

High energy scattering in $N=4$ SYM

- Introduction: Motivation
- Multiregge limit
- Results

Introduction: Motivation

Conjecture (Maldacena 1997):

correspondence between quantum field theory and String theory

Favorite example: $N=4$ supersymmetric QFT dual to String theory in AdS₅

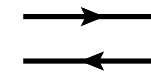
Hope: can be solved exactly (integrable)

Dream: QCD, but dual analogue is not known exactly, needs modelling

$N=4$ SYM : somewhat close to QCD, but has important differences
particle content, scale invariant, no running coupling --> 'Laboratory'

N=4 correspondence:
large N_c

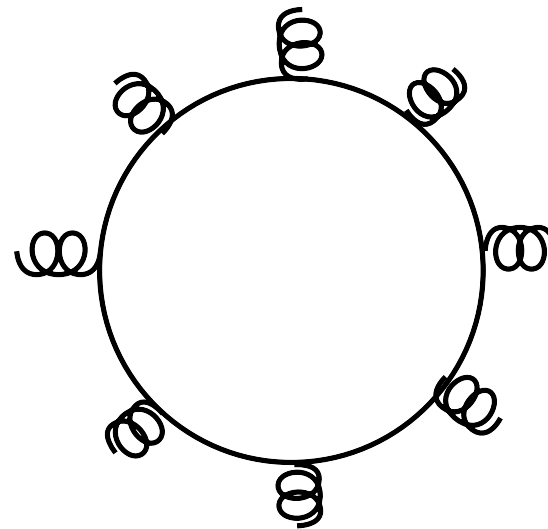
$N=4$ SYM



$AdS_5 \times S^5$

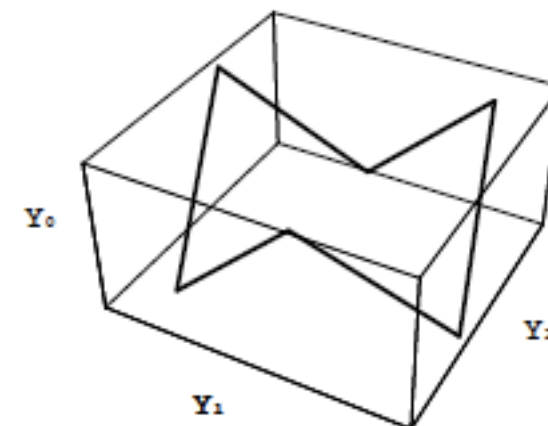
$$g, N_c, \lambda = g^2 N_c$$

$$g^2 N_c = \frac{R^4}{l_s^4}$$



λ small

weak coupling:
perturbation theory



λ large

strong coupling:
semiclassical approximation,
minimal area of polygon

Perturbation theory and beyond:

exact scattering amplitudes in two, three,... loops, strong coupling limit
scattering amplitudes in multiregge limit

BDS conjecture for MHV amplitudes:

Bern, Dixon, Smirnov.

$$\text{tr}(T^{a_1} \dots T^{a_n}) + \text{noncycl.perm}, \quad A_n = A_n^{\text{tree}} \cdot M_n(\epsilon)$$

$$\ln M_n = \sum_l a^l \left[\left(f^{(l)}(\epsilon) I_n(l\epsilon) + F_n(0) \right) + C^{(l)} + E_n^{(l)}[\epsilon] \right]$$

$$a = \frac{N_c \alpha}{2\pi} (4\pi e^{-\gamma})^\epsilon, \quad d = 4 - 2\epsilon$$

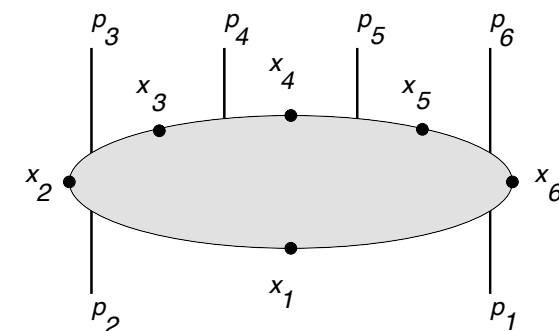
Corrections needed for $n \geq 6$: conformal invariant remainder function $R^{(n)}$

$$M_n \sim \exp(\ln M_n^{\text{BDS}} + R^{(n)})$$

Correspondence: amplitude - Wilson loop (color non-singlet!).

