

Measurement of the weak charge of the proton with Q_{weak}

Ciprian Gal
University of Virginia

The Qweak Collaboration

101 collaborators 26 grad students
11 post docs 27 institutions

Institutions:

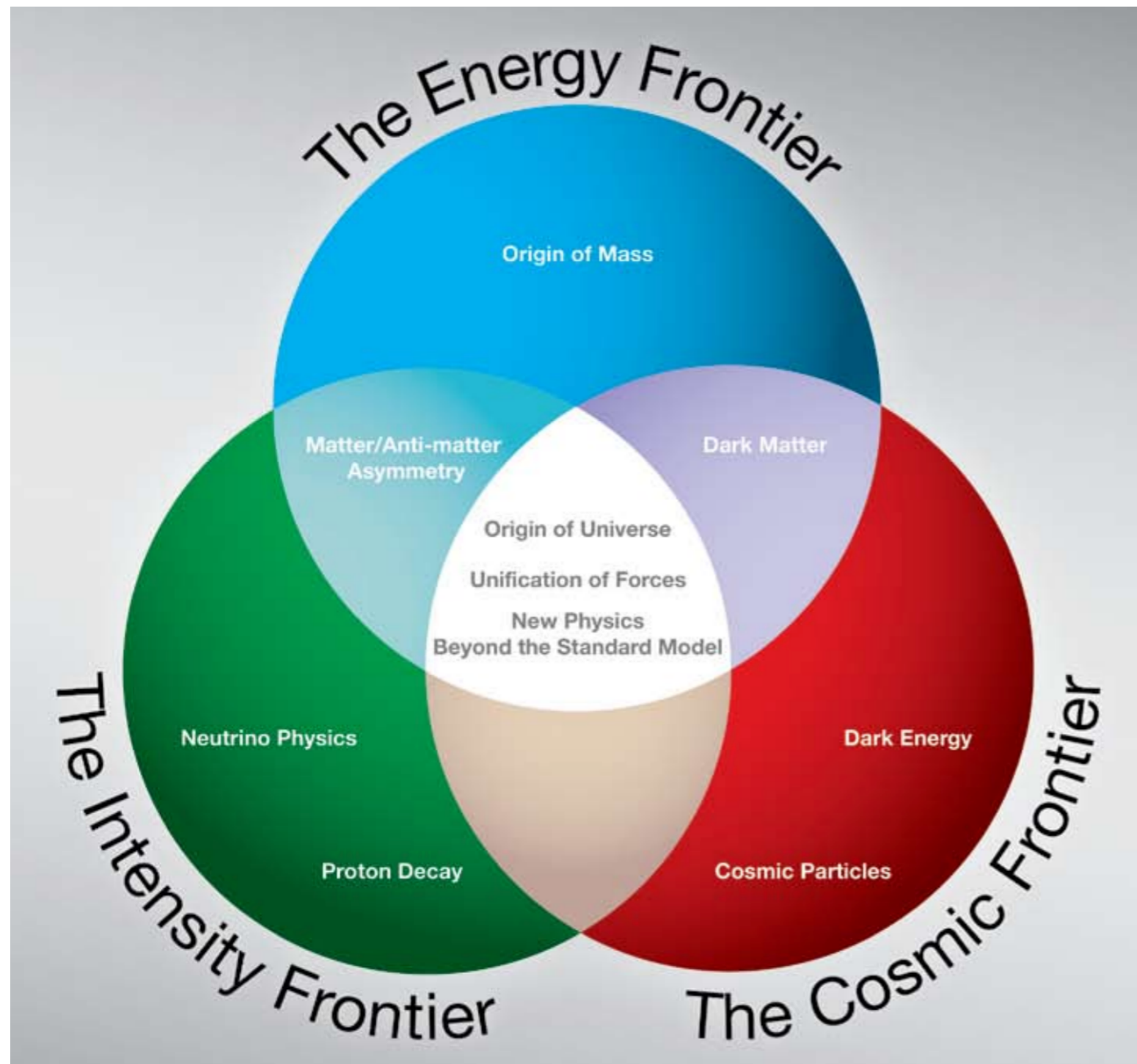
- 1 University of Zagreb
- 2 College of William and Mary
- 3 A. I. Alikhanyan National Science Laboratory
- 4 Massachusetts Institute of Technology
- 5 Thomas Jefferson National Accelerator Facility
- 6 Ohio University
- 7 Christopher Newport University
- 8 University of Manitoba,
- 9 University of Virginia
- 10 TRIUMF
- 11 Hampton University
- 12 Mississippi State University
- 13 Virginia Polytechnic Institute
- 14 Southern University at New Orleans
- 15 Idaho State University
- 16 Louisiana Tech University
- 17 University of Connecticut
- 18 University of Northern British Columbia
- 19 University of Winnipeg
- 20 George Washington University
- 21 University of New Hampshire
- 22 Hendrix College, Conway
- 23 University of Adelaide
- 24 Syracuse University
- 25 Duquesne University



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Spokespersons Project Manager Grad Students *deceased

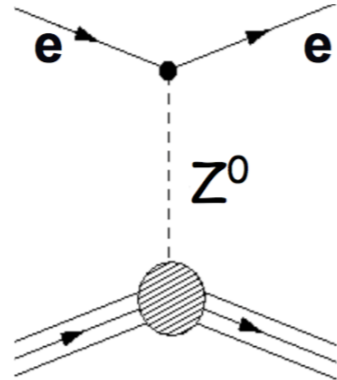
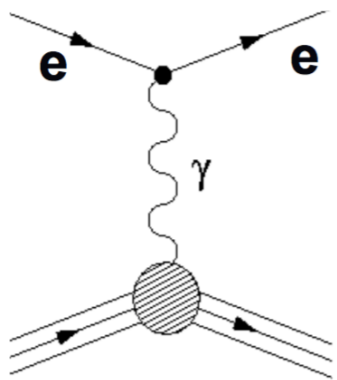
The path to discovery



https://science.energy.gov/~media/hep/pdf/files/pdfs/p5_report_06022008.pdf

- In order to reveal the way nature works a multi-pronged approach is needed
- From astronomical observations (cosmic frontier) to direct measurements (energy frontier) to indirect measurements (intensity frontier)
 - some of our most tantalizing results have come from indirect searches where we get hints at what could lie ahead
- Each of these paths comes with their own pros and cons

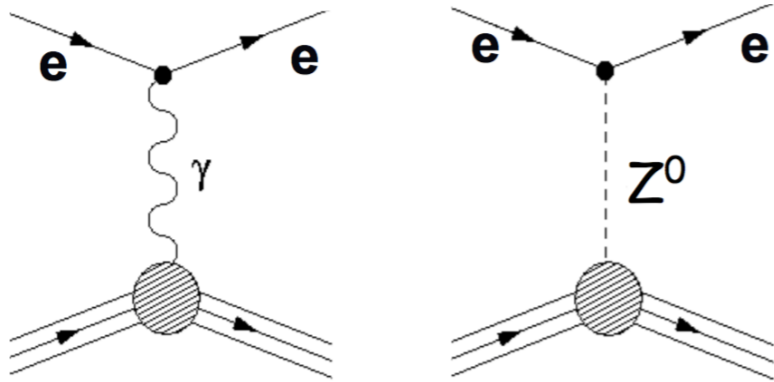
Particle physicists' intro to PVES



$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{\frac{\text{diagram with } \gamma \text{ and } Z^0}{\text{diagram with } \gamma}}{2} \propto \frac{|\mathcal{M}_Z|}{|\mathcal{M}_\gamma|}$$

- uses longitudinally polarized electron beams
- measures asymmetries that are generally on the level of ppm or less

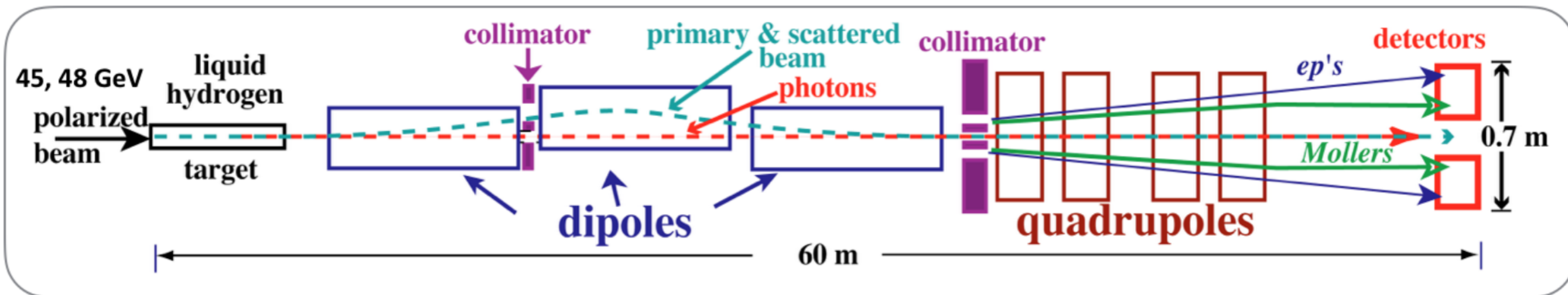
Particle physicists' intro to PVES



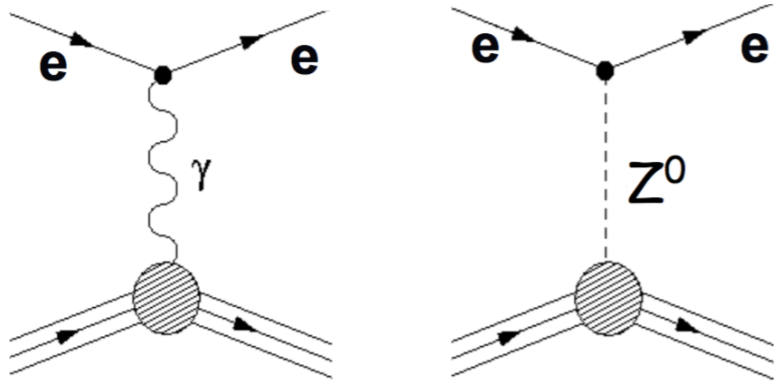
$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{\frac{\text{diagram with } \gamma \text{ and } Z^0}{\left| \text{diagram with } \gamma \right|^2}}{\propto} \frac{|\mathcal{M}_Z|}{|\mathcal{M}_\gamma|}$$

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- measures asymmetries that are generally on the level of ppm or less

Kinematic cuts are already made when the experiment starts:



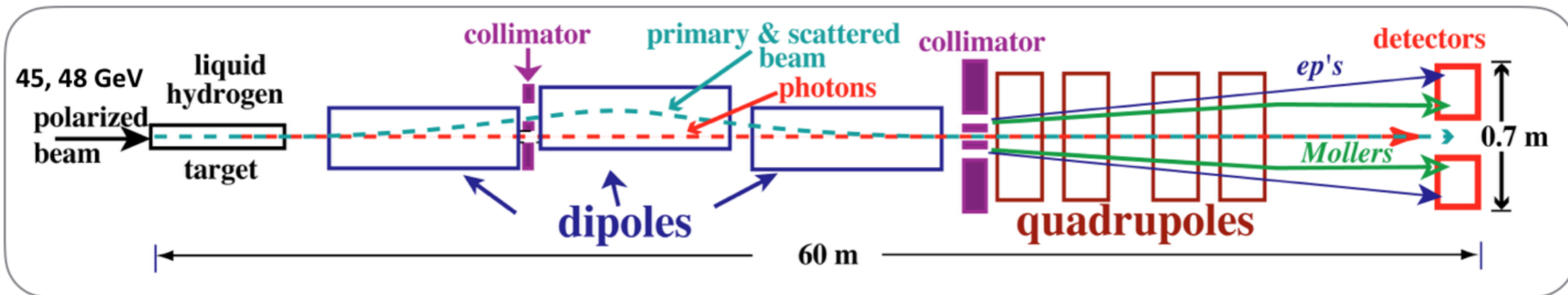
Particle physicists' intro to PVES



$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L} \sim \frac{\frac{\text{diagram with } \gamma}{\text{diagram with } \gamma} + \frac{\text{diagram with } Z^0}{\text{diagram with } \gamma}}{2} \propto \frac{|\mathcal{M}_Z|}{|\mathcal{M}_\gamma|}$$

- uses longitudinally polarized electron beams
- measures asymmetries that are generally on the level of ppm or less

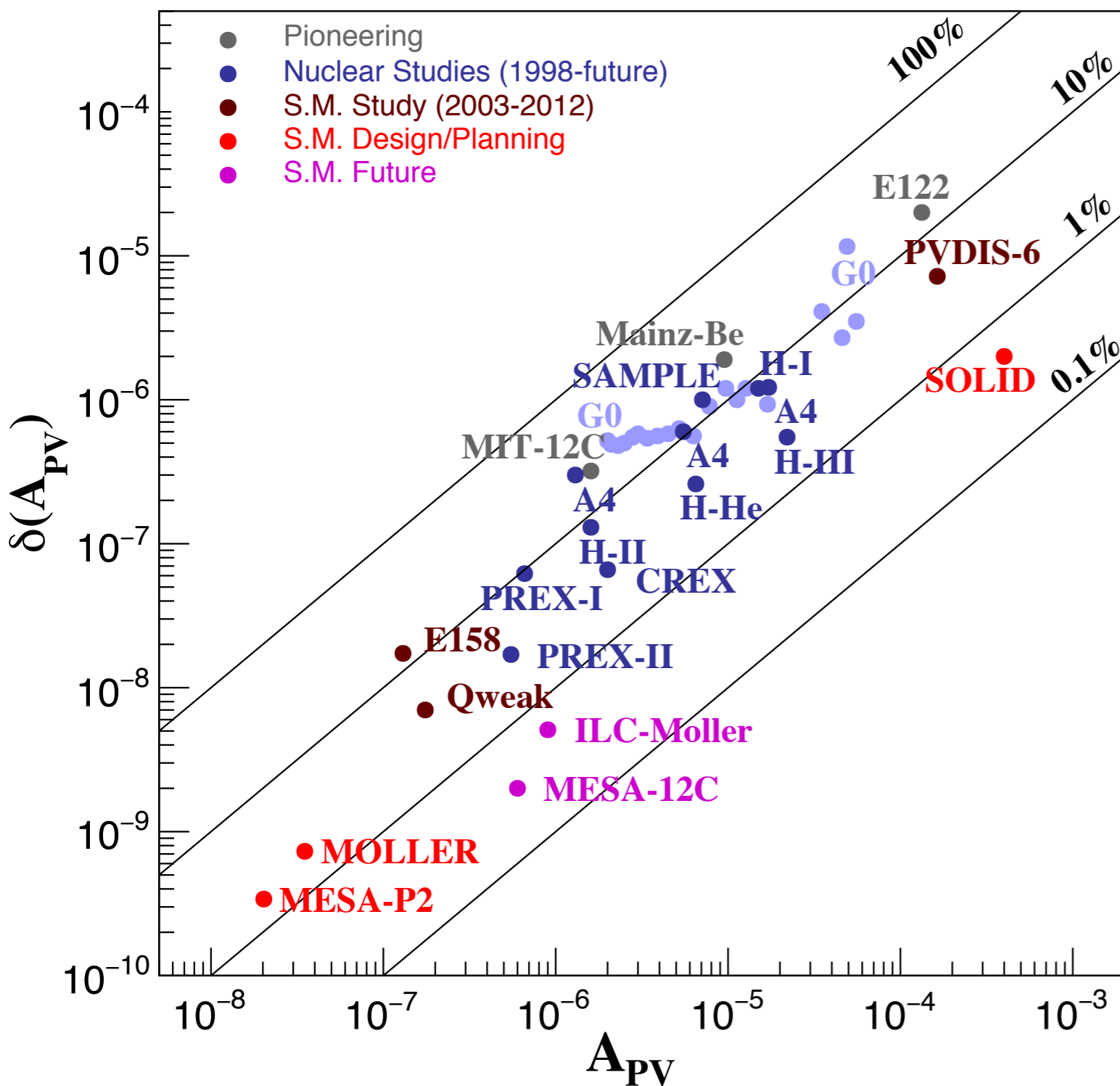
Kinematic cuts are already made when the experiment starts:



