



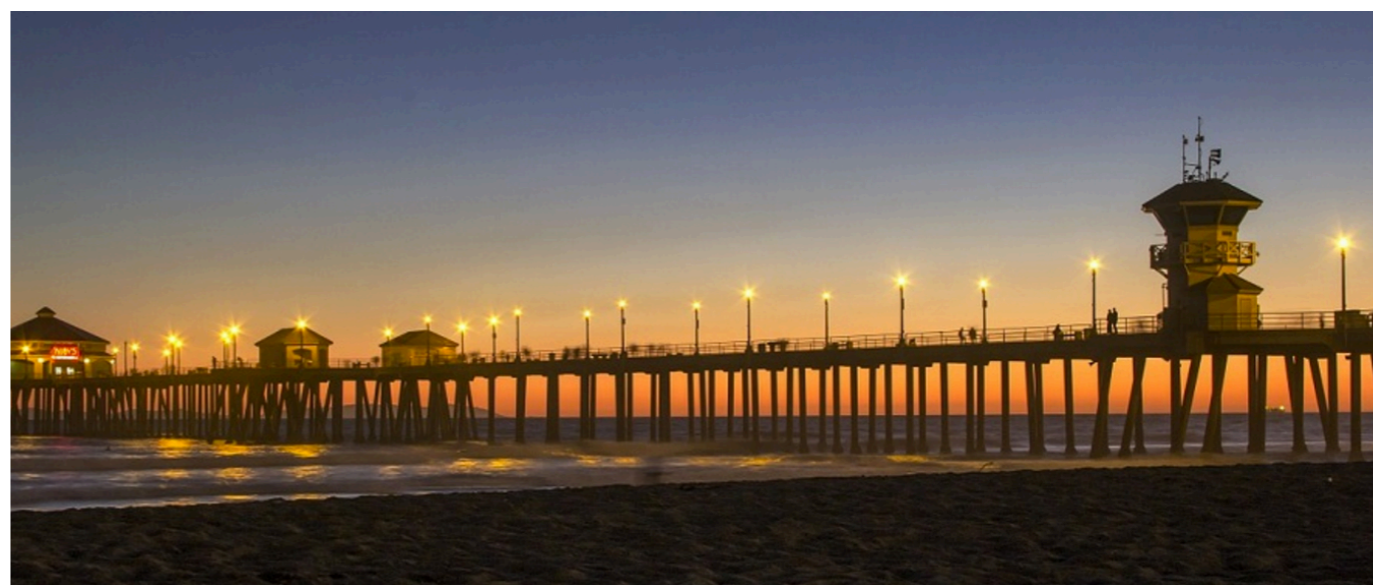
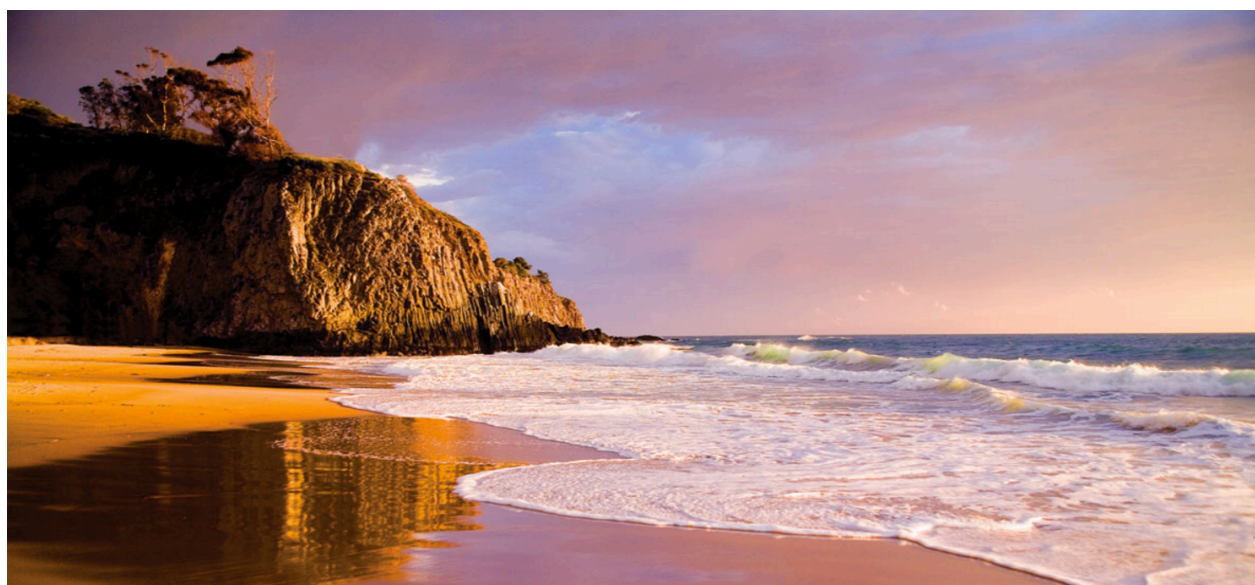
# Loop-induced DM production via ZH



DM@LHC 2017

Irvine - 04.04.2017

Dorival Gonçalves



# Motivation

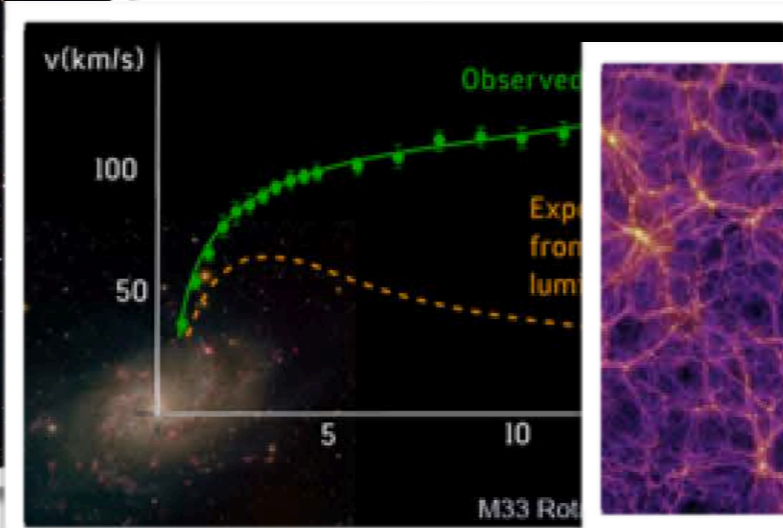
- After the Higgs discovery at the LHC  
Dark Matter is the next most important physics problem to tackle at the LHC  
**DM@LHC?!**

- The evidence just keeps piling up

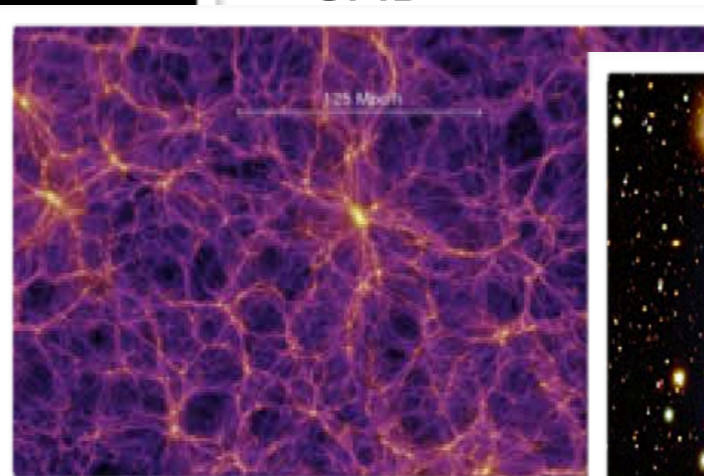
Galaxy clusters



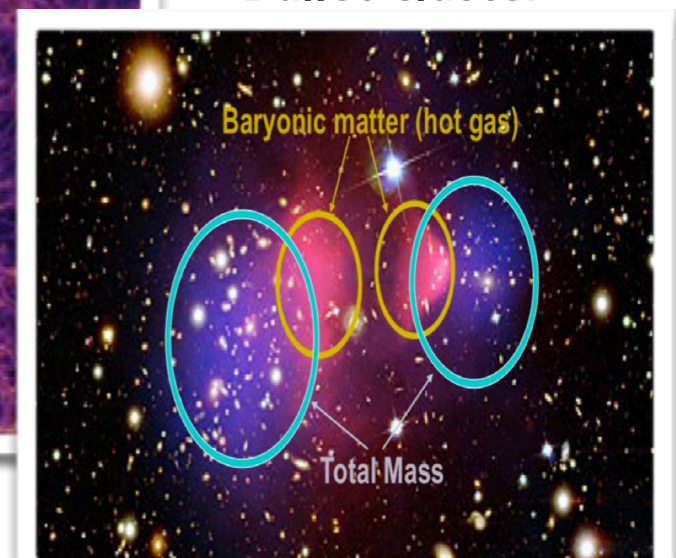
Galactic Rotation Curves



CMB



Bullet cluster

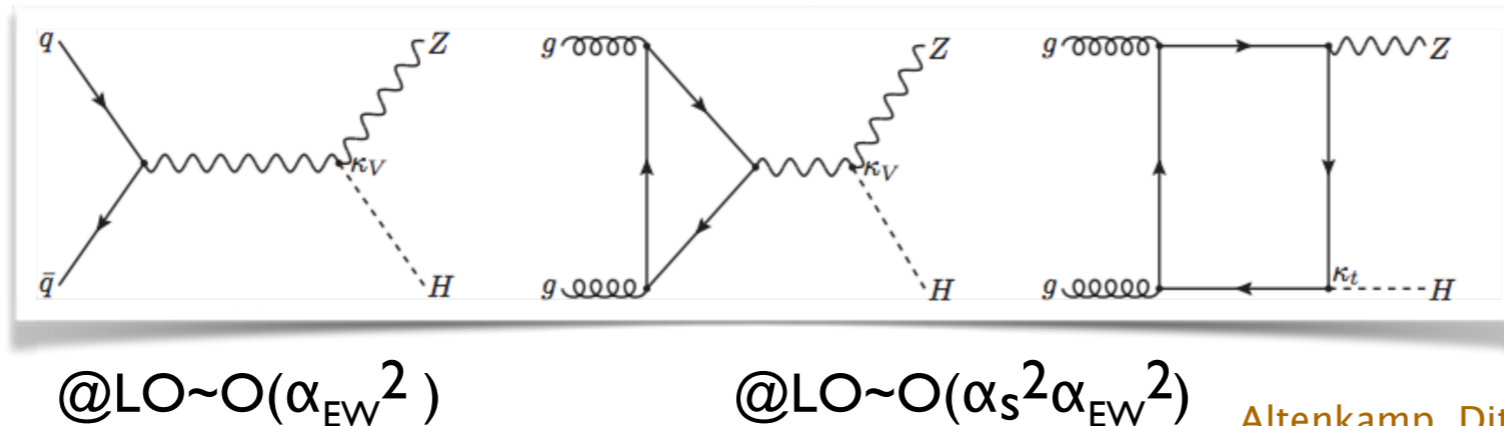


- New paths to address this problem at the LHC are very welcome  
We will focus on ZH channel and show that loop-induced open new possibilities
- Invisible Higgs searches
- Simplified Models for Dark Matter
- Simple Models for Dark Matter: Pseudoscalar Portal



