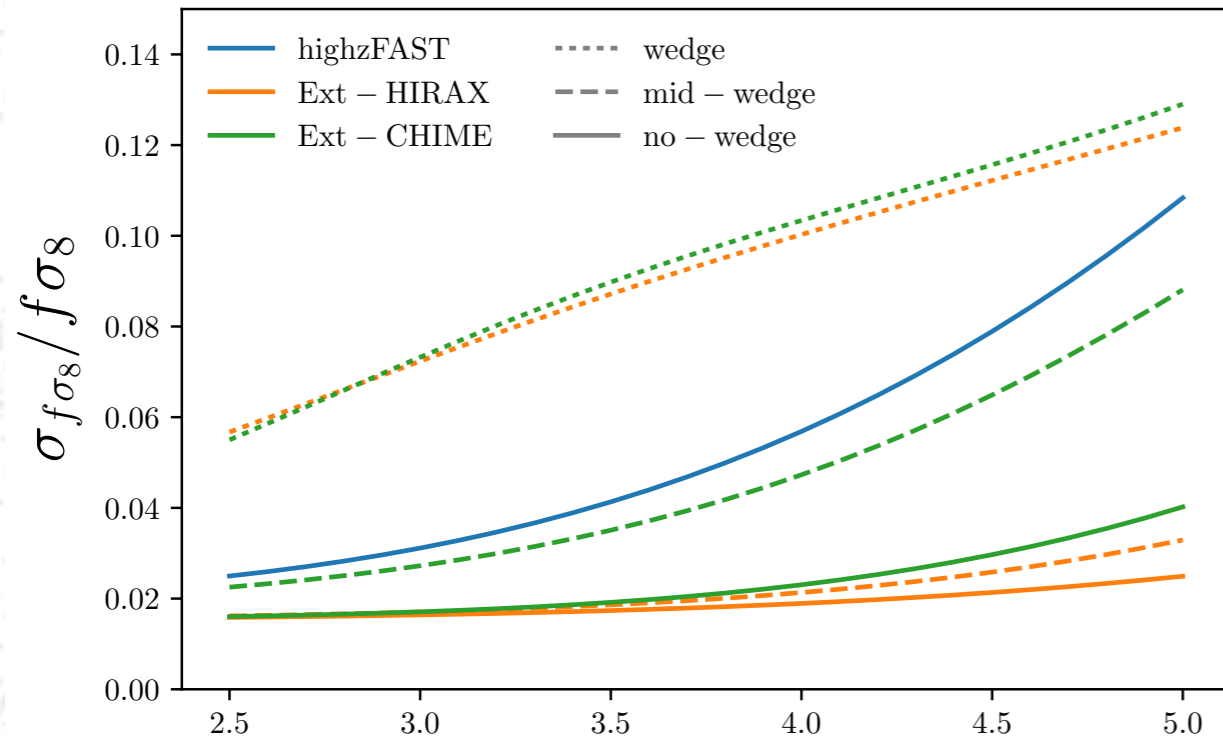
$$V_{\text{eff}}(k, \mu) = V_{\text{sur}} \left( \frac{P_{21}(k, \mu) W(k, \mu)}{P_{21}(k, \mu) W(k, \mu) + P_{\text{N}}^{\text{tot}}(k, \mu)} \right)^2$$

We use two different  $k_{\text{max}}$  scales — 0.2 &  $k_{\text{nl}}(z) = 0.2(1+z)^{2/3} h/\text{Mpc}$

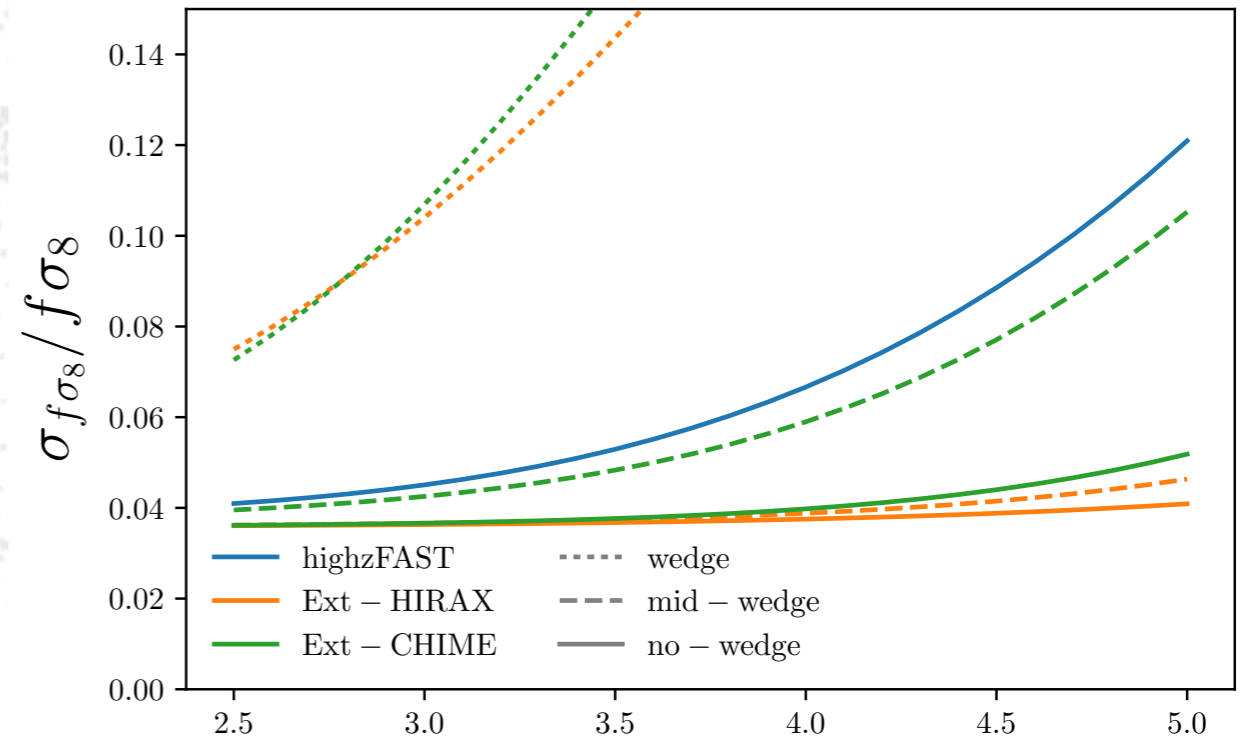
# Growth of structures: $f\sigma_8(z)$

2% priors

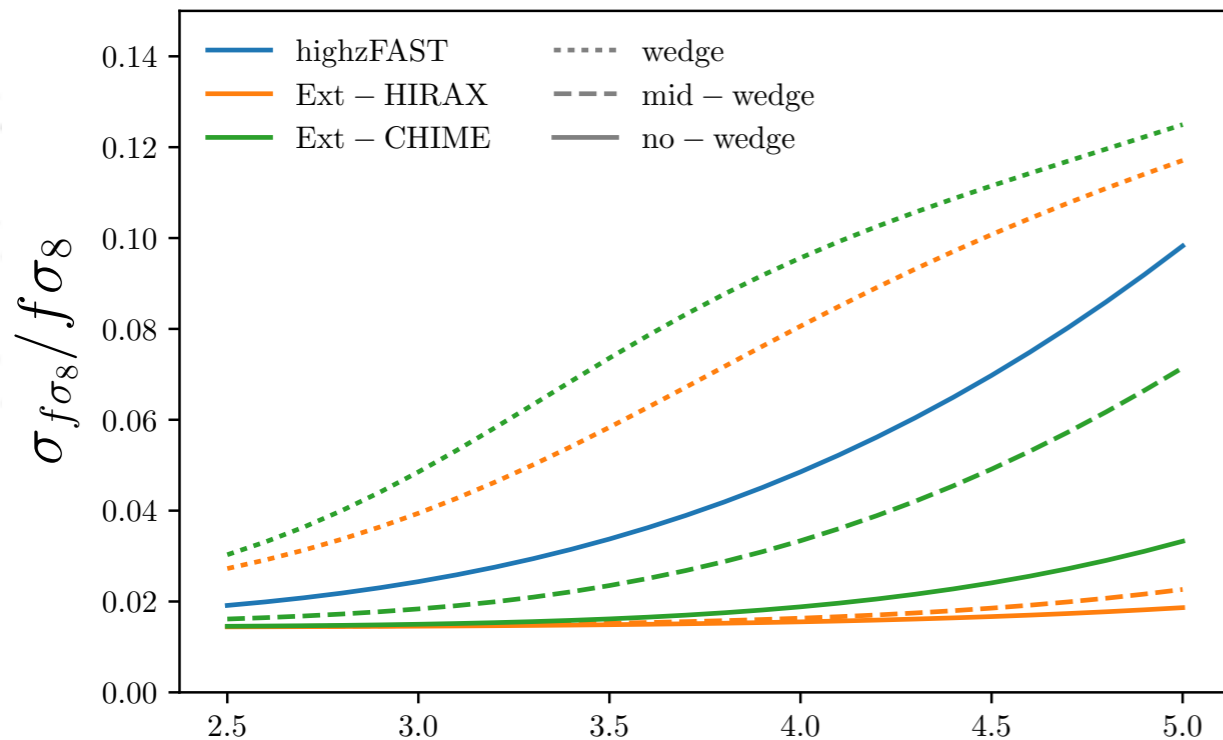
$k_{\max} = 0.2 h\text{Mpc}^{-1}$ , 2% priors –  $b_{\text{HI}}\&\Omega_{\text{HI}}$