Allows for the collection of large amounts of data (100s of MHz) needed to resolve small asymmetries:

E158:
 $A_{PV} = (-131 \pm 14 \pm 10) \text{ ppb}$

History of PVES



- PVES has a long history of pushing the limits of precision and discovery

- E122: ($\Delta A = 10$ ppm)

- G0, A4, HAPPEX ($\Delta A = 0.25$ to 2 ppm)

- E158 ($\Delta A = 17$ ppb)

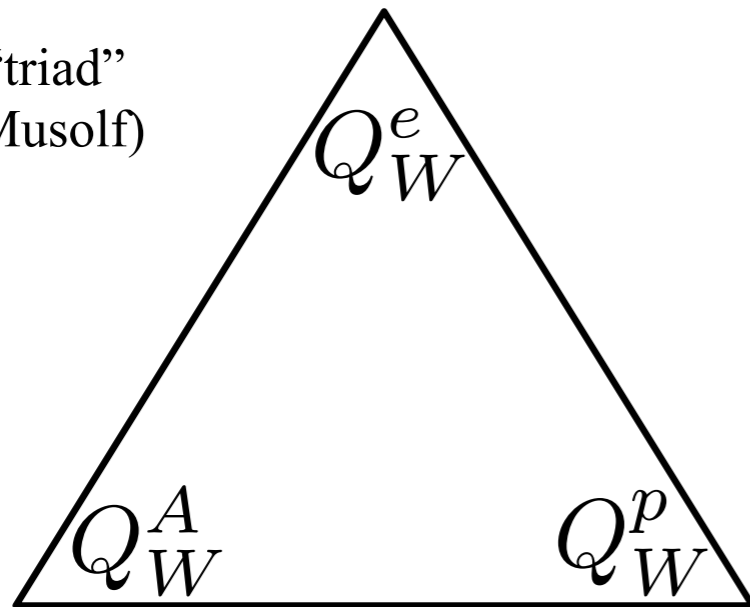
- Qweak ($\Delta A = 9$ ppb)

- Moller ($\Delta A = 0.8$ ppb)

- P2 ($\Delta A = 0.34$ ppb)

Electroweak measurements

Weak charge “triad”
(M. Ramsey-Musolf)



- In the early 2000s E158 made the first measurement of electron weak charge Q_W^e
- Atomic Parity Violation measurements on ^{133}Cs gave unique insights into d-quark weak vector charge
- Finally Qweak directly measures the proton weak vector charge Q_W^p

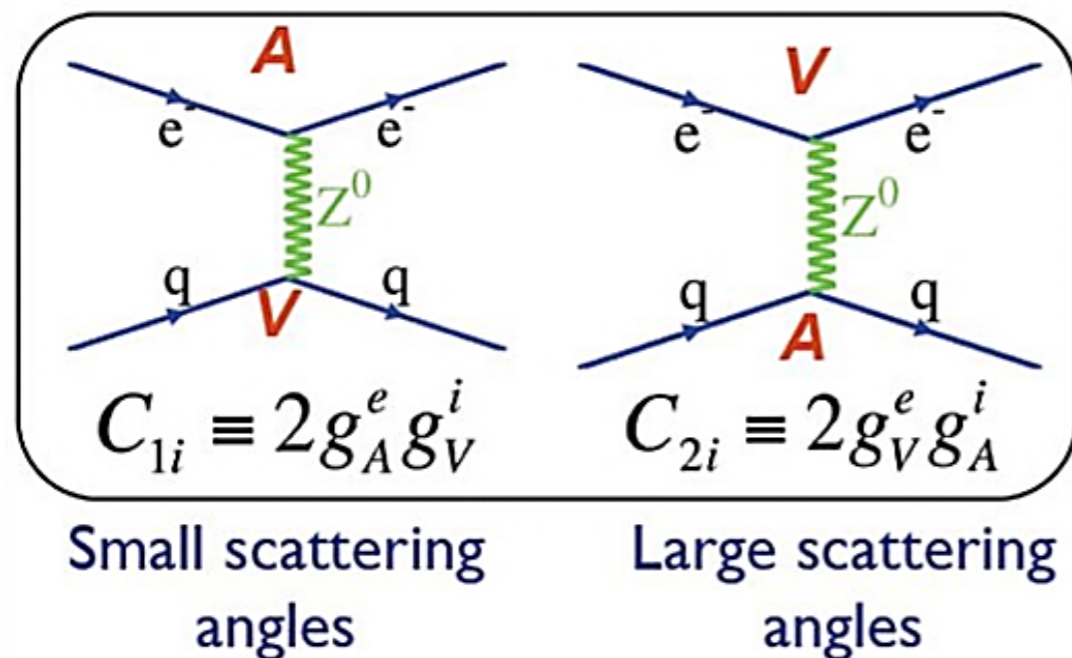
- Weak charge is the analog to the electric charge:

Particle	Electric charge	Weak vector charge ($\sin^2 \theta_W \approx \frac{1}{4}$)
e	-1	$Q_W^e = -1 + 4 \sin^2 \theta_W \approx 0$
u	$+\frac{2}{3}$	$-2C_{1u} = +1 - \frac{8}{3} \sin^2 \theta_W \approx +\frac{1}{3}$
d	$-\frac{1}{3}$	$-2C_{1d} = -1 + \frac{4}{3} \sin^2 \theta_W \approx -\frac{2}{3}$
p(uud)	+1	$Q_W^p = 1 - 4 \sin^2 \theta_W \approx 0$
n(udd)	0	$Q_W^n = -1$

- also defined as $Q^2 \rightarrow 0$ (intrinsic property of particle)
- proton and electron have nearly 0 weak charge
- combined with the very well defined SM prediction makes it a good place to look for deviations (and new physics)

Quark Vector couplings \sim contact interaction

$$\mathcal{L}_{eq}^{PV} = -\frac{G_F}{\sqrt{2}} \sum_i [C_{1i} \bar{e} \gamma_\mu \gamma_5 e \bar{q} \gamma^\mu q + C_{2i} \bar{e} \gamma_\mu e \bar{q} \gamma^\mu \gamma_5 q] + \mathcal{L}_{new}^{PV}$$



- At low Q^2 ($Q^2 \ll M_Z^2$) the SM Lagrangian is effectively a 4-fermion contact interaction
- Qweak is sensitive to quark vector couplings C_{1u} and C_{1d}

Tree-level Qweak asymmetry

$$A_{PV} = \frac{-G_F Q^2}{4\pi\alpha\sqrt{2}} \left[\frac{\varepsilon G_E^\gamma G_E^Z + \tau G_M^\gamma G_M^Z - (1 - 4 \sin^2 \theta_W) \varepsilon' G_M^\gamma G_A^e}{\varepsilon (G_E^\gamma)^2 + \tau (G_M^\gamma)^2} \right]$$

At forward scattering angles and low 4-momentum transfer (Q^2):

$$A_{PV} = \underbrace{\frac{-G_F Q^2}{4\pi\alpha\sqrt{2}}}_{A_0} [Q_W^p + Q^2 B(Q^2, \theta)]$$

- Unlike measurements on Q_W^e , Qweak asymmetry needs to take care of the hadronic part of the interaction
 - small Q^2 makes the contributions smaller compared to previous experiments (proton looks like a point particle)
- The hadronic contributions can be determined from the previous PVES experiments

Qweak experimental setup



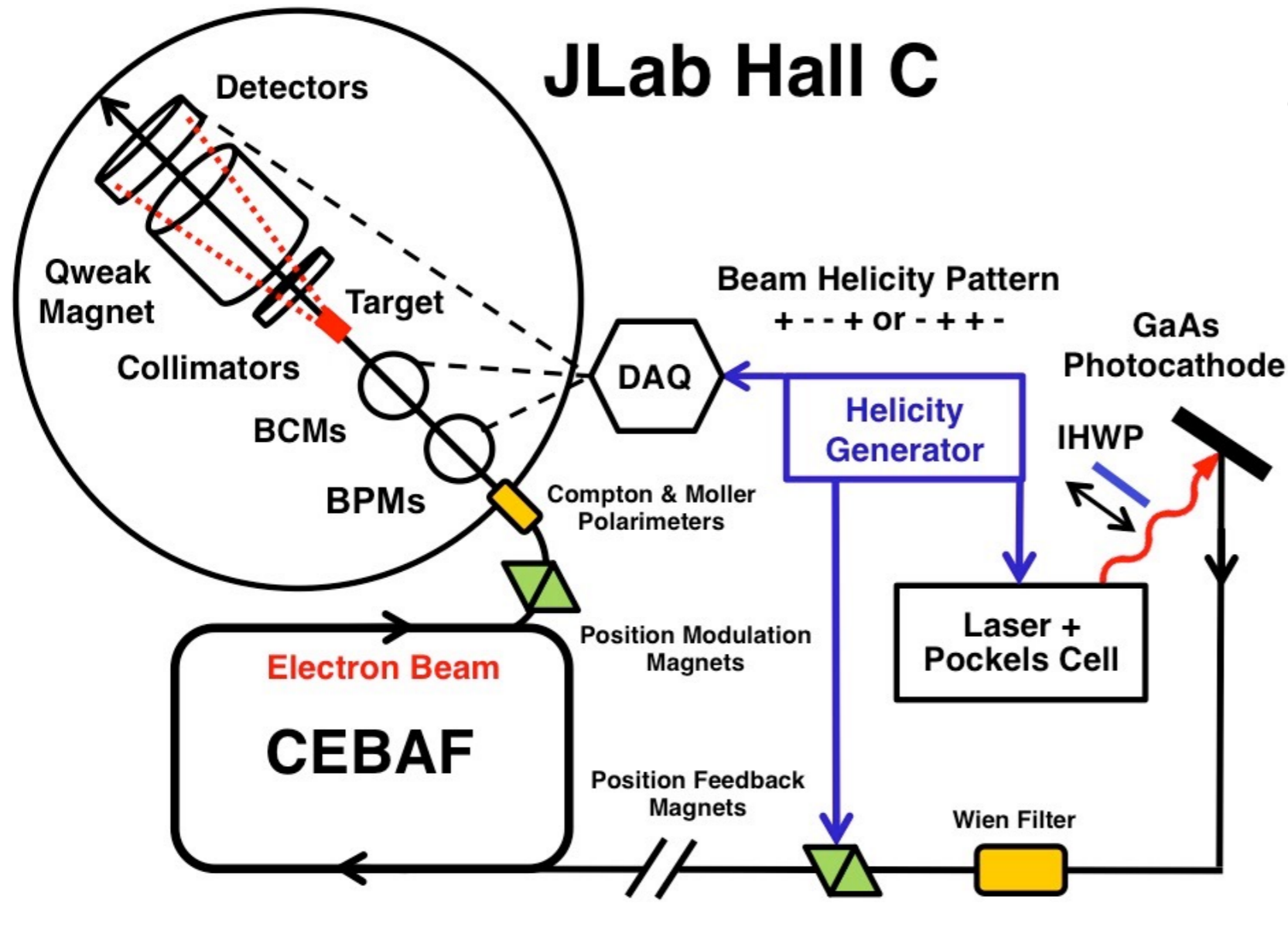
- Ran in three periods between 2010 and 2012 at the Continuous Electron Beam Accelerating facility at DOE's Jefferson Lab
- Commissioning data published in **PRL 111, 141803 (2013)**

$$Q_w^p(SM) = 0.0708 \pm 0.0003$$

initial result: $Q_w^p(PVES) = 0.064 \pm 0.012$

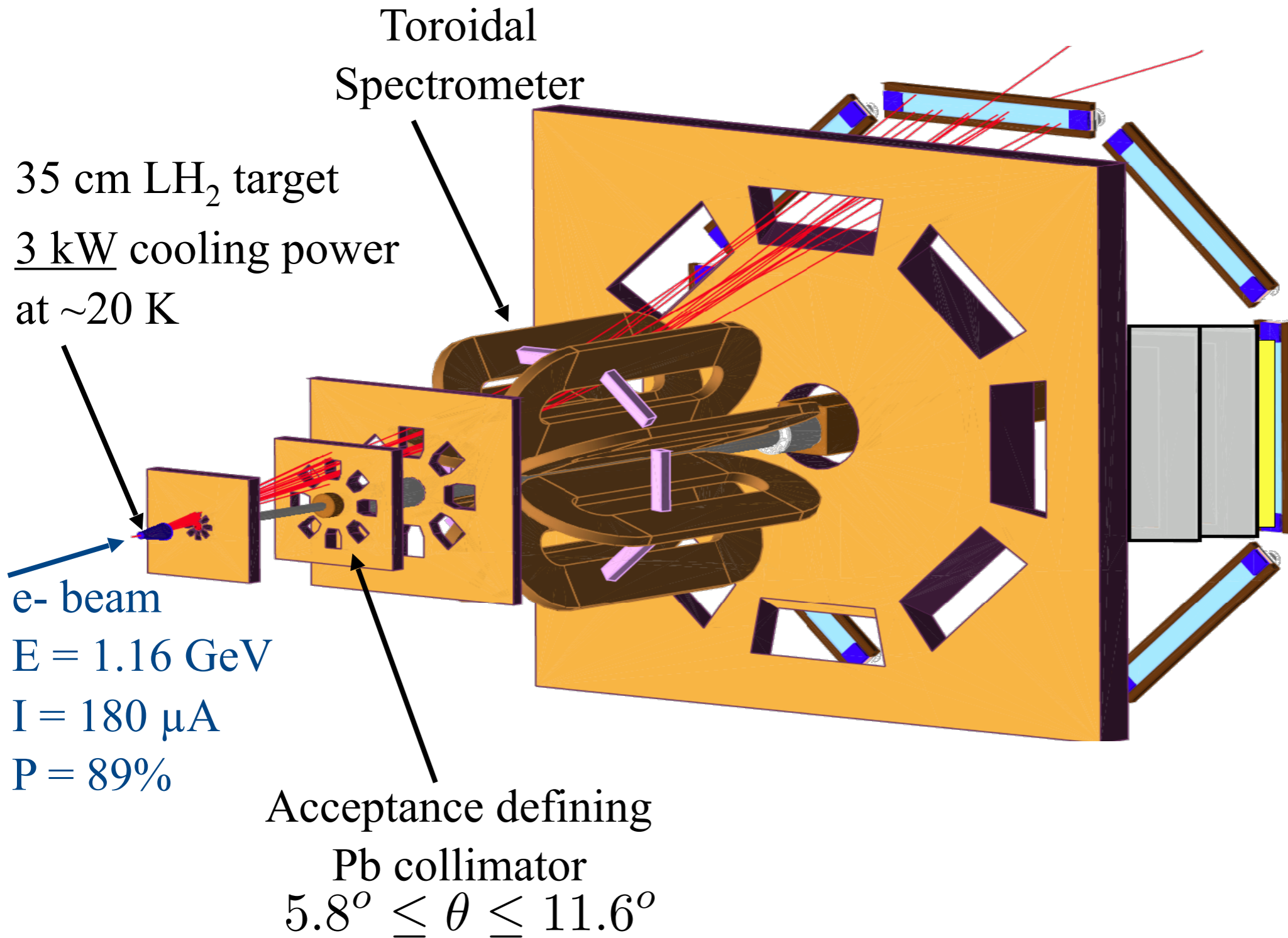
final result: $Q_w^p(PVES) = \quad \pm 0.0045$

