

#### UNIVERSIDAD TECNICA FEDERICO SANTA MARIA







# The Pomeron spin-flip and its measurements

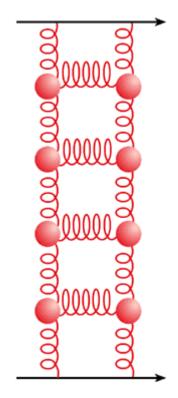
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#### Outline

- Motivation & introduction
- Pomeron spin-flip in *pp* collisions
- Nuclear target case
- Gold nucleus puzzle
- Conclusions







### Pomeron and its spin?

- **Pomeron** and differential cross section prediction is one of triumph of **Regge theory** in 1973.
- However, there is not *general agreement* about the pomeron spin.
- In community is a general agreement about the dominant spin non-flip hadronic interaction.
- But, there is no theoretical reason to think that the spinflip component of the **Pomeron** is zero.
  - If we assume **Pomeron** as Regge pole, we get zero spin asymmetry since the same phase
  - Even if we assume different phase (no Regge pole), we need spin-flip interaction to vary steeply with energy to distinguish them



#### **Spin-flip Pomeron interaction?**

Our goal:

Study the **spin-flip Pomeron** interaction.

#### Our method:

Study of the single spin asymmetry  $A_N(t)$  in the CNI region.

$$A_N \frac{d\sigma}{dt} = 2 \operatorname{Im}[\phi_{++}\phi_{+-}^*]$$
$$\frac{d\sigma}{dt} = |\phi_{++}|^2 + |\phi_{+-}|^2$$

 $\phi_{++}$  - Non-flip amplitude  $\phi_{+-}$  - Spin-flip amplitude



#### Why CNI region?

Let's assume that **Pomeron** can flip the spin.

Then, due to the same phase factor the hadronic single spin asymmetry will be zero anyway.

Solution is the interference with EM amplitude.

CNI (Coulomb-nuclear interference) region = a kinematical region of very low 4-momentum transfer squared, -t, where the interference electromagnetic-hadron terms dominates B.Z.Kopeliovich, B.G.Zakharov, Phys.Lett. B226 (1989) 156



#### How to calculate it?

Coulomb spin-flip and non-flip amplitude are known, as well as no-flip hadronic amplitude from data.

$$\phi^h = \phi_{++} \left( 1 + i \frac{\sqrt{-t}}{m_N} \vec{\sigma} \cdot \vec{n} r_5 \right)$$

Spin-flip hadron amplitude can be parametrized by factor

$$r_5 = \frac{m_p \phi_{+-}}{\sqrt{-t} \operatorname{Im} \phi_{++}}$$

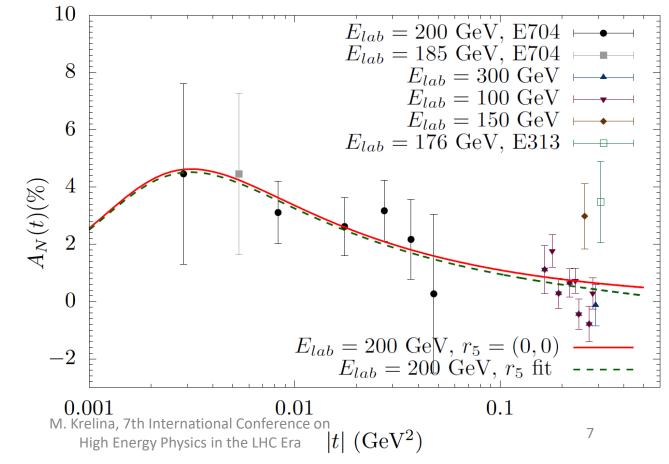
Assuming  $r_5 = 0$  the asymmetry  $A_N(t)$  can be fully predicted.

L.I.Lapidus & B.Kopeliovich Sov. J. Nucl. Phys. 19(1974) 114



#### Problems with spin-flip Pomeron in pp at low energies

- 1) Contribution from Reggeons,  $r_5 = r_{5,\mathbb{P}} + r_{5,\mathbb{R}} \Rightarrow \text{small } r_5$ does not mean small  $r_{5,\mathbb{P}}$ ...
- 2) Large experimental uncertainties

