

11 September 2017
TOOLS workshop, Corfu

Tools for Dark Matter Indirect Detection

Marco Cirelli
(CNRS LPTHE Jussieu)

Based on:

Cirelli, Corcella, Hektor, Hutsi, Kadastik, Panci, Raidal, Sala, Strumia,
JCAP 1103 (2011) 051 [1012.4515]

Baratella, Cirelli, Hektor, Pata, Piibeleht, Strumia
JCAP 1403 (2014) 053 [1312.6408]

Boudaud, Cirelli, Giesen, Salati
JCAP 1505 (2015) no.05, 013 [1412.5695]

Buch, Cirelli, Giesen, Taoso
JCAP 1509 (2015) no.09, 037 [1505.01049]



NewDark

11 September 2017
TOOLS workshop, Corfu

Tools for Dark Matter Indirect Detection

Marco Cirelli
(CNRS LPTHE Jussieu)

Based on:

Cirelli, Corcella, Hektor, Hutsi, Kadastik, Panci, Raidal, Sala, Strumia,
JCAP 1103 (2011) 051 [1012.4515]

Baratella, Cirelli, Hektor, Pata, Piibeleht, Strumia
JCAP 1403 (2014) 053 [1312.6408]

Boudaud, Cirelli, Giesen, Salati
JCAP 1505 (2015) no.05, 013 [1412.5695]

Buch, Cirelli, Giesen, Taoso
JCAP 1509 (2015) no.09, 037 [1505.01049]



NewDark

Outline

1. arouse your interest in DM ID

Outline

1. arouse your interest in DM ID
2. talk to you about the gory details of DM ID for 3 hours...

Outline

1. arouse your interest in DM ID
2. talk to you about the gory details of DM ID for 3/4 hour...

Outline

1. arouse your interest in **DM ID**
2. talk to you about the
gory **details** of DM ID
for 3/4 hour...
3. ...in order to convince you that
you can forget everything and
trust **PPPC4DMID**

DM detection

direct detection

production at colliders

- indirect
 - γ from annihil in galactic center or halo
and from synchrotron emission
Fermi, ICT, radio telescopes...
 - e^+ from annihil in galactic halo or center
PAMELA, Fermi, HESS, AMS, balloons...
 - \bar{p} from annihil in galactic halo or center
 - \bar{d} from annihil in galactic halo or center
GAPS
 - $\nu, \bar{\nu}$ from annihil in galaxy or massive bodies
SK, Icecube, Km3Net

DM detection

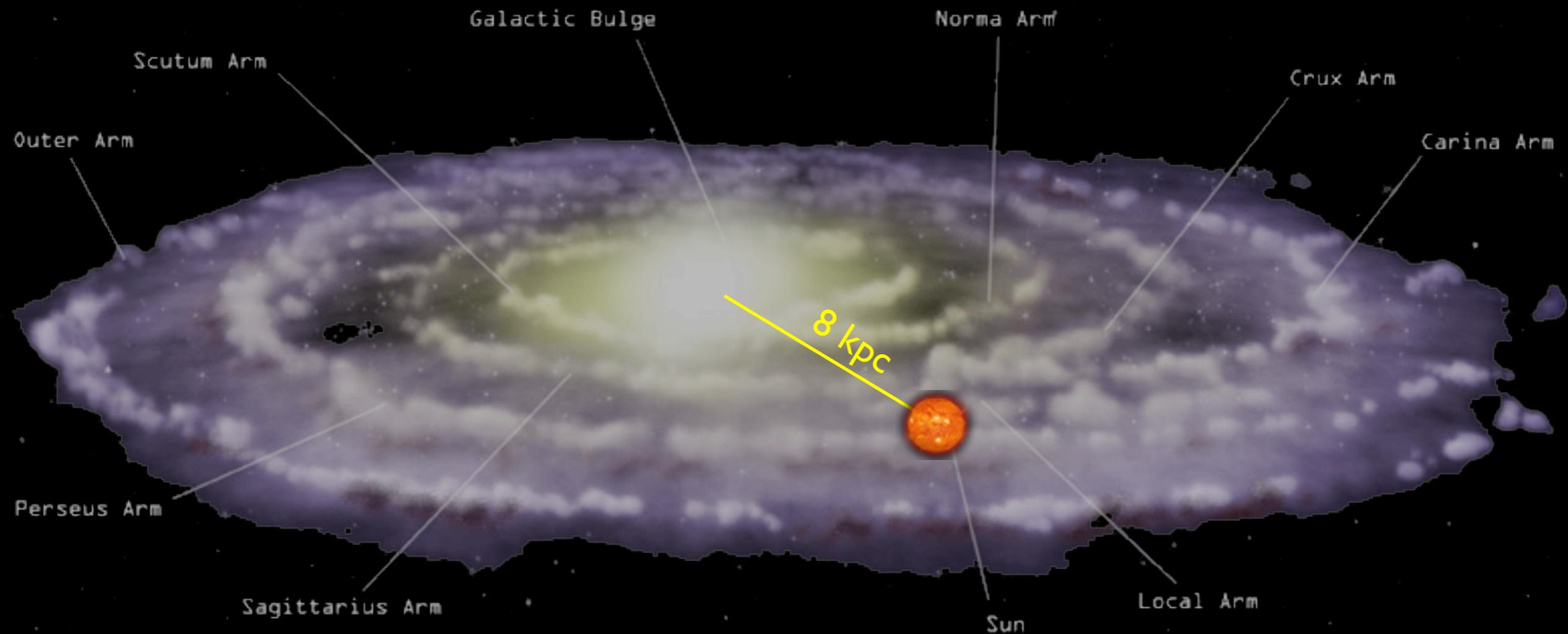
direct detection

production at colliders

- indirect
 - γ from annihil in galactic center or halo
and from synchrotron emission
Fermi, ICT, radio telescopes...
 - e^+ from annihil in galactic halo or center
PAMELA, Fermi, HESS, AMS, balloons...
 - \bar{p} from annihil in galactic halo or center
 - \bar{d} from annihil in galactic halo or center
GAPS
 - $\nu, \bar{\nu}$ from annihil in galaxy or massive bodies
SK, Icecube, Km3Net

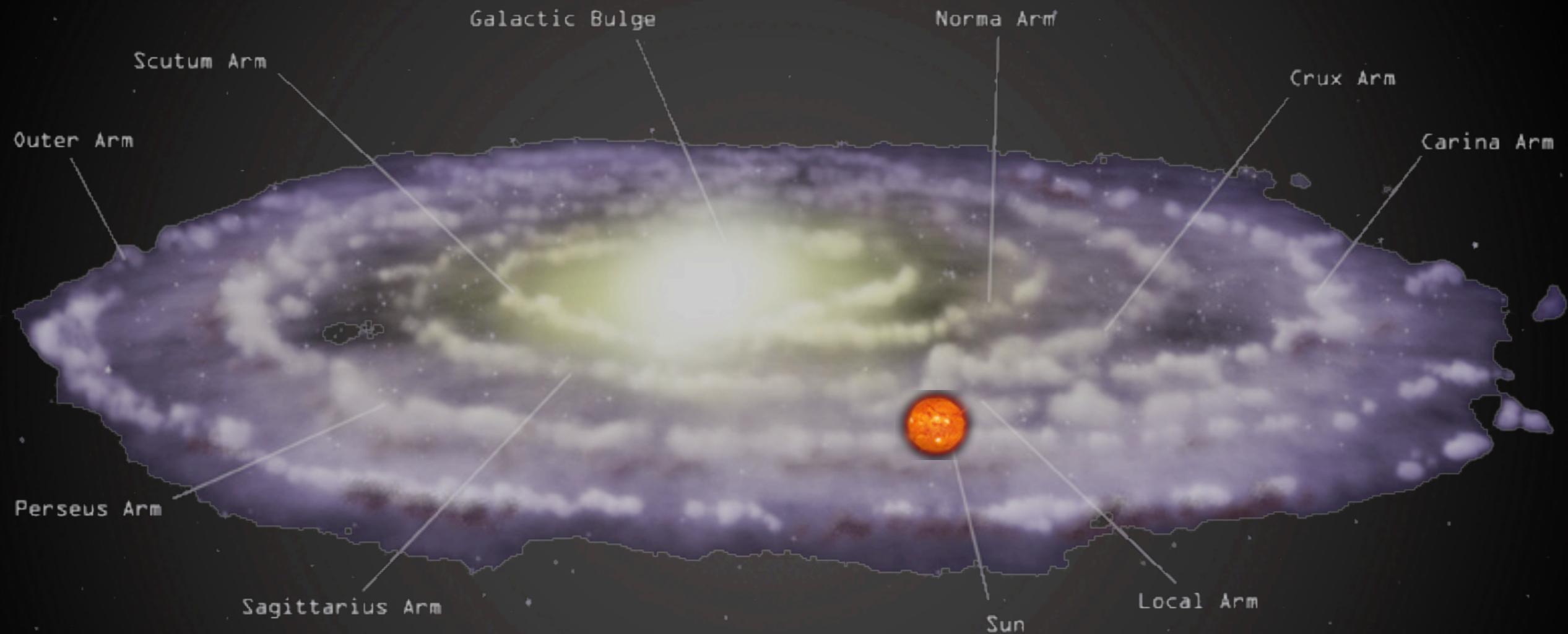
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



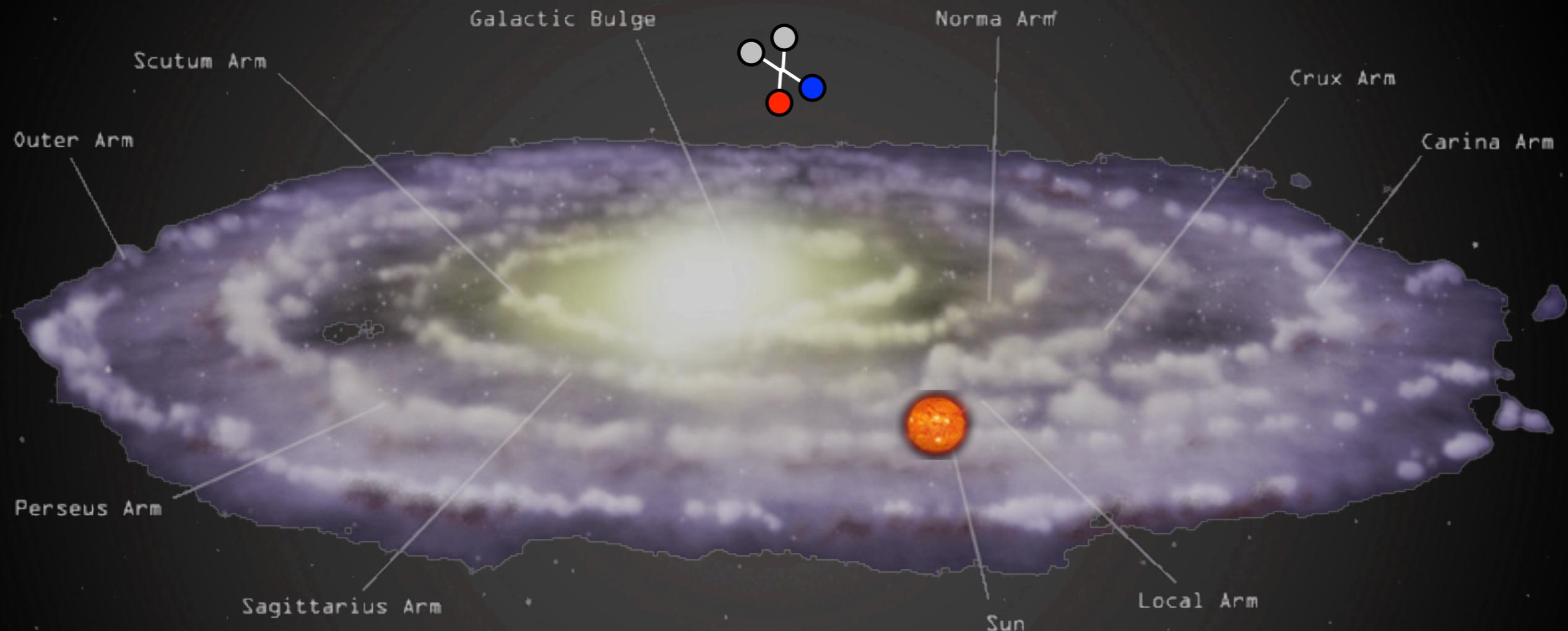
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo

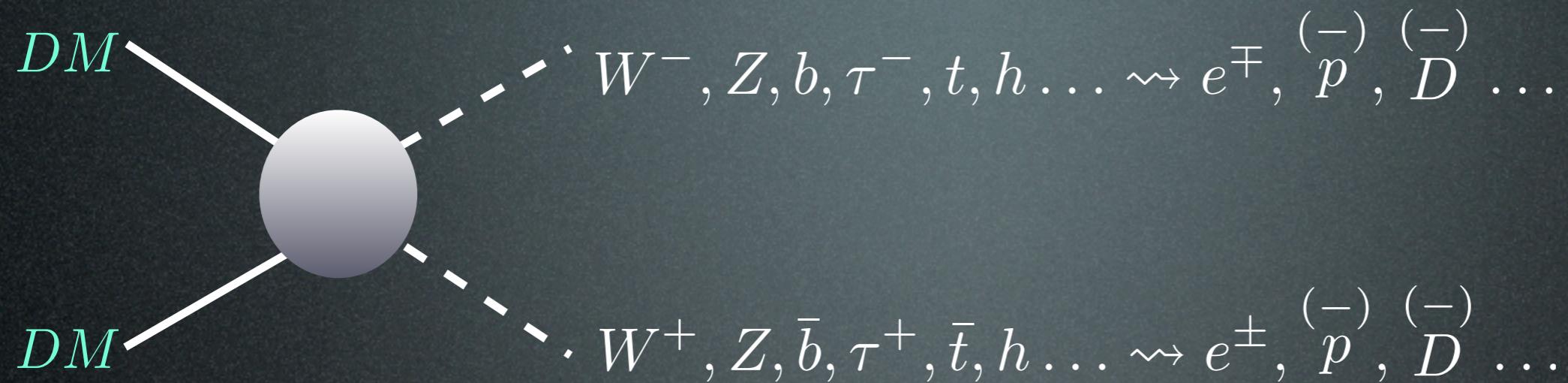


Indirect Detection: basics

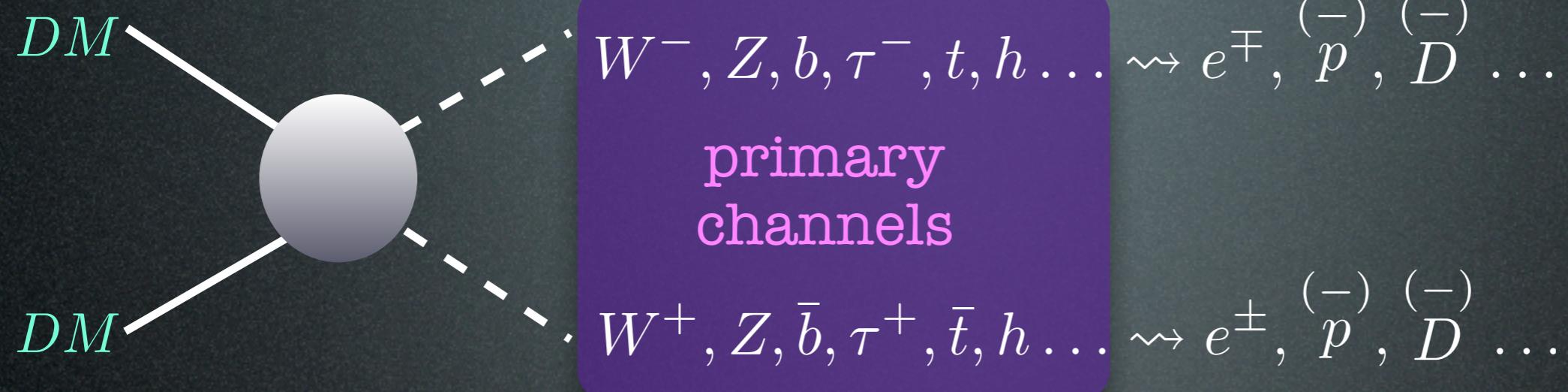
\bar{p} and e^+ from DM annihilations in halo



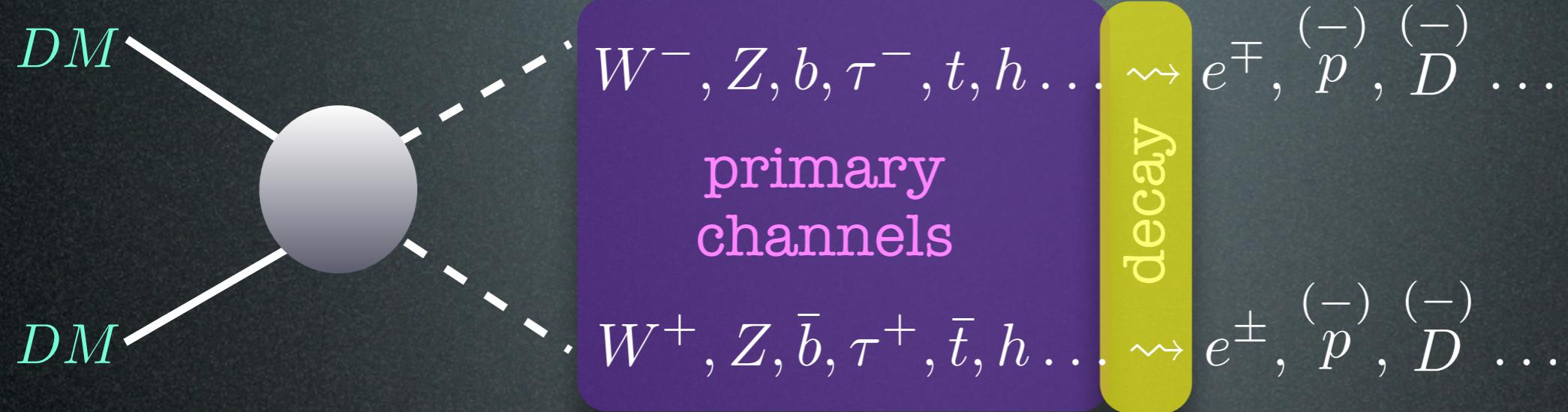
Fluxes at production



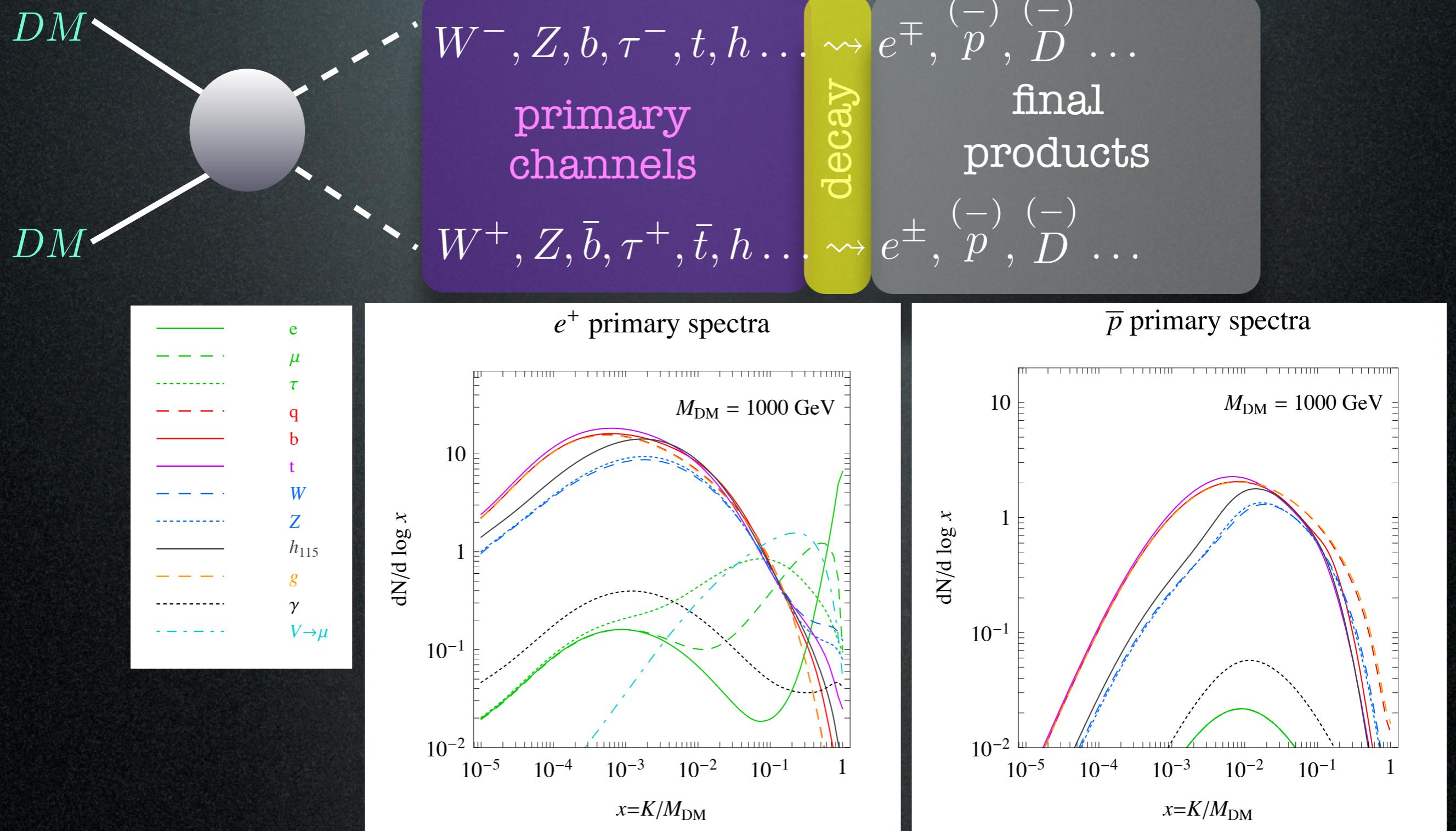
Fluxes at production



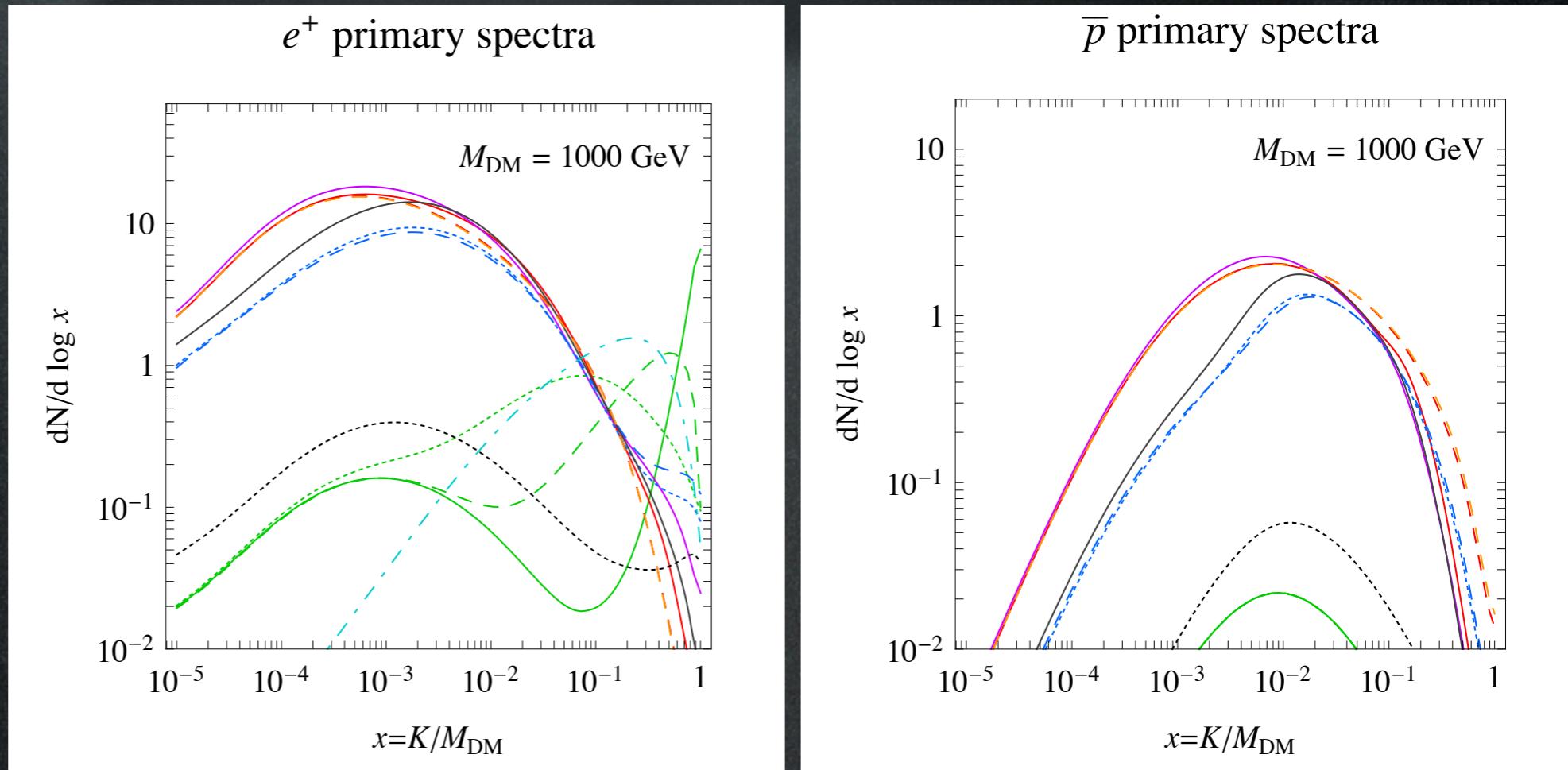
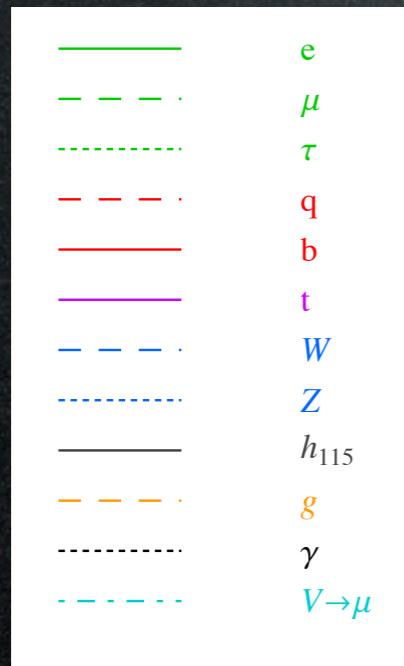
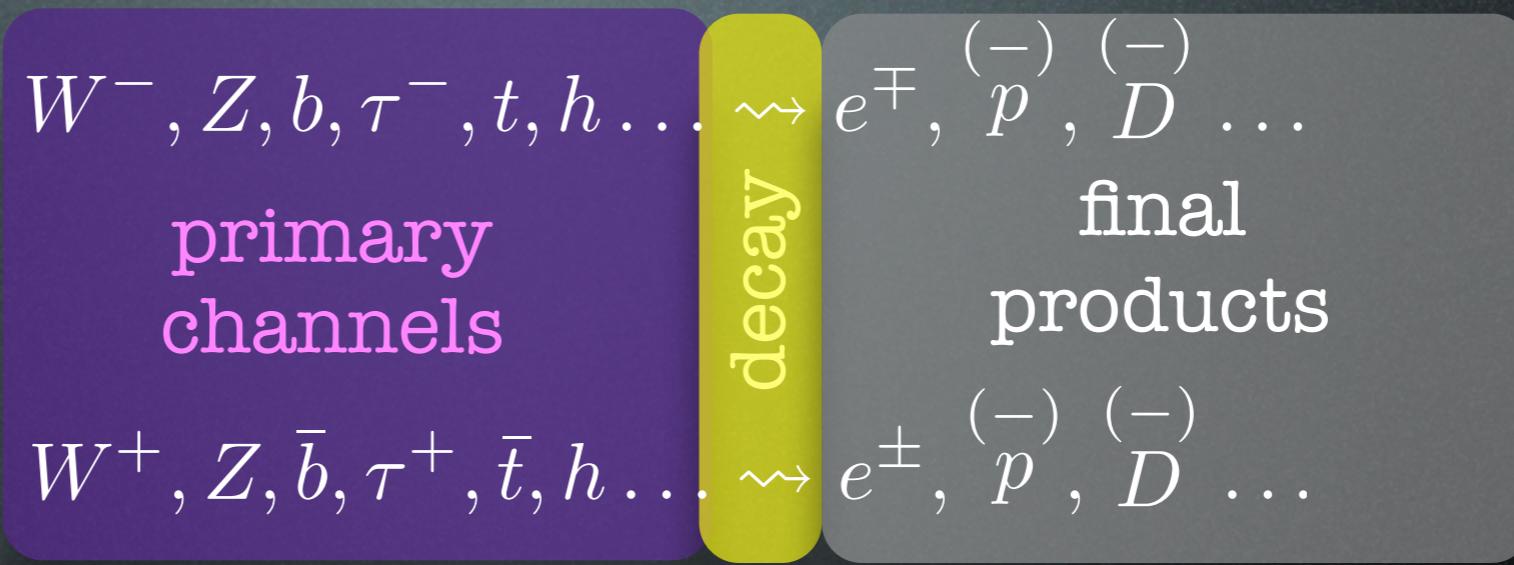
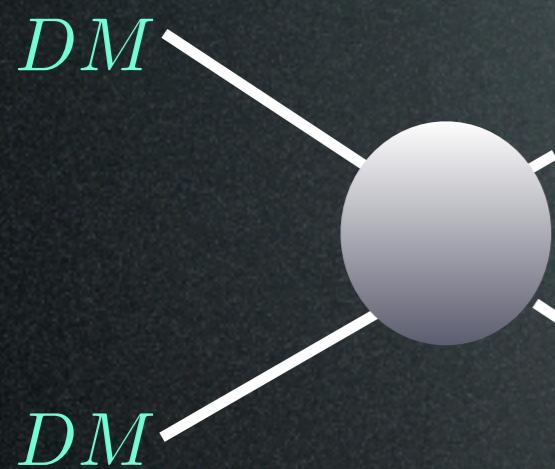
Fluxes at production



Fluxes at production



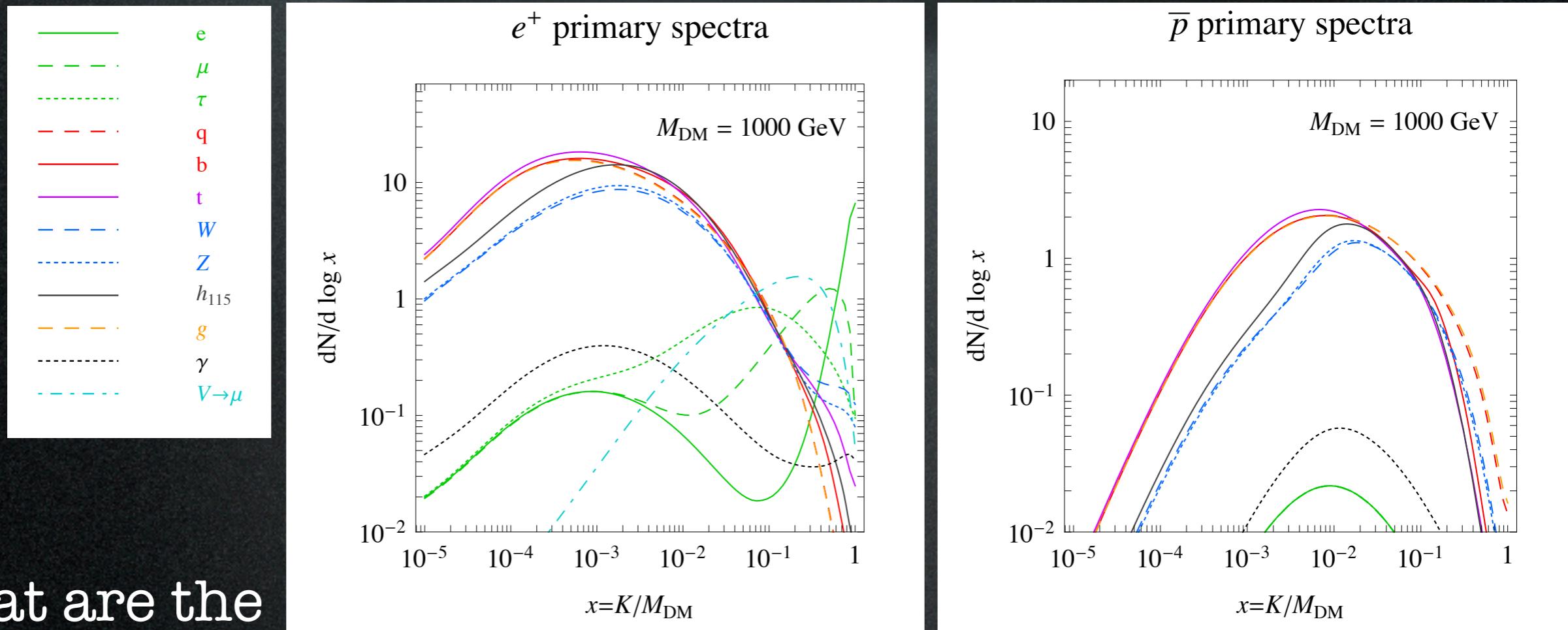
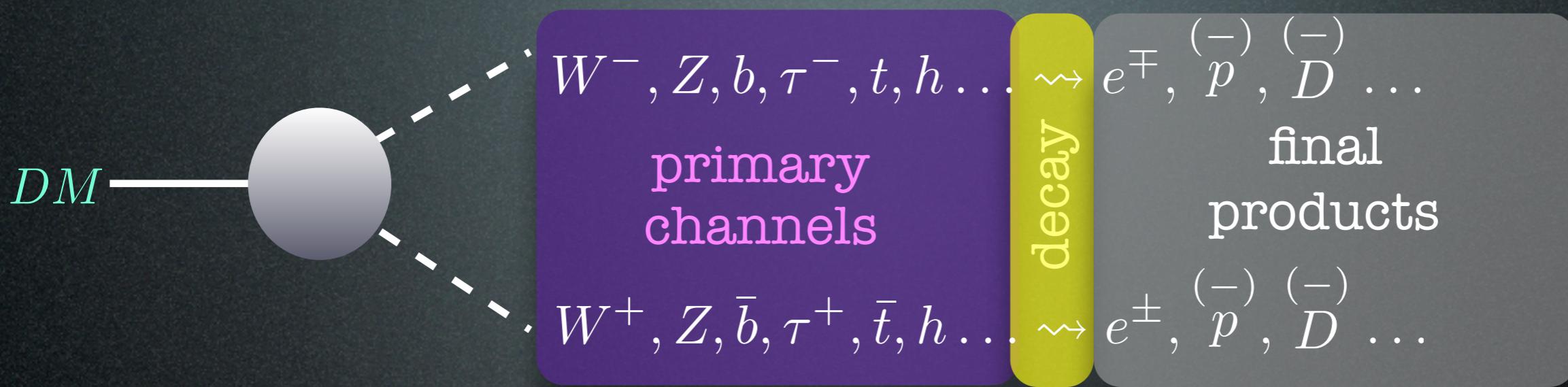
Fluxes at production



So what are the
particle physics
parameters?

1. Dark Matter mass
2. primary channel(s)

Fluxes at production



So what are the particle physics parameters?

1. Dark Matter mass
2. primary channel(s)

Fluxes at production

Different hadronic MonteCarlos could give different products

Or: what is the ‘systematic’ uncertainty?

Fluxes at production

Different hadronic MonteCarlos could give different products

Or: what is the ‘systematic’ uncertainty?

PYTHIA (8.135) vs **HERWIG** (6.510)



e.g. lacks γ radiation
from W^+W^- states
(added)



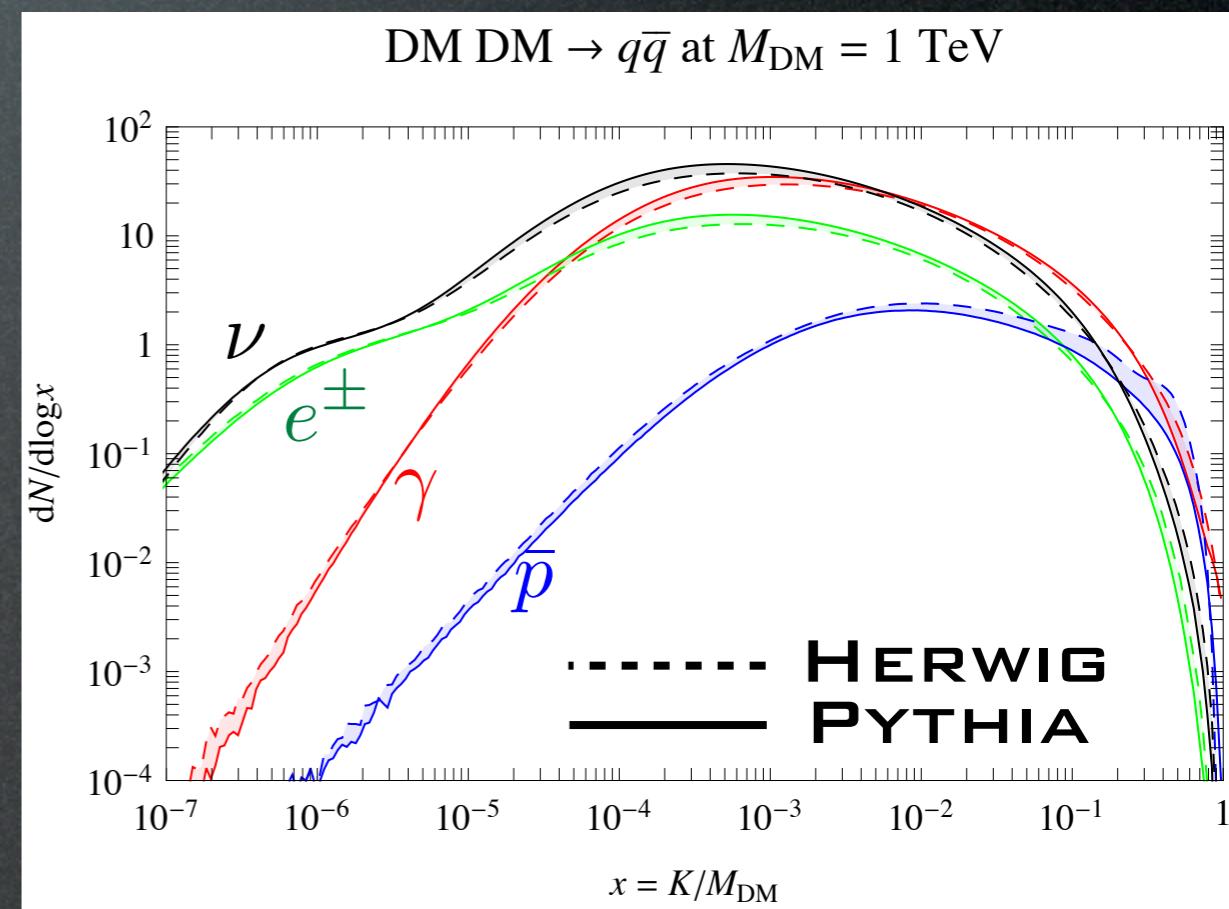
e.g. lacks $\ell \rightarrow \ell\gamma$
and $\gamma \rightarrow f\bar{f}$
branchings

Fluxes at production

Different hadronic MonteCarlos could give different products
Or: what is the ‘systematic’ uncertainty?

PYTHIA (8.135) VS **HERWIG** (6.510)

e.g. lacks γ radiation from W^+W^- states (added)
e.g. lacks $\ell \rightarrow \ell\gamma$ and $\gamma \rightarrow f\bar{f}$ branchings



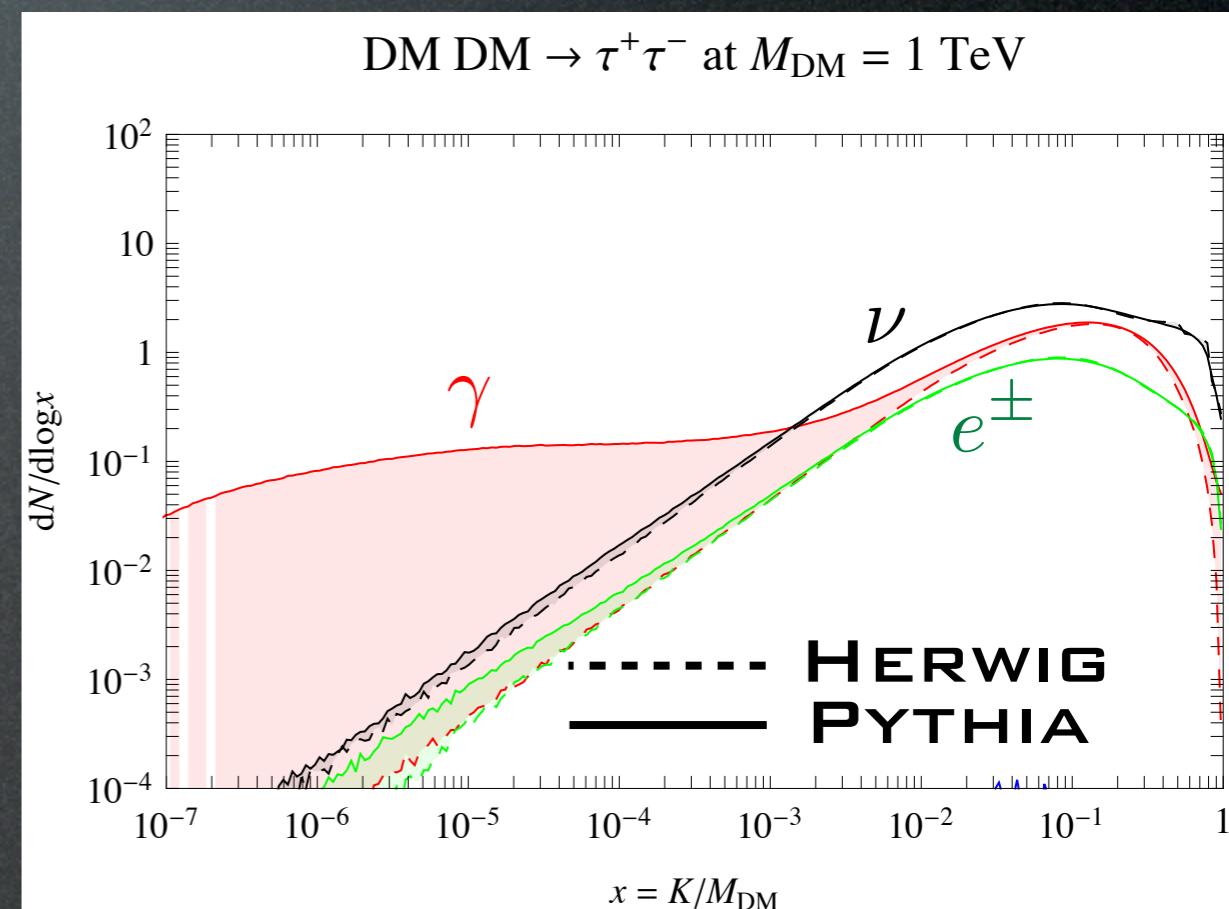
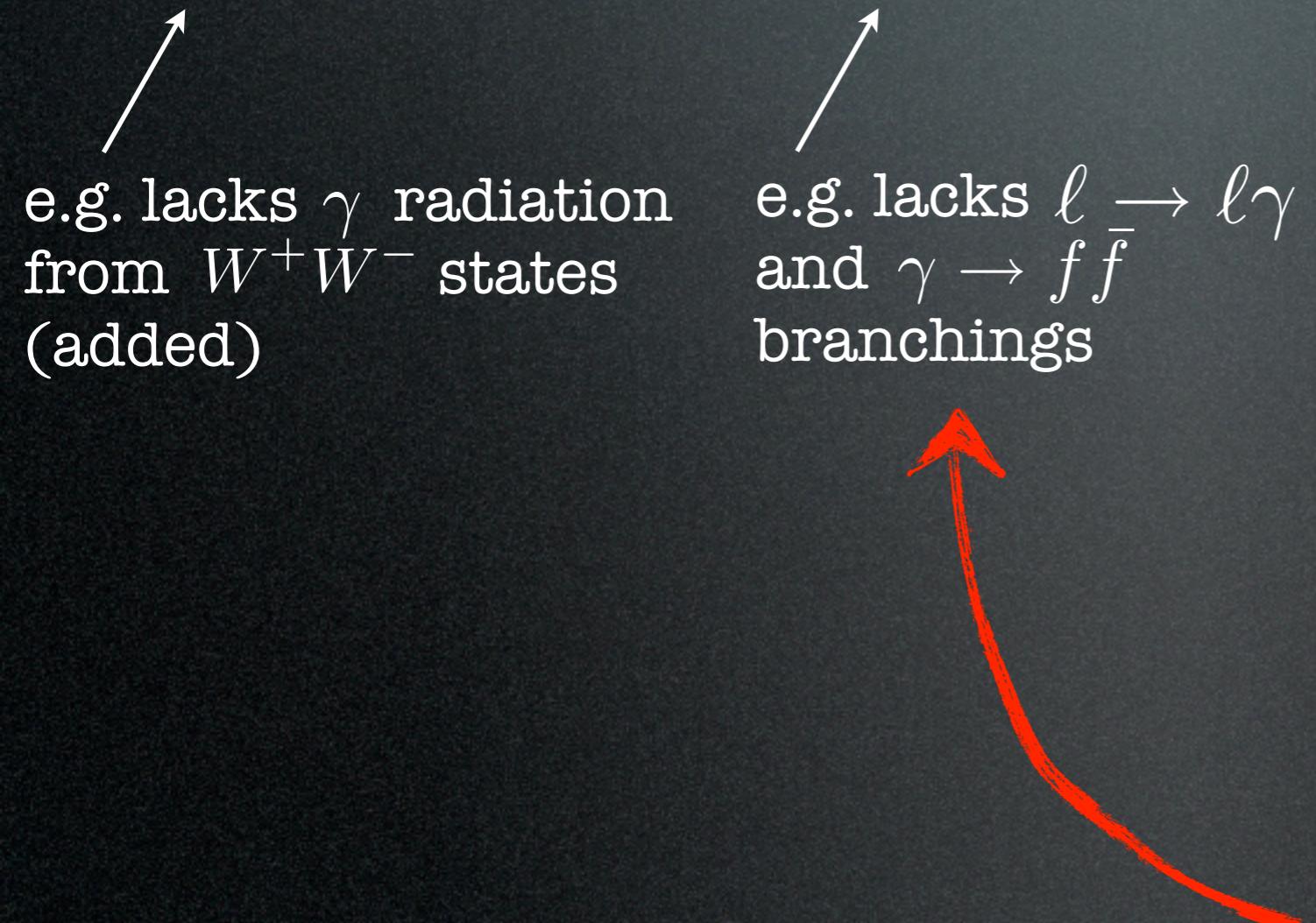
calibrated on LEP processes,
good agreement, overall 20%

Fluxes at production

Different hadronic MonteCarlos could give different products

Or: what is the ‘systematic’ uncertainty?

PYTHIA (8.135) VS **HERWIG** (6.510)



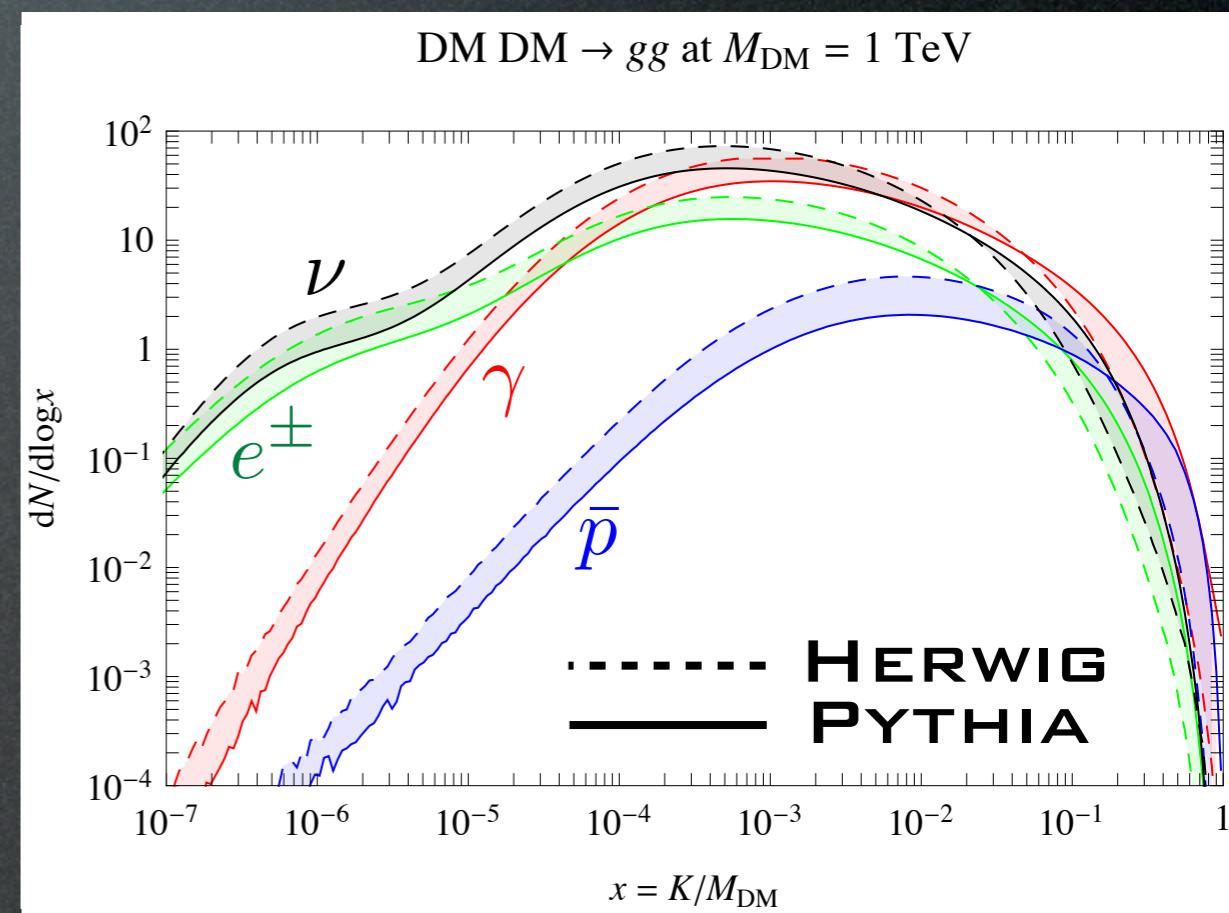
big discrepancy in γ (e^\pm) fluxes at low energy

Fluxes at production

Different hadronic MonteCarlos could give different products
Or: what is the ‘systematic’ uncertainty?

PYTHIA (8.135) VS **HERWIG** (6.510)

e.g. lacks γ radiation from W^+W^- states (added)
e.g. lacks $\ell \rightarrow \ell\gamma$ and $\gamma \rightarrow f\bar{f}$ branchings



Factor 2: not calibrated on LEP?
Anyway not central for DM

Fluxes at production

Different hadronic MonteCarlos could give different products

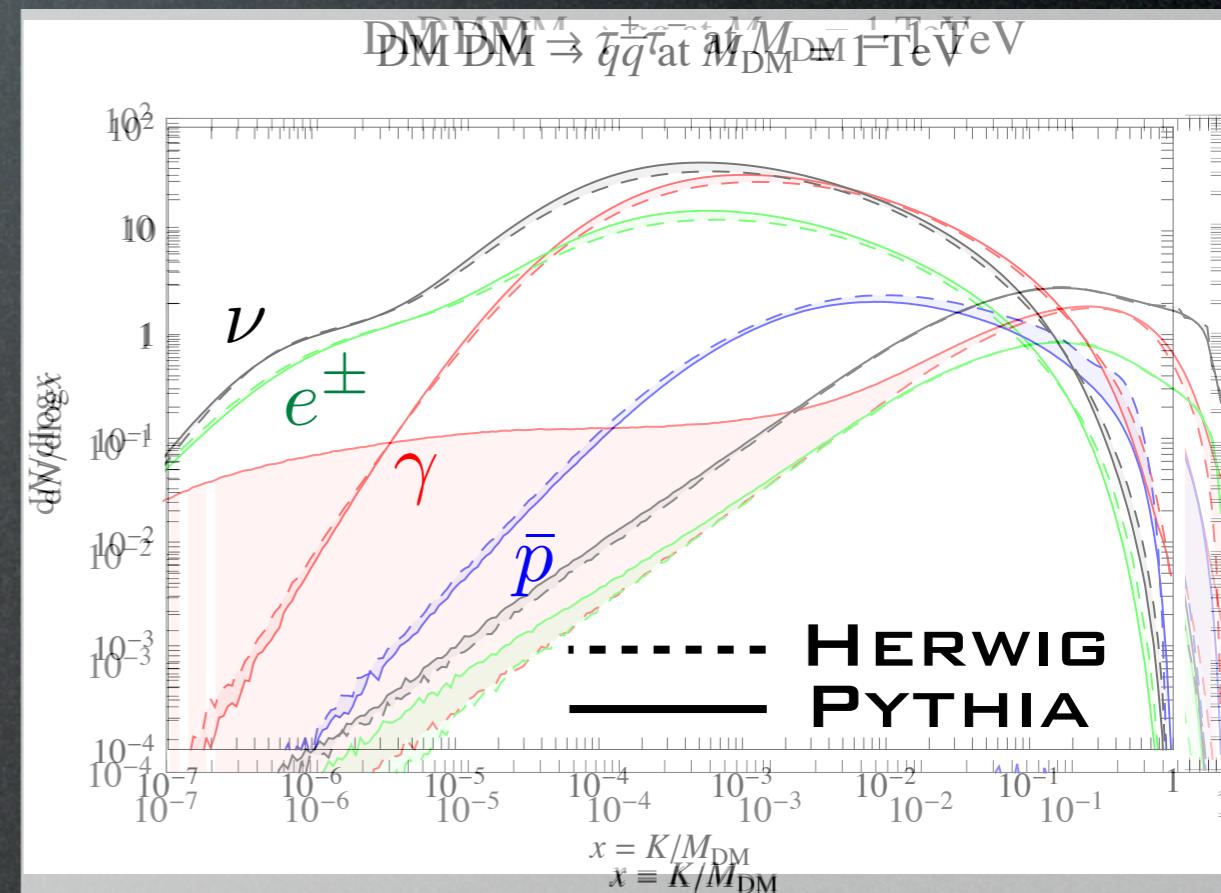
Or: what is the ‘systematic’ uncertainty?

