



# **QCD Phase Transitions and Baryon Number Fluctuations within FRG**

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**Dalian University of Technology**

EMMI Workshop on Critical Fluctuations Near the QCD Phase Boundary in  
Relativistic Nuclear Collisions  
October 10-13, 2017, Wuhan

Collaborators:

fQCD collaboration (J. Braun, L. Corell, A. Cyrol, WF, M. Leonhardt, M. Mitter,  
J.M. Pawlowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink)

Yu-xin Liu (Peking U.), Yue-liang Wu (ITP).....

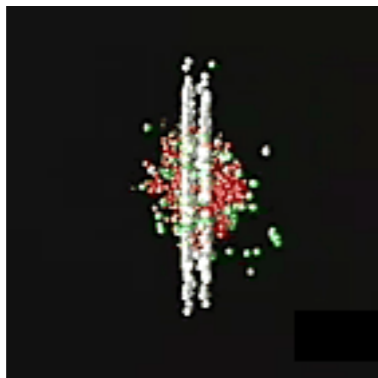
# Outline

- \* **Introduction**
- \* **Fluctuations of Conserved Charges with Mean Field Approximations (MFA)**
- \* **Quantum Fluctuations beyond MFA**
- \* **Preliminary Results from QCD within FRG**
- \* **Summary and outlook**

