

Coherence Effects at Strong Coupling

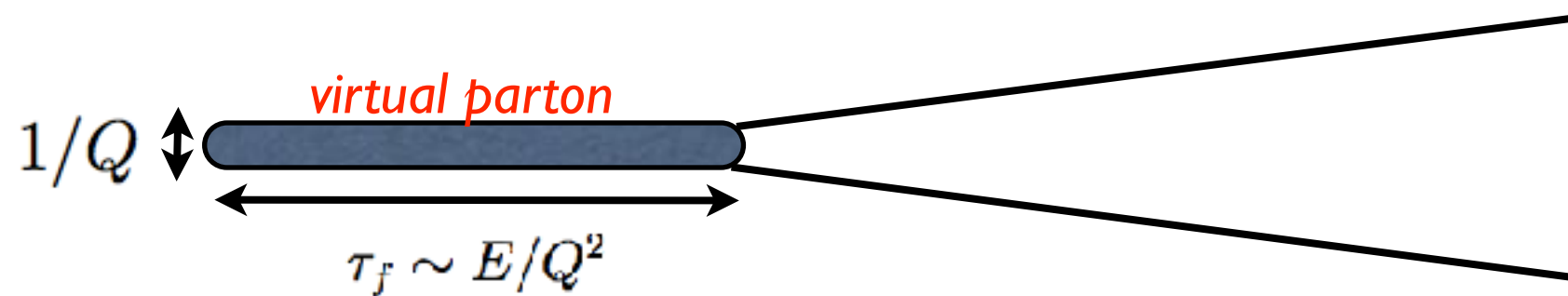
Jorge Casalderrey-Solana



THE ROYAL
SOCIETY

Jets in Medium

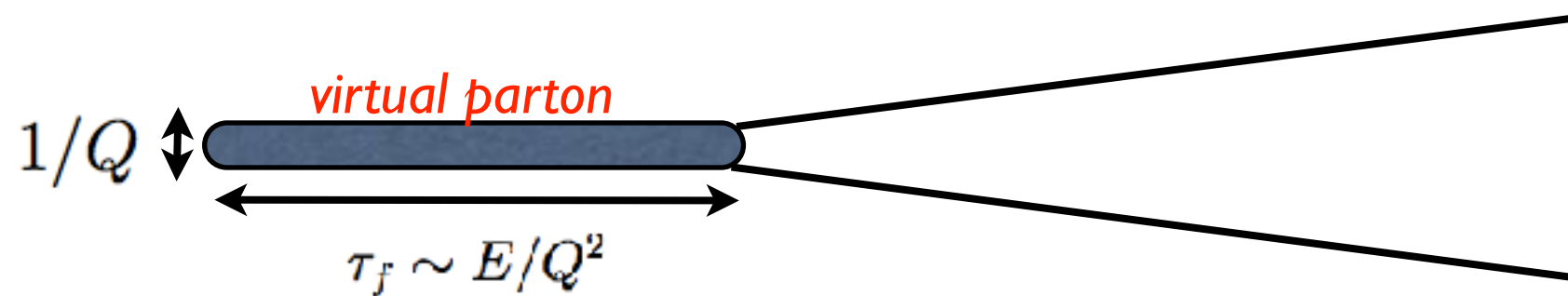
- Jets probe different scales



Jets in Medium

- ⊙ Jets probe different scales

$$T \ll Q$$

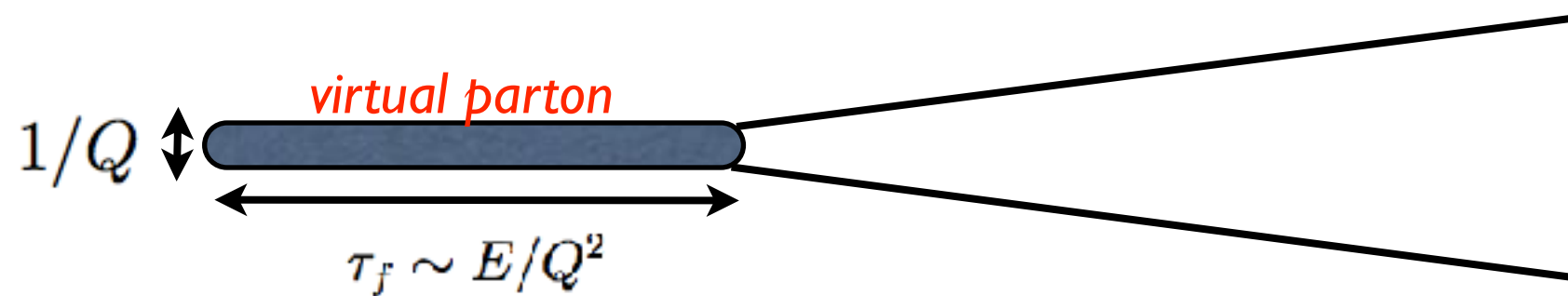


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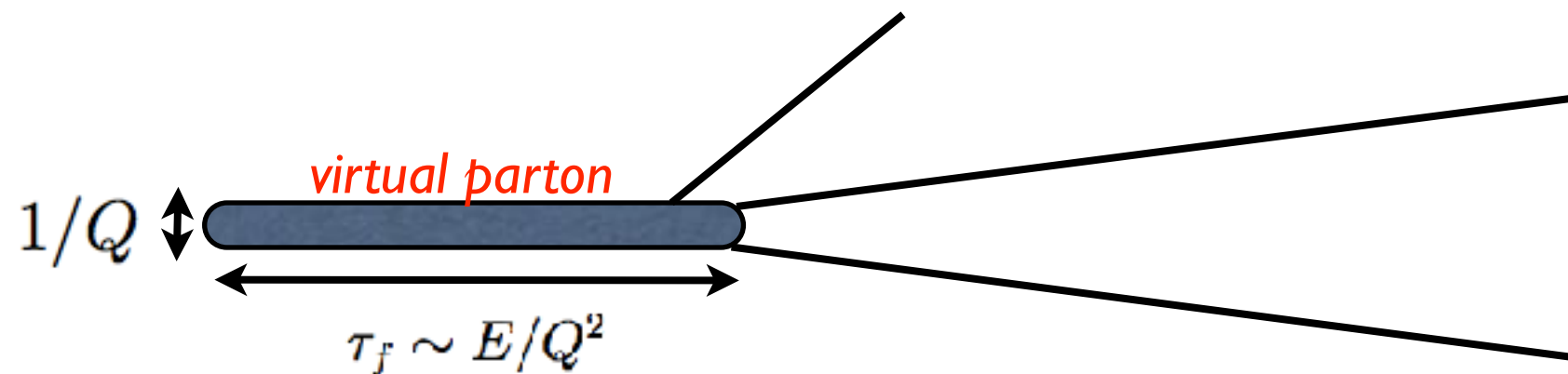
- In medium

- Small size \Rightarrow Only rare fluctuations alter splittings. $\mathcal{P} \sim \alpha^2 \frac{T^2}{Q^2}$

Jets in Medium

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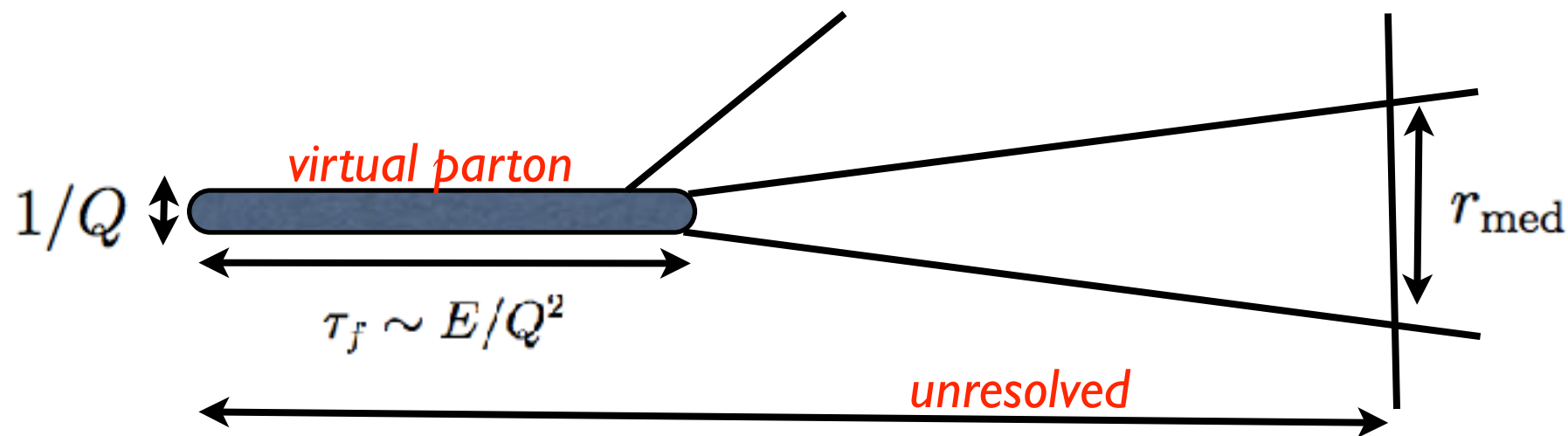
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- Virtual partons carry a colour charge \Rightarrow

Induced radiation is still possible with $k_{\perp} \ll Q$

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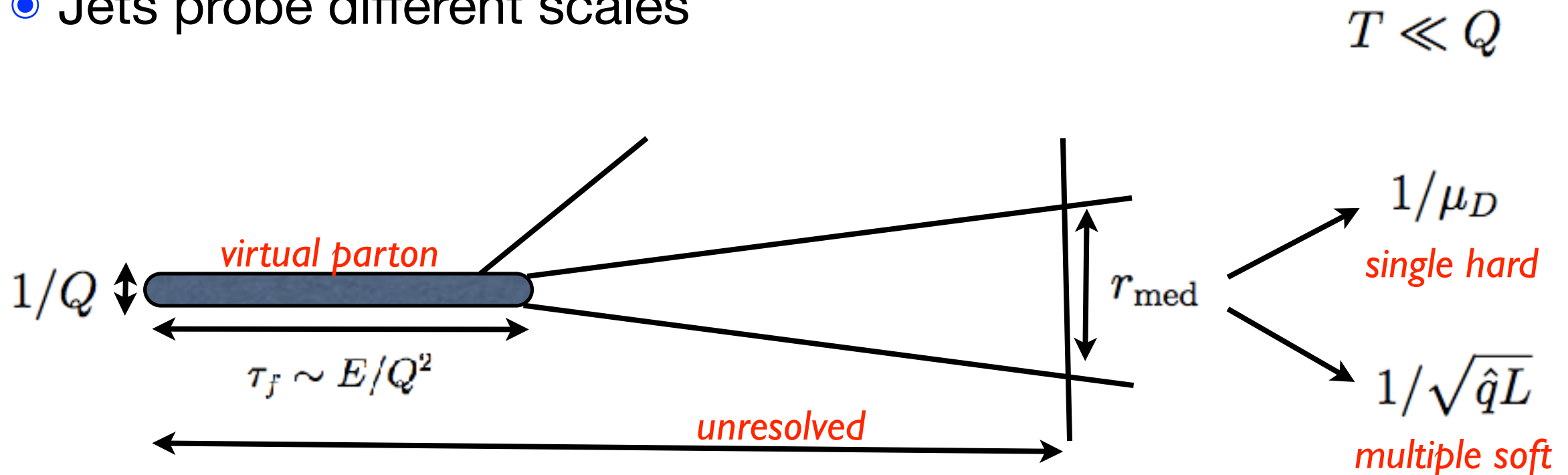
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- After formation, partons may be not resolved by the medium

Radiation as single objects

Jets in Medium

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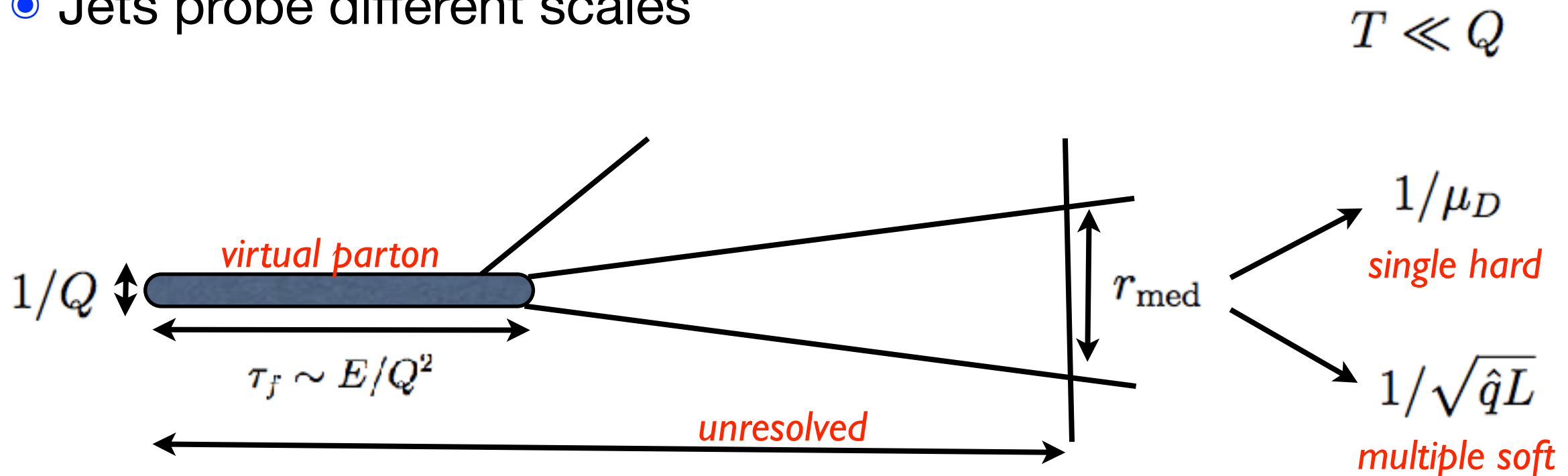
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Radiation as single objects

Why strong coupling for jets?

- Jets probe different scales



- What are the typical scales?

