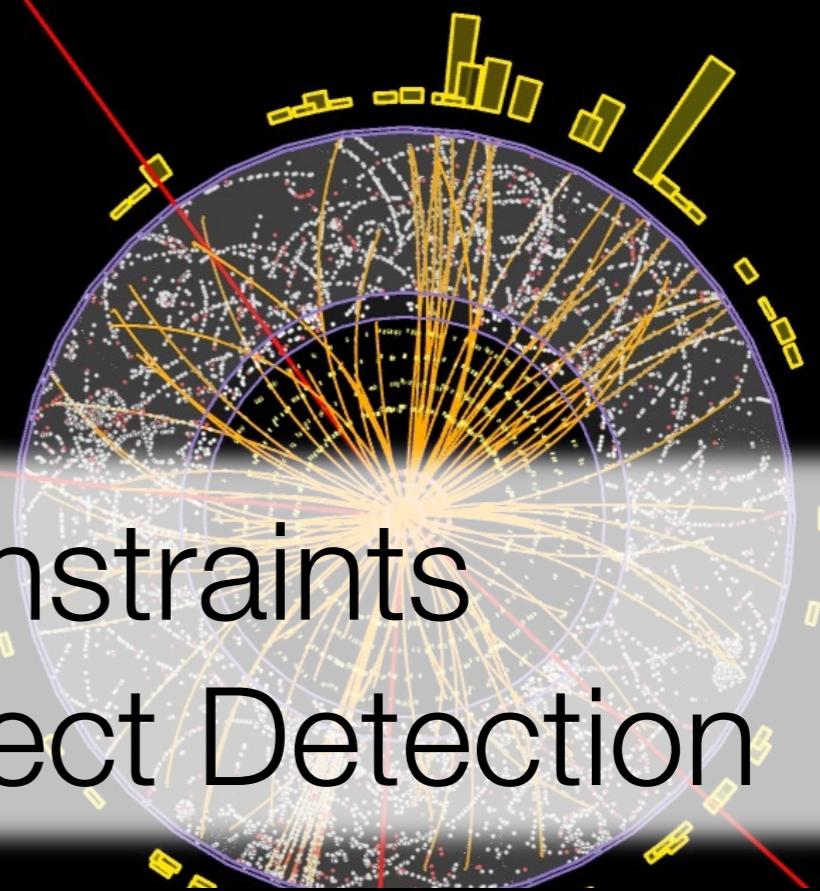


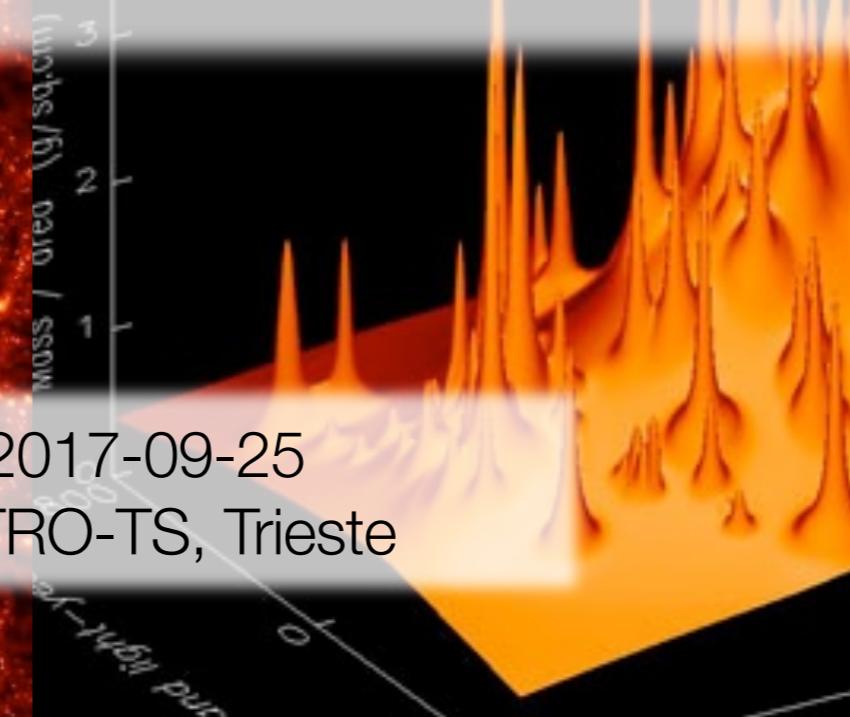
# Future CTA constraints vs Colliders and Direct Detection



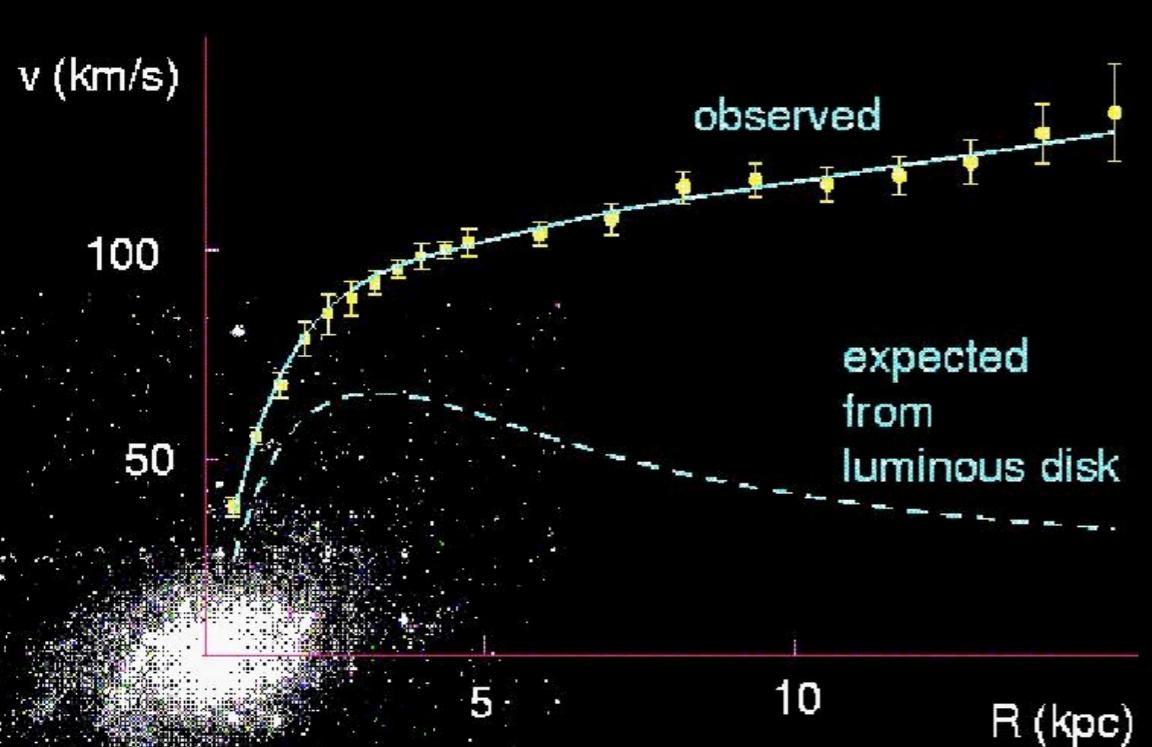
Thomas Jacques

with C. Balázs, J. Conrad, B. Farmer, T. Li, M. Meyer,  
F. S. Queiroz and M. A. Sánchez-Conde

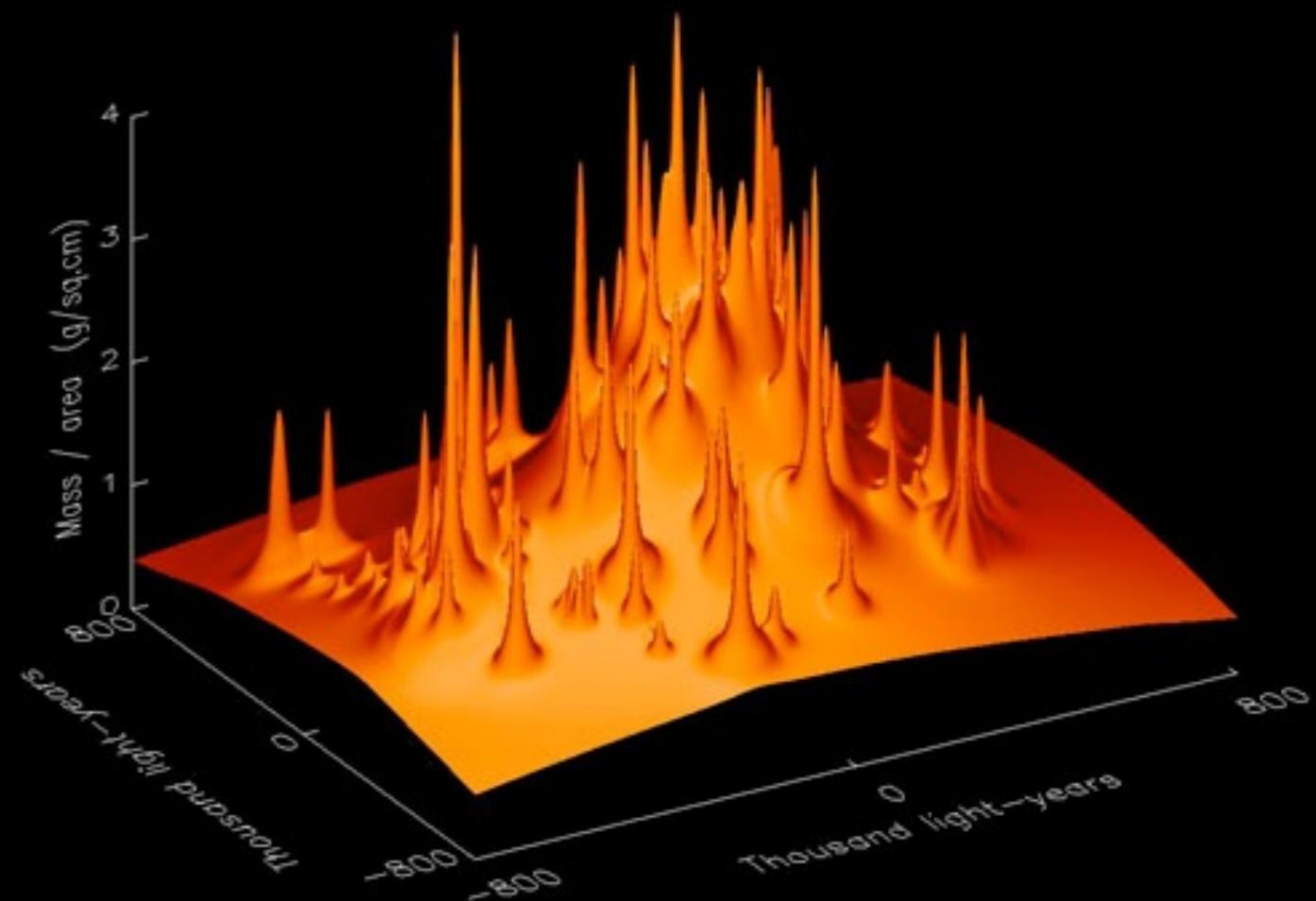
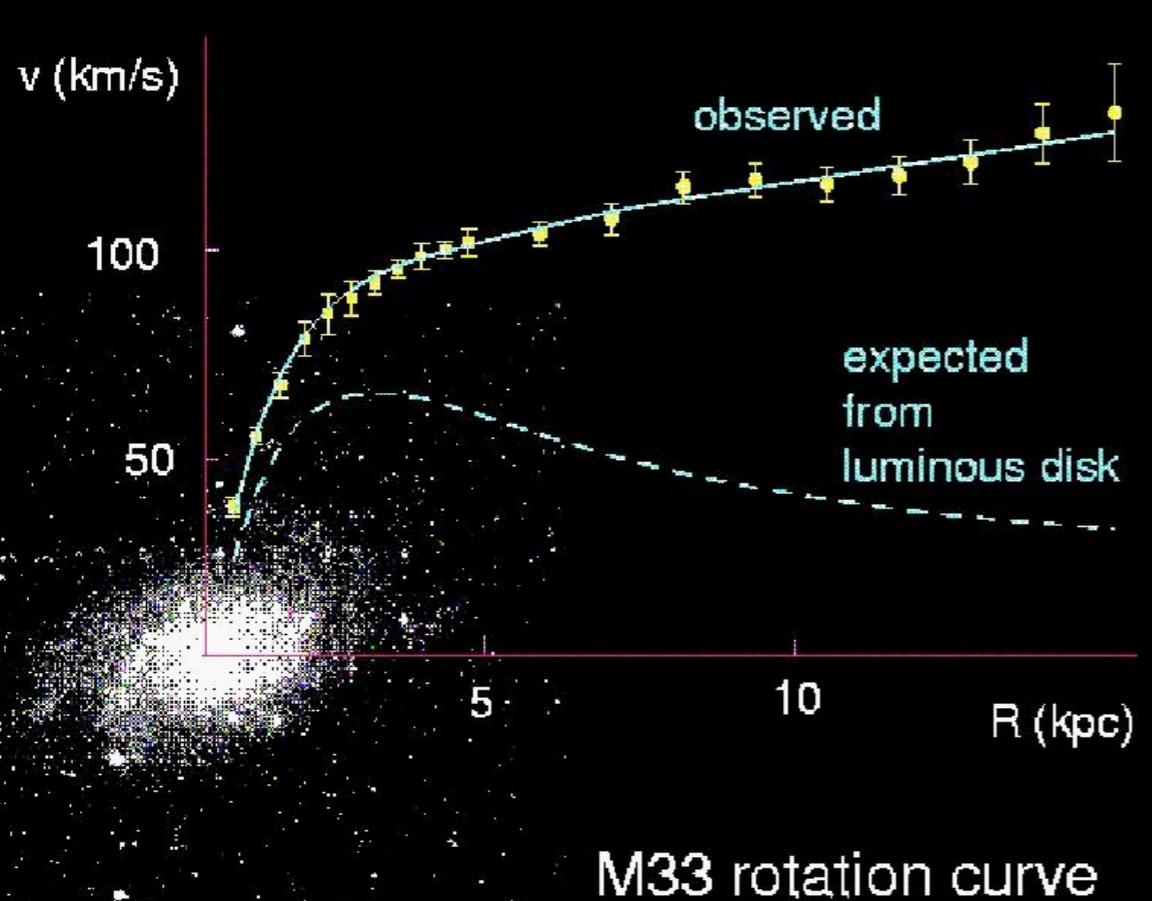
arXiv:1706.01505

