

EW corrections for V+jets

Jonas M. Lindert

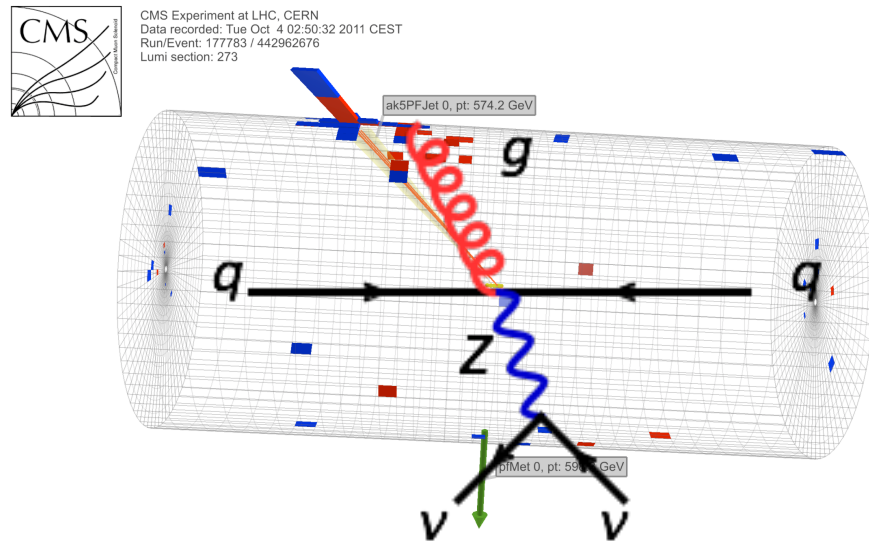


work in collaboration with:

*R. Boughezal, A. Denner, S. Dittmaier, A. Huss, A. Gehrmann-De Ridder,
T. Gehrmann, N. Glover, S. Kallweit, P. Maierhöfer, M. L. Mangano,
T.A. Morgan, A. Mück, M. Schönherr, F. Petriello, S. Pozzorini, G. P. Salam*

illuminating standard candles at the LHC - V+jets
Imperial, London, 25.04.2017

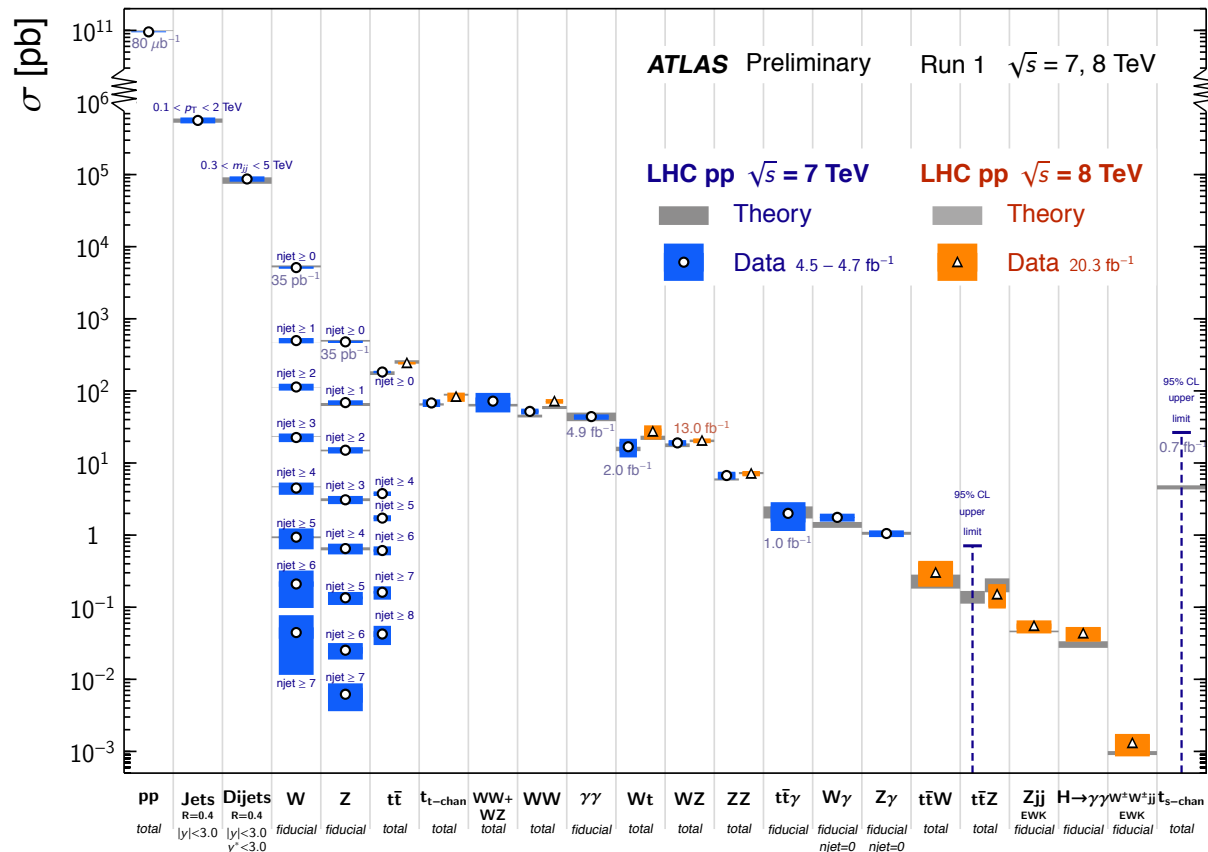
V + multijet production



- ▶ Dominant backgrounds for monojet **DM searches**
- ▶ Important/dominant backgrounds for various **BSM searches** (lepton + missing E_T + ets)
- ▶ Dominant backgrounds for **top physics**
- ▶ Dominant backgrounds for **Higgs physics**, e.g. $VH(\rightarrow bb)$, $H\rightarrow WW$

Standard Model Production Cross Section Measurements

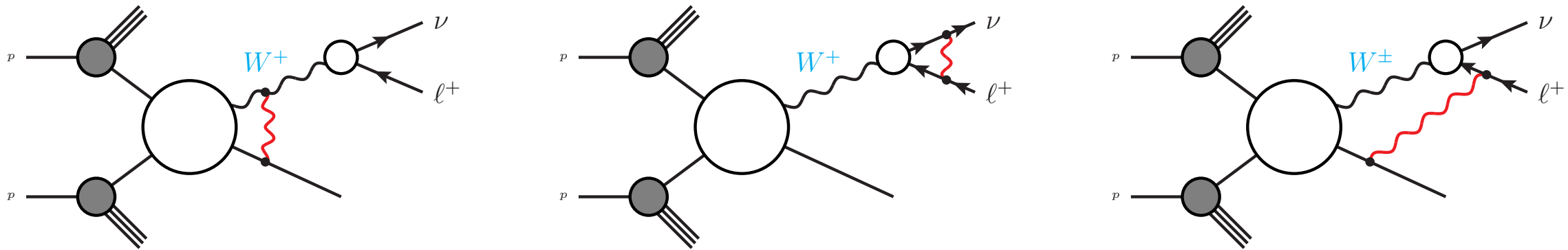
Status: July 2014



- ▶ Large cross-sections and clean leptonic signatures
- ▶ **V+jets**: Precision **QCD** at LHC
- ▶ Playground to probe different aspects of higher-order calculations (LO+PS, NLO+PS, NLO-Merging, NLO EW,...)
- ▶ Probe and constrain PDFs

Decays of heavy particles @ NLO EW

- ▶ Leptonic decays of gauge bosons are trivial at NLO QCD. At NLO EW corrections in production, decay and non-factorizable contributions have to be considered.



- ▶ Scheme of choice: **complex-mass-scheme** [Denner, Dittmaier]
 - gauge invariant and exact NLO
 - **computationally very expensive**: one extra leg per two-body decay
- ▶ Pragmatic choice: **Narrow-width-approximation (NWA)**
 - gauge invariant in strict on-shell limit of NWA
 - **allows to capture all Sudakov effects** (not present in decay)
 - **allows to go to higher jet multiplicities**
 - not applicable to all processes at all perturbative orders

Combination of NLO QCD and EW & Setup

Two alternatives:

$$\sigma_{\text{QCD+EW}}^{\text{NLO}} = \sigma^{\text{LO}} + \delta\sigma_{\text{QCD}}^{\text{NLO}} + \delta\sigma_{\text{EW}}^{\text{NLO}}$$

$$\sigma_{\text{QCD}\times\text{EW}}^{\text{NLO}} = \sigma_{\text{QCD}}^{\text{NLO}} \left(1 + \frac{\delta\sigma_{\text{EW}}^{\text{NLO}}}{\sigma^{\text{LO}}} \right) = \sigma_{\text{EW}}^{\text{NLO}} \left(1 + \frac{\delta\sigma_{\text{QCD}}^{\text{NLO}}}{\sigma^{\text{LO}}} \right)$$

Difference between the two approaches indicates uncertainties due to missing EW-QCD corrections of $\mathcal{O}(\alpha\alpha_s)$

Relative corrections w.r.t. NLO QCD:

$$\frac{\sigma_{\text{QCD+EW}}^{\text{NLO}}}{\sigma_{\text{QCD}}^{\text{NLO}}} = \left(1 + \frac{\delta\sigma_{\text{EW}}^{\text{NLO}}}{\sigma_{\text{QCD}}^{\text{NLO}}} \right)$$

suppressed by large NLO QCD corrections

$$\frac{\sigma_{\text{QCD}\times\text{EW}}^{\text{NLO}}}{\sigma_{\text{QCD}}^{\text{NLO}}} = \left(1 + \frac{\delta\sigma_{\text{EW}}^{\text{NLO}}}{\sigma^{\text{LO}}} \right)$$

“usual” NLO EW w.r.t. LO

V + l jet

$l\nu + 1 \text{ jet}$: inclusive

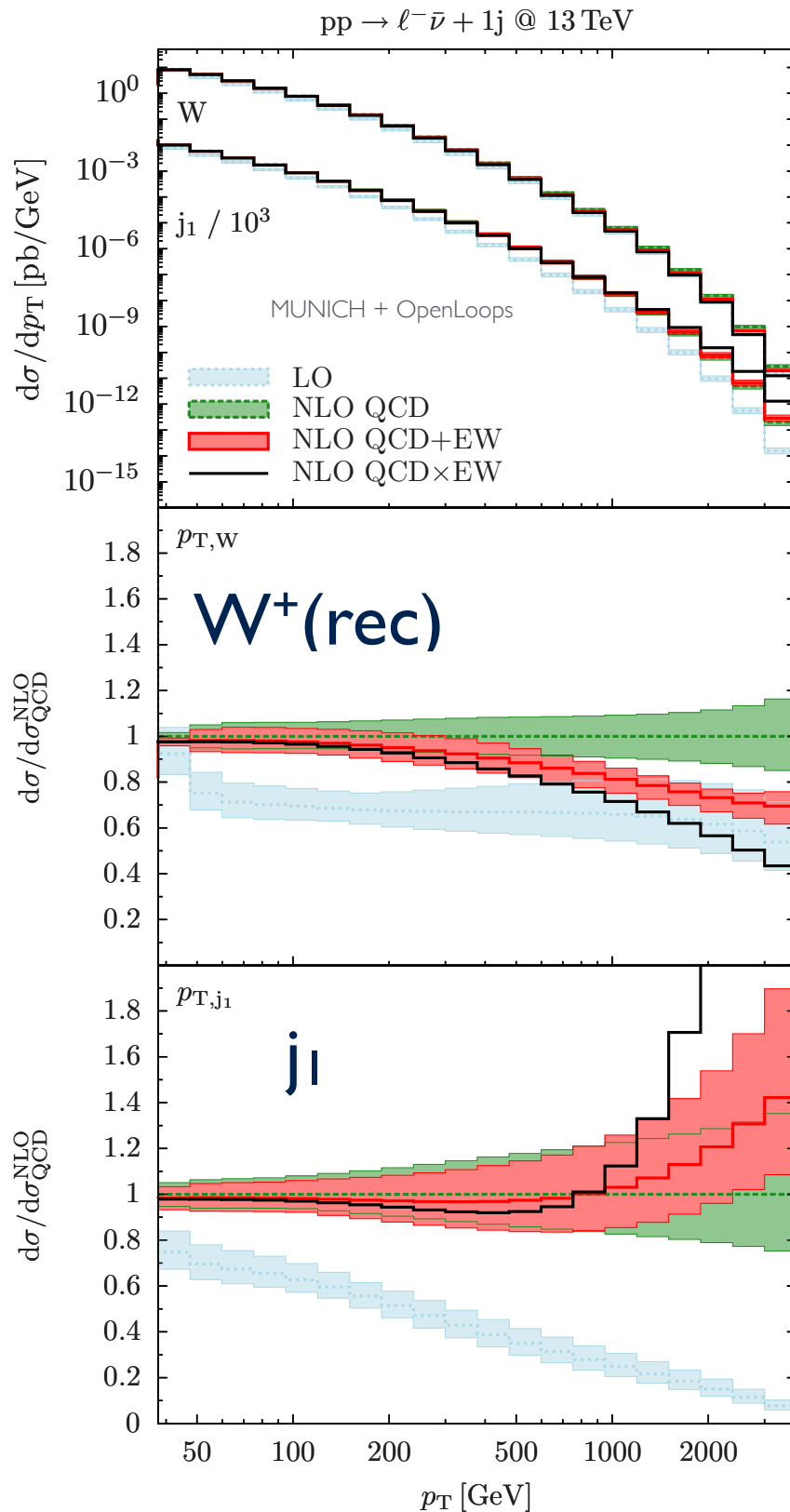
Setup:

$$\sqrt{S} = 13 \text{ TeV}$$

$$p_{T,j} > 30 \text{ GeV}, \quad |\eta_j| < 4.5$$

$$p_{T,1} > 20 \text{ GeV}, \quad |\eta_1| < 2.5, \quad E_T^{\text{miss}} > 25$$

$$\mu_0 = \hat{H}_T/2 \text{ (+ 7-pt. variation)}$$



inclusive

$\approx 1\%$ EW corrections

p_T of W-boson

- ▶ +100% QCD corrections in the tail
- ▶ large negative EW corrections due to **Sudakov behaviour**: -20–35% corrections at 1–4 TeV
- ▶ sizeable difference between QCD+EW and QCDxEW!

p_T of jet

- ▶ “giant QCD K-factors” in the tail [Rubin, Salam, Sapeta '10]
- ▶ dominated by **dijet configurations** (effectively LO, no EW)
- ▶ positive 10–50% EW corrections from quark bremsstrahlung

[S. Kallweit, JML, P. Maierhöfer, M. Schönherr, S. Pozzorini, '14+'15]

$l\nu + 1 \text{ jet}$: inclusive

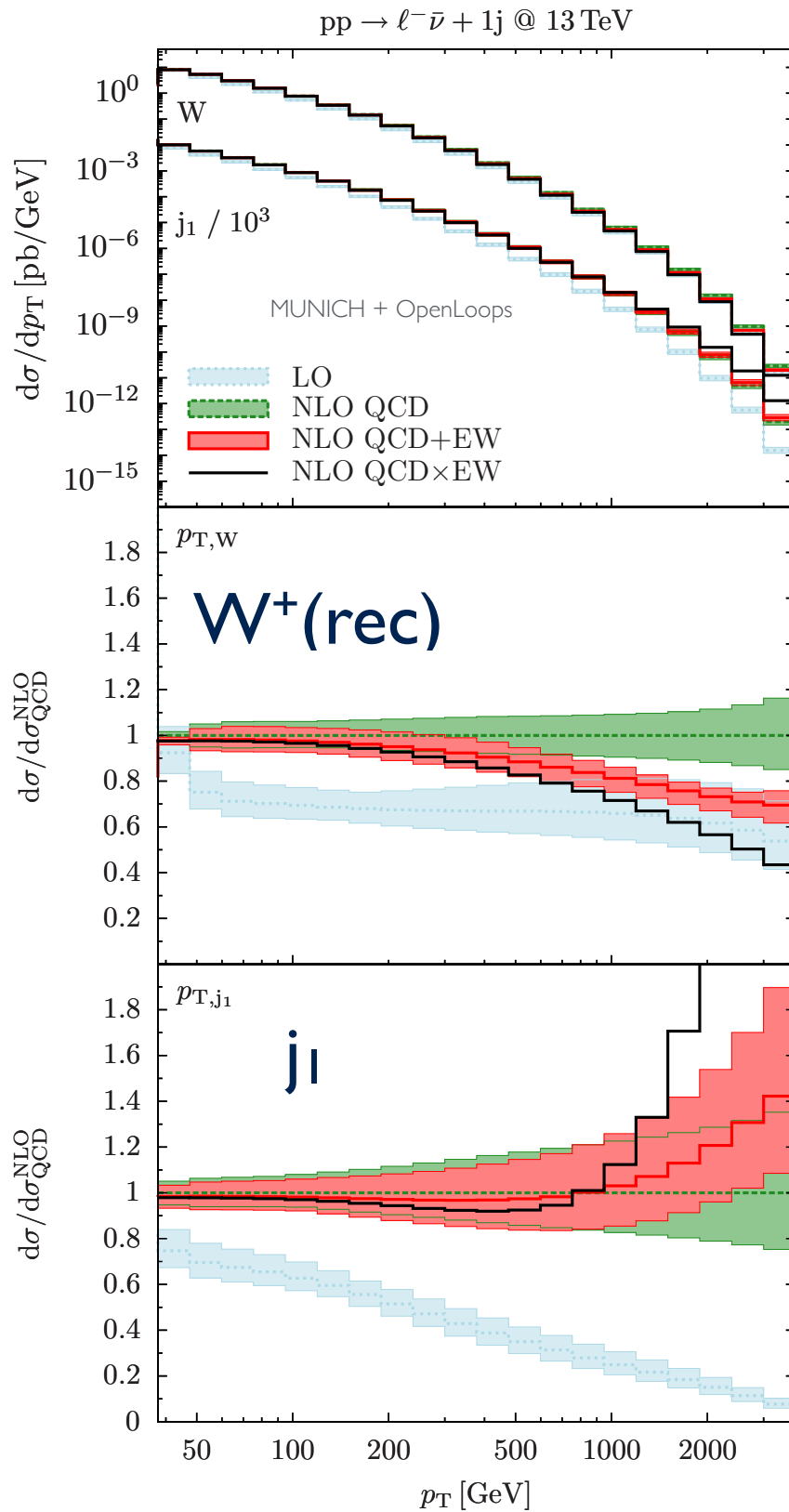
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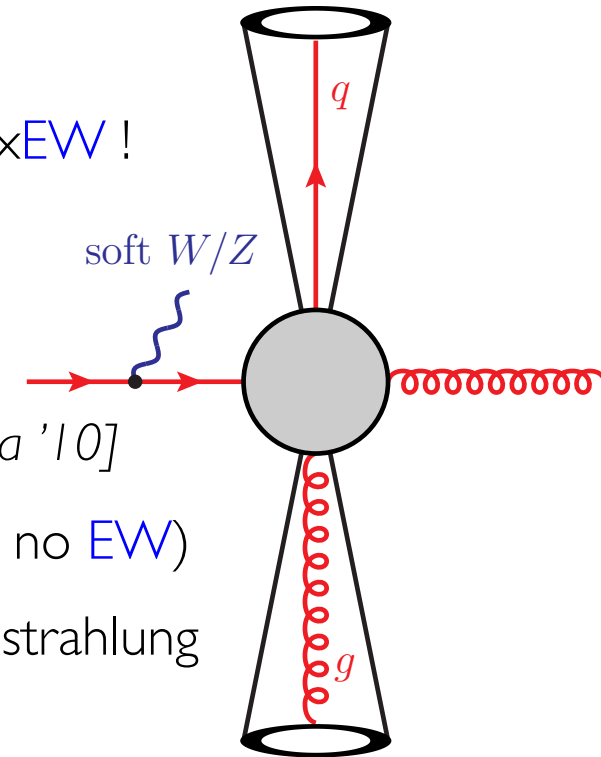
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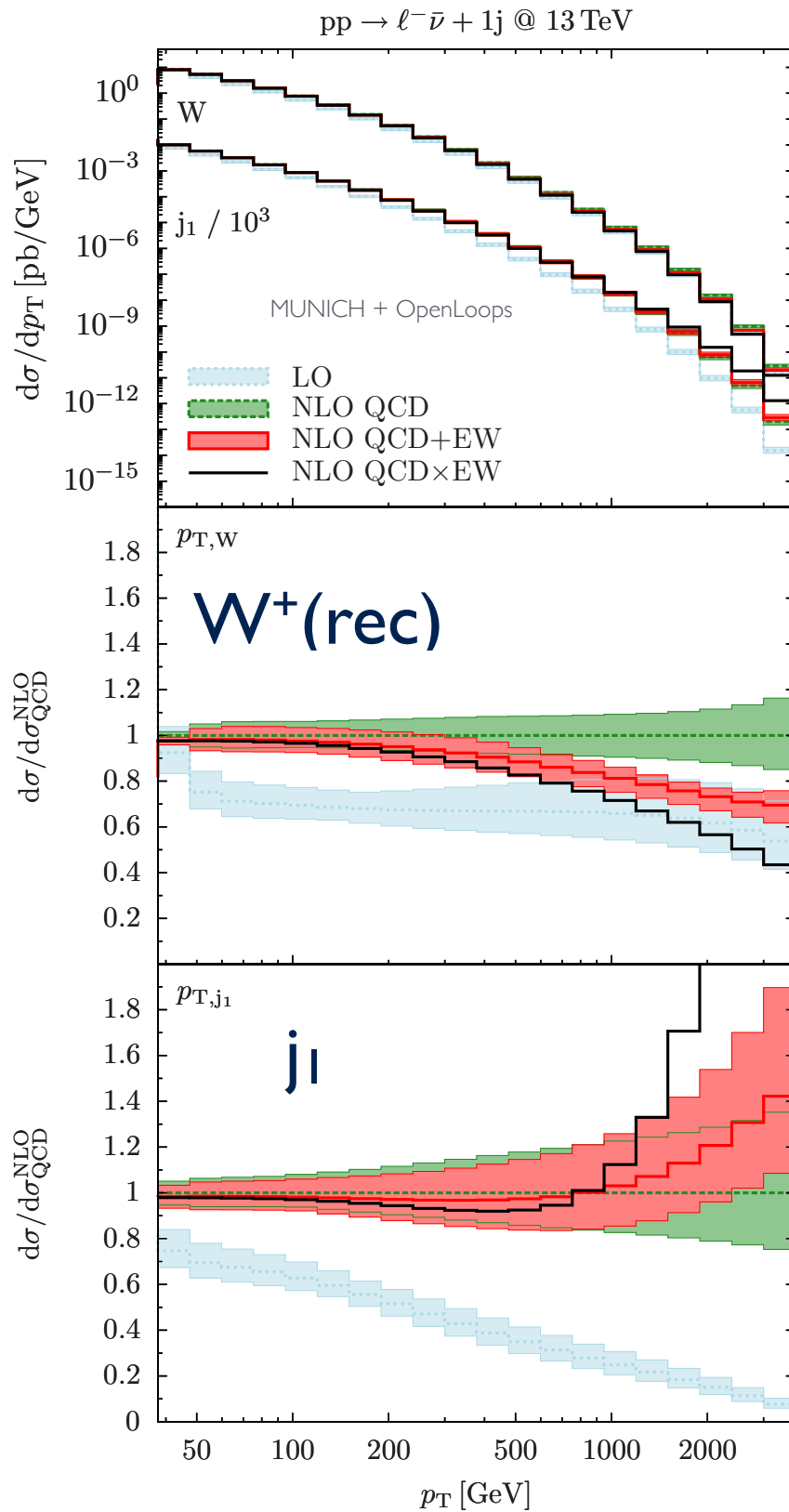
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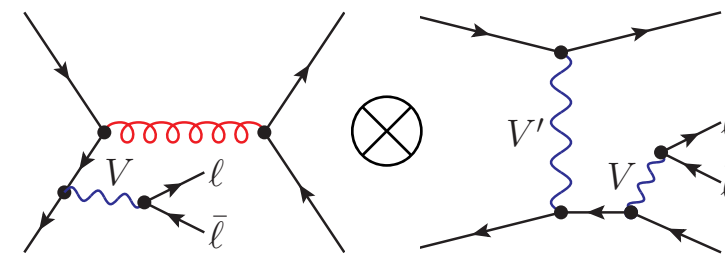
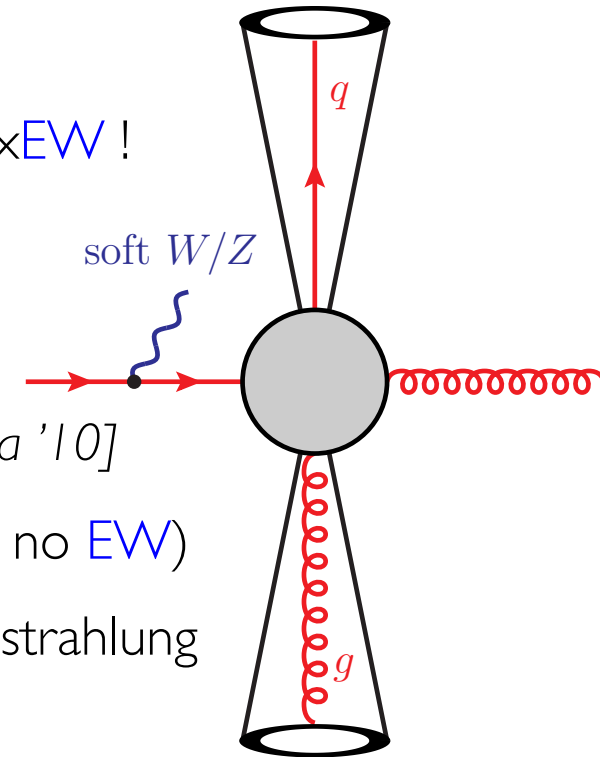
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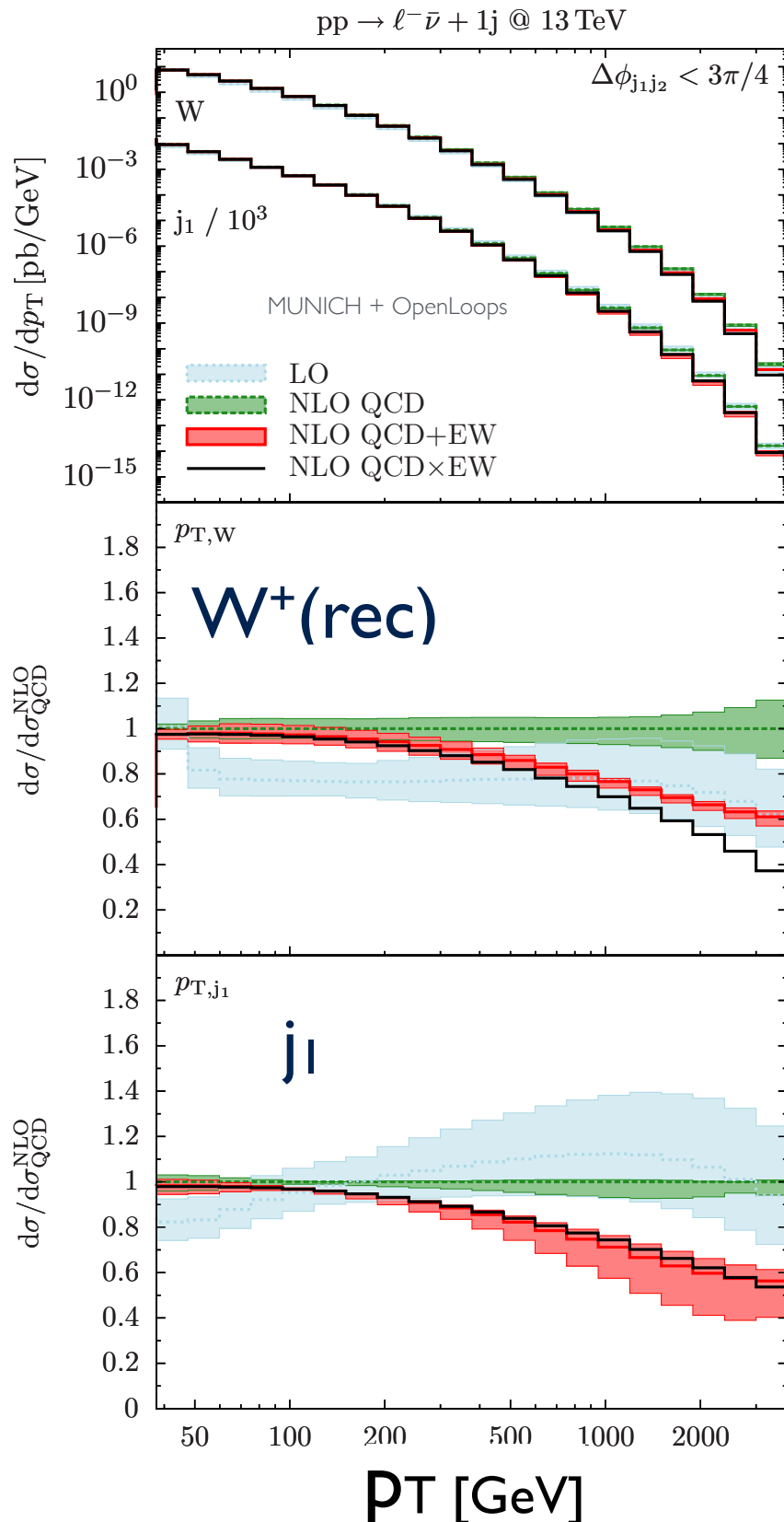
⇒ pathologic with large uncertainties!

