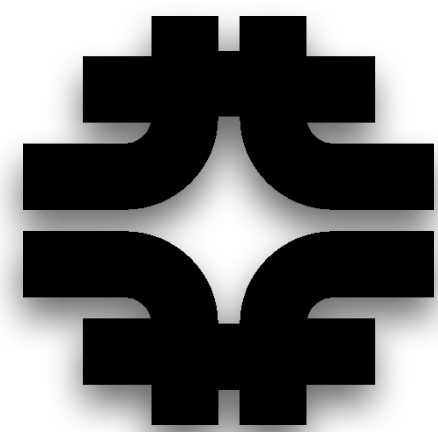


Higgs to bb and Pixel-Phase-II R&D



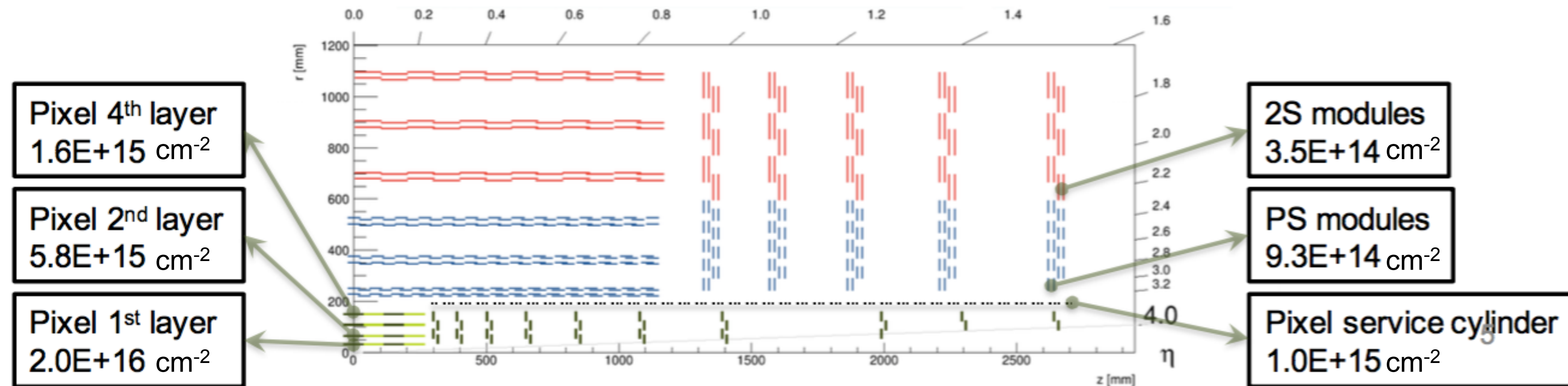
Caterina Vernieri, FNAL and LPC Distinguished Researcher

7.27.2017

Phase-II Pixel upgrade

Ongoing R&D plan to develop a new pixel tracking system to operate at HL-LHC

- To maintain a high tracking and b-jet identification efficiency at luminosities up to $7 \cdot 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ and operate up to 200 $\langle \text{PU} \rangle$
- It will extend the η coverage from the present $\eta = 2.5$ to $\eta = 4$
- Radiation tolerance up to 3000 fb^{-1}
- Different lines of investigation for Phase-II:
 - **Better acceptance per module**/hermeticity \Rightarrow **slim edge sensors**
 - Higher instantaneous luminosity \Rightarrow **high granularity** \Rightarrow **small pitch sensor**
 - Higher integrated luminosity \Rightarrow **high radiation tolerance** \Rightarrow **thinner sensors**

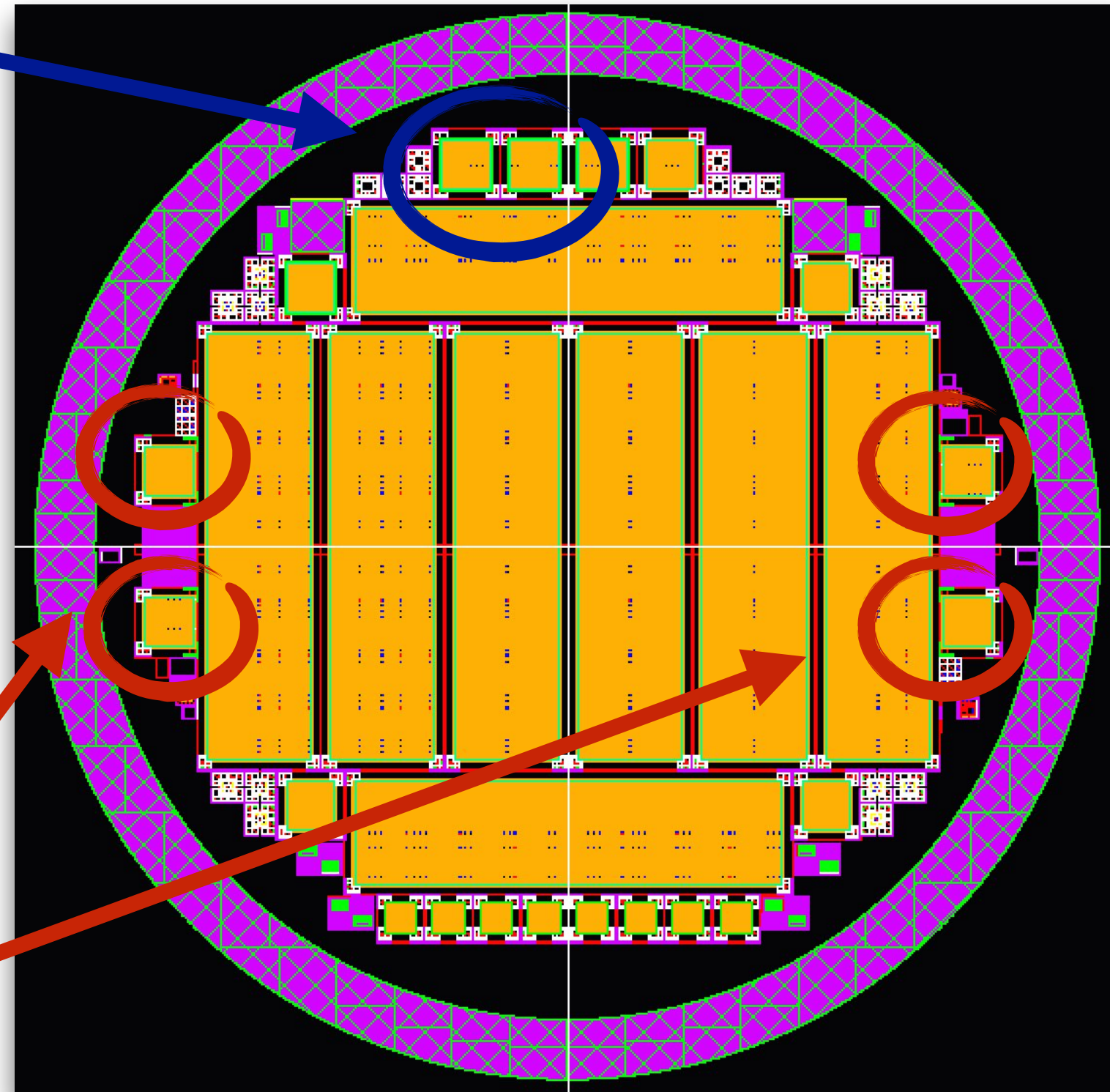


n-in-n

Prototypes were produced along with the pixel sensors of the CMS Forward Pixel detector and tested at Fermilab at the Test Beam Facility

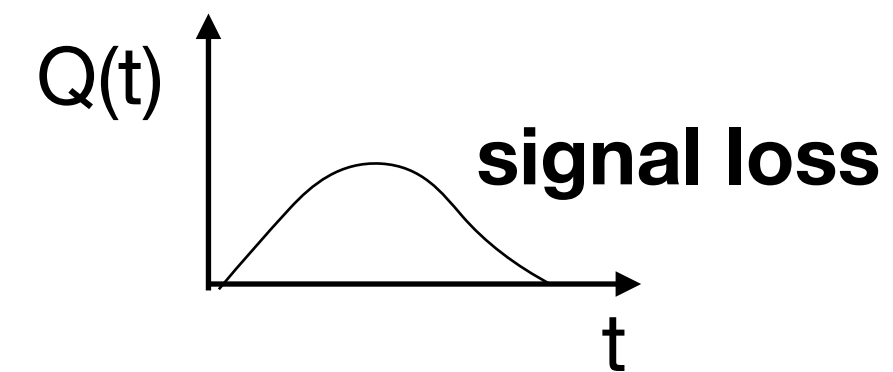
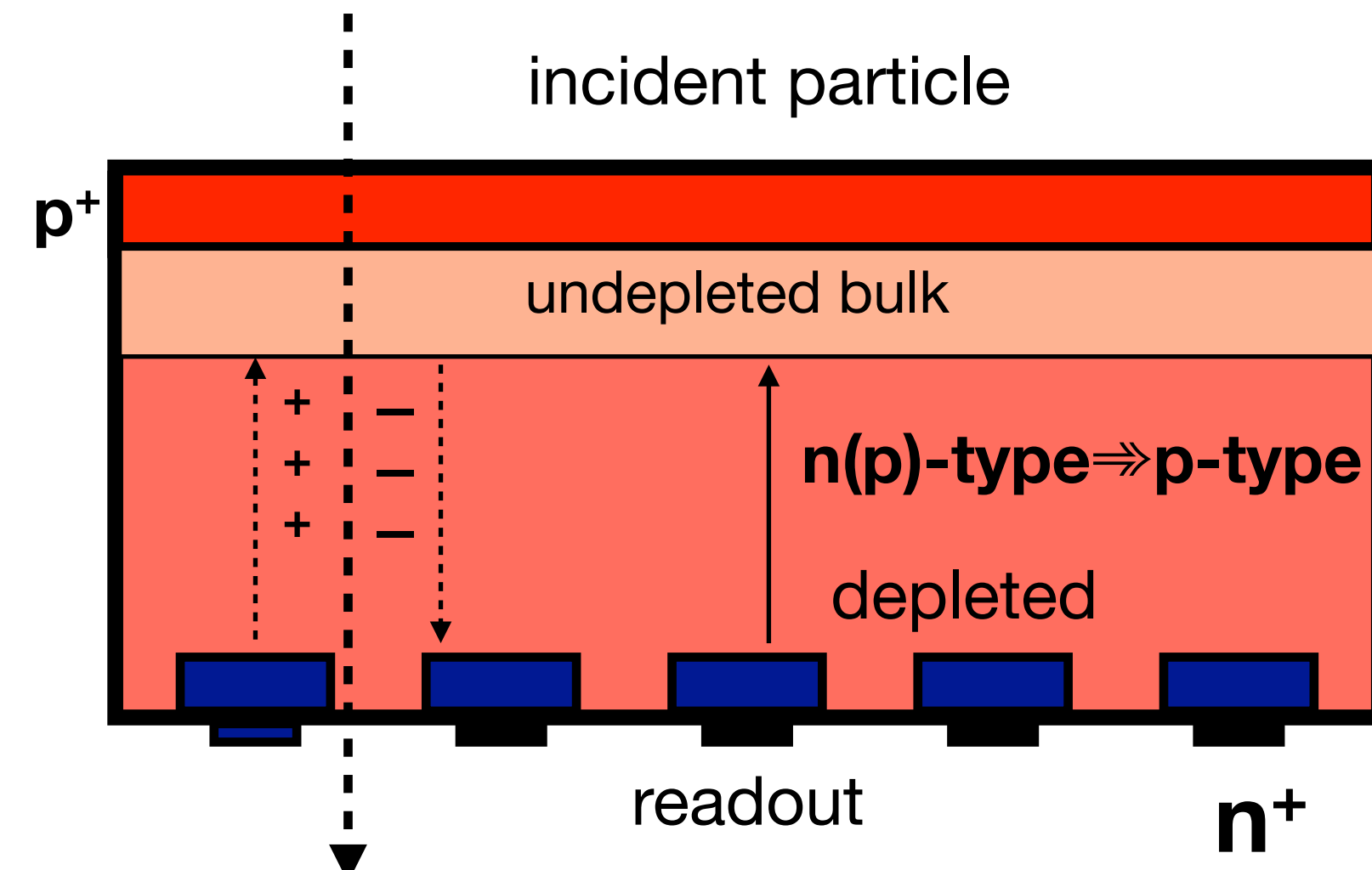
Small pitches

- n-in-n (double sided)
- 300 μm thick
- FZ, 3-5 $\text{k}\Omega\text{cm}$



Slim edge

n-in-p or n-in-n



- faster charge collection (drift for electron is larger)
- CCE degradation

The approach saves money and relies on an established process

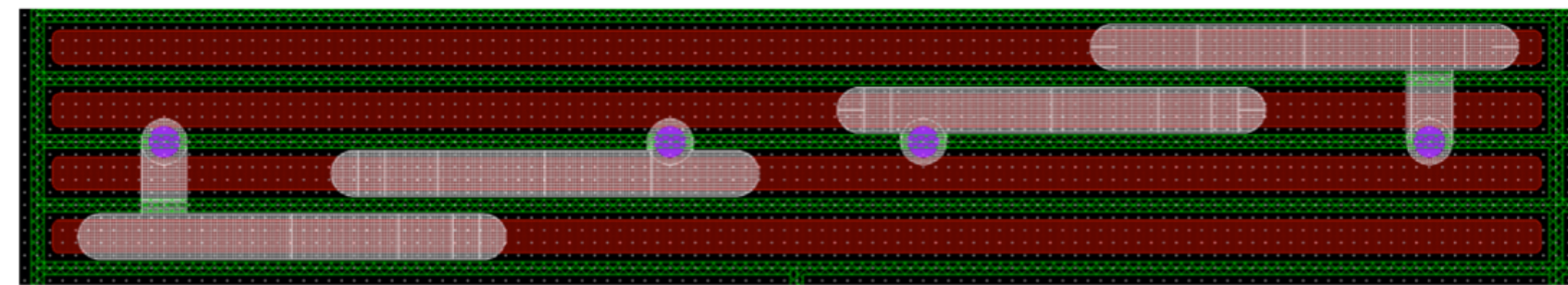
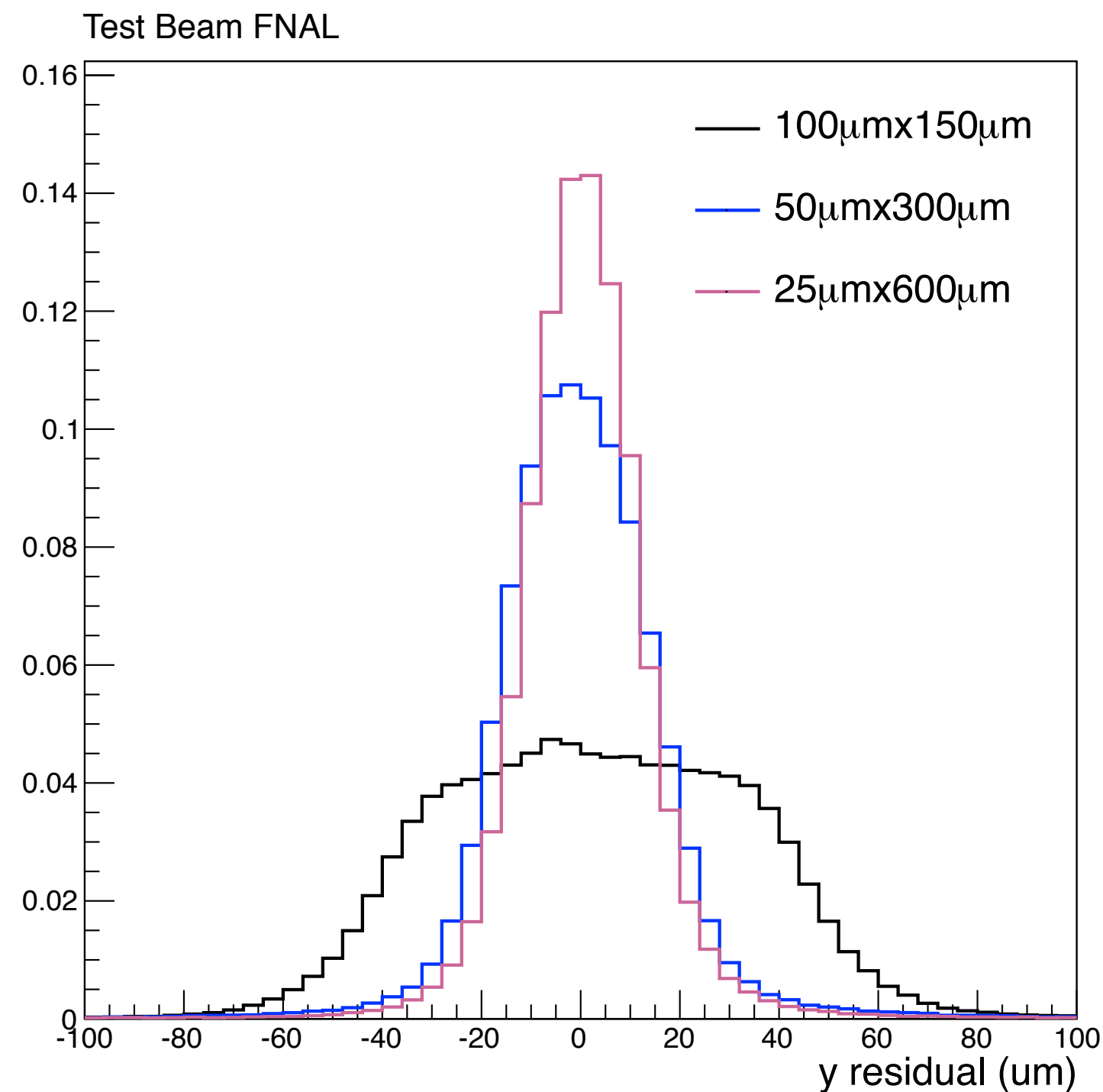
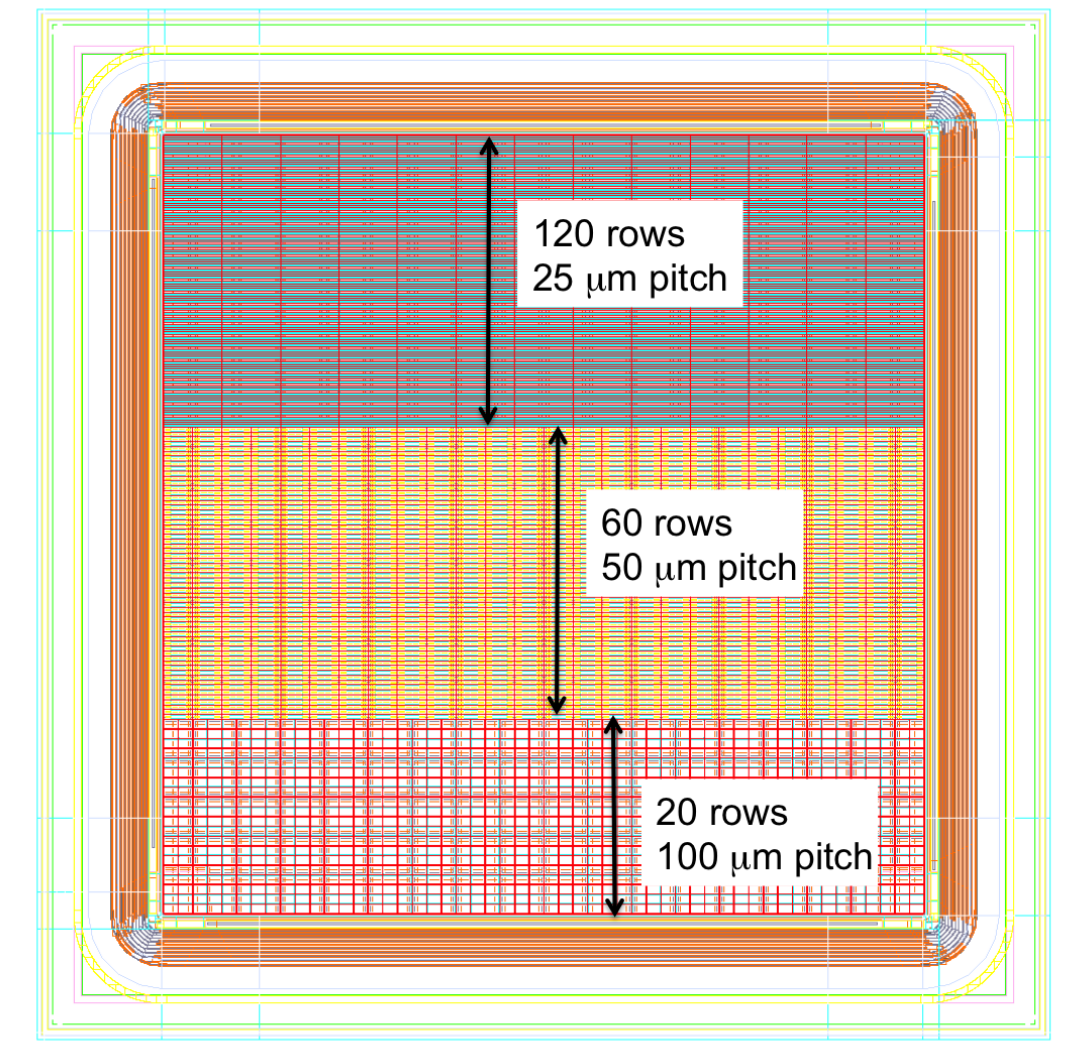
Small pitch sensors