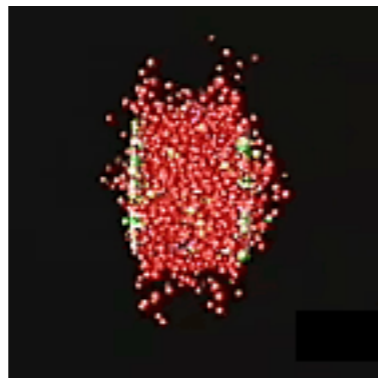
# Heavy-ion collision



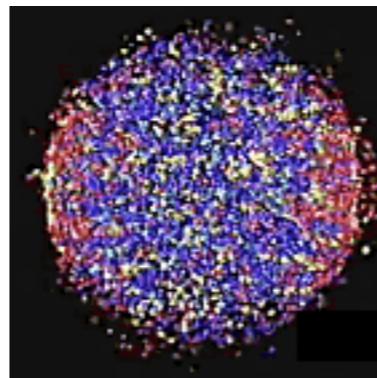
Ions about to collide



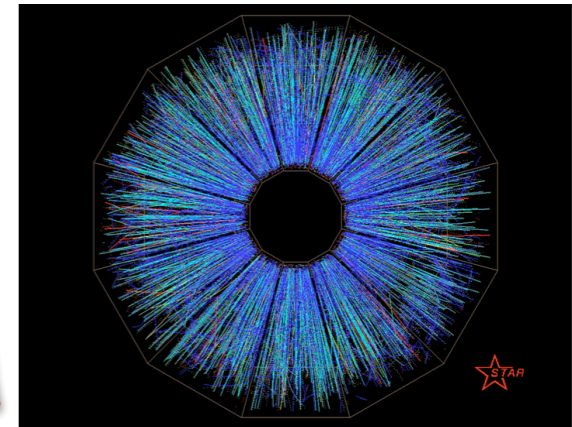
Ion collision



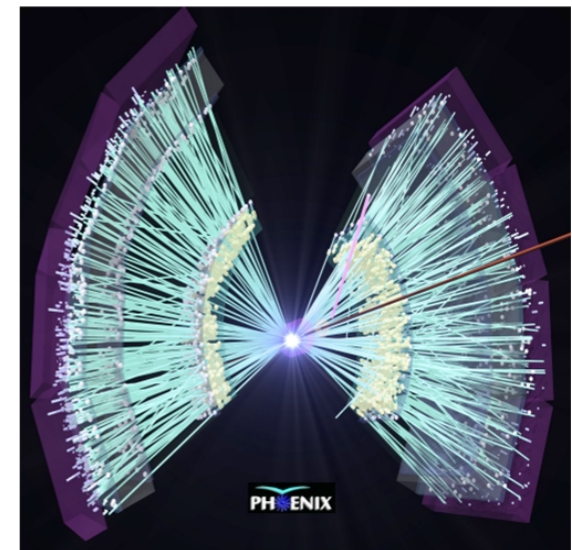
Plasma formation



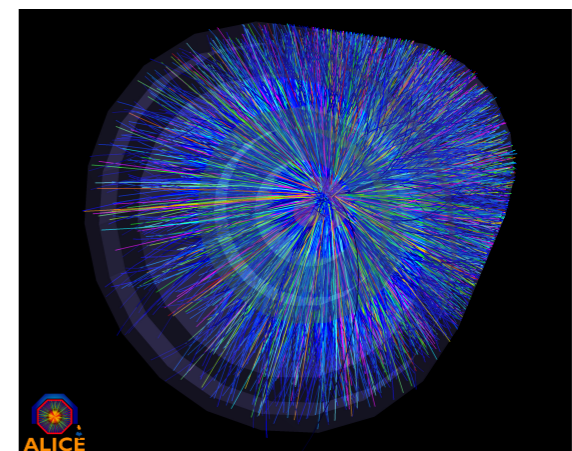
Freeze out



STAR

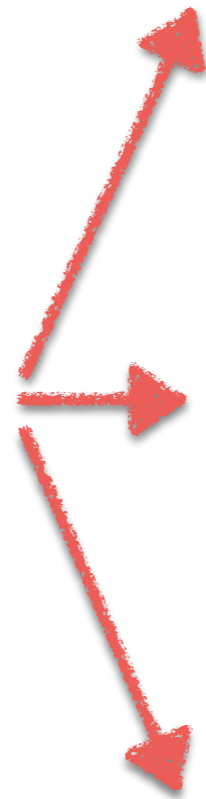


PHENIX

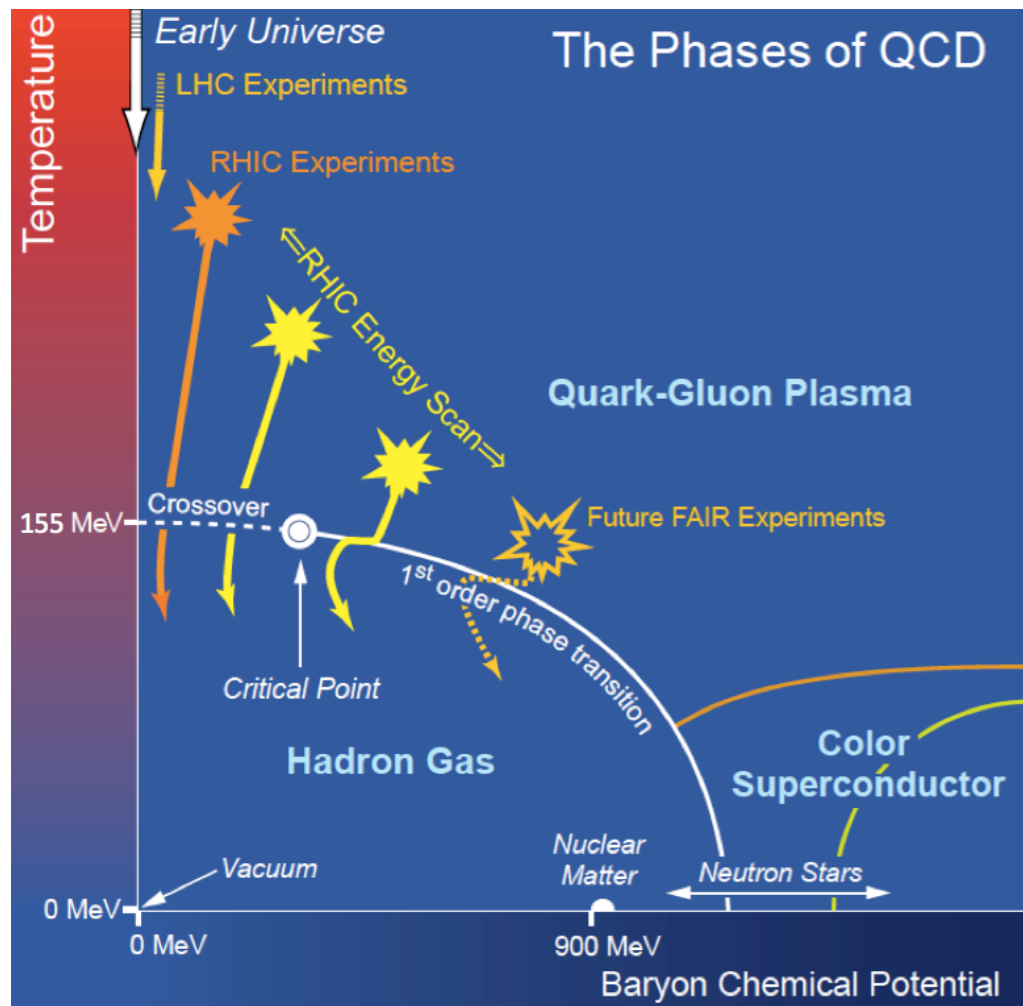


ALICE

What we see

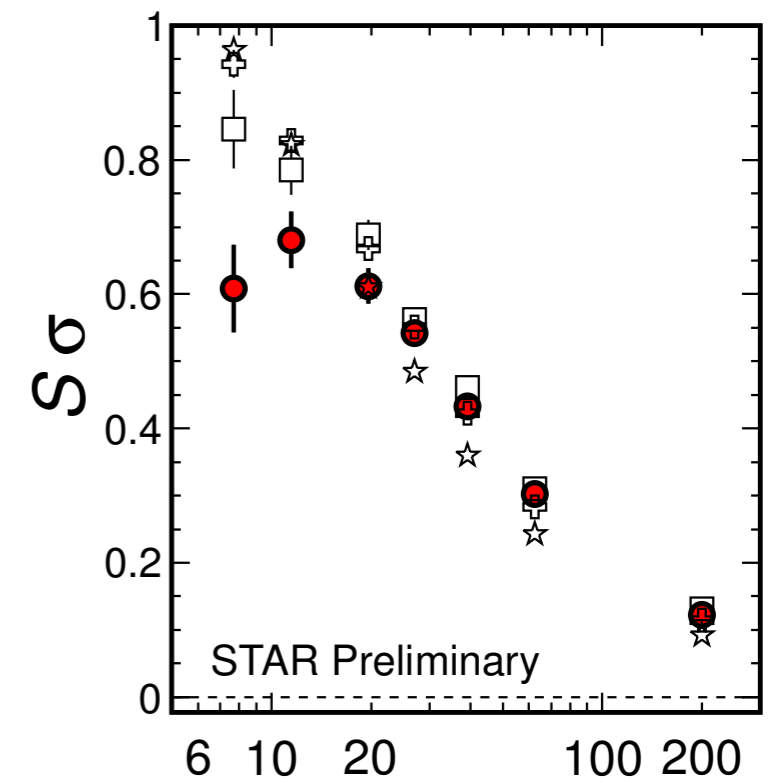
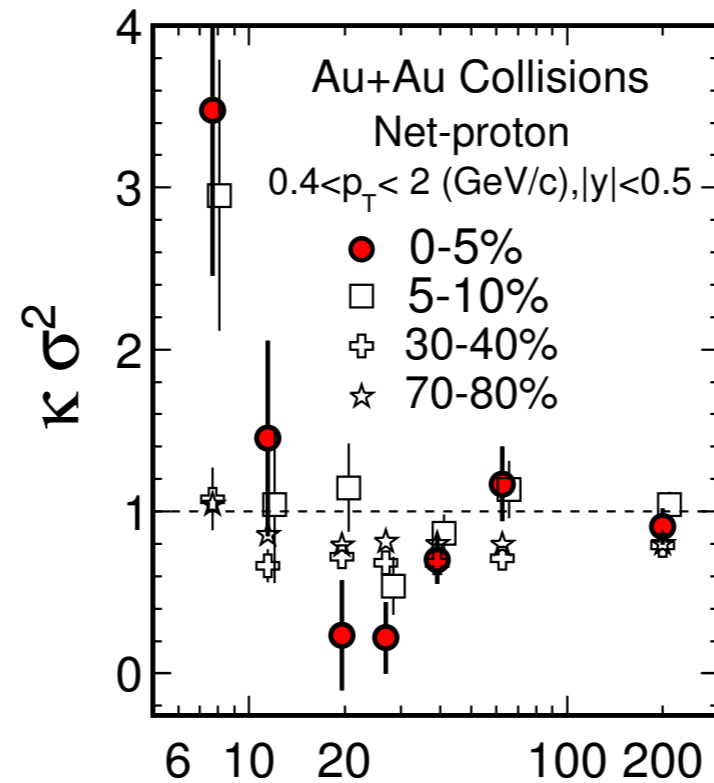


# QCD Phase Structure



The Hot QCD White Paper (2015)

RHIC:

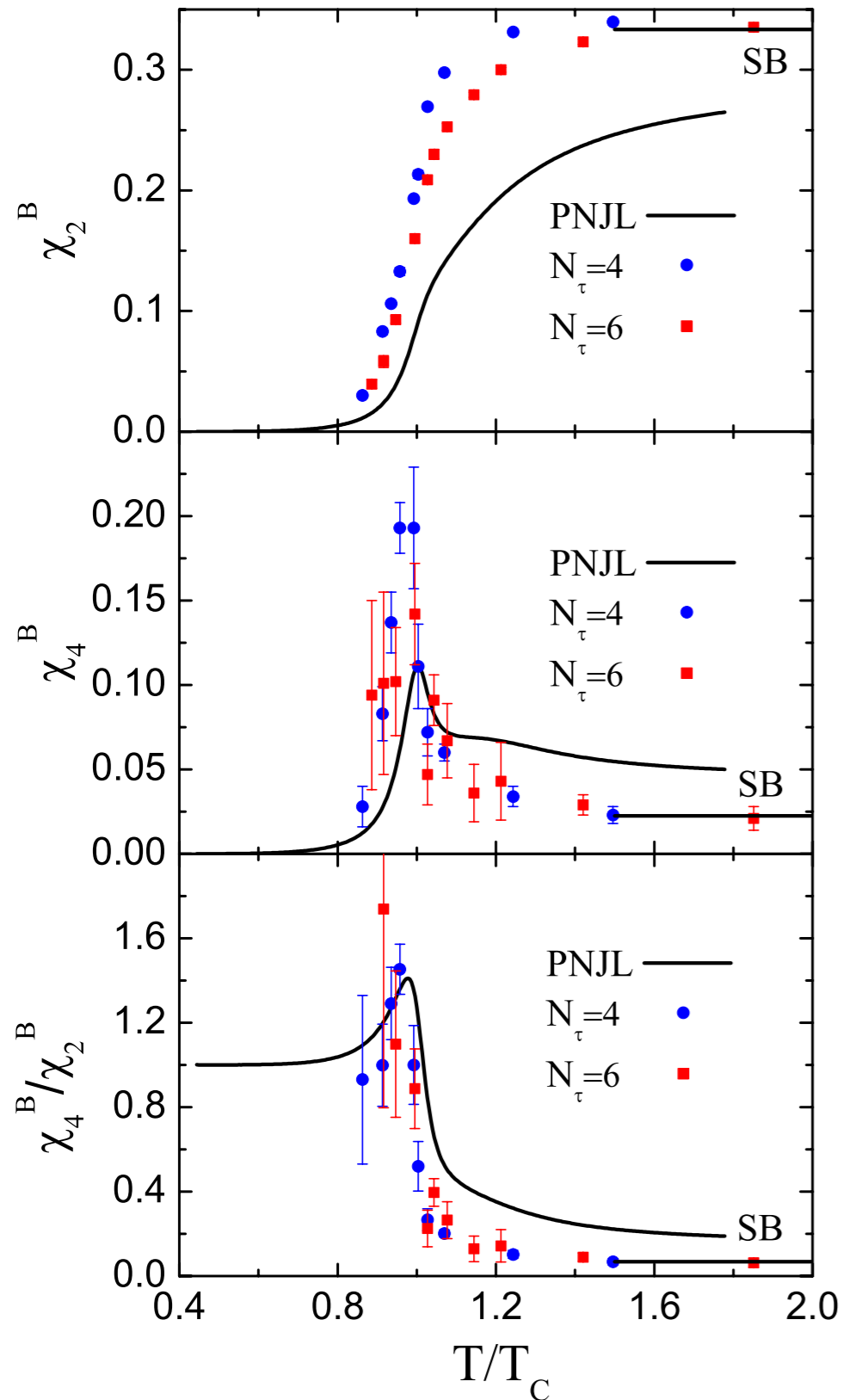


X.Luo (STAR), PoS CPOD2014, 019 (2014)

see talks by Xiao-feng Luo



# 2+1 Polyakov-loop NJL model



Model:

$$\begin{aligned} \mathcal{L}_{PNJL} = & \bar{\psi} (i\gamma_\mu D^\mu + \gamma_0 \hat{\mu} - \hat{m}_0) \psi \\ & + G \sum_{a=0}^8 \left[ (\bar{\psi} \tau_a \psi)^2 + (\bar{\psi} i\gamma_5 \tau_a \psi)^2 \right] \\ & - K \left[ \det_f (\bar{\psi} (1 + \gamma_5) \psi) + \det_f (\bar{\psi} (1 - \gamma_5) \psi) \right] \\ & - \mathcal{U}(\Phi, \Phi^*, T) \end{aligned}$$

Generalized susceptibilities :

$$\chi_{ijk}^{BQS} = \left. \frac{\partial^{i+j+k} (P/T^4)}{\partial(\mu_B/T)^i \partial(\mu_Q/T)^j \partial(\mu_S/T)^k} \right|_{\mu_B, Q, S=0}$$

$$\chi_2^X = \frac{1}{VT^3} \langle \delta N_X^2 \rangle$$

$$\chi_{11}^{XY} = \frac{1}{VT^3} \langle \delta N_X \delta N_Y \rangle$$

$$\chi_3^X = \frac{1}{VT^3} \langle \delta N_X^3 \rangle$$

$$\chi_{12}^{XY} = \frac{1}{VT^3} \langle \delta N_X \delta N_Y^2 \rangle$$

$$\chi_4^X = \frac{1}{VT^3} \left( \langle \delta N_X^4 \rangle - 3 \langle \delta N_X^2 \rangle^2 \right)$$

