

# Jet structure in integrated EPOS3-HQ approach

**Iurii KARPENKO**

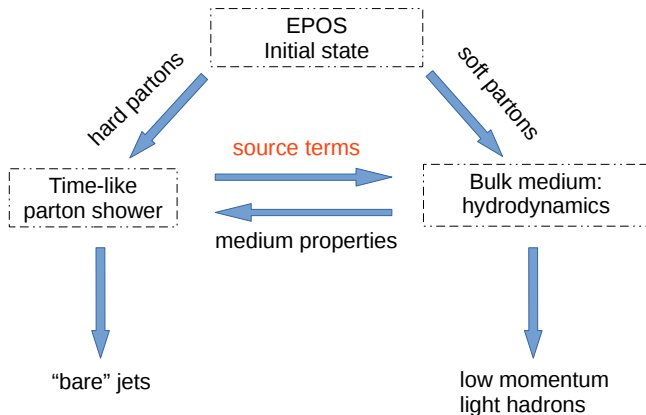
with Martin Rohrmoser, Joerg Aichelin, Pol Gossiaux, Klaus Werner

CNRS/SUBATECH Nantes  
Jan Kochanowski University

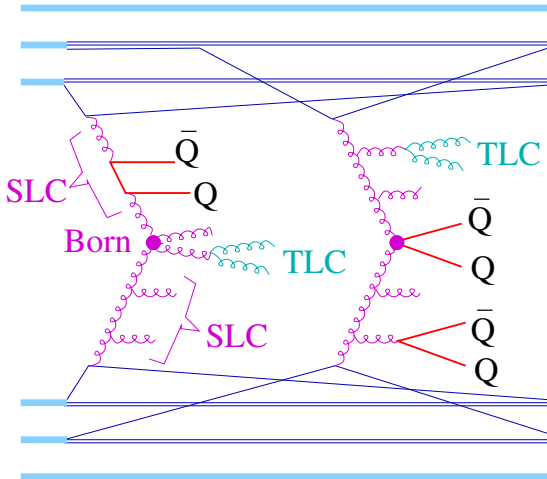


## The goal

To fully integrate jets into the EPOS3 model.



## EPOS initial state



### Parton-Based Gribov-Regge Theory

H. J. Drescher, M. Hladik, S. Ostapchenko, T. Pierog, K. Werner, Phys. Rept. 350, 93, 2001

Pomeron = parton ladder, treated as a kinky string.

Spacelike cascades including Born process in the EPOS IS provide partons with all  $p_T$  which are further separated into core and corona.

## Hydrodynamic background

**For this study** and in order to explore the effects in a clear way:

Averaged(smooth) hydrodynamic initial state compatible with EPOS3.

Equation of state: Laine & Schroeder '06, compatible with s95p-v1.2 EoS.

M. Laine, Y. Schroeder Phys. Rev. D73 (2006) 085009

3+1 dimensional viscous hydrodynamics:

$$T^{\mu\nu} = (\varepsilon + p)u^\mu u^\nu - p \cdot g^{\mu\nu} + \pi^{\mu\nu}$$

$$\partial_{;\nu} T^{\mu\nu} = 0, \quad \partial_{;\nu} N^\nu = 0$$

$$\langle u^\gamma \partial_{;\gamma} \pi^{\mu\nu} \rangle = -\frac{\pi^{\mu\nu} - \pi_{NS}^{\mu\nu}}{\tau_\pi} - \frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma$$

solved with vHLL code, Comput. Phys. Commun. 185 (2014), 3016

<https://github.com/yukarpenko/vhll>

## Time-like parton shower

Cascade made by **Martin Rohrmoser**:

- Monte Carlo simulation of DGLAP equations for a parton shower between virtuality scales  $Q_{\uparrow}$  (from Born process in EPOS) and  $Q_{\downarrow} = 0.6$  GeV.

- Radiative energy loss (virtuality gain) a la YaJEM:

$$\frac{dQ^2}{dt} = \hat{q}_R(t,x), \quad \hat{q}_R(t,x) = \frac{210}{1+53 \cdot T} T^3(t,x)$$

- Collisional energy loss: longitudinal drag

$$\frac{dp_{\parallel}}{dt} = -A(t,x), \quad A = \frac{\hat{q}_R}{0.09+0.715 \cdot T(t,x)/0.16}$$

- Mean lifetime of a parton between the branchings is  $\Delta t = E/Q^2$ .