Require bump-bonding patterns compatible to PSI46 ReadOutChip (Phase-I)

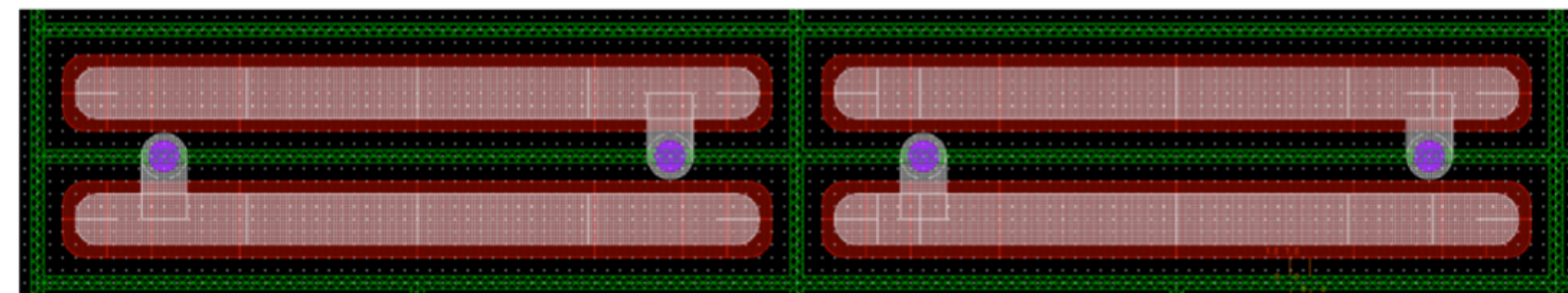
Maintained the same pixel area $100 \times 150 \mu\text{m}^2$ that is implemented in the Phase-I design

Single ROC sensors split in 3 regions with 3 different pitches

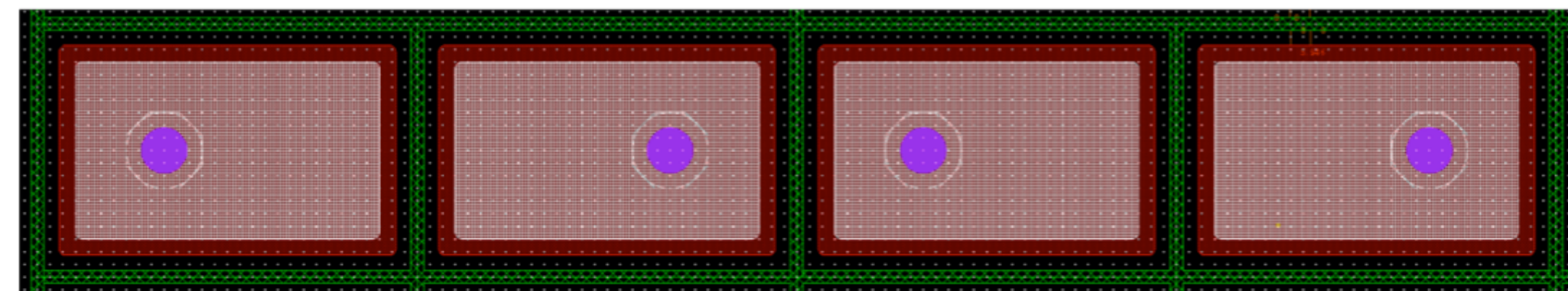
* The hit **resolution is $5.8 \mu\text{m}$** for **$25 \mu\text{m} \times 600 \mu\text{m}$** pixels



n+	→ 13 μm
gap	→ 3.5 μm
p+	→ 5 μm
gap	→ 3.5 μm
Pitch	→ 25 μm



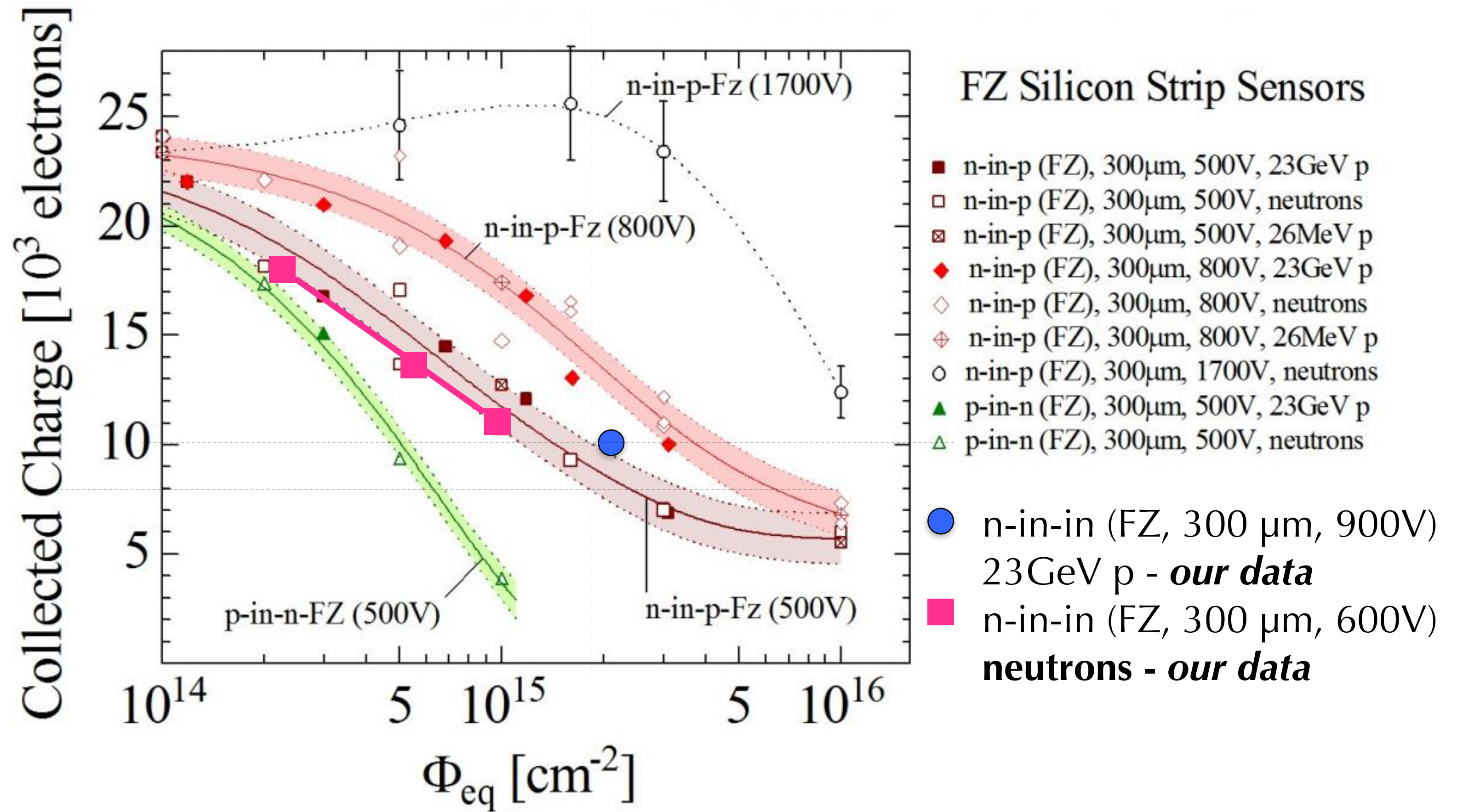
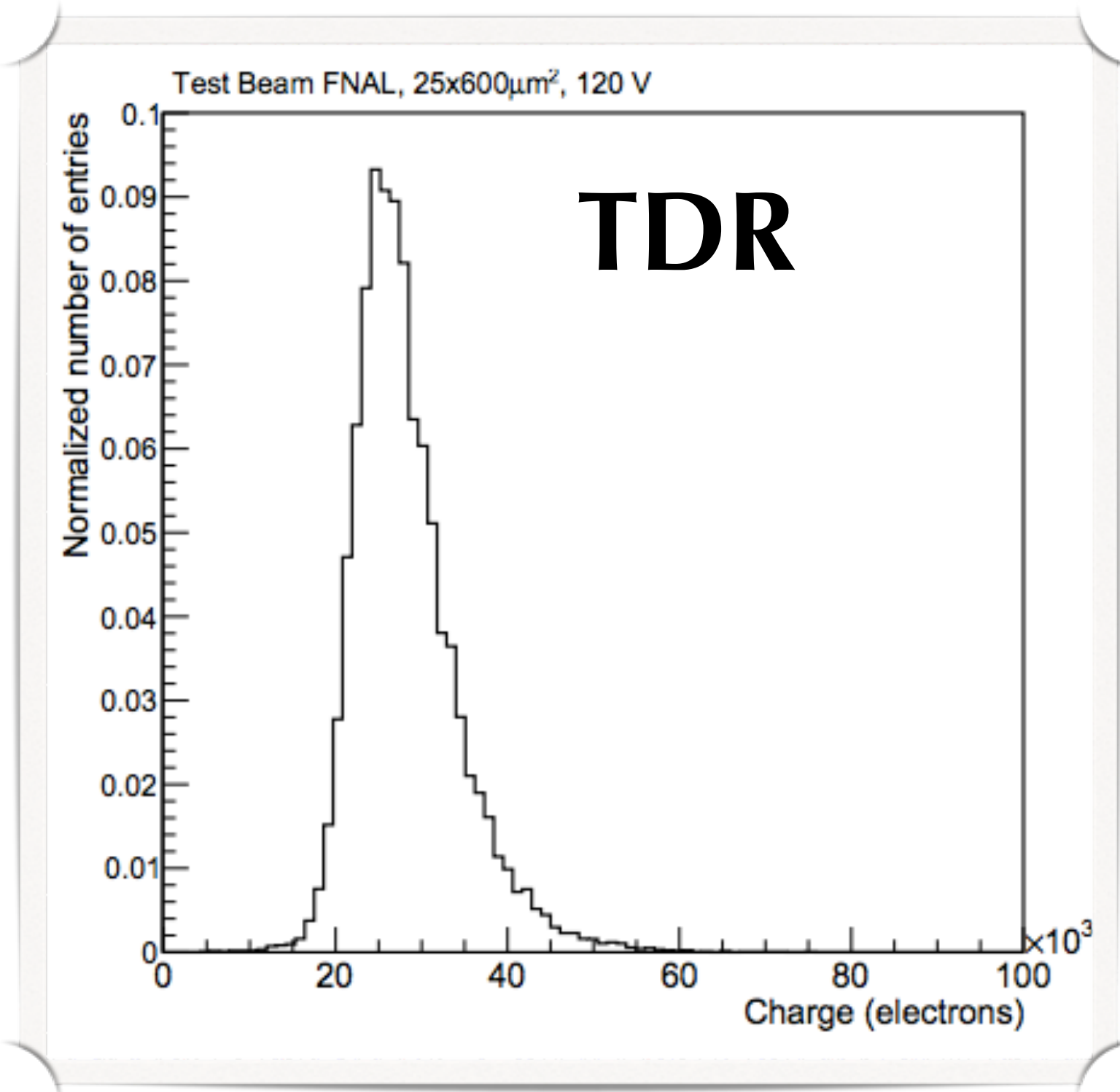
n+	→ 30 μm
gap	→ 7.5 μm
p+	→ 5 μm
gap	→ 7.5 μm
total	→ 50 μm



n+	→ 83 μm
gap	→ 6 μm
p+	→ 5 μm
gap	→ 6 μm
total	→ 100 μm

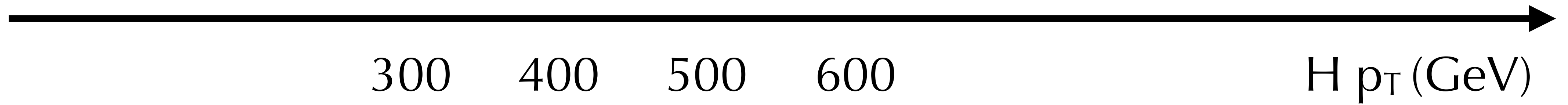
Collected charge vs Fluence

<http://ph-news.web.cern.ch/content/rd50-radiation-tolerant-silicon-detectors-1>



Results have been published in **NIM 2016 06, 020** and included in the Tracker **TDR**

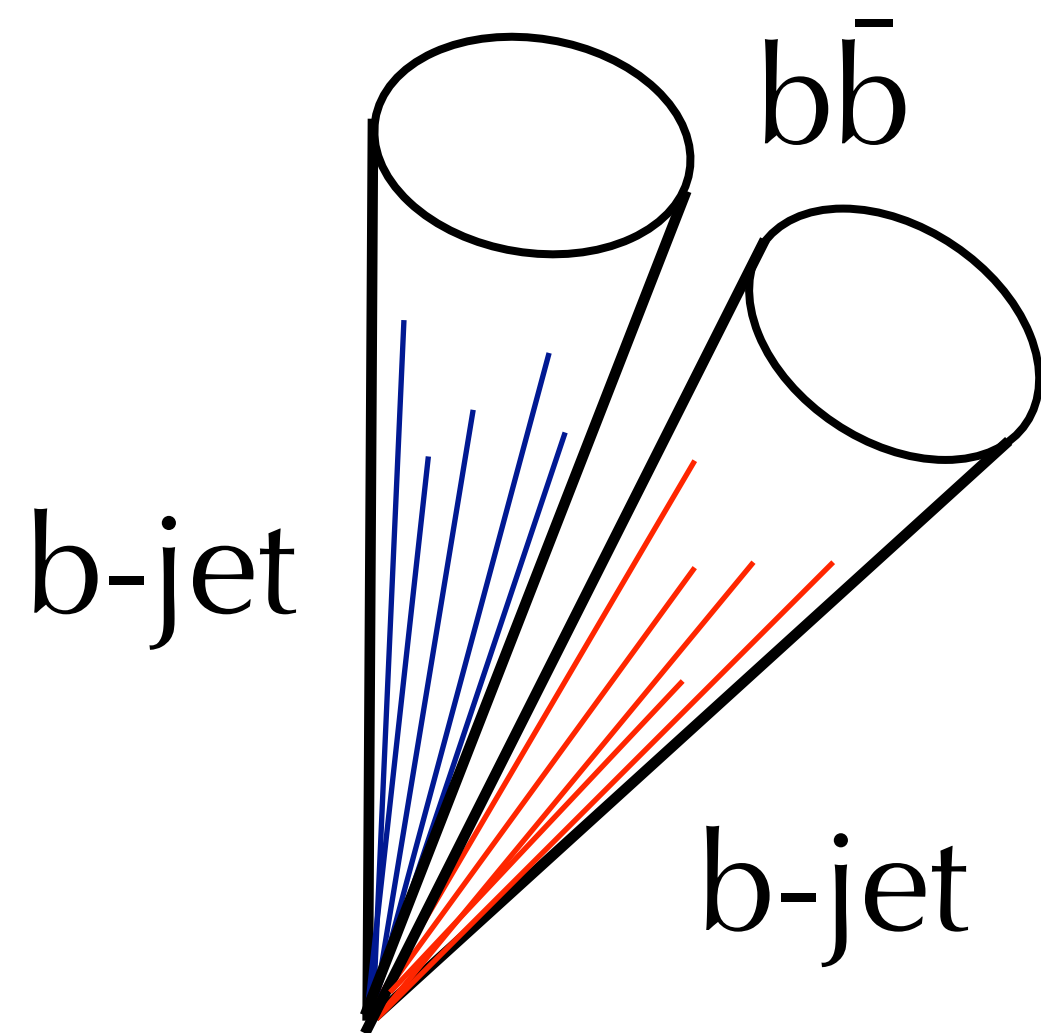
H(bb̄)



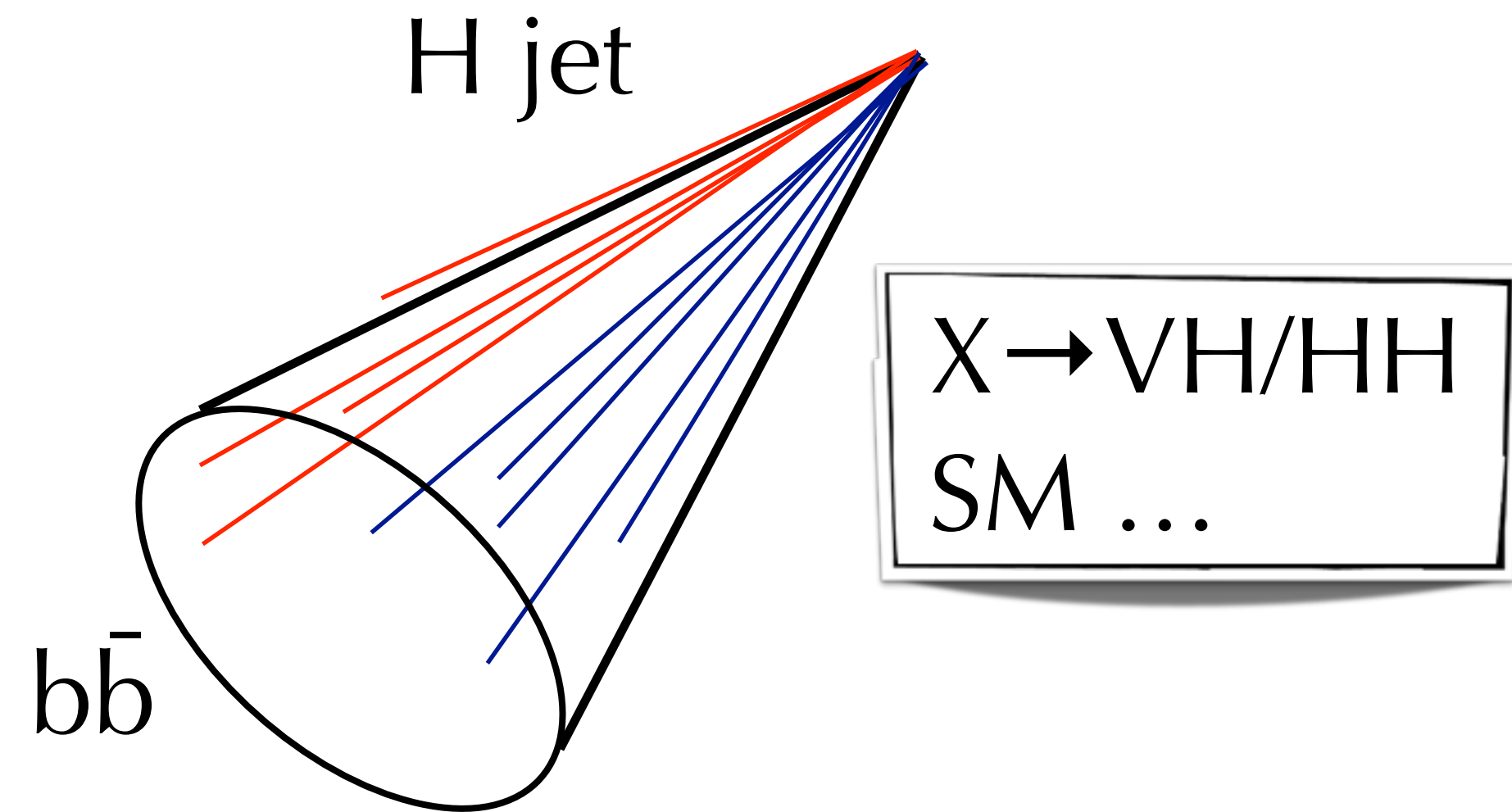
$$dR(bb̄) \sim 2m_H/p_T$$

two-separated b-jets
(R = 0.4)

one single large-cone (fat) jet
(R = 0.8)



VH(bb̄)
VBF H
X → HH(4b)

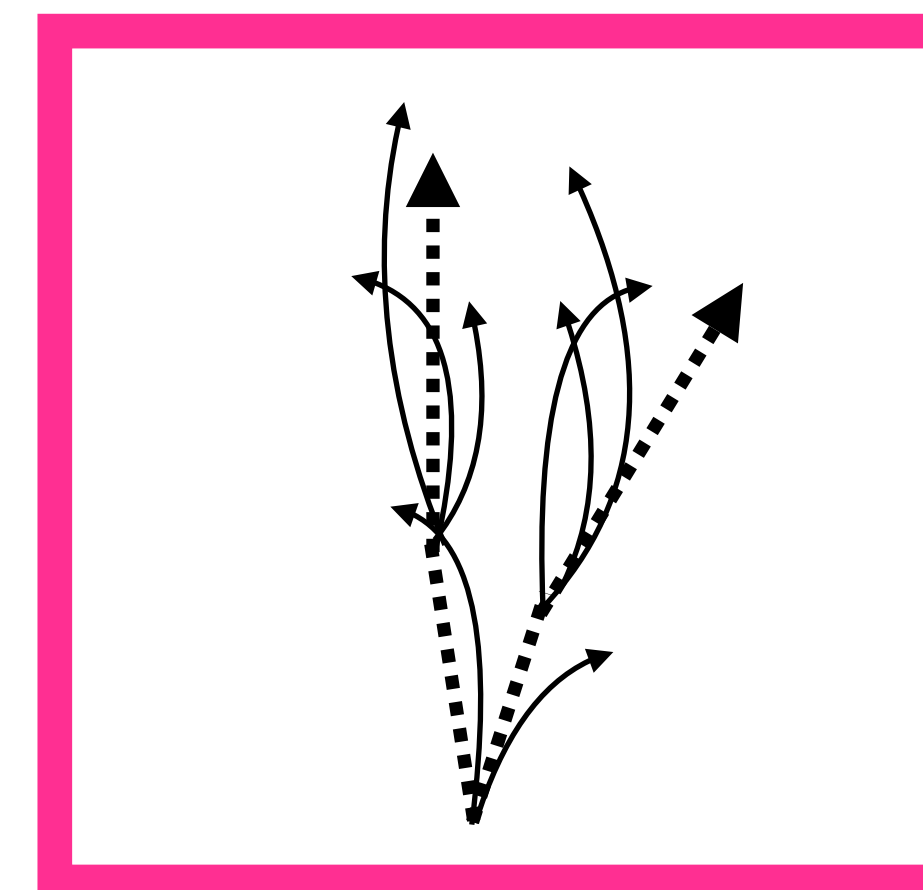


X → VH/HH
SM ...