- overall around 20%
- with some surprises

PYTHIA (8.135) vs **HERWIG** (6.510)

e.g. lacks γ radiation from W^+W^- states (added)

e.g. lacks $\ell \rightarrow \ell\gamma$ and $\gamma \rightarrow f\bar{f}$ branchings



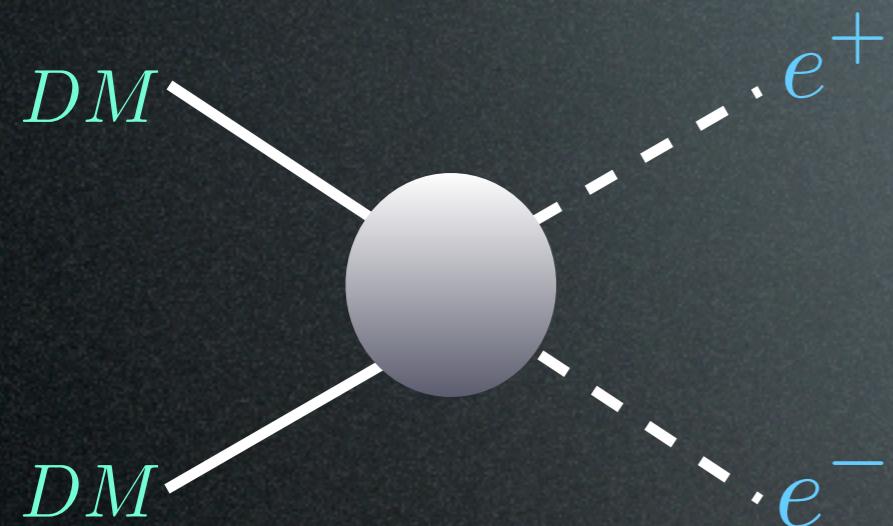
We use (modified) **PYTHIA 8** for all computations.

Fluxes at production

ElectroWeak corrections are important!

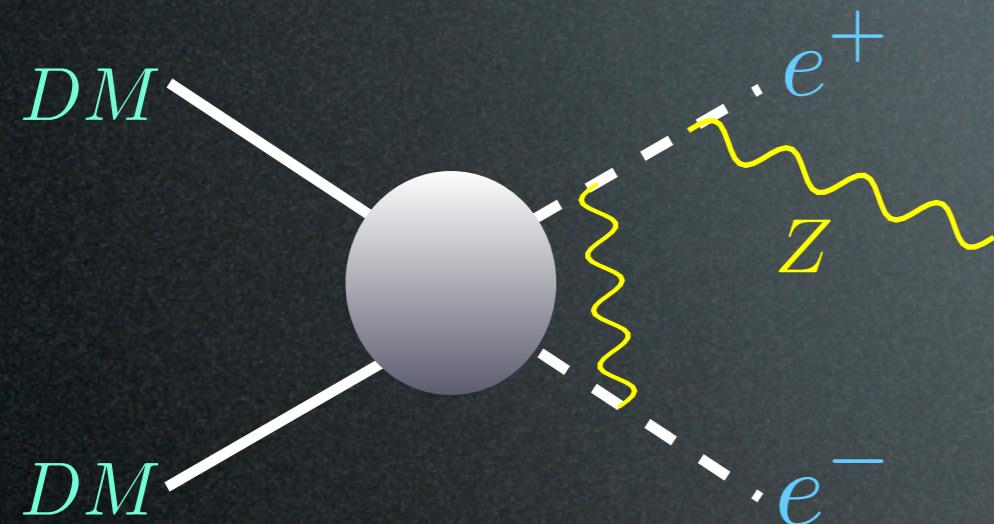
Fluxes at production

ElectroWeak corrections are important!



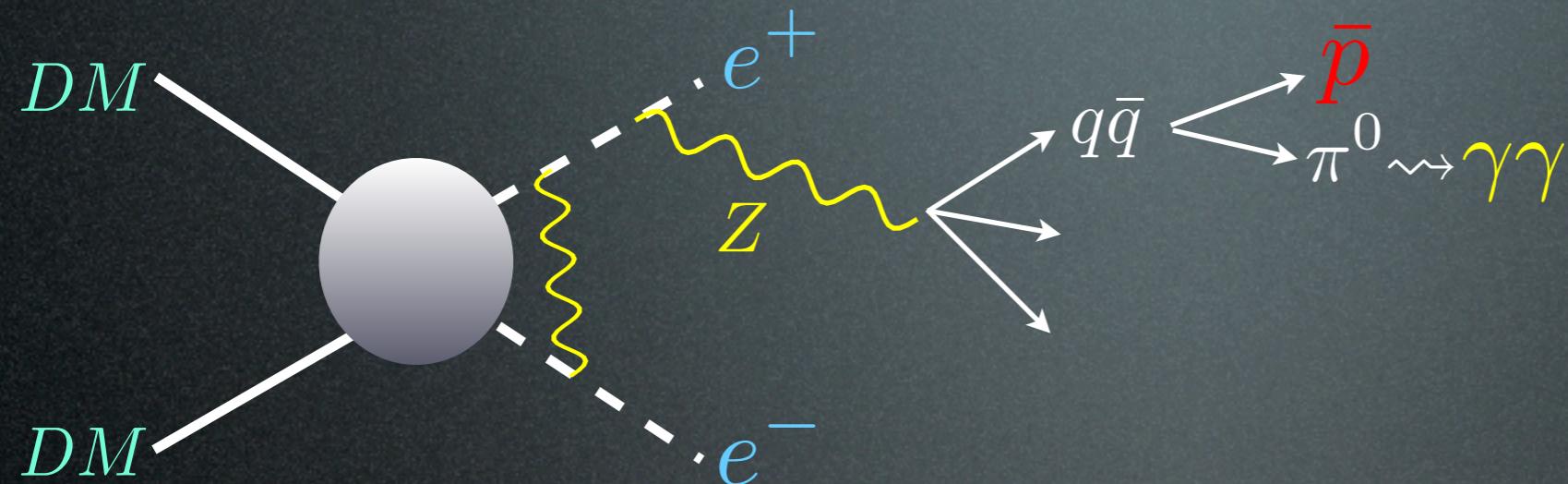
Fluxes at production

ElectroWeak corrections are important!



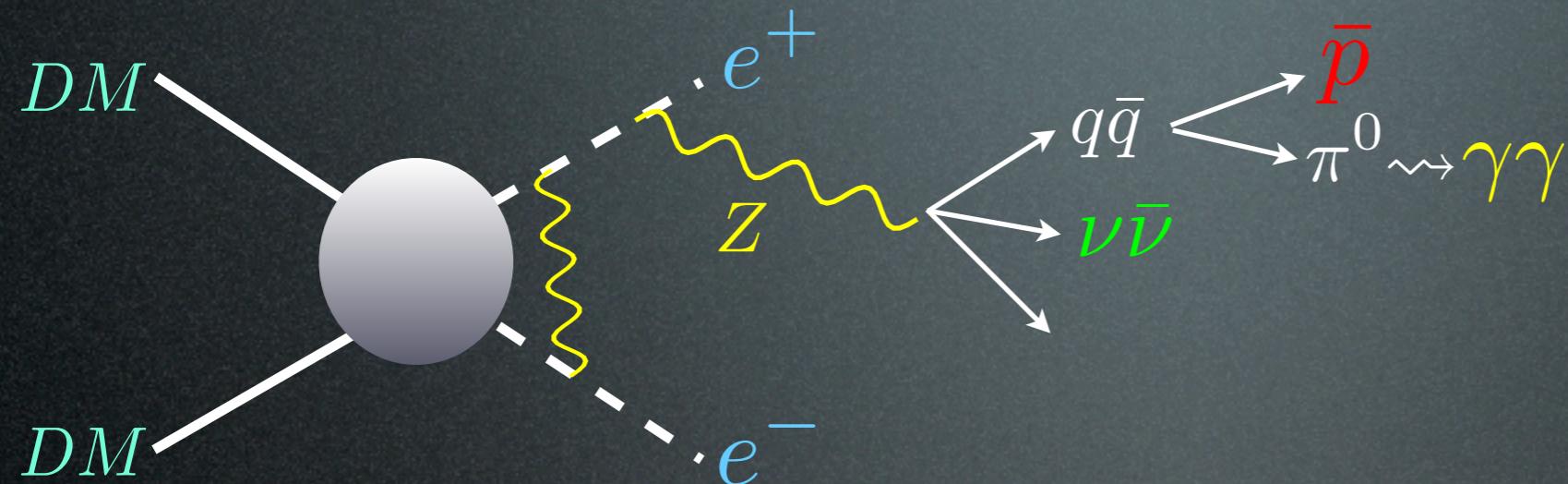
Fluxes at production

ElectroWeak corrections are important!



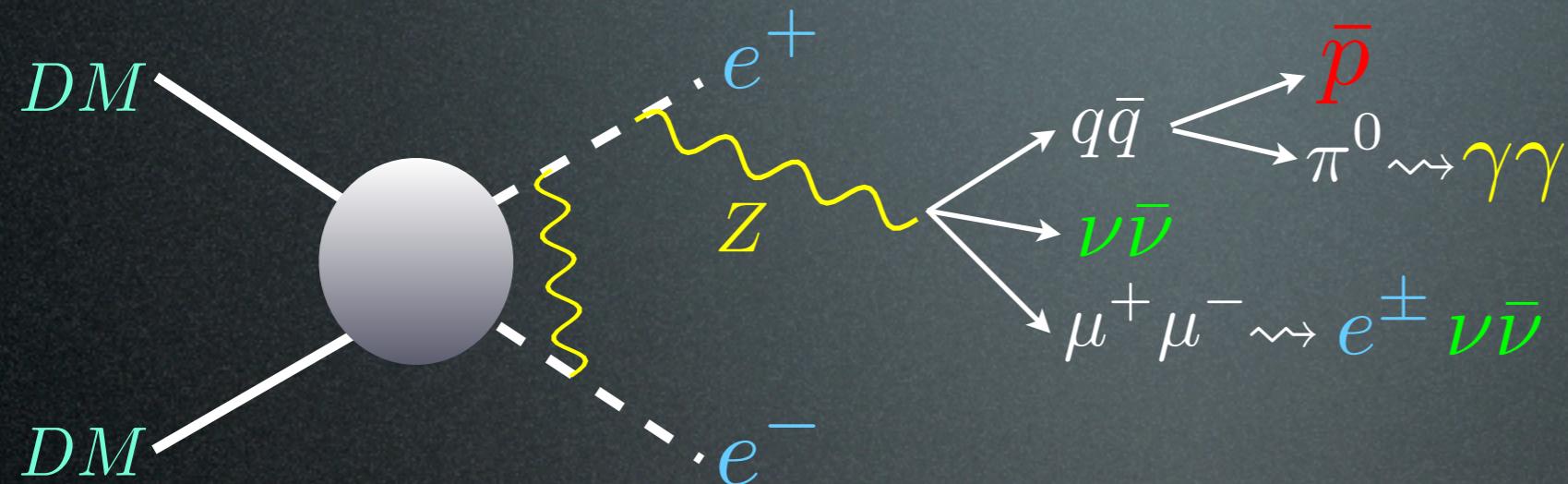
Fluxes at production

ElectroWeak corrections are important!



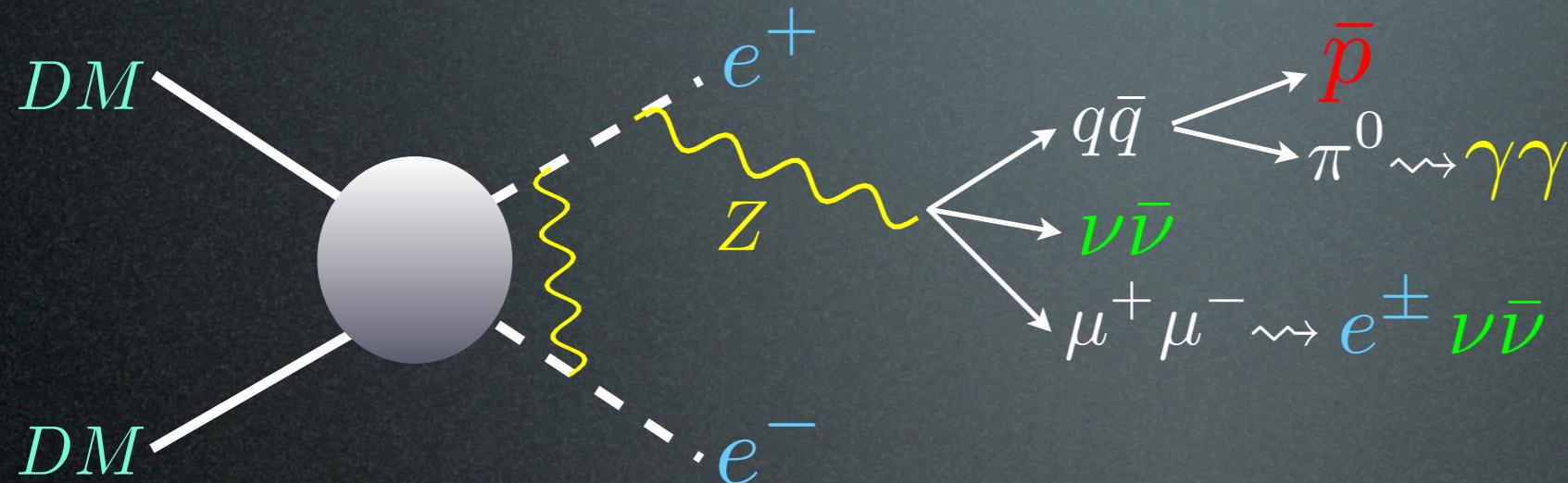
Fluxes at production

ElectroWeak corrections are important!



Fluxes at production

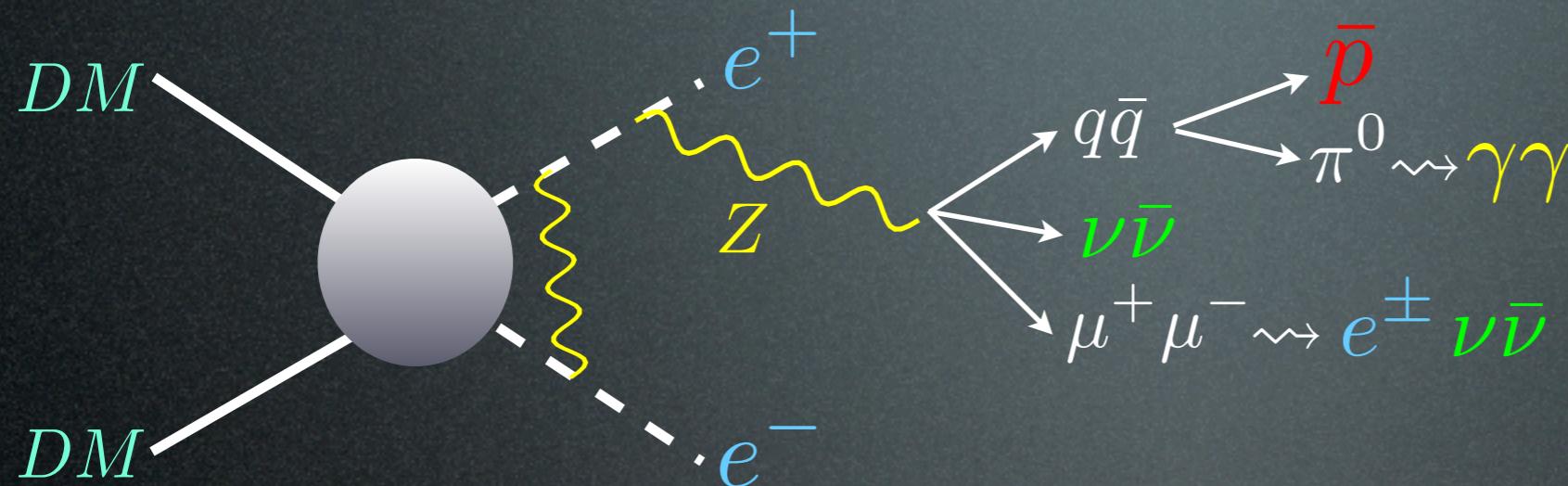
ElectroWeak corrections are important!



$$\frac{\Delta\sigma}{\sigma} \propto \alpha_{\text{weak}} \ln^2 \left(\frac{M_{\text{DM}}^2}{M_Z^2} \right)$$

Fluxes at production

ElectroWeak corrections are important!



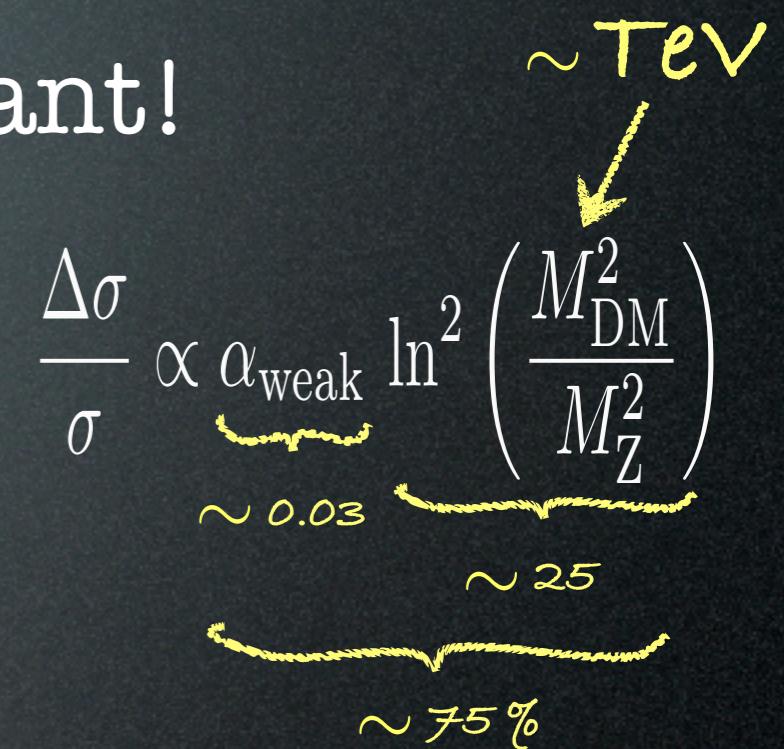
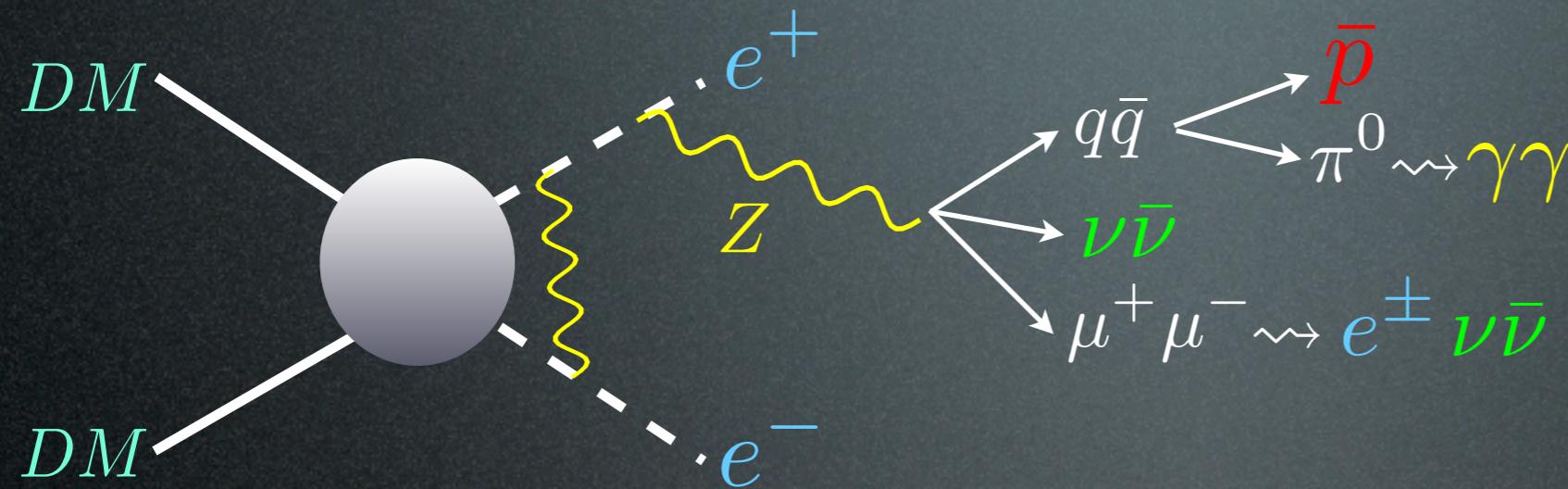
$$\frac{\Delta\sigma}{\sigma} \propto \alpha_{\text{weak}} \ln^2 \left(\frac{M_{\text{DM}}^2}{M_Z^2} \right)$$

Annotations on the right side of the equation:

- $\sim \text{Tev}$ (yellow arrow pointing up)
- ~ 0.03 (yellow wavy line under α_{weak})
- ~ 25 (yellow wavy line under M_Z^2)

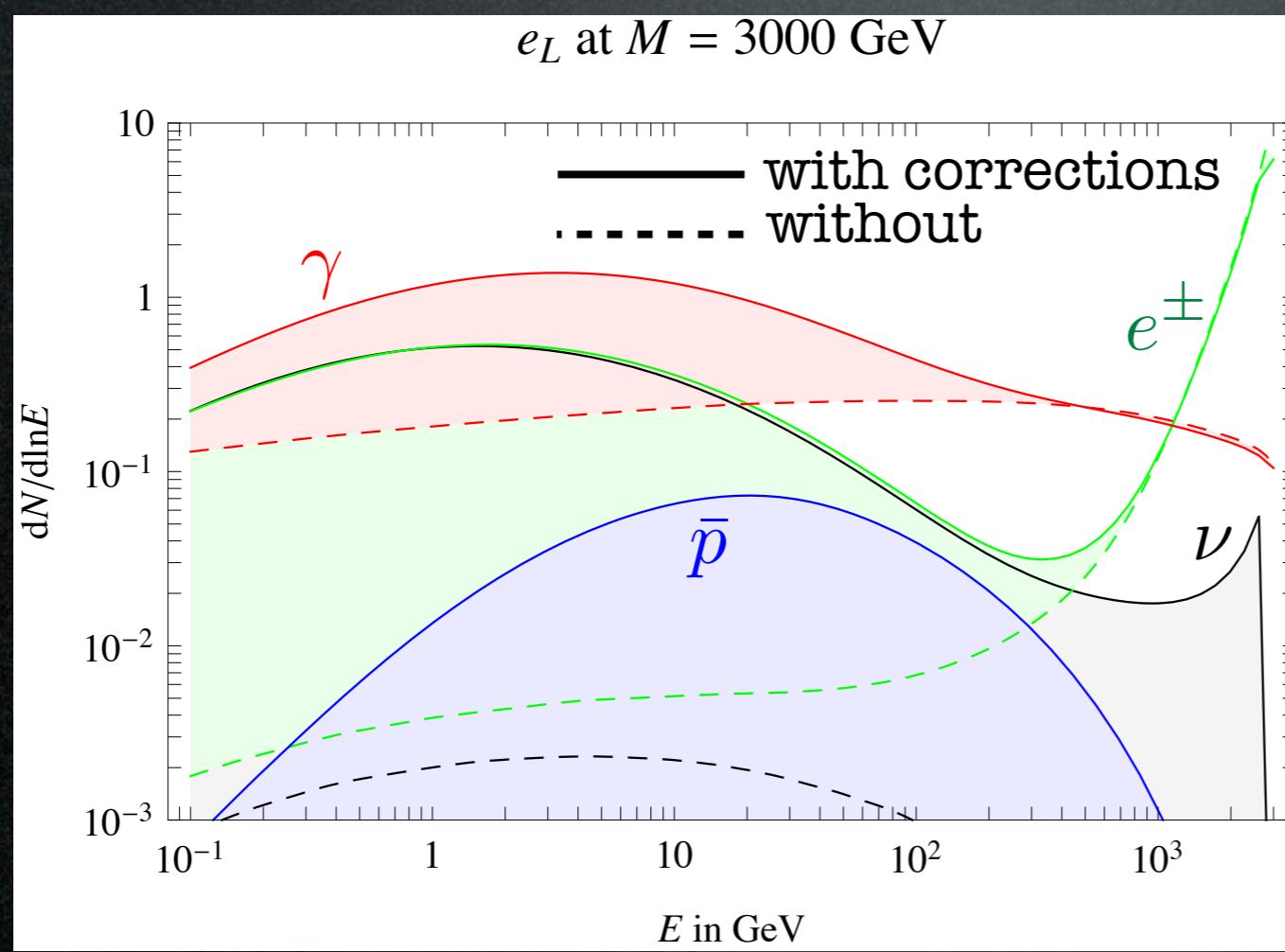
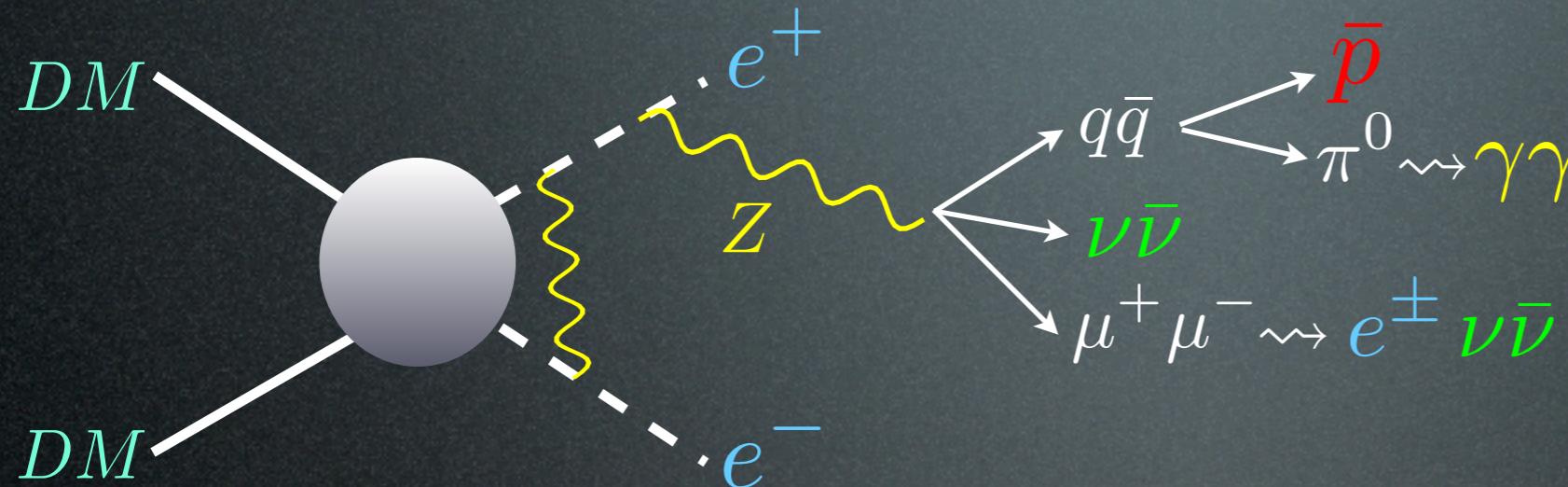
Fluxes at production

ElectroWeak corrections are important!



Fluxes at production

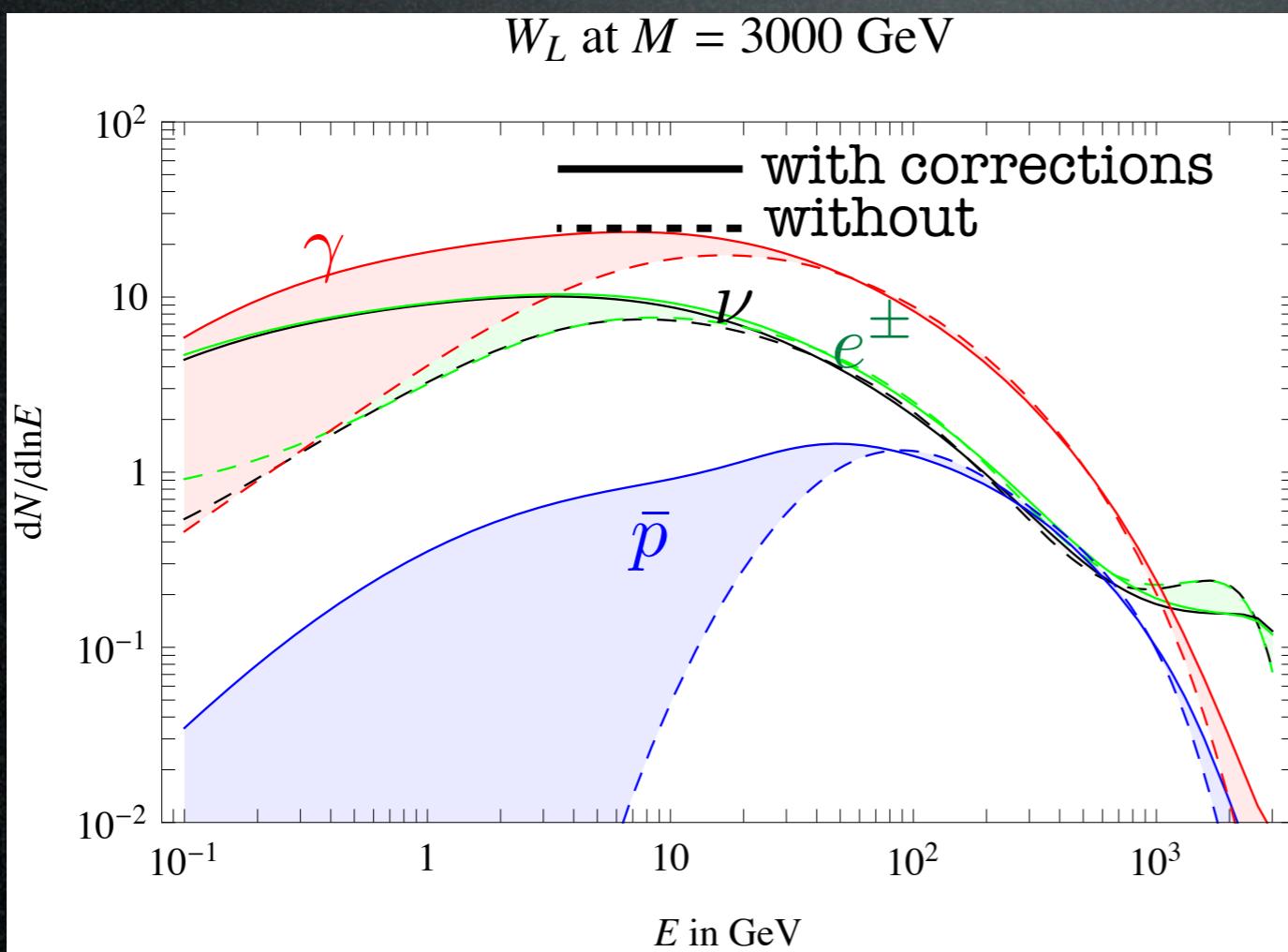
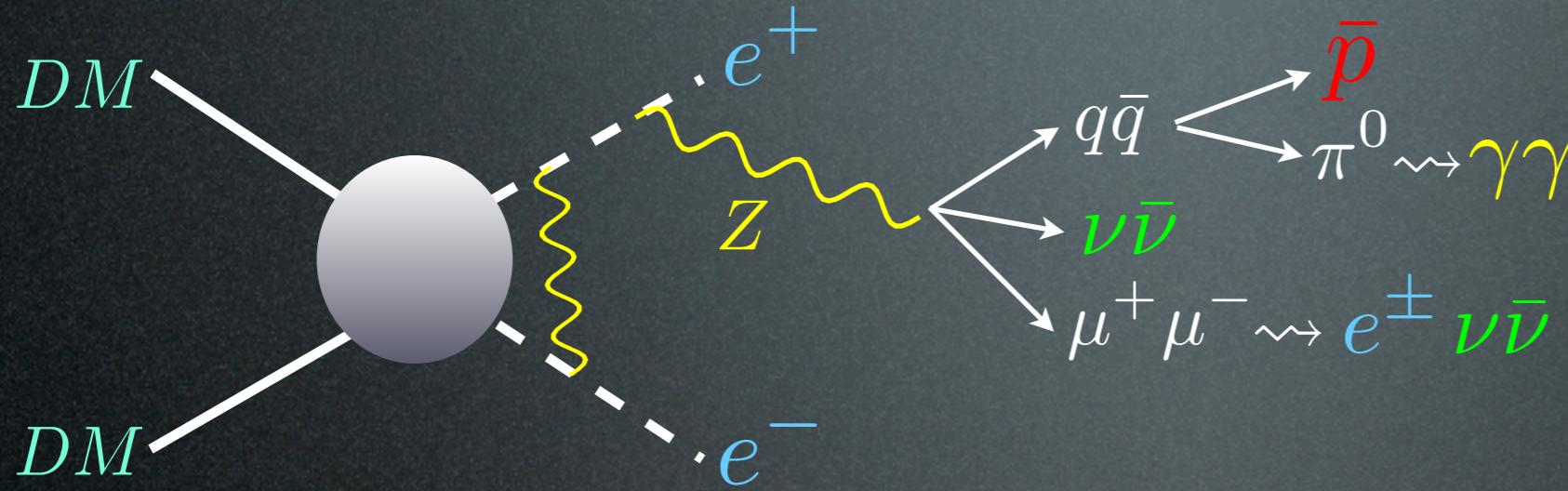
ElectroWeak corrections are important!



- unexpected species
- different spectra
(especially at low energy, but not only)

Fluxes at production

ElectroWeak corrections are important!



- unexpected species
- different spectra
(especially at low energy, but not only)

Fluxes at production

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Cocella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection",
arXiv 1012.4515, JCAP 1103 (2011) 051.
Erratum: JCAP 1210 (2012) E01.

Fluxes at production

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Cocella, A.Hektor, G.Huet
"PPPC 4 DM ID: A Poor Particle Physicist's Guide to Dark Matter Indirect Detection"
arXiv 1012.4515, JCAP 1103 (2011) 053
Erratum: JCAP 1210 (2012) E01.

Fluxes at production:

Complete fluxes at production, including EW corrections as computed in 1009.0224:

Mathematica functions: The file `dINdxEW.m` provides the spectra $\log_{10} [dN/d \log_{10} x]$. The notebook `Sample.nb` shows how to load and use it.

Numerical tables:

Each table provides the spectra $dN/d \log_{10} x$ of stable SM particles (positrons, antiprotons...), normalized per one annihilation.
The columns are: [m_{DM} , $\log_{10} x$, $dN/d \log_{10} x$ for 28 primary channels].

The primary channels are:

$DM \text{ DM} \rightarrow e_L^+ e_L^-$, $e_R^+ e_R^-$, $e^+ e^-$, $\mu_L^+ \mu_L^-$, $\mu_R^+ \mu_R^-$, $\mu^+ \mu^-$, $\tau_L^+ \tau_L^-$, $\tau_R^+ \tau_R^-$, $\tau^+ \tau^-$, $q\bar{q}$, $c\bar{c}$, $b\bar{b}$, $t\bar{t}$, $W_L^+ W_L^-$, $W_T^+ W_T^-$, $Z_L^+ Z_L^-$, $Z_T^+ Z_T^-$, gg , $\gamma\gamma$, hh , $\nu_e \bar{\nu}_e$, $\nu_\mu \bar{\nu}_\mu$, $\nu_\tau \bar{\nu}_\tau$, $VV \rightarrow 4e$, $VV \rightarrow 4\mu$, $VV \rightarrow 4\tau$.

The non-polarized fluxes are just obtained as the appropriate average of the Left and Right or Longitudinal and Transverse ones.
The channel into Higgs bosons assumes a Higgs mass of 125 GeV.

[Positrons](#)

[Antiprotons](#)

[Gamma rays](#)

[Electron Neutrinos](#)

[Muon Neutrinos](#)

[Tau Neutrinos](#)

[Antideuterons](#)

[all of the above](#) All the 7 tables, in a single zipped file.

Fluxes at production without EW corrections (for comparison with previous calculations):

Mathematica functions: The file `dINdxPythia.m` provides the spectra $\log_{10} [dN/d \log_{10} x]$.
The notebook `Sample.nb` shows how to load and use it.

Numerical tables:

Same as above, but here the primary channels are the following 12:

$DM \text{ DM} \rightarrow e^+ e^-$, $\mu^+ \mu^-$, $\tau^+ \tau^-$, $q\bar{q}$, $c\bar{c}$, $b\bar{b}$, $t\bar{t}$, $W^+ W^-$, $Z^+ Z^-$, gg , $\gamma\gamma$, hh .

[Positrons](#)

[Antiprotons](#)

[Gamma rays](#)

[Electron Neutrinos](#)

[Muon Neutrinos](#)

[Tau Neutrinos](#)

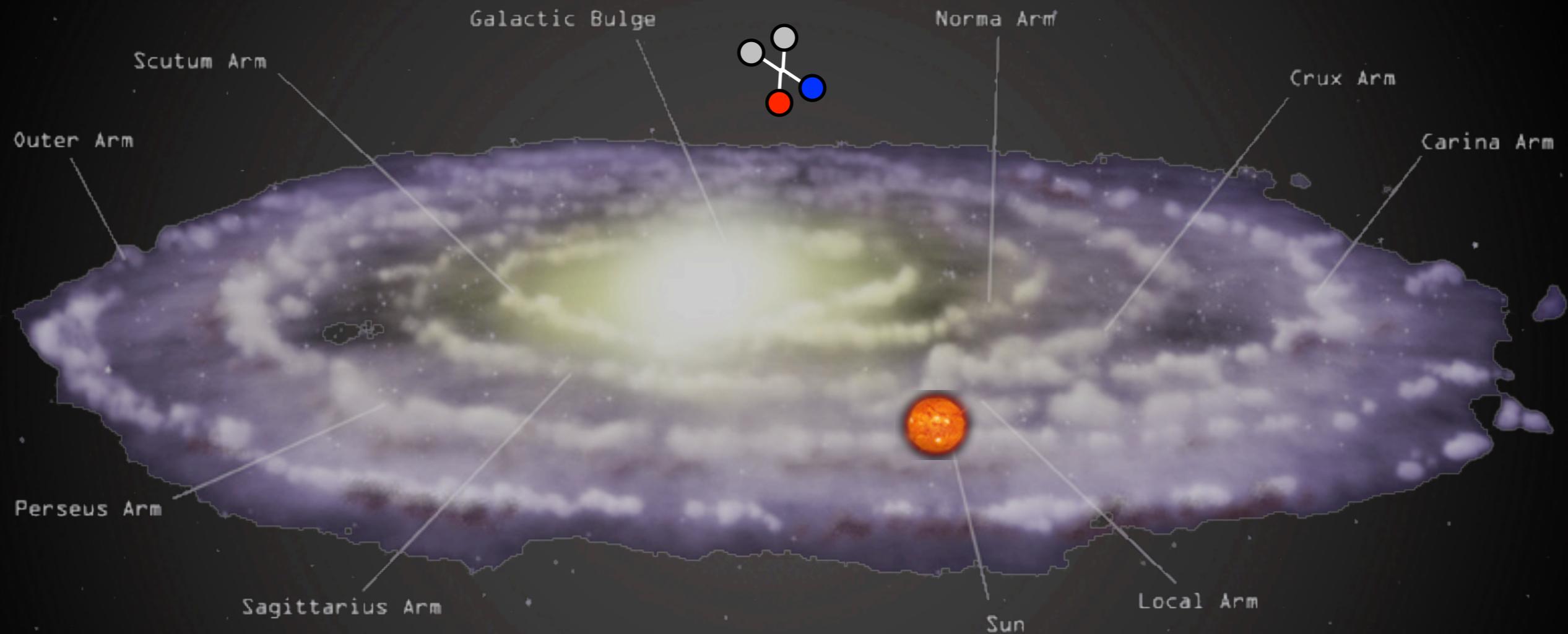
[Antideuterons](#)

[all of the above](#) All the 7 tables, in a single zipped file.

Fluxes at production in models with cascade decays in the hidden sector:

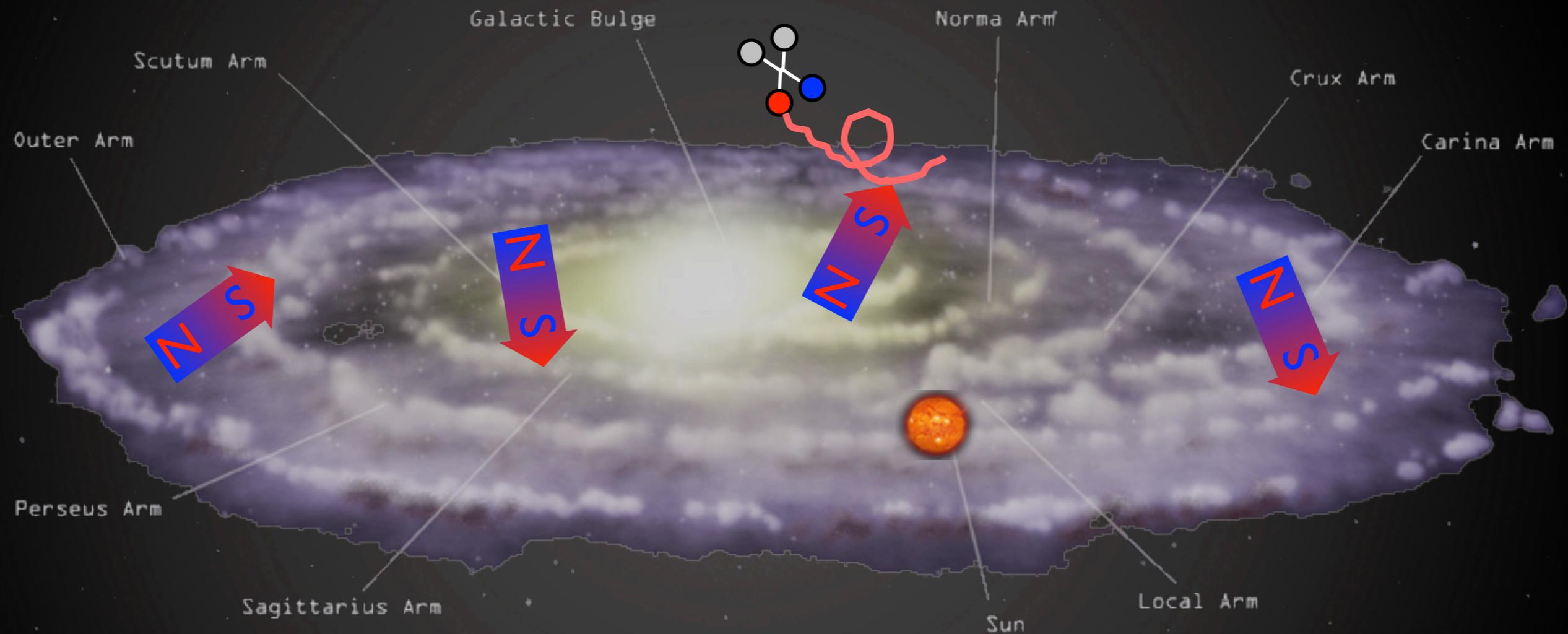
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



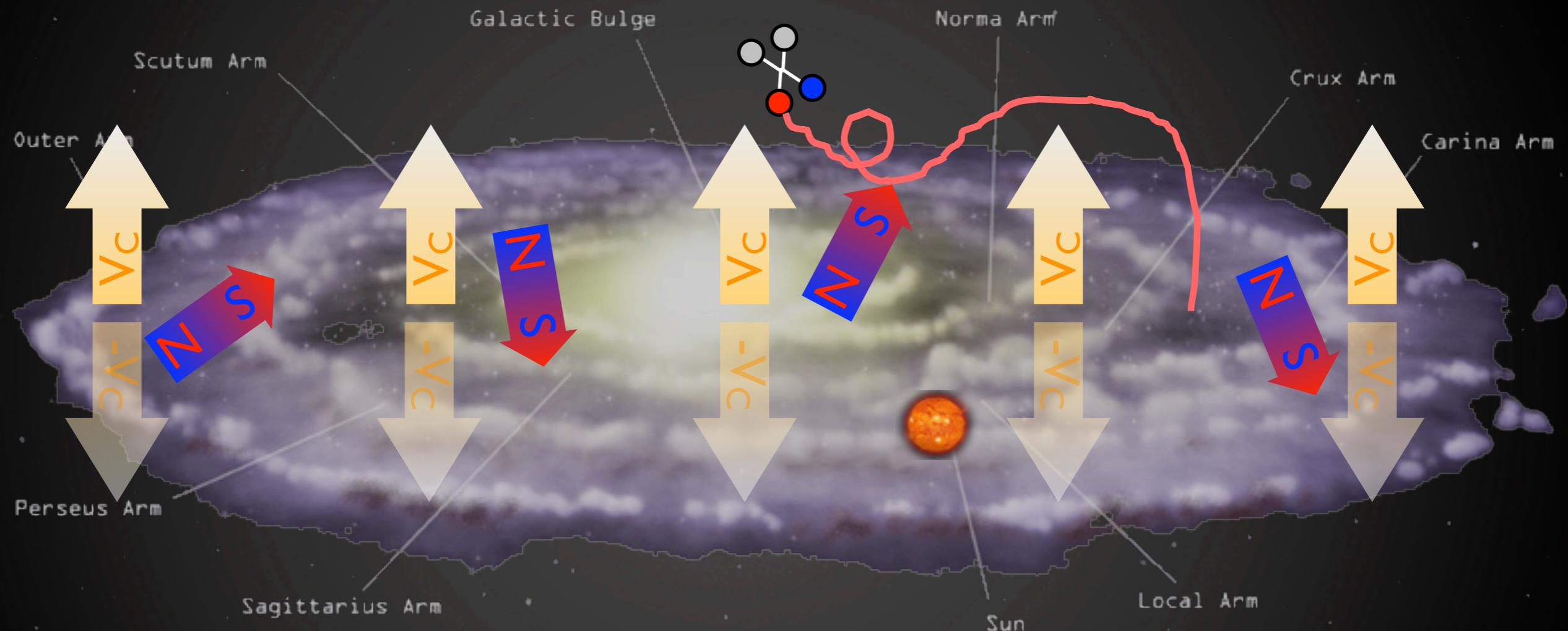
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



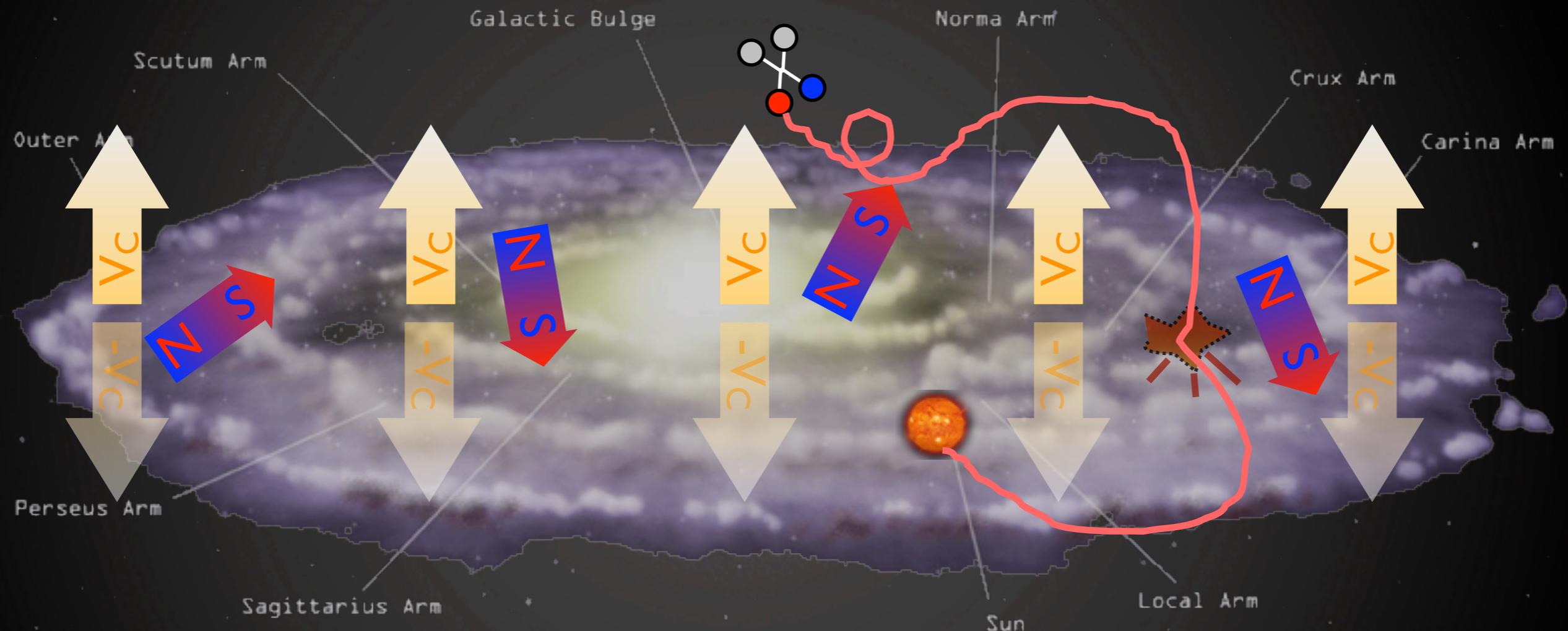
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



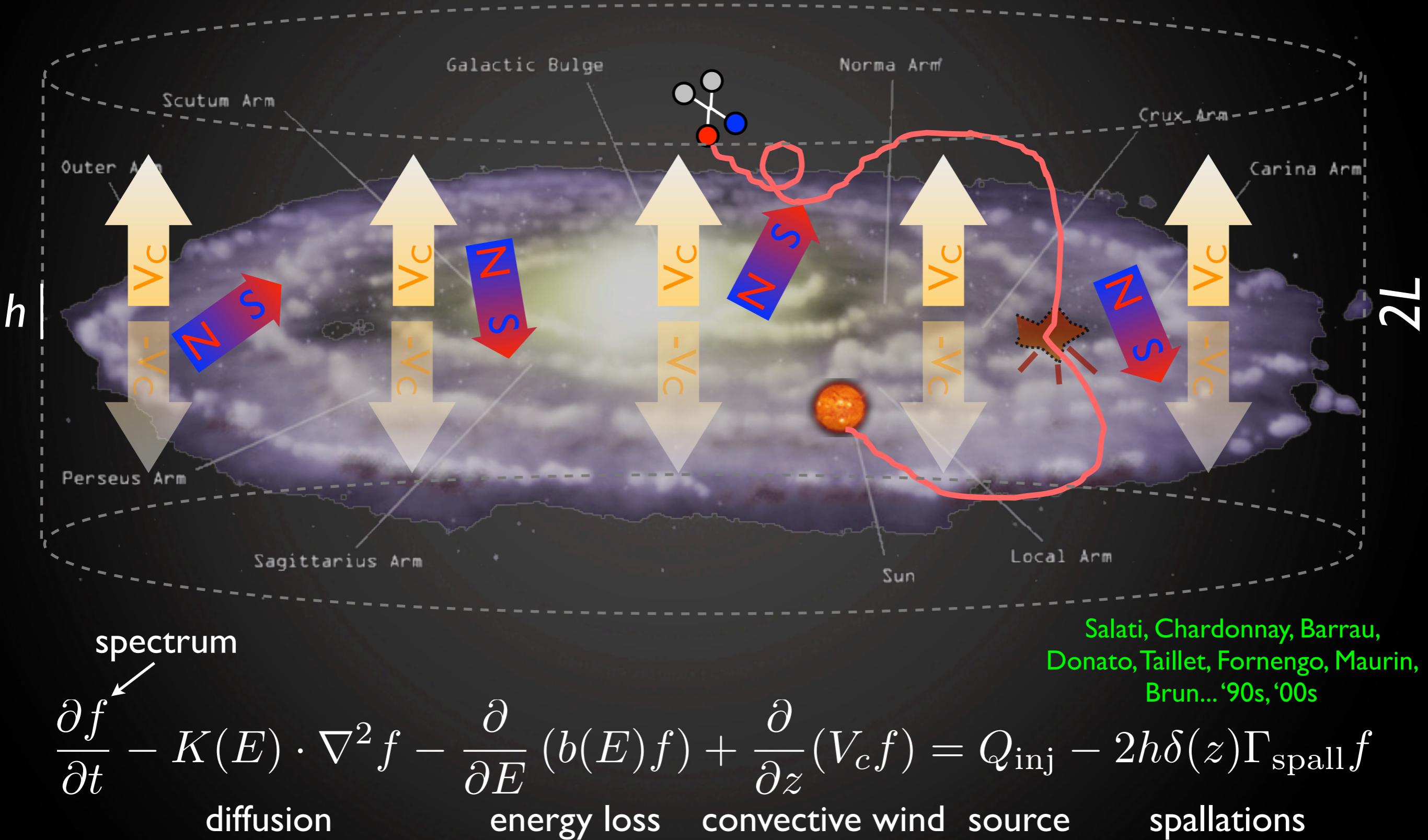
Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



DM halo profiles

From N-body numerical simulations:

$$\text{NFW : } \rho_{\text{NFW}}(r) = \rho_s \frac{r_s}{r} \left(1 + \frac{r}{r_s}\right)^{-2}$$

$$\text{Einasto : } \rho_{\text{Ein}}(r) = \rho_s \exp \left\{ -\frac{2}{\alpha} \left[\left(\frac{r}{r_s}\right)^\alpha - 1 \right] \right\}$$

$$\text{Isothermal : } \rho_{\text{Iso}}(r) = \frac{\rho_s}{1 + (r/r_s)^2}$$

$$\text{Burkert : } \rho_{\text{Bur}}(r) = \frac{\rho_s}{(1 + r/r_s)(1 + (r/r_s)^2)}$$

$$\text{Moore : } \rho_{\text{Moo}}(r) = \rho_s \left(\frac{r_s}{r}\right)^{1.16} \left(1 + \frac{r}{r_s}\right)^{-1.84}$$

At small r: $\rho(r) \propto 1/r^\gamma$

6 profiles:

cuspy: **NFW, Moore**

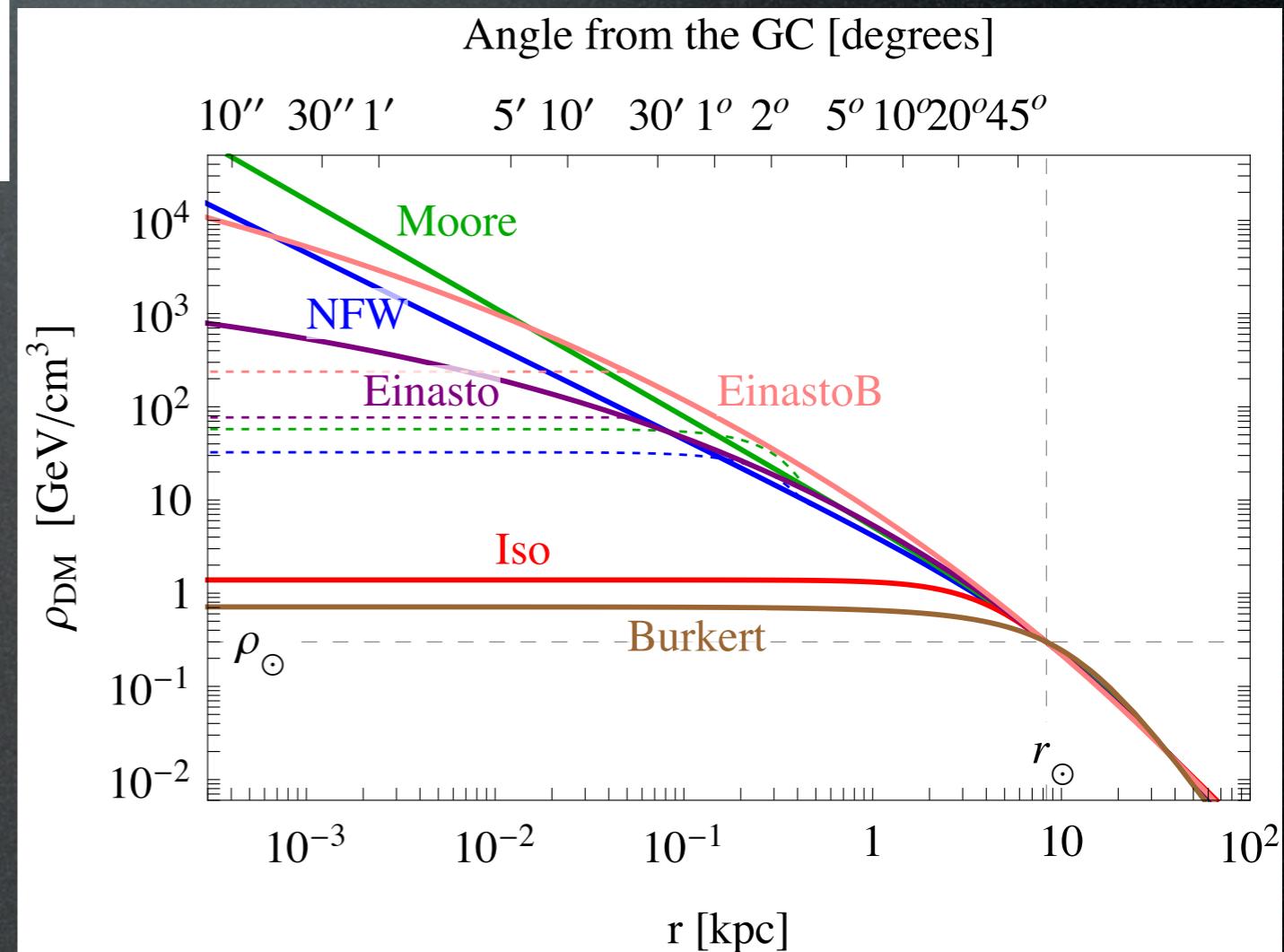
mild: **Einasto**

smooth: **isothermal, Burkert**

EinastoB = steepened Einasto

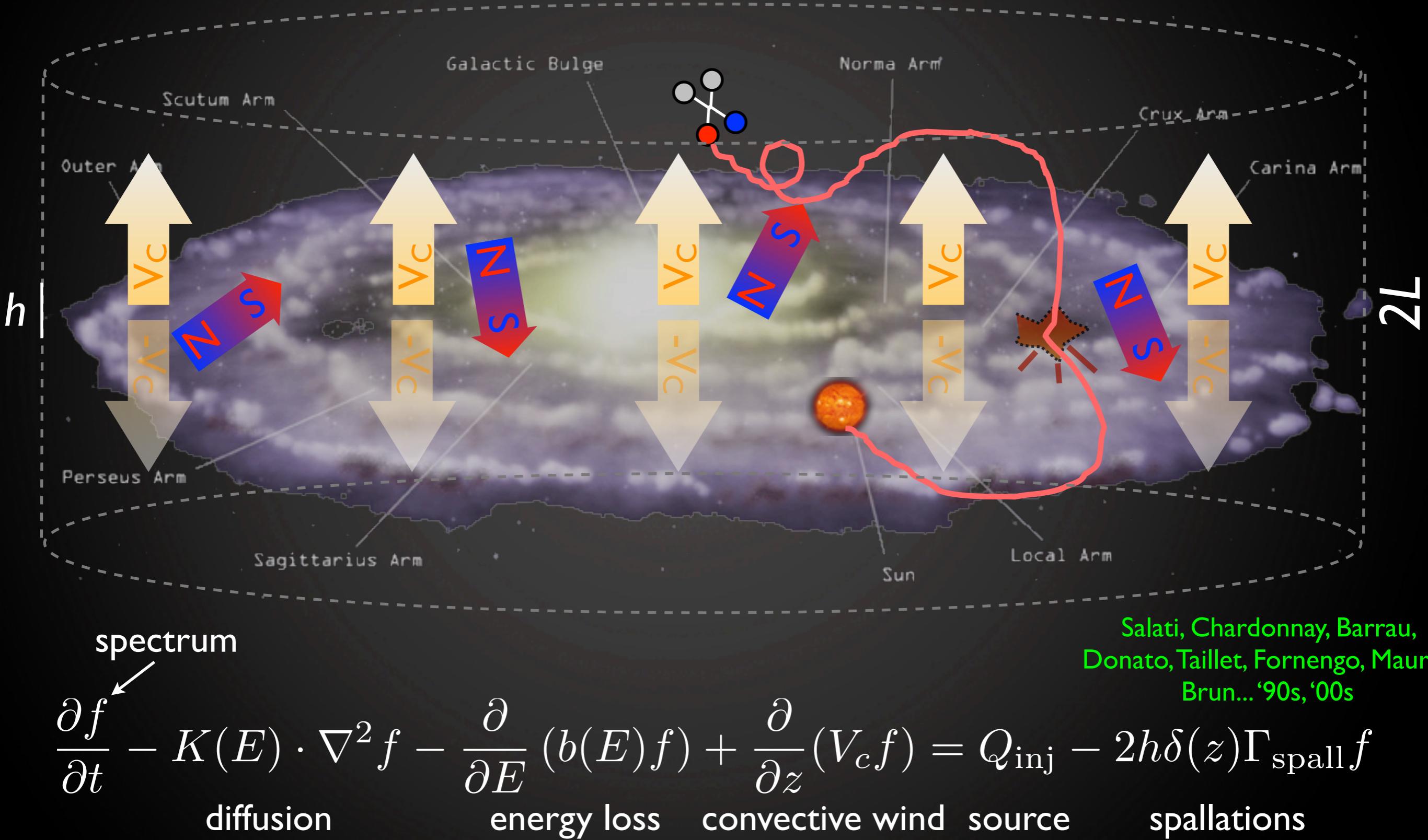
(effect of baryons?)