- Average energy loss controlled by hardest splittings

$$Q_H \sim \sqrt{\hat{q}L} \sim 3.2 \text{ GeV} \left(\frac{\hat{q}}{2 \text{ GeV}^2/\text{fm}} \frac{L}{5 \text{ fm}} \right)^{1/2}$$

multiple soft

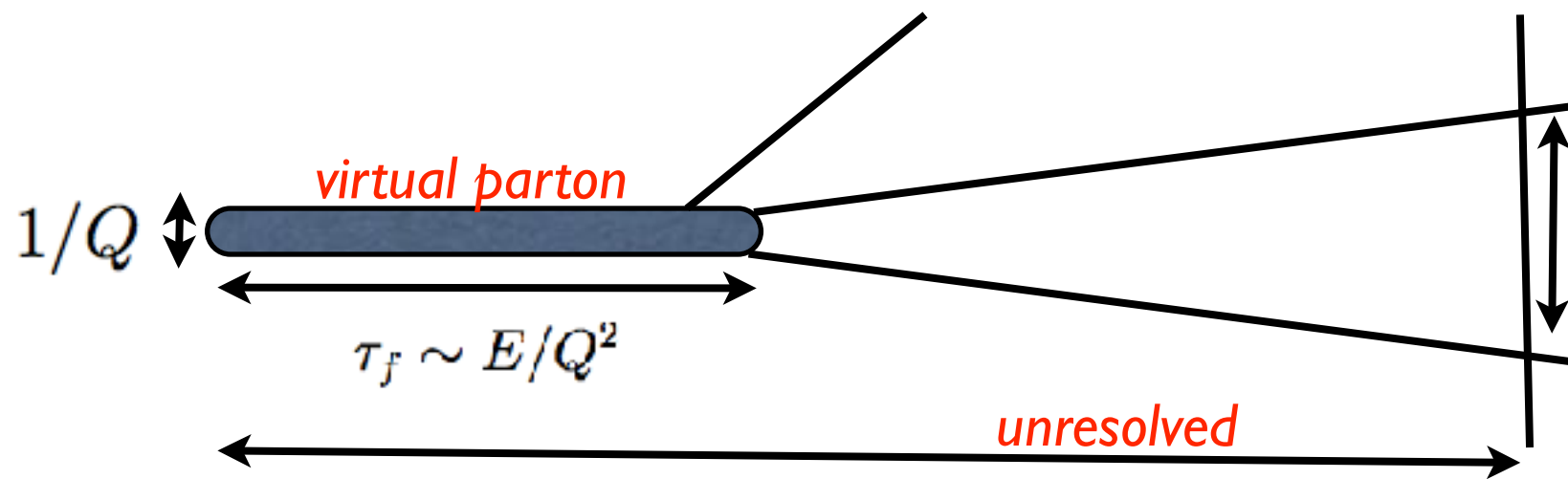
$$Q_H \sim \mu_D \sim 1.2 \text{ GeV} \left(\frac{g}{2} \frac{T}{0.47 \text{ GeV}} \right)$$

single hard

Not terribly perturbative

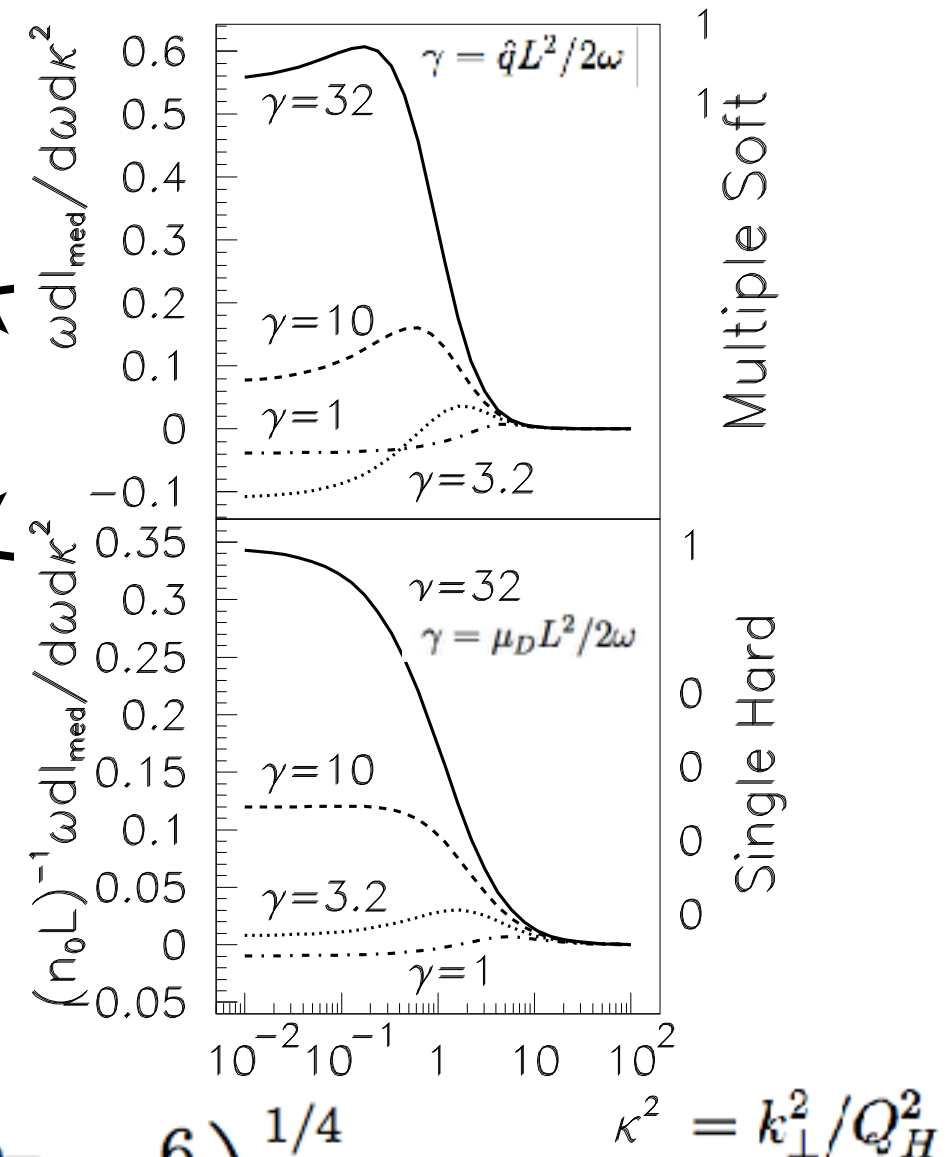
Why strong coupling for jets?

- Jets probe different scales



- Most gluons radiated at smaller scale
 - In a steeply falling spectrum losses are dominated by softer emissions

Salgado & Wiedemann 03

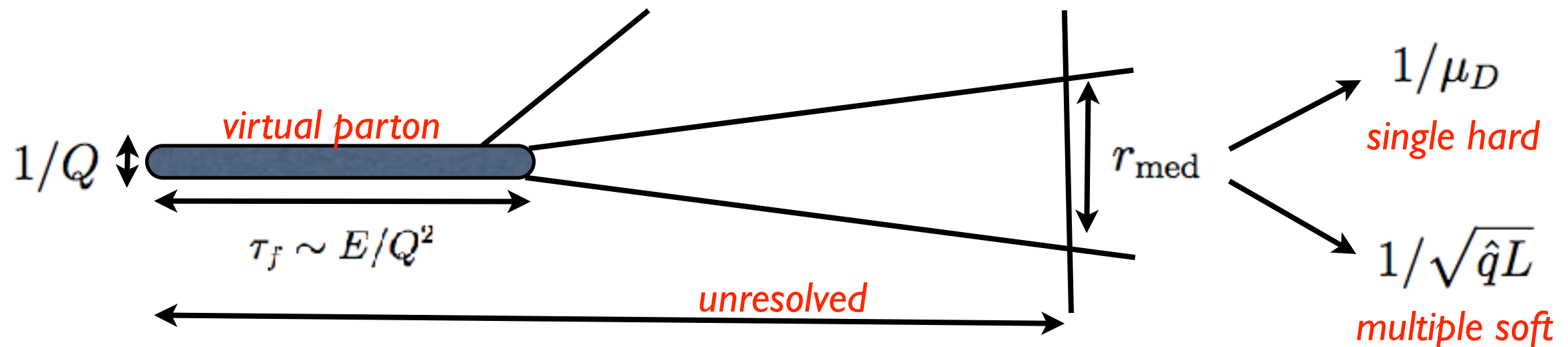


$$Q_{\text{dom}} < \left(\frac{p_T}{n} \hat{q} \right)^{1/4} = 1.6 \text{ GeV} \left(\frac{\hat{q}}{2 \text{ GeV}^2/\text{fm}} \frac{p_T}{100 \text{ GeV}} \frac{6}{n} \right)^{1/4}$$

BDMS 01

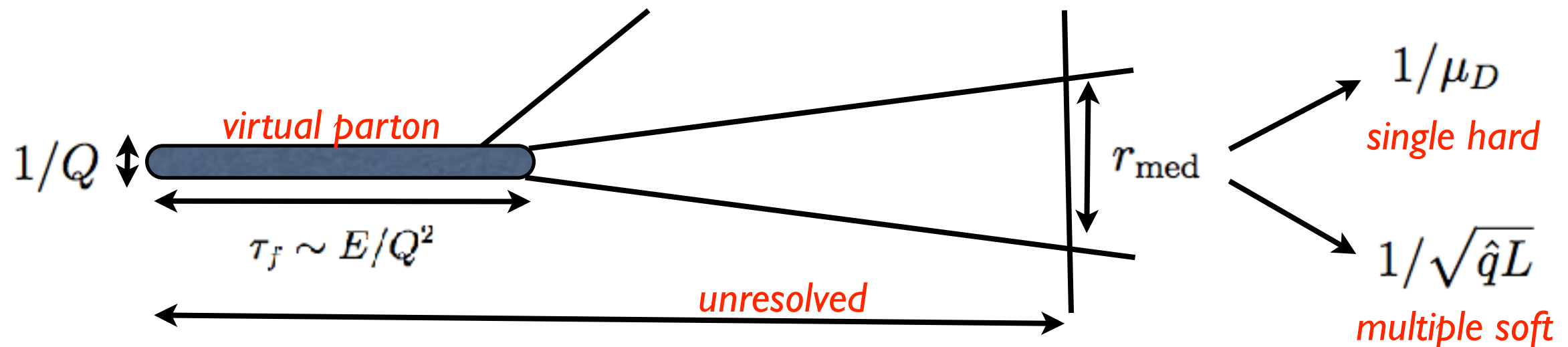
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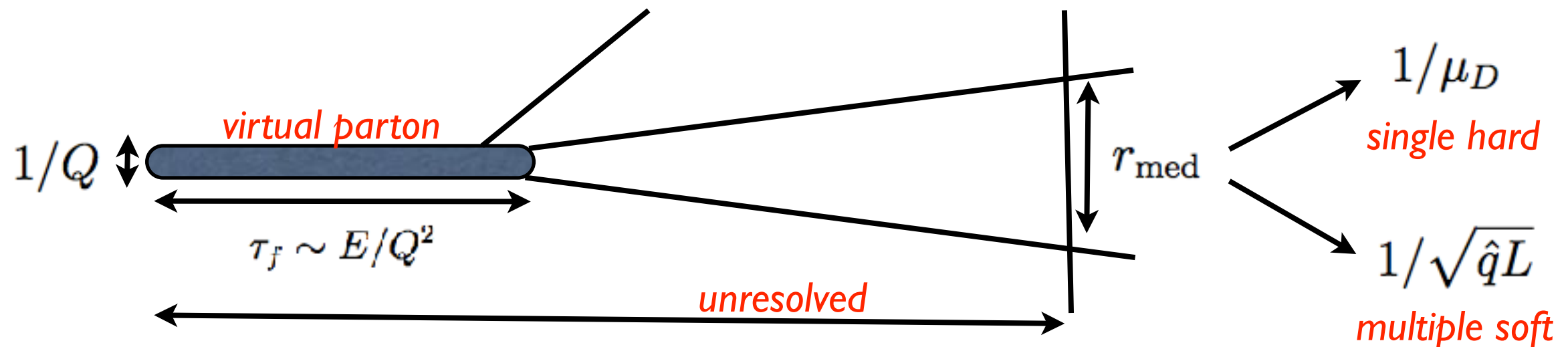
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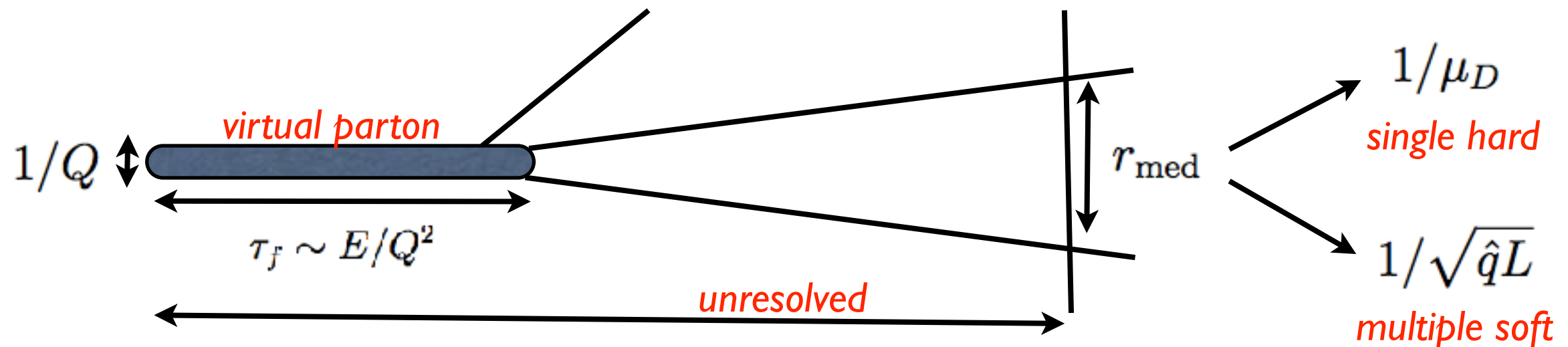
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- Scales involved in quenching are not large
 - Nevertheless, LO perturbative computations describe many data
- It is worth exploring the opposite extreme assumption
 - Infinite coupling limit (only accessible via gauge/gravity duality)

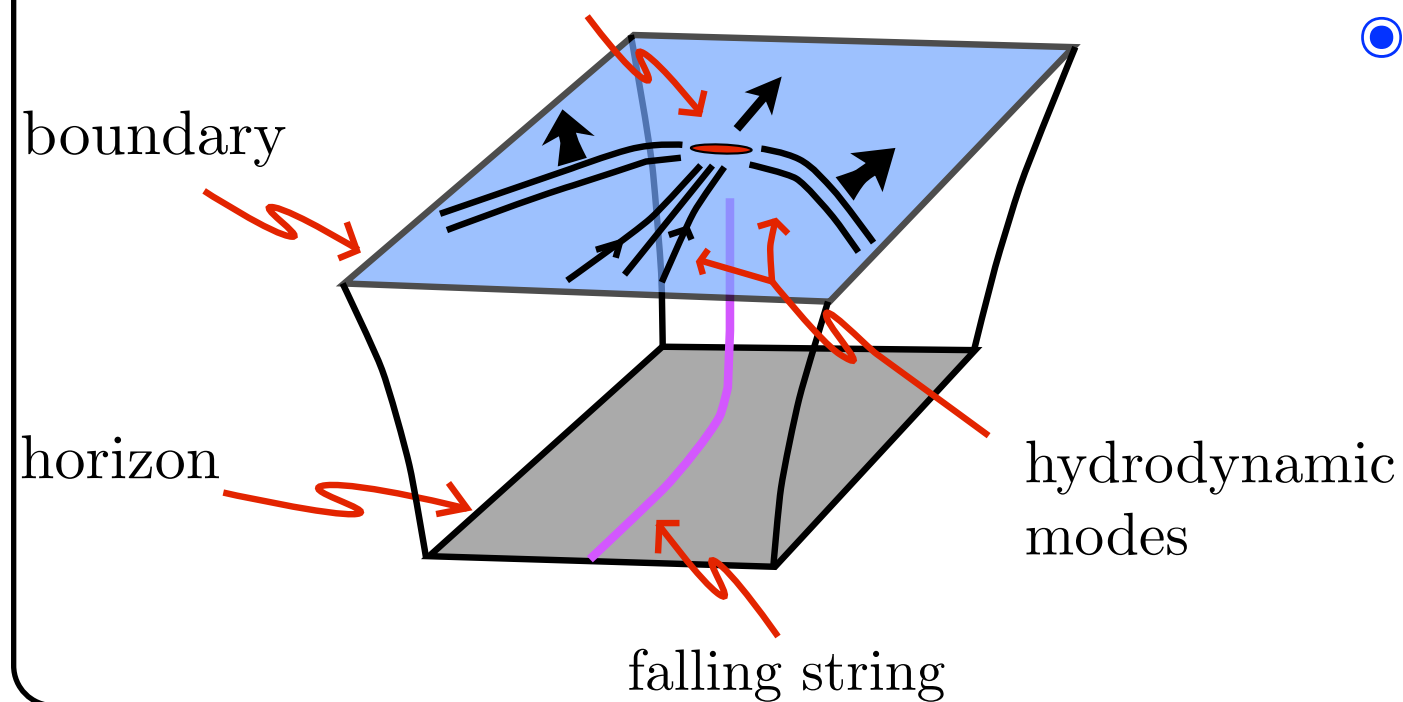
Why strong coupling for jets?



- Scales involved in quenching are not large
 - Nevertheless, LO perturbative computations describe many data
- It is worth exploring the opposite extreme assumption
Infinite coupling limit (only accessible via gauge/gravity duality)
 - Can we get an equally good description?
 - Can we find different characteristic features?

