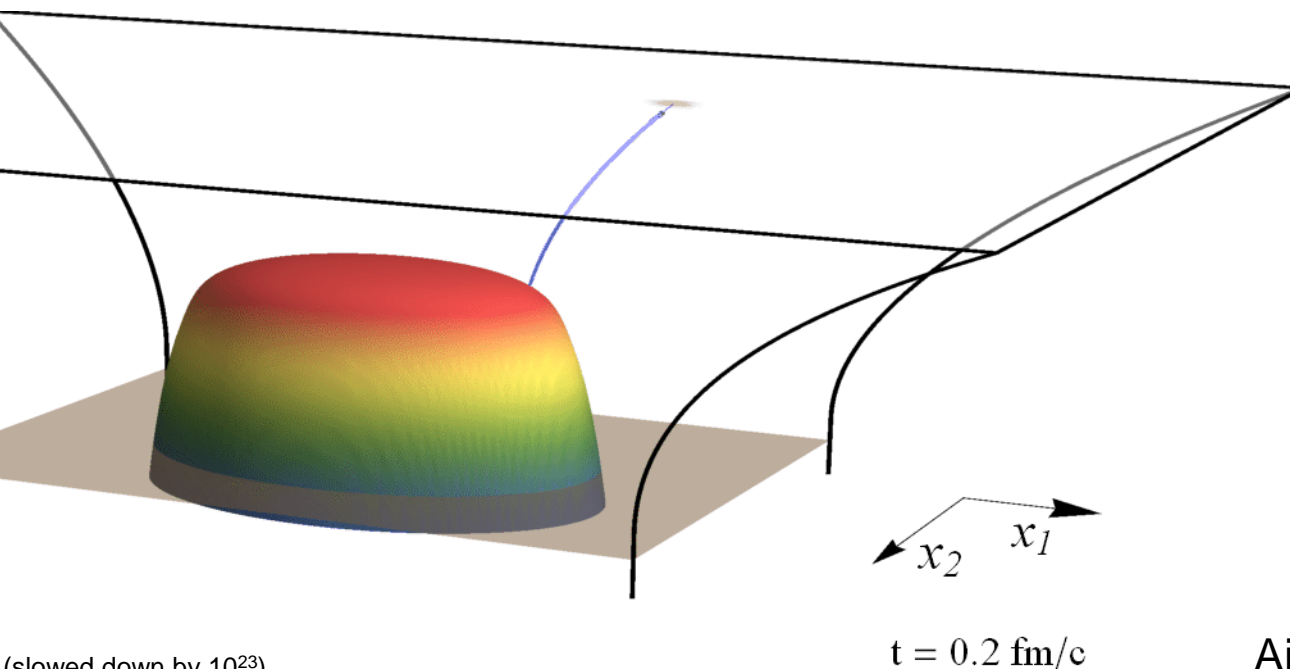


JET ENERGY LOSS IN A FLOWING PLASMA

JET SHAPE MODIFICATIONS IN HOLOGRAPHIC DIJET SYSTEMS

With Jasmine Brewer, Krishna Rajagopal and Andrey Sadofyev
1602.04187 (PRL), 1710.03237 (JHEP) and 1809.10695