Qweak experimental setup

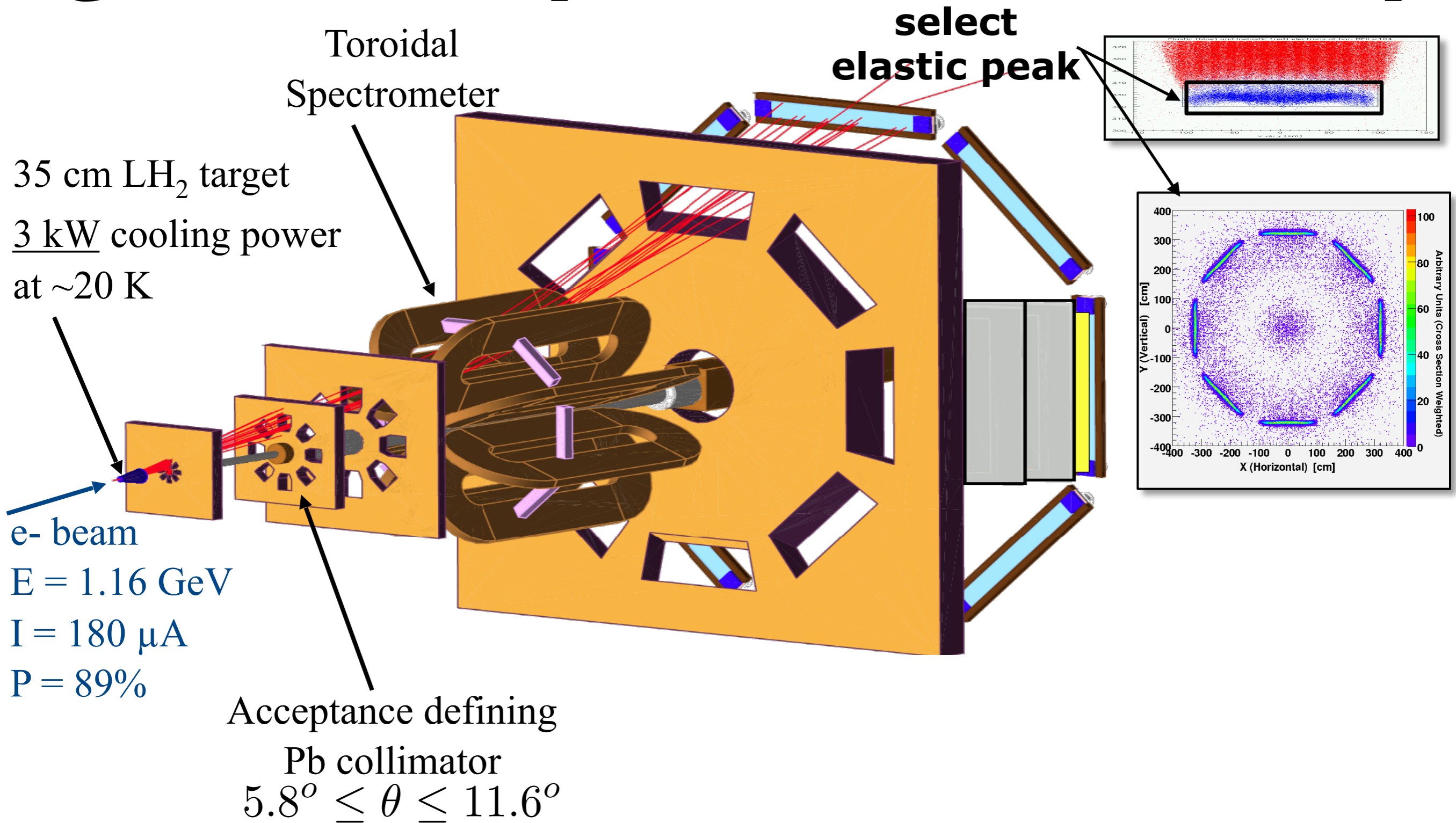


- Careful preparation of the entire machine (from injector all the way to the experiment) has to be made
- The injector was setup to have a fast helicity reversal at 1kHz
- High polarization of the electron beam had to be maintained
- Special care was taken to avoid helicity correlated beam asymmetries (intensity, beam positions)
- Checks were done throughout the machine until the target to make sure beam quality was maintained

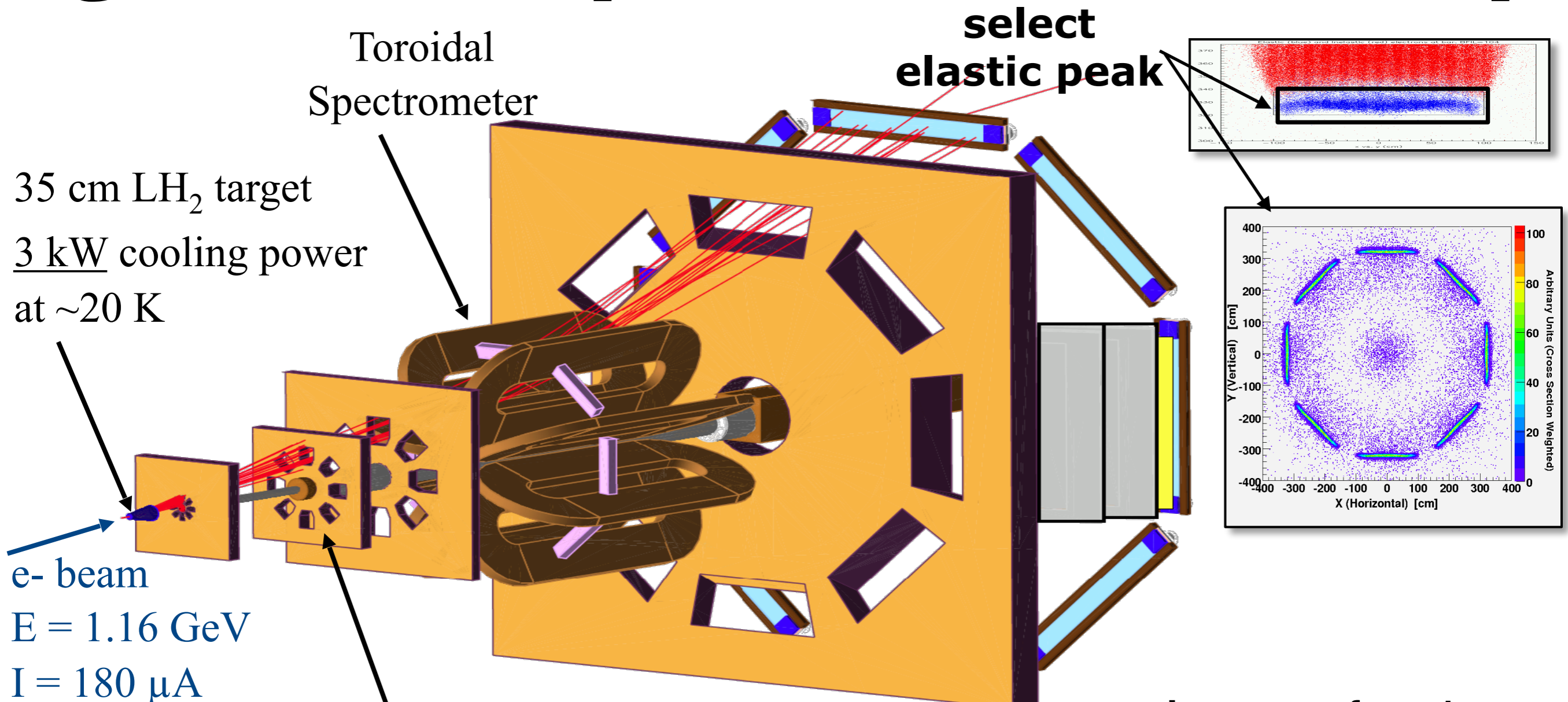
Qweak experimental setup



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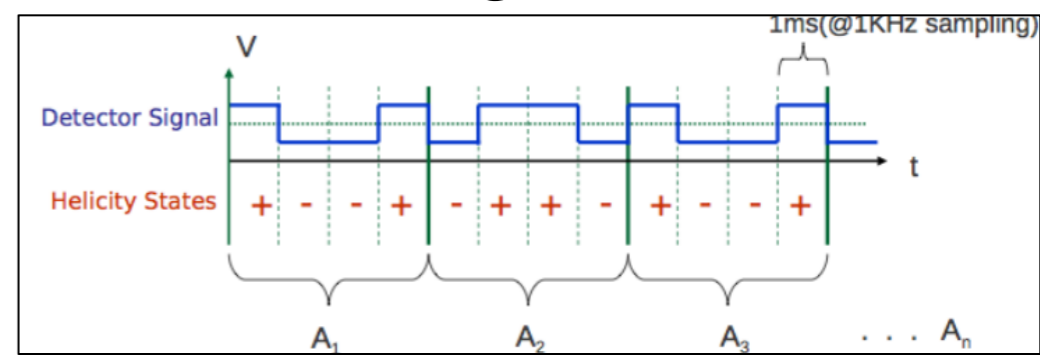


Qweak experimental setup



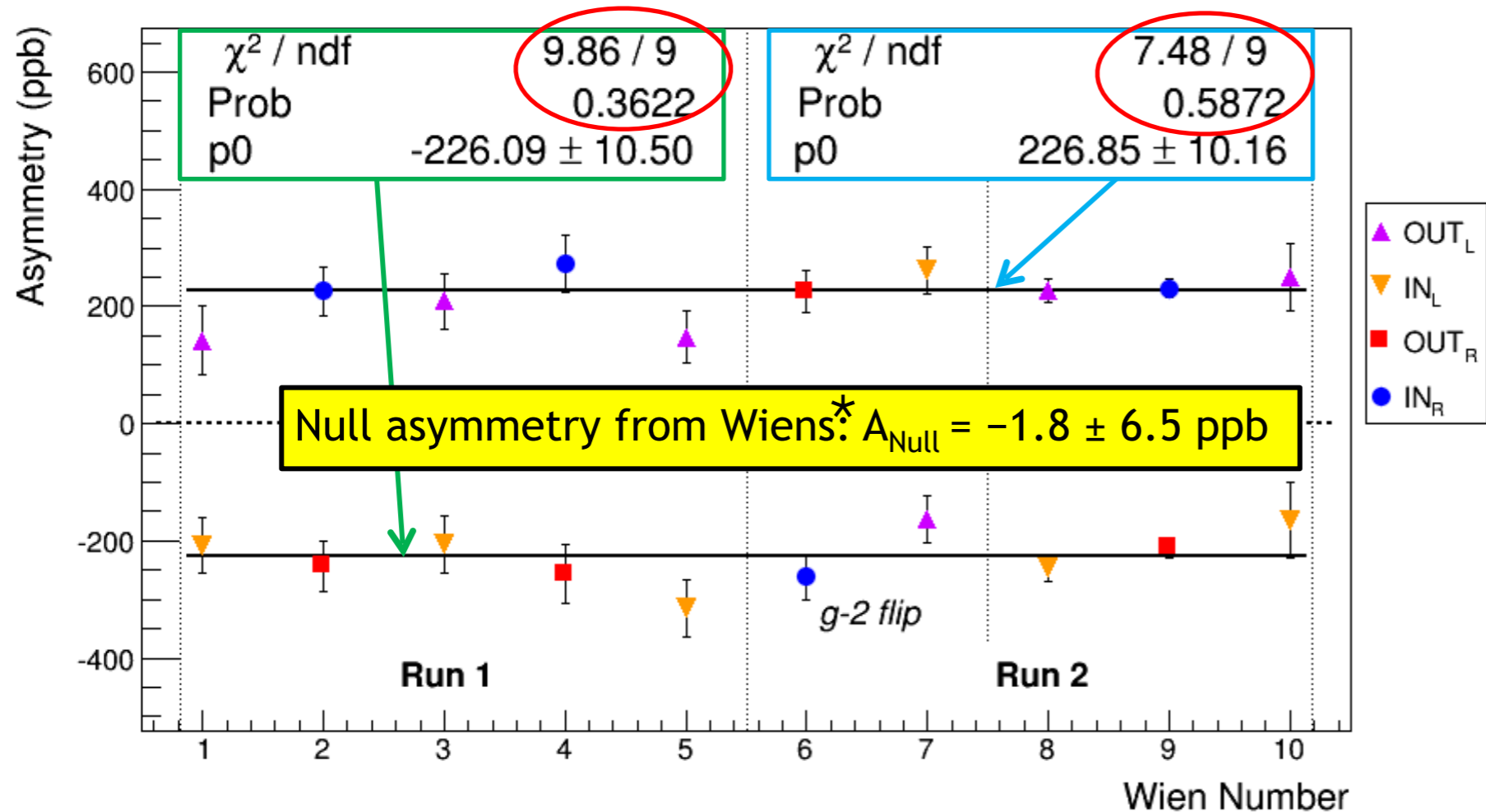
Acceptance defining
 Pb collimator
 $5.8^\circ \leq \theta \leq 11.6^\circ$

integrate for ~1msec



calculate asymmetry from 4 windows

NULL asymmetry



- insertable laser optics flip the polarization of the laser and electron beam
- magnetic spin manipulation (Wein filter) allowed for direct electron beam polarization flip
- Energy variation through the accelerator ($g-2$ flip)

*1 wien ~ 2-6 weeks of data

- The physics asymmetry was consistent through these slow spin reversals

Corrections and systematics

$$A_{\text{msr}} = A_{\text{raw}} + A_T + A_L + A_{\text{BCM}} + A_{\text{BB}} + A_{\text{beam}} + A_{\text{bias}}$$

$$A_{ep} = R_{\text{tot}} \frac{A_{\text{msr}}/P - \sum_{i=1,3,4} f_i A_i}{1 - \sum_{i=1}^4 f_i}$$

- Correct the raw asymmetry for measured experimental factors (detector non-linearity, beam asymmetries)
- Apply additional corrections (polarization, acceptance, backgrounds)

Corrections and systematics

$$A_{\text{msr}} = A_{\text{raw}} + A_T + A_L + A_{\text{BCM}} + A_{\text{BB}} + A_{\text{beam}} + A_{\text{bias}}$$

$$A_{ep} = R_{\text{tot}} \frac{A_{\text{msr}}/P - \sum_{i=1,3,4} f_i A_i}{1 - \sum_{i=1}^4 f_i}$$

- The polarization measurement used both Compton and MOLLER measurements to reach σP of 1.05/0.73% (Run1/Run2)
- Target Al windows caused the largest correction (~ 38 ppb)
- Using vertical drift chambers we benchmarked detailed Q^2 simulations with data

