Open questions/Prospects for collider DM searches

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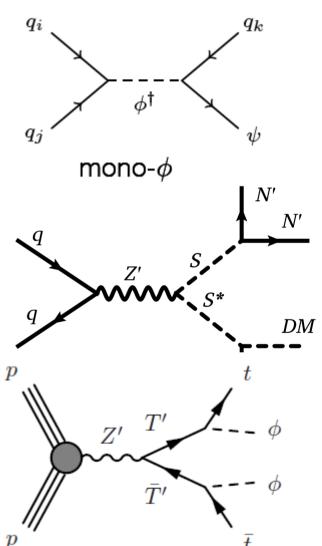
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Future of collider searches for Dark Matter at the LPC FNAL 27-28 July 2017

New Signatures/Topologies/Models

Which we did not study yet and should be ready for Just a few examples

- Monojet from scalar resonance (see talk by Matt Buckley)
- MET recoiling Displaced vertices and/or heavy boosted objects from Z'→ S S*→ (N',N') + (DM,DM) chain (AB, Eung Jin Chun, Dan Locke, Priyotosh Bandyopadhyay, Rusa Mandal)
- tt+MET from Z'→ T'T'→ tt+DMDM
 1707.07000 (AB, Patrick Schaefers, Thomas Flacke, Bithika Jain)



Do we need new tools/improvement of the existing ones?

- Better BG prediction/simulation?
 - ▶ NLO QCD+ NLO EW SM BG simulation can be done with Sherpa
- Common repository for BSM model implementation?
 - FeynRules website, HEPMDB
- Improvements on DM relic density and DM DD and ID signal predictions?
 - ▶ Freeze-in scenario will be in micrOMEGAs in ~ two weeks

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Common framework/agreement on the recasting of the experimental limits:

- (Re)interpreting the results of new physics searches at the LHC Workshop (CERN)
- Supplemental Material section would be very useful, see e.g.

arXiv:1508.04094 (LHCb)

Search for hidden-sector bosons in $B^0 \to K^{*0} \mu^+ \mu^-$ decays

Supplemental Material

The limits reported in the Letter assume a spin-zero hidden-sector boson. To convert these into limits for a spin-one boson, the ratio of efficiencies for the spin-one to spin-zero cases must be accounted for. Determining this ratio involves integrals of the form

$$\frac{\int f_j(\vec{\Omega})\epsilon(\vec{\Omega}, m^2(\mu^+\mu^-))d\vec{\Omega}}{\int f_{1c}(\vec{\Omega})\epsilon(\vec{\Omega}, m^2(\mu^+\mu^-))d\vec{\Omega}},$$

Combining collider searches, relic density and DM detection (DD and ID) limits

- Should experimental papers do this?
- How generic the statement is
 - Non-applicable for non-standard cosmology with low reheating temperature
- Importance of the running of the operator
 - In case of axial operators, e.g.

$$c_A^{(q)} c_\chi \overline{\chi} \gamma^\mu \chi \overline{q} \gamma_\mu \gamma_5 q$$
 (D7) or $c_A^{(q)} c_\phi \phi^\dagger \overleftrightarrow{\partial}_\mu \phi \overline{q} \gamma^\mu \gamma_5 q$ (C4)

couplings $\mathbf{c_v}^{(q)}$ arise due to the running of the wilson coeffcient $\mathbf{c_A}^{(q)}$ leading to sizable constraints on the DM DD constraints

One can use runDM program (github.com/bradkav/runDM) by
 F. D'Eramo, B. J. Kavanagh & P. Panci

$$c_A^{(u)}, c_A^{(d)}, c_V^{(u)}, c_V^{(d)} = (1,1,0,0)[5\text{TeV}] \rightarrow (1.1, 1.1, 0.04, -0.07)[1\text{GeV}]$$

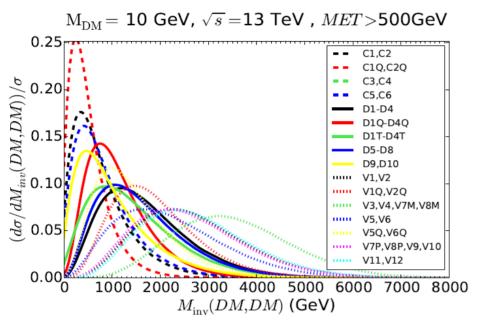
AB, Bertuzzo, Caniu, Eboli, di Cortona

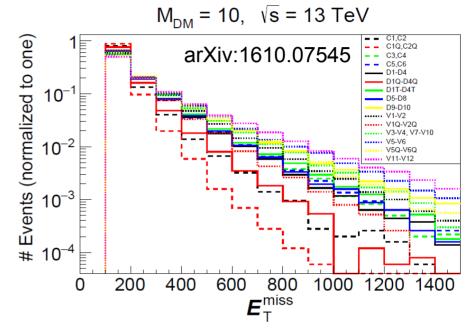
Distinguishing DM operators/theories

M(DM,DM) distributions



are correlated with Different MET shapes





- energy dependence of the DM operator $\rightarrow M_{\text{DMDM}}$ distributions \rightarrow slopes of MET
- projection for 300 fb⁻¹: some operators C1-C2,C5-C6,D9-D10,V1-V2,V3-V4,V5-V6 and V11-12 can be distinguished from each other
- Application beyond EFT: when the DM mediator is not produced on-the-mass-shell and M_{DMDM} is not fixed: t-channel mediator or mediators with mass below 2M_{DM}