



NLO EW corrections in VH production in the PO framework

Andrea Patteri

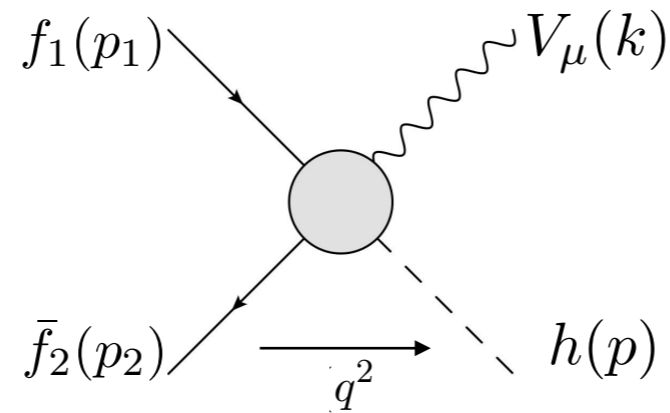
WG2 post-YR4 kickoff
CERN, 08.05.2017

Outline:

- Generalities and motivations
- Sketch of computation
- Results
- Future perspectives and conclusions

POs in Associated Higgs Production

Parametrisation of Associated Production (AP) processes with POs:

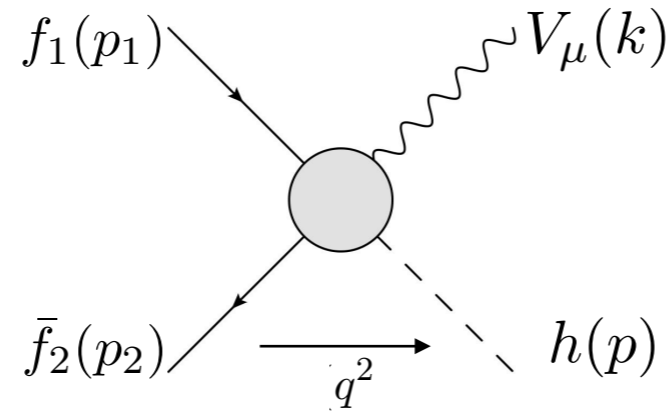


$$V_\mu = W_\mu^\pm, Z_\mu$$

$$\begin{aligned} \mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) &= 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ &\times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right] \end{aligned}$$

POs in Associated Higgs Production

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- ▶ **Longitudinal:** $F_L^{f_1 f_2 V}(q^2) = \kappa_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \frac{\epsilon_V^{f_1 \bar{f}_2}}{M_V^2},$
- ▶ **Transverse:** $F_T^{f_1 f_2 V}(q^2) = \epsilon_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \delta_{VZ} \epsilon_{Z\gamma} \frac{e Q_{f_1}}{q^2},$
- ▶ **CP-odd:** $F_{CP}^{f_1 f_2 V}(q^2) = \epsilon_{VV}^{CP} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} - \delta_{VZ} \epsilon_{Z\gamma}^{CP} \frac{e Q_{f_1}}{q^2},$

Pseudo
Observables

Purposes

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

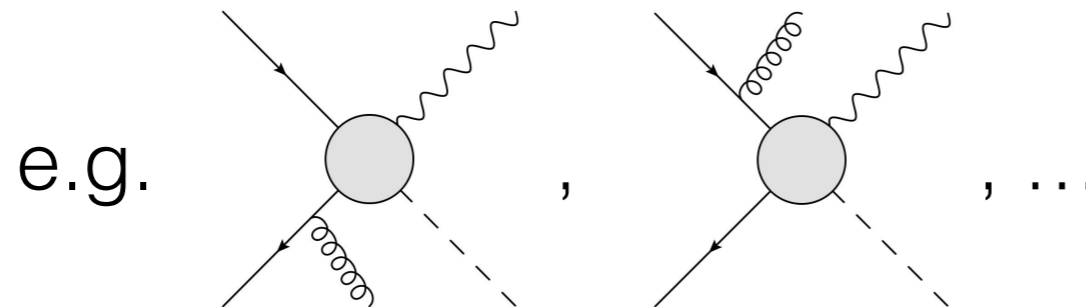
Such analytical decomposition is well motivated at threshold.