H($b\bar{b}$) tagging

The boosted H($b\bar{b}$) signal is identified as large cone size jets:

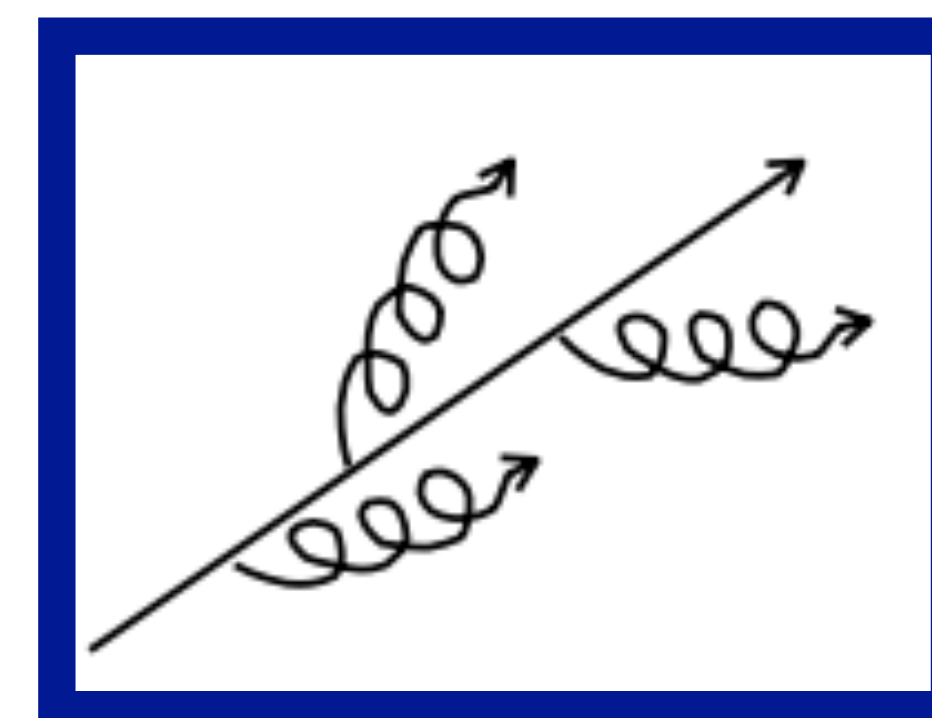
- $R=0.8$
- **PUPPI** (PileUp Per Particle Id) is used to mitigate **pile up effects**



H/Z($b\bar{b}$)

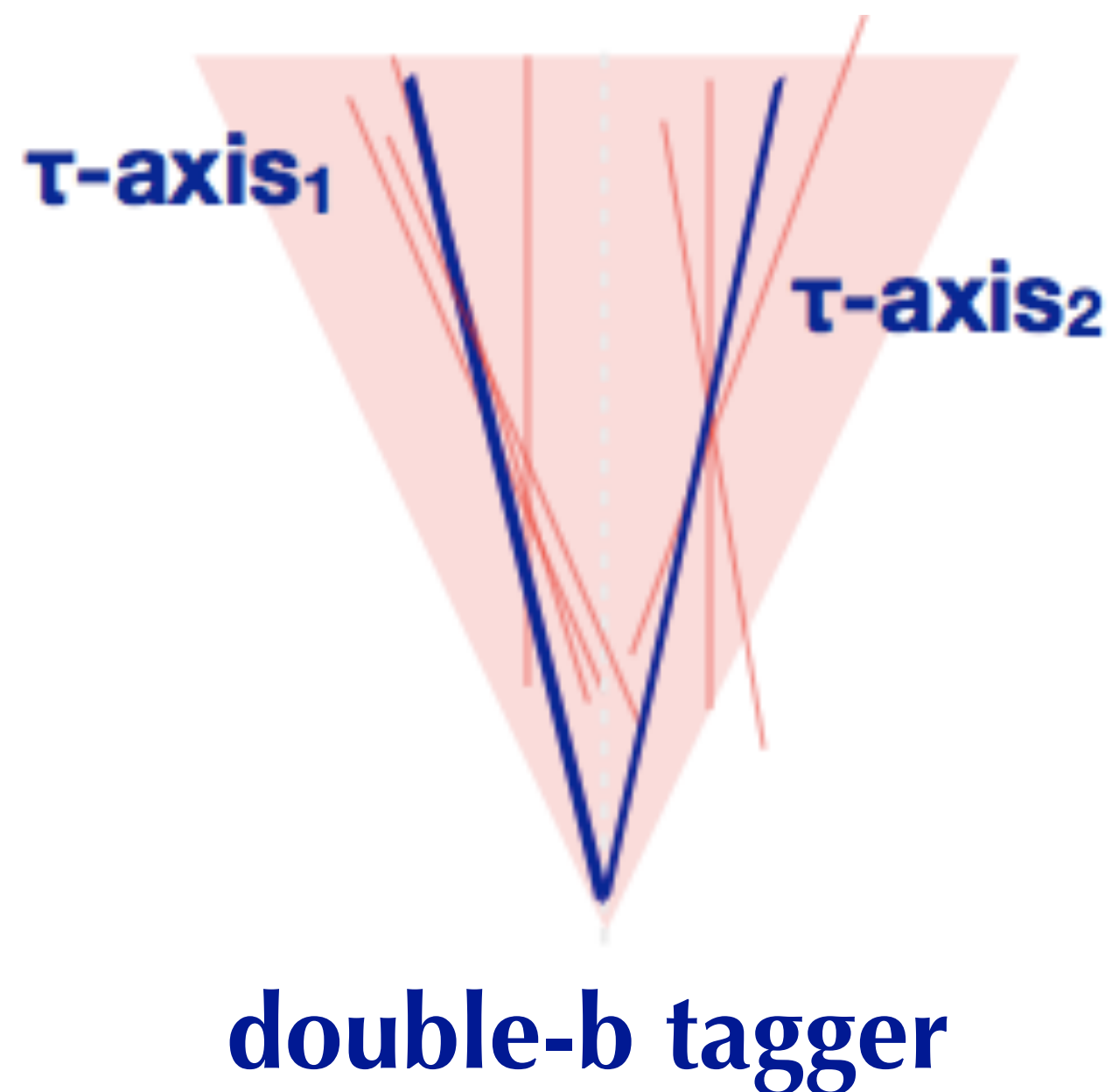
Our tools:

- **b-tagging** to reconstruct the two B hadrons from the b and \bar{b} within the same fat jet
- jet **mass** compatibility with the Higgs
- the composite nature of the jet using **substructure**

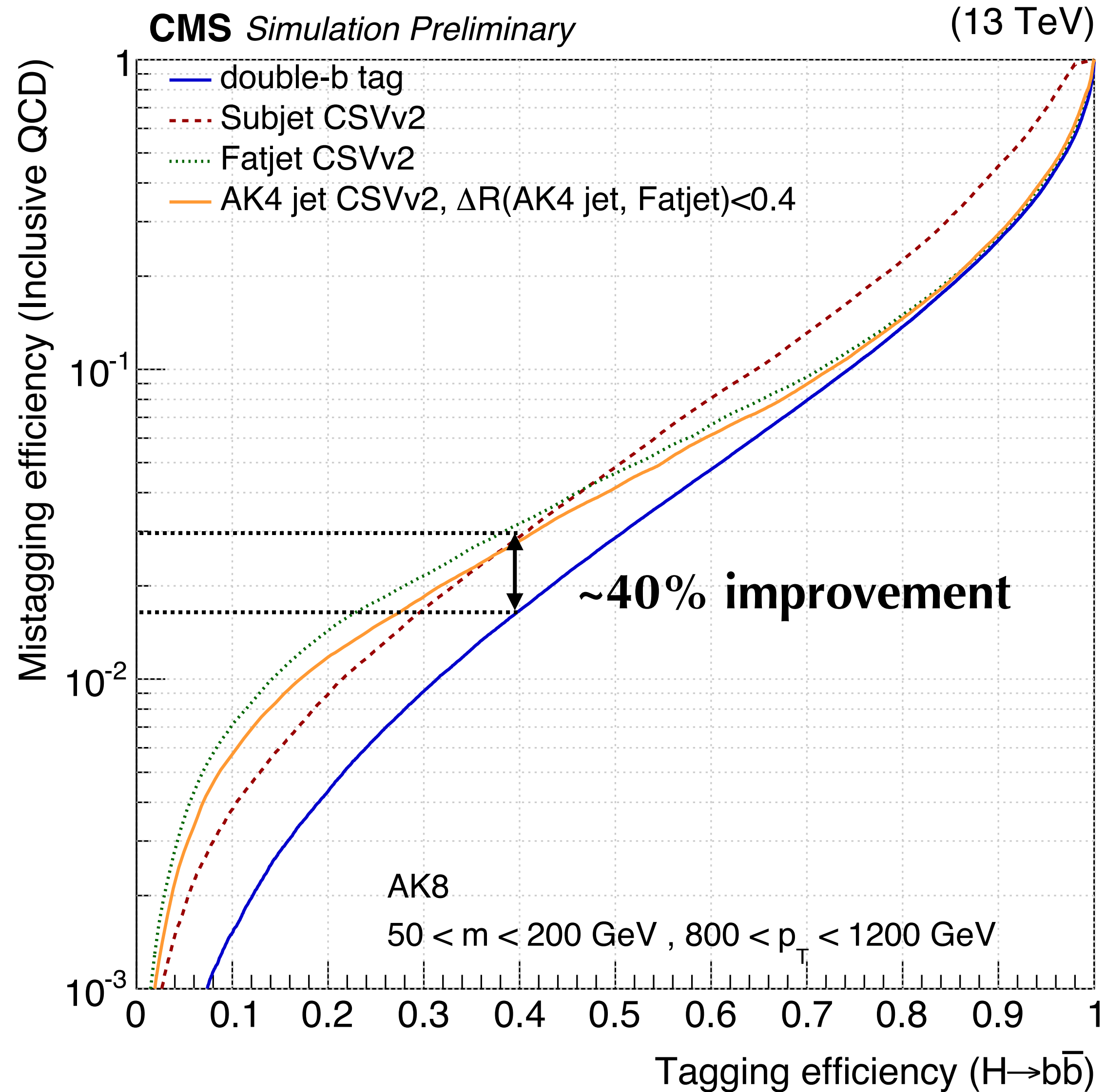


background q/g

double-b tagger

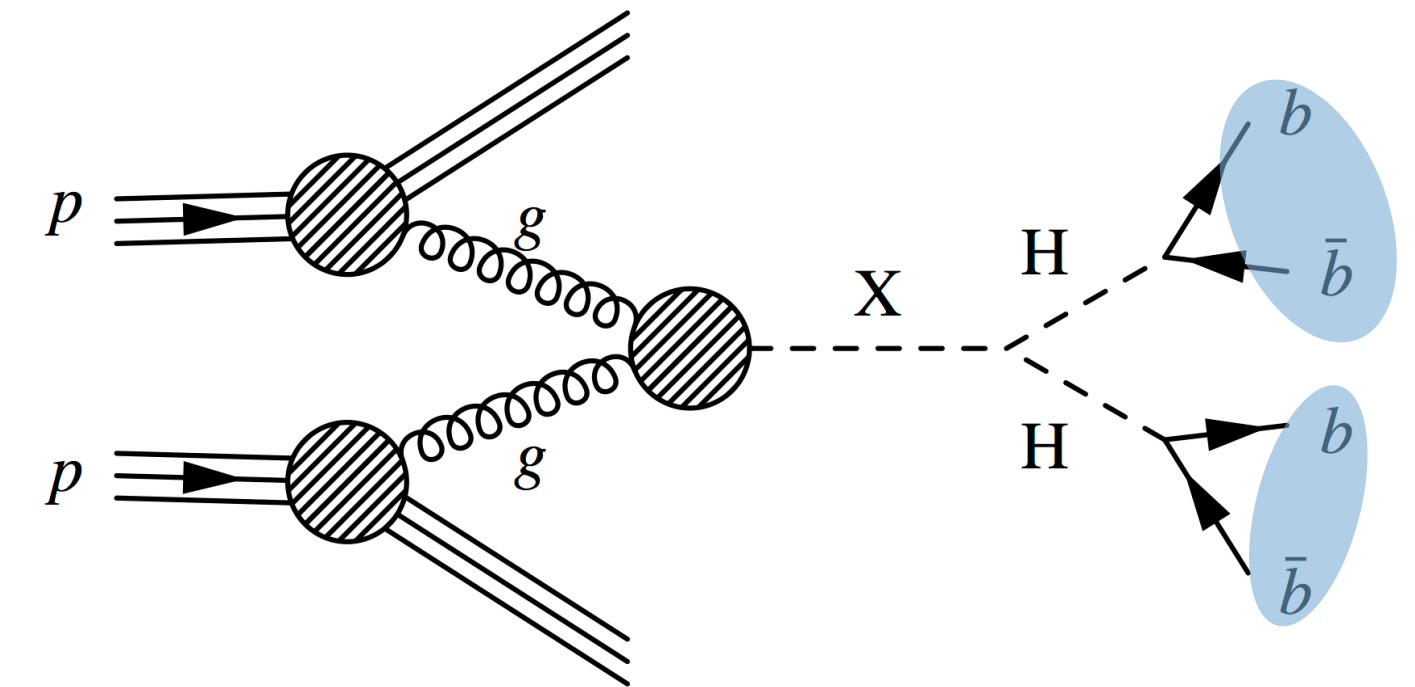


- Tool used in **HIG/B2G/EXO/SUS**
Machine learning approach is being investigated for 2017



Resonant Di-Higgs at CMS

FNAL, Florida U, JHU, Rutgers U, Colorado Boulder U, Kansas U



SM predicts an extremely low rate for hh production (30 fb at 13 TeV)

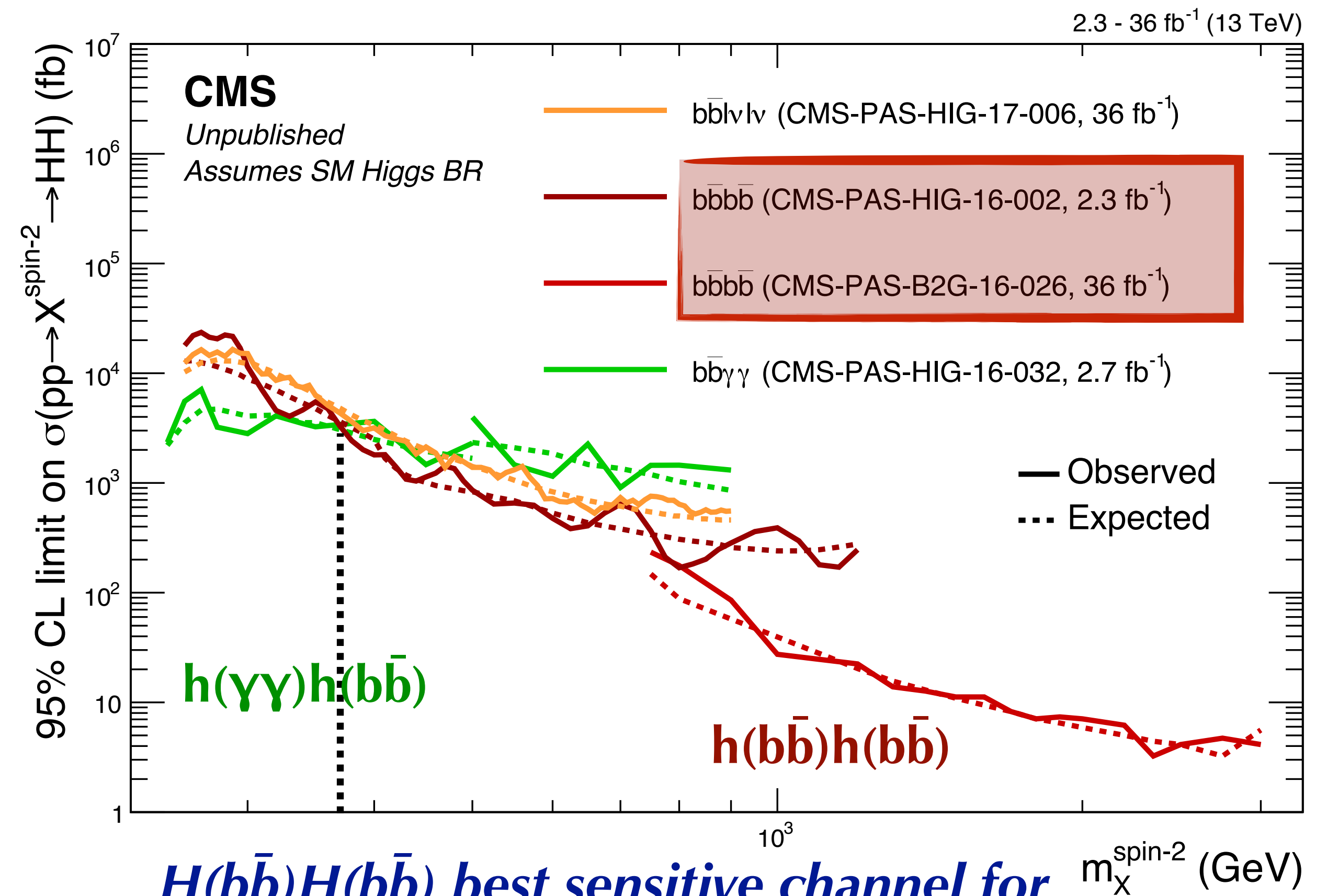
1000 times smaller than the single Higgs boson production

Significantly enhanced in many BSM scenarios

A natural choice is to exploit $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ (highest BR)