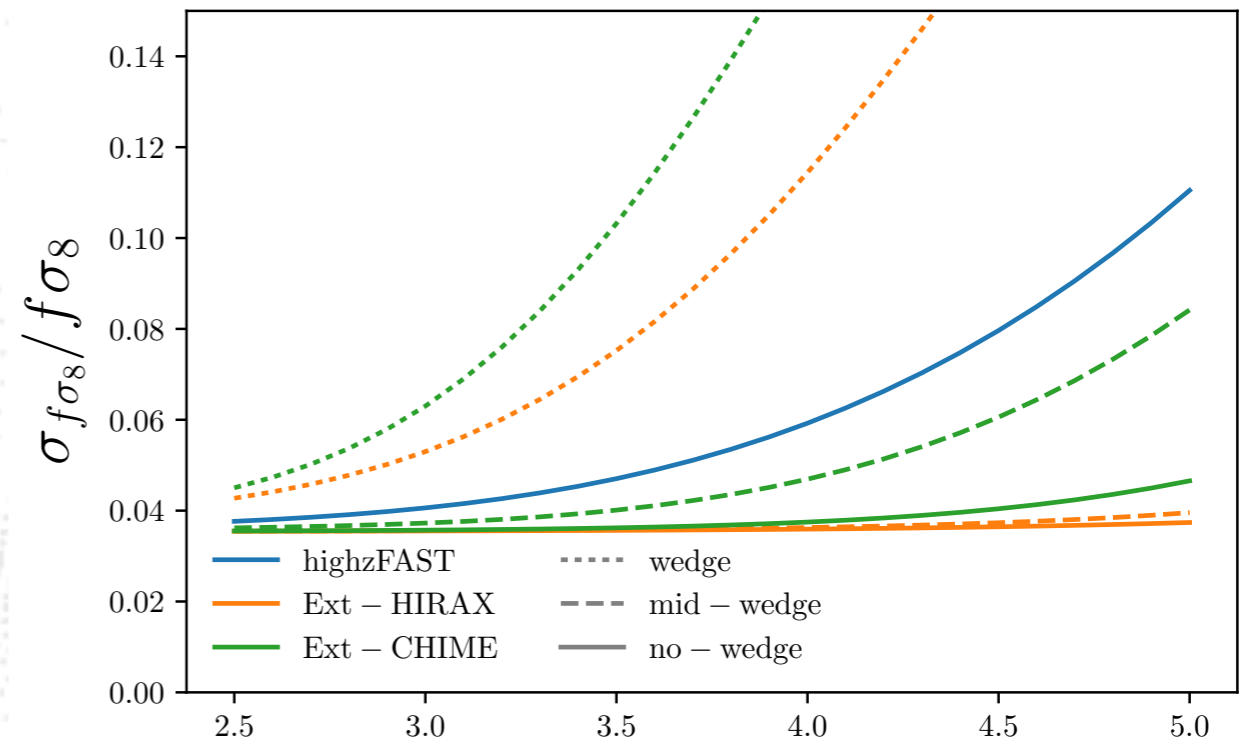
$k_{\max} = 0.2 h\text{Mpc}^{-1}$ , 5% priors –  $b_{\text{HI}}\&\Omega_{\text{HI}}$



$k_{\max} = k_{\text{nl}}(z)$ , 2% priors –  $b_{\text{HI}}\&\Omega_{\text{HI}}$



$k_{\max} = k_{\text{nl}}(z)$ , 5% priors –  $b_{\text{HI}}\&\Omega_{\text{HI}}$

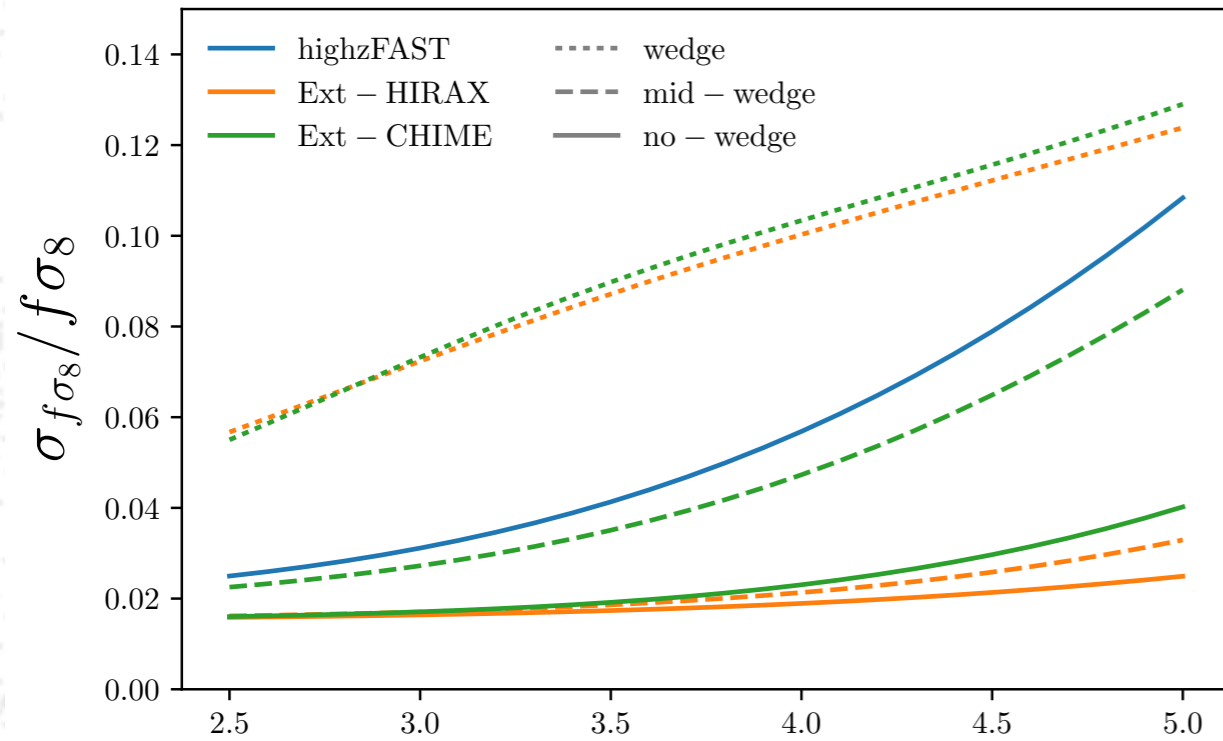


# Growth of structures: $f\sigma_8(z)$

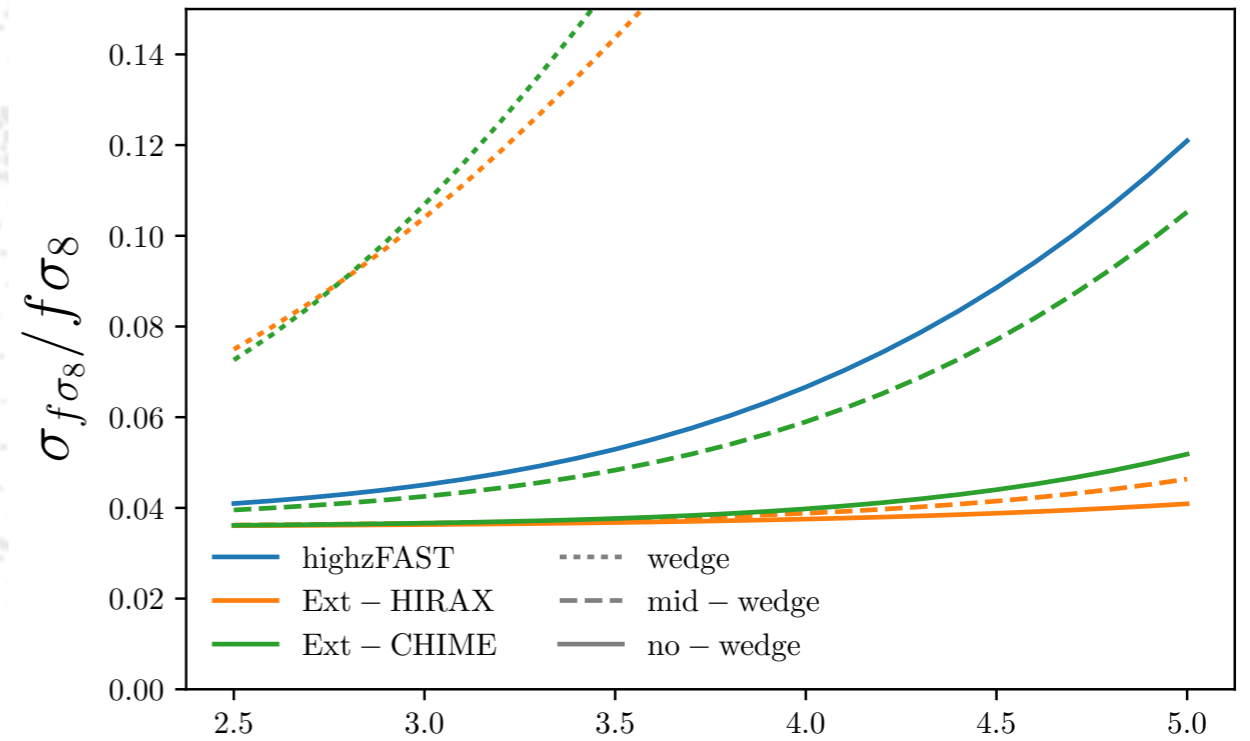
2% priors

5% priors

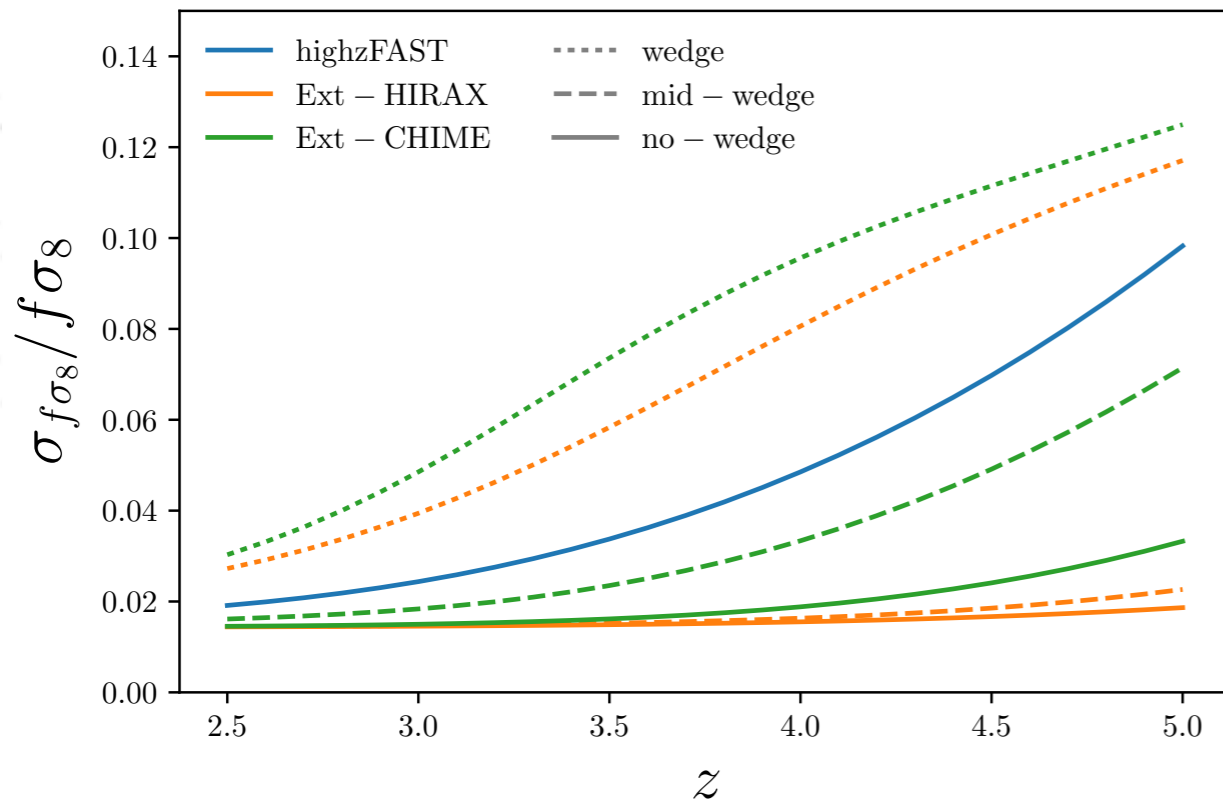
$k_{\max} = 0.2 h\text{Mpc}^{-1}$ , 2% priors -  $b_{\text{HI}}\&\Omega_{\text{HI}}$