DM halo	α	r_s [kpc]	ρ_s [GeV/cm ³]
NFW	—	24.42	0.184
Einasto	0.17	28.44	0.033
EinastoB	0.11	35.24	0.021
Isothermal	—	4.38	1.387
Burkert	—	12.67	0.712
Moore	—	30.28	0.105



Indirect Detection: basics

\bar{p} and e^+ from DM annihilations in halo



Propagation

Propagation for antiprotons:

$$\frac{\partial f}{\partial t} - K(T) \cdot \nabla^2 f + \frac{\partial}{\partial z} (\text{sign}(z) f V_{\text{conv}}) = Q - 2h \delta(z) \Gamma_{\text{ann}} f$$

diffusion

$$K(T) = K_0 \beta (p/\text{GeV})^\delta$$

T kinetic energy

convective wind

spallations

Propagation

Propagation for antiprotons:

$$\frac{\partial f}{\partial t} - K(T) \cdot \nabla^2 f + \frac{\partial}{\partial z} (\text{sign}(z) f V_{\text{conv}}) = Q - 2h \delta(z) \Gamma_{\text{ann}} f$$

diffusion

$$K(T) = K_0 \beta (p/\text{GeV})^\delta$$

T kinetic energy

convective wind

spallations

Model	δ	K_0 in kpc^2/Myr	L in kpc	V_{conv} in km/s
min	0.85	0.0016	1	13.5
med	0.70	0.0112	4	12
max	0.46	0.0765	15	5

Propagation

Propagation for antiprotons:

$$\frac{\partial f}{\partial t} - K(T) \cdot \nabla^2 f + \frac{\partial}{\partial z} (\text{sign}(z) f V_{\text{conv}}) = Q - 2h \delta(z) \Gamma_{\text{ann}} f$$

diffusion

convective wind

spallations

$$K(T) = K_0 \beta (p/\text{GeV})^\delta$$

T kinetic energy

Model	δ	K_0 in kpc^2/Myr	L in kpc	V_{conv} in km/s
min	0.85	0.0016	1	13.5
med	0.70	0.0112	4	12
max	0.46	0.0765	15	5

Solution:

$$\Phi_{\bar{p}}(T, \vec{r}_\odot) = B \frac{v_{\bar{p}}}{4\pi} \left(\frac{\rho_\odot}{M_{\text{DM}}} \right)^2 R(T) \sum_k \frac{1}{2} \langle \sigma v \rangle_k \frac{dN_{\bar{p}}^k}{dT}$$

Propagation

Propagation for antiprotons:

$$\frac{\partial f}{\partial t} - K(T) \cdot \nabla^2 f + \frac{\partial}{\partial z} (\text{sign}(z) f V_{\text{conv}}) = Q - 2h \delta(z) \Gamma_{\text{ann}} f$$

diffusion

$$K(T) = K_0 \beta (p/\text{GeV})^\delta$$

T kinetic energy

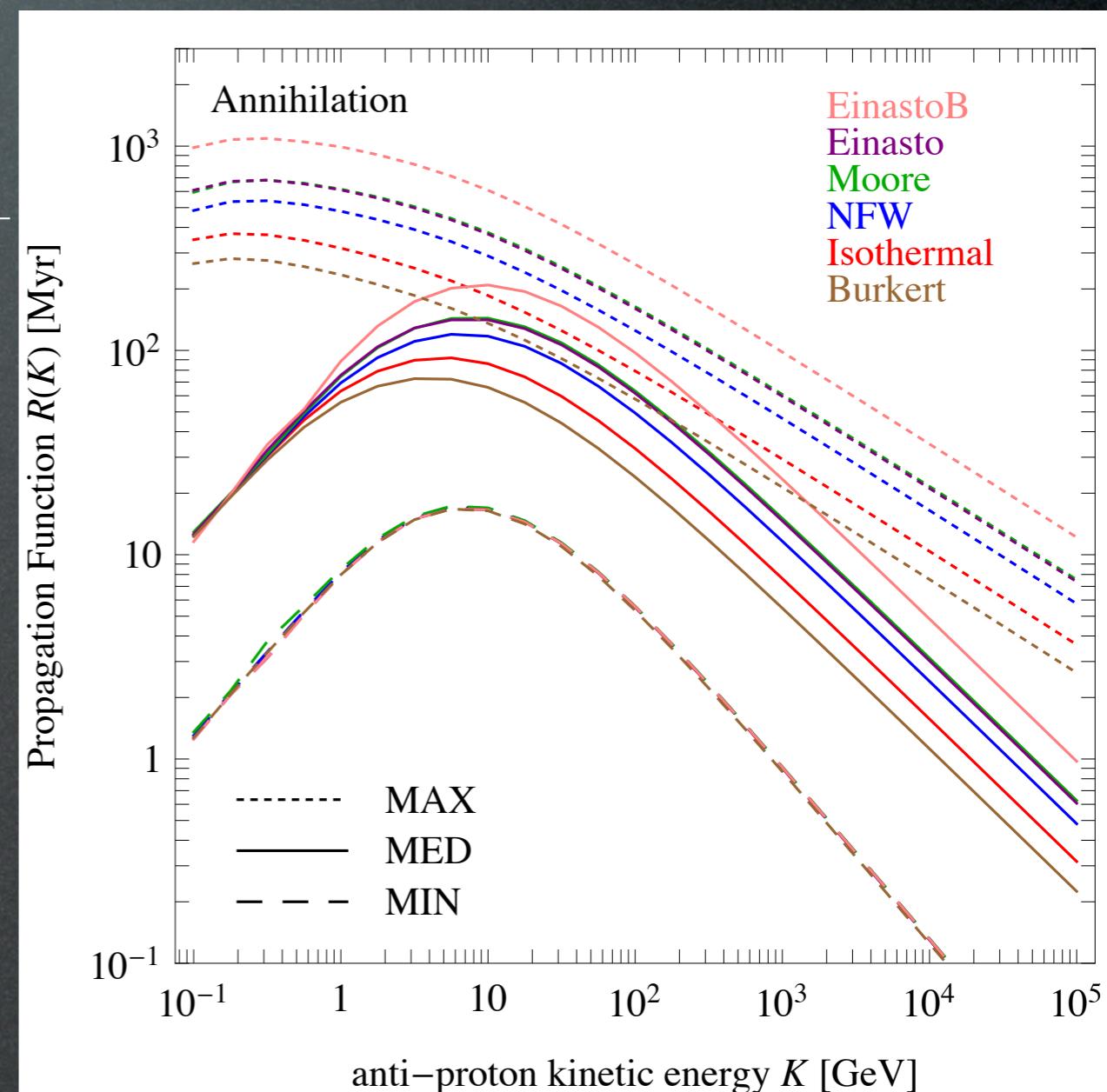
convective wind

spallations

Model	δ	K_0 in kpc^2/Myr	L in kpc	V_{conv} in km/s
min	0.85	0.0016	1	13.5
med	0.70	0.0112	4	12
max	0.46	0.0765	15	5

Solution:

$$\Phi_{\bar{p}}(T, \vec{r}_\odot) = B \frac{v_{\bar{p}}}{4\pi} \left(\frac{\rho_\odot}{M_{\text{DM}}} \right)^2 R(T) \sum_k \frac{1}{2} \langle \sigma v \rangle_k \frac{dN_{\bar{p}}^k}{dT}$$



Propagation

Propagation for antiprotons:

$$\frac{\partial f}{\partial t} - K(T) \cdot \nabla^2 f + \frac{\partial}{\partial z} (\text{sign}(z) f V_{\text{conv}}) = Q - 2h \delta(z) \Gamma_{\text{ann}} f$$

diffusion

$$K(T) = K_0 \beta (p/\text{GeV})^\delta$$

T kinetic energy

convective wind

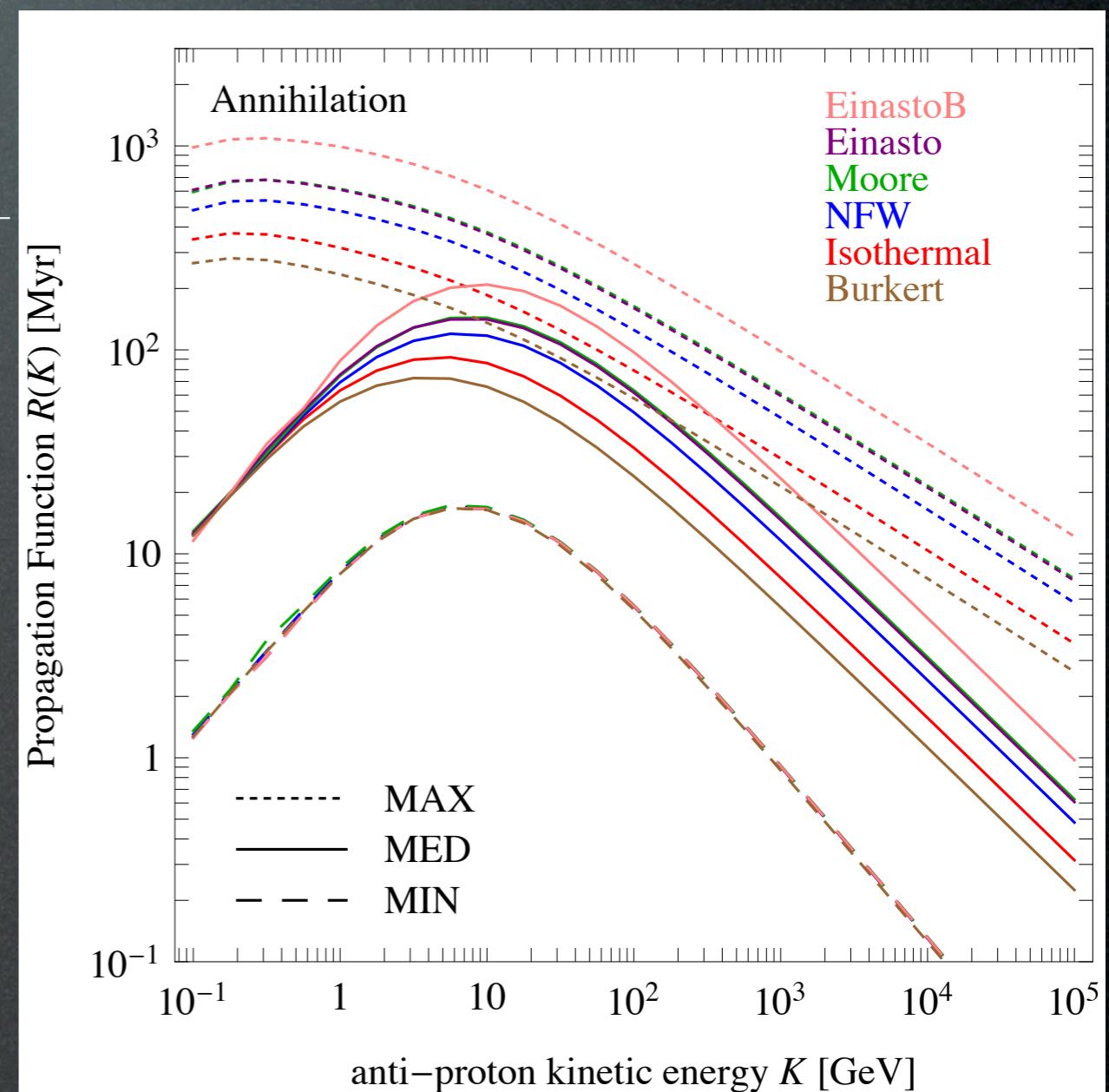
spallations

Model	δ	K_0 in kpc^2/Myr	L in kpc	V_{conv} in km/s
min	0.85	0.0016	1	13.5
med	0.70	0.0112	4	12
max	0.46	0.0765	15	5

Solution:

$$\Phi_{\bar{p}}(T, \vec{r}_\odot) = B \frac{v_{\bar{p}}}{4\pi} \left(\frac{\rho_\odot}{M_{\text{DM}}} \right)^2 R(T) \sum_k \frac{1}{2} \langle \sigma v \rangle_k \frac{dN_{\bar{p}}^k}{dT}$$

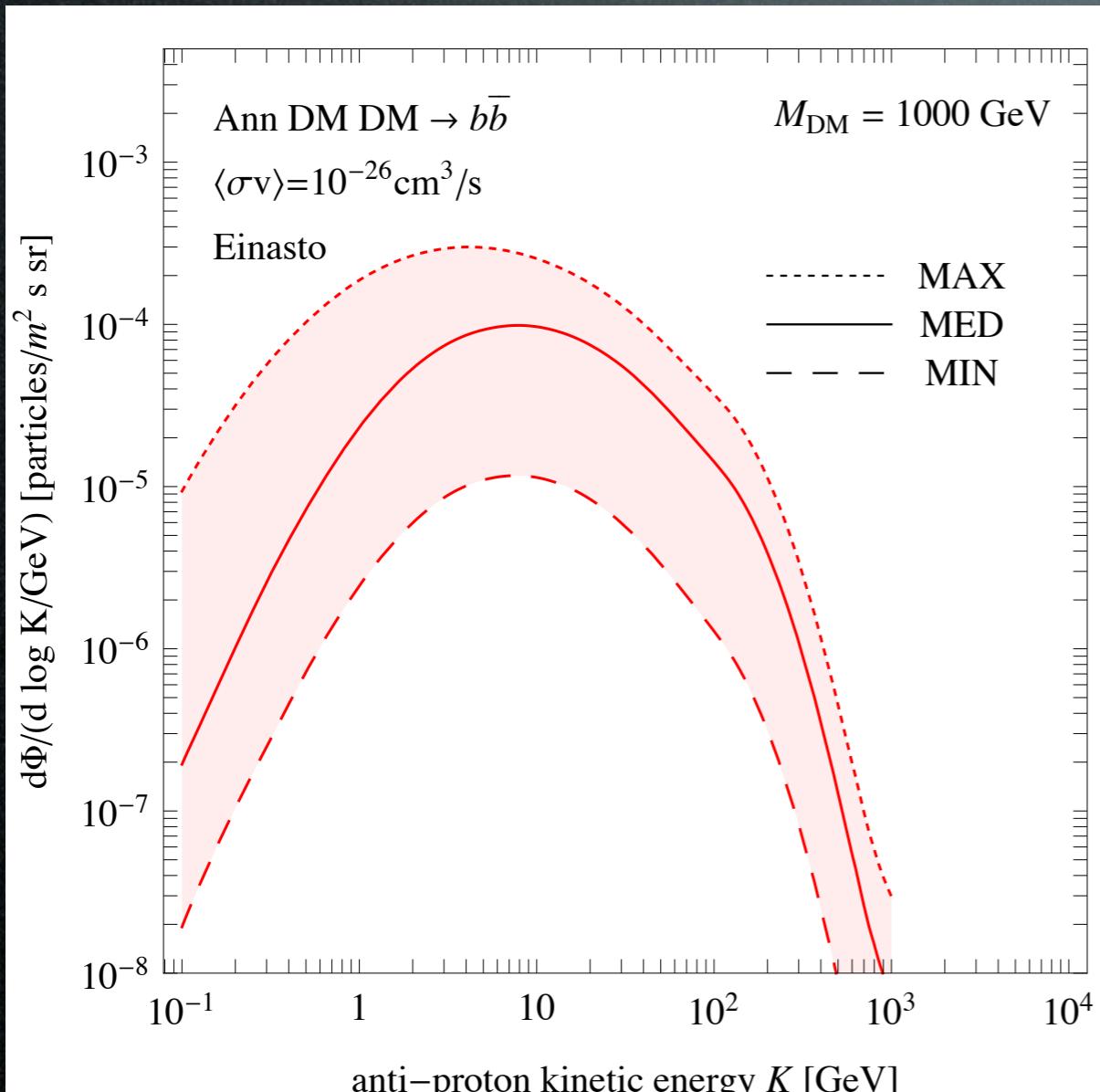
Improvement: added Energy Losses & Diffusive Reacceleration



Propagated fluxes

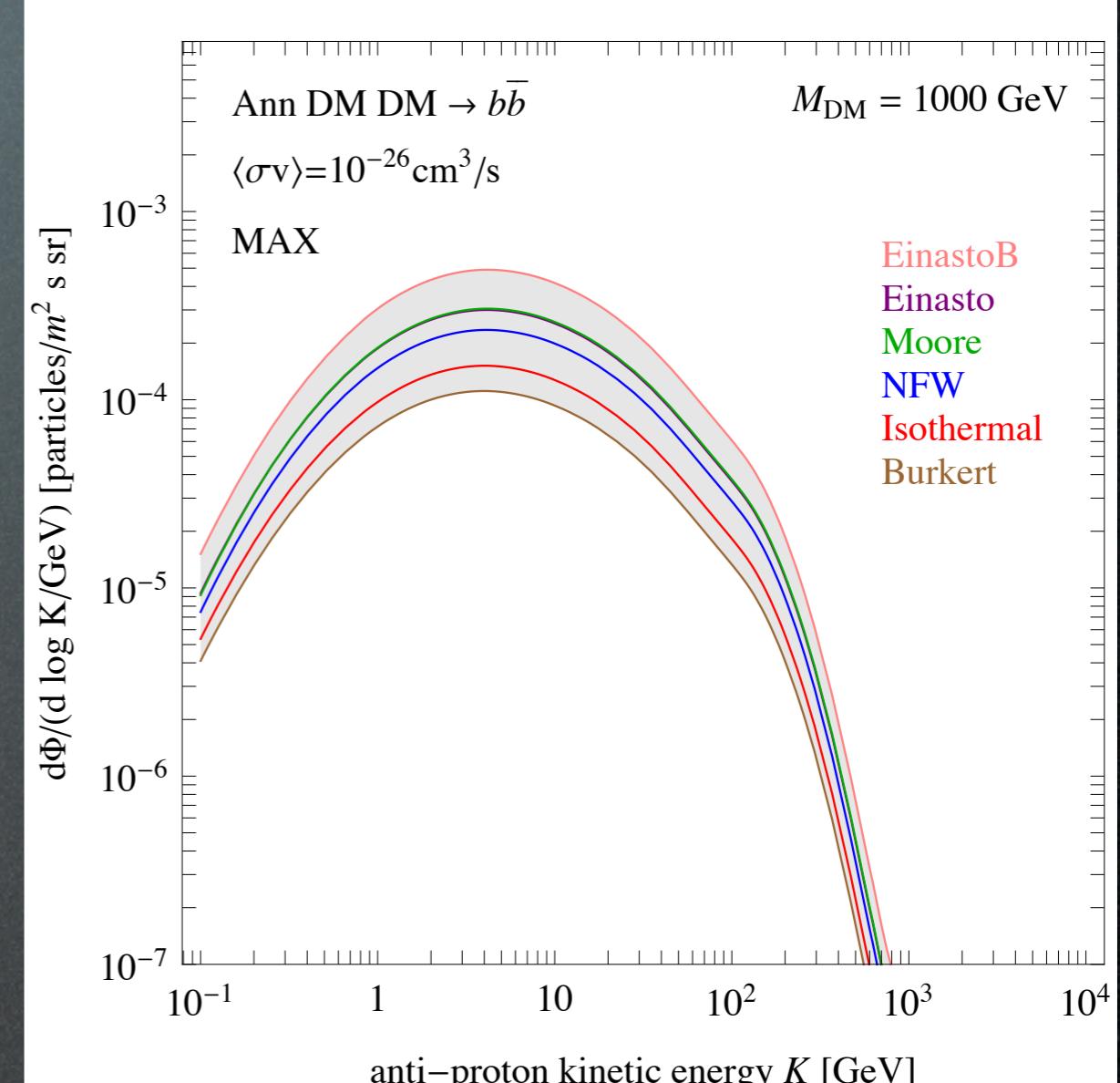
Antiprotons

Varying prop parameters



Almost 2 orders of magnitude

Varying halo profile



Almost 1 order of magnitude

Bottom line: Antiprotons are quite affected by propagation, but spectral shape somewhat preserved

Propagated fluxes

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M. Cirelli,
"PPPC 4
arXiv 10
Erratum:

Propagation functions for charged cosmic rays at the location of the Earth:

Annihilation

Positrons: The file [ElectronHaloFunctEarthAnn.m](#) provides the halo functions $I(x, E_3 r_{\text{Earth}})$ at the location of the Earth.
The notebook [Sample.nb](#) shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ [in the approximated formalism - see paper].

Antiprotons: [Table](#) of fit coefficients for the propagation function $R(T)$.

Antideuterons: [Table](#) of fit coefficients for the propagation function $R(T)$.

Decay

Positrons: The file [ElectronHaloFunctEarthDec.m](#) provides the halo functions $I(x, E_3 r_{\text{Earth}})$ at the location of the Earth.
The notebook [Sample.nb](#) shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ [in the approximated formalism - see paper].

Antiprotons: [Table](#) of fit coefficients for the propagation function $R(T)$.

Antideuterons: [Table](#) of fit coefficients for the propagation function $R(T)$.

Fluxes of charged cosmic rays at the Earth, after propagation:

Annihilation

Antiprotons: Mathematica function: the file [ProtonFluxELDRAnn.m](#) provides the spectra $\log_{10} [d\Phi/d \log_{10} K]$.
Refer to the notebook [Sample.nb](#) for usage.

Numerical [table](#): provides the spectra $d\Phi/d \log_{10} K$.
The columns are: [m_{DM} , halo, propagation, $\log_{10} K$, $d\Phi/d \log_{10} K$ for 28 primary channels].
The spectra are computed at a benchmark annihilation cross of $\langle \sigma v \rangle = 10^{-26} \text{ cm}^3/\text{s}$ and renormalized multiplying by $(m_{\text{DM}}/\text{GeV})^2$.
Units are particles/ $(\text{m}^2 \text{ s} \text{ sr})$.

The file [ProtonFluxAnn.m](#) provides the spectra without ELDR (Energy Losses including tertiaries and Diffusive Reacceleration, see [1412.5696](#) for details); this is now superseded and provided here only for the purpose of comparison with previous calculations.

Decay

Antiprotons: Mathematica function: the file [ProtonFluxELDRDec.m](#) provides the spectra $\log_{10} [d\Phi/d \log_{10} K]$.
Refer to the notebook [Sample.nb](#) for usage.

Numerical [table](#): provides the spectra $d\Phi/d \log_{10} K$.
The columns are: [m_{DM} , halo, propagation, $\log_{10} K$, $d\Phi/d \log_{10} K$ for 28 primary channels].
The spectra are computed at a benchmark decay life of $\tau = 10^{26} \text{ s}$ and renormalized multiplying by $[m_{\text{DM}}/\text{GeV}]$.
Units are particles/ $(\text{m}^2 \text{ s} \text{ sr})$.

The file [ProtonFluxDec.m](#) provides the spectra without ELDR (Energy Losses including tertiaries and Diffusive Reacceleration, see [1412.5696](#) for details); this is now superseded and provided here only for the purpose of comparison with previous calculations.

Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion energy loss (assumed space indep.)
(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.) $b(E) = (E/\text{GeV})^2 / \tau_E$
 $K(E) = K_0(E/\text{GeV})^\delta$ $\tau_E = 10^{16} \text{ s}$

$$Q = \frac{1}{2} \left(\frac{\rho}{M_{\text{DM}}} \right)^2 f_{\text{inj}}, \quad f_{\text{inj}} = \sum_k \langle \sigma v \rangle_k \frac{dN_{e^+}^k}{dE}$$

Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion
(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

energy loss (assumed space indep.)
 $b(E) = (E/\text{GeV})^2 / \tau_E$
 $K(E) = K_0 (E/\text{GeV})^\delta$ $\tau_E = 10^{16} \text{ s}$

$$Q = \frac{1}{2} \left(\frac{\rho}{M_{\text{DM}}} \right)^2 f_{\text{inj}}, \quad f_{\text{inj}} = \sum_k \langle \sigma v \rangle_k \frac{dN_{e^+}^k}{dE}$$

Model	δ	K_0 in kpc^2/Myr	L in kpc
min (M2)	0.55	0.00595	1
med	0.70	0.0112	4
max (M1)	0.46	0.0765	15

Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion
(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

energy loss (assumed space indep.)
 $b(E) = (E/\text{GeV})^2 / \tau_E$
 $\tau_E = 10^{16} \text{ s}$

$K(E) = K_0 (E/\text{GeV})^\delta$

$$Q = \frac{1}{2} \left(\frac{\rho}{M_{\text{DM}}} \right)^2 f_{\text{inj}}, \quad f_{\text{inj}} = \sum_k \langle \sigma v \rangle_k \frac{dN_{e^+}^k}{dE}$$

Model	δ	K_0 in kpc^2/Myr	L in kpc
min (M2)	0.55	0.00595	1
med	0.70	0.0112	4
max (M1)	0.46	0.0765	15

Solution:

$$\Phi_{e^+}(E, \vec{r}_\odot) = B \frac{v_{e^+} \tau_E}{4\pi} \int_E^{M_{\text{DM}}} dE' Q(E') \cdot I(\lambda_D(E, E'))$$

$$\lambda_D^2 = 4K_0 \tau_E \left[\frac{(E/\text{GeV})^{\delta-1} - (E'/\text{GeV})^{\delta-1}}{\delta - 1} \right]$$

Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion

(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

$$K(E) = K_0(E/\text{GeV})^\delta$$

energy loss (assumed space indep.)

$$b(E) = (E/\text{GeV})^2/\tau_E$$

$$\tau_E = 10^{16} \text{ s}$$

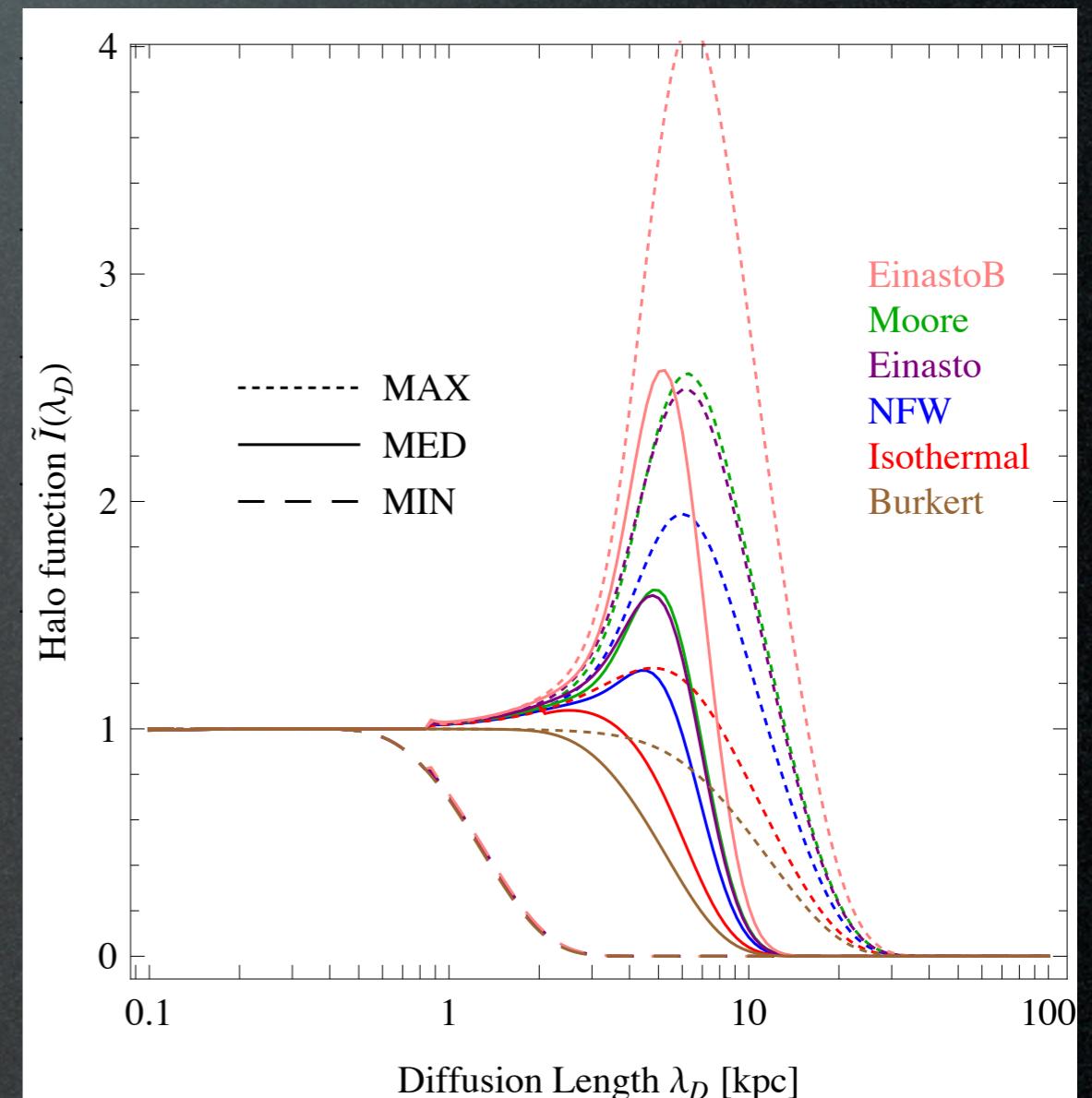
$$Q = \frac{1}{2} \left(\frac{\rho}{M_{\text{DM}}} \right)^2 f_{\text{inj}}, \quad f_{\text{inj}} = \sum_k \langle \sigma v \rangle \frac{dN_{e^+}^k}{dE}$$

Model	δ	K_0 in kpc^2/Myr	L in kpc
min (M2)	0.55	0.00595	1
med	0.70	0.0112	4
max (M1)	0.46	0.0765	15

Solution:

$$\Phi_{e^+}(E, \vec{r}_\odot) = B \frac{v_{e^+} \tau_E}{4\pi} \int_E^{M_{\text{DM}}} dE' Q(E') \cdot I(\lambda_D(E, E'))$$

$$\lambda_D^2 = 4K_0\tau_E \left[\frac{(E/\text{GeV})^{\delta-1} - (E'/\text{GeV})^{\delta-1}}{\delta-1} \right]$$



Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion

(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

$$K(E) = K_0(E/\text{GeV})^\delta$$

energy loss (assumed space indep.)

$$b(E) = (E/\text{GeV})^2 / \tau_E$$

$$\tau_E = 10^{16} \text{ s}$$

!

Propagation

Propagation for positrons: conventional treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) = Q$$

diffusion

(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

$$K(E) = K_0(E/\text{GeV})^\delta$$

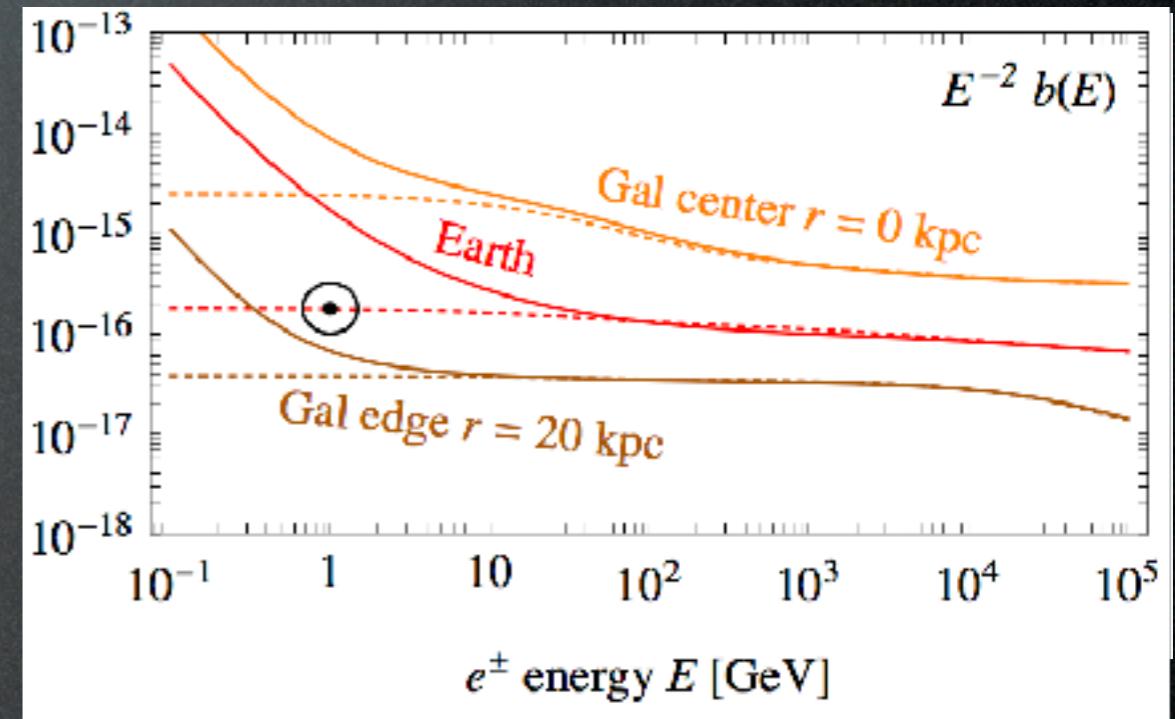
energy loss (assumed space indep.)

$$b(E) = (E/\text{GeV})^2 / \tau_E$$

$$\tau_E = 10^{16} \text{ s}$$

!

- it is **not** space independent
- it is **not** simple E^2



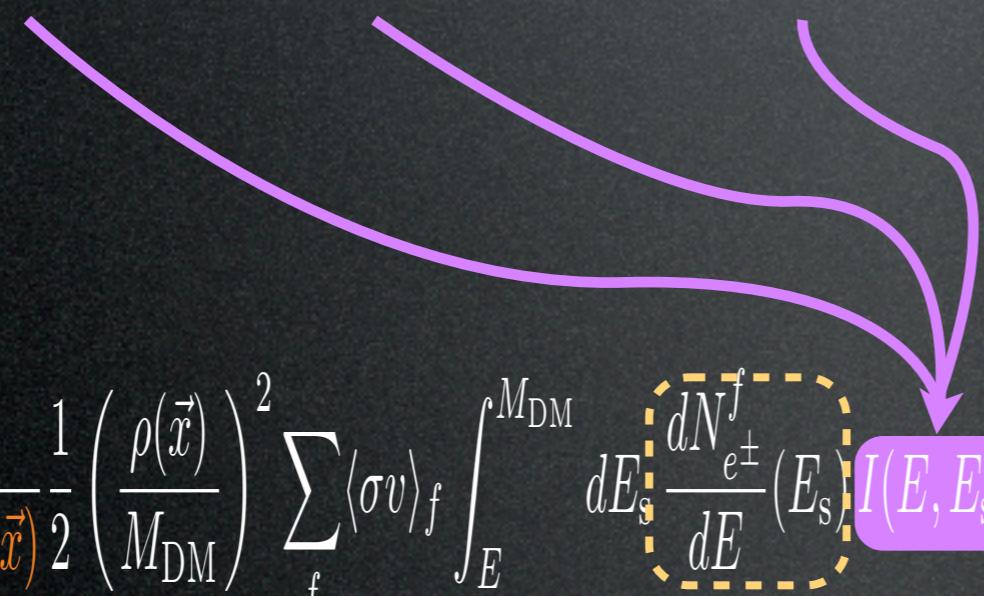
Propagation

Propagation for positrons: improved treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (\mathbf{b}(E, \vec{x}) f) = Q$$

diffusion
 (in turbulent $\bar{B} \approx \mu\text{G}$,
 assumed space indep.) energy loss function
 $K(E) = K_0(E/\text{GeV})^\delta$
 extracted using GalProp simulations

Model	δ	K_0 in kpc^2/Myr	L in kpc
min (M2)	0.55	0.00595	1
med	0.70	0.0112	4
max (M1)	0.46	0.0765	15



Solution:

$$\Phi_{e^\pm}(E, \vec{r}_\odot) = \frac{v_{e^\pm}}{4\pi b(E, \vec{x})} \frac{1}{2} \left(\frac{\rho(\vec{x})}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \int_E^{M_{\text{DM}}} dE_s \frac{dN_{e^\pm}^f}{dE}(E_s) I(E, E_s, \vec{x})$$

Propagation

Propagation for positrons: improved treatment

$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (\mathbf{b}(E, \vec{x}) f) = Q$$

diffusion

(in turbulent $\bar{B} \approx \mu\text{G}$,
assumed space indep.)

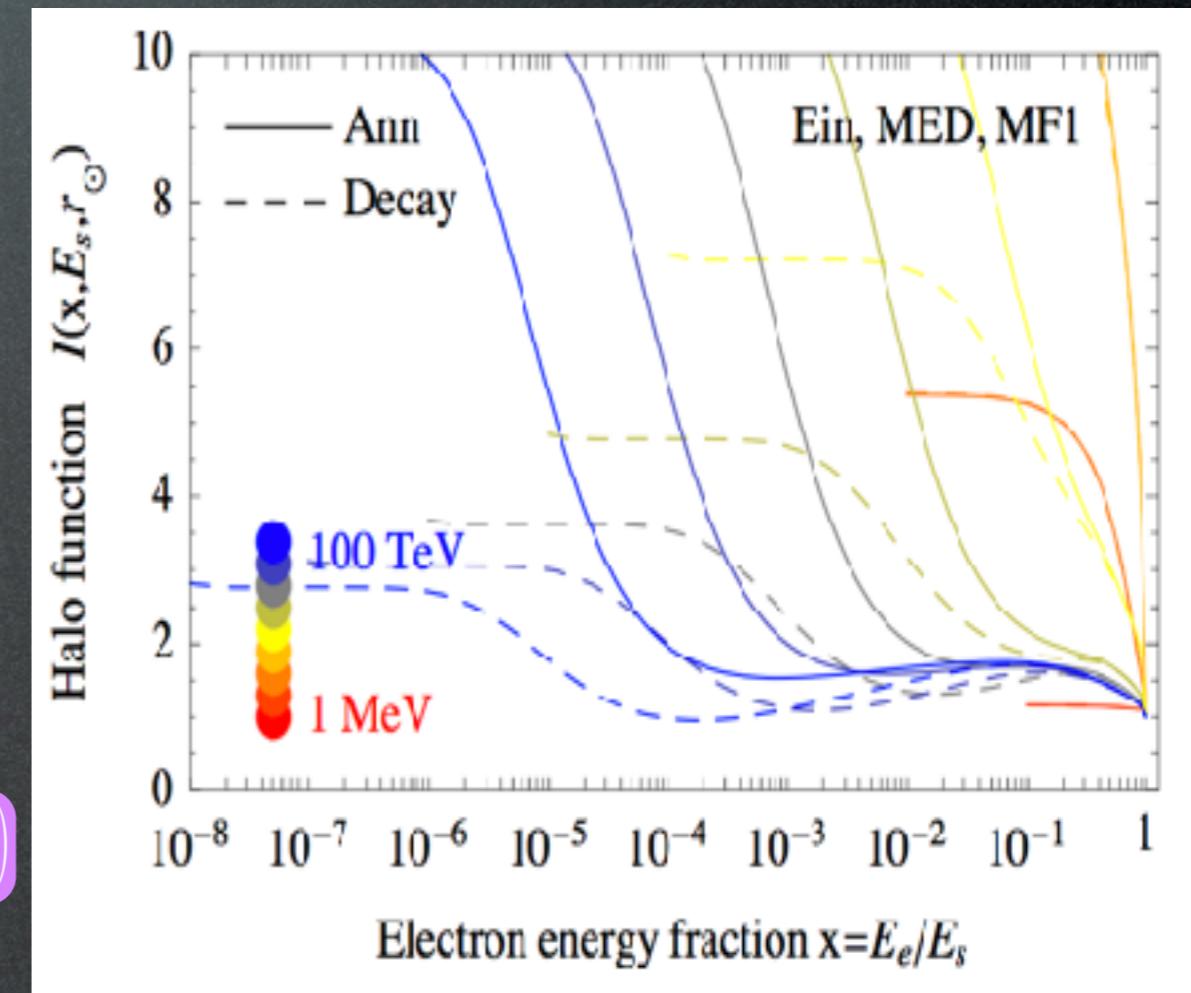
$$K(E) = K_0(E/\text{GeV})^\delta$$

energy loss function
extracted using GalProp simulations

Model	δ	K_0 in kpc^2/Myr	L in kpc
min (M2)	0.55	0.00595	1
med	0.70	0.0112	4
max (M1)	0.46	0.0765	15

Solution:

$$\Phi_{e^\pm}(E, \vec{r}_\odot) = \frac{v_{e^\pm}}{4\pi b(E, \vec{x})} \frac{1}{2} \left(\frac{\rho(\vec{x})}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \int_E^{M_{\text{DM}}} dE_s \frac{dN_{e^\pm}^f}{dE}(E_s) I(E, E_s, \vec{x})$$

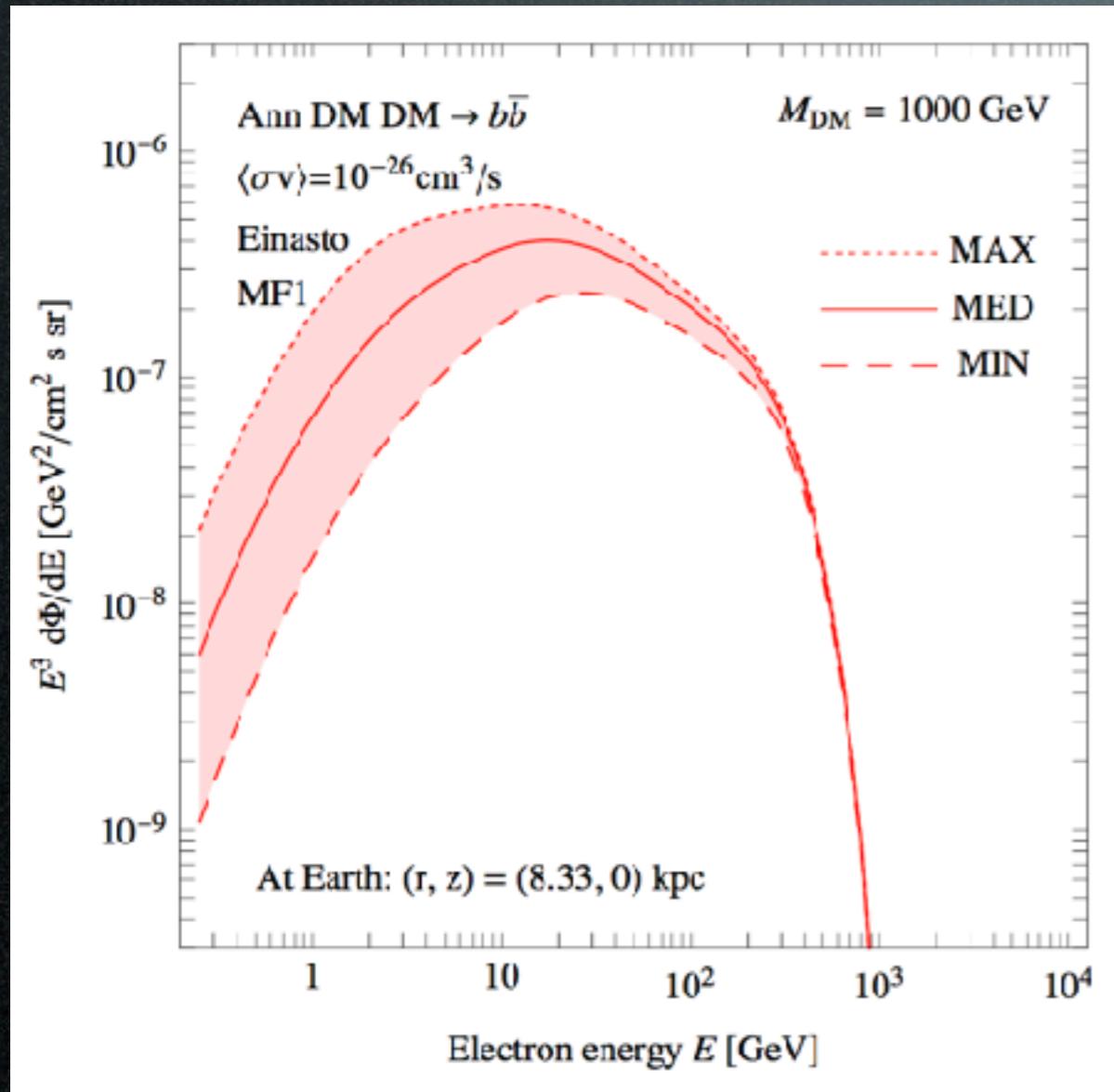


There is now a ‘halo function’ per each positron injection energy

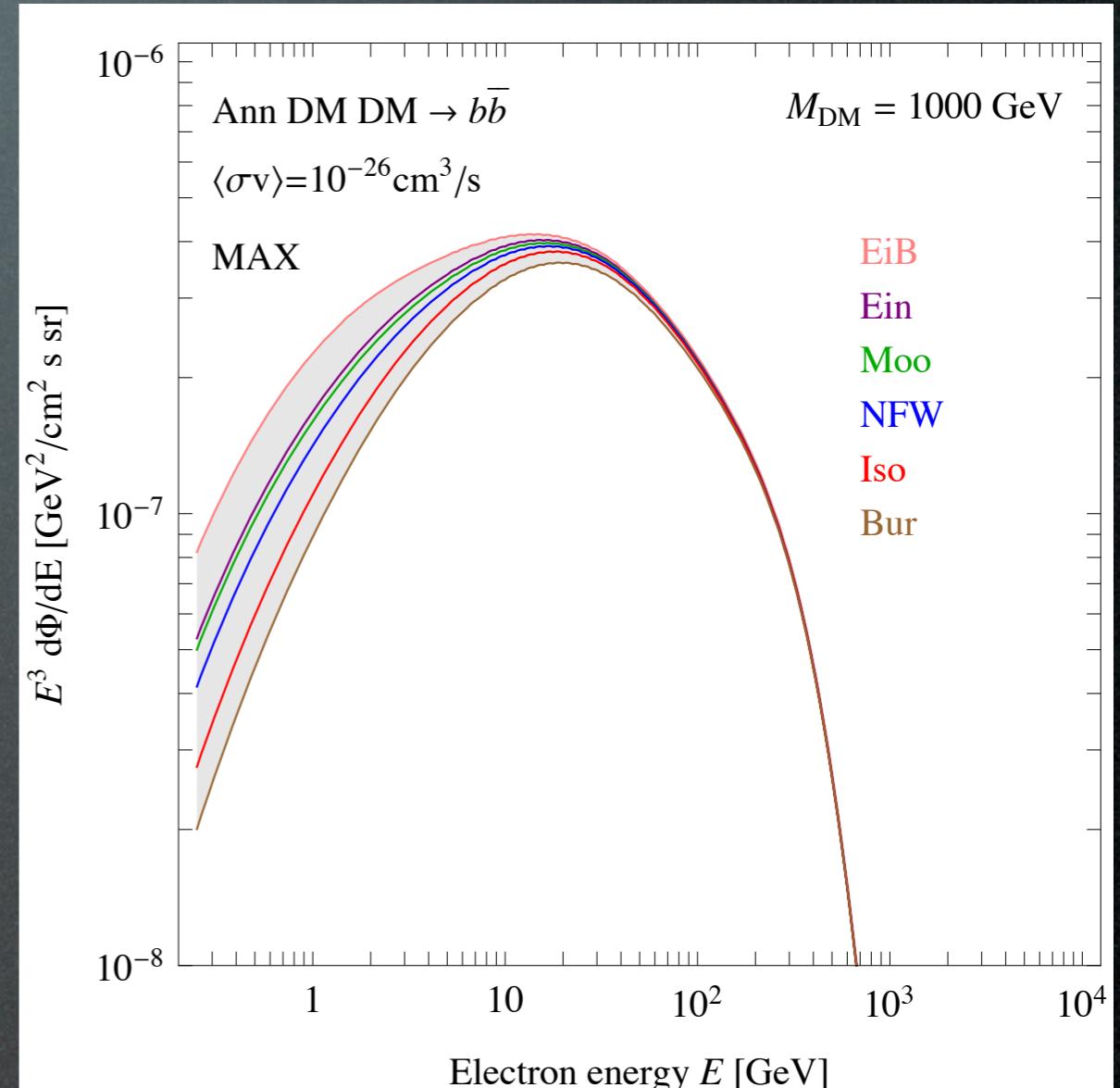
Propagated fluxes

Positrons

Varying prop parameters



From factor 10 to no effect



From factor 10 to no effect

Bottom line: Positrons are affected by propagation,
mainly at low energy

Propagated fluxes

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Cocella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4
arXiv 10
Erratum:

Propagation functions for electrons and positrons everywhere in the Galaxy:

Energy loss coefficient function $b[E, r, z]$ for electrons and positrons in the Galaxy: Mathematica function [bhol.m](#), refer to the notebook [Sample.nb](#) for usage.