Energy Distribution to a Jet

jet Chesler & Rajagopal 16



○ $T^{\mu\nu}$ associated to falling string

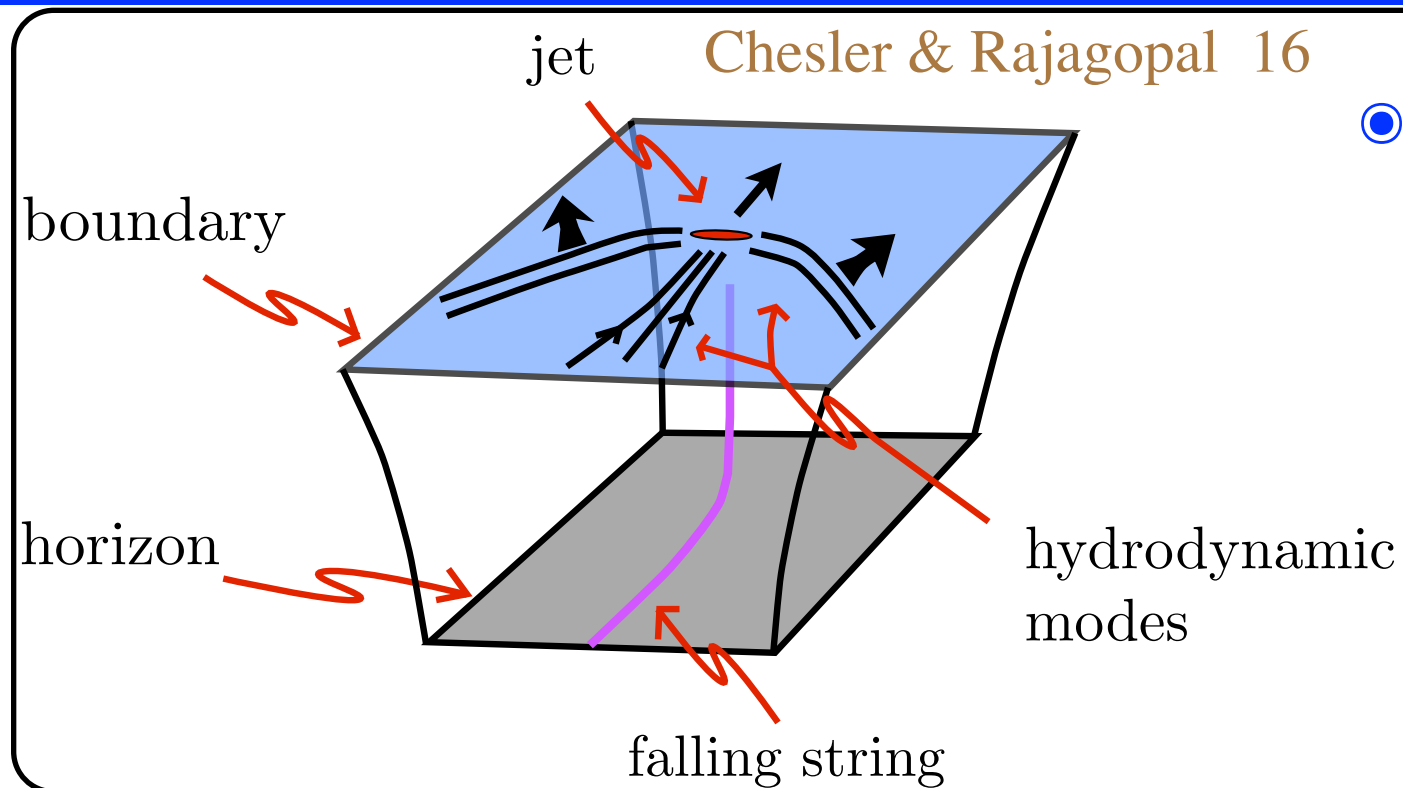
- High energy core

Jet

- Hydro excitations

Thermalised lost energy

Energy Distribution to a Jet



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Thermalised lost energy

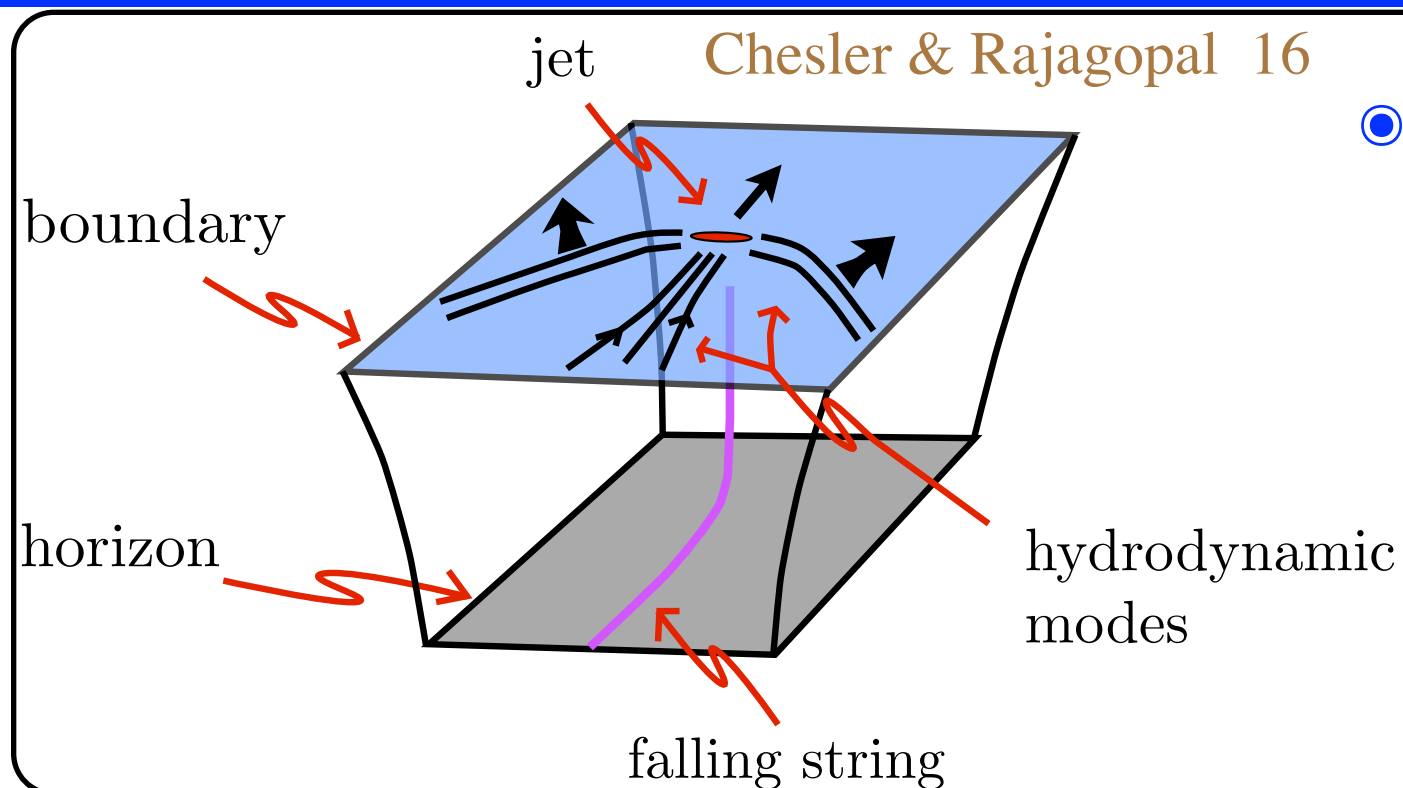
• Dumping of energy into hydro modes at a rate

$$\frac{1}{E_{\text{init}}} \frac{dE_{\text{jet}}}{dx} = - \frac{4x^2}{\pi x_{\text{therm}}^2 \sqrt{x_{\text{therm}}^2 - x^2}} \xrightarrow{\text{depend on jet energy}} x_{\text{stop}} = \frac{1}{2 \kappa_{\text{sc}}} \frac{E_{\text{in}}^{1/3}}{T^{4/3}},$$

Gubser et al 08, Chesler et al. 08,
Ficnar and Gubser 13

number

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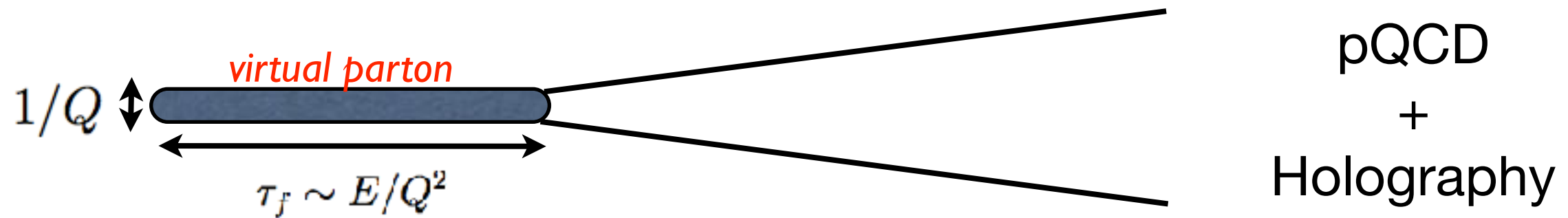
number

• Short distance physics remain strongly coupled

⇒ use as a pheno tool requires additional modelling

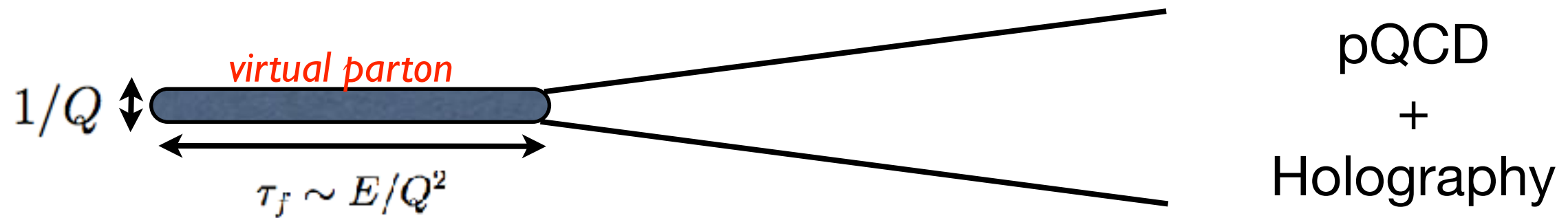
A Hybrid Model

JCS, Gulhan, Milhano, Pablos and Rajagopal 14,15,16



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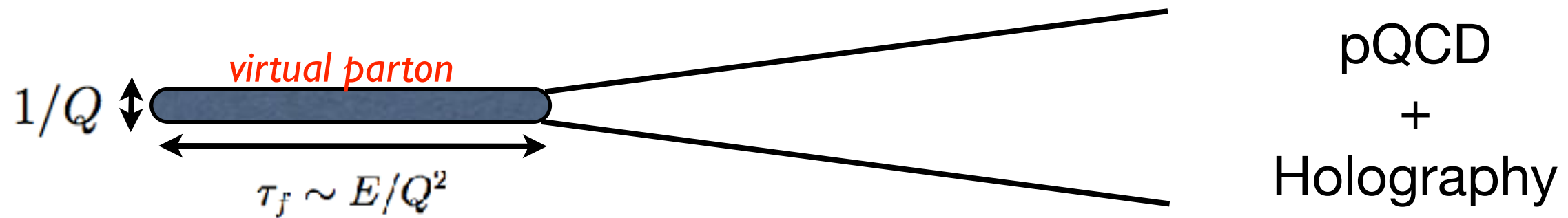
JCS, Gulhan, Milhano, Pablos and Rajagopal 14,15,16



- Keep the DGLAP structure inherited from QCD

A Hybrid Model

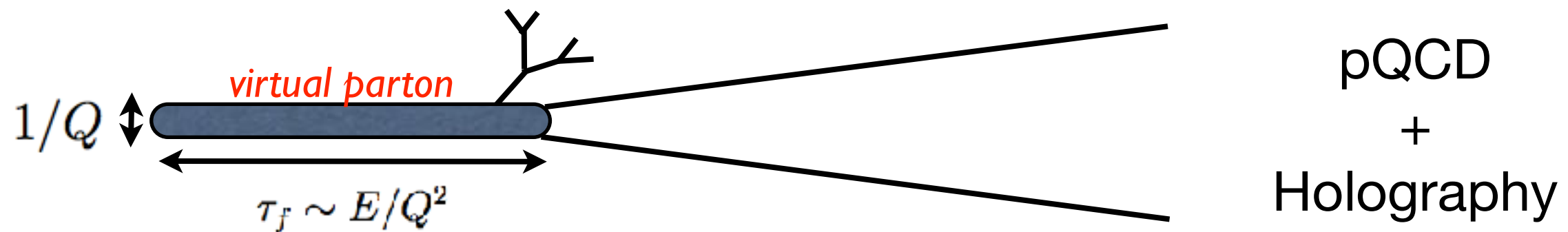
JCS, Gulhan, Milhano, Pablos and Rajagopal 14,15,16



- ⊙ Keep the DGLAP structure inherited from QCD
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A Hybrid Model

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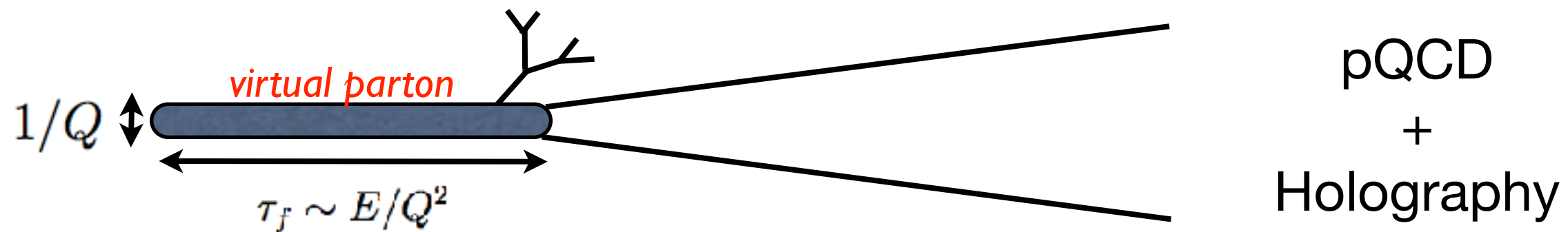


- Keep the DGLAP structure inherited from QCD
- Neglect rare fluctuations (evolution unchanged)
- Treat small scale processes at strong coupling
 - Each parton is considered as an energetic excitation in plasma
 - Energy loss proceeds as in strong coupling

$$x_{\text{stop}} = \frac{1}{2\kappa_{\text{sc}}} \frac{E_{\text{in}}^{1/3}}{T^{4/3}} \quad \text{(unique) fitting parameter}$$