No significant excess in the range 1-3 TeV

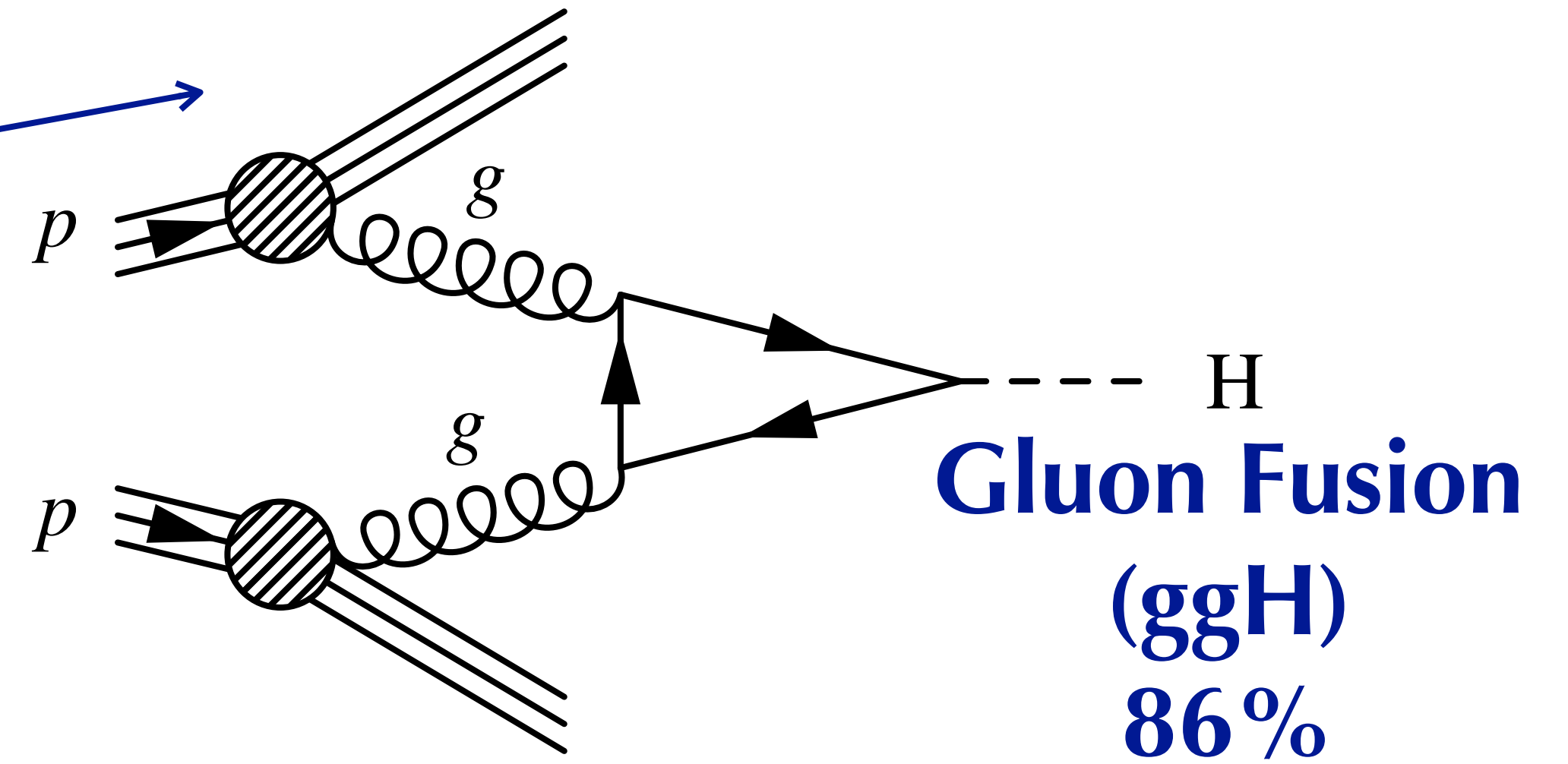
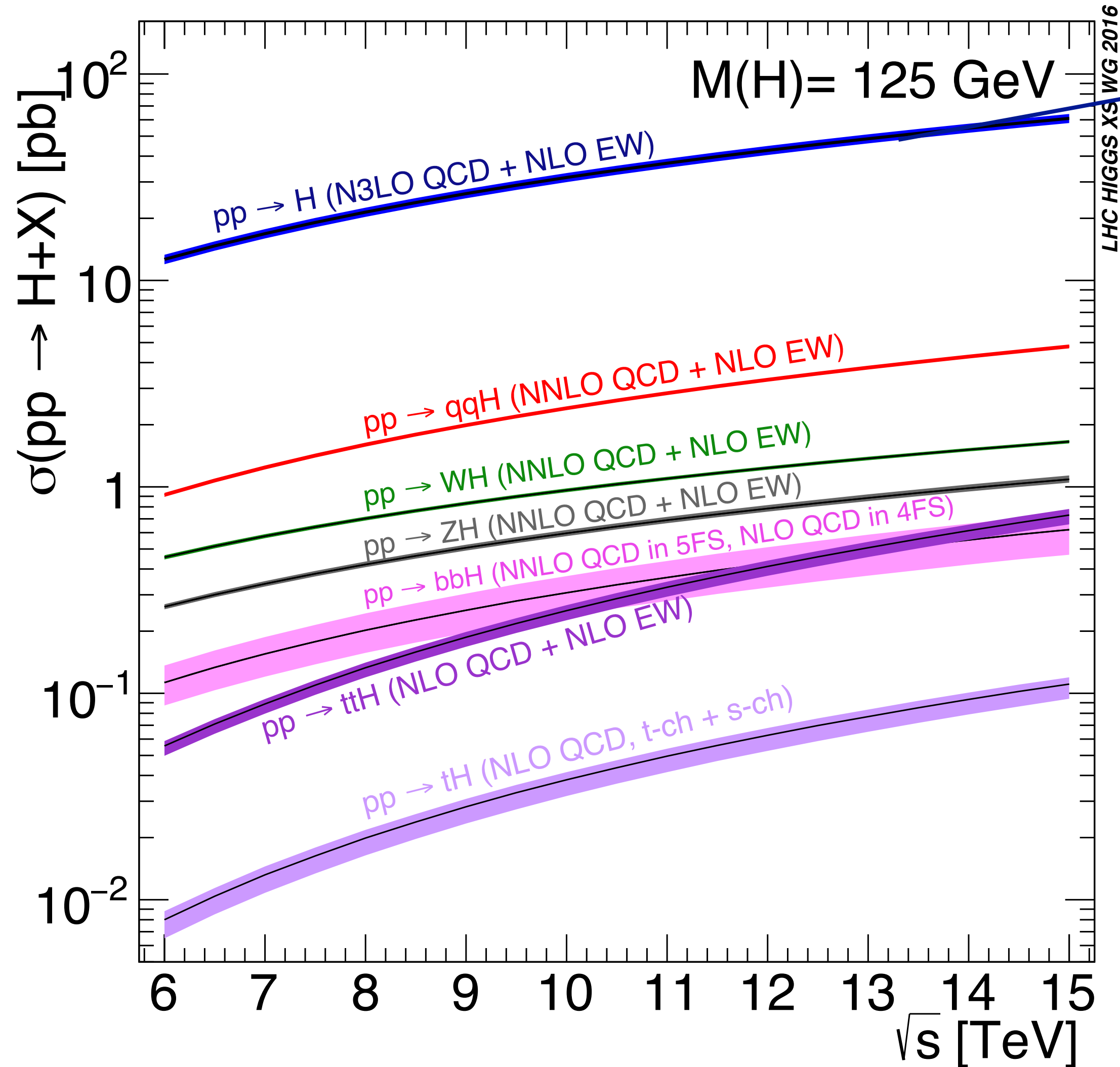
- Results are being published in **B2G-16-026** (just approved) and **HIG-17-009** (being approved)
- **Phase II studies for the TDR**



$H(b\bar{b})H(b\bar{b})$ best sensitive channel for heavy di-higgs resonances

H(b \bar{b}) at LHC

Javier Duarte, Nhan Tran, **C.V. (FNAL)**, Phil Harris (MIT),
Cristina Suarez, Petar Maksimovic (JHU), Michael Krohn
(Colorado Boulder U)

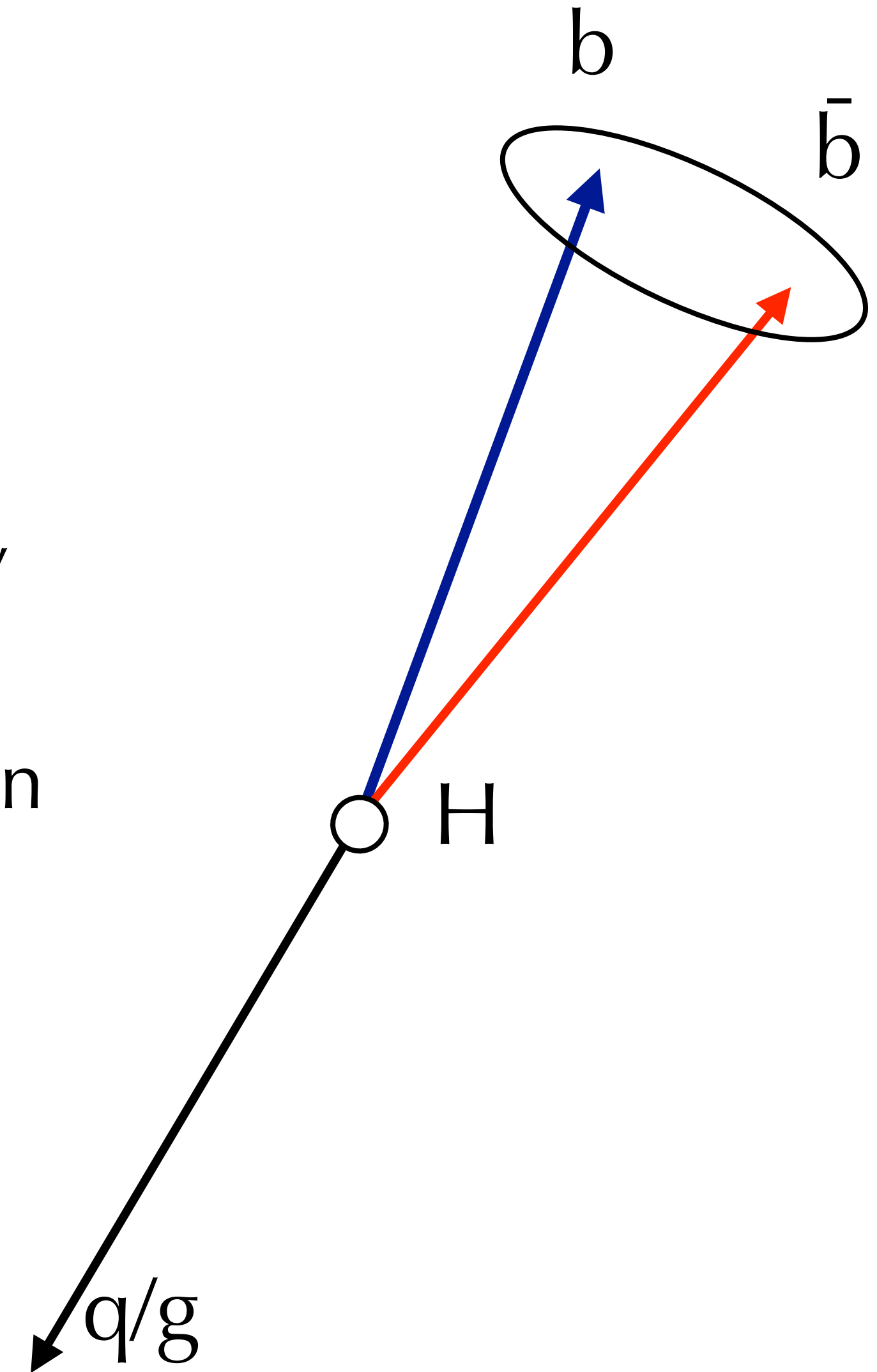


- Two-jets final state
- Overwhelming background from QCD production of b quarks
 - 10^7 larger

Search for $gg \rightarrow H \rightarrow b\bar{b}$ historically **deemed impossible**

Search for inclusive H to $b\bar{b}$

- We can access this process in the **boosted dijet** topology
- Use initial state jet to get above the trigger threshold
- Look for boosted **H boson in a single jet mass** distribution
 - Use the **Z boson as Standard Model** candle
 - b-tagging to disentangle W/Z

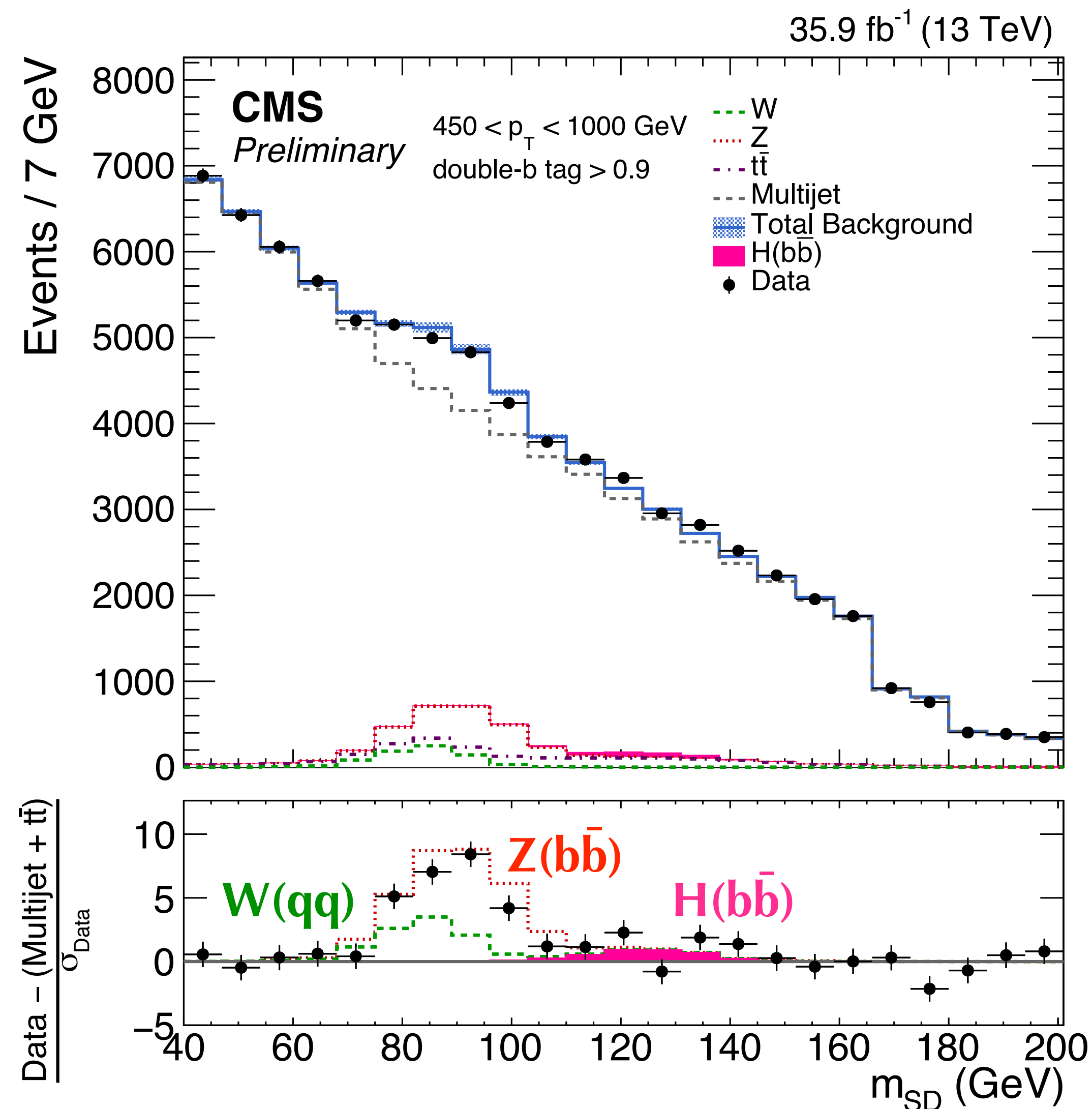


Simultaneous fit of the Z and H signals

- **First Observation of the $Z(b\bar{b})$ in the one-jet topology** 5.1σ (5.8σ)
- The observed significance for the **$H(b\bar{b})$** is **1.5σ** (0.7σ)
- The measured **cross sections** for Z+jets and Higgs for jet $p_T > 450$ GeV are:

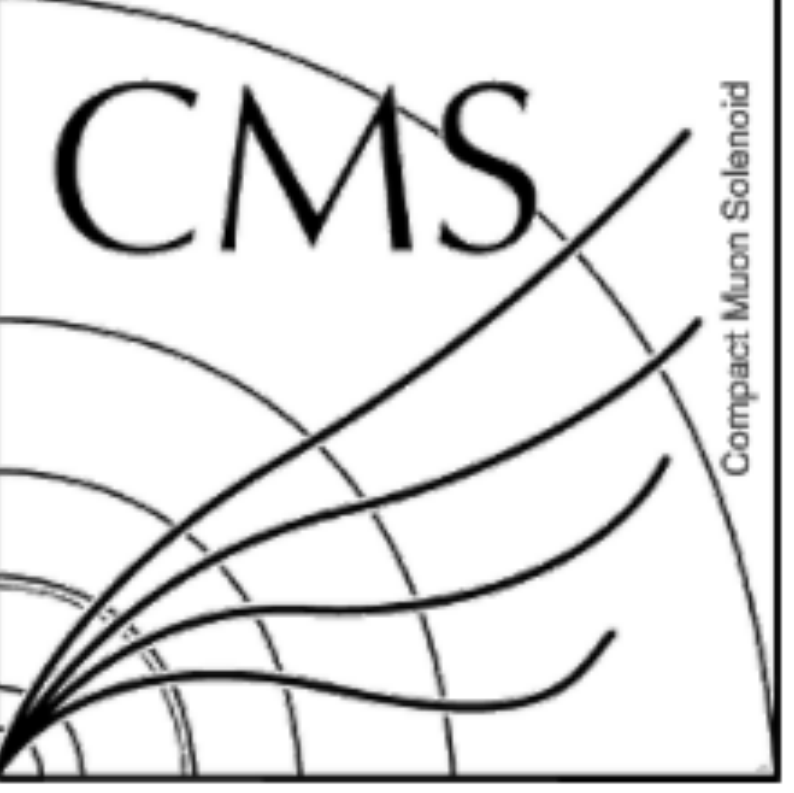
$$\sigma_Z = 849 +155/-155 \text{ (stat.)} +140/-205 \text{ (syst.)}$$

$$\sigma_H = 74 +48/-48 \text{ (stat.)} +10/-17 \text{ (syst.)}$$

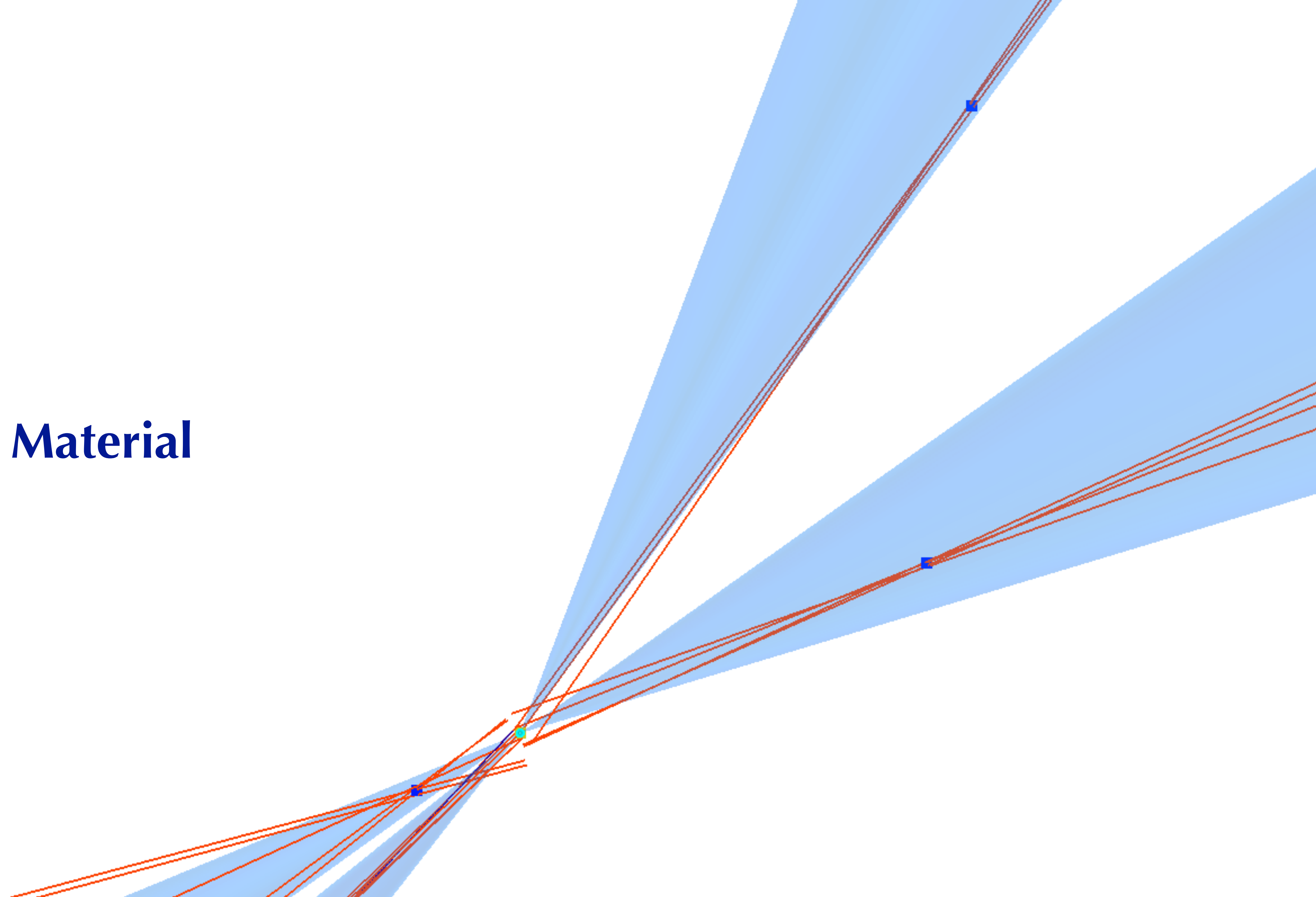


Summary

- Being at LPC allows one to effectively take part in an R&D program as well as have face-to-face discussions with local experts in physics object reconstruction, algorithms and other analyzers to improve the quality of the physics results
- The LPC effort of the Phase-2 pixels R&D has contributed to the CMS Tracker TDR
 - taking advantage of the Fermilab test beam facility and expertise
- Development of a new tagging approach is relying entirely on local expertise
 - continuing developing to make use of new deepNN techniques
 - local tutorials to build new expertise at LPC
- The most sensitive result on Dihiggs searches and the first search for inclusive Higgs boson decay to bb are mostly an LPC effort