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OUTLINE

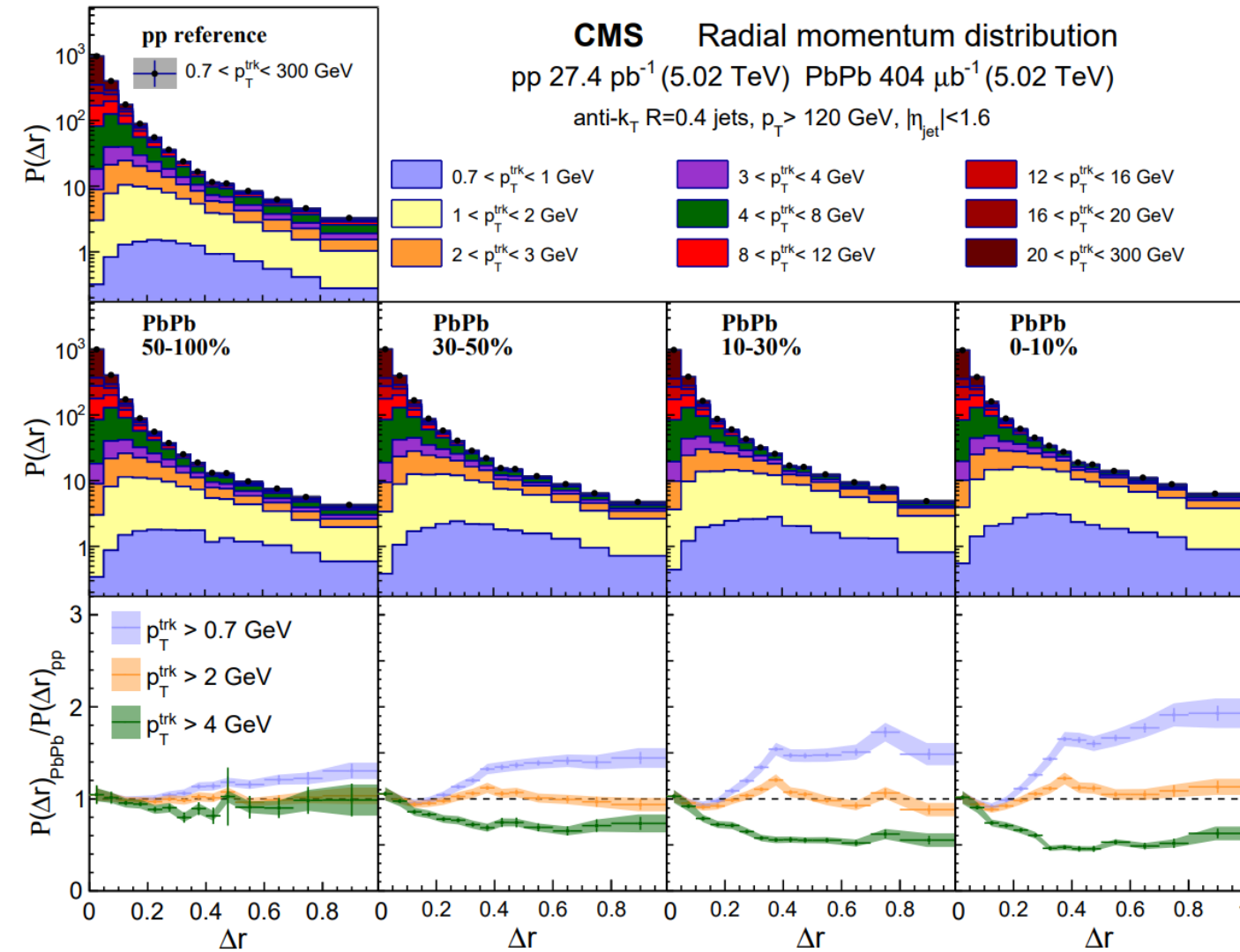
Motivation and a simple model

- (weakly coupled) jet energy interacts strongly with QGP
- From production (pQCD) to strings to null geodesics in AdS
- Goal is to provide energy loss formalism valid for flowing plasma

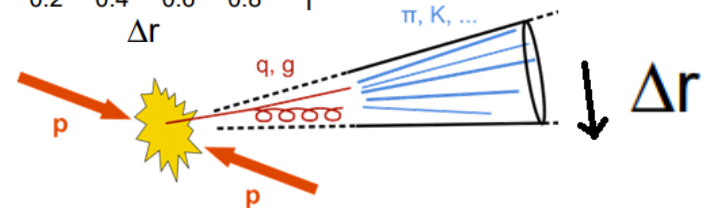
Jet shapes and (holographic) dijet systems

- Dijet asymmetry interesting: path length & fluctuations essential
- Study leading and subleading shape modifications
- Observable: dijet asymmetry for different subleading jet width