[S. Kallweit, JML, P. Maierhöfer, M. Schönherr, S. Pozzorini, '14+'15]

$l\nu + 1 \text{ jet}$: exclusive

$$\Delta\phi_{j1j2} < 3\pi/4$$

(veto on dijet configurations)



QCD corrections

- ▶ mostly moderate and stable QCD corrections

EW corrections

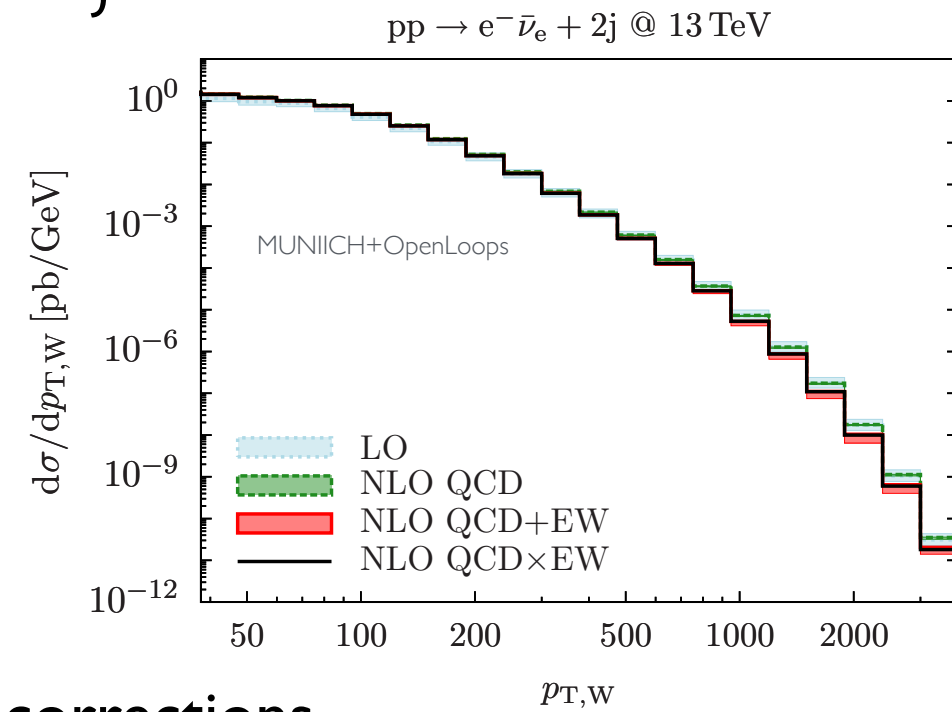
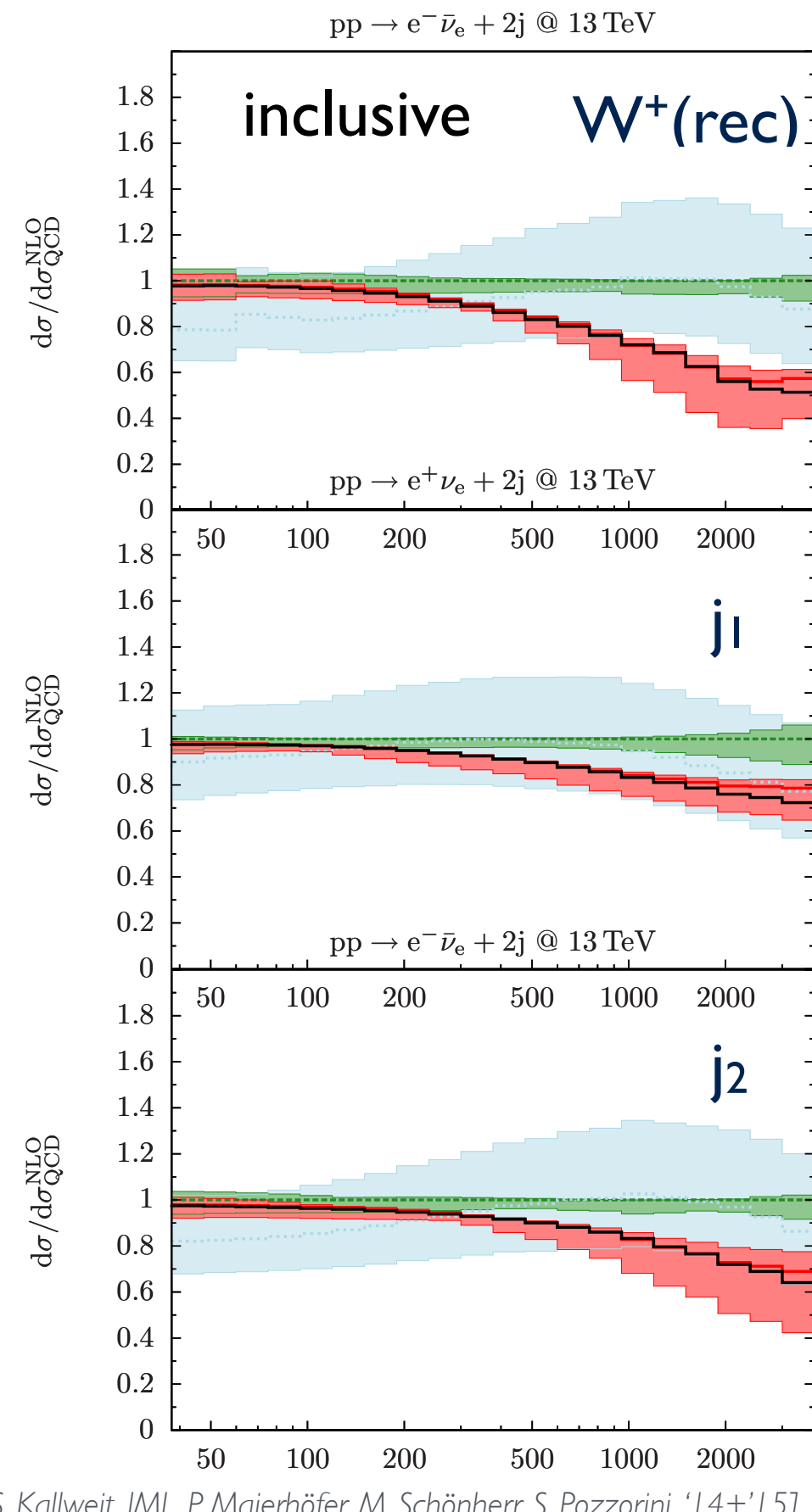
- ▶ **Sudakov behaviour** in both tails:
-20–50% EW corrections at 1–4 TeV
- ▶ EW corrections larger than QCD uncertainties for $p_{T,W^+} > 300$ GeV

\Rightarrow for jet-observables *inclusive* $W^+ 1 \text{ jet}$ requires merging with $W^+ 2 \text{ jets}$ at NLO QCD+EW!

[S. Kallweit, JML, P. Maierhöfer, M. Schönherr, S. Pozzorini, '14+'15]

$V + 2 \text{ jet}$

$l\nu + 2 \text{ jets}$



QCD corrections

- ▶ small and very stable
- ▶ $\approx 10\%$ scale uncertainties

EW corrections

- ▶ **Sudakov behaviour** in all p_T tails:

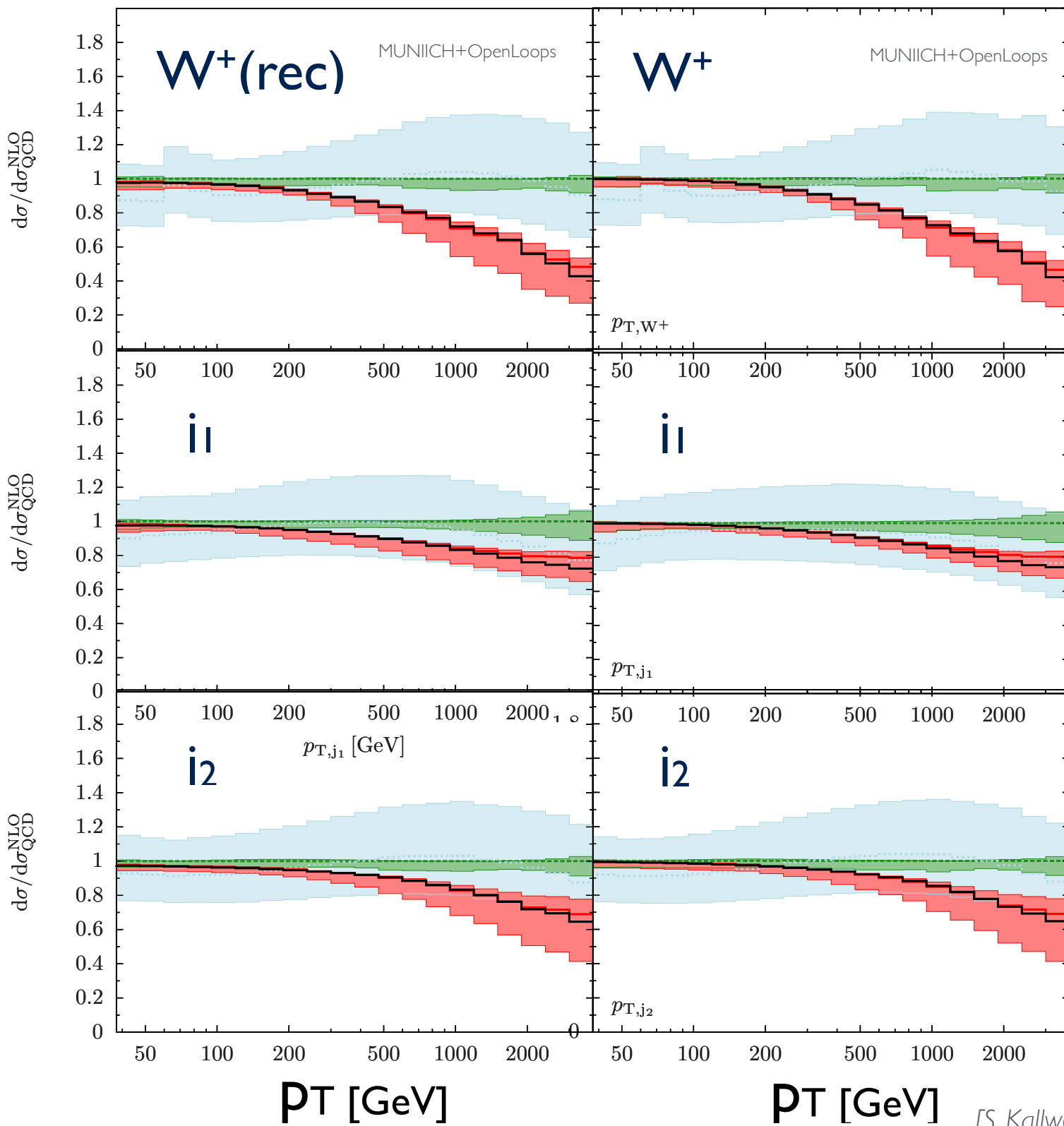
- -30–60% for W-boson at 1-4 TeV
- -15–25% for 1st and 2nd jet at 1-4 TeV

different!

- ▶ Might need resummation of leading EW Sudakov logs

off-shell vs. on-shell production

$l^+ \nu + 2 \text{ jets}$ vs. $W^+ + 2 \text{ jets}$

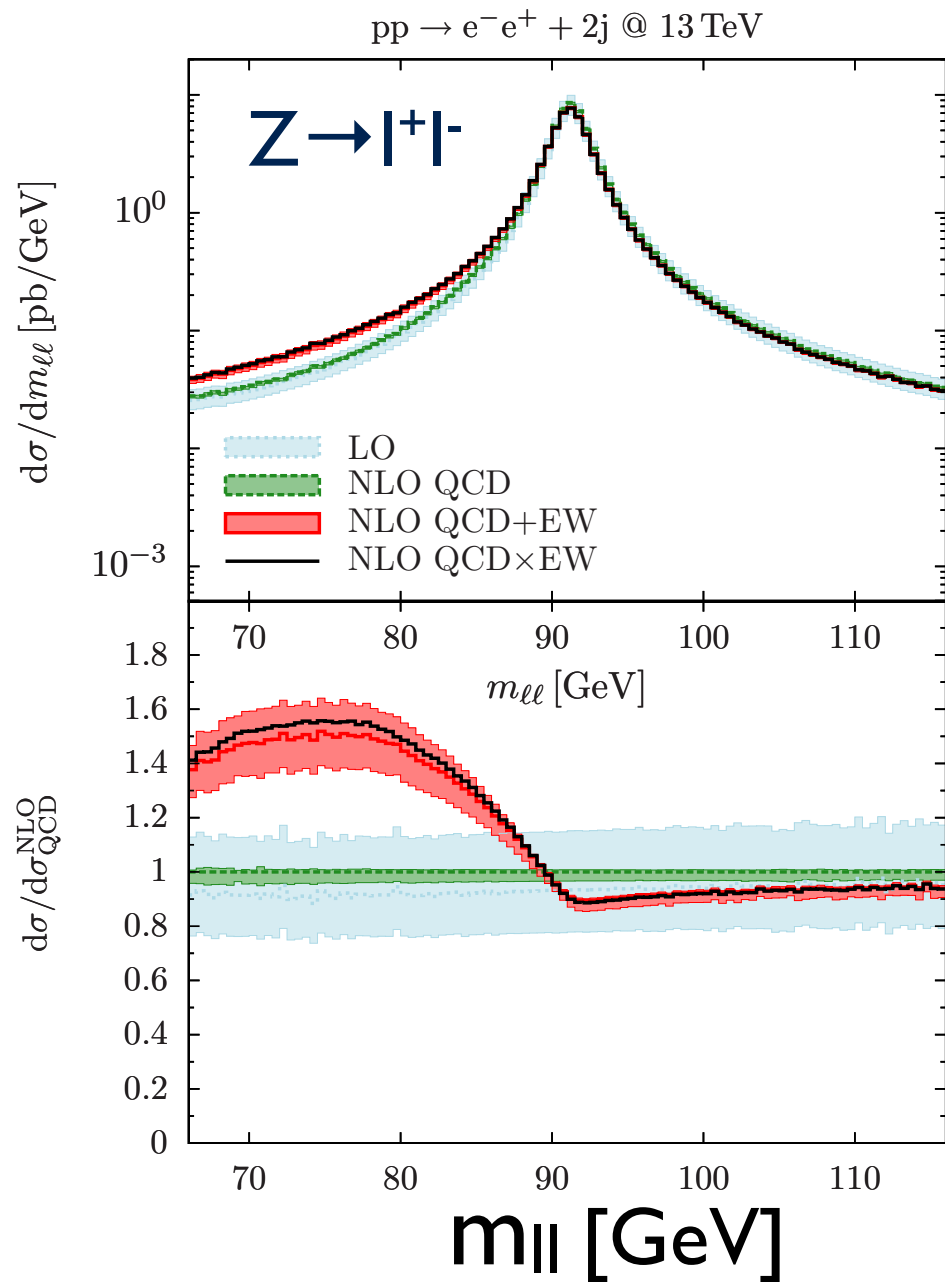


Effect of decays:

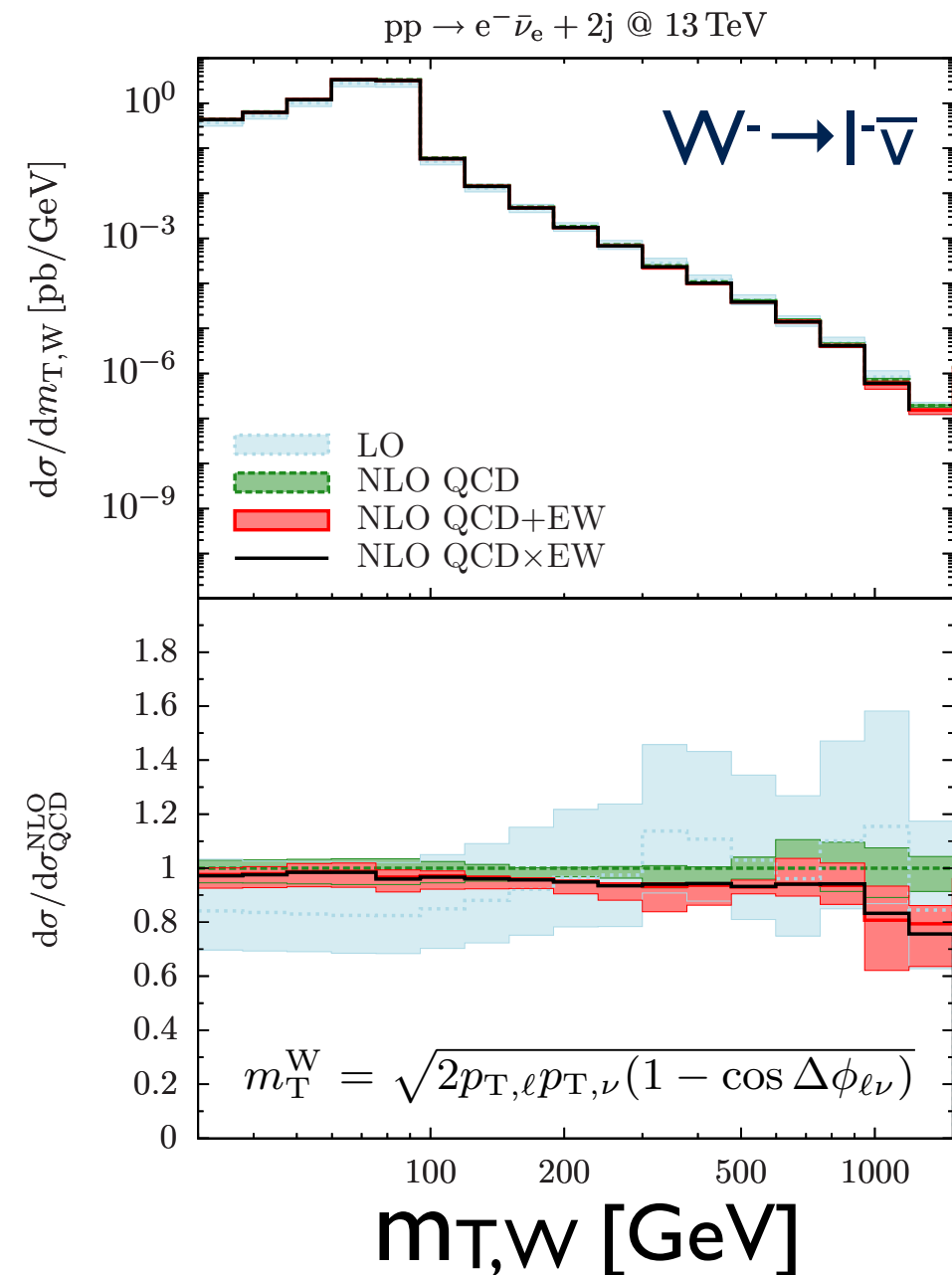
- ▶ Large Sudakov corrections unaffected
- ▶ However: needed for realistic experimental cuts

[S. Kallweit, JML, P. Maierhöfer, M. Schönherr, S. Pozzorini, '14+'15]

Leptonic observables: only in off-shell calculation



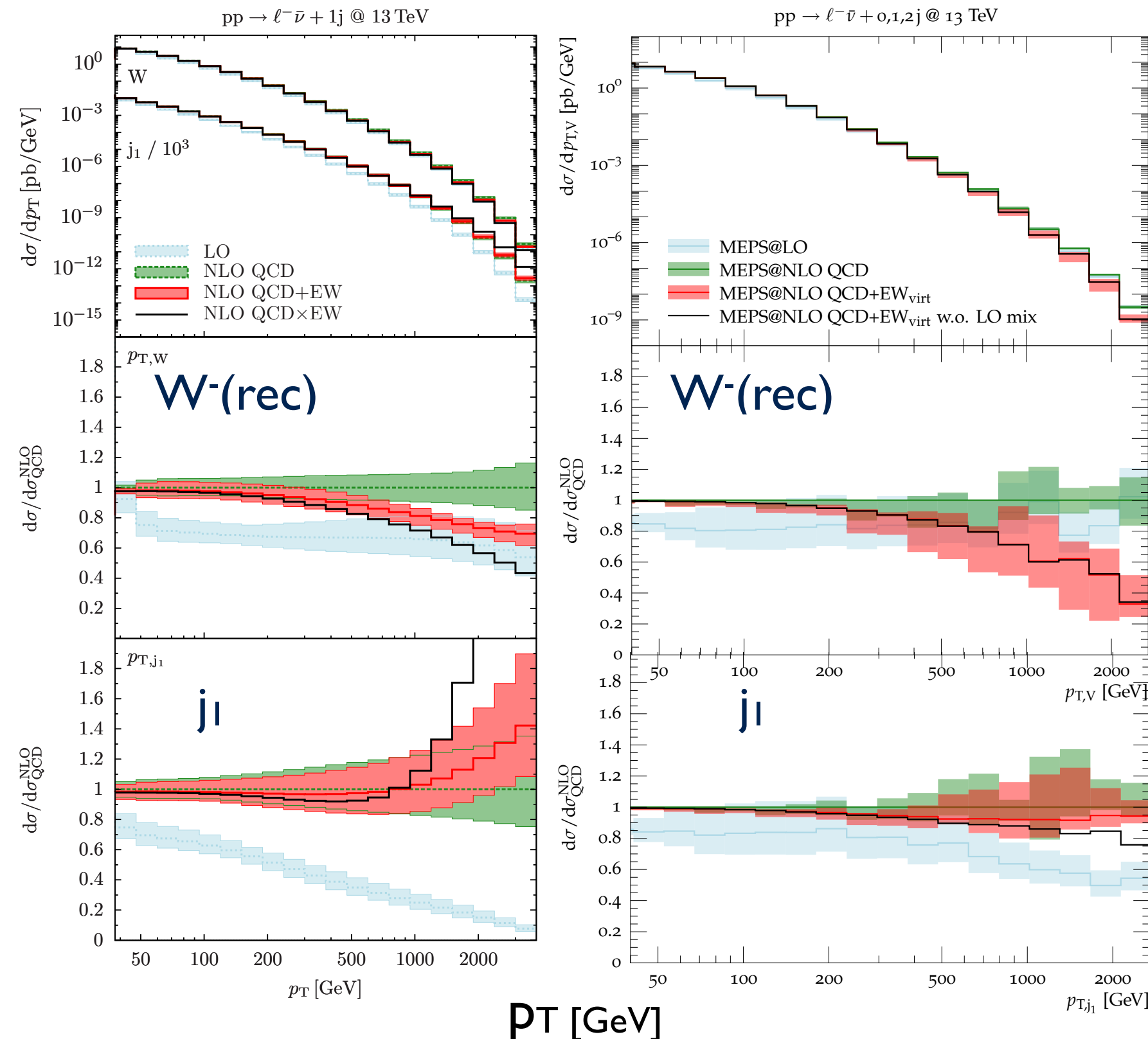
- ▶ up to 50% from QED Bremsstrahlung.
- ▶ Similar shape as for NC DY



- ▶ moderate EW corrections at large $m_{T,W}$
- ▶ no (strong) Sudakov enhancement

$$V + 1 \text{ jet} \otimes V + 2 \text{ jets}$$

inclusive V+1jet: MEPS@NLO QCD+EW_{virt}



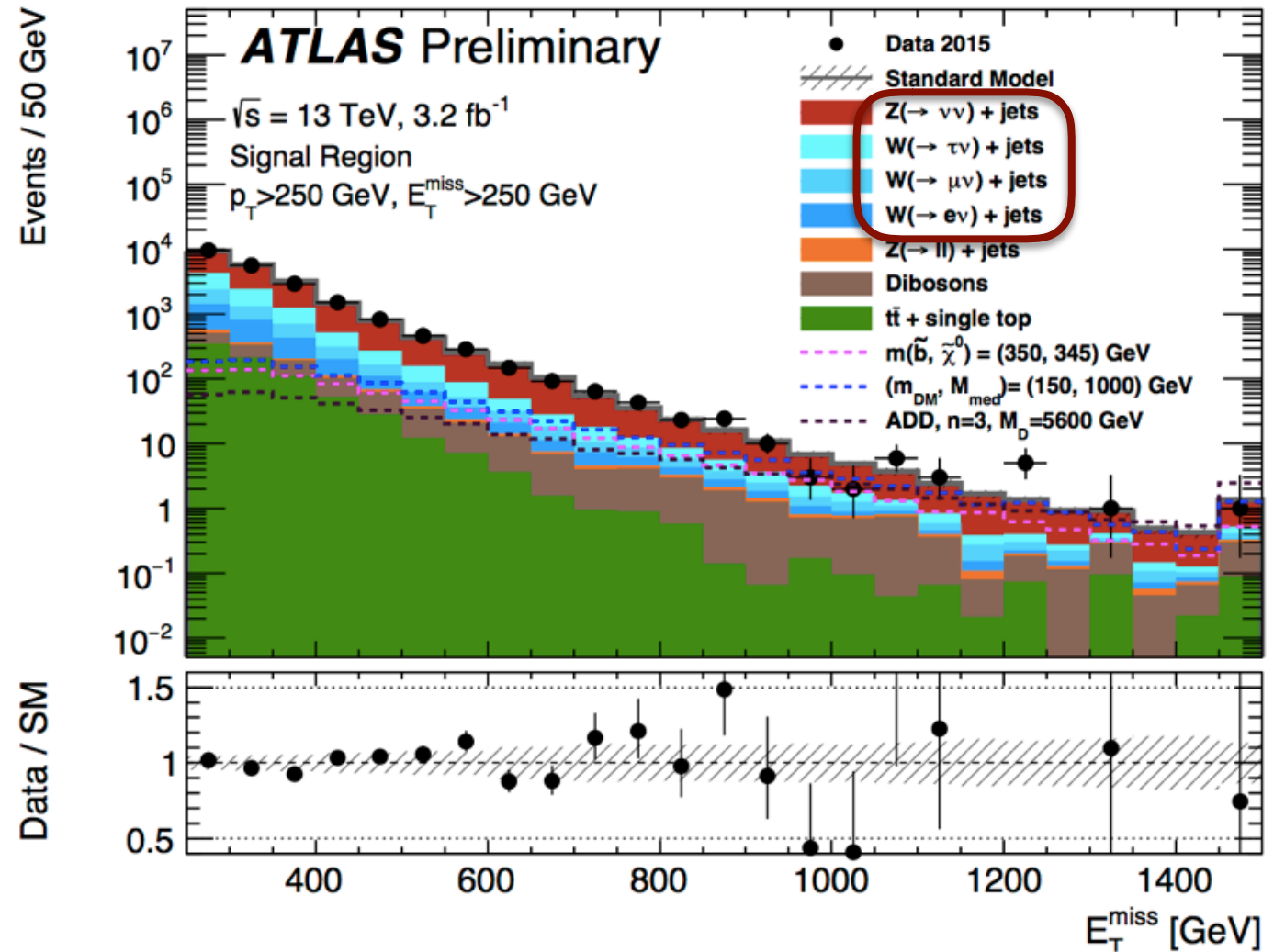
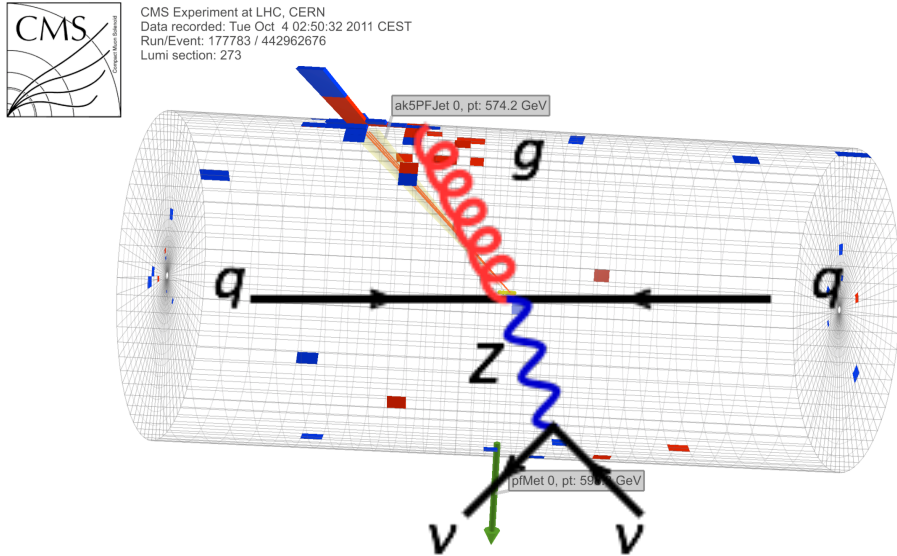
- ▶ Stable NLO QCD+EW predictions in all of the phase-space...
- ▶ ...including Parton-Shower effects.
- ▶ Can directly be used by the experimental collaborations
- ▶ $p_{T,V}$: MEPS@NLO QCD+EW in agreement with QCD \times EW (fixed-order)
- ▶ p_{T,j_1} : compensation between negative Sudakov and LO mix

[see Marek's talk]

[S. Kallweit, JML, P. Maierhöfer, M. Schönherr, S. Pozzorini, '15]

Uncertainties

V+jets backgrounds in monojet/MET + jets searches

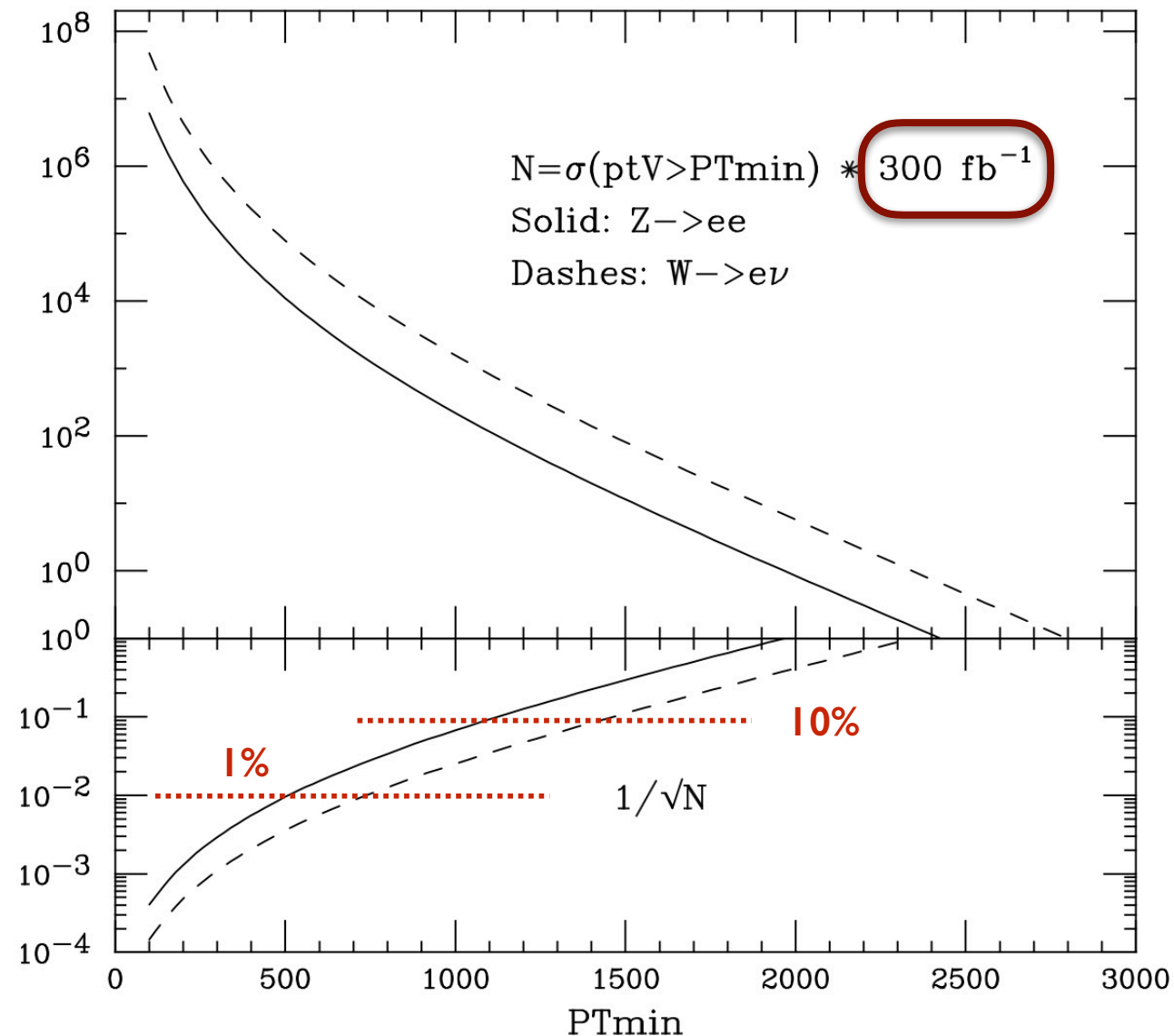


irreducible backgrounds:

$$pp \rightarrow Z(\rightarrow \nu\bar{\nu}) + \text{jets} \Rightarrow \text{MET} + \text{jets}$$

$$pp \rightarrow W(\rightarrow l\nu) + \text{jets} \Rightarrow \text{MET} + \text{jets} \quad (\text{lepton lost})$$