Additional Material



ongoing R&D

Three different lines of the investigations for the inner parts of the pixel detector:

1. planar pixel sensors in n-in-n

2. 3D pixel sensors

- small pitch sensors investigated by CNM (Spain) and FBK/INFN (Italy)
- thinner wafers (200 μm)

3. thin planar pixel sensors in n-in-p

• single sided process favored

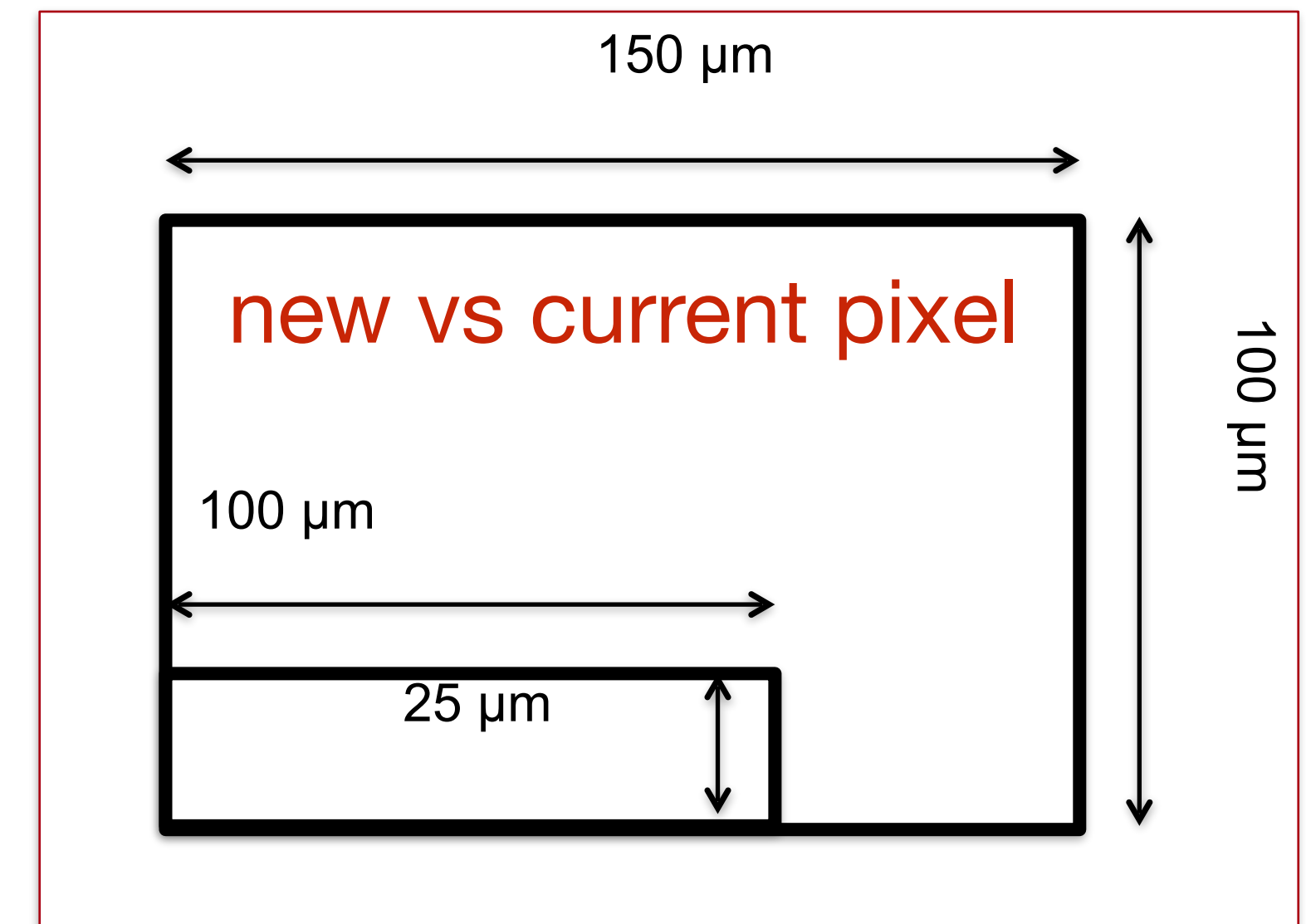
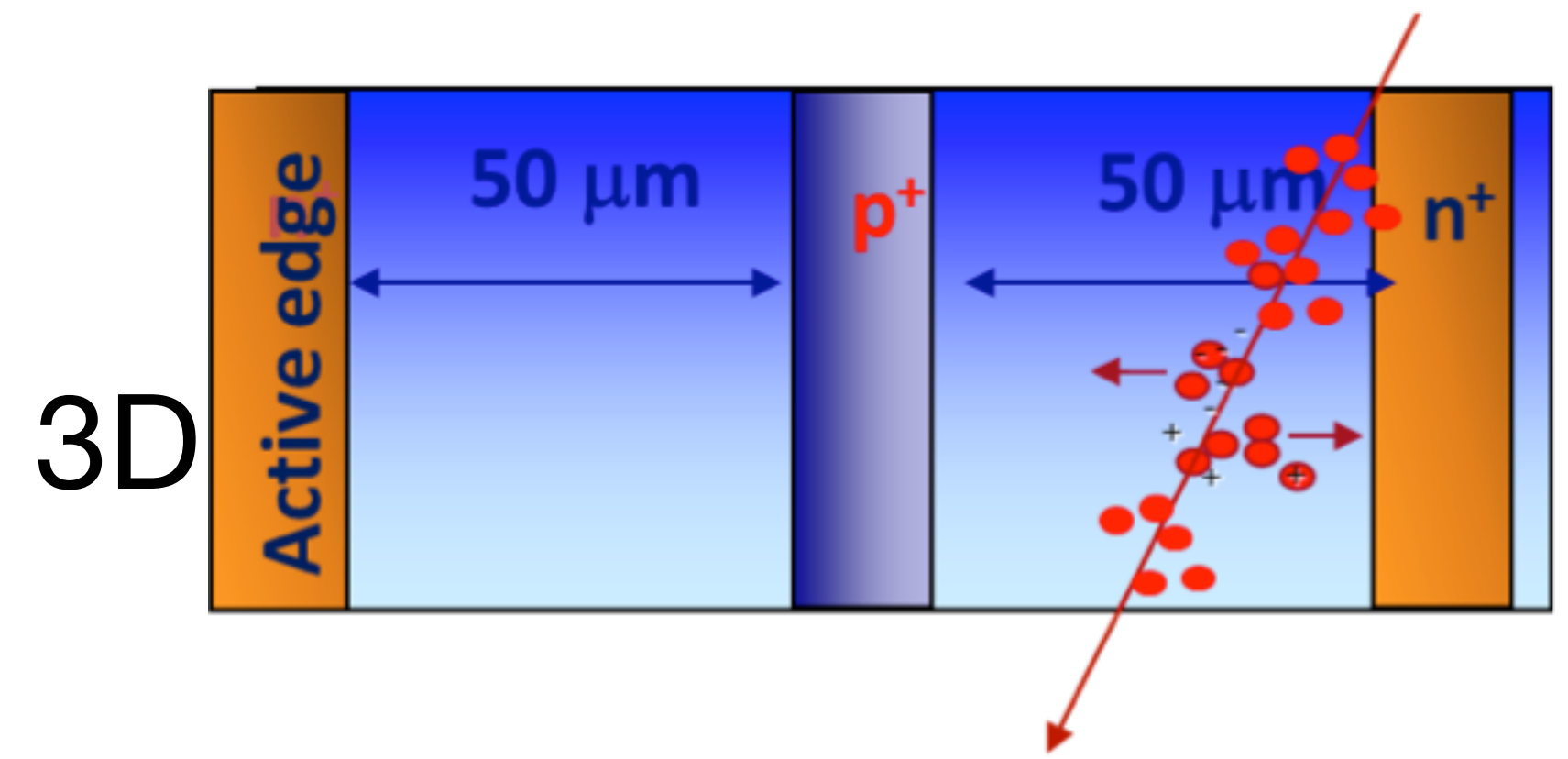
- More vendors, **cost effective**

- 6" n-in-p FZ with 150 μm thickness submission of small pitch pixel sensors

- pixel sizes of 25 μm x100 μm and 50 μm x50 μm under investigations

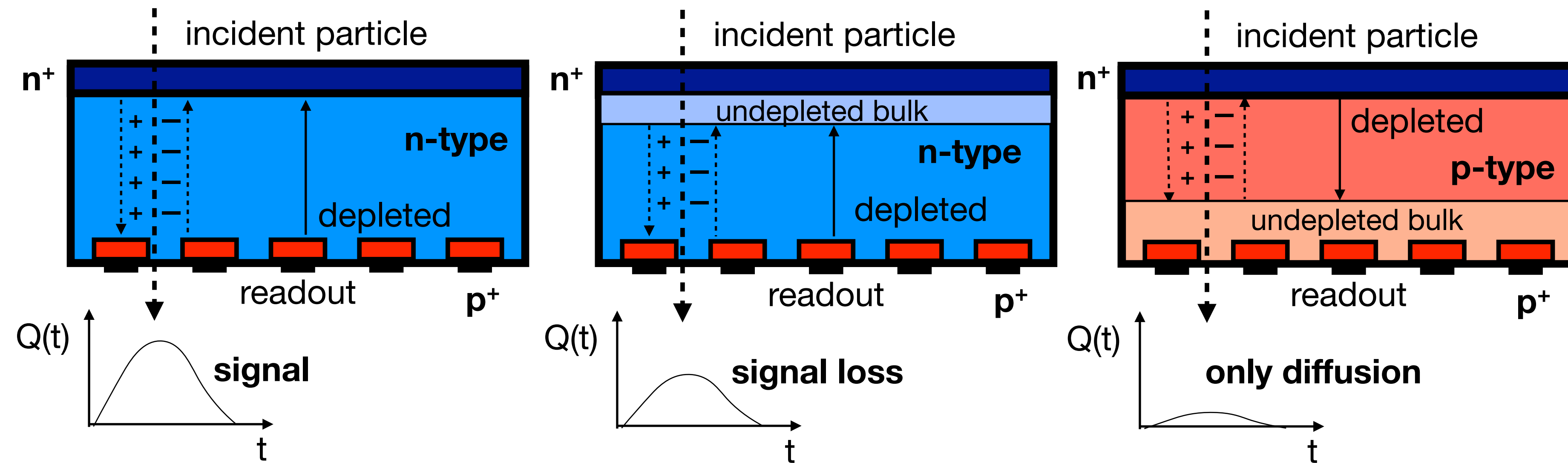
Common advantages:

- Short drift path
- Higher field at same V_{bias}
- Lower operation voltage



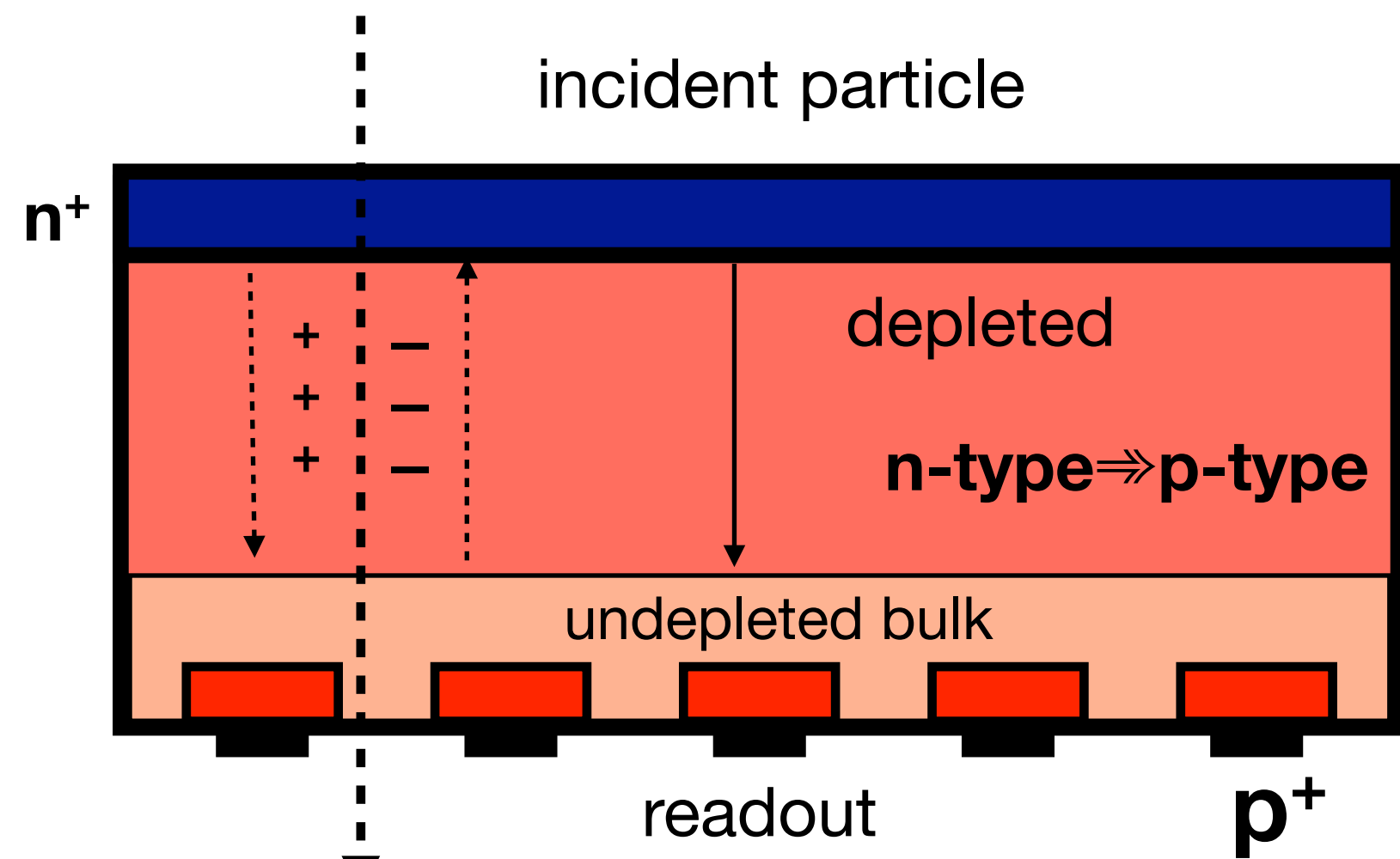
p-in-n

After irradiation (to a level implying type inversion) p-in-n does not function well if not fully depleted.