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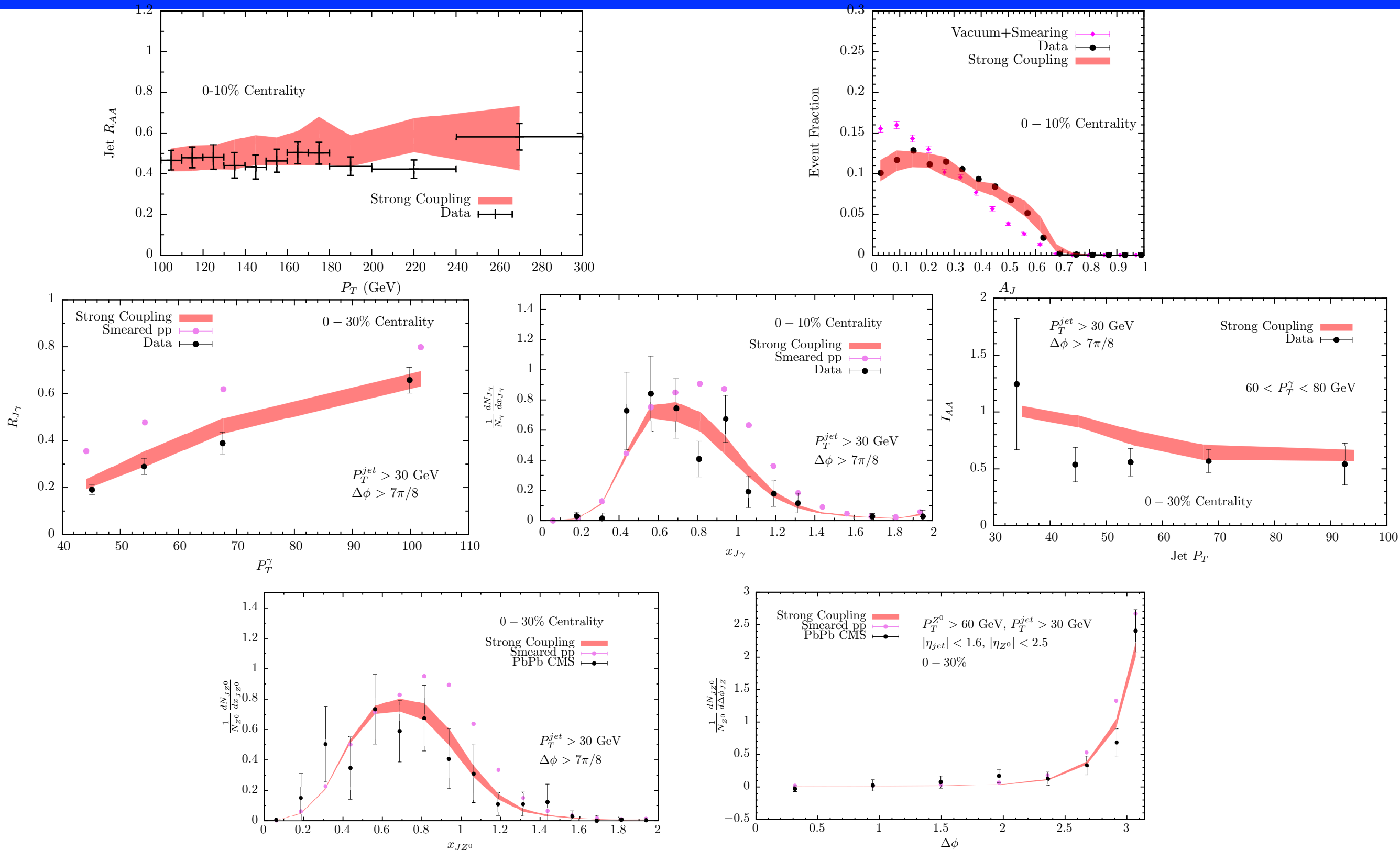


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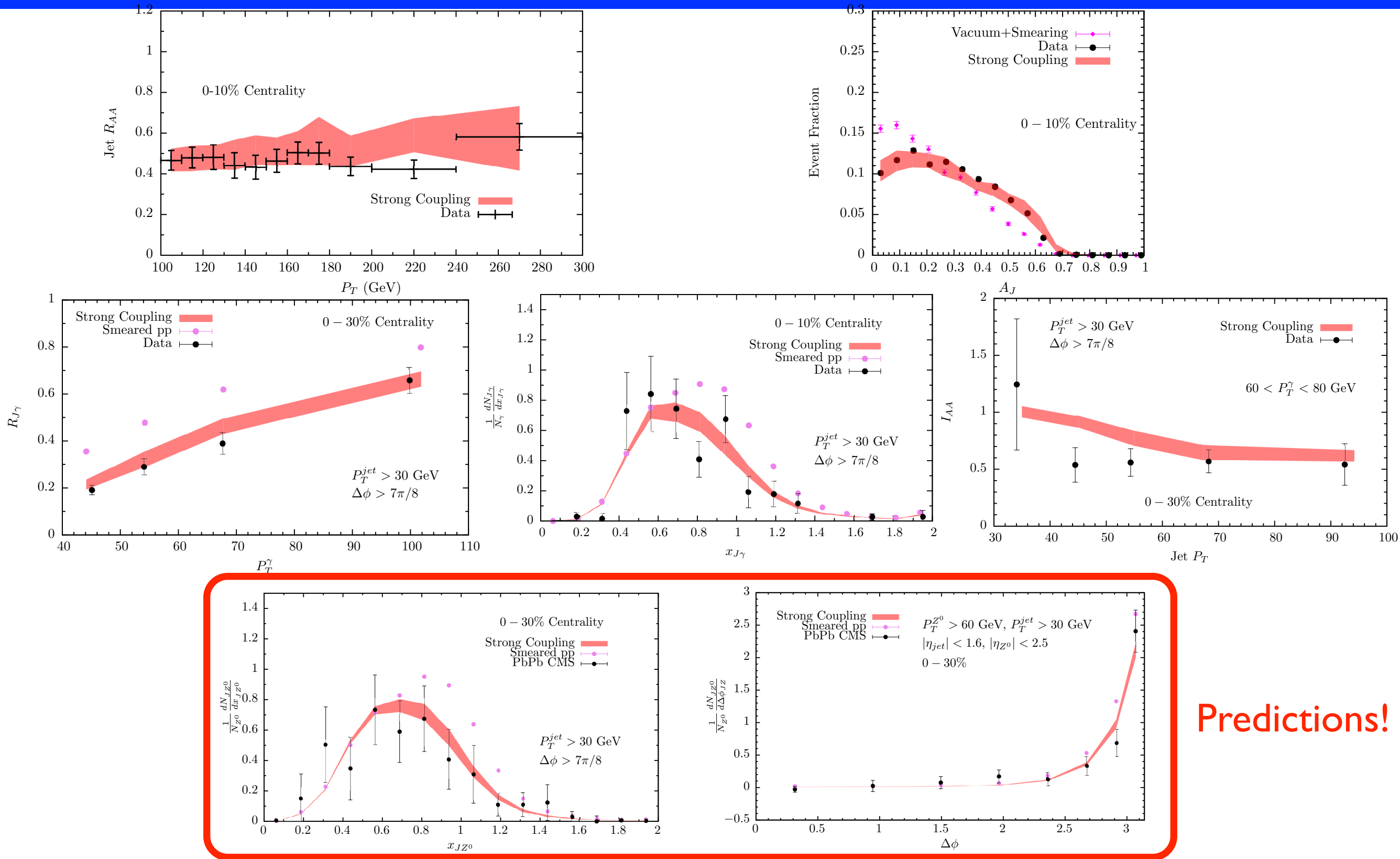
$$x_{\text{stop}} = \frac{1}{2\kappa_{\text{sc}}} \frac{E_{\text{in}}^{1/3}}{T^{4/3}} \quad \text{(unique fitting parameter)}$$

- Neglect coherence effects (for the moment)

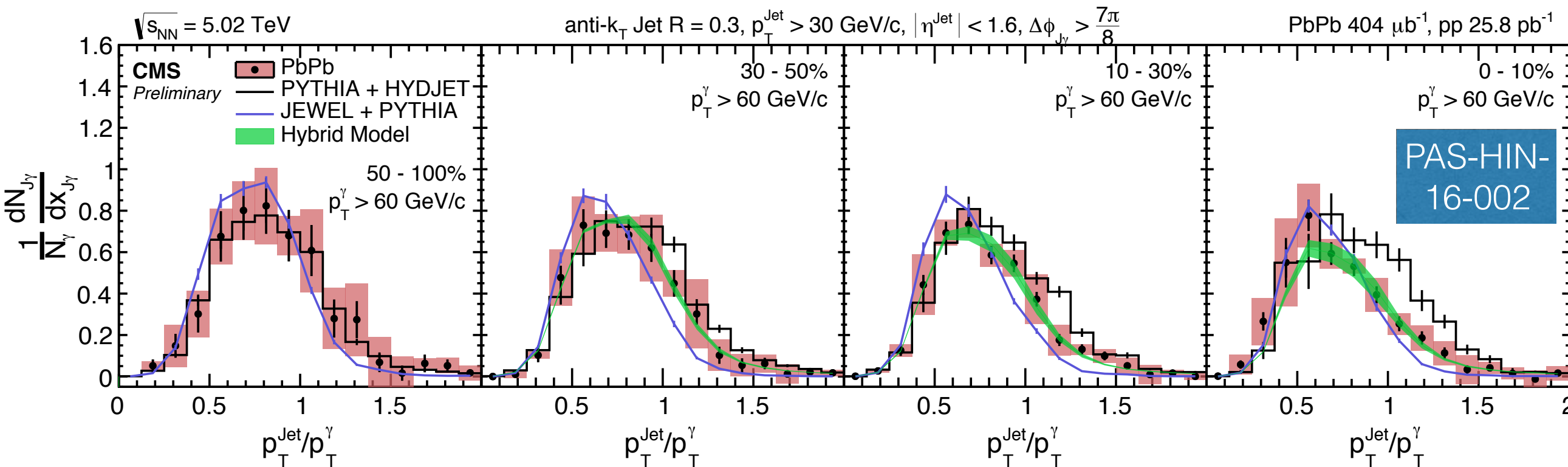
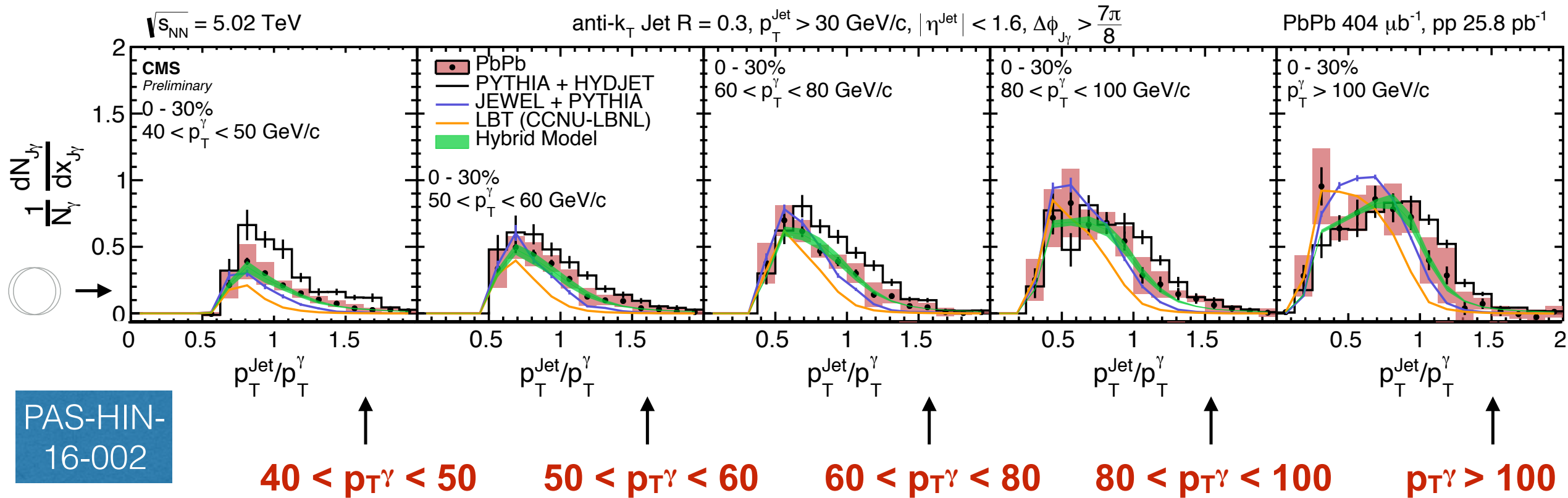
Success of the Hybrid Model

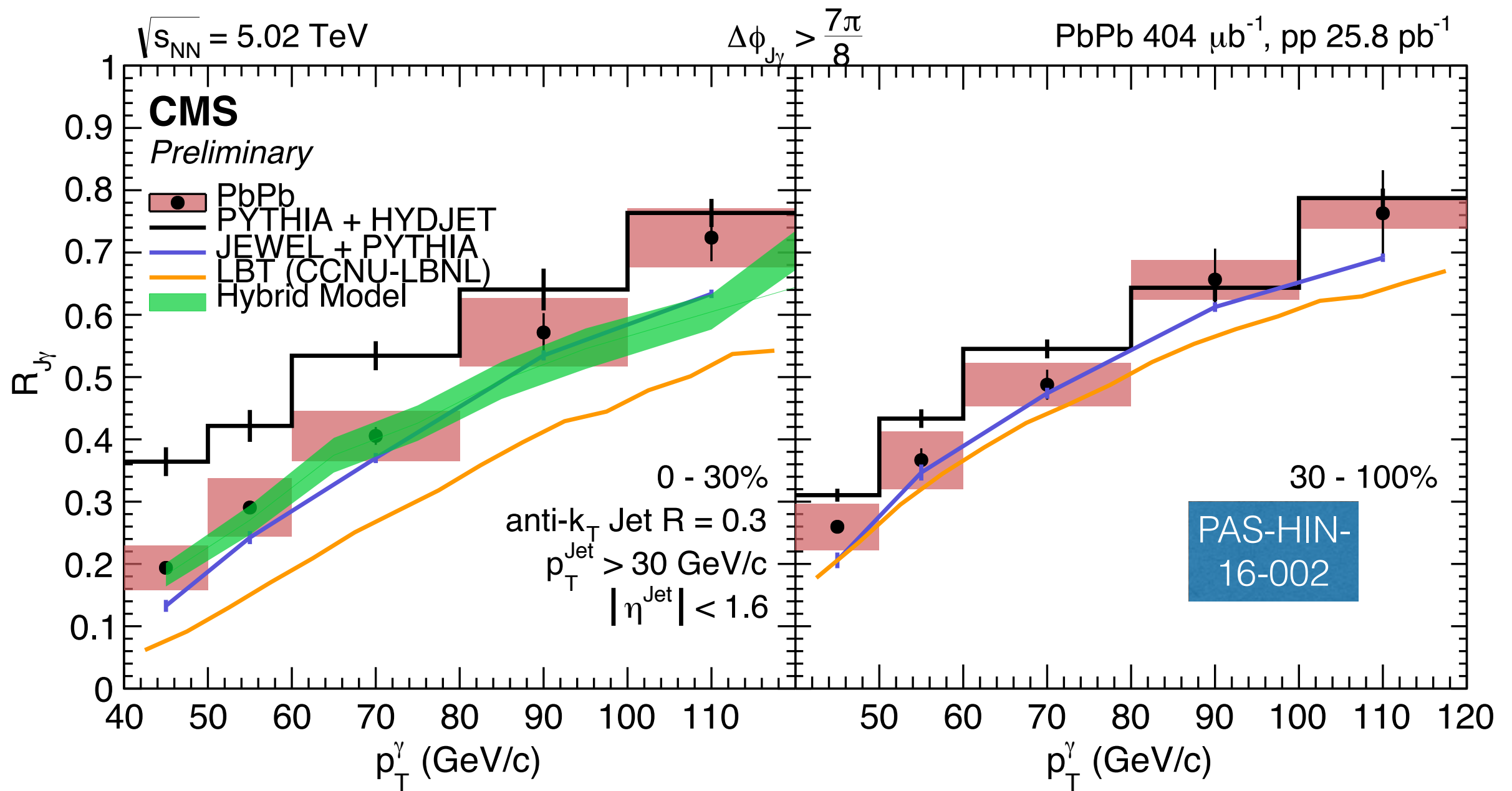


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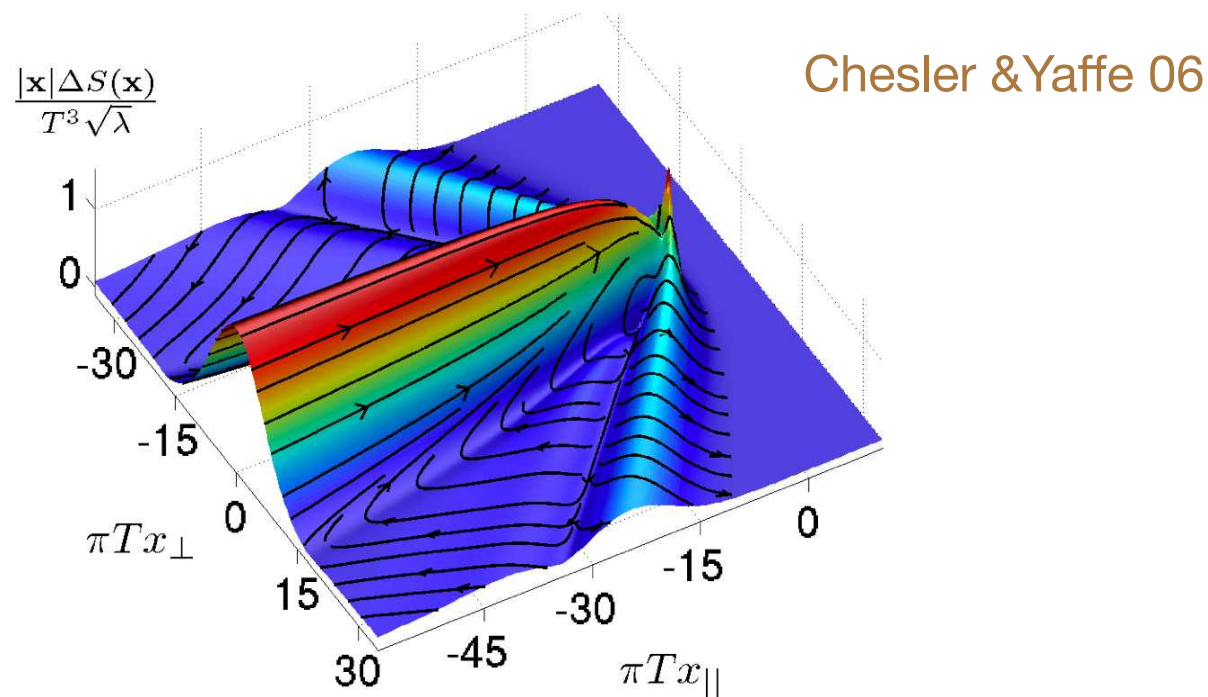


Predictions!