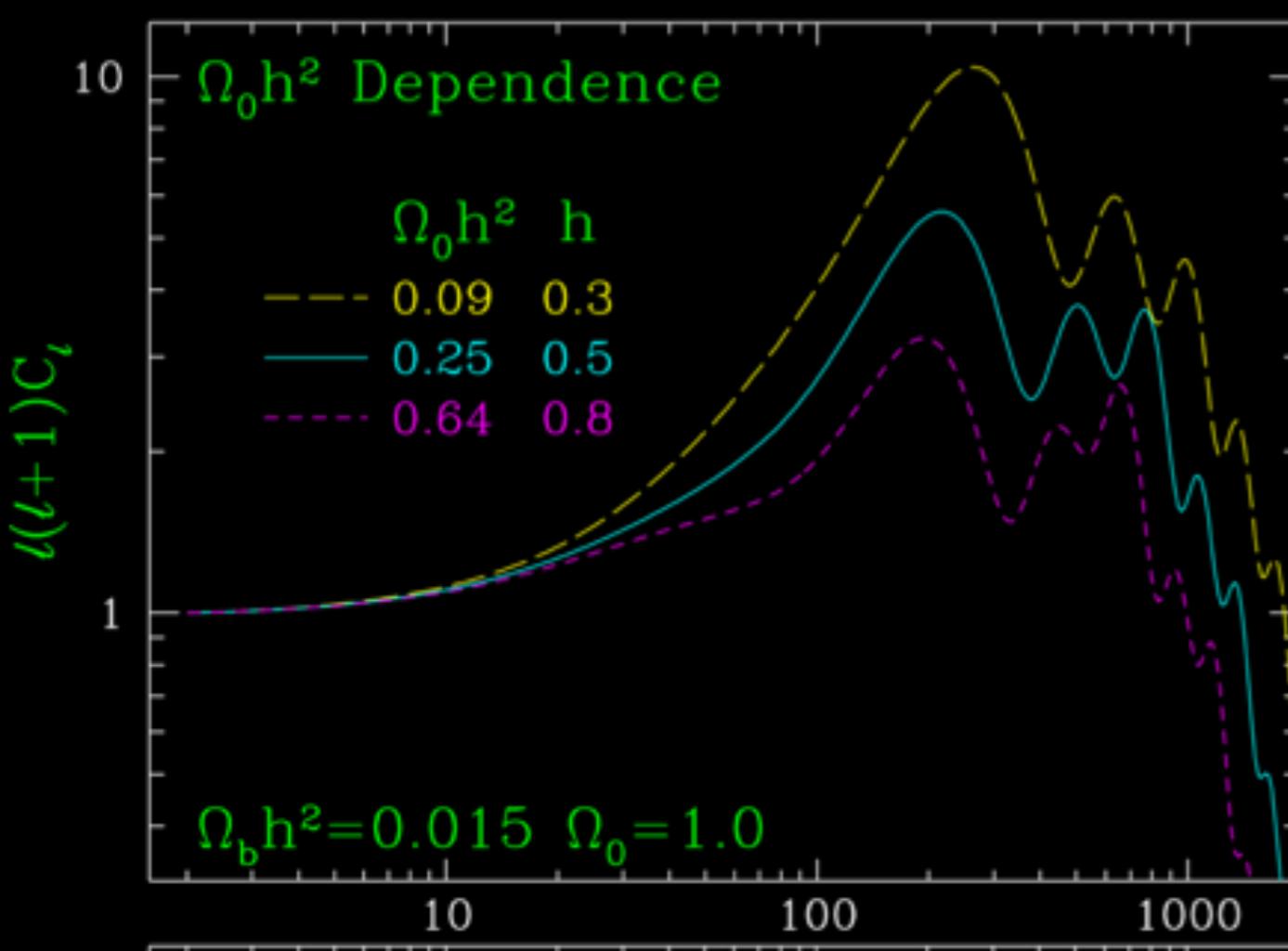
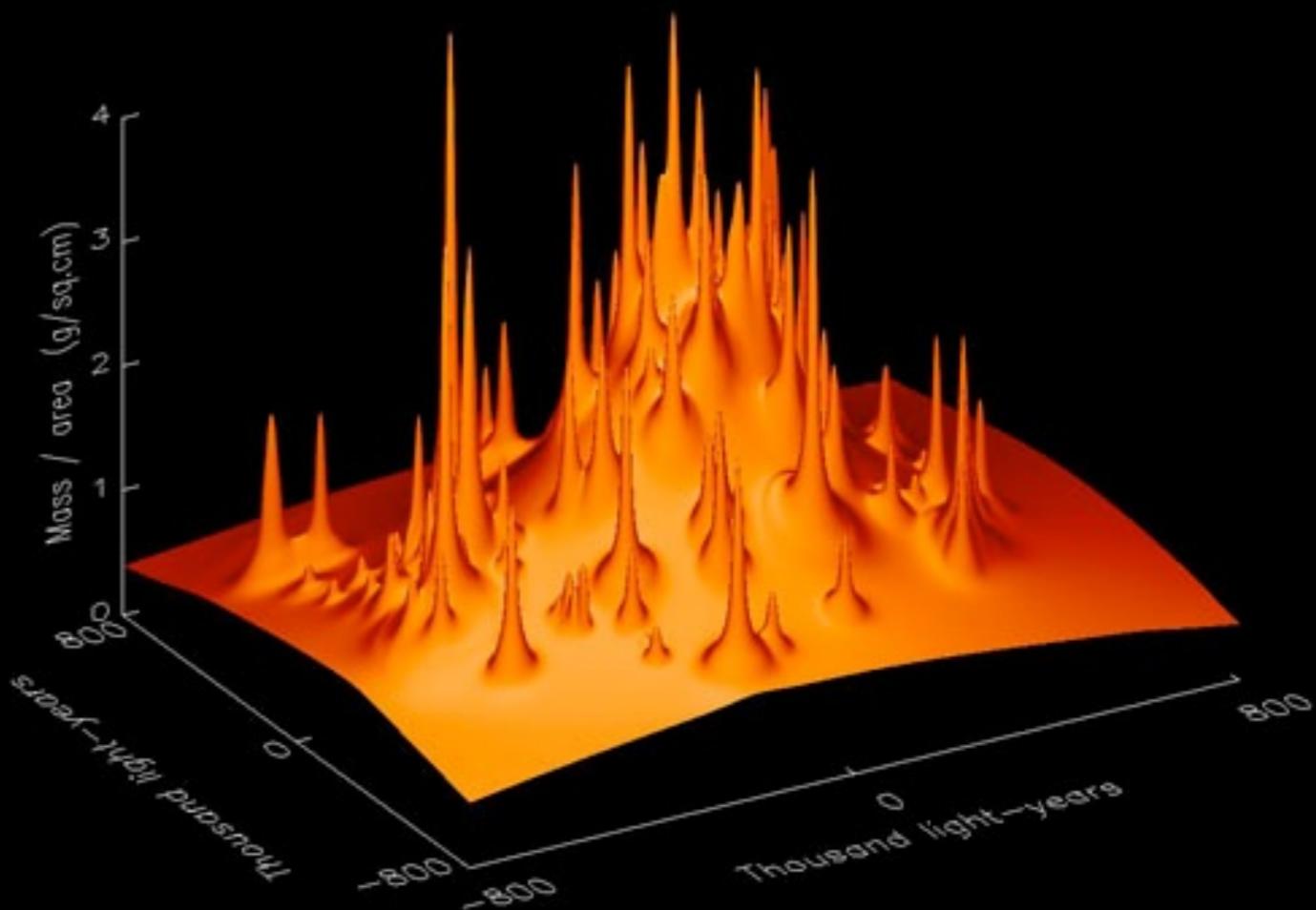
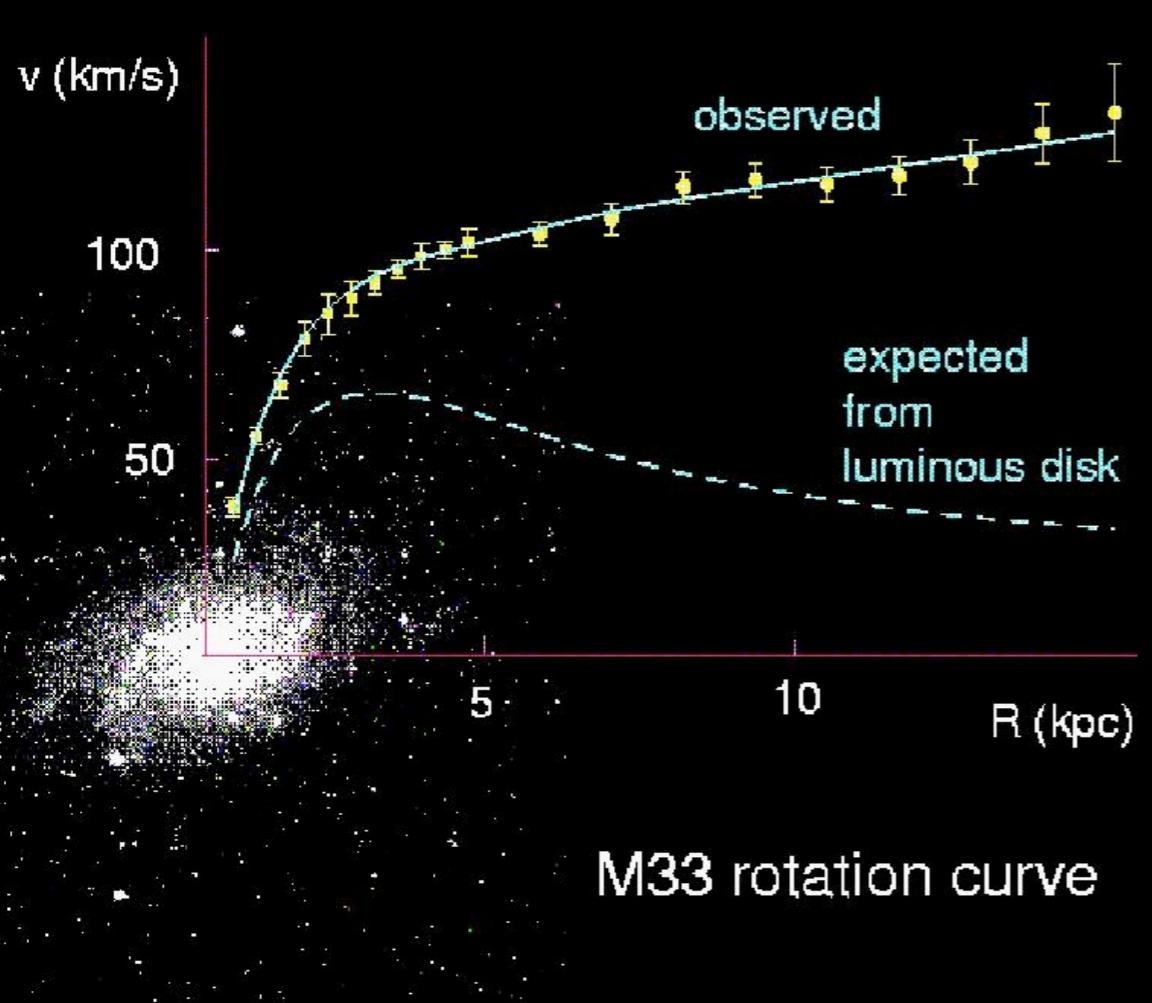