TRENDS IN JET ANALYSIS

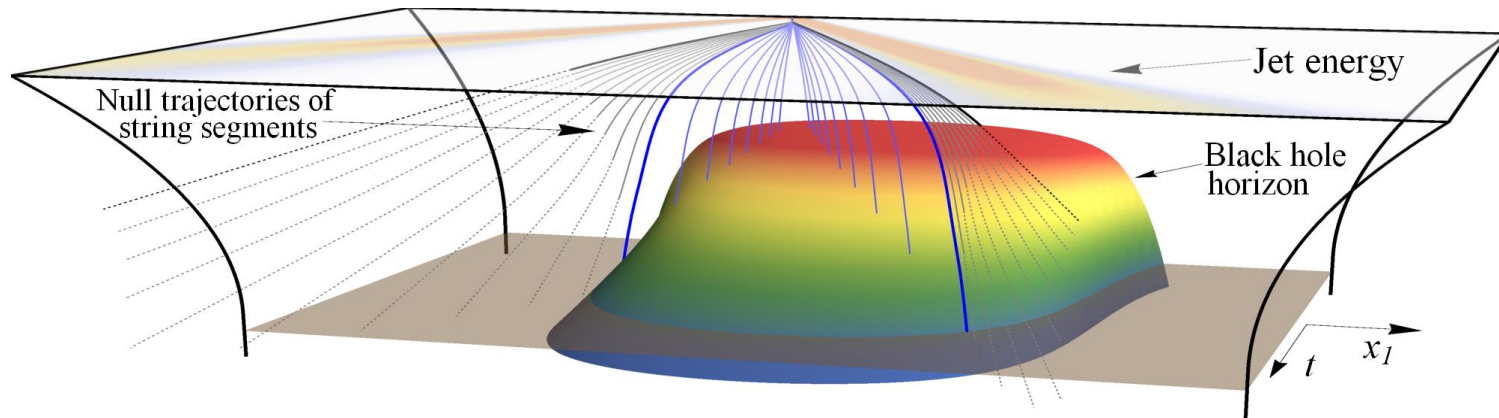


p_T cut separates
medium effects
from jet physics



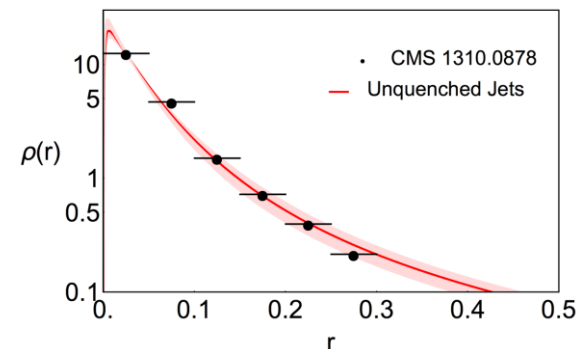
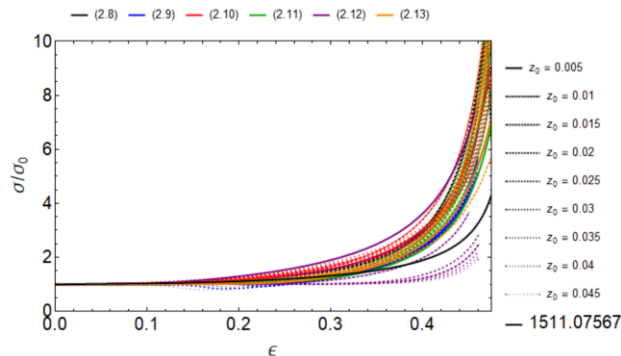
JET ENERGY LOSS IN ADS/CFT

Quark-antiquark pairs \sim strings in AdS geometry



Leads to (simplified) model for jet evolution

- String segments quickly follow null geodesics in semi-universal way
- Possible to track null geodesics falling in: determines energy loss
- Final shape depends on specific evolution



FLOW AND IDEAL HYDRODYNAMICS

Final formula for geodesics in ideal hydrodynamics, including flow:

$$Z''(t) = -2\pi^4 T^4 (1 - v_{||}) Z(t)^3 \left(\frac{1 - v_{||}}{v_{||}^2 + v_{\perp}^2 - 1} \right)$$

- Z is holographic coordinate string segment, T and v depend on (τ, x_{\perp})
- Initial position close to zero, initial Z' related to jet width

Flow had not yet been included

- Formula can have analytic solutions for constant temperature
- Varying temperature is non-trivial (even when gradients are small)
- Effects of gradients is relatively small (upcoming work)
- (full formula including Z' and viscous terms is known but longer)

$$ds^2 = -2 u_{\mu} dx^{\mu} dr - r^2 f(b r) u_{\mu} u_{\nu} dx^{\mu} dx^{\nu} + r^2 P_{\mu\nu} dx^{\mu} dx^{\nu} \\ + 2 r^2 b F(b r) \sigma_{\mu\nu} dx^{\mu} dx^{\nu} + \frac{2}{3} r u_{\mu} u_{\nu} \partial_{\lambda} u^{\lambda} dx^{\mu} dx^{\nu} - r u^{\lambda} \partial_{\lambda} (u_{\nu} u_{\mu}) dx^{\mu} dx^{\nu}$$

A SIMPLE ALGORITHM

Energy loss from AdS/CFT in a dynamic setting

- Start with several string segments at boundary (~ 20), with different Z'
 - Z' of endpoint is determined by pQCD opening angle of q/g
 - Z' of other segments is taken from semi-universal curve (slide 4)
- Evolve $Z(t)$ according to simple differential equation
- Straightforward to determine energy outside horizon

Main difference with current dE/dx approaches:

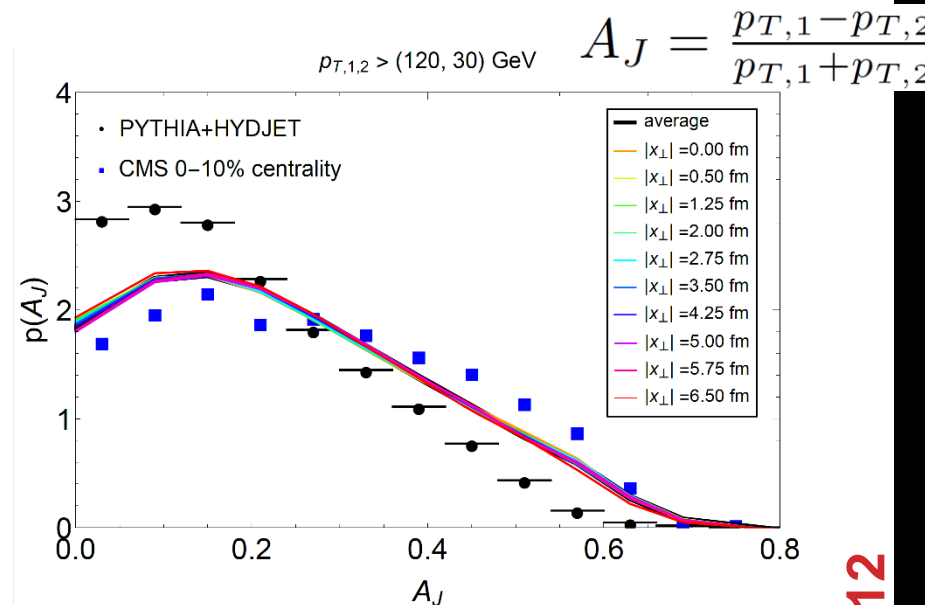
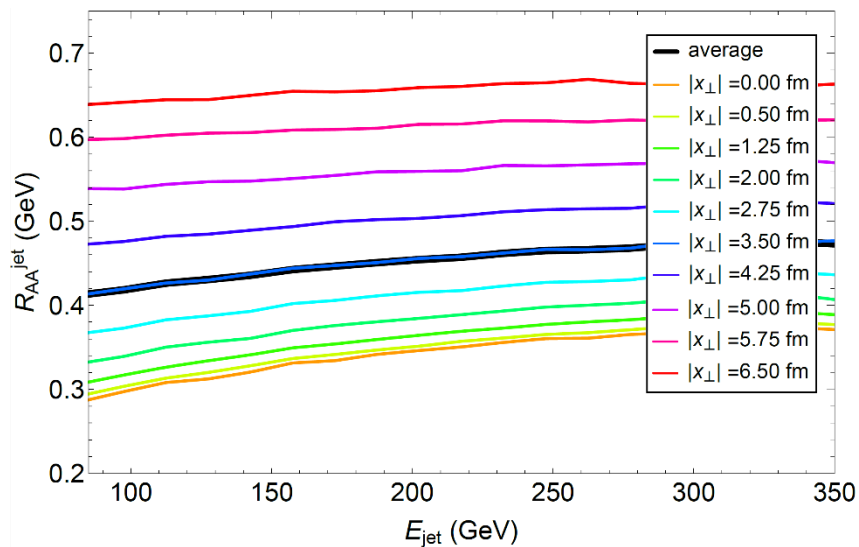
- Need to keep track of ~ 20 variables per parton, i.e. parton wave function more complicated than just energy
- 2nd order equation: memory effect
 - Perhaps similar to L^2 or L^3 scaling of current approaches
- Relatively non-linear interplay of $E(x)$ versus $T(t)$ and $v(t)$