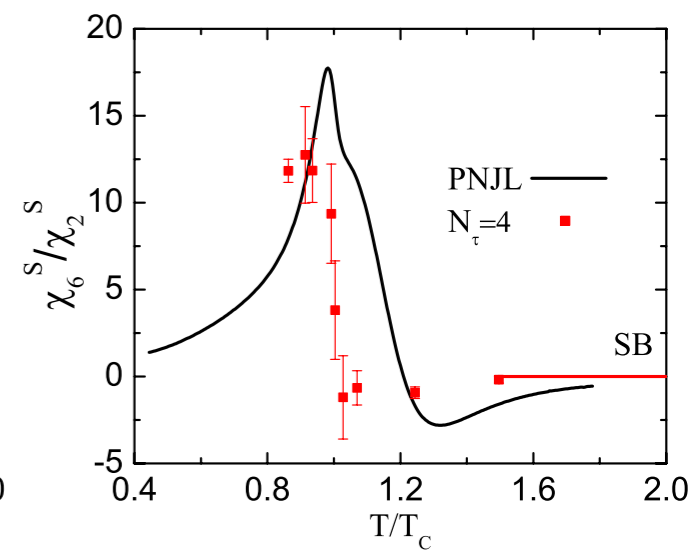
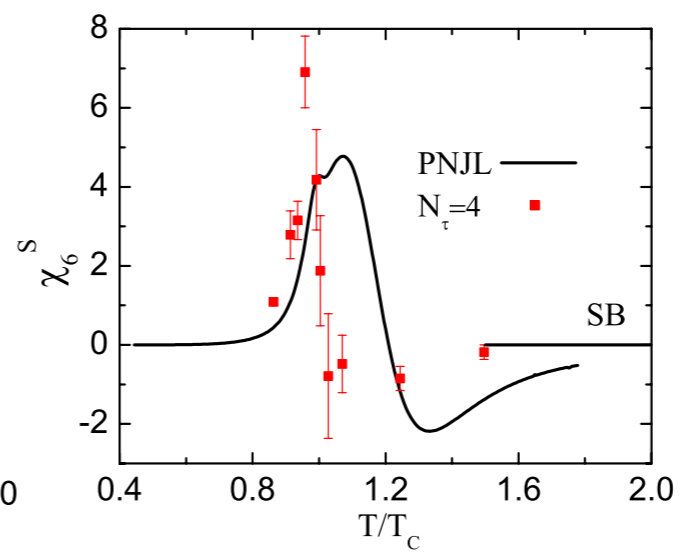
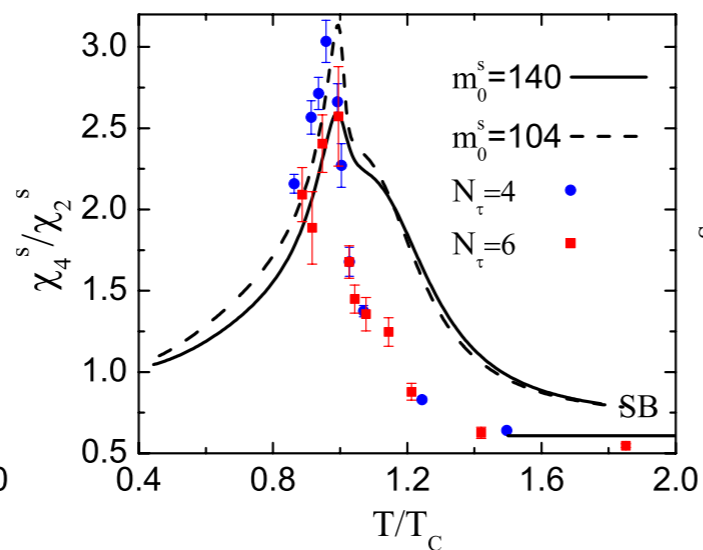
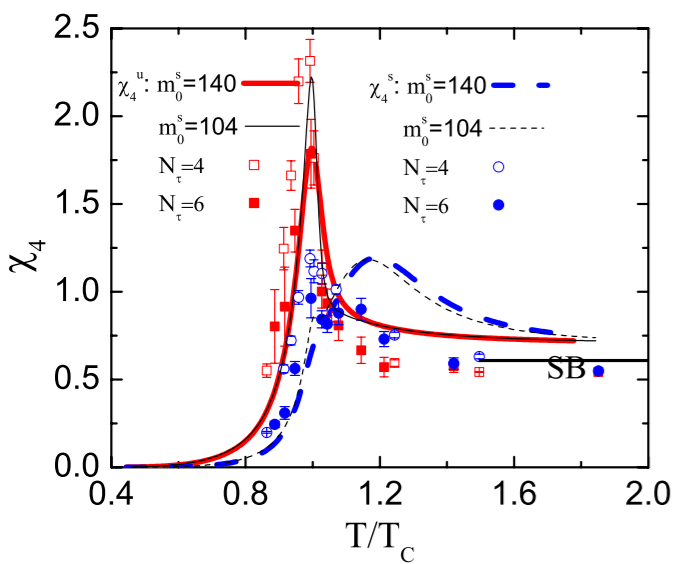
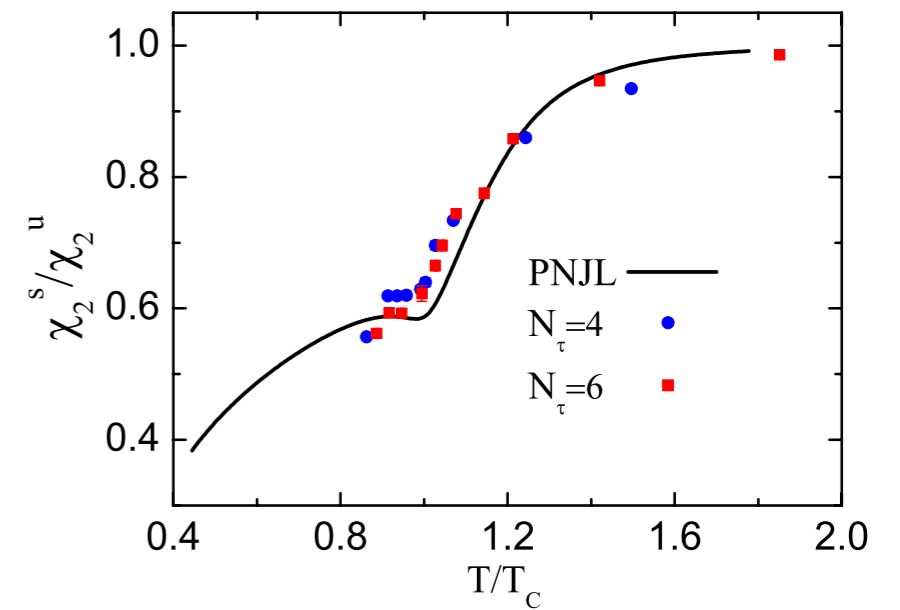
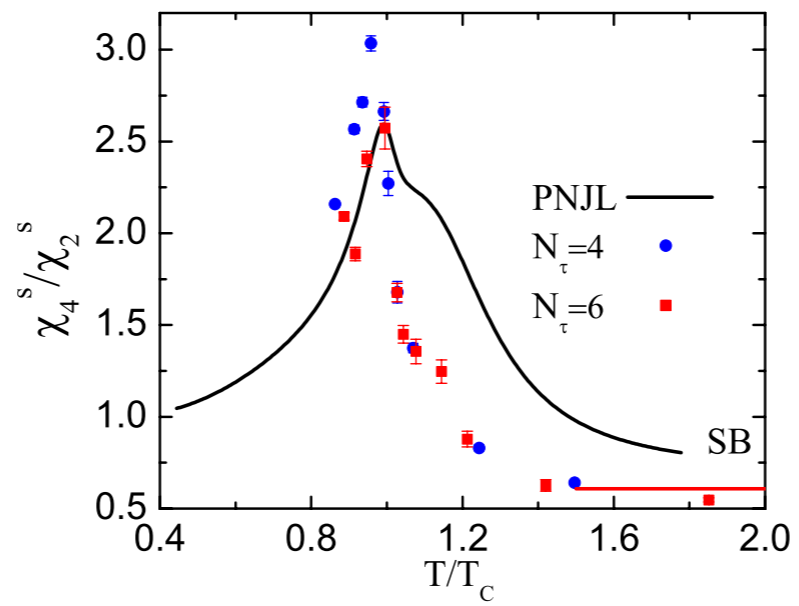
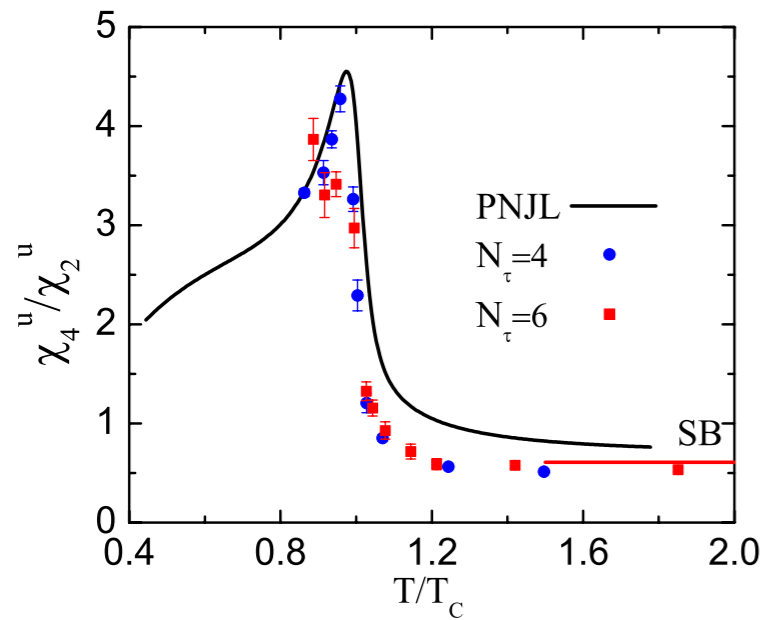
$$\chi_{111}^{XYZ} = \frac{1}{VT^3} \langle \delta N_X \delta N_Y \delta N_Z \rangle$$

$$\begin{aligned} \chi_6^X = & \frac{1}{VT^3} \left( \langle \delta N_X^6 \rangle - 15 \langle \delta N_X^4 \rangle \langle \delta N_X^2 \rangle \right. \\ & \left. - 10 \langle \delta N_X^3 \rangle^2 + 30 \langle \delta N_X^2 \rangle^3 \right) \end{aligned}$$