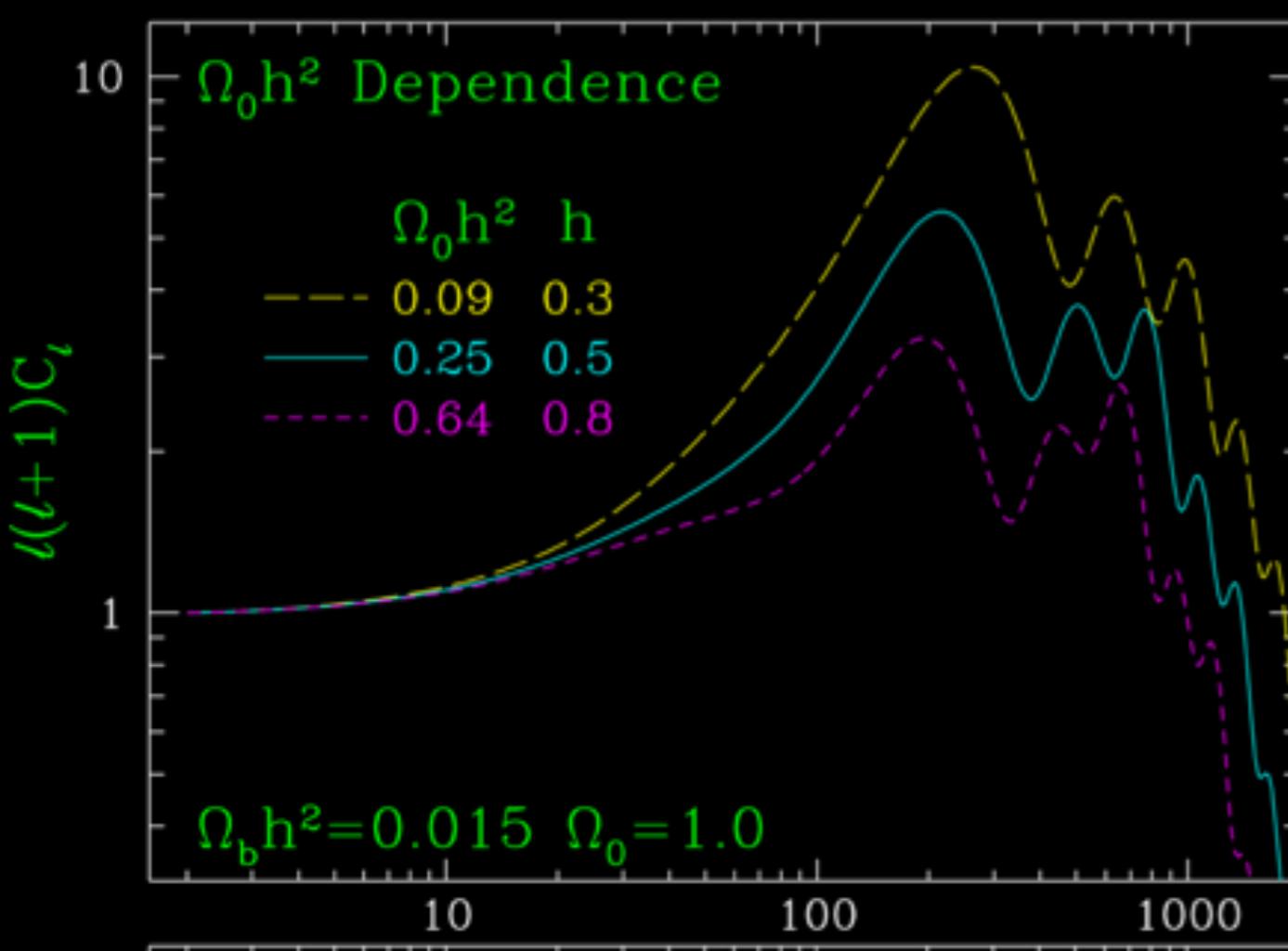
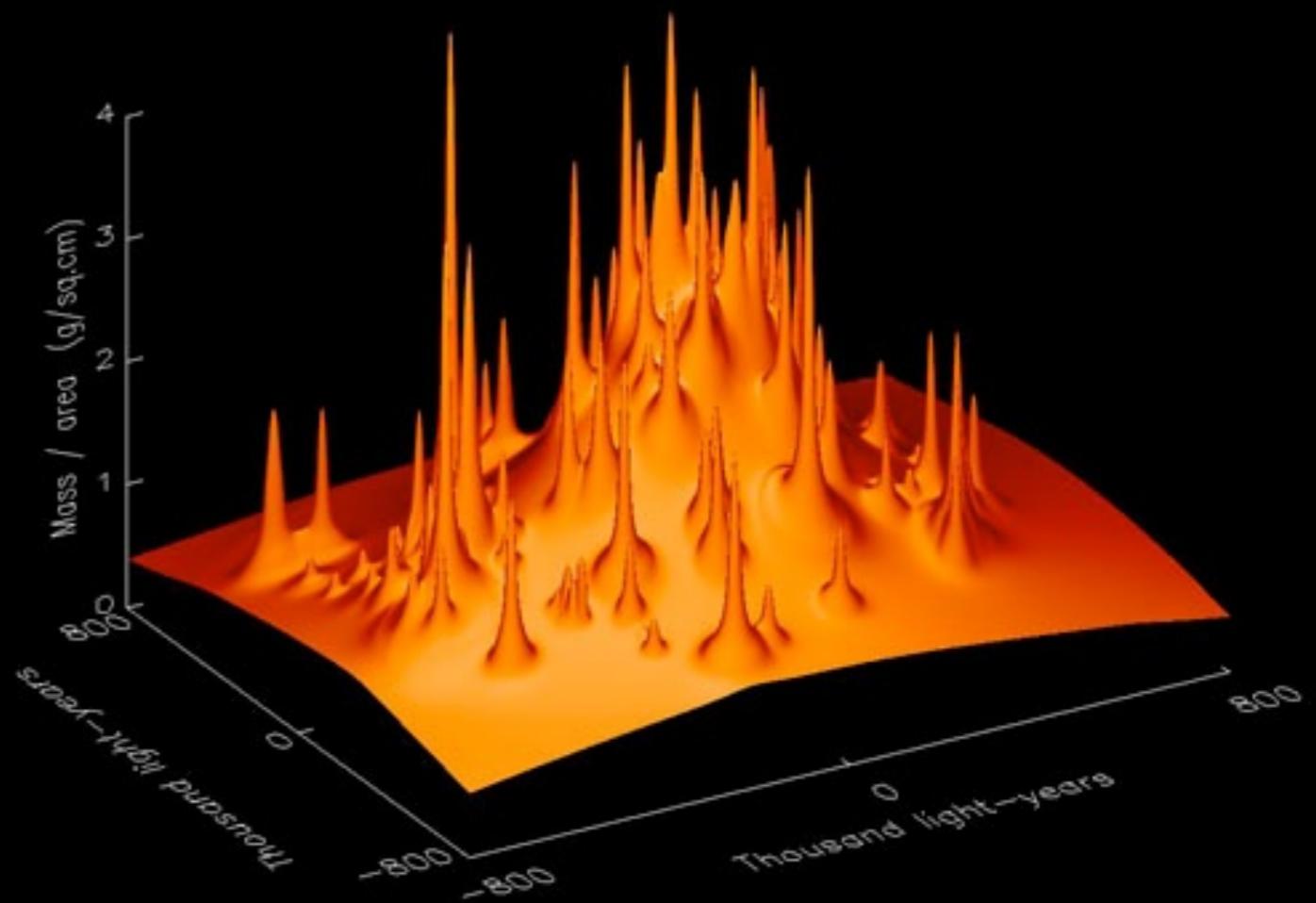
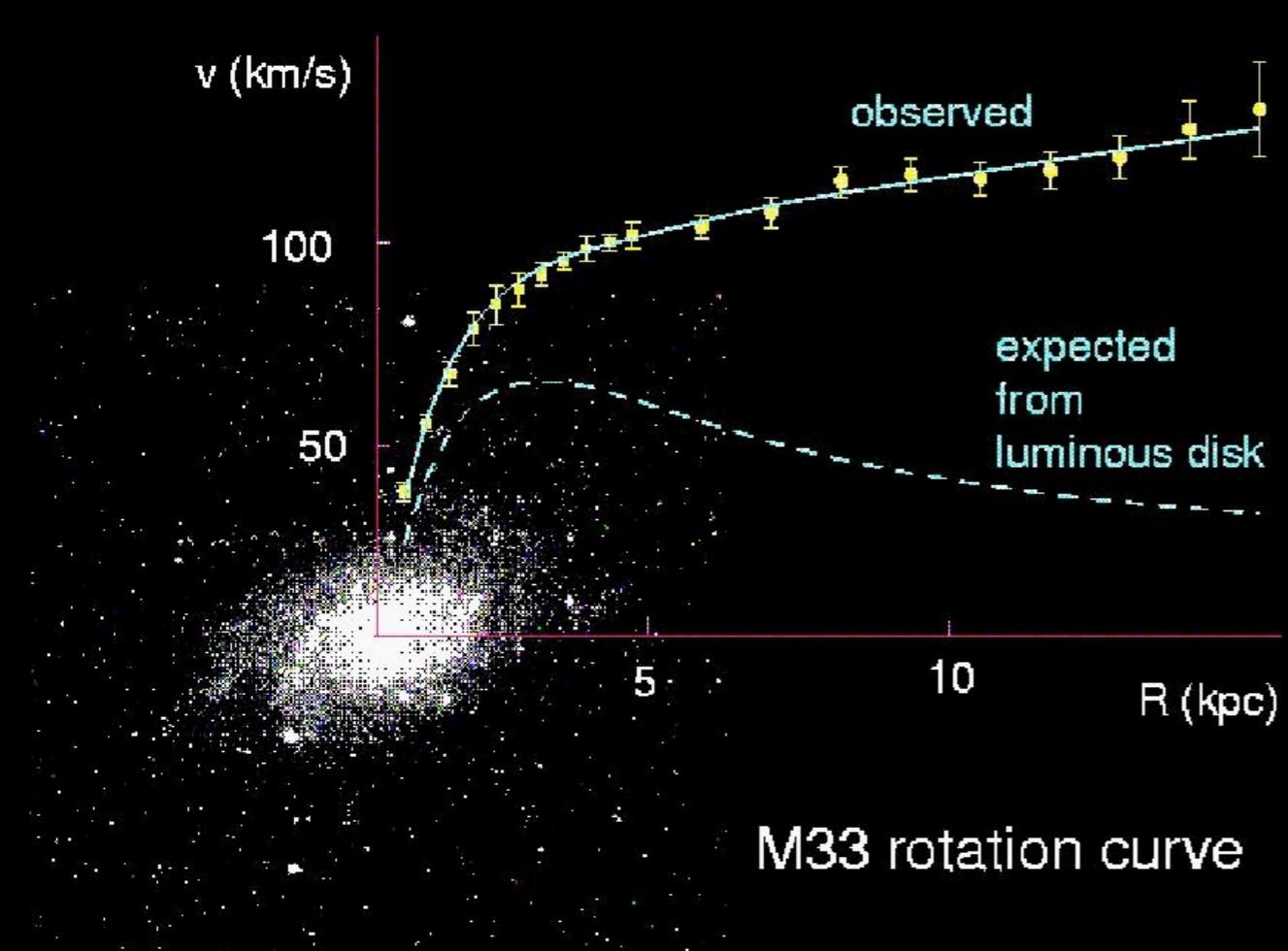
2017-09-25  
ASTRO-TS, Trieste



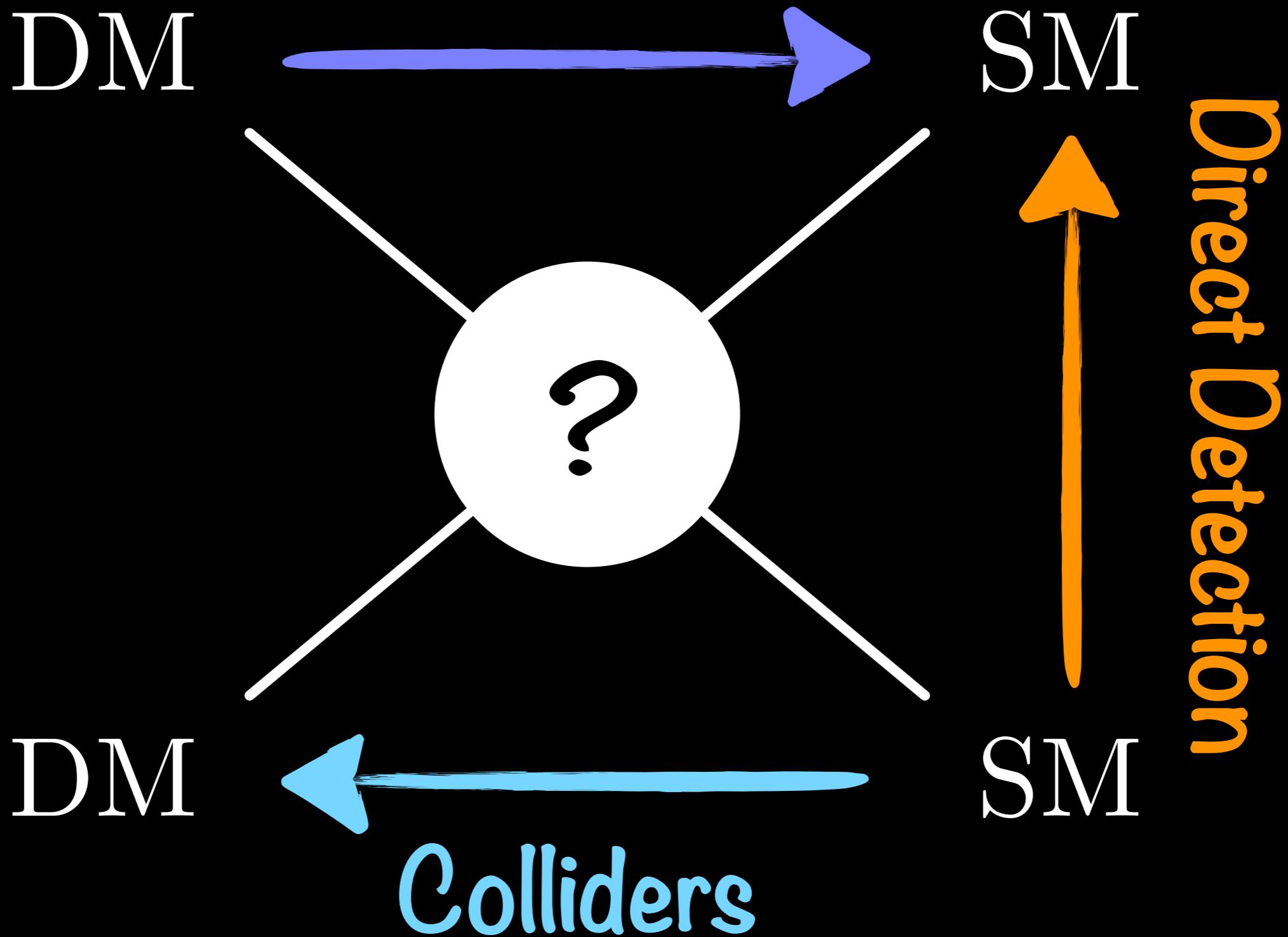
M33 rotation curve





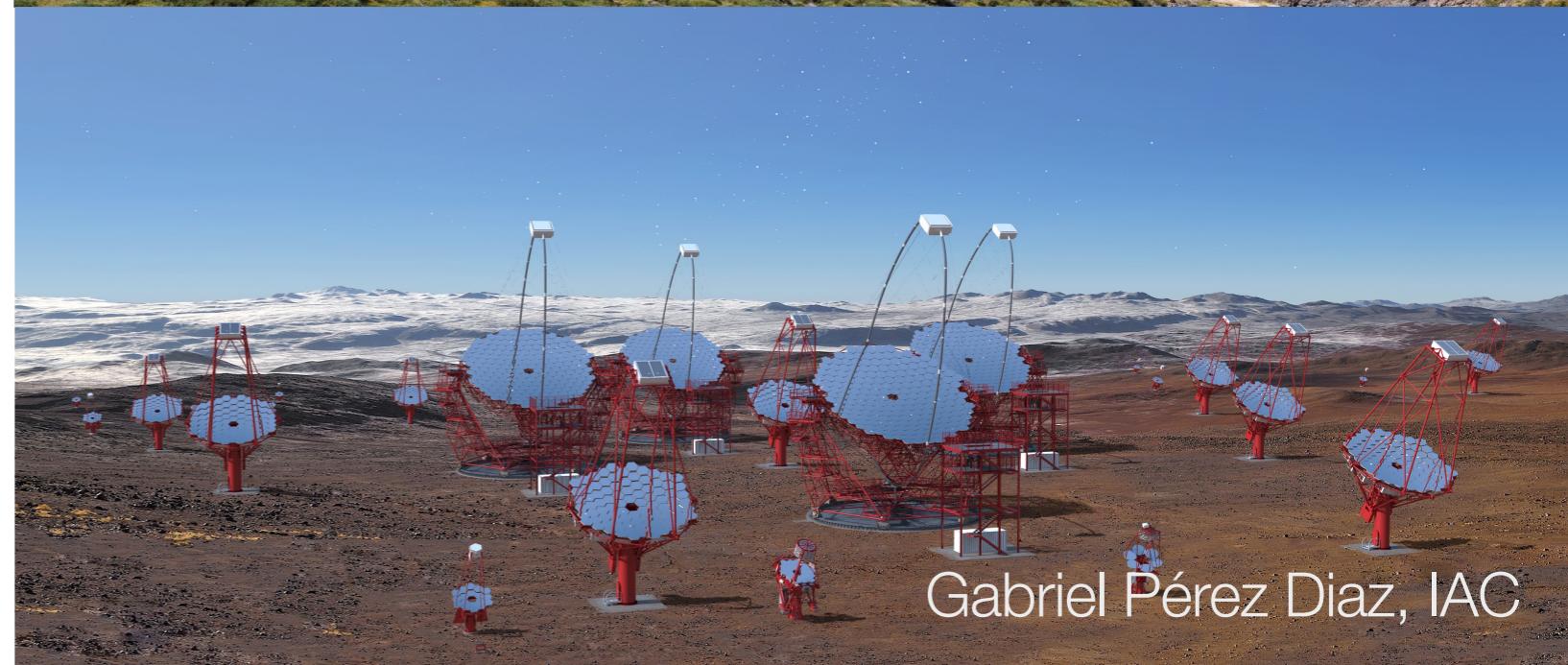


# Indirect Detection



# Cherenkov Telescope Array

- Planned imaging atmospheric cherenkov telescope array
- Will observe gamma-rays in energy range  $10 \text{ GeV}$  to  $\sim 300 \text{ TeV}$
- Order of magnitude more sensitive than existing observatories
- Two planned sites: Atacama desert & Canary Islands
- CTA Consortium includes 1,350 members from 210 institutes in 32 countries
- Not me! - Done in collaboration with CTA members, not an official CTA paper