### Caveats

- No hadronization model plugged in yet.
- No jet reconstruction algorithm plugged in yet.

⇒ **I do not draw any comparison to experimental data.**

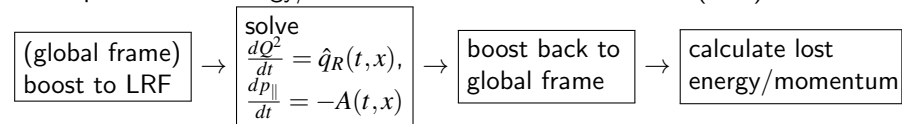
## Jet-medium interaction

- Fluid and jet evolutions run in parallel:

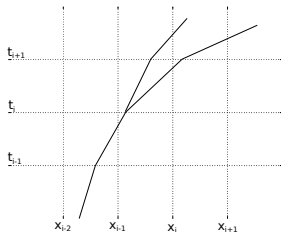


- The temperature and flow velocity are taken from the hydrodynamic evolution

- Jet partons lose energy/momentum in the local rest frame (LRF) of the fluid:

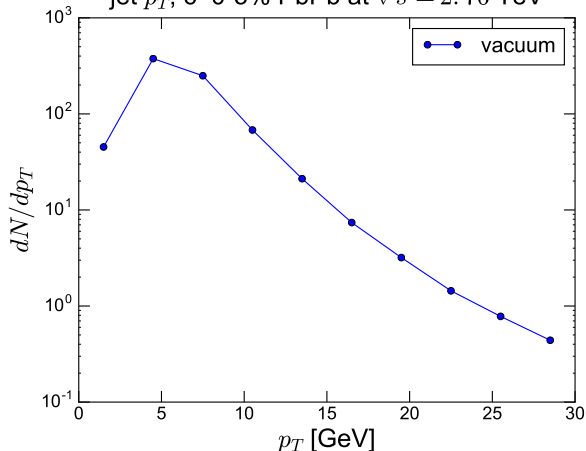


- Once the energy of a parton in the fluid rest frame drops below  $\alpha \cdot T(t, x)$ , the parton is melted into the fluid: its energy/momentum is distributed around nearby fluid cells, and the parton is removed from the parton cascade.
- The fluid acquires the lost energy/momentum (**absorption**) via the source terms:  $\partial_{; \nu} T^{\mu \nu} = J^{\mu}$



## Results: jet $p_T$

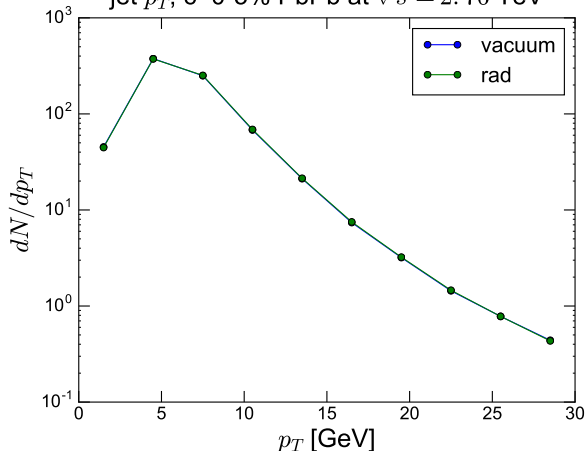
jet  $p_T$ , c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV



● vacuum

## Results: jet $p_T$

jet  $p_T$ , c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV

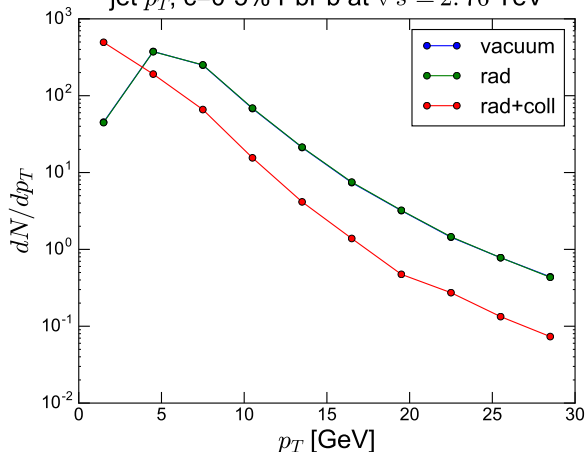


- vacuum
- radiative



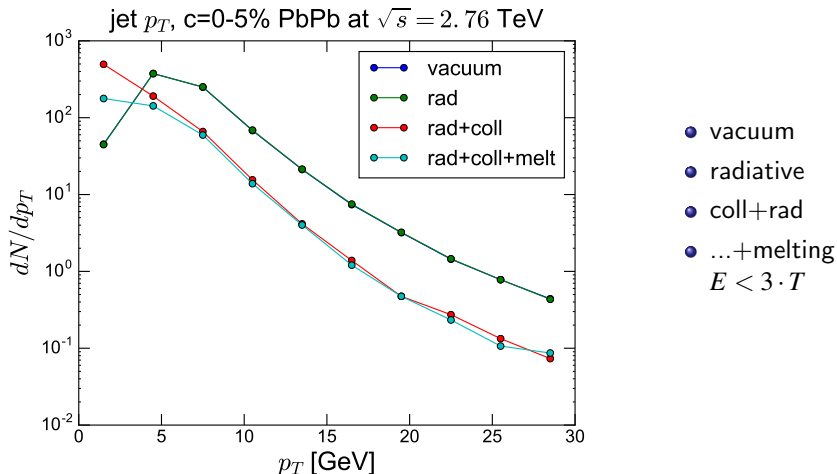
## Results: jet $p_T$

jet  $p_T$ , c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV

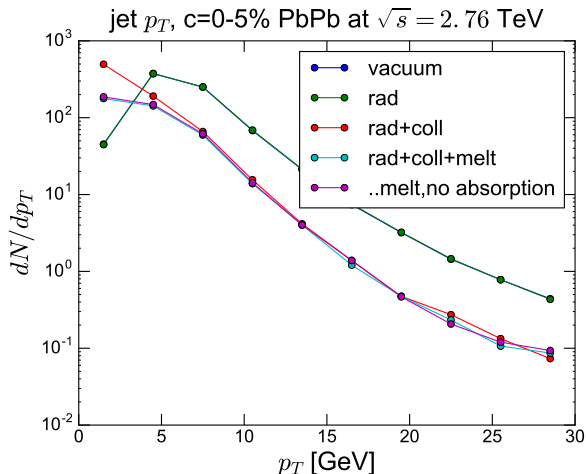


- vacuum
- radiative
- coll+rad

# Results: jet $p_T$

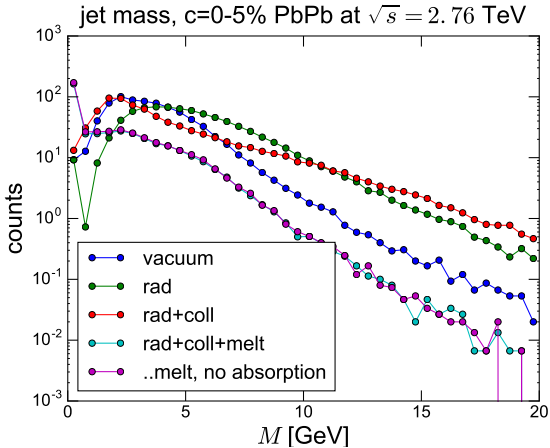


# Results: jet $p_T$



- vacuum
- radiative
- coll+rad
- ...+melting  
 $E < 3 \cdot T$
- ...+no absorption  
by the fluid

