

Probing the **chirality** of dark matter at colliders with dark photon showering

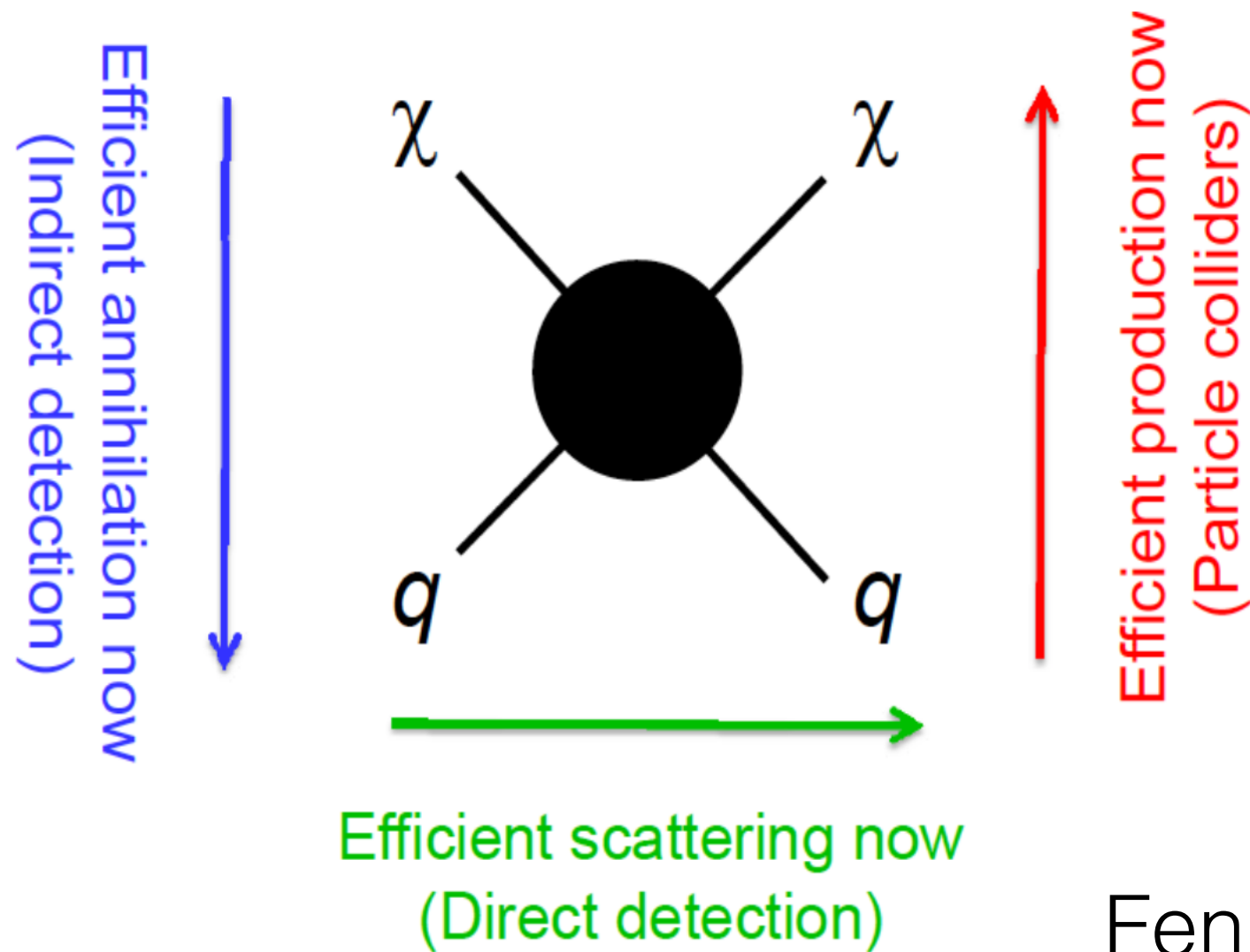
Myeonghun Park



Based on arXiv:1612:02850
with Mengchao Zhang, Minho Kim and Hye-Sung Lee

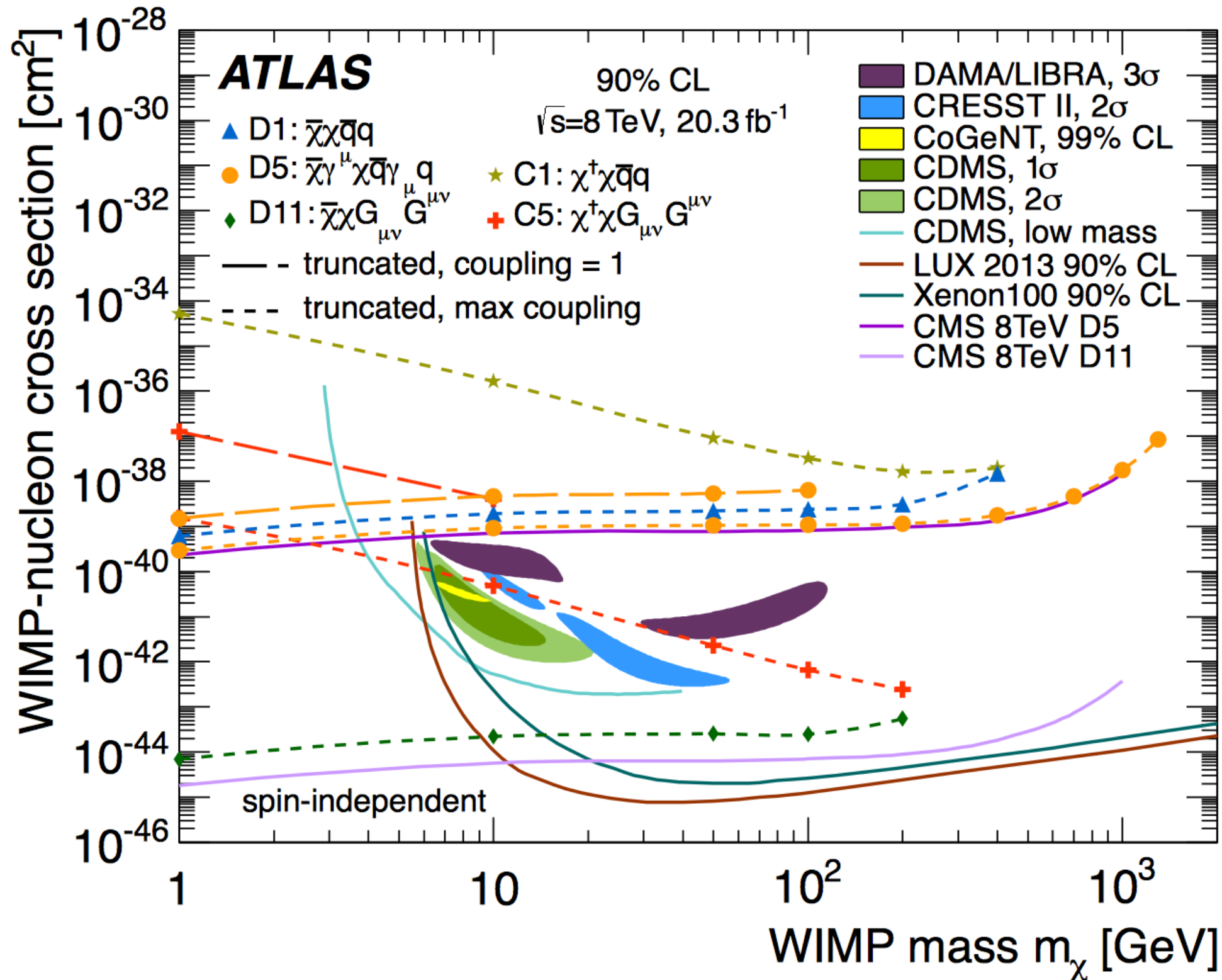
DM @ LHC 2017

All directional efforts!

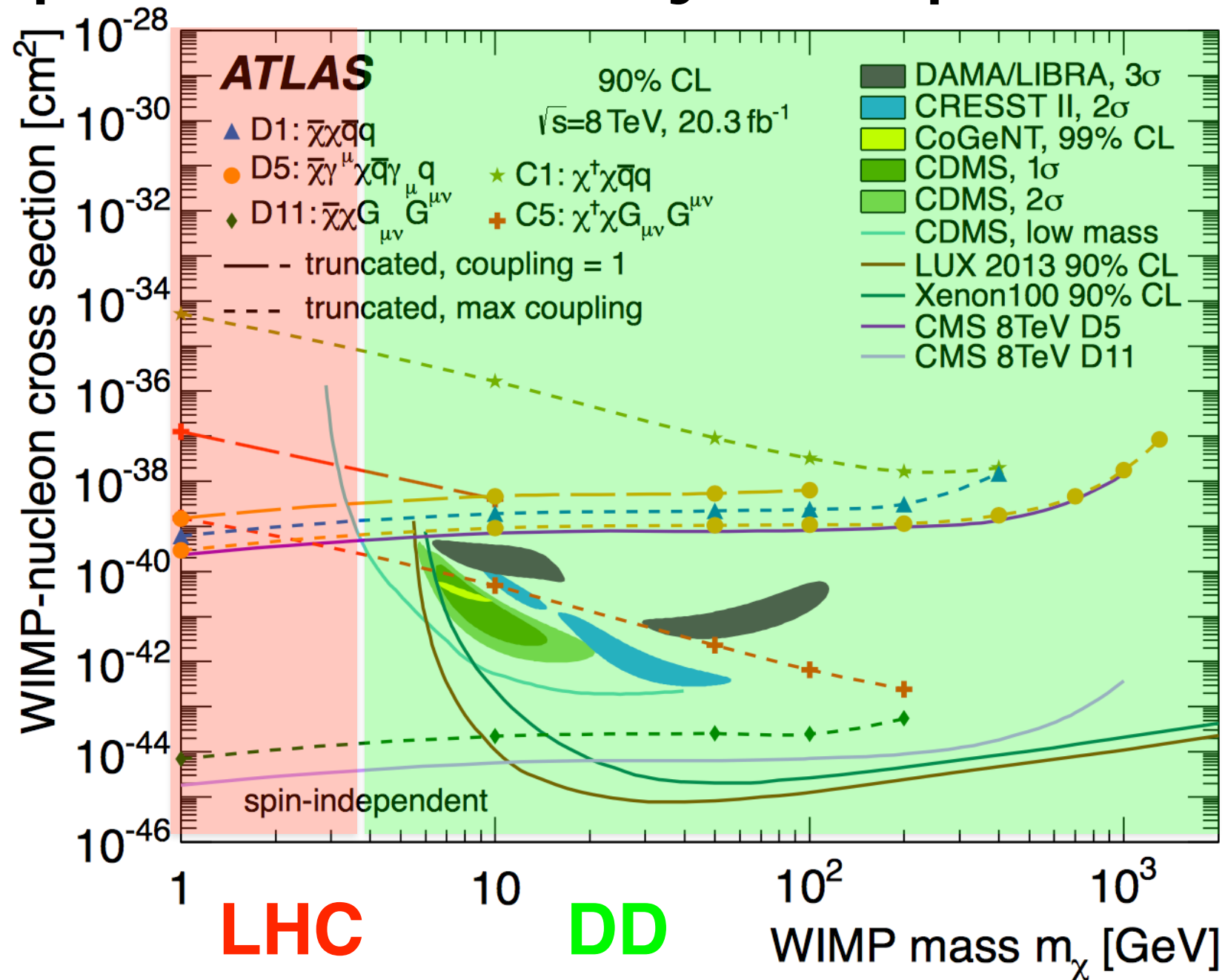


Feng (2008)