Purposes

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

Such analytical decomposition is well motivated at threshold.

- QCD (& QED) gives distortions due to IR physics (Greljo)

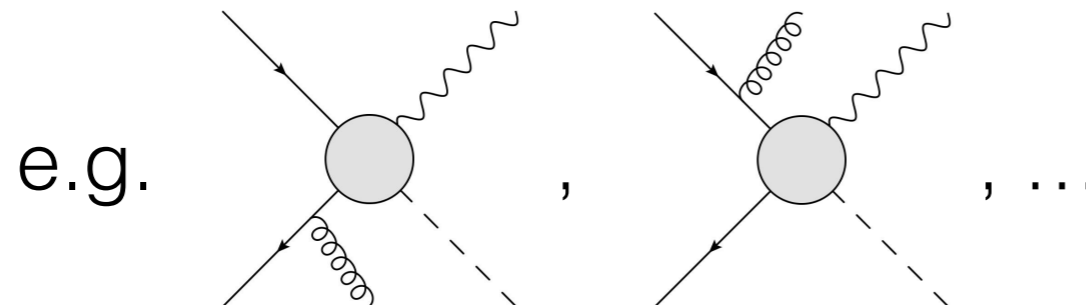


Purposes

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

Such analytical decomposition is well motivated at threshold.

- QCD (& QED) gives distortions due to IR physics (Greljo)



- At high energy \Rightarrow kinematic deformation due to EW sector

Assumptions and approximations

Frame of the computation:

- Kinematic region:
 - on-shell external momenta
 - $S \sim T \sim U \sim E^2 \gg M_W^2$
- Retaining only double-logarithmic (DL) corrections (Sudakov double logarithms):

$$\log^2 \left(\frac{S}{M^2} \right), \quad \log^2 \left(\frac{T}{M^2} \right), \quad \log^2 \left(\frac{U}{M^2} \right)$$

- Effective expansion parameter: $v/E \ll 1$

Neglecting terms order $\mathcal{O}(v^2/E^2)$

↳ drawback: $\left(\frac{1}{P_Z(q^2)} \sim \frac{1}{P_W(q^2)} \sim \frac{1}{P_A(q^2)} \right) + \mathcal{O}(v^2/E^2)$

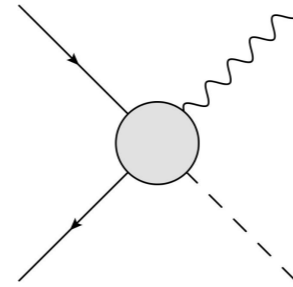
Brief outline of the computation

Computation performed in a diagrammatic fashion:

- I. Determination of Feynman diagrams giving DL contributions
- II. Derivation of a master formula for those Feynman diagrams
- III. Careful computation of all contributions

computation of DL contributions

Tree-level
process:



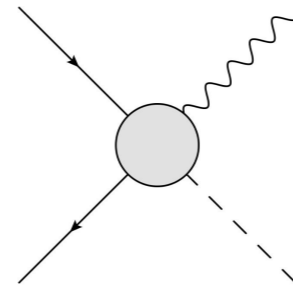
I. Which 1-loop diagrams give DLs?

Pozzorini, hep-ph/0201077

Feynman gauge computation:

computation of DL contributions

Tree-level
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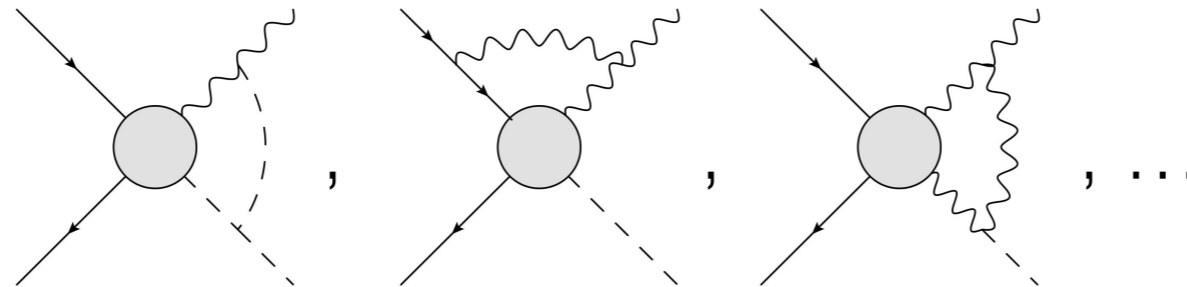


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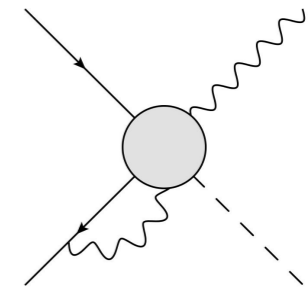
Feynman gauge computation:

1. Topologies \Rightarrow particle exchange between external legs

e.g.

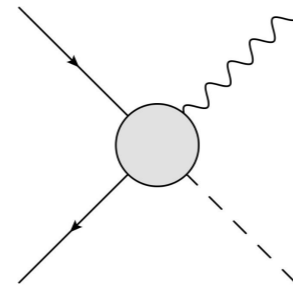


no:



computation of DL contributions

Tree-level
process:

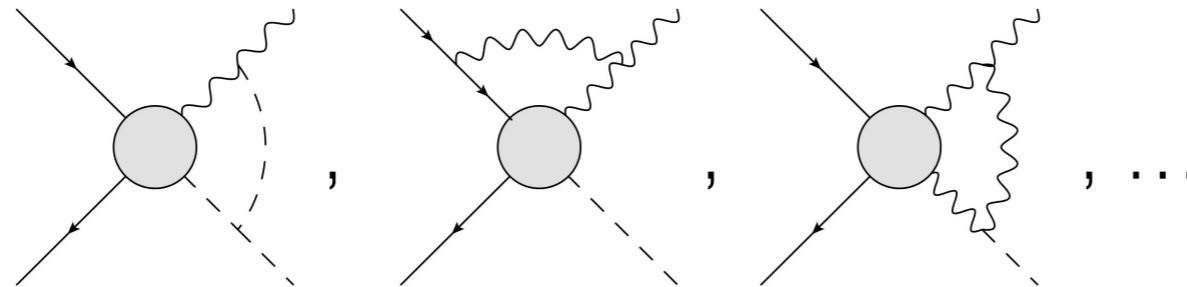


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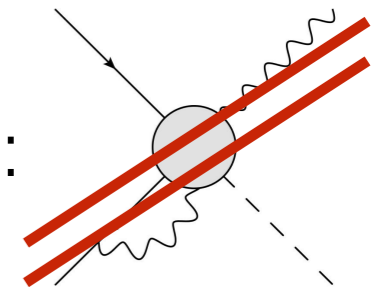
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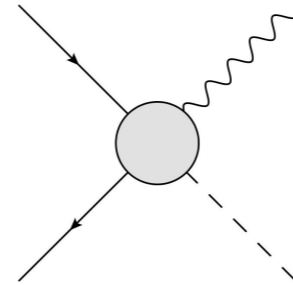


no:



computation of DL contributions

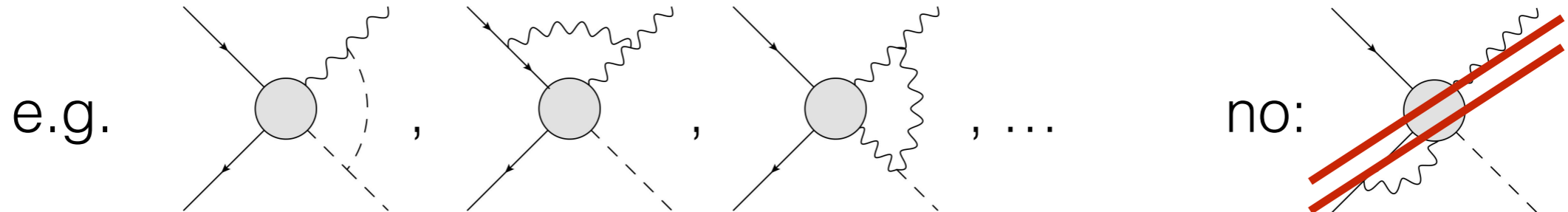
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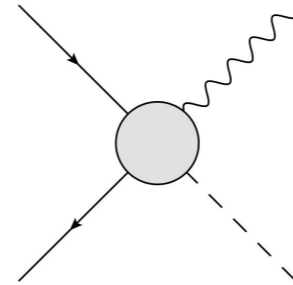
1. Topologies \Rightarrow particle exchange between external legs



2. Loop momentum \Rightarrow soft and collinear exchanged particle

computation of DL contributions

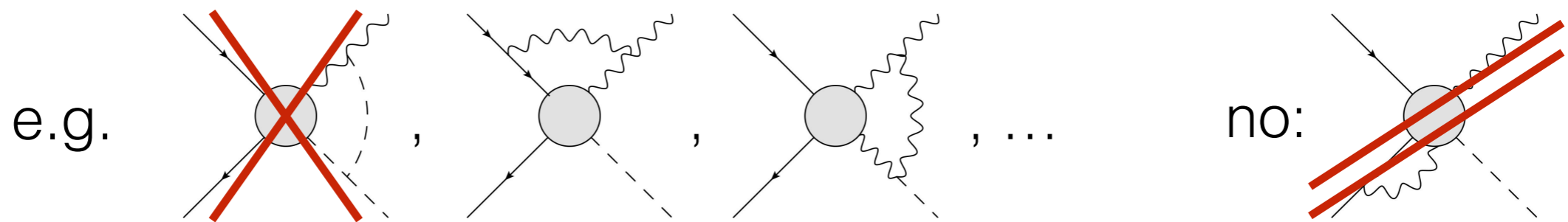
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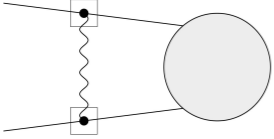
3. v/E expansion \Rightarrow exchanged particle: vector boson

computation of DL contributions

All possible topologies giving rise to DL mass singularities:

	A	B
I		\emptyset
II	1: , 2:	1: , 2:
III	1: , 2:	1: , 2:
IV		1: , 2:

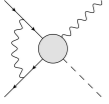
II. Master formula for DL contributions

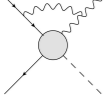


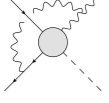
$$= \frac{1}{(4\pi)^2} \eta_{\mu\nu} R_i^\mu R_j^\nu \mathcal{M}_0 \left[-i (4\pi)^2 \int \frac{d^d q}{(2\pi)^d} \frac{1}{D_i D_j D_V} \right]$$

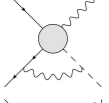
$$= \frac{1}{(4\pi)^2} \eta_{\mu\nu} R_i^\mu R_j^\nu \mathcal{M}_0 \cdot \frac{1}{4 p_i \cdot p_j} \text{DL}(2 p_i \cdot p_j, V, \varphi_i, \varphi_j) .$$

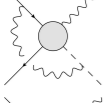
III. Careful computation of all contributions

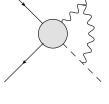


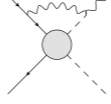
$$= \frac{\alpha}{4\pi} I_{f_1' f_1}^{\bar{V}_s} I_{f_2' f_2}^{V_s} \cdot \text{DL}(S, V_s, f_1, f_2) \cdot \mathcal{M}_0^{f_1' f_2' \bar{V} h} ,$$