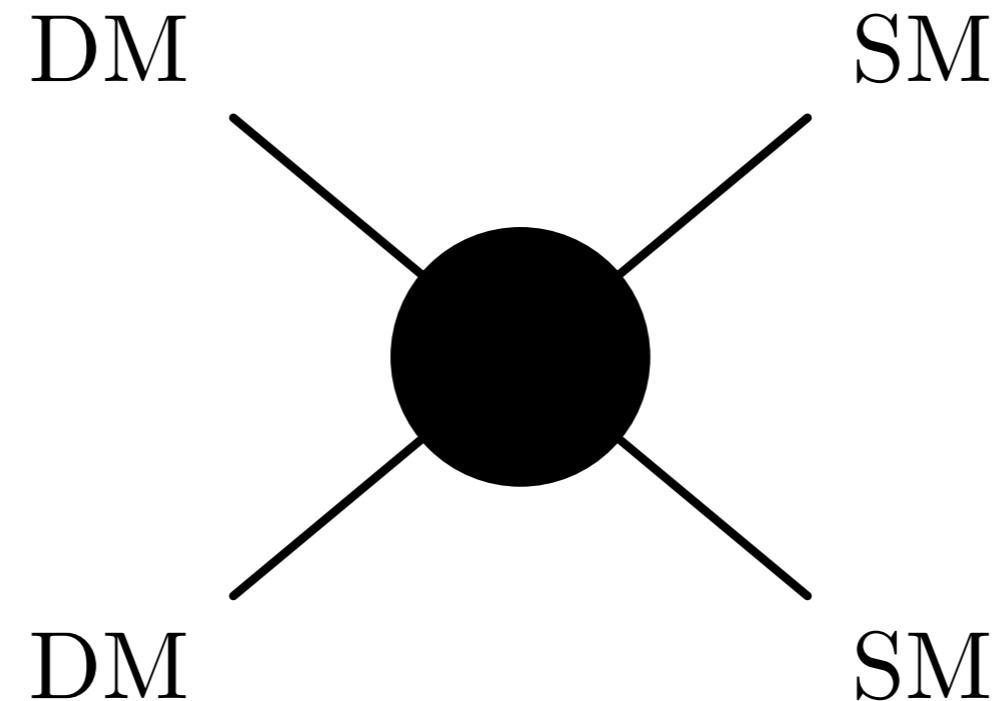


Gabriel Pérez Diaz, IAC

# Indirect Detection

Thomas Jacques

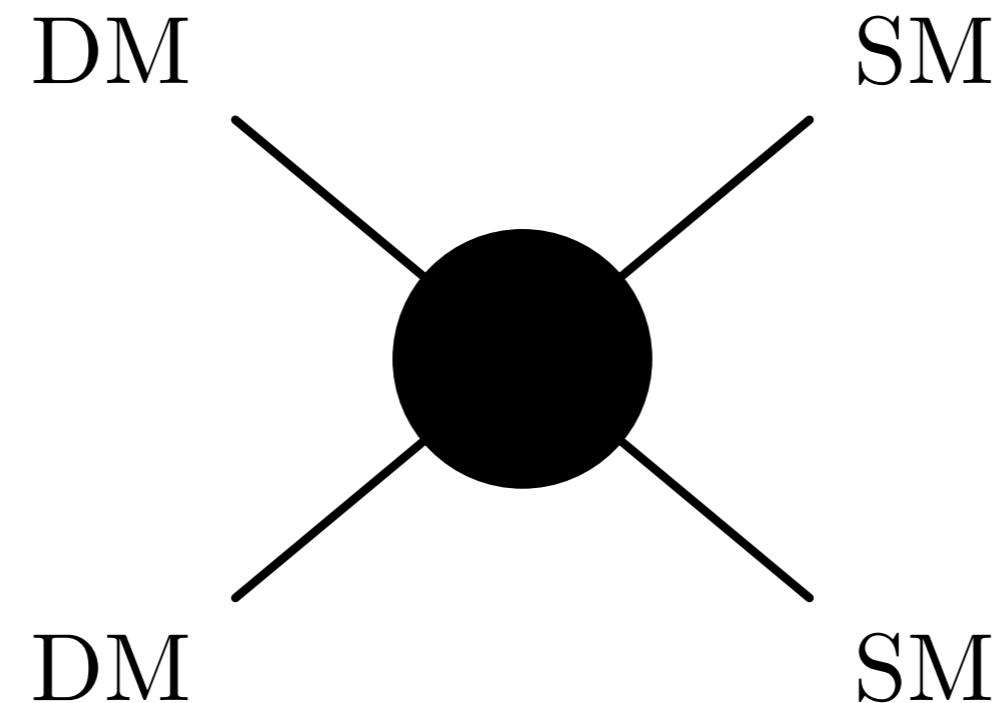
$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



# Indirect Detection

Thomas Jacques

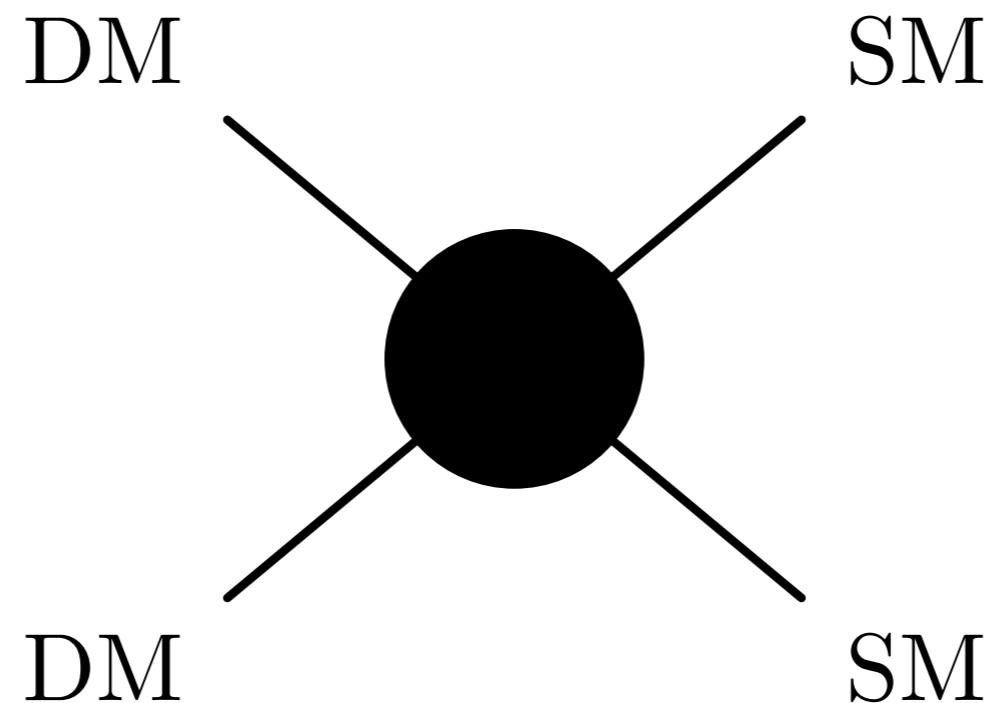
$$\frac{d\phi}{dE} = \boxed{\frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE}} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



# Indirect Detection

Thomas Jacques

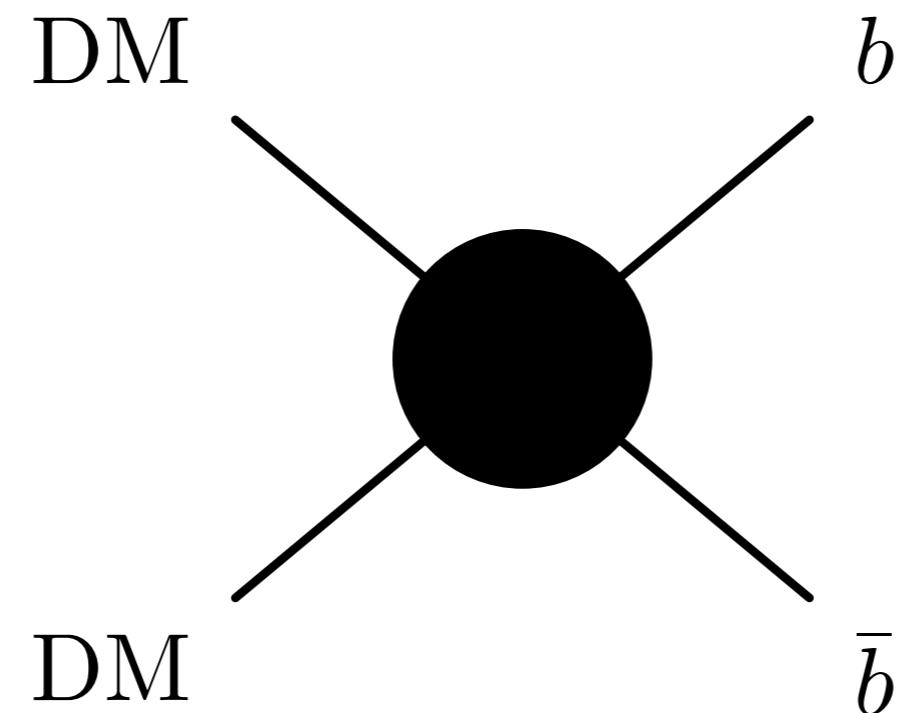
$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



# Indirect Detection

Thomas Jacques

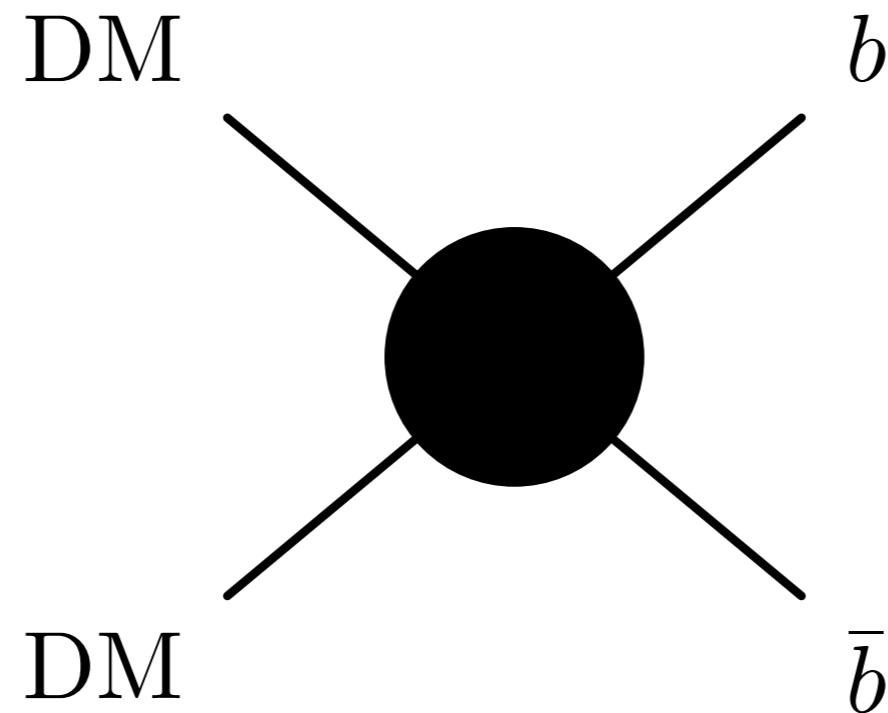
$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



# Indirect Detection

Thomas Jacques

$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$

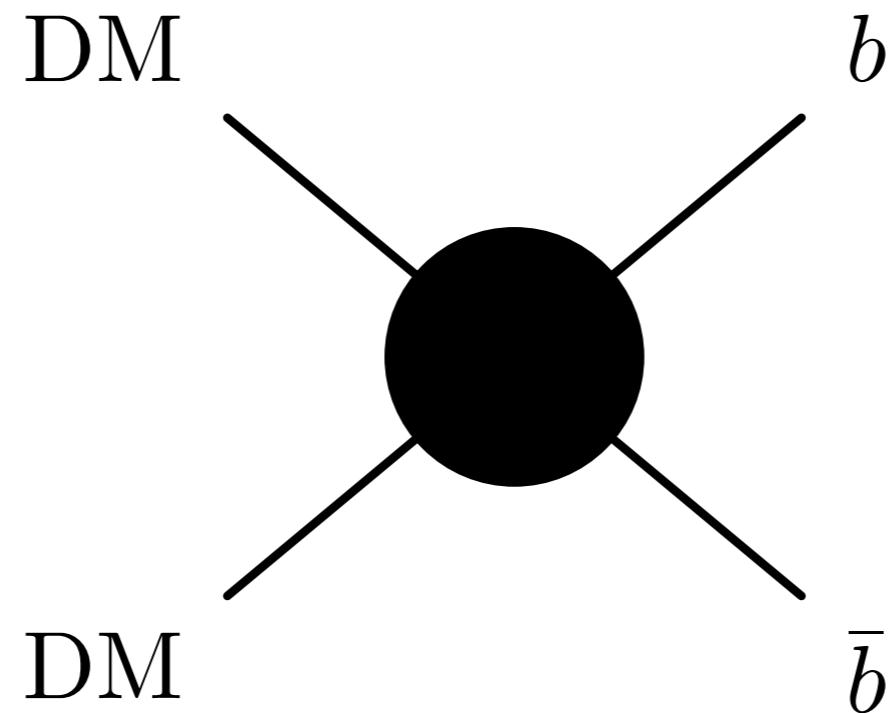


- Usually assume 100% branching to  $b$ -quarks

# Indirect Detection

Thomas Jacques

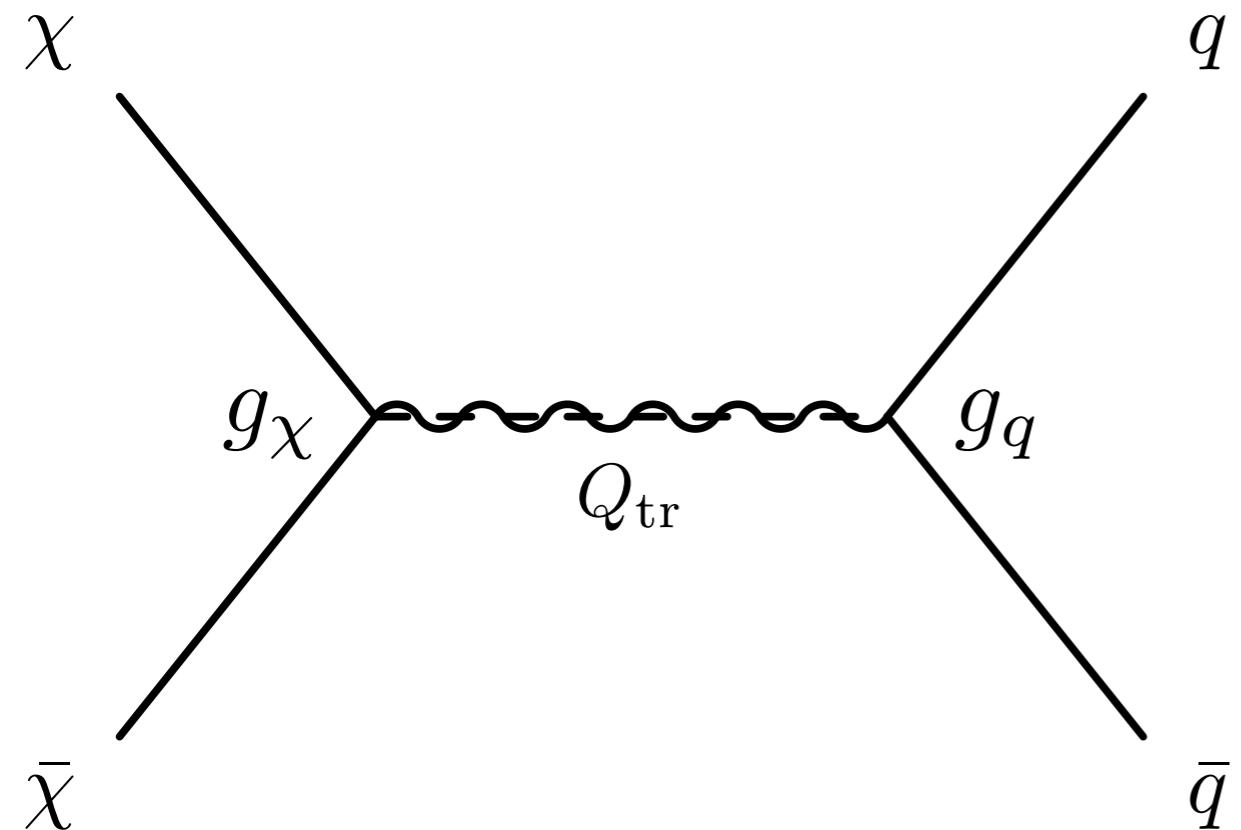
$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