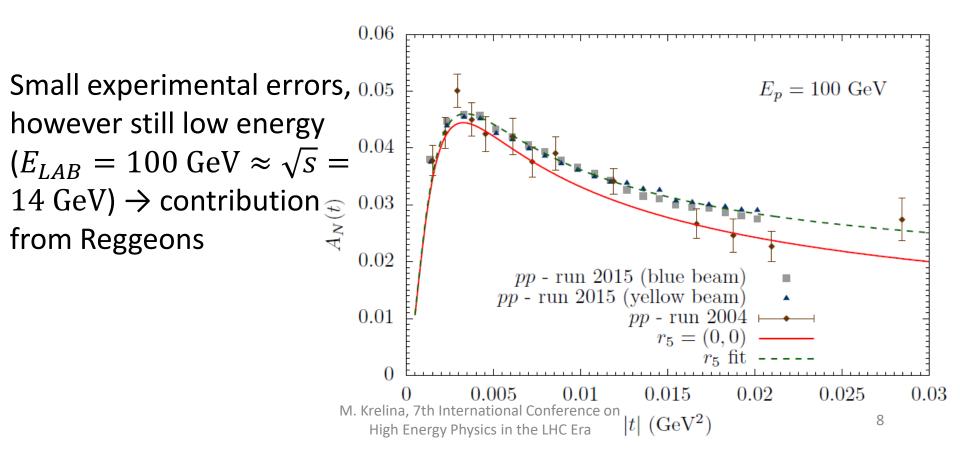




#### pp data from H-JET

Combined  $r_5$  fit result

#### $r_5 = -0.0073 \pm 0.0032 - i0.0289 \pm 0.0128$





#### What next?

#### We want to suppress Reggeons

# ⇒ Higher energy⇒ Nuclear targets



### Why nuclear target?

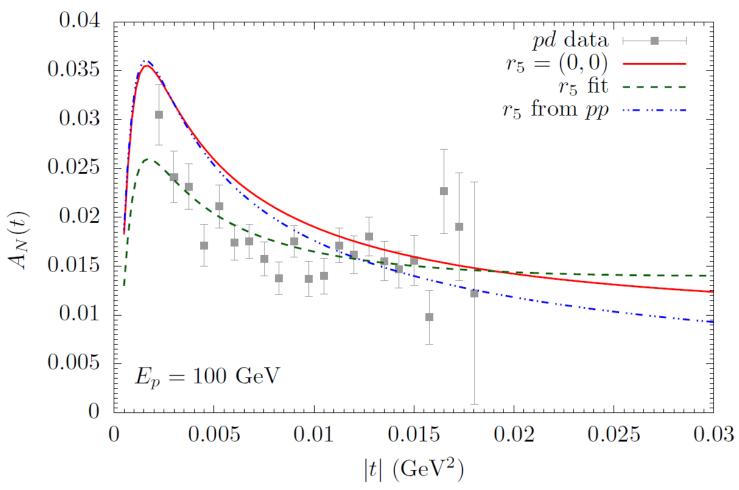
Two motivations:

**Polarimetry** – was actual 10 years ago, expected smaller errors at *pA* elastic scattering.

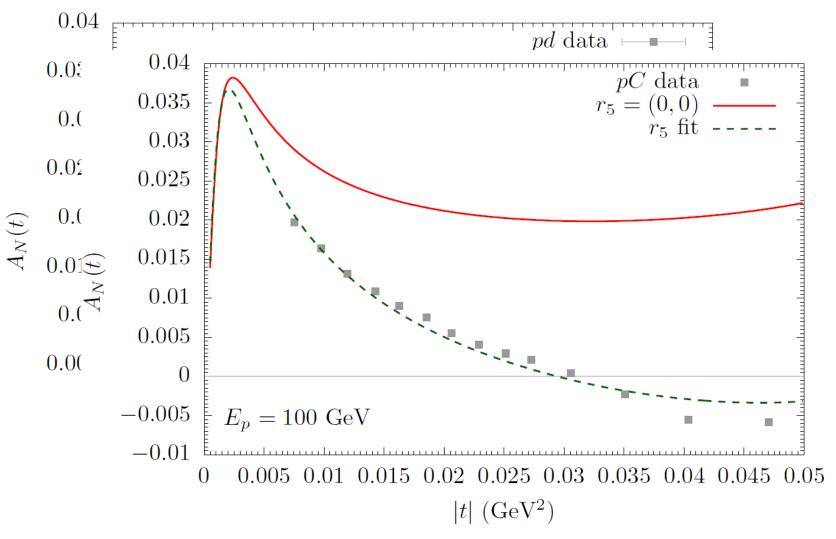
**Reggeons** – experimental data mostly from RHIC ( $E_{LAB} = 100 \text{ GeV} \approx \sqrt{s} = 14 \text{ GeV}$ ). Expected a contribution mainly from the iso-vector Reggeons.