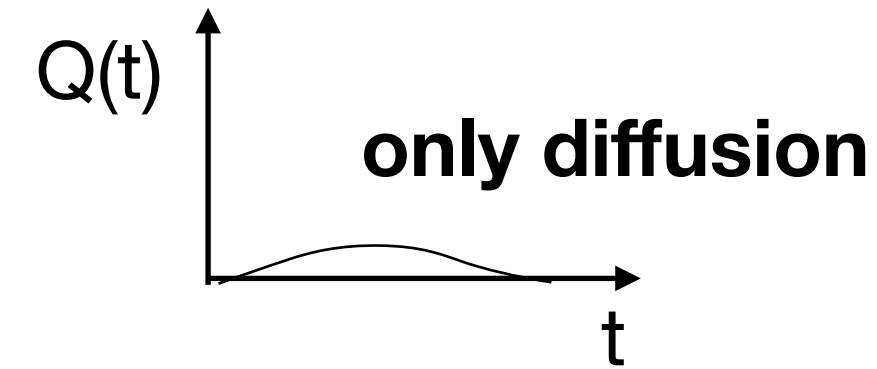


p-in-n vs. n-in-n

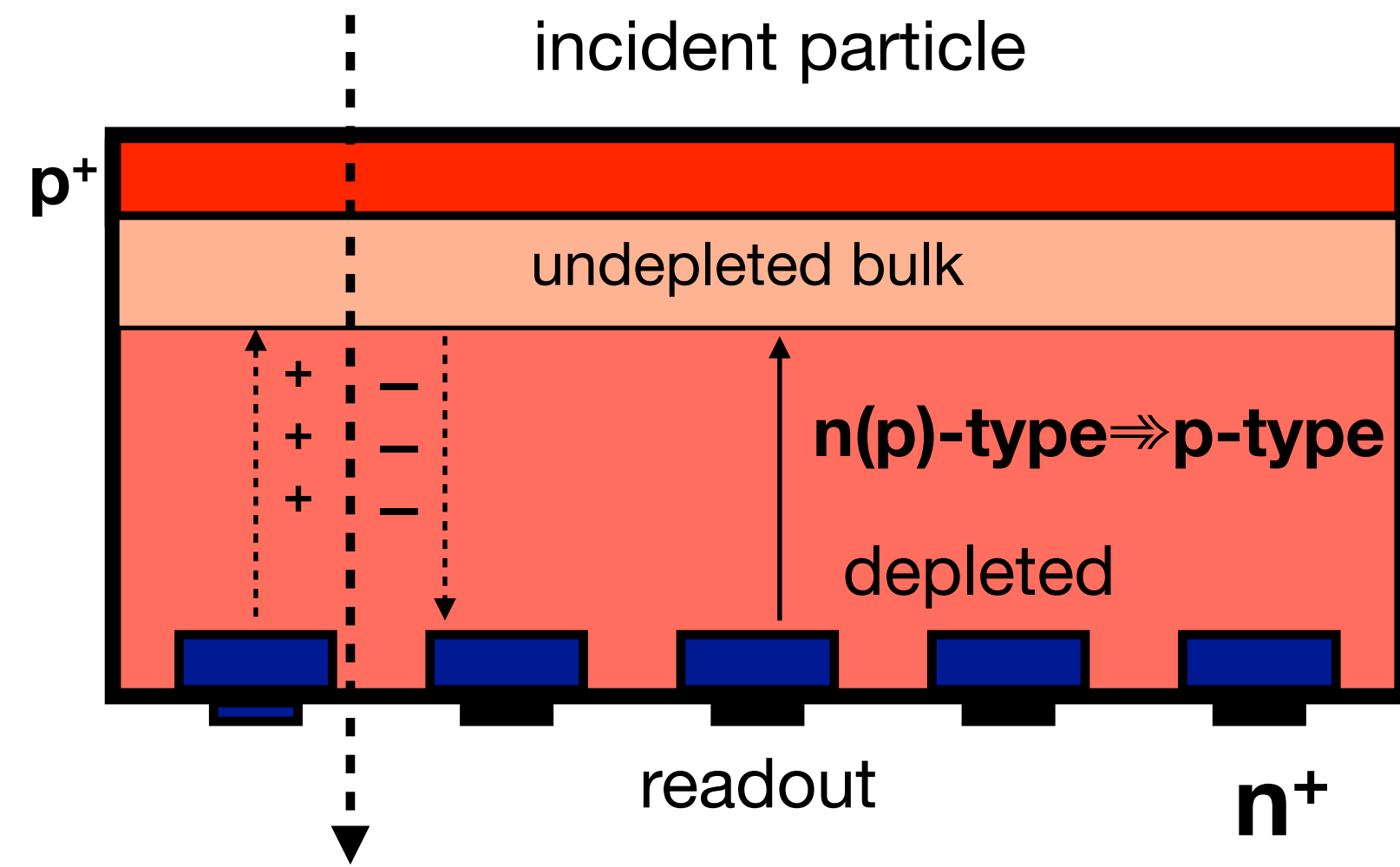
after irradiation and type-inversion



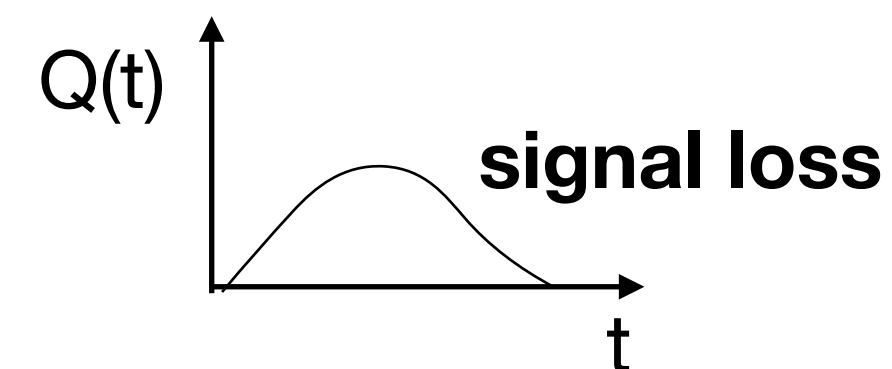
p-in-n



- signal loss
- resolution degradation due to charge spreading



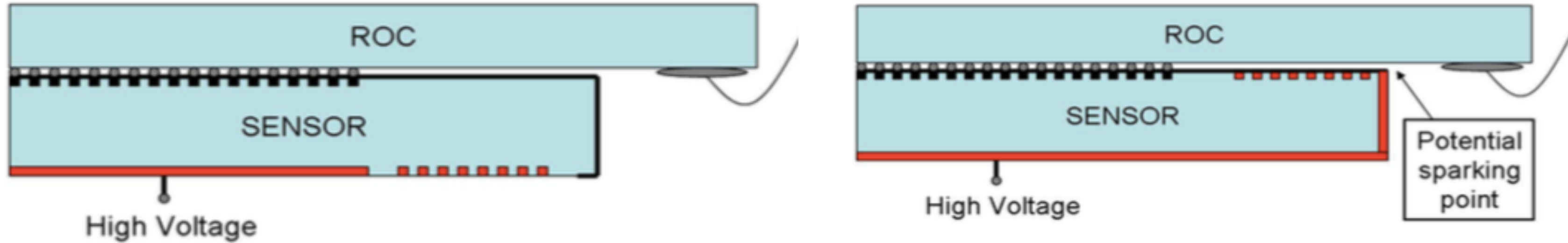
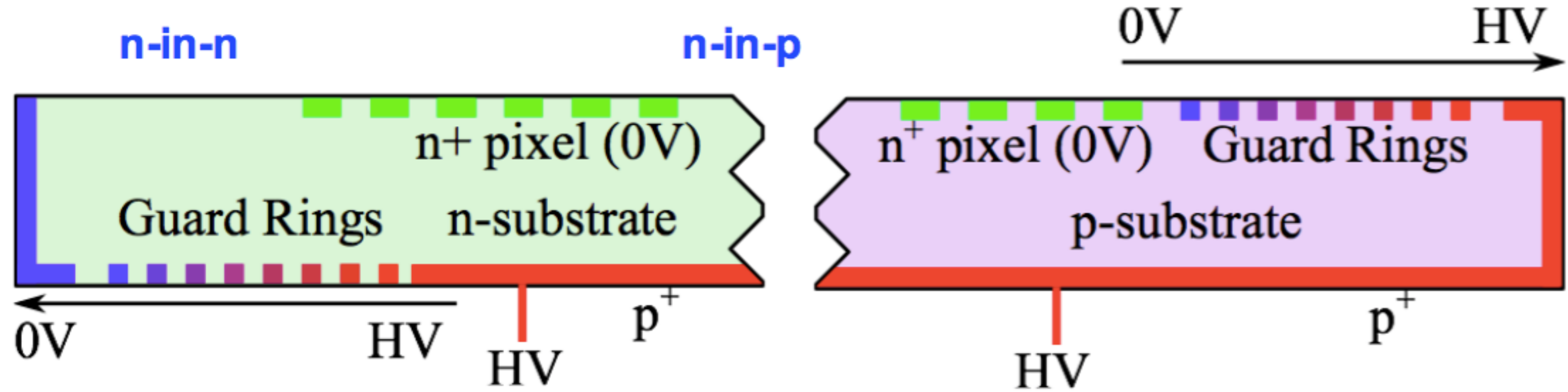
n-in-p or n-in-n



- faster charge collection (drift for electron is larger)
- CCE degradation

n-in-n vs n-in-p

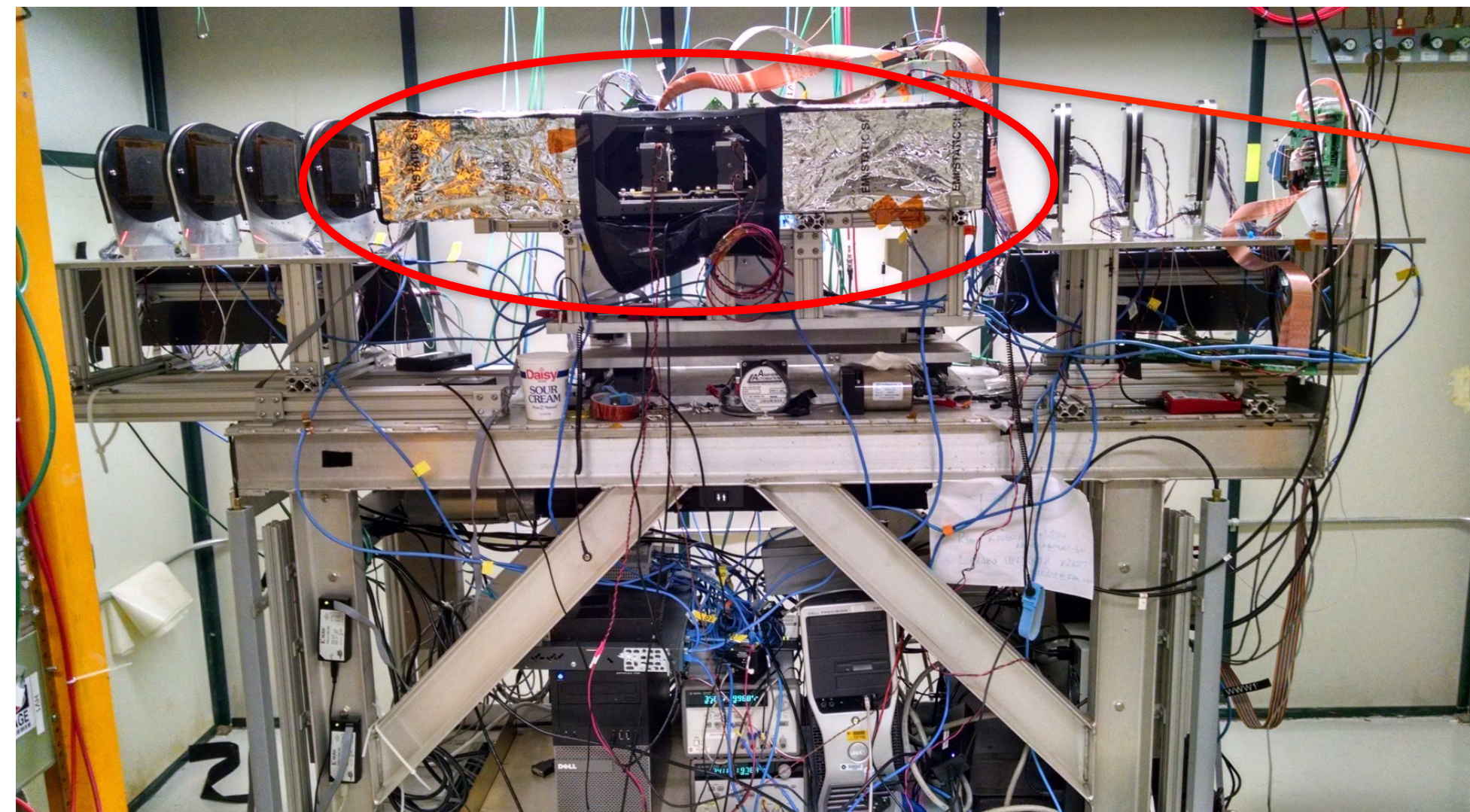
double sided vs single sided



The Fermilab MTest beam facility

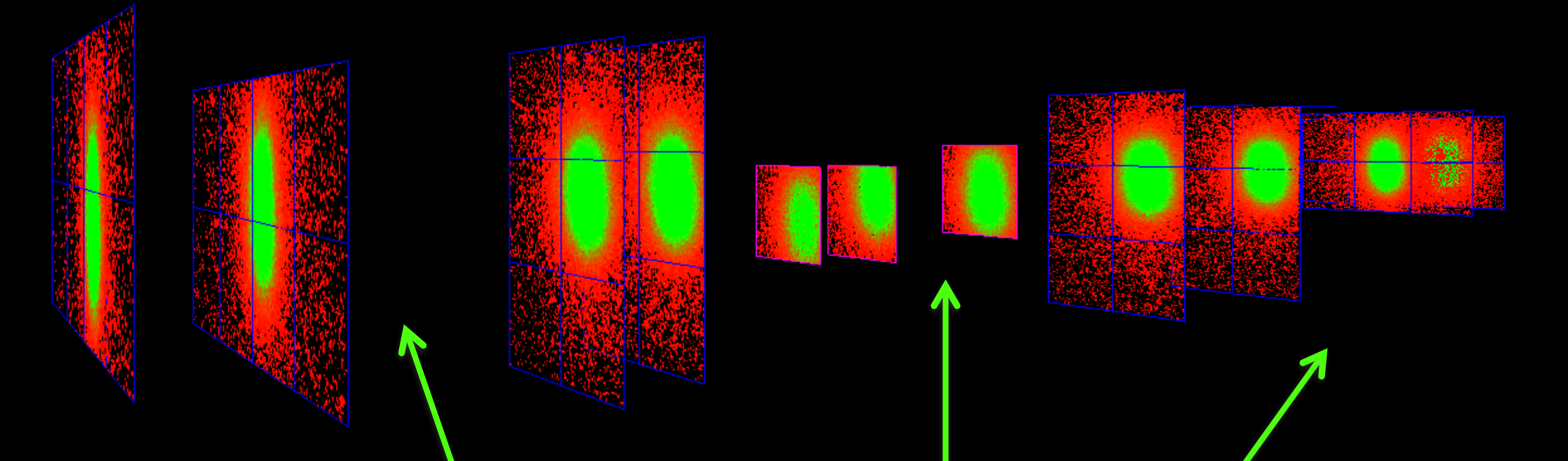
Fermilab Test Beam Facility provides **120 GeV protons** from Main Injector

- At present 2 independent tracking telescopes are installed:
 - The legacy **pixel telescope** built using leftover CMS modules
8 pixel planes
readout based on PSI46 analog chip (Phase 0)
(100x150 μm^2 pixel cell for 80 rows and 52 columns)
~8 μm resolution on each coordinate
 - The new strip based telescope (being commissioned now)