- Usually assume 100% branching to  $b$ -quarks
- Need an underlying model to get the branching ratios

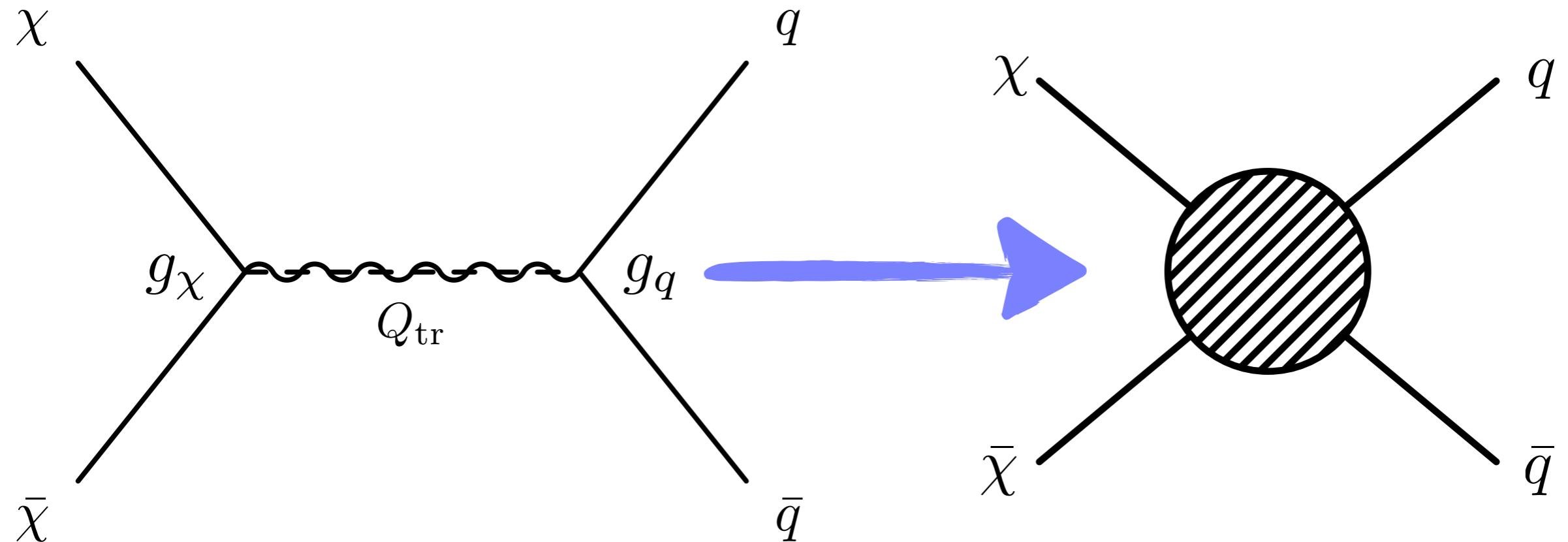
# Use Simplified Models and EFTs!

Thomas Jacques



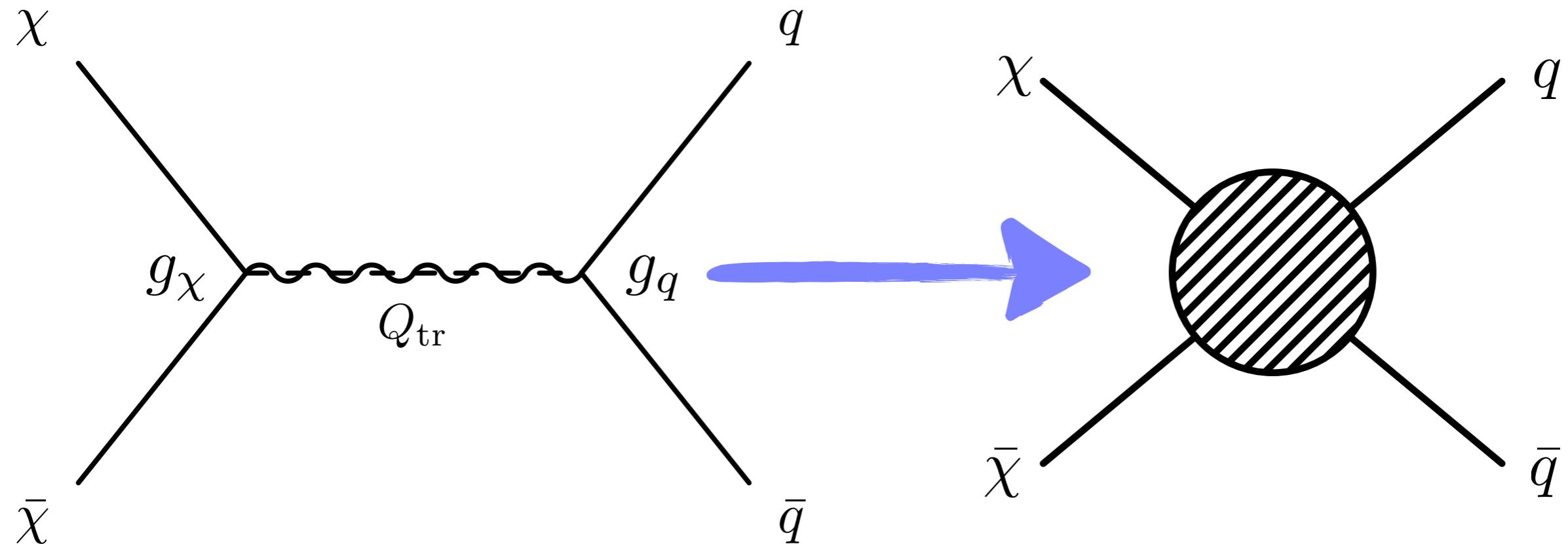
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Thomas Jacques



# Use Simplified Models and EFTs!

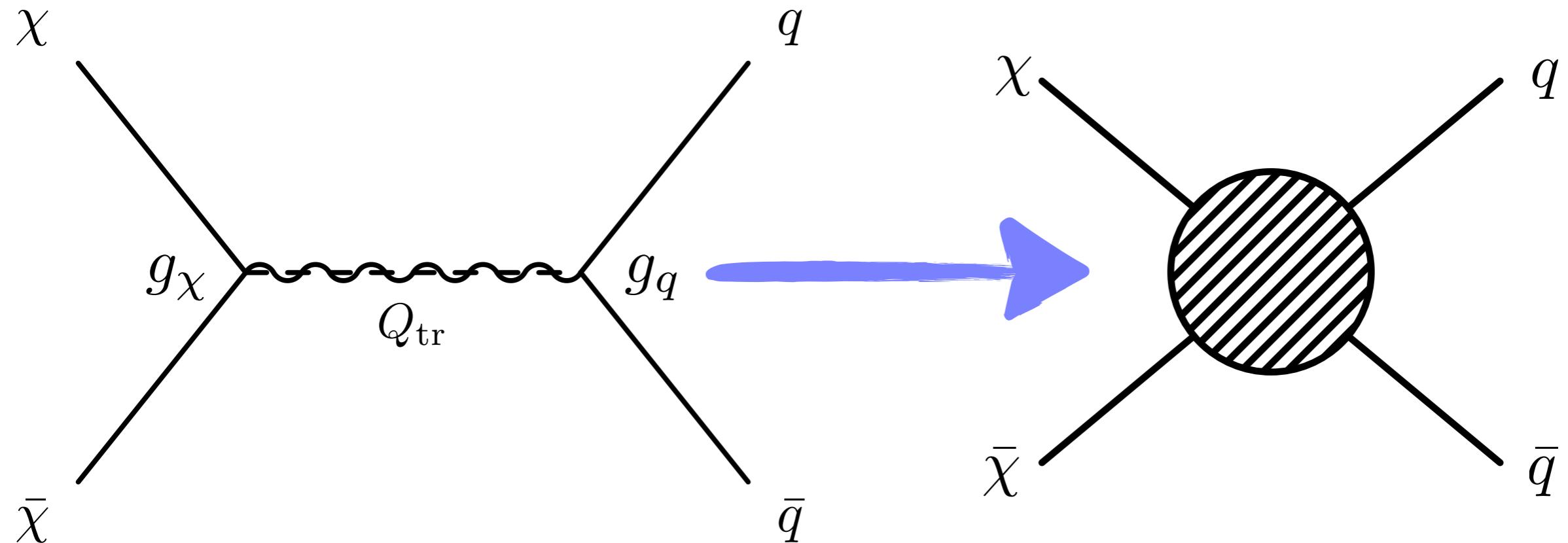
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$$\frac{g_q g_\chi}{M^2 - Q_{\text{tr}}^2} \simeq \frac{g_q g_\chi}{M^2} \equiv \frac{1}{M_\star^2}$$

# Use Simplified Models and EFTs!

Thomas Jacques

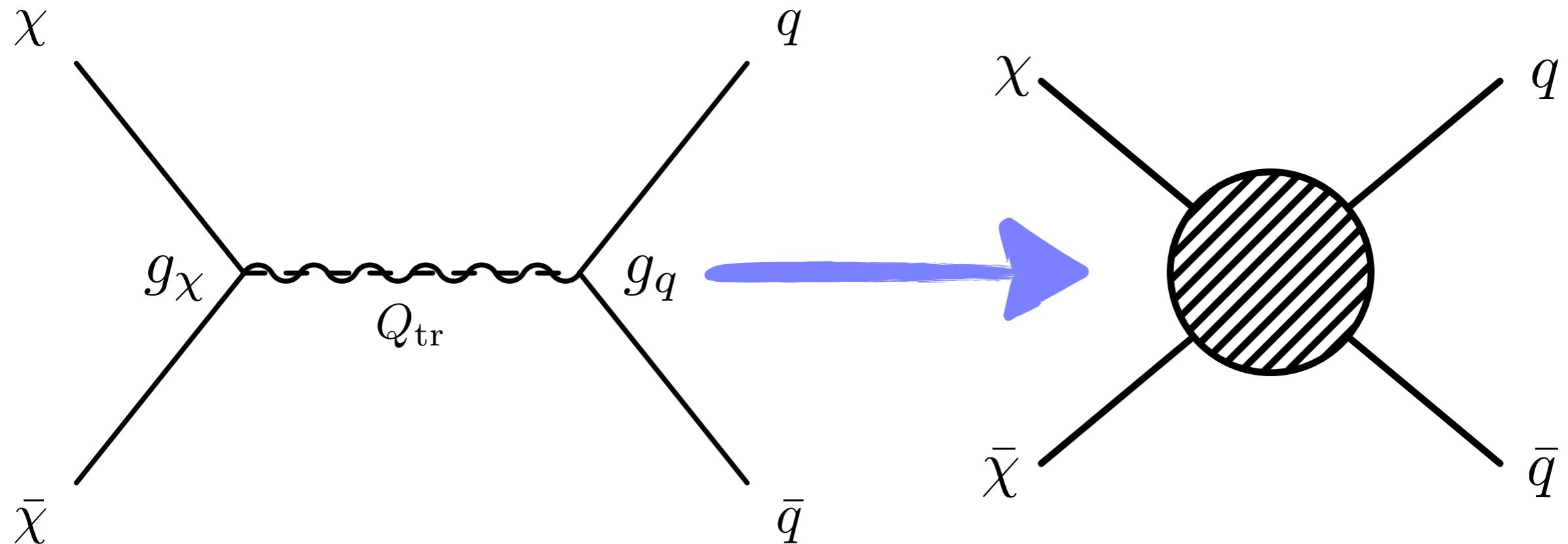


$$\frac{g_q g_\chi}{M^2 - Q_{\text{tr}}^2} \simeq \frac{g_q g_\chi}{M^2} \equiv \frac{1}{M_\star^2}$$

- Effective Field Theories have fewer parameters, just  $m_{\text{DM}}$  and  $M_\star$

# Use Simplified Models and EFTs!

Thomas Jacques



$$\frac{g_q g_\chi}{M^2 - Q_{\text{tr}}^2} \simeq \frac{g_q g_\chi}{M^2} \equiv \frac{1}{M_*^2}$$

- Effective Field Theories have fewer parameters, just  $m_{\text{DM}}$  and  $M_*$
- Simplified models have wider range of validity