PNJL: WF, Yu-xin Liu, Yue-liang Wu, PRD 81 (2010) 014028

Lattice: M. Cheng *et al.*, PRD 79 (2009) 074505

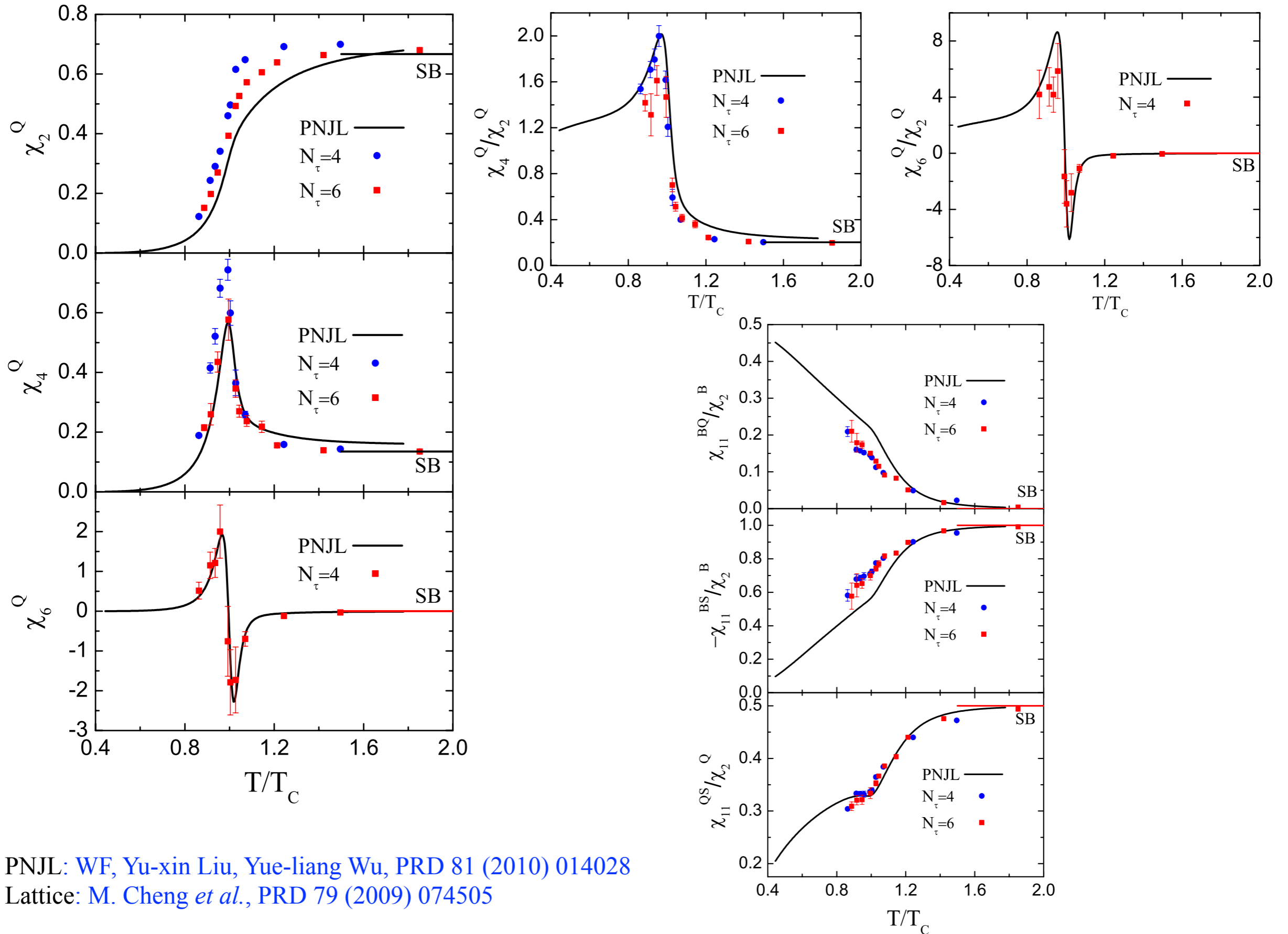
# Quark Number Fluctuations



PNJL: WF, Yu-xin Liu, Yue-liang Wu, PRD 81 (2010) 014028

Lattice: M. Cheng *et al.*, PRD 79 (2009) 074505

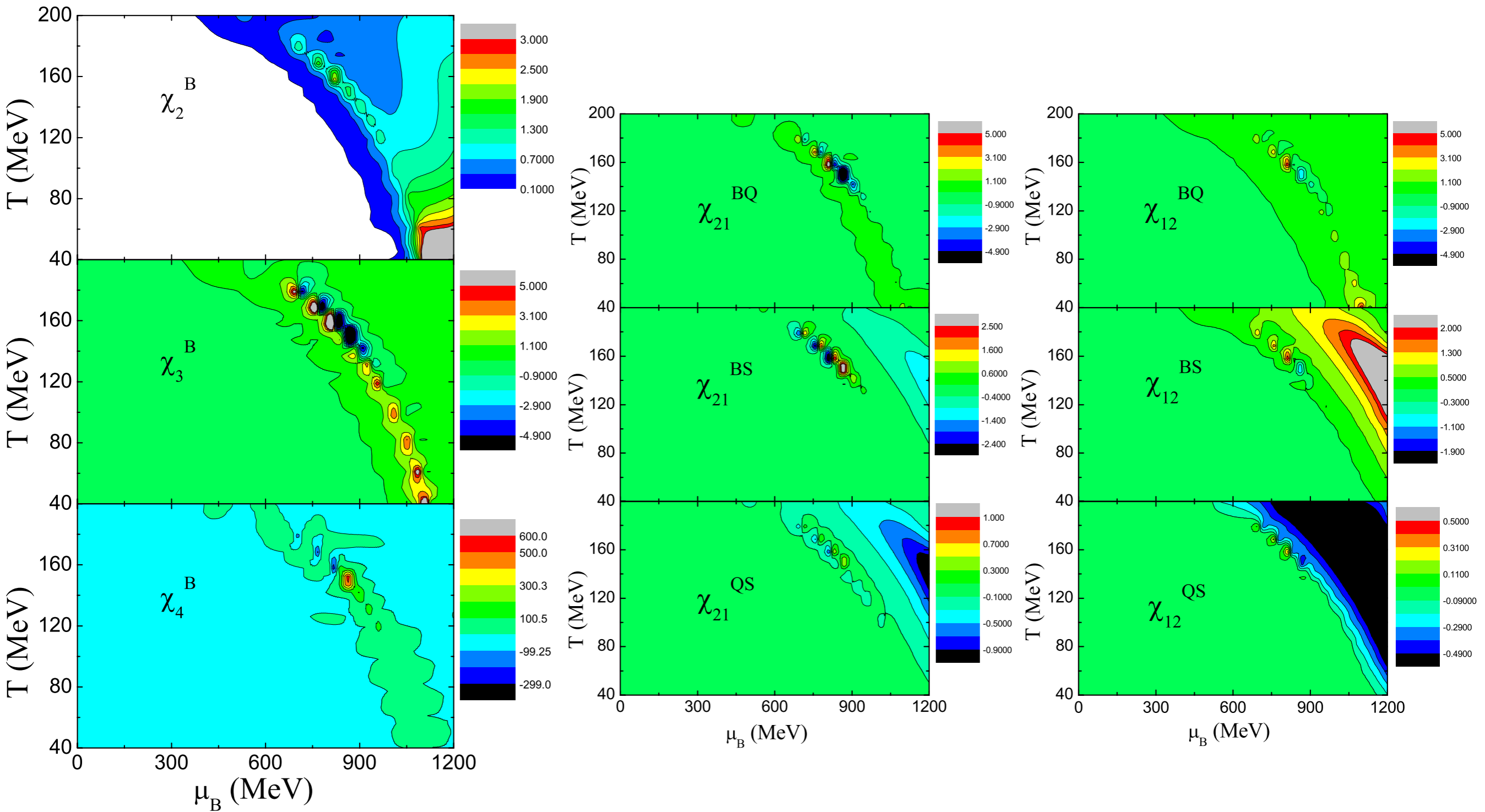
# Fluctuations of Electric Charge and Correlations



PNJL: WF, Yu-xin Liu, Yue-liang Wu, PRD 81 (2010) 014028

Lattice: M. Cheng *et al.*, PRD 79 (2009) 074505

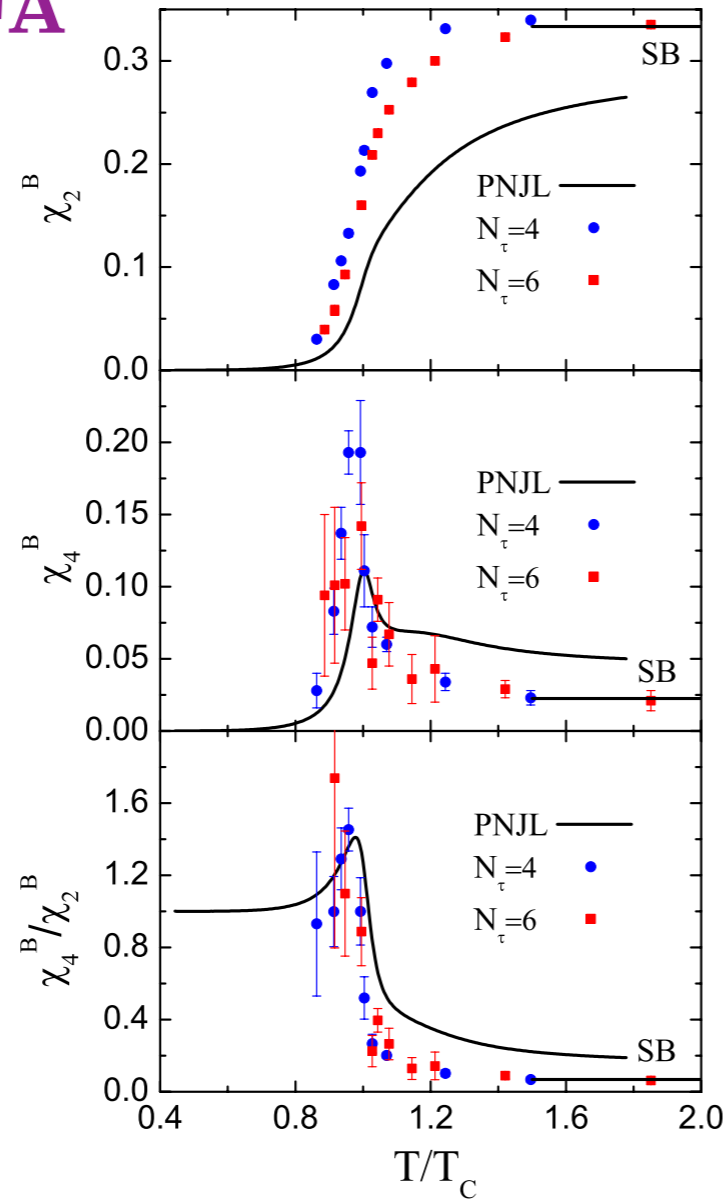
# Fluctuations on the Phase Diagram



# Quantum Fluctuations beyond MFA

## Summary on MFA

- Easy to calculate
- Agree with lattice results only qualitatively
- Overestimate the phase transition strength in the MFA
- Thus, one has to go beyond the MFA