# ZH Production

## ZH signal components:



Altenkamp, Dittmaier, Harlander, Rzehak, Zirke (2013)  
 Englert, McCullough, Spannowsky (2013)  
 Hespel, Maltoni, Vryonidou (2015)  
 Campbell, Ellis, Williams (2016)

There are four major factors that guarantee GF larger than the anticipated naive  $\alpha_s^2 \approx 1\%$ :

➔ Larger gluon PDF

➔ Larger initial state colour factor

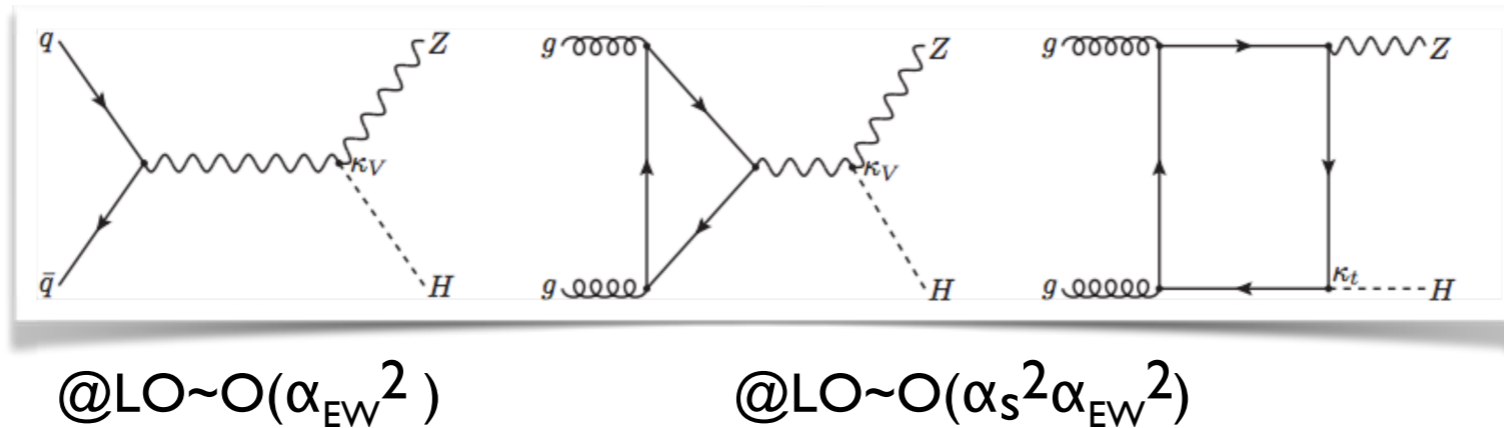
➔ Top Yukawa coupling appears in the place  $\alpha_{EW}$  factors:  $y_t \sim O(1)$

➔ Threshold enhancement at  $m_{ZH} \sim 2m_t$ , which gives rise to relevant rates at the boosted regime  $p_{TH} \sim m_t$

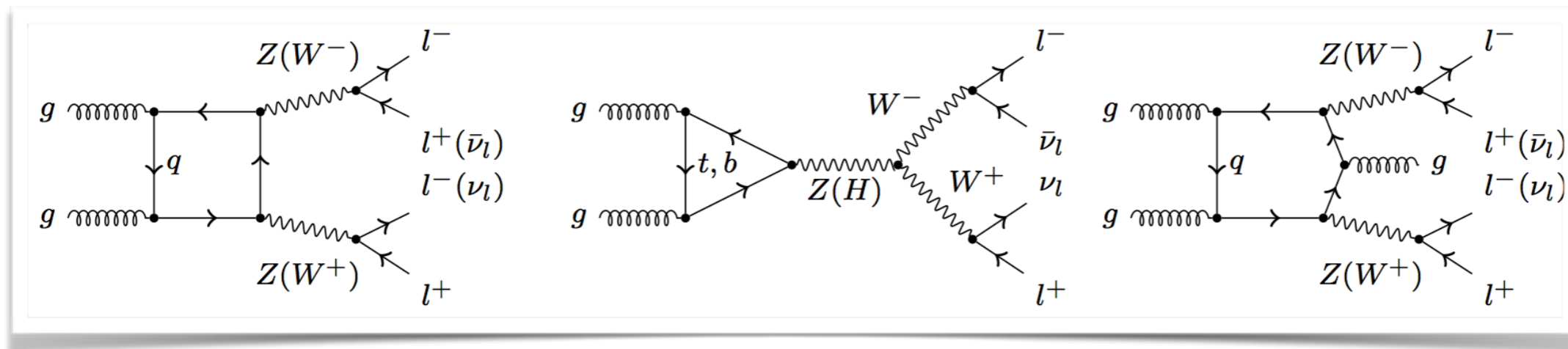
DG, Krauss, Kuttimalai, Maierhoefer (2015,2016)

# ZH Production

## ZH signal components:



By the same arguments, we should also consider the loop-induced background contributions  
 E.g., the major backgrounds for Invisible searches  $VV'$  also present relevant loop-induced components



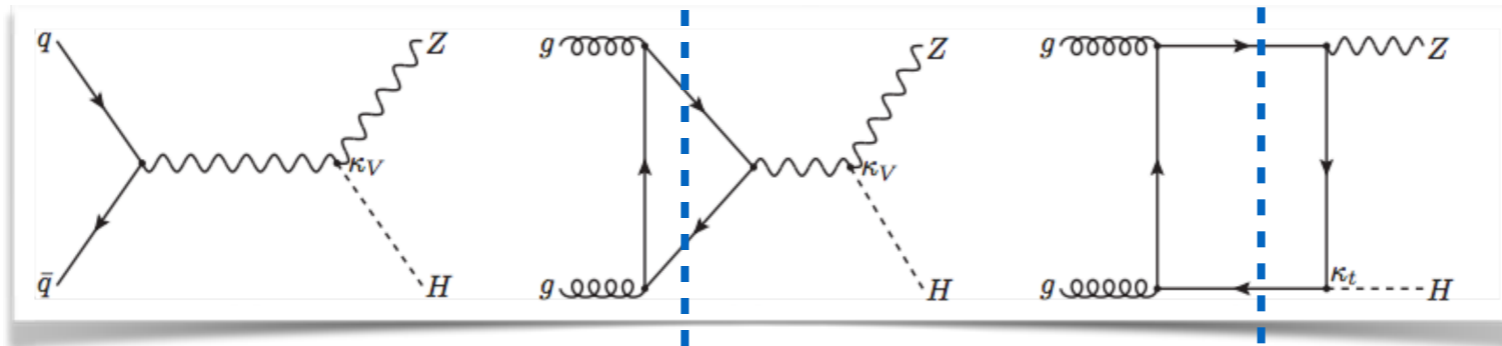
DG, Krauss, Kuttimalai, Maierhoefer (2015,2016)

The studies presented here were performed with Sherpa+OpenLoops  
 See Stefano's talk for details on OpenLoops

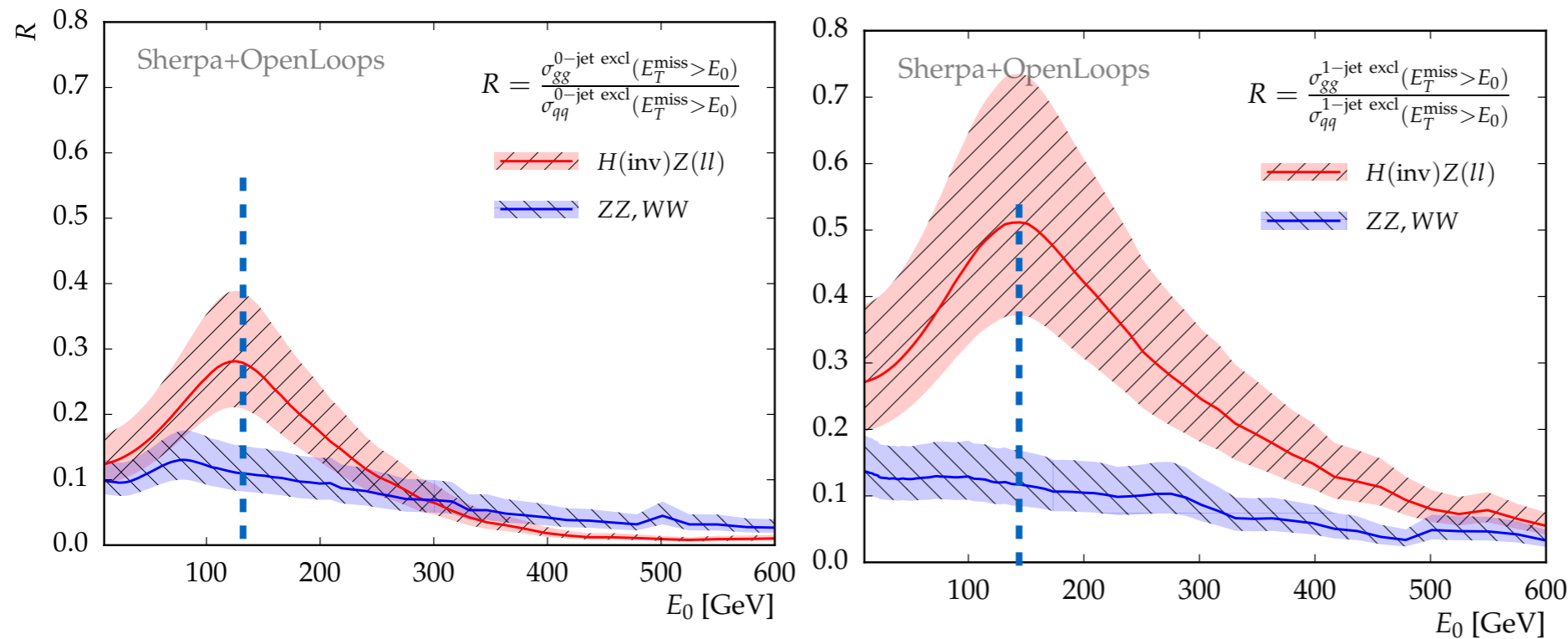


# ZH Production

ZH signal components:



Boosted kinematics enhances loop-induced component



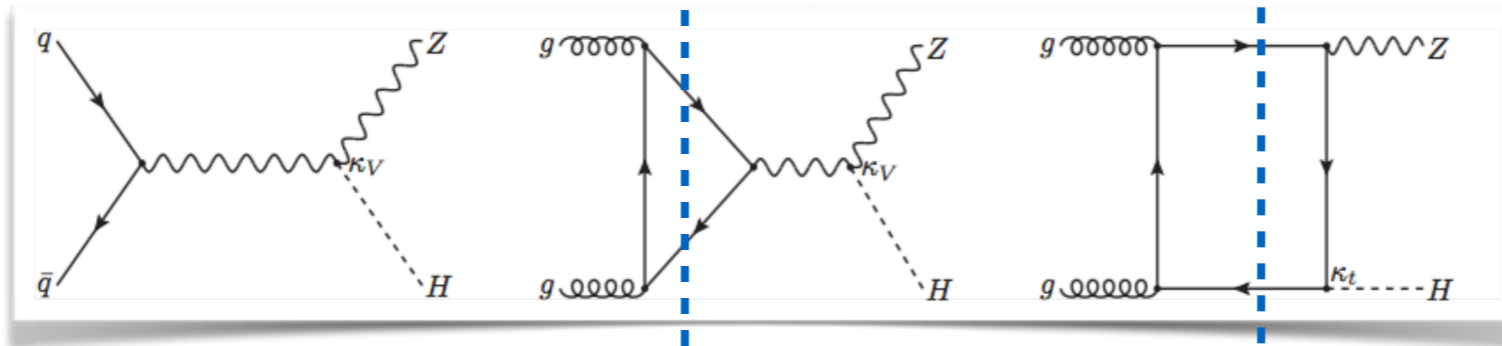
➡ DY signal/back via MEPS@NLO (0 and 1-jet at NLO): it works as towers of MC@NLO

➡ Loop-induced signal/back via MEPS@Loop<sup>2</sup> (0 and 1-jet loop<sup>2</sup>): CKKW merging

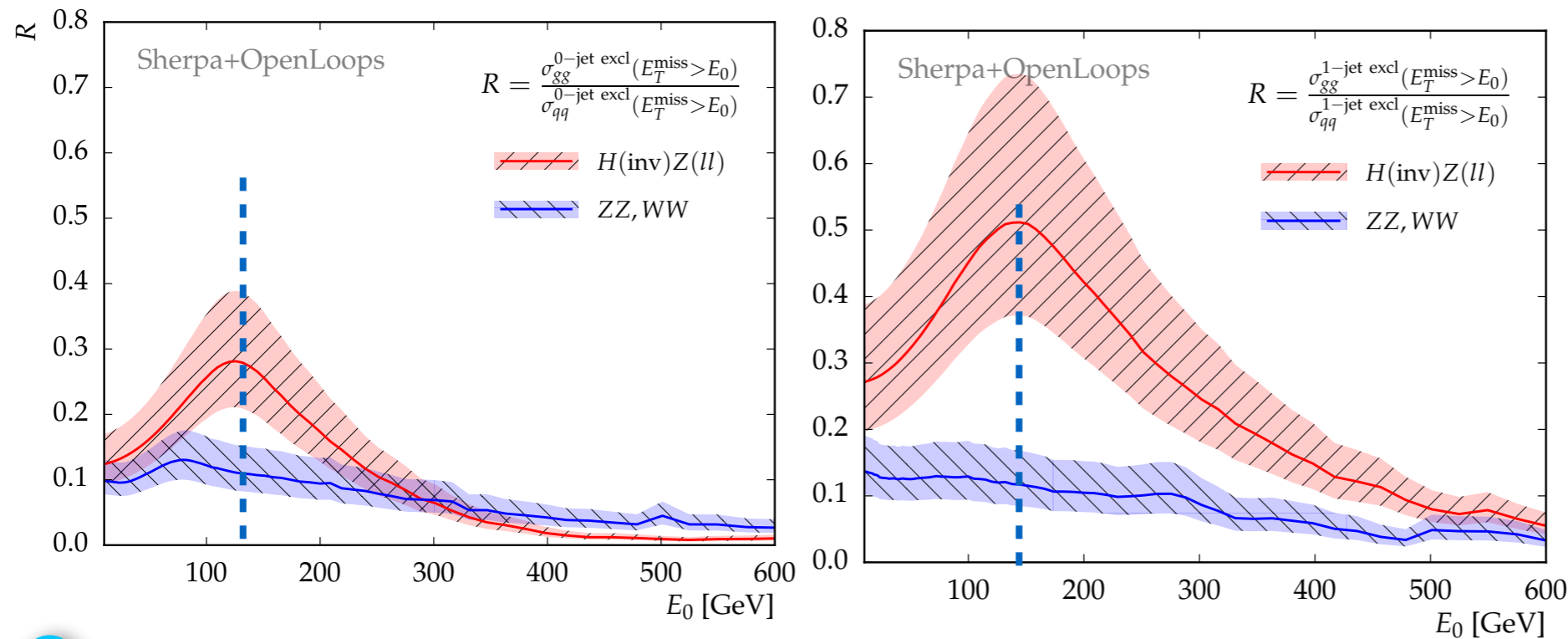
DG, Krauss, Kuttimalai, Maierhoefer (2015,2016)

# ZH Production

ZH signal components:



Boosted kinematics enhances loop-induced component



Despite looking at 0 and 1-jet bins, CMS Run I analyses neglect this component. This impacts on:

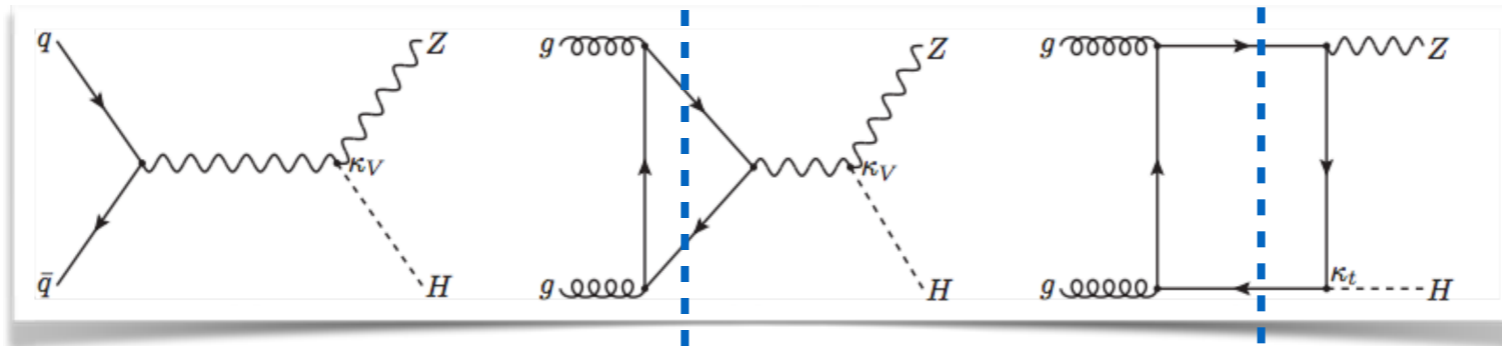
➔ Invisible bounds from  $Z(\text{ll})H(\text{inv})$

➔ And  $y_b$  from  $Z(\text{ll})H(\text{bb})$

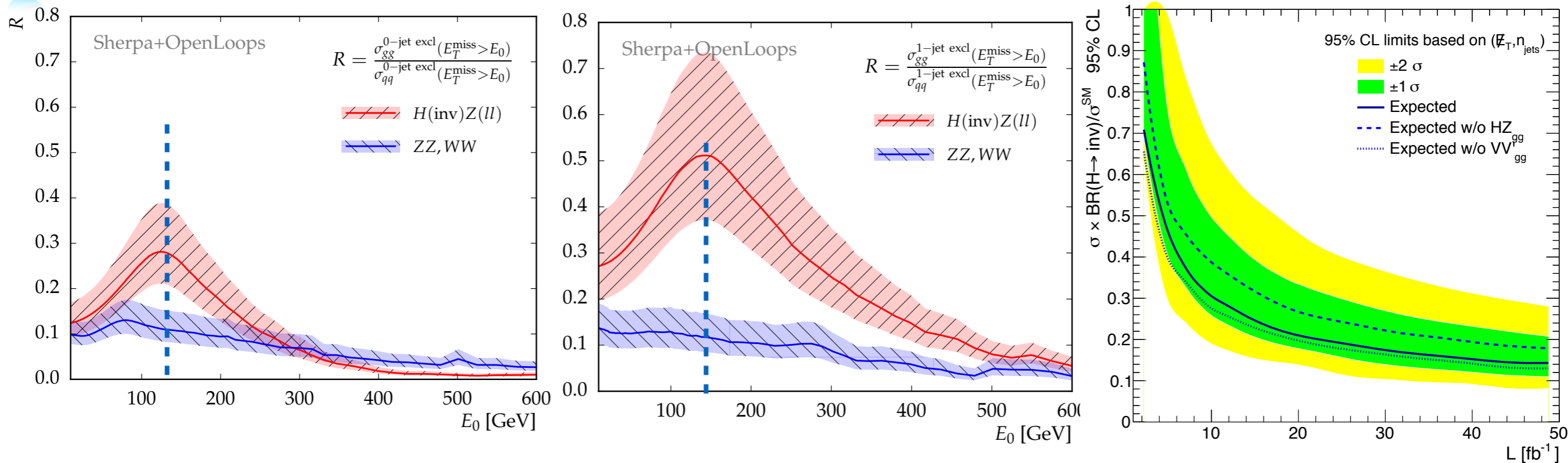
DG, Krauss, Kuttimalai, Maierhoefer (2015,2016)

# ZH Production

ZH signal components:



Boosted kinematics enhances loop-induced component



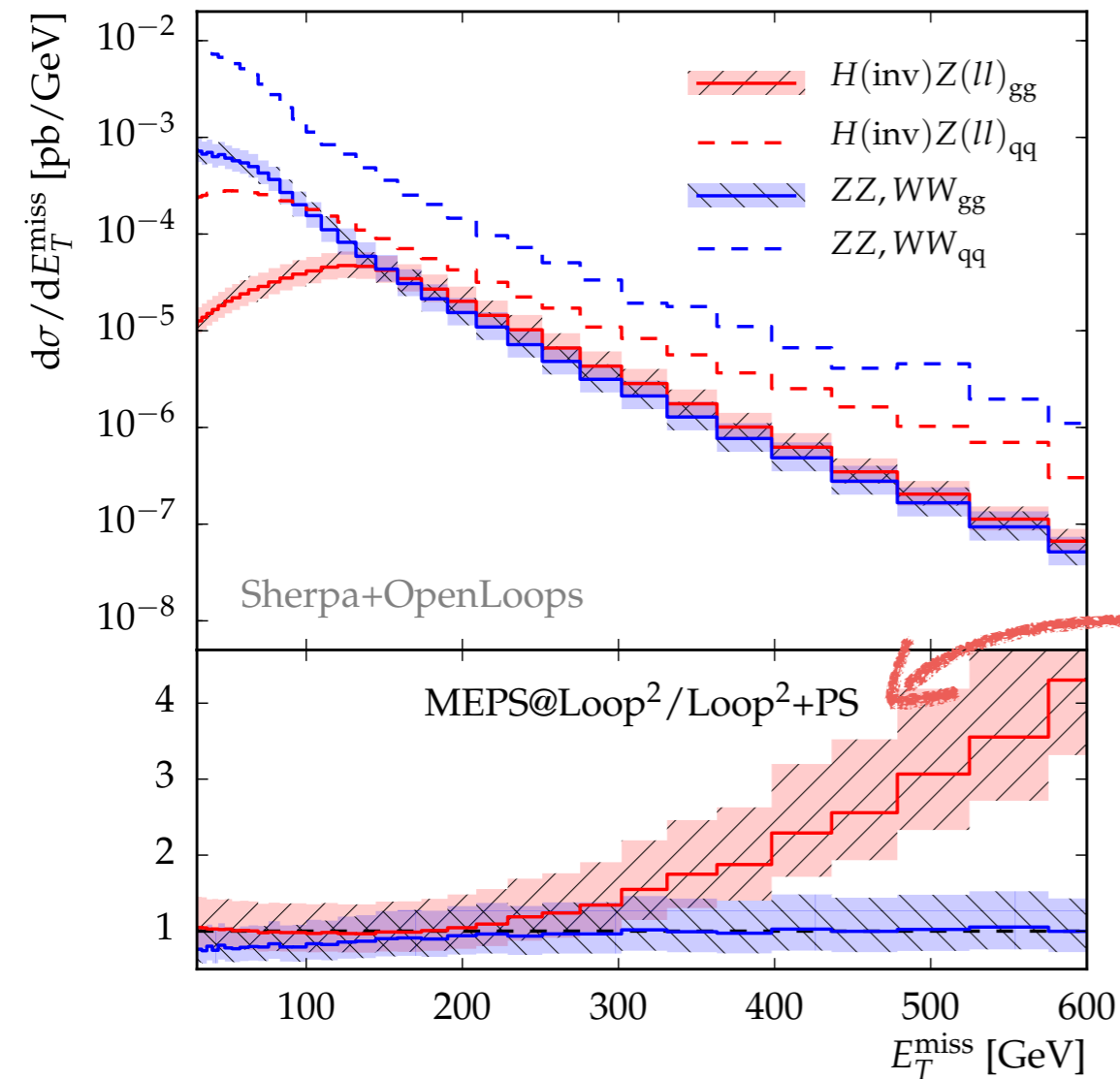
Loop-induced ZH makes the  $\text{BR}_{\text{inv}}$  bound stronger by more than 30%

DG, Krauss, Kuttimalai, Maierhoefer (2015,2016)

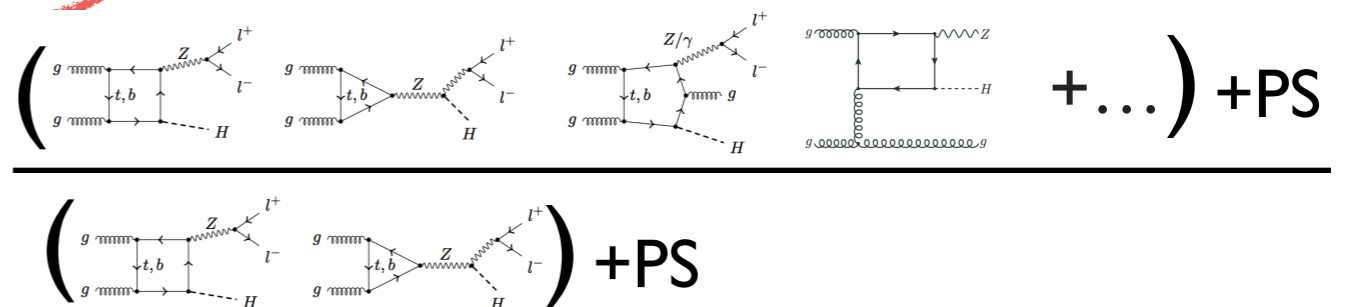


# Multi-jet Merging effects

As the GF becomes a significant player at boosted regimes, a proper modelling is of vital importance



- Effects induced by higher jet multiplicity ME beyond the scope of conventional PS alone
- Multi-jet merging correctly fill these phase space regions



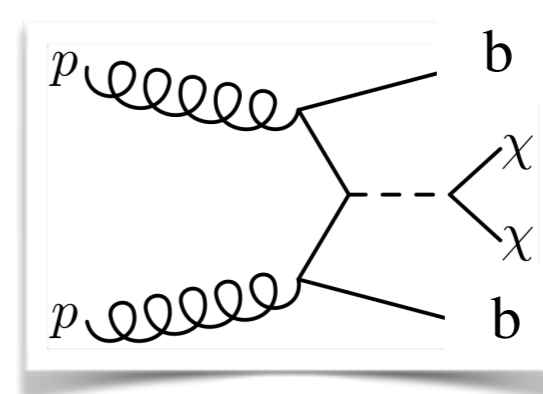
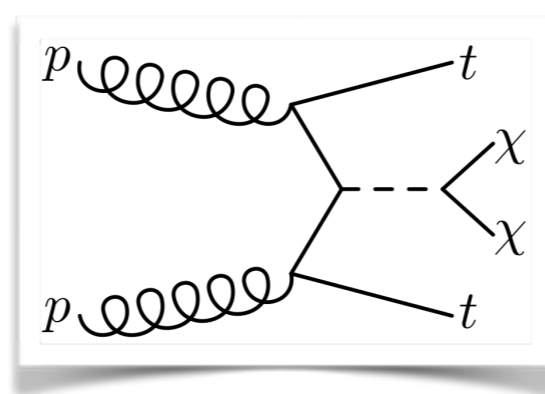
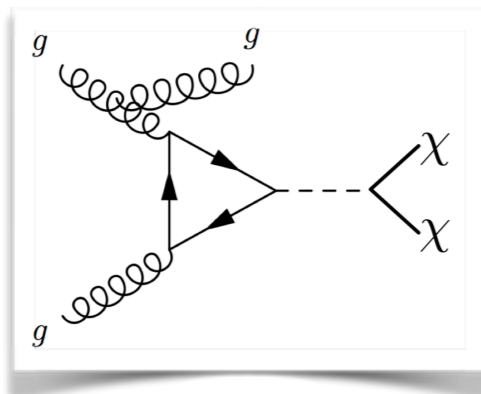
DG, Krauss, Kuttimalai, Maierhoefer (2015, 2016)

# Scalar Simplified Models for DM

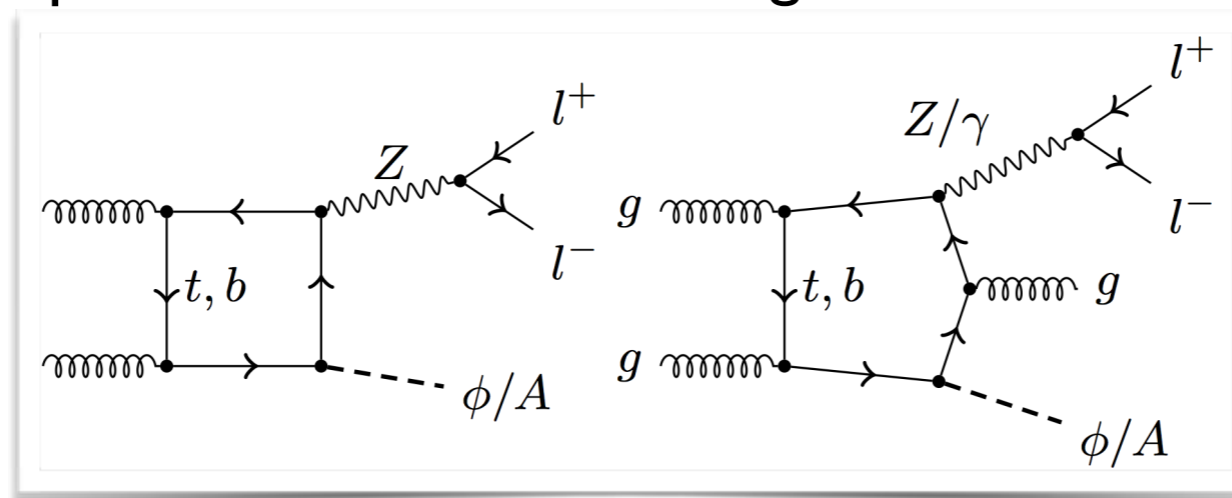
- Let's focus on Dirac fermion DM with **scalar/pseudoscalar** mediator
- Fermionic couplings proportional to the SM Yukawas ( $y_f$ ) - MFV avoids Flavor constraints

$$\mathcal{L} \supset - \sum_f \frac{y_f}{\sqrt{2}} (g_v^\phi \phi \bar{f} f + i g_v^A A \bar{f} \gamma_5 f) - g_\chi^\phi \phi \bar{\chi} \chi - i g_\chi^A A \bar{\chi} \gamma_5 \chi$$

- Three well studied signatures at Run-I LHC:



- Loop-induced ZH produce a relevant new signature:

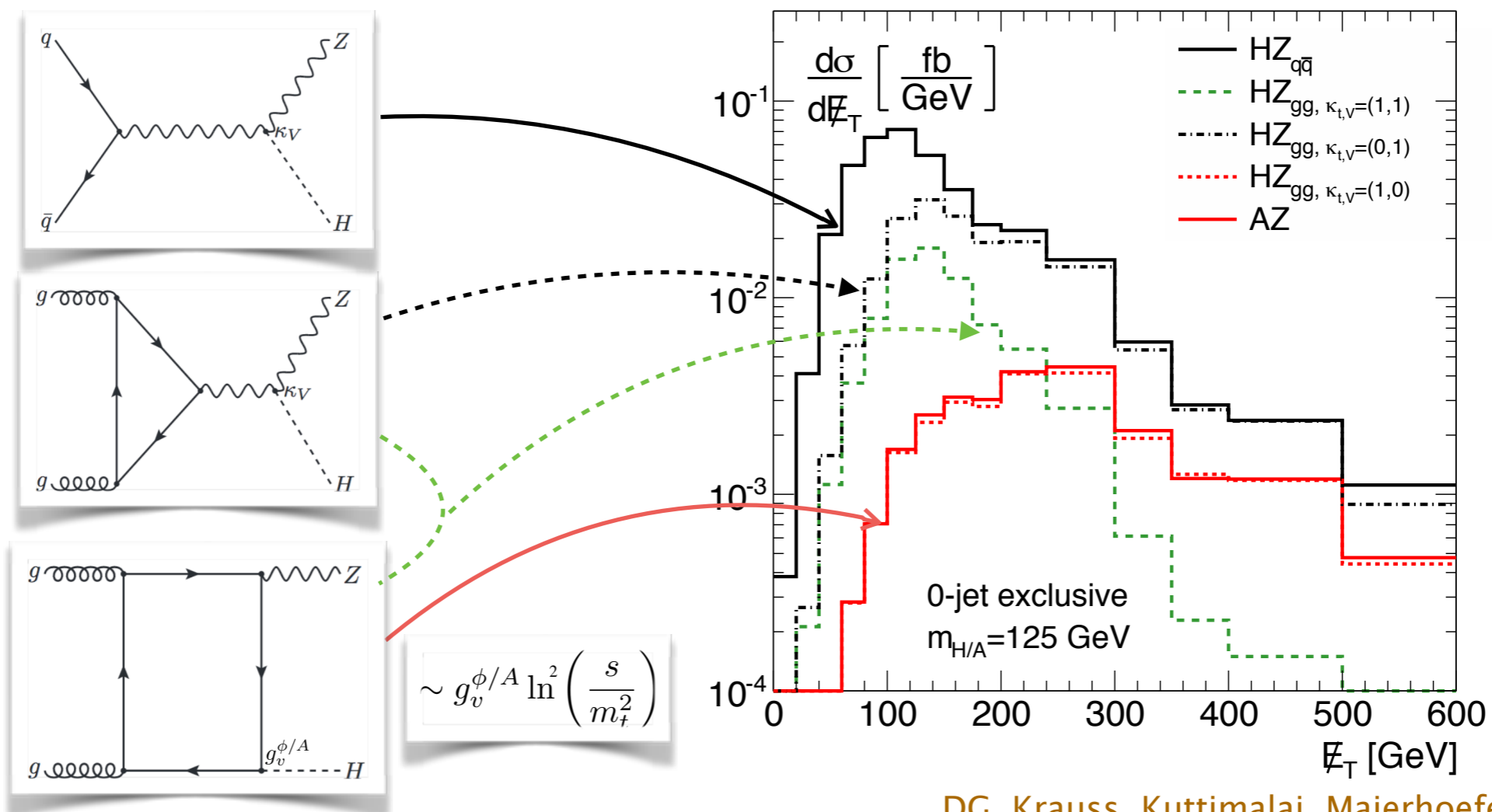


DG, Krauss, Kuttimalai, Maierhoefer (2016)

# Scalar Simplified Models for DM

- Let's focus on Dirac fermion DM with **scalar/pseudoscalar** mediator
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DG, Krauss, Kuttimalai, Maierhoefer (2016)



# UV completions

→ Can we improve these searches and/or find another sensitive channel(s)?

→ Which kind of UV completions are we mapping? ~~Simplified~~ **Simple Models**

Since Scalar case is mostly bound by DD, let's focus on Pseudoscalar scenario where DD is suppressed & LHC is more sensitive

$$\mathcal{L}_s = \bar{\chi}(i\not{\partial} - m_\chi)\chi + \frac{1}{2}(\partial_\mu a)^2 - \frac{m_a^2}{2}a^2 - g_\chi a \bar{\chi}i\gamma^5\chi - g_{SM} a \sum_f \frac{y_f}{\sqrt{2}} \bar{f}i\gamma^5 f$$

if  $\chi$  is SM gauge singlet  $\mathcal{L}$  is not gauge invariant

"a" needs  $SU(2)_L \times U(1)_Y$  charge to couple to SM fermions

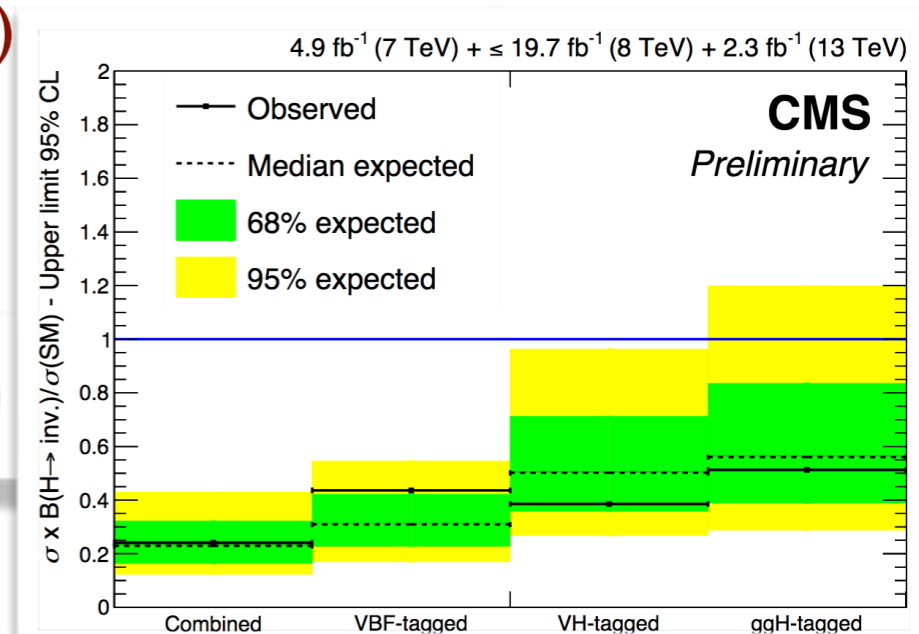
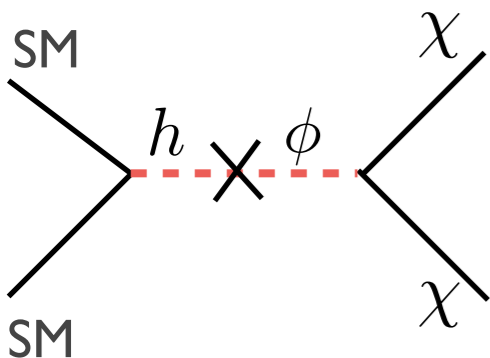
Needs mixing between "a" & scalar EW multiplet (new states)

→ Ex: Higgs mixing  $\mathcal{L} \supset -y_f \bar{f}Hf - \mu_\phi \phi |H|^2 - y_\chi \phi \bar{\chi}\chi$

O'Connell, Ramsey-Musolf, Wise (2006)

$$\mathcal{L} \supset \sin\theta\phi \left( \frac{m_f}{v} \bar{f}f + \frac{m_l}{v} \bar{l}l + \frac{2m_W^2}{v} W^2 + \frac{m_Z^2}{v} Z^2 \right)$$

Craig, Lou, McCullough, Thalapillil (2014)



# Pseudoscalar Portal

**Pseudoscalar Portal:** simple renormalizable completion with **EW singlet DM**

Ipek, McKeen, Nelson (2014)

It works as a generalization of the Higgs Mixing framework to 2HDM

$$\mathcal{L} \supset \underbrace{-V_{2\text{HDM}}}_{\text{visible sector}} - \underbrace{\frac{m_{a_0}^2}{2} a_0 - ig_\chi a_0 \bar{\chi} \gamma_5 \chi}_{\text{dark sector}} + \underbrace{(i\kappa a_0 H_1^\dagger H_2 + h.c.)}_{\text{portal}}$$

with  $y_\chi, \kappa \in \text{Real}$  & no CP-violation in 2HDM

$$V_{\text{dark}} \supset g_\chi (c_\theta a + s_\theta A) \bar{\chi} i \gamma^5 \chi$$

**mixing**

$$\begin{aligned} A &= c_\theta A_0 + s_\theta a_0 \\ a &= c_\theta a_0 - s_\theta A_0 \end{aligned}$$

$$\begin{aligned} V_{\text{portal}} &= \frac{(m_A^2 - m_a^2) s_{2\theta}}{2v} (c_{\beta-\alpha} H - s_{\beta-\alpha} h) \\ &\times [aA (s_\theta^2 - c_\theta^2) + (a^2 - A^2) s_\theta c_\theta] . \end{aligned}$$

→ 2HDM Yukawas:

Type-I

$$i(c_\theta A - s_\theta a) t_\beta^{-1} \sum_q \frac{y_q}{\sqrt{2}} \bar{q} \gamma^5 q$$

Type-II

$$i(c_\theta A - s_\theta a) \left( t_\beta^{-1} \sum_u \frac{y_q}{\sqrt{2}} \bar{q} \gamma^5 q + t_\beta \sum_d \frac{y_q}{\sqrt{2}} \bar{q} \gamma^5 q \right)$$

DG, Machado, No (2016)

Bauer, Haisch, Kahlhoefer (2017)

# Pseudoscalar Portal

SM + DM + “a” mediator + new states



A: new (heavier) mediator

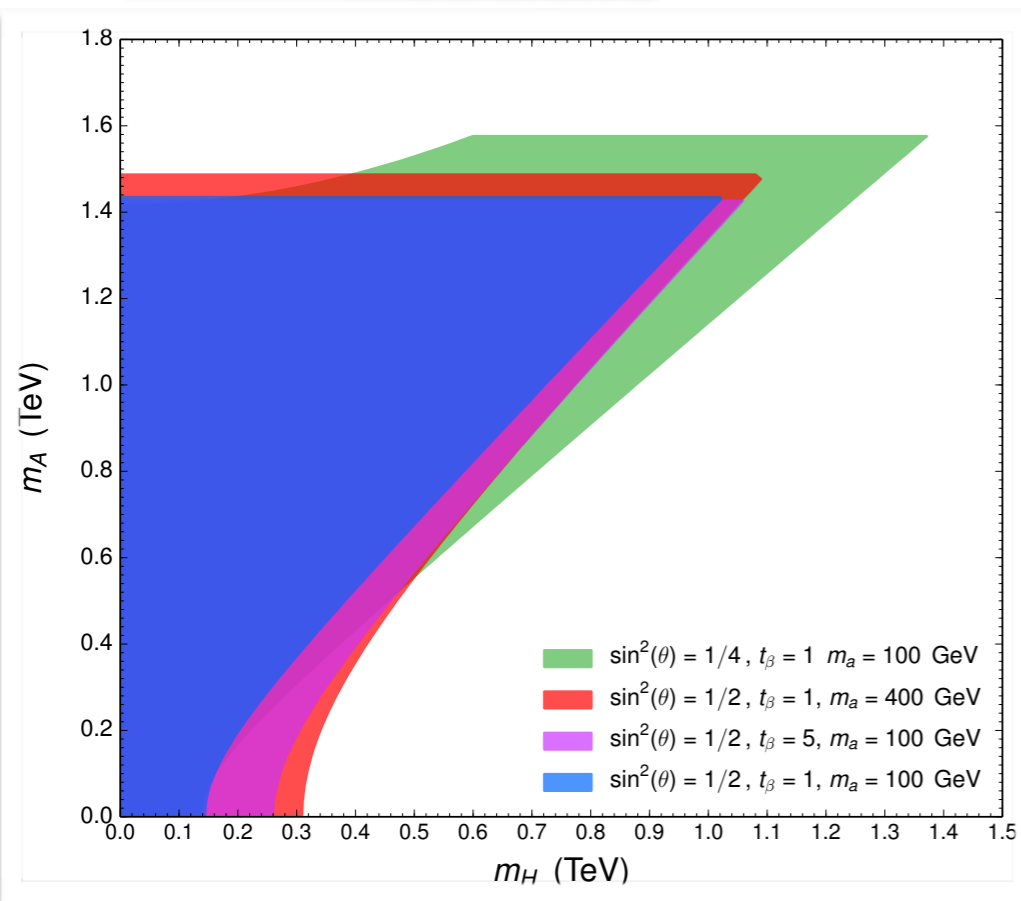
$H^\pm, H_0$ : New states (mediator gauge partners)