DIJET ASYMMETRY

Two sources of (extra) imbalance in heavy ion dijet asymmetry:

- Path length: different path length \rightarrow less balanced
- Fluctuations of energy loss: jet width fluctuations (our model)
- Nice trick to turn off path length imbalance: start all jets at center
 - Turns out that in our model (and in JEWEL) A_J is insensitive to position

Nuclear modification factor

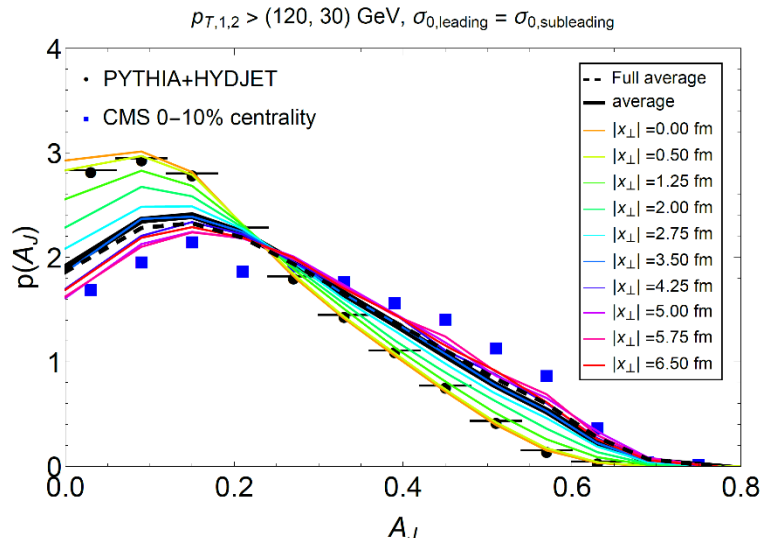


DIJET ASYMMETRY

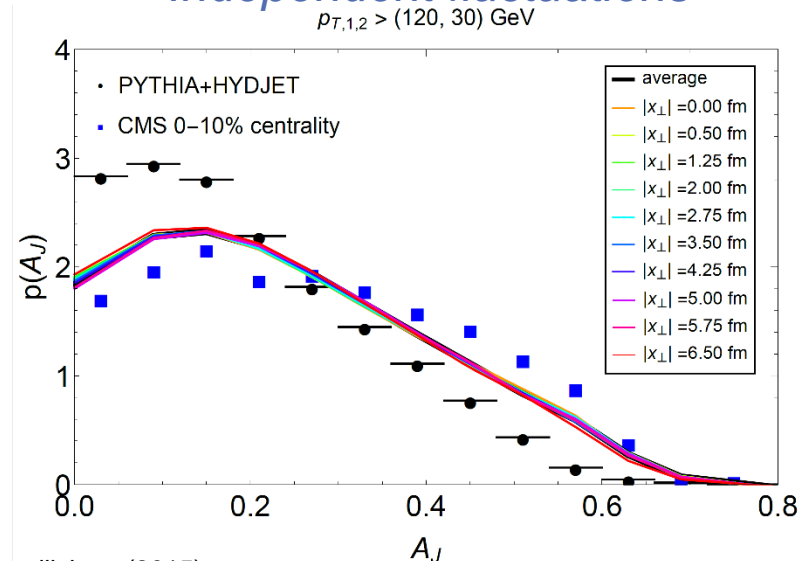
(Artificial) way of turning off independent energy loss fluctuations

- Select subset of dijets in sample that have equal jet widths
- Relative energy loss fully determined by path length imbalance
- Now indeed dijet asymmetry of jets produced at center is unmodified
 - Curiously average dijet asymmetry agrees with full result