Target precision

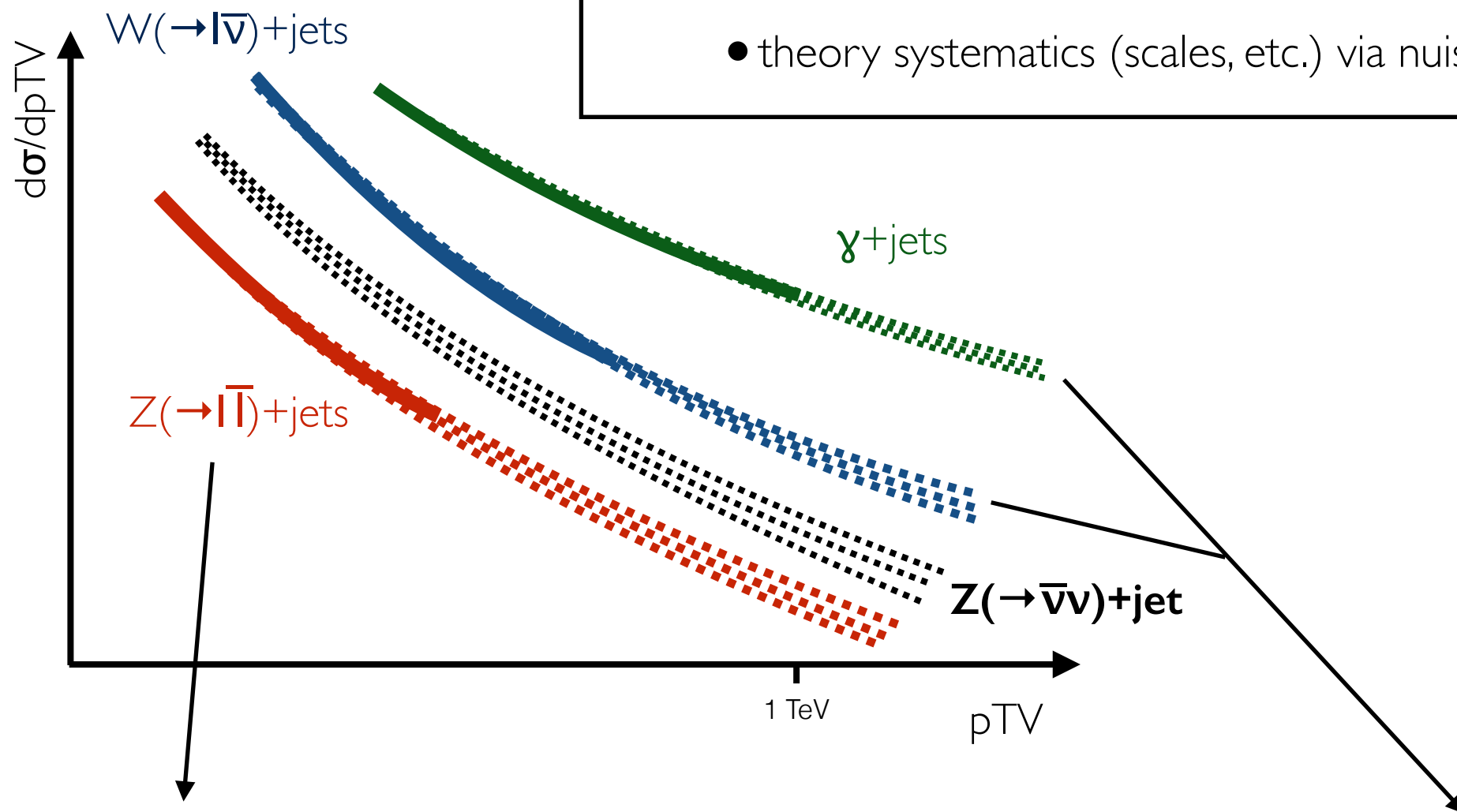


- for $500 \text{ GeV} < p_{TV} < 1000 \text{ GeV}$: background statistics will be at **1% level**
- understanding of V +jets backgrounds at this level increases sensitivity in DM searches
- this level of precision is theoretically possible @ **NNLO QCD + NNLO EW**
- requires solid understanding of **uncertainties!**

Determine V +jets backgrounds

global fit of $Z(\rightarrow l\bar{l})$ +jets, $W(\rightarrow l\bar{\nu})$ +jets and γ +jets measurements

- to determine $Z(\rightarrow \bar{\nu}\nu)$ +jet
- and the visible channels at high- p_T
- theory systematics (scales, etc.) via nuisance parameters in fit



- hardly any systematics (just QED dressing)
- very precise at low p_T
- but: limited statistics at large p_T

- fairly large data samples at large p_T
- systematics from transfer factors

Goal of the ongoing study [to be published soon, already available to ATLAS & CMS]

work in collaboration with:

R. Boughezal, A. Denner, S. Dittmaier, A. Huss, A. Gehrmann-De Ridder, T. Gehrmann, N. Glover,
S. Kallweit, M. L. Mangano, T.A. Morgan, A. Mück, M. Schönherr, F. Petriello, S. Pozzorini, G. P. Salam

- Combination of state-of-the-art predictions: (N)NLO QCD + (N)NLO EW in order to match (future) experimental sensitivities (1-10% accuracy in the few hundred GeV-TeV range)

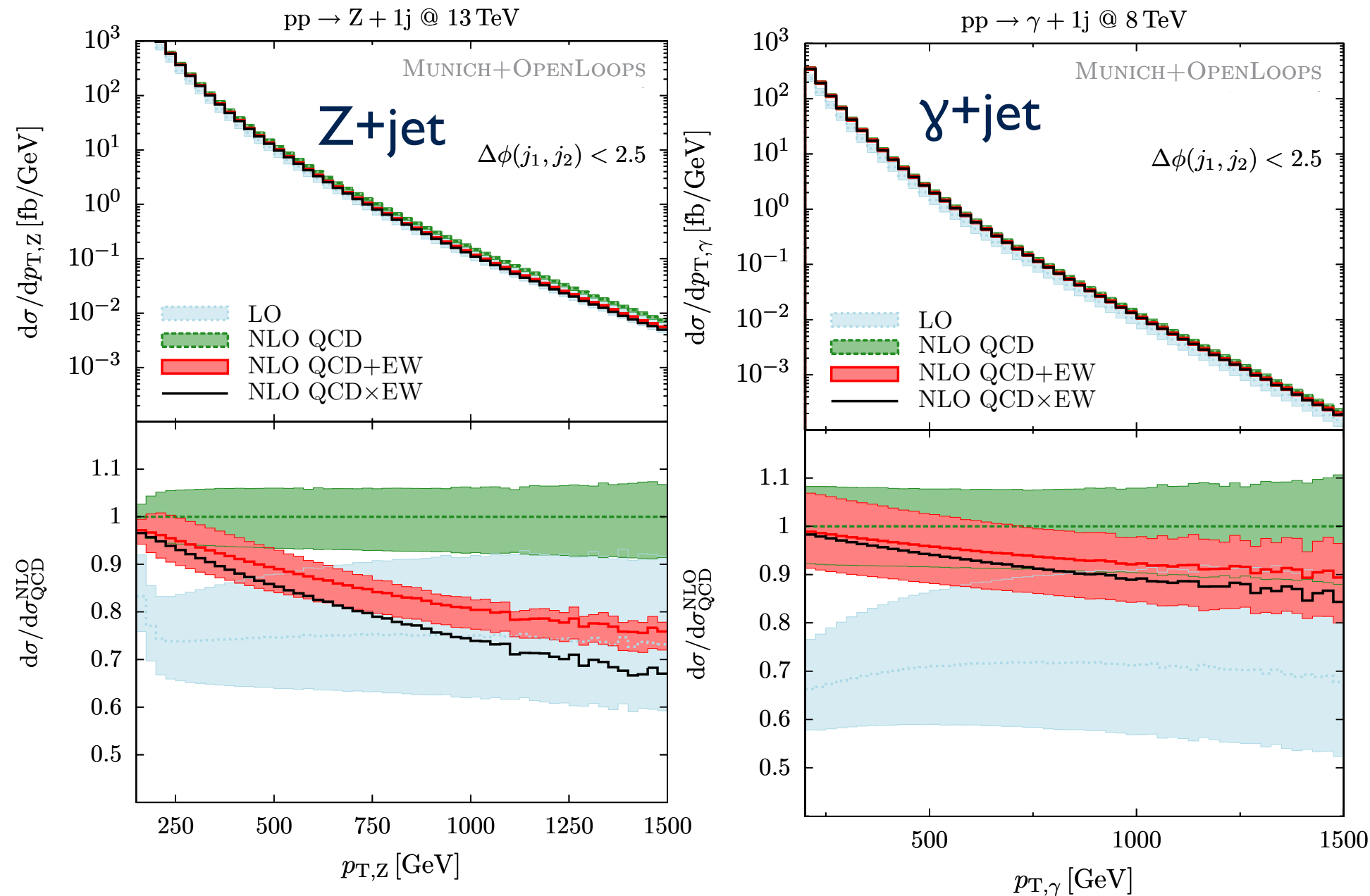
$$\frac{d}{dx} \frac{d}{dy} \sigma^{(V)}(\vec{\epsilon}_{\text{MC}}, \vec{\epsilon}_{\text{TH}}) := \frac{d}{dx} \frac{d}{dy} \sigma_{\text{MC}}^{(V)}(\vec{\epsilon}_{\text{MC}}) \left[\begin{array}{c} \frac{d}{dx} \sigma_{\text{TH}}^{(V)}(\vec{\epsilon}_{\text{TH}}) \\ \frac{d}{dx} \sigma_{\text{MC}}^{(V)}(\vec{\epsilon}_{\text{MC}}) \end{array} \right]$$

one-dimensional reweighting of MC samples in $x = p_{\text{T}}^{(V)}$

with
$$\frac{d}{dx} \sigma_{\text{TH}}^{(V)} = \frac{d}{dx} \sigma_{\text{QCD}}^{(V)} + \frac{d}{dx} \sigma_{\text{mix}}^{(V)} + \frac{d}{dx} \Delta \sigma_{\text{EW}}^{(V)} + \frac{d}{dx} \sigma_{\gamma\text{-ind.}}^{(V)}$$

- Robust **uncertainty estimates** including
 1. Pure QCD uncertainties
 2. Pure EW uncertainties
 3. Mixed QCD-EW uncertainties
 4. PDF, γ -induced uncertainties
- Prescription for **correlation** of these uncertainties
 - ▶ within a process (between low-pT and high-pT)
 - ▶ across processes

Prelude: Z+jet vs. γ + 1 jet



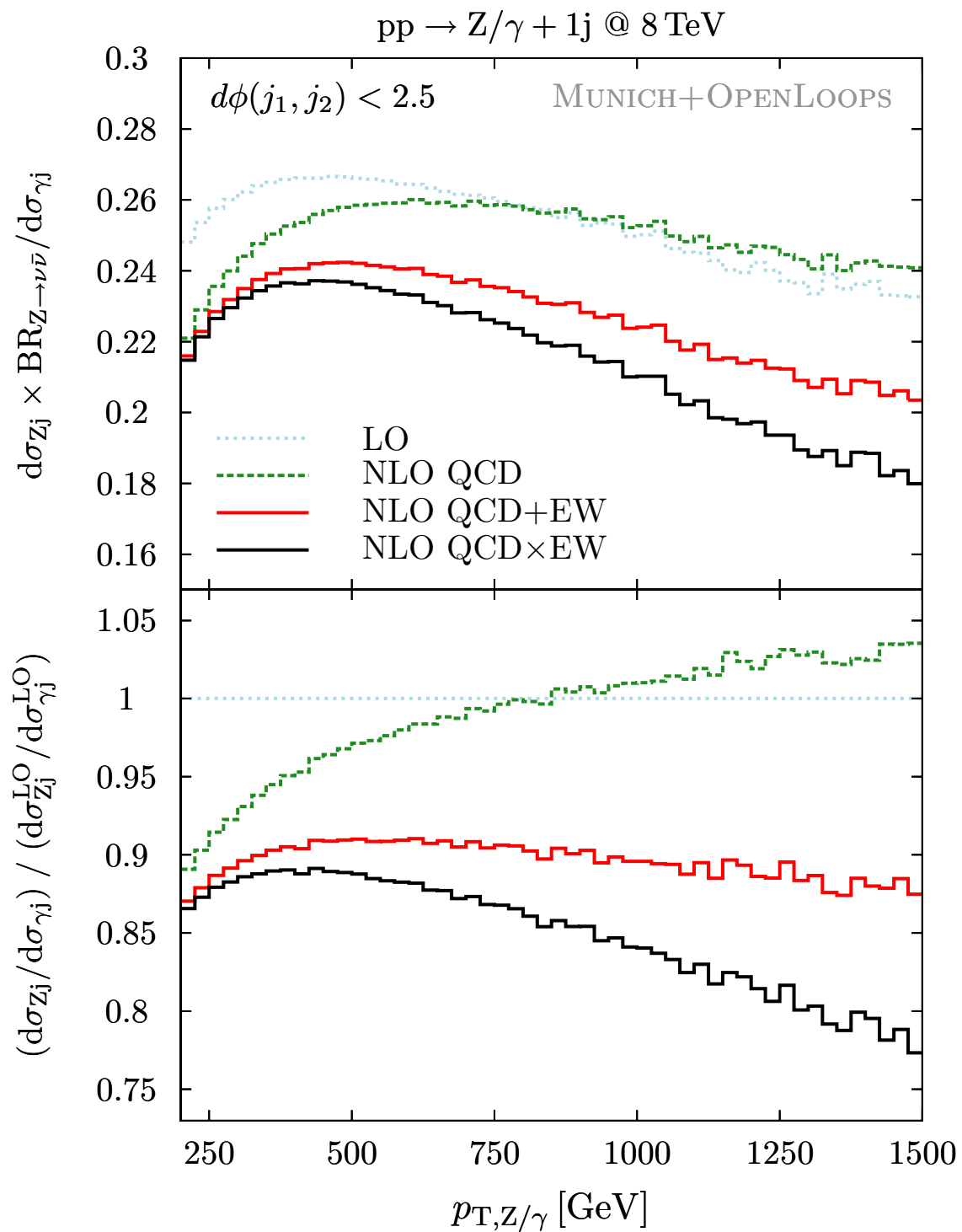
QCD corrections

- ▶ mostly moderate and stable QCD corrections
- ▶ (almost) **identical QCD corrections in the tail**, sizeable differences for small p_T

EW corrections

- ▶ **correction in $p_T(Z) >$ correction in $p_T(\gamma)$**
- ▶ **-20/-8%** for Z/ γ at 1 TeV
- ▶ EW corrections $>$ QCD uncertainties for $p_{T,Z} > 350$ GeV

Prelude: Z/γ pT-ratio



Overall

- ▶ mild dependence on the boson pT

QCD corrections

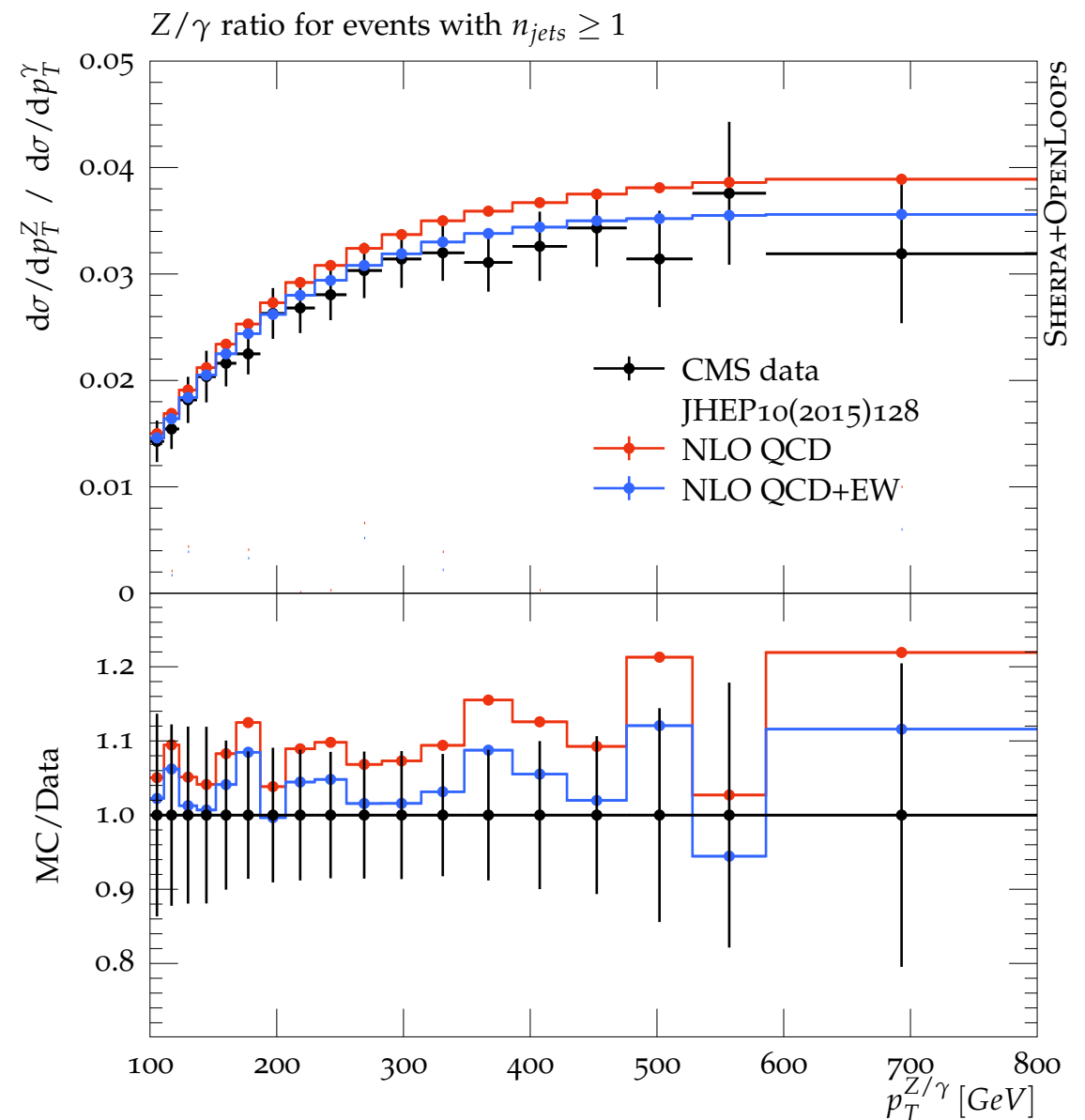
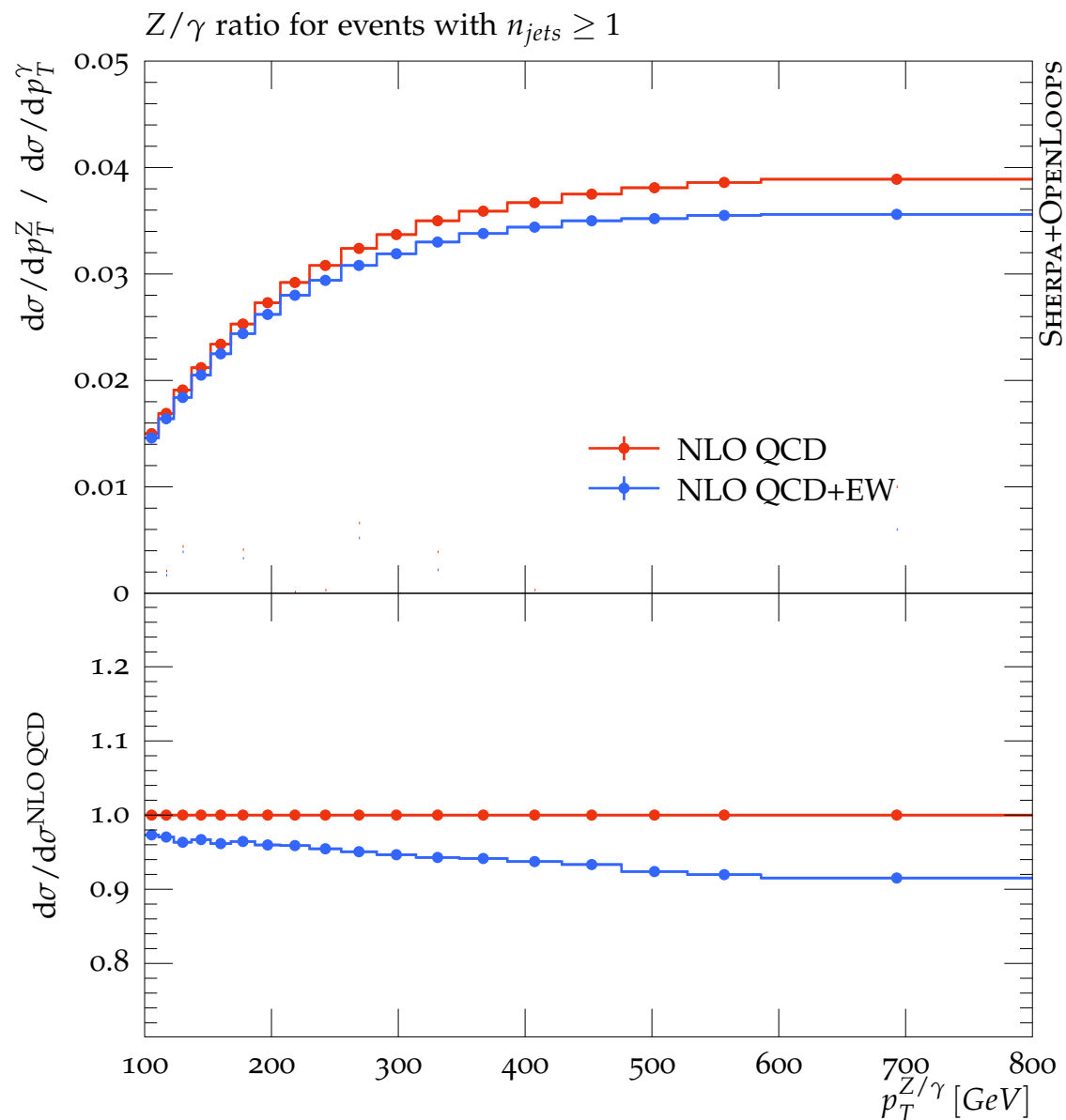
- ▶ 10-15% below 250 GeV
- ▶ $\approx 5\%$ above 350 GeV

EW corrections

- ▶ sizeable difference in EW corrections results in 10-15% corrections at several hundred GeV
- ▶ $\sim 5\%$ difference between NLO QCD+EW and NLO QCD×EW

Prelude: compare against Z/γ -data

[JHEP10(2015)128]



[Ciulli, Kallweit, JML, Pozzorini, Schönherr for **LH'15**]

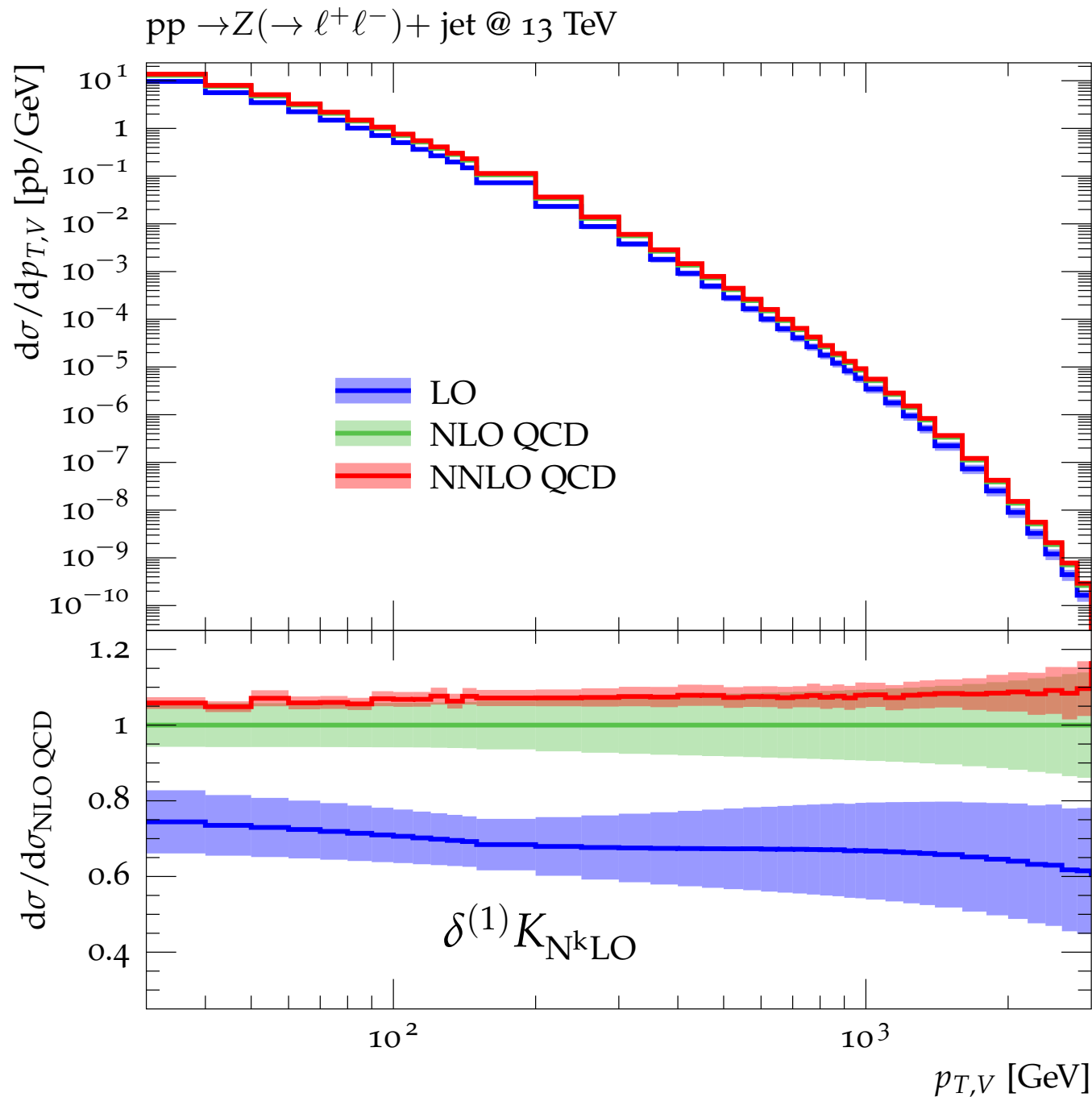
- ▶ remarkable agreement with data at @ NLO **QCD+EW**!