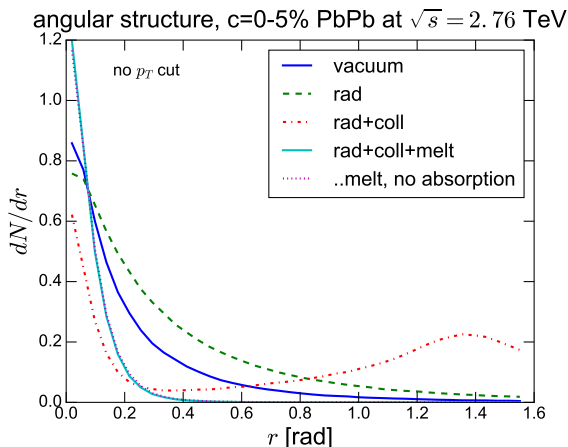
# Jet mass



$$M^2 = \left( \sum_{i \in jet} p_i^\mu \right)^2$$

- Jet mass distribution at high  $M$  is very sensitive to the melting scenario.
- Turning the absorption on/off does not make difference for the jet itself (which is naively expectable).

## Angular structure of a jet



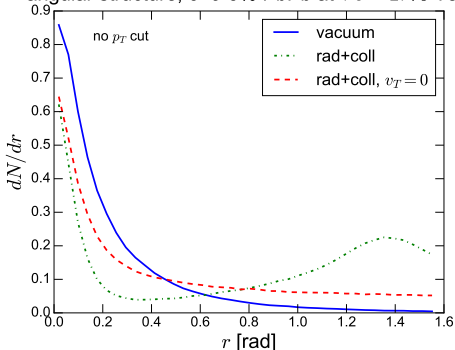
- Radiative EL  $\rightarrow$  broadening (more secondary splittings)
- Collisional EL  $\rightarrow$  1) grooming at small  $\rho$  and 2) a wide peak around  $\approx \pi/2$
- Parton melting kills the peak around  $\approx \pi/2$
- Switching the absorption off does not influence the jet structure

## Effects of radial flow

The structure around  $\rho \approx \pi/2$  is an effect of radial flow.

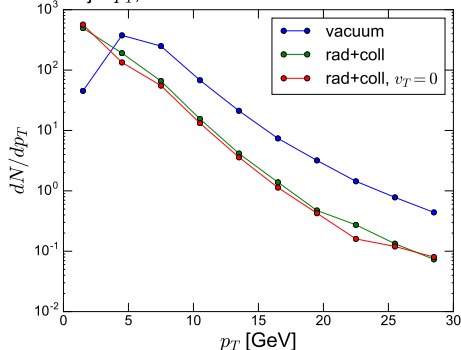
### Effect for $dN/dr$

angular structure, c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV



### No effect for jet $p_T$

jet  $p_T$ , c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV

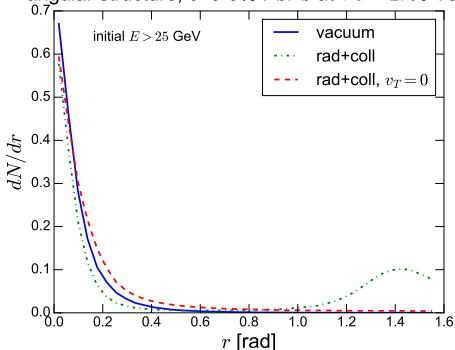


- Switching the transverse expansion off kills the peak in the  $\rho$  distribution.
- Influence on the  $p_T$  distribution is tiny: the main effect of radial flow (switched off here) is faster system cooling, and so smaller energy loss.