Quantity	Run 2 error (ppb)	Run 2 fractional
BCM Normalization: A_{BCM}	2.3	17%
Beamline Background: A_{BB}	1.2	5%
Beam Asymmetries: A_{beam}	1.2	5%
Rescattering bias: A_{bias}	3.4	37%
Beam Polarization: P	1.2	4%
Target windows: A_{b1}	1.9	12%
Kinematics: R_{Q^2}	1.3	5%
Total of others	2.2	15%
Combined in quadrature	5.6	

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- Cross checks between redundant sets of high precision beam monitoring systems allowed us to obtain a small uncertainty for our determination of beam properties (charge, position)
- Using these monitors we could determine remaining asymmetry in beam properties under the fast helicity flip

Quantity	Run 2 error (ppb)	Run 2 fractional
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$$A_{ep} = R_{\text{tot}} \frac{A_{\text{msr}}/P - \sum_{i=1,3,4} f_i A_i}{1 - \sum_{i=1}^4 f_i}$$

- In the early stages of the experiment we observed a large amount of background coming from beam line elements
 - We introduced a tungsten plug to reduce this background (not fully taken care of)
- A small contribution to the total signal (0.19%) remained after the introduction of the W-plug
- Further studies were done to determine and correct this background contribution
 - background detectors were used to account for possible helicity correlated asymmetries from this background

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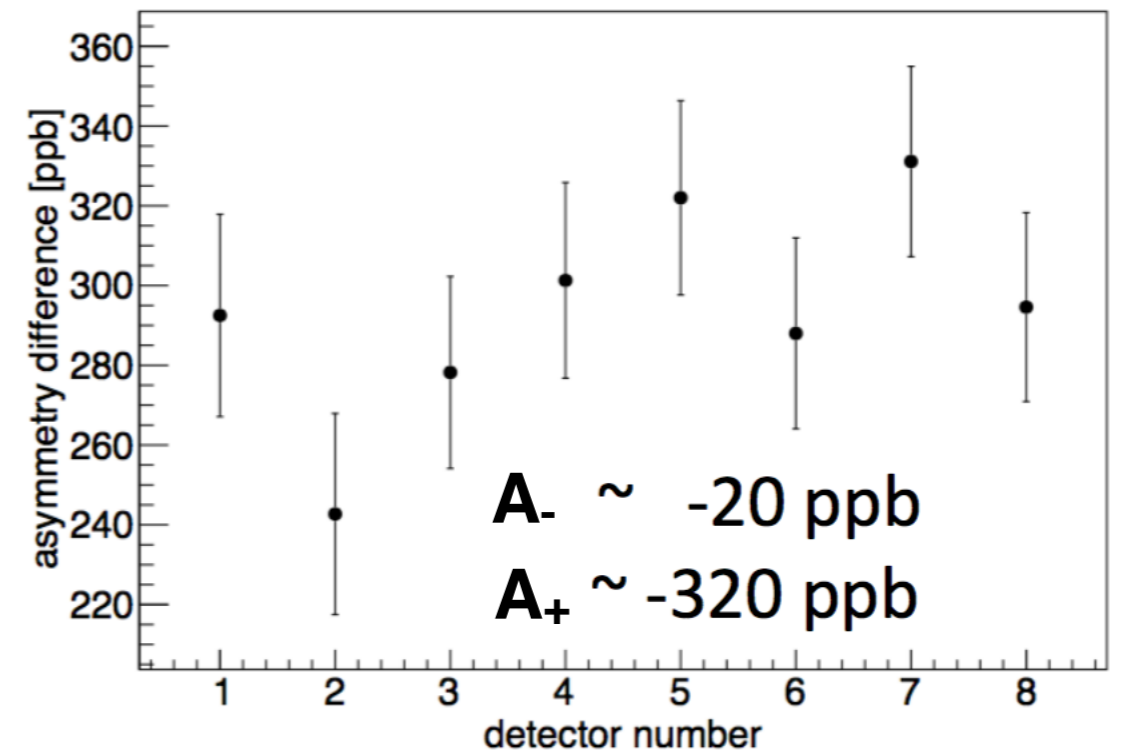
Corrections and systematics

$$A_{\text{msr}} = A_{\text{raw}} + A_T + A_L + A_{\text{BCM}} + A_{\text{BB}} + A_{\text{beam}} + A_{\text{bias}}$$

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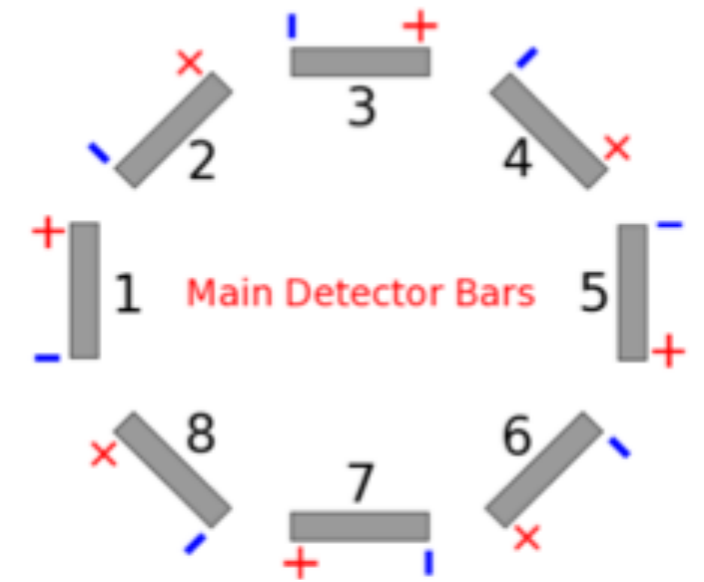
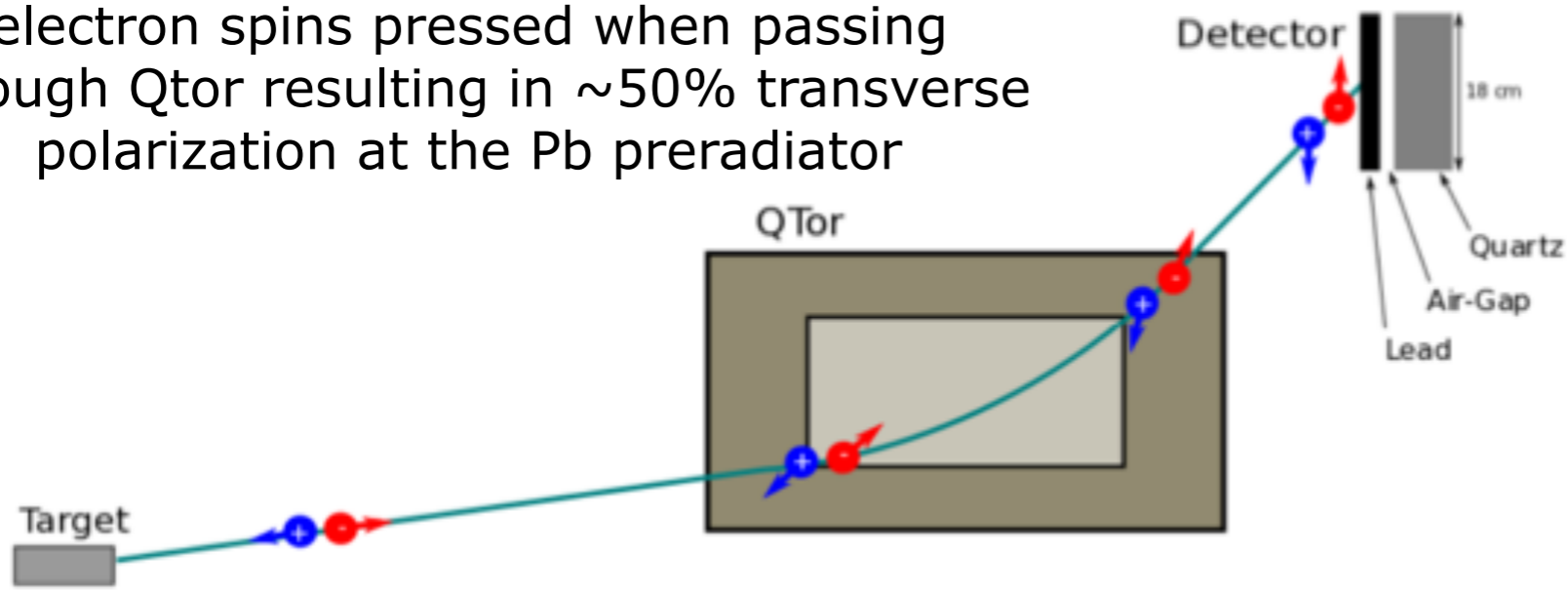
- During the analysis we observed a systematic difference between the asymmetries measured with the positive and negative PMTs

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Polarized scattering in preradiator

electron spins pressed when passing through Qtor resulting in $\sim 50\%$ transverse polarization at the Pb preradiator



produced significant asymmetry difference between + and - PMTs

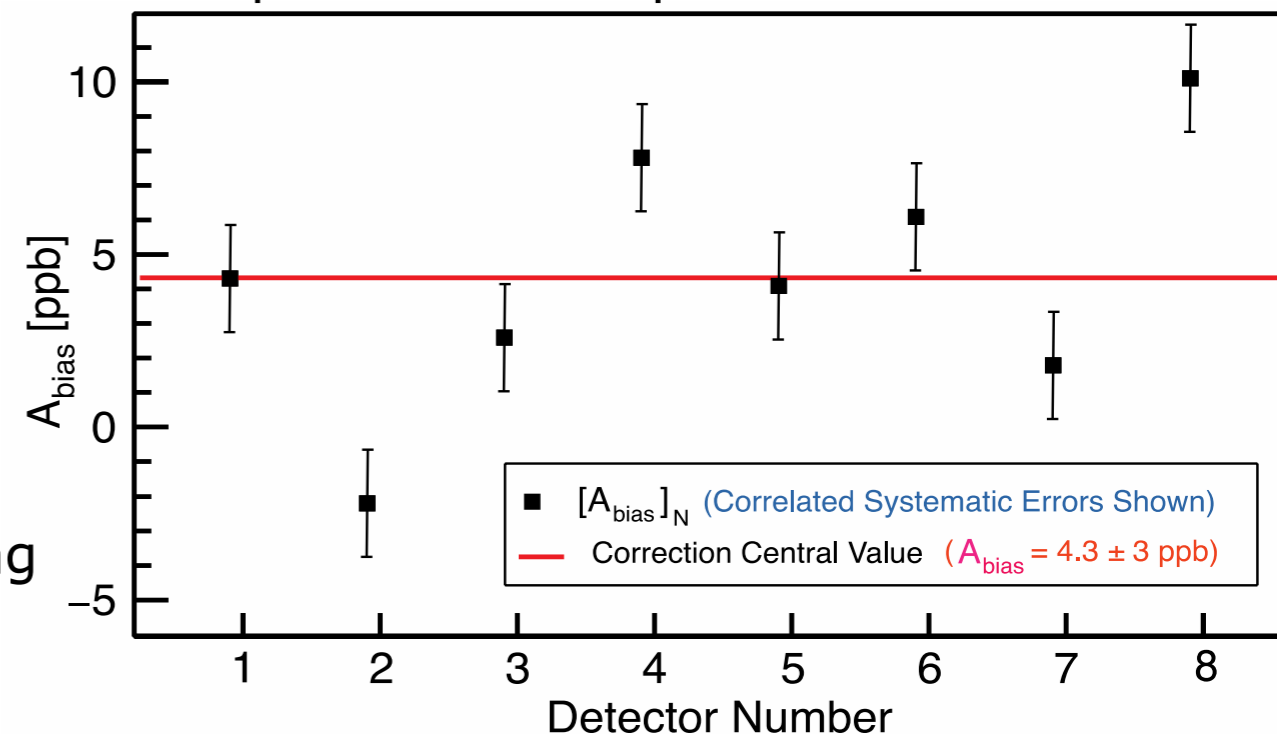
$$A_{PMTDD} = A_- - A_+$$

this effect cancels at first order for our measurement

$$A_{PV} = (A_- + A_+)/2$$

Effect was evaluated using several MCs implementing low energy (few MeV) Mott Scattering and optical properties of the Quartz bars.

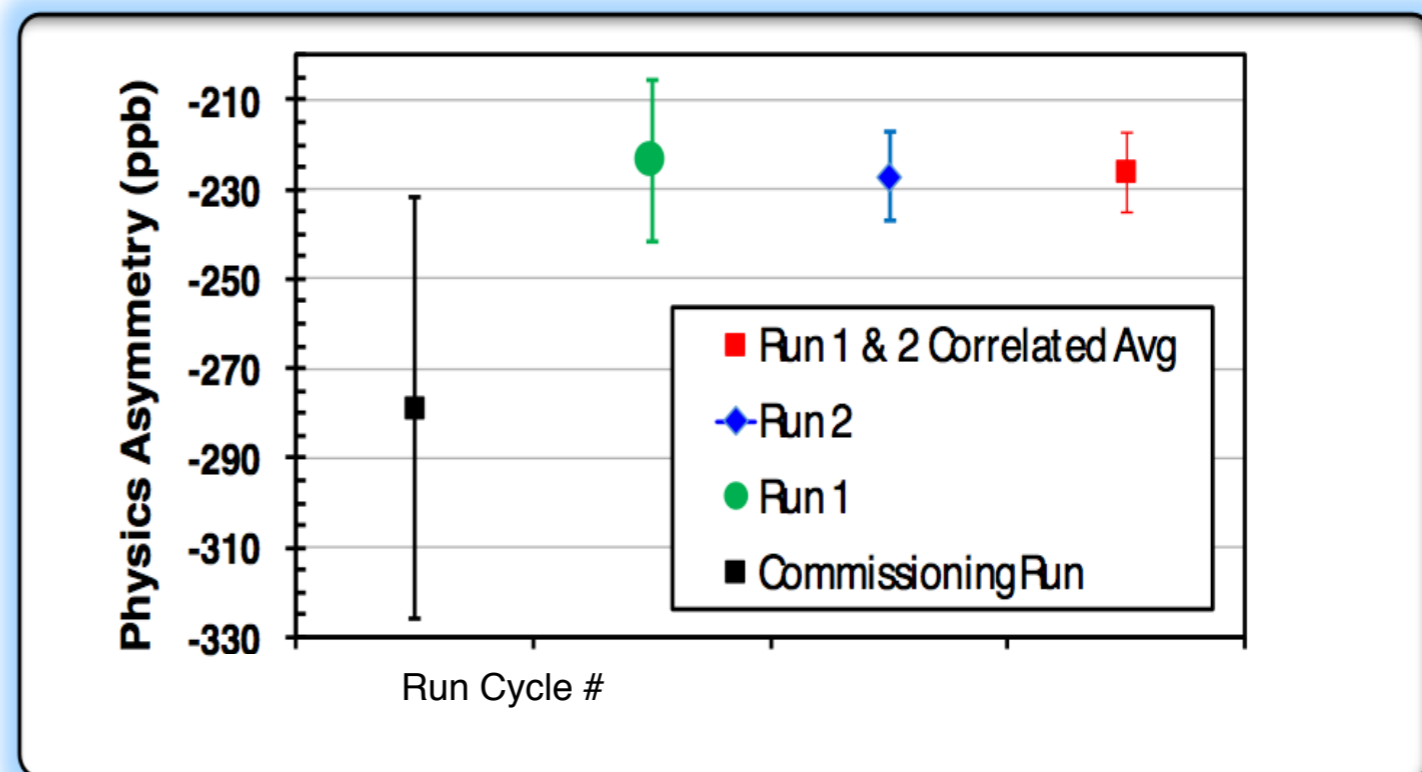
Corrected asymmetries were due to imperfections in quartz bar construction