Complementary experiments ¹

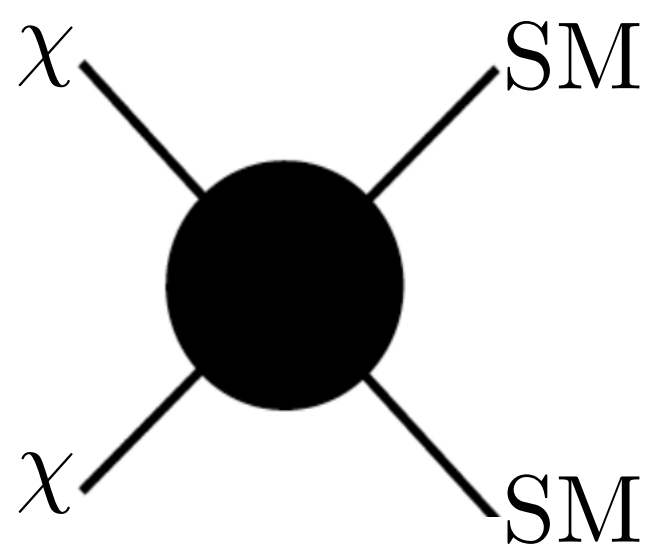


Complementary experiments ²

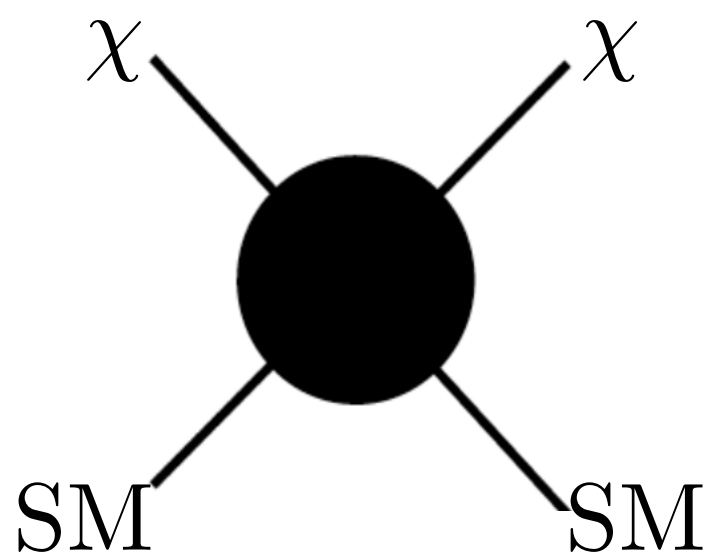


- A collider has a sensitivity in probing a light dark matter particle!
- There is a big difference among Collider / Indirect / Direct exps.

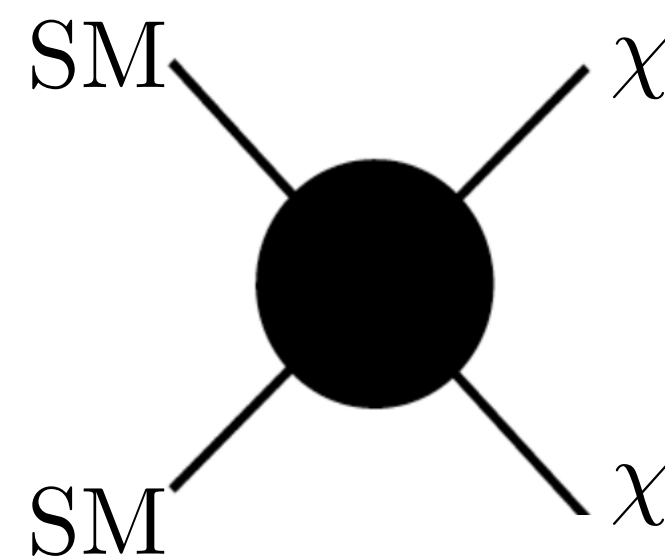
“active”



**Dark matter
Indirect
searches**



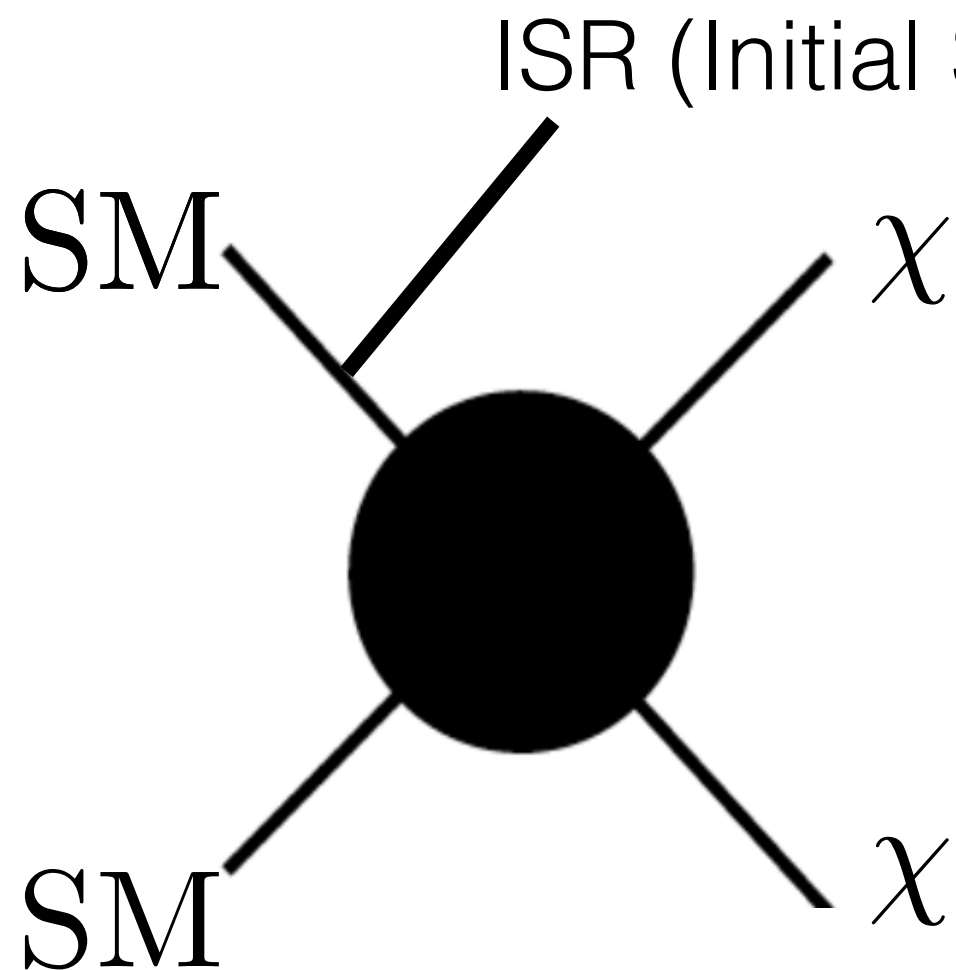
**Dark matter
Direct searches**



Collider

Understanding details

- Once we tag “dark matter” events over backgrounds with ISR, ISR can talk about the details...

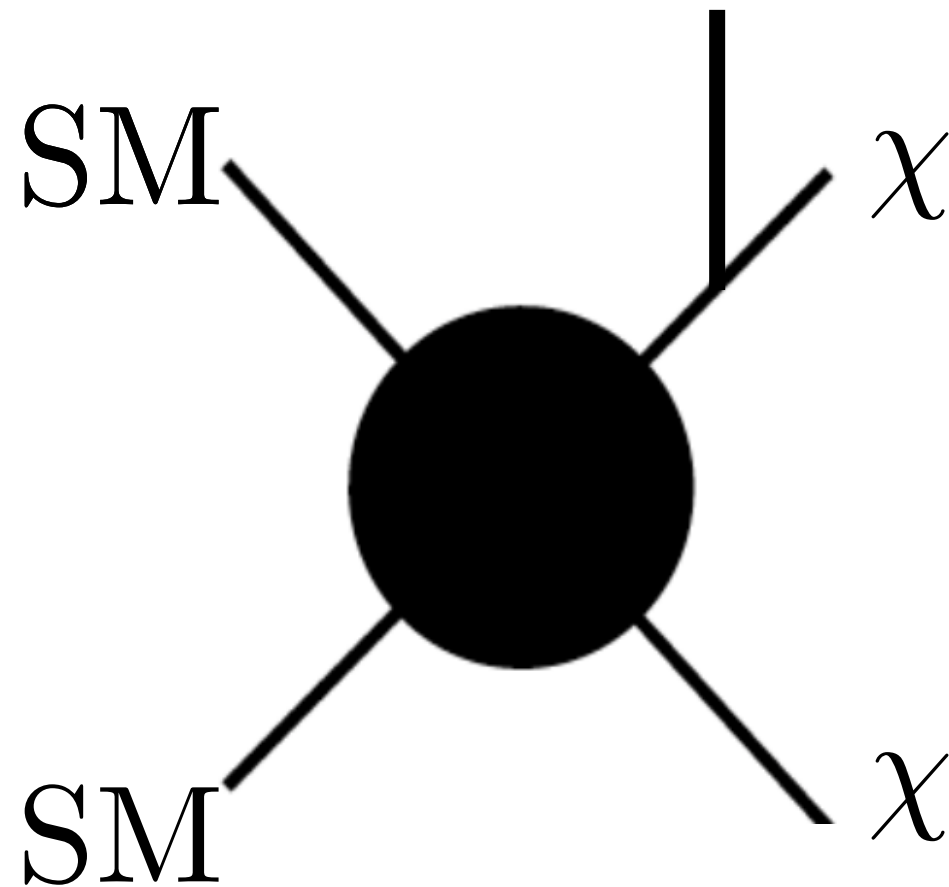


- Determine the properties of dark matter
 - **Mass** spectrum (from kinematics)
 - **Spin** of dark matter / **Interaction** between SM and Dark (from MET, talk [by Alexander Belyaev](#)
~ ISR (photon, jet) PT, talk [by Jonathan Feng](#))

Understanding details

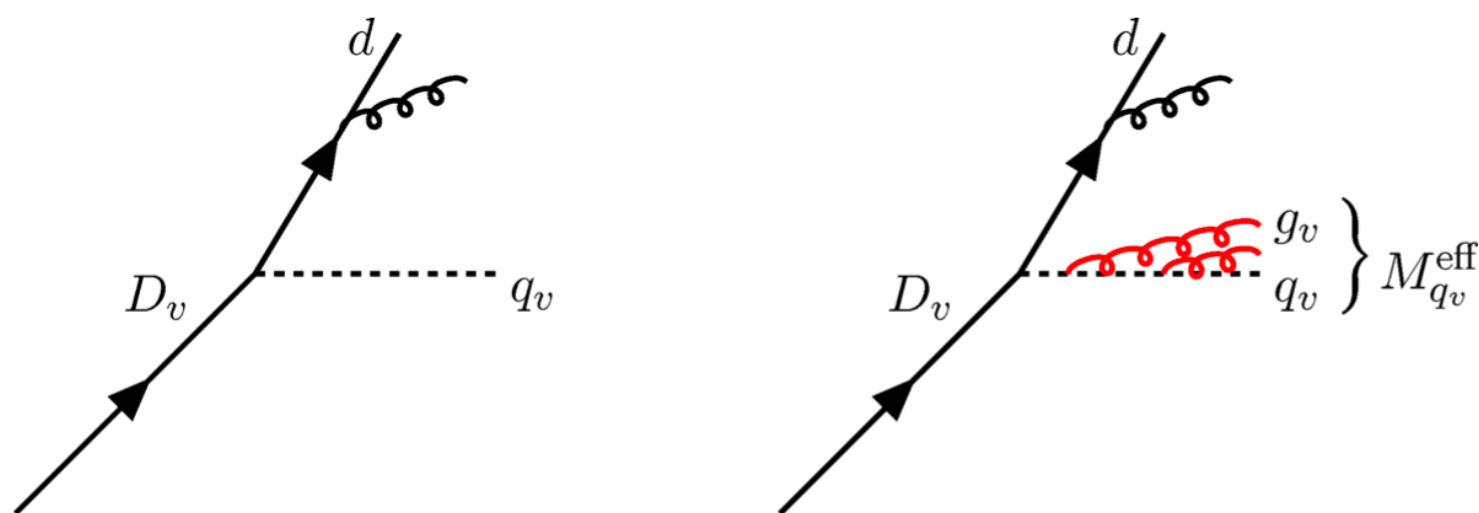
- FSR can talk about a gauge structure of Dark matter!