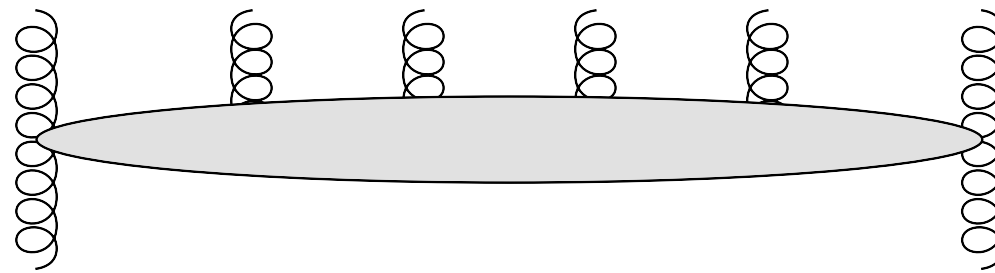
Remainder function depends upon anharmonic ratios, e.g. for $n=6$:

$$u_1 = \frac{x_{13}^2 x_{46}^2}{x_{14}^2 x_{36}^2}, \quad u_2 = \frac{x_{24}^2 x_{15}^2}{x_{25}^2 x_{41}^2}, \quad u_3 = \frac{x_{26}^2 x_{35}^2}{x_{25}^2 x_{36}^2} = \frac{s_{25} s_{36}}{s_{345} s_{456}}$$



Motivation: why multiregge limit?

- Historically, the Regge limit has played an important role
- structural insight (Regge factorization): test of 2,3, ... loop results
- Regge limit is of practical interest: want to know multi-leg partonic amplitudes at high energies:



$2 \rightarrow n$ gluon production

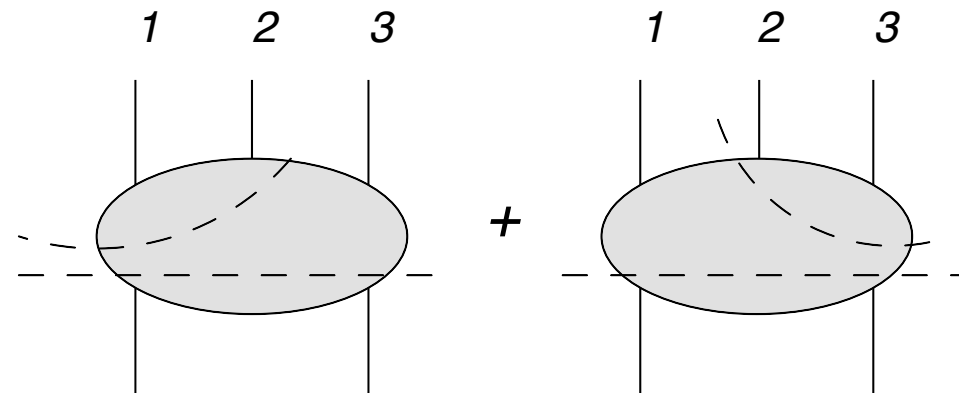
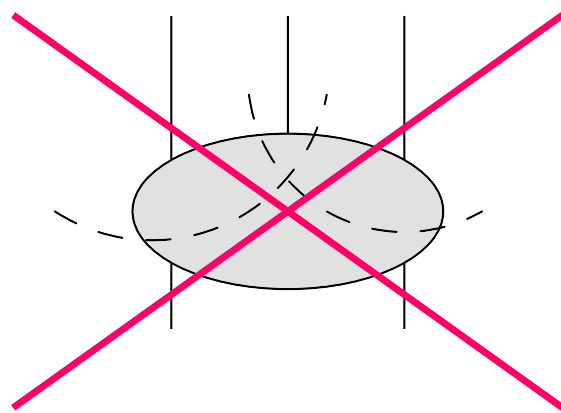
In the following:
compute scattering amplitudes for