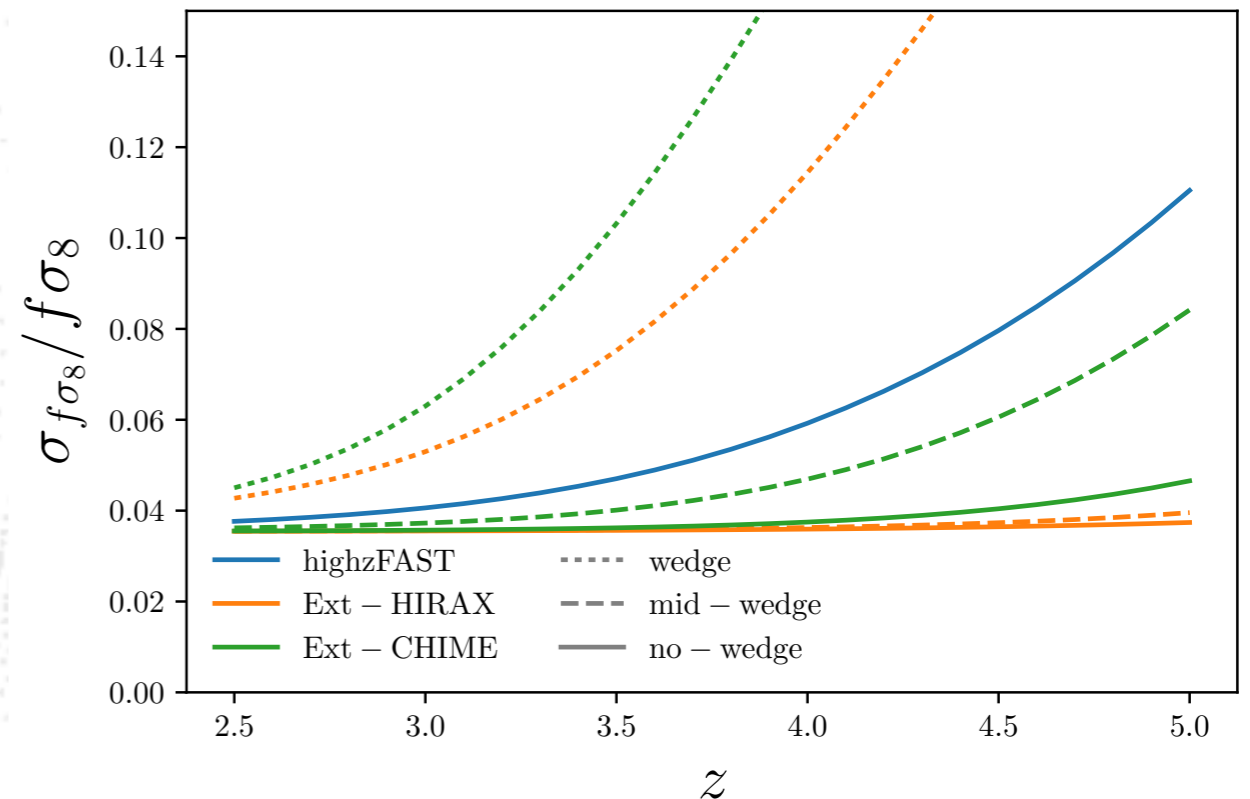
$k_{\max} = 0.2 h\text{Mpc}^{-1}$ , 5% priors -  $b_{\text{HI}}\&\Omega_{\text{HI}}$