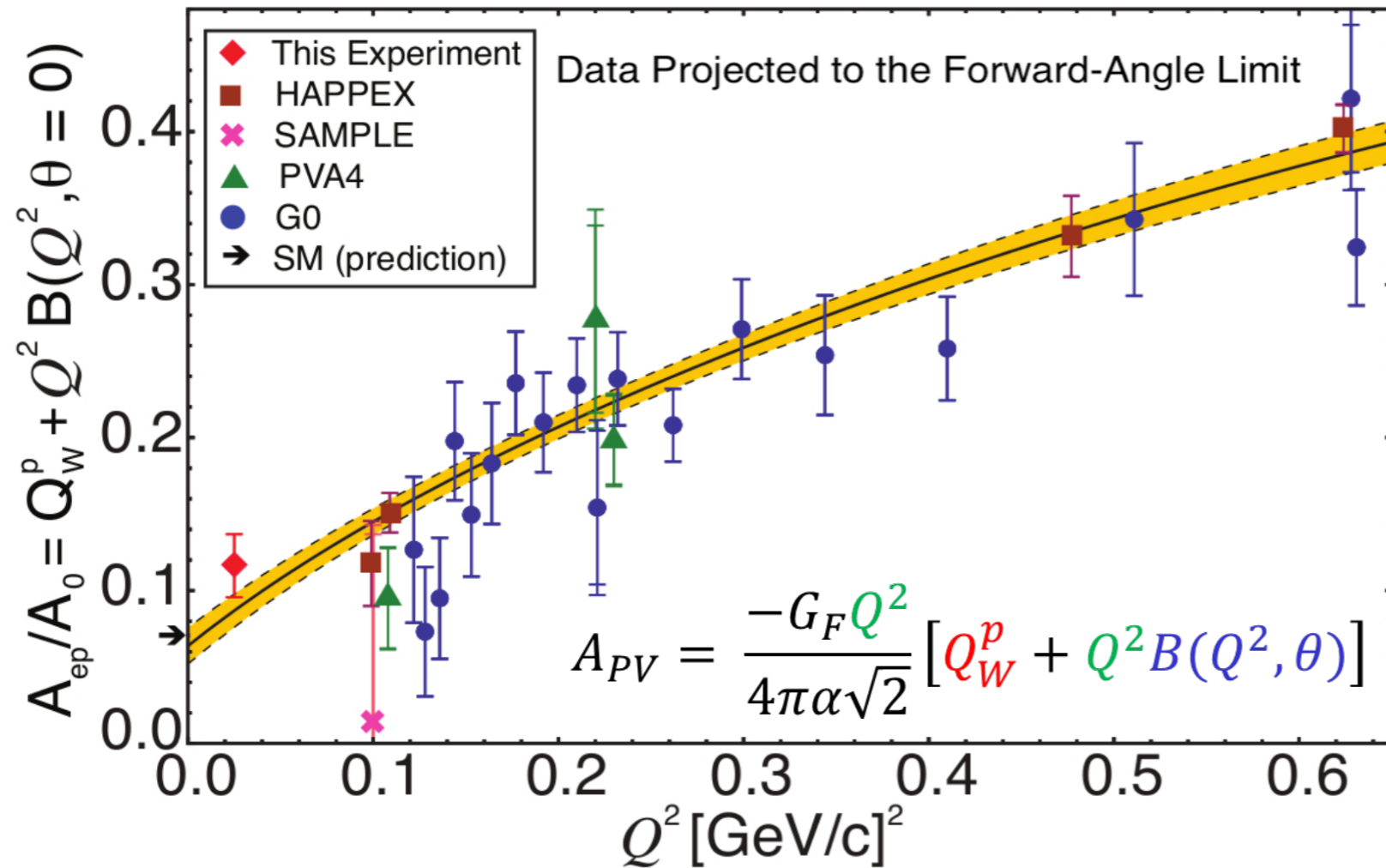
Corrections and systematics

Period	Asymmetry (ppb)	Stat. Unc. (ppb)	Syst. Unc. (ppb)	Tot. Uncertainty (ppb)
Run 1	-223.5	15.0	10.1	18.0
Run 2	-227.2	8.3	5.6	10.0
Run 1 and 2 combined with correlations	-226.5	7.3	5.8	9.3

- Our final result is still statistically limited and dominated by the results from our second running period

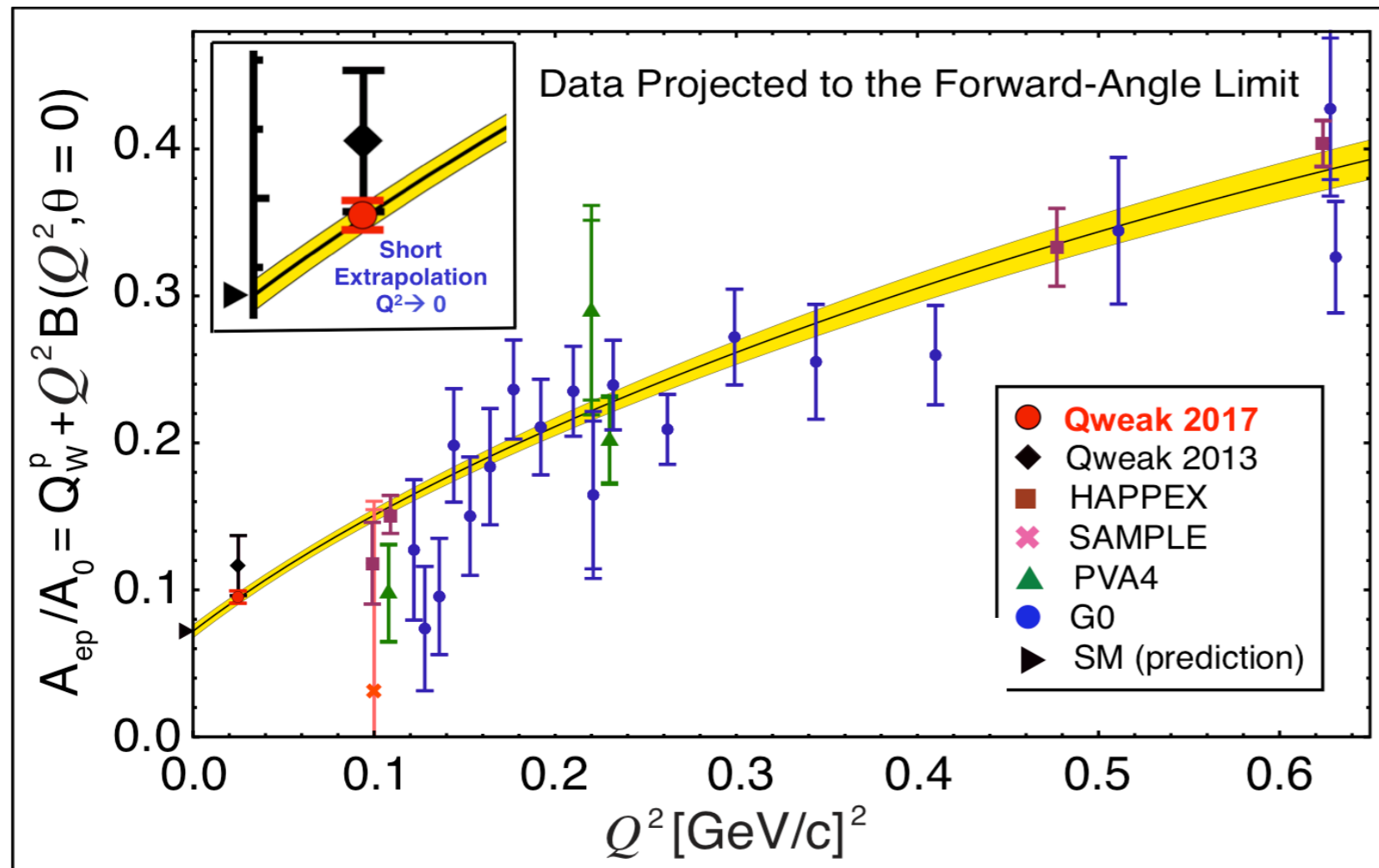


Determining Q_W^p



- Global fit:
 - 5 parameters: C_{1u} , C_{1d} , ρ_s , μ_s and G_A^Z
- Using all PVES data up to $Q^2 = 0.63 \text{ GeV}^2$

Determining Q_w^p



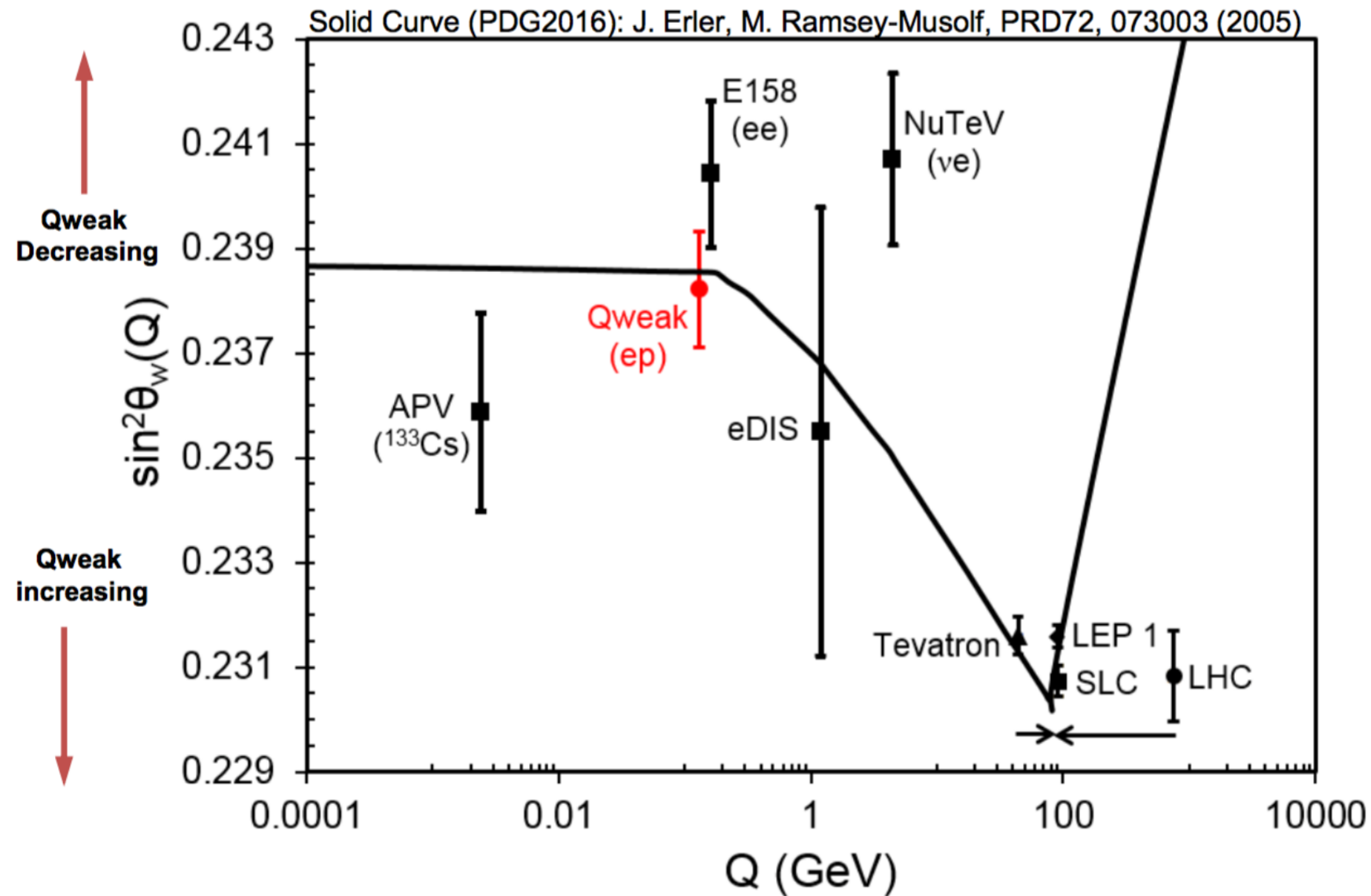
- Global fit:
 - 5 parameters: C_{1u} , C_{1d} , ρ_s , μ_s and G_A^Z
- Using all PVES data up to $Q^2 = 0.63 \text{ GeV}^2$

- Q_w^p is determined from a global fit of PVES data
 - our measurement is the closest to the extrapolation point and anchors the fit while the rest of the data determine the hadronic contributions

$$Q_w^p(SM) = 0.0708 \pm 0.0003$$

$$Q_w^p(PVES) = 0.0719 \pm 0.0045$$

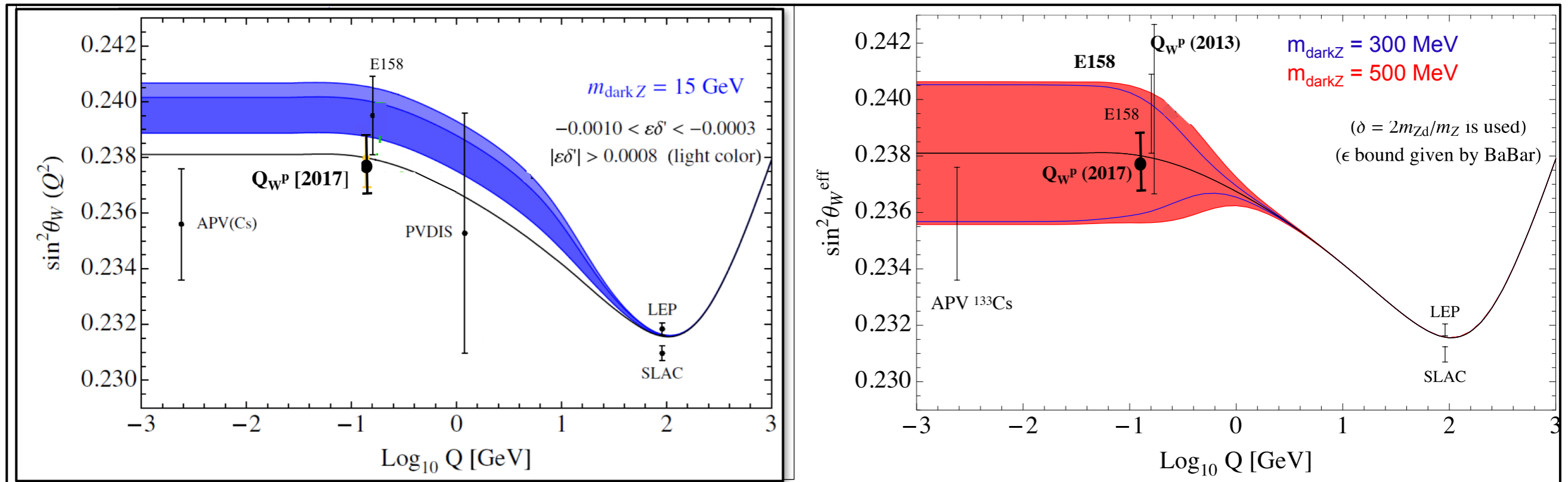
Weak mixing angle



- E158 and Qweak are sensitive to different types of new physics
 - strong consistency with SM for Qweak should put a stronger limit on scalar lepto-quarks (E158 insensitive)