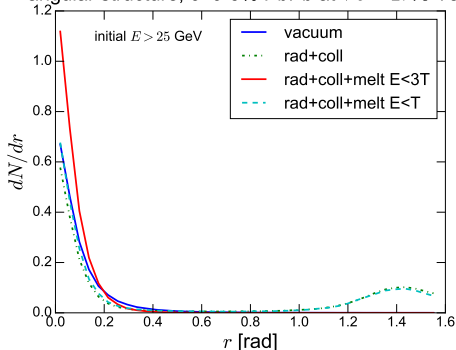
## Other effects

- larger initial jet energy selection
- different parton melting criterion

angular structure, c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV



angular structure, c=0-5% PbPb at  $\sqrt{s} = 2.76$  TeV

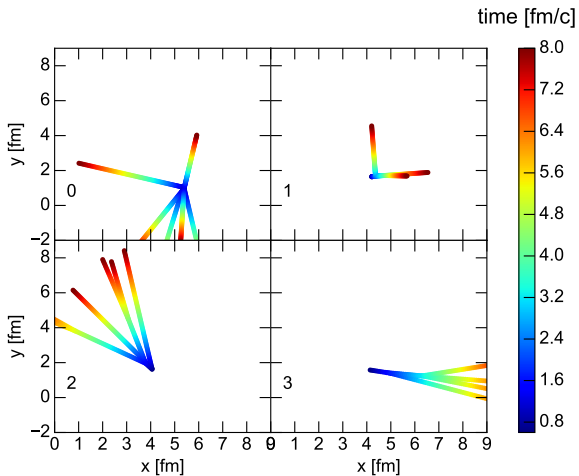


- cut on initial  $E_{\text{jet}}$ : less differences at small  $r$ , the large  $r$  peak survives.
- Parton melting at  $E < 3T$  destroys the peak, however milder melting criterion  $E < T$  preserves it.

## Effects of radial flow (2)

Space trajectories of 4 randomly chosen jets.

No medium interaction:

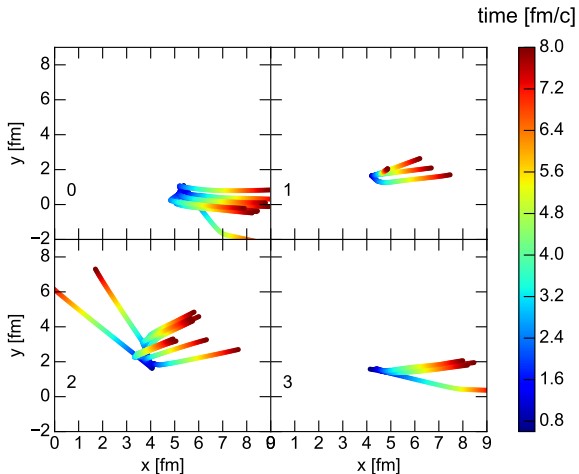




## Effects of radial flow (2)

Space trajectories of 4 randomly chosen jets.

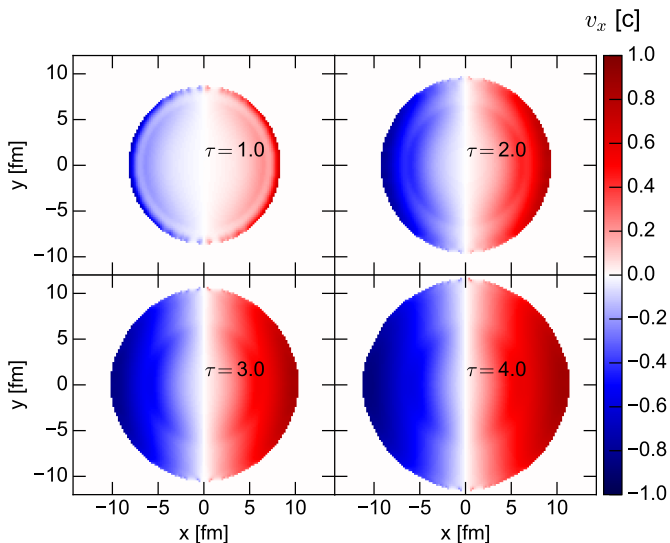
radiative+collisional energy loss, no melting:



## Back reaction on the fluid

Snapshots of the  $x$  component of fluid velocity at different times.

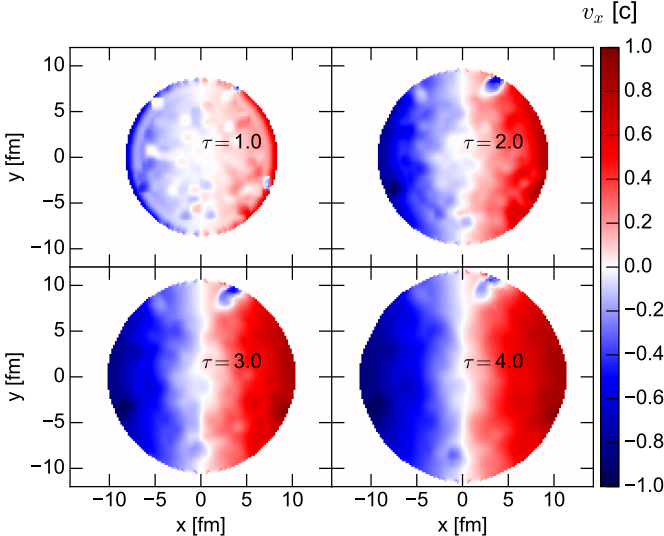
On this slide: **fluid with no absorption (benchmark/to guide the eye).**