Annihilation

Positrons: The file [ElectronHaloFunctGalaxyAnn.m](#) provides the halo functions $I(x, E_s, r, z)$ at a point (r, z) in the Galaxy. [Note: very large file, almost 4GB.]
The notebook [Sample.nb](#) shows how to load and use it.

Decay

Positrons: The file [ElectronHaloFunctGalaxyDec.m](#) provides the halo functions $I(x, E_s, r, z)$ at a point (r, z) in the Galaxy. [Note: very large file, almost 4GB.]
The notebook [Sample.nb](#) shows how to load and use it.

Propagation functions for charged cosmic rays at the location of the Earth:

Annihilation

Positrons: The file [ElectronHaloFunctEarthAnn.m](#) provides the halo functions $I(x, E_s, r_{\text{Earth}})$ at the location of the Earth.
The notebook [Sample.nb](#) shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ [in the approximated formalism - see paper].

Decay

Positrons: The file [ElectronHaloFunctEarthDec.m](#) provides the halo functions $I(x, E_s, r_{\text{Earth}})$ at the location of the Earth.
The notebook [Sample.nb](#) shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ [in the approximated formalism - see paper].

Fluxes of charged cosmic rays at the Earth, after propagation:

Annihilation

Positrons: Mathematica function: the file [ElectronFluxAnn.m](#) provides the spectra $\log_{10} [d\Phi/d \log_{10} E]$.
Refer to the notebook [Sample.nb](#) for usage.

Numerical [table](#): provides the spectra $d\Phi/d \log_{10} E$.
The columns are: [m_{DM} , halo, propagation, $\log_{10} E$, $d\Phi/d \log_{10} E$ for 28 primary channels].
The spectra are computed at a benchmark annihilation cross of $\langle\sigma v\rangle = 10^{-26} \text{ cm}^3/\text{s}$ and renormalized multiplying by $(m_{\text{DM}}/\text{GeV})^2$.
Units are particles/(cm² s sr).

Decay

Positrons: Mathematica function: the file [ElectronFluxDec.m](#) provides the spectra $\log_{10} [d\Phi/d \log_{10} E]$.
Refer to the notebook [Sample.nb](#) for usage.

Numerical [table](#): provides the spectra $d\Phi/d \log_{10} E$.
The columns are: [m_{DM} , halo, propagation, $\log_{10} E$, $d\Phi/d \log_{10} E$ for 28 primary channels].
The spectra are computed at a benchmark decay life of $\tau = 10^{26} \text{ s}$ and renormalized multiplying by $(m_{\text{DM}}/\text{GeV})^2$.
Units are particles/(cm² s sr).

DM detection

direct detection

production at colliders

γ from annihil in galactic center or halo
and from synchrotron emission

Fermi, ICT, radio telescopes...

indirect e^+ from annihil in galactic halo or center

PAMELA, Fermi, HESS, AMS, balloons...

\bar{p} from annihil in galactic halo or center

\bar{d} from annihil in galactic halo or center

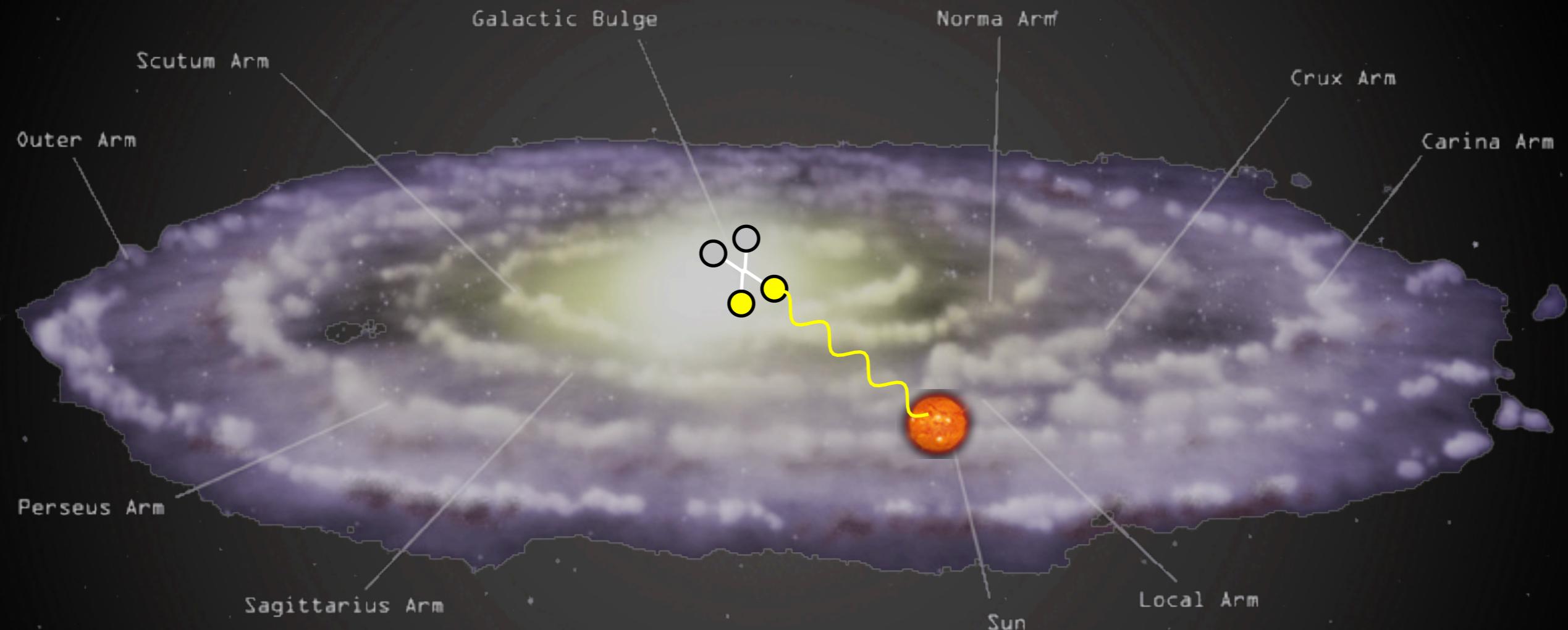
GAPS

$\nu, \bar{\nu}$ from annihil in massive bodies

SK, Icecube, Km3Net

Indirect Detection: basics

γ from DM annihilations in galactic center

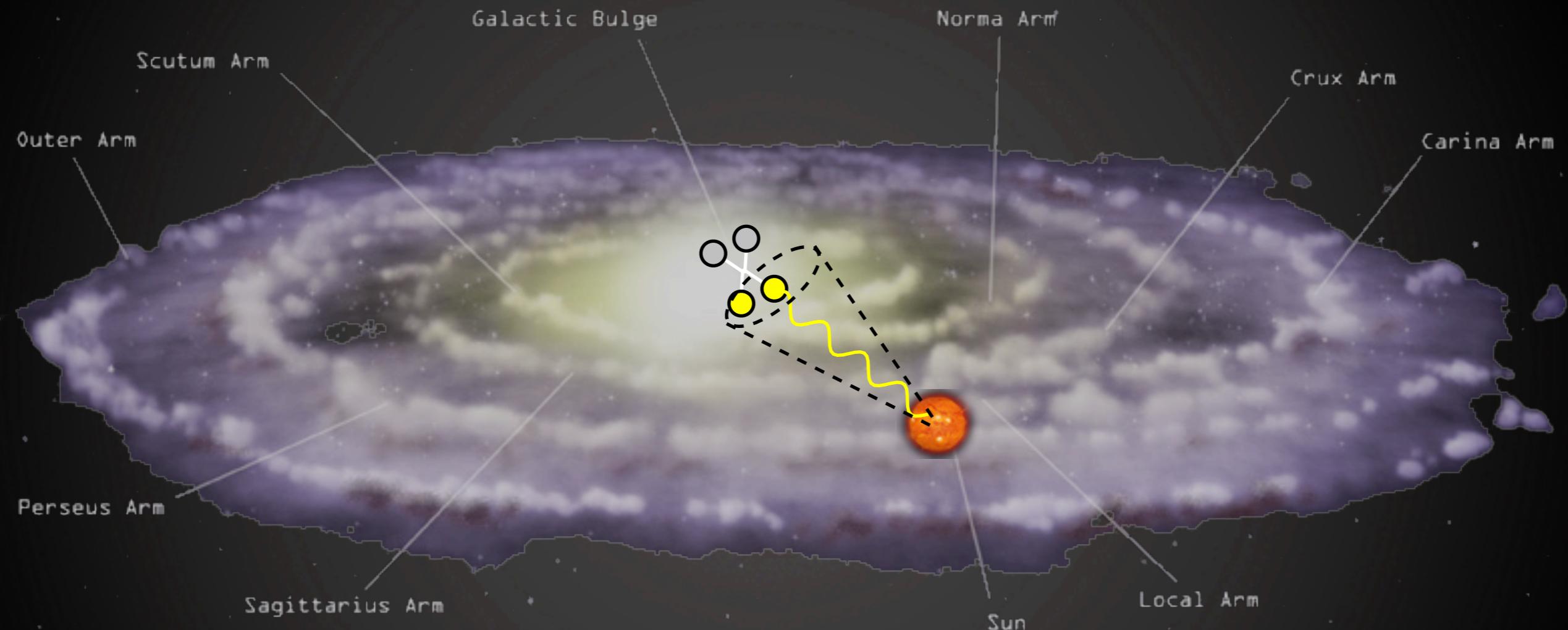


$DM \rightarrow W^-, Z, b, \tau^-, t, h \dots \rightsquigarrow e^\mp, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma$

$DM \rightarrow W^+, Z, \bar{b}, \tau^+, \bar{t}, h \dots \rightsquigarrow e^\pm, \overset{(-)}{p}, \overset{(-)}{D} \dots \text{ and } \gamma$

Indirect Detection: basics

γ from DM annihilations in galactic center

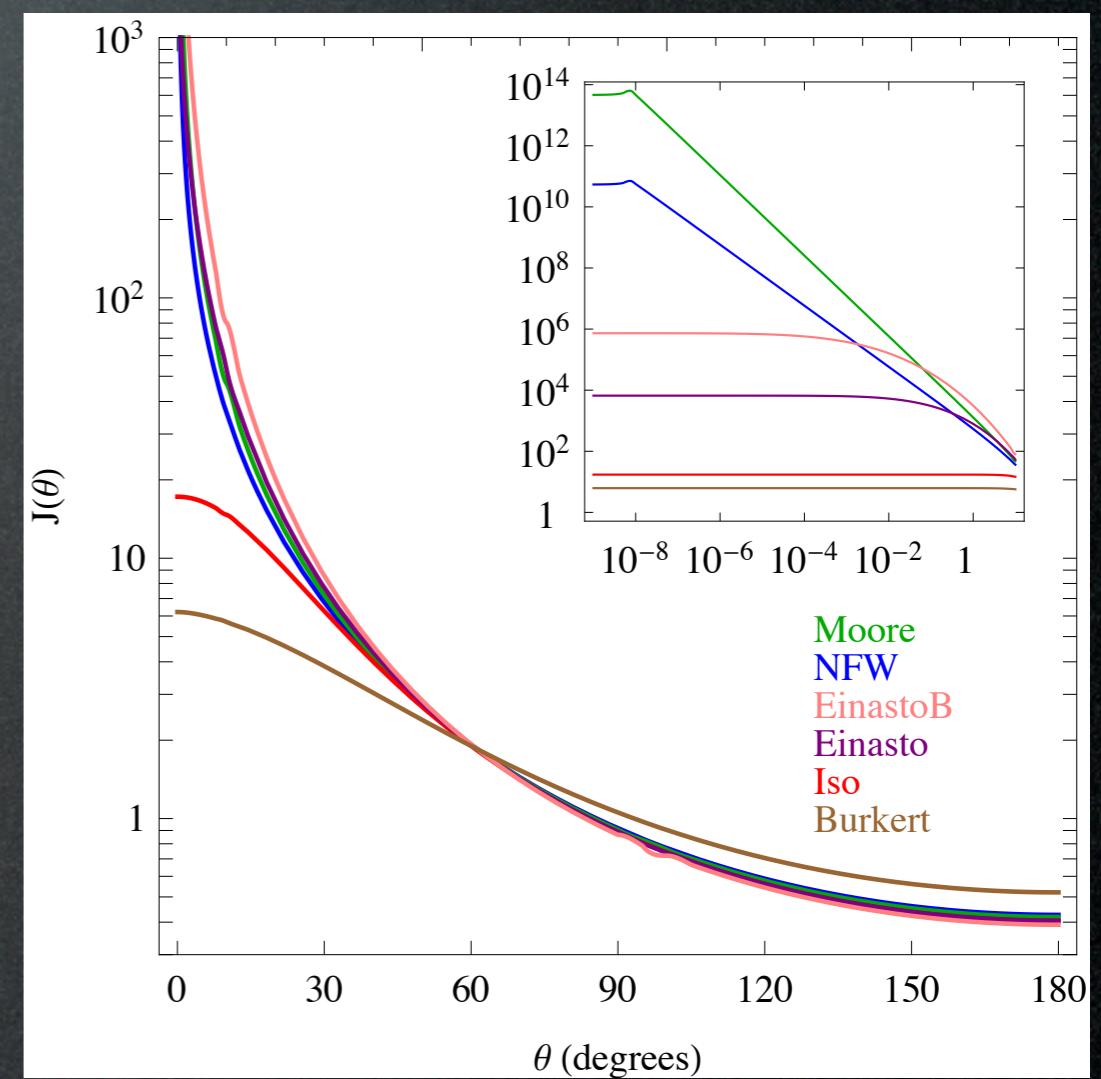
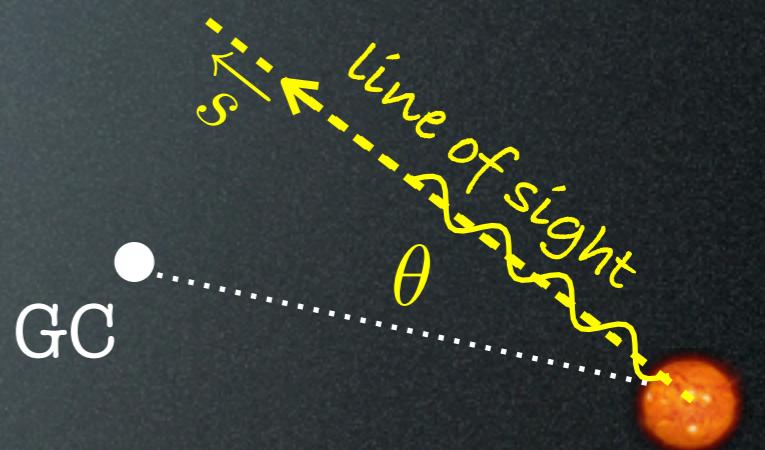


$DM \rightarrow W^-, Z, b, \tau^-, t, h \dots \rightsquigarrow e^\mp, \overset{(-)}{p}, \overset{(-)}{D} \dots$ and γ

$DM \rightarrow W^+, Z, \bar{b}, \tau^+, \bar{t}, h \dots \rightsquigarrow e^\pm, \overset{(-)}{p}, \overset{(-)}{D} \dots$ and γ

$$\frac{d\Phi_\gamma}{d\Omega \, dE} = \frac{1}{2} \frac{r_\odot}{4\pi} \left(\frac{\rho_\odot}{M_{\text{DM}}} \right)^2 J \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE}$$

$$J = \int_{\text{l.o.s.}} \frac{ds}{r_\odot} \left(\frac{\rho(r(s, \theta))}{\rho_\odot} \right)^2$$



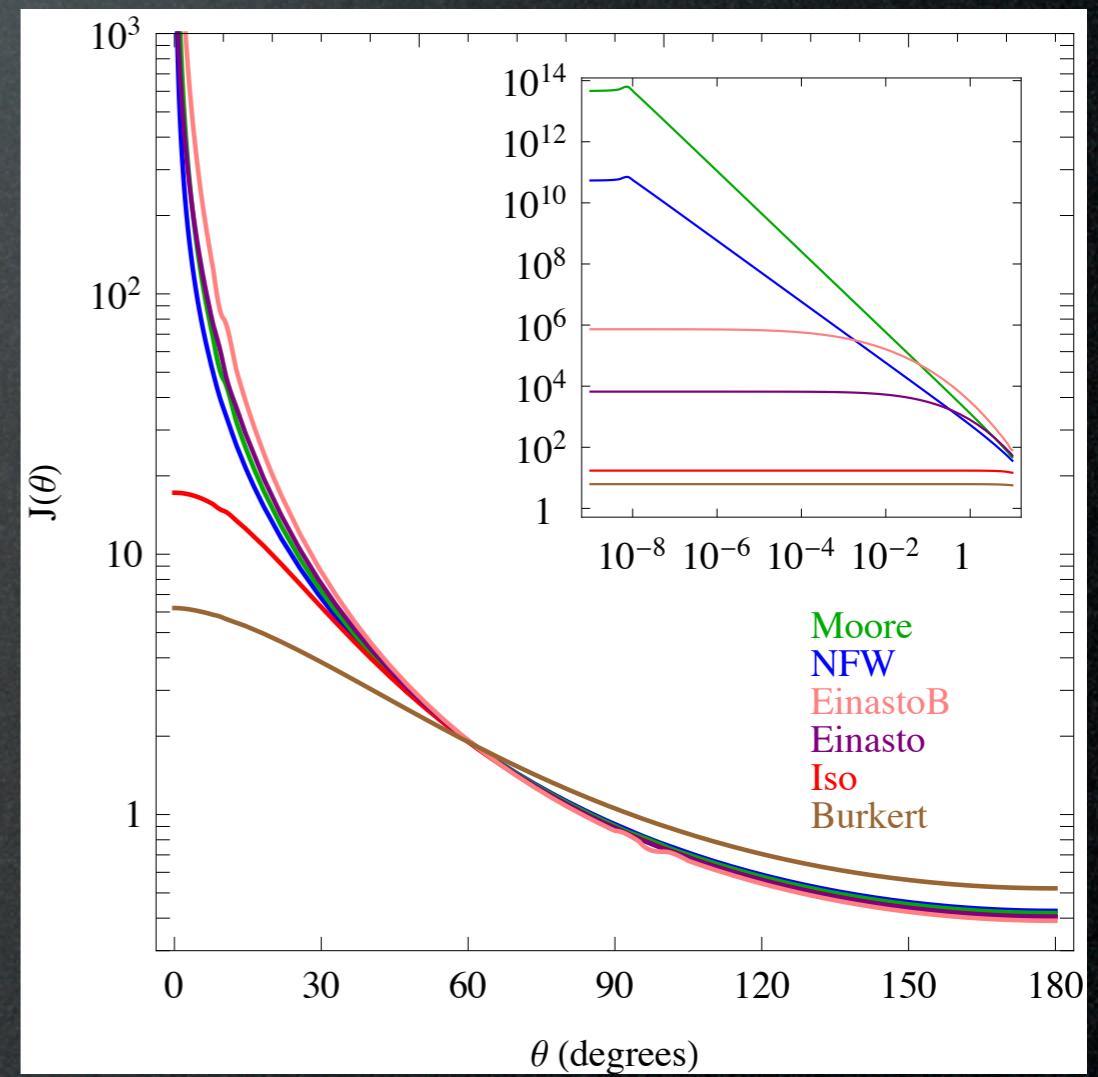
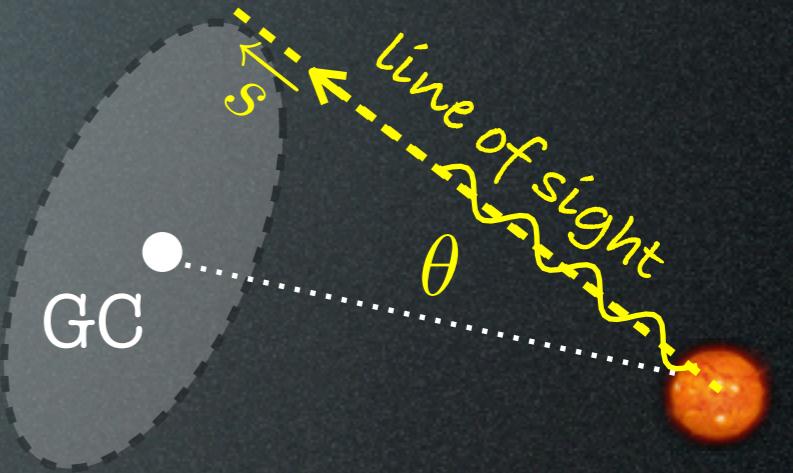
‘Prompt’ gamma rays

$$\frac{d\Phi_\gamma}{dE} = \frac{1}{2} \frac{r_\odot}{4\pi} \left(\frac{\rho_\odot}{M_{\text{DM}}} \right)^2 \bar{J} \Delta\Omega \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE_\gamma}$$

$$\bar{J} = \frac{1}{\Delta\Omega} \int_{\text{l.o.s.}} \int ds \frac{1}{r_\odot} \left(\frac{\rho(r(s, \theta))}{\rho_\odot} \right)^2$$

Region	$\Delta\Omega$ [steradians]	NFW	Ein	EinB	Iso	Bur	Moore
‘GC 0.1°’	$0.96 \cdot 10^{-5}$	11579	3579	55665	17.2	6.21	81751
‘GC 0.14°’	$0.19 \cdot 10^{-4}$	8255	3206	43306	17.2	6.21	52395
‘GC 1°’	$0.96 \cdot 10^{-3}$	1118	1196	6945	17.2	6.21	3855
‘GC 2°’	0.004	542	711	3103	17.2	6.19	1521
‘Gal Ridge’	$0.29 \cdot 10^{-3}$	1904	1605	11828	17.2	6.21	7927
‘3 × 3’	0.011	306	443	1577	17.1	6.16	741
‘5 × 5’	0.030	174	264	783	16.8	6.10	367
‘5 × 30’	0.183	47.7	70.5	170	12.1	5.16	84.8
‘10 × 10’	0.121	77.7	118	280	15.5	5.85	138
‘10 × 30’	0.364	35.5	51.8	109	11.7	5.09	57.2

Spread is very large
for small regions close to GC



'Prompt' gamma rays

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Corcella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection",
arXiv 1012.4515, JCAP 1103 (2011) 051.
Erratum: JCAP 1210 (2012) E01.

J factors for prompt gamma rays:

Annihilation

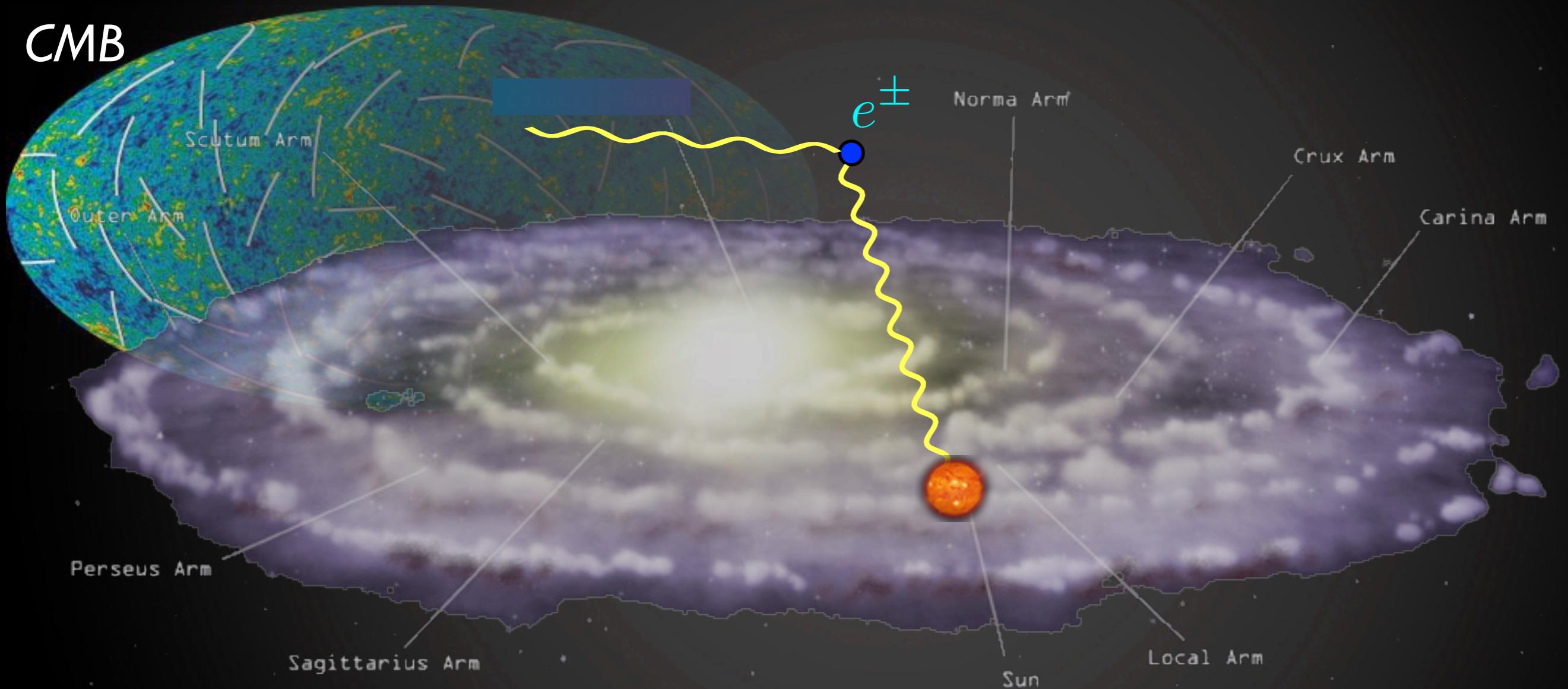
Mathematica function: the file [JAnn.m](#) provides $\text{Log}_{10}[J(\theta)]$.
Refer to the notebook [Sample.nb](#) for usage.

Decay

Mathematica function: the file [JDec.m](#) provides $\text{Log}_{10}[J(\theta)]$.
Refer to the notebook [Sample.nb](#) for usage.

Indirect Detection: basics

γ from Inverse Compton on e^\pm in halo

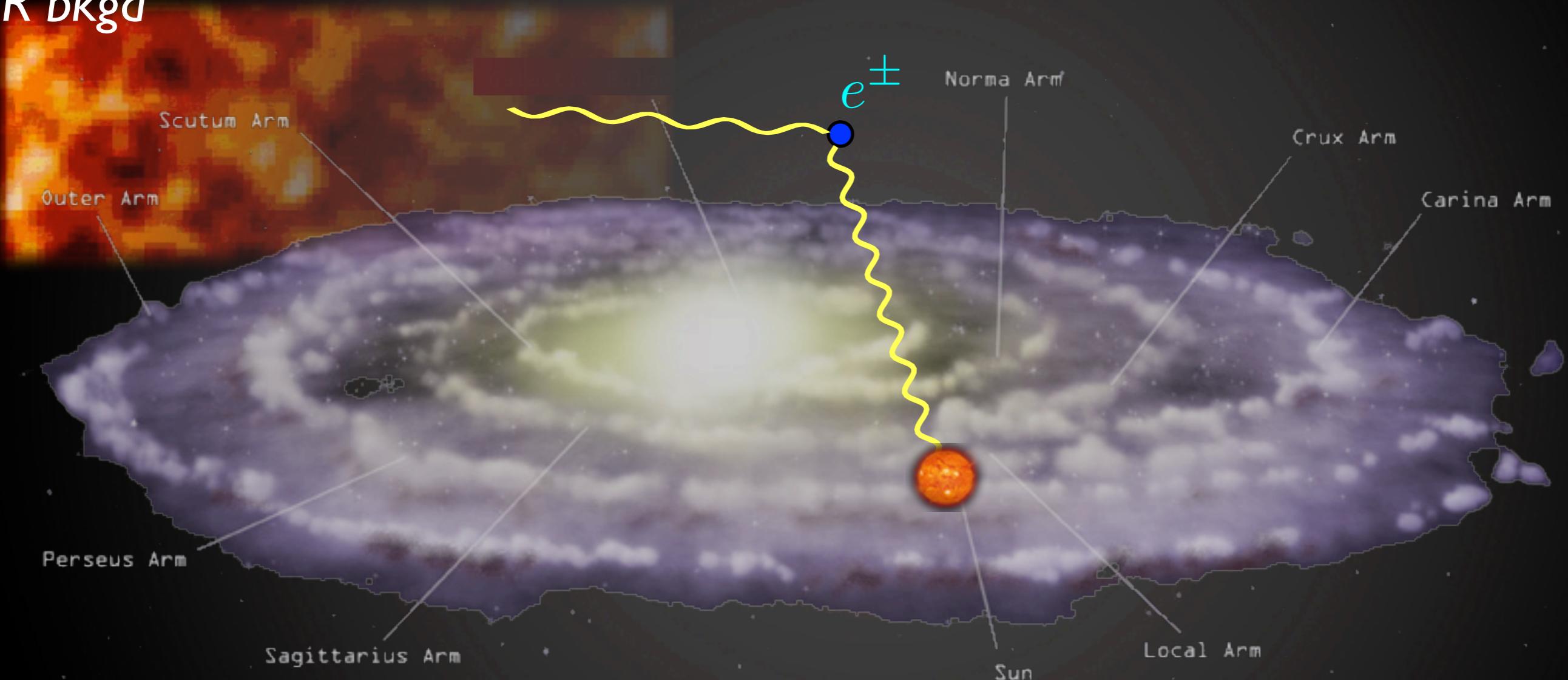


- upscatter of CMB, infrared and starlight photons on energetic e^\pm
- probes regions outside of Galactic Center

Indirect Detection: basics

γ from Inverse Compton on e^\pm in halo

IR bkgd

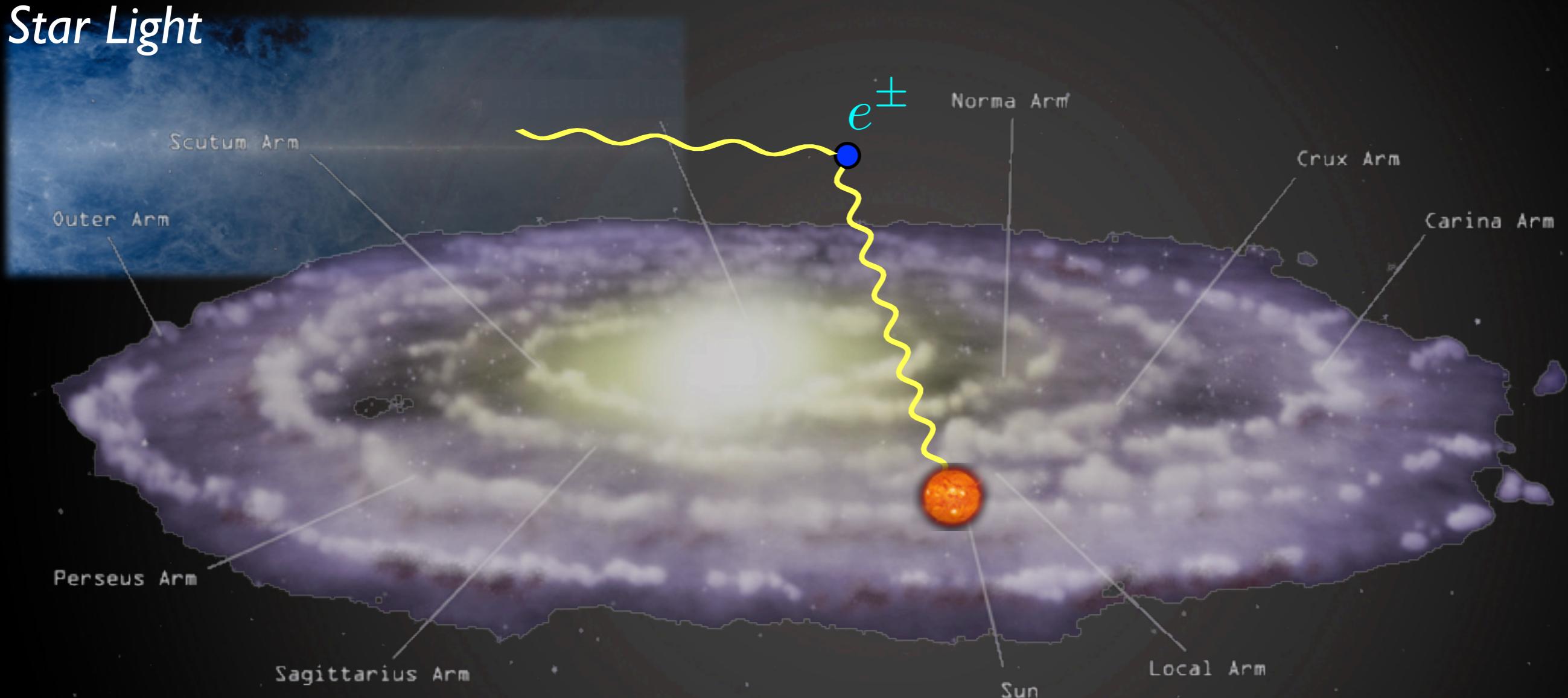


- upscatter of CMB, infrared and starlight photons on energetic e^\pm
- probes regions outside of Galactic Center

Indirect Detection: basics

γ from Inverse Compton on e^\pm in halo

Star Light



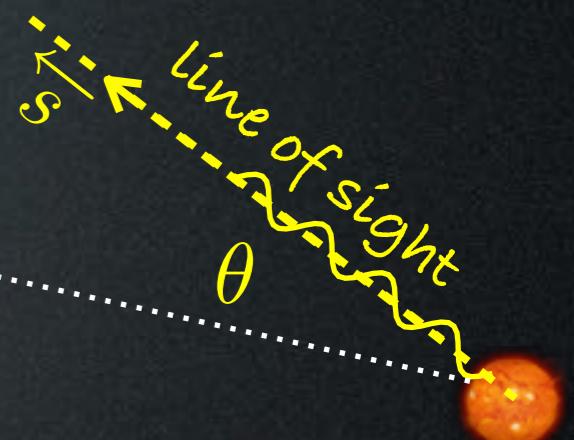
- upscatter of CMB, infrared and starlight photons on energetic e^\pm
- probes regions outside of Galactic Center

ICS gamma rays

Flux of ICS γ
from direction θ :

$$\frac{d\Phi_{\text{IC}\gamma}}{dE_\gamma d\theta} = \frac{1}{E_\gamma} \int_{\text{l.o.s.}} \frac{ds}{4\pi} 2 \int_{m_e}^{M_{\text{DM}}} dE_e \sum_i \mathcal{P}_{\text{IC}}^i \frac{dn_{e^\pm}}{dE_e}$$

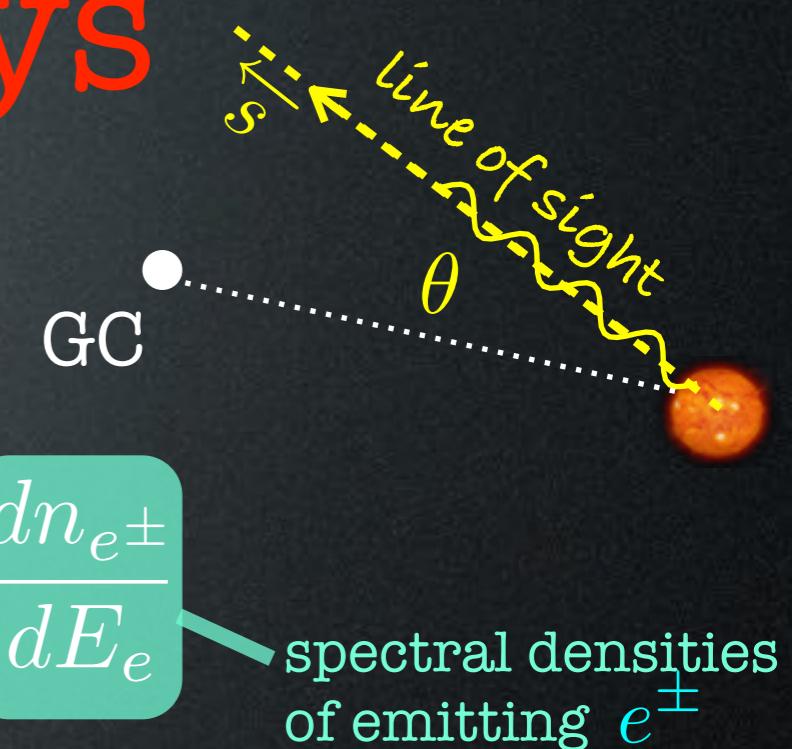
integral
over s emission from a 'cell' in \vec{s}



ICS gamma rays

Flux of ICS γ
from direction θ :

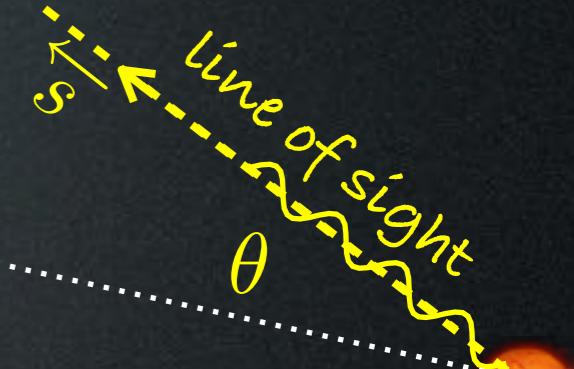
$$\frac{d\Phi_{\text{IC}\gamma}}{dE_\gamma d\theta} = \frac{1}{E_\gamma} \int_{\text{l.o.s.}} \frac{ds}{4\pi} 2 \int_{m_e}^{M_{\text{DM}}} dE_e \sum_i \mathcal{P}_{\text{IC}}^i \frac{dn_{e^\pm}}{dE_e}$$



spectral densities
of emitting e^\pm

ICS gamma rays

Flux of ICS γ
from direction θ :

$$\frac{d\Phi_{\text{IC}\gamma}}{dE_\gamma d\theta} = \frac{1}{E_\gamma} \int_{\text{l.o.s.}} \frac{ds}{4\pi} 2 \int_{m_e}^{M_{\text{DM}}} dE_e \sum_i \mathcal{P}_{\text{IC}}^i \frac{dn_{e^\pm}}{dE_e}$$


spectral densities
of emitting e^\pm

power emitted by individual ‘cell’:

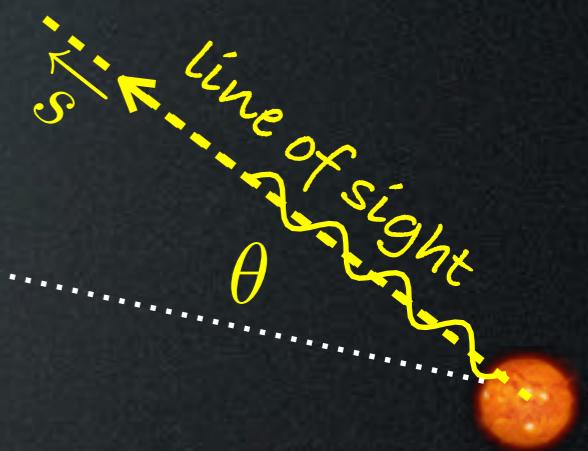
$$\mathcal{P}_{\text{IC}}^i(E_\gamma, E_e, \vec{x}) = \frac{3\sigma_T}{4\gamma^2} \int_{1/4\gamma^2}^1 dq \left(E_\gamma - \frac{E_\gamma}{4q\gamma^2(1-\epsilon)} \right) \frac{n_i(E_\gamma^0(q), \vec{x})}{q} \left[2q \ln q + q + 1 - 2q^2 + \frac{1}{2} \frac{\epsilon^2}{1-\epsilon} (1-q) \right]$$

$$q = \frac{\epsilon}{\Gamma_E(1-\epsilon)}, \quad \Gamma_E = \frac{4E_\gamma^0 E_e}{m_e^2}, \quad \epsilon = \frac{E_\gamma}{E_e}$$

ICS gamma rays

Flux of ICS γ
from direction θ :

$$\frac{d\Phi_{\text{IC}\gamma}}{dE_\gamma d\theta} = \frac{1}{E_\gamma} \int_{\text{l.o.s.}} \frac{ds}{4\pi} 2 \int_{m_e}^{M_{\text{DM}}} dE_e \sum_i \mathcal{P}_{\text{IC}}^i \frac{dn_{e^\pm}}{dE_e}$$



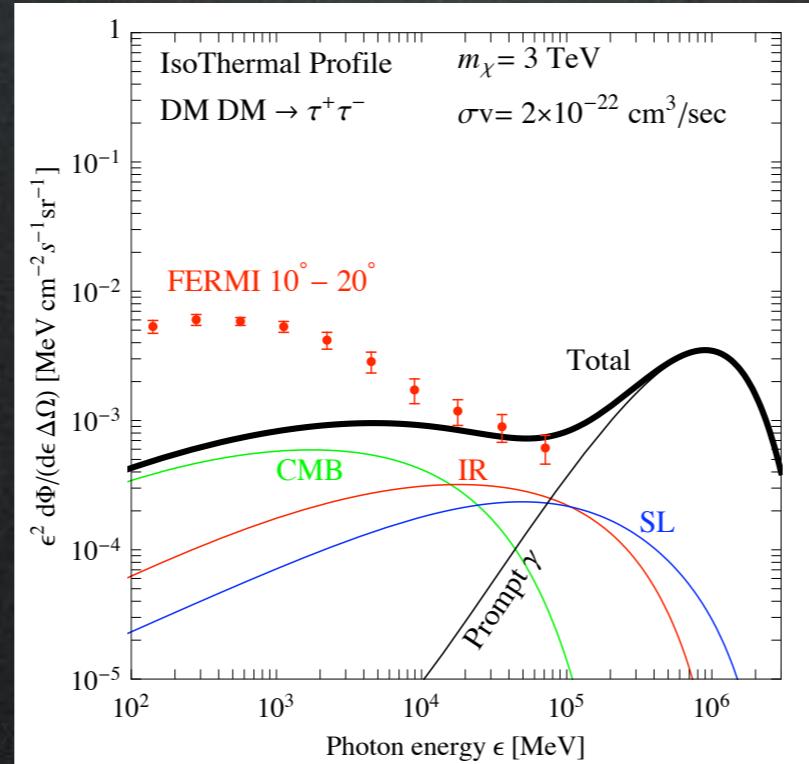
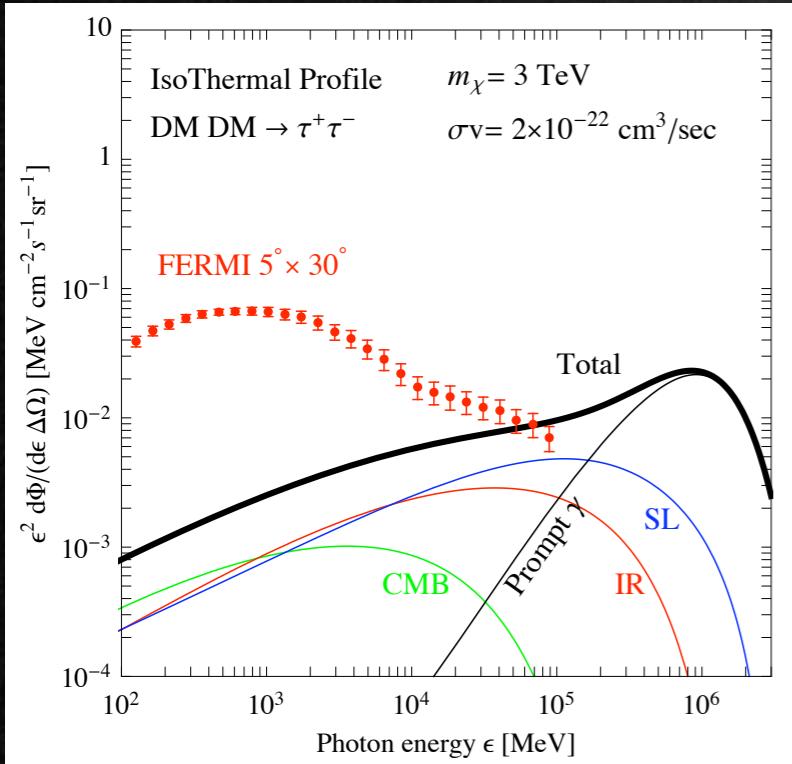
spectral densities
of emitting e^\pm

power emitted by individual ‘cell’:

$$\mathcal{P}_{\text{IC}}^i(E_\gamma, E_e, \vec{x}) = \frac{3\sigma_T}{4\gamma^2} \int_{1/4\gamma^2}^1 dq \left(E_\gamma - \frac{E_\gamma}{4q\gamma^2(1-\epsilon)} \right) \frac{n_i(E_\gamma^0(q), \vec{x})}{q} \left[2q \ln q + q + 1 - 2q^2 + \frac{1}{2} \frac{\epsilon^2}{1-\epsilon} (1-q) \right]$$

$$q = \frac{\epsilon}{\Gamma_E(1-\epsilon)}, \quad \Gamma_E = \frac{4E_\gamma^0 E_e}{m_e^2}, \quad \epsilon = \frac{E_\gamma}{E_e}$$

Different windows = different weights of target lights:



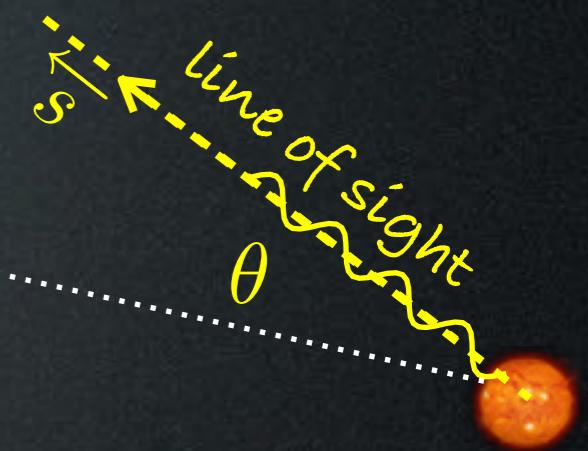
Cirelli, Pani, Serpico 0912.0663

spectral densities
of target lights

ICS gamma rays

Flux of ICS γ
from direction θ :

$$\frac{d\Phi_{\text{IC}\gamma}}{dE_\gamma d\theta} = \frac{1}{E_\gamma} \int_{\text{l.o.s.}} \frac{ds}{4\pi} 2 \int_{m_e}^{M_{\text{DM}}} dE_e \sum_i \mathcal{P}_{\text{IC}}^i \frac{dn_{e^\pm}}{dE_e}$$

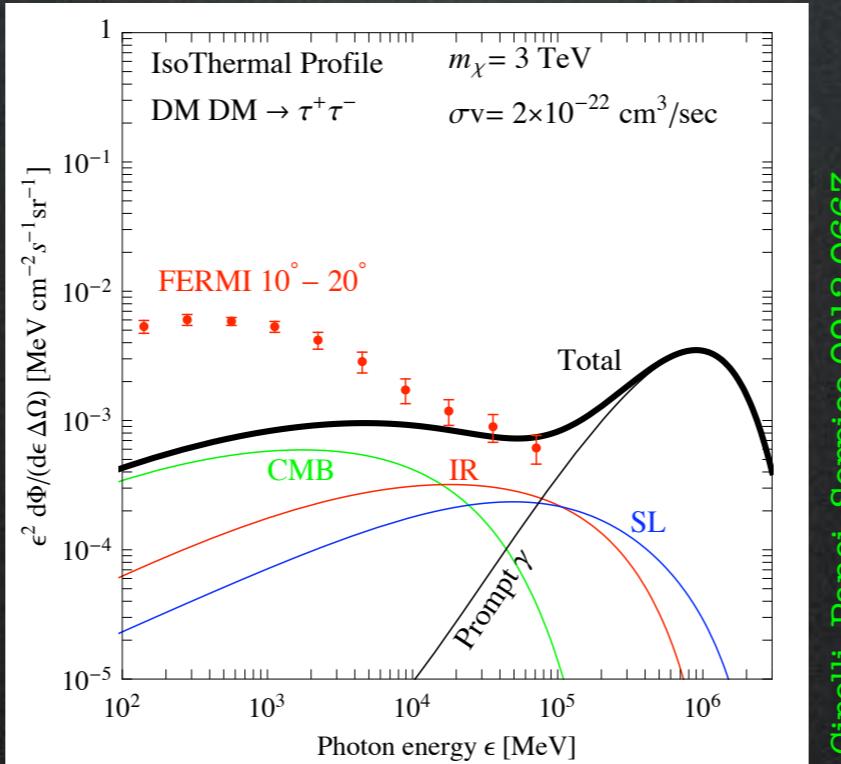
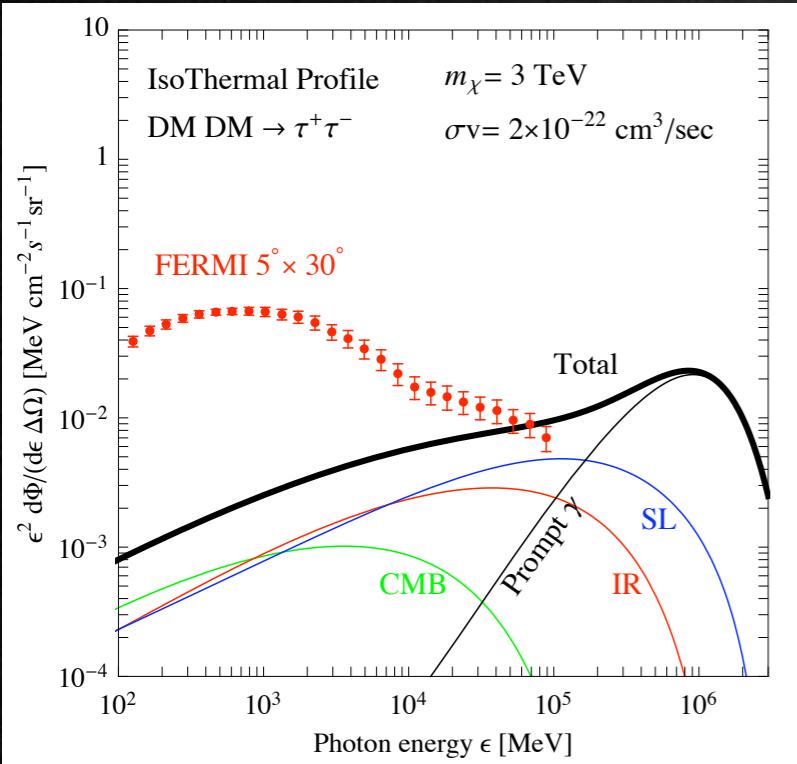


power emitted by individual ‘cell’:

$$\mathcal{P}_{\text{IC}}^i(E_\gamma, E_e, \vec{x}) = \frac{3\sigma_T}{4\gamma^2} \int_{1/4\gamma^2}^1 dq \left(E_\gamma - \frac{E_\gamma}{4q\gamma^2(1-\epsilon)} \right) \frac{n_i(E_\gamma^0(q), \vec{x})}{q} \left[2q \ln q + q + 1 - 2q^2 + \frac{1}{2} \frac{\epsilon^2}{1-\epsilon} (1-q) \right]$$

$$q = \frac{\epsilon}{\Gamma_E(1-\epsilon)}, \quad \Gamma_E = \frac{4E_\gamma^0 E_e}{m_e^2}, \quad \epsilon = \frac{E_\gamma}{E_e}$$

Different windows = different weights of target lights:



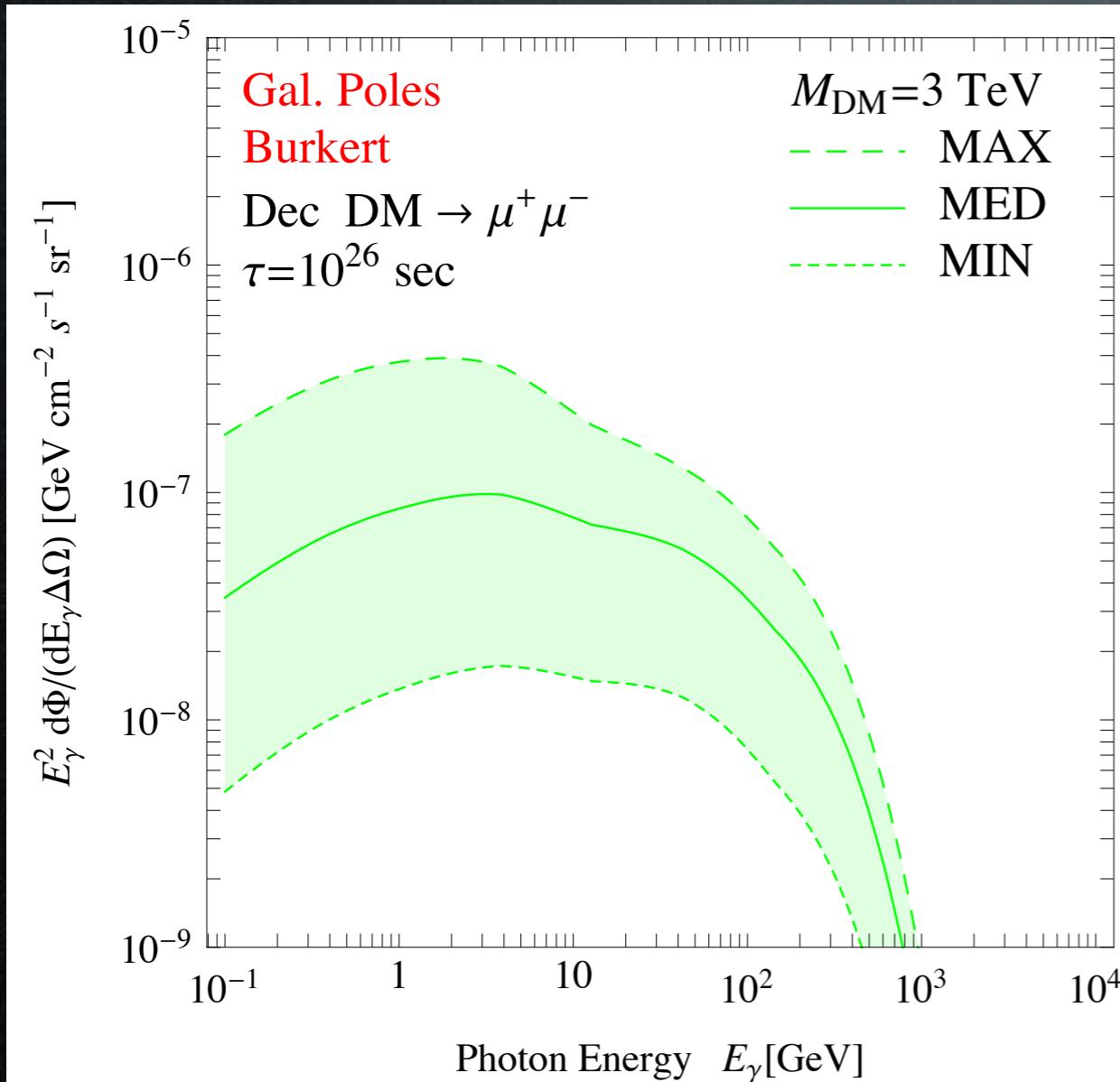
Cirelli, Pani, Serpico 0912.0663

spectral densities
of target lights

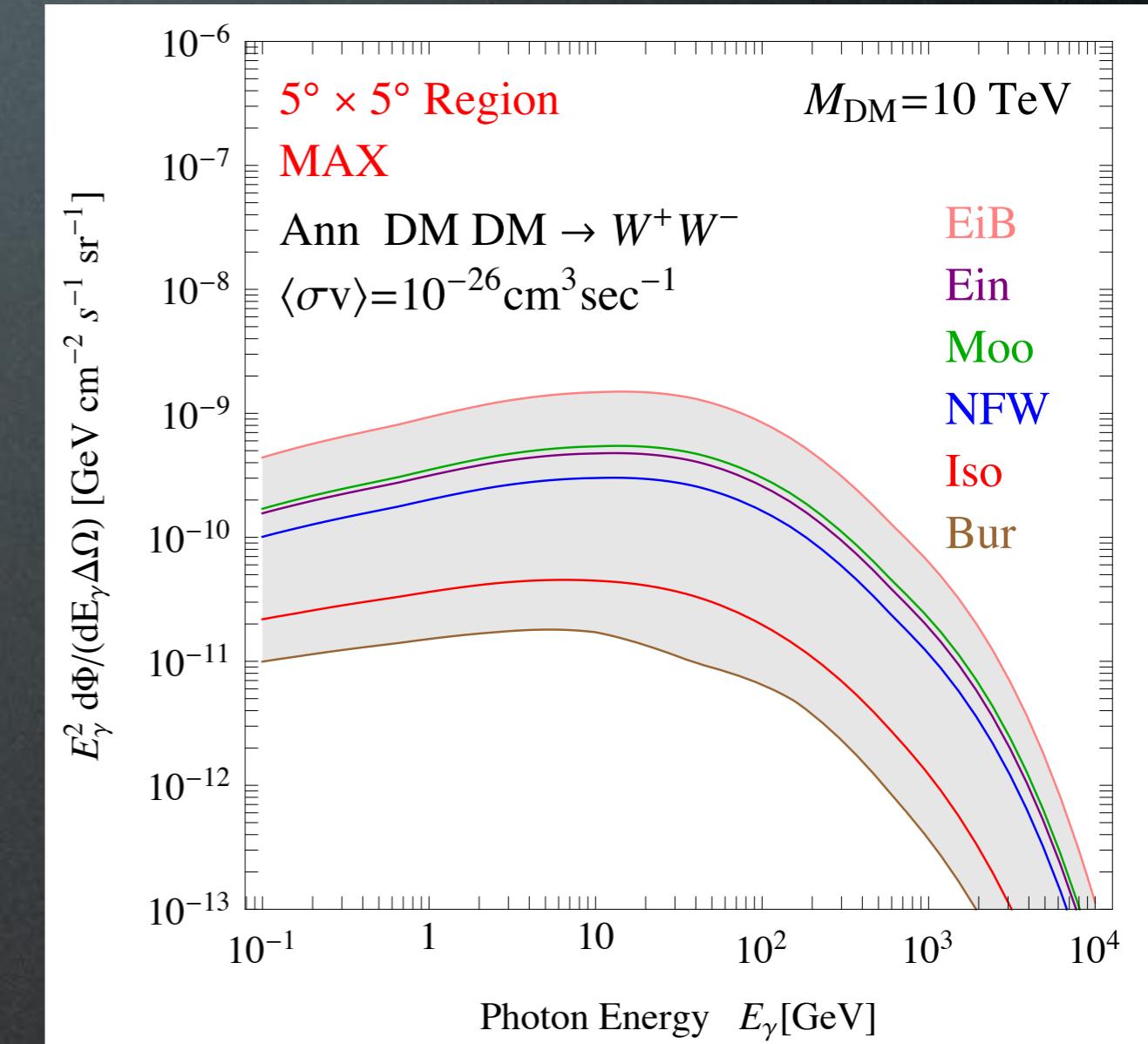
Everything can
be rearranged
in terms of ICS
halo functions
(see 1012.4515 for details)

ICS gamma rays

Varying prop parameters

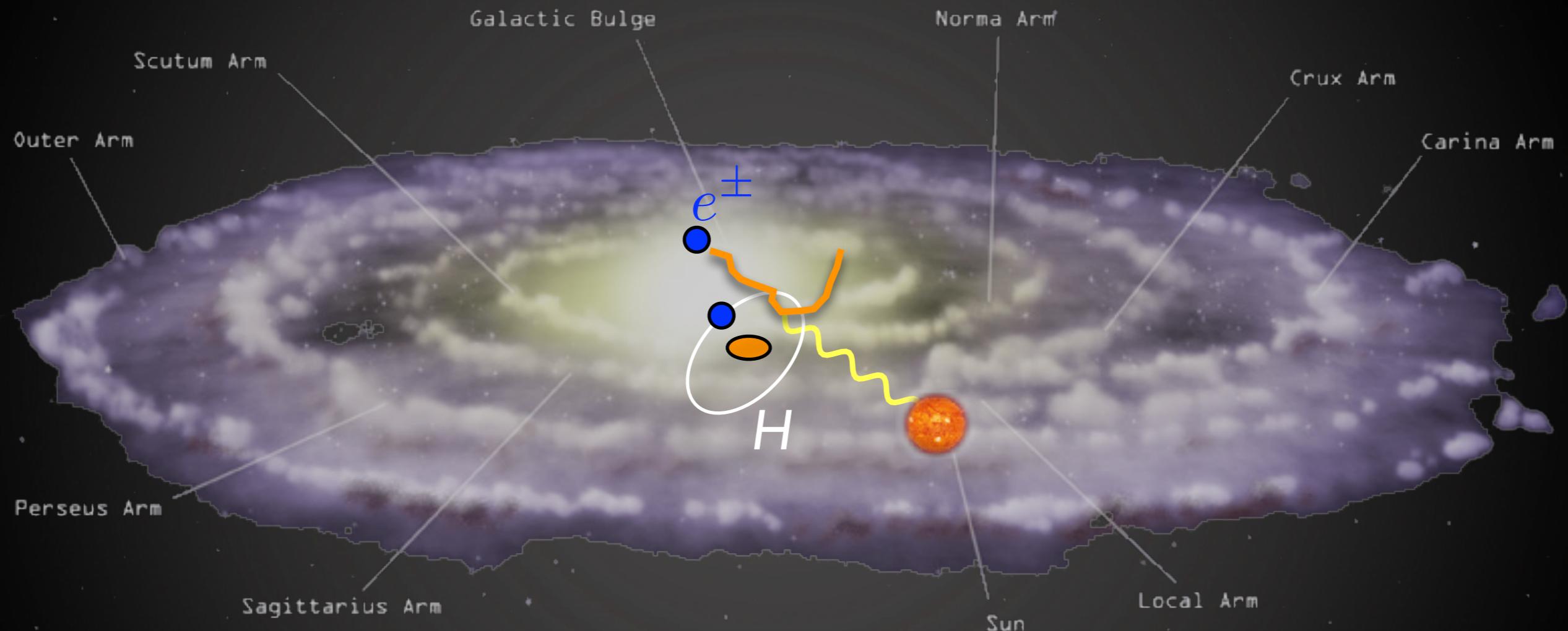


Varying halo profile



Secondary emission

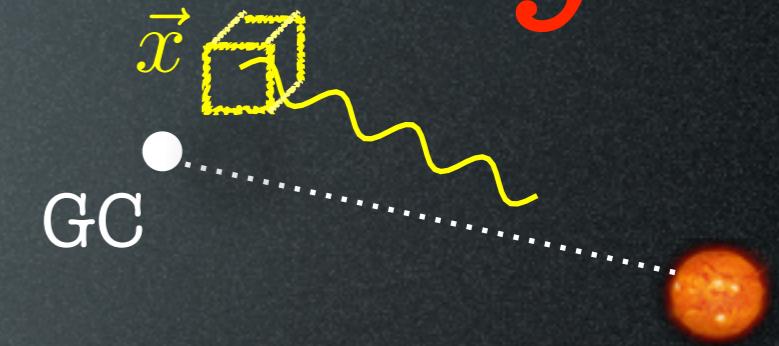
soft gammas from bremsstrahlung of e^\pm on ISM



Bremsstrahlung gamma rays

Emission of bremsstrahlung γ
from a cell located in \vec{x} :

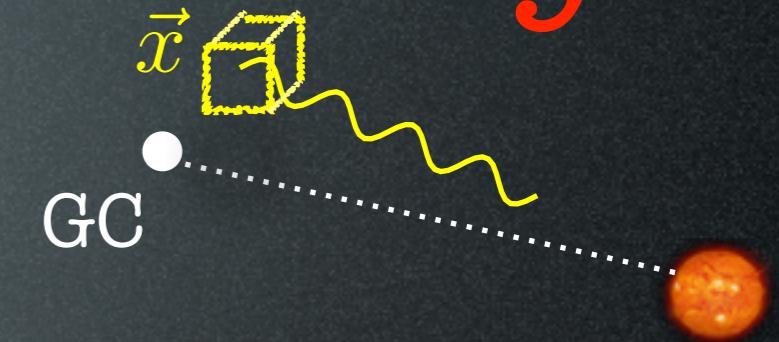
$$\frac{d\mathcal{E}_{\gamma, \text{brem}}(\vec{x})}{dE_\gamma} = \sum_i n_i(\vec{x}) \int_{E_L} dE_{e^\pm} 2 \frac{d\Phi_{e^\pm}(\vec{x})}{dE_{e^\pm}} \cdot \frac{d\sigma_i}{dE_\gamma}$$



Bremsstrahlung gamma rays

Emission of bremsstrahlung γ
from a cell located in \vec{x} :

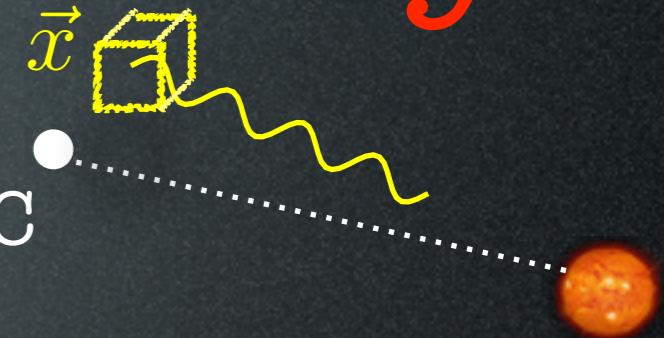
$$\frac{d\mathcal{E}_{\gamma, \text{brem}}(\vec{x})}{dE_\gamma} = \underbrace{\sum_i n_i(\vec{x})}_{\text{gas densities}} \int_{E_L} dE_{e^\pm} \underbrace{2 \frac{d\Phi_{e^\pm}(\vec{x})}{dE_{e^\pm}} \cdot \frac{d\sigma_i}{dE_\gamma}}_{e^\pm \text{ input spectrum (after propagation!)}}$$



Bremsstrahlung gamma rays

Emission of bremsstrahlung γ
from a cell located in \vec{x} :

$$\frac{d\mathcal{E}_{\gamma, \text{brem}}(\vec{x})}{dE_\gamma} = \sum_i n_i(\vec{x}) \int_{E_L} dE_{e^\pm} 2 \frac{d\Phi_{e^\pm}(\vec{x})}{dE_{e^\pm}} \cdot \boxed{\frac{d\sigma_i}{dE_\gamma}}$$



bremsstrahlung
differential cross section:

$$\boxed{\frac{d\sigma_i(E_{e^\pm}, E_\gamma)}{dE_\gamma}} = \frac{3 \alpha_{\text{em}} \sigma_T}{8\pi E_\gamma} \left\{ \left[1 + \left(1 - \frac{E_\gamma}{E_{e^\pm}} \right)^2 \right] \phi_1^i - \frac{2}{3} \left(1 - \frac{E_\gamma}{E_{e^\pm}} \right) \phi_2^i \right\}$$

Bremsstrahlung gamma rays



Emission of bremsstrahlung γ
from a cell located in \vec{x} :

$$\frac{d\mathcal{E}_{\gamma, \text{brem}}(\vec{x})}{dE_\gamma} = \sum_i n_i(\vec{x}) \int_{E_L} dE_{e^\pm} 2 \frac{d\Phi_{e^\pm}(\vec{x})}{dE_{e^\pm}} \cdot \boxed{\frac{d\sigma_i}{dE_\gamma}}$$

bremsstrahlung
differential cross section:

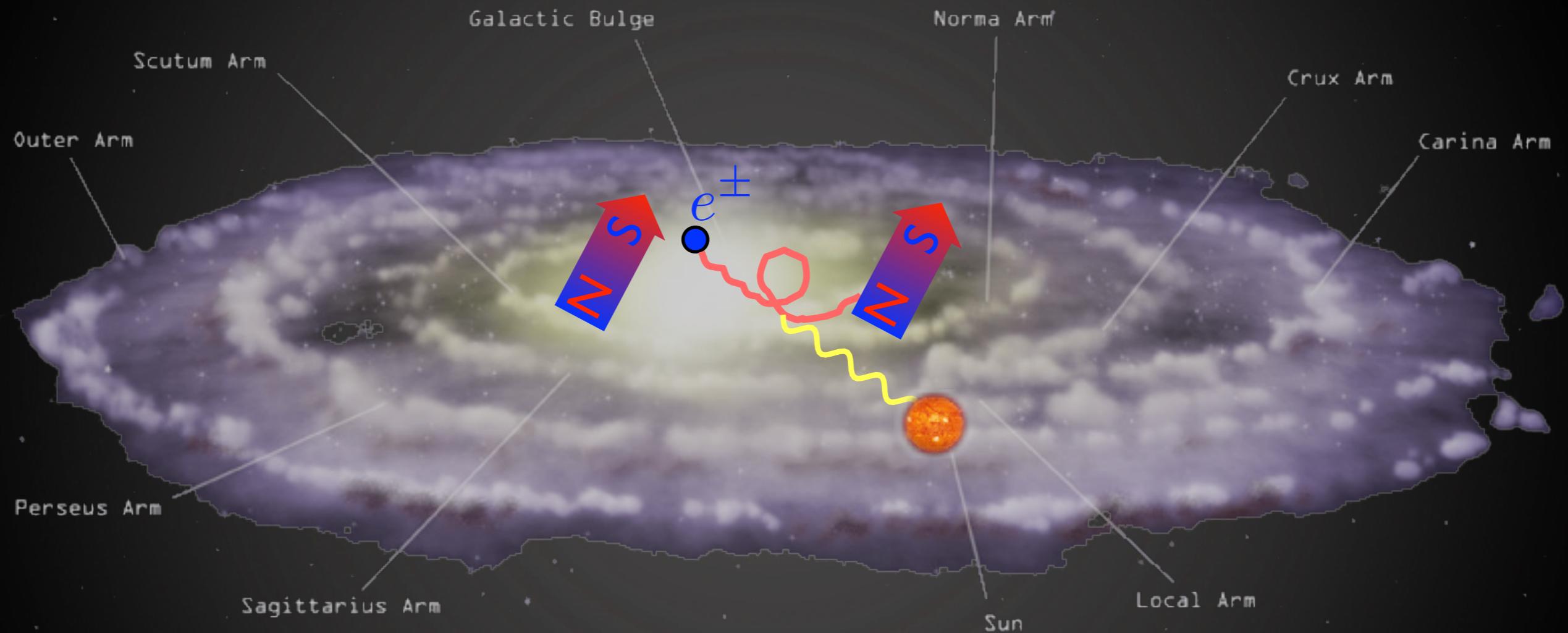
$$\boxed{\frac{d\sigma_i(E_{e^\pm}, E_\gamma)}{dE_\gamma}} = \frac{3 \alpha_{\text{em}} \sigma_T}{8\pi E_\gamma} \left\{ \left[1 + \left(1 - \frac{E_\gamma}{E_{e^\pm}} \right)^2 \right] \phi_1^i - \frac{2}{3} \left(1 - \frac{E_\gamma}{E_{e^\pm}} \right) \phi_2^i \right\}$$

on ionized gas: $\phi_1^{\text{ion}}(E_{e^\pm}, E_\gamma) = \phi_2^{\text{ion}}(E_{e^\pm}, E_\gamma) = 4(Z^2 + Z) \left\{ \log \left[\frac{2E_{e^\pm}}{mc^2} \left(\frac{E_{e^\pm} - E_\gamma}{E_\gamma} \right) \right] - \frac{1}{2} \right\}$

on neutral gas: $\phi_1^{\text{H}}(\Delta = 0) \equiv \phi_{1,\text{ss}}^{\text{H}} = 45.79, \phi_2^{\text{H}}(\Delta = 0) \equiv \phi_{2,\text{ss}}^{\text{H}} = 44.46, \phi_1^{\text{He}}(\Delta = 0) \equiv \phi_{1,\text{ss}}^{\text{He}} = 134.60,$
 $\phi_2^{\text{He}}(\Delta = 0) \equiv \phi_{2,\text{ss}}^{\text{He}} = 131.40, \phi_{(1,2)}^{\text{H}_2}(\Delta = 0) \simeq 2 \phi_{(1,2),\text{ss}}^{\text{H}}$
 $\left(\Delta = \frac{E_\gamma m_e}{4\alpha_{\text{em}} E_{e^\pm} (E_{e^\pm} - E_\gamma)} \right)$

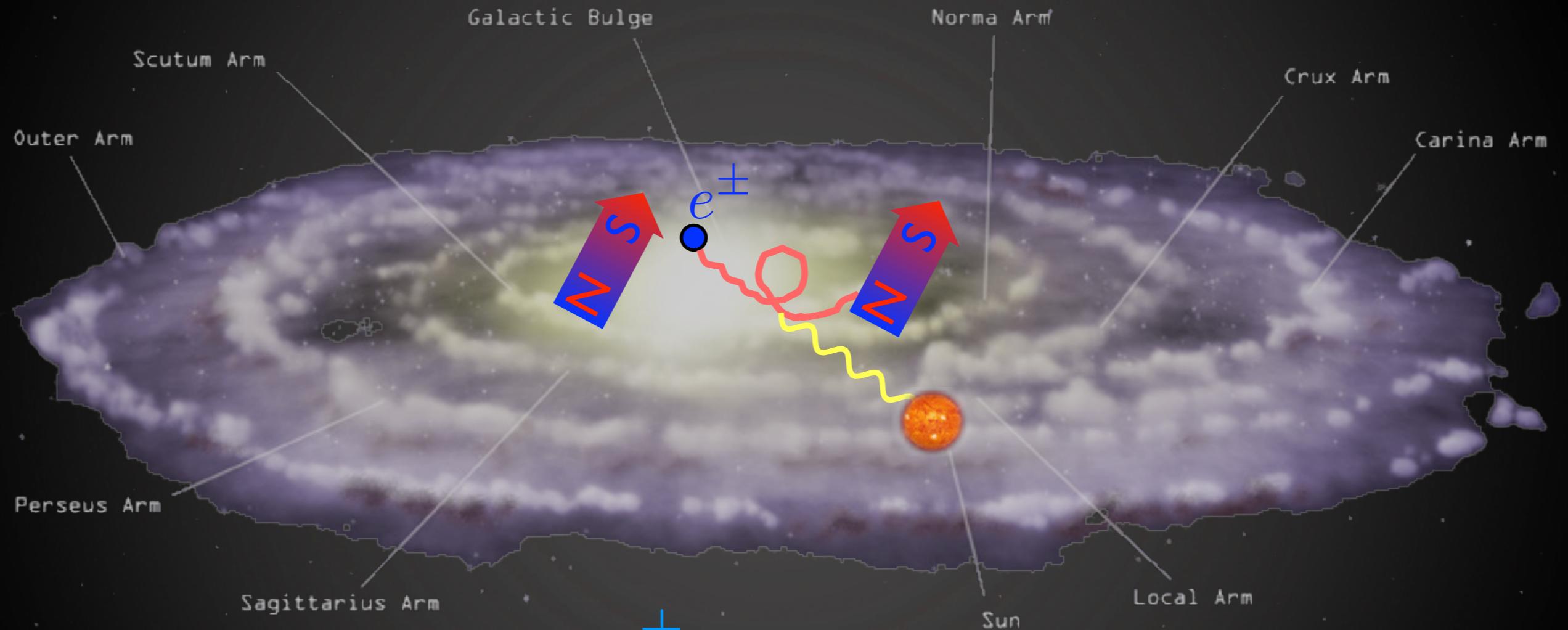
Secondary emission

radio-waves from synchro radiation of e^\pm in GC



Secondary emission

radio-waves from synchro radiation of e^\pm in GC



- compute the population of e^\pm from DM annihilations in the GC
- compute the synchrotron emitted power for different configurations of galactic \vec{B}

(assuming ‘scrambled’ B; in principle, directionality could focus emission, lift bounds by O(some))

Synchrotron radiation

Spectrum of synchrotron emission
from region $\Delta\Omega$:

$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

Synchrotron radiation

Spectrum of synchrotron emission
from region $\Delta\Omega$:

$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

spectral densities
of emitting e^\pm

Synchrotron radiation

Spectrum of synchrotron emission
from region $\Delta\Omega$:

$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

spectral densities
of emitting e^\pm

synchr power emitted by individual electron:

$$\frac{dW_{\text{syn}}}{d\nu} = \frac{\sqrt{3}}{6\pi} \frac{e^3 B}{m_e} F\left(\frac{\nu}{\nu_{\text{syn}}}\right) \quad F(x) = x \int_x^\infty K_{5/3}(\xi) d\xi \approx \frac{8\pi}{9\sqrt{3}} \delta(x - 1/3)$$
$$\nu_{\text{syn}} = \frac{3eBp^2}{4\pi m_e^3}$$

Synchrotron radiation

Spectrum of synchrotron emission from region $\Delta\Omega$:

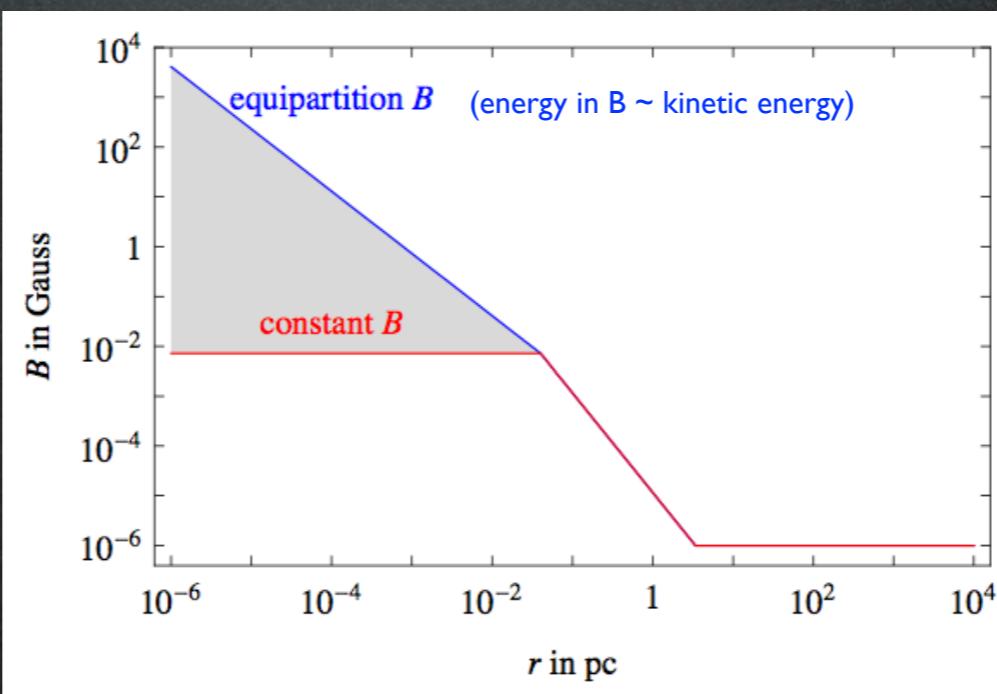
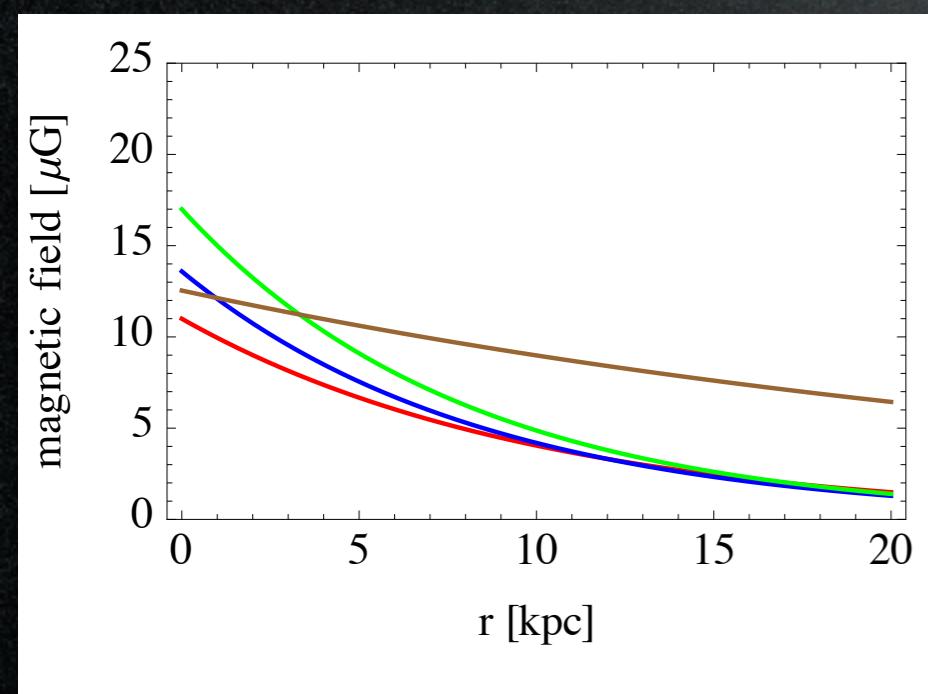
$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

spectral densities
of emitting e^\pm

synchr power emitted by individual electron:

$$\frac{dW_{\text{syn}}}{d\nu} = \frac{\sqrt{3}}{6\pi} \frac{e^3 B}{m_e} F\left(\frac{\nu}{\nu_{\text{syn}}}\right) \quad F(x) = x \int_x^\infty K_{5/3}(\xi) d\xi \approx \frac{8\pi}{9\sqrt{3}} \delta(x - 1/3)$$

$$\nu_{\text{syn}} = \frac{3eBp^2}{4\pi m_e^3}$$



Synchrotron radiation

Spectrum of synchrotron emission from region $\Delta\Omega$:

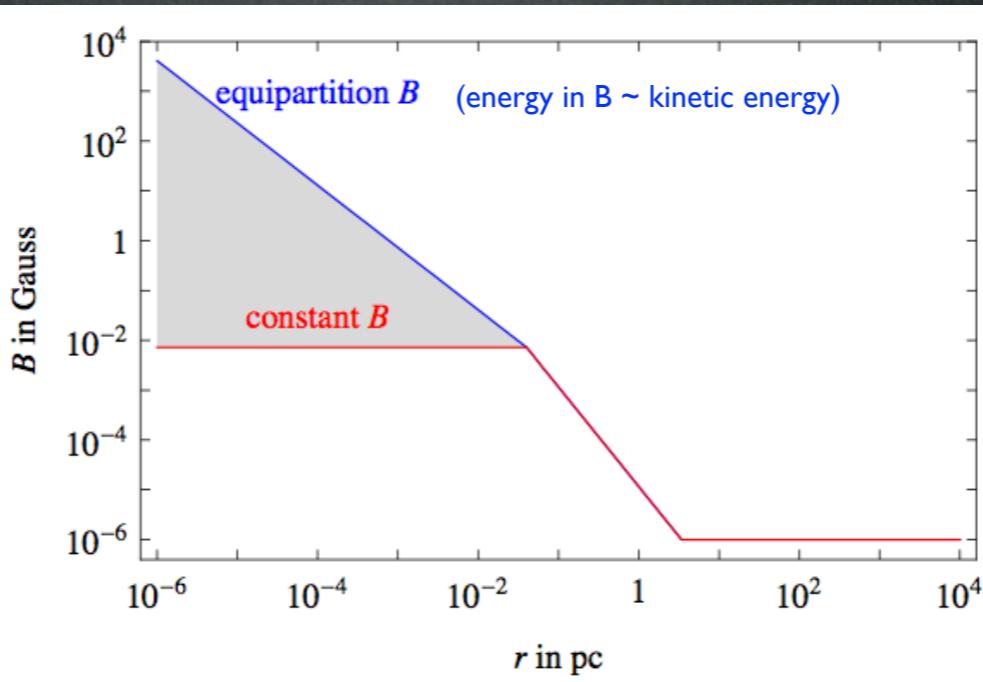
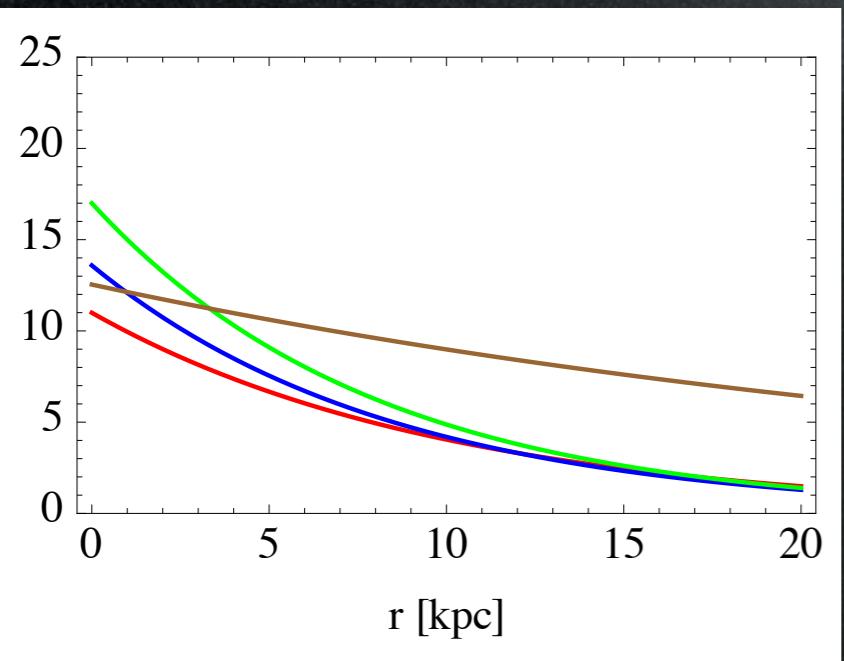
$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

spectral densities
of emitting e^\pm

synchr power emitted by individual electron:

$$\frac{dW_{\text{syn}}}{d\nu} = \frac{\sqrt{3}}{6\pi} \frac{e^3 B}{m_e} F\left(\frac{\nu}{\nu_{\text{syn}}}\right) \quad F(x) = x \int_x^\infty K_{5/3}(\xi) d\xi \approx \frac{8\pi}{9\sqrt{3}} \delta(x - 1/3)$$

$$\nu_{\text{syn}} = \frac{3eBp^2}{4\pi m_e^3}$$



Rule of thumb:

$$p_{e^\pm} \longrightarrow \nu_{\text{syn}}/3$$

$$\nu_{\text{syn}} = 4.2 \text{ MHz} \frac{B}{\text{G}} \left(\frac{p}{m_e} \right)^2$$

Synchrotron radiation

Spectrum of synchrotron emission from region $\Delta\Omega$:

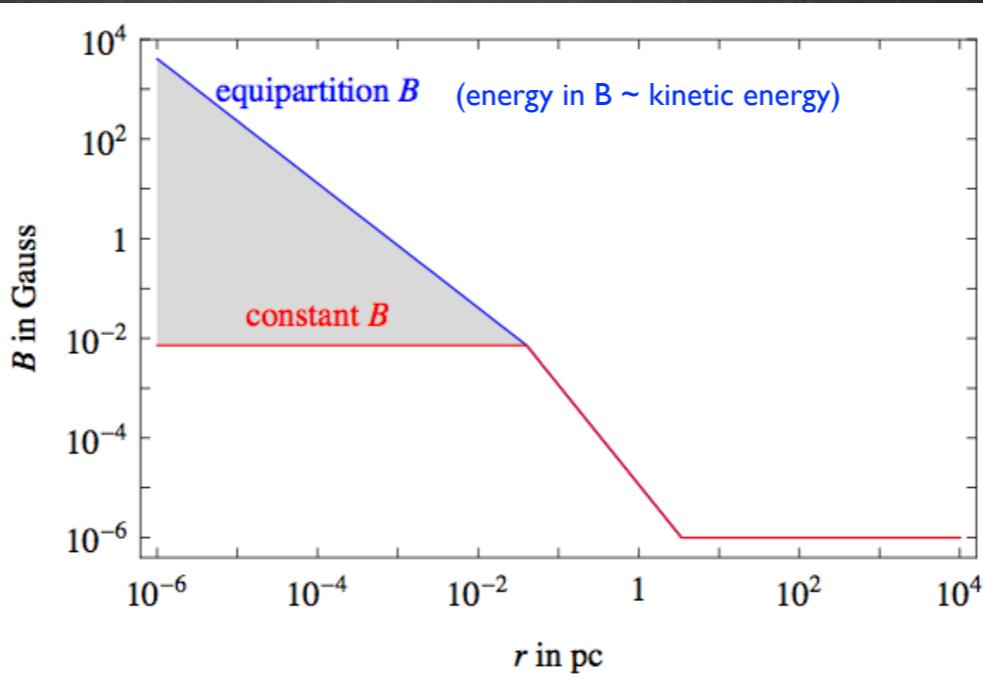
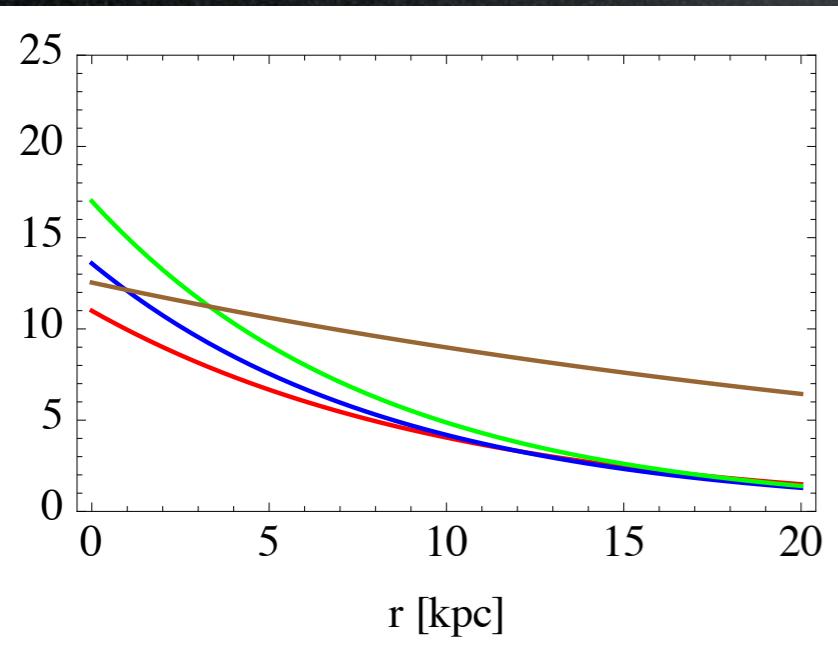
$$\frac{dW_{\text{syn}}}{d\nu d\Omega} = \frac{2}{4\pi} \int_{\text{l.o.s.}} ds \int dE_e \frac{dn_{e^\pm}}{dE_e} \frac{dW_{\text{syn}}}{d\nu}$$

spectral densities
of emitting e^\pm

synchr power emitted by individual electron:

$$\frac{dW_{\text{syn}}}{d\nu} = \frac{\sqrt{3}}{6\pi} \frac{e^3 B}{m_e} F\left(\frac{\nu}{\nu_{\text{syn}}}\right) \quad F(x) = x \int_x^\infty K_{5/3}(\xi) d\xi \approx \frac{8\pi}{9\sqrt{3}} \delta(x - 1/3)$$

$$\nu_{\text{syn}} = \frac{3eBp^2}{4\pi m_e^3}$$



Rule of thumb:

$$p_{e^\pm} \longrightarrow \nu_{\text{syn}}/3$$

$$\nu_{\text{syn}} = 4.2 \text{ MHz} \frac{B}{\text{G}} \left(\frac{p}{m_e} \right)^2$$

'Prompt' gamma rays

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Cocella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection".

arXiv 10
Erratum:

Fluxes of Inverse Compton gamma rays:

Annihilation

Mathematica function: the file [ICAnn.m](#) provides the Inverse Compton halo functions $I_{\text{IC}}(E_g, E_\gamma, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

Such notebook contains also a Mathematica code bite to compute the resulting IC flux.

Decay

Mathematica function: the file [ICDec.m](#) provides the Inverse Compton halo functions $I_{\text{IC}}(E_g, E_\gamma, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

Such notebook contains also a Mathematica code bite to compute the resulting IC flux.

Fluxes of bremsstrahlung gamma rays:

Annihilation

Mathematica function: the file [BremAnn.m](#) provides the bremsstrahlung halo functions $I_{\text{brem}}(E_g, E_\gamma, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

Decay

Mathematica function: the file [BremDec.m](#) provides the bremsstrahlung halo functions $I_{\text{brem}}(E_g, E_\gamma, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

Fluxes of synchrotron radiation:

Annihilation

Mathematica function: the file [SynAnn.m](#) provides the synchrotron halo functions $I_{\text{syn}}(E_g, v, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

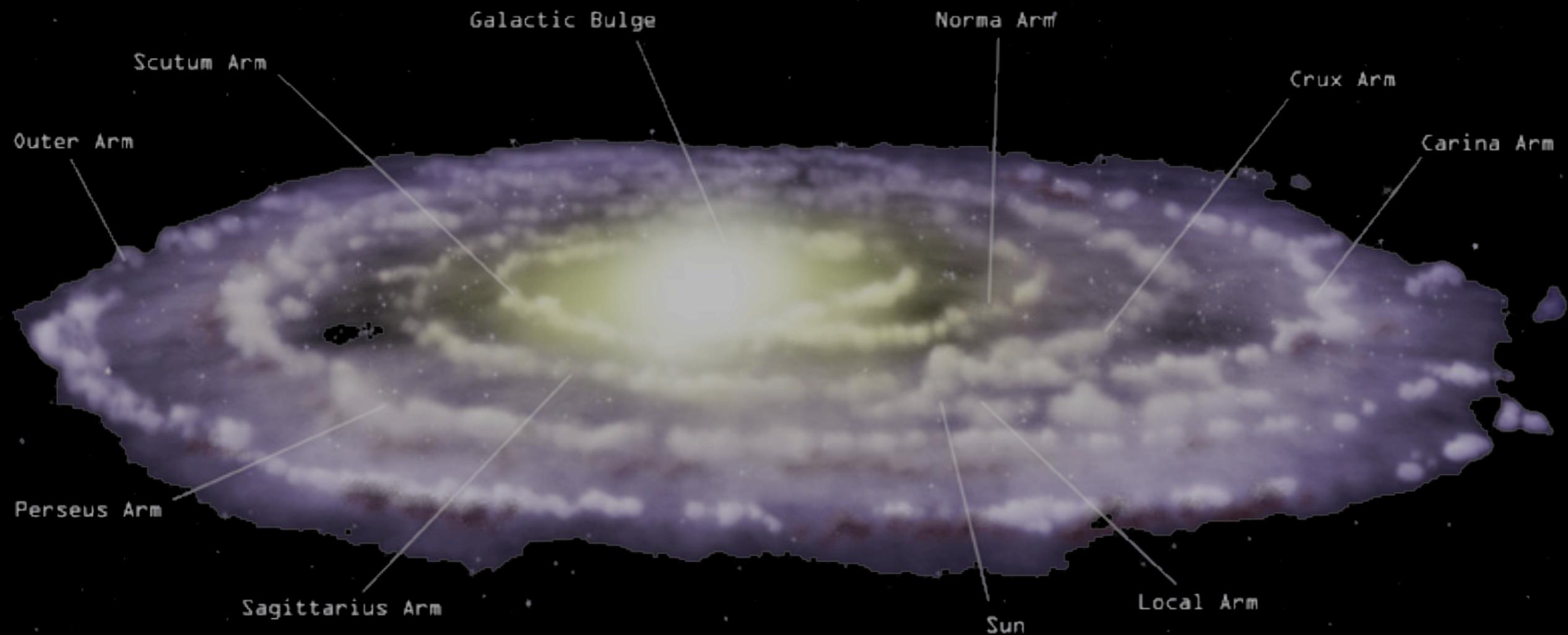
Decay

Mathematica function: the file [SynDec.m](#) provides the synchrotron halo functions $I_{\text{syn}}(E_g, v, l, b)$.

Refer to the notebook [Sample.nb](#) for usage.

Indirect Detection: basics

γ from outside the Galaxy



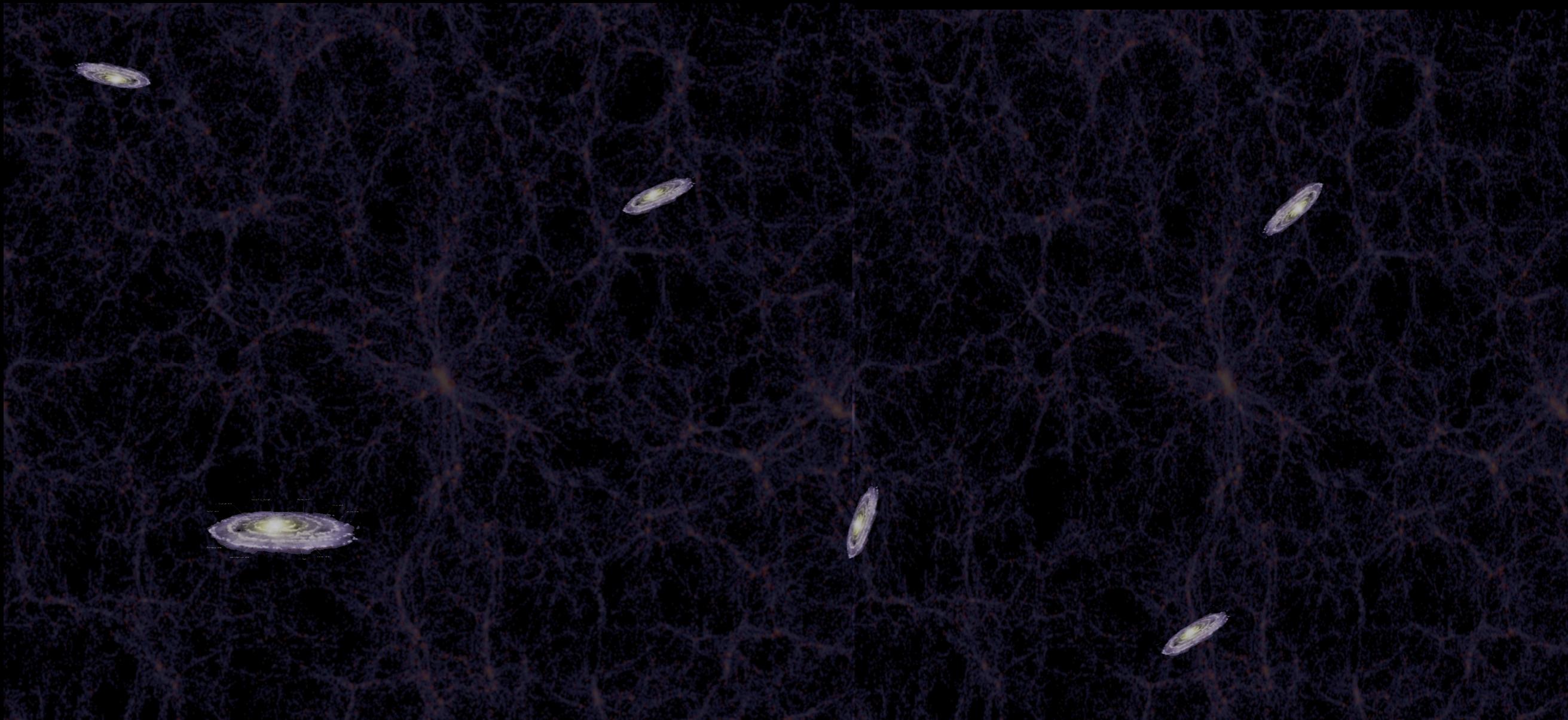
Indirect Detection: basics

γ from outside the Galaxy



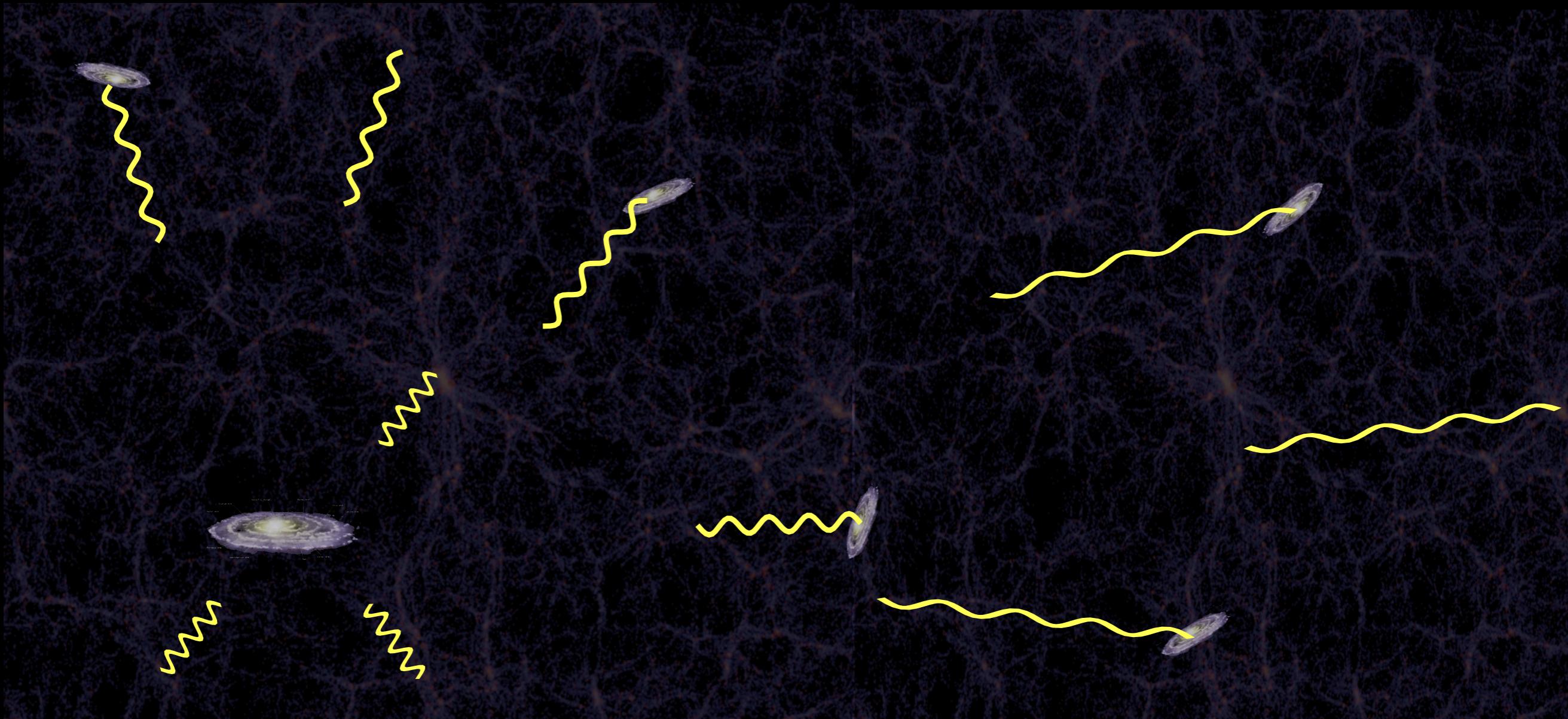
Indirect Detection: basics

γ from outside the Galaxy



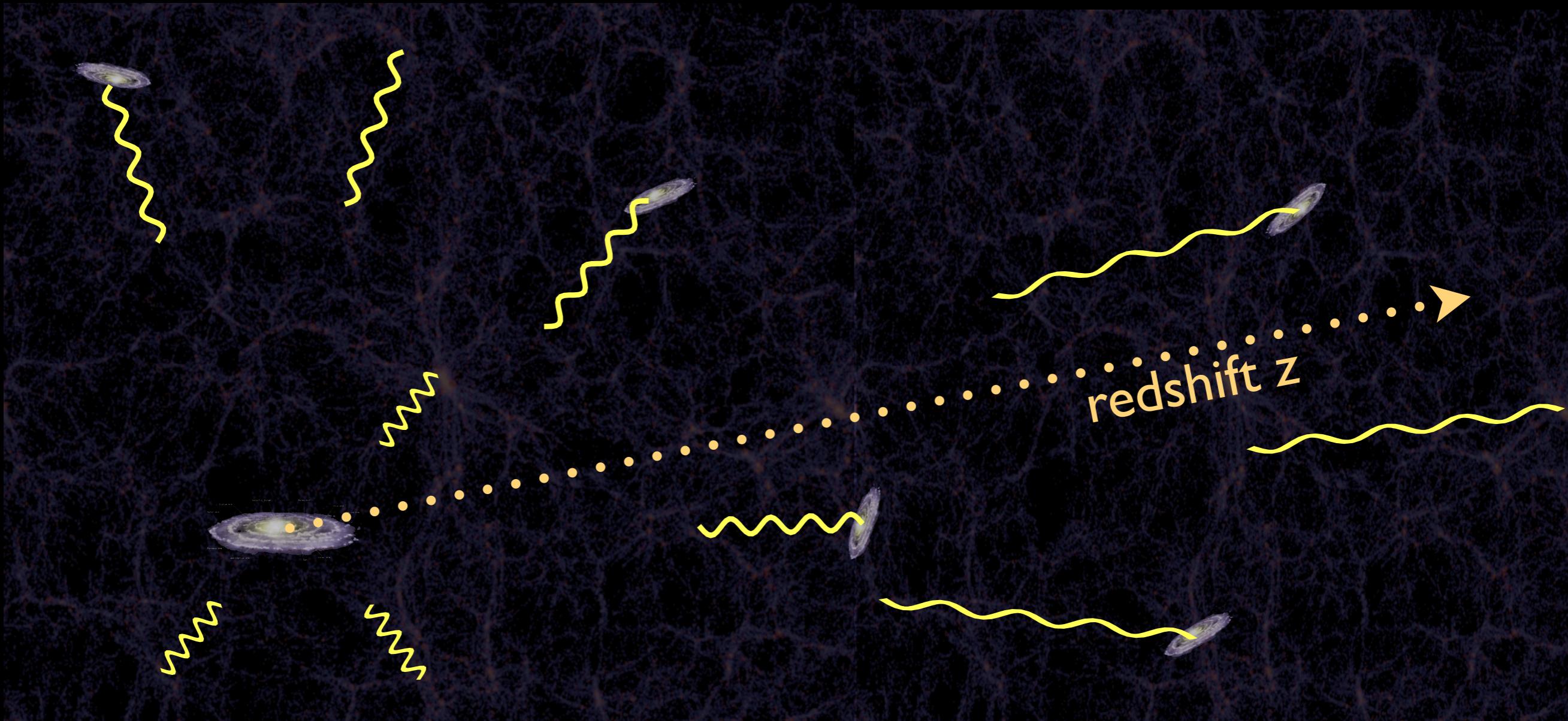
Indirect Detection: basics

γ from outside the Galaxy



Indirect Detection: basics

γ from outside the Galaxy



- isotropic flux of prompt and ICS gamma rays, integrated over z and r
- depends strongly on halo formation details and history

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

*flux emitted
at z'*

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

redshift due to the expansion of the Universe

flux emitted at z'

intervening absorption

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

cluster boost

$$B(z, M_{\min}) = 1 + \frac{\Delta_c}{3\bar{\rho}_{m,0}} \int_{M_{\min}}^\infty dM M \frac{dn}{dM}(M, z) f[c(M, z)]$$

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

cluster boost

$$B(z, M_{\min}) = 1 + \frac{\Delta_c}{3\bar{\rho}_{m,0}} \int_{M_{\min}}^\infty dM M \frac{dn}{dM}(M, z) f[c(M, z)]$$

minimal
halo
mass

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

cluster boost

$$B(z, M_{\min}) = 1 + \frac{\Delta_c}{3\bar{\rho}_{m,0}} \int_{M_{\min}}^\infty dM M \frac{dn}{dM}(M, z) f [c(M, z)]$$

minimal
halo
mass

halo mass fct
(number of halos
with mass M
at redshift z)