Dark photon / "Z"

(Davoudiasl, Lee, Marciano, Phys. Rev. D89, 095006 (2014), & Marciano (private communication))



- Some models point to a heavy dark photon (Z_d) could be detected at low Q through it's mixing with between Z_0 and Z_d
- Complementary with direct searches
 - in this scenario the Z_d would not have any coupling with SM

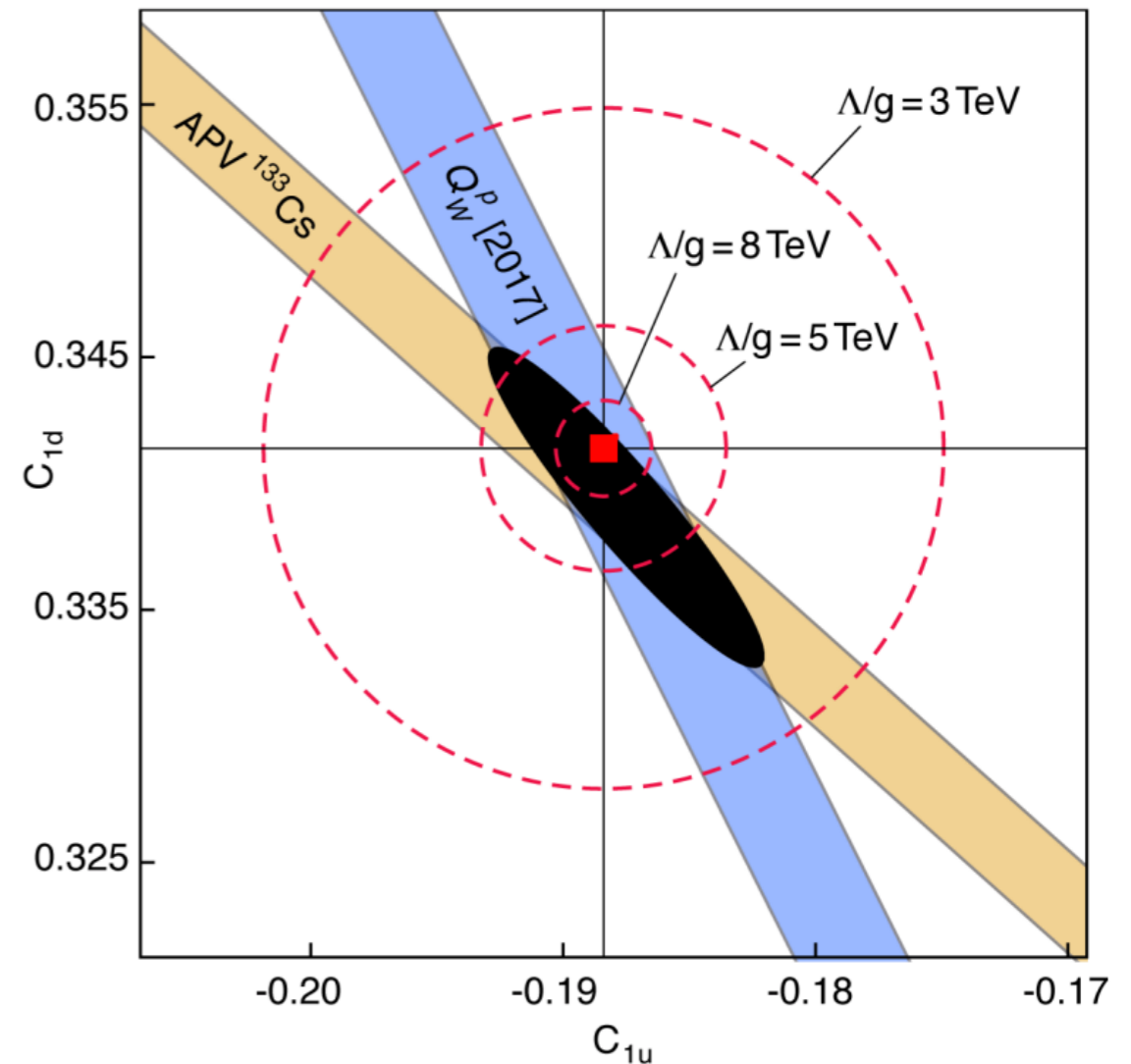
Semi-leptonic PV Physics

Following prescription for new physics from contact interactions in **PhysRevD.68.016006**:

$$\mathcal{L} = \mathcal{L}_{\text{SM}}^{\text{PV}} + \mathcal{L}_{\text{NEW}}^{\text{PV}}$$

$$\mathcal{L}_{\text{SM}}^{\text{PV}} = -\frac{G_{F-}}{\sqrt{2}} \bar{e} \gamma_{\mu} \gamma_5 e \sum_q C_{1q} \bar{q} \gamma^{\mu} q,$$

$$\mathcal{L}_{\text{NEW}}^{\text{PV}} = \frac{g^2}{4\Lambda^2} \bar{e} \gamma_{\mu} \gamma_5 e \sum_f h_V^q \bar{q} \gamma^{\mu} q,$$

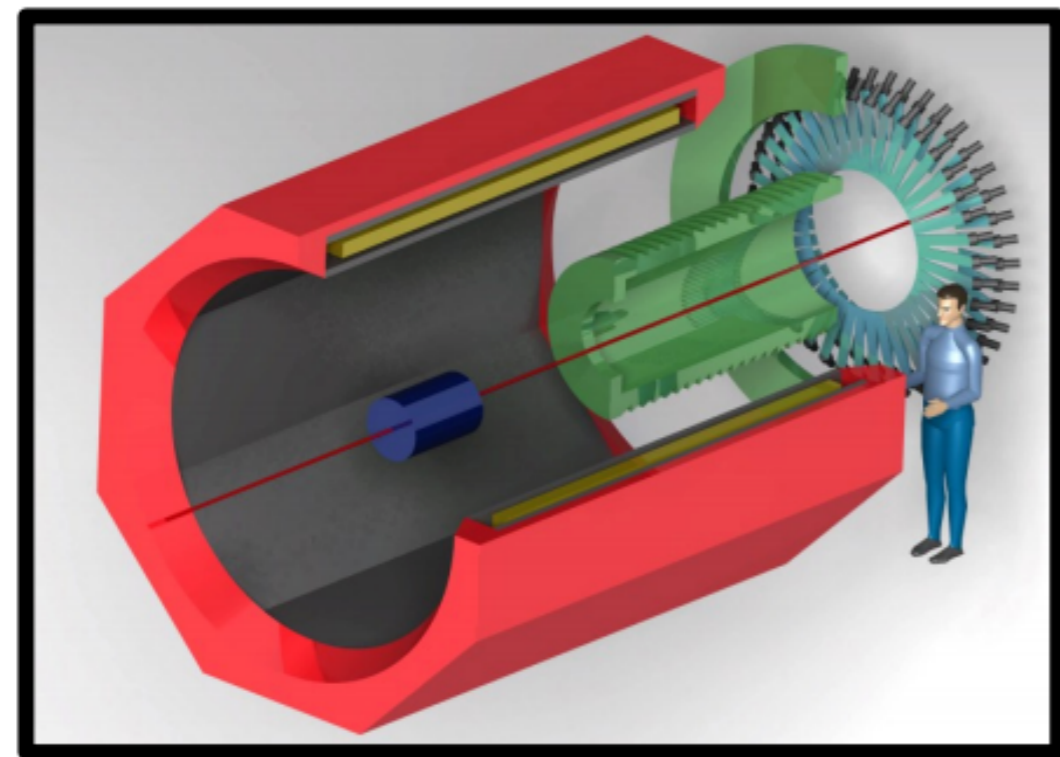


- Our measurement of vector-proton weak neutral current charge constraints new physics which are comparable with current limits from LHC

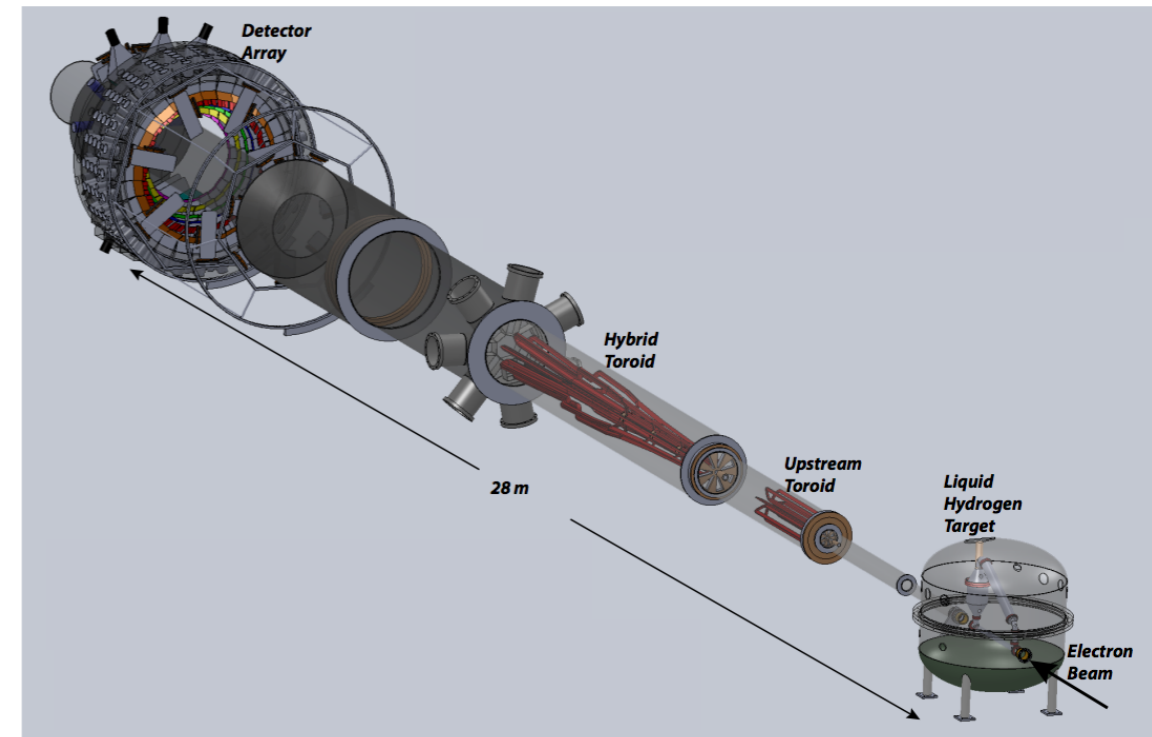
$$g^2 = 4\pi \text{ limit is } 26.3 \text{ TeV}$$