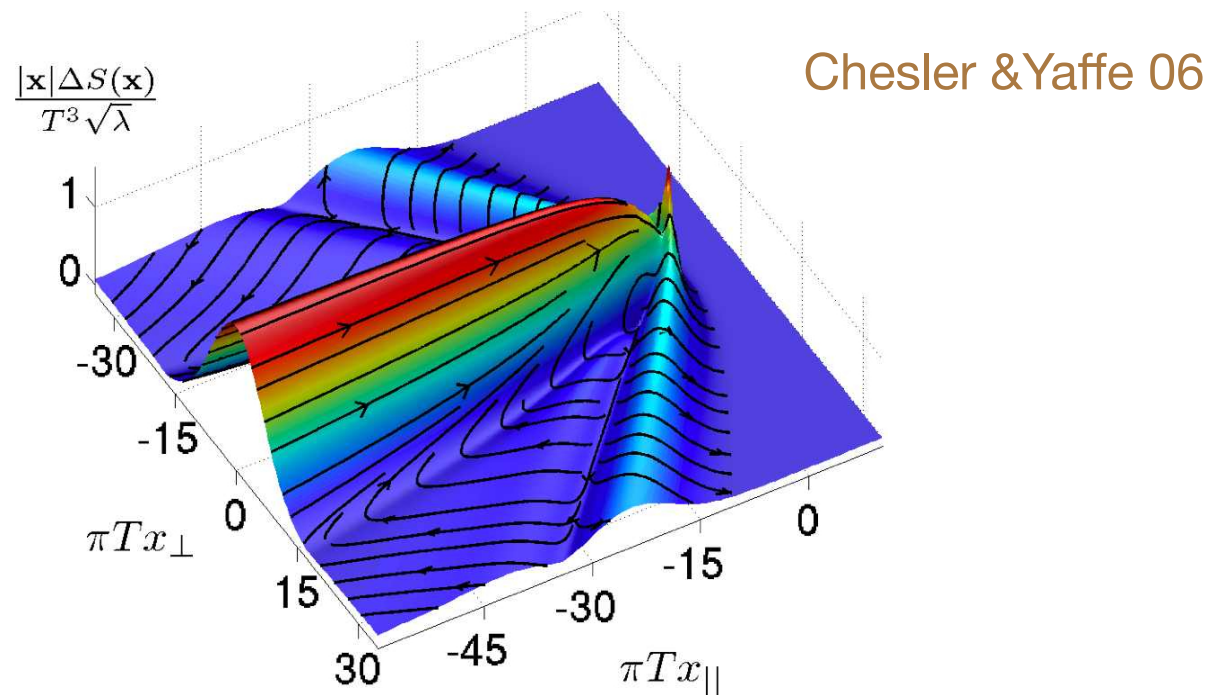
Back-Reaction and Wake



- The QGP is an extremely good fluid
 - Medium response to Eloss must be collective
 - Strong coupling computations provide an explicit example
 - Collectivity starts at short distance $1/T$ from the jet
 - There is a strong momentum flux along the jet direction

JCS, Shuryak & Teaney 06

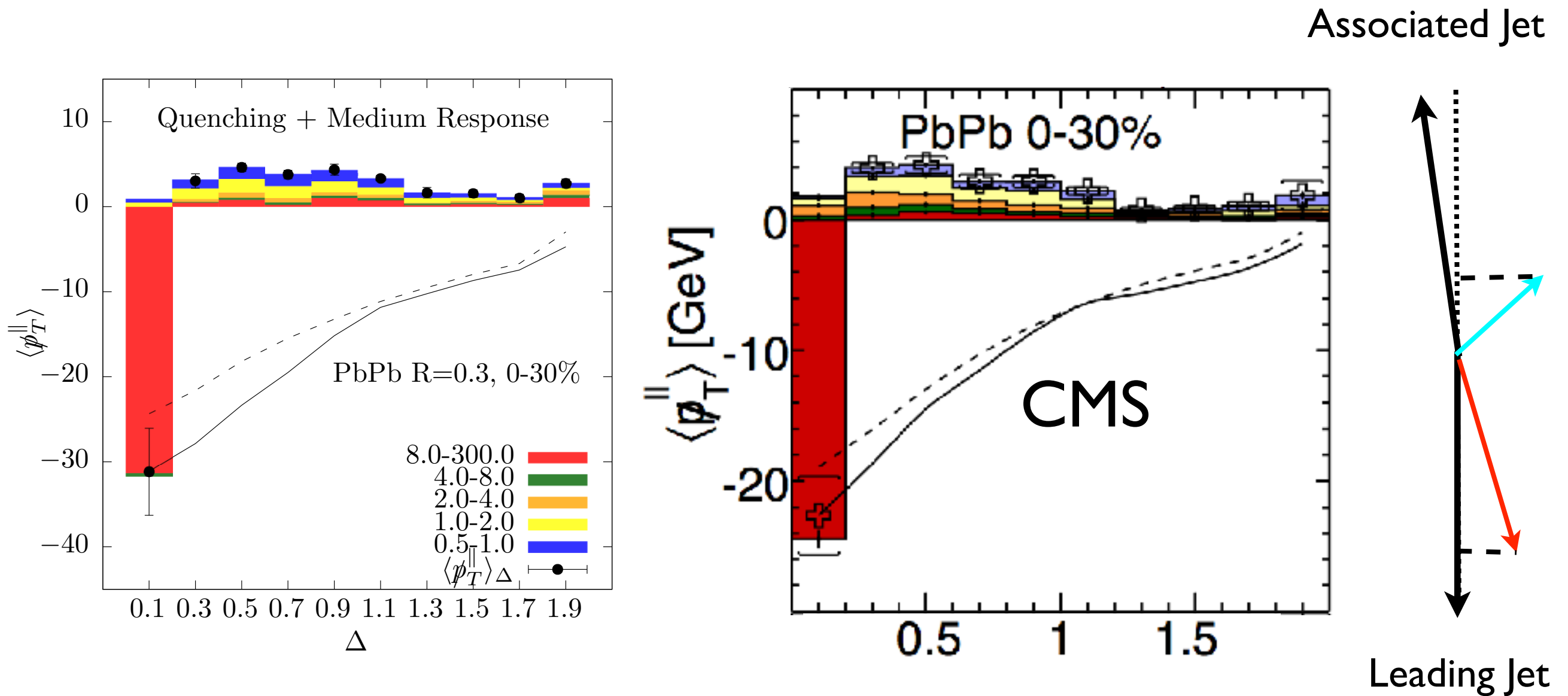
Back-Reaction and Wake



- The QGP is an extremely good fluid
 - Medium response to Eloss must be collective
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 - Collectivity starts at short distance $1/T$ from the jet
 - There is a strong momentum flux along the jet direction
- We only model the generic contribution to (soft) particles from E&M conservation
 - Underestimates production at $p_T \gg T$ (model dependent)

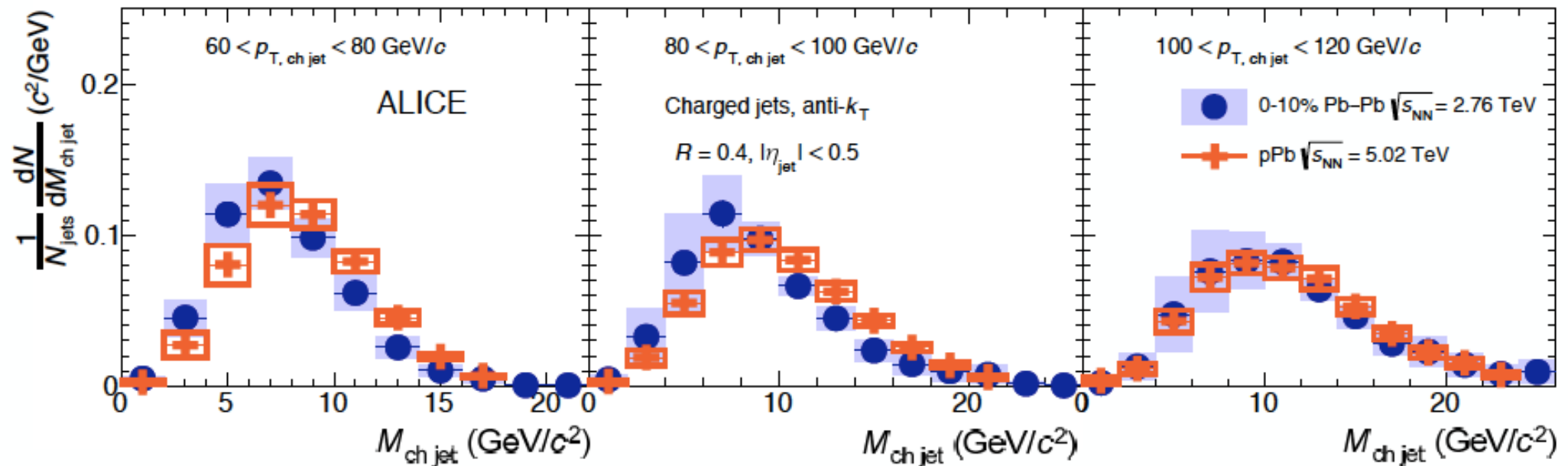
JCS, Shuryak & Teaney 06

Recovering Jet Energy



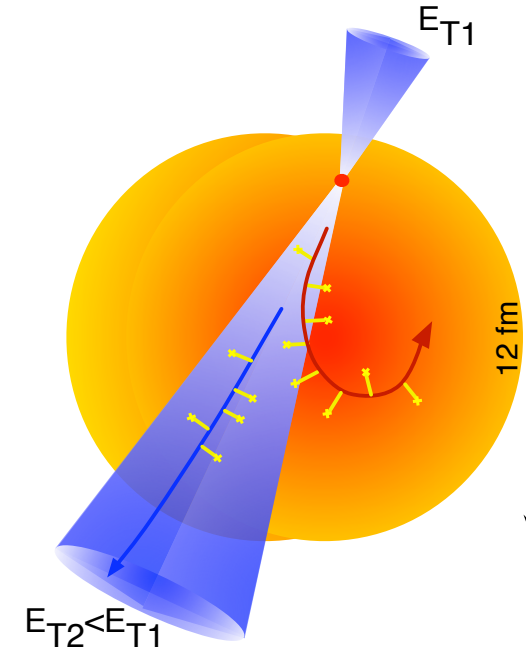
- Medium response completely fixed by Eloss
 - No additional parameters

Jet Masses



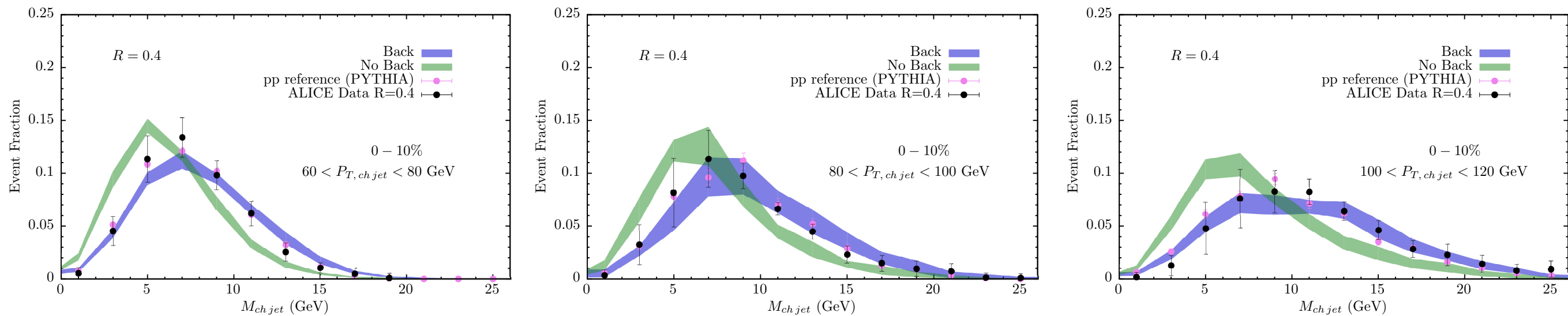
- Little sensitivity to strong quenching!