$k_{\max} = k_{\text{nl}}(z)$ , 2% priors -  $b_{\text{HI}}\&\Omega_{\text{HI}}$

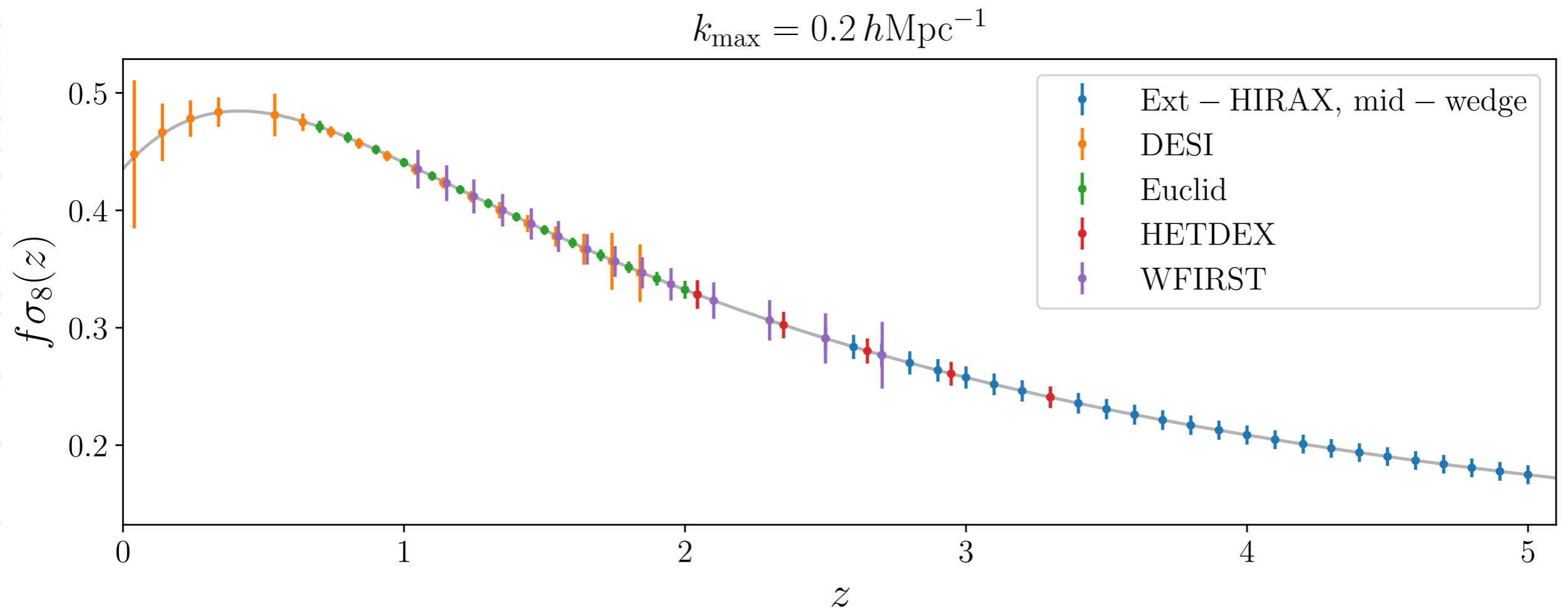


$k_{\max} = k_{\text{nl}}(z)$ , 5% priors -  $b_{\text{HI}}\&\Omega_{\text{HI}}$

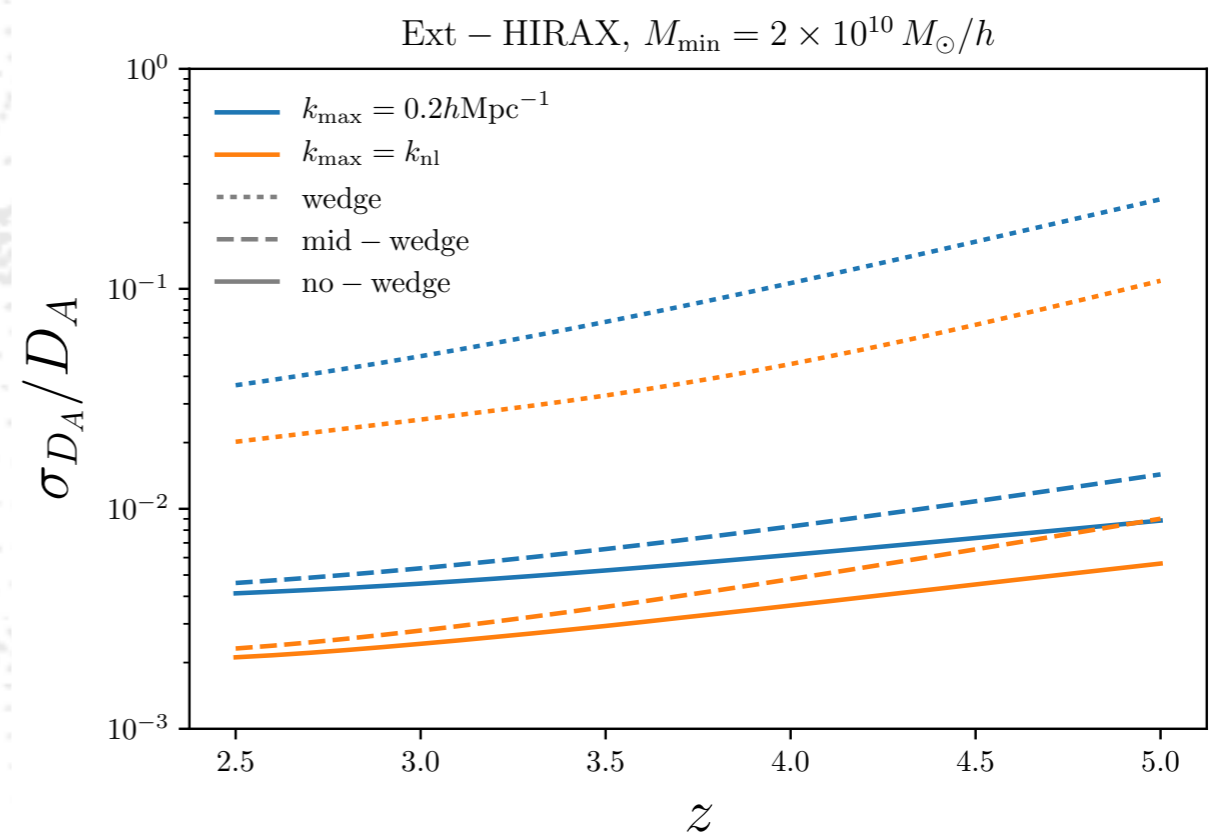
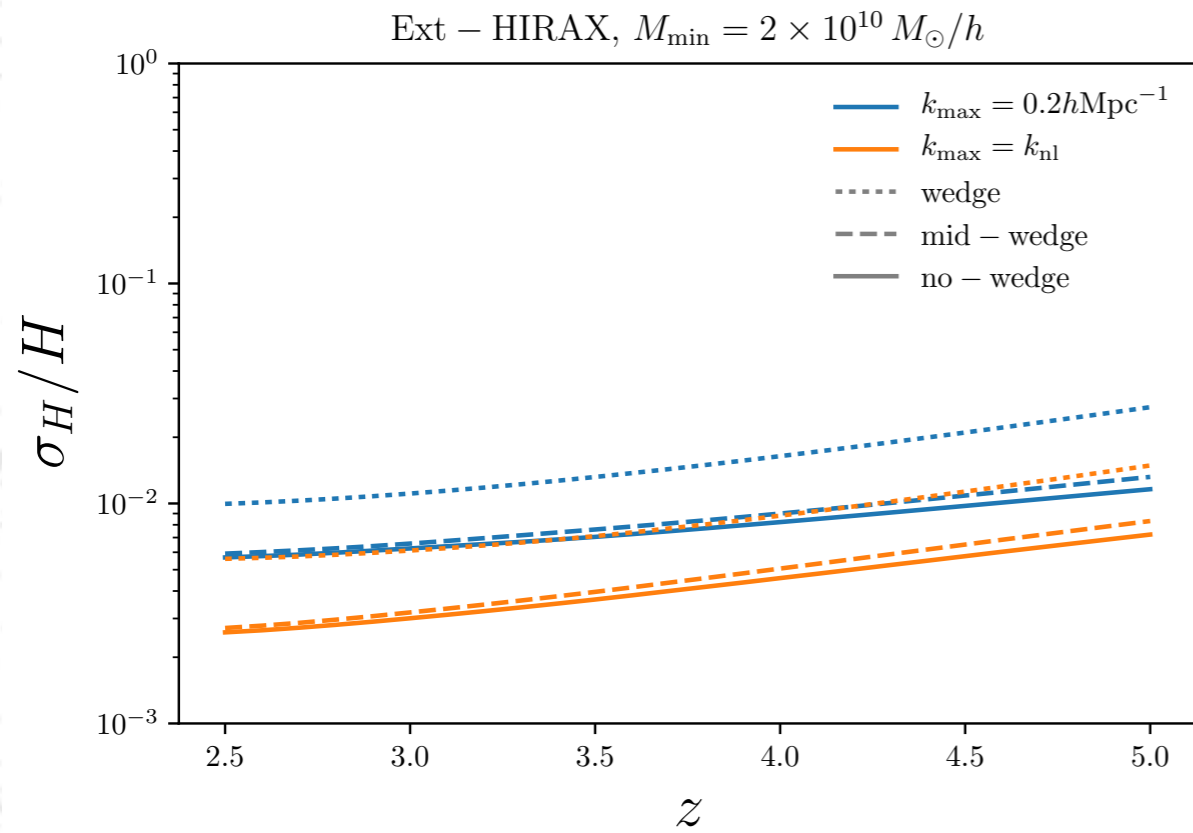


# Growth of structures: $f\sigma_8(z)$

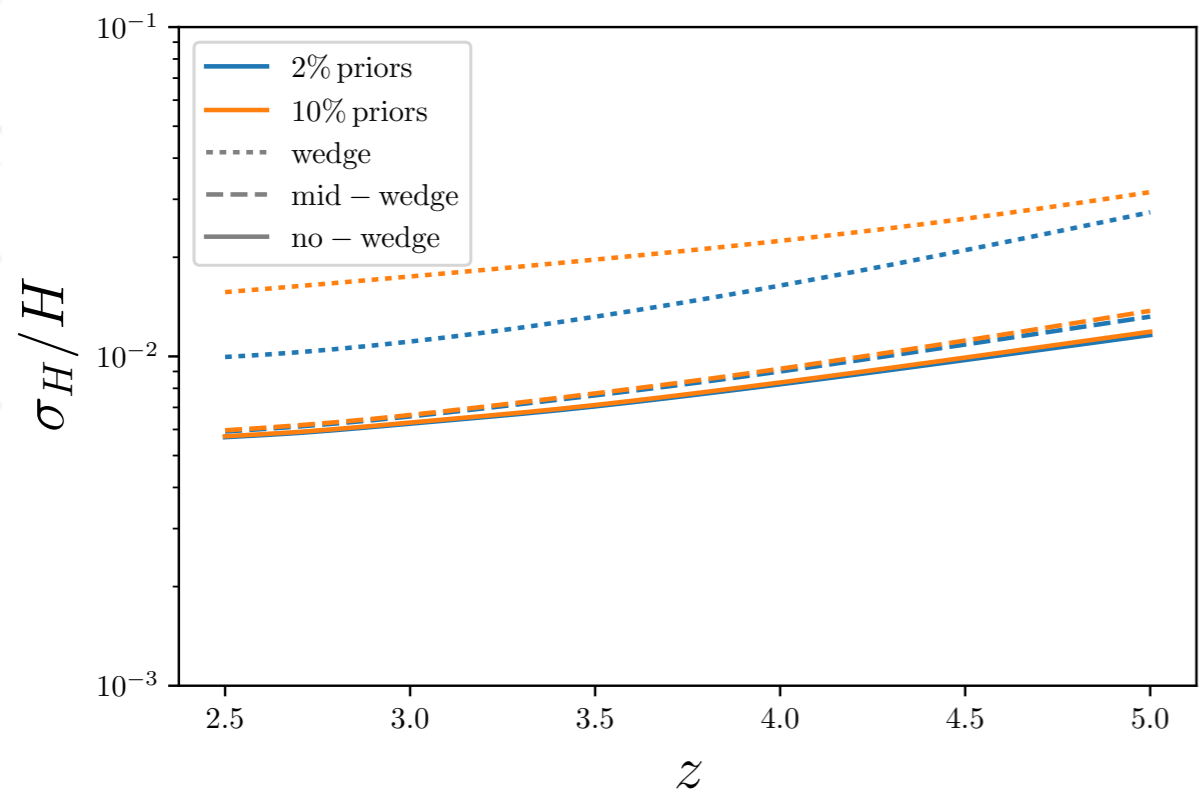
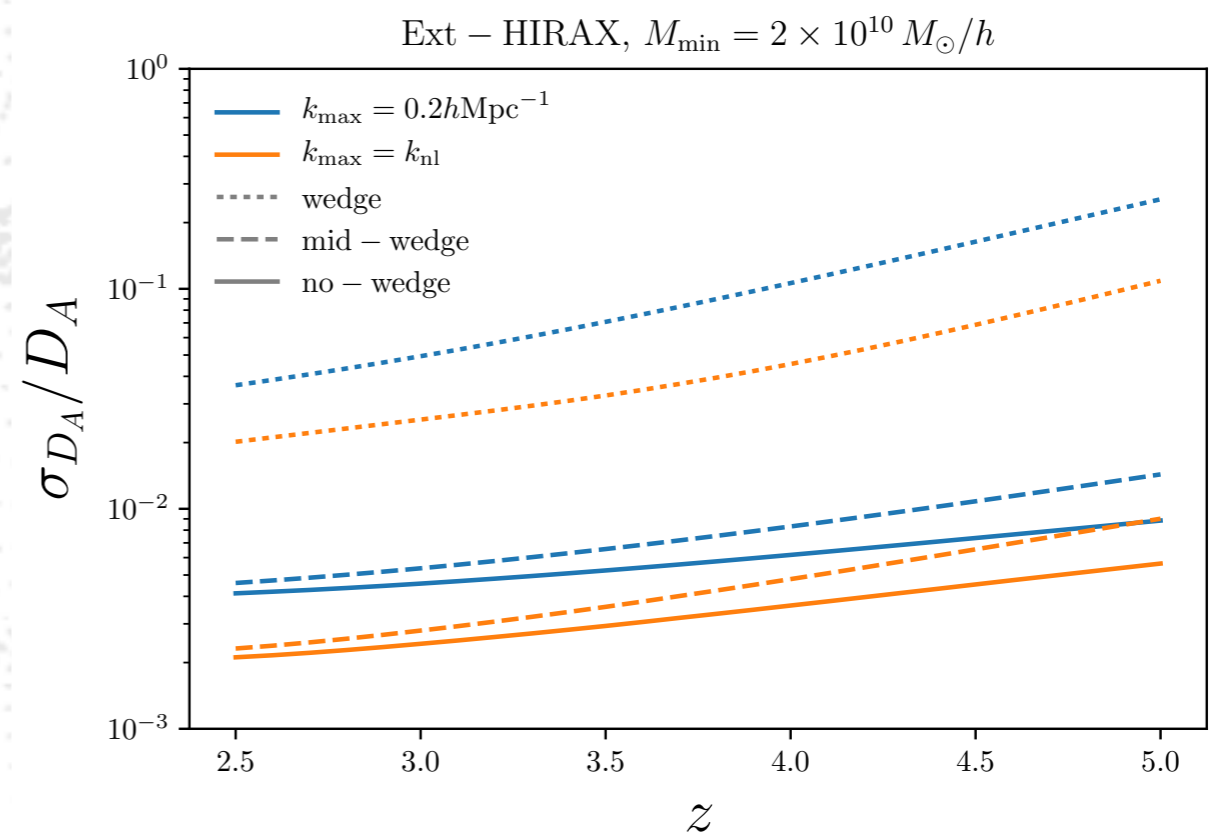
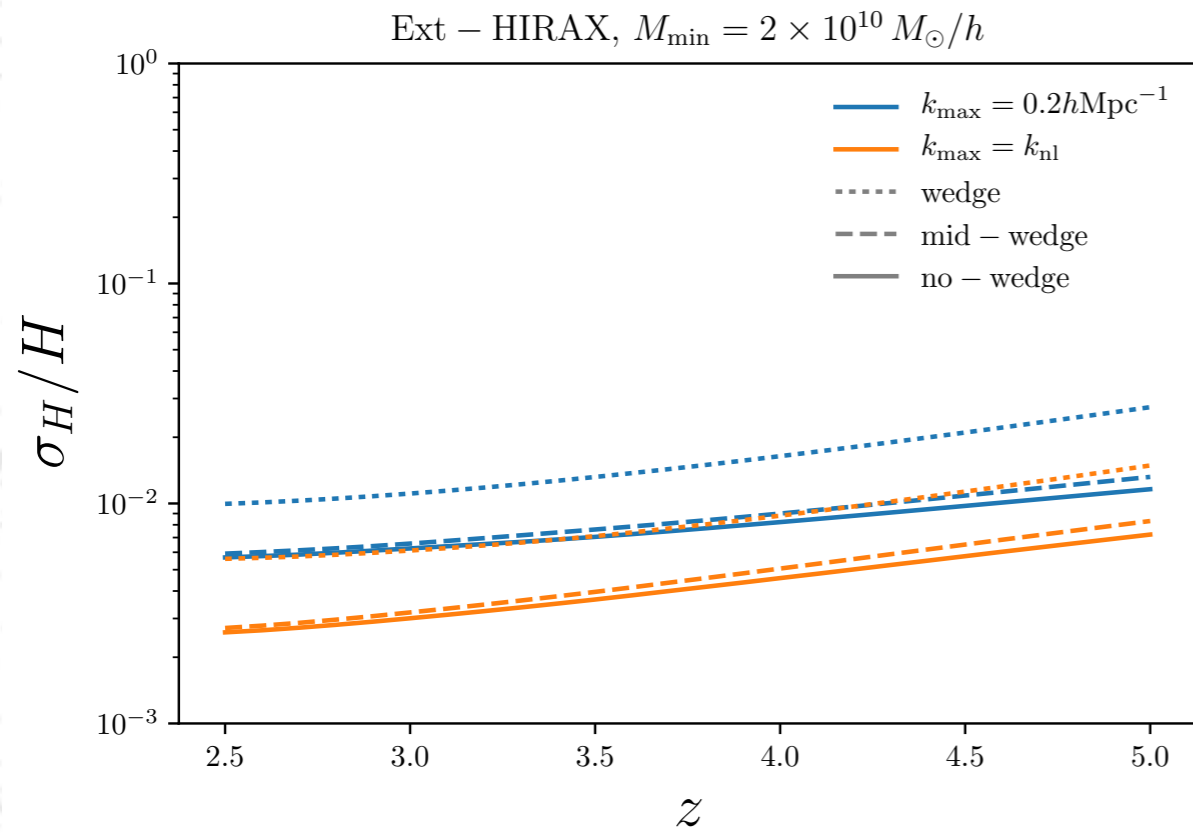
Redshift coverage comparison with other surveys