Future PVES experiments

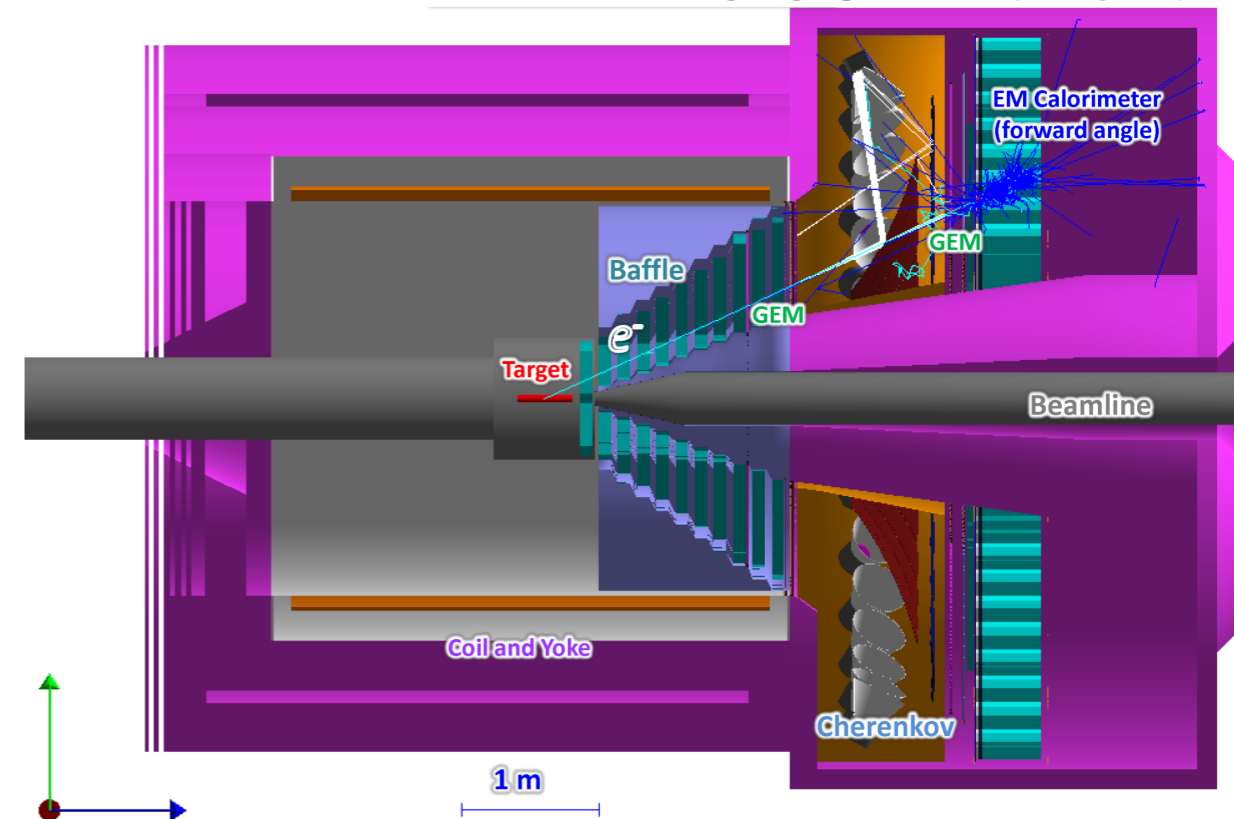
MESA/P2 at Mainz



MOLLER at JLab



PVDIS SOLID at JLab

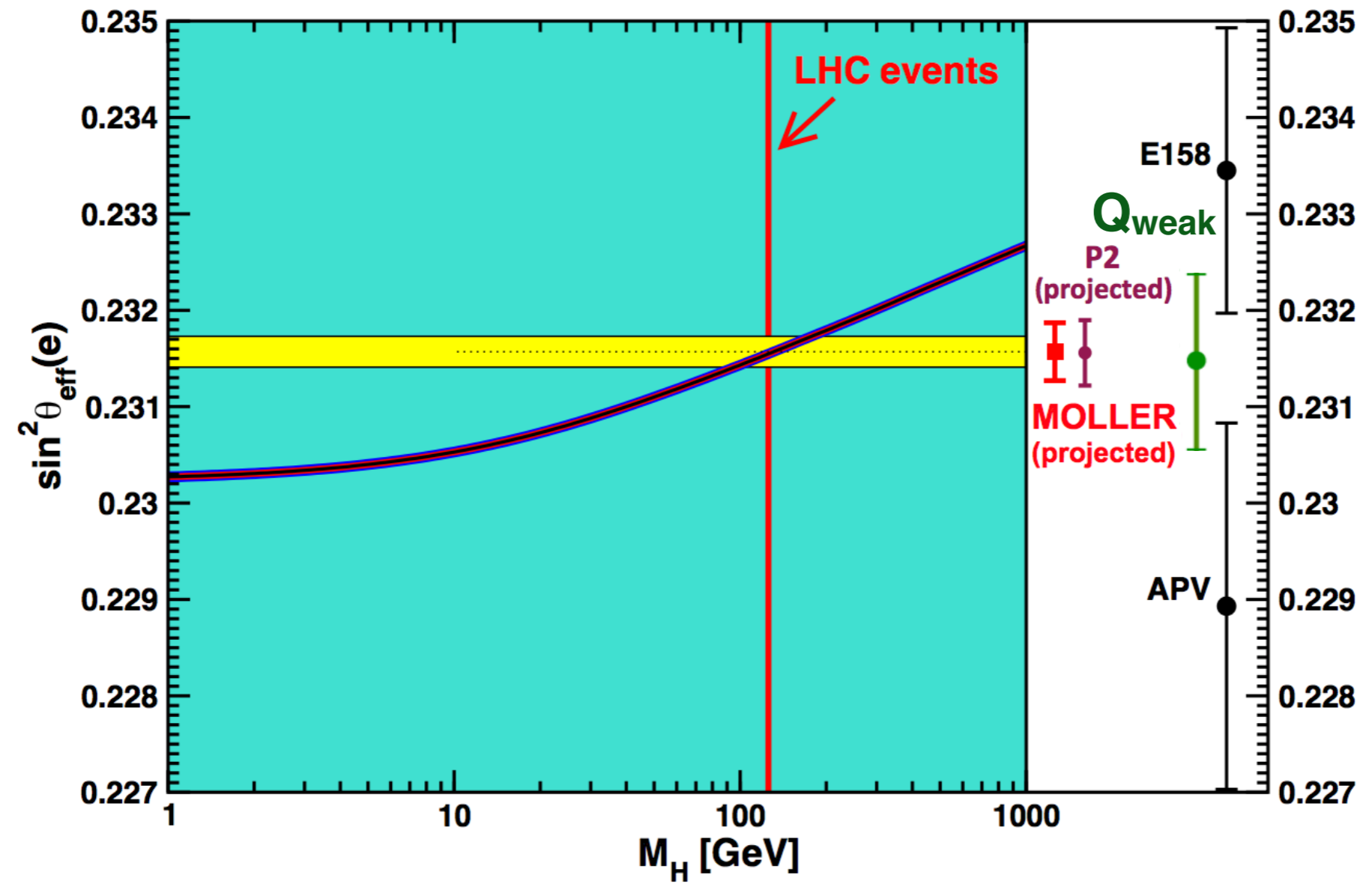
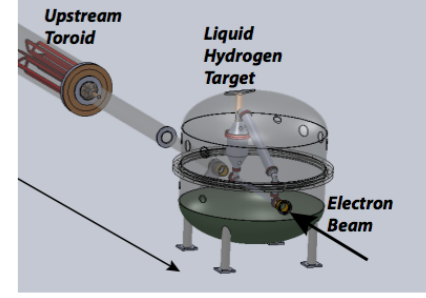
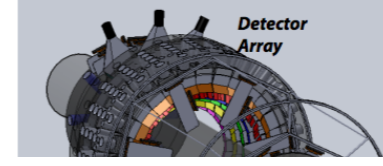
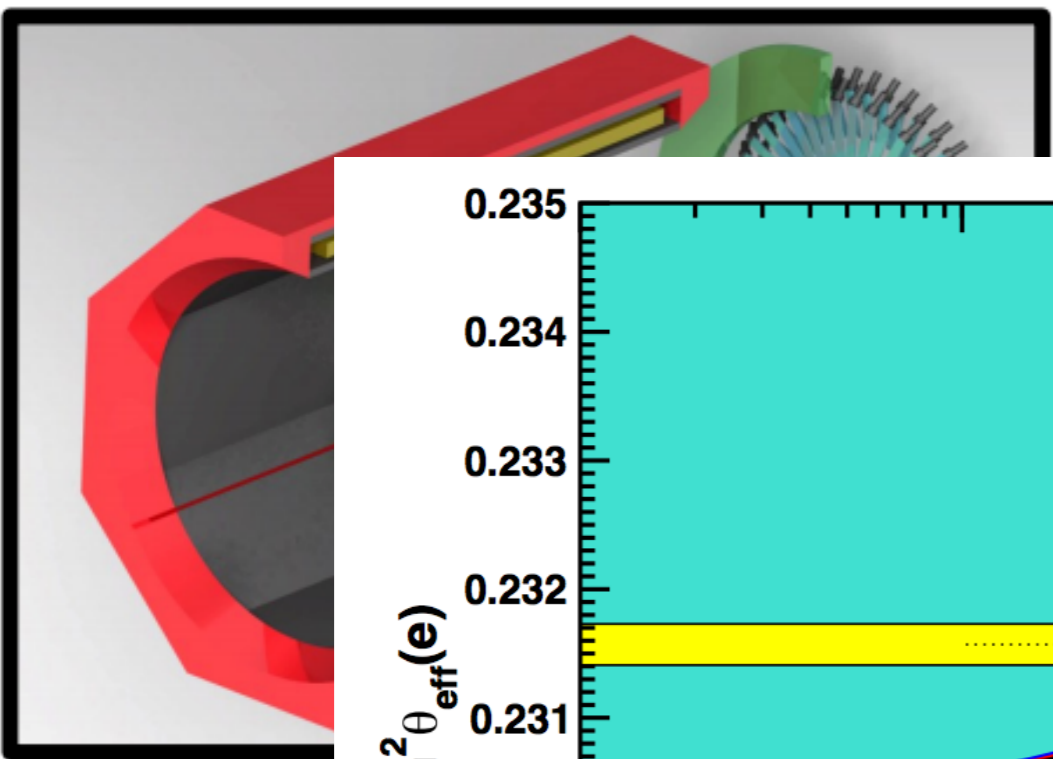


- P2 will improve on the weak charge of the proton
- MOLLER will improve on the weak charge of the electron
- SOLID will make measurement over a wide kinematic range that will include test SM prediction for the axial proton weak charge

Future PVES experiments

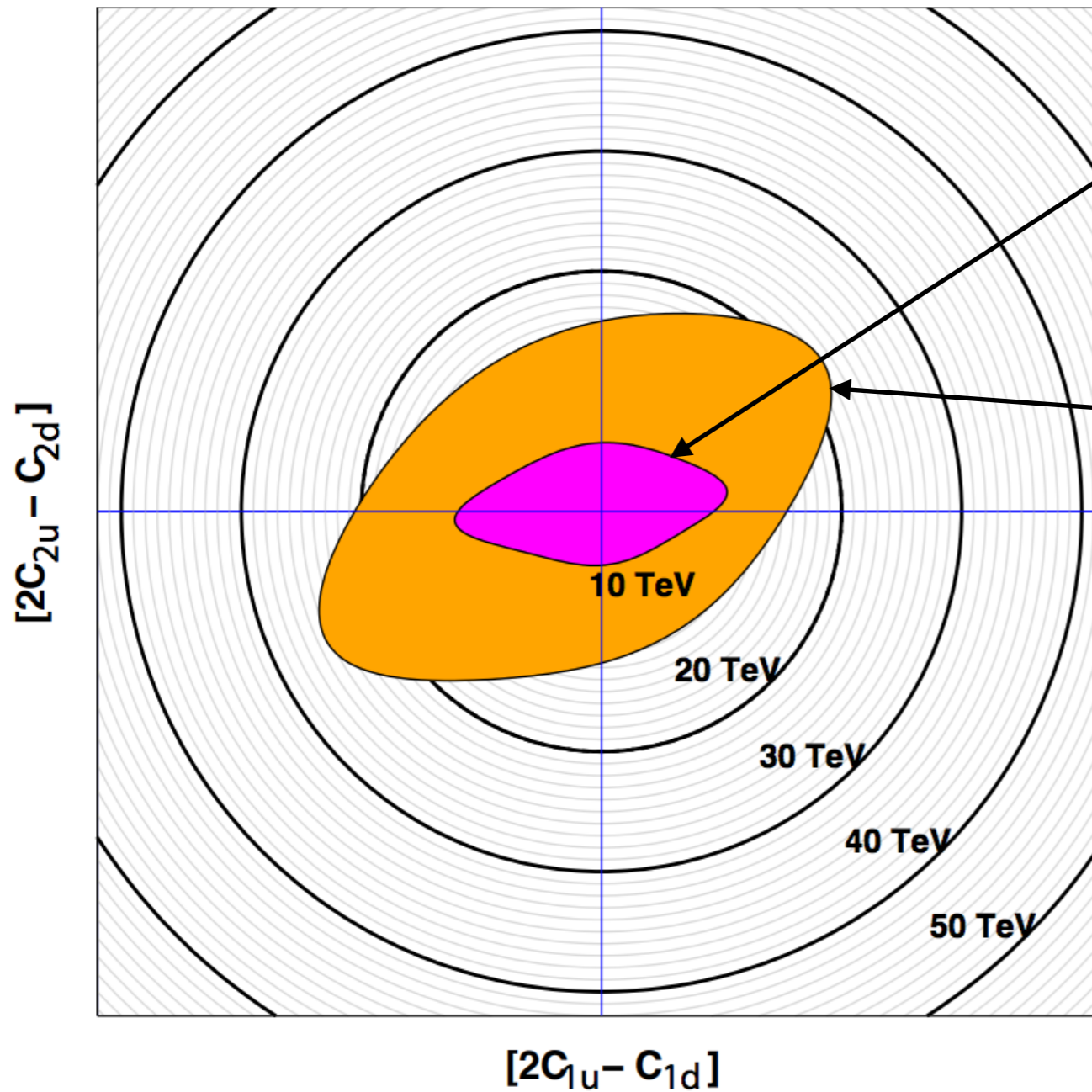
MESA/P2 at Mainz

MOLLER at JLab



- P2 and MOLLER will have weak mixing angle determinations as precise as Z-pole measurement but at very low Q^2

Future PVES experiments

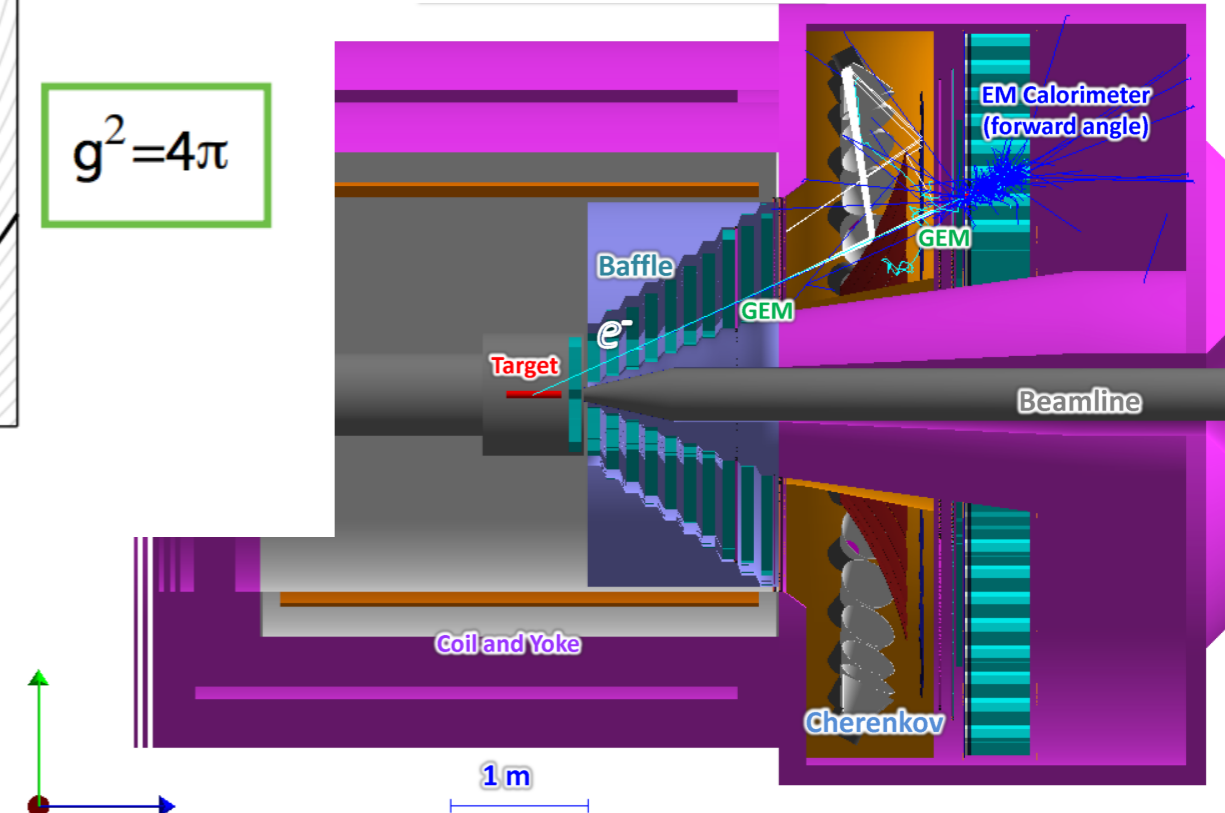


Currently published constraints
(6GeV PVDIS, Qweak, APV)

Constraints after SOLID comparable
with future LHC running.

PVDIS SOLID at JLab

$$g^2 = 4\pi$$



Summary

- Qweak has obtained the most precise determination of a PVES asymmetry and extracted the vector weak charge of the proton Q_w^p
- Sensitive measurement of weak mixing angle limits the phase space of dark Zs
- Consistency with SM further constrains semi-leptonic PV physics above 26.3 TeV
- Future PVES experiments will bring us even closer to the answers we all seek

Backup

Qweak ~ 2.5 kW LH₂ Target

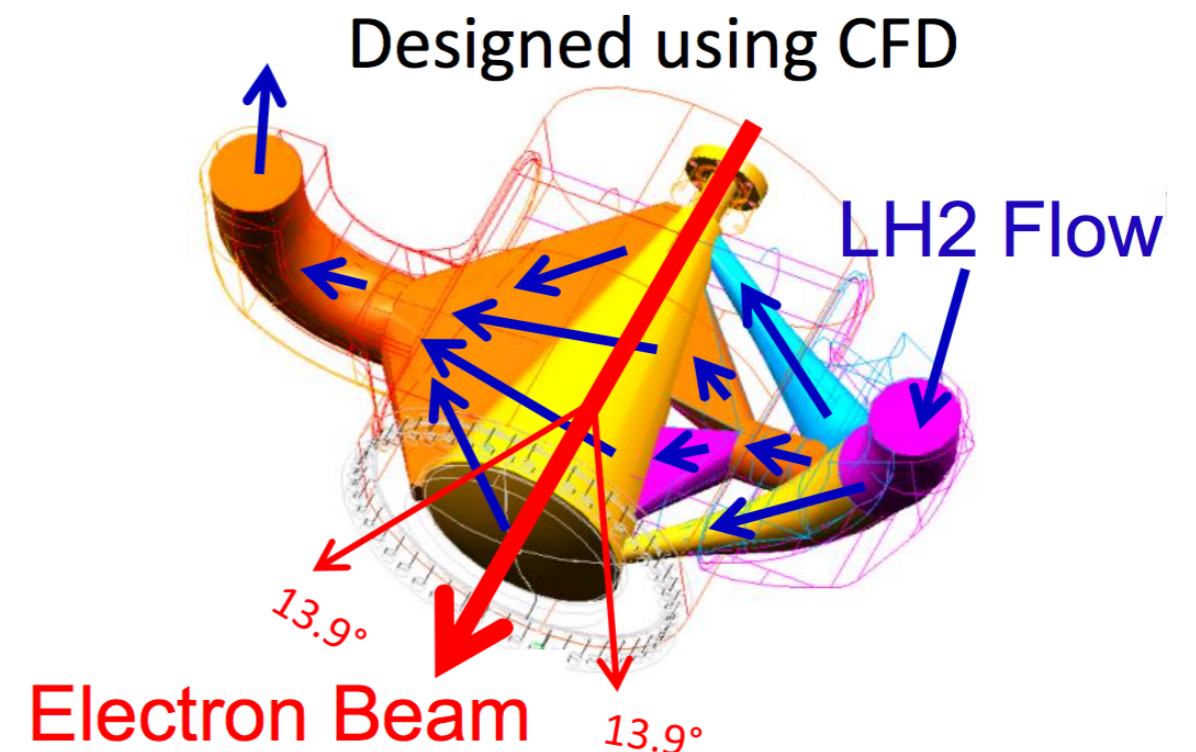
The highest power cryo-target ever built!
35 cm long liquid hydrogen (LH₂)



Target density fluctuations must be small compared to statistical uncertainty

This was achieved by:

- First use of fluid *dynamics simulation* in design to minimize “density changes”, in liquid or at windows.
- *Fast helicity reversal* – up to ~ 1 ms flip rate allows common mode rejection “boiling” noise, line noise and undesired helicity correlated beam properties.
- *Additional safeguards*: large raster size $\sim (3\text{mm} \times 3\text{mm})$, faster pump speed, and more cooling directed onto windows....