- Puzzling result
- Removing soft fragments \Rightarrow
Jet mass narrowing



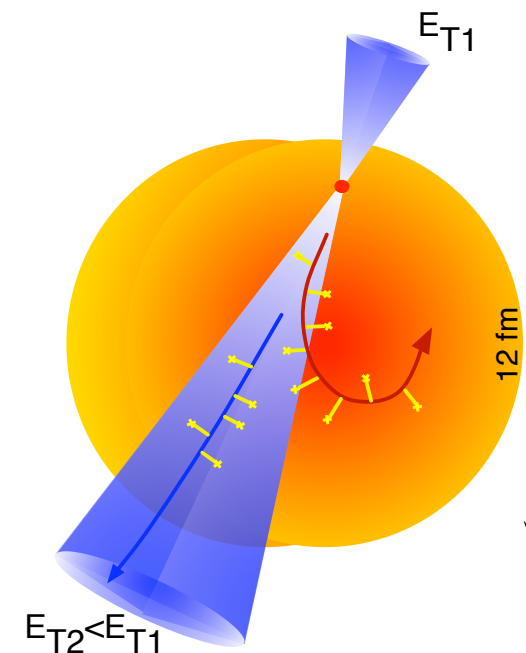
Jet Masses

JCS, Pablos, Hulcher, Milano, Rajagopal (in preparation)



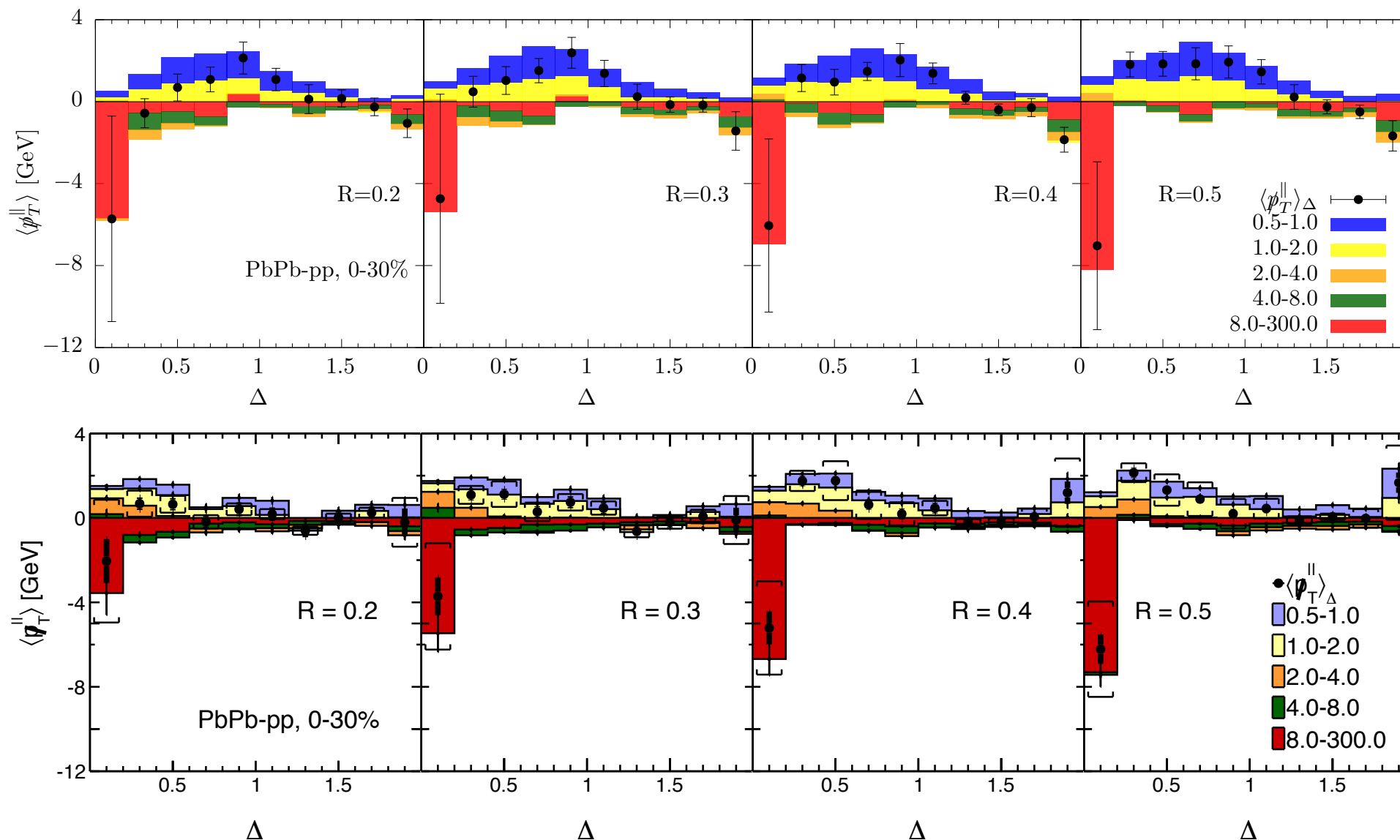
- Little sensitivity to strong quenching!

- Puzzling result
- Removing soft fragments ⇒
Jet mass narrowing



- Medium response regenerates the missing mass

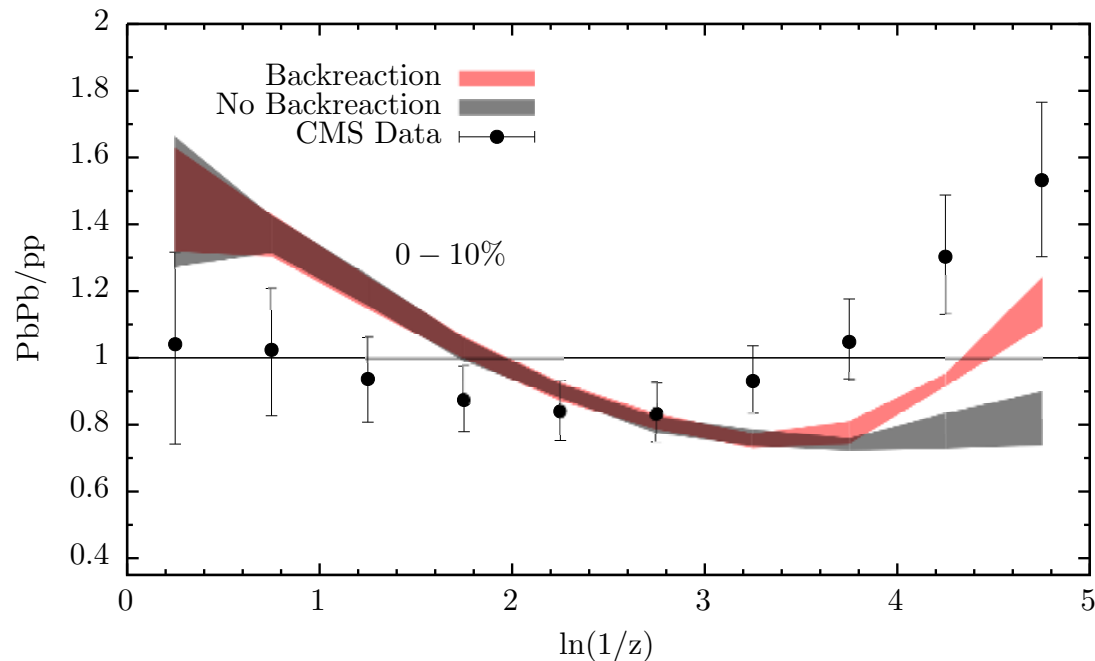
Soft Back-reaction



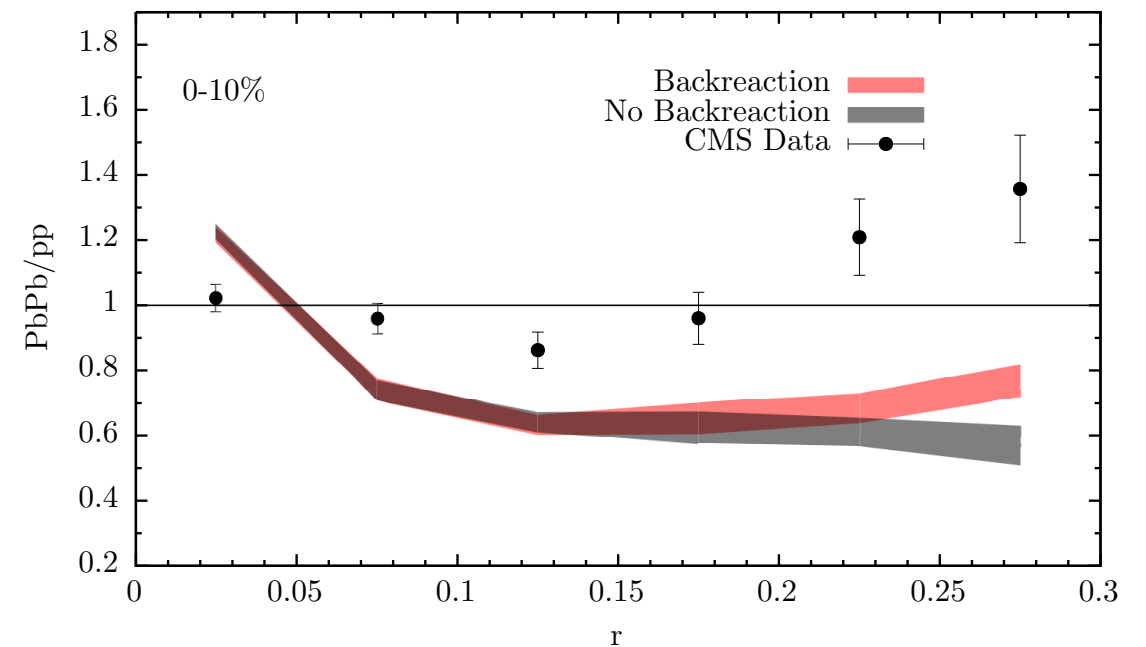
- Energy is recovered in soft ($\sim T$) particles
 - Expected deficiency of the treatment
 - But also in the region where incomplete thermalisation should appear (e.g. radiative processes)

Not Everything Works

Ratio of FF



Ratio of jet shapes



What is the origin of the discrepancy?

- Shape modification may be sensitive to perturbative emissions
- Treatment of back-reaction may be too crude
 - Other model implementations show sensitivity to recoil/back-reaction in these and other observables

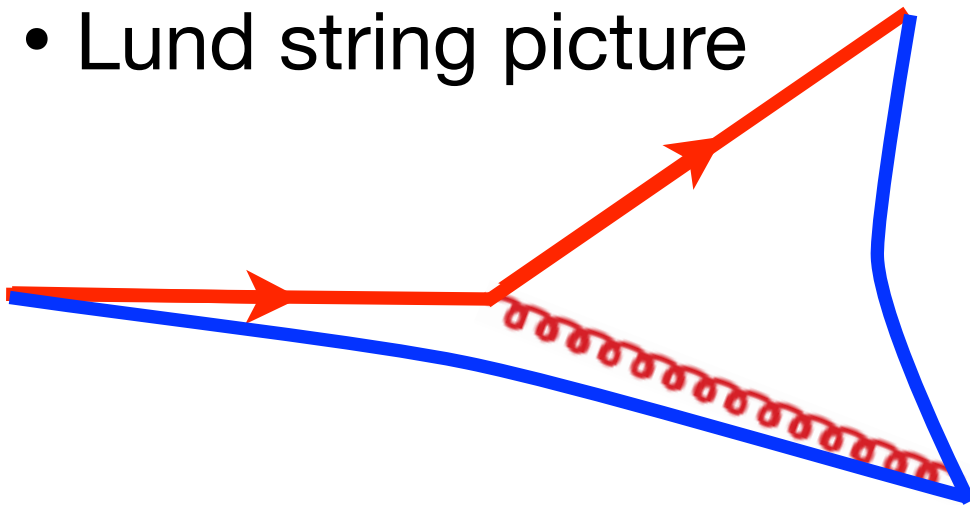
Tachibana, Chang and Quin, 17, Kunnawalkam, Elayavalli and Zapp 17, Milhano, Wiedemann and Zapp 17