# BAO: AP parameters $H(z)$ & $D_A(z)$

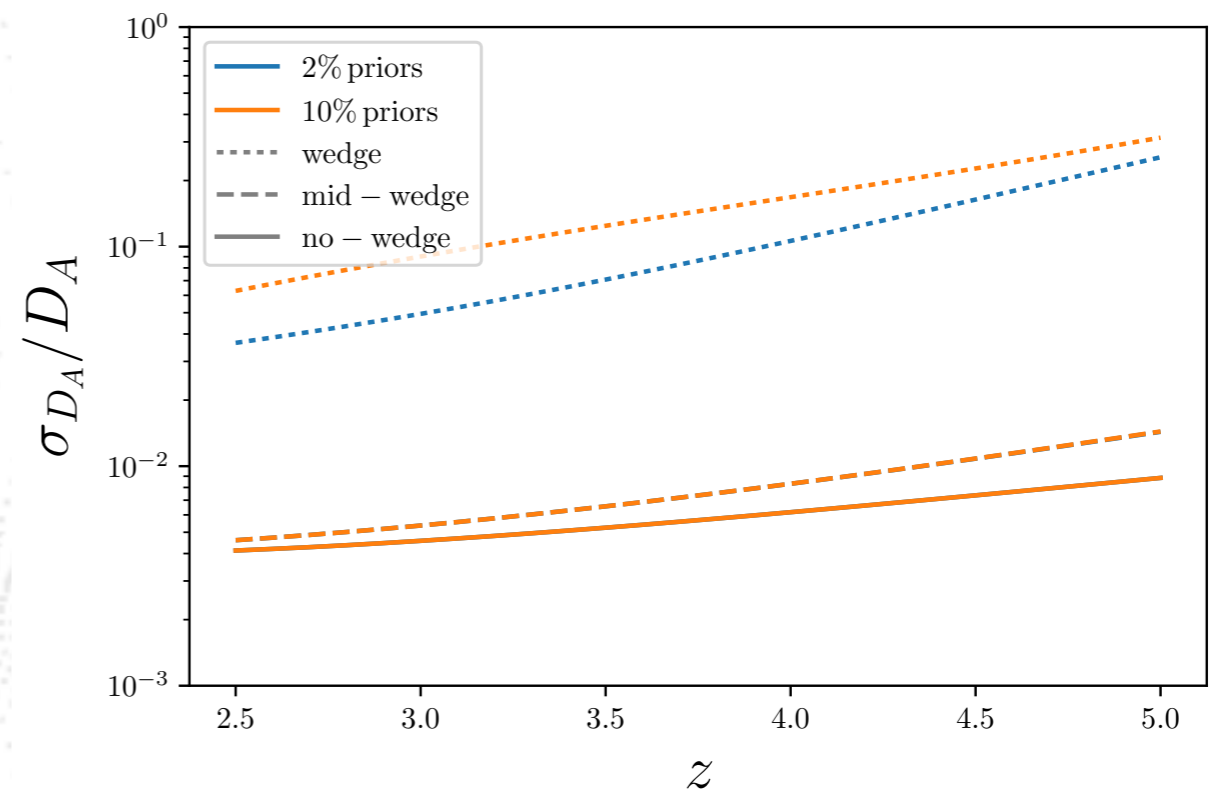
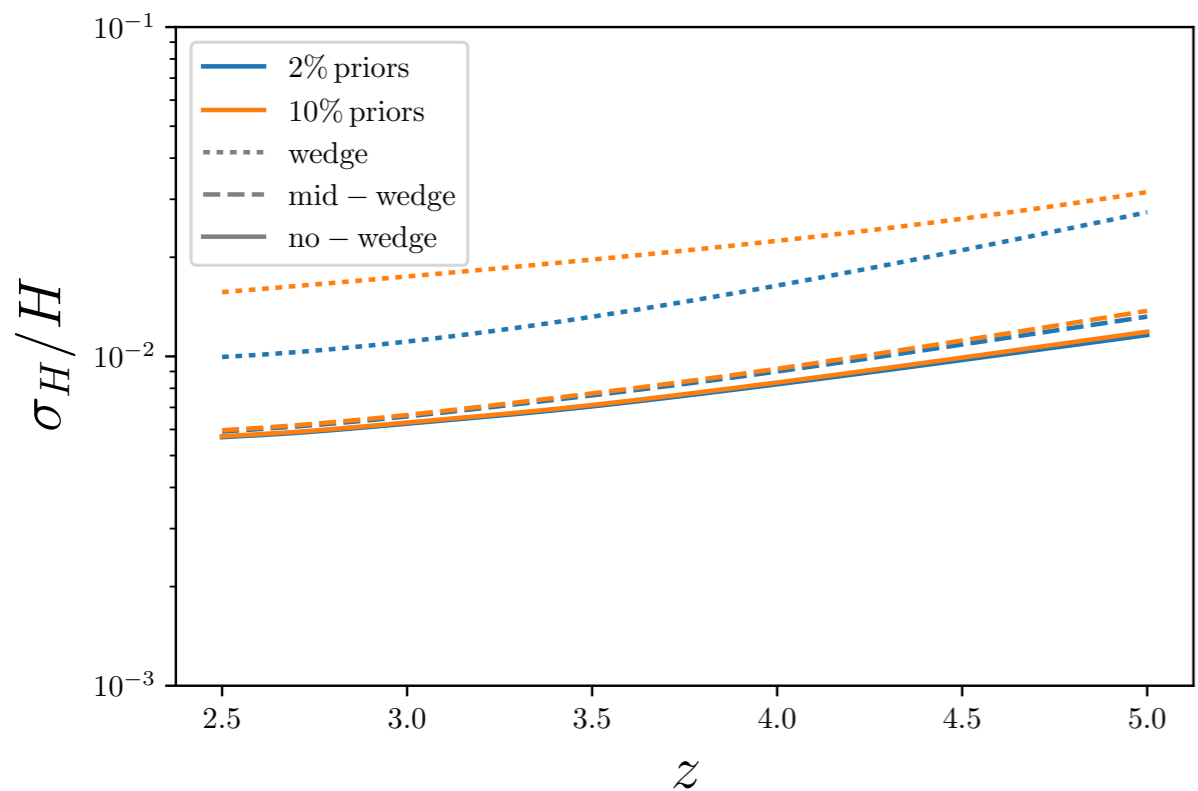
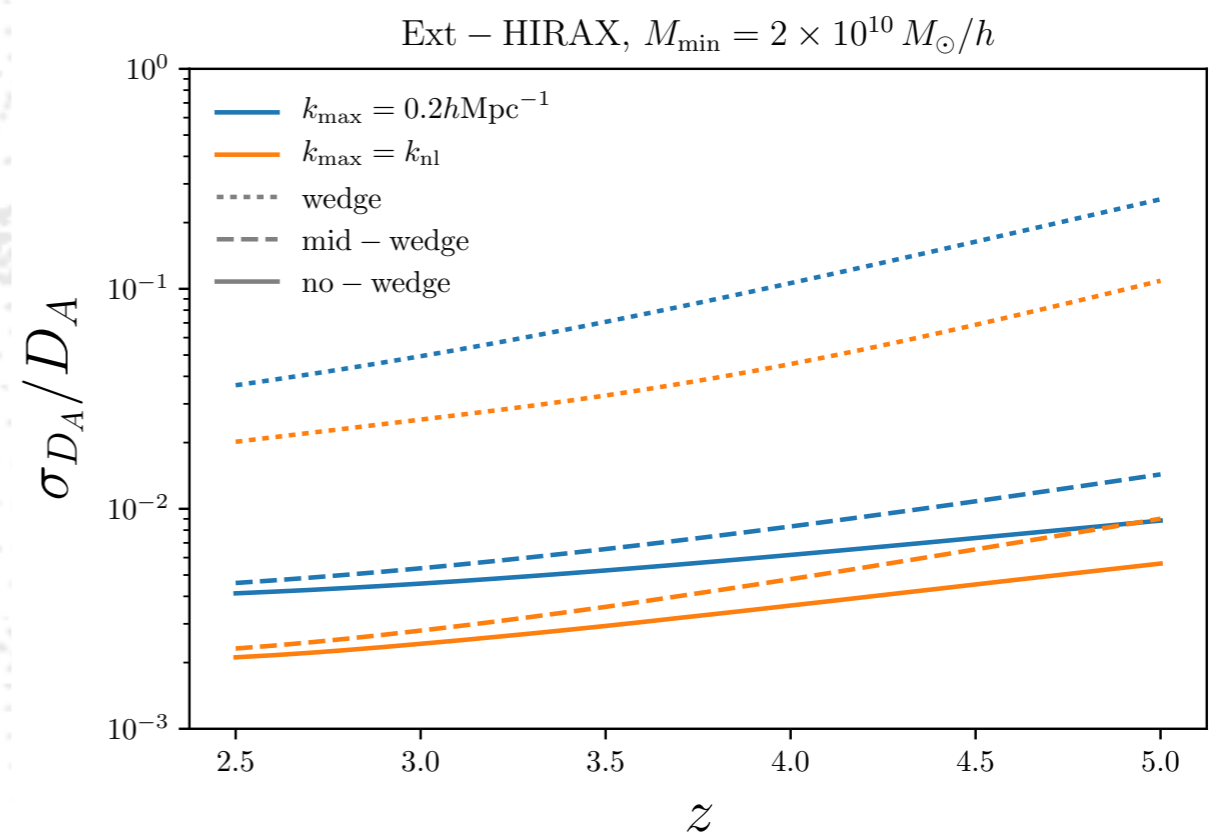
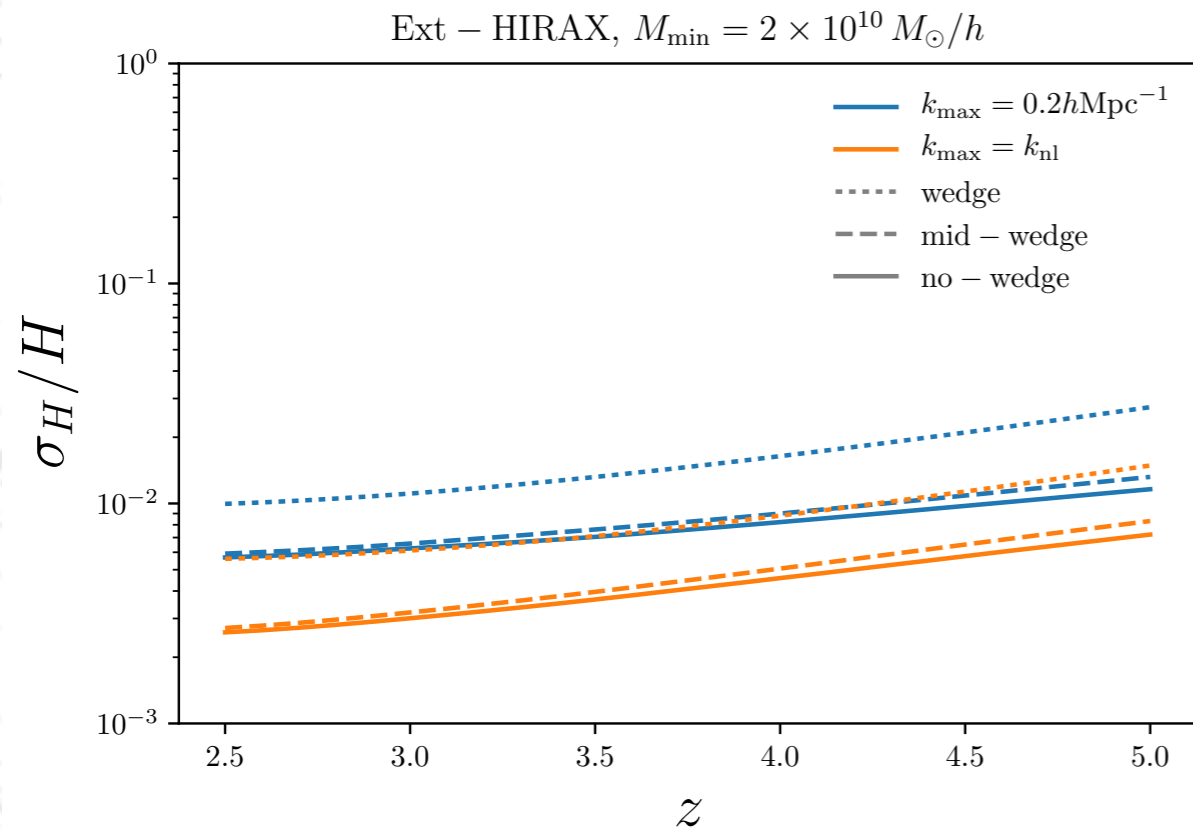