If we use the nucleus with zero isospin (e.g. Carbon), these Reggeons are excluded. For other nuclei are suppressed as 1/A. B. Kopeliovich, hep-ph/9801414

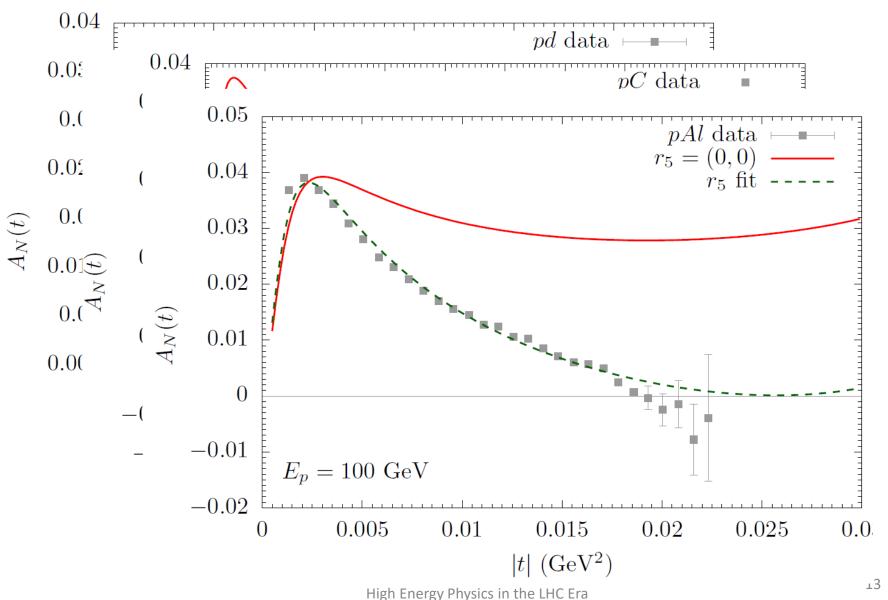














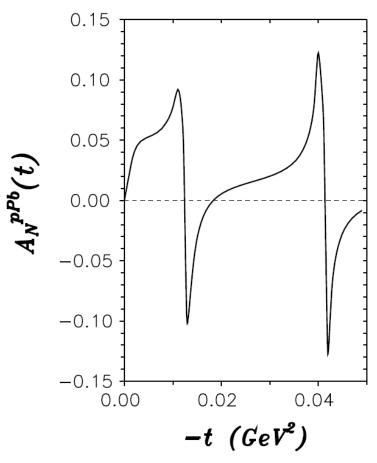
- *pd*:  $r_5 = -0.005 \pm 0.003 + i0.267 \pm 0.056$
- *pC*:  $r_5 = -0.051 \pm 0.001 i0.014 \pm 0.014$
- *pAI*:  $r_5 = -0.100 \pm 0.003 i0.183 \pm 0.096$
- *pC* and *pAI* indicates higher  $r_5$  then  $pp \Rightarrow$  significant contribution from Reggeons in *pp*.

Remember result from *pp*:

 $r_5 = -0.0073 \pm 0.0032 - i0.0289 \pm 0.0128$ 

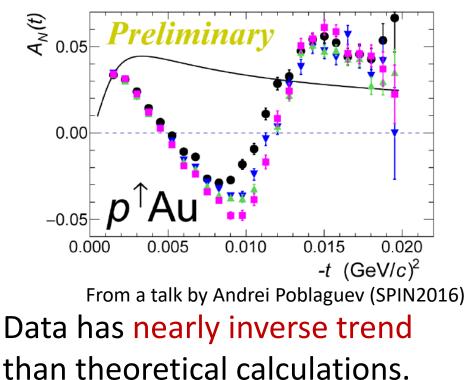


#### ...but the Gold is the challenge



B. Kopeliovich, hep-ph/9801414

Estimation of  $r_{5,\mathbb{P}}$  form Carbon is sufficient, for Gold the situation is more complicated. However, take a look at it...





#### Mistakes in formulas

We want to understand the *t* dependence  $\Rightarrow$  need to revise the calculation. We start from the Glauber. We found some mistakes in signs in original formulas.