$2 \rightarrow 4, 2 \rightarrow 5, 2 \rightarrow 6, \dots$

Multiregge limit

Method: unitarity, analytic properties

Steinmann relations: no simultaneous singularities in overlapping channels



$$(-s_{12})^{\omega_1 - \omega_2} (-s)^{\omega_2} F_L + (-s_{23})^{\omega_2 - \omega_1} (-s)^{\omega_1} F_R$$

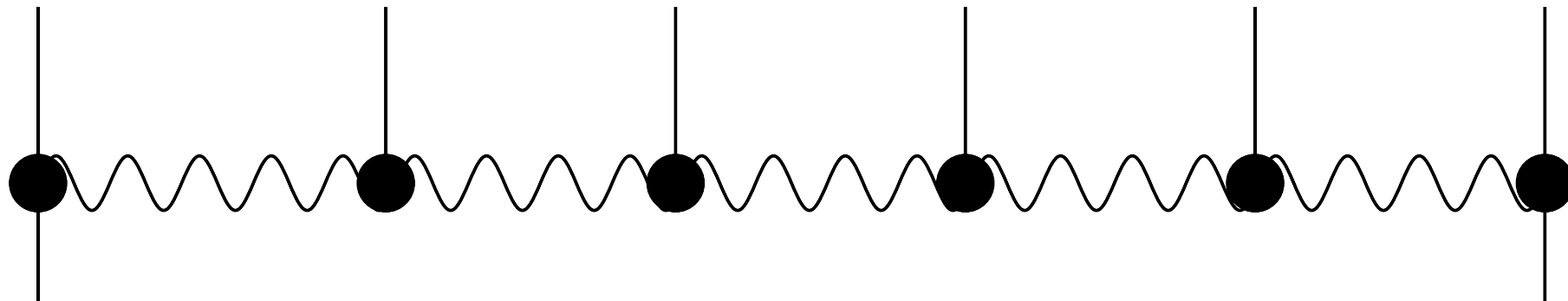
Each term: energy factors times real partial wave.

Allows analytic continuation into different kinematic regions

- 2 \rightarrow 4 five terms
- 2 \rightarrow 5 fourteen terms
- 2 \rightarrow 6 forty two terms
- ... (Catalan numbers)

Results (Summary)

Regge poles: (BDS)



Signatured amplitudes: factorization

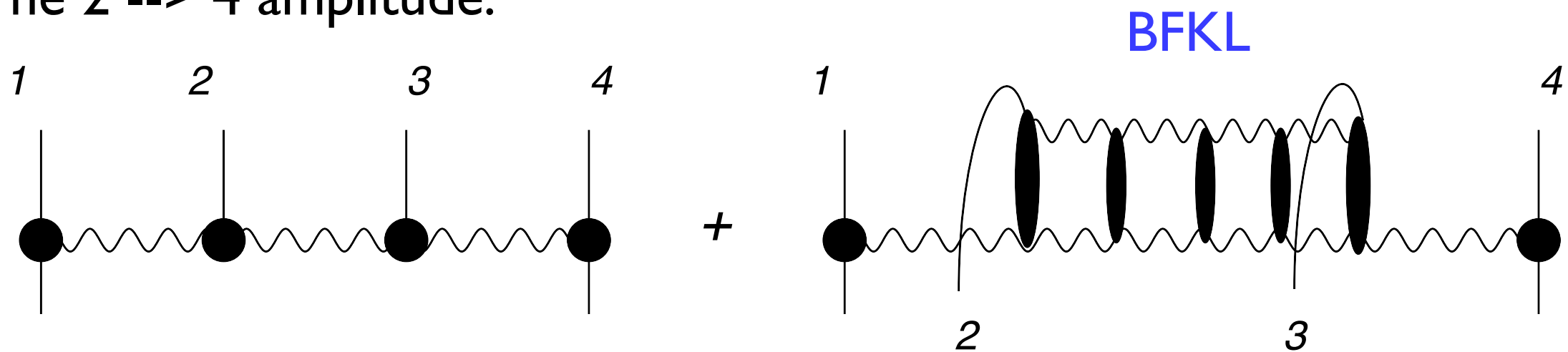
Large N_c (planar approximation):

