

# Standard Model Prediction of $R(D^*)$

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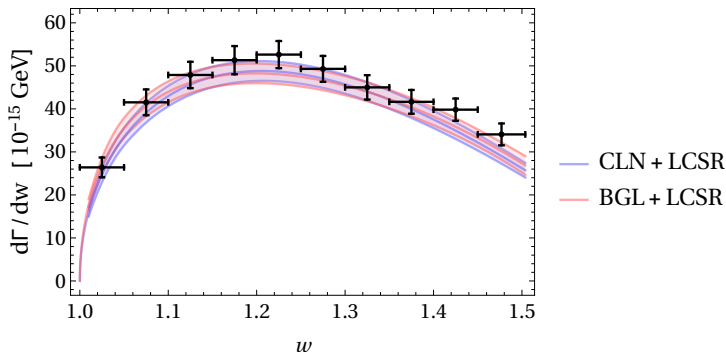
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based on work in progress with Dante Bigi and Paolo Gambino  
and Phys.Lett. B769 (2017) 441-445 [1703.06124]

## Belle has new (preliminary) data

- First time  $w$  and angular **deconvoluted distributions independent** of parametrization.  
↳ Possible to use **different parametrizations**.



$$w = \frac{m_B^2 + m_{D^*}^2 - q^2}{2m_B m_{D^*}}, \quad q^2 = (p_B - p_{D^*})^2$$

## Form factor parametrizations

Boyd Grinstein Lebed: **Unitarity**, crossing symmetry, analyticity

$$f_i(z) = \frac{1}{P_i(z)\phi_i(z)} \sum_{n=0}^{\infty} a_n^i z^n, \quad z = \frac{\sqrt{1+w} - \sqrt{2}}{\sqrt{1+w} + \sqrt{2}}, \quad w = \frac{m_B^2 + m_{D^*}^2 - q^2}{2m_B m_{D^*}}.$$

- Theory input: **unitarity constraints**  $\sum_{n=0}^{\infty} (a_n^i)^2 < 1$ .
- $0 < z < 0.056$  for  $B \rightarrow D^* l \nu \Rightarrow$  truncation at  $N = 2$  enough,  $z^3 \sim 10^{-4}$ .
- $P_i(z)$ : “Blaschke factor”: removes poles,  $\phi_i(z)$ : phase space factors.

Caprini Lellouch Neubert: Use **HQET** relations between form factors

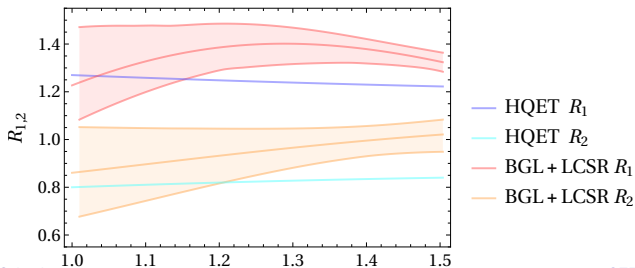
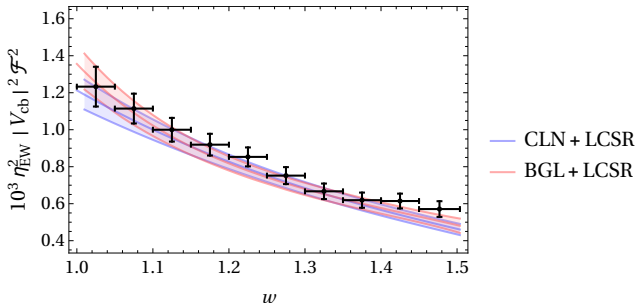
- **Less** parameters, **slope** of form factor ratios  $R_i$  fixed.

$$h_{A_1}(w) = h_{A_1}(1) \left( 1 - 8\rho^2 z + (53\rho^2 - 15)z^2 - (231\rho^2 - 91)z^3 \right),$$

$$R_1(w) = R_1(1) - 0.12(w - 1) + 0.05(w - 1)^2,$$

$$R_2(w) = R_2(1) + 0.11(w - 1) - 0.06(w - 1)^2$$

## Fit results for $B \rightarrow D^* l \nu$



## Different results for $V_{cb}$

[Bigi Gambino Schacht, 1703.06124, agreeing with Grinstein Kobach, 1703.08170]

BGL Fit:	Data + lattice	Data + lattice + LCSR
$\chi^2/\text{dof}$	27.9/32	31.4/35
$ V_{cb} $	$0.0417 \left( \begin{smallmatrix} +20 \\ -21 \end{smallmatrix} \right)$	$0.0404 \left( \begin{smallmatrix} +16 \\ -17 \end{smallmatrix} \right)$