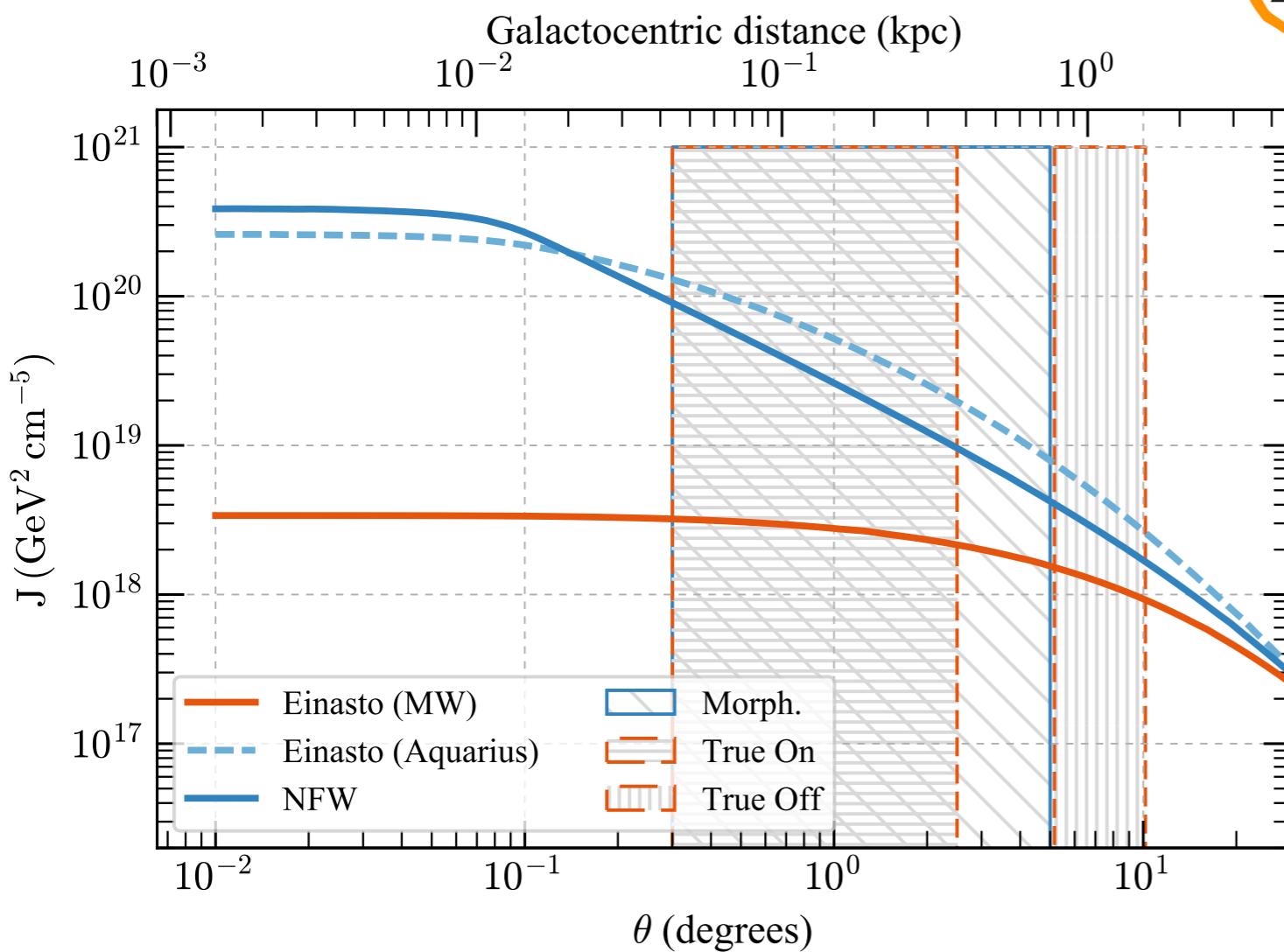
# Choice of Models

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Operator / Mediator	Operator	Annihilation Suppression	Scattering Suppression
Scalar	$\frac{m_q}{M_\star^3} (\bar{\chi}\chi)(\bar{q}q)$	$v^2$	1
Pseudo-scalar	$\frac{m_q}{M_\star^3} (\bar{\chi}\gamma^5\chi)(\bar{q}\gamma^5q)$	1	$(\vec{s}_\chi \cdot \vec{q})(\vec{s}_N \cdot \vec{q})$
Vector	$\frac{1}{M_\star^2} (\bar{\chi}\gamma^\mu\chi)(\bar{q}\gamma_\mu q)$	1	1
Axial-Vector	$\frac{1}{M_\star^2} (\bar{\chi}\gamma^\mu\gamma^5\chi)(\bar{q}\gamma_\mu\gamma^5q)$	$m_q^2, v^2$	$\vec{s}_\chi \cdot \vec{s}_N$

# Astrophysics

$$\frac{d\phi}{dE} = \frac{1}{8\pi} \frac{\langle \sigma v \rangle}{2m_\chi^2} \sum_f B_f \frac{dN_f}{dE} \int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_\chi^2(r) dl d\Omega$$



$$\rho_{\text{NFW}}(r) = \frac{\rho_0}{\frac{r}{r_s} \left[ 1 + \frac{r}{r_s} \right]^2}$$

$$\rho_{\text{Ein}}(r) = \rho_0 \exp \left( -\frac{2}{\alpha} \left[ \left( \frac{r}{r_s} \right)^\alpha - 1 \right] \right)$$

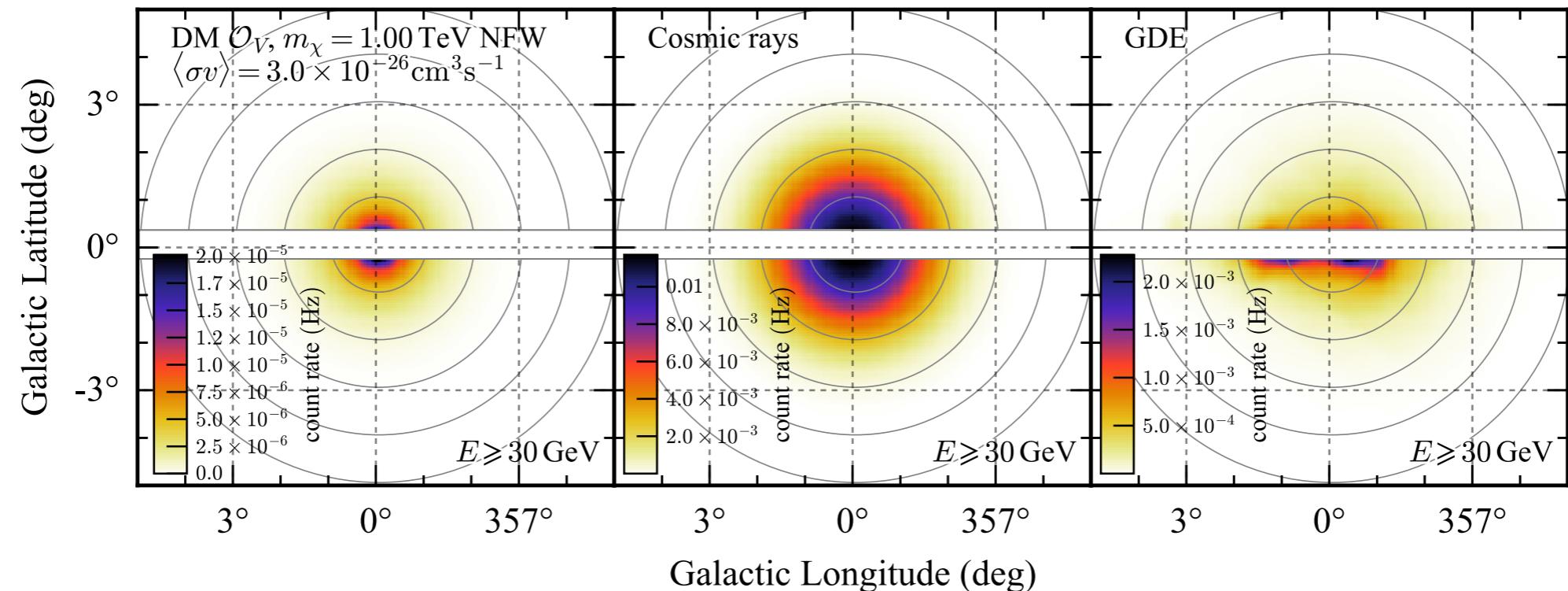
$$\rho(r_{\text{sun}}) = 0.42 \text{ GeV/cm}^3$$

$$r_s = 20 \text{ kpc} \quad \alpha = 0.36$$

- Cuspyness of inner core is uncertain:  
Use two density profiles, Einasto and NFW
- Large uncertainties in fit parameters, use recent best fit for MW

# Analysis

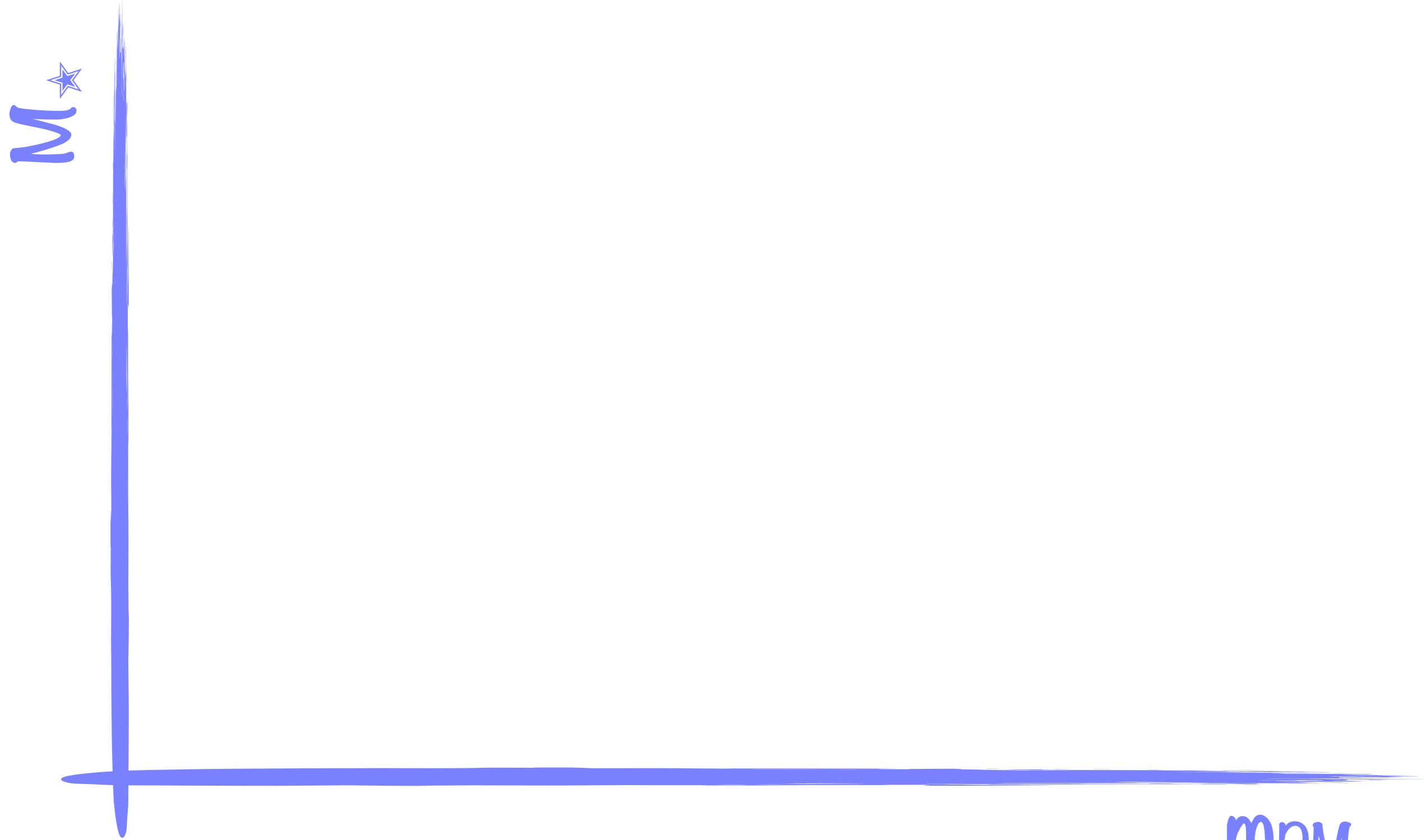
- Galactic Diffuse background taken from Fermi-LAT Template
- Cosmic Ray Background taken from CTA Monte-Carlo simulations



- Limit on  $\langle \sigma v \rangle$  converted to limit on model parameters (masses and couplings) using known analytic relationships