1)  $m_{A, H^\pm, H_0} \sim m_a$

2 mediator scenario +  $H^\pm, H_0$  pheno as in usual 2HDM

2)  $m_{A, H^\pm, H_0} \gg m_a$

Closer to Simplified Model



DG, Machado, No (2016)

If  $\sin \theta$  is fixed new states do not fully decouple

Mass Splittings among  $A, H^\pm, H_0$  bounded by 2HDM unitarity

Ginzburg, Ivanov, *Phys. Rev D* **72** (2005) 115010

$$m_i - m_j \leq \mathcal{O}(\text{few}) \times v$$

Assume  $[c_{\beta-\alpha} = 0 \quad m_{H^\pm} = m_{H_0}]$

Perturbativity unitarity





# Pseudoscalar Portal

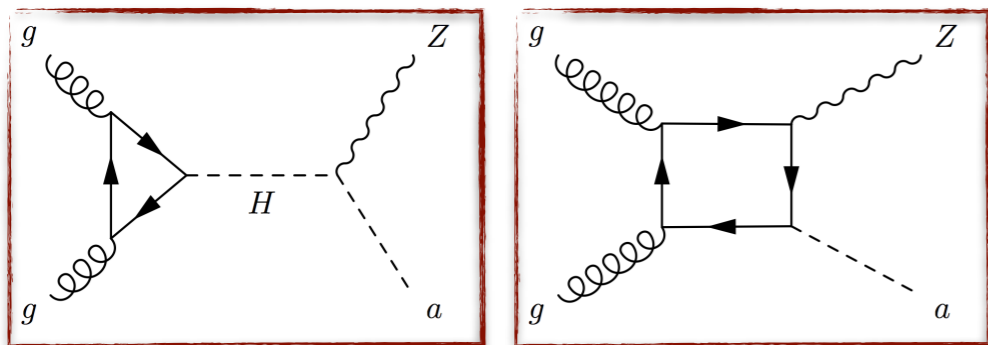
SM + DM + “a” mediator + new states

A: new (heavier) mediator

$H^\pm, H_0$ : New states (mediator gauge partners)

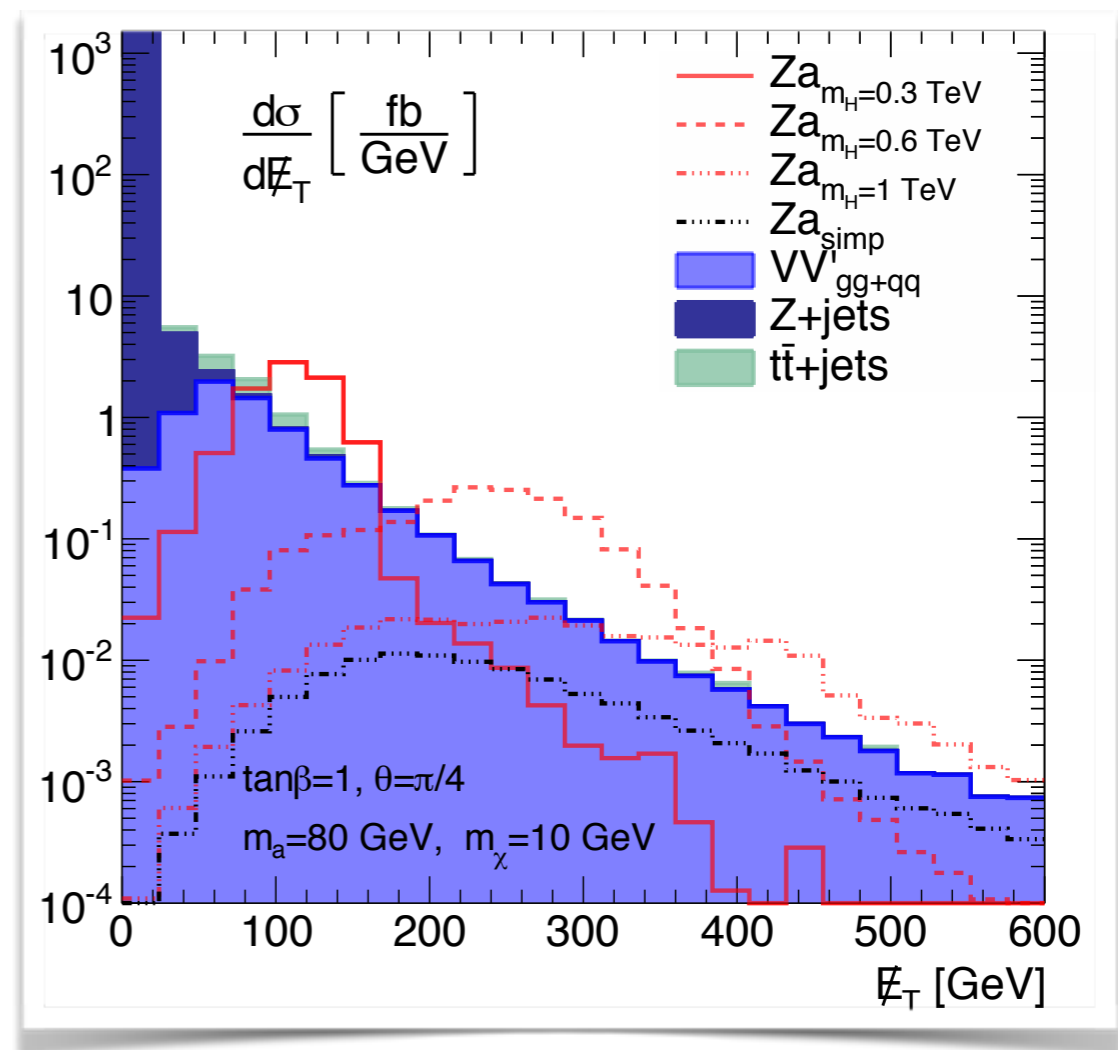
In the presence of mixing new states do not (fully) decouple

present in Simplified Model



Resonante mono-Z for  $m_H > m_a + m_Z$

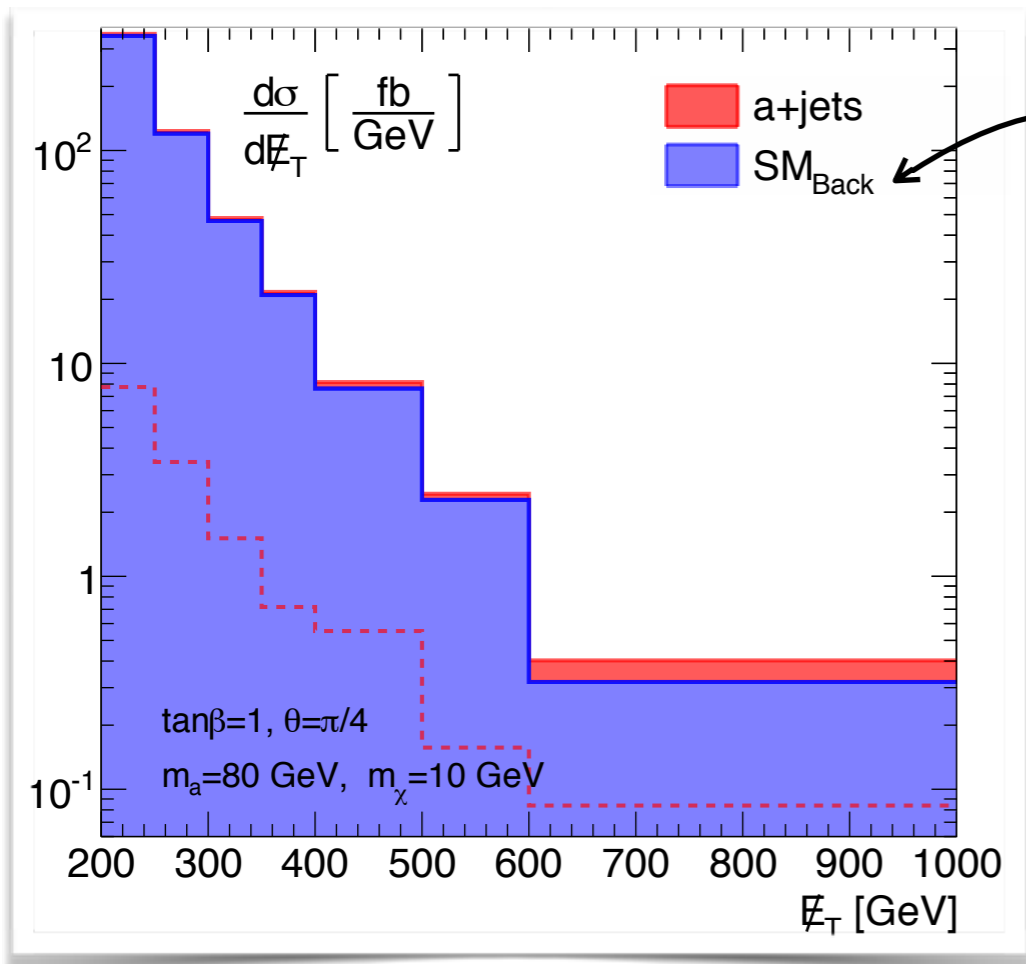
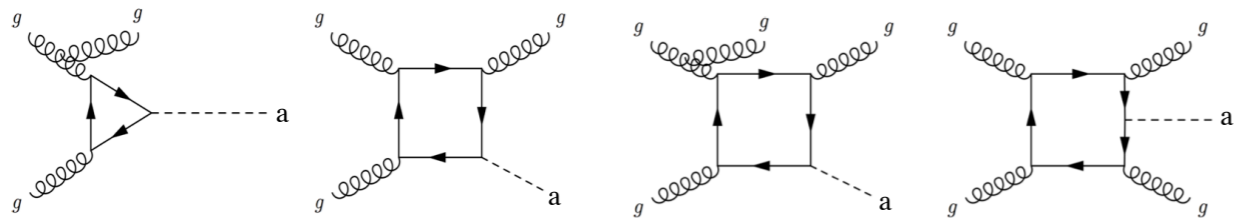
$$\cancel{E}_T^{\max} \sim \frac{1}{2m_H} \sqrt{(m_H^2 - m_a^2 - m_Z^2)^2 - 4m_Z^2 m_a^2}$$