- Additional physics processes may be required

Coherence at Strong Coupling

- Preparing “multi partonic” excitations in holography

- Lund string picture

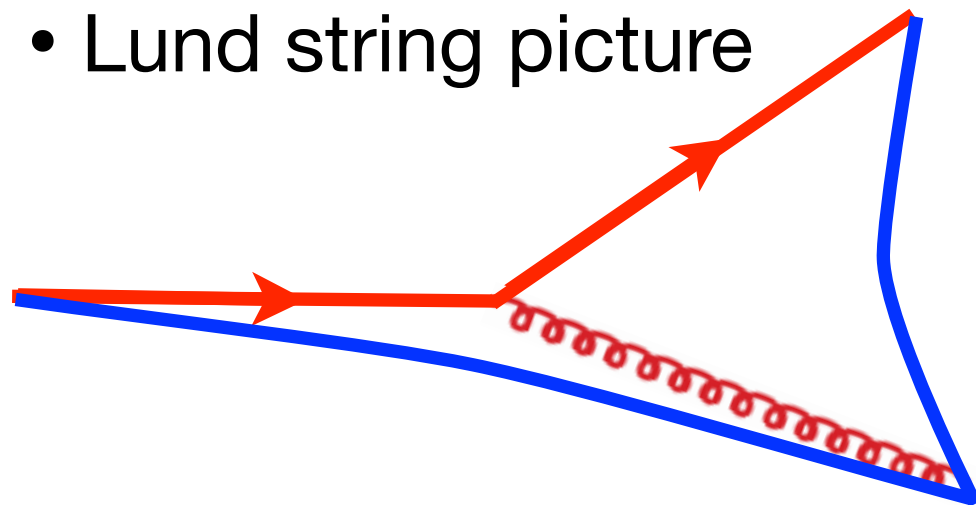


gluons \Rightarrow string with kinks

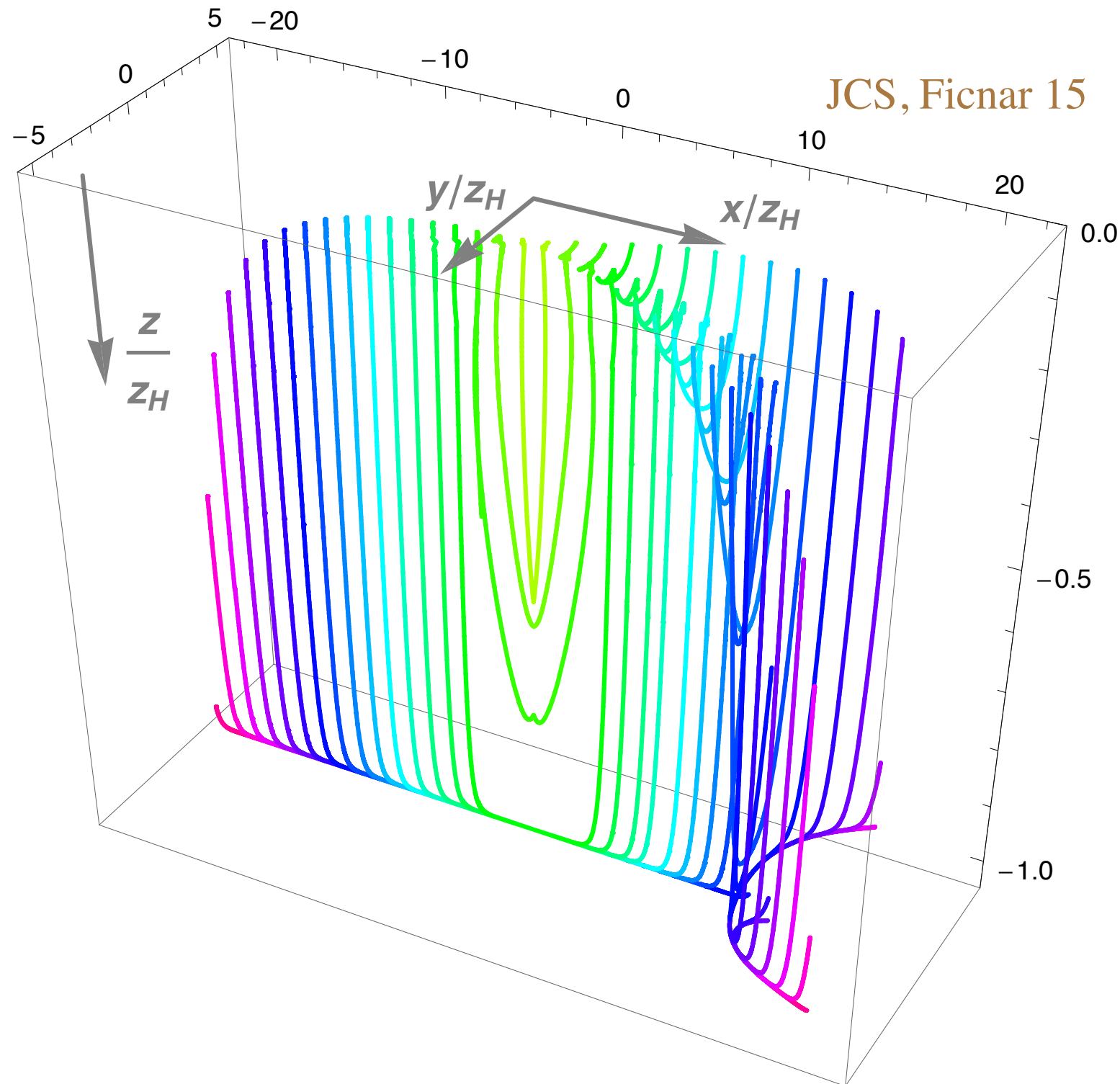
Coherence at Strong Coupling

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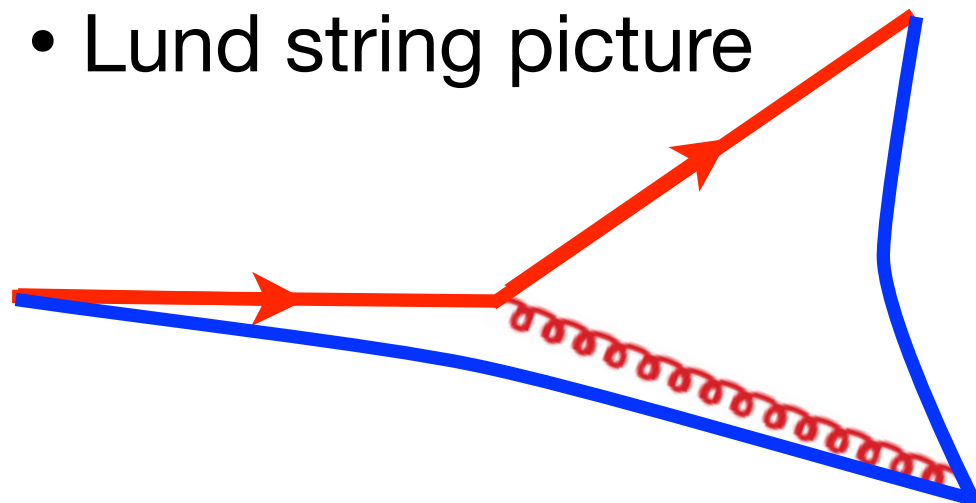
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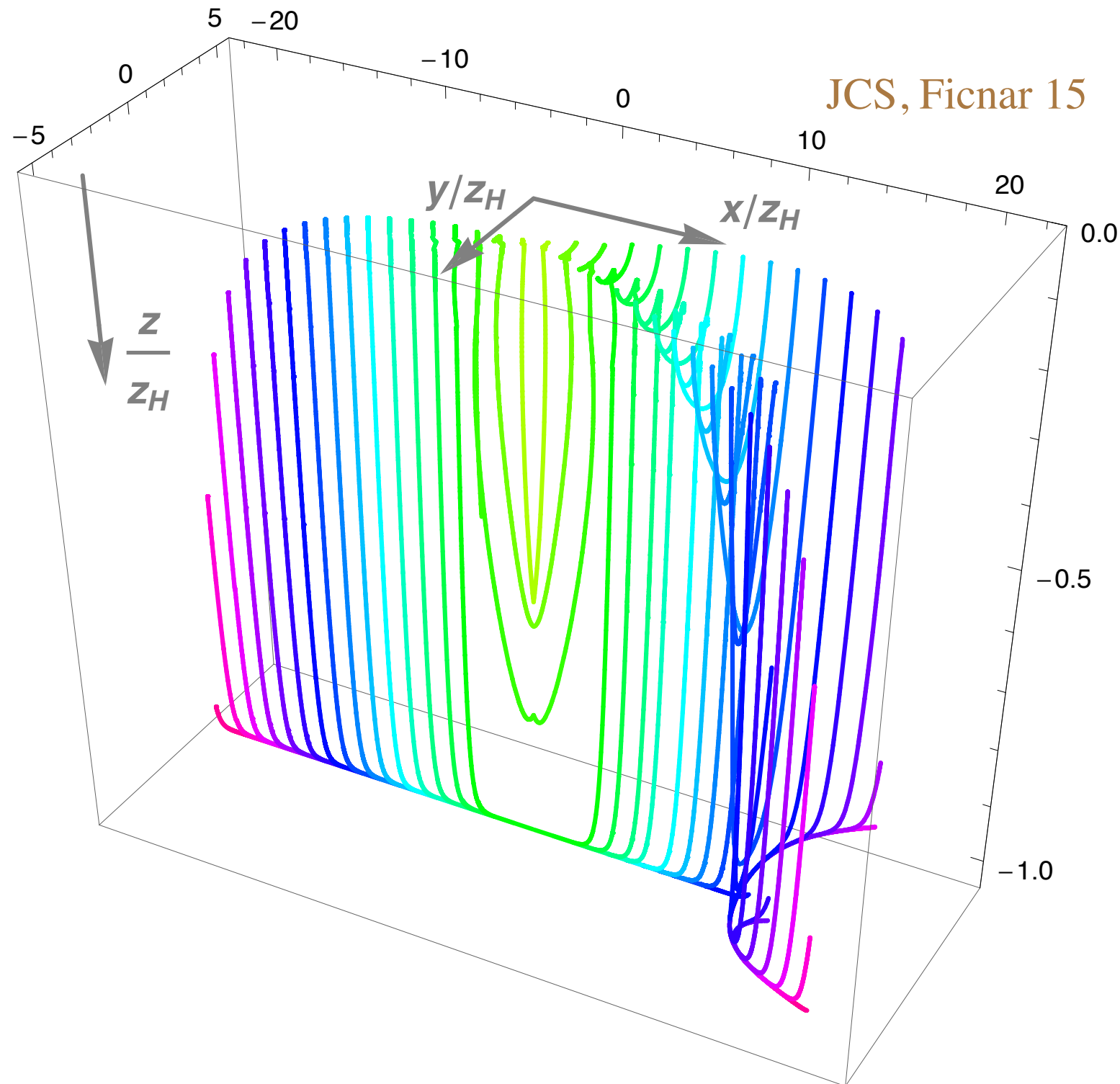
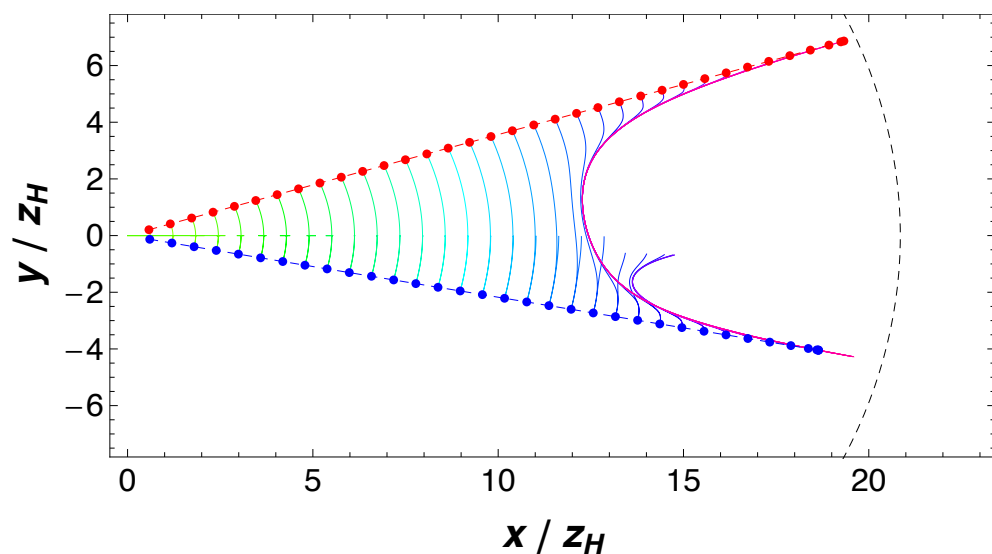
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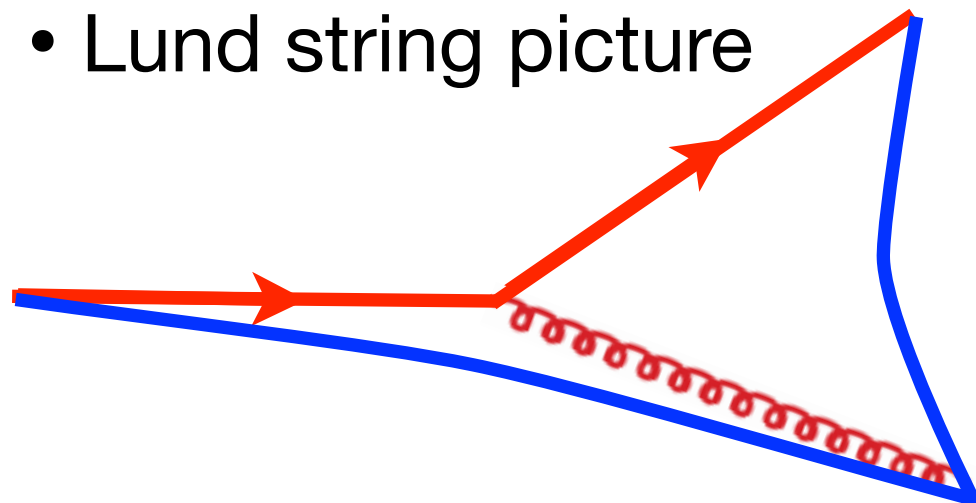
resolved jets



Coherence at Strong Coupling

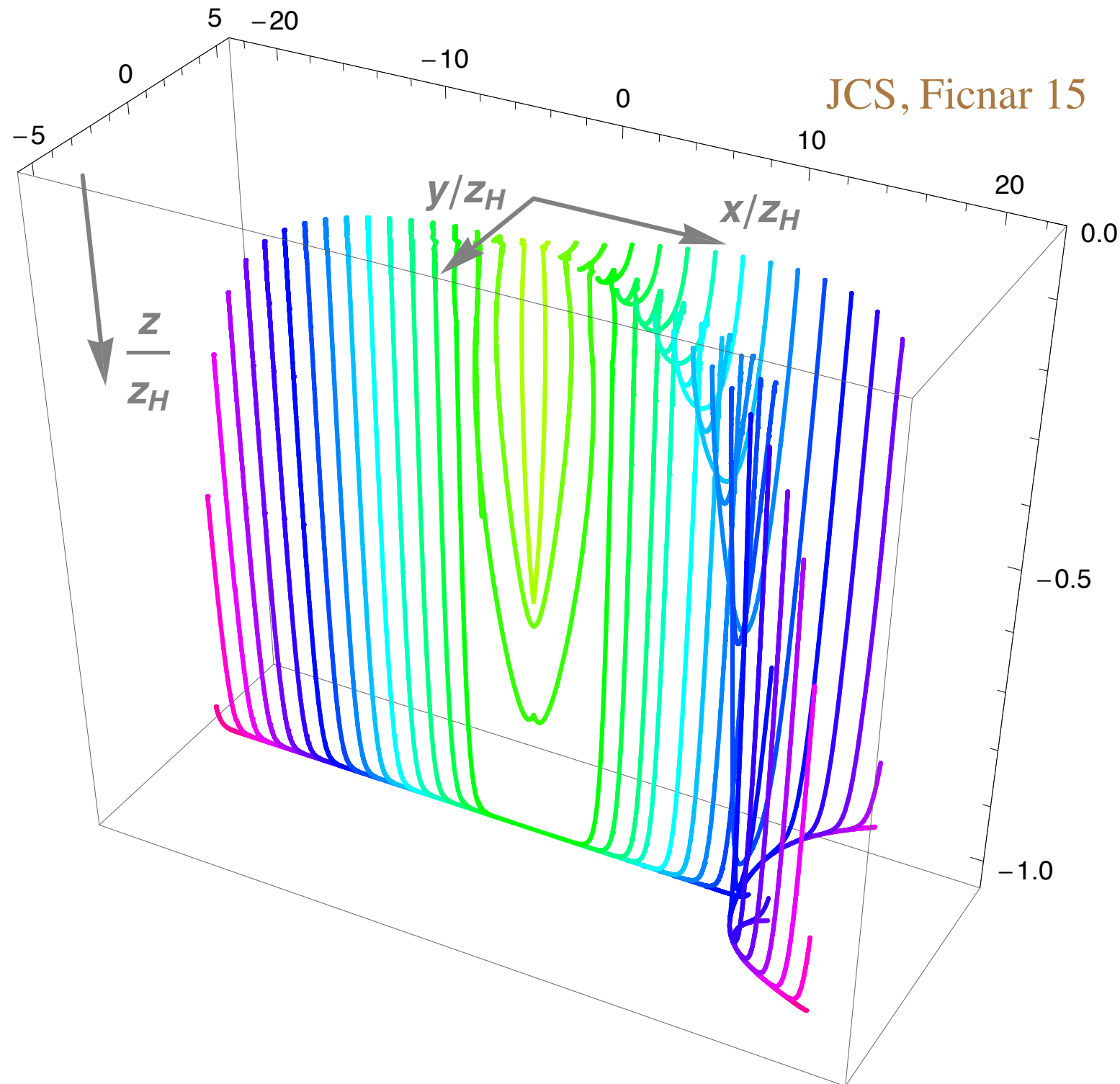
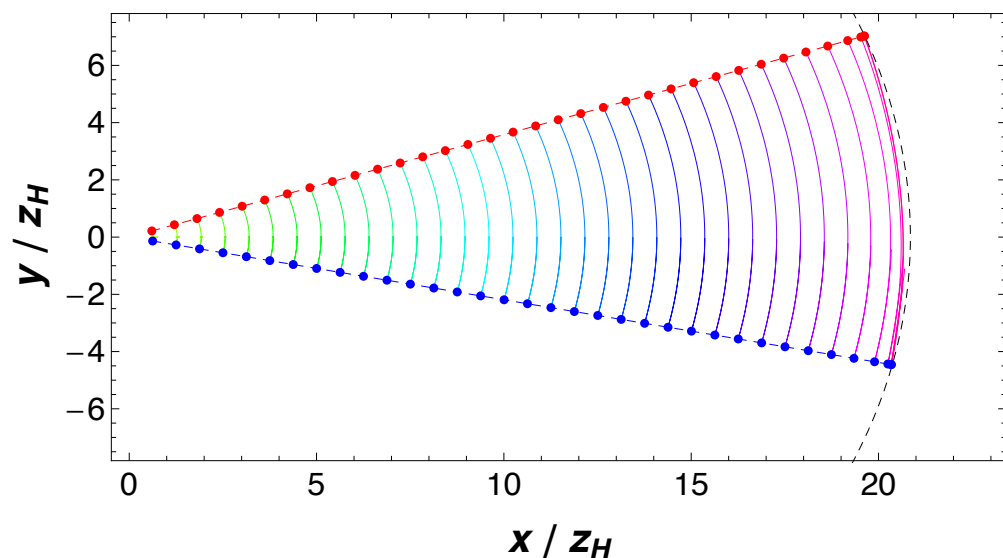
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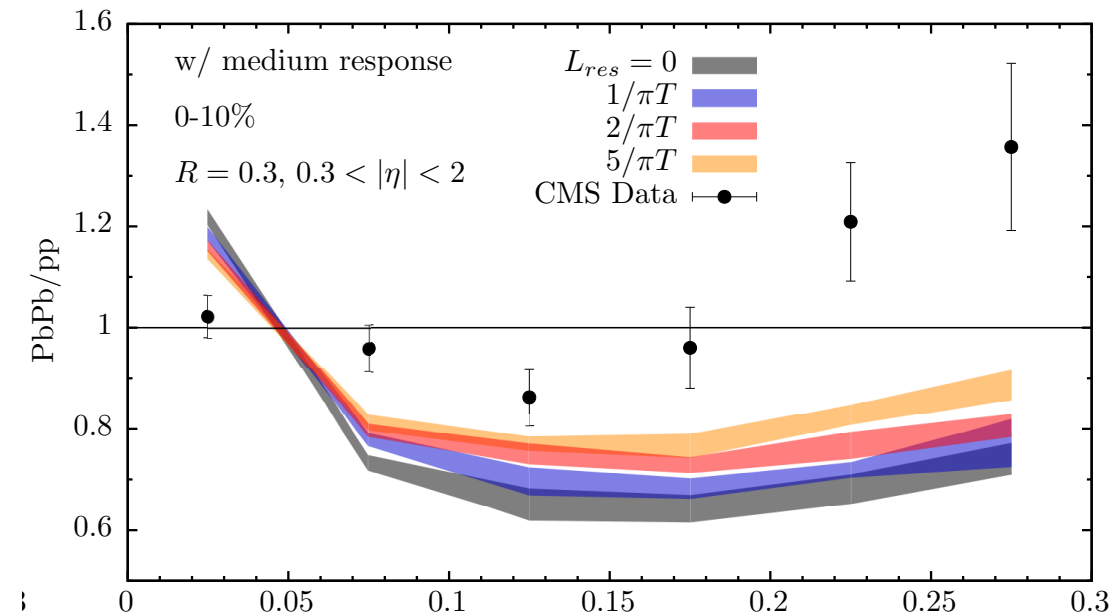
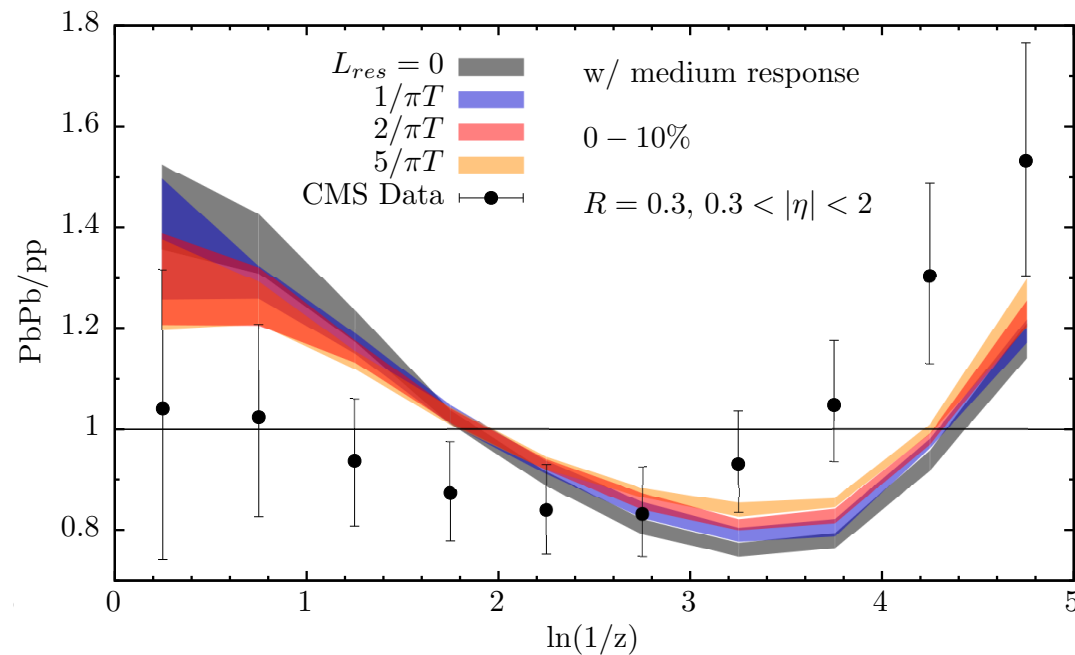
gluons \Rightarrow string with kinks