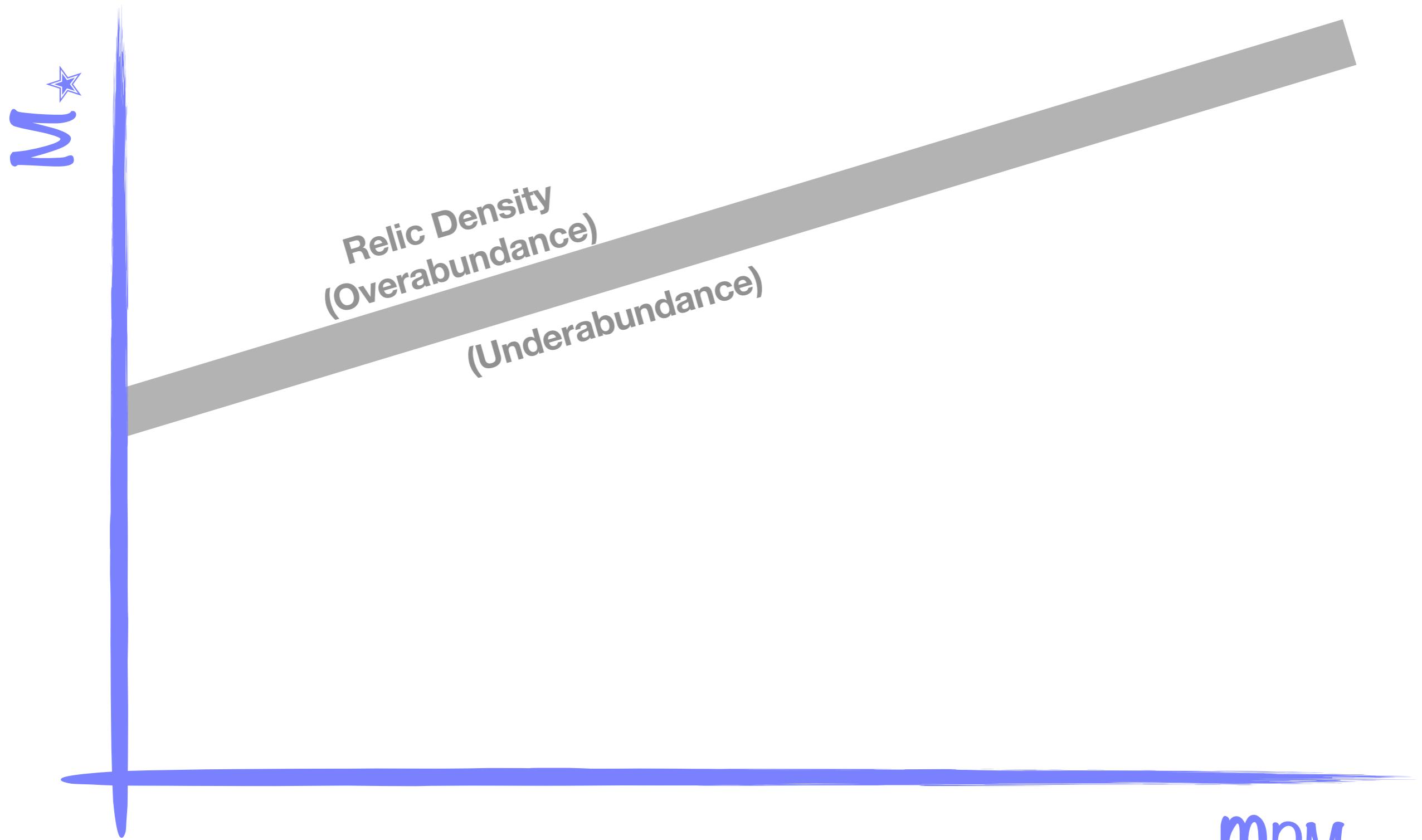
# Limits

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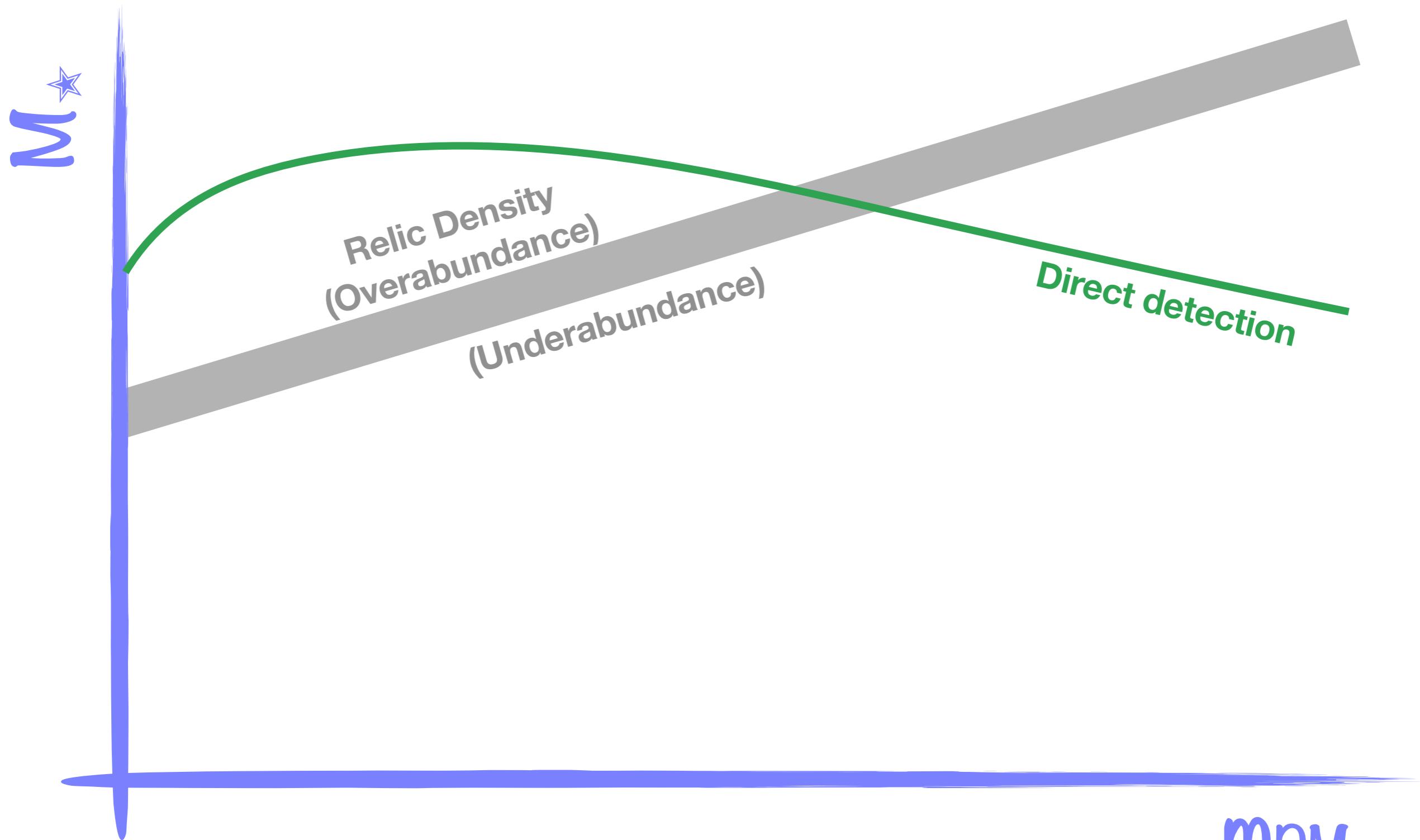
# Limits