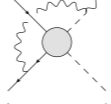
$$= \frac{\alpha}{4\pi} I_{f_1' f_1}^{\bar{V}_s} I_{\bar{V}' \bar{V}}^{V_s} \cdot \text{DL}(T, V_s, f_1, V) \cdot \left(\mathcal{M}_0^{f_1' f_2' \bar{V} h} + \frac{p_1 \cdot \epsilon_V}{T} k^\rho \mathcal{G}_{0\rho}^{f_1' f_2' \bar{V} h} \right) ,$$


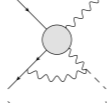
$$= \frac{\alpha}{4\pi} I_{f_2' f_2}^{\bar{V}_s} I_{\bar{V}' \bar{V}}^{V_s} \cdot \text{DL}(U, V_s, f_2, V) \cdot \left(\mathcal{M}_0^{f_1' f_2' \bar{V} h} + \frac{p_2 \cdot \epsilon_V}{U} k^\rho \mathcal{G}_{0\rho}^{f_1' f_2' \bar{V} h} \right) ,$$


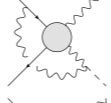
$$= \frac{\alpha}{4\pi} \kappa_{VV} I_{f_2' f_2}^{\bar{V}_s} I_{\phi_{V_s} h}^V \cdot \text{DL}(T, V_s, f_2, h) \cdot \mathcal{M}_0^{f_1' f_2' \bar{V} \phi_{V_s}} ,$$


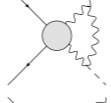
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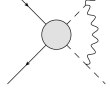
$$= \frac{\alpha}{4\pi} \kappa_{VV} I_{\bar{V}' \bar{V}}^{\bar{V}_s} I_{\phi_{V_s} h}^{V_s} \cdot \text{DL}(S, V_s, V, h) \cdot \left(\mathcal{M}_0^{f_1' f_2' \bar{V} \phi_{V_s}} + \frac{p_h \cdot \epsilon_V}{S} k^\rho \mathcal{G}_{0\rho}^{f_1' f_2' \bar{V} \phi_{V_s}} \right) ,$$


$$= -\frac{\alpha}{4\pi} e \Upsilon_{VV_s}^{\phi'} I_{f_1' f_1}^{\bar{V}_s} \frac{v}{T} \cdot \text{DL}(T, V_s, f_1, V) \cdot p_1 \cdot \epsilon_V \mathcal{M}_0^{f_1' f_2' \bar{\phi}' h} ,$$


$$= -\frac{\alpha}{4\pi} e \Upsilon_{VV_s}^{\phi'} I_{f_1' f_1}^{\bar{V}_s} \frac{v}{U} \cdot \text{DL}(U, V_s, f_2, V) \cdot p_2 \cdot \epsilon_V \mathcal{M}_0^{f_1' f_2' \bar{\phi}' h} ,$$


$$= \frac{\alpha}{4\pi} e \kappa_{V_s V_s} \Upsilon_{V_s V_s}^h I_{f_2' f_2}^{\bar{V}_s} \frac{v}{T} \cdot \text{DL}(T, V_s, f_2, h) \cdot \epsilon_V^\mu p_2^\nu \mathcal{G}_{0\mu\nu}^{f_1' f_2' \bar{V} V_s} ,$$


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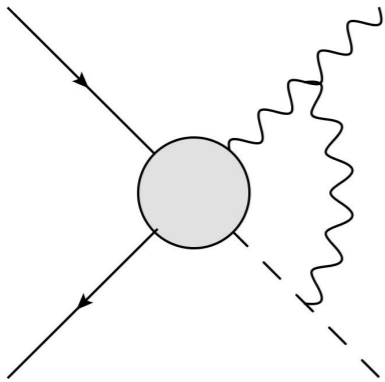
$$= \frac{\alpha}{4\pi} e \kappa_{V_s V_s} \Upsilon_{V_s V_s}^h I_{VV'}^{\bar{V}_s} \frac{v}{S} \cdot \text{DL}(S, V_s, V, h) \cdot \left(\epsilon_V^\mu k^\nu - \frac{1}{2} \epsilon_V^\nu k^\mu \right) \mathcal{G}_{0\mu\nu}^{f_1' f_2' \bar{V} V_s} ,$$


$$= -\frac{\alpha}{4\pi} e \kappa_{V_s V_s} \Upsilon_{V_s V_s}^{\phi'} I_{\bar{\phi}'' h}^{V_s} \frac{v}{S} \cdot \text{DL}(S, V_s, V, h) \cdot p_h \cdot \epsilon_V \mathcal{M}_0^{f_1' f_2' \bar{\phi}''} .$$

