

Probing pseud Nambu-Goldstone Boson by stimulated photon colliders in the mass range $0.1\text{eV} \sim 10\text{keV}$

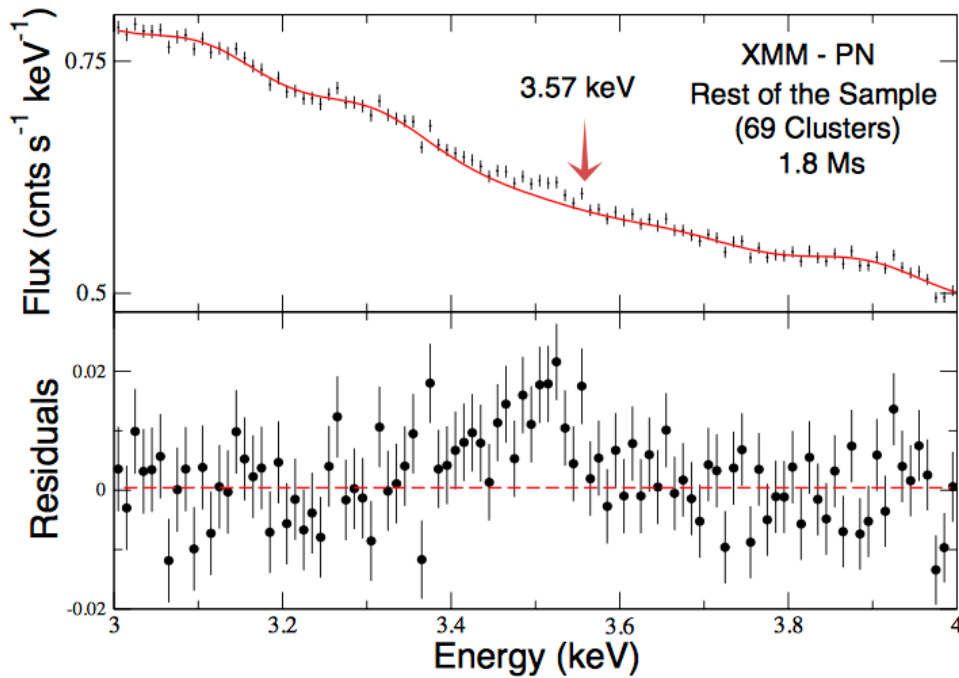
arXiv:1701.04282, (2017)

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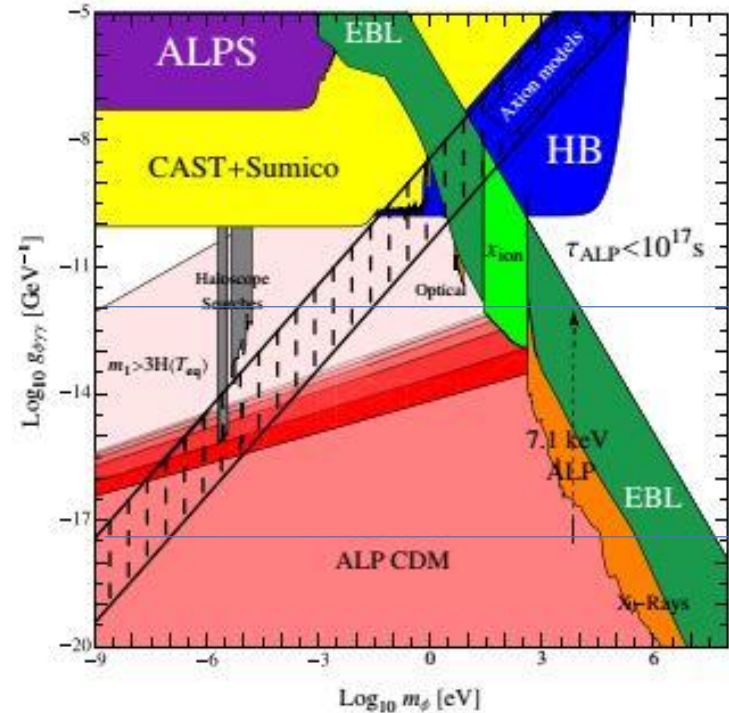
Yuichi Toyota and Kensuke Homma

4/19/2017

Unidentified emission in 3.5keV



arXiv:1402.2301v2



PHYSICAL REVIEW D 89, 103511 (2014)

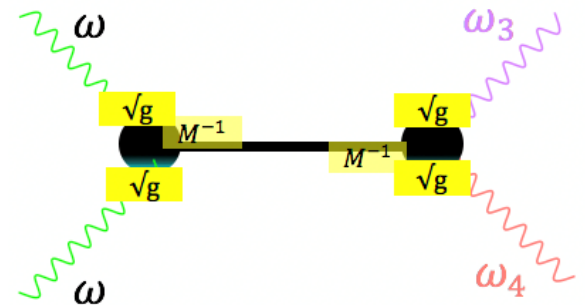
- Recently, an unidentified emission line, $\omega \sim 3.5$ keV, was reported in the photon energy spectra from a single galaxy and galaxy clusters.
- If we assume a resonance state, the resonance state of 7 keV decay into two photons which has 3.5 keV.