in some kinematic regions factorization is broken, unphysical singularities appear

BDS: not correct in just these regions

---> need remainder function, correction to BDS formula: Regge cuts

The 2 --> 4 amplitude:



Regge cut cancels the unphysical singularity of the Regge pole.
Appears only in the kinematic region where

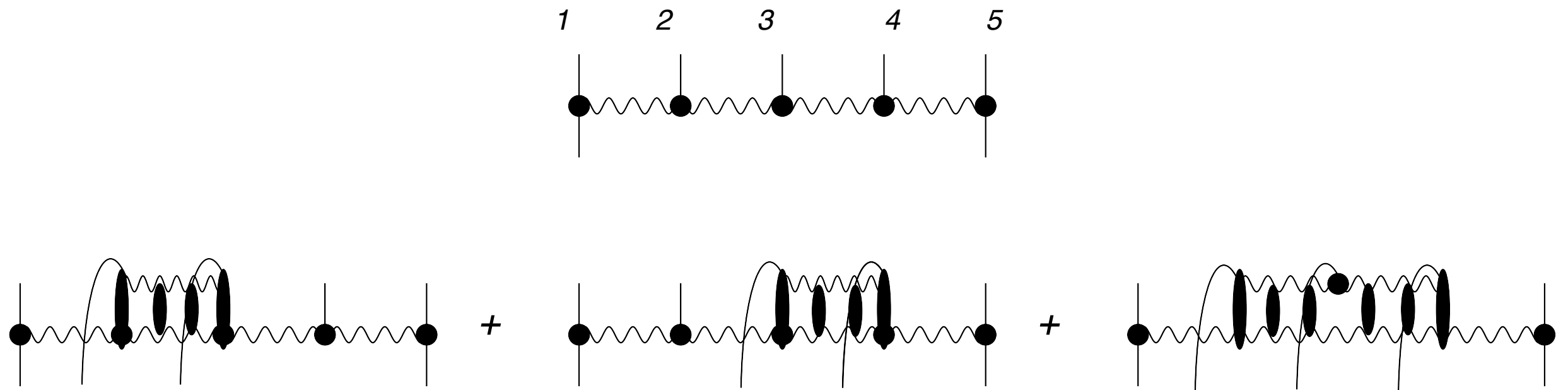
$$s, s_{23} > 0, s_{123} < 0, s_{234} < 0$$

BDS formula contains one loop part of the Regge cut.

In the large N_c (planar) approximation:

Regge cuts are needed for consistency reasons.