FSR (Final State Radiation)



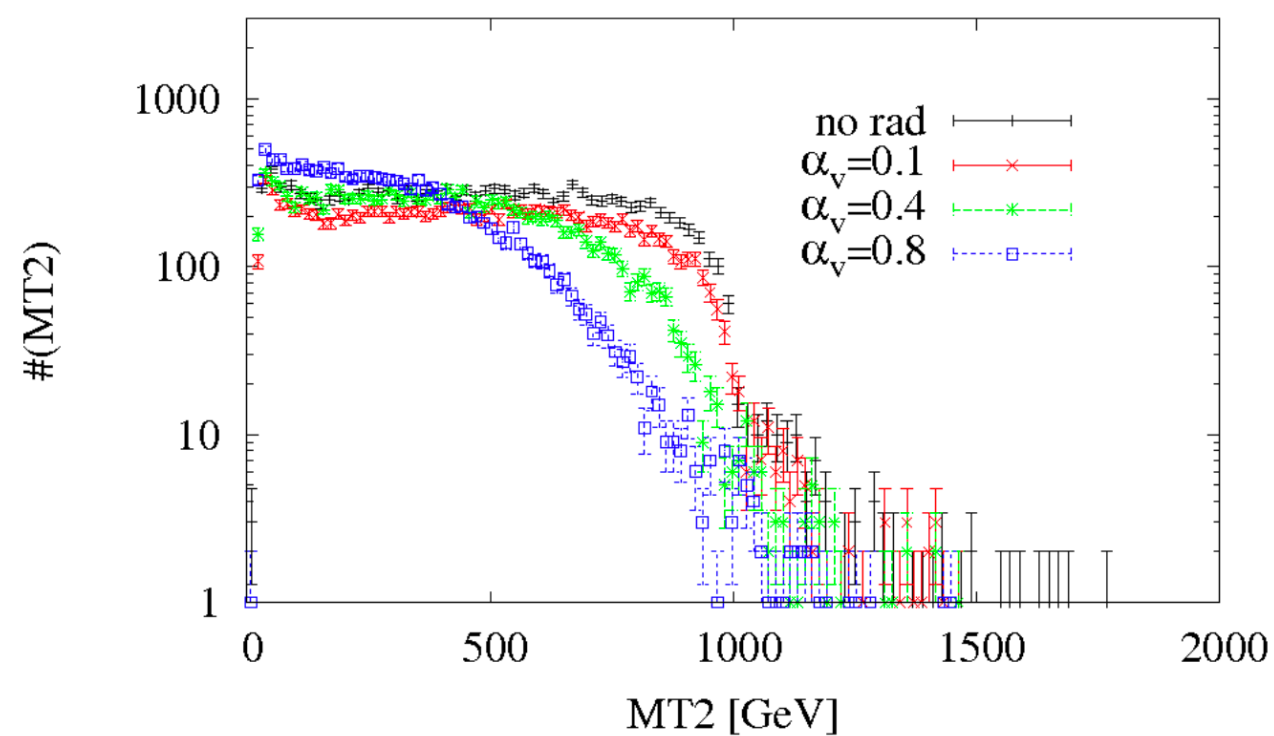
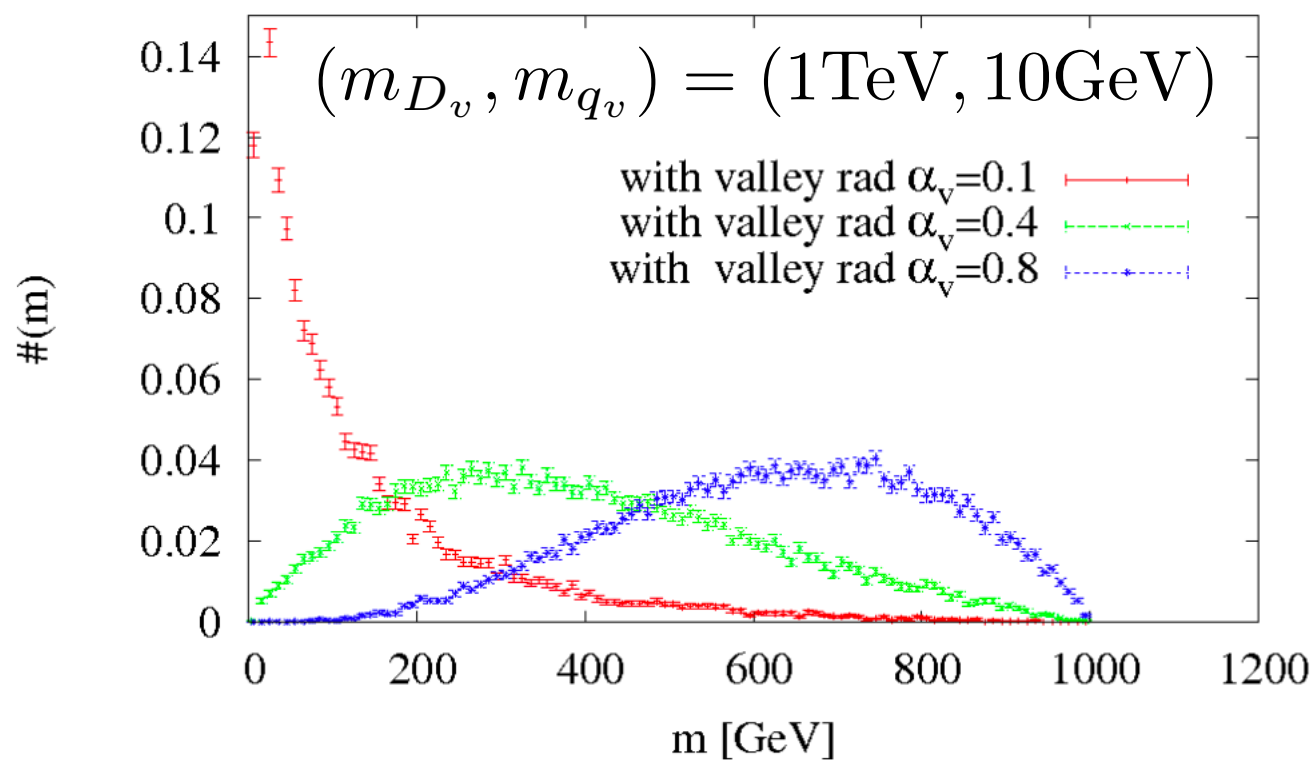
Understanding details

Lisa Carloni, Torbjorn Sjostrand (2010)

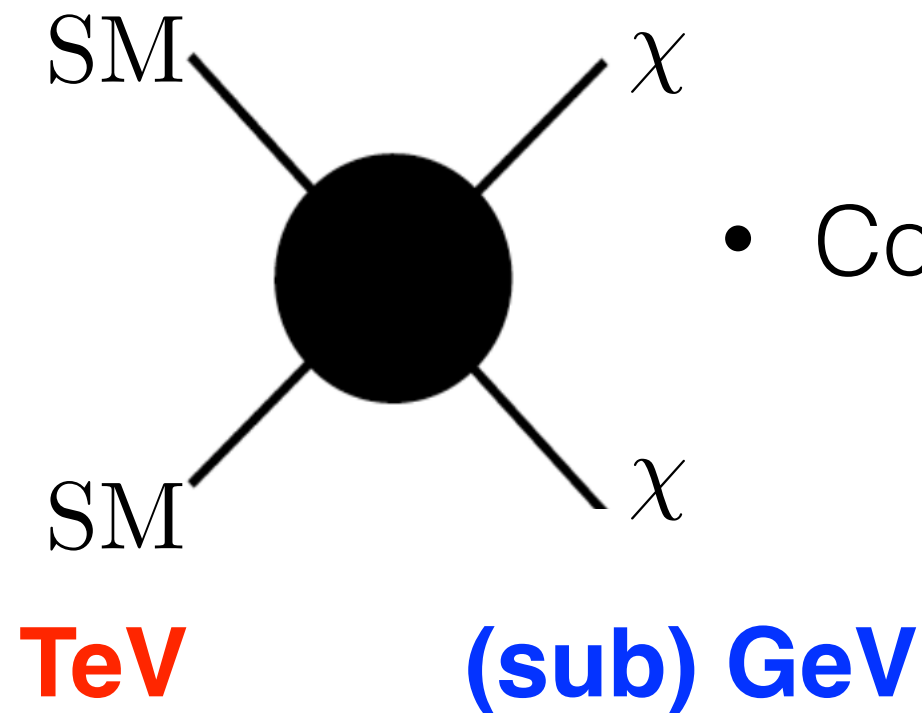


$$SU(3)_c \times SU(3)_v$$

| | | |
|-------|---|---|
| Q_v | 3 | 3 |
| q_v | 1 | 3 |



DM @ Colliders



- Collider can **BOOST** dark matter particles!