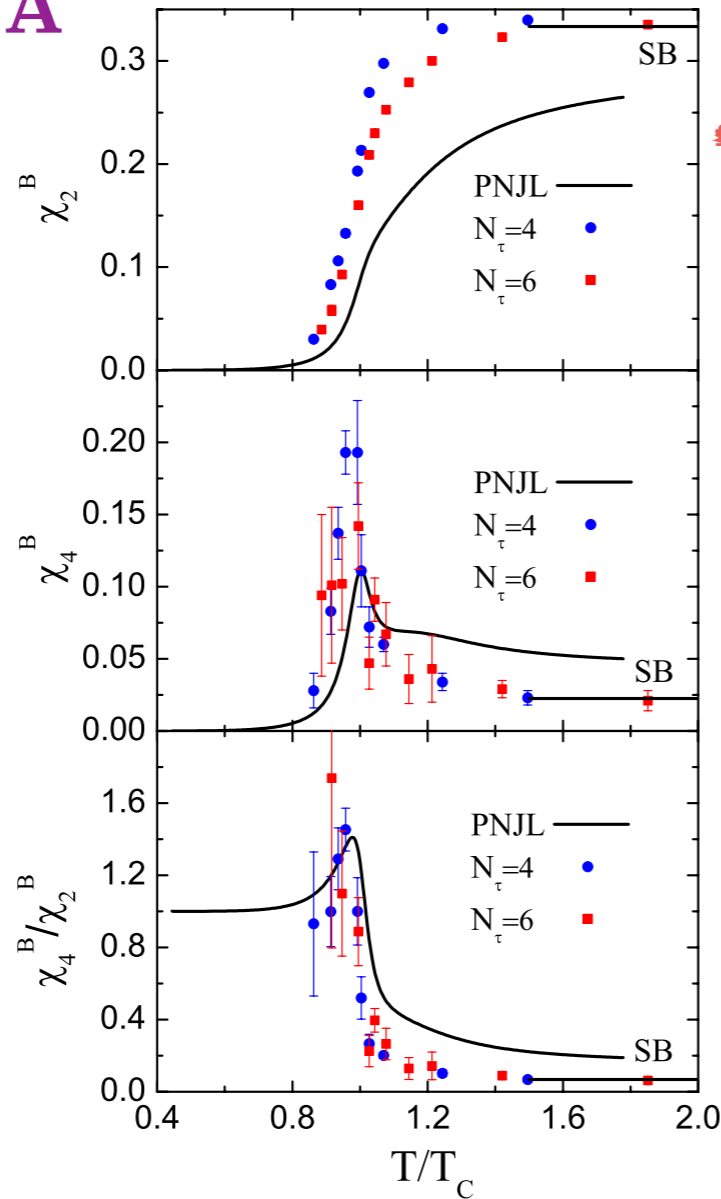




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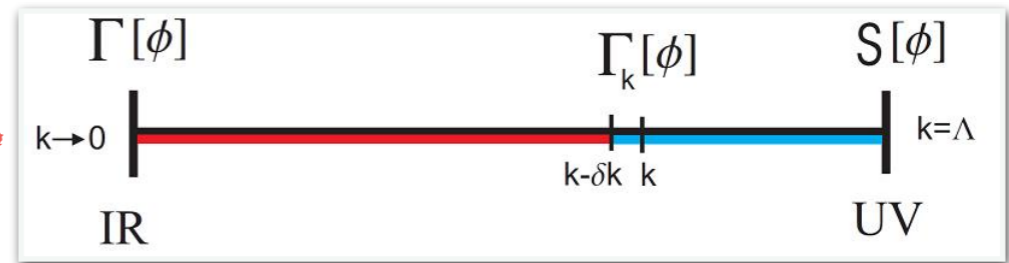
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## FRG

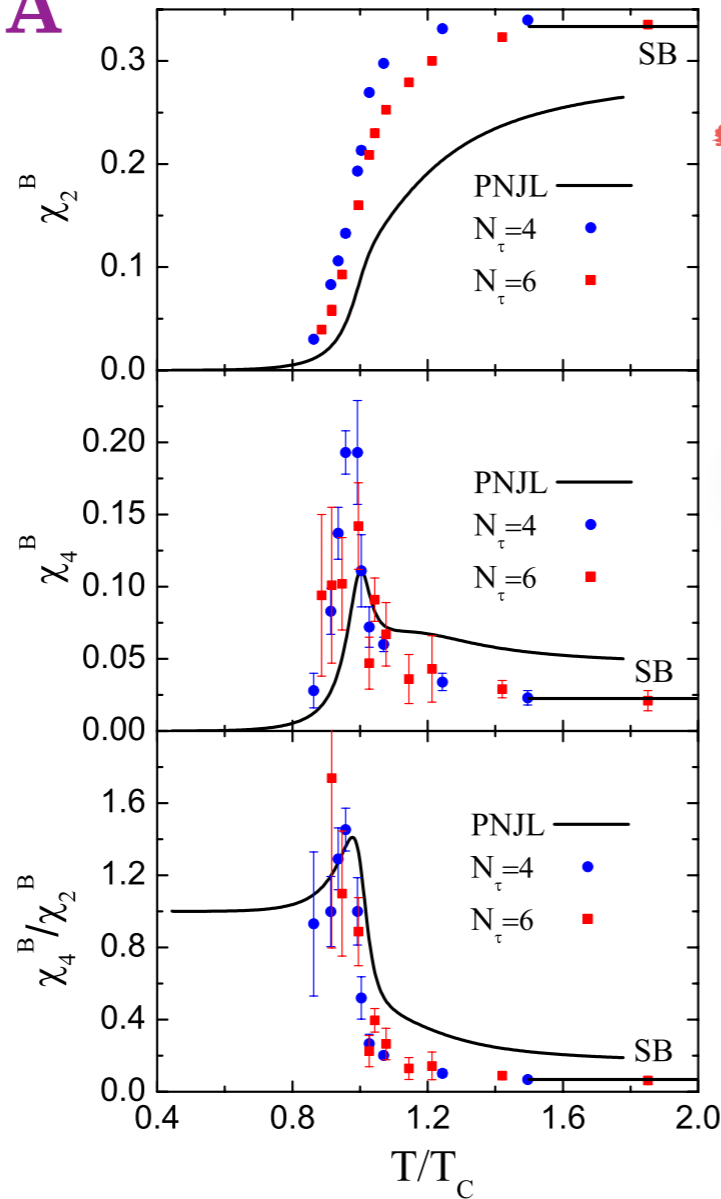
see talks by Jan, Defu



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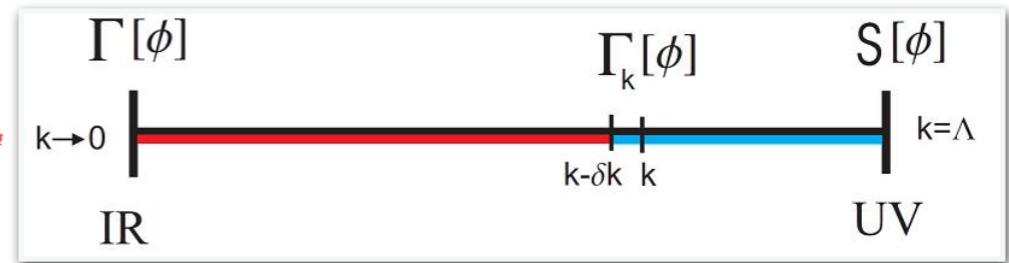
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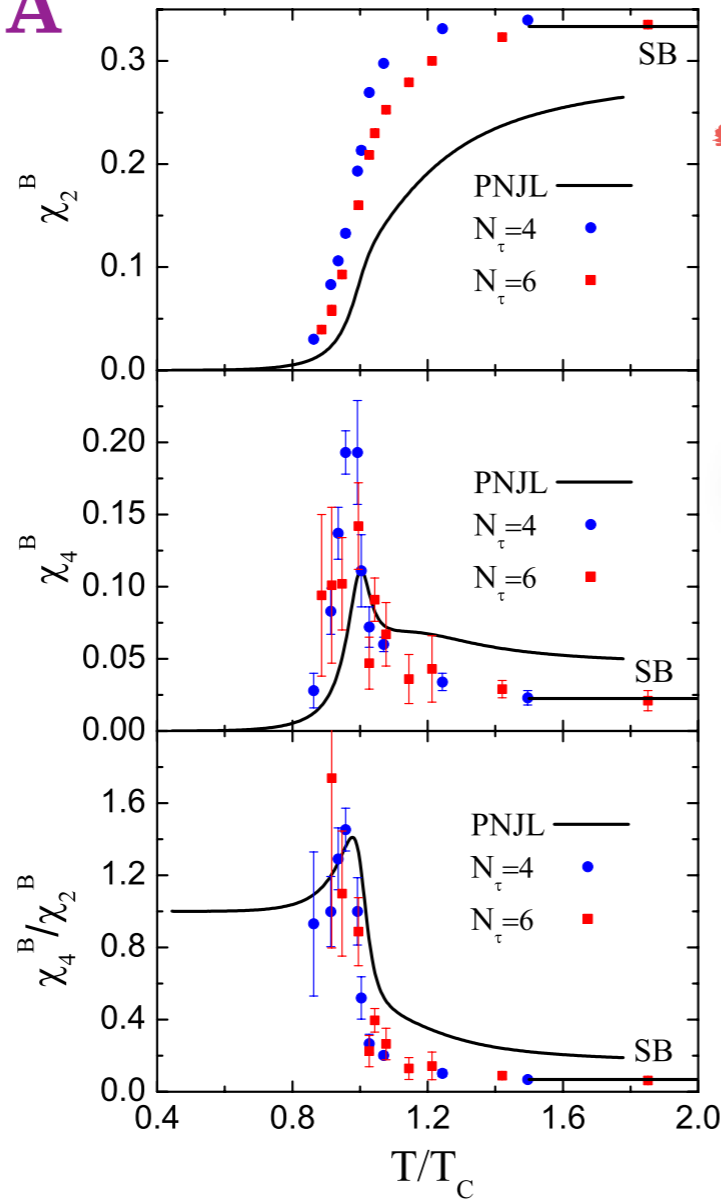


$$\partial_t \Gamma_k = \frac{1}{2} \left( \text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} + \frac{1}{2} \text{Diagram 4} \right)$$