1. pure QCD uncertainties

QCD effects

[see Nigel's talk]

$$\frac{d}{dx}\sigma_{\text{QCD}}^{(V)} = \frac{d}{dx}\sigma_{\text{LO QCD}}^{(V)} + \frac{d}{dx}\sigma_{\text{NLO QCD}}^{(V)} + \frac{d}{dx}\sigma_{\text{NNLO QCD}}^{(V)}$$



$$\mu_0 = \frac{1}{2} \left(\sqrt{p_{T,\ell+\ell^-}^2 + m_{\ell+\ell^-}^2} + \sum_{i \in \{q,g,\gamma\}} |p_{T,i}| \right)$$

this is a 'good' scale for V+jets

- at large $p_{T,V}$: $HT'/2 \approx p_{T,V}$
- modest higher-order corrections
- sufficient convergence

scale uncertainties due to 7-pt variations

$$\mu_{R,F} = \xi_{R,F} \mu_0$$

$$(\xi_R, \xi_F) = (2, 2), (2, 1), (1, 2), (1, 1), (1, 0.5), (0.5, 1), (0.5, 0.5)$$

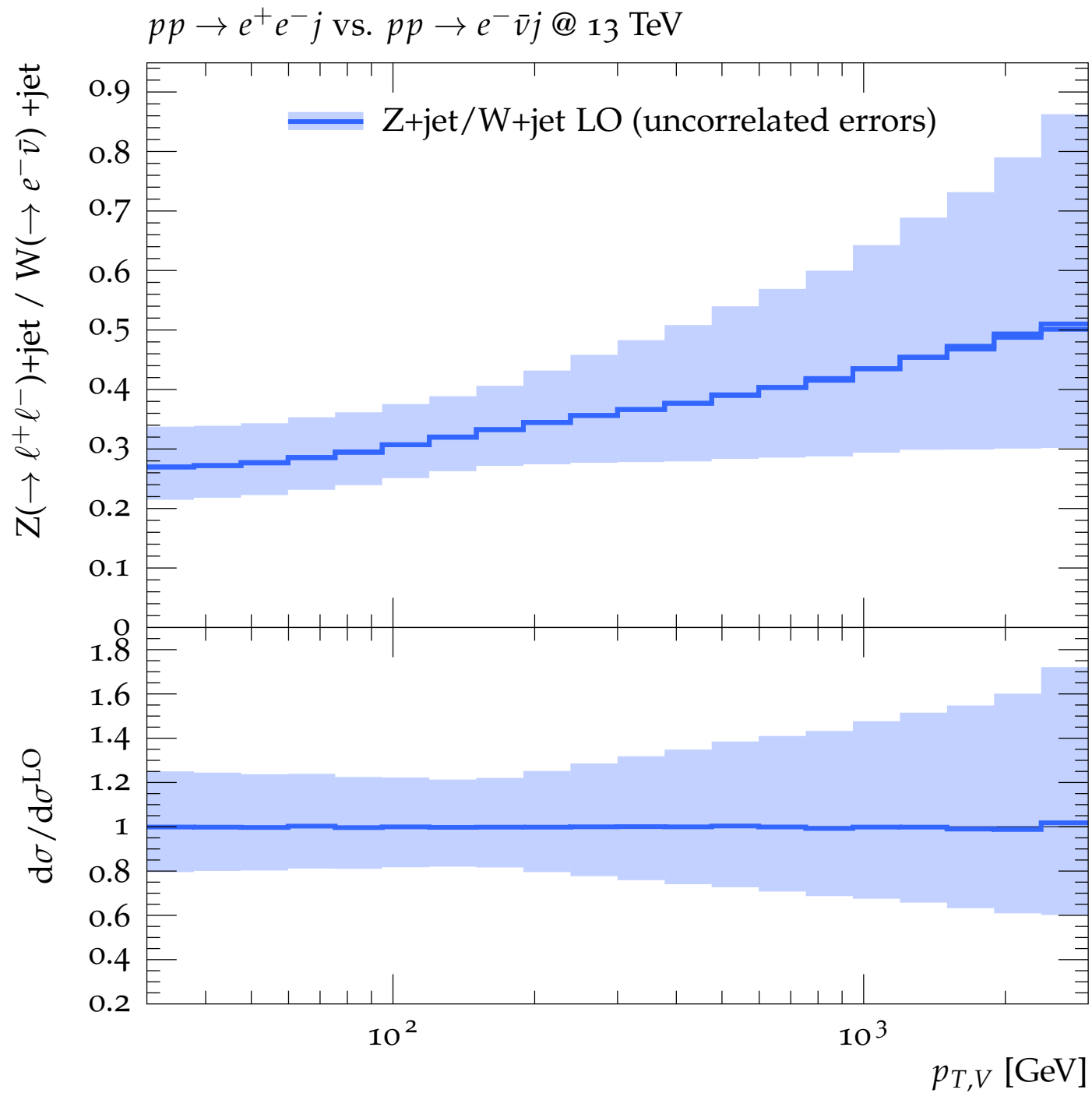
yields

- $\mathcal{O}(20\%)$ uncertainties at LO
- $\mathcal{O}(10\%)$ uncertainties at NLO
- $\mathcal{O}(5\%)$ uncertainties at NNLO

with minor shape variations

NNLO from [A. Huss, A. Gehrmann-De Ridder,
T. Gehrmann, N. Glover, T.A. Morgan]

Correlation of scale variations

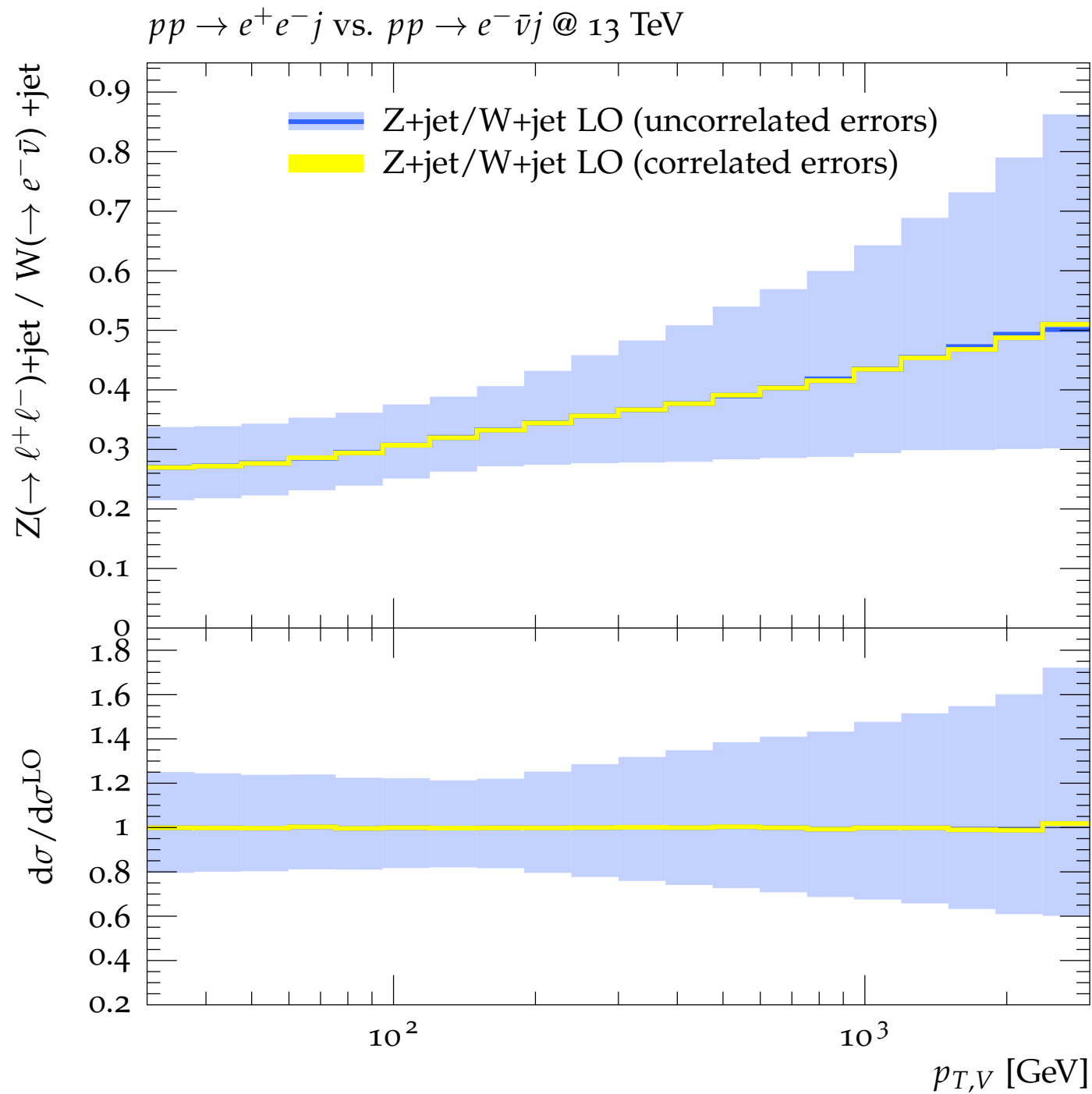


consider Z+jet / W+jet $p_{T,V}$ -ratio @ LO

uncorrelated treatment yields

O(40%) uncertainties

Correlation of scale variations



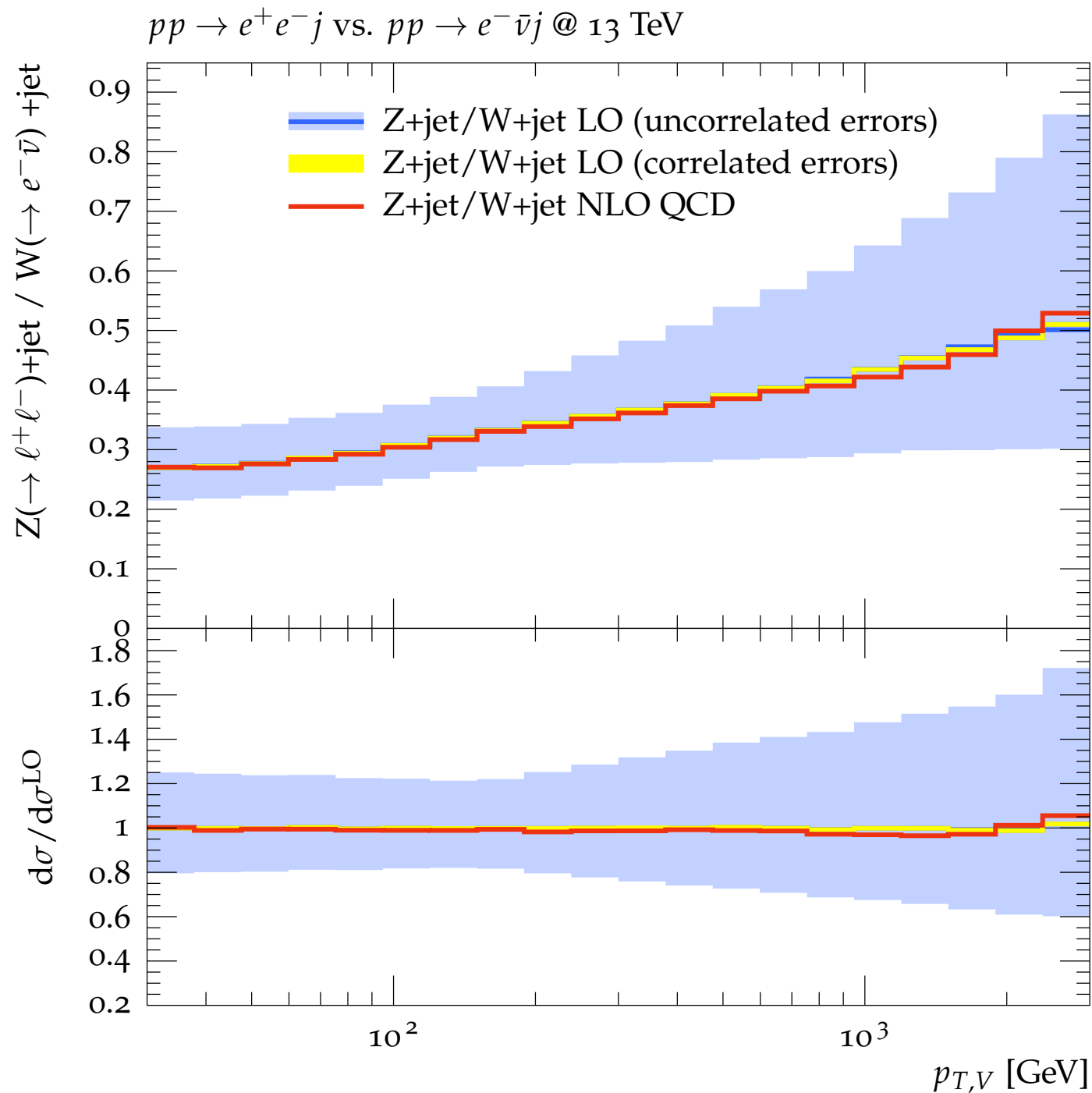
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uncorrelated treatment yields
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correlated treatment yields tiny
O($< \sim 1\%$) uncertainties

check against NLO QCD!

Correlation of scale variations



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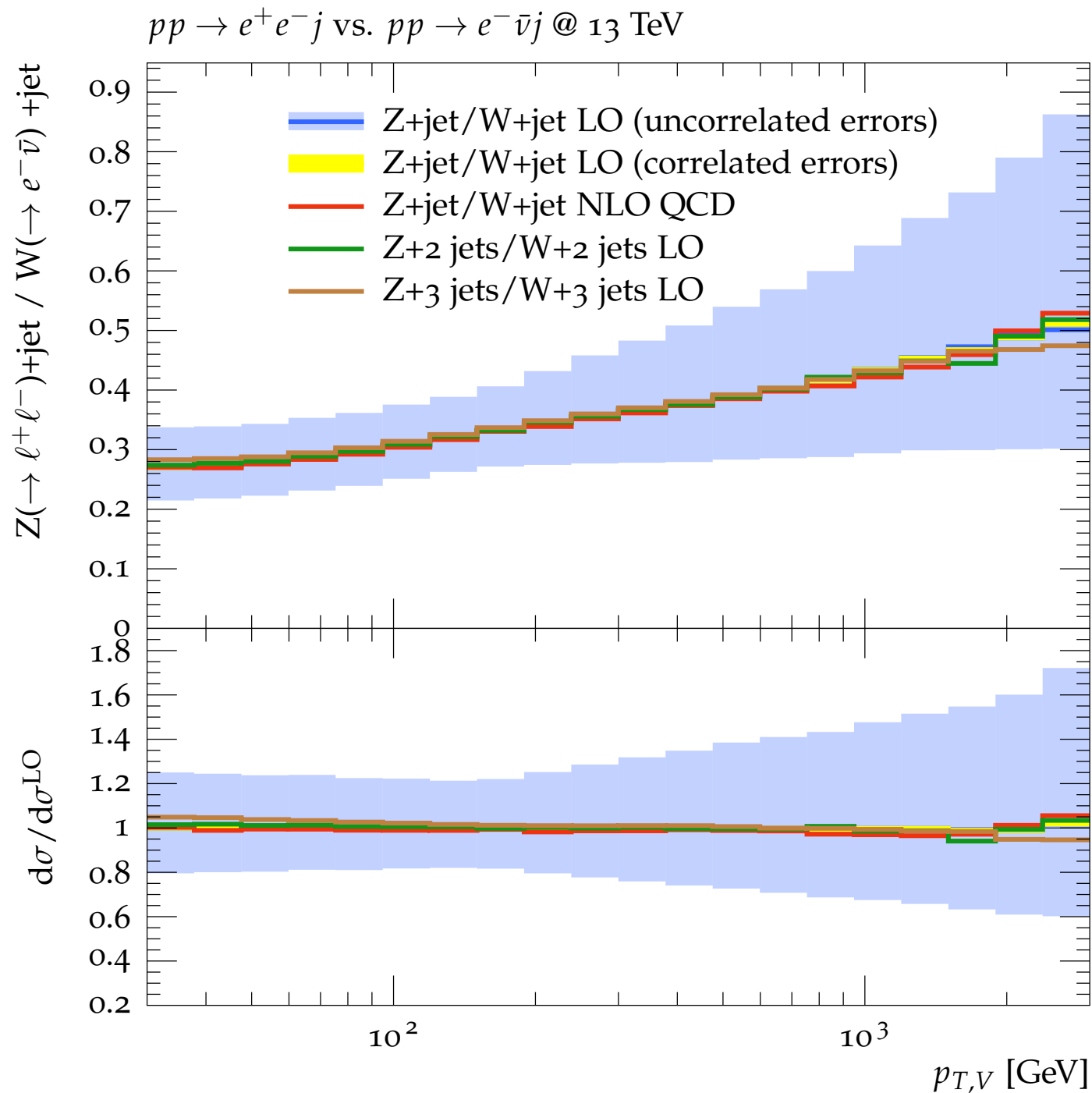
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NLO QCD corrections remarkably flat
in Z+jet / W+jet ratio!

→ supports correlated treatment of
uncertainties!

Correlation of scale variations



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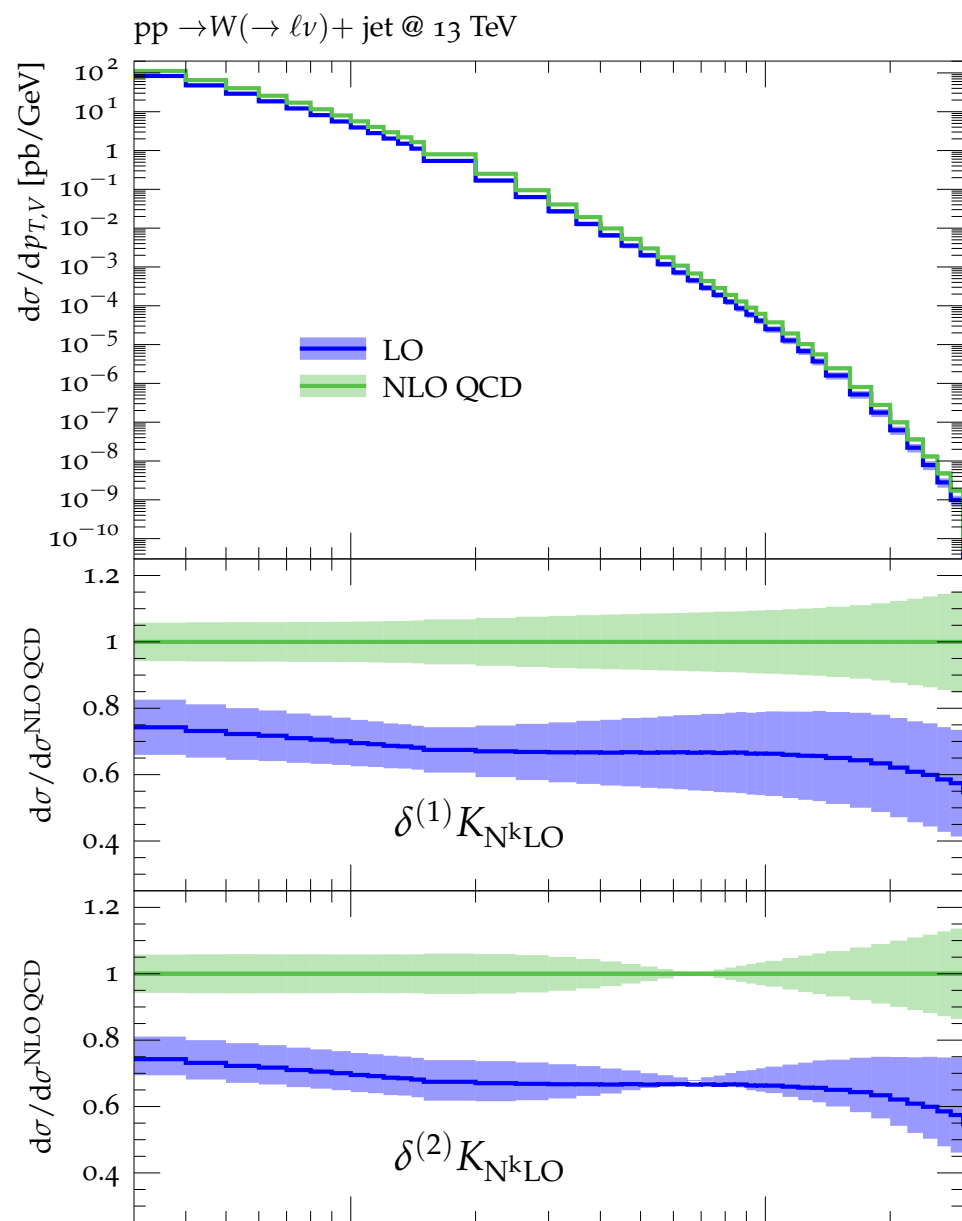
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NLO QCD corrections remarkably flat
in Z+jet / W+jet ratio!

→ supports correlated treatment of
uncertainties!

Also holds for higher jet-multiplicities
→ indication of correlation also in
higher-order corrections beyond NLO!

QCD uncertainties



$$\frac{d}{dx} \sigma_{N^k \text{LO QCD}}^{(V)}(\vec{\epsilon}_{\text{QCD}}) = \left[K_{N^k \text{LO}}^{(V)}(x) + \sum_{i=1}^3 \epsilon_{\text{QCD},i} \delta^{(i)} K_{N^k \text{LO}}^{(V)}(x) \right] \times \frac{d}{dx} \sigma_{\text{LO QCD}}^{(V)}(\vec{\mu}_0).$$

$$\epsilon_{\text{QCD},i}^{(Z)} = \epsilon_{\text{QCD},i}^{(W^\pm)} = \epsilon_{\text{QCD},i}^{(\gamma)} = \epsilon_{\text{QCD},i}$$

- correlated across processes
- correlated across pT bins

nuisance parameters:
interpreted as 1σ Gaussian

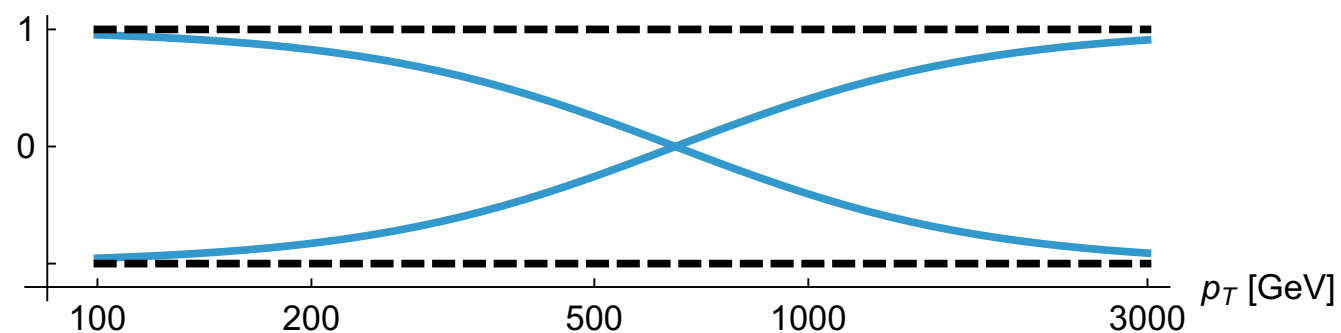
- $\delta^{(1)} K_{N^k \text{LO}}^V = \frac{1}{2} \left[K_{N^k \text{LO}}^{V,\text{max}} - K_{N^k \text{LO}}^{V,\text{min}} \right]$

symmetrized **scale uncertainty**

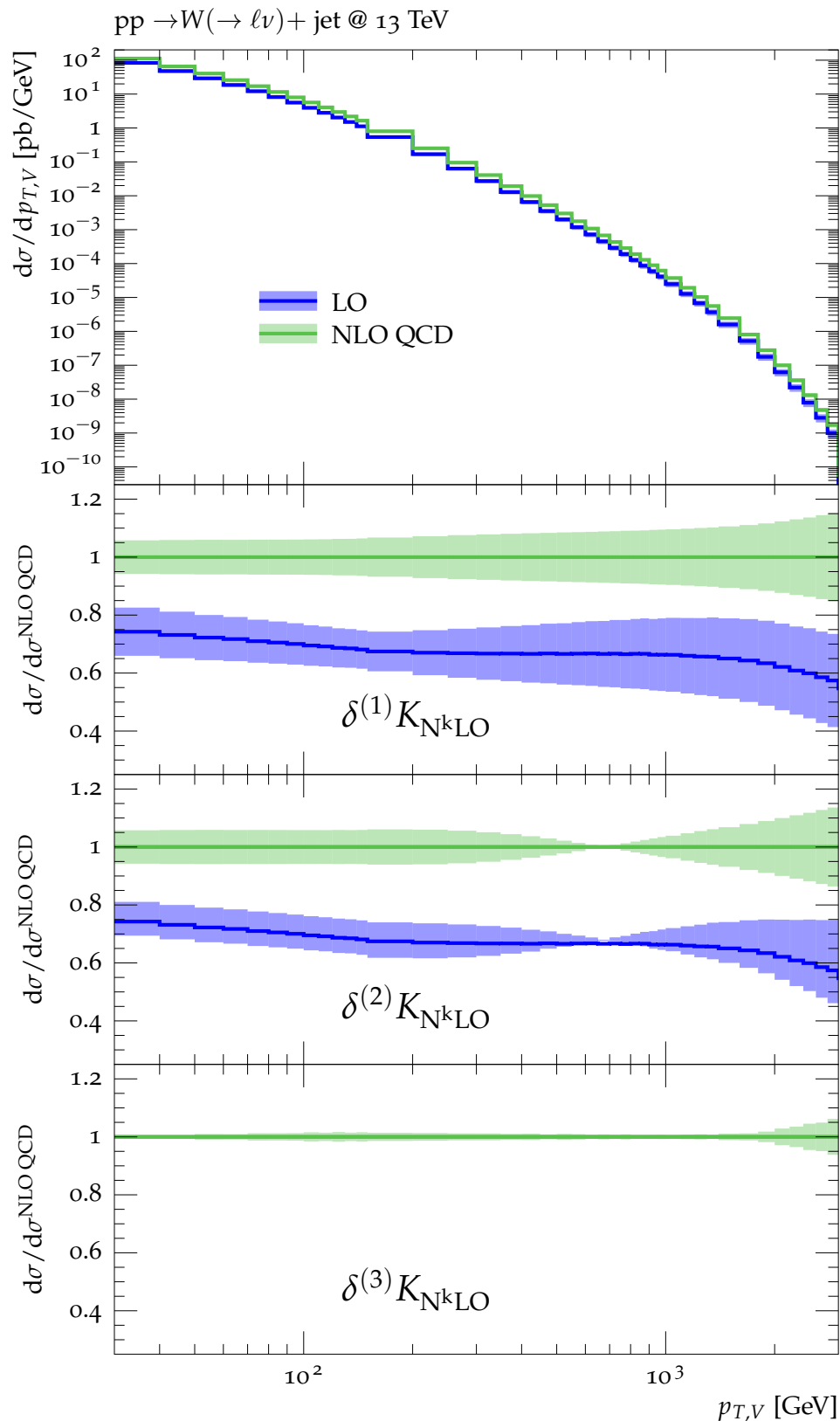
- $\delta^{(2)} K_{N^k \text{LO}}^V = \frac{p_T^2 - 650 \text{ GeV}}{p_T^2 + 650 \text{ GeV}} \delta^{(1)} K_{N^k \text{LO}}^V$

yields max **shape distortion** within scale variation band

$\pm \omega_{\text{shape}}(p_T)$



QCD uncertainties



$$\frac{d}{dx} \sigma_{N^k \text{LO QCD}}^{(V)}(\vec{\epsilon}_{\text{QCD}}) = \left[K_{N^k \text{LO}}^{(V)}(x) + \sum_{i=1}^3 \epsilon_{\text{QCD},i} \delta^{(i)} K_{N^k \text{LO}}^{(V)}(x) \right] \times \frac{d}{dx} \sigma_{\text{LO QCD}}^{(V)}(\vec{\mu}_0).$$

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symmetrized **scale uncertainty**

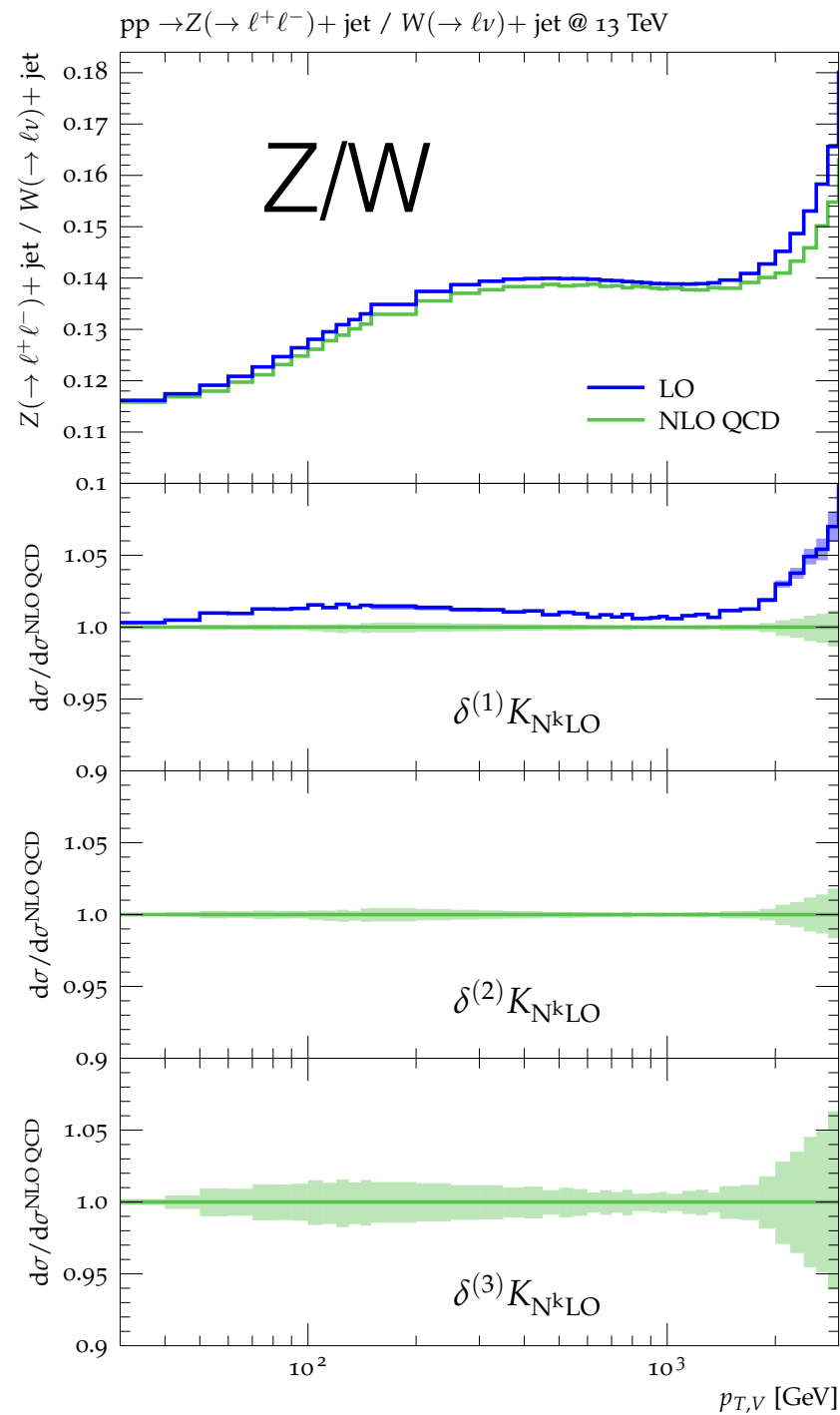
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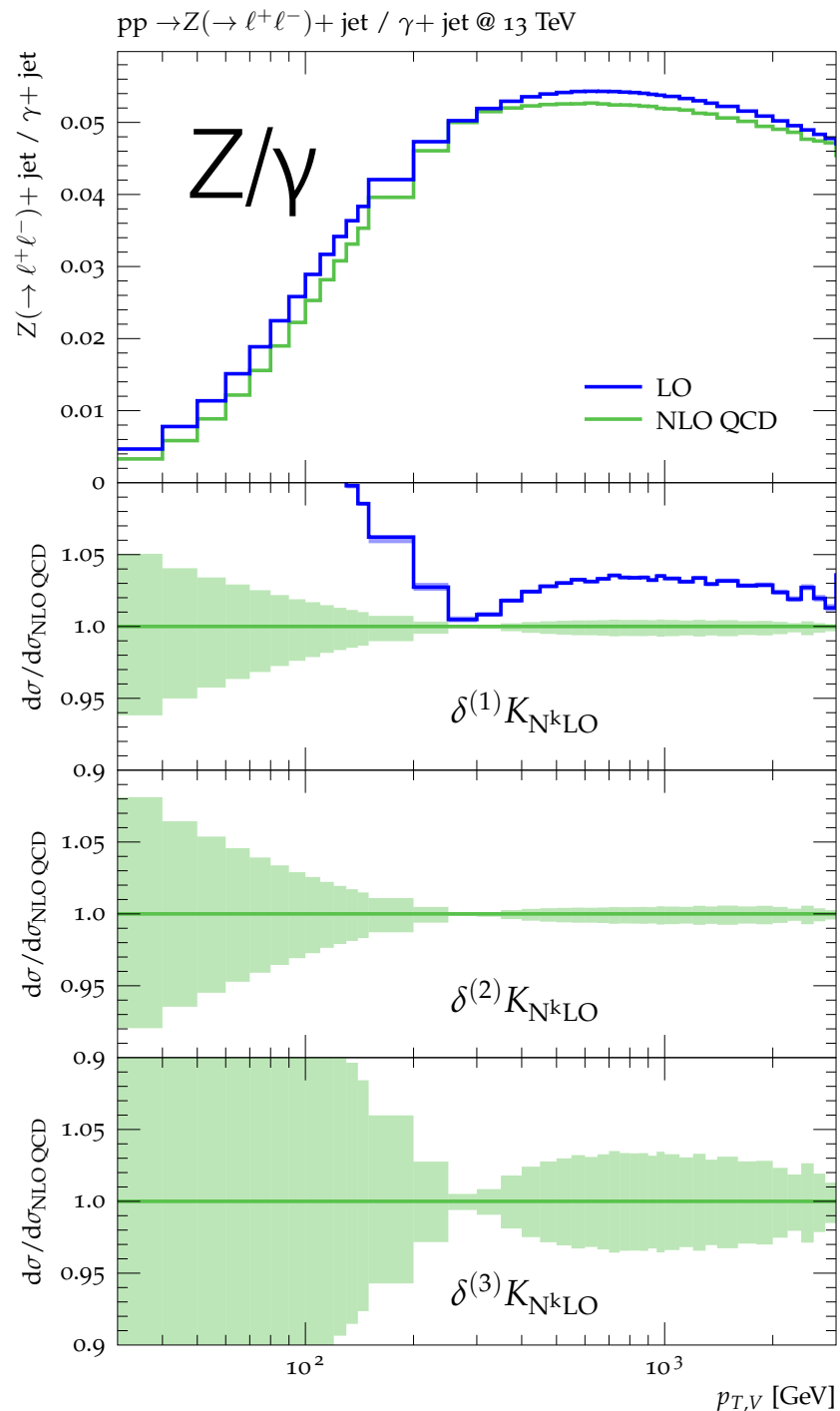
- $\delta^{(3)} K_{N^k \text{LO}}^V = \frac{K_{N^k \text{LO}}^V}{K_{N^{k-1} \text{LO}}^V} - \frac{K_{N^k \text{LO}}^Z}{K_{N^{k-1} \text{LO}}^Z}$