- **Radiations** from **BOOSTED** “dark charged” particle will acquire certain level of energy, enough to be “tagged”
(detectable effect)

DM with an abelian charge

- Dark matter may have a dark-U(1) charge
(good to have proper relic, see a talk by Tongyan Lin)
- dark-U(1) can mix with SM U(1)-hyper through a Gauge-kinetic mixing : Gauge-invariant term

$$\mathcal{L} \ni \epsilon F'_{\mu\nu} F_Y^{\mu\nu}$$

- If dark-U(1) is massless (unbroken), then a dark matter can have a milli-charged under SM U(1)
- If dark-U(1) is massive, dark matter would be totally neutral under SM U(1) (Holdom 1986)

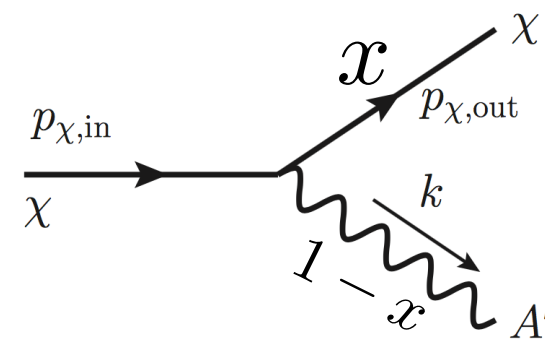
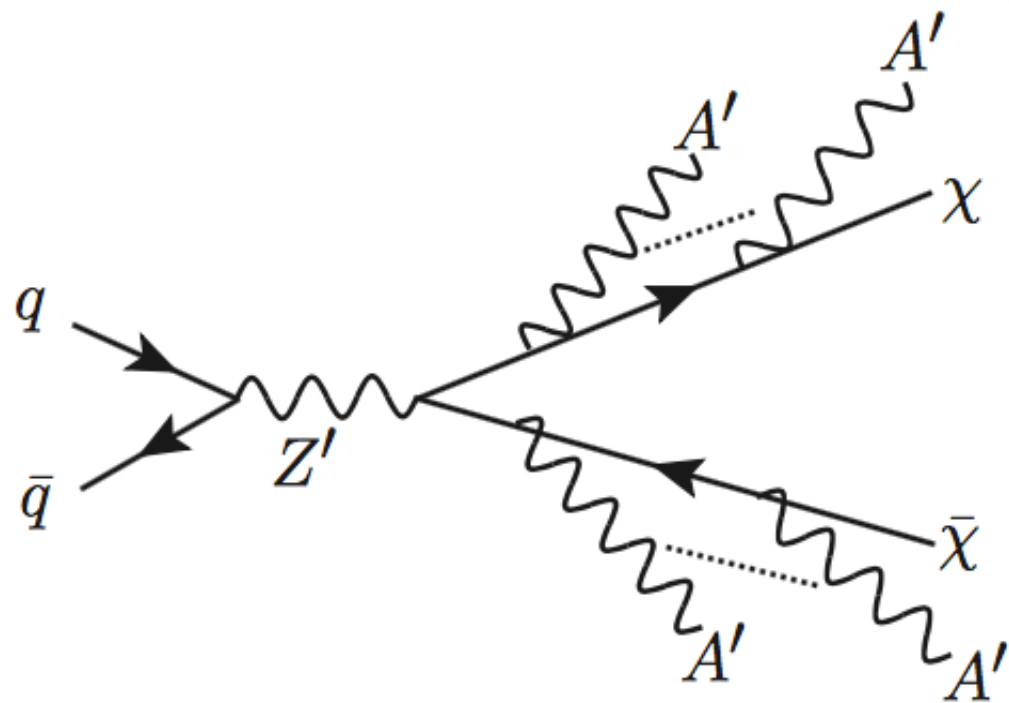
(see also J. Feng, J. Smolingsky and P. Tanedo 2016)

Highly boosted DM@collider

- For a vector-like Dark matter case:

M. Buschmann et.al arXiv:1505.07459

$$\mathcal{L}_{\text{dark}} \equiv \bar{\chi}(i\not{\partial} - m_{\chi} + ig_{A'}\not{A}')\chi - \frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} - \frac{1}{2}m_{A'}^2 A'_{\mu}A'^{\mu} - \frac{\epsilon}{2}F'_{\mu\nu}F^{\mu\nu}$$



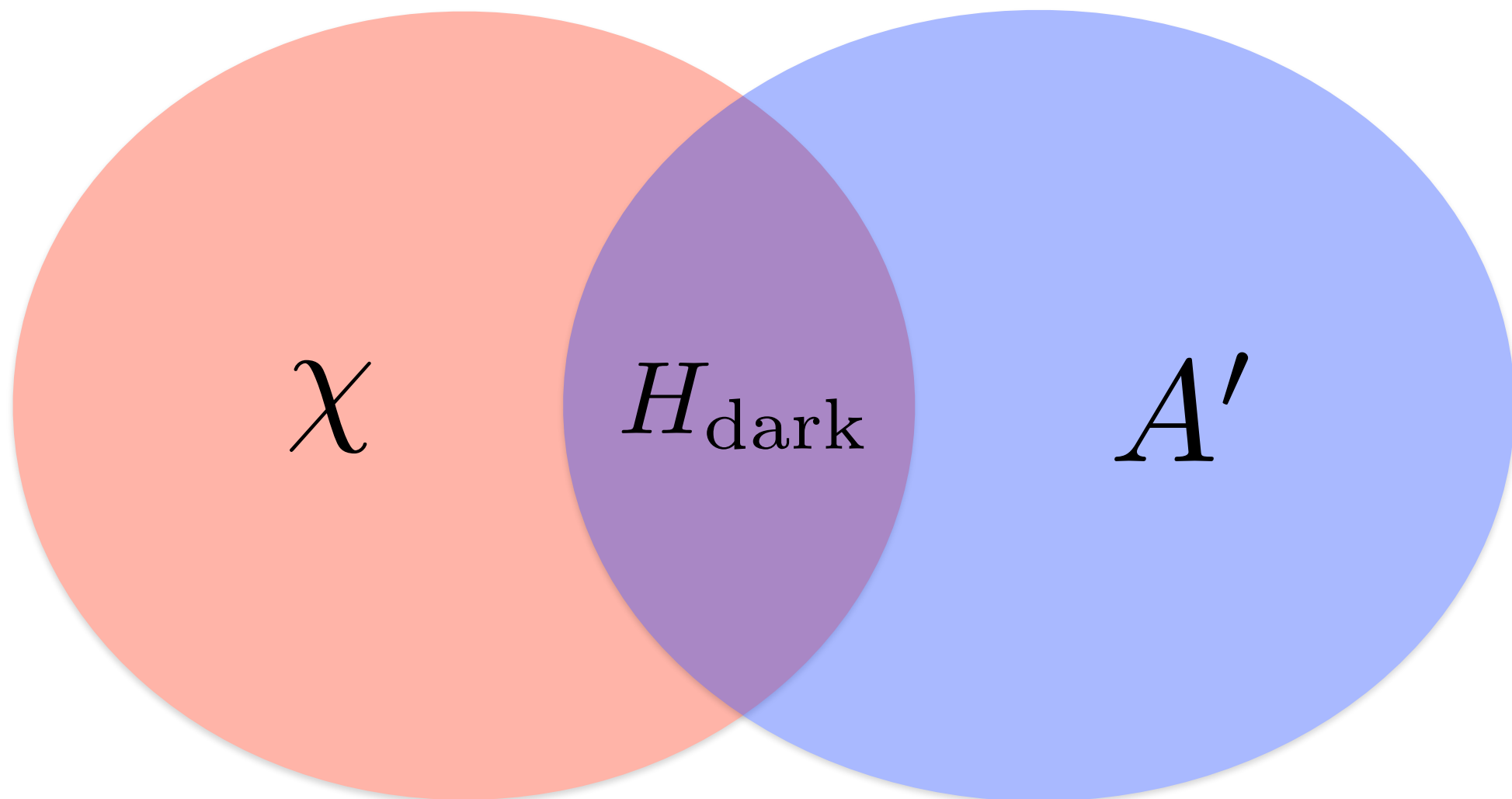
x is an energy fraction

Splitting function

$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq Q_V'^2 \frac{1+x^2}{1-x}$$

What if

a dark matter and a dark-photon
share the same **origin** for their **mass**?



DM, dark gauge boson and a Dark Higgs

| | | | | | |
|----------|-----------|---------------|---------------|---------------|---------------|
| A'_μ | Φ | χ_L | χ_R | ψ_L | ψ_R |
| | Q'_Φ | Q'_{χ_L} | Q'_{χ_R} | Q'_{ψ_L} | Q'_{ψ_R} |

$$Q'_\Phi = Q'_{\chi_R} - Q'_{\chi_L} = -(Q'_{\psi_R} - Q'_{\psi_L})$$

$$\mathcal{L}_{\text{vector+scalar}} = -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{\varepsilon}{2}F_{\mu\nu}F'^{\mu\nu} + |D_\mu\Phi|^2$$

$$\begin{aligned} \mathcal{L}_{\text{matter}} = & \bar{\chi}_L i\gamma^\mu D_\mu \chi_L + \bar{\chi}_R i\gamma^\mu D_\mu \chi_R + \bar{\psi}_L i\gamma^\mu D_\mu \psi_L \\ & + \bar{\psi}_R i\gamma^\mu D_\mu \psi_R - y_\chi \bar{\chi}_L \Phi^* \chi_R - y_\chi \bar{\chi}_R \Phi \chi_L \\ & - y_\psi \bar{\psi}_L \Phi \psi_R - y_\psi \bar{\psi}_R \Phi^* \psi_L \end{aligned}$$

DM, dark gauge boson and a Dark Higgs

| | | | | | |
|----------|-----------|---------------|---------------|---------------|---------------|
| A'_μ | Φ | χ_L | χ_R | ψ_L | ψ_R |
| | Q'_Φ | Q'_{χ_L} | Q'_{χ_R} | Q'_{ψ_L} | Q'_{ψ_R} |

$$\mathcal{L}_{\text{matter}} \ni -g' Q'_V A'_\mu \bar{\chi} \gamma^\mu \chi - g' Q'_A A'_\mu \bar{\chi} \gamma^\mu \gamma_5 \chi$$

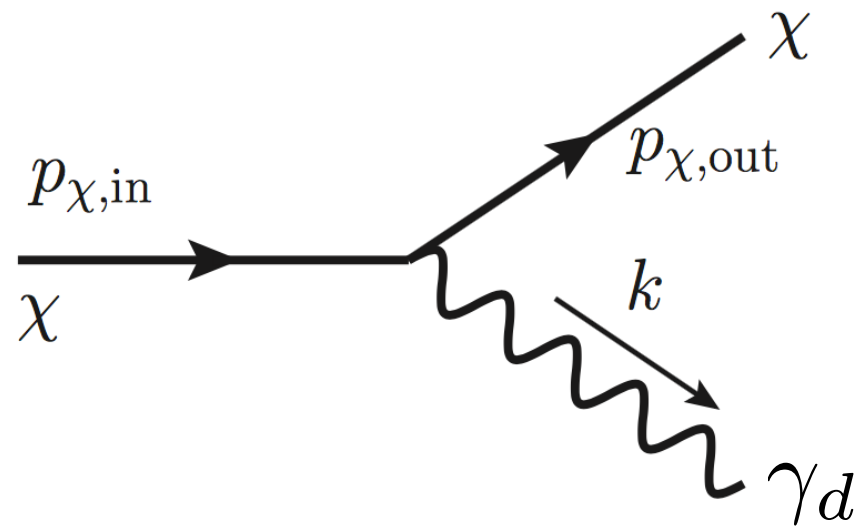
$$Q'_A = \frac{1}{2} (Q'_{\chi_R} - Q'_{\chi_L}) = \frac{Q'_\Phi}{2}$$

$$Q'_V = \frac{1}{2} (Q'_{\chi_R} + Q'_{\chi_L}) = \frac{Q'_\Phi}{2} + Q'_{\chi_L}$$