# BAO: AP parameters $H(z)$ & $D_A(z)$



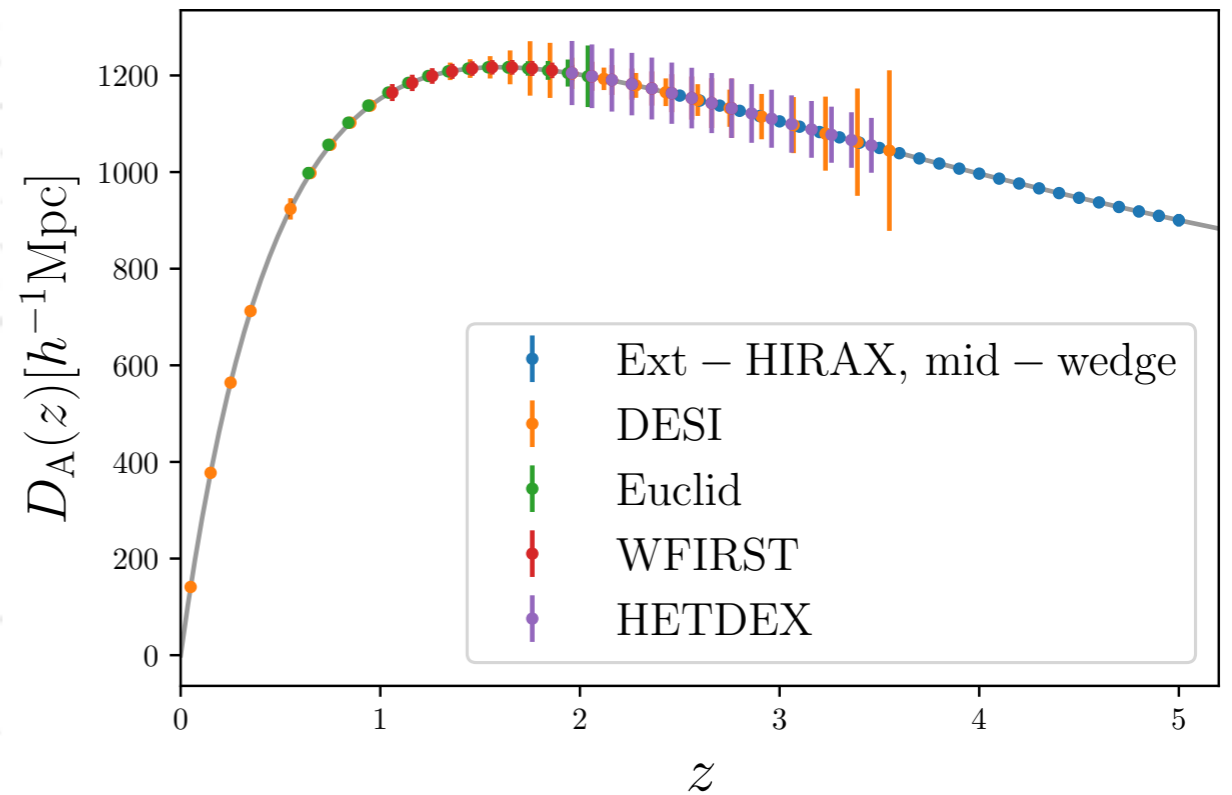
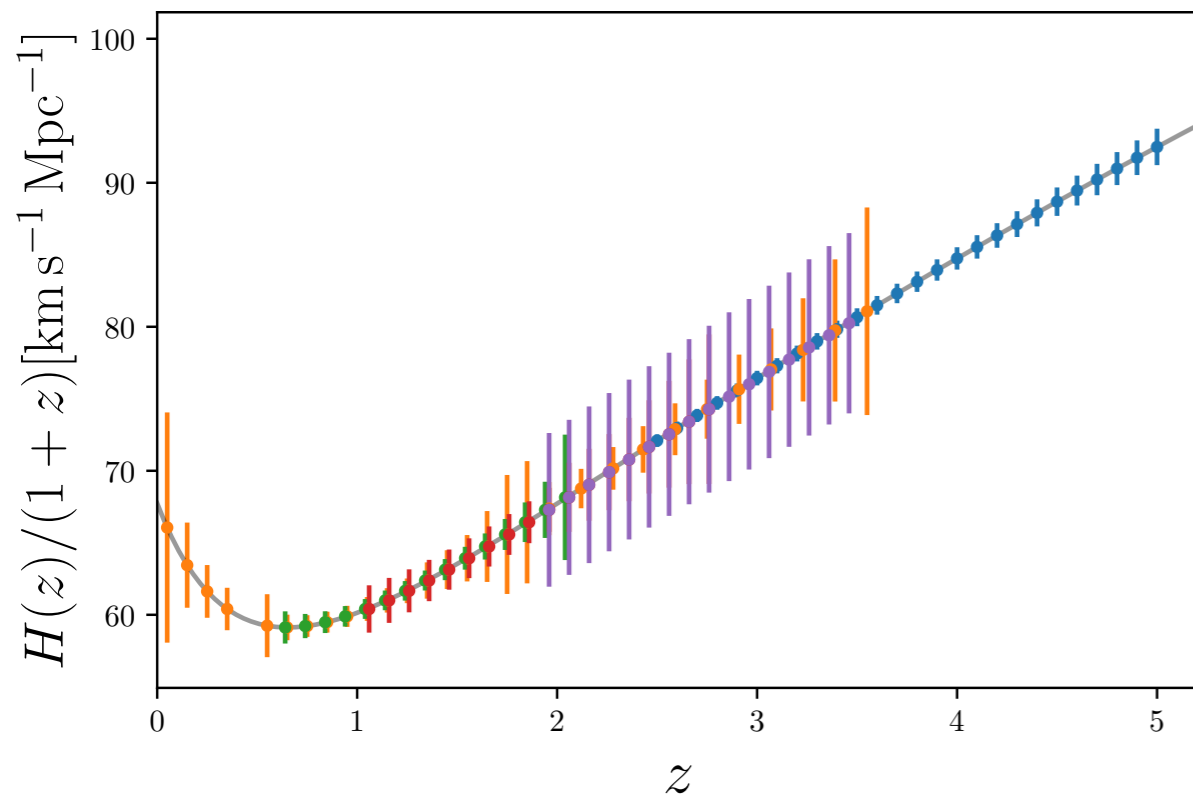