POs and loop computations



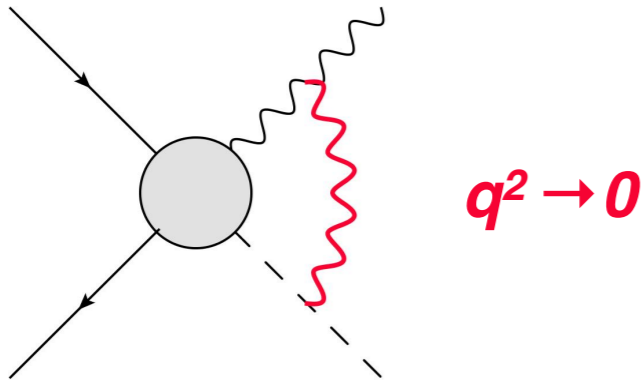
use of PO parametrisation into a loop computation?
is it consistent?



POs and loop computations



use of PO parametrisation into a loop computation?
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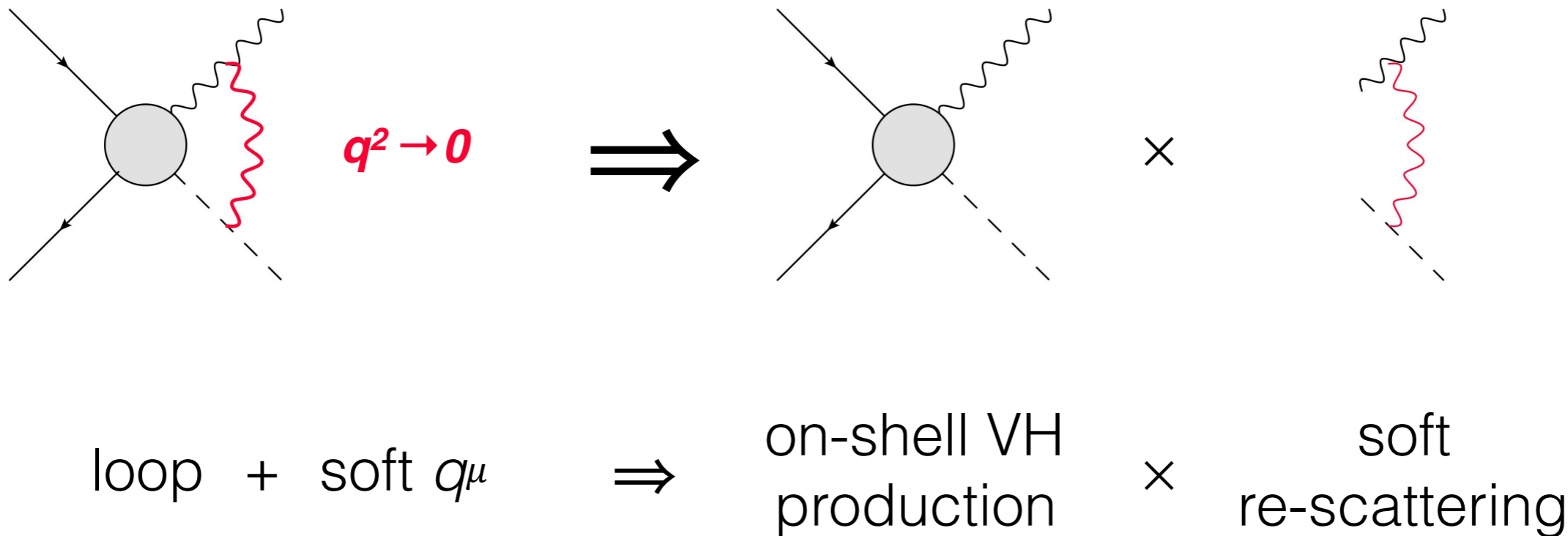