Thus we always have the axial coupling
between DM and a Dark photon

if a dark photon and dark matter share the origin of mass

Showering process



Splitting function

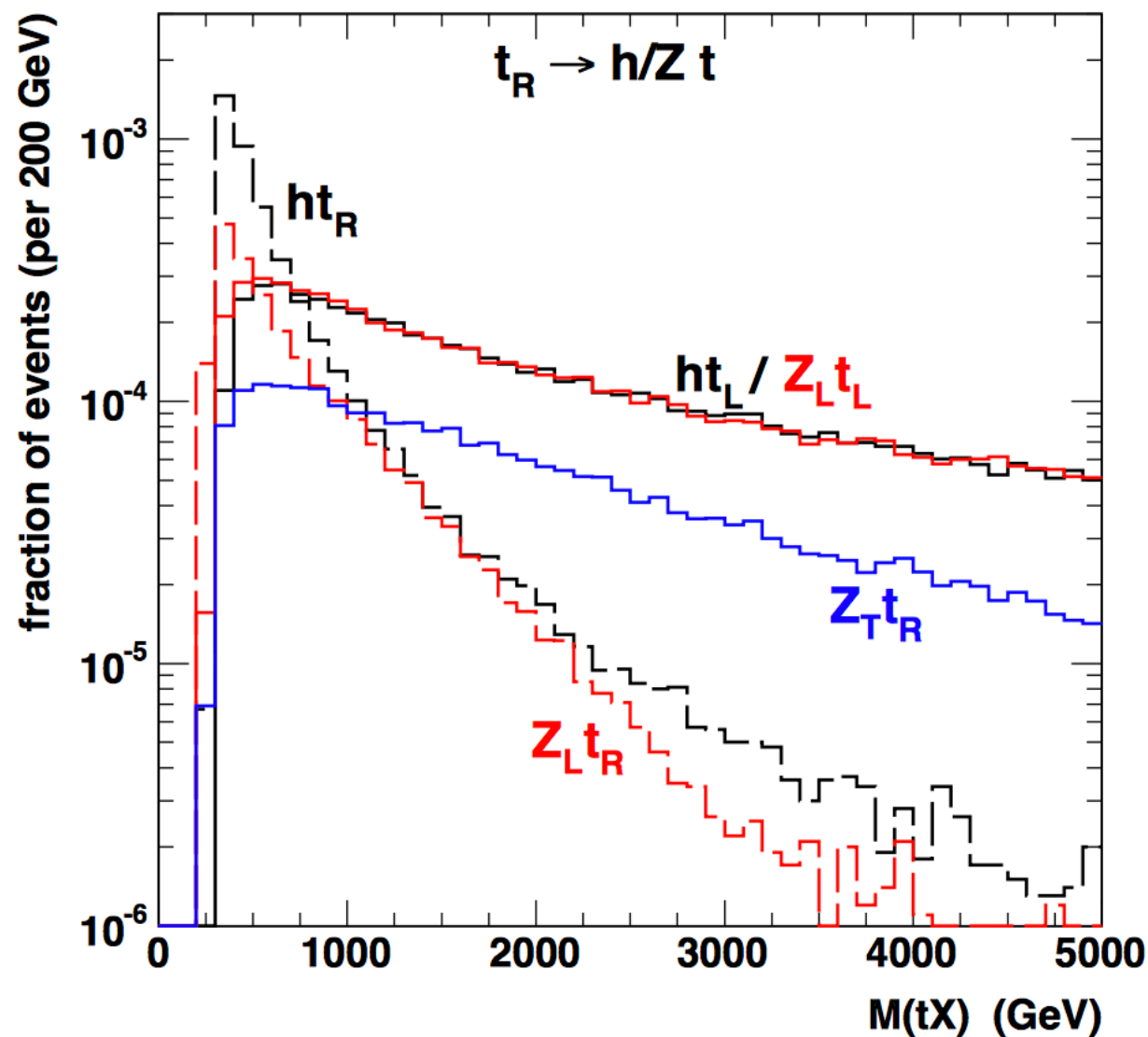
$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq Q_V'^2 \frac{1 + x^2}{1 - x}$$

- In a chiral case, the longitudinal component of a dark photon couples to a dark matter

$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq (Q_V'^2 + Q_A'^2) \frac{1 + x^2}{1 - x} + 2Q_A'^2 \frac{m_\chi^2}{m_{\gamma_d}^2}$$

- We implemented this shower profile in PYTHIA 8

Lesson from SM



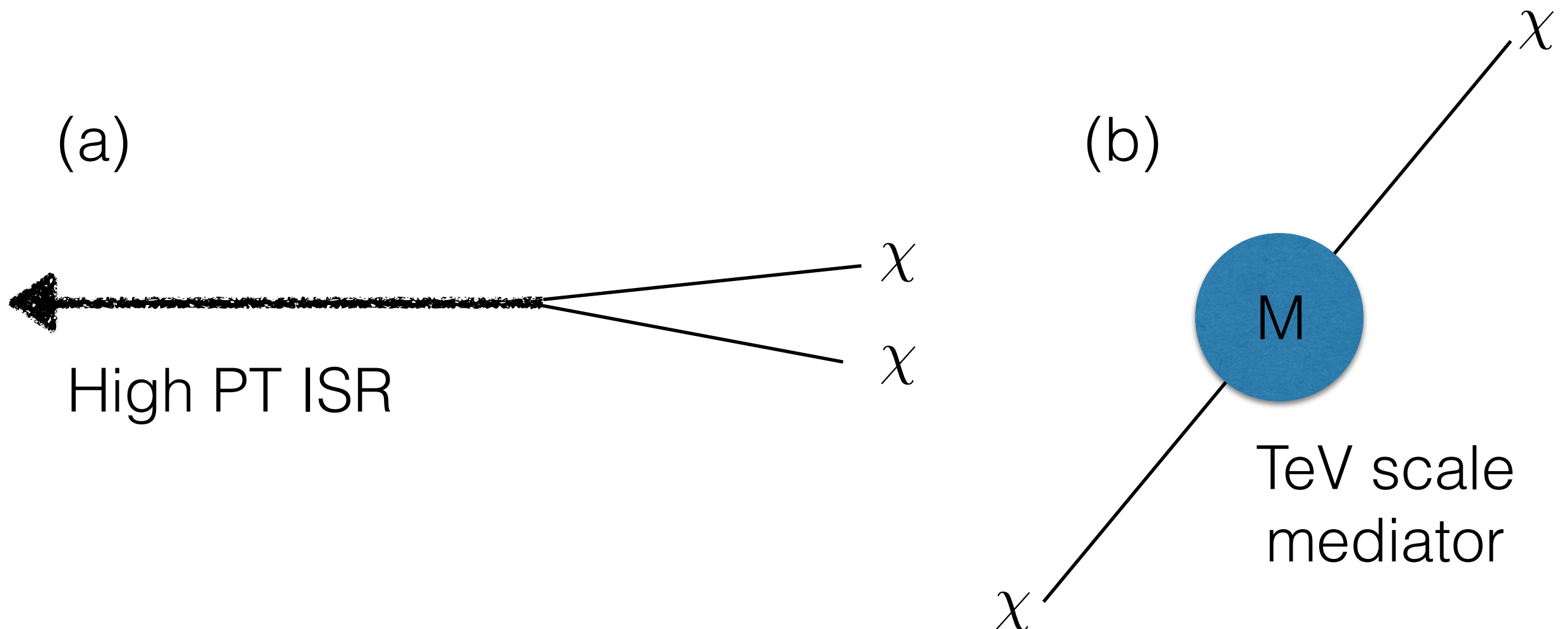
In High Energetic Top-quark case:

Goldston boson Equivalent (GET) show the growth single-logarithmically with energy

$$\mathcal{P}(t_R \rightarrow ht_L) \simeq \mathcal{P}(t_R \rightarrow Z_L t_L)$$

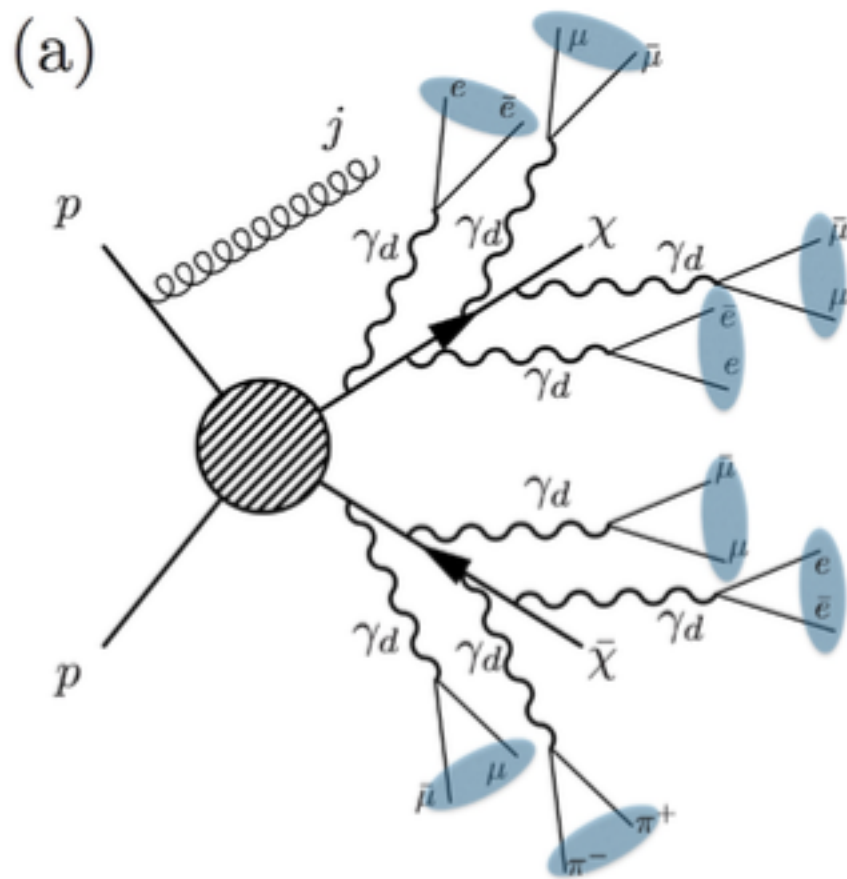
DM production @ collider

- To be more generic, we simulated “boosted” dark matter via
 - (a). Hard recoil from High PT ISR jet
 - (b). Hard back-to-back boost from a heavy mediator

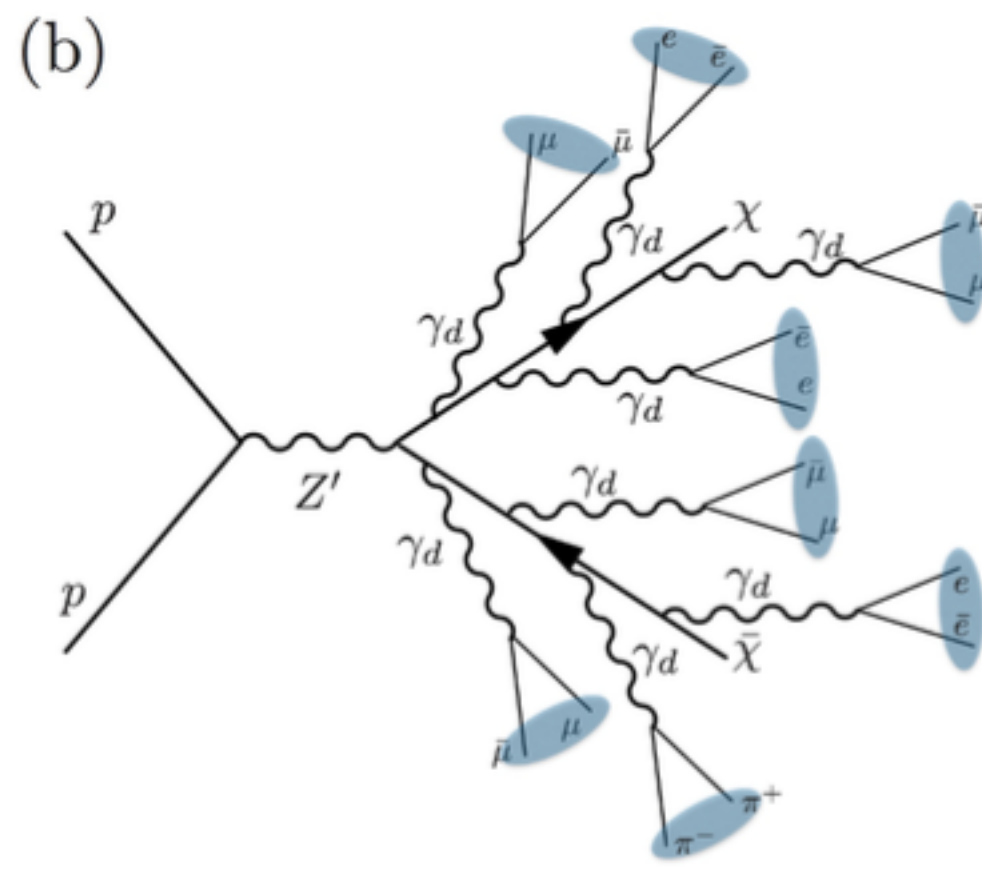


DM production @ collider

- To be more generic, we simulated “boosted” dark matter via
 - Hard recoil from High PT ISR jet
 - Hard back-to-back boost from a heavy mediator