# BAO: AP parameters $H(z)$ & $D_A(z)$



# BAO: AP parameters

Redshift coverage comparison with other surveys



# **LambdaCDM extensions using the broadband shape**

- The sum of neutrino masses
- The effective number of relativistic degrees of freedom

# LambdaCDM extensions using the broadband shape

- The sum of neutrino masses
- The effective number of relativistic degrees of freedom

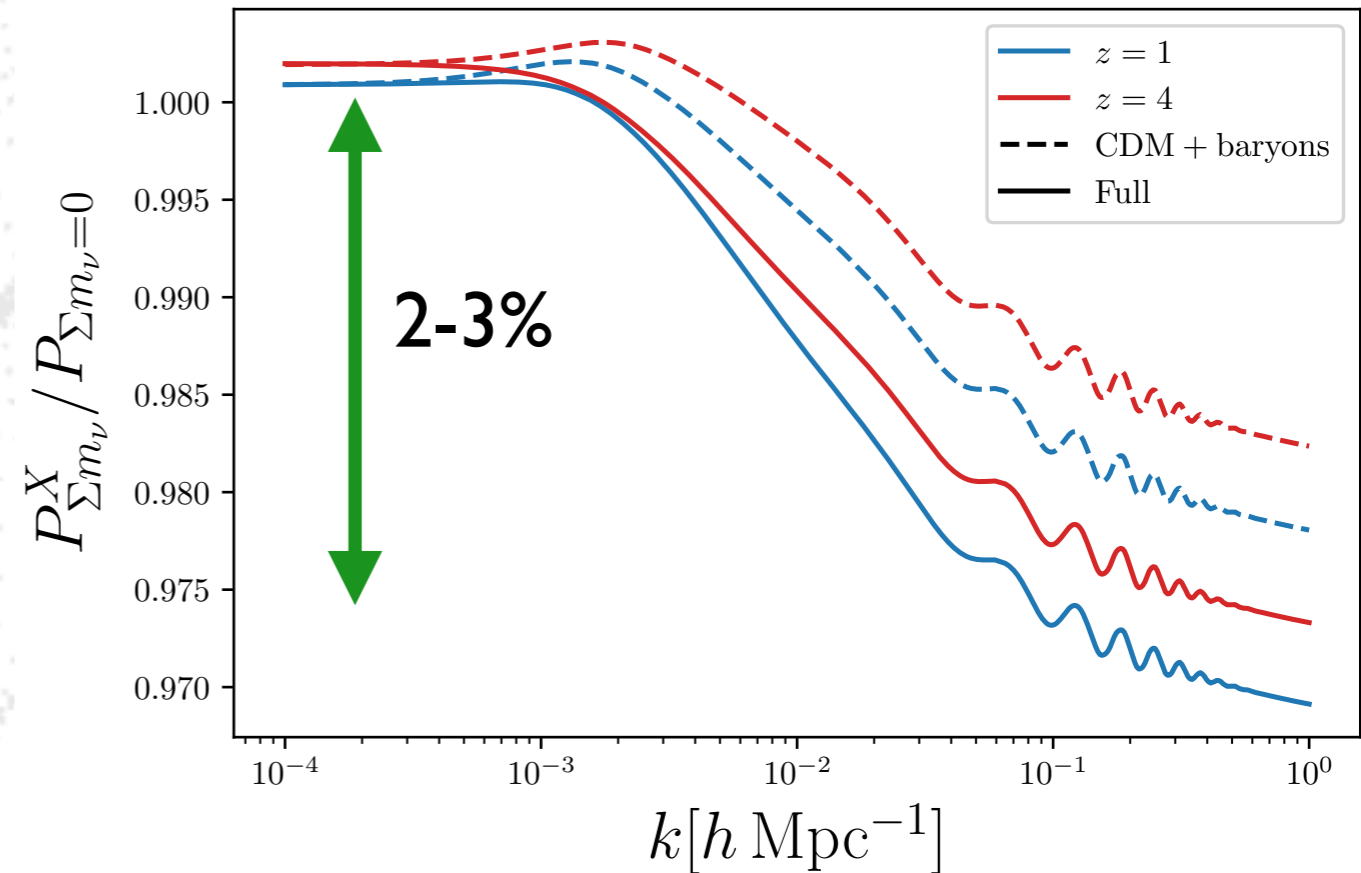
We make use of the synergy with current and future datasets and forecasts:

- Planck 2015 + BOSS BAO
- Future galaxy redshift survey — Euclid
- Future CMB Stage 4

# Broadband shape + massive neutrinos

The effect of massive neutrinos on the power spectrum

CDM+baryons power spectrum



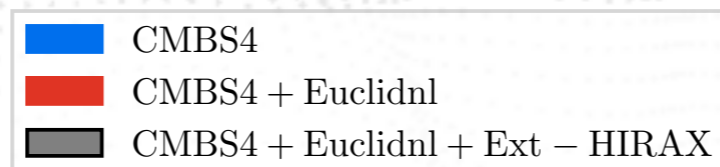
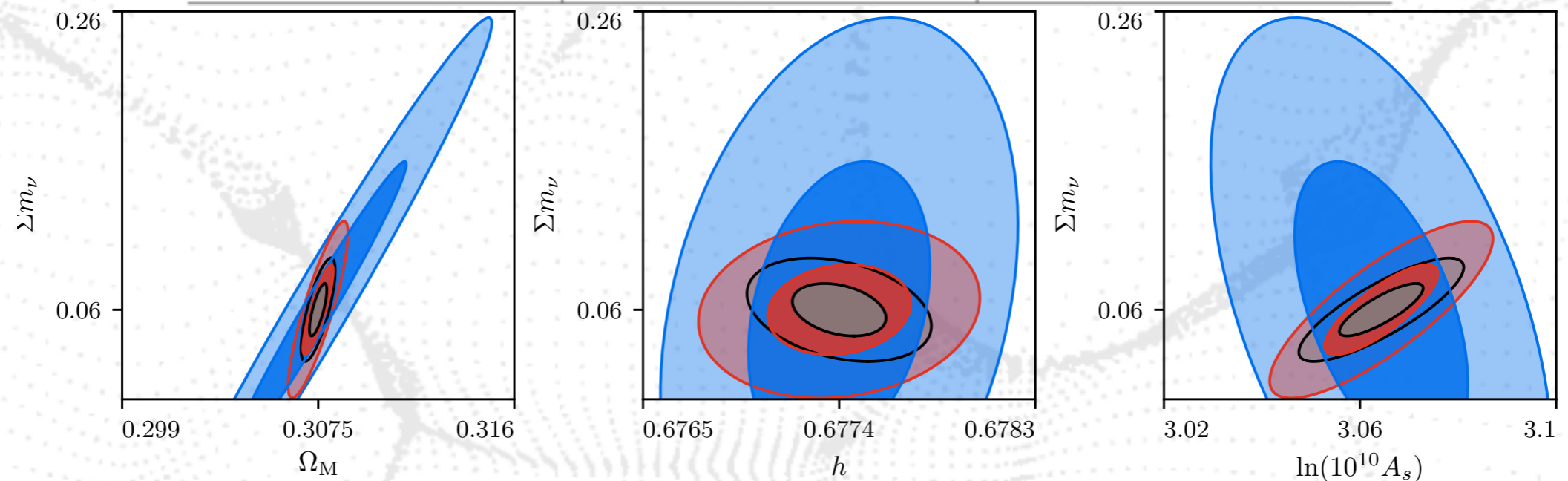
## Ext-HIRAX alone

Ext-HIRAX $\Sigma m_\nu$					
	$\Omega_M$	$h$	$\Sigma m_\nu [\text{eV}]$	$\ln(10^{10} A_s)$	$n_s$
Fiducial values	0.3075	0.6774	0.060	3.064	0.9667
No wedge $k_{\text{max}} = 0.2h\text{Mpc}^{-1}$					
2% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.0016	0.0010	0.059	0.026	0.0036
2% $b_{\text{HI}} \& \Omega_{\text{HI}}$ , diff $M_{\text{min}}$	0.0015	0.0009	0.056	0.025	0.0033
2% $b_{\text{HI}} \& \Omega_{\text{HI}}$ , 1-loop	0.0020	0.0010	0.081	0.038	0.014
5% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.0024	0.0015	0.093	0.047	0.0038
10% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.0029	0.0018	0.11	0.065	0.0040