Feature of this method

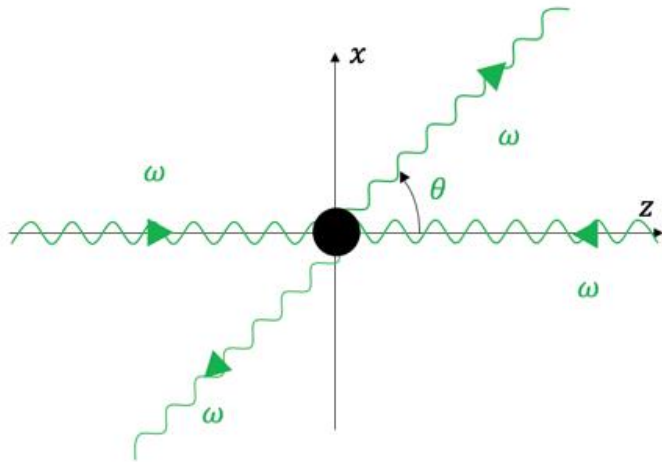
Effective Interaction Lagrangian

$$-L = gM^{-1} \frac{1}{4} F_{\mu\nu} F^{\mu\nu} \phi$$

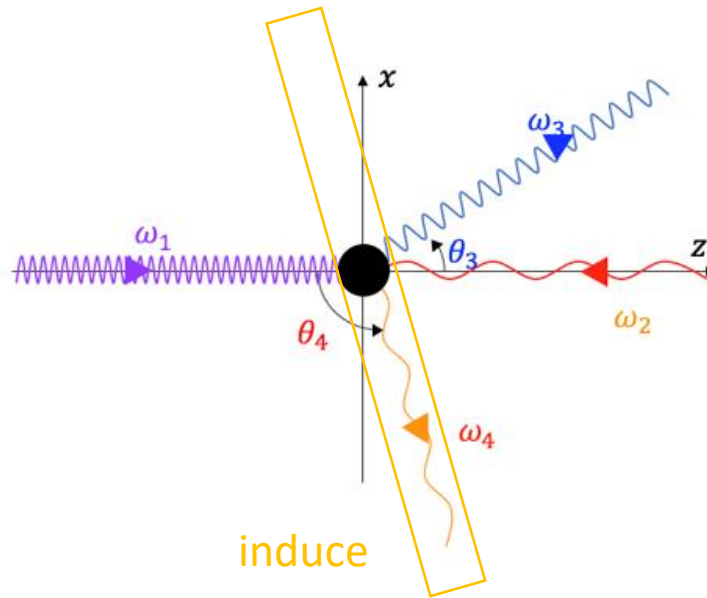
- s-channel resonant pNGB exchange.
- Utilizing the coherent nature of induction laser.
- head-on-collision



Asymmetric Collision System(ACS)



Center of Mass system



Asymmetric collision system

u, u^+, u^- are constant related β

$$\begin{aligned} \text{O-axis: } & \omega_1 + \omega_2 = \omega_3 + \omega_4 \\ \text{z-axis: } & \omega_3 \sin \theta_3 = \omega_4 \sin \theta_4 \\ \text{x-axis: } & \omega_1 - \omega_2 = \omega_3 \cos \theta_3 - \omega_4 \cos \theta_4 \end{aligned}$$

$$\begin{aligned} \omega_1 &= u\omega, & \omega_2 &= u^{-1}\omega \\ \omega_3 &= \frac{2\omega}{u^+ - u^- \cos \theta_3} \end{aligned}$$

Short pulsing effect

$$|\mathcal{M}| = -(gM^{-1})^2 \frac{\omega^4 (\cos 2\vartheta - 1)^2}{2\omega^2 (\cos 2\vartheta - 1) + m^2},$$

Then we introduce the imaginary part,

$$m^2 \rightarrow (m - i\Gamma/2)^2 \approx m^2 - im\Gamma \quad , \Gamma = (16\pi)^{-1} (gM^{-1})^2 m^3,$$