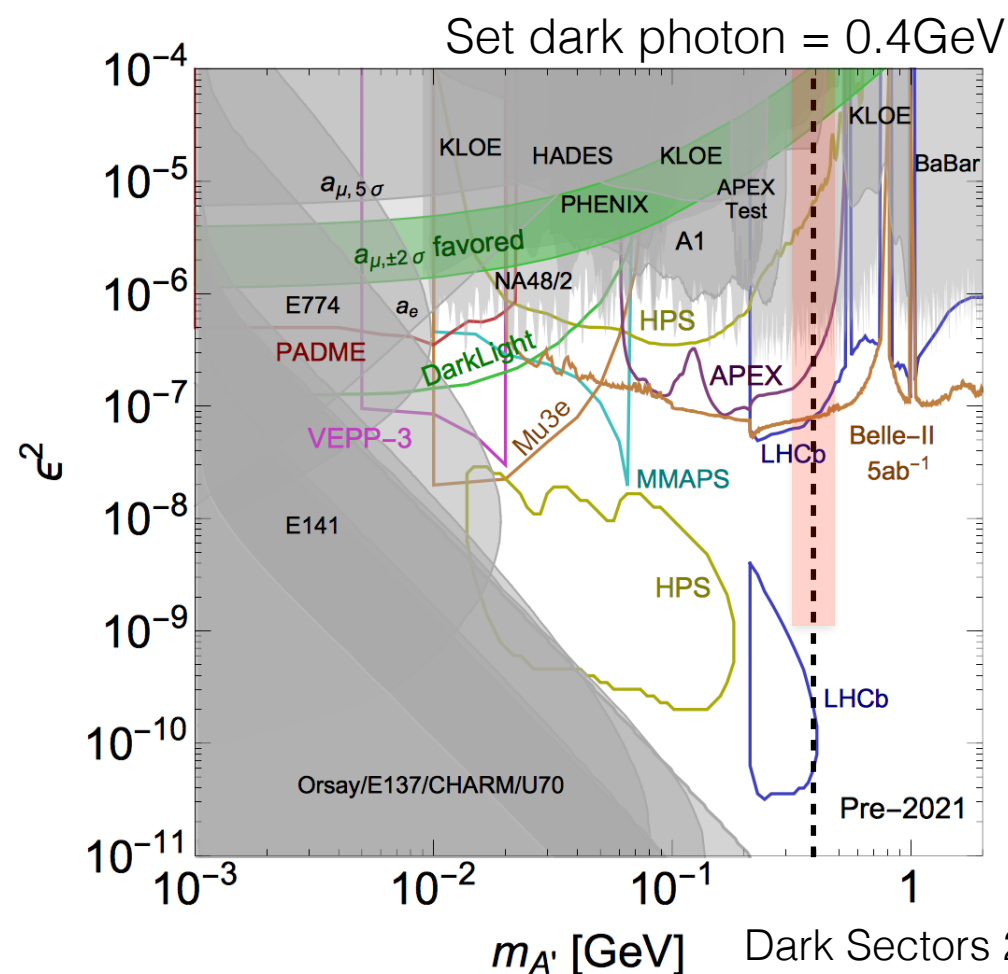
LHC can not produce a mediator directly (Effective operator)



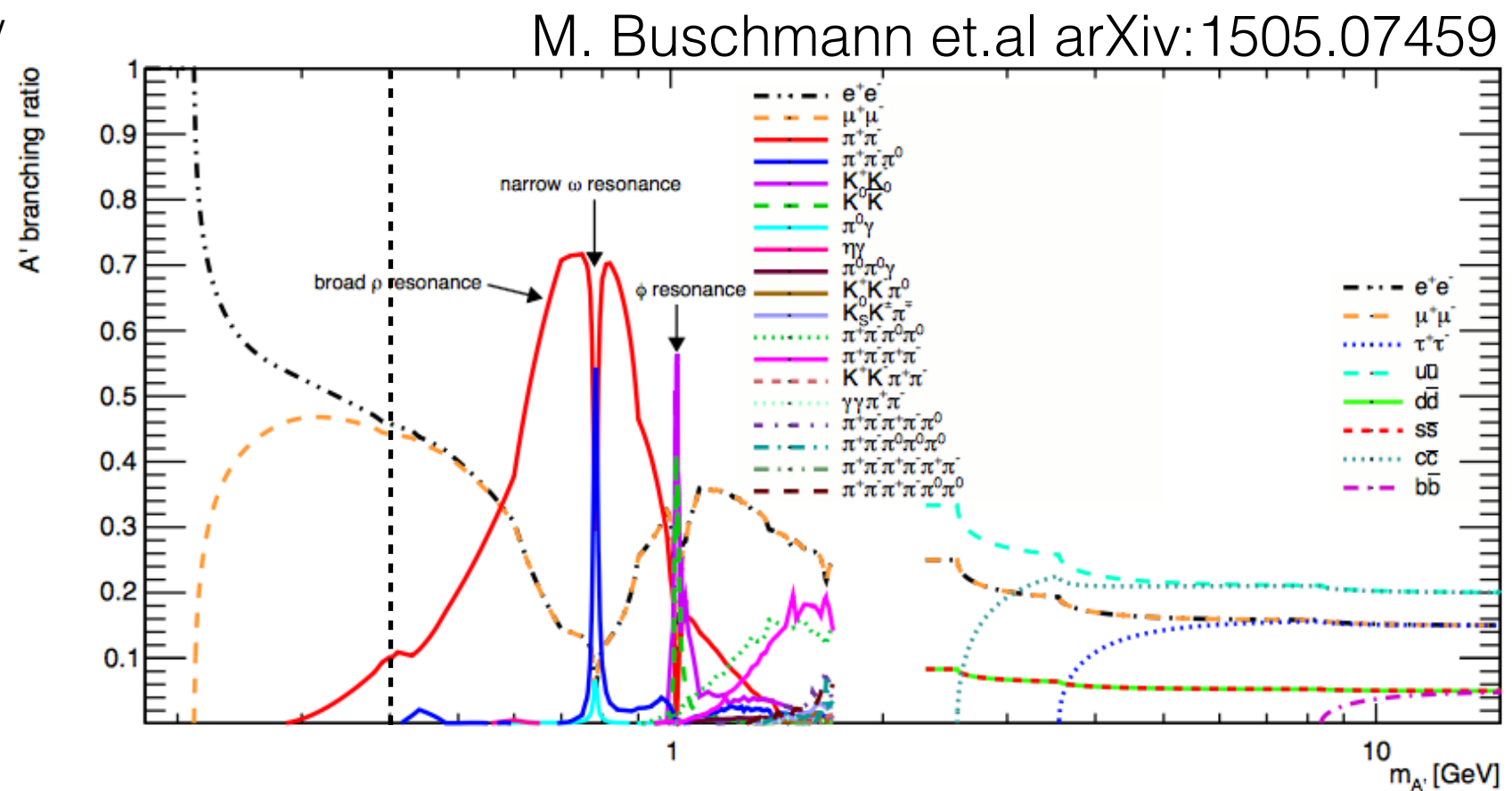
LHC can produce a massive mediator (here Z')

Benchmark points

- We choose a bench mark point for
 - the prompt decays of a dark photon
 - Non-negligible decay mode into muons-pair to tag!



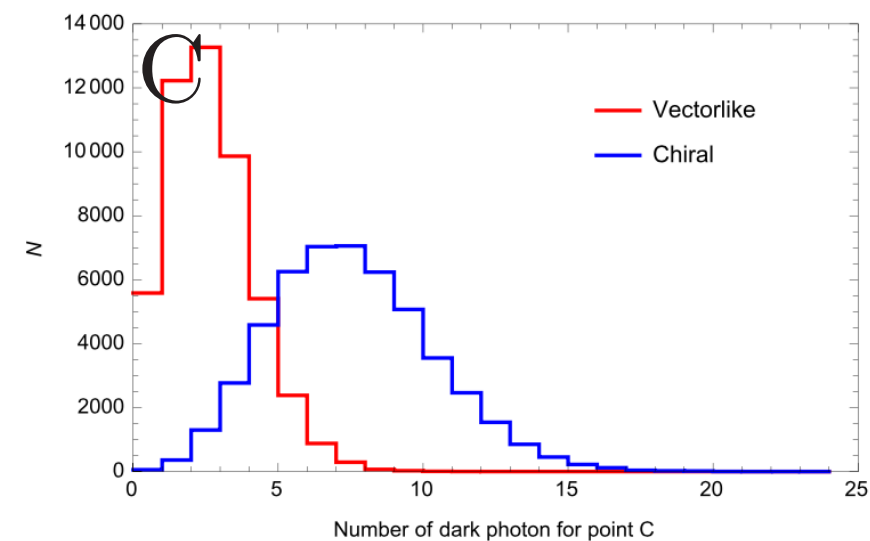
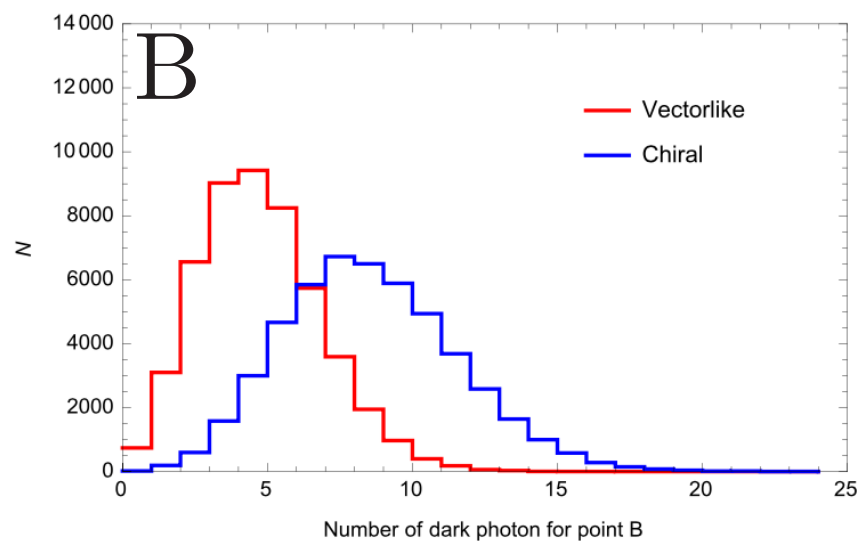
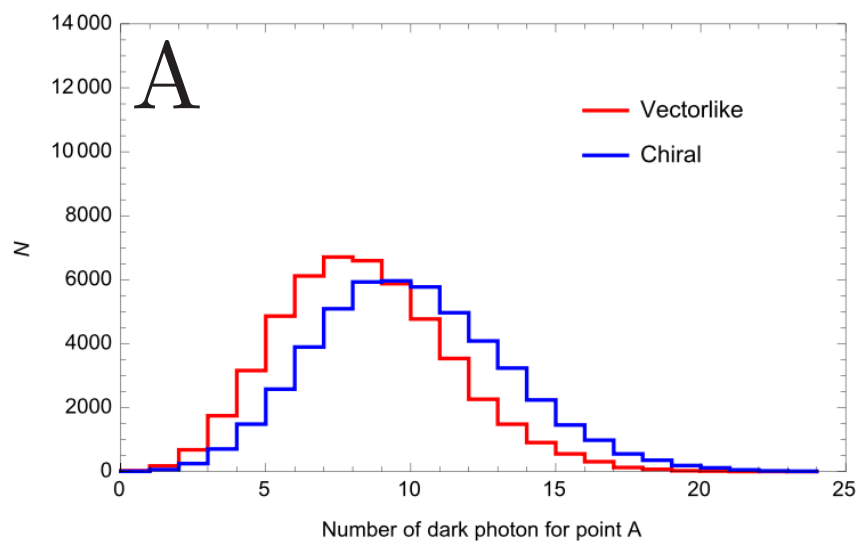
Dark Sectors 2016 Workshop
arXiv:1608.08632