CLN Fit:	Data + lattice	Data + lattice + LCSR
$\chi^2/\text{dof}$	34.3/36	34.8/39
$ V_{cb} $	$0.0382 (15)$	$0.0382 (14)$

- $|V_{cb}|$  central values deviate by 9% and 6% (with LCSR).
- **LCSR**: Light Cone Sum Rule results [Faller, Khodjamirian, Klein, Mannel 0809.0222]  
 $h_{A_1}(w_{\max}) = 0.65(18)$ ,  $R_1(w_{\max}) = 1.32(4)$ ,  $R_2(w_{\max}) = 0.91(17)$ .
- **Lattice**:  $h_{A_1}(1) = 0.906 \pm 0.013$ . [FNAL/MILC 1403.0635]

$$\text{Anatomy of } R(D^*) \equiv \frac{\int_1^{w_{\tau,\max}} dw d\Gamma_{\tau}/dw}{\int_1^{w_{\max}} dw d\Gamma/dw}$$

Differential decay rate for  $B \rightarrow D^* \tau \nu_{\tau}$

[BGL, hep-ph/9705252]

$$\begin{aligned} \frac{d\Gamma_{\tau}}{dw} &= \frac{d\Gamma_{\tau,1}}{dw} + \frac{d\Gamma_{\tau,2}}{dw} \\ \frac{d\Gamma_{\tau,1}}{dw} &= \left(1 - m_{\tau}^2/q^2\right)^2 \left(1 + m_{\tau}^2/(2q^2)\right) \frac{d\Gamma}{dw} \\ \frac{d\Gamma_{\tau,2}}{dw} &= |V_{cb}|^2 m_{\tau}^2 \times \text{kinematics} \times \mathcal{F}_2(z)^2 \end{aligned}$$

- $d\Gamma/dw$ : **Measured** differential decay rate of  $B \rightarrow D^* l \nu$  with  $m_l = 0$ , depends on axial vector form factors  $f$ ,  $\mathcal{F}_1$  and vector form factor  $g$ .
- $\mathcal{F}_2$ : Additional **unconstrained** pseudoscalar **form factor**.
- $d\Gamma_{\tau,2}/dw$  contributes  $\sim 10\%$  to  $R(D^*)$ .
- Common normalization/notation:

$$R_0 = \frac{P_1}{A_1} = m_{D^*} \left( \frac{1+w}{1+r} \right) \frac{\mathcal{F}_2}{f}, \quad r = m_{D^*}/m_B$$

## Calculating $R_0(w)$

Heavy quark limit

[BGL, hep-ph/9705252]

$$R_0(w) = 1 \quad \forall w.$$

HQET at  $O(\Lambda/m_{c,b})$

[Bernlochner Ligeti Papucci Robinson (BLPR), 1703.05330,

Neubert, Phys. Lett. B264 (1991) 455; hep-ph/9408290, hep-ph/9306320]

$$R_0(w) = R_0(1) + R'_0(1)(w - 1),$$

$$R_0(1) = 1.09 + 0.25\eta(1),$$

$$R'_0(1) \equiv \left. \frac{d}{dw} R_0(w) \right|_{w=1} = -0.18 + 0.87\hat{\chi}_2(1) + 0.06\eta(1) + 0.25\eta'(1),$$

Sum rule parameters

$$\eta(1) = 0.62 \pm 0.02$$

$$\eta'(1) = 0 \pm 0.2$$

$$\hat{\chi}_2(1) = -0.06 \pm 0.02$$

How large could **higher order corrections** to  $R_0(w)$  beyond  $\mathcal{O}(\Lambda/m_{c,b}, \alpha_s)$  be?

**Rough dimensional analysis of higher order corrections**

$$\Lambda^2/m_c^2 \sim (0.3)^2 \simeq 10\%$$

$$\alpha_s(m_c)^2 \sim (0.4)^2 \simeq 16\%$$

$$\alpha_s(m_c) \times \Lambda/m_c \sim 0.3 \times 0.4 \simeq 12\%$$



## Direct Comparison of HQET and Lattice QCD Results

$A_1/V_1$  at  $w = 1$ : Central values deviate by up to 12%.