un-resolved jets



Resolution effects

Pablos, Hulcher, Rajagopal 17



Phenomenological implementation

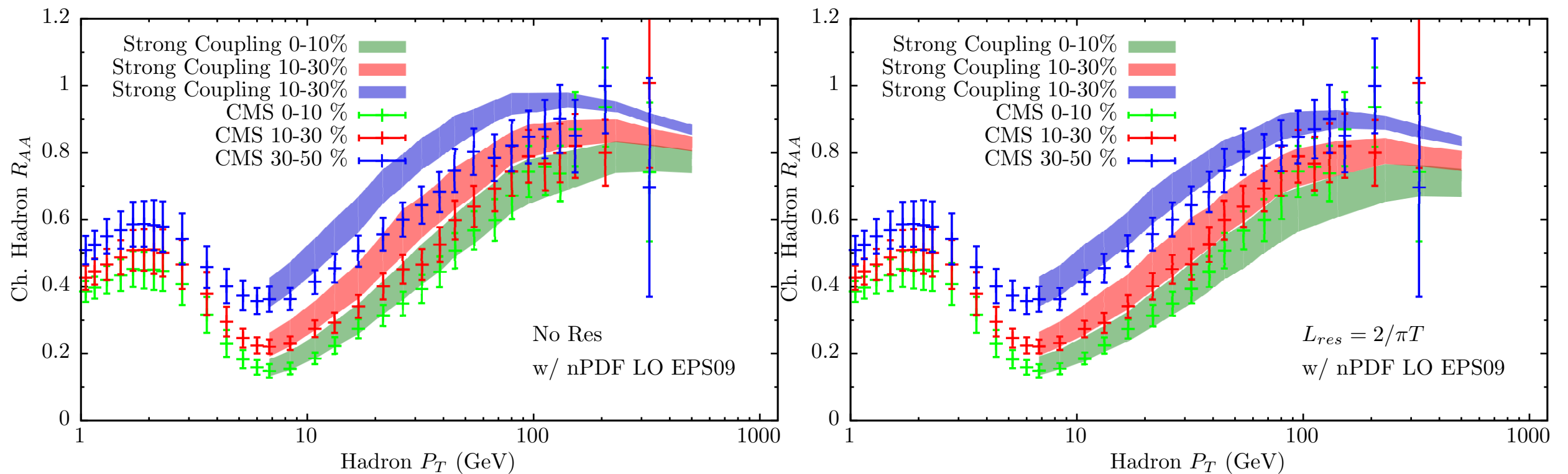
- Introduce a transverse resolution parameter $\pi T L_{res} \sim \mathcal{O}(1)$
- Partons in shower loose independently if $L > L_{res}$

Combination of resolution and back-reaction

- Pushes distributions in the right direction
- But still not enough...

Hadron RAA

JCS, Pablos, Hulcher, Milano, Rajagopal (in preparation)



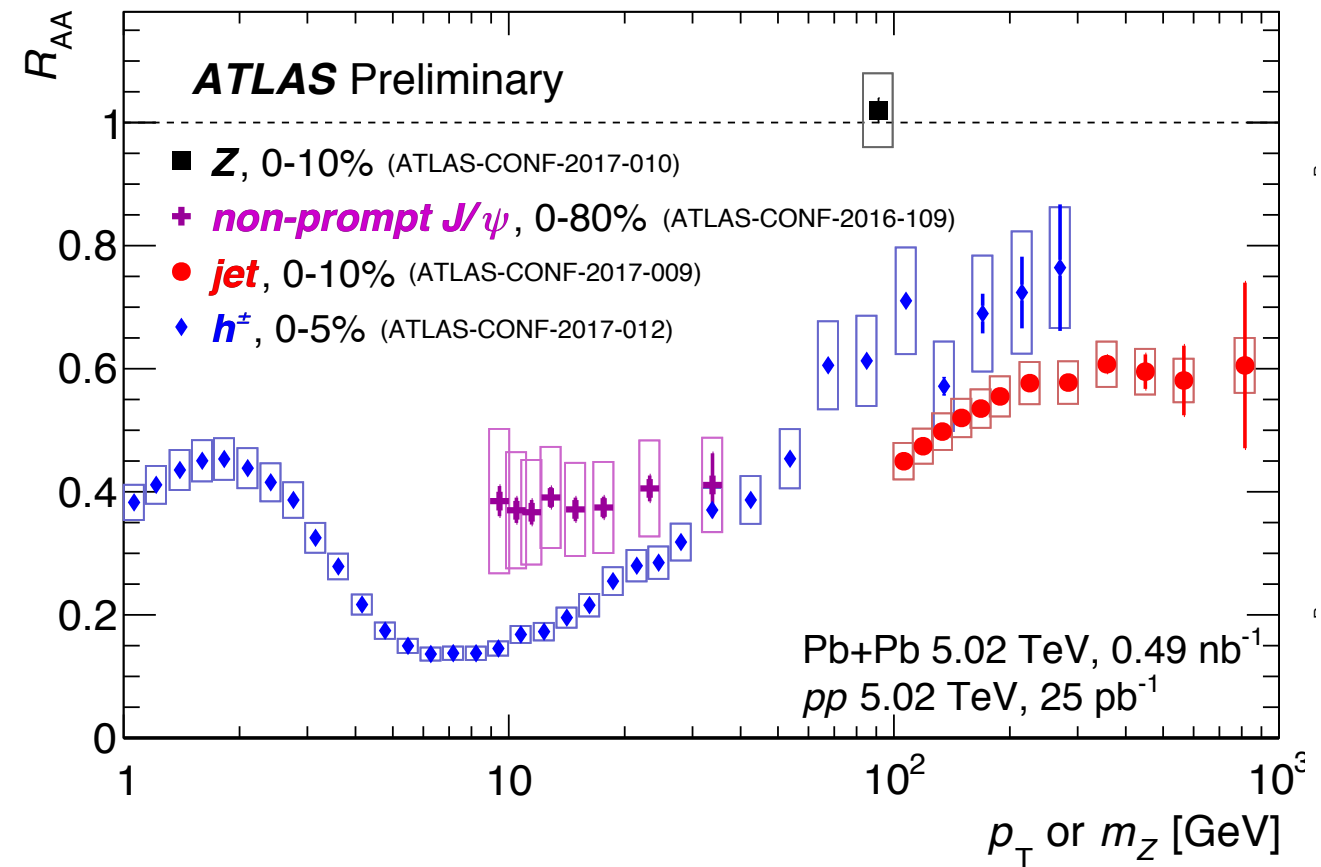
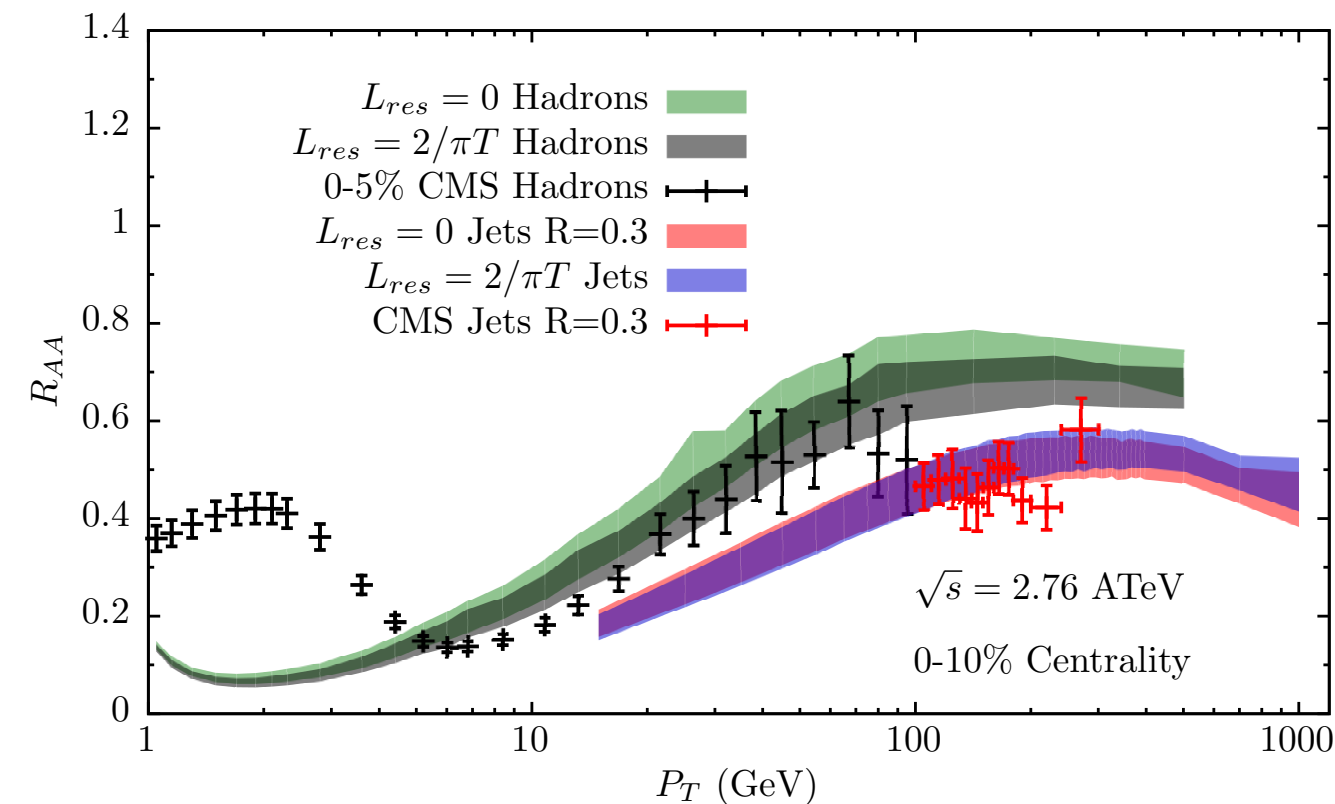
⊙ Resolution effects have an impact in the description of charged particle R_{AA}

⊙ Jet and charged particle R_{AA} show different sensitivity to resolution

also noted in Mehtar-Tani, Tywoniuk, 17

Particle vs Jets

JCS, Pablos, Hulcher, Milano, Rajagopal (in preparation)



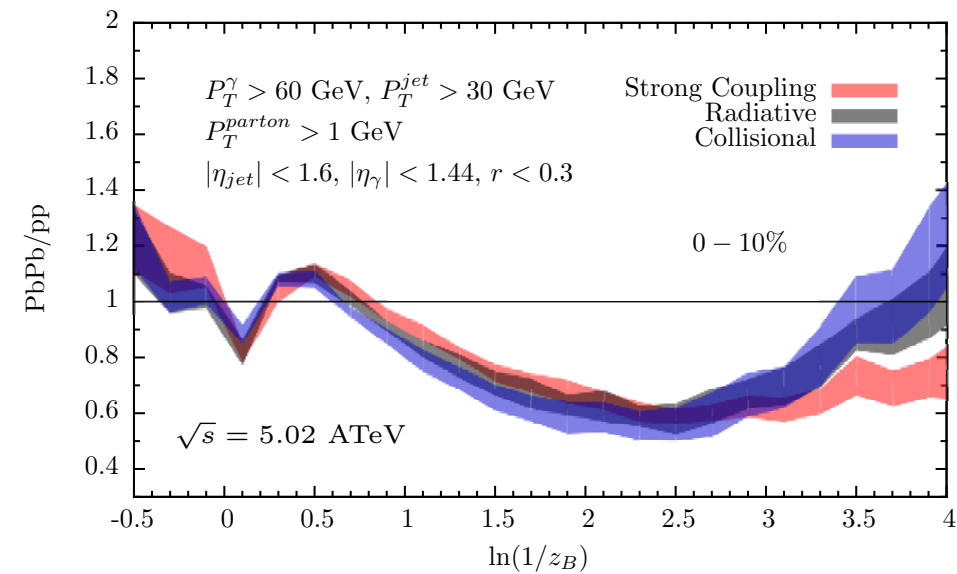
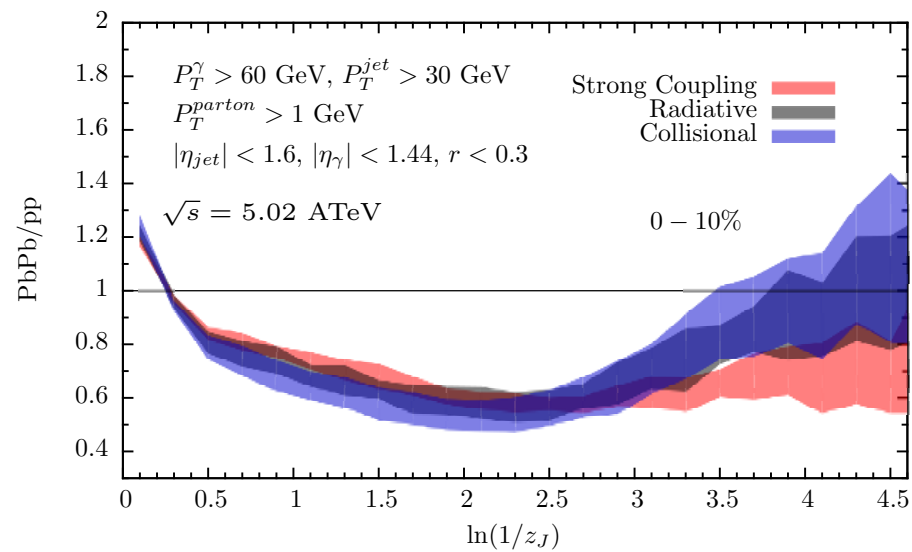
- Simultaneous description of Jet and hadron RAA
 - Including correct spectrum and flavour
 - NPDF
 - Fluctuation in jet structure (i.e. not all jets loose the same energy)

Conclusion

- ⦿ How can we discern the nature of the d.o.f from hard probes?
 - Not clear yet
 - We are exploring the consequences of a strongly coupled physics.
- ⦿ A simple model
 - Incorporates relevant physics from strong coupling
 - We can implement new physics processes
 - Is testable and predictive

Photon Fragmentation Functions

•Parton level



Inensitivity to Broadening