loop + soft q_μ

POs and loop computations



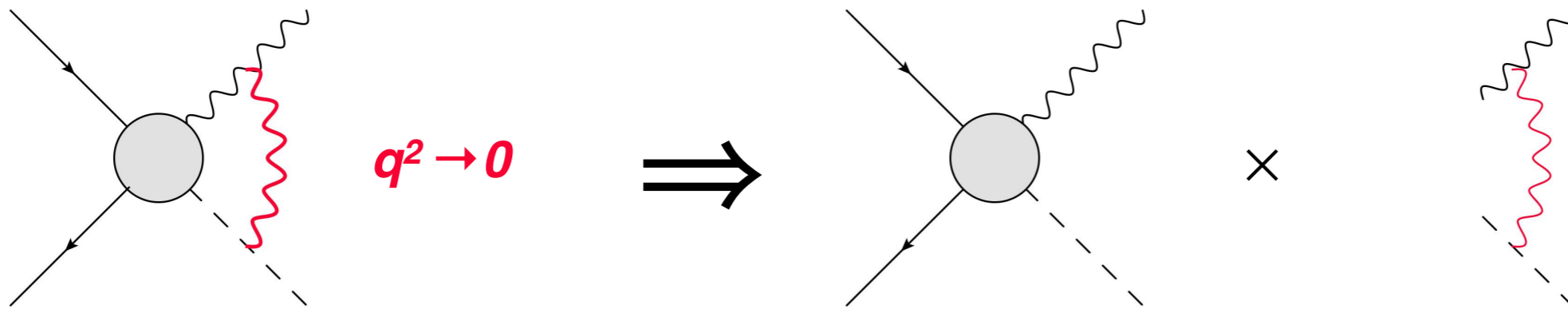
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POs and loop computations



use of PO parametrisation into a loop computation?
is it consistent?



loop + soft q_μ



on-shell VH
production

×

soft
re-scattering



PO
parametrisation

master formula
for DLs

Results

results

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

▶ **Longitudinal:** $F_L^{f_1 f_2 V}(q^2) = \kappa_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \frac{\epsilon_V^{f_1 f_2}}{M_V^2},$

▶ **Transverse:** $F_T^{f_1 f_2 V}(q^2) = \epsilon_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \delta_{VZ} \epsilon_{Z\gamma} \frac{e Q_{f_1}}{q^2},$

▶ **CP-odd:** $F_{CP}^{f_1 f_2 V}(q^2) = \epsilon_{VV}^{CP} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} - \delta_{VZ} \epsilon_{Z\gamma}^{CP} \frac{e Q_{f_1}}{q^2},$

{ κ_{ZZ}, κ_{WW} $\epsilon_Z^{f_1 f_2}, \epsilon_W^{f_1 f_2}$
 $\epsilon_{ZZ}, \epsilon_{WW}, \epsilon_{Z\gamma}$
 $\epsilon_{ZZ}^{CP}, \epsilon_{WW}^{CP}, \epsilon_{Z\gamma}^{CP}$

results

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

▶ **Longitudinal:** $F_L^{f_1 f_2 V}(q^2) = \kappa_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \frac{\epsilon_V^{f_1 f_2}}{M_V^2},$

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▶ **CP-odd:** $F_{CP}^{f_1 f_2 V}(q^2) = \epsilon_{VV}^{CP} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} - \delta_{VZ} \epsilon_{Z\gamma}^{CP} \frac{e Q_{f_1}}{q^2},$