No independent fluctuations



Independent fluctuations

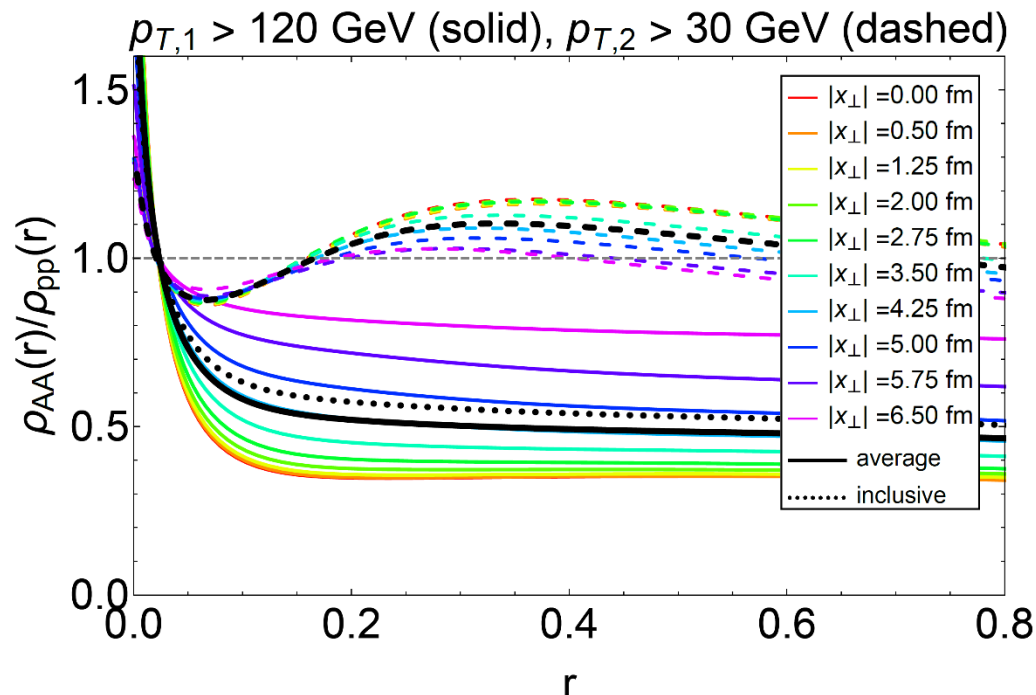


JET SHAPE MODIFICATIONS

Two clear lessons:

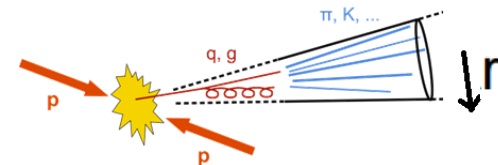
- Influence of path length imbalance and fluctuations 'cancels'
- Dijet asymmetry is uninformative on where a jet was produced
- Jets from center lose more energy, hence larger fluctuations

Jets at center are also more heavily modified:



Leading and subleading (dashed) modifications

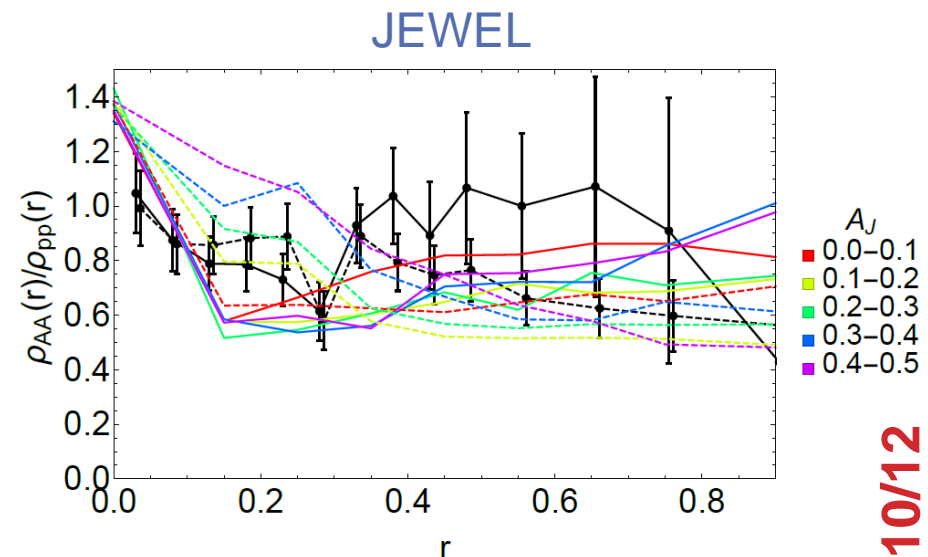
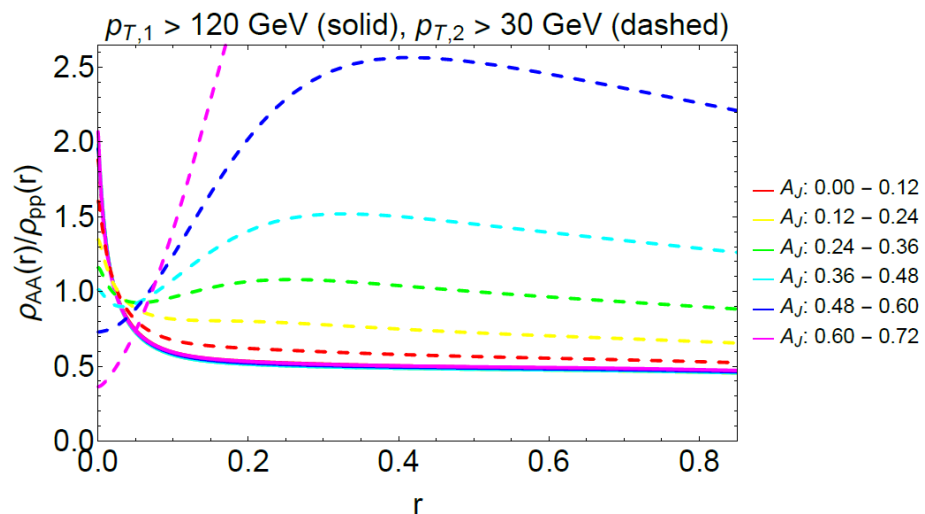
- Jets from center more strongly modified
- Leading jet much narrower
- Subleading jets seem wider (but subtle in our model)



JET SHAPE MODIFICATIONS

Can these differences be probed experimentally?