Extragalactic gamma rays

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

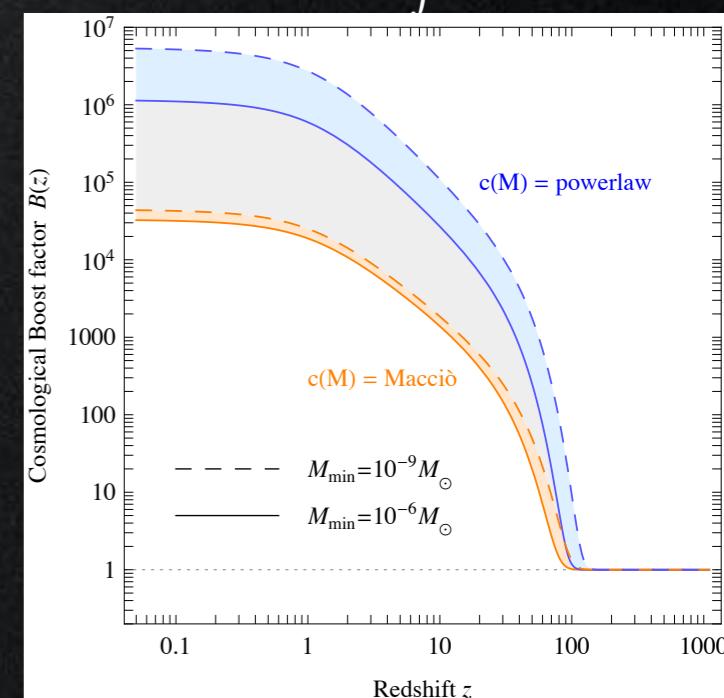
cluster boost

$$B(z, M_{\min}) = 1 + \frac{\Delta_c}{3\bar{\rho}_{m,0}} \int_{M_{\min}}^\infty dM M \frac{dn}{dM}(M, z) f[c(M, z)]$$

minimal halo mass

halo mass fct
(number of halos with mass M at redshift z)

concen-
tra-tion
function



Extragalactic gamma rays

absorption

Flux of ExGal γ :

$$\frac{d\Phi_{\text{EG}\gamma}}{dE_\gamma} = \frac{c}{E_\gamma} \int_0^\infty dz' \frac{1}{H(z')(1+z')^4} \frac{j_{\text{EG}\gamma}(E'_\gamma, z')}{4\pi} e^{-\tau(E'_\gamma, z')}$$

emissivities:

$$j_{\text{EG}\gamma}^{\text{prompt}}(E', z') = \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f \frac{dN_\gamma^f}{dE'_\gamma}$$

$$j_{\text{EG}\gamma}^{\text{IC}}(E', z') = 2 \int_{m_e}^{M_{\text{DM}}} dE_e \frac{\mathcal{P}_{\text{IC}}^{\text{CMB}}(E'_\gamma, E_e, z')}{b_{\text{IC}}^{\text{CMB}}(E_e, z')} \int_{E_e}^{M_{\text{DM}}} d\tilde{E}_e \frac{dN_e}{d\tilde{E}_e} \frac{1}{2} B(z') \left(\frac{\bar{\rho}(z')}{M_{\text{DM}}} \right)^2 \sum_f \langle \sigma v \rangle_f$$

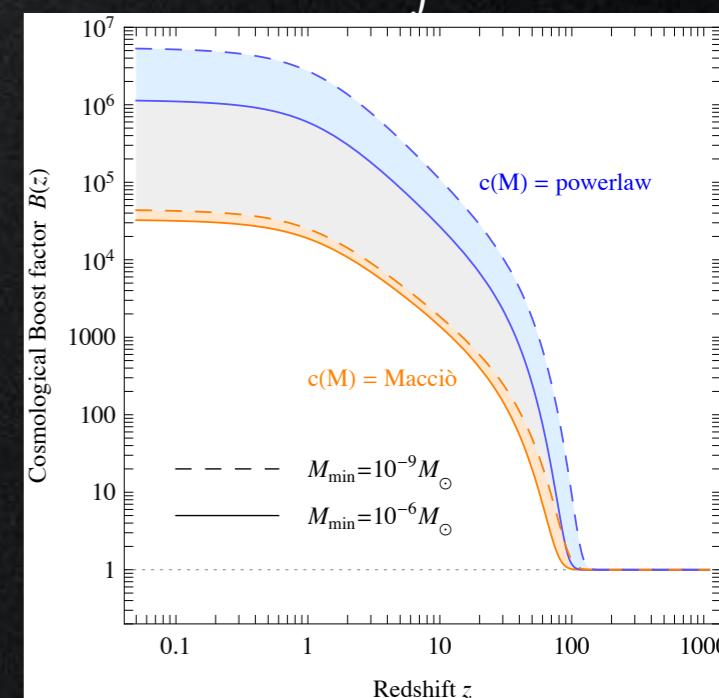
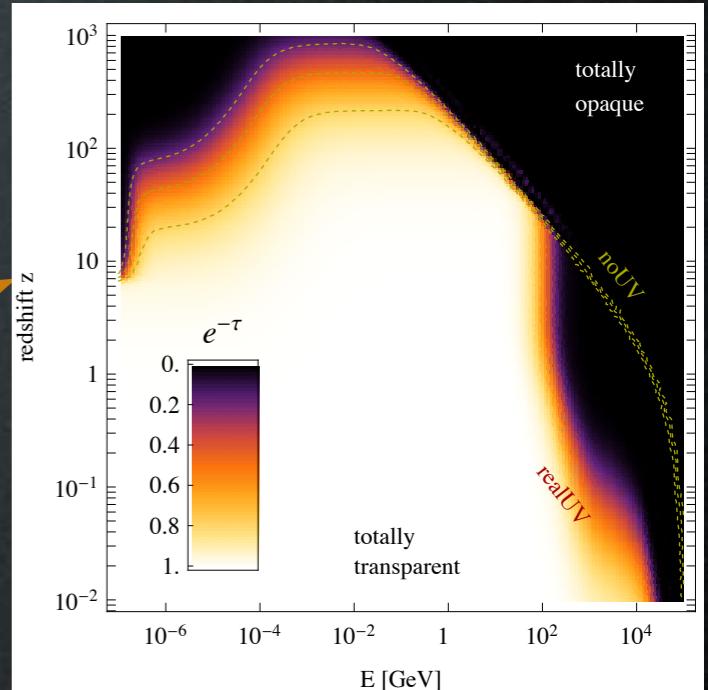
cluster boost

$$B(z, M_{\min}) = 1 + \frac{\Delta_c}{3\bar{\rho}_{m,0}} \int_{M_{\min}}^\infty dM M \frac{dn}{dM}(M, z) f[c(M, z)]$$

minimal halo mass

halo mass fct
(number of halos with mass M at redshift z)

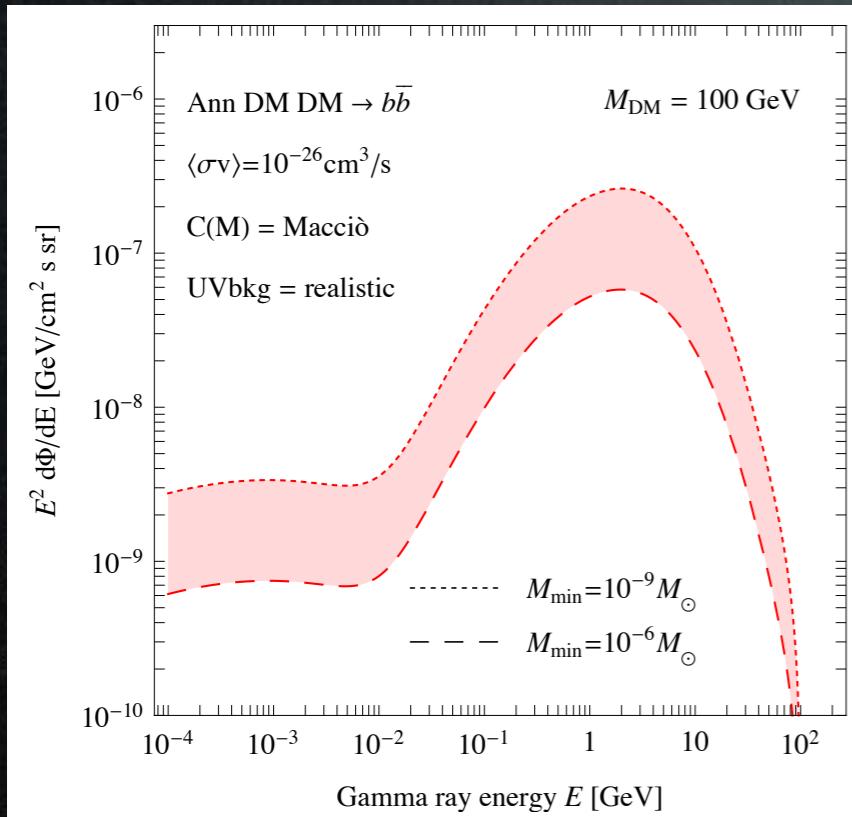
concen-
tra-tion
function



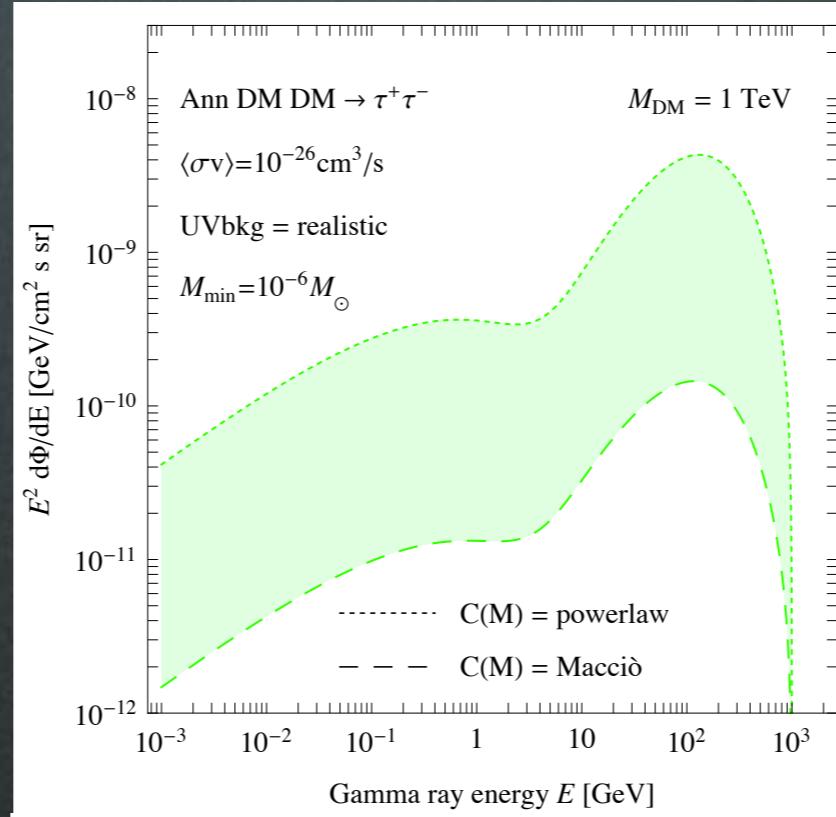
Extragalactic gamma rays

Varying:

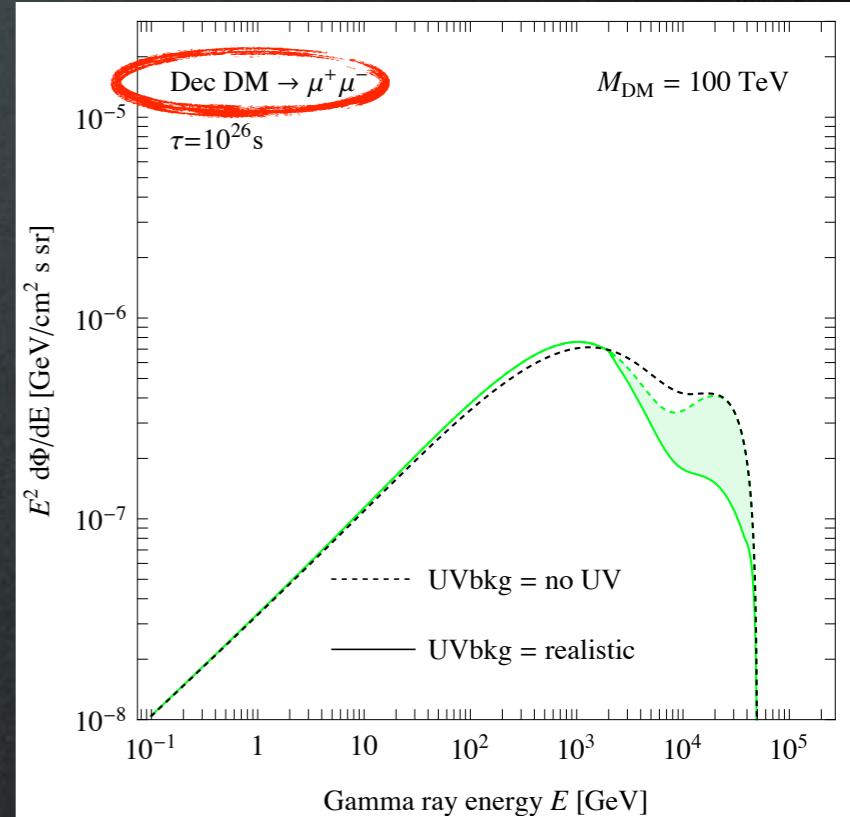
minimal halo mass



concentration param



UV absorption



Bottom line: ExGal gamma rays are affected by uncertainties of a few orders of magnitude due to cosmic history, for annihilations.
Much less for decay.

'Prompt' gamma rays

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Cocella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection",
arXiv 1012.4515, JCAP 1103 (2011) 051.
Erratum: JCAP 1210 (2012) E01.

Fluxes of extragalactic gamma rays:

Cosmological Boost factor B: Mathematica function `BoostF.m`, refer to the notebook [Sample.nb](#) for usage.

Optical Depth of the Universe: transparency factor $\text{Exp}[-\tau]$: Mathematica function `OpticalDepth.m`, refer to the notebook [Sample.nb](#) for usage.

Annihilation

Mathematica function: the file `EGgammaFluxAnn.m` provides the spectra $\log_{10} [\frac{d\Phi}{d\log_{10} E}]$.

Refer to the notebook [Sample.nb](#) for usage.

Decay

Mathematica function: the file `EGgammaFluxDec.m` provides the spectra $\log_{10} [\frac{d\Phi}{d\log_{10} E}]$.

Refer to the notebook [Sample.nb](#) for usage.

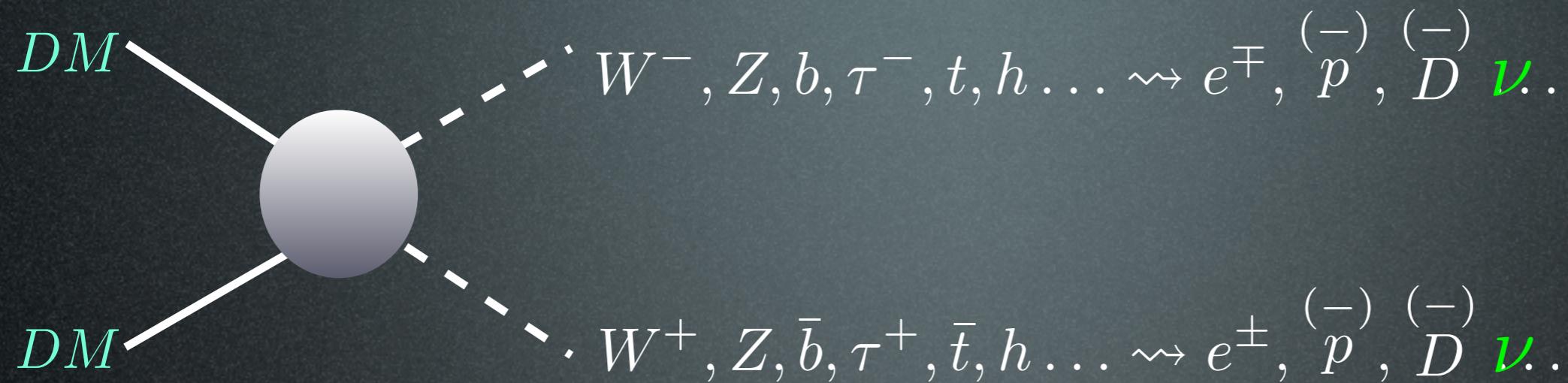
DM detection

direct detection

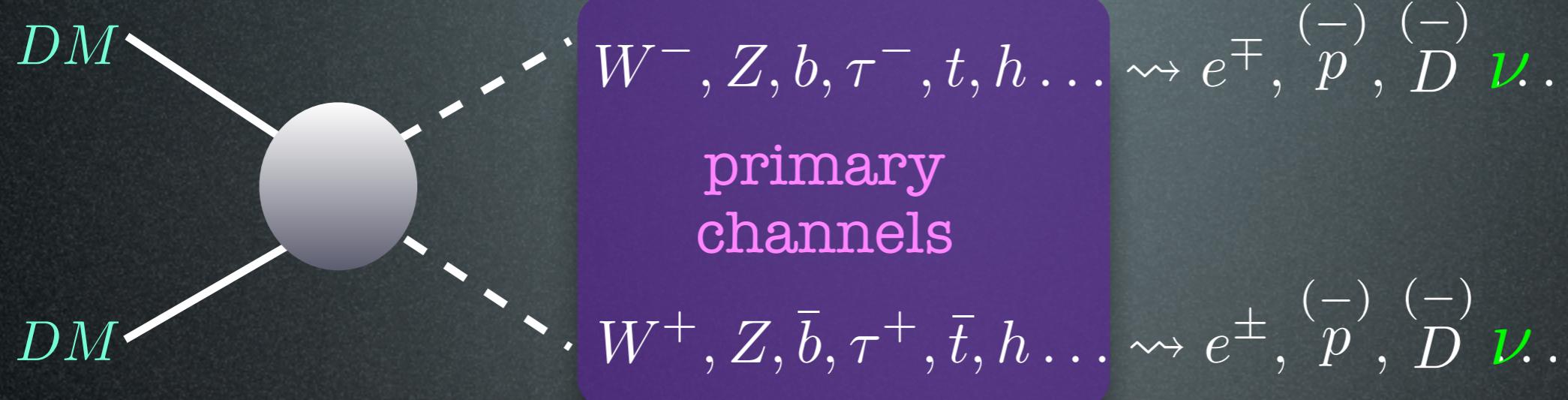
production at colliders

- indirect
 - γ from annihil in galactic center or halo
and from synchrotron emission Fermi, HESS, radio telescopes
 - e^+ from annihil in galactic halo or center PAMELA, ATIC, Fermi
 - \bar{p} from annihil in galactic halo or center
 - \bar{d} from annihil in galactic halo or center
 - $\nu, \bar{\nu}$ from annihil in galactic center/halo/Sun

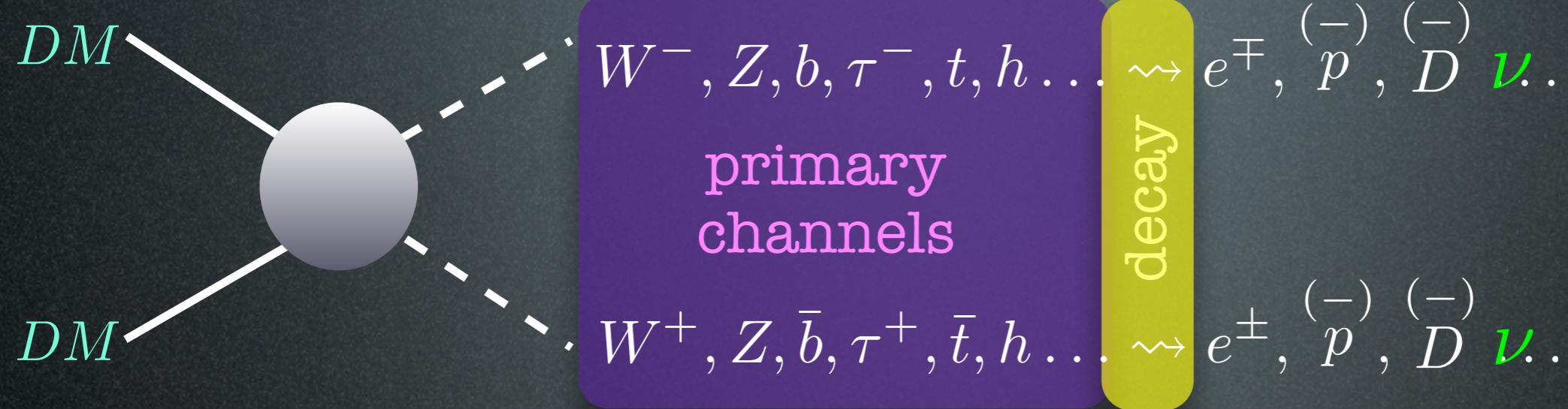
ID with neutrinos



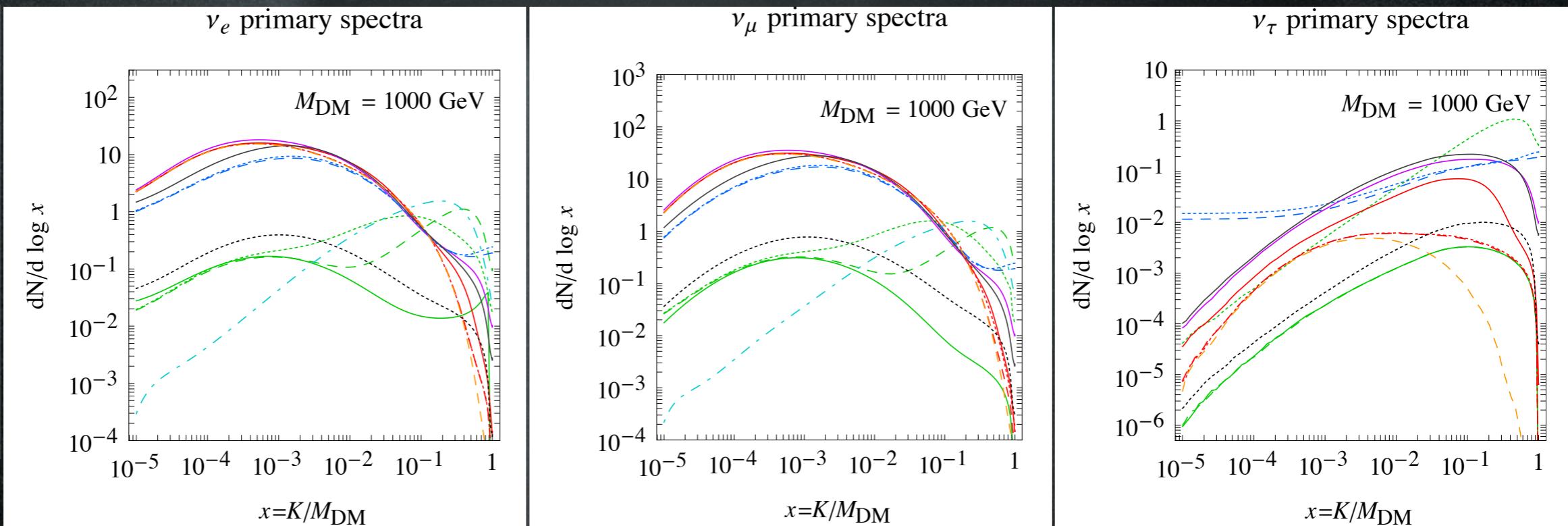
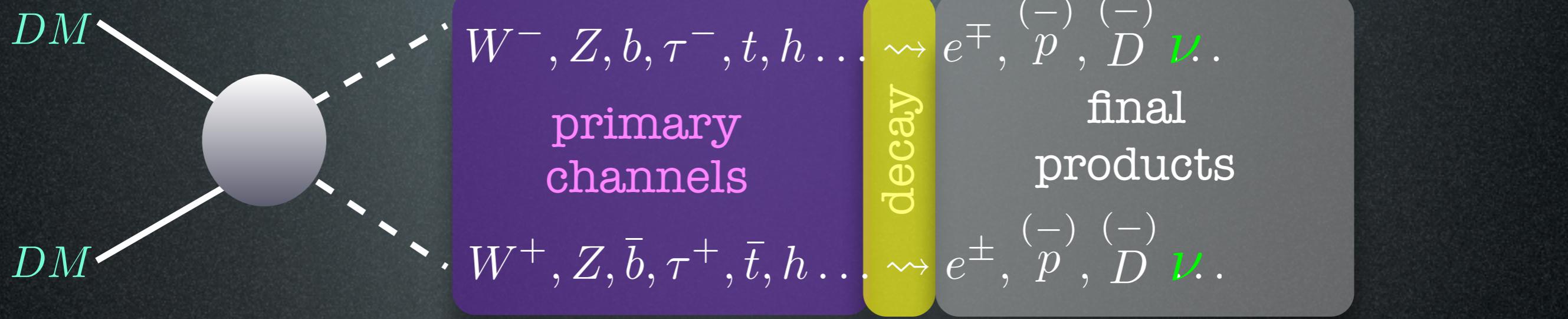
ID with neutrinos



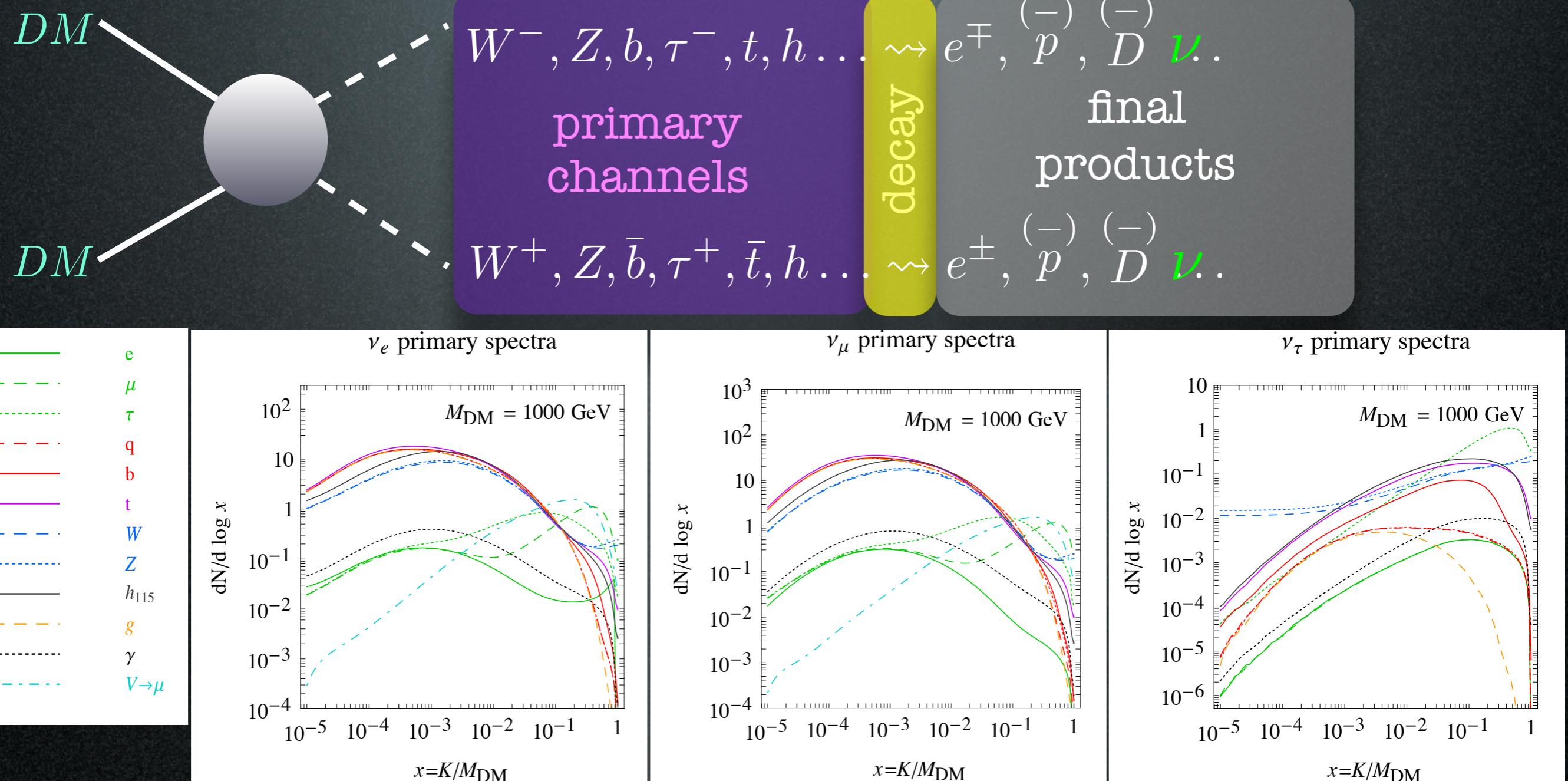
ID with neutrinos



ID with neutrinos



ID with neutrinos



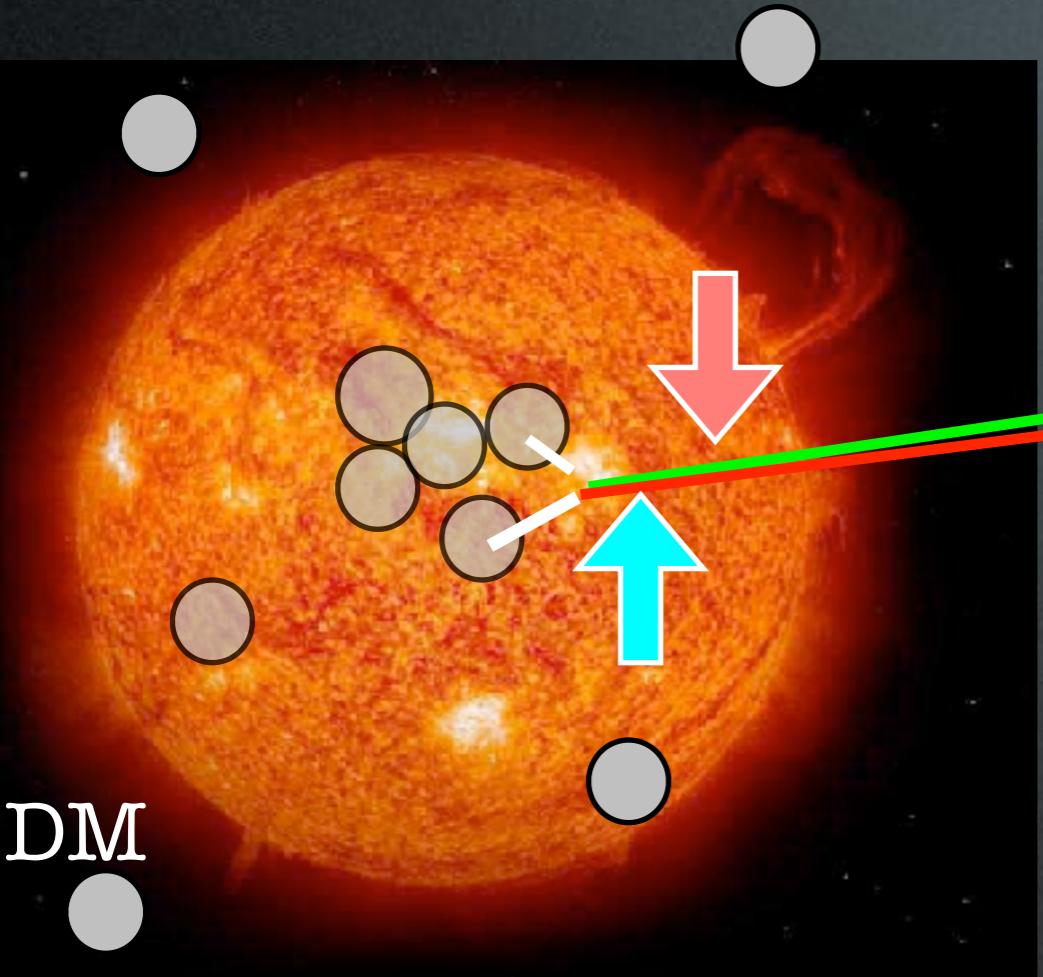
So what are the particle physics parameters?

1. Dark Matter mass
2. primary channel(s)
3. annihilation cross section σ_{ann}

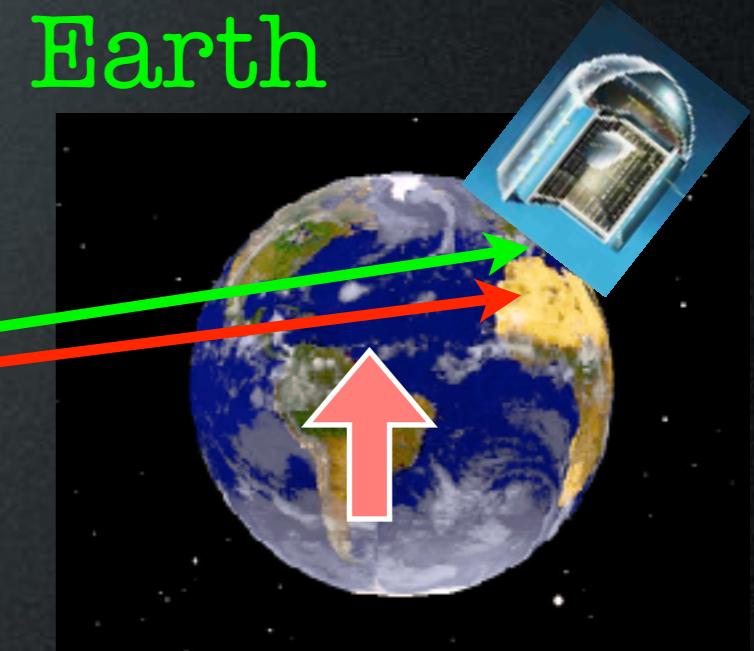
ID with neutrinos

ν from DM annihilations in the Sun

Sun



Earth

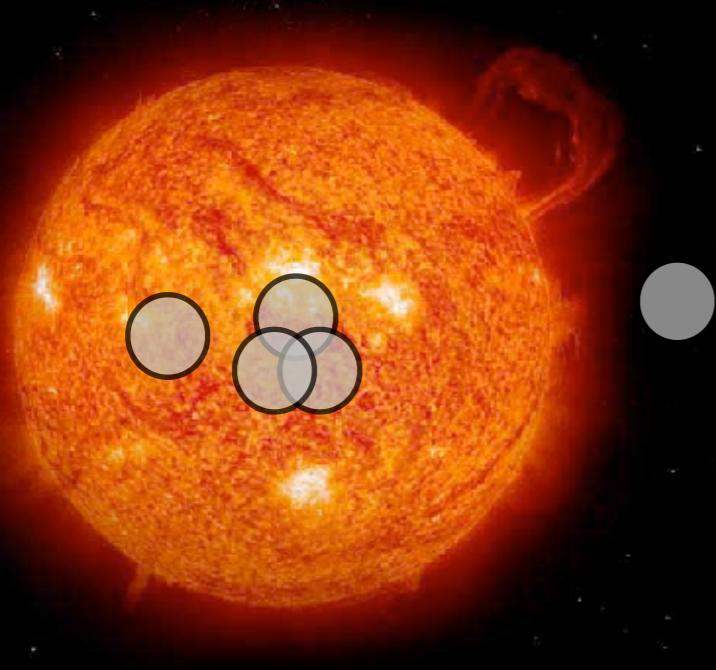


DM

Include oscillations + interactions:

- reshuffling of the 3 flavors
- distortions the spectra
- attenuations of the fluxes

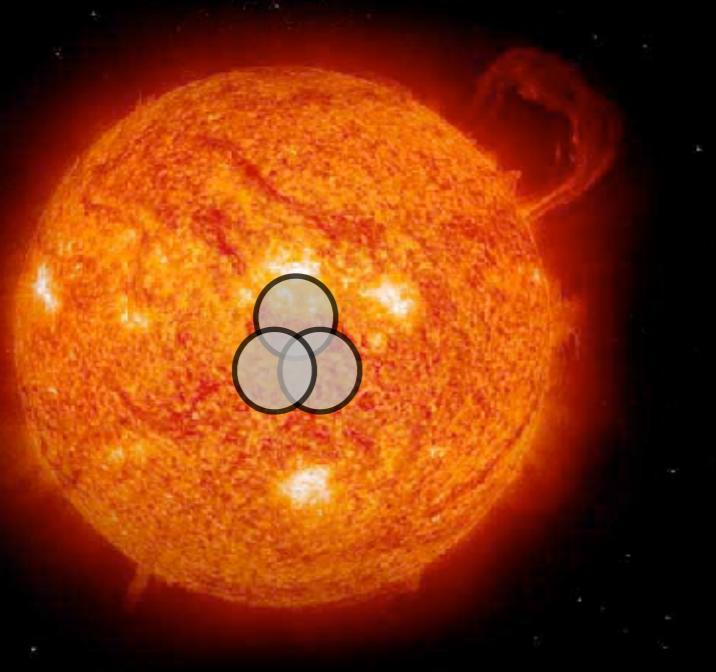
1. Capture & annihilation



basics: DM particle scatters with nuclei and loses energy
if $v_f < v_{\text{esc}}$ particle is gravitationally trapped
it spirals to center of body and accumulates
annihilates

$$\begin{aligned}v_{\text{halo}} &\simeq 270 \text{ km/s} \\v_{\text{esc}, \odot} &\simeq 620 \text{ km/s} \\v_{\text{esc}, \oplus} &\simeq 12 \text{ km/s}\end{aligned}$$

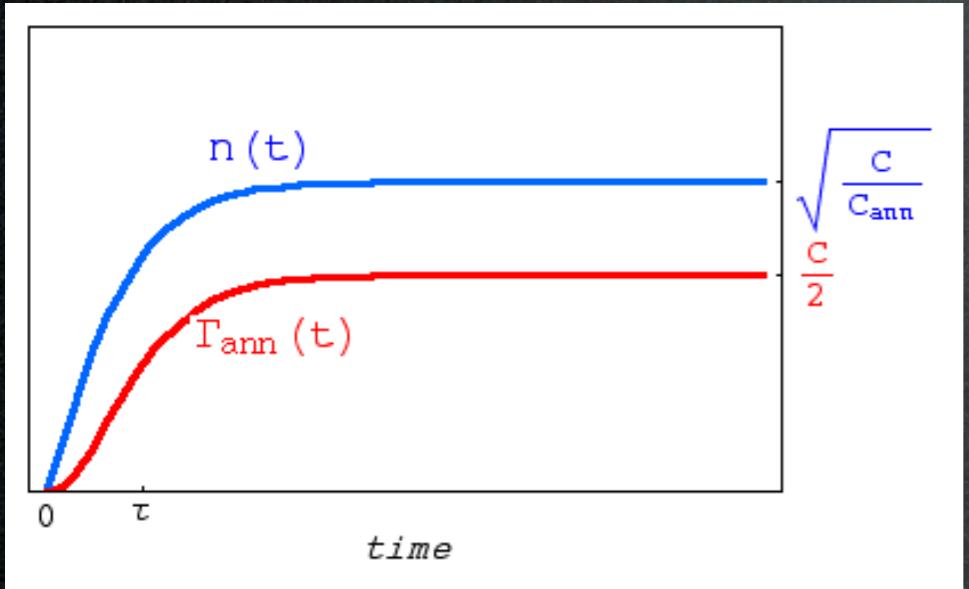
1. Capture & annihilation



basics: DM particle scatters with nuclei and loses energy
if $v_f < v_{\text{esc}}$ particle is gravitationally trapped
it spirals to center of body and accumulates
annihilates

$$\begin{aligned}v_{\text{halo}} &\simeq 270 \text{ km/s} \\v_{\text{esc},\odot} &\simeq 620 \text{ km/s} \\v_{\text{esc},\oplus} &\simeq 12 \text{ km/s}\end{aligned}$$

equilibrium attained:

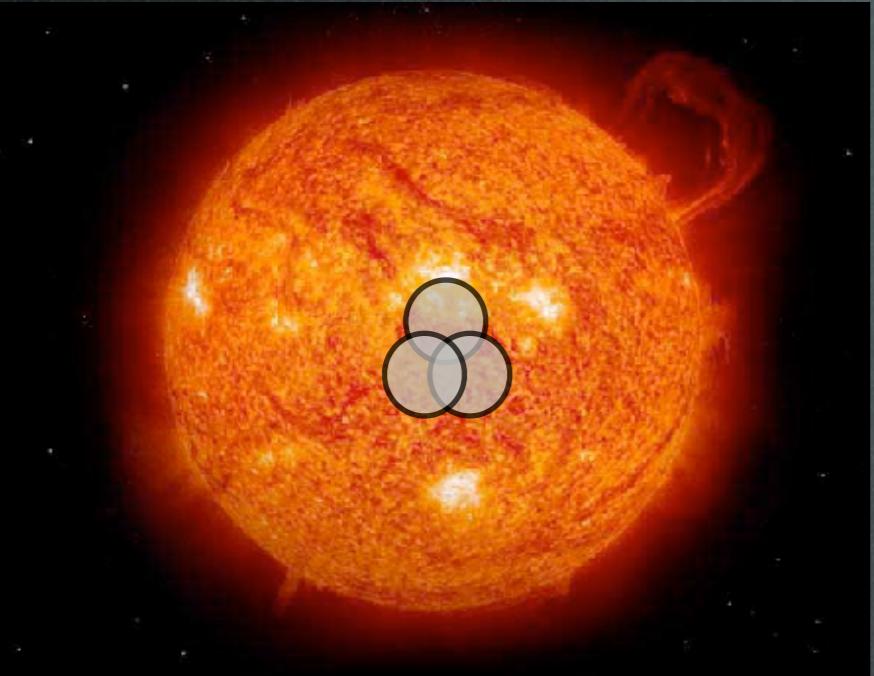


$$\dot{n} = \Gamma_{\text{capt}} - C_{\text{ann}} n^2$$

σ_N $\langle \sigma_{\text{ann}} v \rangle$

$$C_{\text{ann}} = \langle \sigma v \rangle \left(\frac{G_N M_{\text{DM}} \rho_{\odot}}{3 T_{\odot}} \right)^{3/2}$$
$$n(t) = \sqrt{\frac{\Gamma_{\text{capt}}}{C_{\text{ann}}}} \tanh \left(\frac{t}{\tau} \right)$$
$$\tau = \frac{1}{\sqrt{\Gamma_{\text{capt}} C_{\text{ann}}}}$$
$$\Gamma_{\text{ann}}(t) = \frac{\Gamma_{\text{capt}}}{2} \tanh^2 \left(\frac{t}{\tau} \right)$$

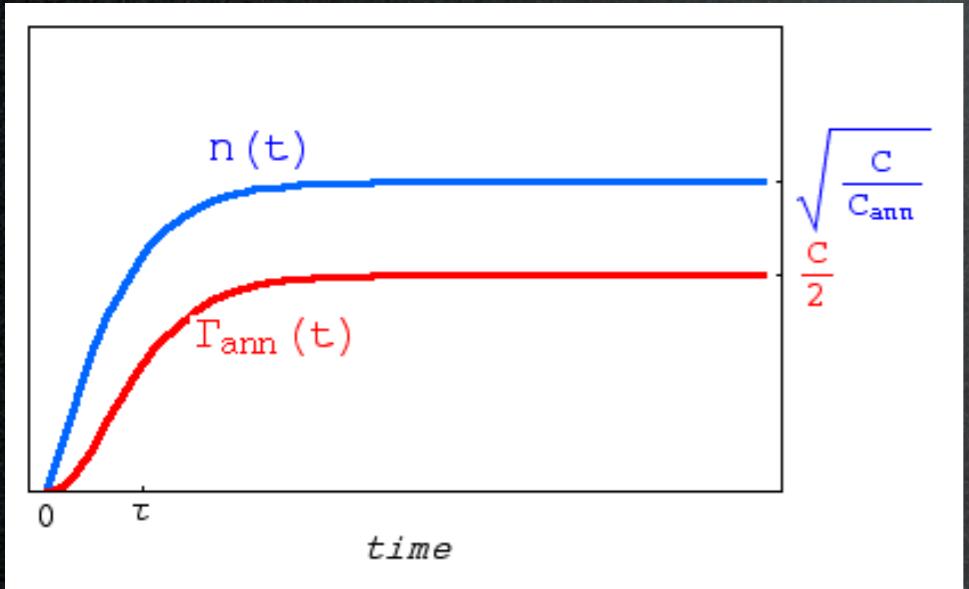
1. Capture & annihilation



basics: DM particle scatters with nuclei and loses energy
if $v_f < v_{\text{esc}}$ particle is gravitationally trapped
it spirals to center of body and accumulates
annihilates

$$\begin{aligned}v_{\text{halo}} &\simeq 270 \text{ km/s} \\v_{\text{esc},\odot} &\simeq 620 \text{ km/s} \\v_{\text{esc},\oplus} &\simeq 12 \text{ km/s}\end{aligned}$$

equilibrium attained:



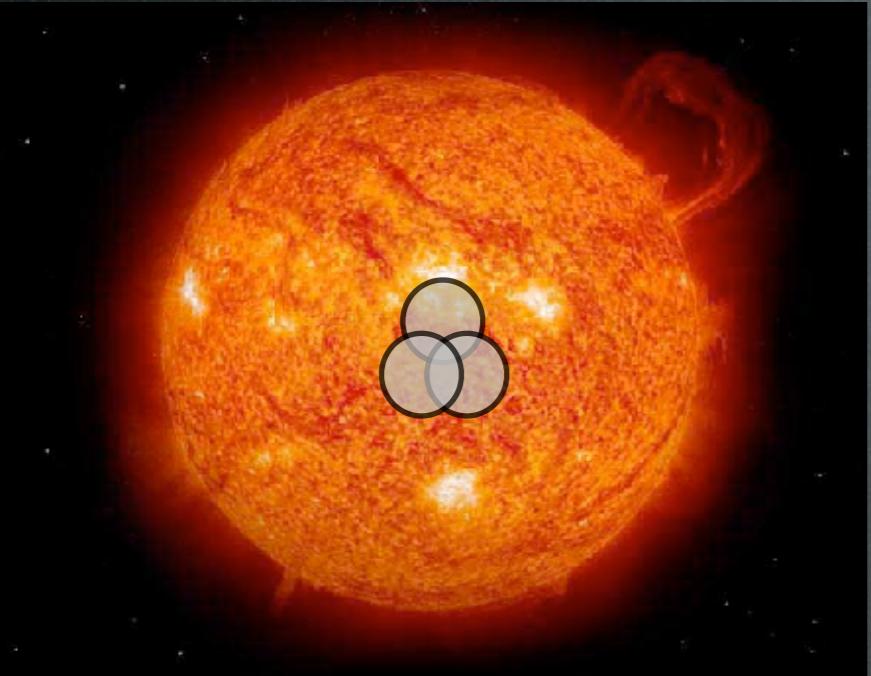
$$\dot{n} = \Gamma_{\text{capt}} - C_{\text{ann}} n^2$$

σ_N $\langle \sigma_{\text{ann}} v \rangle$

$$C_{\text{ann}} = \langle \sigma v \rangle \left(\frac{G_N M_{\text{DM}} \rho_{\odot}}{3 T_{\odot}} \right)^{3/2}$$
$$n(t) = \sqrt{\frac{\Gamma_{\text{capt}}}{C_{\text{ann}}}} \tanh \left(\frac{t}{\tau} \right)$$
$$\tau = \frac{1}{\sqrt{\Gamma_{\text{capt}} C_{\text{ann}}}}$$
$$\Gamma_{\text{ann}}(t) = \frac{\Gamma_{\text{capt}}}{2} \tanh^2 \left(\frac{t}{\tau} \right)$$

The main physical parameter is: σ_N (DM-nucleon scattering cross section)

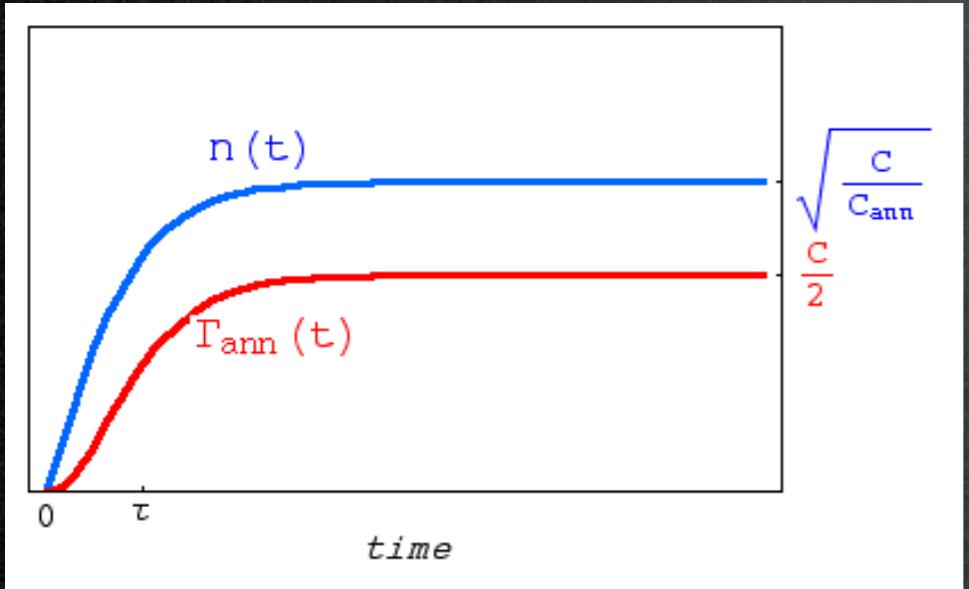
1. Capture & annihilation



basics: DM particle scatters with nuclei and loses energy
if $v_f < v_{\text{esc}}$ particle is gravitationally trapped
it spirals to center of body and accumulates
annihilates

$$\begin{aligned}v_{\text{halo}} &\simeq 270 \text{ km/s} \\v_{\text{esc},\odot} &\simeq 620 \text{ km/s} \\v_{\text{esc},\oplus} &\simeq 12 \text{ km/s}\end{aligned}$$

equilibrium attained:



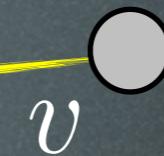
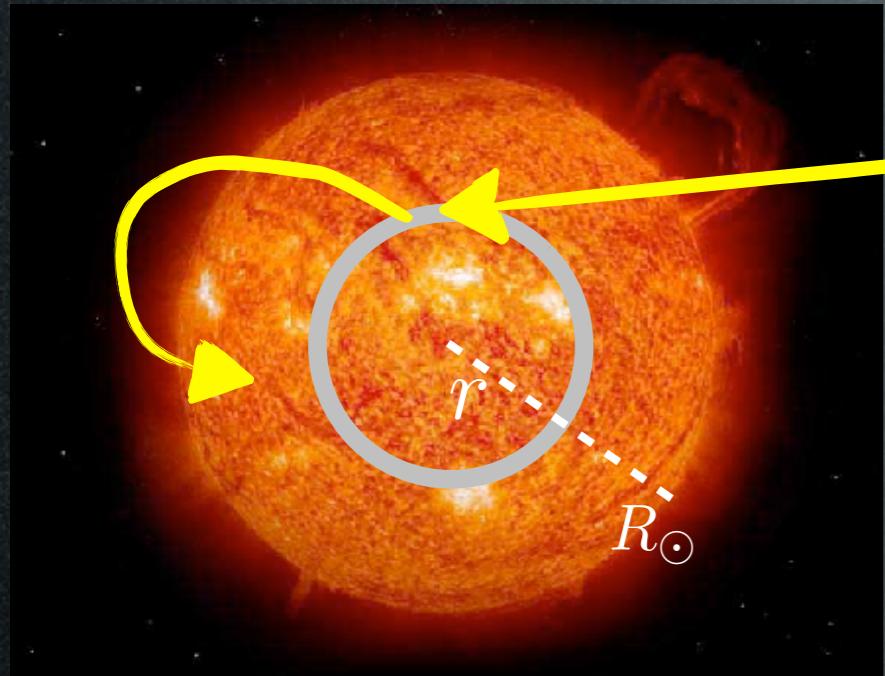
$$\dot{n} = \Gamma_{\text{capt}} - C_{\text{ann}} n^2$$

σ_N $\langle \sigma_{\text{ann}} v \rangle$

$$C_{\text{ann}} = \langle \sigma v \rangle \left(\frac{G_N M_{\text{DM}} \rho_{\odot}}{3 T_{\odot}} \right)^{3/2}$$
$$n(t) = \sqrt{\frac{\Gamma_{\text{capt}}}{C_{\text{ann}}}} \tanh \left(\frac{t}{\tau} \right) \quad \tau = \frac{1}{\sqrt{\Gamma_{\text{capt}} C_{\text{ann}}}}$$
$$\Gamma_{\text{ann}}(t) = \frac{\Gamma_{\text{capt}}}{2} \tanh^2 \left(\frac{t}{\tau} \right) \rightarrow \frac{\Gamma_{\text{capt}}}{2}$$

The main physical parameter is: σ_N (DM-nucleon scattering cross section)

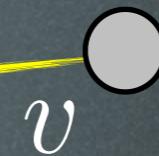
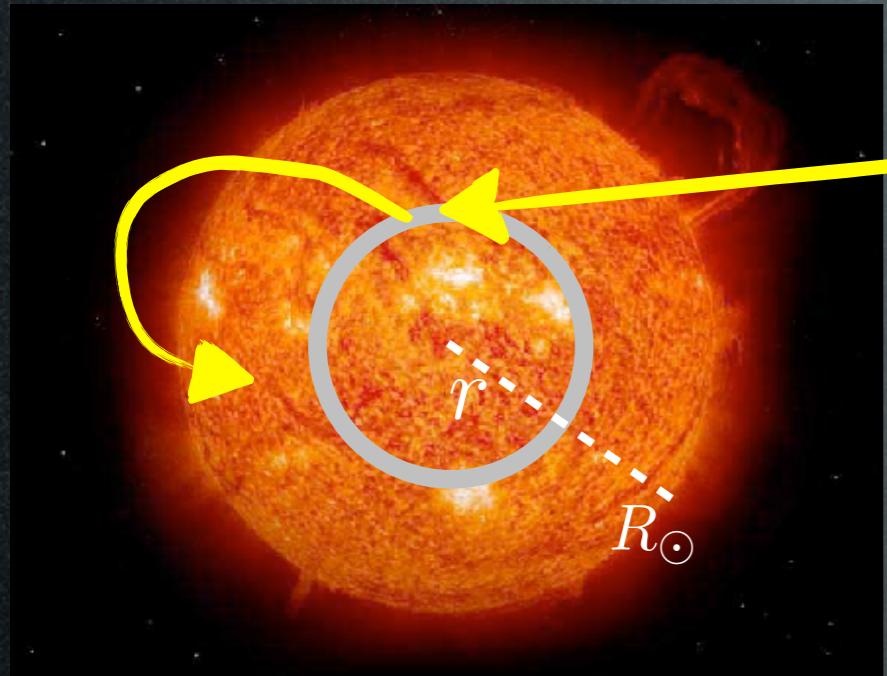
1. Capture & annihilation



A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_{\odot}} dr \ 4\pi r^2 \ n_i(r) \int_0^{\infty} dv \ 4\pi v^2 f_{\odot}(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

