

A closer look into the peak shape in nuclear A closer look into the peak shape in nuclear modification factor in relativistic heavy ion collisions modification factor in relativistic heavy ion collisions Rathijit Biswas¹, Sumit K. Saha², Souvik P. Adhya³, Sidharth K. Prasad¹, Supriya Das¹, Tapan K. Nayak⁴ and Sumit Basu⁵ ¹Department of Physics, Bose Institute, India ¹Department of Physics, Bose Institute, India ²Experimental High Energy Physics and Applications Group, Variable Energy Cyclotron Center, India ²Experimental High Energy Physics and Applications Group, Variable Energy Cyclotron Center, India 3 Institute of Particle and Nuclear Physics, Charles University, Czech Republic 3 Institute of Particle and Nuclear Physics, Charles University, Czech Republic ⁴CERN, Switzerland ⁴CERN, Switzerland

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Physics Motivation:

At very high temperature and energy density, a phase transition occurs from hadronic \gtrapprox matter to de-confined quarks and gluons, known as Quark Gluon Plasma(QGP). Partons produced in the early stages of the Hadrons collisions, while passing through the medium, loose energy via multiple scatterings and gluon radiations [1-3]. This leads to modified p_T spectra of final state hadrons for A-A collisions with respect to p-p collisions, the Phase Diagram for Hadronic key physical quantity being the Nuclear matter to QGP. modification factor (R_{AA}) and the ratio of

Summary

Where $d^2N_{AA}/dp_Td\eta$ represents the differential particle yield in nucleus-nucleus collisions and $d^2\sigma_{\text{on}}/dp_{\text{T}}dp$ is the cross-section in proton-proton collisions. In the above expression, nuclear overlap function $(T_{AA}$ is obtained from Glauber model and is proportional to the number of binary collisions $(\langle N_{\text{coll}}\rangle)$. $\langle N_{\text{cell}}\rangle$ and <N^{peri}_{coll}> are the average number of binary collisions in central and peripheral A-A collisions, respectively.

Energy species dependence of R_{CP} and R_{AA} as a function of p_{T} with different beam energies for available central collisions at ALICE [5] and STAR [1] experiments with EPOS comparison.

> beam energies at different $p_T^{}$ range in available central collisions for inclusive charged particles and for identified charged particles from EPOS. The left plot stands for ISF measured from R_{AA} and the right one from R_{CP}

- \triangleright We observe a shift of the peak position with energy as well as with different identified species for RAA. But for RCP, we only observe species dependency or in other words mass dependency.
- \triangleright The mismatch in the trends of the peak shifting suggests that for RCP it is not coming from collectivity. Whether this is coming from the radial boost or the mass hierarchy in the final states or both, is yet to conclude as the analysis is going on.
- \triangleright For all species, ISF follows a particular trend at different beam energy both from RAA and RCP measurements. In addition at a particular beam energy, the ISF value is mass dependent, for example, pions have a lower mass than kaon and proton and show a larger ISF value.

$$
R_{AA} \equiv \frac{d^2 N_{AA} / dp_T d\eta}{\langle T_{AA} \rangle d^2 \sigma_{pp} / dp_T d\eta}
$$

Based on the available results of $R_{AA}(p_T)$ and $R_{CP}(p_T)$ measurements from various experiments and that obtained from EPOS we have systematically studied shape of the Cronin peak [4] for all charged as well as identified charged particles for different energies. For the quantitative understanding of the Cronin peak, we have defined a quantity called the Integrated Suppression Fraction (ISF) and we studied it for different identified charged particles at different √s obtained from EPOS and inclusive charged particles from available experimental data.

$$
R_{CP} \equiv \frac{\langle N_{coll}^{peri} \rangle d^2 N_{cent} / dp_T d\eta}{\langle N_{coll}^{cent} \rangle d^2 N_{peri} / dp_T d\eta}
$$

ISF as a function of

We define the ISF as the ratio of two quantities integrated over a range of $\bm{{\mathsf{p}}}_\text{\tiny T},$

Definition:

The R_{AA} and R_{CP} are defined as:

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[1] J. Adams et al. [STAR Collaboration], Nucl. Phys. A 757 (2005) 102.

$$
ISF = \frac{\int_{\rho_T^{\text{min}}}^{\rho_T^{\text{max}}} [R_{AA} = 1] d\rho_T - \int_{\rho_T^{\text{min}}}^{\rho_T^{\text{max}}} R_{AA} d\rho_T}{\int_{\rho_T^{\text{min}}}^{\rho_T^{\text{max}}} [R_{AA} = 1] d\rho_T}
$$

we choose two p_T ranges, (a) 1 GeV/c to 2 GeV/c and (b) 1 GeV/c to 4 GeV/c to study the p_T dependence of the ISF. We show ISF as a function of collision energy from both R_{AA} and R_{CP} values

Evolution of Cronin peak

transventum (p_T) differentiated spectra from central over peripheral collisions, studied through the quantity R_{CP} .