B.Z. Kopeliovich and T.L. Trueman, Phys. Rev. D64, 034004 (2001)

$$\frac{16\pi}{(\sigma_{tot}^{pA})^2} \frac{d\sigma_{pA}}{dt} = \left[\frac{t_c}{t} F_A^{em}(t)\right]^2 + \left[\operatorname{Im} F_A^h(t)\right]^2 \left(\rho_{pA}^2(t) + 1 - \frac{t}{m_N^2} |r_5|^2\right) \\ + 2\frac{t_c}{t} F_A^{em}(t) F_A^h(t) \left(\rho_{pA}(t) + \delta_{pA}\right), \\ \frac{16\pi}{(\sigma_{tot}^{pA})^2} \frac{d\sigma_{pA}}{dt} A_N^{pA}(t) = \frac{\sqrt{-t}}{m_N} \operatorname{Im} F_A^h(t) \left\{ 2\operatorname{Im} F_A^h(t) (\operatorname{Rer}_5 - \operatorname{Im} r_5 \rho_{pA}(t)) \\ + F_A^{em}(t) \frac{t_c}{t} \left[ (\mu_p - 1)(1 - \delta_{pA} \rho_{pA}(t)) - 2(\operatorname{Im} r_5 - \delta_{pA} \operatorname{Rer}_5) \right] \right\}$$

However, wrong signs were at not dominant terms and have no effect for the Gold description.



### Wrong EM form factor

We found that the source of the trouble is the incorrect electromagnetic form factor, where we discovered the importance of the absorption

$$\phi_{em}(q) = \sqrt{\pi} Z \alpha_{em} \left(\frac{2}{q^2} + \frac{\mu_p - 1}{q}\right) F_A^{em}(q^2) e^{i\delta_{pA}} \otimes e^{-\frac{1}{2}\sigma_{tot}^{pp}T_A(b)}$$

$$\overset{0.15}{\underset{0}{}_{0}} \overset{0.15}{\underset{0}{}_{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0.15}{\underset{0}}{\underset{0}} \overset{0.15}{\underset{0}} \overset{0$$



#### Other corrections

To have a full description we should add other correction such as Gribov correction or nucleon-nucleon correlations.

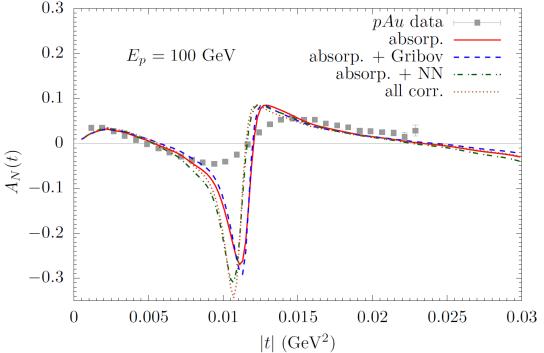
#### Gribov corrections – effectively increase the pA cross section

B. Z. Kopeliovich, Int. J. Mod. Phys. A31 no. 28n29, (2016) 1645021, arXiv:1602.00298 [hep-ph].

B. Z. Kopeliovich, I. K. Potashnikova, and I. Schmidt, Phys. Rev. C73 (2006) 034901, arXiv:hep-ph/0508277 [hep-ph].

#### NN correlations – effectively reduce the nuclear thickness function

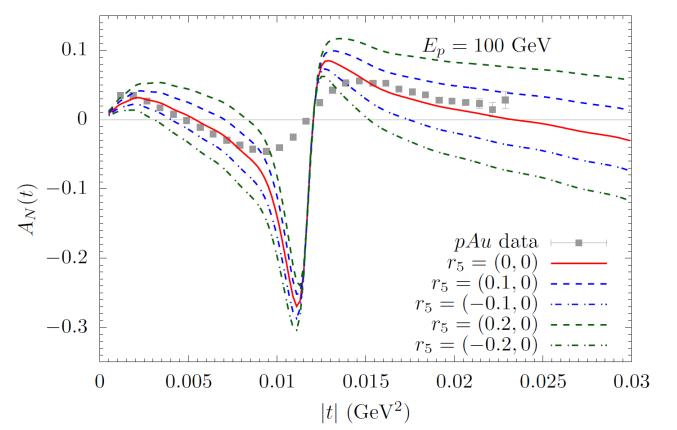
M. Alvioli, C. Ciofi degli Atti, B. Z. Kopeliovich, I. K. Potashnikova, and I. Schmidt, Phys. Rev. C81 (2010) 025204, arXiv:0911.1382 [nucl-th].