Difference of (N)NLO corrections as **process correlation uncertainty**

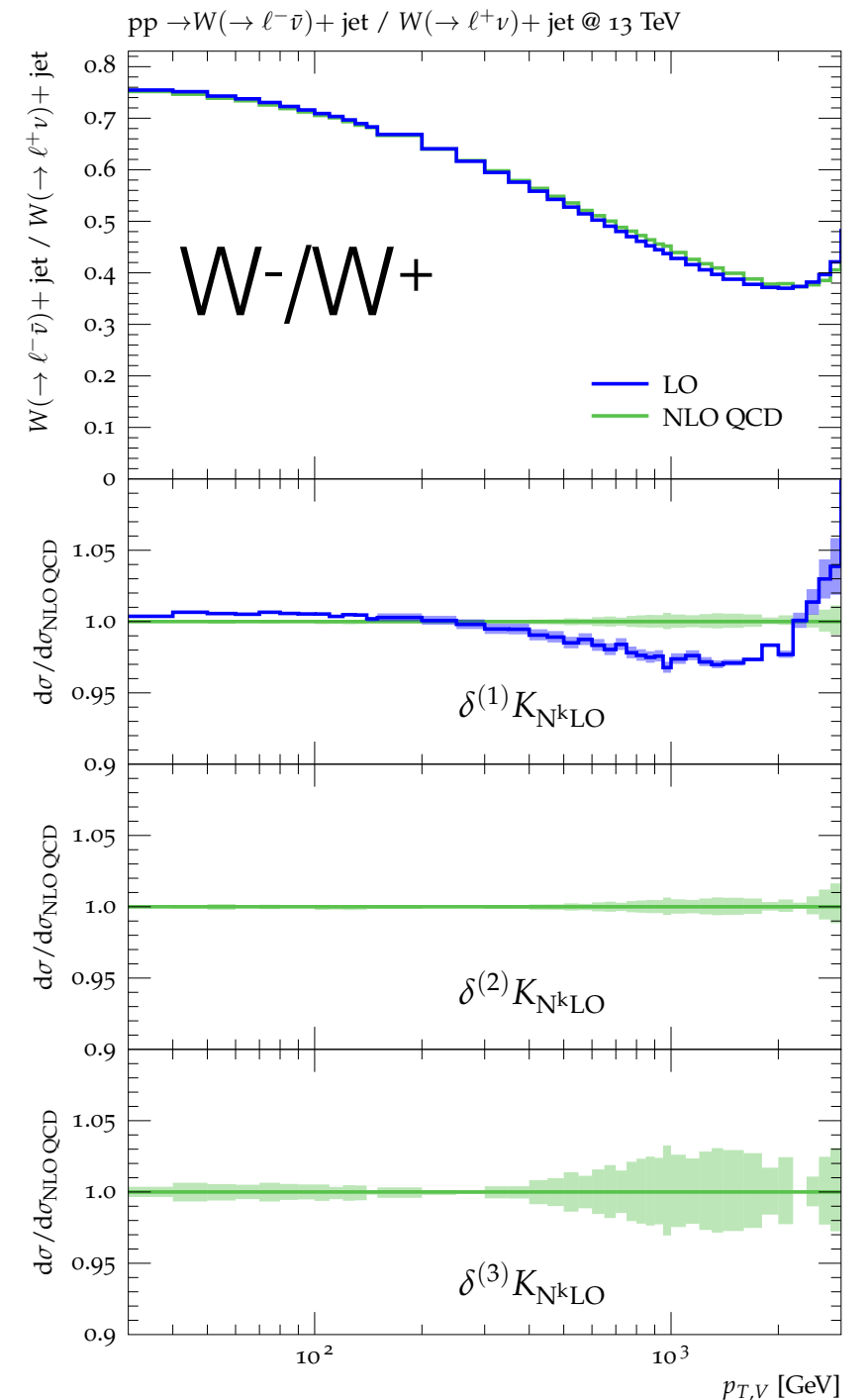
QCD uncertainties in pT-ratios



$Z/W^\pm \approx 1\%$



$Z/\gamma \approx 3\%$



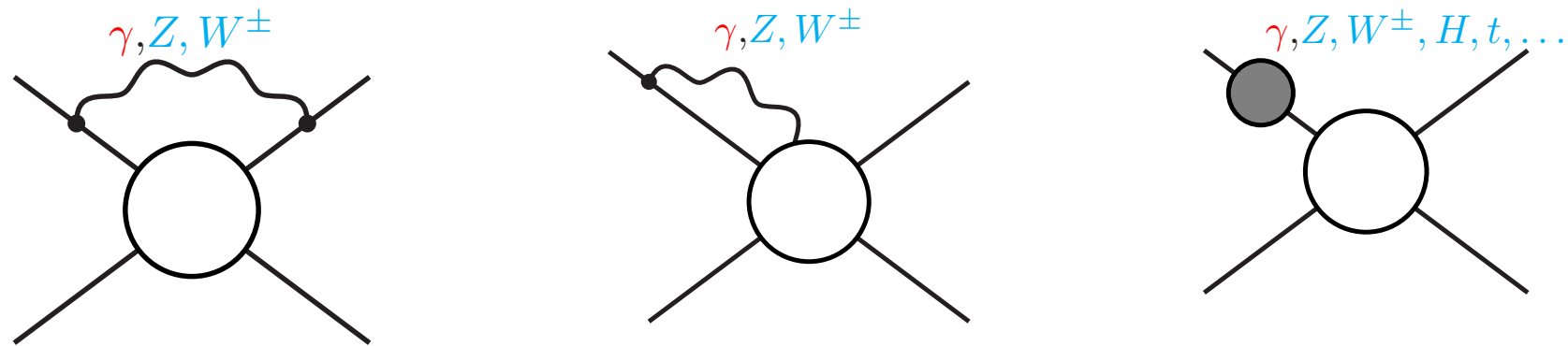
$W^+/W^- \approx 1-3\%$

08.07.06.05.04

2. pure EW effects uncertainties

Virtual EW Sudakov logarithms

Originate from soft/collinear virtual EW bosons coupling to on-shell legs



Universality and factorisation similar as in QCD [Denner, Pozzorini; '01]

$$\delta\mathcal{M}_{\text{LL+NLL}}^{1\text{-loop}} = \frac{\alpha}{4\pi} \sum_{k=1}^n \left\{ \frac{1}{2} \sum_{l \neq k} \sum_{a=\gamma, Z, W^\pm} I^a(k) I^{\bar{a}}(l) \ln^2 \left(\frac{\hat{s}_{kl}}{M_W^2} \right) + \gamma^{\text{ew}}(k) \ln \left(\frac{\hat{s}}{M_W^2} \right) \right\} \mathcal{M}_0$$

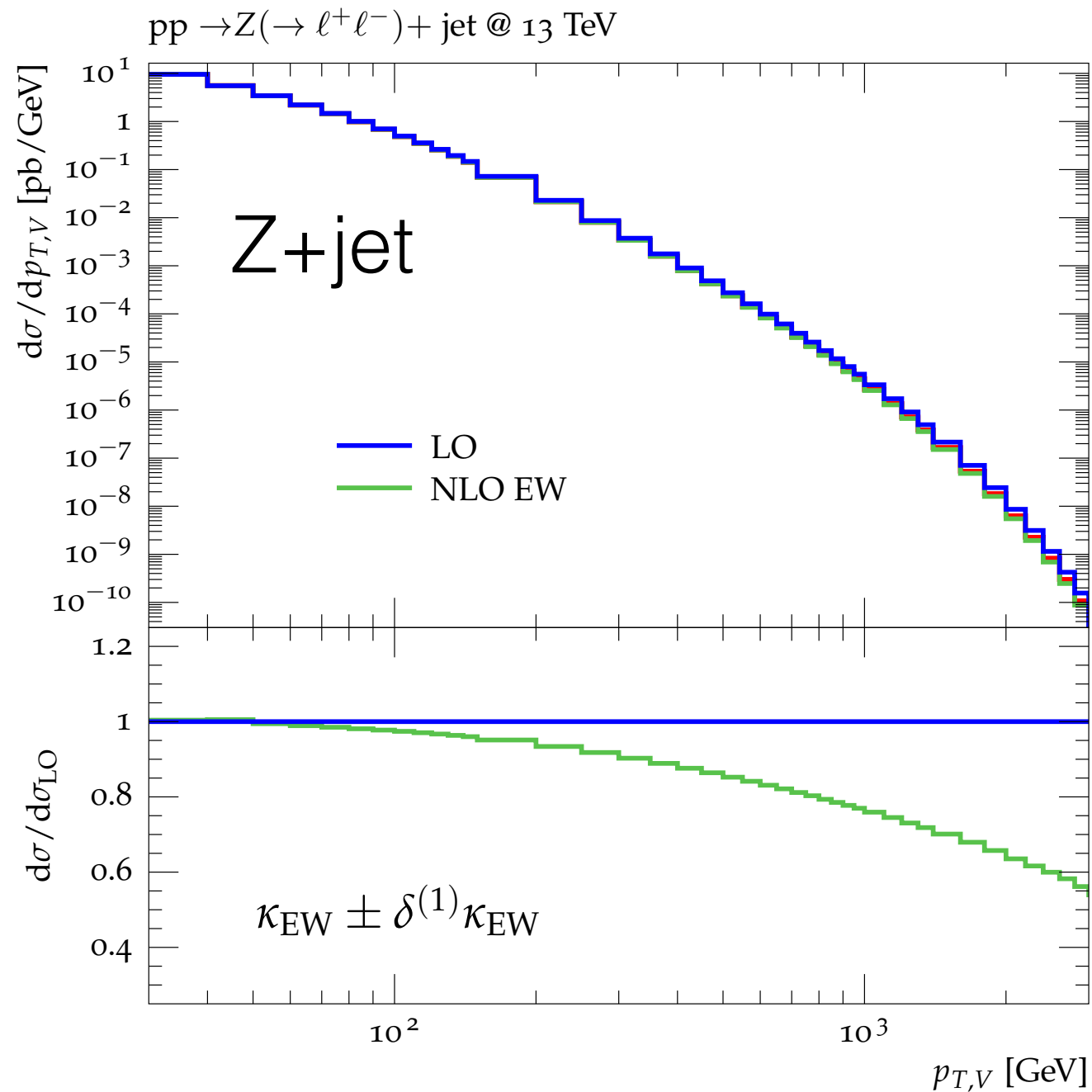
- process-independent, simple structure
- typical size at $\sqrt{\hat{s}} = 1, 5, 10 \text{ TeV}$:

$$\delta_{\text{LL}} \sim -\frac{\alpha}{\pi s_W^2} \log^2 \frac{\hat{s}}{M_W^2} \simeq -28, -76, -104\%,$$

$$\delta_{\text{NLL}} \sim +\frac{3\alpha}{\pi s_W^4} \log \frac{\hat{s}}{M_W^2} \simeq +16, +28, +32\%$$

- ➔ large (negative) corrections at high energies (pT, MET, HT, Minv)
- ➔ sizeable cancellations between leading and subleading terms possible

Pure EW uncertainties



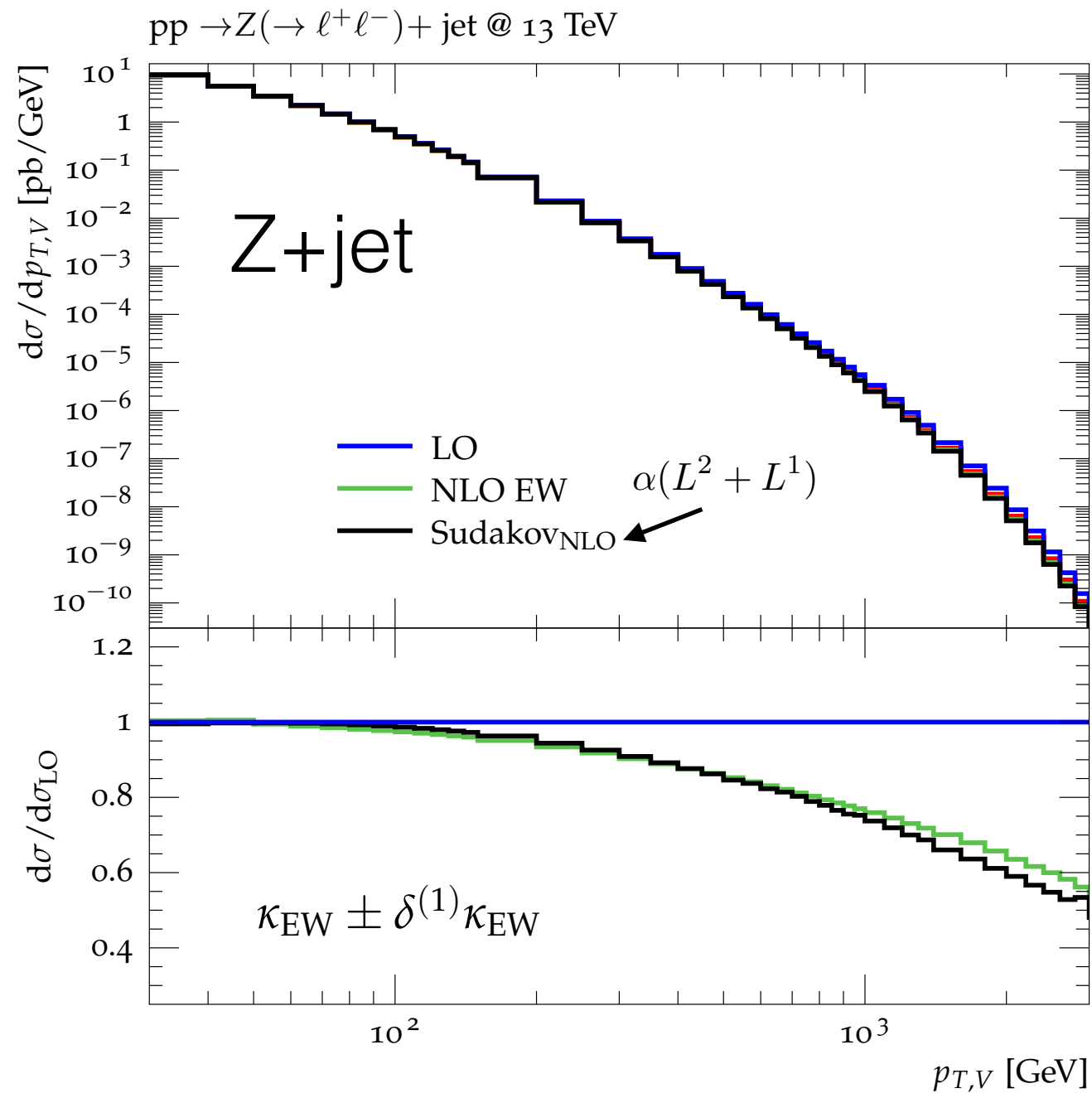
EW corrections become sizeable at large $p_{T,V}$

Origin: virtual EW Sudakov logarithms

Note: real EW Sudakov logarithms included as separate $VV(+jets)$ backgrounds

How to estimate corresponding pure EW uncertainties of relative $\mathcal{O}(\alpha^2)$?

Pure EW uncertainties



Large EW corrections
dominated by Sudakov logs

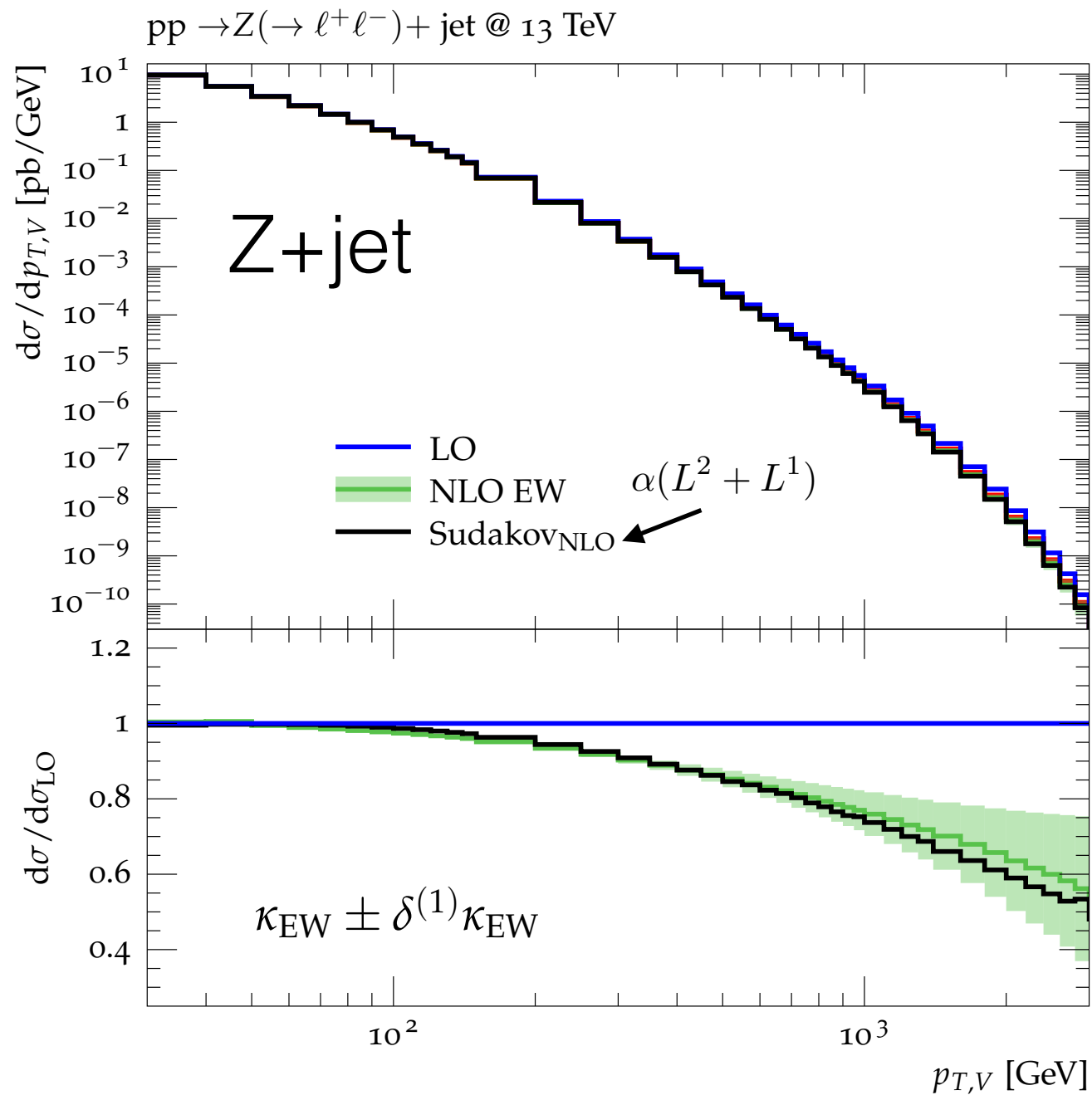


Uncertainty estimate of NLO EW
from naive exponentiation $\times 2$:

$$\delta^{(1)}\kappa_{EW} \simeq \frac{2}{k!} \left(\kappa_{NLO,EW} \right)^k$$

$$\kappa_{NLO,EW}(\hat{s}, \hat{t}) = \frac{\alpha}{\pi} \left[\delta_{\text{hard}}^{(1)} + \delta_{\text{Sud}}^{(1)} \right]$$

Pure EW uncertainties



Large EW corrections dominated by Sudakov logs

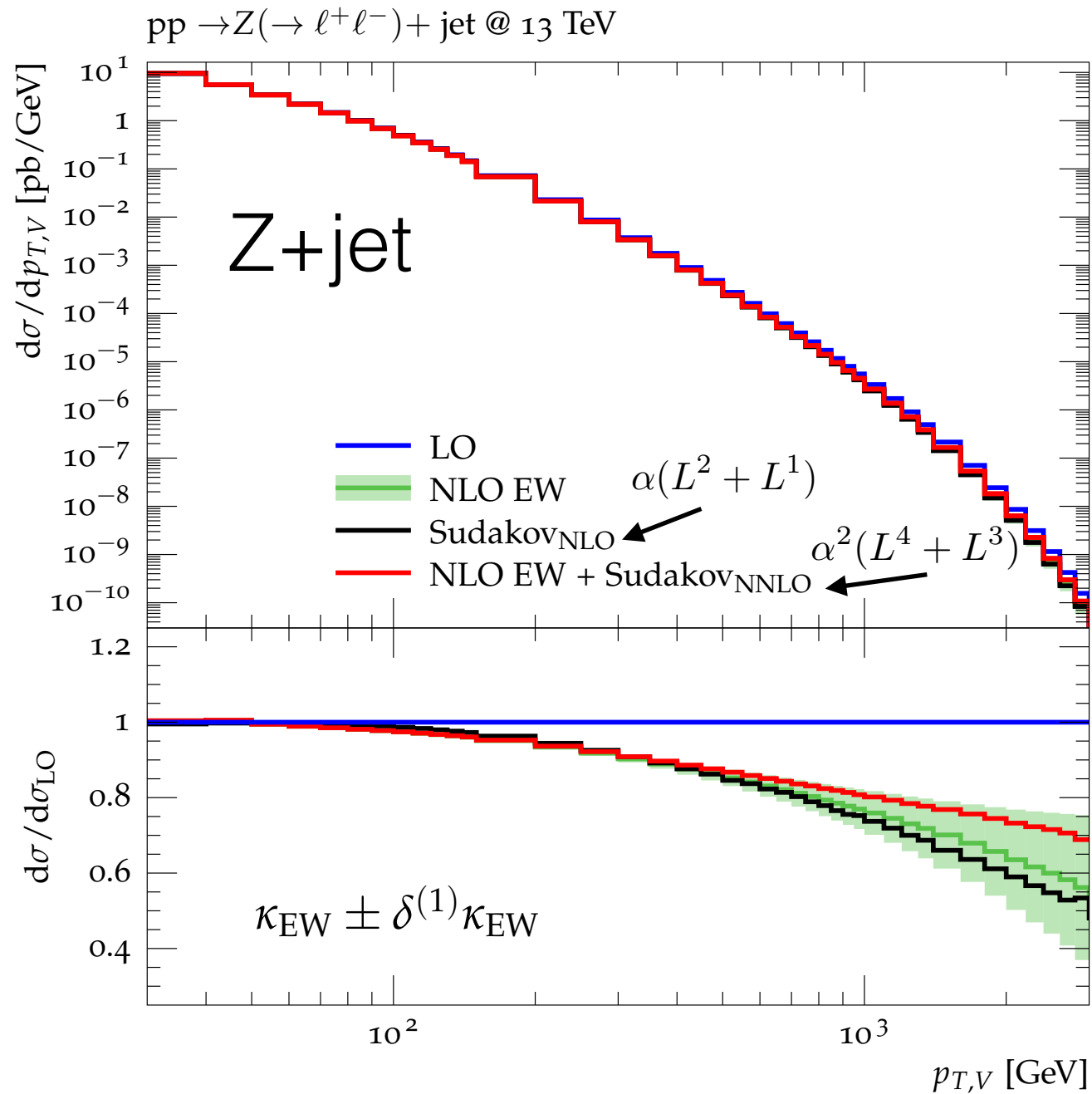


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Pure EW uncertainties



Large EW corrections dominated by Sudakov logs

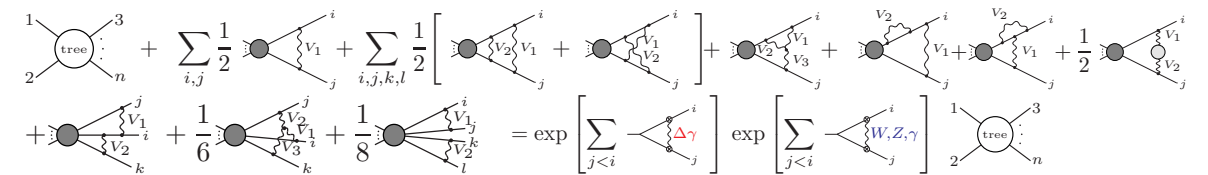


Uncertainty estimate of NLO EW from naive exponentiation $\times 2$:

$$\delta^{(1)}\kappa_{EW} \simeq \frac{2}{k!} \left(\kappa_{NLO,EW} \right)^k$$



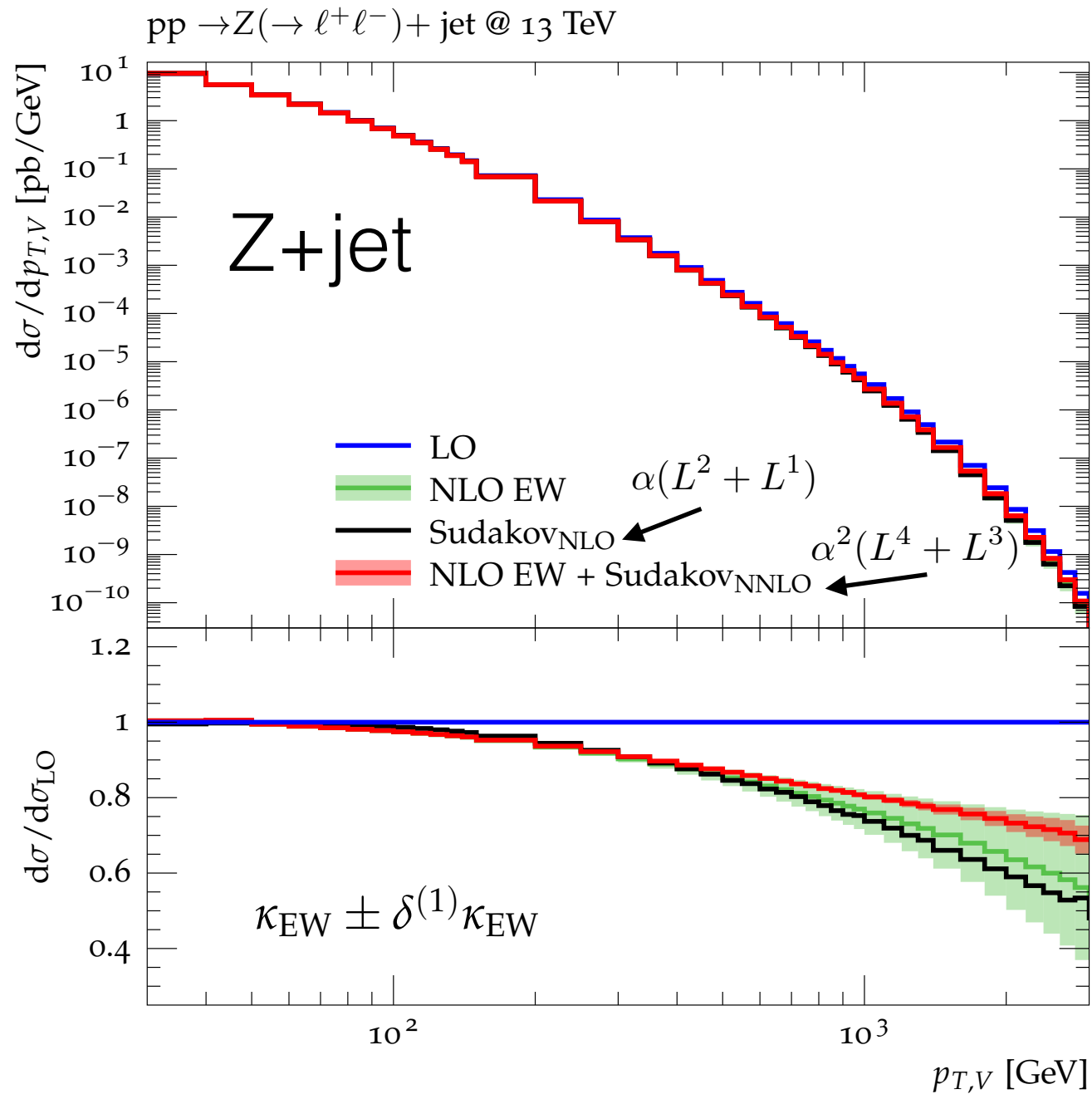
check against two-loop Sudakov logs
[Kühn, Kulesza, Pozzorini, Schulze; 05-07]



$$\kappa_{NLO,EW}(\hat{s}, \hat{t}) = \frac{\alpha}{\pi} \left[\delta_{\text{hard}}^{(1)} + \delta_{\text{Sud}}^{(1)} \right]$$

$$\kappa_{NNLO,Sud}(\hat{s}, \hat{t}) = \left(\frac{\alpha}{\pi} \right)^2 \delta_{\text{Sud}}^{(2)}$$

Pure EW uncertainties



Large EW corrections
dominated by Sudakov logs



Uncertainty estimate of NLO EW
from naive exponentiation $\times 2$:

$$\delta^{(1)}\kappa_{EW} \simeq \frac{2}{k!} \left(\kappa_{\text{NLO,EW}} \right)^k$$



check against two-loop Sudakov logs
[Kühn, Kulesza, Pozzorini, Schulze; 05-07]