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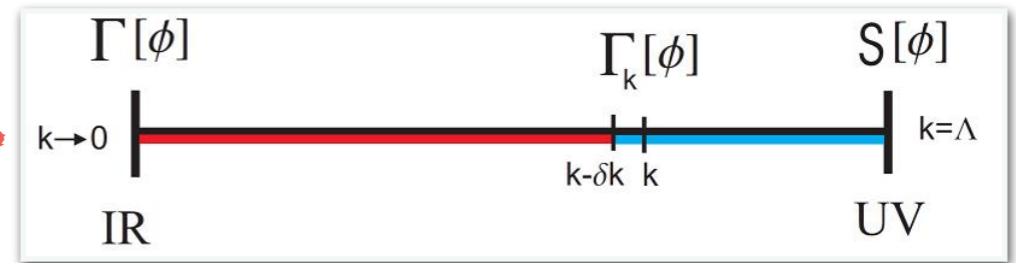
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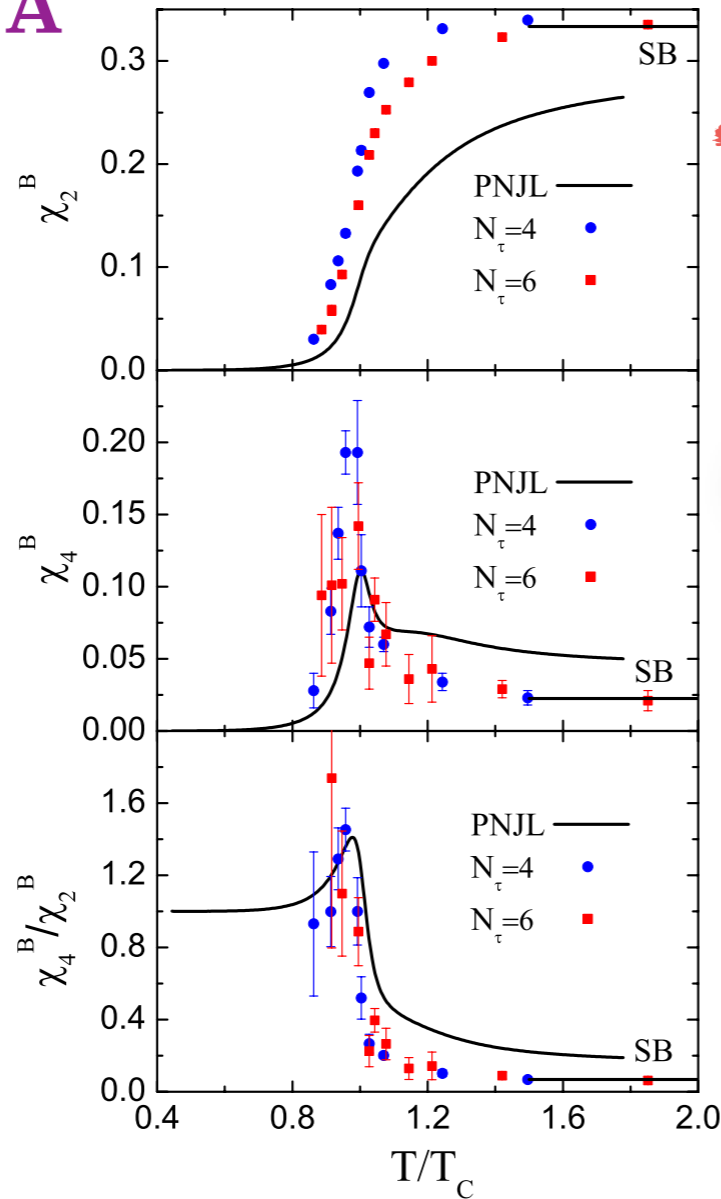
$$\partial_t \Gamma_k = \frac{1}{2} \left( \text{Diagram 1} - \text{Diagram 2} \right) - \text{Diagram 3} + \frac{1}{2} \text{Diagram 4}$$

glue sector

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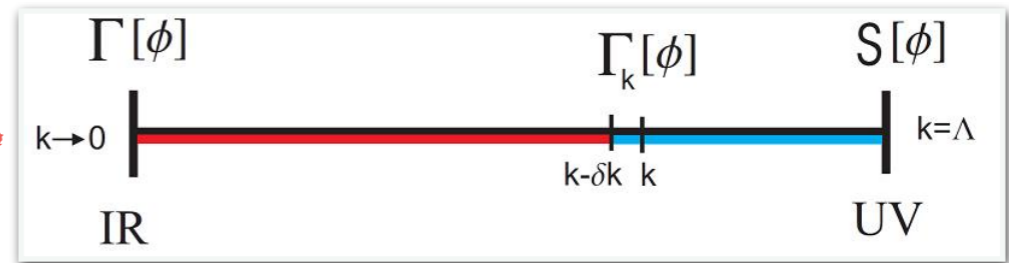
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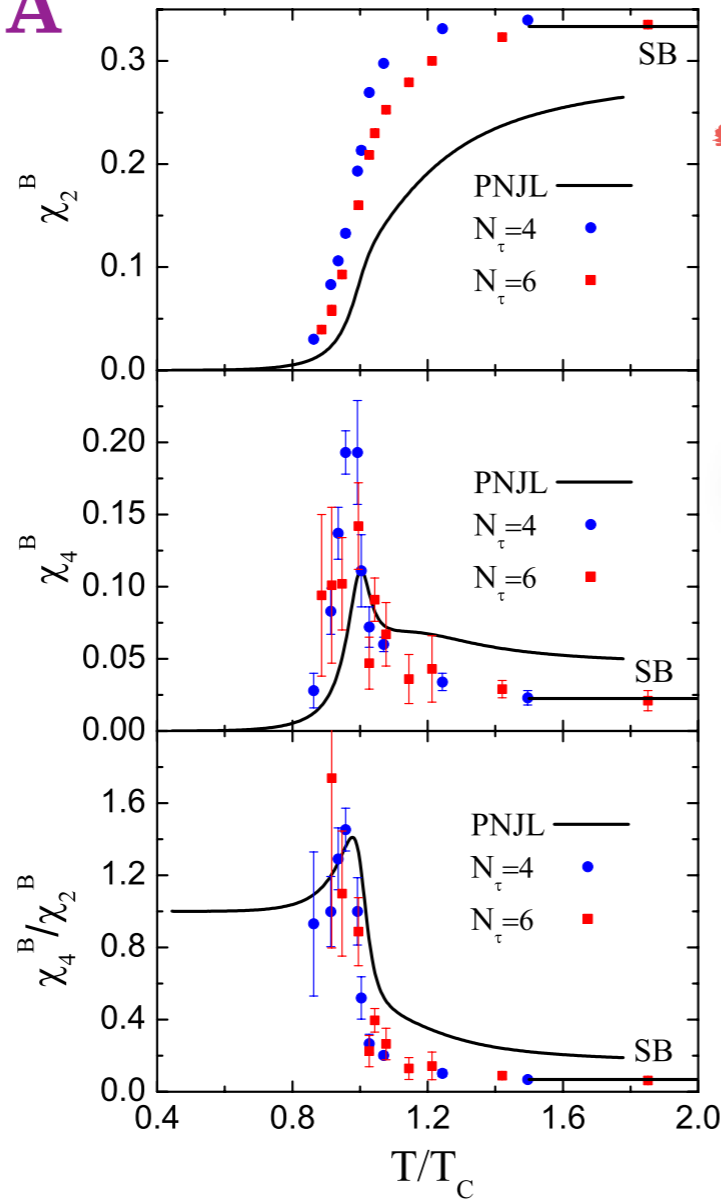


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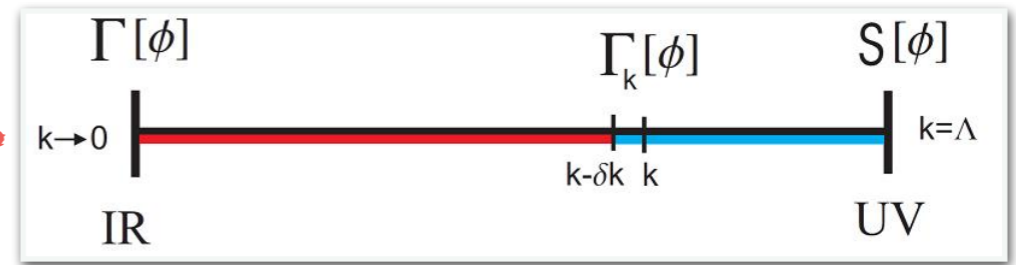
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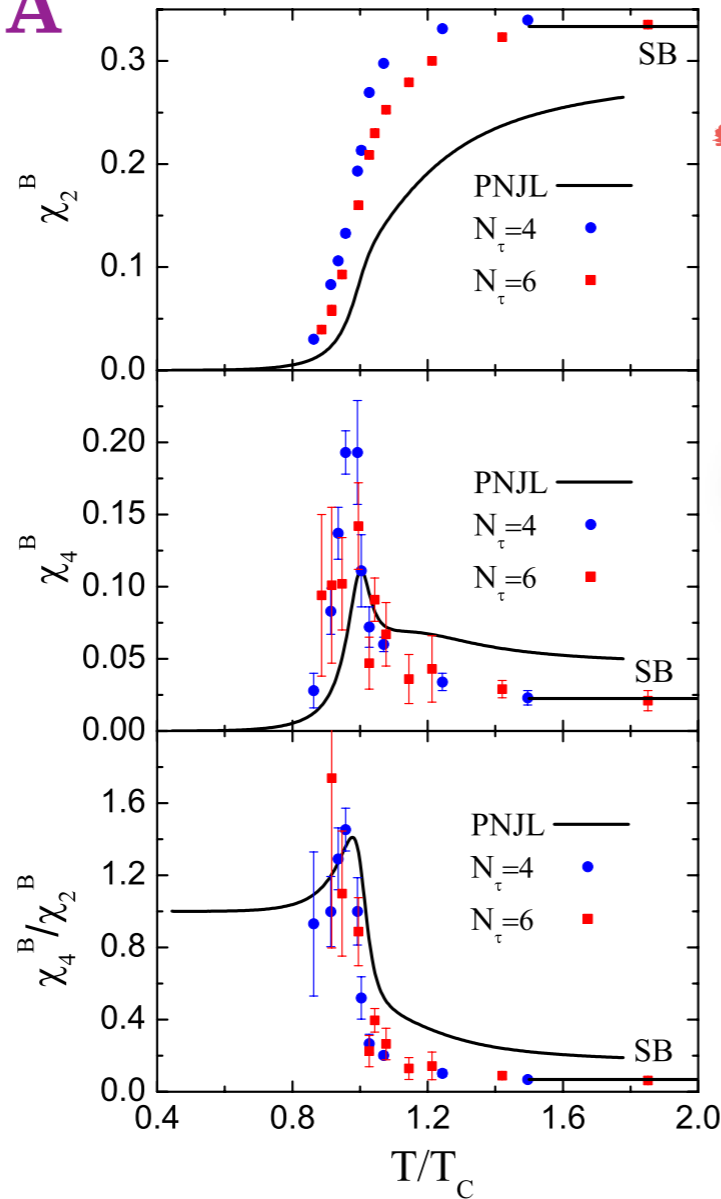
MFA