1. Capture & annihilation

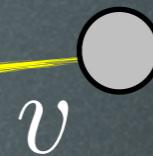
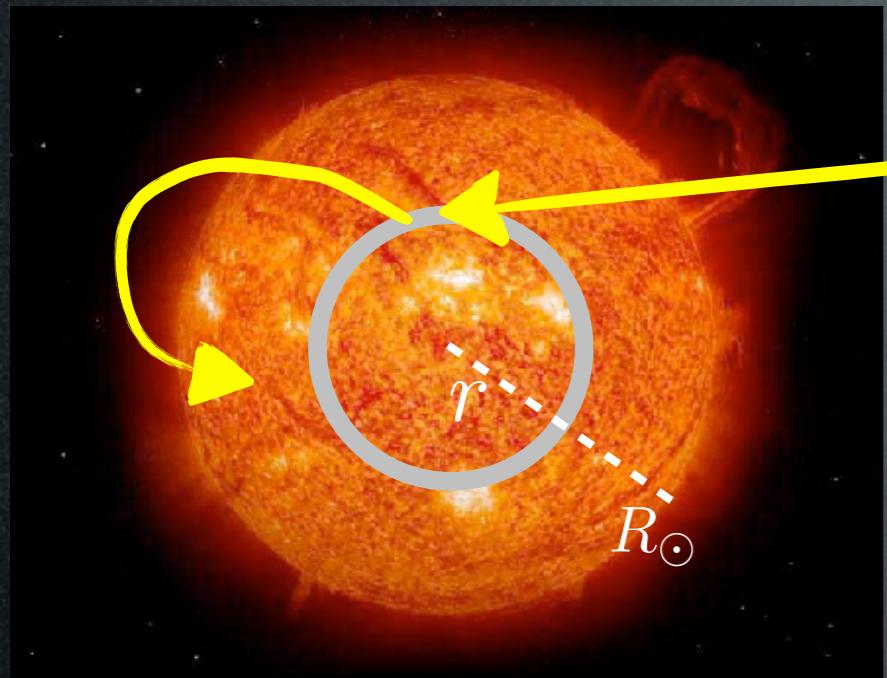


A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 \ n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM
number
density

1. Capture & annihilation



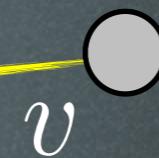
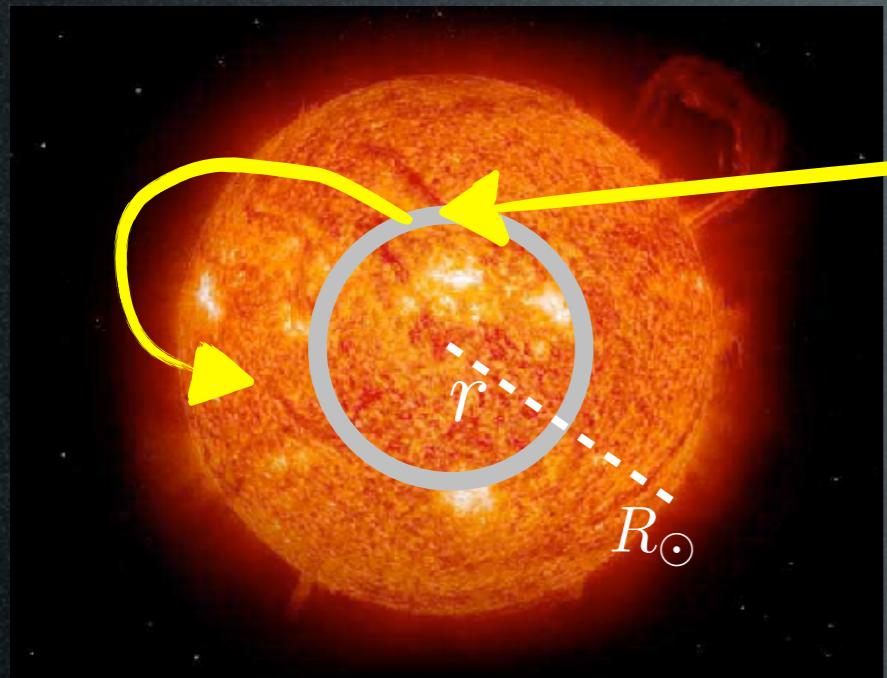
A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 \ n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM number density

scattering cross section on element **i**

1. Capture & annihilation



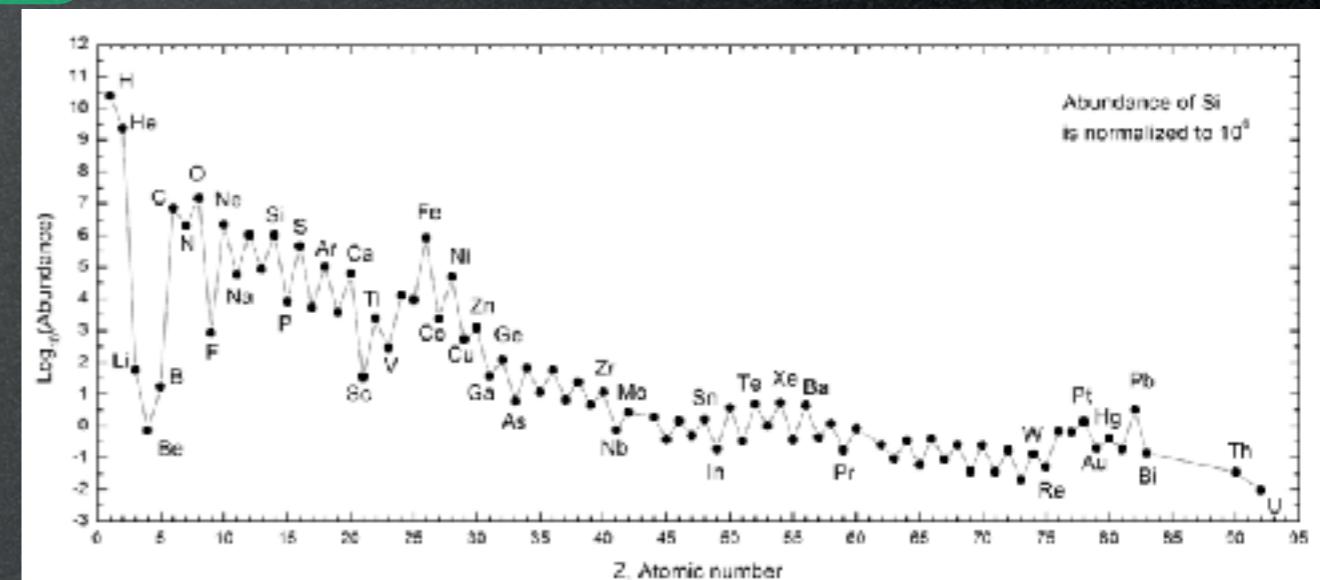
A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

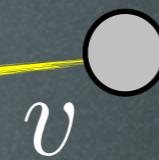
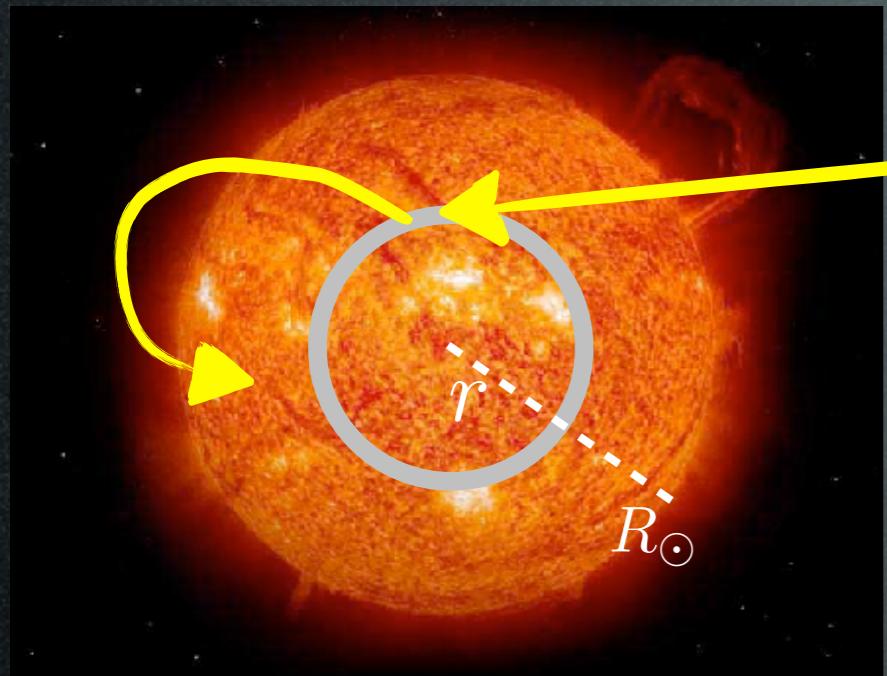
DM number density

scattering cross section on element **i**

number density of element **i**



1. Capture & annihilation



A.Gould 1987, 1988, 1990

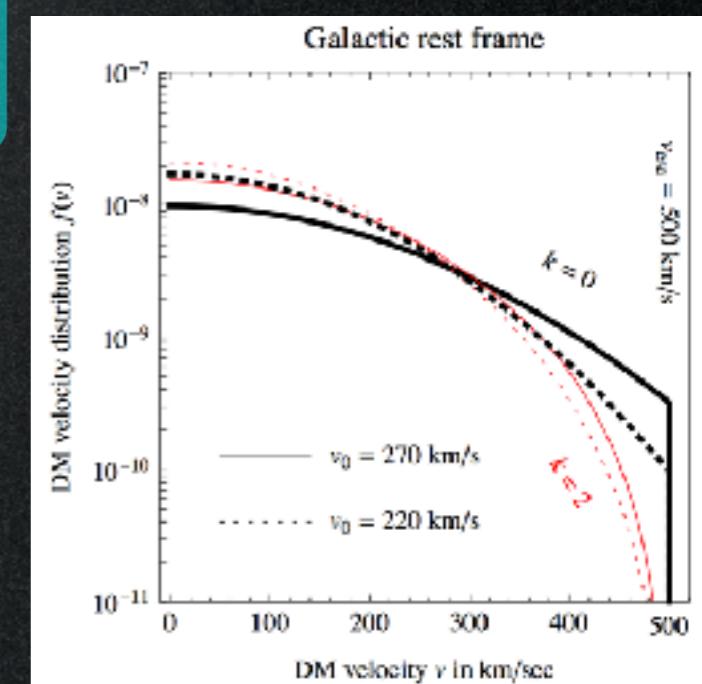
$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM number density

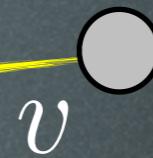
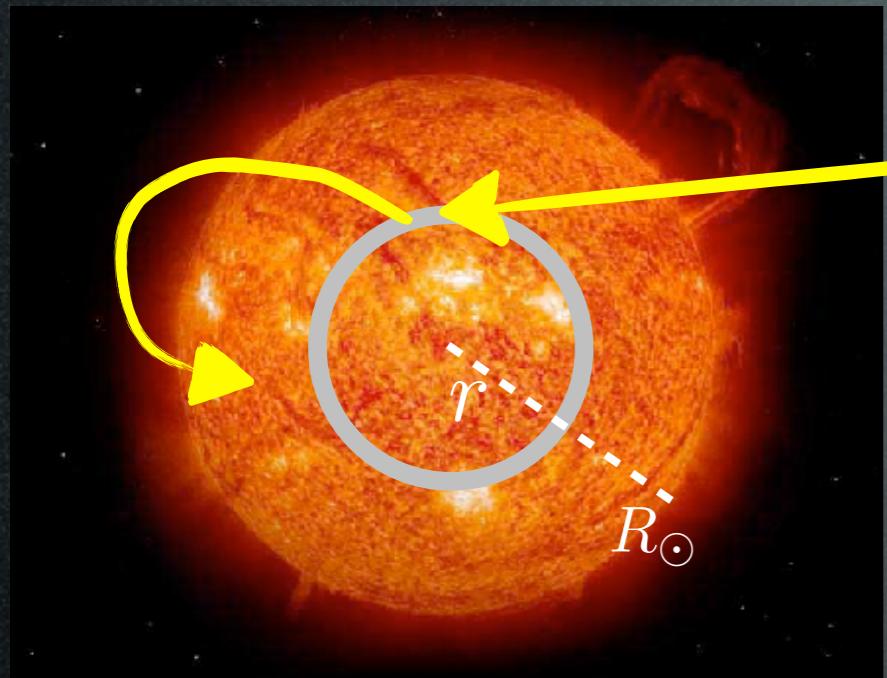
scattering cross section on element **i**

number density of element **i**

velocity distribution
(in solar frame,
without Sun's gravity)



1. Capture & annihilation



A.Gould 1987, 1988, 1990

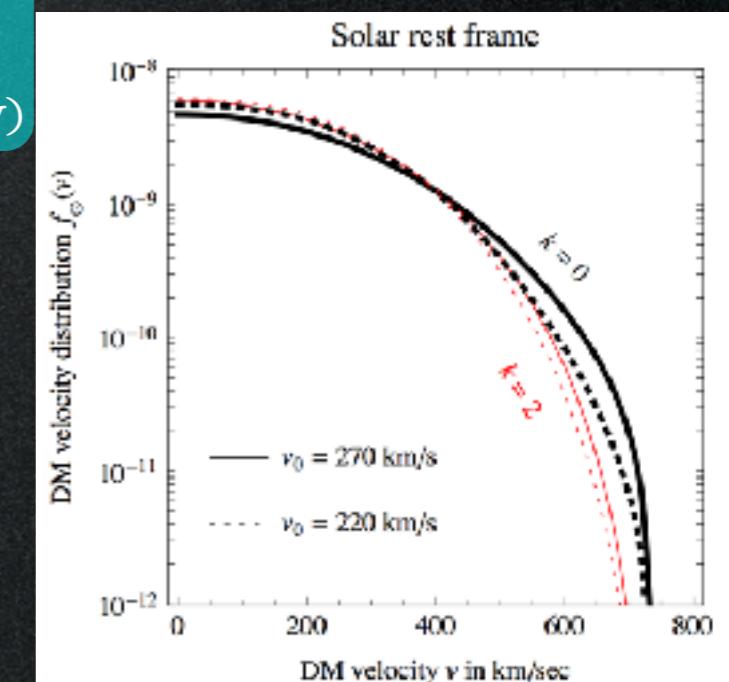
$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM number density

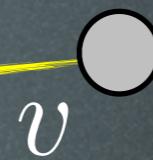
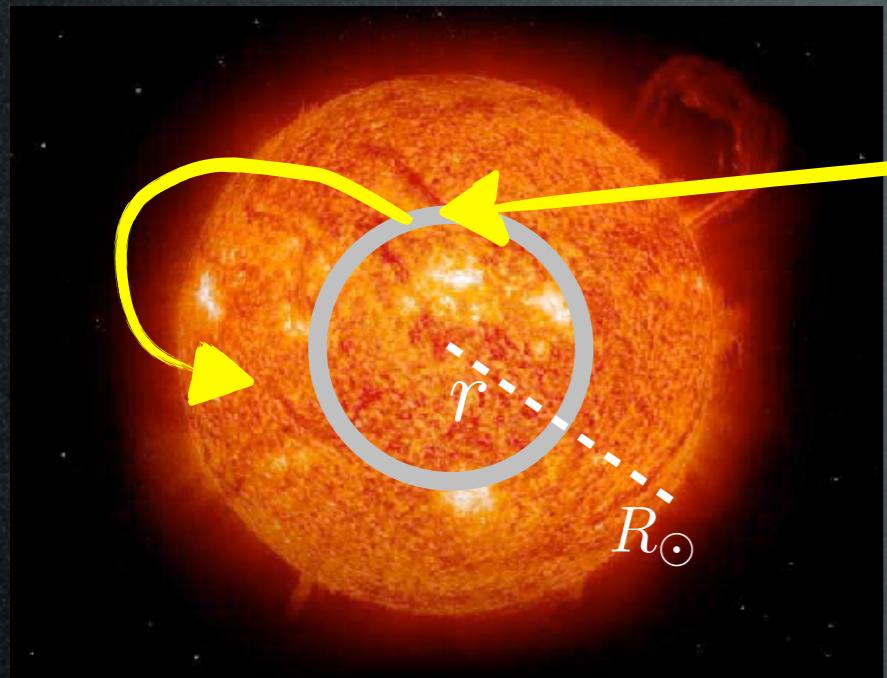
scattering cross section on element **i**

number density of element **i**

velocity distribution (in solar frame, without Sun's gravity)



1. Capture & annihilation



A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM number density

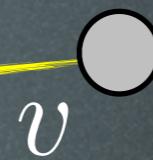
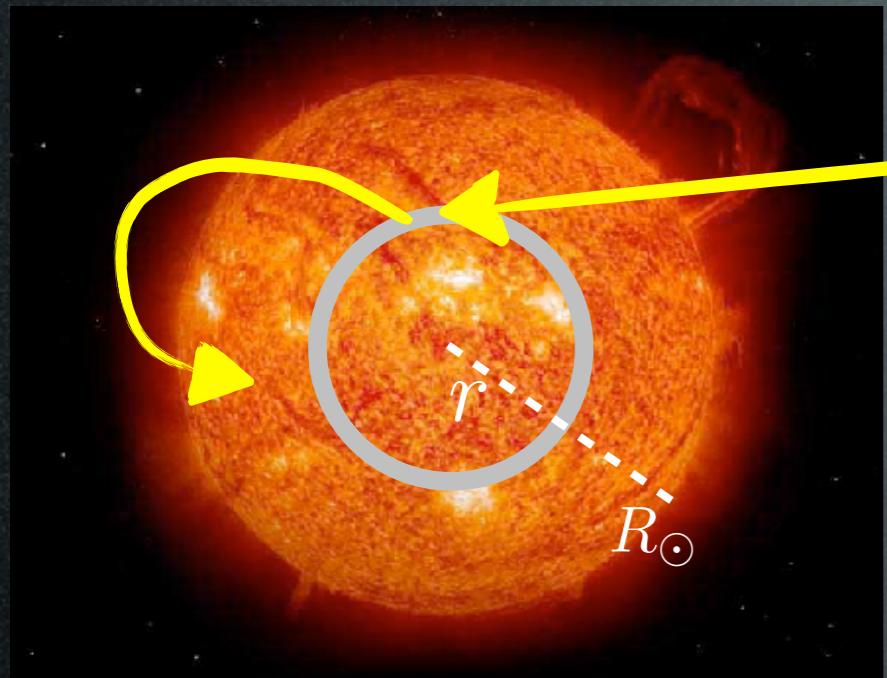
scattering cross section on element **i**

number density of element **i**

velocity distribution
(in solar frame,
without Sun's gravity)

effect of solar gravity

1. Capture & annihilation



A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

DM number density

scattering cross section on element **i**

number density of element **i**

velocity distribution (in solar frame, without Sun's gravity)

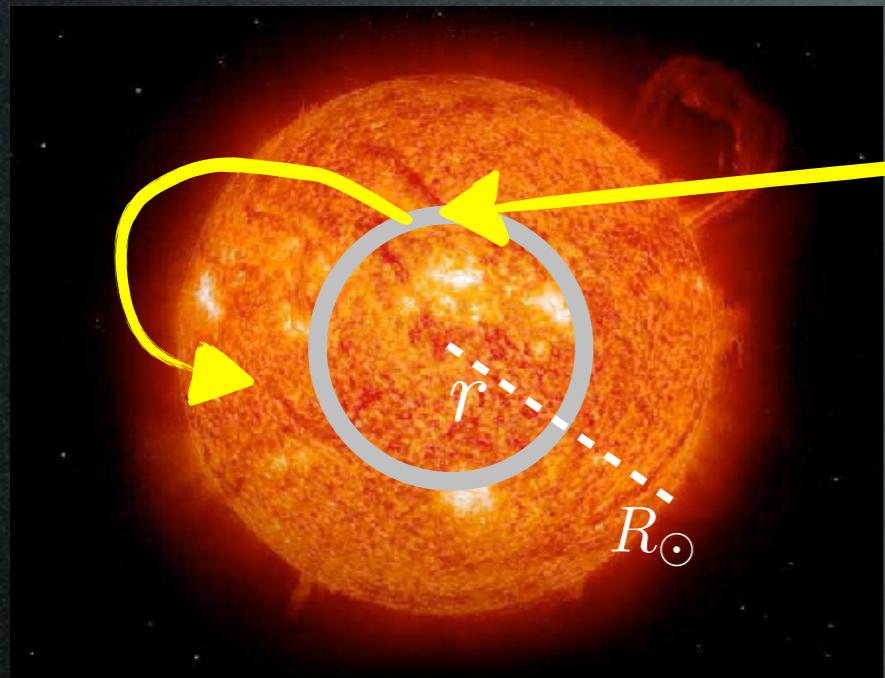
effect of solar gravity

scattering probability:

$$\wp_i(v, v_{\odot \text{esc}}) = \max \left(0, 1 - \frac{\Delta_{\min}}{\Delta_{\max}} \right)$$

$$\Delta_{\max} = \frac{4 m_i M_{\text{DM}}}{(M_{\text{DM}} + m_i)^2} \quad \Delta_{\min} = \frac{v^2}{v^2 + v_{\odot \text{esc}}^2}$$

1. Capture & annihilation



A.Gould 1987, 1988, 1990

$$\Gamma_{\text{capt}} = \frac{\rho_{\text{DM}}}{M_{\text{DM}}} \sum_i \sigma_i \int_0^{R_\odot} dr \ 4\pi r^2 n_i(r) \int_0^\infty dv \ 4\pi v^2 f_\odot(v) \frac{v^2 + v_{\odot \text{esc}}^2}{v} \wp_i(v, v_{\odot \text{esc}})$$

ρ_{DM}

DM number density

σ_i

scattering cross section on element **i**

$n_i(r)$

number density of element **i**

$f_\odot(v)$

velocity distribution
(in solar frame,
without Sun's gravity)

$\wp_i(v, v_{\odot \text{esc}})$

effect of solar gravity

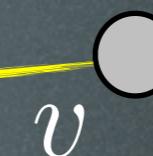
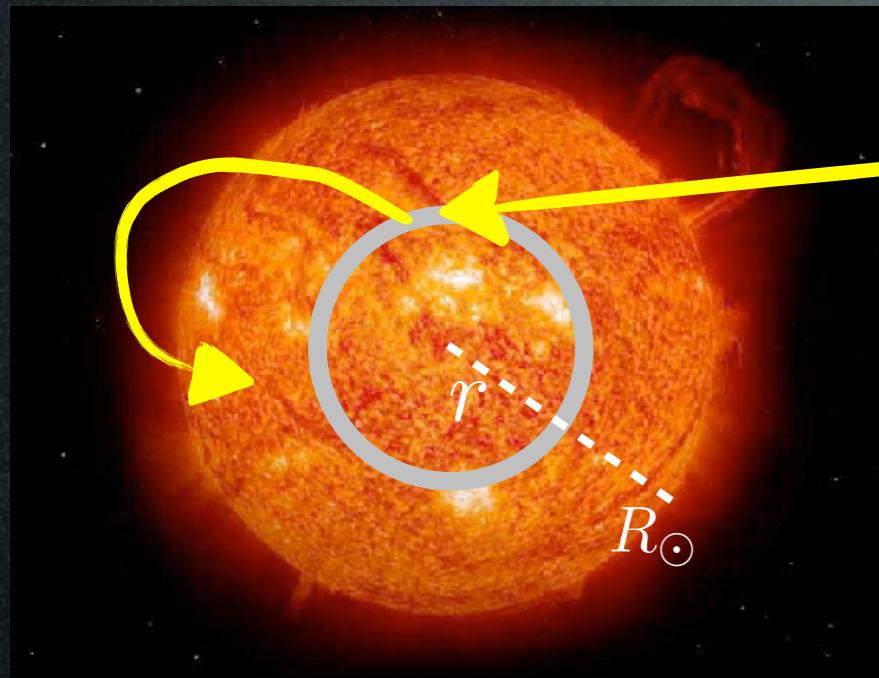
scattering probability:

$$\wp_i(v, v_{\odot \text{esc}}) = \max \left(0, 1 - \frac{\Delta_{\min}}{\Delta_{\max}} \right)$$

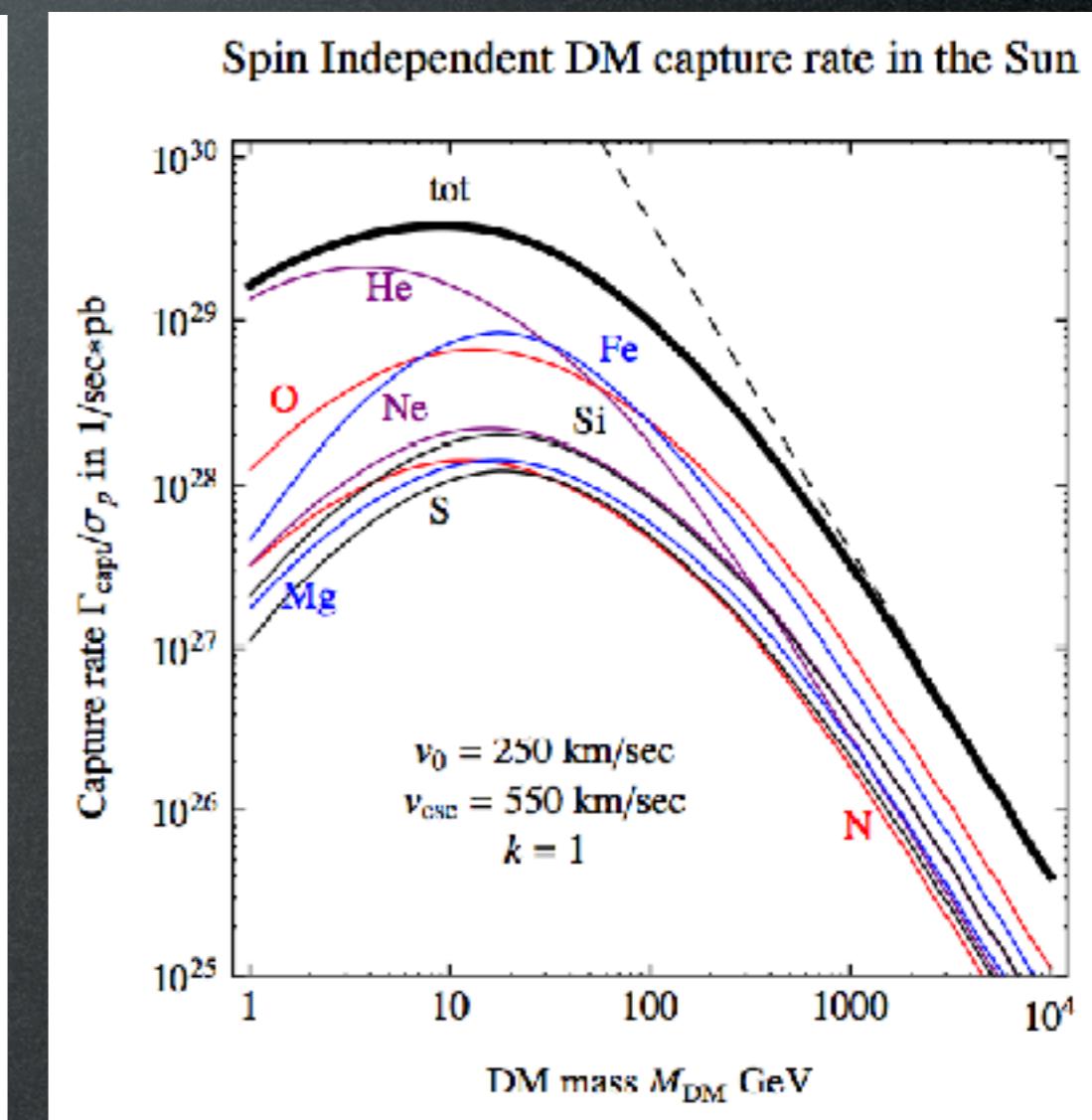
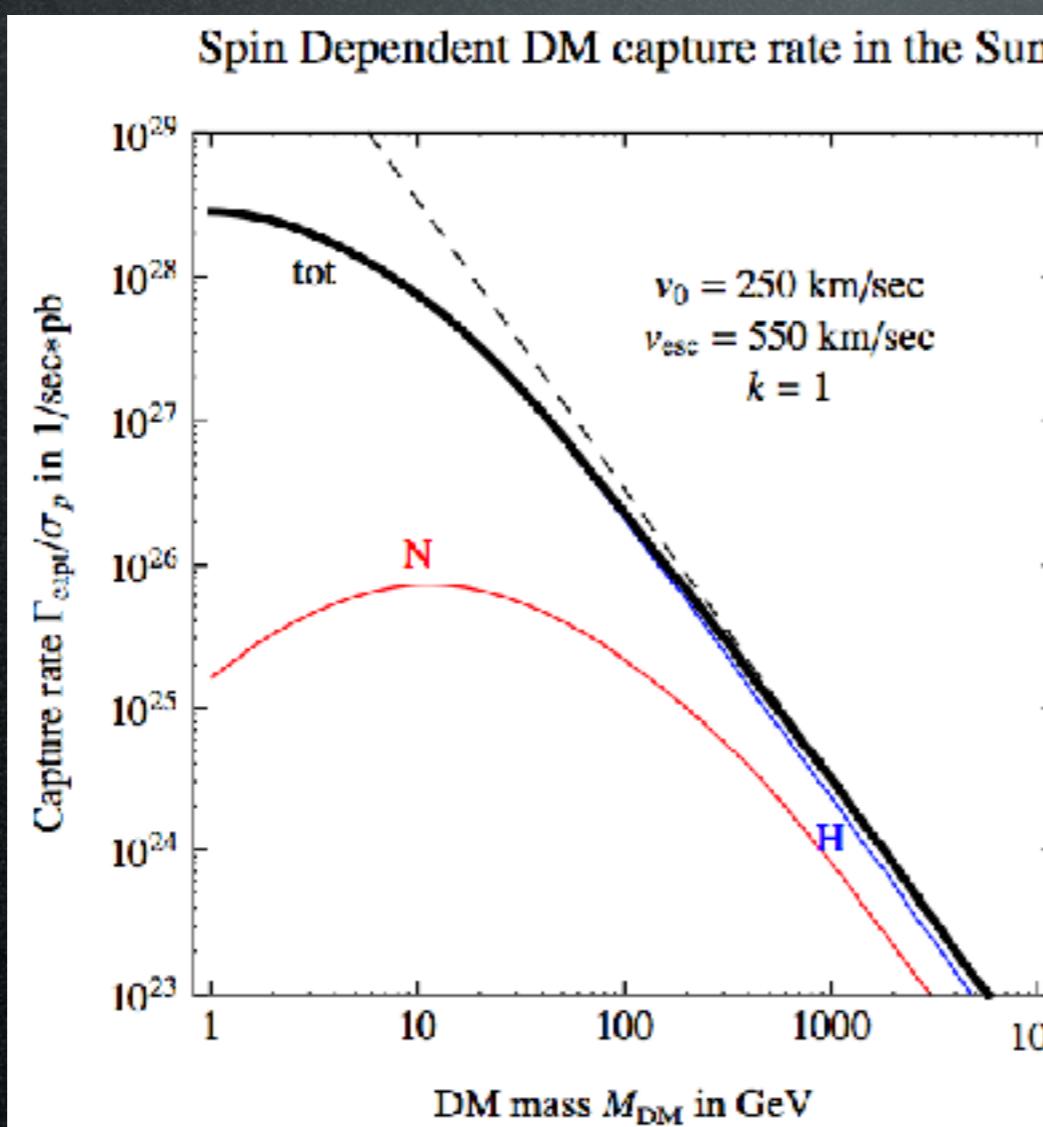
$$\Delta_{\max} = \frac{4 m_i M_{\text{DM}}}{(M_{\text{DM}} + m_i)^2} \quad \Delta_{\min} = \frac{v^2}{v^2 + v_{\odot \text{esc}}^2}$$

$$\wp_i(v, v_{\odot \text{esc}}) = \frac{1}{E \Delta_{\max}} \int_{E \Delta_{\min}}^{E \Delta_{\max}} d(\Delta E) |F_i(\Delta E)|^2 \quad |F_i(\Delta E)|^2 = e^{-\Delta E/E_0} \quad E_0^{\text{SI}} = 5/2 m_i r_i^2 \\ E_0^{\text{SD}} = 3/2 m_i r_i^2$$

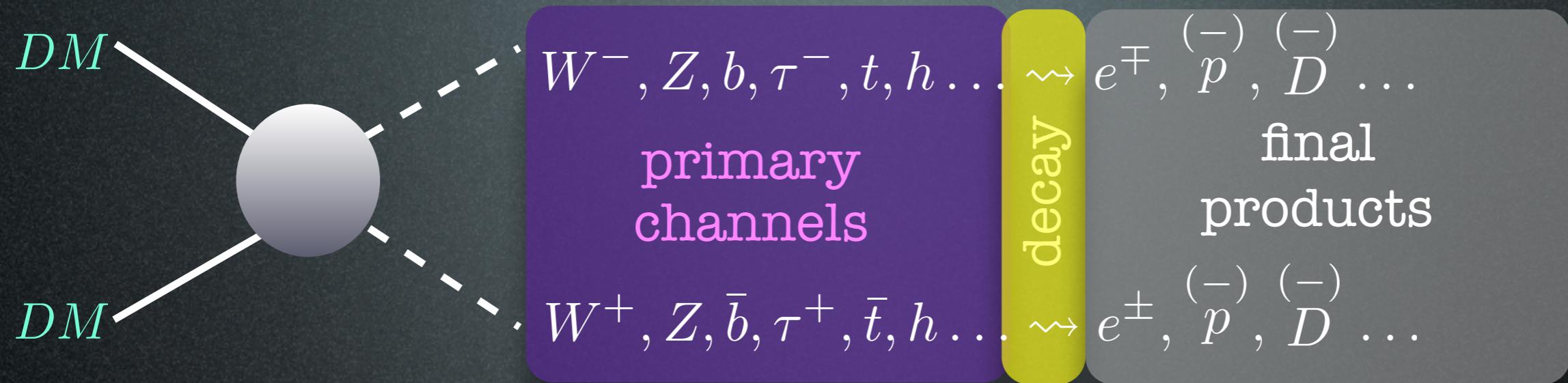
1. Capture & annihilation



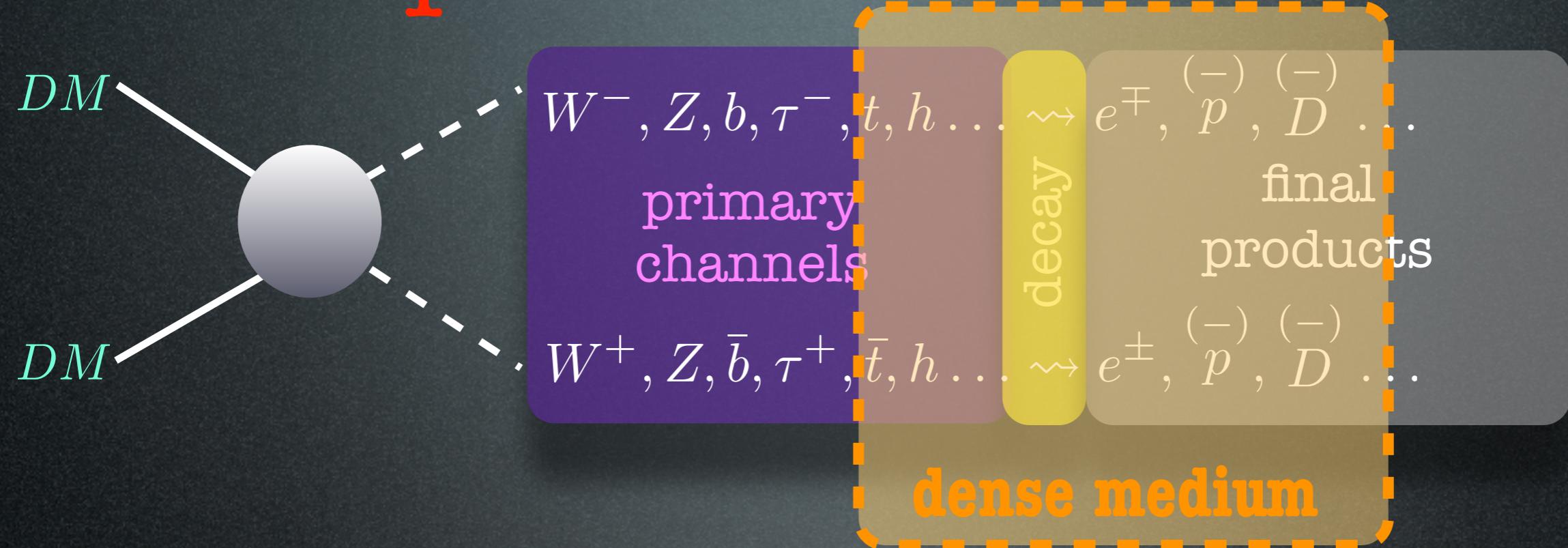
A.Gould 1987, 1988, 1990



1. Capture & annihilation



1. Capture & annihilation

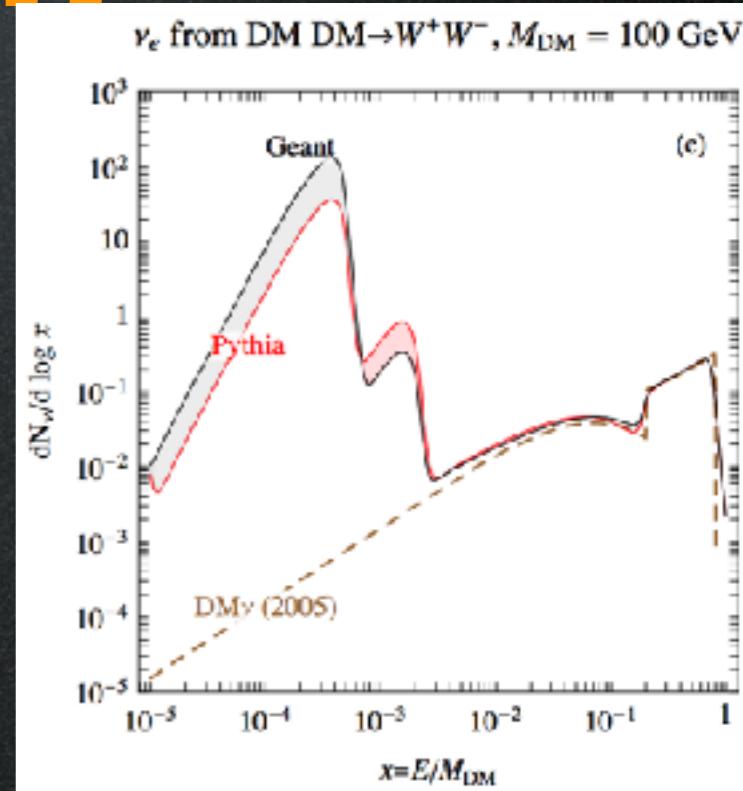
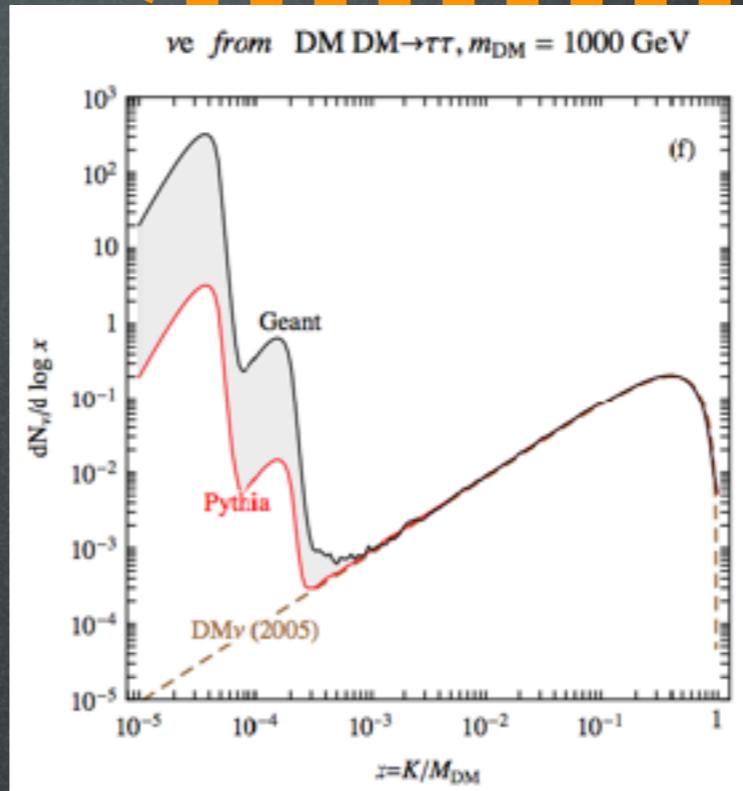
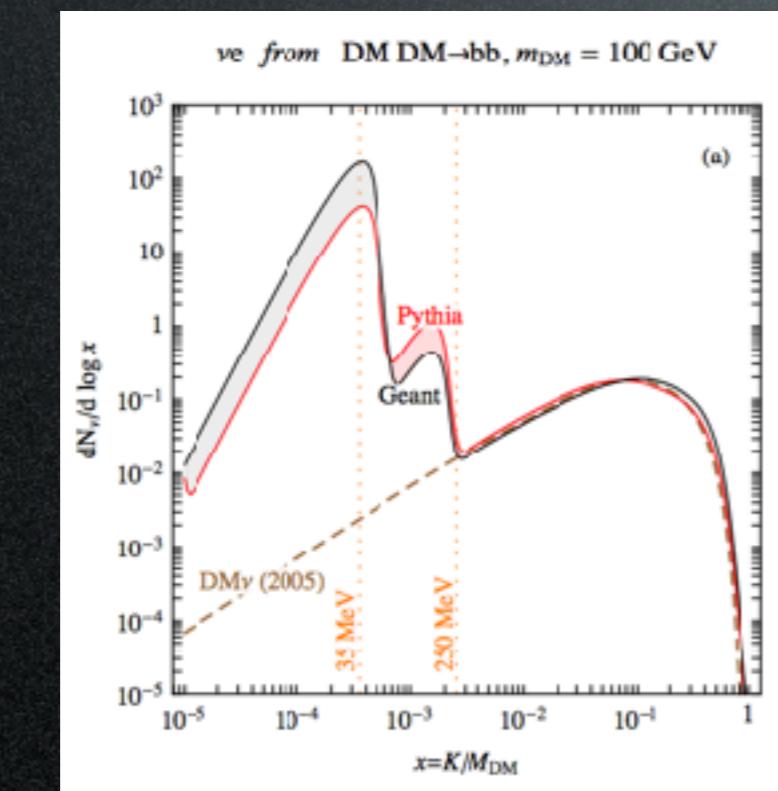


Effects of the medium:

- 1) light hadrons ($\pi, K\dots$) and leptons (μ) are stopped and decay at rest
- 2) heavy hadrons/leptons lose some energy before decaying

1. Capture & annihilation

DM



$W^-, Z, b, \tau^-, t, h \dots \rightsquigarrow e^\mp, p^\pm, D^\pm \dots$
primary channels

$e^\mp, p^\pm, D^\pm \dots$
final products

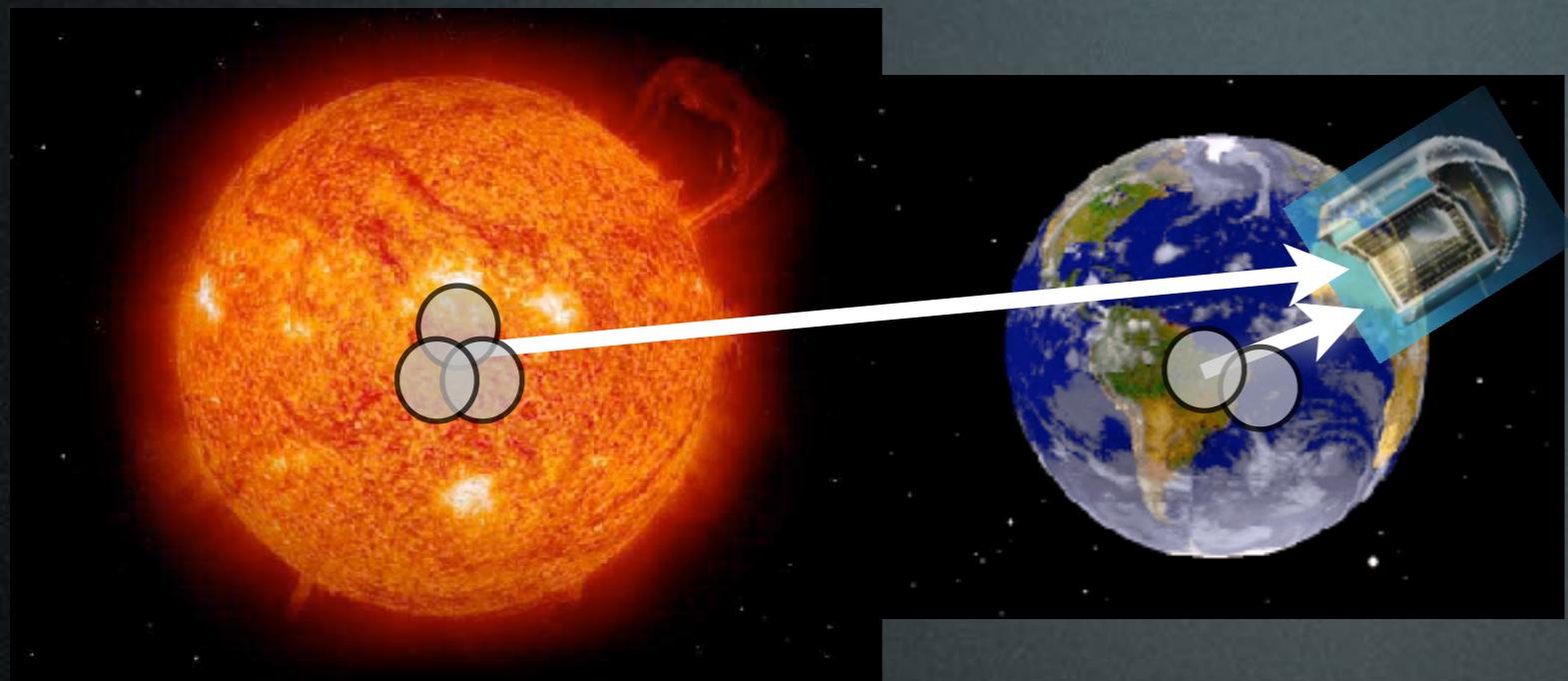
decay

dense medium

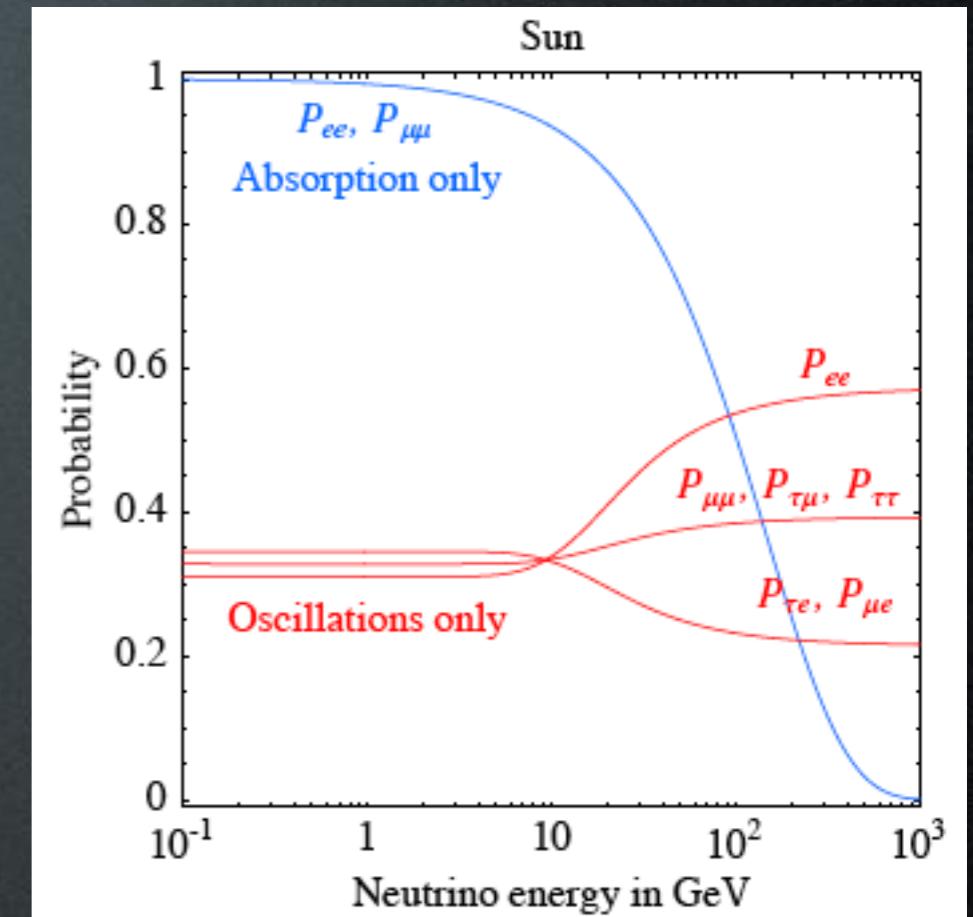
Effects of the medium:

- 1) light hadrons ($\pi, K\dots$) and leptons (μ) are stopped and decay at rest
- 2) heavy hadrons/leptons lose some energy before decaying