Thomas Jacques

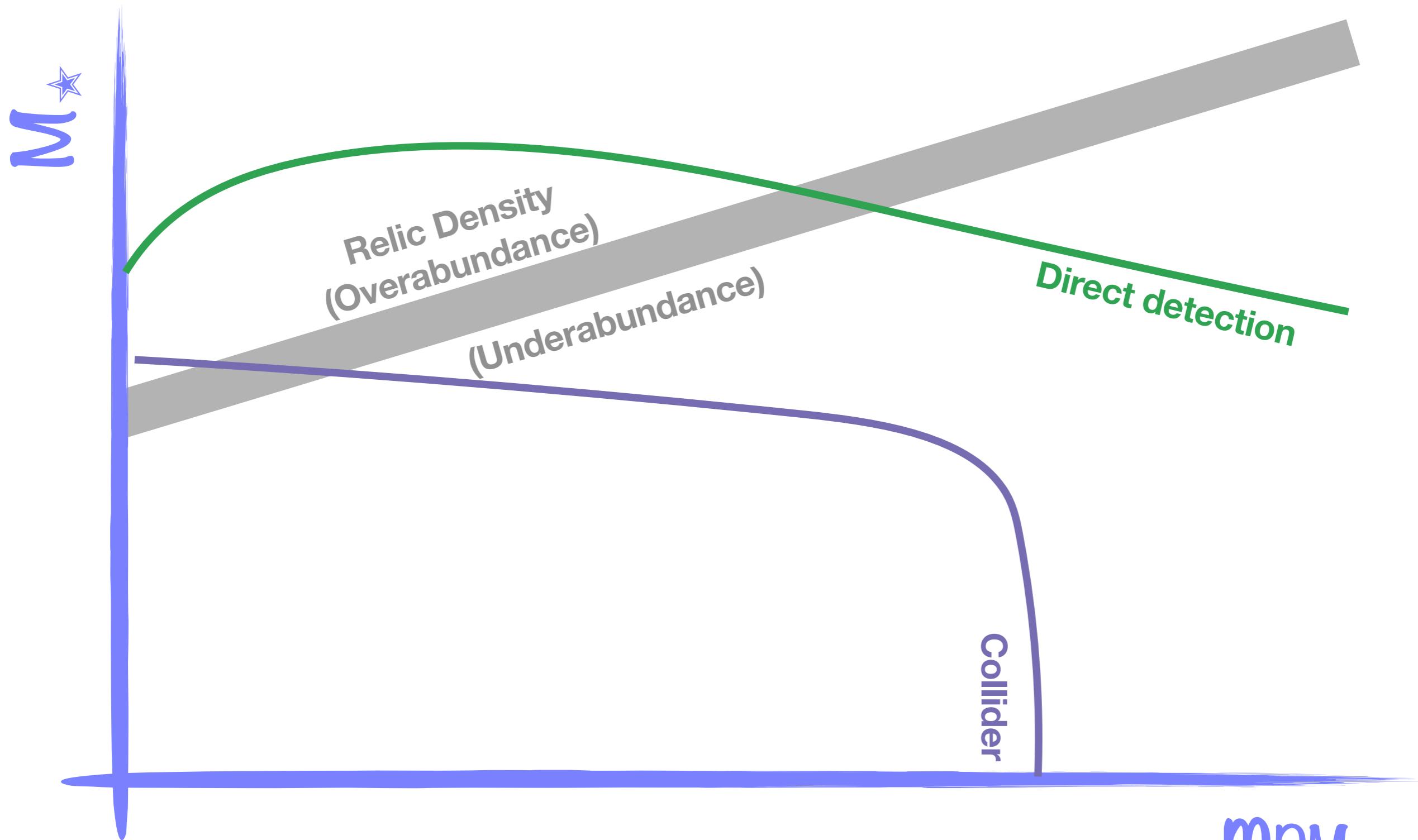


# Limits

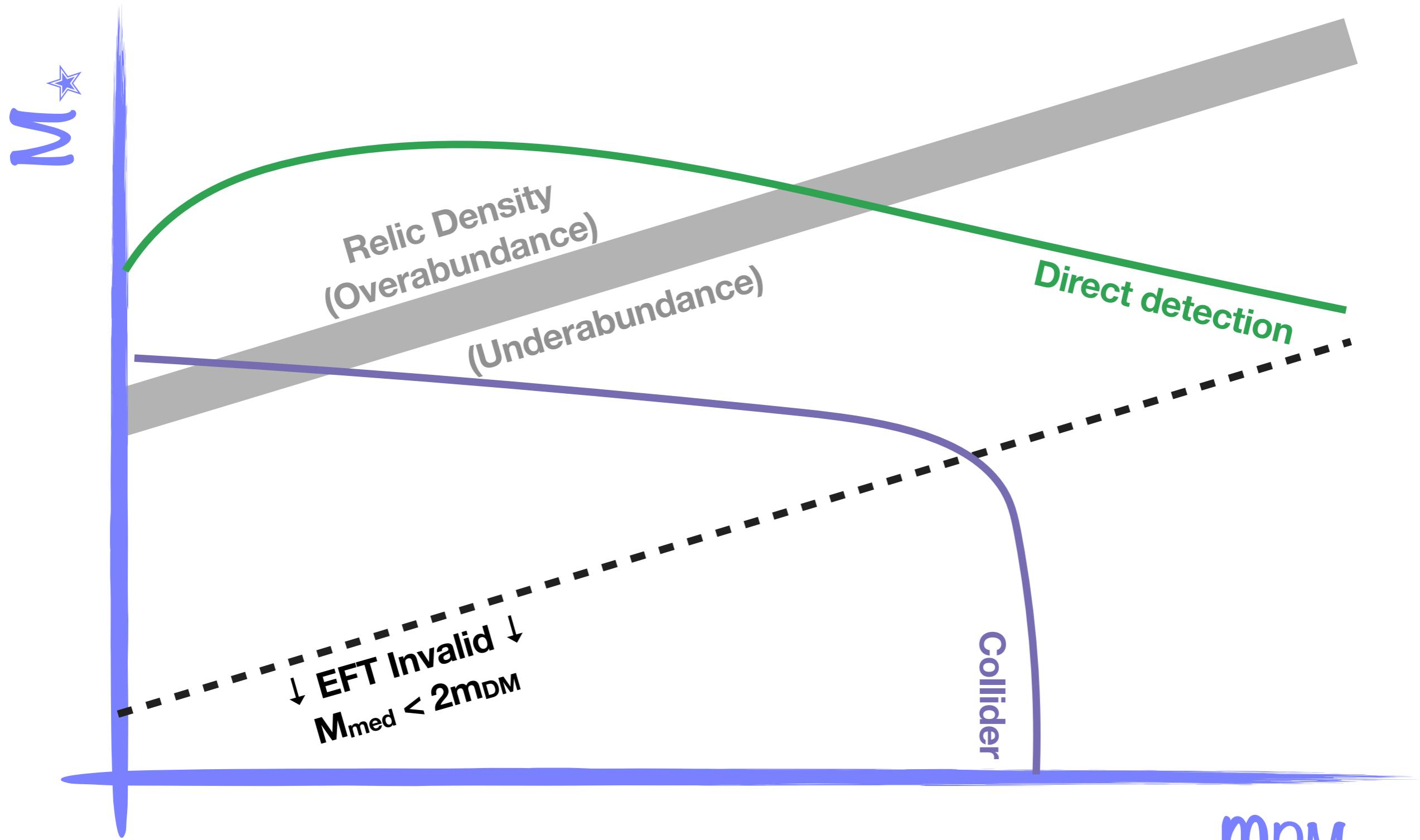
Thomas Jacques



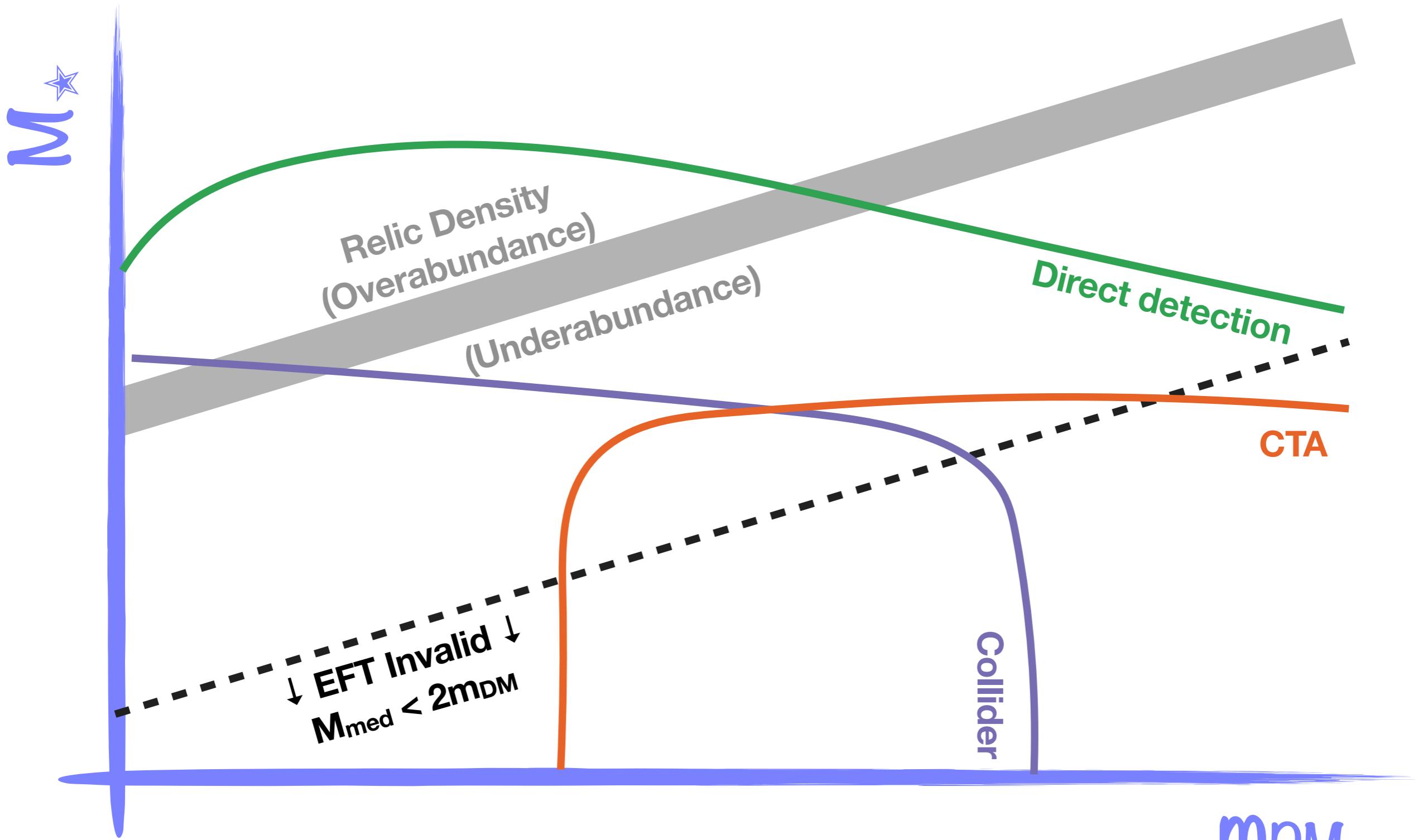
# Limits

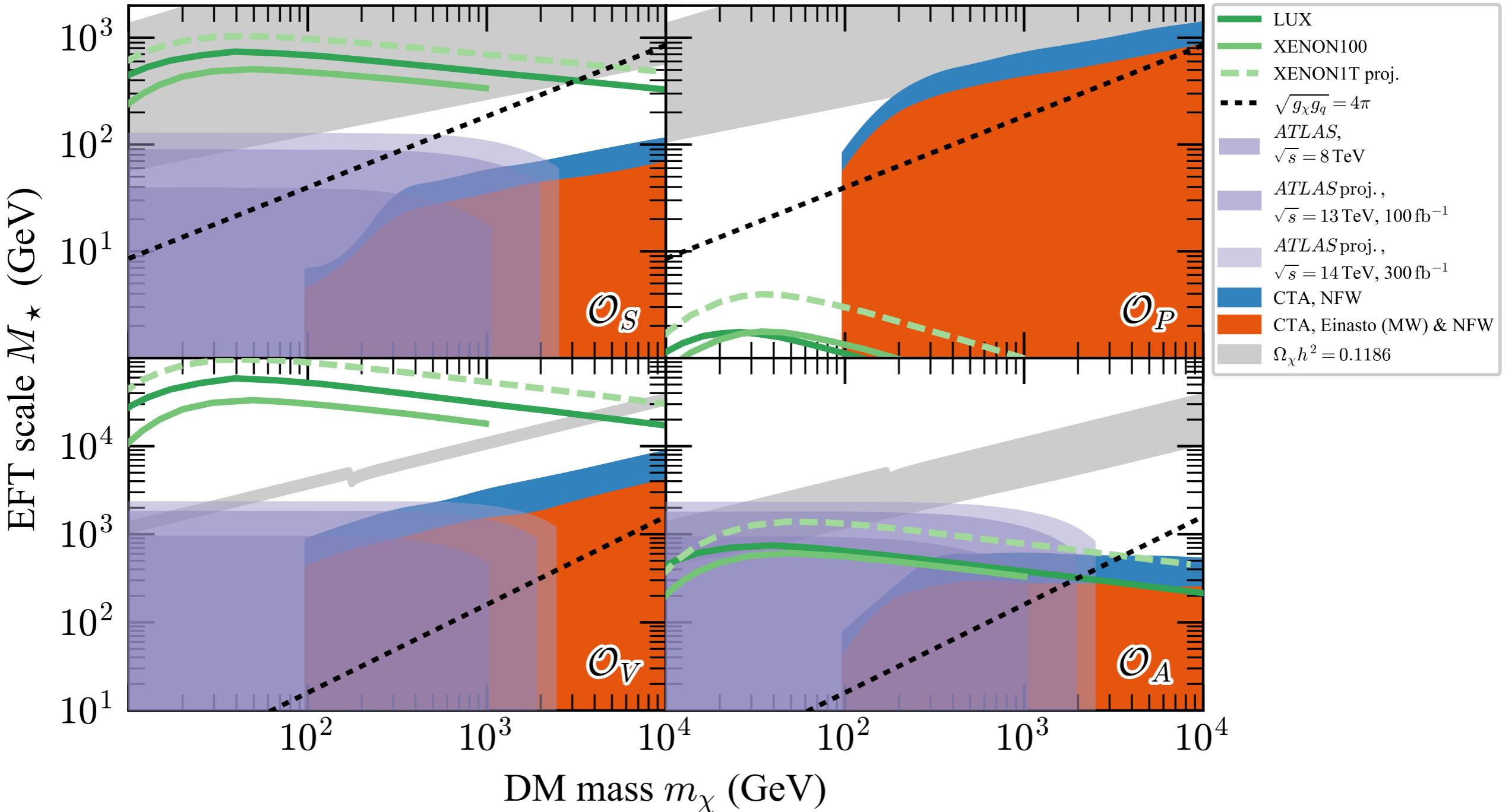


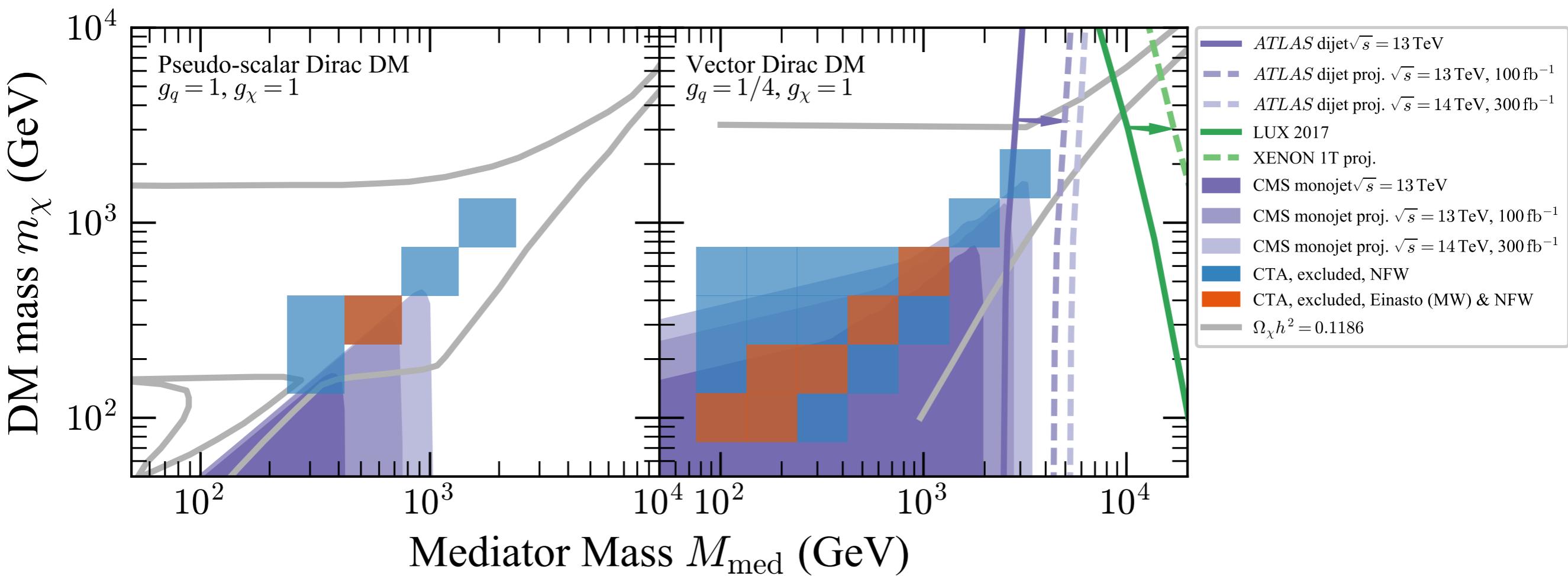
# Limits



# Limits







# Conclusion

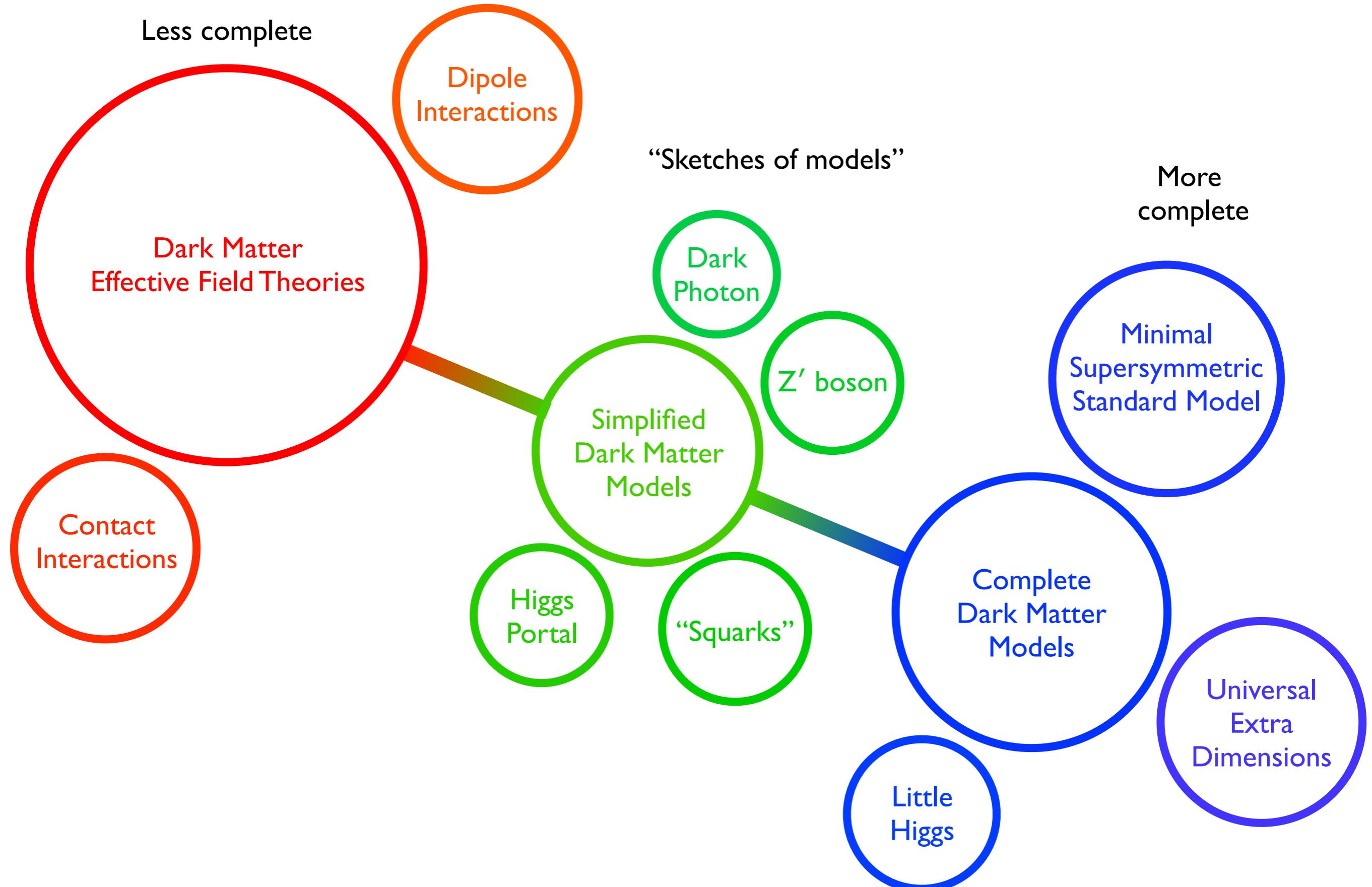
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- Though we have strong evidence for the existence of dark matter, its particle nature remains unknown
- A wide array of complementary experimental efforts are underway to probe the nature of dark matter
- Effective Field Theories and simplified models allow for an elegant comparison of constraints from a wide range of experiments
- CTA will place strong constraints on dark matter
- In particular, CTA will provide strong constraints on Vector mediated dark matter, and will lead the field in constraints on Pseudoscalar models



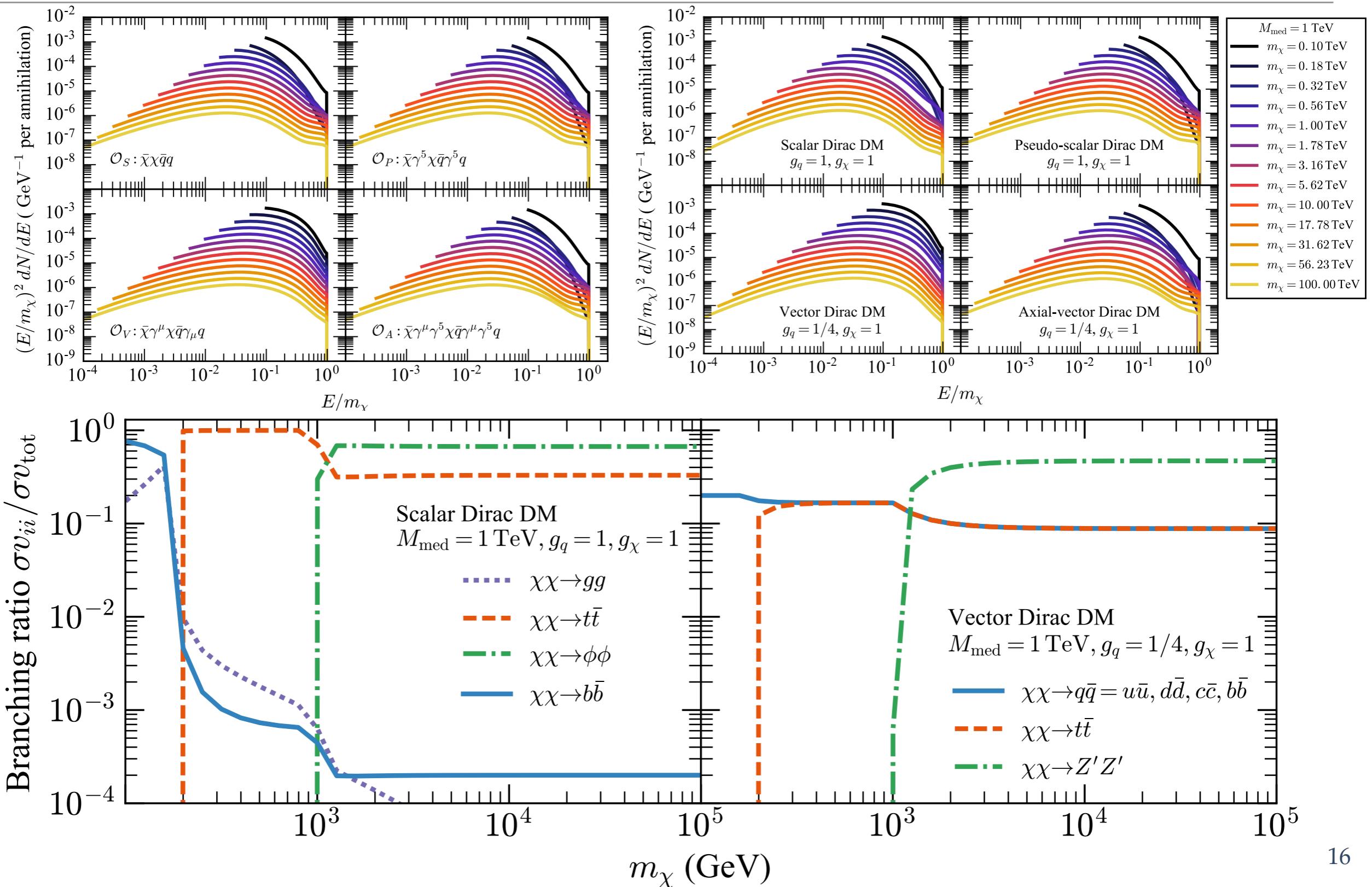
# Backup

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# Spectra and branching ratios

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# Limits

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