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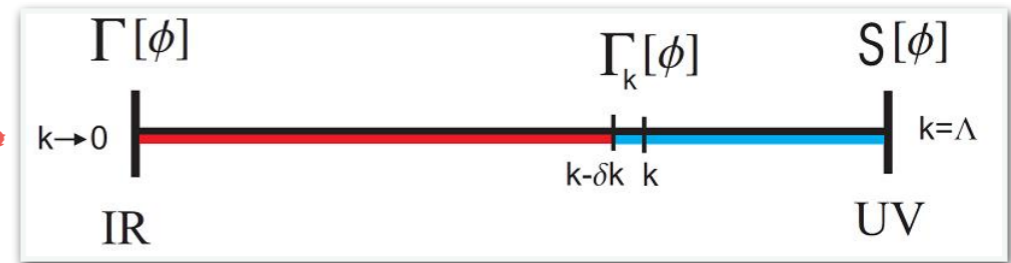
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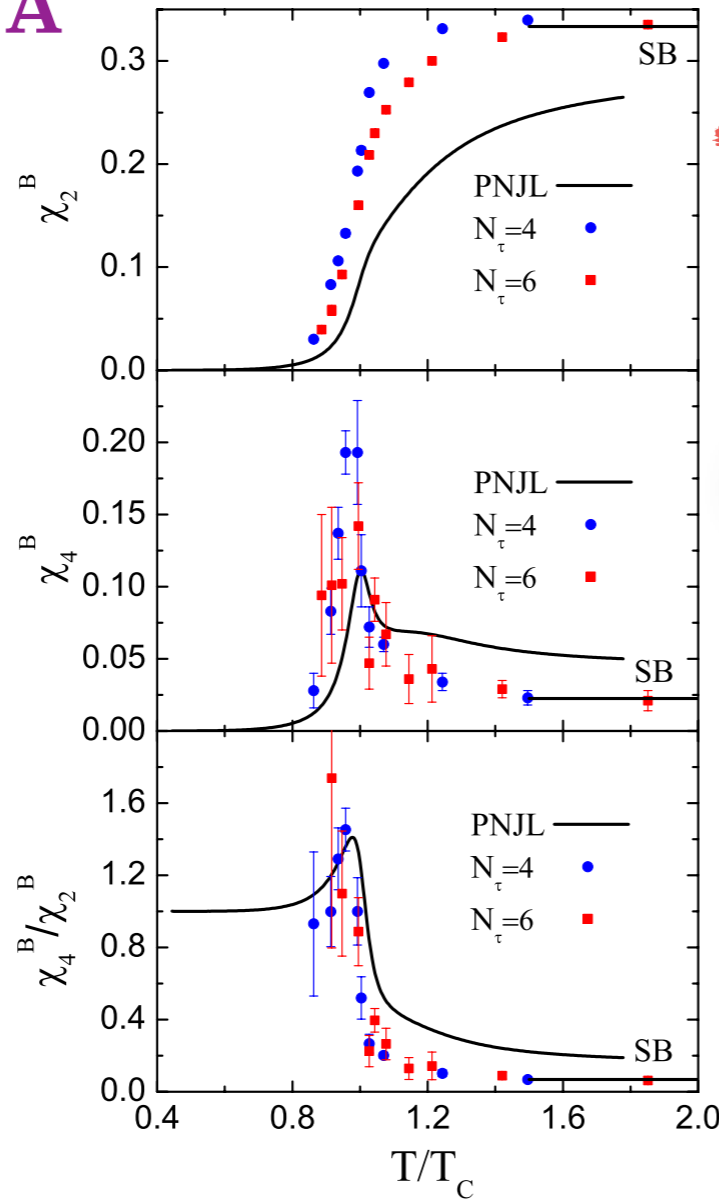


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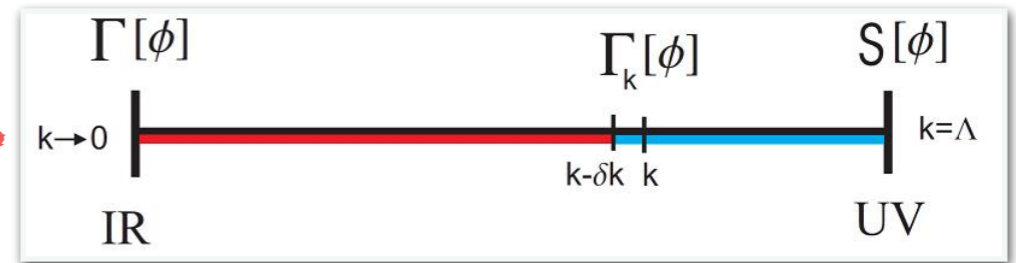
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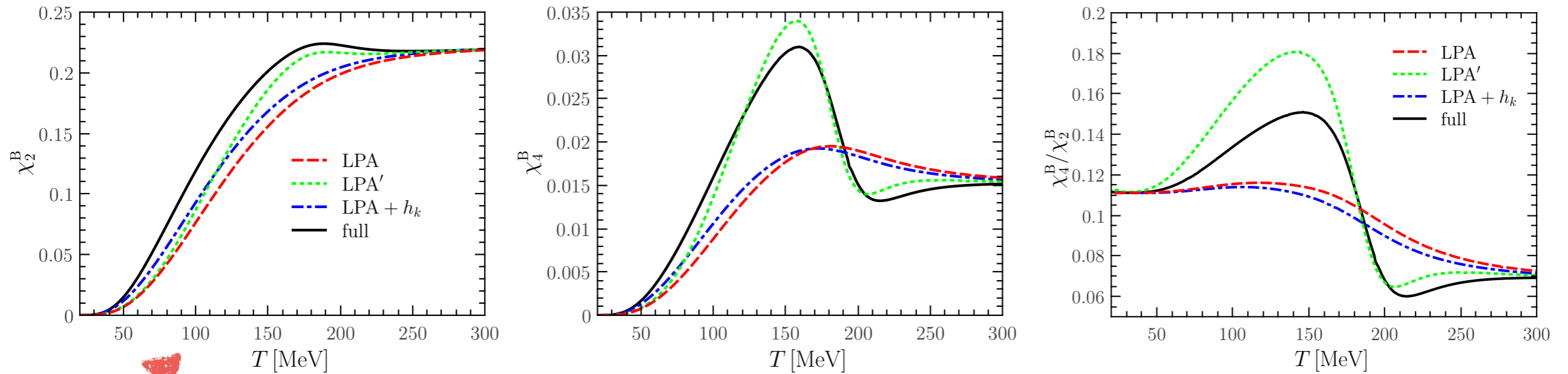


$$\partial_t \Gamma_k = \frac{1}{2} \left( \text{Loop 1} - \text{Loop 2} - \text{Loop 3} + \frac{1}{2} \text{Loop 4} \right)$$

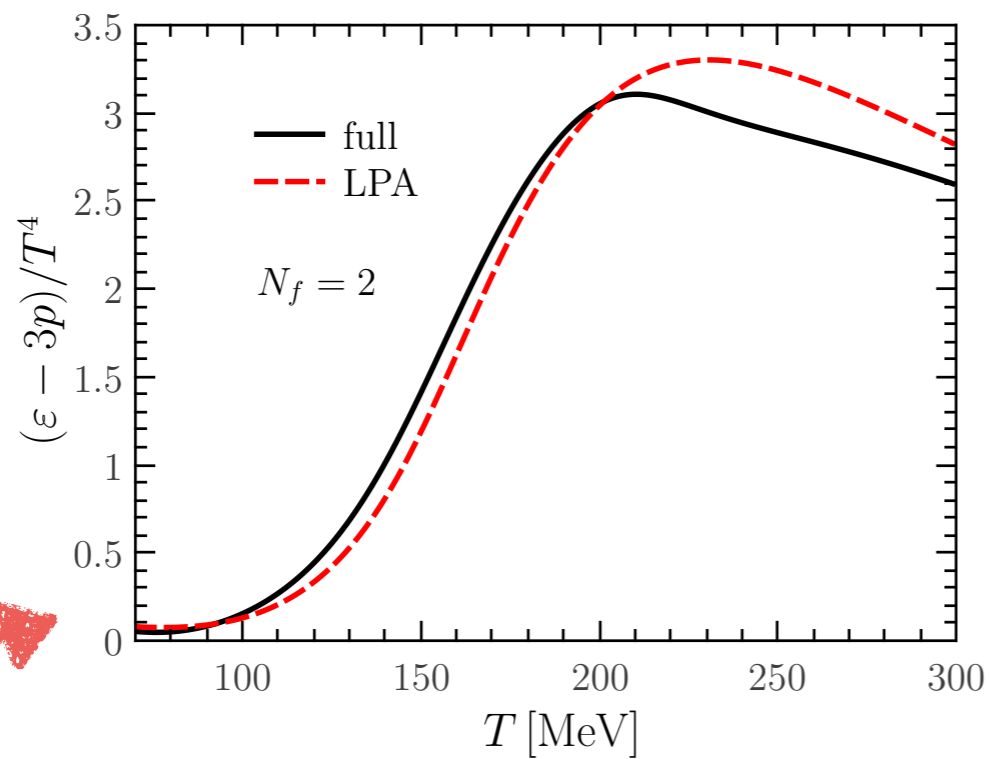
Matter part, the effective model:

$$\Gamma_k = \int_x \left\{ Z_{q,k} \bar{q} (\gamma_\mu \partial_\mu - \gamma_0 \mu) q + \frac{1}{2} Z_{\phi,k} (\partial_\mu \phi)^2 + h_k \bar{q} \left( T^0 \sigma + i \gamma_5 \vec{T} \cdot \vec{\pi} \right) q + V_k(\rho) - c\sigma \right\} + \dots$$