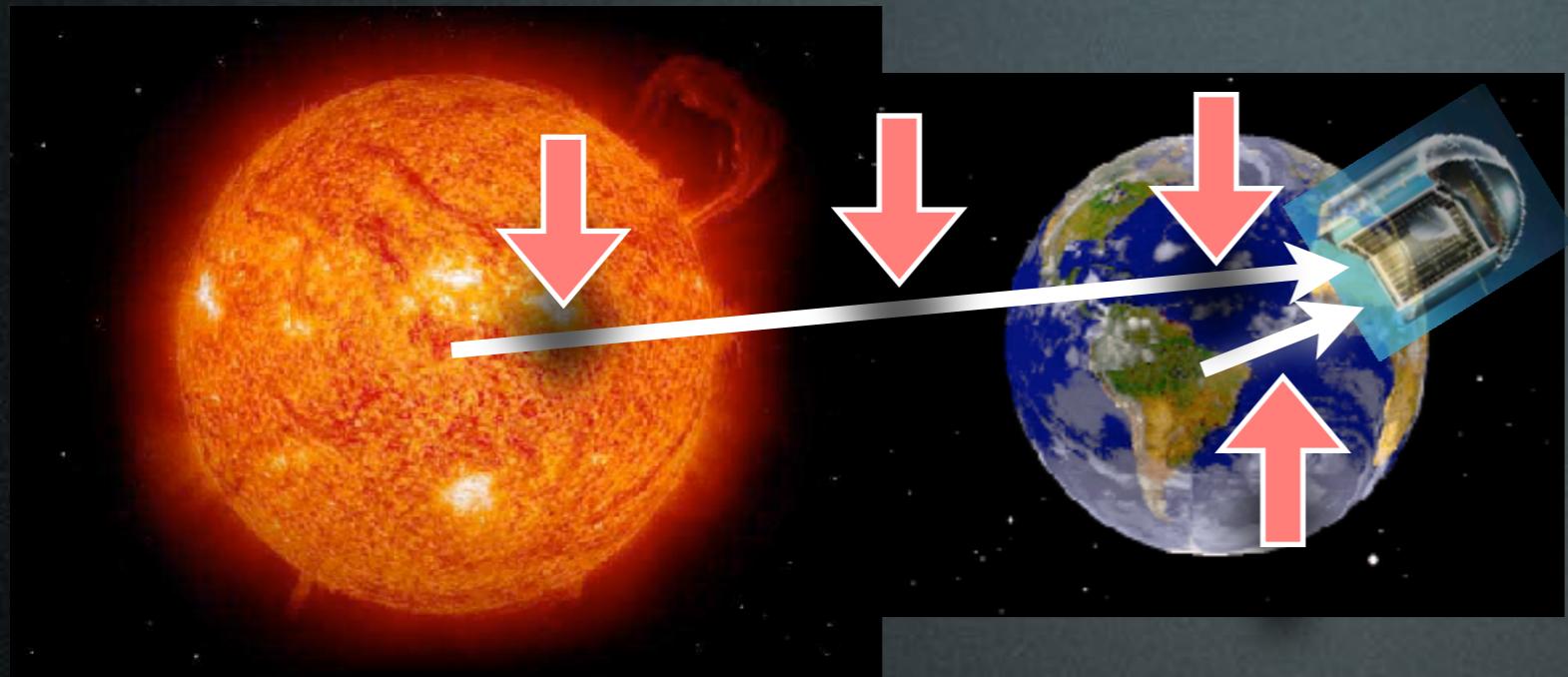
2. Propagation



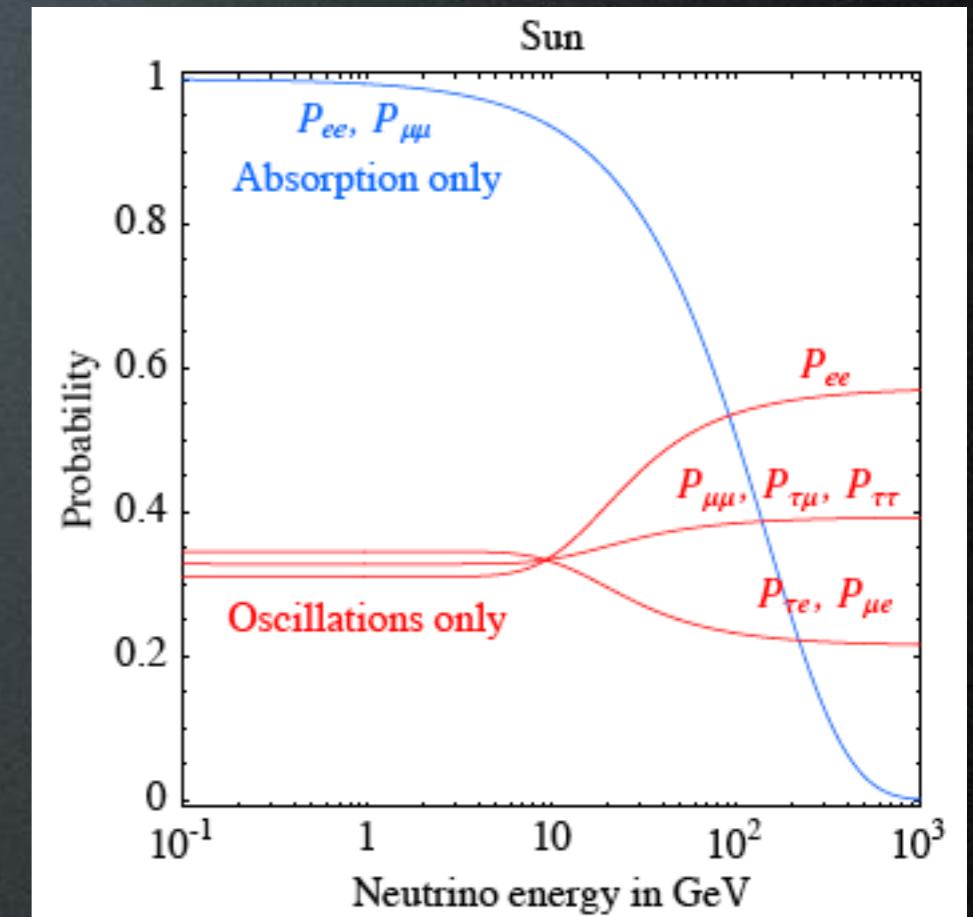
oscillations + interactions



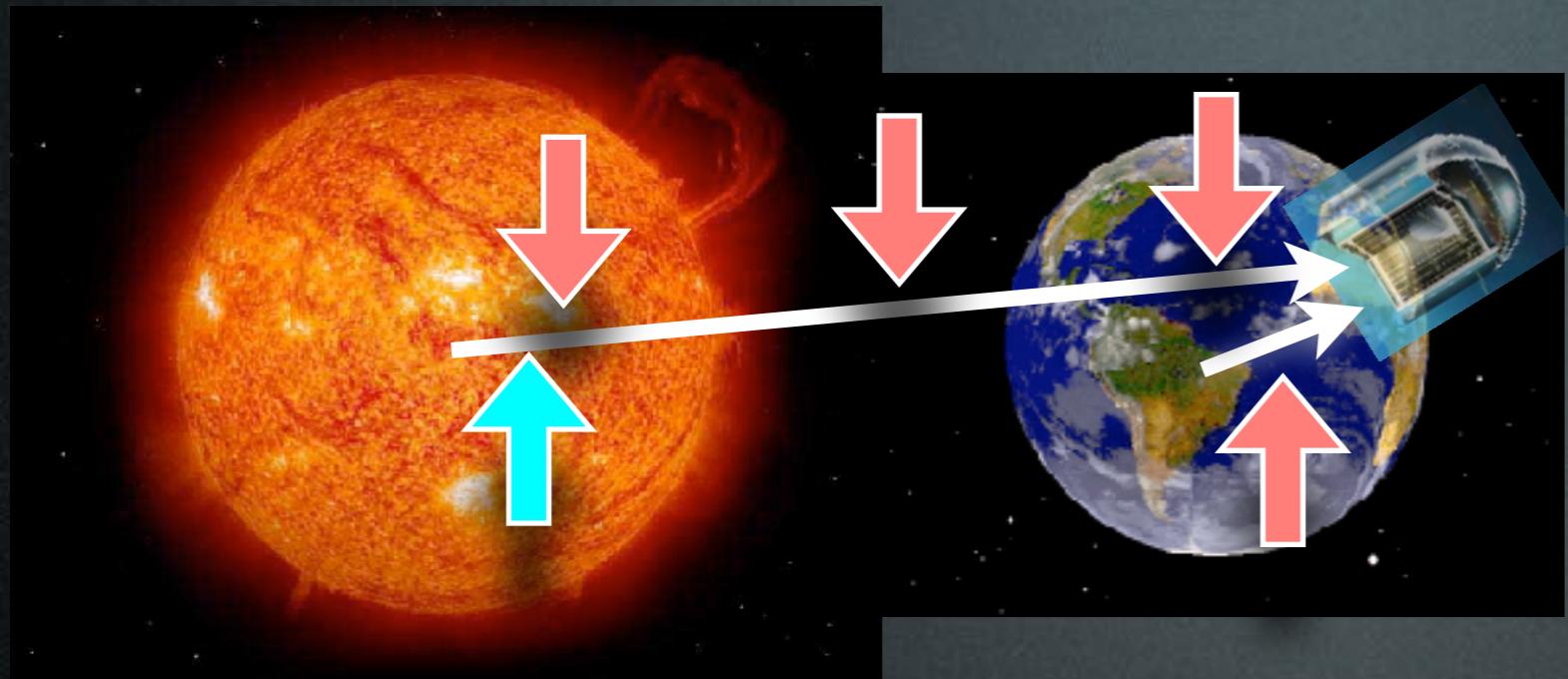
2. Propagation



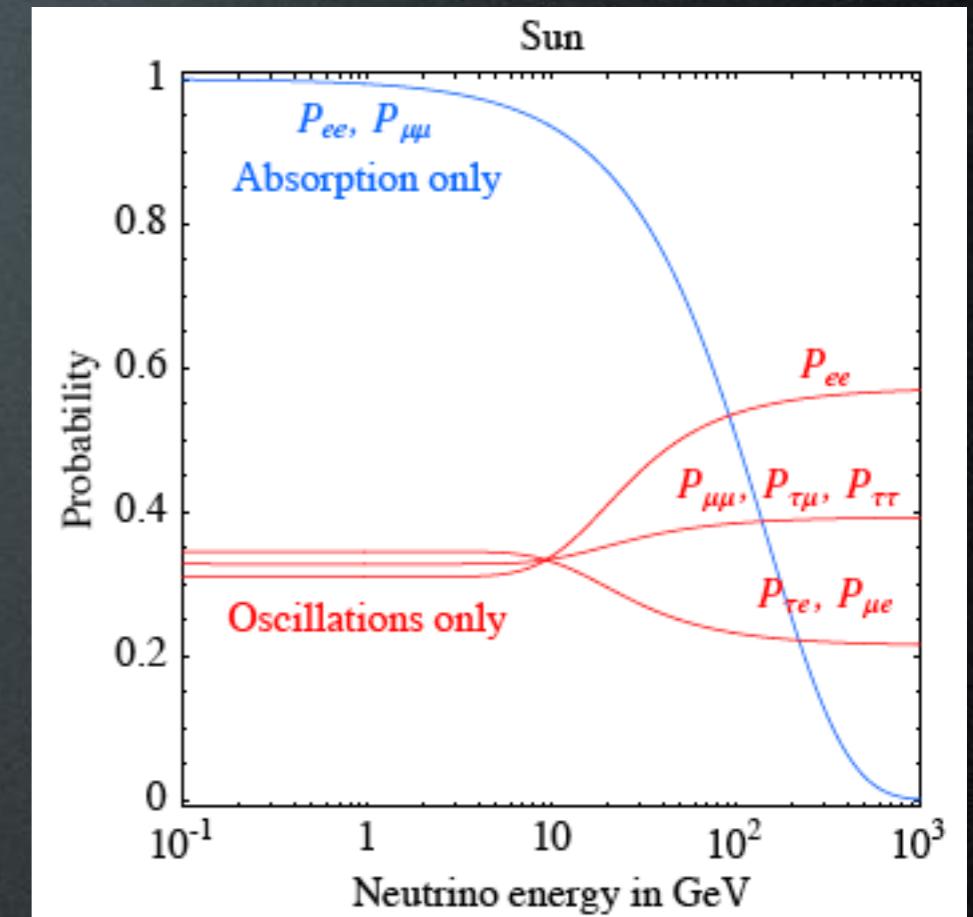
oscillations + interactions



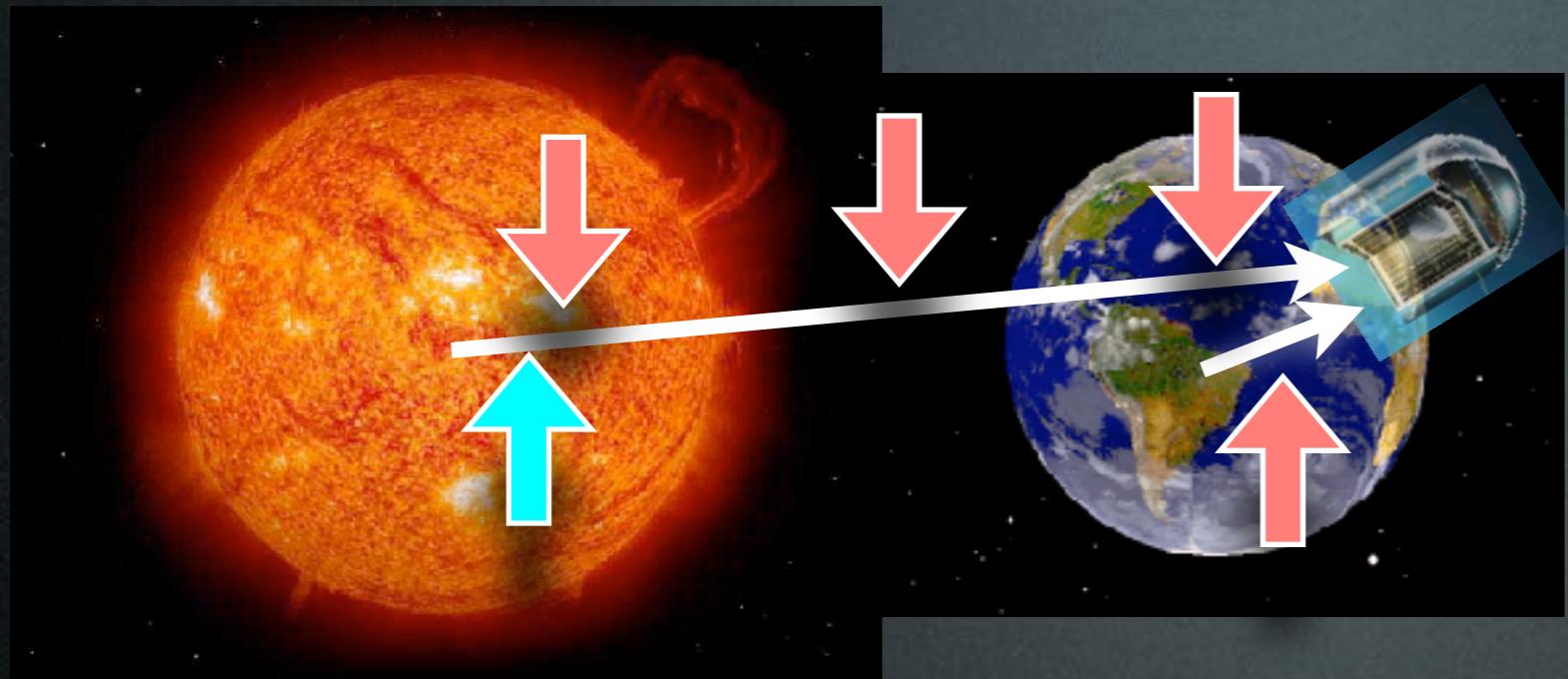
2. Propagation



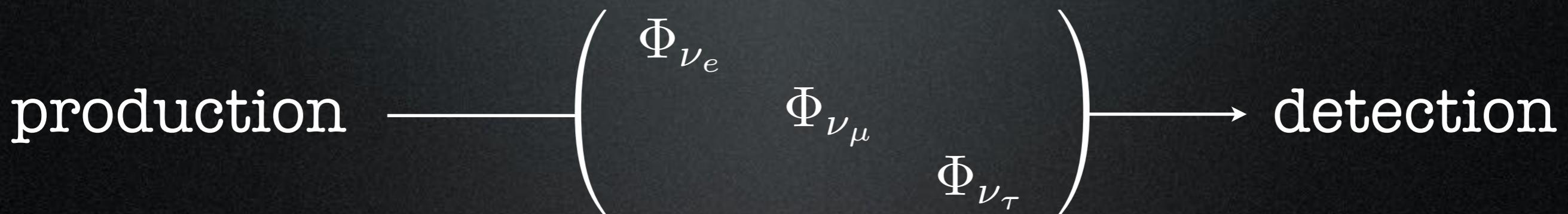
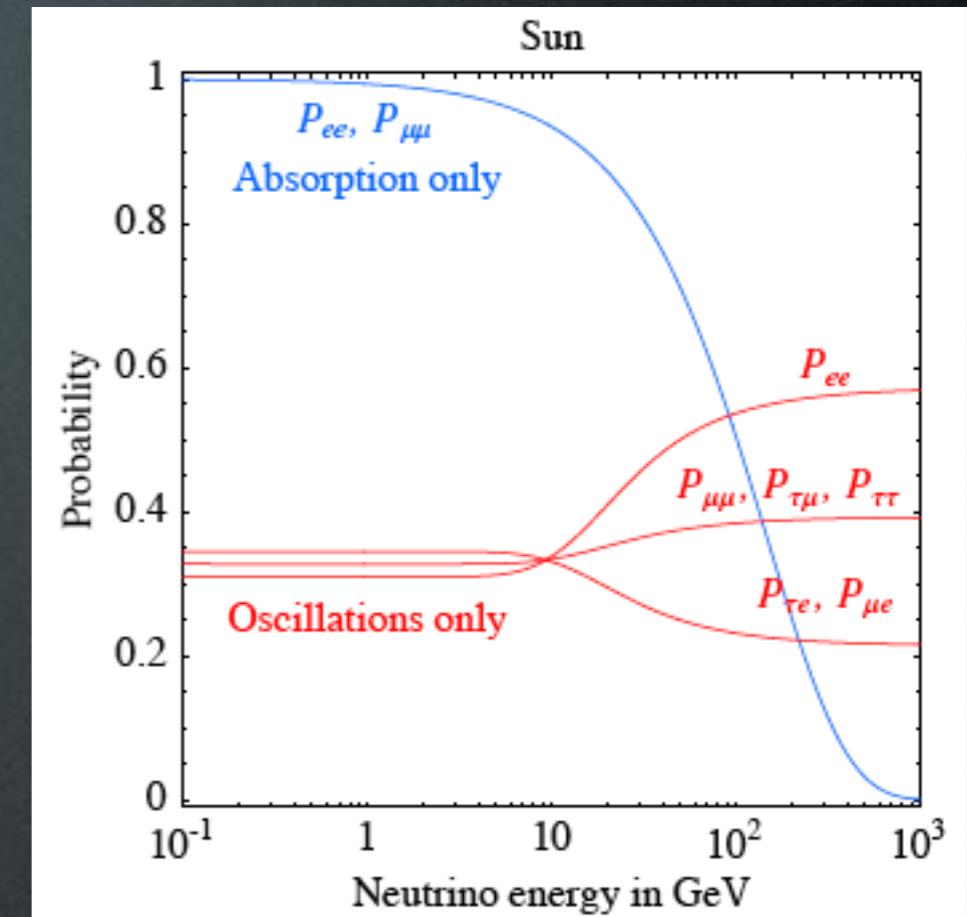
oscillations + interactions



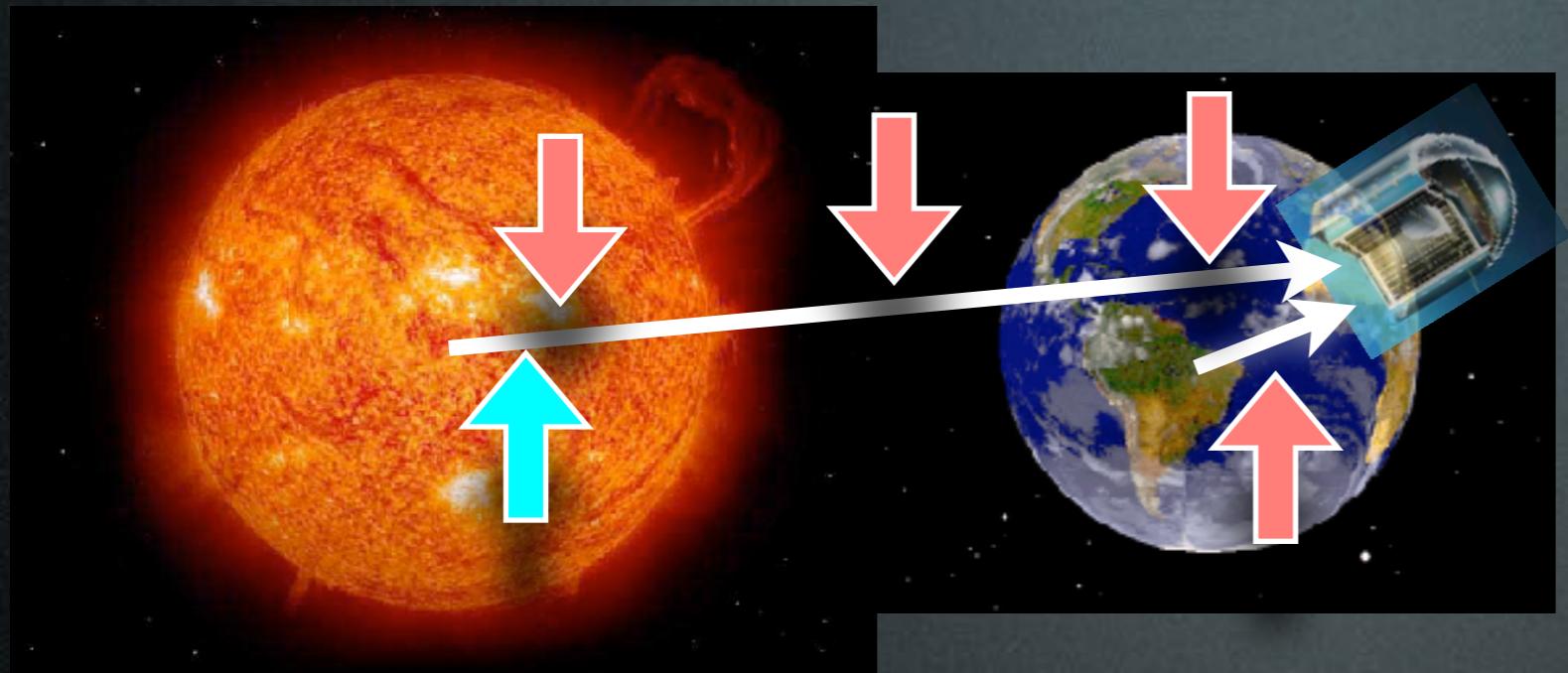
2. Propagation



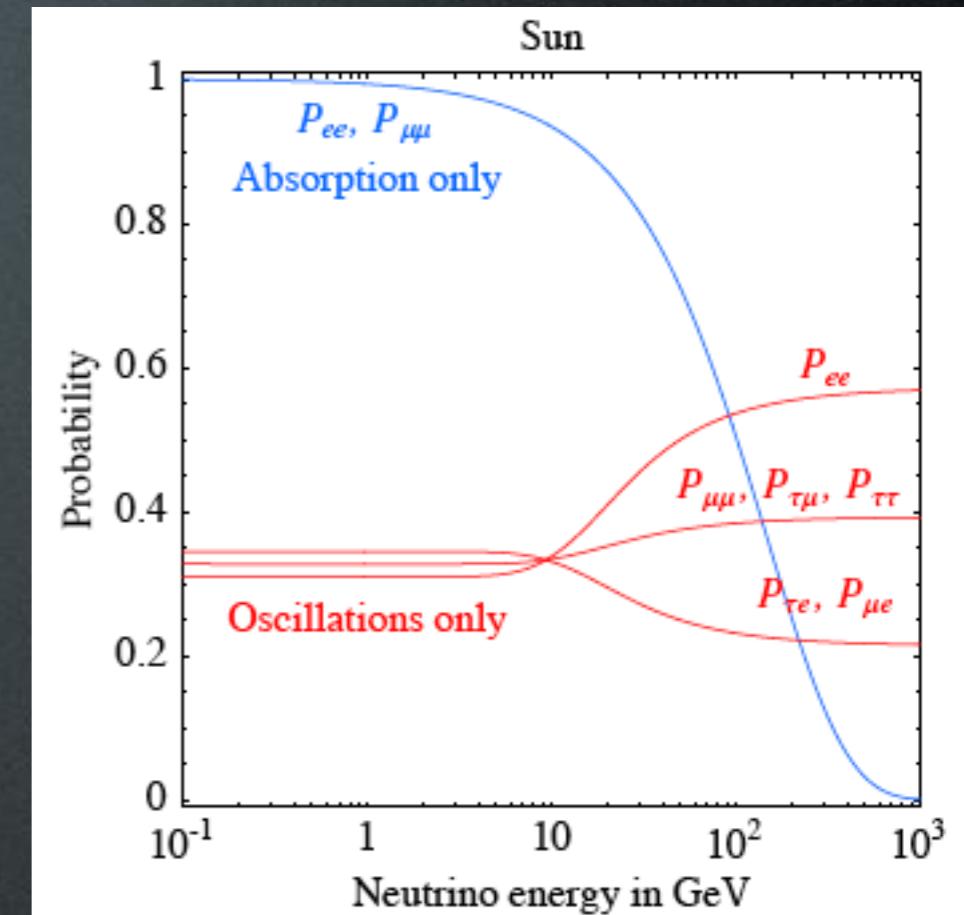
oscillations + interactions



2. Propagation



oscillations + interactions



density matrix

$$\rho = \begin{pmatrix} \rho_{ee} & \rho_{e\mu} & \rho_{e\tau} \\ \rho_{\mu e} & \rho_{\mu\mu} & \rho_{\mu\tau} \\ \rho_{\tau e} & \rho_{\tau\mu} & \rho_{\tau\tau} \end{pmatrix}$$

full evolution equation:

$$\frac{d\rho}{dr} = -i[H, \rho] + \left. \frac{d\rho}{dr} \right|_{CC} + \left. \frac{d\rho}{dr} \right|_{NC} + \left. \frac{d\rho}{dr} \right|_{in}$$

2. Propagation: oscillations

$$\frac{d\rho}{dr} = -i[H, \rho]$$

$$H = \frac{\mathbf{m}^\dagger \mathbf{m}}{2E_\nu} + \sqrt{2}G_F \left[N_e \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix} - \frac{N_n}{2} \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix} \right]$$

2. Propagation: oscillations

$$\frac{d\rho}{dr} = -i[H, \rho]$$

$$H = \frac{\mathbf{m}^\dagger \mathbf{m}}{2E_\nu} + \sqrt{2}G_F \left[N_e \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix} - \frac{N_n}{2} \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix} \right]$$



vacuum mixing:

$$\mathbf{m}^\dagger \mathbf{m} = V \cdot \begin{pmatrix} m_1^2 & & \\ & m_2^2 & \\ & & m_3^2 \end{pmatrix} \cdot V^\dagger$$

$$\theta_{\text{sun}} = 32^\circ$$

$$\theta_{\text{atm}} = 45^\circ$$

$$\theta_{13} = 8.8^\circ$$

$$\Delta m_{\text{sun}}^2 = 8.0 \cdot 10^{-5} \text{ eV}^2$$

$$|\Delta m_{\text{atm}}^2| = 2.5 \cdot 10^{-3} \text{ eV}^2$$

2. Propagation: oscillations

$$\frac{d\rho}{dr} = -i[H, \rho]$$

$$H = \frac{m^\dagger m}{2E_\nu} + \sqrt{2}G_F \left[N_e \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix} - \frac{N_n}{2} \begin{pmatrix} 1 & & \\ & 1 & \\ & & 1 \end{pmatrix} \right]$$

vacuum mixing:

$$m^\dagger m = V \cdot \begin{pmatrix} m_1^2 & & \\ & m_2^2 & \\ & & m_3^2 \end{pmatrix} \cdot V^\dagger$$

$$\theta_{\text{sun}} = 32^\circ$$

$$\theta_{\text{atm}} = 45^\circ$$

$$\theta_{13} = 8.8^\circ$$

$$\Delta m_{\text{sun}}^2 = 8.0 \cdot 10^{-5} \text{ eV}^2$$

$$|\Delta m_{\text{atm}}^2| = 2.5 \cdot 10^{-3} \text{ eV}^2$$

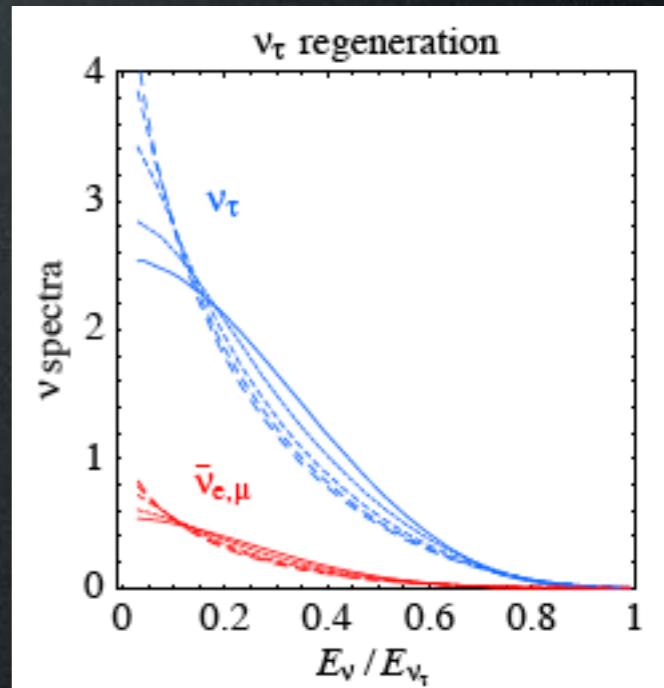
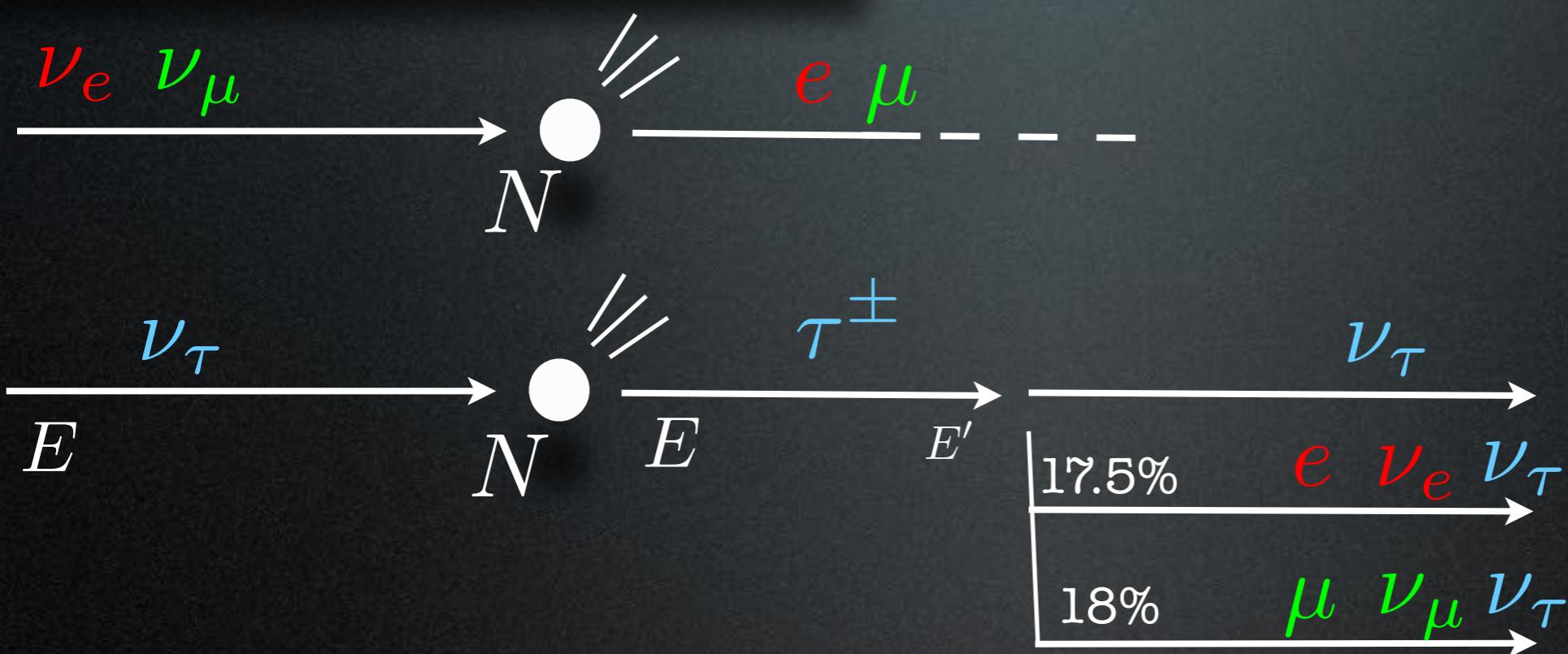
matter effect (MSW):

$$N_e(r), N_n(r)$$

from solar/
Earth models

2. Propagation: CC absorption & tau

$$\frac{d\rho}{dr} = -i[H, \rho] + \left. \frac{d\rho}{dr} \right|_{\text{CC}}$$

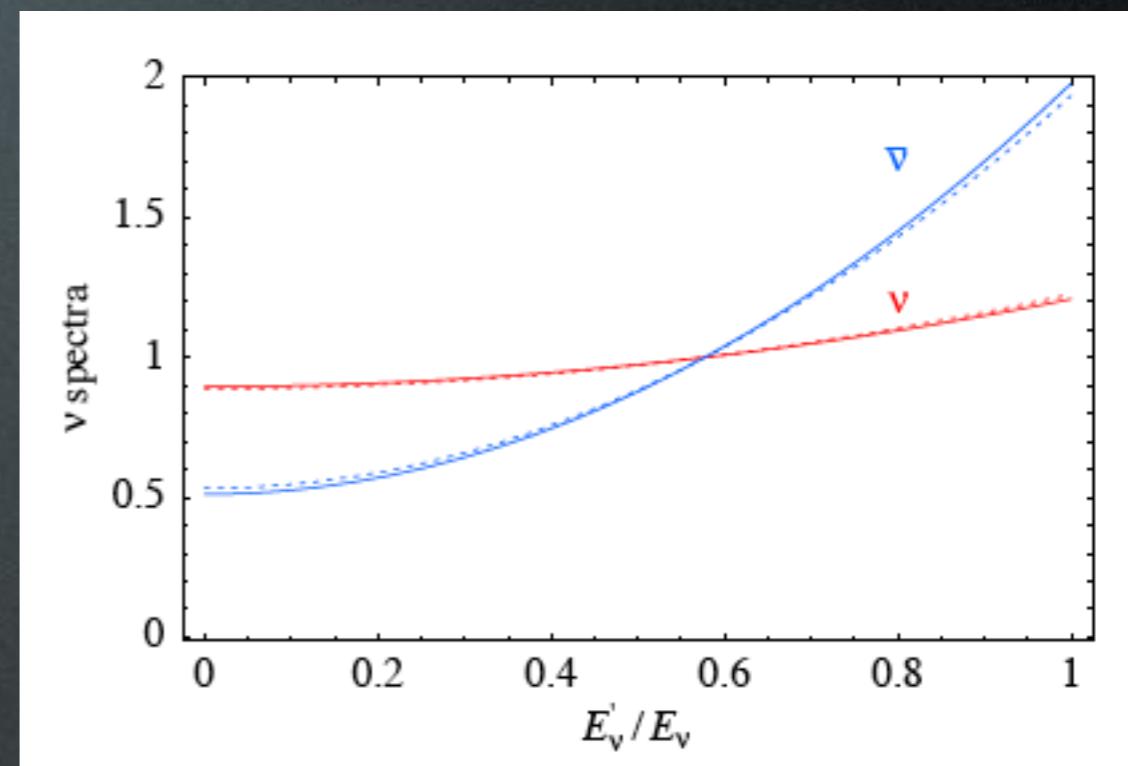
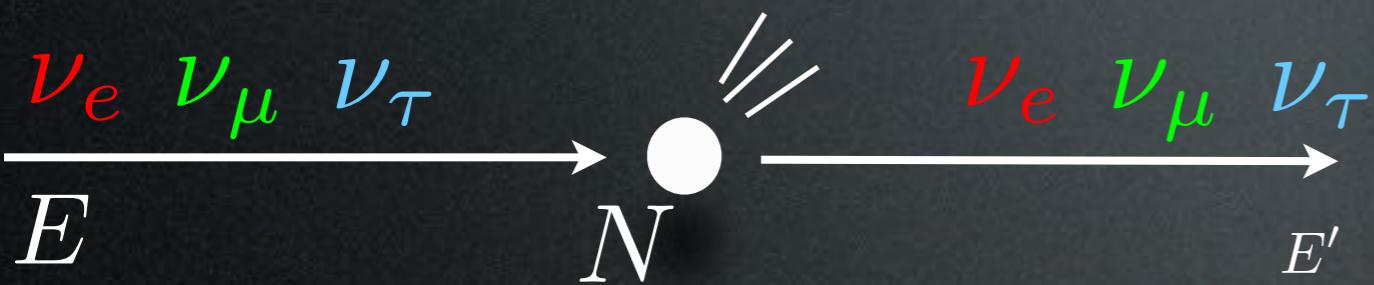


(re)generation

$$\begin{aligned} \left. \frac{d\rho}{dr} \right|_{\text{CC}} = & -\frac{\{\Gamma_{\text{CC}}, \rho\}}{2} + \int \frac{dE_\nu^{\text{in}}}{E_\nu^{\text{in}}} \left[\Pi_\tau \rho_{\tau\tau}(E_\nu^{\text{in}}) \Gamma_{\text{CC}}^\tau(E_\nu^{\text{in}}) f_{\tau \rightarrow \tau}(E_\nu^{\text{in}}, E_\nu) \right. \\ & \left. + \Pi_{e,\mu} \bar{\rho}_{\tau\tau}(E_\nu^{\text{in}}) \bar{\Gamma}_{\text{CC}}^\tau(E_\nu^{\text{in}}) f_{\bar{\tau} \rightarrow e,\mu}(E_\nu^{\text{in}}, E_\nu) \right] \end{aligned}$$

2. Propagation: NC scatterings

$$\frac{d\rho}{dr} = -i[H, \rho] + \left. \frac{d\rho}{dr} \right|_{\text{CC}} + \left. \frac{d\rho}{dr} \right|_{\text{NC}}$$



$$\left. \frac{d\rho}{dr} \right|_{\text{NC}} = - \int_0^{E_\nu} dE'_\nu \frac{d\Gamma_{\text{NC}}}{dE'_\nu}(E_\nu, E'_\nu) \rho(E_\nu) + \int_{E_\nu}^{\infty} dE'_\nu \frac{d\Gamma_{\text{NC}}}{dE'_\nu}(E'_\nu, E_\nu) \rho(E'_\nu)$$

'Prompt' gamma rays

www.marcocirelli.net/PPPC4DMID.html

PPPC 4 DM ID - A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection

We provide ingredients and recipes for computing signals of TeV-scale Dark Matter annihilations and decays.

Data and Results from [1012.4515 \[hep-ph\]](#) (and [1009.0224 \[hep-ph\]](#)), from [1312.6408 \[hep-ph\]](#), [1412.5696 \[astro-ph.HE\]](#), from [1505.01049 \[hep-ph\]](#) and from [1511.08787 \[hep-ph\]](#).

If you use the data provided on this site, please cite:

M.Cirelli, G.Corcella, A.Hektor, G.Hütsi, M.Kadastik, P.Panci, M.Raidal, F.Sala, A.Strumia,
"PPPC 4 DM ID: A Poor Particle Physicist Cookbook for Dark Matter Indirect Detection",
arXiv 1012.4515, JCAP 1103 (2011) 051.
Erratum: JCAP 1210 (2012) E01.

DMν : Neutrinos from the Sun:

DM annihilation rate in the Sun: Mathematica function [GammaAnn.m](#), refer to the notebook [Sample.nb](#) for usage.

Neutrino energy spectra at production: Mathematica function [d1NnuclxEW.m](#), refer to the notebook [Sample.nb](#) for usage.

[03 jun 2015] Warning: some bugs in these files have been brought to our attention, we are working to fix them. Sorry for the inconvenience.

Neutrino energy spectra at detection: Mathematica function [d1NnuclxEarth.m](#), refer to the notebook [Sample.nb](#) for usage.

[03 jun 2015] Warning: some bugs in these files have been brought to our attention, we are working to fix them. Sorry for the inconvenience.

DM detection

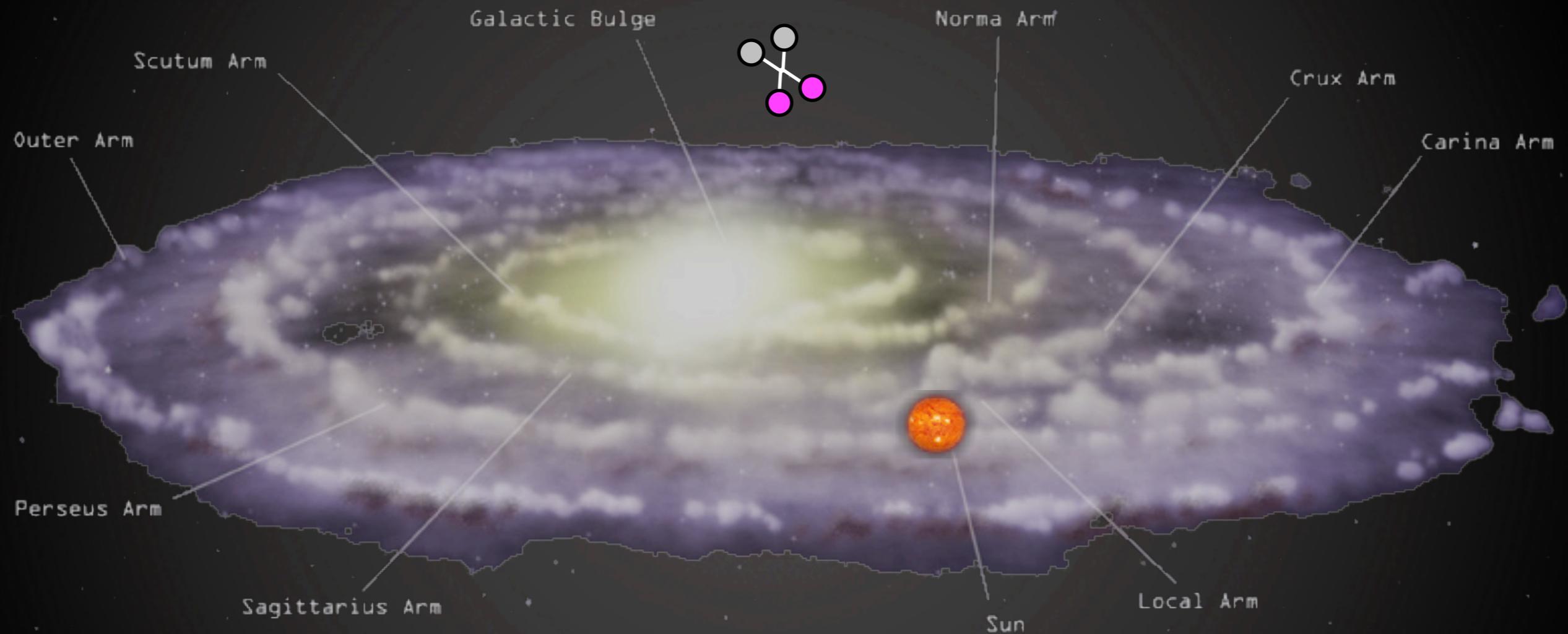
direct detection

production at colliders

- indirect
 - γ from annihil in galactic center or halo
and from synchrotron emission Fermi, HESS, radio telescopes
 - e^+ from annihil in galactic halo or center PAMELA, ATIC, Fermi
 - \bar{p} from annihil in galactic halo or center
 - \bar{d} from annihil in galactic halo or center
 - $\nu, \bar{\nu}$ from annihil in galactic center

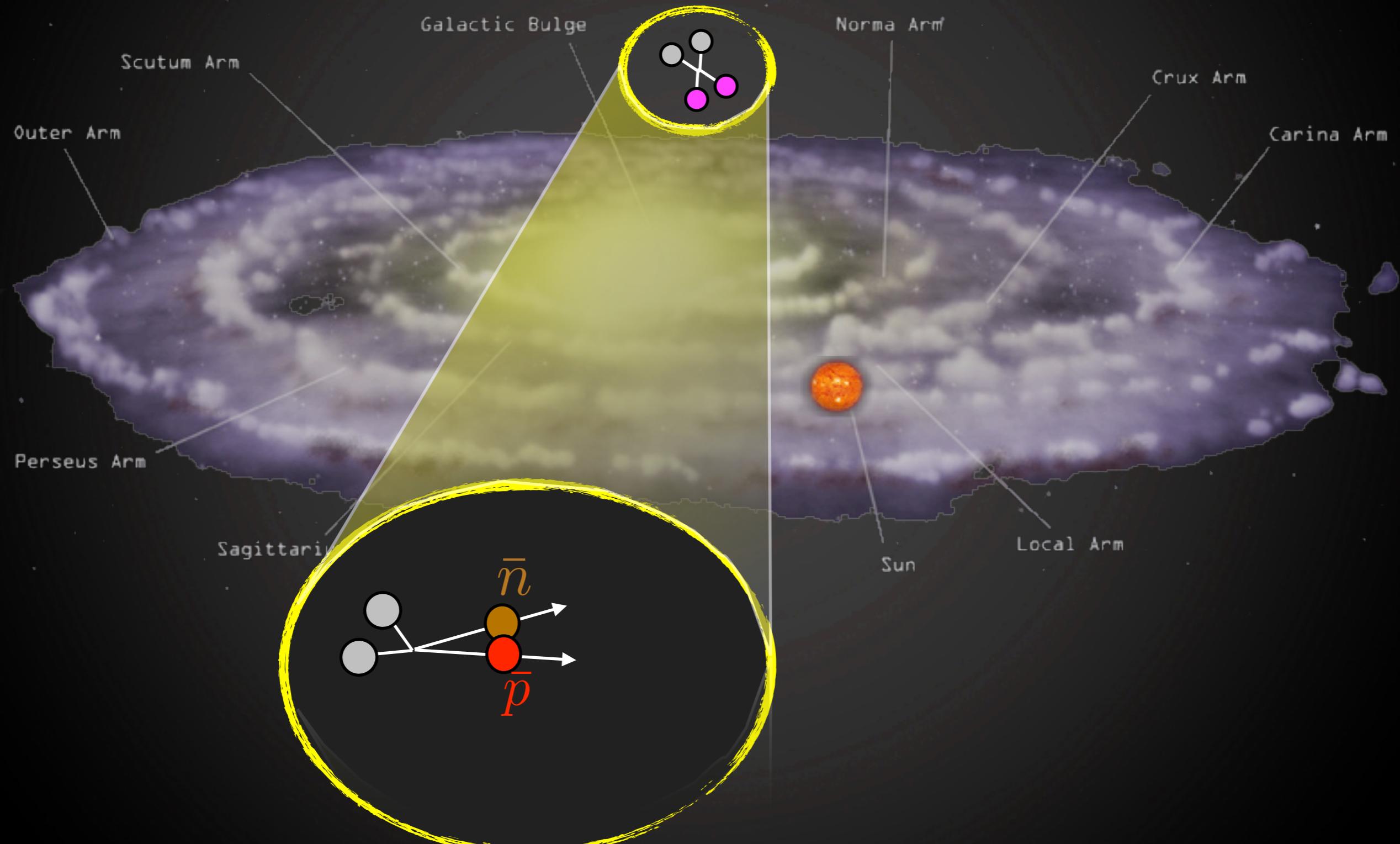
Indirect Detection: basics

\bar{d} from DM annihilations in halo



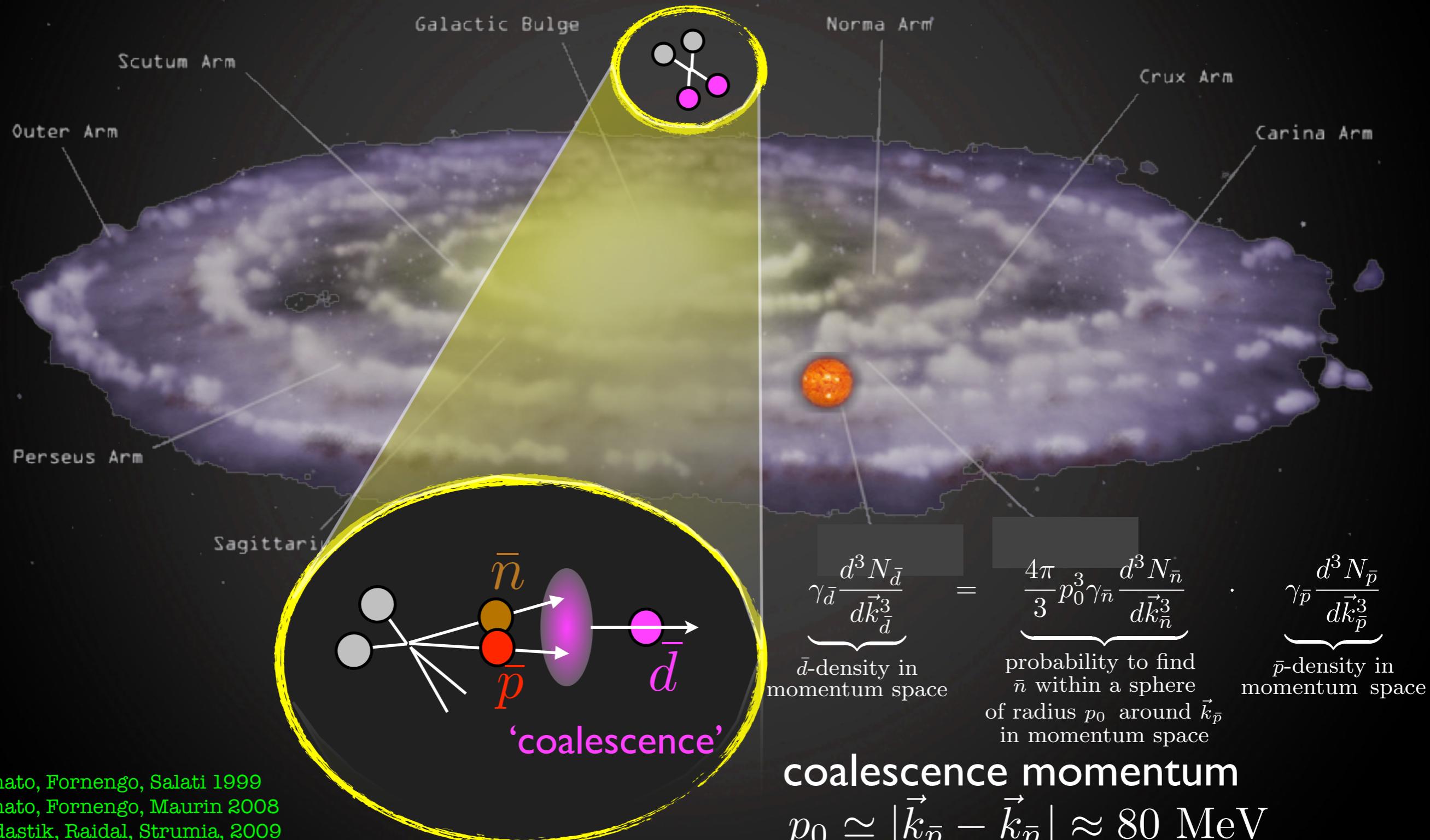
Indirect Detection: basics

\bar{d} from DM annihilations in halo



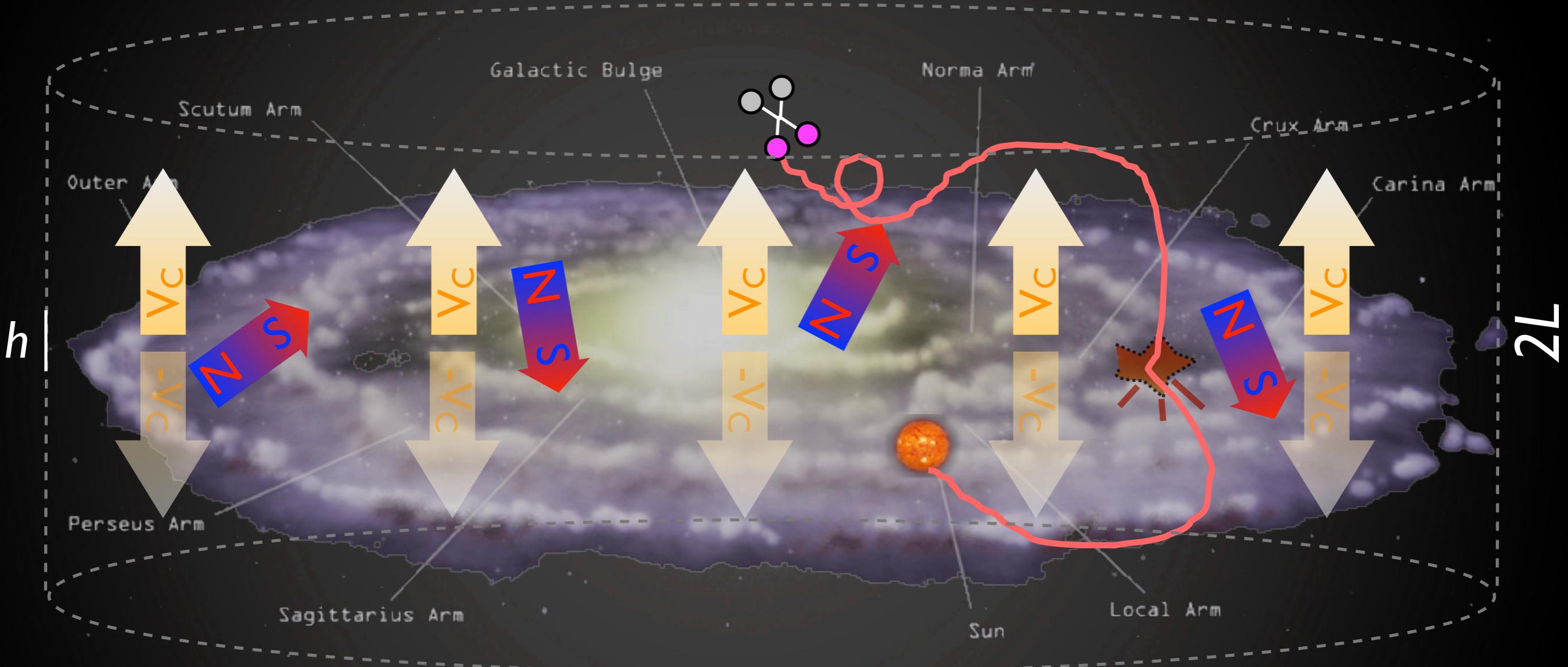
Indirect Detection: basics

\bar{d} from DM annihilations in halo



Indirect Detection: basics

\bar{d} from DM annihilations in halo



$$\frac{\partial f}{\partial t} - K(E) \cdot \nabla^2 f - \frac{\partial}{\partial E} (b(E)f) + \frac{\partial}{\partial z} (V_c f) = Q_{\text{inj}} - 2h\delta(z)\Gamma_{\text{spall}}f$$

diffusion energy loss convective wind source spallations

Advertisement

You need a quick **reference** for formulæ and methods
to compute indirect detection signals?

You want to compute all **signatures** of your DM model in
positrons, electrons, neutrinos, gamma rays...
but you don't want to mess around with astrophysics?

Advertisement

You want to compute all **signatures** of your DM model in positrons, electrons, neutrinos, gamma rays...
but you don't want to mess around with astrophysics?

‘The Poor Particle Physicist Cookbook
for Dark Matter Indirect Detection’

PPPC 4 DM ID

We provide ingredients and recipes for computing signals of TeV-scale
Dark Matter annihilations and decays in the Galaxy and beyond.

Cirelli, Corcella, Hektor,
Hütsi, Kadastik, Panci,
Raidal, Sala, Strumia

1012.4515 [hep-ph]

www.marcocirelli.net/PPPC4DMID.html



Advertisement

You want to compute all **signatures** of your DM model in positrons, electrons, neutrinos, gamma rays...
but you don't want to mess around with astrophysics?

Propagation functions for electrons and positrons everywhere in the Galaxy:

Energy loss coefficient function $b[E, r, z]$ for electrons and positrons in the Galaxy: Mathematica function `b.m`, refer to the notebook `Sample.nb` for usage.

Annihilation

Positrons: The file `ElectronHaloFunctGalaxyAnn.m` provides the halo functions $I(x, E_p, r, z)$ at a point (r, z) in the Galaxy.
The notebook `Sample.nb` shows how to load and use it.

Decay

Positrons: The file `ElectronHaloFunctGalaxyDec.m` provides the halo functions $I(x, E_p, r, z)$ at a point (r, z) in the Galaxy
The notebook `Sample.nb` shows how to load and use it.

Propagation functions for charged cosmic rays at the location of the Earth:

Annihilation

Positrons: The file `ElectronHaloFunctEarthAnn.m` provides the halo functions $I(x, E_p, r_{\text{Earth}})$ at the location of the Earth.
The notebook `Sample.nb` shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ (in the approximated formalism - see paper).

Antiprotons: [Table](#) of fit coefficients for the propagation function $R[T]$.

Antideuterons: [Table](#) of fit coefficients for the propagation function $R[T]$.

Decay

Positrons: The file `ElectronHaloFunctEarthDec.m` provides the halo functions $I(x, E_p, r_{\text{Earth}})$ at the location of the Earth.
The notebook `Sample.nb` shows how to load and use it.

[Table](#) of fit coefficients for the reduced halo function $I(\lambda)$ (in the approximated formalism - see paper).

Antiprotons: [Table](#) of fit coefficients for the propagation function $R[T]$.

Antideuterons: [Table](#) of fit coefficients for the propagation function $R[T]$.

Fluxes of charged cosmic rays at the Earth, after propagation:

Annihilation

Positrons: Mathematica function: the file `ElectronFluxAnn.m` provides the

Decay

Positrons: Mathematica function: the file `ElectronFluxDec.m` provides the

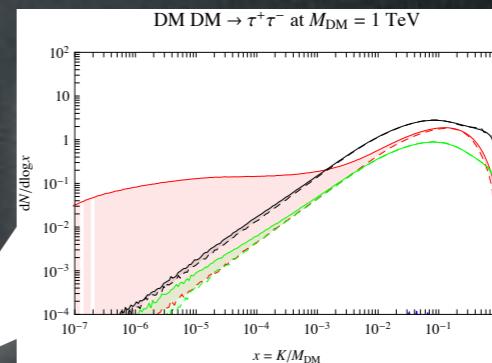
Advertisement

You want to compute all **signatures** of your DM model in positrons, electrons, neutrinos, gamma rays...
but you don't want to mess around with astrophysics?

Main added value features:

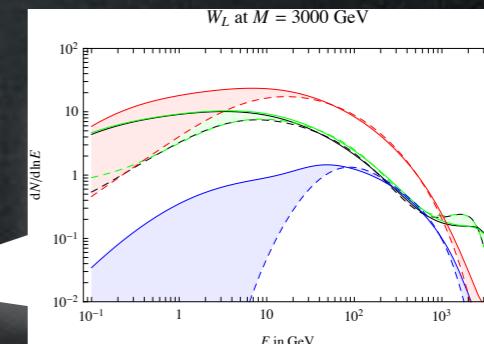


compare different MCs

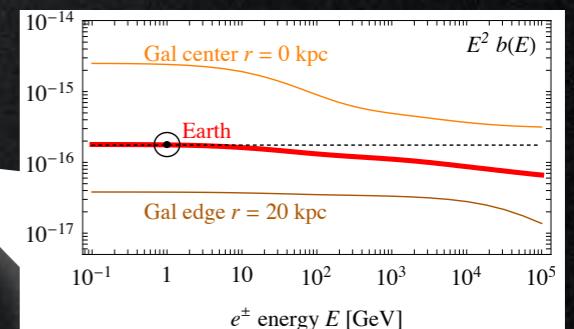


include EW corrections

Ciafaloni, Riotto et al., 1009.0224



improved e^\pm propagation



improved ICS γ -ray computation