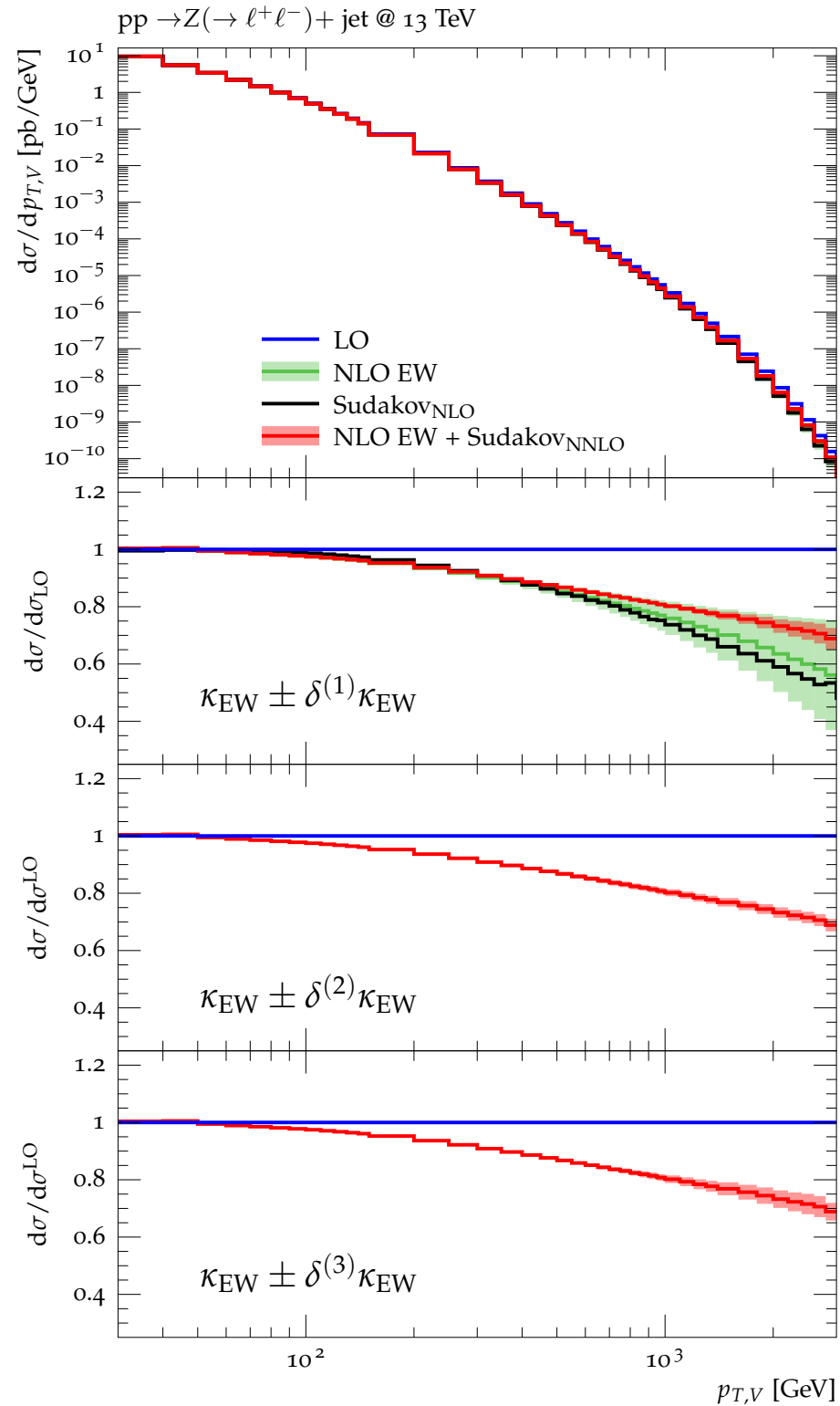
Uncertainty estimate of NNLO EW:

$$\delta^{(1)}\kappa_{EW}^{(V)}(x) = \frac{2}{3} \kappa_{\text{NLO,EW}}^{(V)}(x) \kappa_{\text{NNLO,Sud}}^{(V)}(x)$$

$$\kappa_{\text{NLO,EW}}(\hat{s}, \hat{t}) = \frac{\alpha}{\pi} \left[\delta_{\text{hard}}^{(1)} + \delta_{\text{Sud}}^{(1)} \right]$$

$$\kappa_{\text{NNLO,Sud}}(\hat{s}, \hat{t}) = \left(\frac{\alpha}{\pi} \right)^2 \delta_{\text{Sud}}^{(2)}$$

Pure EW uncertainties



- 'higher-order logs'

$$\delta^{(1)} \kappa_{\text{EW}}^{(V)}(x) = \frac{2}{3} \kappa_{\text{NLO EW}}^{(V)}(x) \kappa_{\text{NNLO Sud}}^{(V)}(x)$$

(correlated)

Additional uncorrelated uncertainties:

- 'hard non-log NNLO effects I'

$$\delta^{(2)} \kappa_{\text{EW}}^{(V)}(x) = 0.05 \kappa_{\text{NLO EW}}^{(V)}(x)$$

(uncorrelated)

$$\Leftrightarrow \delta_{\text{hard}}^{(2)} \leq \frac{0.05\pi}{\alpha} \delta_{\text{hard}}^{(1)} \simeq 20 \delta_{\text{hard}}^{(1)}$$

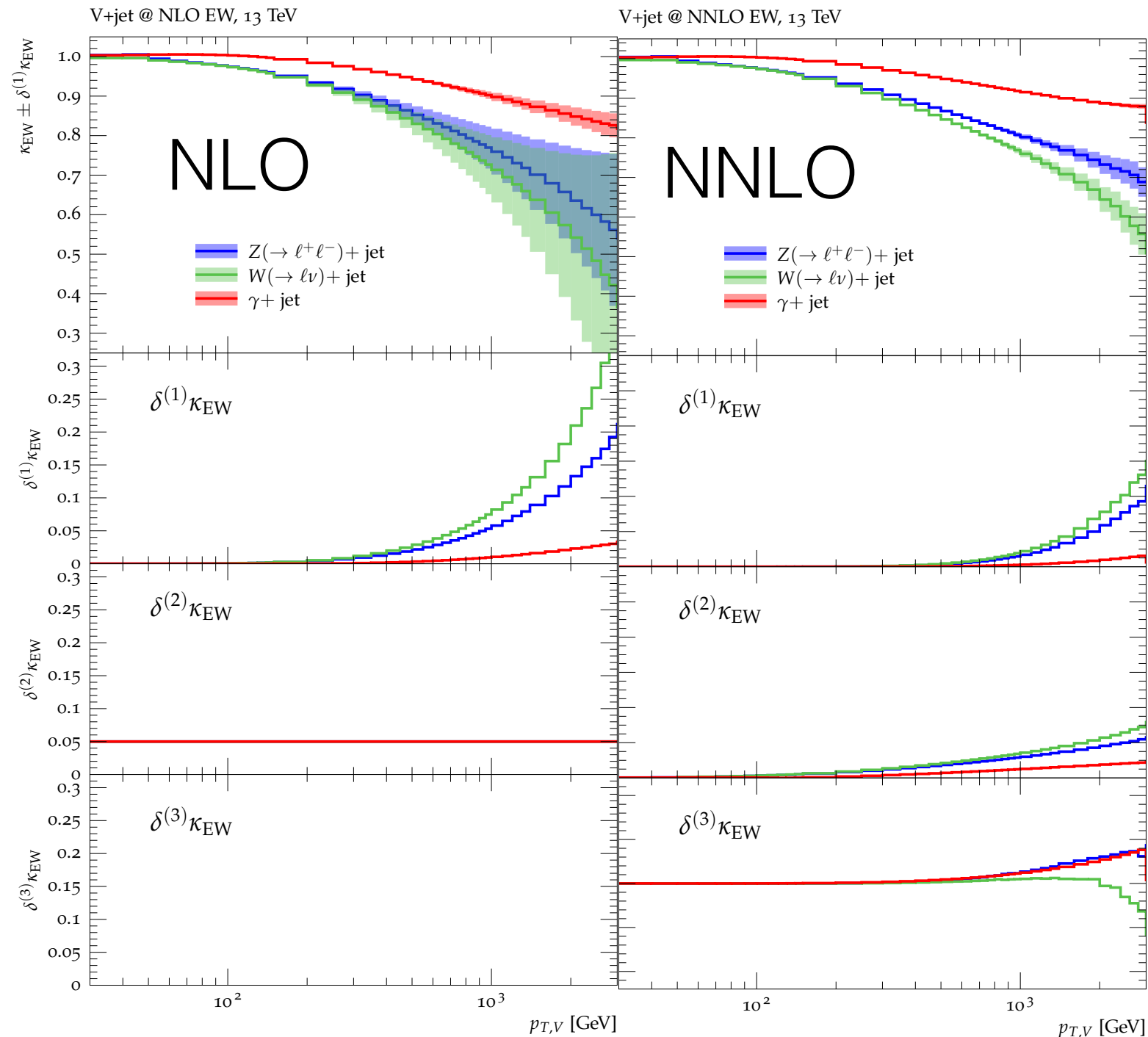
- 'hard non-log NNLO effects II'

$$\delta^{(3)} \kappa_{\text{EW}}^{(V)}(x) = \kappa_{\text{NNLO Sud}}^{(V)}(x) - \frac{1}{2} [\kappa_{\text{NLO EW}}^{(V)}(x)]^2$$

(uncorrelated)

estimate of typical size of $[\delta_{\text{hard}}^{(1)}]^2$ or $\delta_{\text{hard}}^{(1)} \times \delta_{\text{Sud}}^{(1)}$.

Pure EW uncertainties



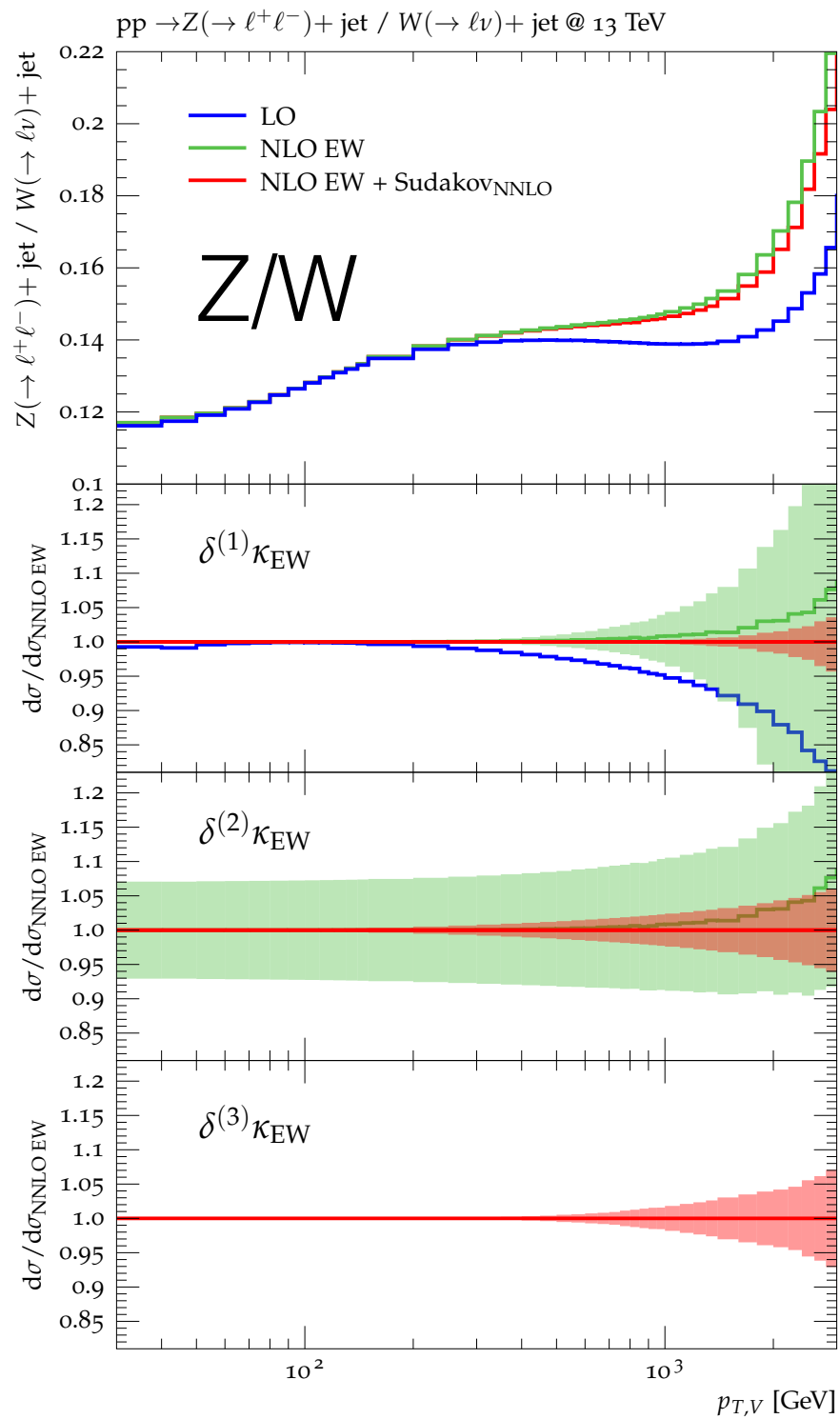
NNLO EW corrections at 1 TeV

- γ +jets: -10%
- Z+jet: -20%
- W+jet: -25%

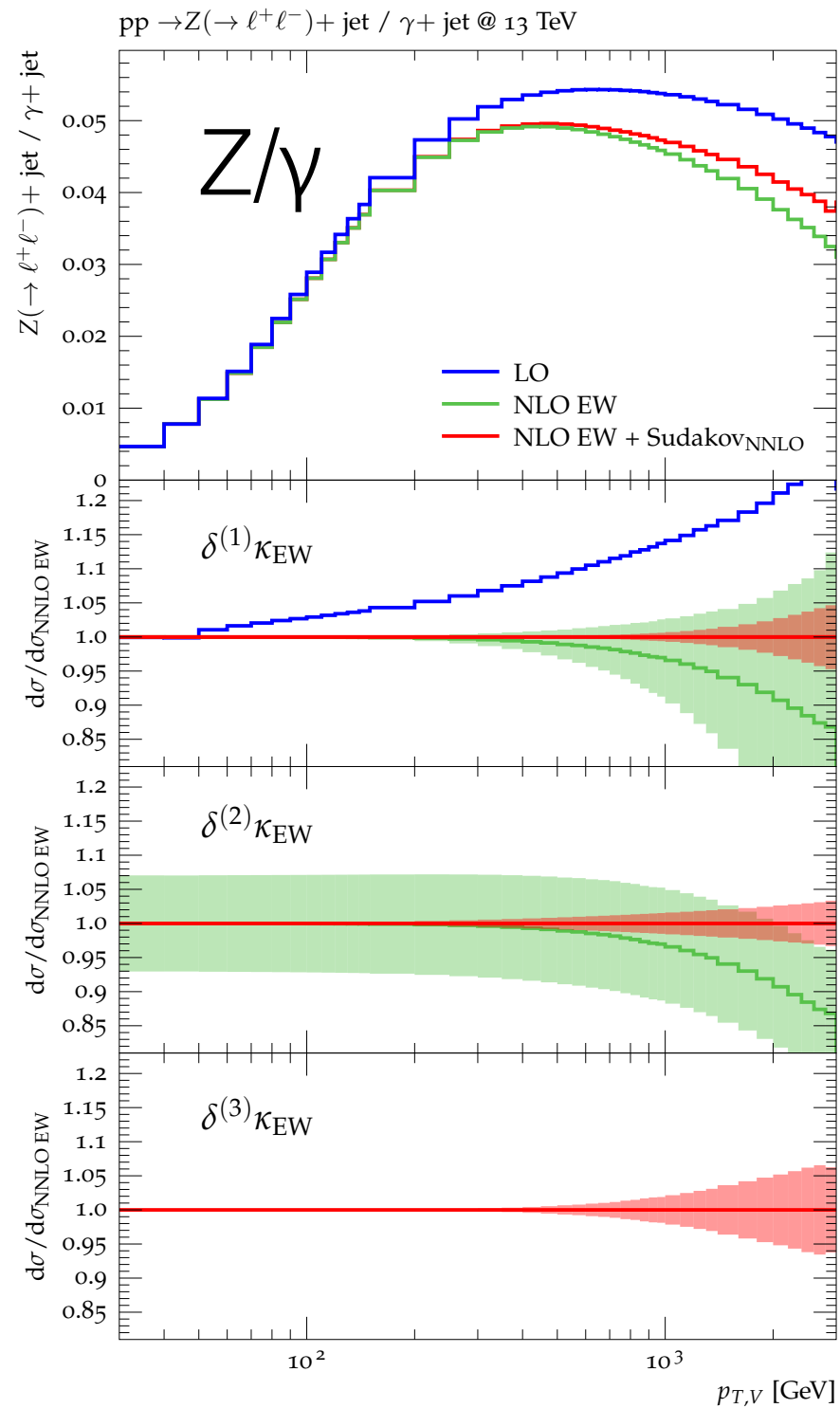
Pure EW uncertainties

- tiny at low p_T and only 1-2% at 1 TeV
- thanks to NNLO Sudakov logs (up to ~ 5%)

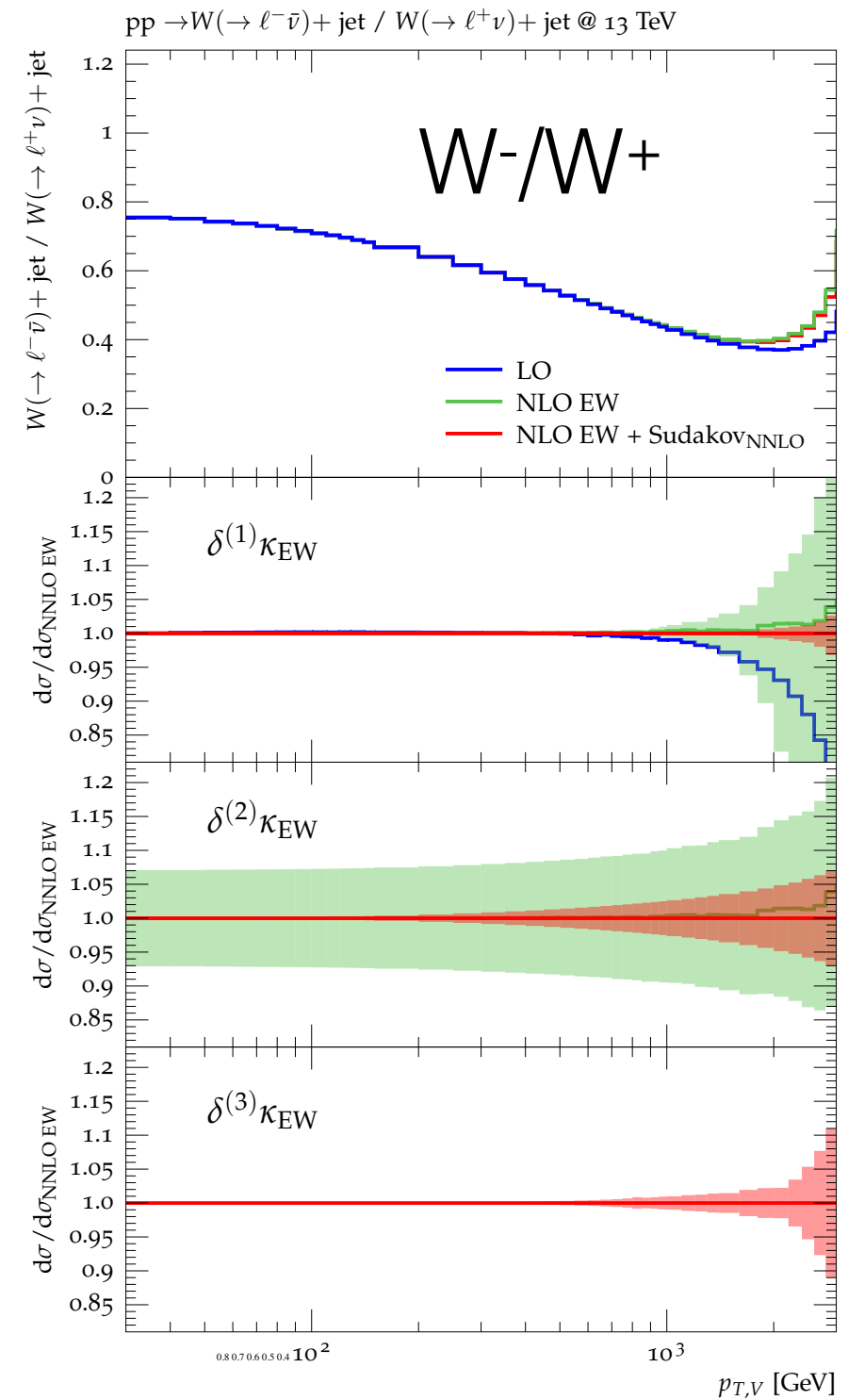
EW uncertainties in pT-ratios



$Z/W^\pm \approx 2-3\%$



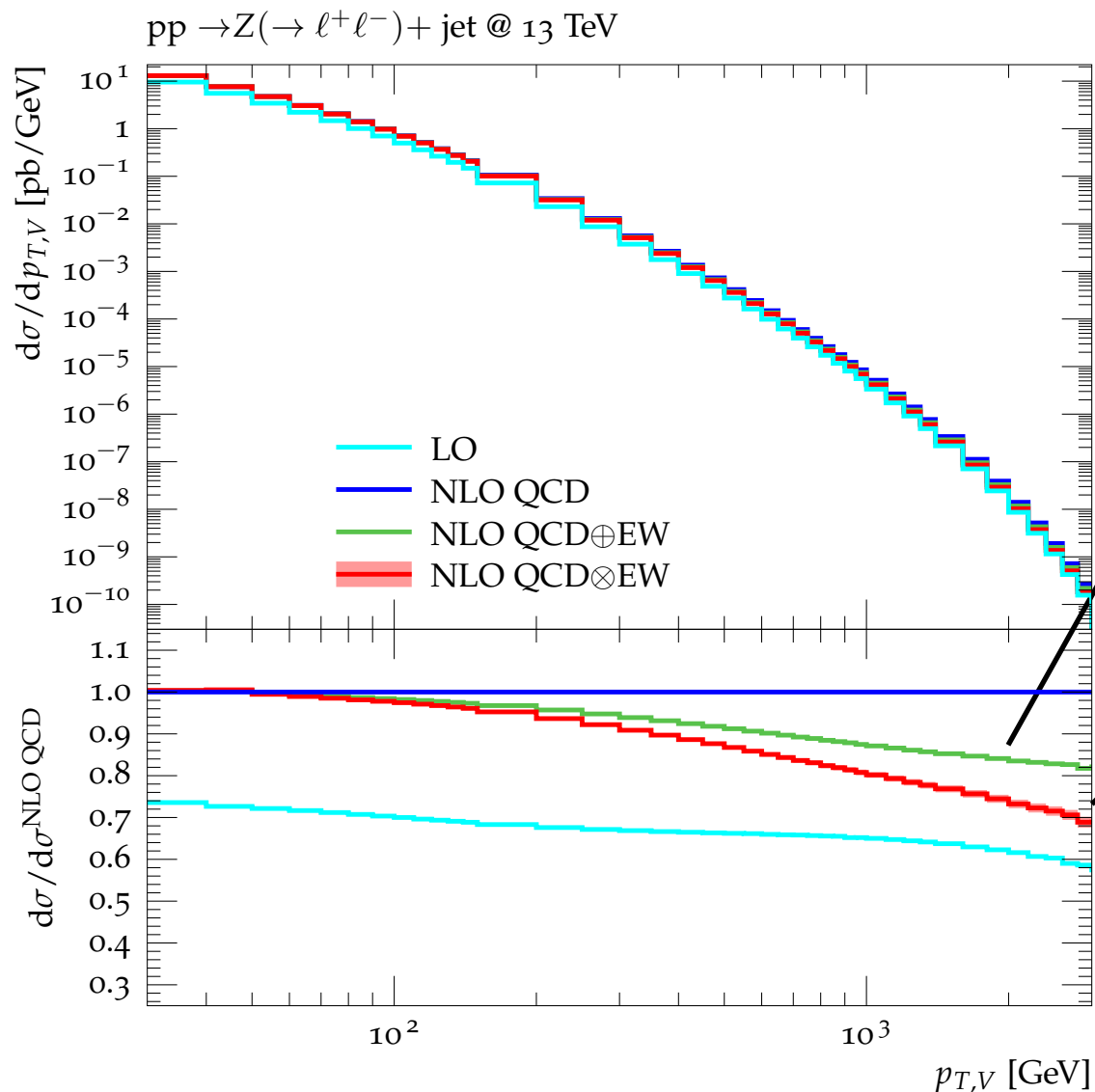
$Z/\gamma \approx 2-3\%$



$W^+/W^- \approx 2-3\%$

3. mixed QCD-EW uncertainties

Mixed QCD-EW uncertainties



Given QCD and EW corrections are sizeable, also mixed QCD-EW uncertainties of relative $\mathcal{O}(\alpha\alpha_s)$ have to be considered.

Additive combination

$$\sigma_{\text{QCD+EW}}^{\text{NLO}} = \sigma^{\text{LO}} + \delta\sigma_{\text{QCD}}^{\text{NLO}} + \delta\sigma_{\text{EW}}^{\text{NLO}}$$

(no $\mathcal{O}(\alpha\alpha_s)$ contributions)

Multiplicative combination

$$\sigma_{\text{QCD}\times\text{EW}}^{\text{NLO}} = \sigma_{\text{QCD}}^{\text{NLO}} \left(1 + \frac{\delta\sigma_{\text{EW}}^{\text{NLO}}}{\sigma^{\text{LO}}} \right) :$$

(some dominant $\mathcal{O}(\alpha\alpha_s)$ contributions, e.g. EW Sudakov logs \times soft QCD)

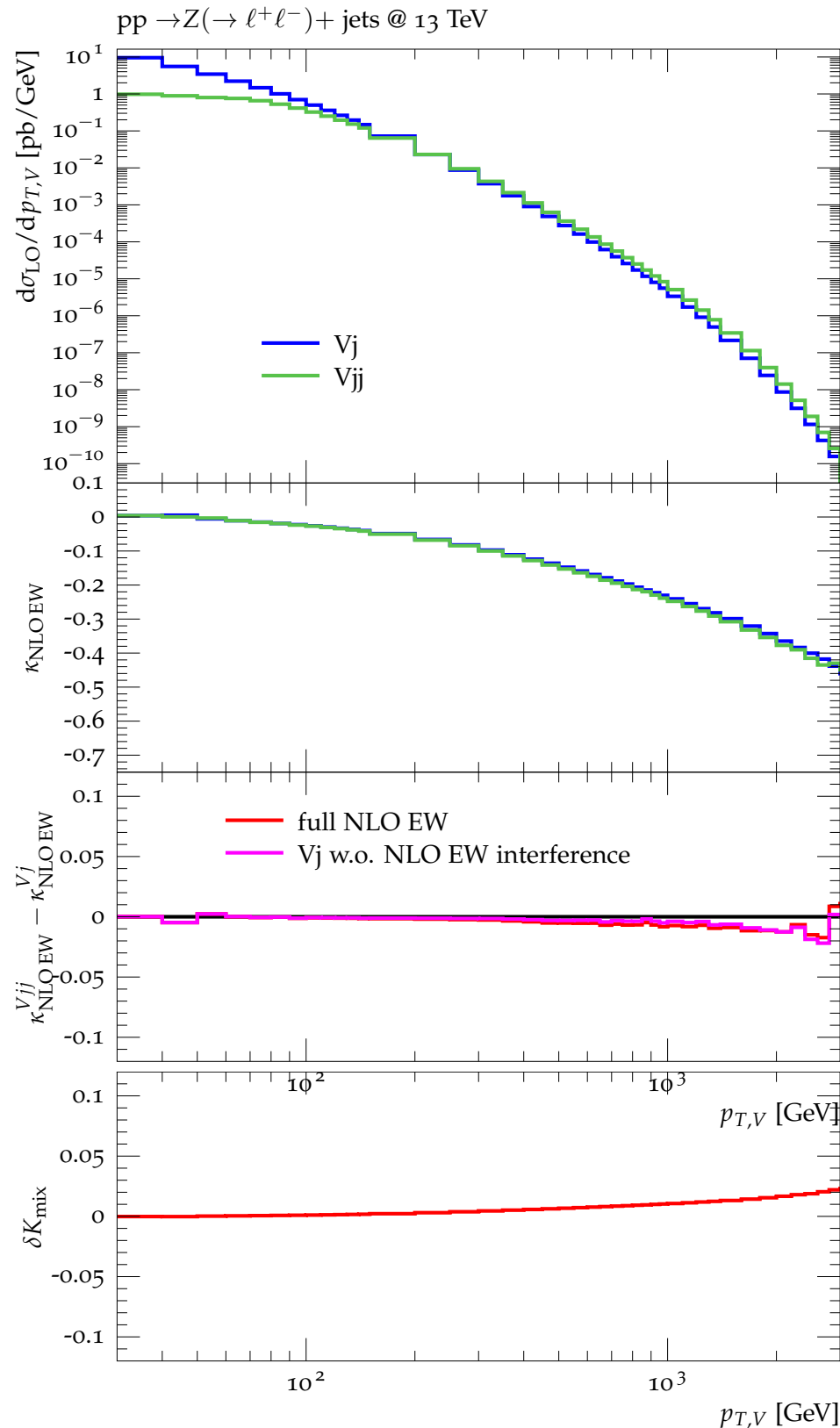
Difference between these two approaches indicates size of missing mixed EW-QCD corrections.

$$K_{\text{QCD}\otimes\text{EW}} - K_{\text{QCD}\oplus\text{EW}} \sim 10\% \quad \text{at 1 TeV}$$

Too conservative!?

For dominant Sudakov EW logarithms factorization should be exact!