Different showering pattern @ collider

Vector : $(Q'L, Q'R) = (1, 1)$

Chiral: $(Q'L, Q'R) = (1, 0)$



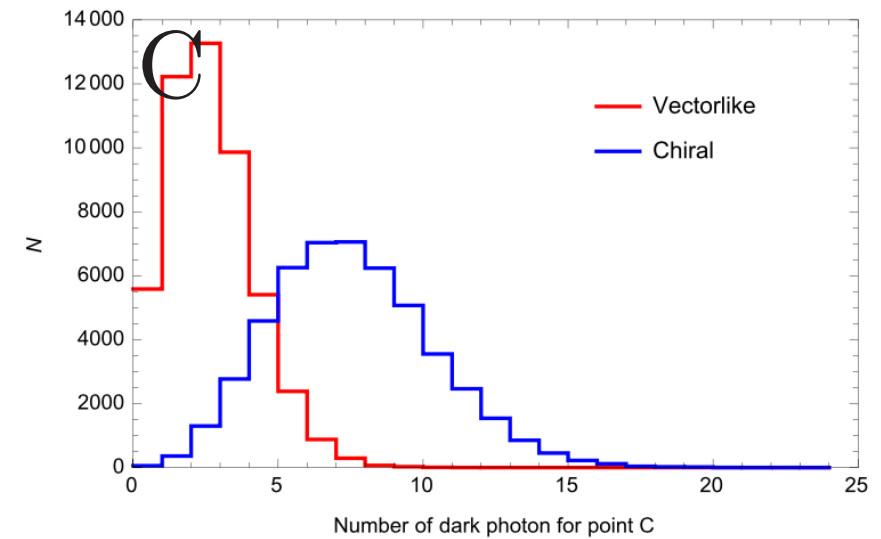
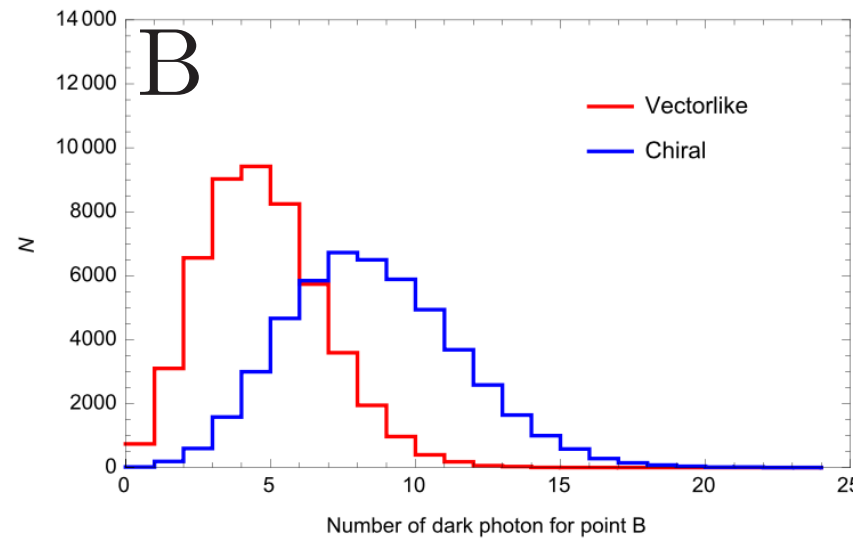
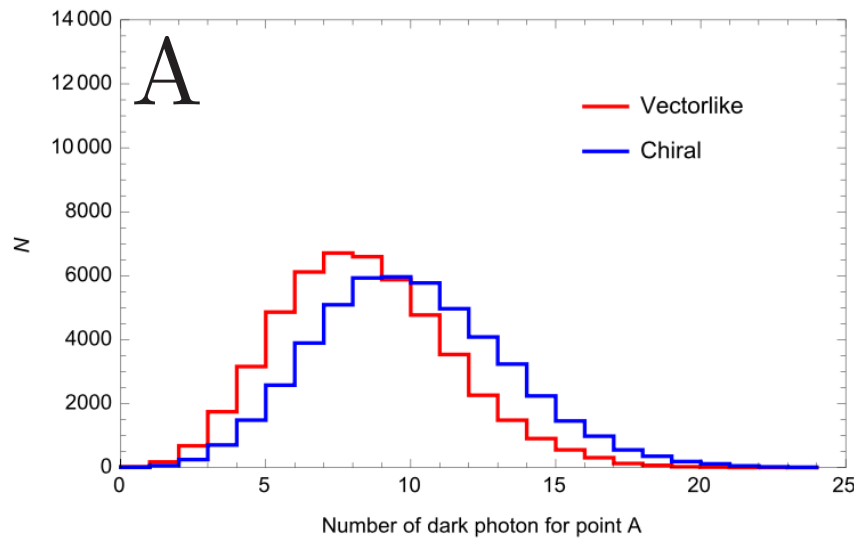
| Benchmark Points (BP) | A | B | C |
|-----------------------|-----|------|-------|
| α' | 0.3 | 0.15 | 0.075 |
| m_χ (GeV) | 0.7 | 1.0 | 1.4 |
| m_{γ_d} (GeV) | | 0.4 | |

$$(y_\chi/\sqrt{2})^2 \lesssim 4\pi \quad \leftarrow \quad \begin{aligned} m_{\gamma_d} &= g' Q'_\phi v_s \\ m_\chi &= y_\chi v_s / \sqrt{2} \end{aligned}$$

TABLE I. Benchmark points we have chosen. They obey the perturbative limit of $\alpha' \frac{m_\chi^2}{m_{\gamma_d}^2} \lesssim 1$.

Different showering pattern @ collider

Vector : $(Q'L, Q'R) = (1, 1)$
 Chiral: $(Q'L, Q'R) = (1, 0)$



| Benchmark Points (BP) | A | B | C |
|-----------------------|-----|------|-------|
| α' | 0.3 | 0.15 | 0.075 |
| m_χ (GeV) | 0.7 | 1.0 | 1.4 |
| m_{γ_d} (GeV) | | 0.4 | |

@LHC, we may see the difference among various mechanism behind the mass of dark matter & a dark-photon

$-(\alpha' \ll 1)$ limits

$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq (Q_V'^2 + Q_A'^2) \frac{1+x^2}{1-x} + 2Q_A'^2 \frac{m_\chi^2}{m_{\gamma_d}^2}$$

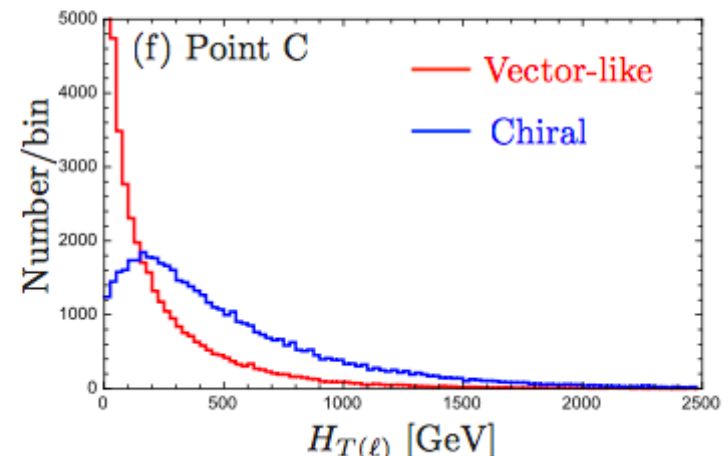
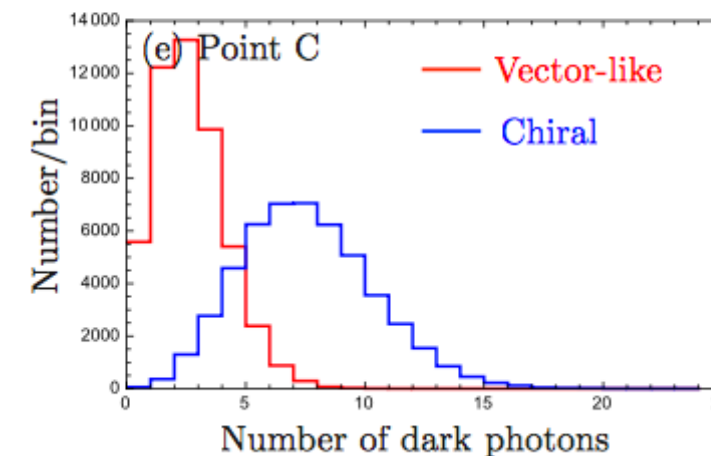
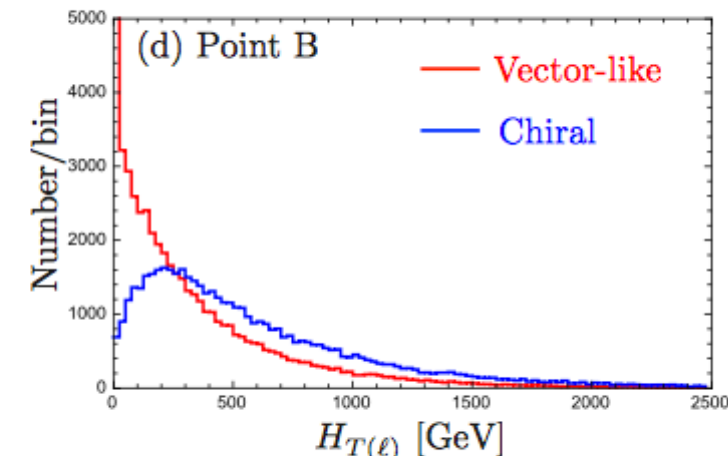
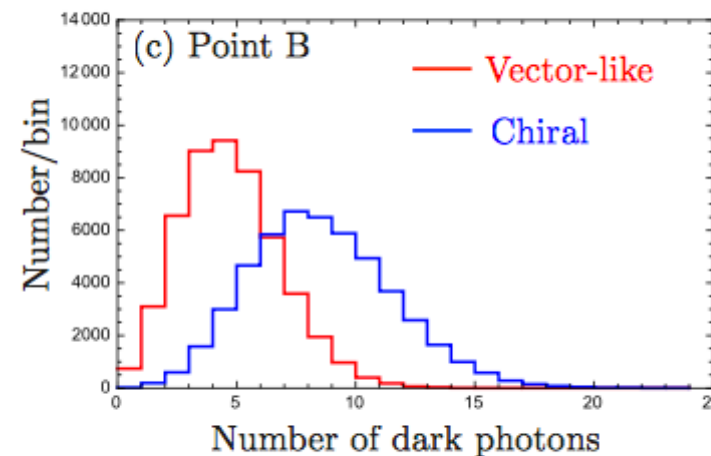
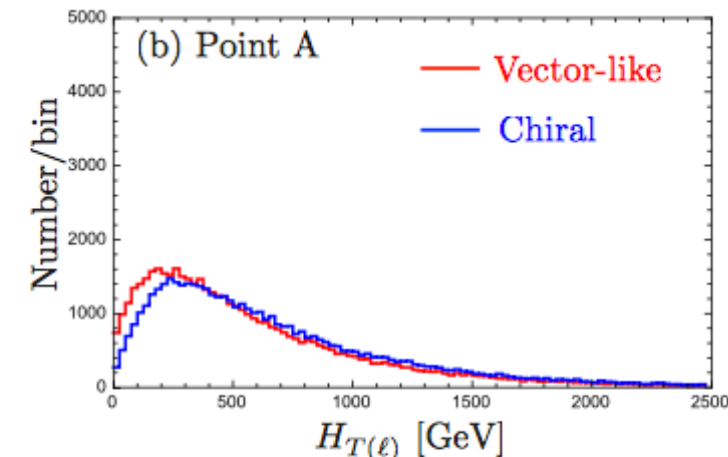
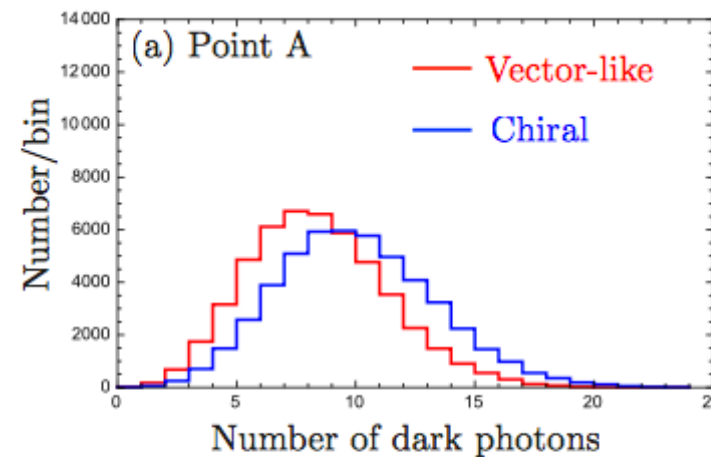
$$\alpha' P_{\chi \rightarrow \chi \gamma_d} \sim \frac{m_{\gamma_d}^2}{v_s^2} \frac{m_\chi^2}{m_{\gamma_d}^2} \sim y_\chi^2$$