Pixel Telescope

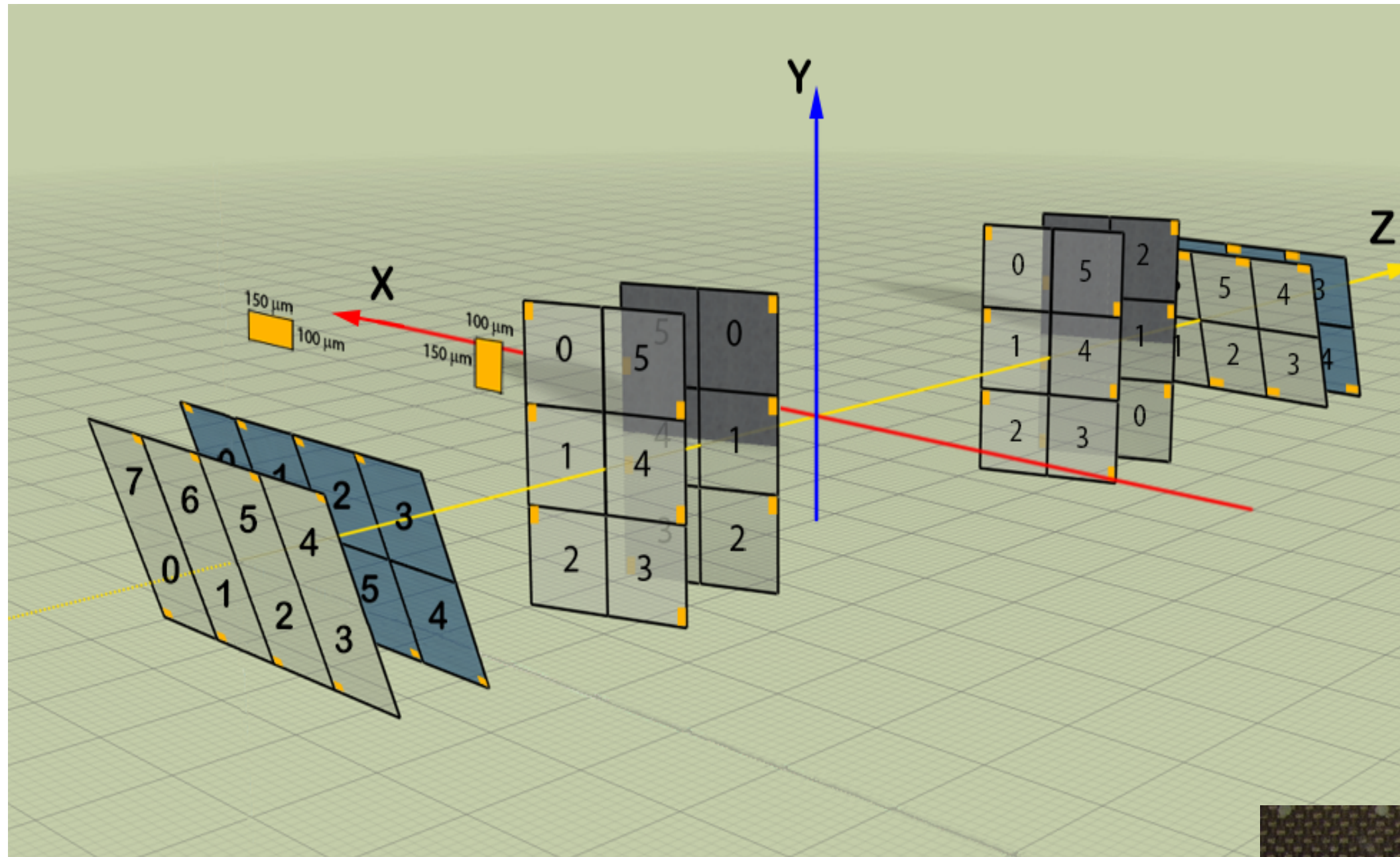
Online data taking



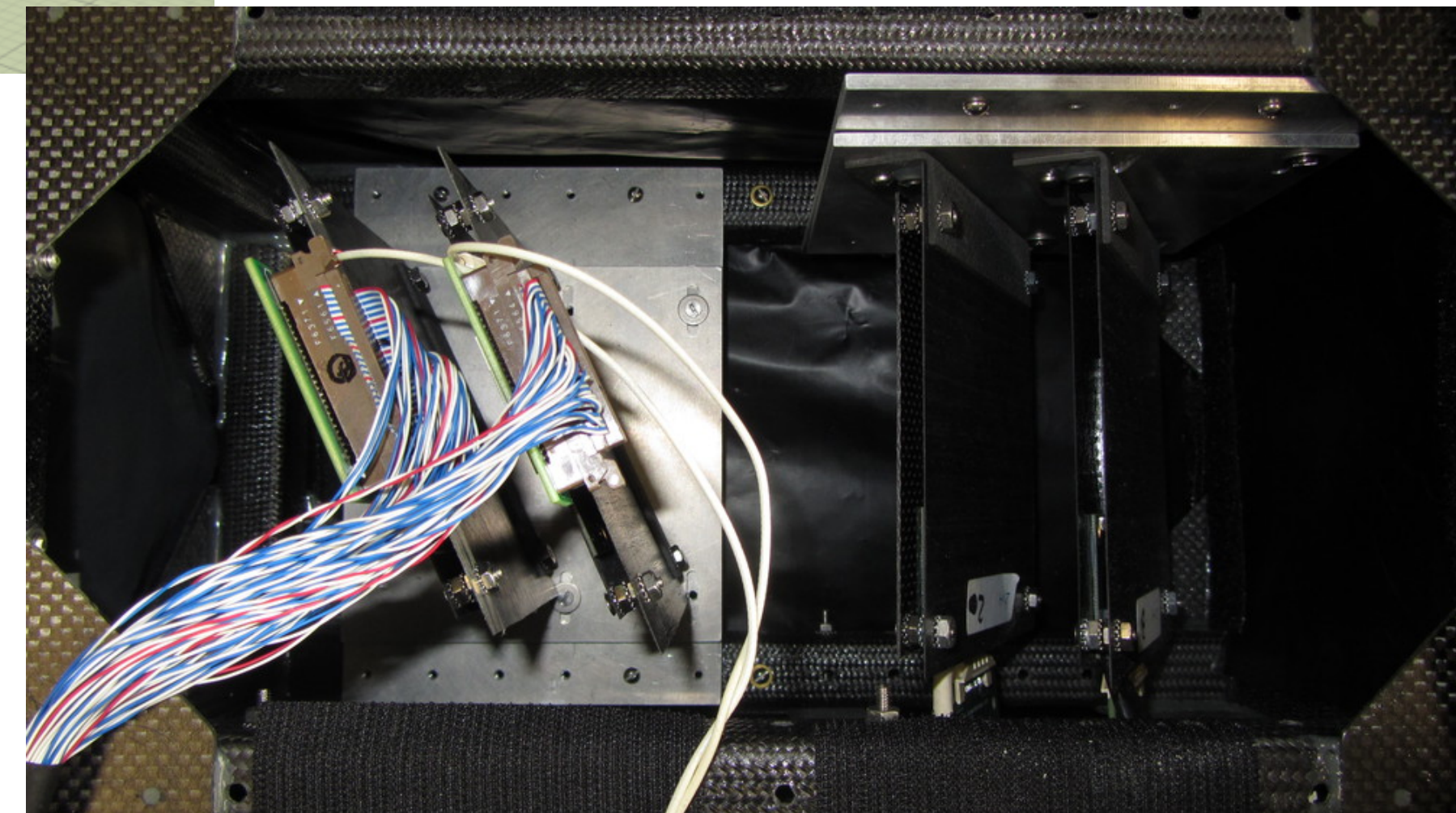
3 detectors under test

Pixel telescope

Pixel telescope planes

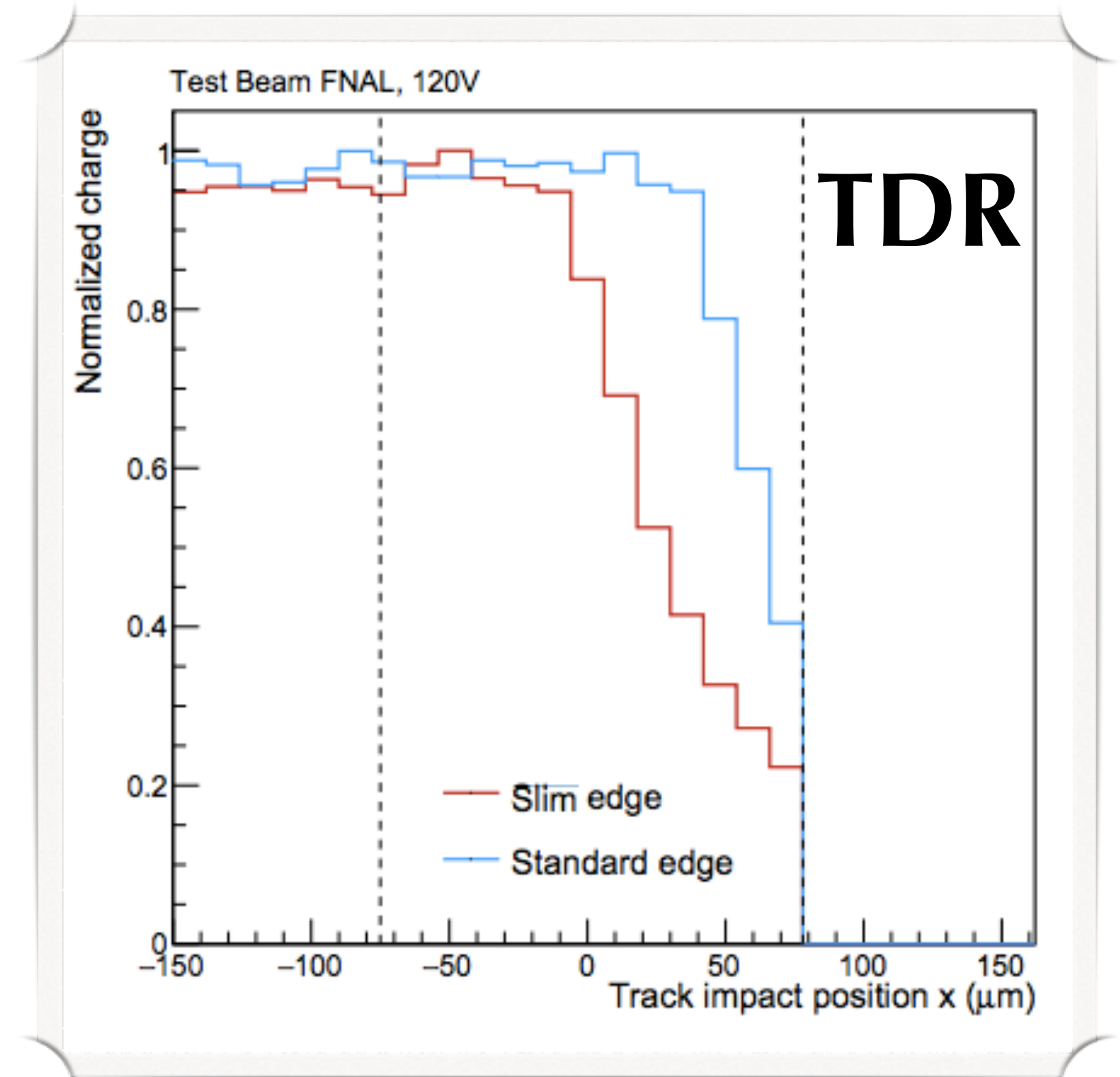
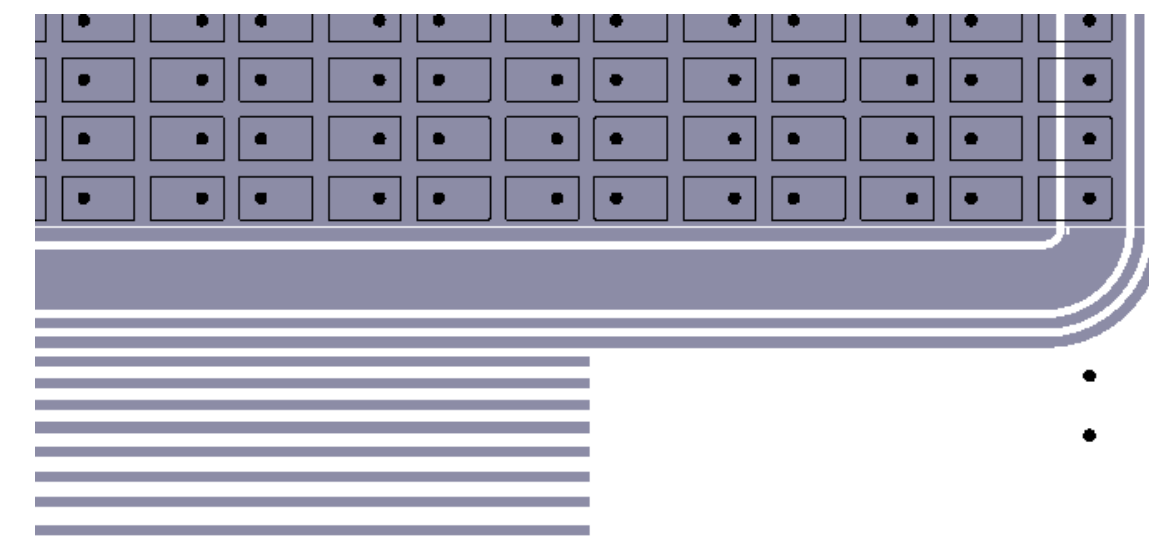
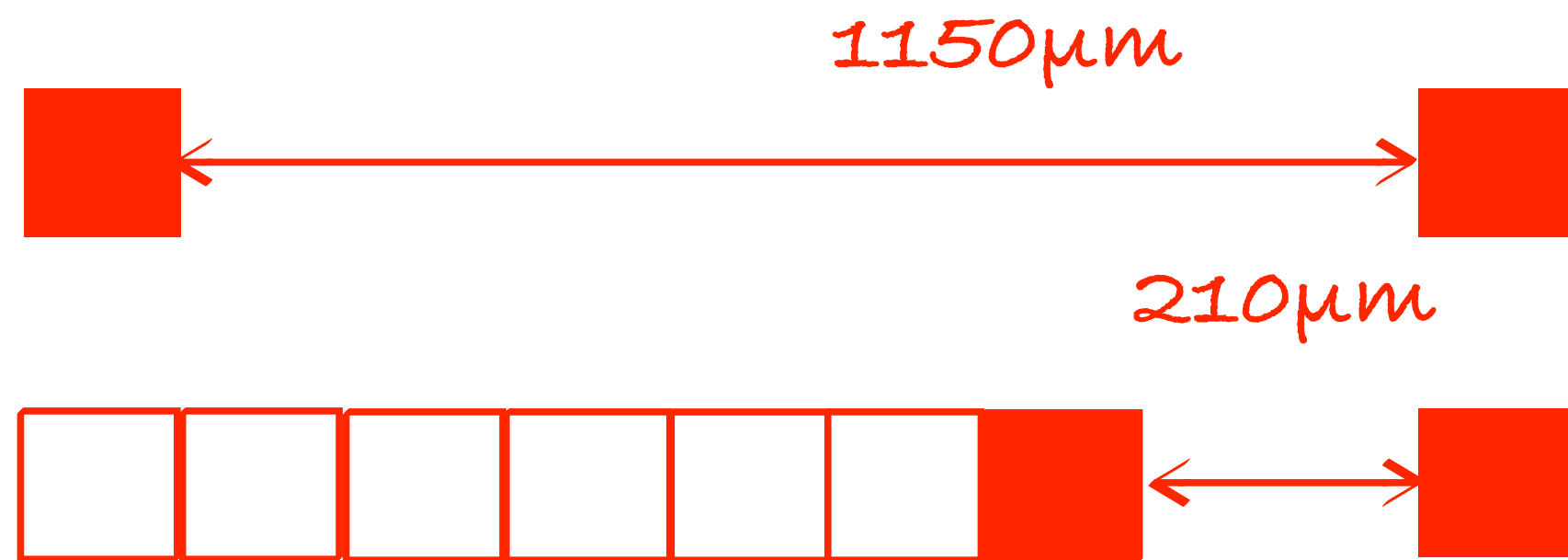


- 8 modules tilted at 25° to maximize charge sharing
- Coverage area only $1.6 \times 1.6 \text{ cm}^2$
- Resolution on the detector under test of $\sim 8 \text{ } \mu\text{m}$



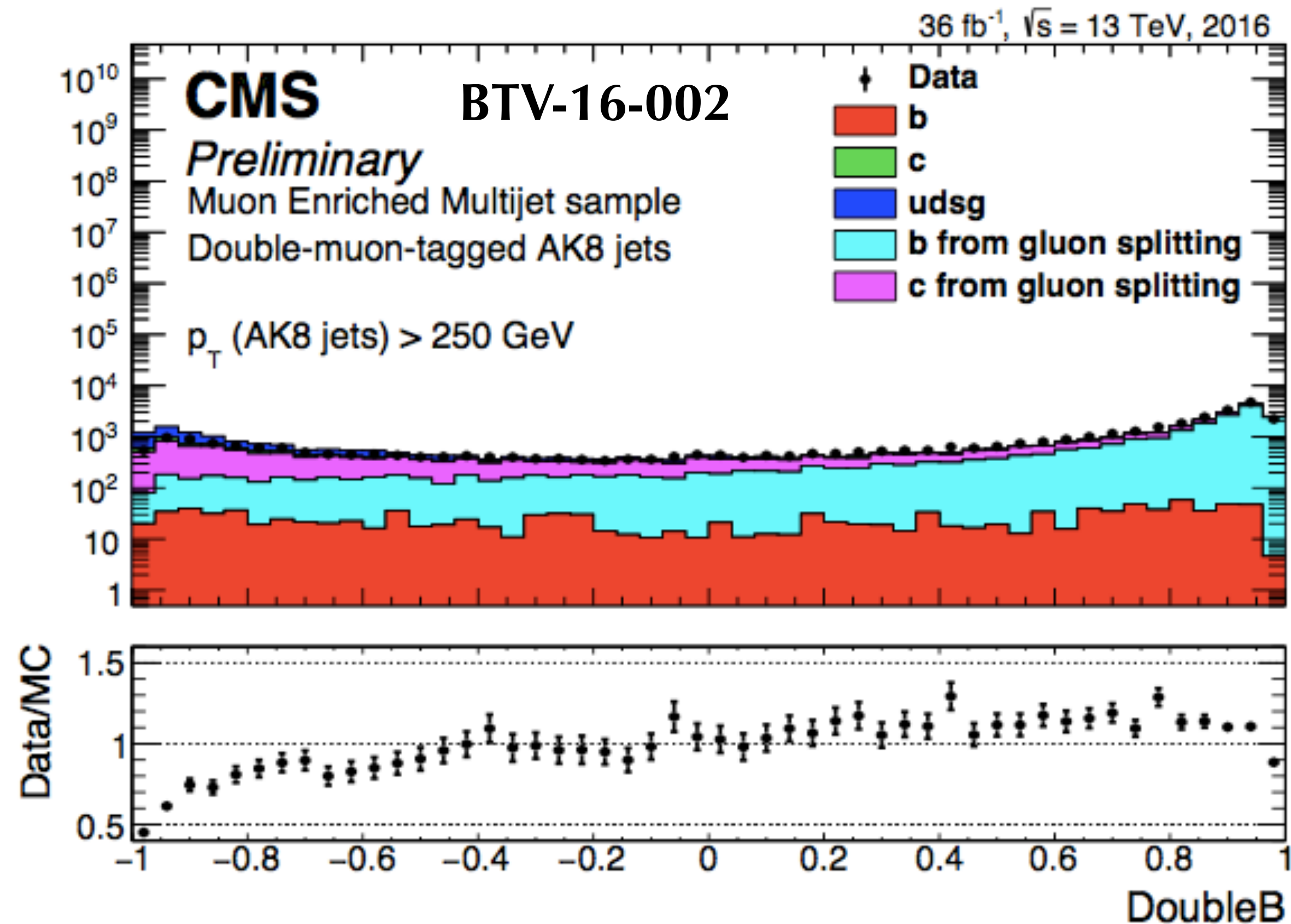
Slim Edge

- Aim at minimizing the dead region at the physical edge of the sensor itself
- The efficiency of the pixels closely located to the edge depends on the geometry of the guard-rings on the opposite side (p-side).
- The pixel array (active area) normally ends at 1.15 mm from the dicing edge.
 - In these prototypes such distance has been reduced down to only 210 μm .
 - The reduction is of $\sim 950 \mu\text{m}$
- * Loss of efficiency on the edge due to error on the track extrapolation
- * Right and left edge behave similarly
- * Slight worse efficiency for the Slim edge
 - * effect on the last 15 μm , but 950 μm are gained (~ 6 pixels)



double-b Efficiency measurement in data

- Since there is no signal (yet!), and signal/background for $Z(b\bar{b})$ is too small to measure now, we use:
 - $g(b\bar{b})$ jets to measure the signal efficiency
- Event selection includes high p_T jet ($p_T > 300$ GeV), double-muon tagged jets, and a soft drop mass cut to ensure that $g(b\bar{b})$ is in a phase space that is signal-like
- Measurement performed using 36/fb collected at 13 TeV (2016)
 - single jet triggers with thresholds such that we are 99% efficient for jet $p_T > 300$ GeV



The Higgs boson

LHC Run I legacy

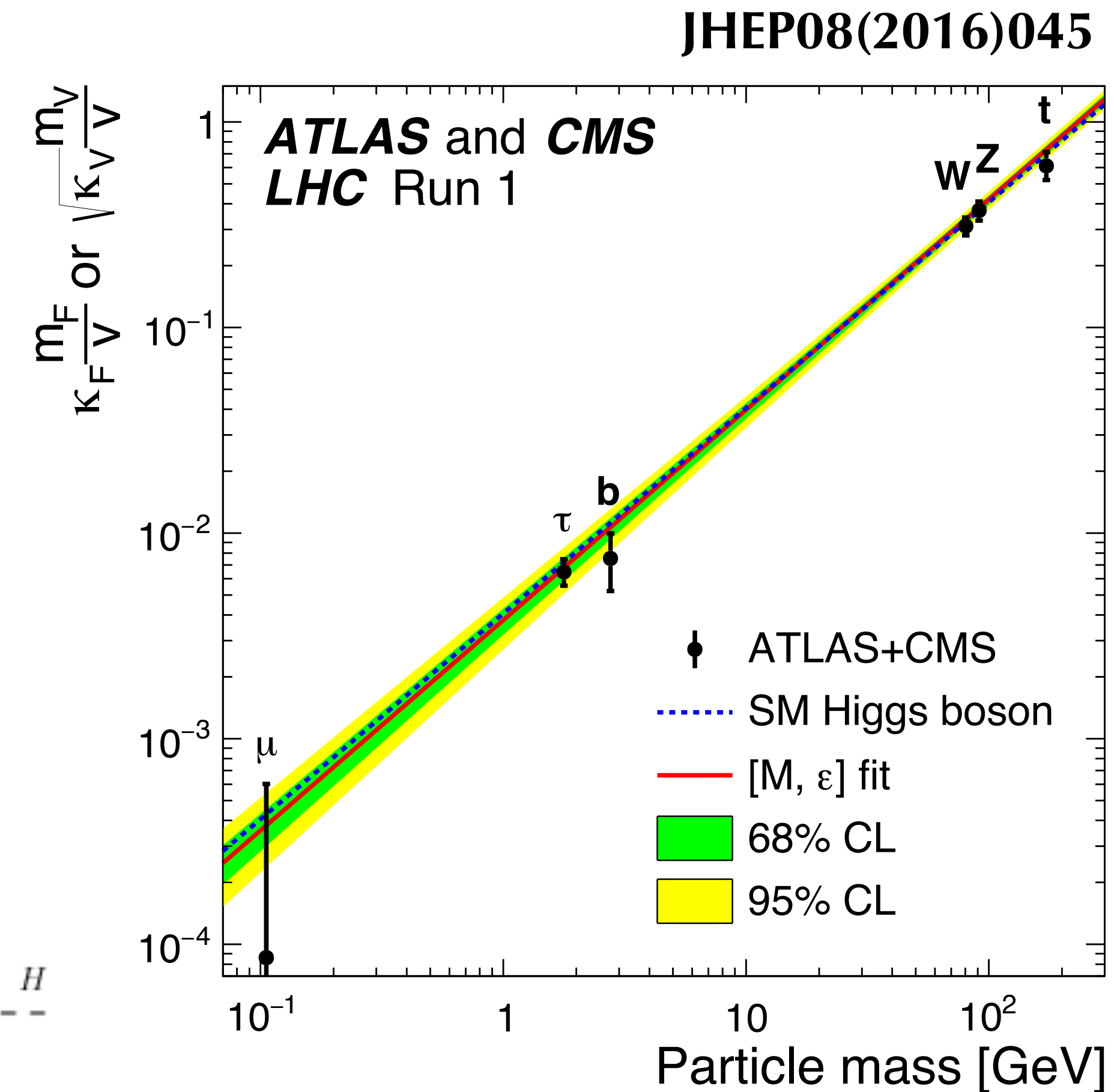
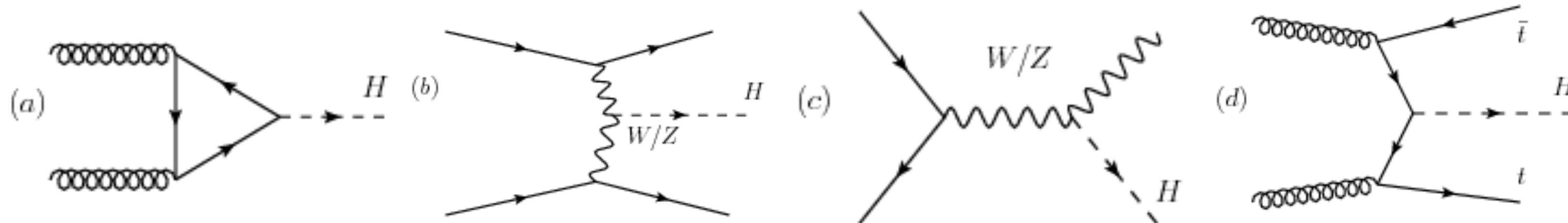
H

$$J^P = 0^+$$

$$m_H = 125.09 \pm 0.24 \text{ GeV}$$

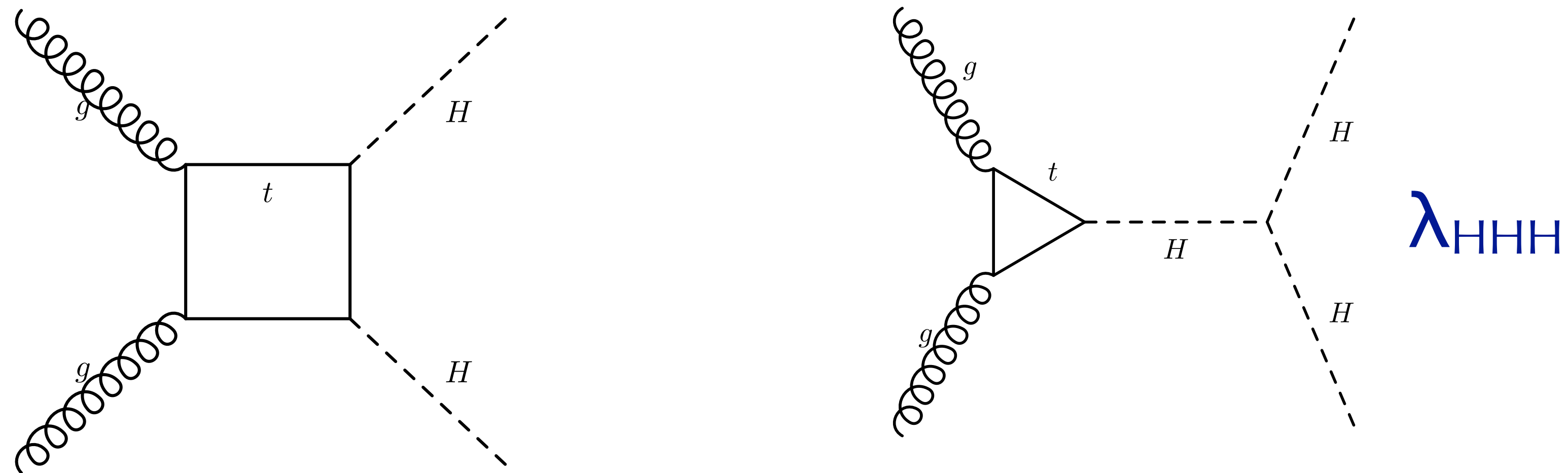
observed production and decay rates are consistent with a SM Higgs boson within uncertainties

- 1st indication of the $H(b\bar{b})$ decay at LHC (2.6σ)
- largest BR for SM H ($\sim 58\%$), dominates total width



Why di-Higgs

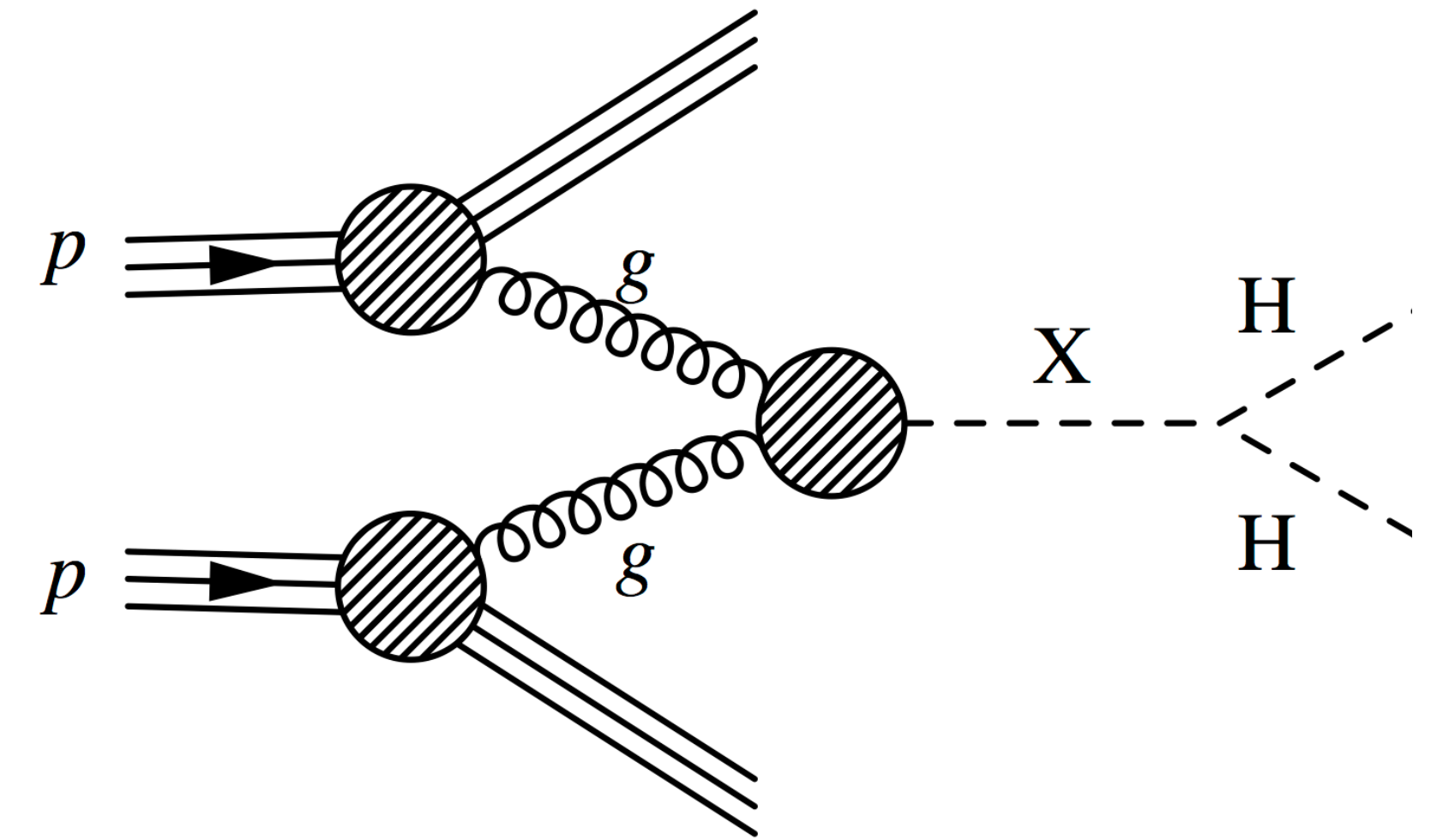
The measurement of the Higgs boson **self coupling** is a **fundamental** test of the SM
SM predicts a **extremely small cross section** for HH production (~ 30 fb at 13 TeV)
1000 times smaller than the single Higgs boson production