Mixed QCD-EW uncertainties



Bold estimate:

Consider real $\mathcal{O}(\alpha\alpha_s)$ correction to V+jet

\simeq NLO EW to V+2jets

and we observe

$$\left. \frac{d\sigma_{\text{NLO EW}}}{d\sigma_{\text{LO}}} \right|_{V+2\text{jet}} - \left. \frac{d\sigma_{\text{NLO EW}}}{d\sigma_{\text{LO}}} \right|_{V+1\text{jet}} \simeq 1\%$$

strong support for

- factorization
- multiplicative QCD \times EW combination

Estimate of non-factorising contributions

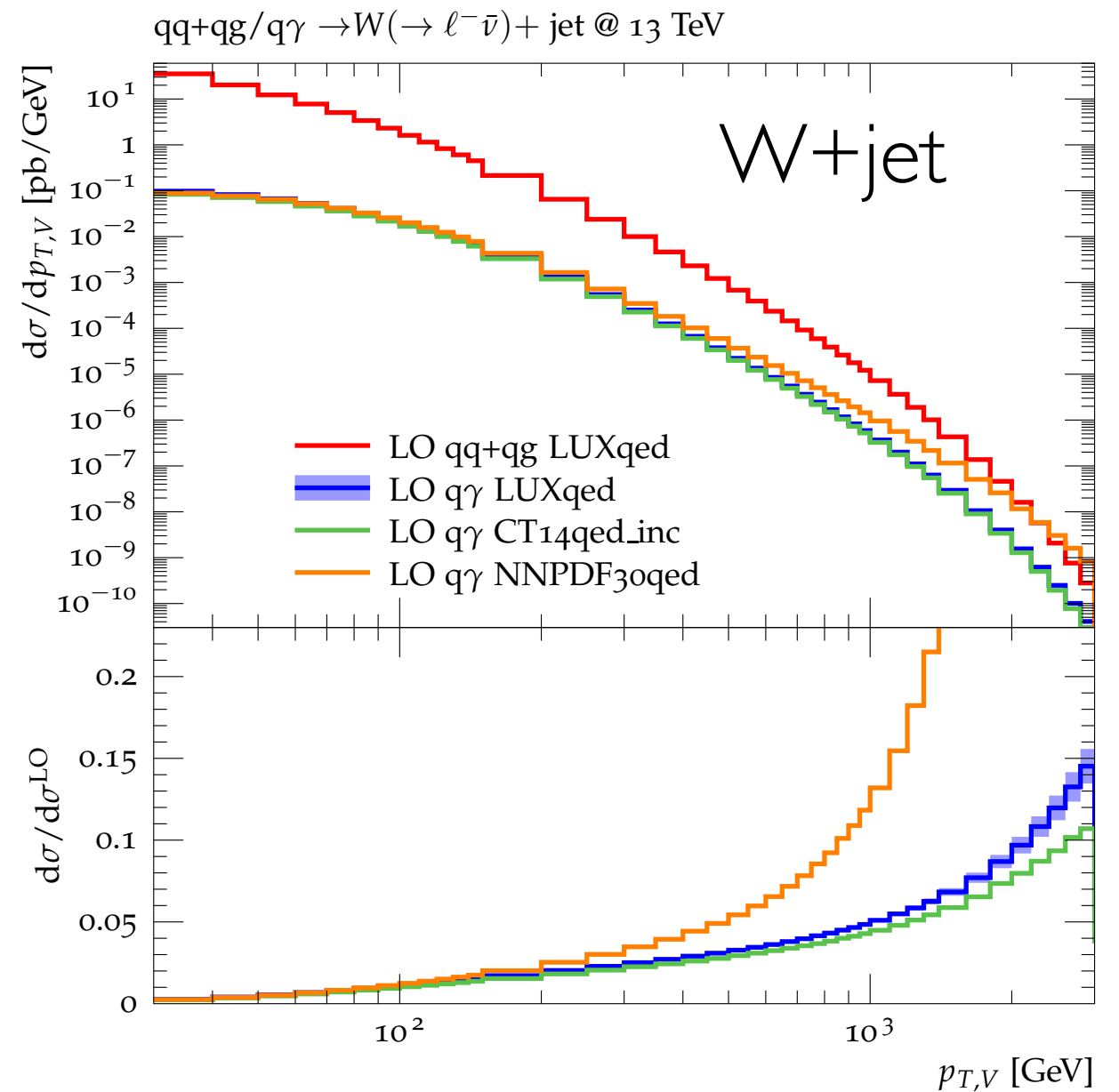
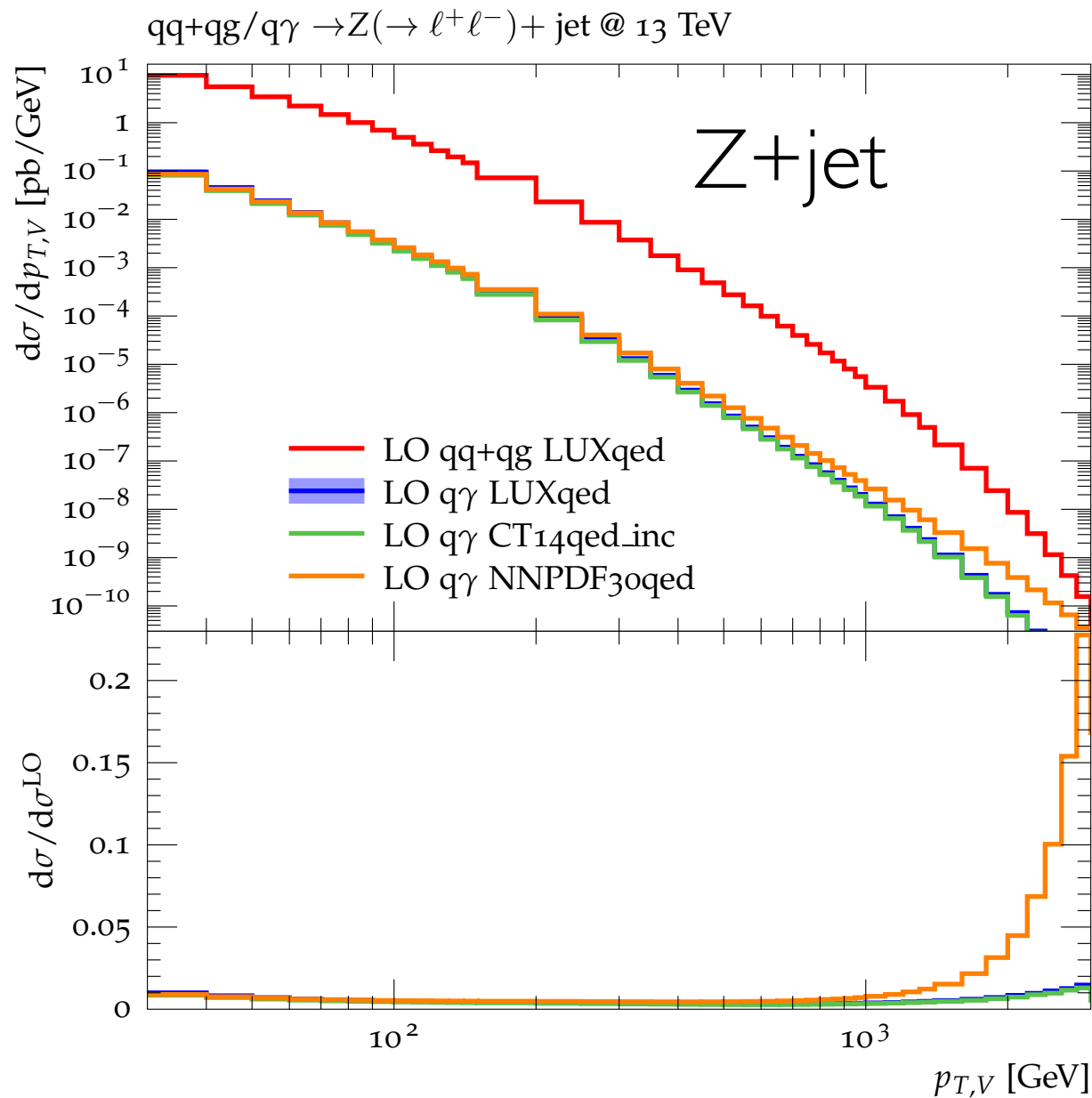
(correlated)

$$\leftarrow \delta K_{\text{mix}}^{(V)}(x) = 0.1 \left[K_{\text{TH},\oplus}^{(V)}(x, \vec{\mu}_0) - K_{\text{TH},\otimes}^{(V)}(x, \vec{\mu}_0) \right]$$

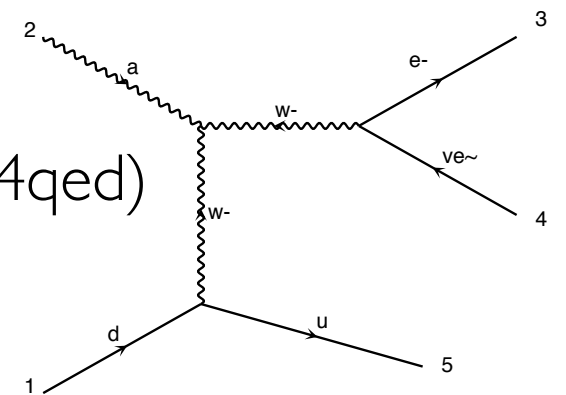
(tuned to cover above difference of EW K-factors)

4. Other issues (PDFs, γ -induced)

Photon-induced production

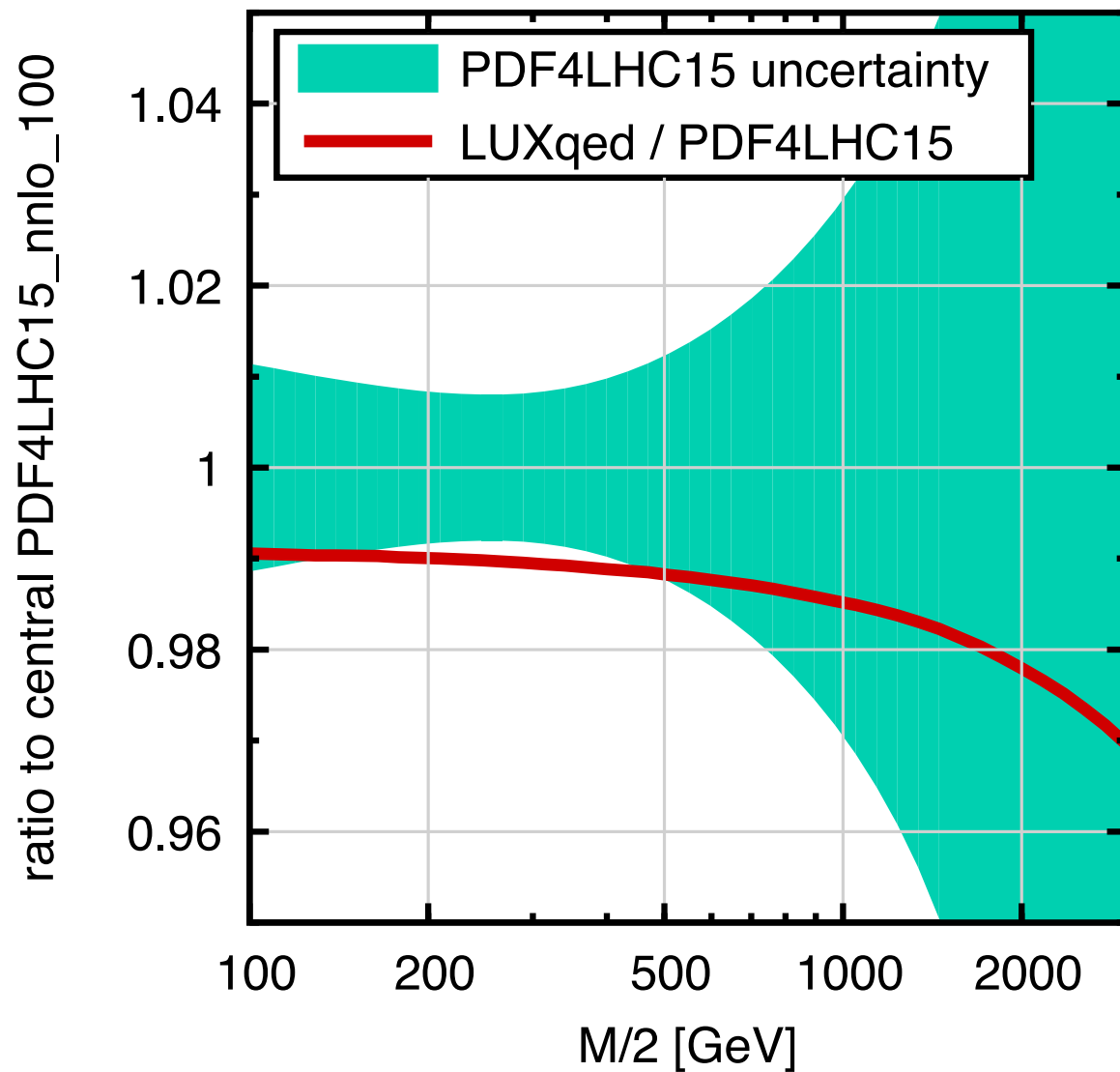


- photon-induced production irrelevant for Z+jet (and γ +jet)
- in W+jet $\mathcal{O}(5\%)$ contribution with LUXqed (consistent with CT14qed) (due to t-channel enhancement)
- $\sim 1\%$ uncertainties in photon PDFs due to LUXqed

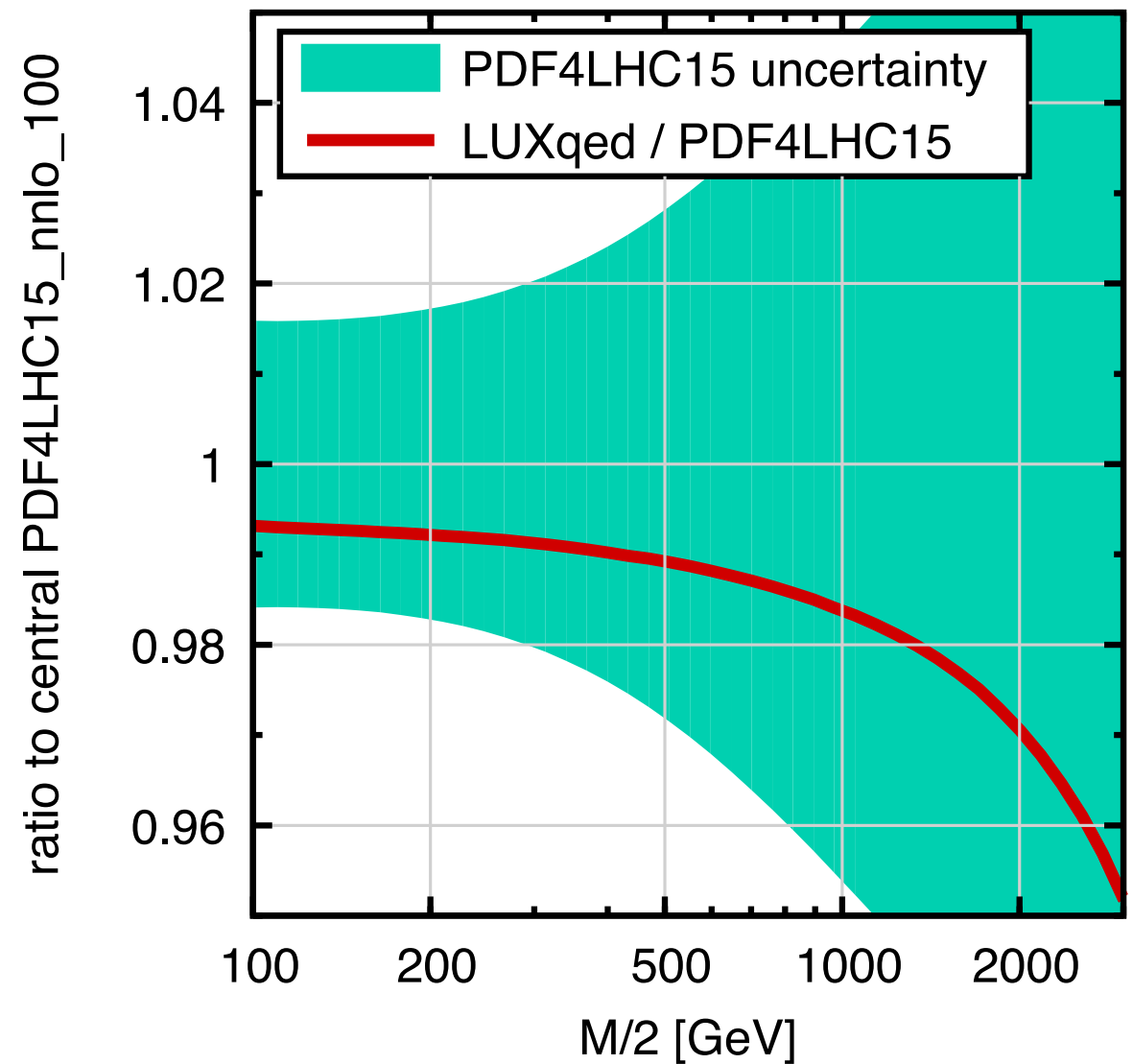


PDFs

QED effects on $(g\Sigma)$ luminosity



QED effects on $(q\bar{q})$ luminosity



- small percent-level QED effects on $qg/q\bar{q}$ luminosities (included via LUXqed)
- 1.5-5% PDF uncertainties

Conclusions

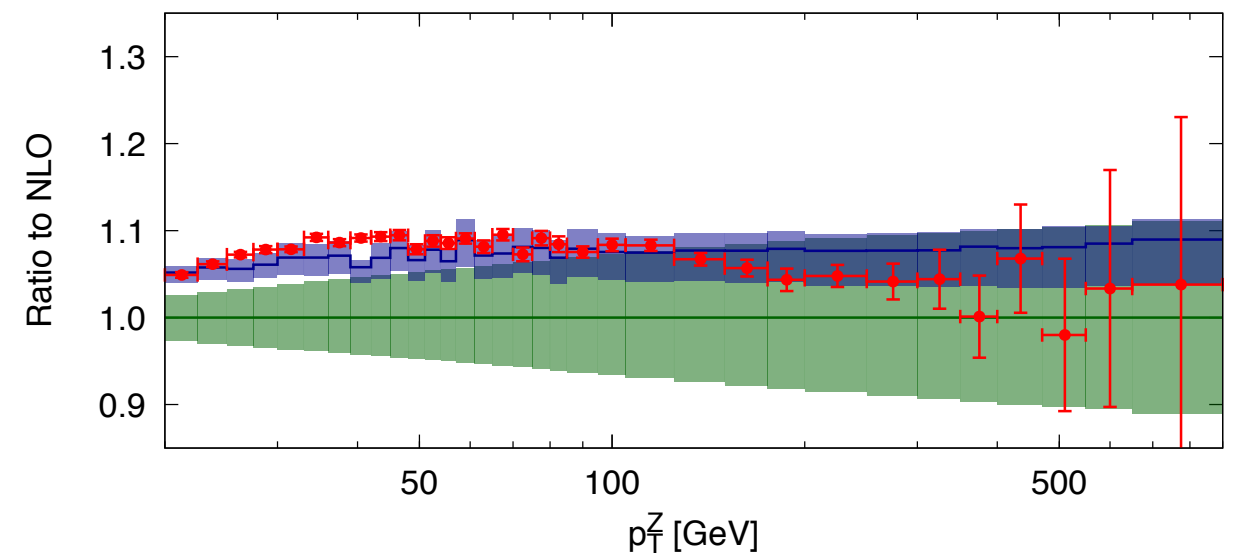
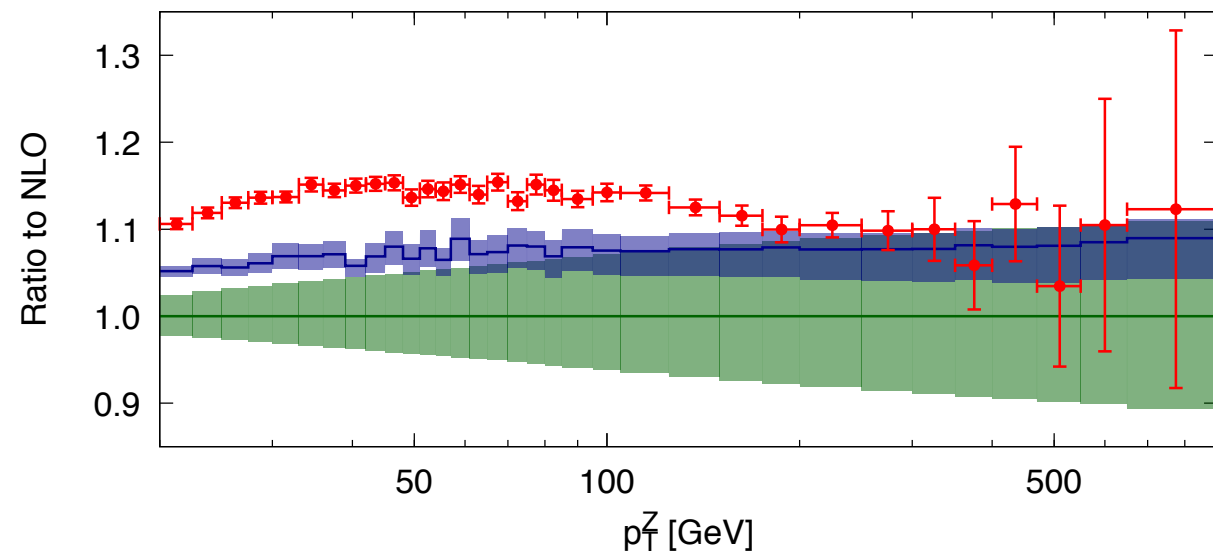
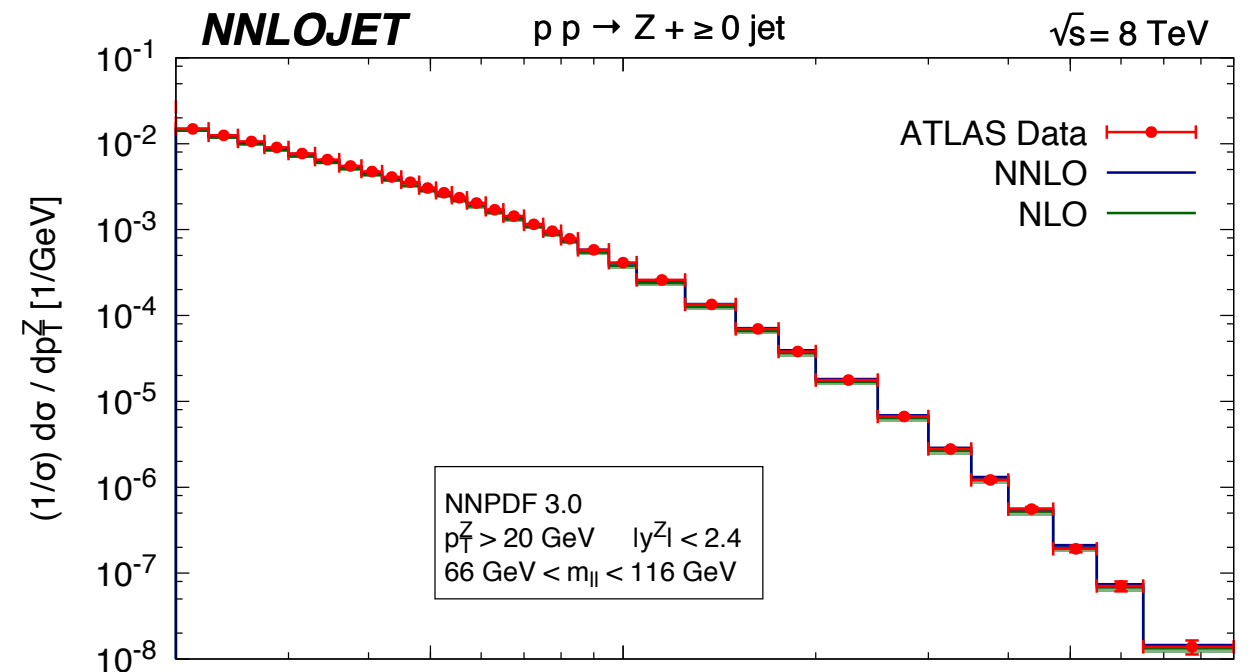
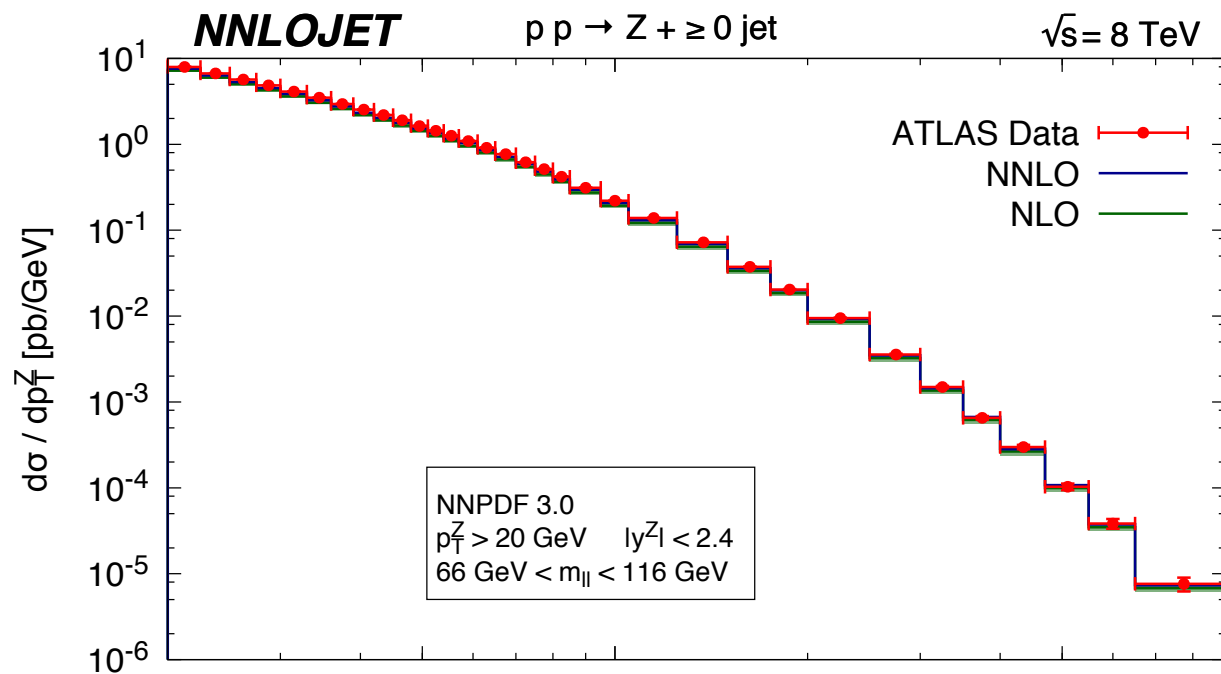
- ▶ monojet / MET+jets searches *soon* limited by V+jets background systematics
- ▶ MC reweighting allows to promote V + jet to NNLO QCD+(N)NLO EW:
 - inclusion of EW corrections *crucial* due to large Sudakov logs
- ▶ Perturbative systematics in pTV under control at the level of 1-10% up to the TeV

Outlook / Questions

- ▶ percent precision requires scrutiny of many subtleties and close TH/EXP interplay
- ▶ Experimental closure tests in control regions
- ▶ Uncertainty estimates applicable to other more exclusive observables?
- ▶ Framework applicable to other process classes?
- ▶ Utilisation in PDF fits?

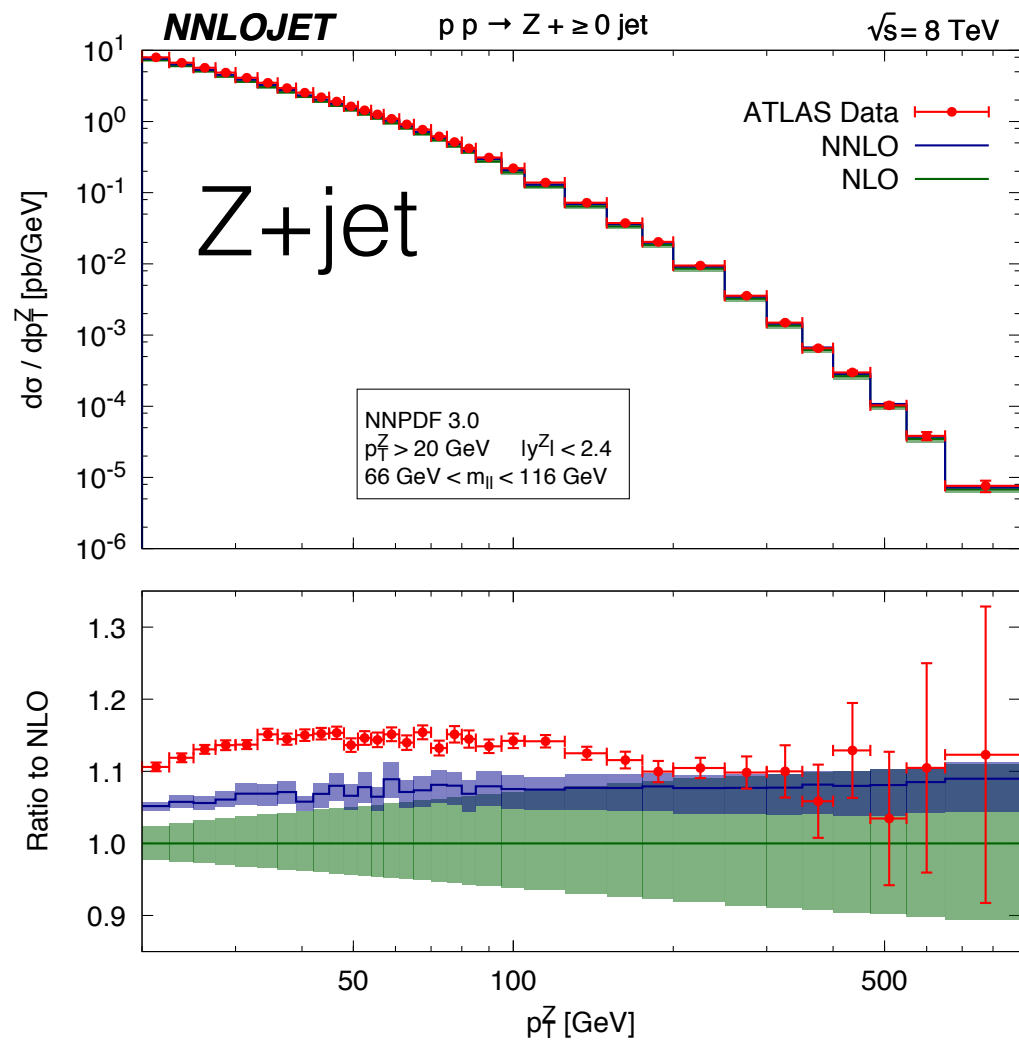
BACKUP

NNLO for Z+jet

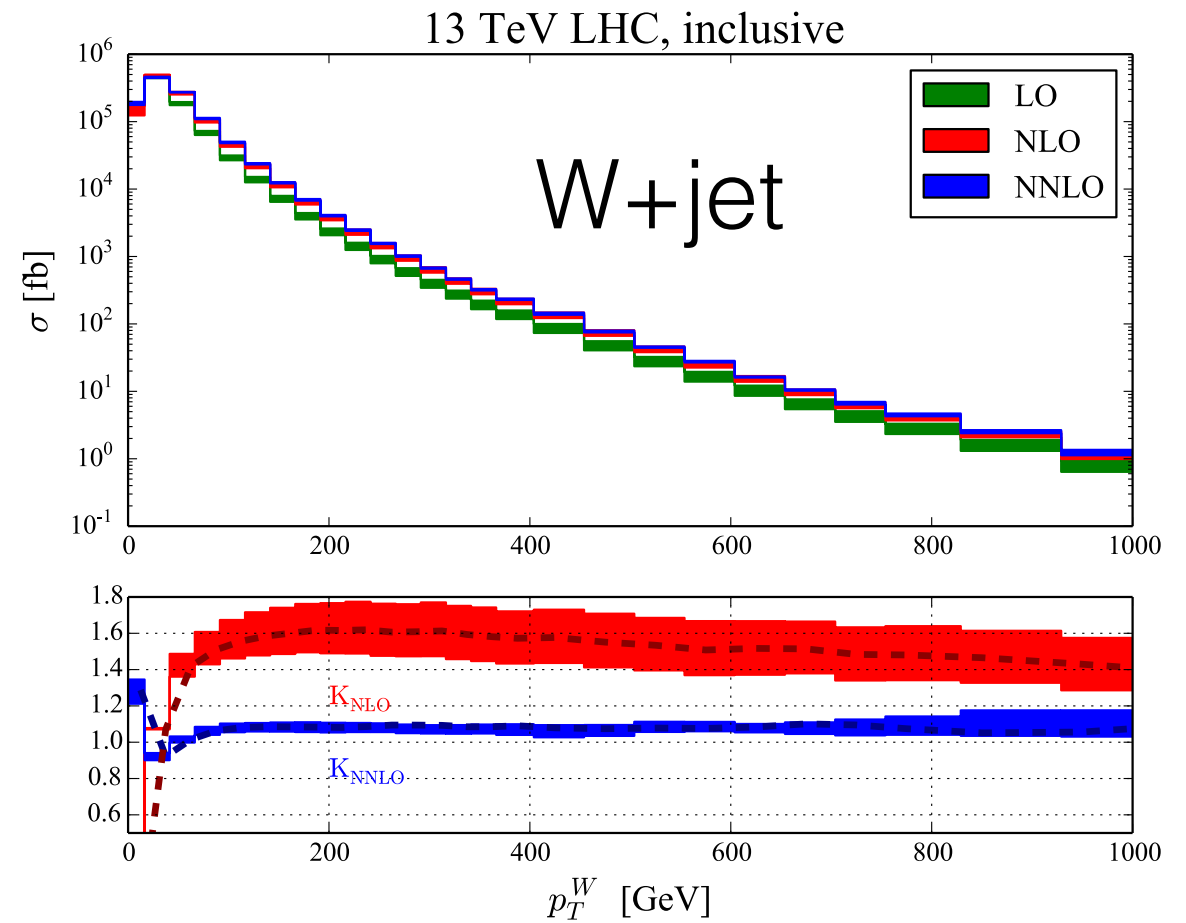


[Gehrmann-De Ridder, Gehrmann, Glover, A. Huss, Morgan; '16]