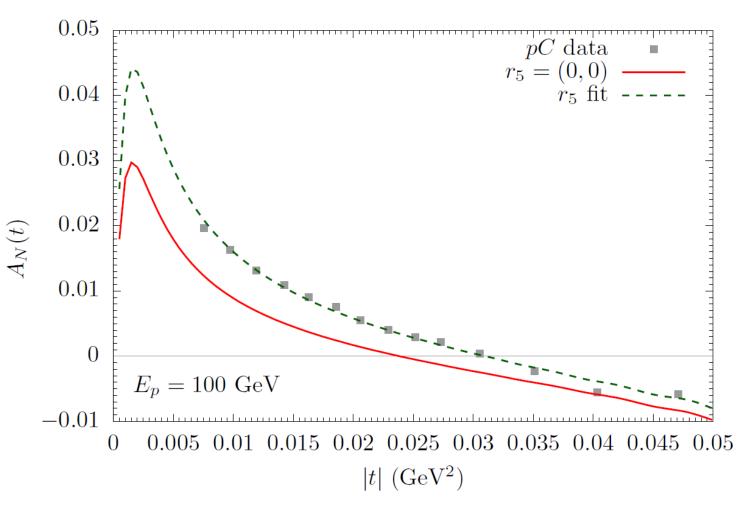
#### Further adjustments

#### Finally, we can make some adjustment by non-zero $r_5$

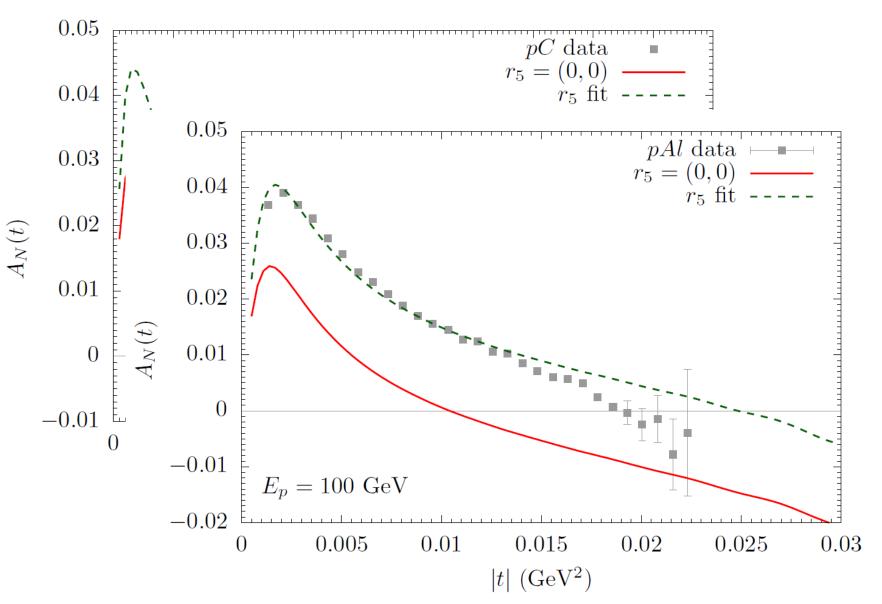


The result looks reasonable, good agreement at low and high *t*, good position of the cross points.

# *pC, pAI* with absorption correction



# *pC, pAI* with absorption correction





## pC, pAl with absorption correction

*pC*:  $r_5 = -0.031 \pm 0.001 - i0.384 \pm 0.017$ 

*pAI*:  $r_5 = -0.074 \pm 0.002 - i0.376 \pm 0.029$ 

- Very different spin asymmetry for zero  $r_5$
- *pC* and *pAI* closer to each other
- High sensitivity for real part of  $r_5$

Remember result from *pp*:

 $r_5 = -0.0073 \pm 0.0032 - i0.0289 \pm 0.0128$ 



#### Conclusions

- We study the CNI region to see the effect of spin-flip hadronic amplitude.
- Indicated small  $r_5$  in pp at RHIC does not report about Pomeron spinflip interaction, it is combination of Pomeron and Reggeon.
- We are interested into the nuclear target because of exclusion or suppression of Reggeons.
- Data for pC indicated higher  $r_5$  then in pp. It can be interpreted that the pomeron  $r_5$  is suppressed in pp. Similar situation for pAl.
- More complex situation in case of Gold target. Unexpected experimentally measured *t* dependence.
- Wrong signs found in the original formulas.
- A novel mechanism of interference of electromagnetic UPC with central hadronic collisions is proposed attempting at explanations of pAu data for CNI generated  $A_N(t)$
- We included other expected correction. Finally we have good agreement at low and high *t*, good position of the crossing points.
- Nevertheless, an accurate determination of  $r_5$  from *pAu* data is not possible so far.



#### Thank you for your attention



#### Acknowledgement

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- This work was funded by Conicyt PIA/Basal FB0821, Chile