$$\text{Lattice QCD: } \left. \frac{A_1(w=1)}{V_1(w=1)} \right|_{\text{FNAL/MILC}} = 0.859(14) \quad [\text{obtained from 1403.0635, 1503.07237}]$$

$$\text{HQET: } \left. \frac{A_1(w=1)}{V_1(w=1)} \right|_{\text{CLN}} = 0.948 \quad [\text{hep-ph/9712417}]$$

$$\text{HQET: } \left. \frac{A_1(w=1)}{V_1(w=1)} \right|_{\text{BLPR}} = 0.966(30) \quad [\text{obtained from 1703.05330}]$$

$f_0/f_+$  at  $w = 1$ : Central values deviate by 3%.

$$\text{Lattice QCD: } \left. \frac{f_0(w=1)}{f_+(w=1)} \right|_{\text{FNAL/MILC}} = 0.753(3) \quad [\text{obtained from 1503.07237}]$$

$$\text{HQET: } \left. \frac{f_0(w=1)}{f_+(w=1)} \right|_{\text{CLN}} = 0.775 \quad [\text{obtained from hep-ph/9712417}]$$

**HPQCD:** Less precise but generally consistent results for the form factors.  
 $A_1/V_1$ : only marginally consistent with FNAL, but even lower result.

# Direct Comparison of HQET and Lattice QCD Results

Slope of  $f_0/f_+$  at  $w = 1$ : Central values deviate by 20%

Lattice QCD:  $\left. \frac{d}{dw} \left( \frac{f_0}{f_+} \right) \right|_{w=1, \text{FNAL/MILC}} = 0.457(35)$  [obtained from 1503.07237]

HQET:  $\left. \frac{d}{dw} \left( \frac{f_0}{f_+} \right) \right|_{w=1, \text{CLN}} = 0.382$  [obtained from hep-ph/9712417]

# Implications of Dimensional Analysis and Comparison HQET $\Leftrightarrow$ Lattice QCD results

## Possible size of higher order corrections of HQET results

- Corrections could modify form factor ratios by  $\sim 12\%$ .  
↳ For prediction of  $R(D^*)$  vary  $R_0(w)$  in a band of  $12\%$ .

Take this into account by variation of additional parameter:

$$R_0(w, E) = E \left( R_0(1) + R'_0(1)(w - 1) \right)$$

$$\text{vary } E = 1.0 \pm 0.12$$

# Overview on Sources of Uncertainty

preliminary results

- Our analysis leads to a **central value**  $R(D^*) = 0.258$ .  
↳ Very good agreement to [BLPR, 1703.05330].

Error due to **experimental** error of measurement of  $B \rightarrow D^* l \nu$ .

$$\delta R(D^*) = 0.005$$

Theory error due to **sum rule** parameters.

- Scan:  $\delta R(D^*) = 0.003$
- Gaussian:  $\delta R(D^*) = 0.002$

Theory error due to **higher order** effects.

- Scan/Gaussian:  $\delta R(D^*) = \begin{matrix} +0.007 \\ -0.006 \end{matrix}$

# Total Uncertainties for $R(D^*)$

preliminary results

## BGL fit

Higher orders	Sum rule parameters	Prediction for $R(D^*)$
Scan	Scan	$0.258^{+15}_{-13}$
Scan	Gaussian	$0.258^{+12}_{-11}$
Gaussian	Gaussian	$0.258^{+9}_{-8}$

## CLN fit

Higher orders	Sum rule parameters	Prediction for $R(D^*)$
Scan	Scan	$0.257^{+15}_{-13}$
Scan	Gaussian	$0.257^{+12}_{-11}$
Gaussian	Gaussian	$0.257^{+9}_{-8}$

Experiment:  $0.310 \pm 0.015 \pm 0.008$  (HFAG average [1612.07233])

# Conclusions

- Belle has **new data**: Deconvoluted, independent of parametrization.
- Different parametrizations give **different results** for  $|V_{cb}|$ .
- $R(D^*)$  depends to an amount of  $\sim 10\%$  on the unconstrained form factor  $\mathcal{F}_2$ , which has to be **estimated by theory**.  
↳ **No lattice** calculation available, use **HQET** input from BLPR.
- Our **central value 0.258 agrees** well with the literature.
- We find a **larger uncertainty**, coming from **three sources**:
  - Experimental error in  $B \rightarrow D^* l \nu$ : **0.005**.
  - Sum rule parameters: **0.003** (scan),  
**0.002** (gaussian).
  - Higher order HQET corrections:  $\begin{matrix} +0.007 \\ -0.006 \end{matrix}$  (scan/gaussian).
- The **anomaly is persistent**.

[numbers: preliminary results]