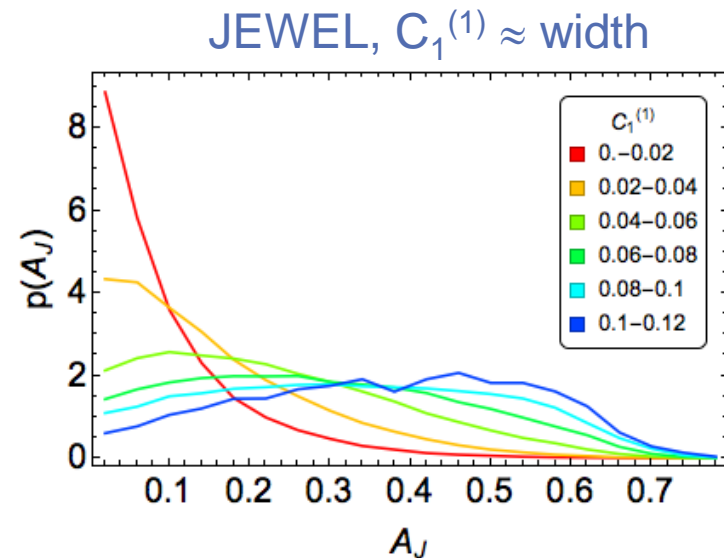
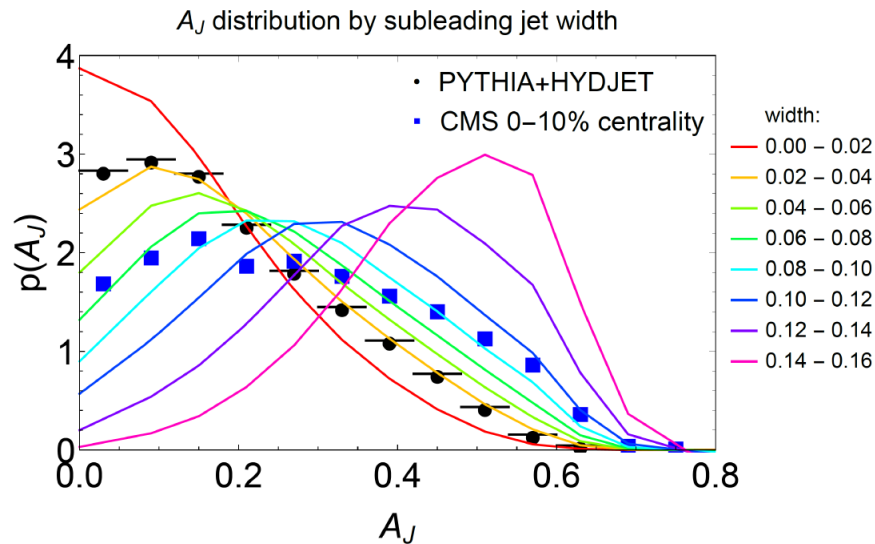
- *Select jet shapes according to dijet imbalance*
- Compare with JEWEL, matches qualitatively
- Includes CMS data with $p_T > 3$ GeV, to exclude thermal particles
- Subtlety for subleading: rather large contribution from 3rd jets and incoherent partons (included in JEWEL, not in our model)



A NEW OBSERVABLE?

Dijet asymmetry is very sensitive to subleading jet width

- Large subleading width leads to imbalanced dijet
- JEWEL effect is a bit stronger for small width, weaker for large width
- Quite intuitive: large subleading width \rightarrow large energy loss \rightarrow large A_J



DISCUSSION

The model

- Take initial conditions of strings from pQCD
- Treat evolution in holography: medium interaction is strongly coupled
- Currently oversimplified: no hadronisation, back-to-back jet, medium etc