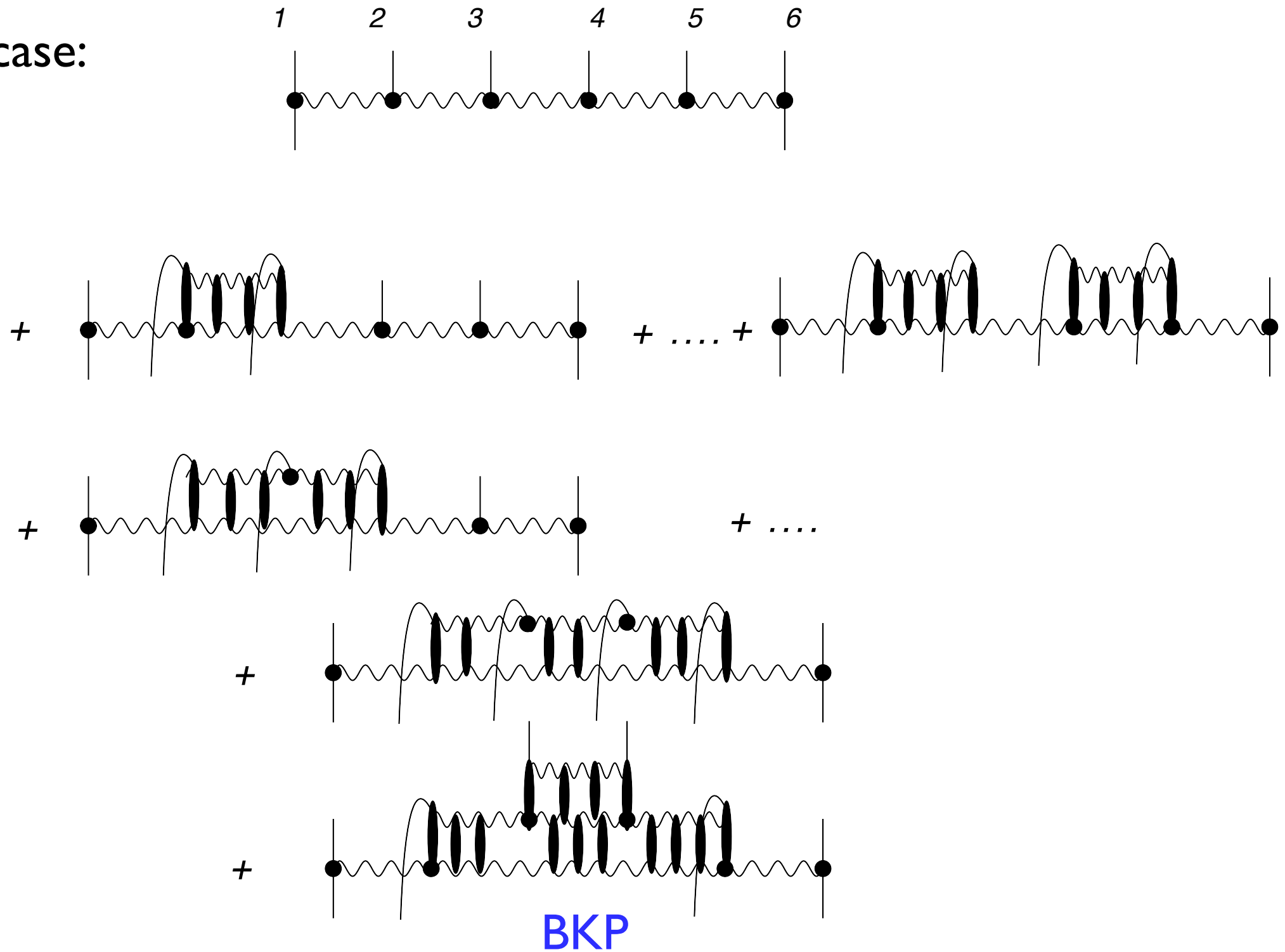
The 2--> 5 amplitude:



The same story:

- in certain kinematic regions Regge pole contribution does not factorize but generates unphysical singularities
- Regge cuts compensate these unwanted terms
- BDS contains Regge cuts in one loop.

The 2 --> 6 case:



Novel features:

- repetition of cuts (beginning of exponentiation?)
- new cut: three reggeon cut --> **spin chain, integrability**

Summary (ongoing work)

n-point scattering amplitudes, remainder functions = corrections to BDS
multi-Regge limit

in large N_c : Regge poles alone are inconsistent (cf. Veneziano model)

remainder functions = Regge cuts (beginning of cuts contained in BDS)

integrable spin chain: $2 \rightarrow 4, 2 \rightarrow 5$ two sites
 $2 \rightarrow 6, 2 \rightarrow 7$: three sites

a step towards an exact solution of $N=4$ SYM