DG, Machado, No (2016)

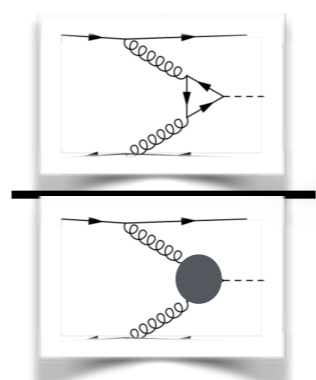
# Pseudoscalar Portal

Impact on usual monojets search:

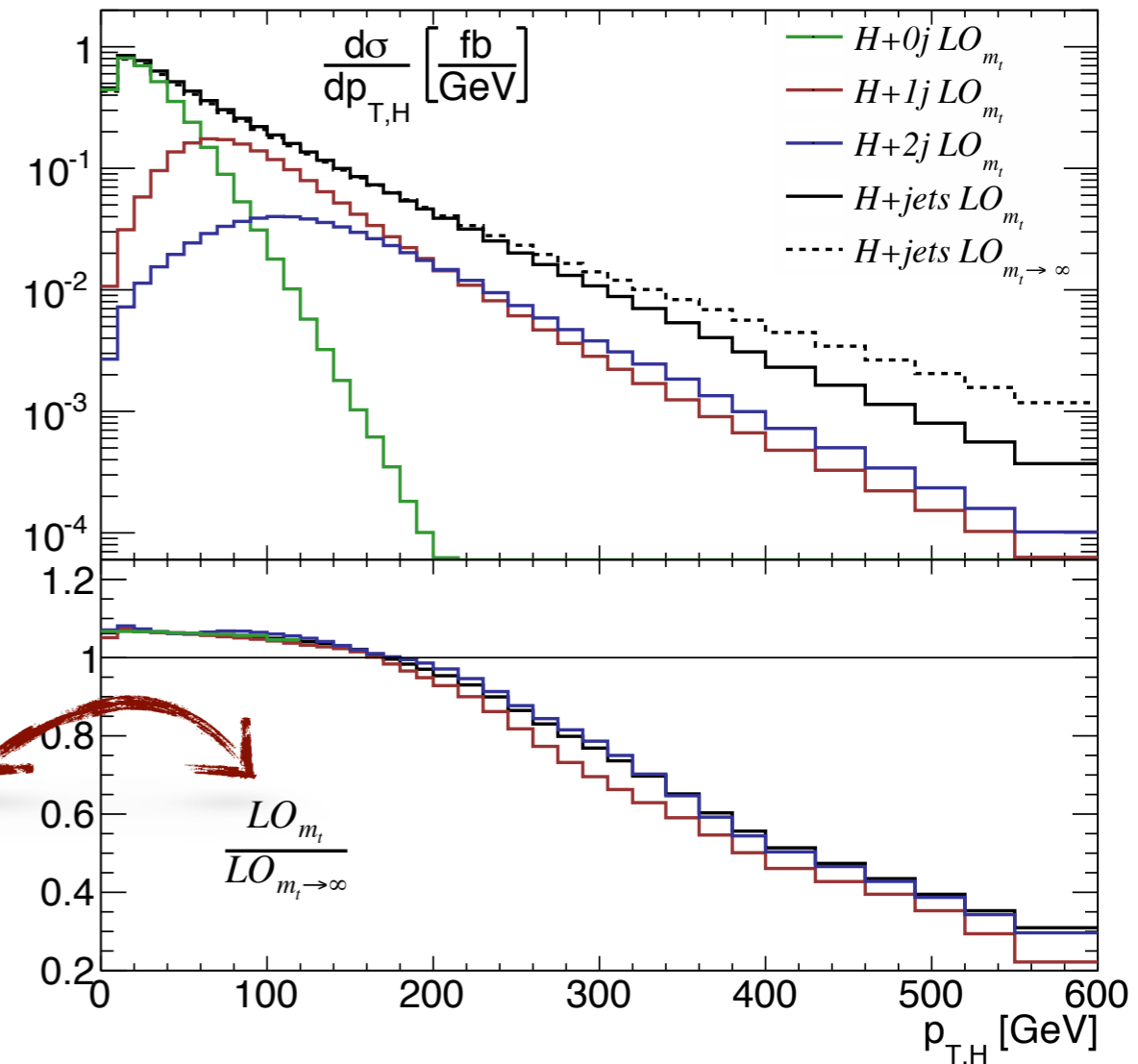


DG, Machado, No (2016)

CMS PAS EXO-15-003



Precision is fundamental: Top mass effects correction of  $O(4)$  at  $p_{T,H} \sim 600$  GeV



Bushman, DG, Krauss, Kuttimalai, Schonherr, Plehn (2014)

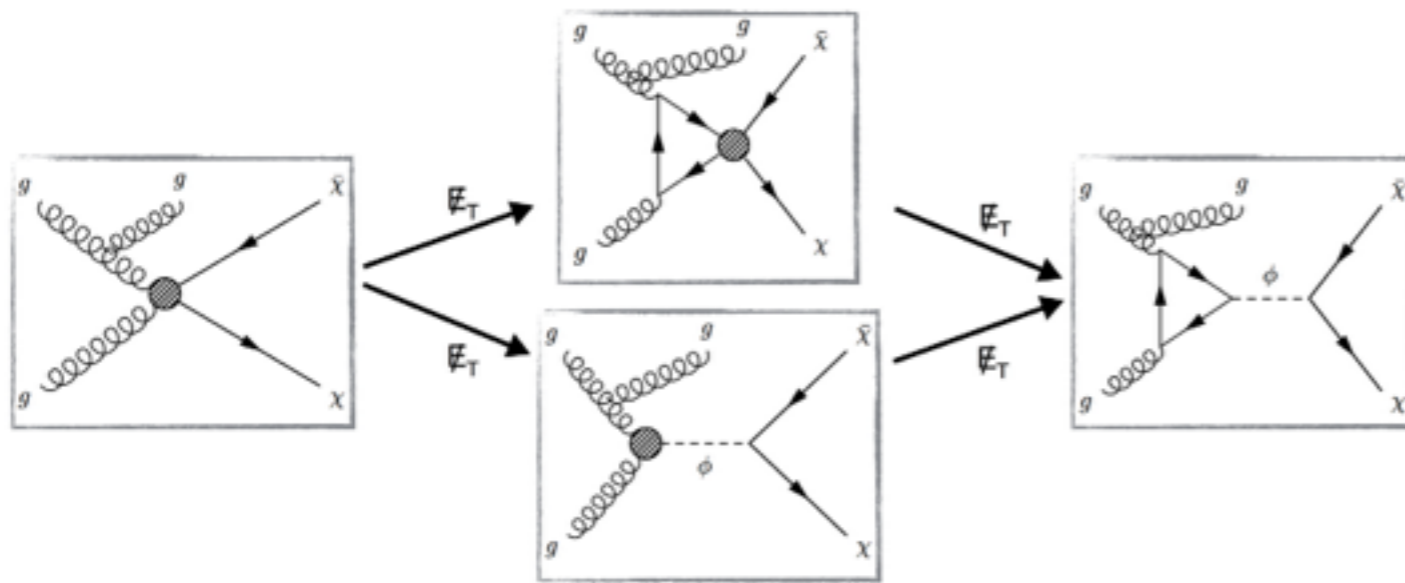
➡ Monojet searches will soon face systematic uncertainty wall at Run-II

➡ DM searches commonly require  $E_T \gg m_t$

# Pseudoscalar Portal

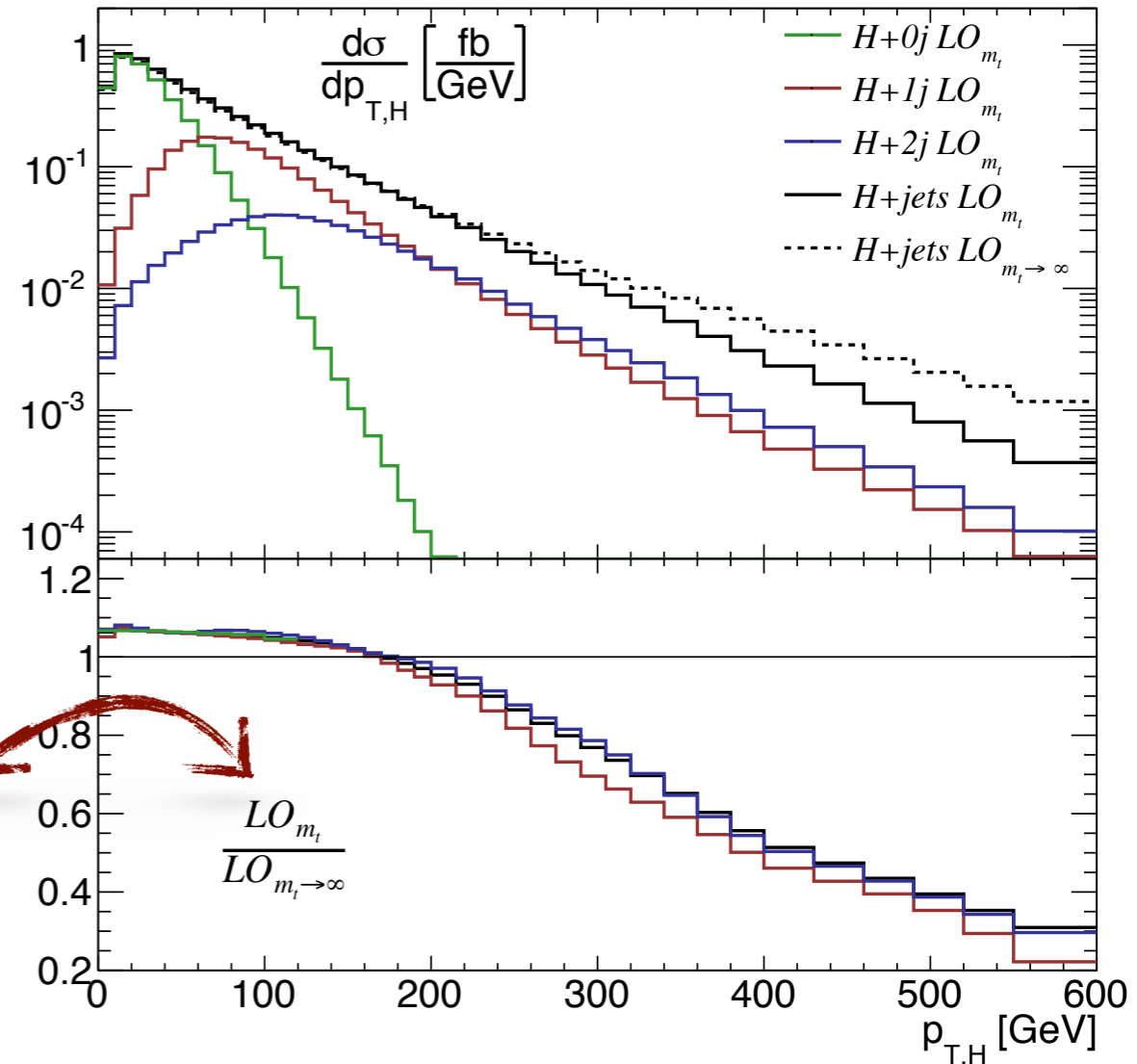
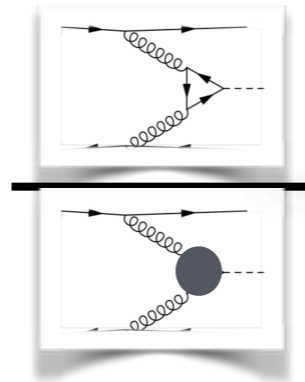
Top mass effects show up in other frameworks too. Take EFT as an example:

Precision is fundamental: Top mass effects correction of  $O(4)$  at  $p_{T,H} \sim 600$  GeV



Buckley, Feld, DG (2014)

Haisch, Kahlhoefer, Unwin (2013)

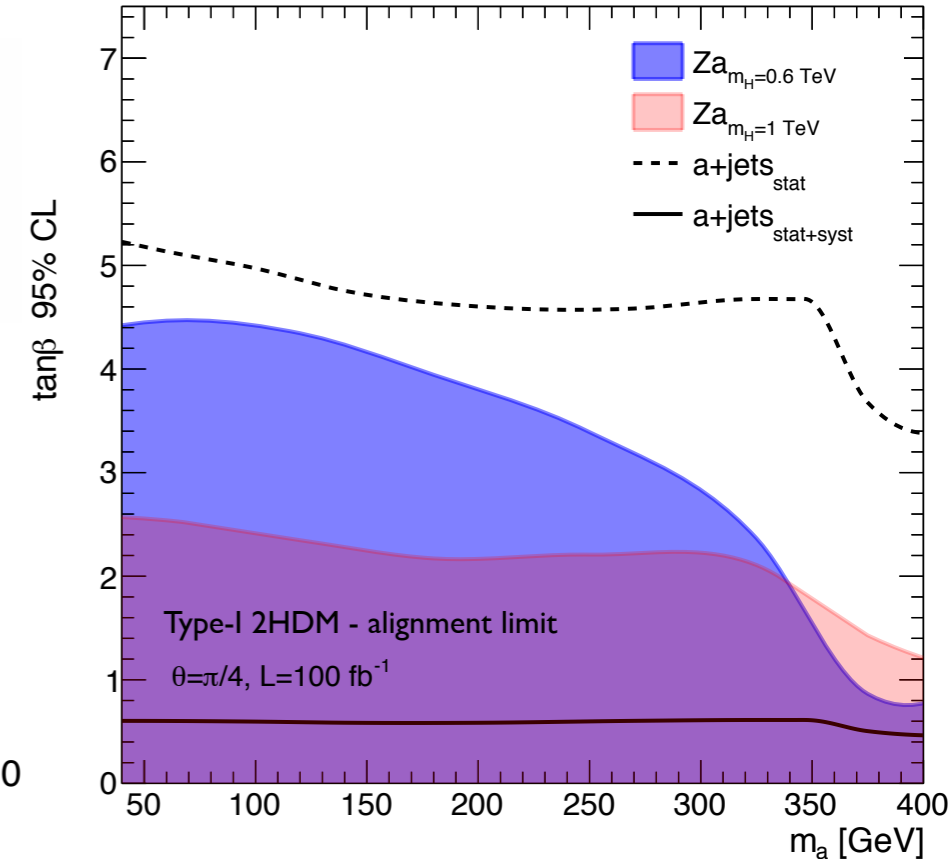
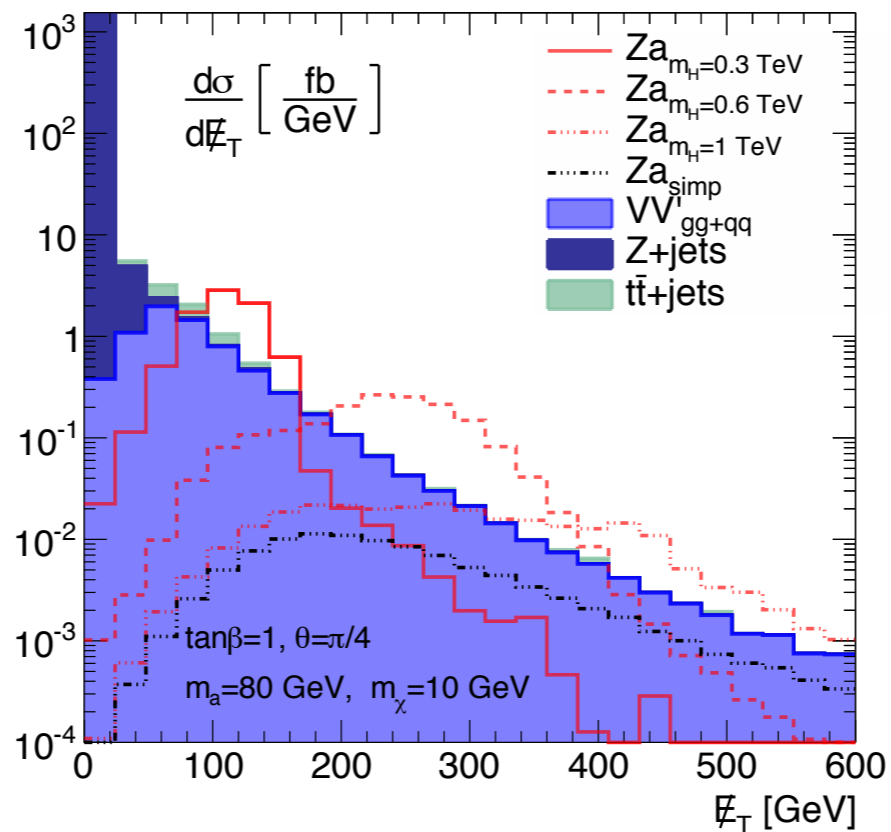
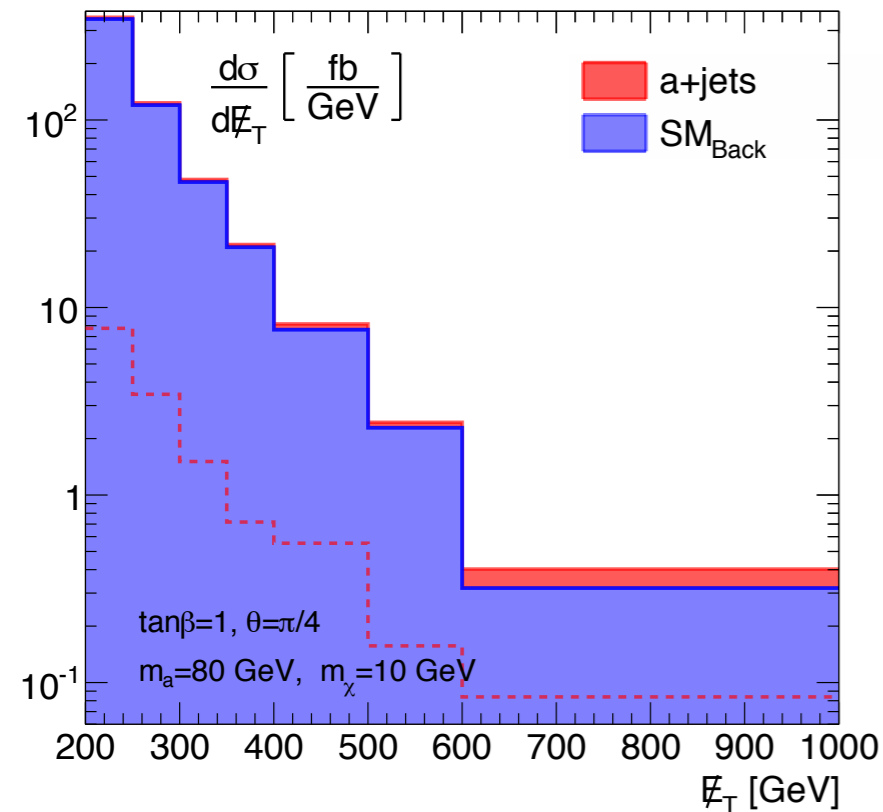


Bushman, DG, Krauss, Kuttimalai, Schonherr, Plehn (2014)

DM searches commonly require  $\cancel{E}_T \gg m_t$

# Pseudoscalar Portal

Sensitivity comparison:



DG, Machado, No (2016)

- ➡ Monojets searches will soon face systematic uncertainty wall at Run-II
- ➡ New channels are welcome. Mostly ones with better control on backgrounds
- ➡ Z(II)A(xx) provides competitive bounds to monojets in the pseudoscalar portal scenario

# Summary

- Loop-induced  $ZH(\text{inv})$  provide relevant contributions: bound stronger by  $\sim 30\%$
- Currently studied Simplified Models channels will severely suffer with systematic uncertainties at Run-II LHC. Alternative channels are welcome
- We showed that  $Za(\chi\chi)$  channel provides competitive bounds in the
  - **Simplified Model** framework (via loop-induced production)
  - **Pseudoscalar Portal** (via resonant heavy Higgs interfering with loop-induced)
- Extra channels can be motivated with ~~Simplified~~ <sup>Simple</sup> Models for DM  
Pseudoscalar Portal matches well with the Higgs(es) search program



**Thank you for your attention!**