}

κ_{ZZ}, κ_{WW}

$\epsilon_Z^{f_1 f_2}, \epsilon_W^{f_1 f_2}$

$\epsilon_{ZZ}, \epsilon_{WW}, \epsilon_{Z\gamma}$

$\epsilon_{ZZ}^{CP}, \epsilon_{WW}^{CP}, \epsilon_{Z\gamma}^{CP}$

EW corrections mix POs among themselves

But only within the same “class” (same tensor and pole structure)

results

$$\mathcal{A}(f_1(p_1)\bar{f}_2(p_2) \rightarrow V(k)h(p)) = 2i \frac{M_V^2}{v} \bar{f}_2(p_2)\gamma_\nu f_1(p_1)\epsilon_\mu^{V*}(k) \\ \times \left[F_L^{f_1 f_2 V}(q^2) \eta^{\mu\nu} + F_T^{f_1 f_2 V}(q^2) \frac{q^\mu k^\nu - (q \cdot k)\eta^{\mu\nu}}{M_V^2} + F_{CP}^{f_1 f_2 V}(q^2) \frac{\epsilon^{\mu\nu\alpha\beta} q_\alpha k_\beta}{M_V^2} \right]$$

▶ **Longitudinal:** $F_L^{f_1 f_2 V}(q^2) = \kappa_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \frac{\epsilon_V^{f_1 f_2}}{M_V^2},$

▶ **Transverse:** $F_T^{f_1 f_2 V}(q^2) = \epsilon_{VV} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} + \delta_{VZ} \epsilon_{Z\gamma} \frac{e Q_{f_1}}{q^2},$

▶ **CP-odd:** $F_{CP}^{f_1 f_2 V}(q^2) = \epsilon_{VV}^{CP} \frac{g_V^{f_1 \bar{f}_2}}{P_V(q^2)} - \delta_{VZ} \epsilon_{Z\gamma}^{CP} \frac{e Q_{f_1}}{q^2},$

{ κ_{ZZ}, κ_{WW} $\epsilon_Z^{f_1 f_2}, \epsilon_W^{f_1 f_2}$
 $\epsilon_{ZZ}, \epsilon_{WW}, \epsilon_{Z\gamma}$
 $\epsilon_{ZZ}^{CP}, \epsilon_{WW}^{CP}, \epsilon_{Z\gamma}^{CP}$

EW corrections mix POs among themselves

But only within the same “class” (same tensor and pole structure)

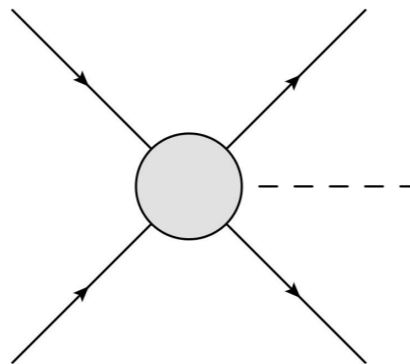
e.g:

$$\kappa_{ZZ} \rightarrow \kappa_{ZZ} + \frac{\alpha}{\pi} \left[\kappa_{ZZ} \left((\dots) \log^2 \frac{S}{M_Z^2} + \dots \right) + \kappa_{WW} \left((\dots) \log^2 \frac{S}{M_W^2} + \dots \right) \right],$$

$$\epsilon_{ZZ} \rightarrow \epsilon_{ZZ} + \frac{\alpha}{\pi} \left[\epsilon_{ZZ} \left((\dots) \log^2 \frac{S}{M_Z^2} + \dots \right) + \epsilon_{WW} \left((\dots) \log^2 \frac{S}{M_W^2} + \dots \right) + \epsilon_{Z\gamma} (\dots) \right],$$

Future perspectives

- Montecarlo simulations of the impact of such DL EW corrections
- Inclusion of these corrections into the PO tool
- Extension of the computation to other interesting processes (e.g. vector boson fusion)



Conclusions

- NLO EW corrections can be sizeable when far from threshold
- Computation of relevant loop EW correction (i.e. DLs) can be consistently performed within the PO framework
- Those corrections introduce a mixing between POs, but a very contained and reasonable one
- Study of EW corrections will be pushed further and could be also extended to other processes