# Back reaction on the fluid

Snapshots of the  $x$  component of fluid velocity at different times.

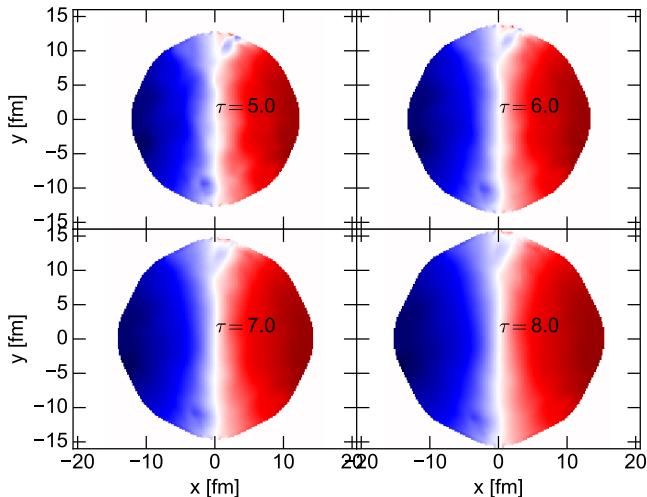
On this slide: **jet energy/momentum absorption at early times.**



## Back reaction on the fluid

Snapshots of the  $x$  component of fluid velocity at different times.

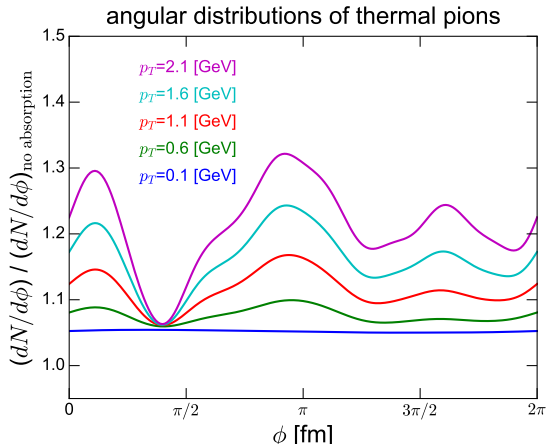
On this slide: **Hydro smears out the perturbations at late times.**



## Corresponding azimuthal distributions of hydro-born hadrons (pions)

Cooper-Frye prescription at  $\varepsilon = \varepsilon_{sw} = 0.5 \text{ GeV}/\text{fm}^3$ :

$$p^0 \frac{d^3 n_i}{d^3 p} = \sum f(x, p) p^\mu \Delta \sigma_\mu, \quad f(x, p) = f_{\text{eq}} \cdot \left( 1 + (1 \mp f_{\text{eq}}) \frac{p_\mu p_\nu \pi^{\mu\nu}}{2T^2(\varepsilon + p)} \right)$$

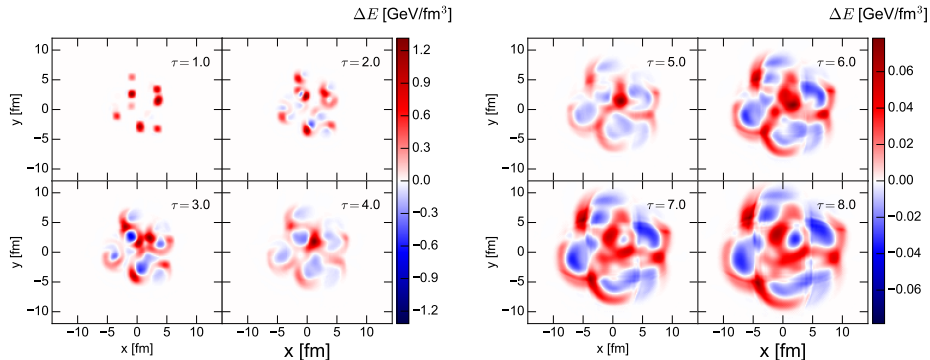


In hydrodynamics:

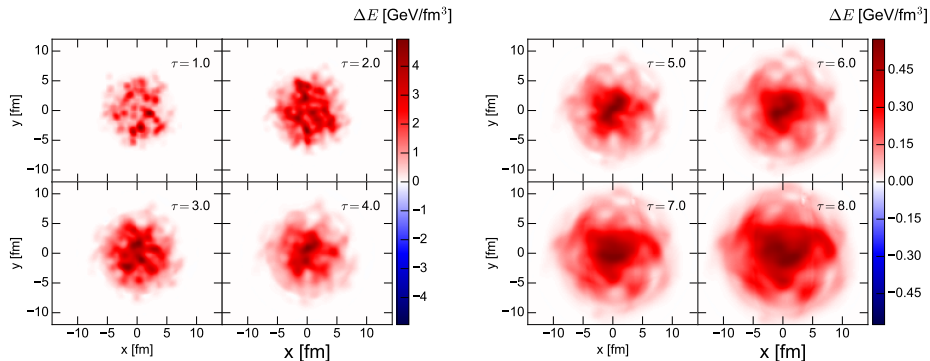
- low- $p_T$  hadrons are dominantly produced at late times
- high- $p_T$  hadrons are dominantly produced at early times from the periphery  
↓
- high- $p_T$  hadrons are sensitive to the perturbations in the medium

# Mach cones!

A simplified setup: few most energetic jet partons with initial  $Q > 20$  GeV.

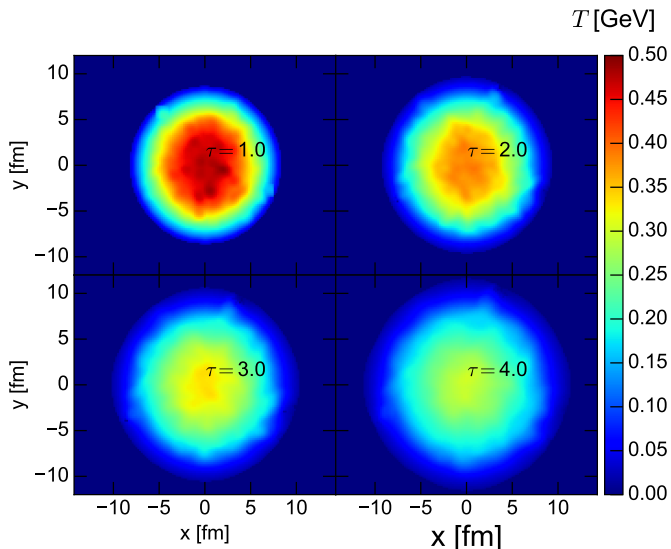


A more realistic case at LHC energy: jet partons with initial  $Q > 2$  GeV.



## Back reaction on the fluid (2)

Corresponding temperature profiles at different times:





## Summary

We have presented a first calculation where jets and bulk hydrodynamic evolution run in parallel mode.

- Initial conditions and initial jet partons: EPOS
- Timelike parton cascade by Martin Rohrmoser
- 3+1 dimensional viscous hydrodynamics for the medium
- Bi-directional interaction between the two

Some lessons:

- Radial flow causes a spread of a fraction of jet energy to a relatively large angle  $\approx \pi/2$ .
- Back reaction of the jet energy loss to the fluid has a negligible effect for the **jet itself** (no jet recoil here!).
- However, the energy absorption causes perturbations in the hydro evolution which are strongest at early times and therefore influence thermal hadron production at “large thermal” momenta  $> 1 - 1.5$  GeV.
- $\rightarrow$  This influences the correlations of high- $p_T$  jet hadrons with their  $p_T > 1.5$  GeV “thermal” colleagues.

Obviously, work in progress.