# Broadband shape + massive neutrinos

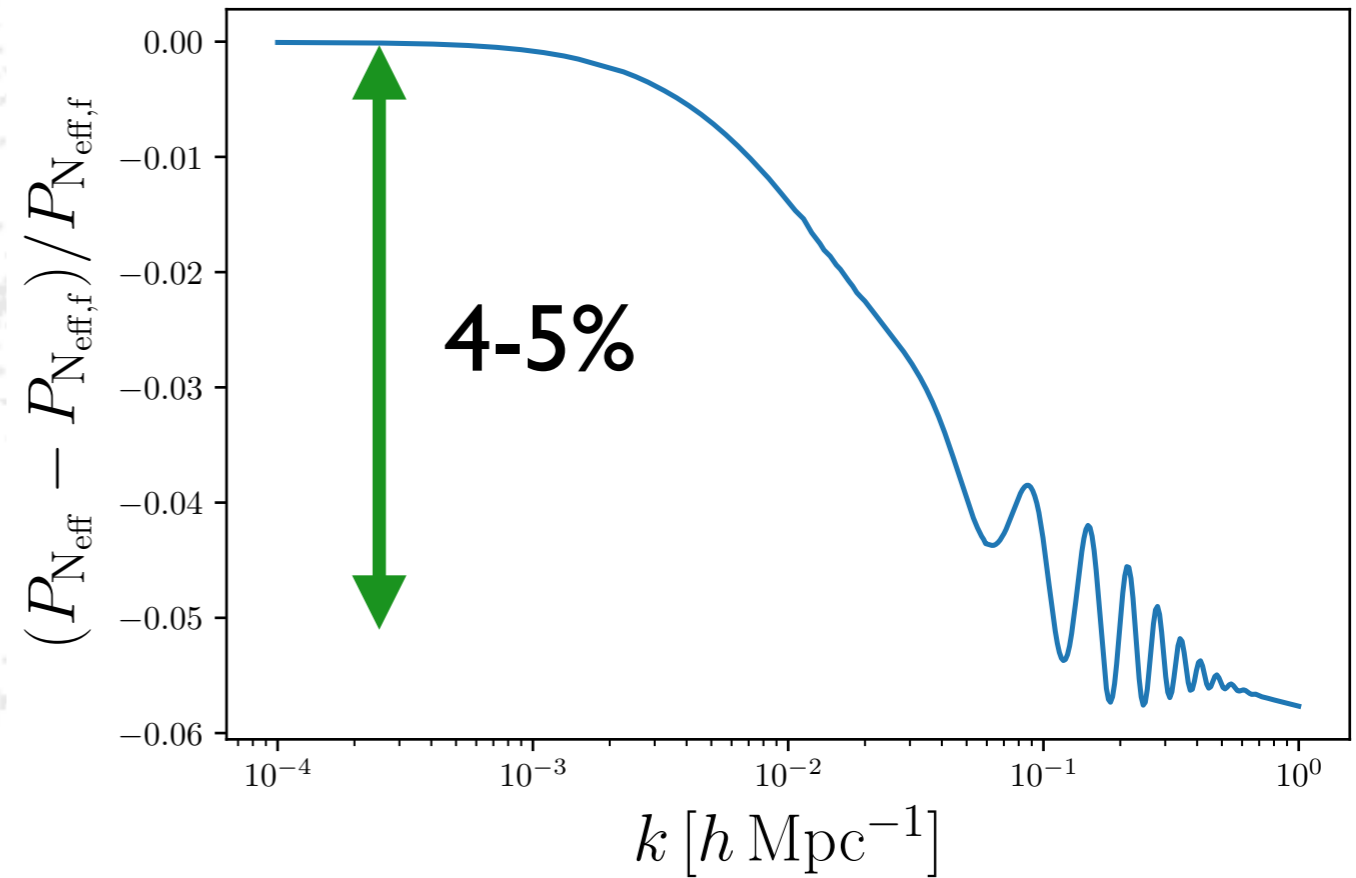
## CMB S4 + Euclid + Ext-HIRAX

Ext-HIRAX $\Sigma m_\nu$ [eV]						
	$k_{\max} = 0.2h\text{Mpc}^{-1}$			$k_{\max} = k_{\text{nl}}(z)$		
Euclid No Rec.+PlanckBAO	0.050			0.049		
+21cm	No wedge	Mid wedge	Wedge	No wedge	Mid wedge	Wedge
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.037	0.038	0.045	0.030	0.031	0.040
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$ , diff $M_{\min}$	0.037	0.038	0.044	0.028	0.029	0.039
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$ , 1-loop	0.038	0.039	0.046	0.035	0.035	0.042
5% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.042	0.043	0.048	0.034	0.035	0.043
10% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.044	0.045	0.049	0.035	0.036	0.044
Euclid No Rec.+CMB-S4	0.031			0.030		
+21cm	No wedge	Mid wedge	Wedge	No wedge	Mid wedge	Wedge
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.022	0.022	0.028	0.018	0.018	0.025
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$ , diff $M_{\min}$	0.021	0.021	0.028	0.017	0.017	0.024
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$ , 1-loop	0.023	0.023	0.028	0.020	0.020	0.025
5% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.023	0.023	0.028	0.018	0.019	0.025
10% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.023	0.023	0.028	0.019	0.019	0.025



# Broadband shape + Neff

The effect of Neff on the power spectrum



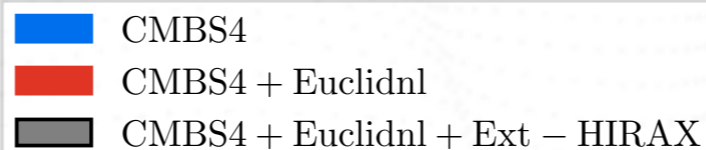
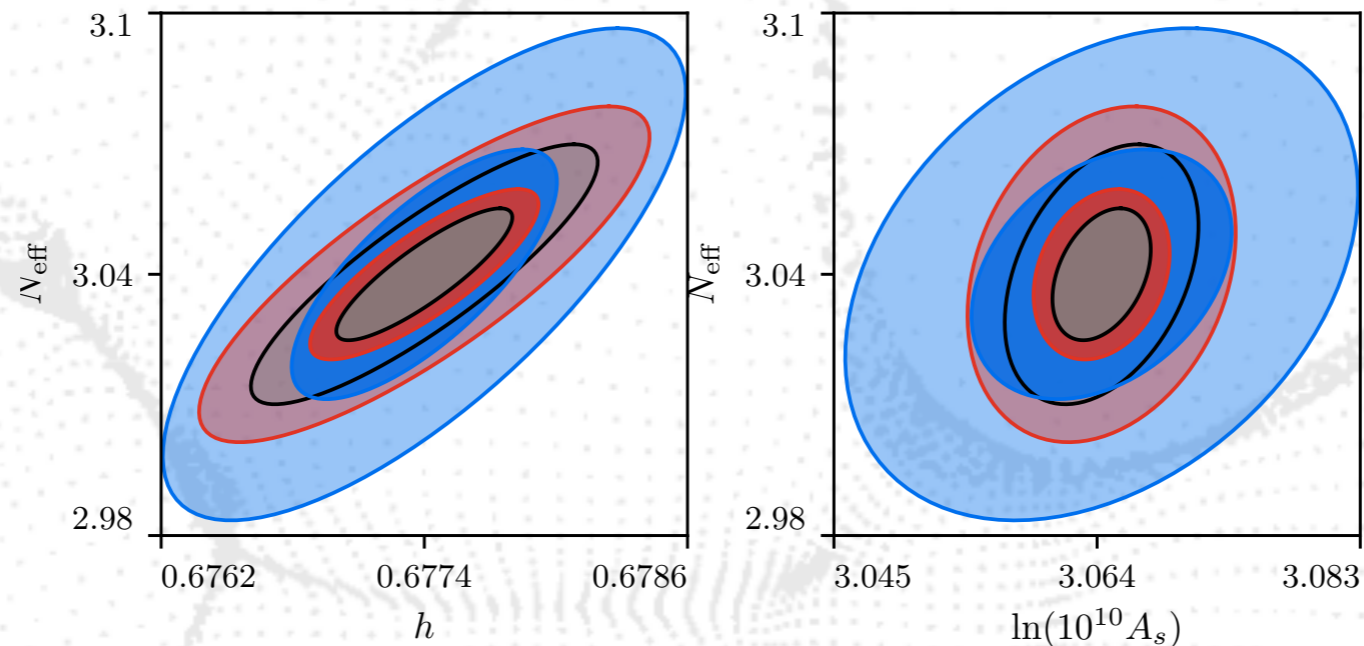
## Ext-HIRAX alone

Ext-HIRAX $N_{\text{eff}}$					
	$\Omega_M$	$h$	$\ln(10^{10} A_s)$	$n_s$	$N_{\text{eff}}$
Fiducial values	0.3075	0.6774	3.064	0.9667	3.04
No wedge $k_{\text{max}} = 0.2h\text{Mpc}^{-1}$					
2% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.0017	0.0036	0.013	0.0061	0.22
5% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.0019	0.0037	0.023	0.0061	0.23
10% $b_{\text{HI}}$ & $\Omega_{\text{HI}}$	0.0020	0.0037	0.041	0.0061	0.23

# Broadband shape + Neff

## CMB S4 + Euclid + Ext-HIRAX

Ext-HIRAX $N_{\text{eff}}$						
	$k_{\text{max}} = 0.2h\text{Mpc}^{-1}$			$k_{\text{max}} = k_{\text{nl}}(z)$		
Euclid Rec.+PlanckBAO	0.067			0.064		
+21cm	No wedge	Mid wedge	Wedge	No wedge	Mid wedge	Wedge
2% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.046	0.047	0.057	0.038	0.039	0.049
5% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.047	0.048	0.057	0.039	0.040	0.049
10% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.047	0.048	0.057	0.039	0.041	0.050
Euclid Rec.+CMB-S4	0.022			0.020		
+21cm	No wedge	Mid wedge	Wedge	No wedge	Mid wedge	Wedge
2% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.019	0.019	0.021	0.015	0.015	0.017
5% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.020	0.020	0.022	0.015	0.016	0.018
10% $b_{\text{HI}} \& \Omega_{\text{HI}}$	0.020	0.020	0.022	0.015	0.016	0.018





# Main conclusions

- We investigate the possibility of performing cosmological studies in the redshift range  $2.5 < z < 5$  through suitable extensions of existing and upcoming radio-telescopes like CHIME, HIRAX and FAST.
- We use Fisher formalism to forecast tight constraints on growth parameter  $f\sigma_8$  (4%) and AP parameters (1%) as a function of redshift in narrow redshift bins  $dz=0.1$ .
- In combination with data from Euclid-like galaxy surveys and CMB S4, the sum of the neutrino masses can be constrained with an error equal to 23 meV ( $1\sigma$ ), while  $N_{\text{eff}}$  can be constrained within 0.02 ( $1\sigma$ ).
- We study in detail the dependence of our results on the instrument, amplitude of the HI bias, the foreground wedge coverage, the non-linear scale used in the analysis, uncertainties in the theoretical modeling and the priors on  $b_{\text{HI}}$  and  $\Omega_{\text{HI}}$ .



**Thank you**