*courtesy of R. Carlini

Corrections and systematics

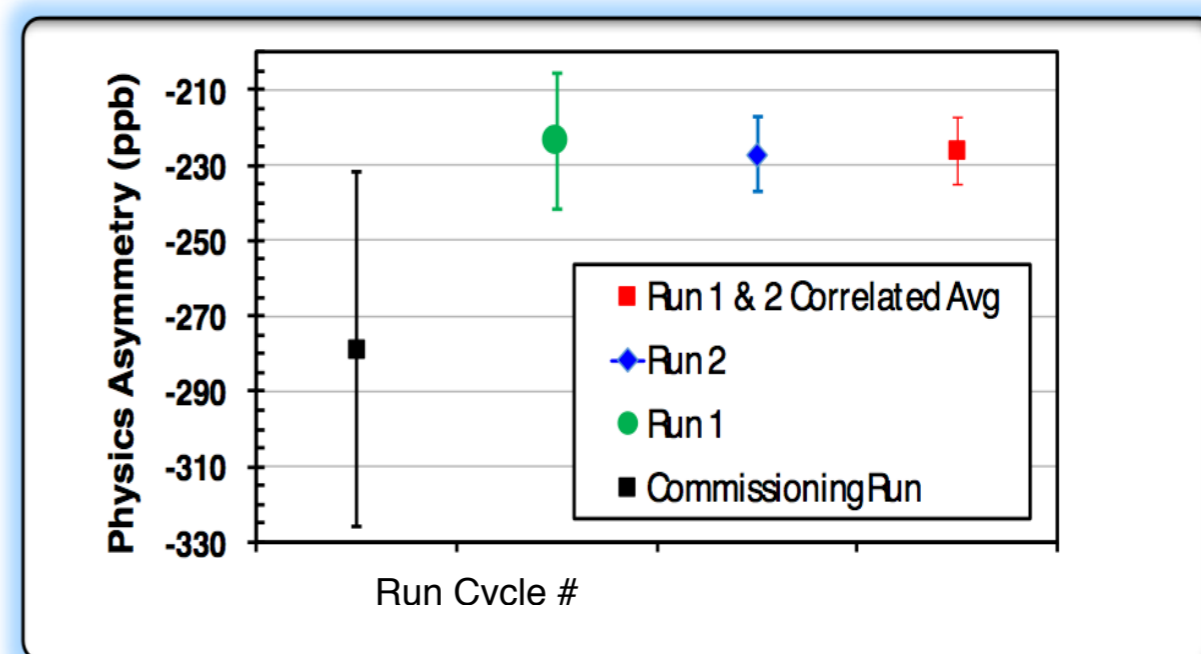
$$A_{\text{msr}} = A_{\text{raw}} + A_T + A_L + A_{\text{BCM}} + A_{\text{BB}} + A_{\text{beam}} + A_{\text{bias}}$$

$$A_{ep} = R_{\text{tot}} \frac{A_{\text{msr}}/P - \sum_{i=1,3,4} f_i A_i}{1 - \sum_{i=1}^4 f_i}$$

Quantity	Run 1 error (ppb)	Run 1 fractional	Run 2 error (ppb)	Run 2 fractional
BCM Normalization: A_{BCM}	5.1	25%	2.3	17%
Beamline Background: A_{BB}	5.1	25%	1.2	5%
Beam Asymmetries: A_{beam}	4.7	22%	1.2	5%
Rescattering bias: A_{bias}	3.4	11%	3.4	37%
Beam Polarization: P	2.2	5%	1.2	4%
Target windows: A_{b1}	1.9	4%	1.9	12%
Kinematics: R_{Q^2}	1.2	2%	1.3	5%
Total of others	2.5	6%	2.2	15%
Combined in quadrature	10.1		5.6	

- Our final result is still statistically limited

Period	Asymmetry (ppb)	Stat. Unc. (ppb)	Syst. Unc. (ppb)	Tot. Uncertainty (ppb)
Run 1	-223.5	15.0	10.1	18.0
Run 2	-227.2	8.3	5.6	10.0
Run 1 and 2 combined with correlations	-226.5	7.3	5.8	9.3



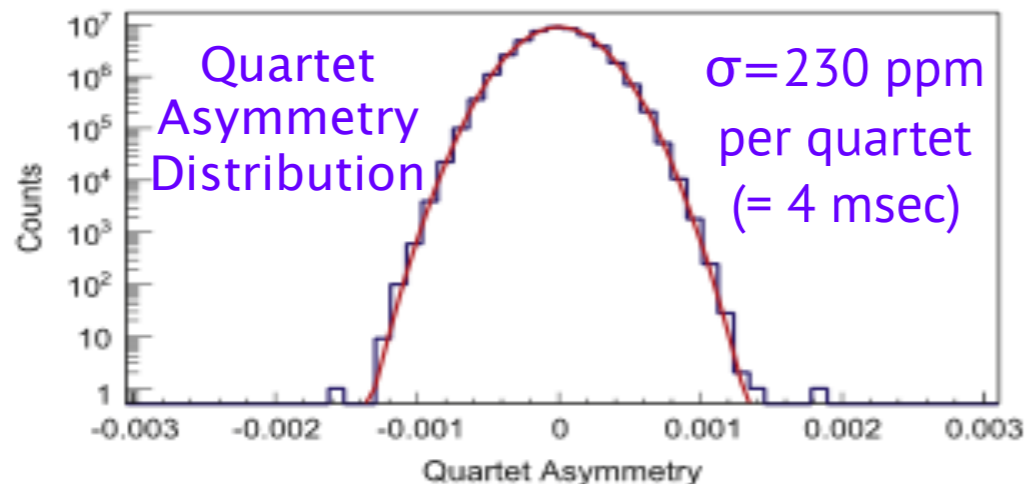
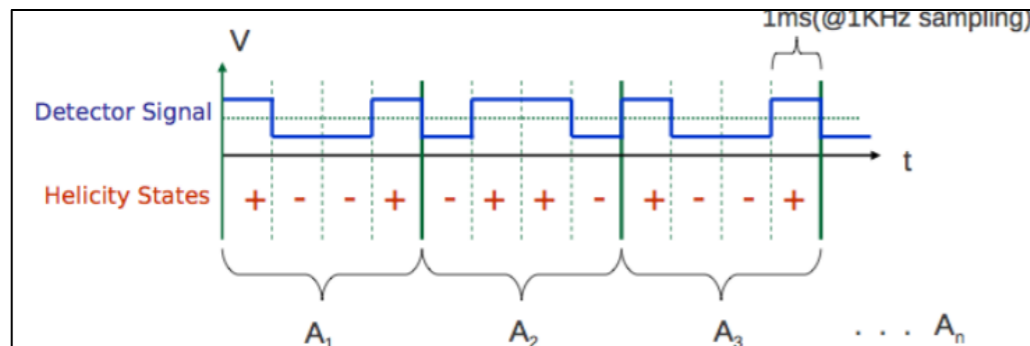
Result stability

Quantity	Value	Error	Method
Q_W^p	0.0719	0.0045	Qweak A_{ep}
ρ_s	0.19	0.11	+
μ_s	-0.18	0.15	PVES data base
$G_A^{Z(T=1)}$	-0.67	0.33	
Q_W^p	0.0718	0.0045	Qweak A_{ep}
Q_W^n	-0.9808	0.0063	+
C_{1u}	-0.1874	0.0022	PVES data base
C_{1d}	0.3389	0.0025	+
C_1 correlation	-0.9317		APV ^{133}Cs
Q_W^p	0.0684	0.0039	Qweak A_{ep}
			+
			PVES data base
			+
			LQCD (strange quarks)
Q_W^p	0.0706	0.0047	Qweak A_{ep}
			+
			EMFF's & theory axial
			+
			LQCD (strange)

Q_W^p (this result)	0.0719 ± 0.0045
Q_W^p (SM)	0.0708 ± 0.0003

- Including the APV data gives discrimination power for C_{1u} and C_{1d} which leads to Q_W^n
- The addition of LQCD data add additional constraints on the Q_W^p determination
- A determination without the rest of the PVES data (but with some ansatz for the form factors) shows the power of the Qweak experiment

PV Measurement



LH2 statistical width (per quartet):

- Counting statistics: 200 ppm
- Main detector resolution: 92 ppm
- Target noise/boiling: 55 ppm
- BCM Resolution: 43 ppm
- Electronic noise: 3 ppm

- Integrate light signal for ~ 1 msec
- Calculate asymmetry for 4 adjacent data samples and add a blinding factor
- Analyzed $\sim 10^9$ quartets which is equivalent to $\sim 10^{17}$ detected electrons

Electroweak radiative corrections

$$Q_W^p = [1 + \Delta\rho + \Delta_e] [(1 - 4\sin^2\theta_W(0)) + \Delta_{e'}] + \square_{WW} + \square_{ZZ} + \square_{\gamma Z}$$

Correction to Q_{Weak}^p	Uncertainty
$\Delta \sin\theta_W (M_Z)$	± 0.0006
$Z\gamma$ box (6.4% \pm 0.6%)	0.00459 ± 0.00044
$\Delta \sin\theta_W (Q)_{hadronic}$	± 0.0003
WW, ZZ box - pQCD	± 0.0001
Charge symmetry	0
Total	± 0.0008

Erlar et al., PRD 68(2003)016006.

Calculations of Two Boson Exchange effects on Q_W^p at our Kinematics:

Recent theory calculations applied to entire data set of PV measurements as appropriate in global analysis.

Our ΔA_{ep} precise enough that corrections to higher Q^2 points make little difference in extrapolation to zero Q^2 .

Energy Dependence γZ correction:

Hall, N.L., Blunden, P.G., Melnitchouk, W., Thomas, A.W., Young, R.D. Quark-hadron duality constraints on γZ box corrections to parity-violating elastic scattering. *Phys. Lett. B* 753, 221-226 (2016).

Axial Vector γZ correction:

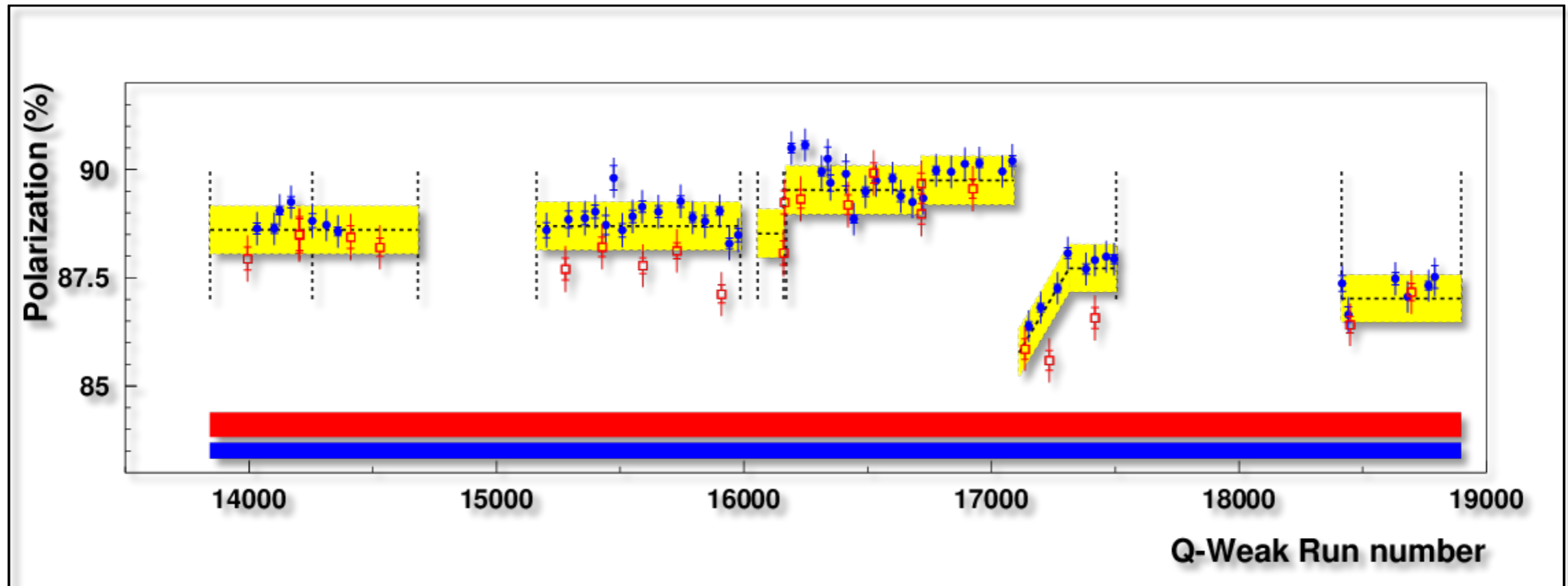
Peter Blunden, P.G., Melnitchouk, W., Thomas, A.W. New Formulation of γZ Box Corrections to the Weak Charge of the Proton. *Phys. Rev. Lett.* 107, 081801 (2011).

Q^2 Dependence γZ :

Gorchtein, M., Horowitz, C.J., Ramsey-Musolf, M.J. Model dependence of the γZ dispersion correction to the parity-violating asymmetry in elastic ep scattering. *Phys. Rev. C* 84, 015502 (2011).

*courtesy of R. Carlini

Polarization measurement



- Inner error bars statistical, outer error bars point-to-point systematic uncertainties added in quadrature with statistical uncertainties.
- Yellow band incorporates overall normalization uncertainties determining by weighted average and total uncertainty.
- Time dependence of reported polarization driven by continuous Compton measurements, with small scale correction (0.21%) determined from uncertainty-weighted global comparison of Compton and Møller polarimeters.

*courtesy of R. Carlini