Jet shape modifications and dijet asymmetry

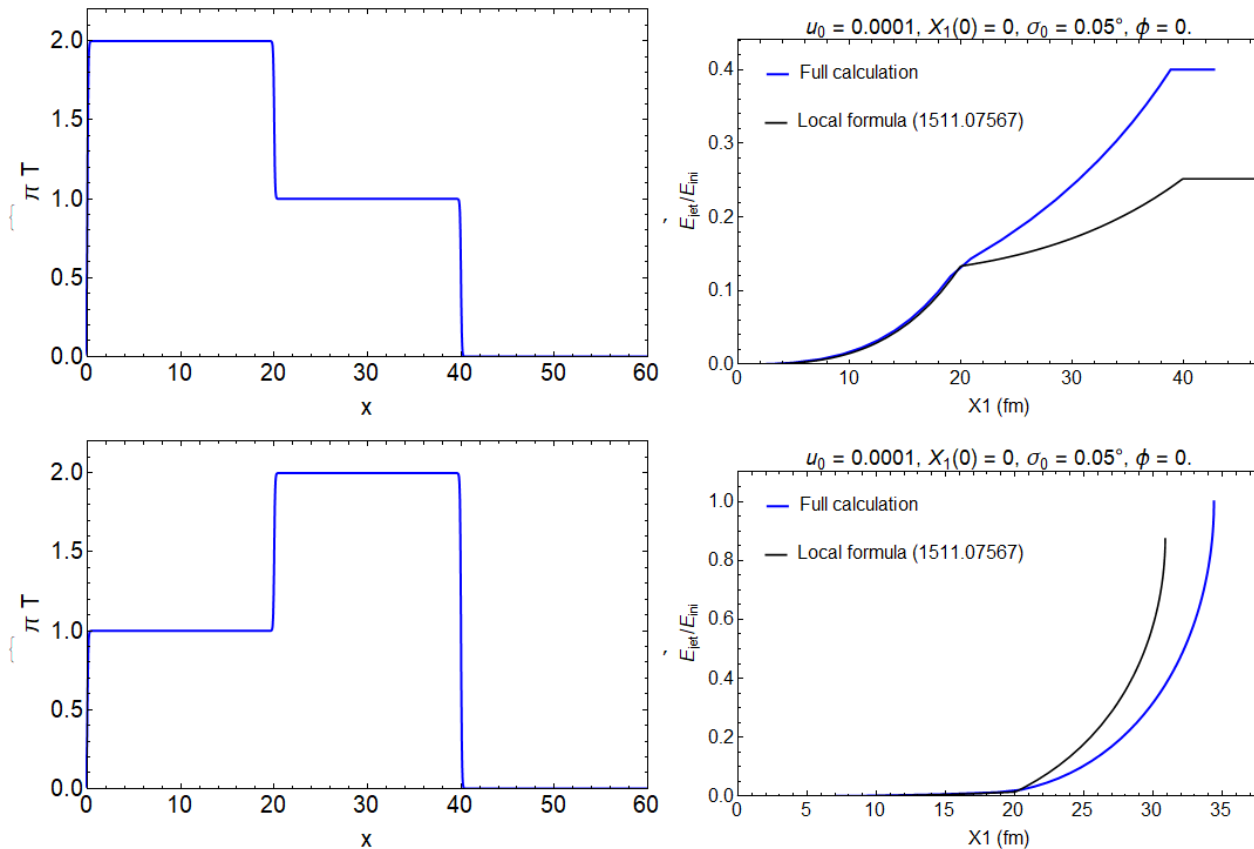
- Confirmed dijet asymmetry does not depend on starting position
- Jet shape modifications strongly depend on position
- Leading jet shape narrower, independent of A_J
- Subleading jet shape increasingly wider for increasing A_J
- Observable: dijet asymmetry for different subleading jet width?

Outlook

- Implementation in Monte Carlo?
- Full understanding initial condition string: 3-jets etc
 - Can have important effects on subleading jet at large r
- Not quite related: back-reaction lost energy on medium: where does E/p go?

VARYING TEMPERATURE

Energy loss depends on temperature evolution non-linearly:

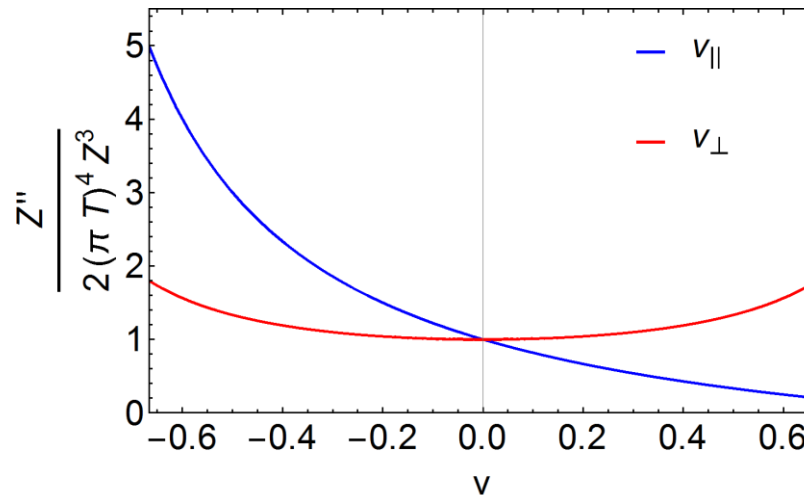


- First phase agrees with (ultra)local formula (Chesler, Rajagopal)
- Interesting: (final) energy loss much bigger for 2nd profile
- Illustrates 'memory-effect': wave function remembers evolution

SECOND EFFECT: FLOW

Final formula for geodesics in ideal hydrodynamics:

$$Z''(t) = -2\pi^4 T^4 (1 - v_{||}) Z(t)^3 \left(\frac{1 - v_{||}}{v_{||}^2 + v_{\perp}^2 - 1} \right)$$



(Z'' essential for energy loss, but not proportional to energy loss)