Quantify the difference

- We use the transverse energy deposits from leptonic decay modes

| Benchmark Points (BP) | A | B | C |
|-----------------------|-----|------|-------|
| α' | 0.3 | 0.15 | 0.075 |
| m_χ (GeV) | 0.7 | 1.0 | 1.4 |
| m_{γ_d} (GeV) | | 0.4 | |

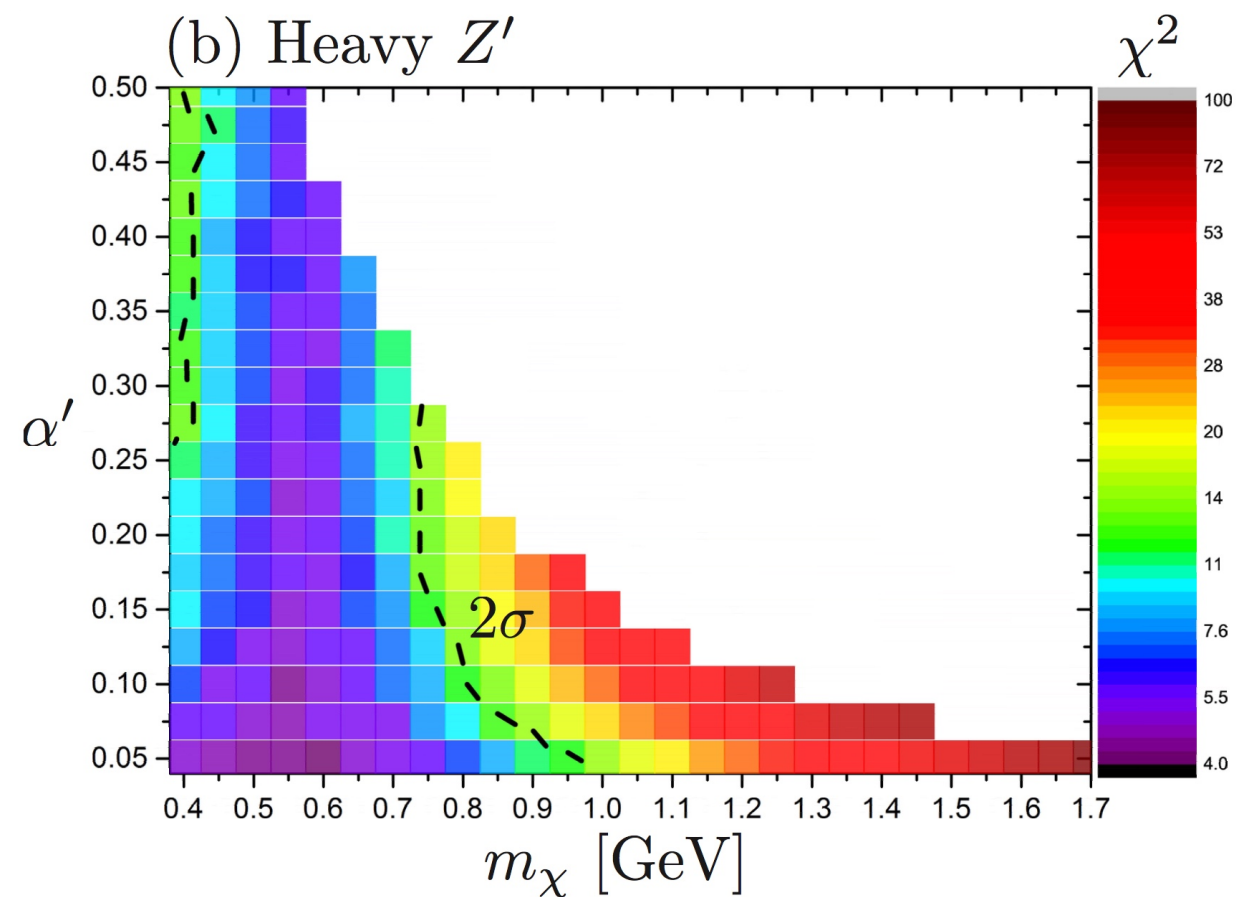
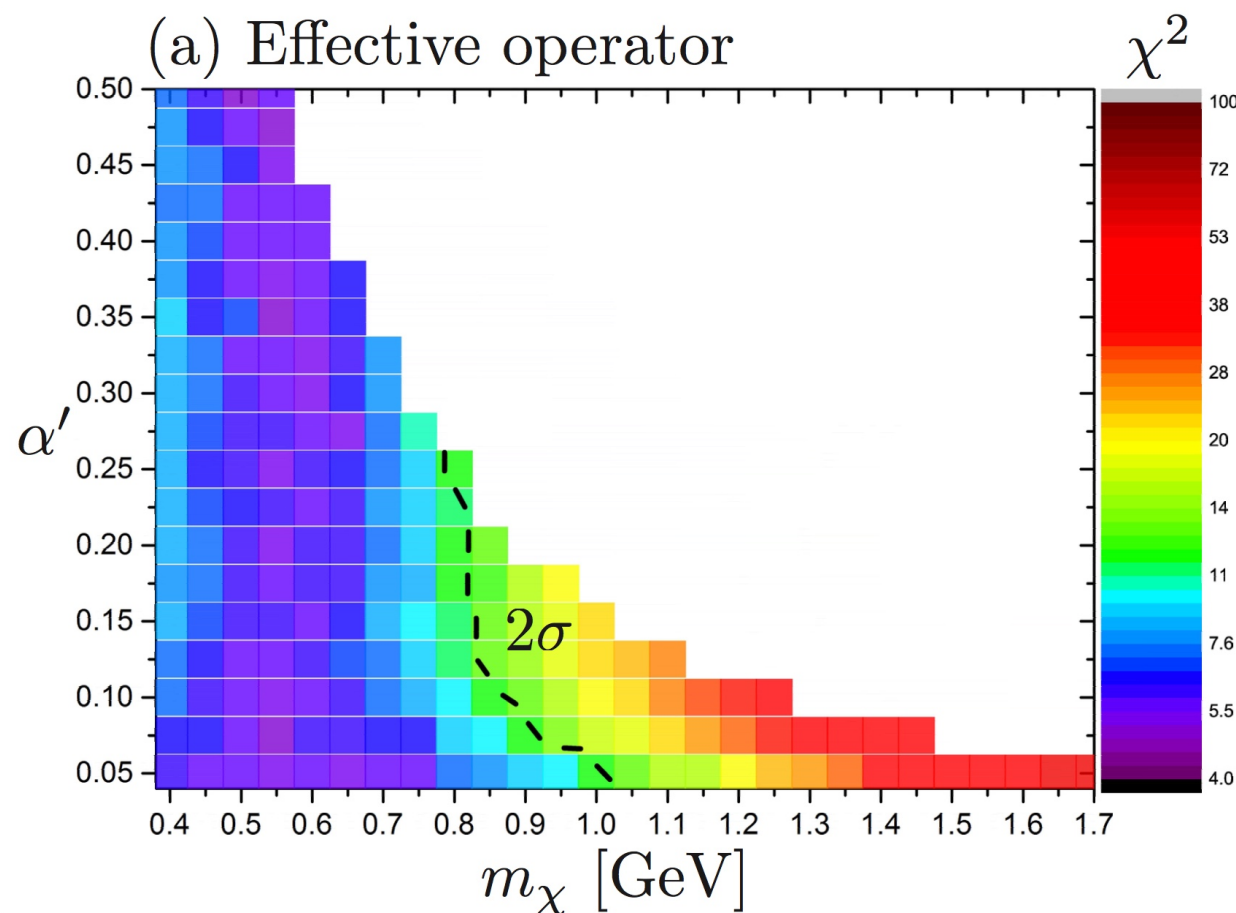
$$H_{T(\ell)} = \sum_{i=\mu^\pm, e^\pm} |p_{Ti}|$$



-Due to GBET, the energy spectrum of leptons from a **longitudinal mode** is larger compared to the case of leptons from a transverse mode

Checking chirality@ Collider

- After triggering signal events by tagging a collimated muon-jet (a jet only with muons)



with 200 signal events after cuts to reduce BKG
 (As BKG does not interfere with signal,
 we can subtract BKG distribution)

Conclusion

- Collider is an **active** experiment
 - to find dark sector (dark matter).
 - to measure properties of a dark sector.
- The **mass origin** in dark sector (like SM-Higgs mechanism) can strongly affect the dark photon showering in “**boosted**” dark matter.
- Collider can probe the nature of dark matter by examine the **pattern** of dark photon **showering!**