Beam energy fluctuation occurs because of short pulsing, so we integrate and take the average value.

$$\begin{aligned} |\bar{M}|^2 &= \frac{1}{\chi_+ - \chi_-} \int_{\chi_-}^{\chi_+} |M|^2 d\chi \\ \omega_{\pm} &= \frac{m}{2} \pm \Delta\omega, \omega = \frac{m}{2} \\ \chi_{\pm} &= \omega_{\pm}^2 - \left(\frac{m}{2}\right) = \Delta\omega^2 \pm m\Delta\omega \sim \pm \Delta\omega \\ \overline{|M|^2} &\sim 8\pi^3 \frac{a}{m\Delta\omega} = \frac{\pi^2}{4\Delta\omega} \left(\frac{g}{M}\right)^2 m^3, \end{aligned}$$

Mean value of the squared amplitude

$$\overline{|M|^2} \sim 8\pi^3 \frac{a}{m\Delta\omega} = \frac{\pi^2}{4\Delta\omega} \left(\frac{g}{M}\right)^2 m^3,$$

Here, energy resolution R defined as,

$$\Delta\omega = \sqrt{\omega_1\omega_2} \sqrt{\left(\frac{\delta\omega_1}{\omega_1}\right)^2 + \left(\frac{\delta\omega_2}{\omega_2}\right)^2} \equiv \frac{mR}{\sqrt{2}},$$

Mean value is,

$$\overline{|M|^2} \sim \frac{\sqrt{2}\pi^2}{8R} \left(\frac{g}{M}\right)^2 m^2$$

Yield

Y_s : spontaneous yield, Y_i : induction yield

$$Y_s = \int dt \int dx^i \rho_1(t, x^i) \rho_2(t, x^i) K \left[\frac{1}{L^2} \right] \tilde{\sigma}_{dm} [L^2]$$

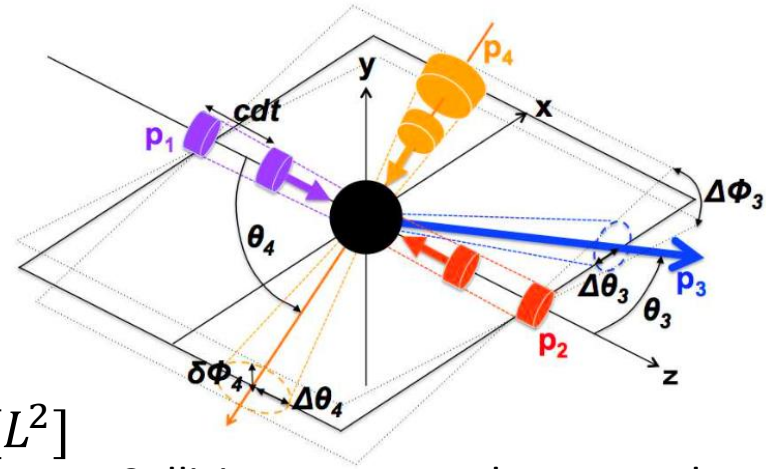
$$Y_i = \int dt \int dx^i \rho_1(t, x^i) \rho_2(t, x^i) K P_4(t, x^i) \tilde{\sigma}_{dm}$$

$$\equiv K \tilde{\sigma}_{dm} N_1 N_2 N_4 G$$

$$G = \left(\frac{2}{\pi} \right)^{\frac{2}{3}} \left(\frac{\omega_{04}}{c\tau} \right)^2 \int_{-\tau L}^{\tau L} dt (\omega_1 \omega_2 \omega_4)^{-2} \times \left\{ A \left(\frac{1}{\omega_1^2} + \frac{1}{\omega_2^2} + \frac{1}{\omega_4^2} \right) B \right\}^{-\frac{1}{2}} e^{-2Dt^2}$$

$$\tilde{\sigma}_{ngb} = \frac{|M|^2}{(8\pi\omega)^2} \int_{\underline{\phi}_3}^{\overline{\phi}_3} d\phi_3 \int_{\underline{\theta}_3}^{\overline{\theta}_3} \left(\frac{\omega_3}{2\omega} \right)^2 \sin \theta_3 d\theta_3 = \frac{\sqrt{2}}{128R} \left(\frac{g}{M} \right)^2 I,$$

$$I \equiv \frac{\Delta\phi_3 (\cos \underline{\theta}_3 - \cos \overline{\theta}_3)}{(u^+ - u^- \cos \underline{\theta}_3)(u^+ - u^- \cos \overline{\theta}_3)}$$



Collision geometry between three pulsed photon beams.