NNLO for W/Z+jet



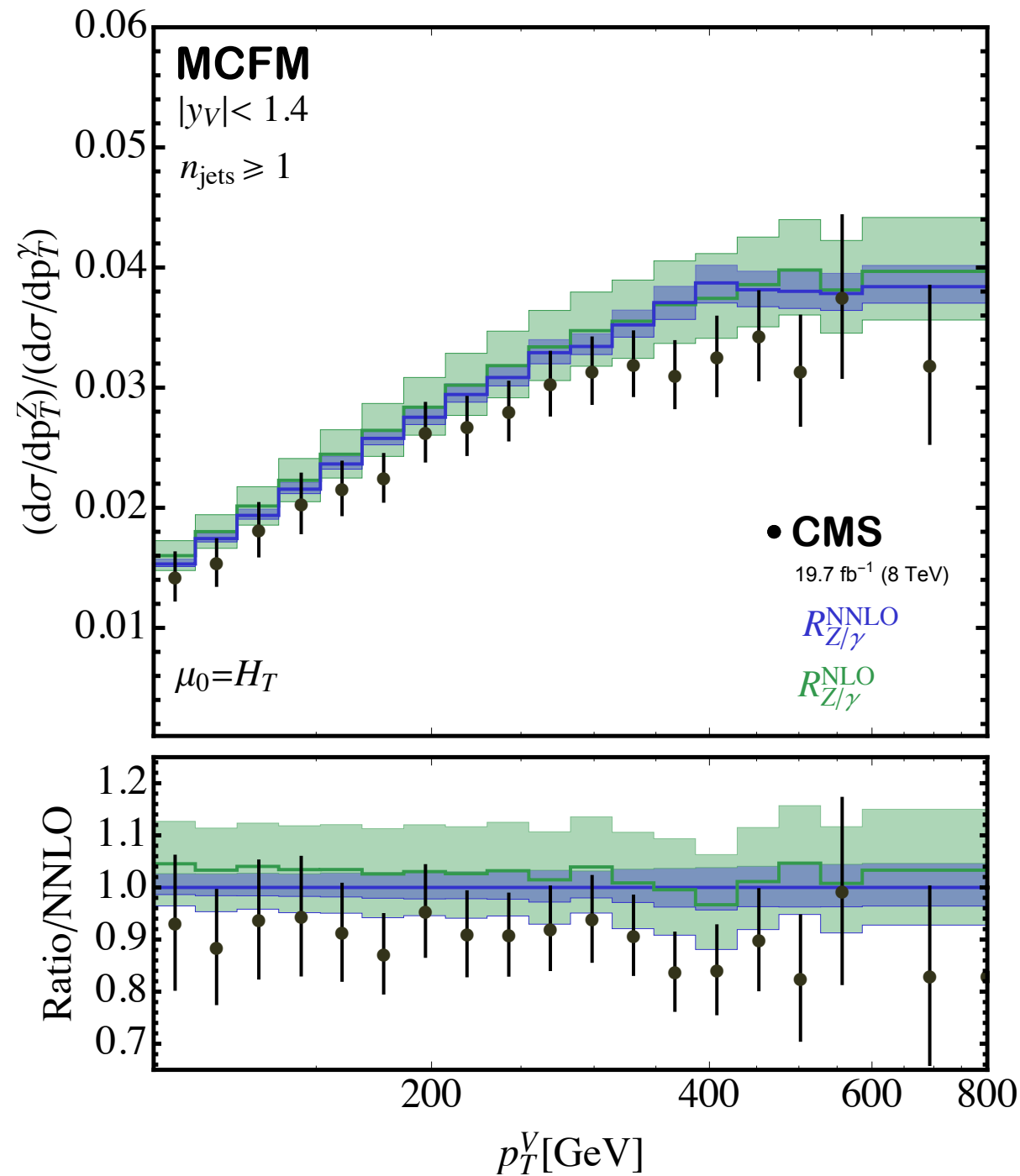
[Gehrmann-De Ridder, Gehrmann, Glover, A. Huss, Morgan; '16]



[Boughezal, Liu, Petriello; '16]

- unprecedented reduction of scale uncertainties at NNLO: $O(\sim 5\%)$
- we can now check the correlation of the uncertainties going from NLO to NNLO

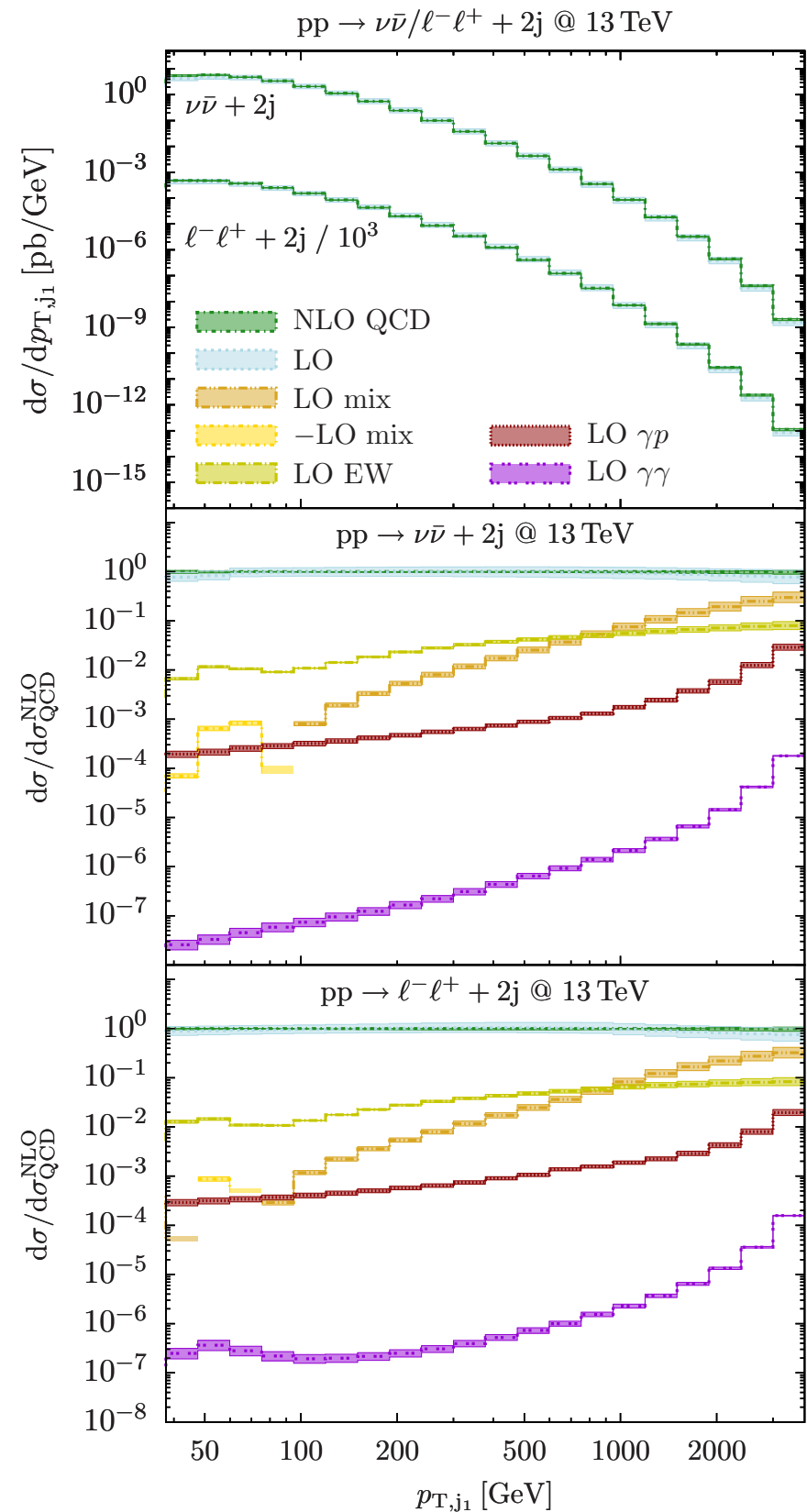
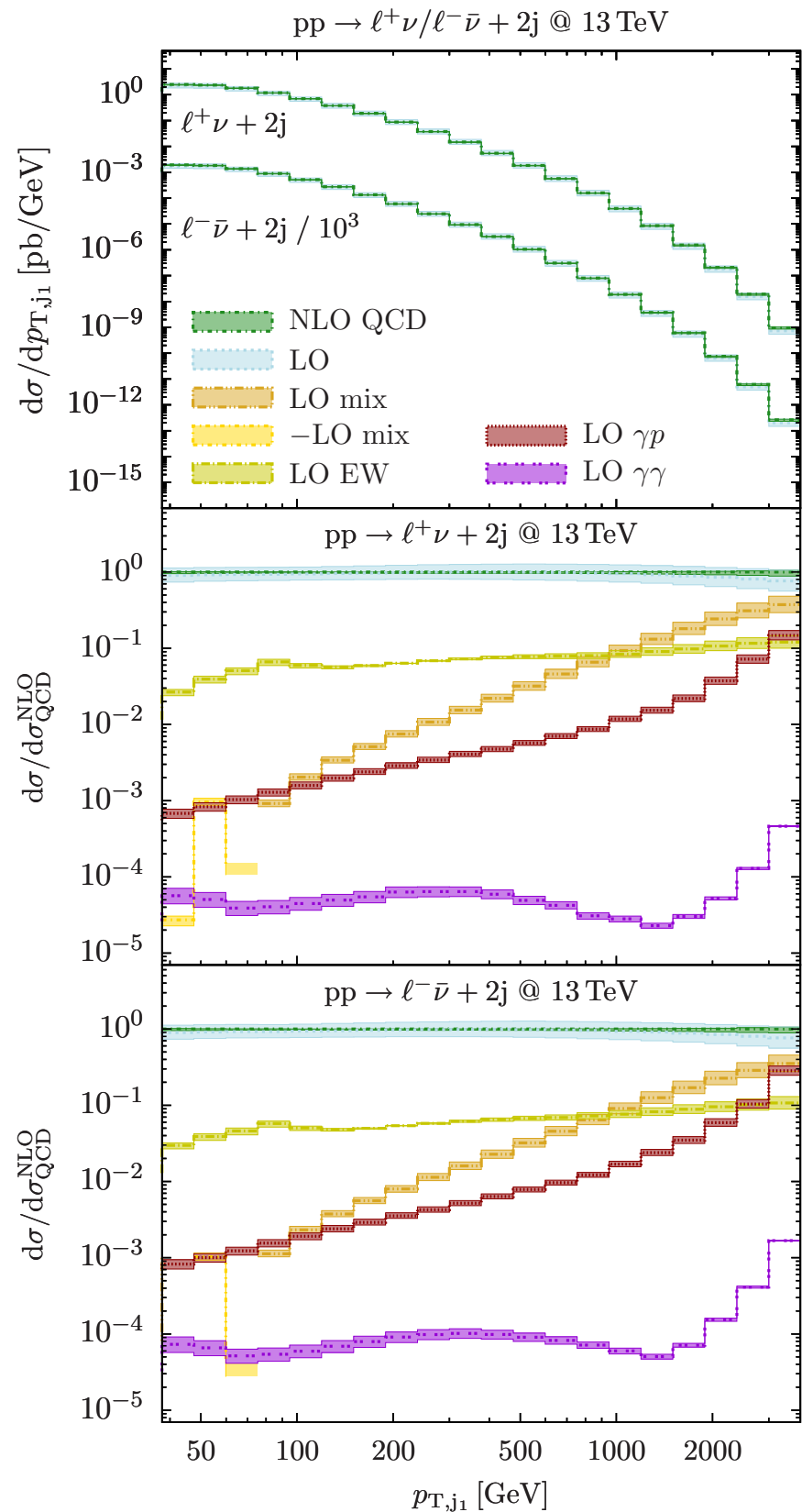
NNLO for Z/γ +jet



[Campbell, Ellis, Williams; '17]

NNLO/NLO ~ 1 for large p_T !

Subleading Born: $p_{T,j1}$



Caveat: γ +jet

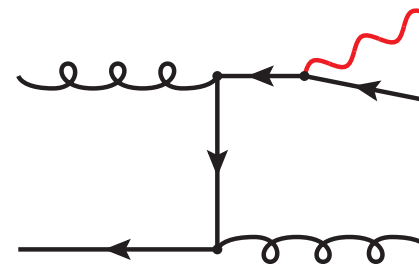
Note: this modelling of process correlations assumes close similarity of QCD effects between different V +jets processes

- apart from PDF effects it is the case for W +jets vs. Z +jets
- at $p_T > 200$ GeV it is also the case for γ +jets vs. Z +jets.

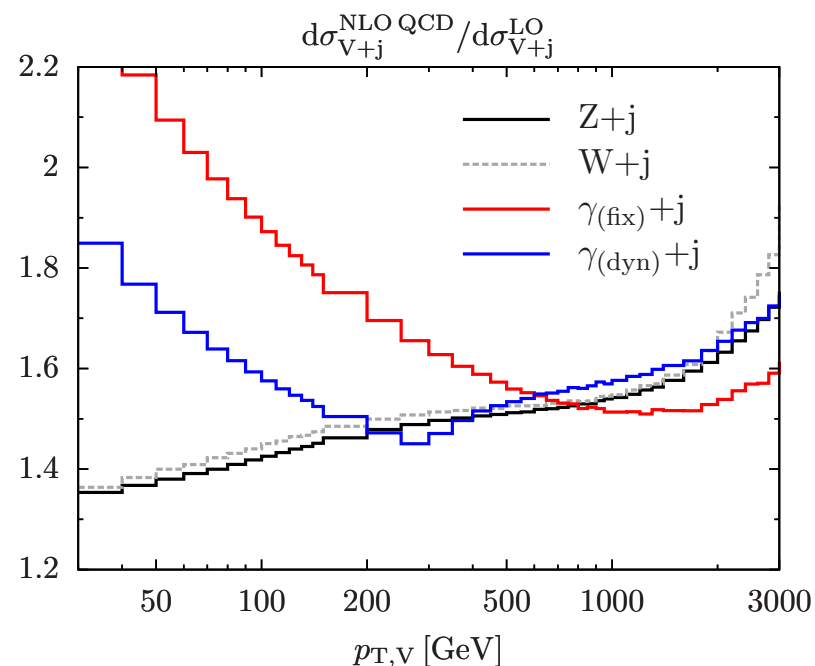
$$\left| \frac{\sigma_{\text{NLO}}^{(V)}}{\sigma_{\text{LO}}^{(V)}} - \frac{\sigma_{\text{NLO}}^{(Z)}}{\sigma_{\text{LO}}^{(Z)}} \right| \ll \left| \frac{\sigma_{\text{NLO}}^{(Z)}}{\sigma_{\text{LO}}^{(Z)}} \right|$$

Different logarithmic effects from fragmentation

- W/Z +jet: masscut-off $M_{Vj} \geq M_V$
- γ + jet: Frixione-isolation cone of radius R_0

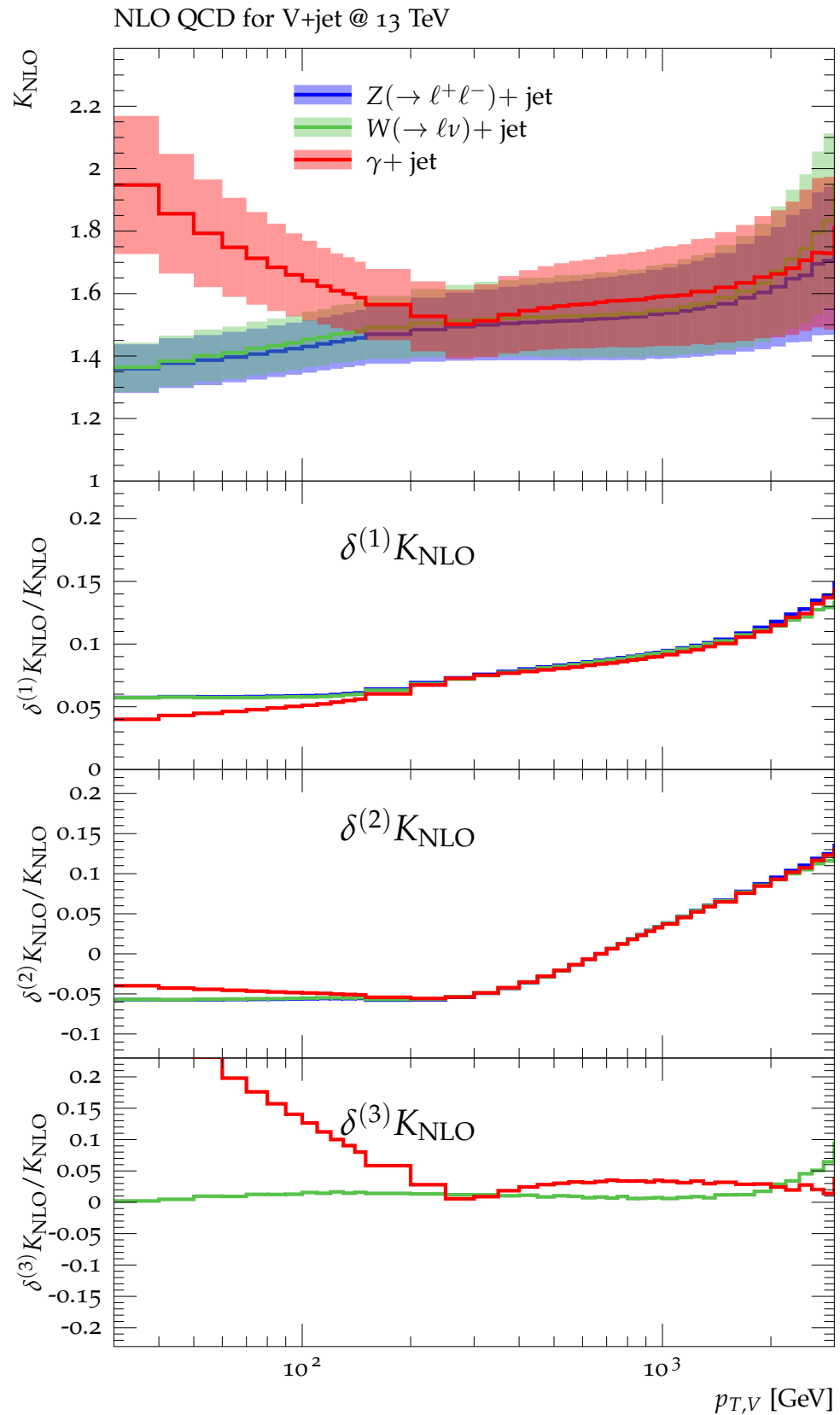


Consider dynamic γ -isolation with $R_{\text{dyn}}(p_T, \gamma) = \min\{1.0, M_Z/p_{T\gamma}\}$



- γ_{dyn} behaves like W/Z at $p_T > M_Z$
 \Rightarrow justifies process-correlation estimate
- remnant part $\gamma_{\text{fix}} - \gamma_{\text{dyn}}$ uncorrelated
 (uncertainty through extra reweighting and MC)

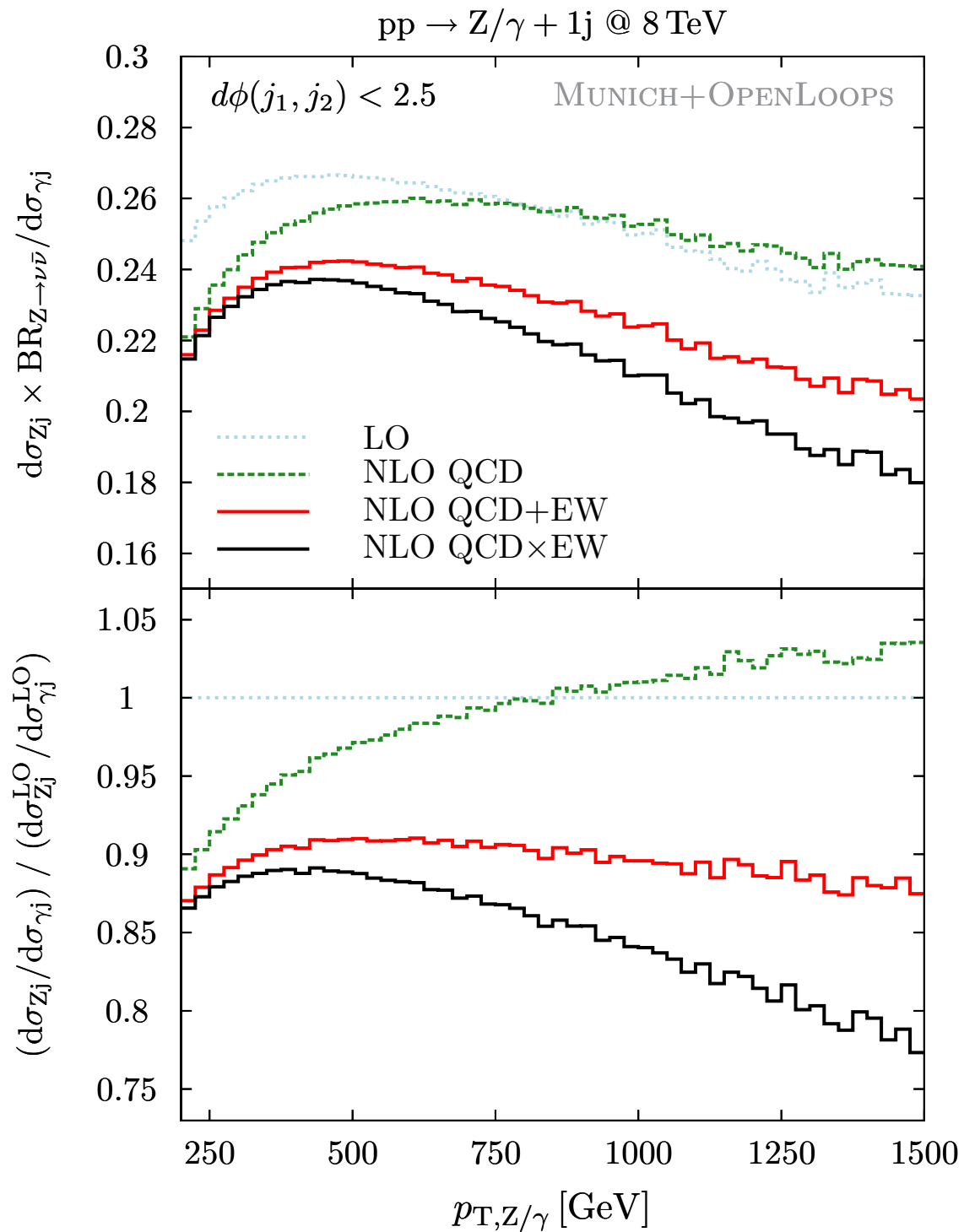
QCD uncertainties



NLO QCD corrections and uncertainties

- almost identical for W/Z/ γ for $p_{T,V} > 200$ GeV
- sizeable γ +jet fragmentation for $p_{T,V} < 200$ GeV

Z/ γ + 1 jet: pT-ratio



Overall

- ▶ mild dependence on the boson pT

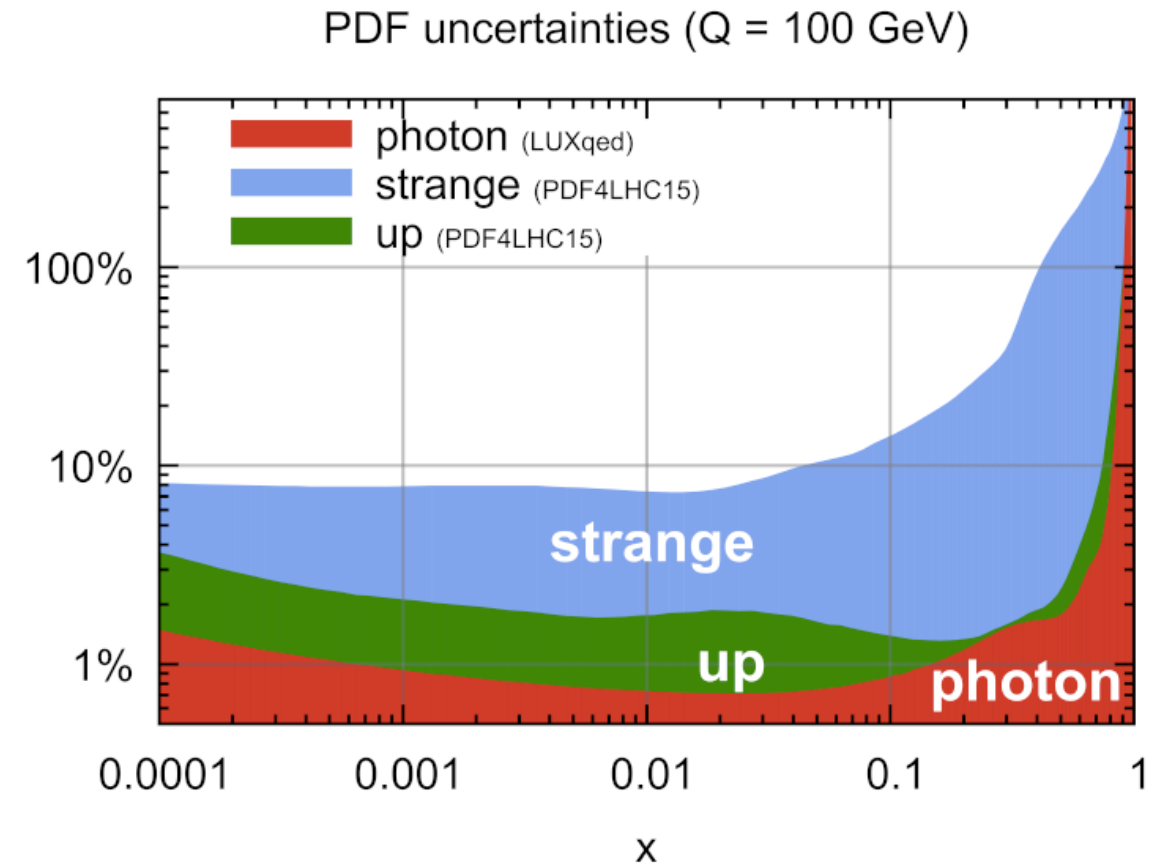
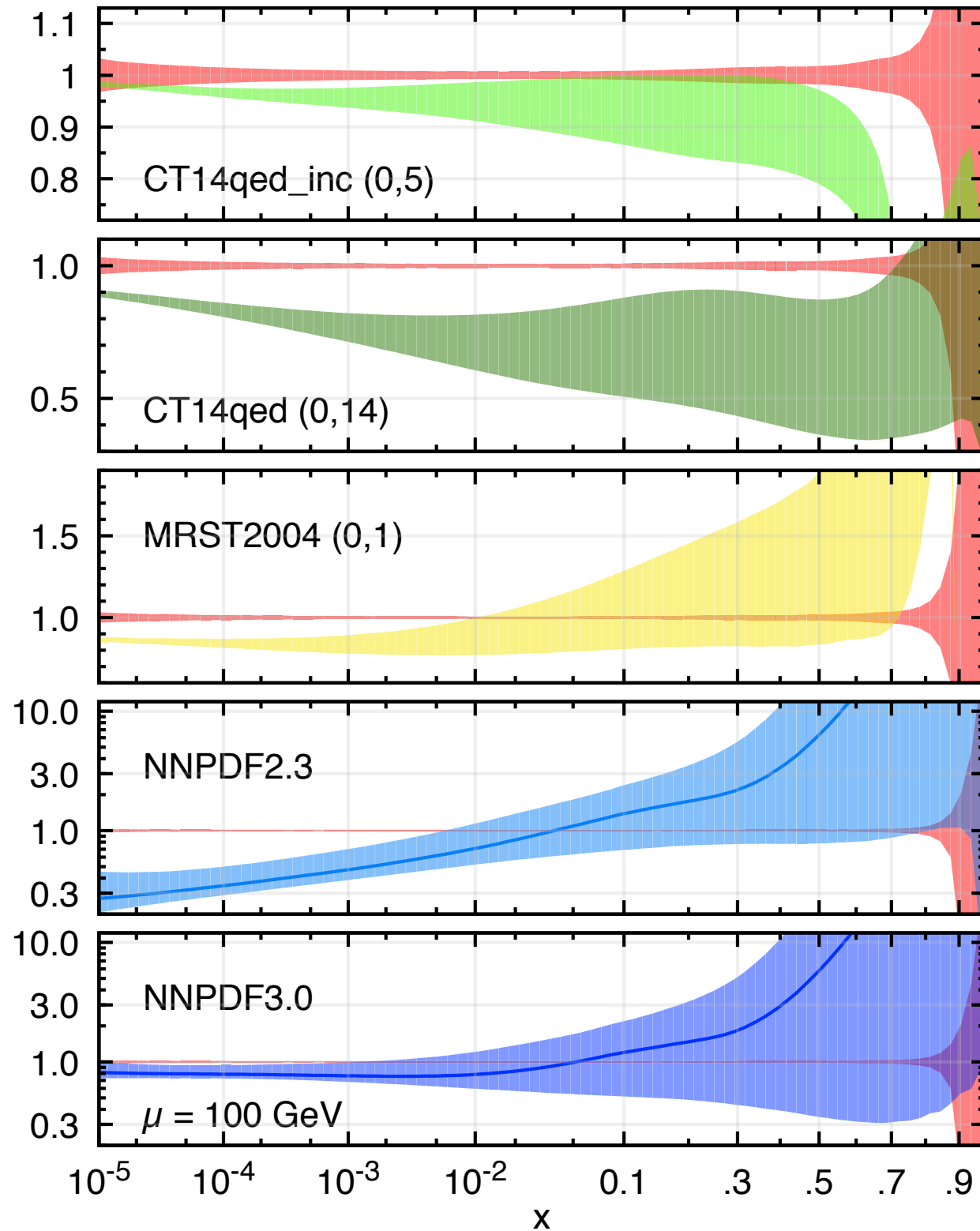
QCD corrections

- ▶ 10-15% below 250 GeV
- ▶ \approx 5% above 350 GeV

EW corrections

- ▶ sizeable difference in EW corrections results in 10-15% corrections at several hundred GeV
- ▶ \sim 5% difference between NLO QCD+EW and NLO QCD \times EW

LUXqed



[Manohar, Nason, Salam, Zanderighi, '16]