But, it is significantly enhanced in many BSM scenarios

Why di-Higgs

- Gluon fusion production of a massive X - **resonant HH** state
 - **small natural width**
- Depending on the resonance mass different models can be probed:
 - The mass range around 300-500 GeV is interesting for **(N)MSSM**
 - From 500 GeV it is interesting for **warped extra dimensions** models
 - spin-0 Radion and spin-2 KK-Graviton



Higgs, as a new powerful tool to search for new physics

hh resonant production

$h(\gamma\gamma)$

simple topology

excellent mass resolution

clean final state for low mass resonances

Limited by small BR

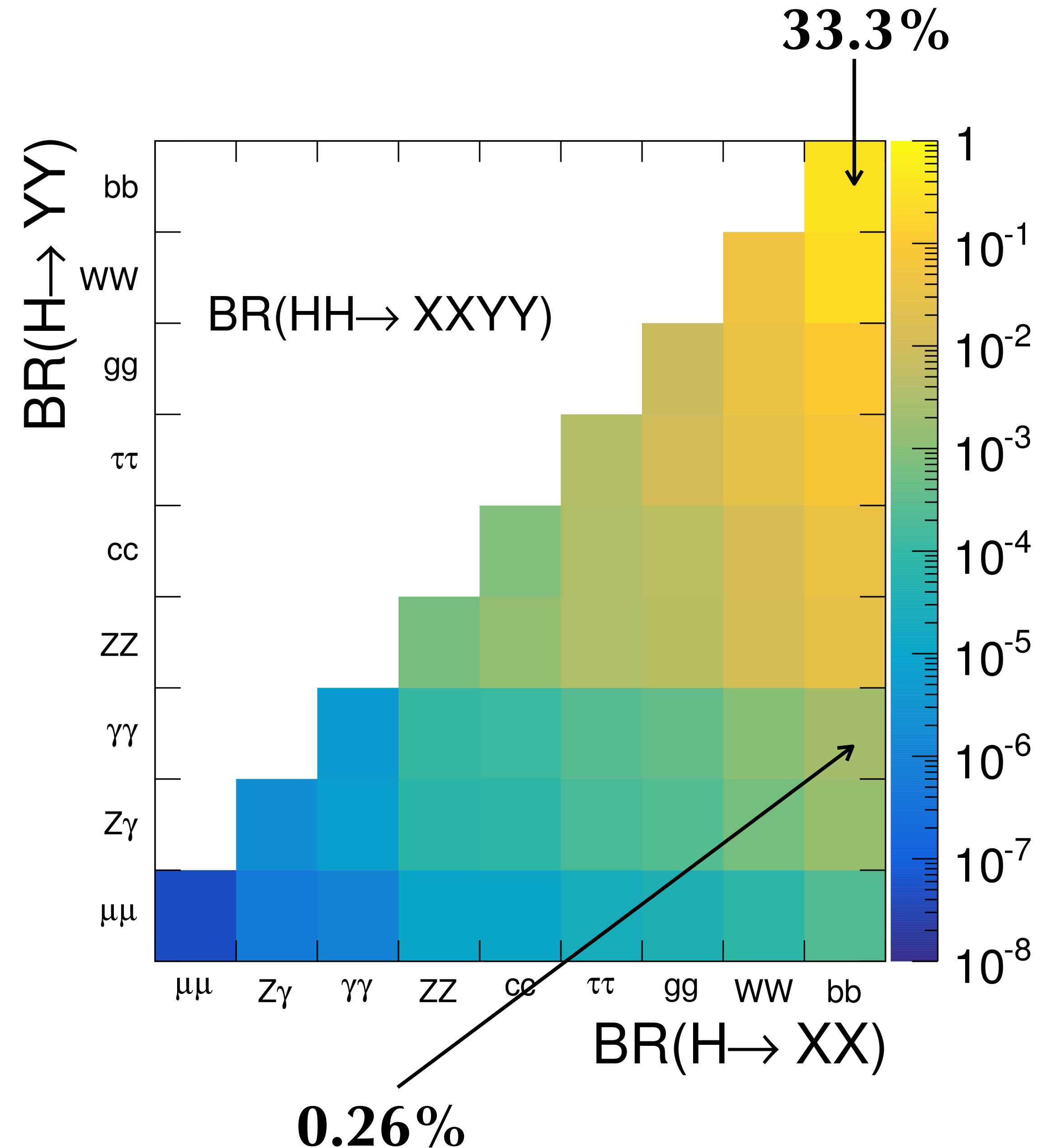
$h(b\bar{b})$

highest BR: larger statistics

high b-tag efficiency and low fake rate

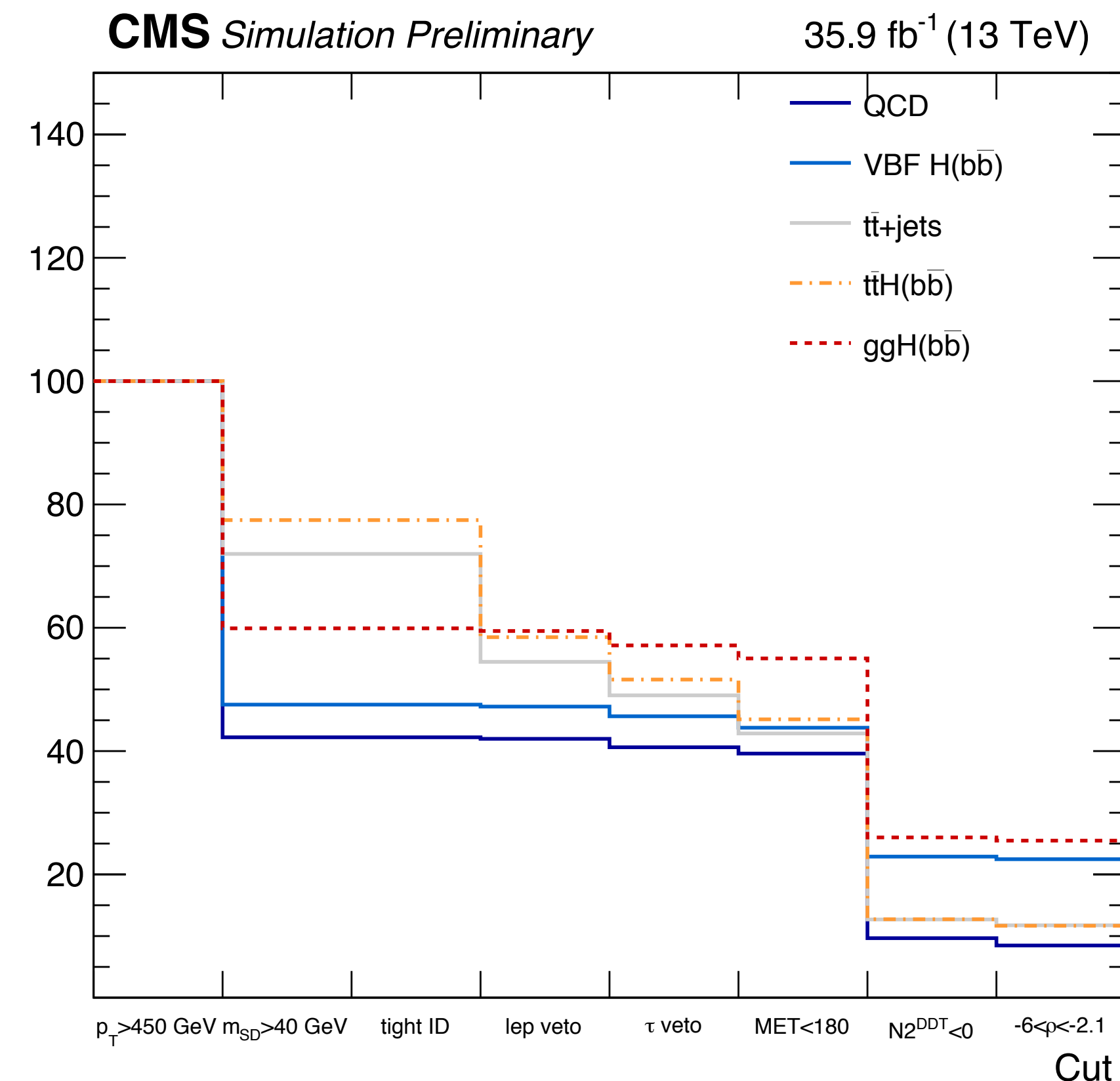
multi-light jets background is highly reduced

Identifying b-quarks plays a critical role



Event Selection

- Highest p_T AK8 jet with Puppi inputs:
 - $p_T > 450$ GeV $|\eta| < 2.5$, tight PF jet ID
 - jet soft-drop mass > 40 GeV ($z = 0.1$, $\beta = 0$)
 - **Substructure**: two prongs discrimination
 - N_2^{DDT} @ 26% background efficiency,
 - **double-b tagger** > 0.9
 - @30% sig. efficiency; 1% bkg efficiency
- **Lepton veto**
- **tt+jets rejection**:
 - Puppi MET < 180 GeV
- **$\rho = \log(m_{SD}^2/p_T^2)$ range $-6.0 < \rho < -2.1$.**
 To avoid instabilities at the edges of the distribution, due to finite cone effects from the AK8 jet clustering (around $\rho \sim -2$), and to avoid the non-perturbative regime of the soft drop mass calculation (below $\rho \sim -6$).
 - At $p_T = 450$ GeV, $\rho = -2.1$ corresponds to $m_{SD} \sim 166$ GeV;
 - $\rho = -6$ corresponds to $m_{SD} \sim 22$ GeV;

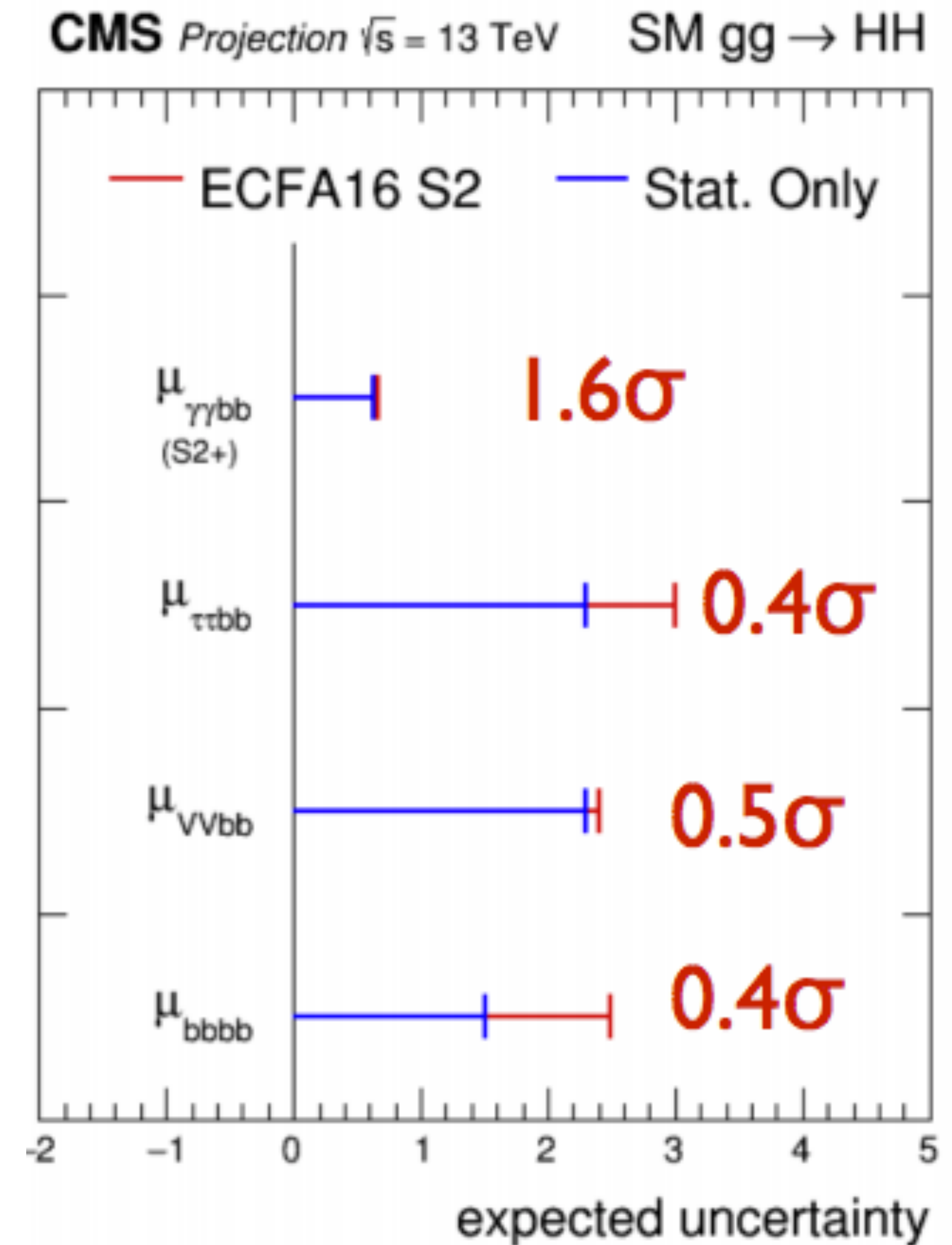


Systematics Uncertainties

Systematic uncertainty source	Type (shape/normalization)	Relative size (or description)
QCD transfer factor	both	float polynomial coefficients $a_{k\ell}$ and QCD normalization
Luminosity	normalization	2.6%
V-tag ($N_2^{1,DDT}$) efficiency	normalization	4.3%
Muon veto efficiency	normalization	0.5%
Electron veto efficiency	normalization	0.5%
Trigger efficiency	normalization	4%
Muon ID efficiency	shape	up to 0.3%
Muon isolation efficiency	shape	up to 0.2%
Muon trigger efficiency	shape	up to 15%
$t\bar{t}$ normalization SF	normalization	from 1μ CR: 8%
$t\bar{t}$ double-b mis-tag SF	normalization	from 1μ CR: 15%
W/Z NLO QCD corrections	normalization	10%
W/Z NLO electroweak corrections per p_T category	normalization	15% – 35%
W/Z NLO electroweak ratio decorrelation per p_T category	normalization	5% – 15%
Double-b tagging efficiency	normalization	4%
Jet energy scale	normalization	up to 10%
Jet energy resolution	normalization	up to 15%
Jet mass scale	shape	shift m_{SD} peak by $\pm 0.4\%$
Jet mass resolution	shape	smear m_{SD} distribution by $\pm 9\%$
Jet mass scale p_T	normalization	0.4%/100 GeV (p_T)
Monte Carlo statistics	normalization	-
H p_T corr (ggH)	normalization	30%

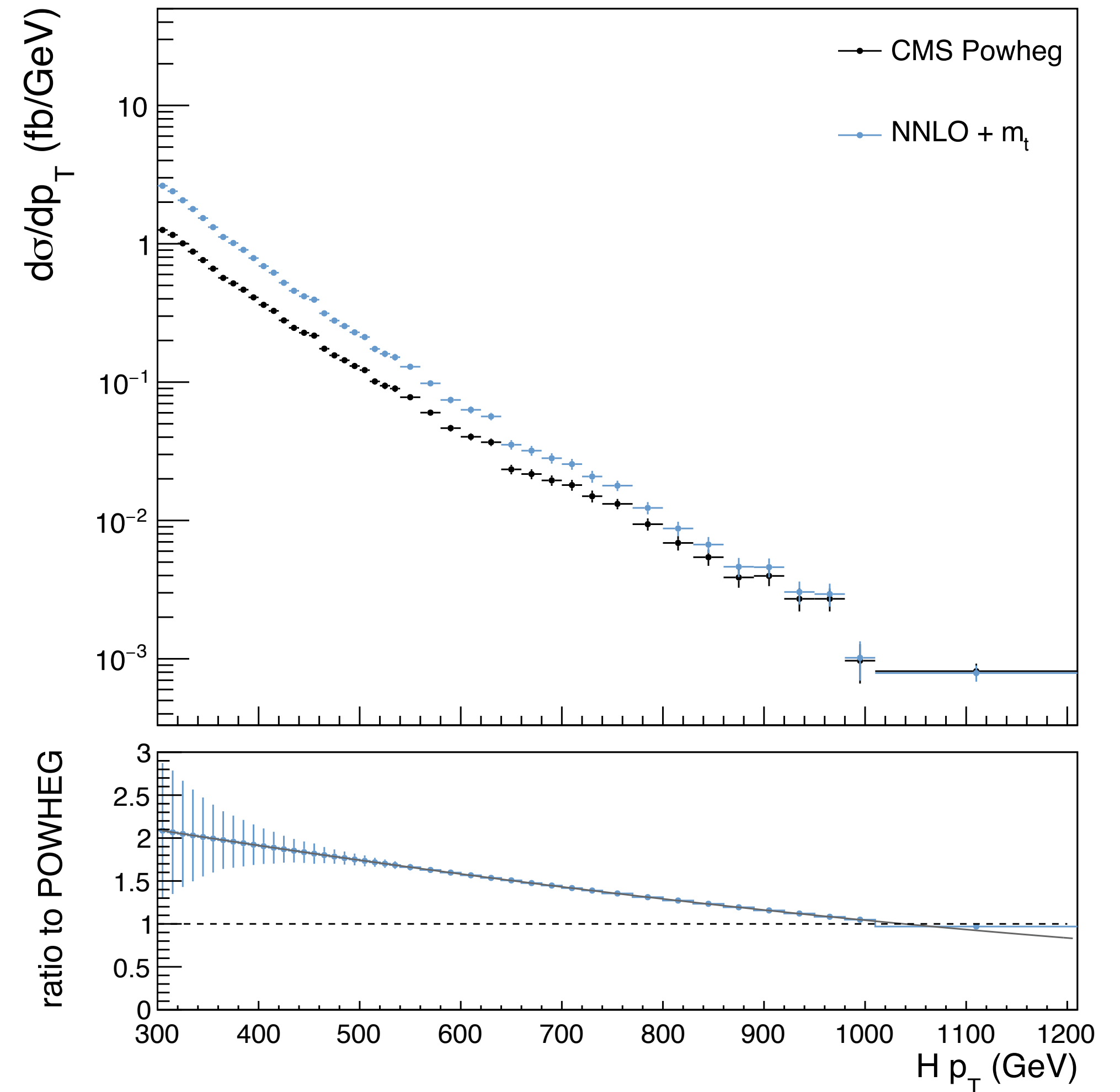
Perspectives

- Extrapolation from Run II to HL-LHC (3000 fb⁻¹)
SM HH production
 - but also analyses improve fast



ggH Higgs p_T reweighting finite m_t + NNLO

- CMS default for ggH is Powheg (*)
- To account for both the effects of higher order corrections and for the finite top mass loop a multi-correction approach is adopted
- Two separate effects: finite top mass corrects **down**, N(N)LO corrects **up**
- Estimate k-factor of ~ 1.6 for Higgs $p_T > 450$ GeV wrt the Powheg sample
- We adopt a 30% uncertainty following the addition in quadrature of the NNLO and NLO* uncertainties.



(*) is generated with Higgs matrix elements up to 1 jet assuming the infinite top mass approximation ($m_{top} \rightarrow \infty$)

Transfer factor for the QCD prediction

- If the double-b tag discriminator value were completely uncorrelated from the jet p_T and m_{SD} , the transfer factor would be equal to unity.
- To account for deviations from unity, we Taylor expand F in ρ and p_T

$$F(\rho, p_T) = ((1 + a_{01} p_T + a_{02} p_T^2 + \dots) + (a_{10} + a_{11} p_T + a_{12} p_T^2 + \dots)\rho + (a_{20} + a_{21} p_T + a_{22} p_T^2 + \dots)\rho^2 + \dots)$$

We performed the F-test on the blinded, full 2016 dataset to decide the degree of the polynomial form:

- $(n_\rho = 2, n_{p_T} = 1)$ polynomial is the default