Coupling

Detected signal photon Number is,

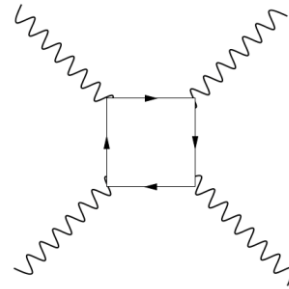
$$Y = fT\epsilon Y_i,$$

f :collision repetition rate(Hz)
 T :accumulation time (s)
 ϵ :efficiency

finally a reachable coupling strength is

$$\frac{g}{M} = 2^{1/4} 8 \sqrt{\frac{RY}{IfT\epsilon KGN_1N_2N_4}}$$

QED Background



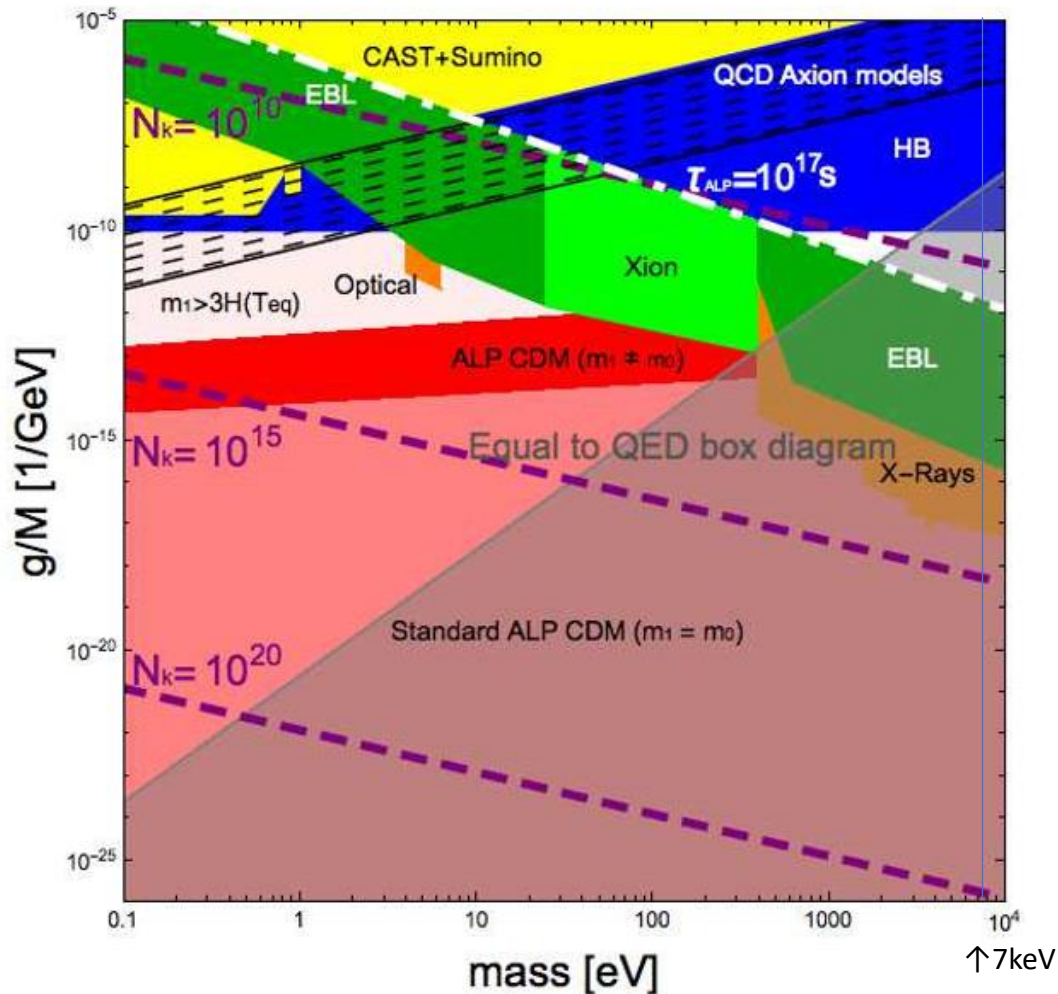
- Above the keV range, we must consider QED Background from the standard model. This calculation are based on Euler-Heisenberg equation.