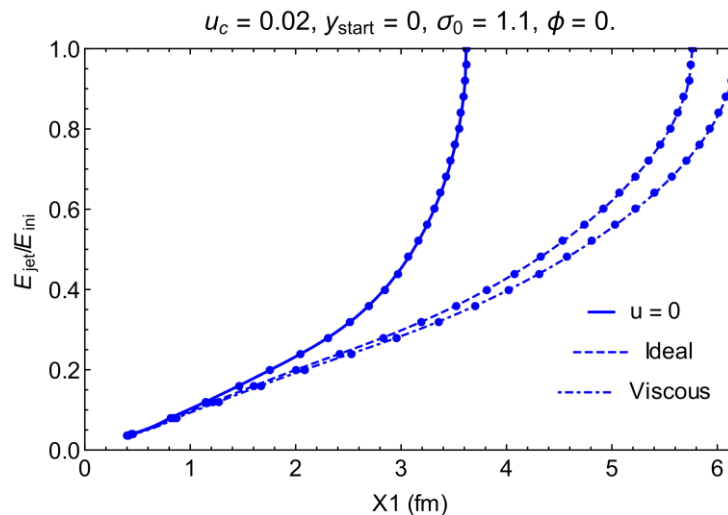
Corrections due to viscosity:

$$\Delta Z''(t) = -2\pi^4 T^3 \left(\frac{2}{3} \frac{\partial v_x}{\partial x} + 2 \left(\frac{\partial v_x}{\partial x} + 2 \frac{\partial v_x}{\partial t} \right) T Z(t) \right) Z(t)^3$$

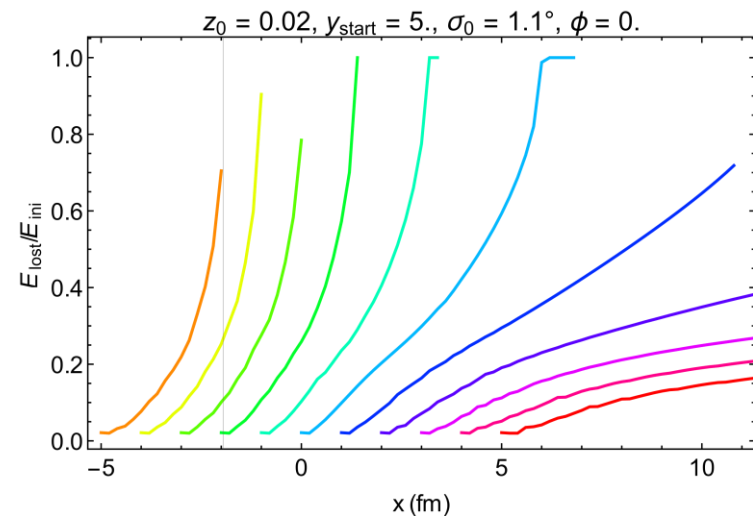
RESULTS IN GUBSER FLOW

Resulting energy loss in (analytic) simple model for central collision

Compare no flow, ideal and viscous



different starting points



- Flow has extremely important effect, doubling stopping distance

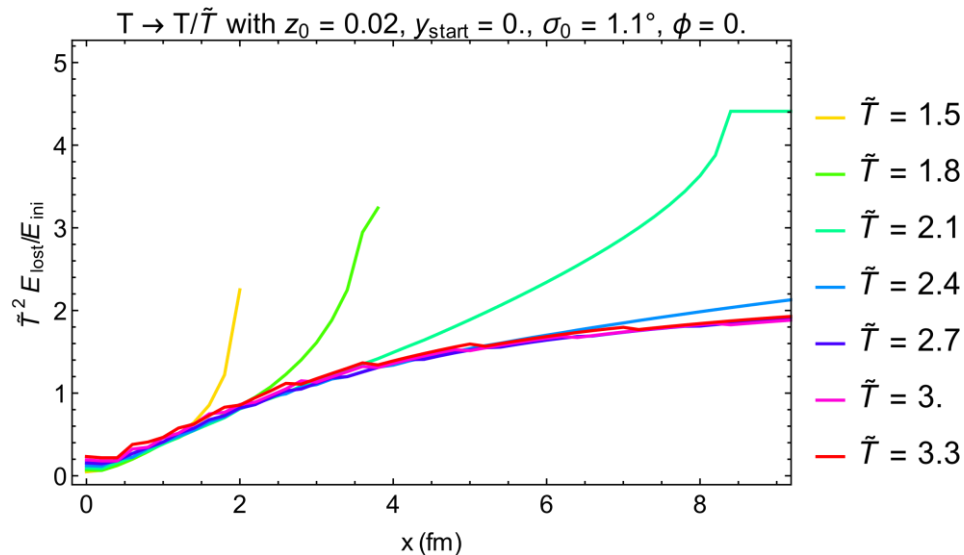
- Recall old result:
$$\Delta x_{\text{stop}} = \left[\frac{2^{1/3} \Gamma(\frac{5}{4})}{\sqrt{\pi} \Gamma(\frac{3}{4})} \right] \frac{1}{T} \left(\frac{E_*}{\sqrt{\lambda} T} \right)^{1/3}$$

- Corrections due to gradients significant, but small

ENERGY LOSS VERSUS TEMPERATURE

Are there semi-general lessons for this energy loss?

- Try extracting dependence dE/dx on temperature: rescale T by \tilde{T}
- Numerical finding: different curves collapse when scaling by T^2
 - Up to point where particular jet loses all energy: early time scaling



- (fairly robust, but scaling somewhat dependent on semi-universal curve, e.g. result from different black line gives T^3)

MODIFIED JET WIDTH DISTRIBUTIONS

Jet shapes become narrower: jet width distribution

- Can also compare with original pp 'truth' jet width
- Even individual jets can become narrower
 - Only for intermediate widths and intermediate energy losses