$$d\sigma_{qed} = \frac{(\alpha r_0)^2}{4\pi^2} \frac{139}{90^2} \omega^6 \left(3 + \frac{\gamma^2 (\cos \theta_3 - \beta)^2}{\gamma^2 (\cos \theta_3 - \beta)^2 + \sin^2 \theta_3} \right)^2$$

$$\times \left(1 + \frac{160}{90^2} \frac{\omega^2 \sin^2 \theta_3}{4\gamma^2 (\cos \theta_3 - \beta)^2 + 3 \sin^2 \theta_3} \right)$$

$$\times \frac{\gamma(1 - \beta \cos \theta_3)}{\gamma^2 (\cos \theta_3 - \beta)^2 + \sin^2 \theta_3} d\theta_3 d\phi_3$$

Sensitivity to mass-coupling domain



- In keV region the signal is hidden by QED.
- In the low energy region the signal is dominate than QED

Light source

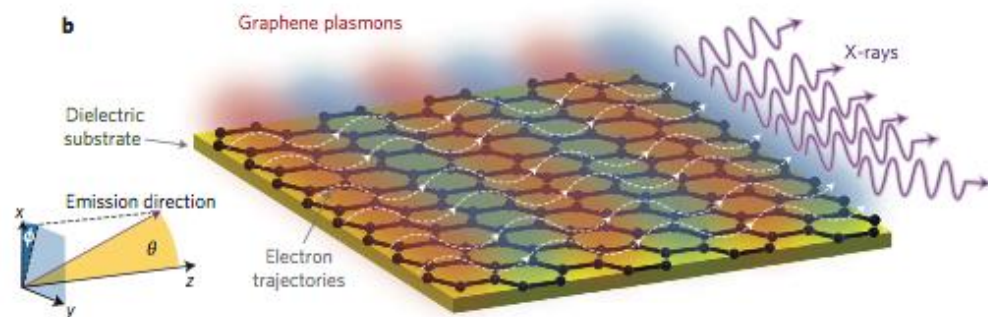
- In the 10 eV–1 keV range
The generation of higher harmonics by shooting high-intensity laser pulses into material targets.
- In the 0.1–10 eV range
variable-wavelength lasers based on optical parametric amplification.

Light sources



http://www.spring8.or.jp/ja/news_publications/press_release/2011/111005/

- X-ray laser(1 – 10 keV)
 - SACLA : $N_k \sim 10^{11}$ at 60 Hz
undulator length of 90 m
- Although introducing three long XFEL lines would likely be difficult.
↓
Incoherent light source at the creation part.
- Graphene based laser



GP-based free-electron source Ref.[4]

Summary

- We formulated a stimulated photon–photon scattering process via an s-channel pNGB exchange in an asymmetric head-on collision system that would be applicable to the mass range of 0.1eV to 10keV.
- In the keV region, the effect of QED is dominant and resonance can not be observed, but there is a prospect of observation in a low region.

Reference

- [1] K. Homma and Y. Toyota, arXiv:1701.04282, (2017).
- [2] K. Homma, T. Hasebe, and K. Kume, Prog. Theor. Exp. Phys. 083C01 (2014).
- [3] F. Moulin and D. Bernard, Opt. Commun. 164, 137 (1999); E. Lundström et al., Phys. Rev. Lett. 96, 083602 (2006); J. Lundin et al., Phys. Rev. A 74, 043821 (2006); D. Bernard et al., Eur. Phys. J. D 10, 141 (2000).
- [4] L. J. Wong, I. Kaminer, O. Ilic, J. D. Joannopoulos, and M. Soljačić, Nature Photonics 10, 46 (2016).

Back up

Setting parameter

Lorentz factor to boost CMS energies	$\gamma = 1.5$
Scattering angle in CMS	$\theta = \pi/4$
Incident angle of induction beam	$\theta_4 = 1.65 \text{ rad}$
Scattering angle of signal photons	$\theta_3 = 0.31 \text{ rad}$
Common f-number of creation beams	$F = 100$
Induction beam f-number	$F_4 = 10$
Common duration time for creation beams	$\tau = Z_{R2}/c$
Integration time from longest Rayleigh length	$\tau_L = Z_{R2}/c$
Common energy uncertainty of creation beams	$R = 5\%$
Collision repetition rate	$f = 1 \text{ Hz}$
Data accumulation time	$T = 10^6 \text{ s}$
Total number of signal photons	$Y = 100$
Detector efficiency to signal photons	$\epsilon = 100\%$

Creation beam energy in CMS	$\omega = 3.50 \text{ keV}$
Higher creation beam energy in ACS	$\omega_1 = 9.16 \text{ keV}$
Lower creation beam energy in ACS	$\omega_2 = 1.34 \text{ keV}$
Induction beam energy in ACS	$\omega_4 = 2.48 \text{ keV}$
Signal photon energy in ACS	$\omega_3 = 8.02 \text{ keV}$