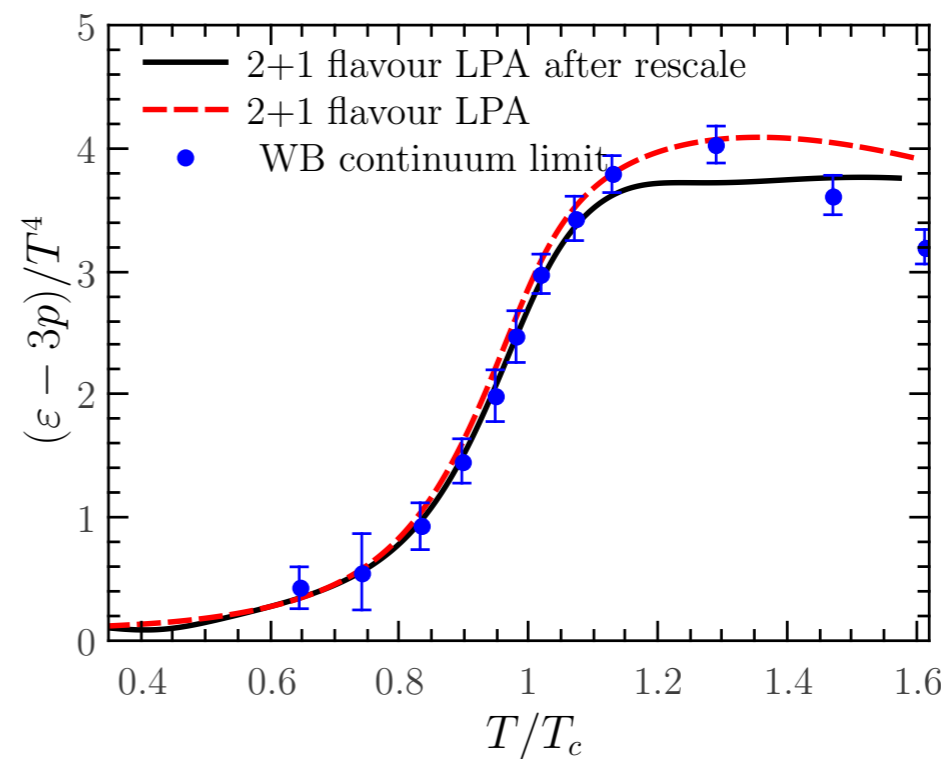
# Thermodynamics of the effective model within FRG



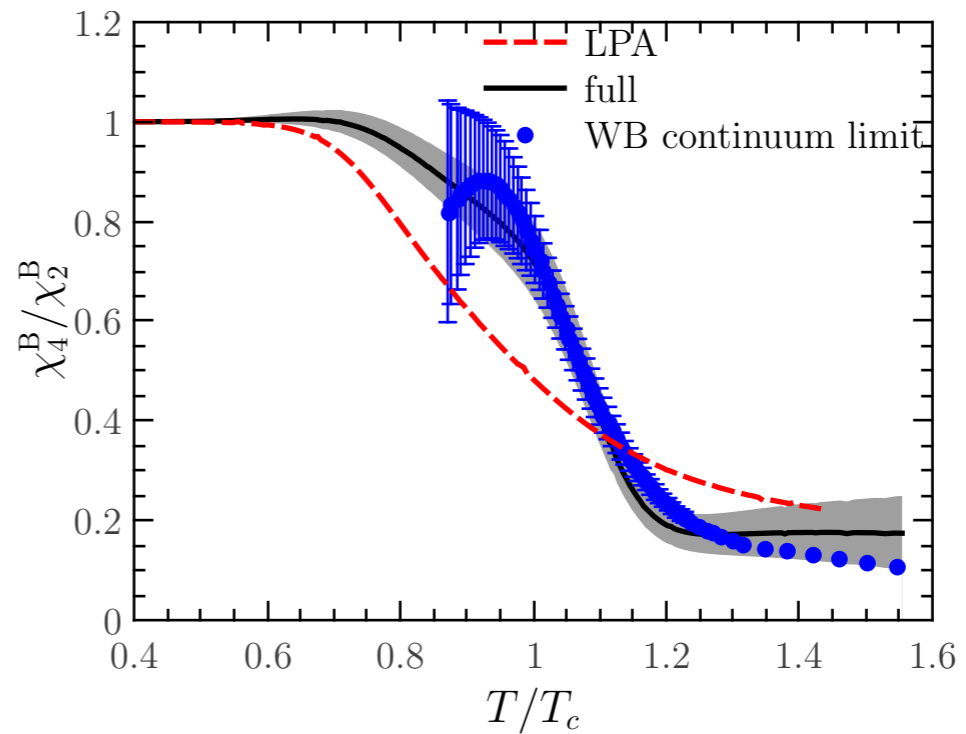
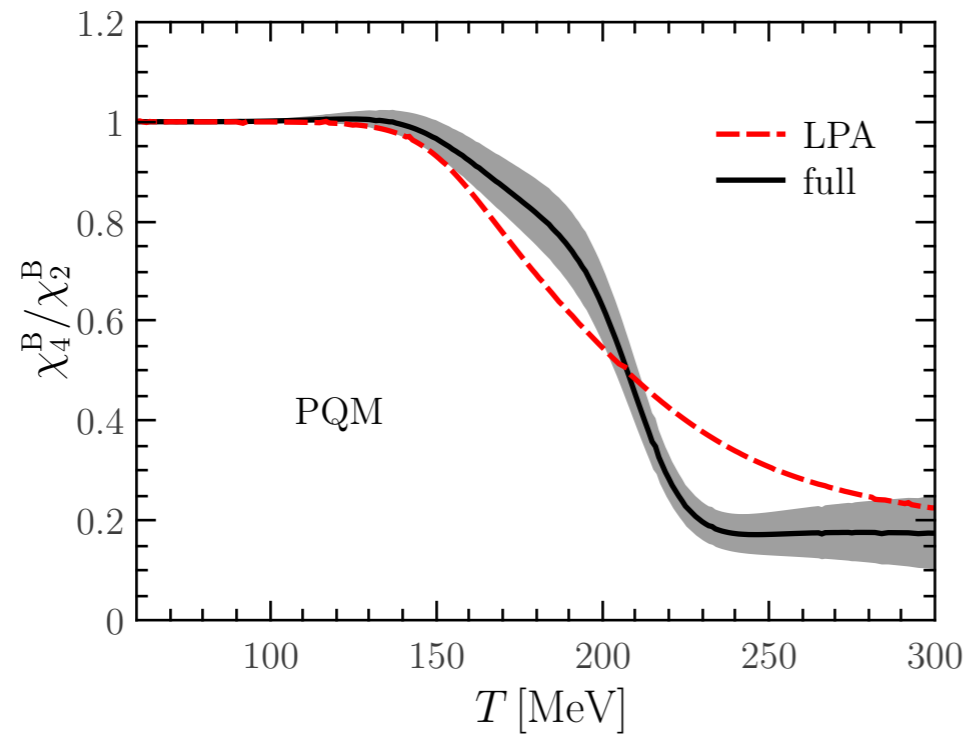
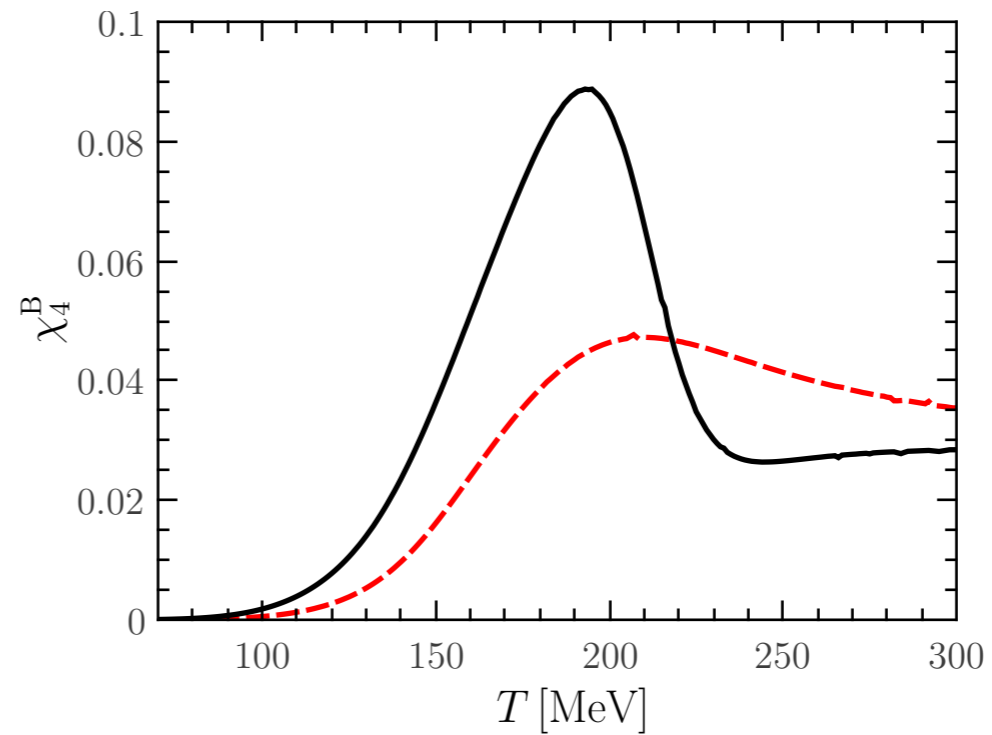
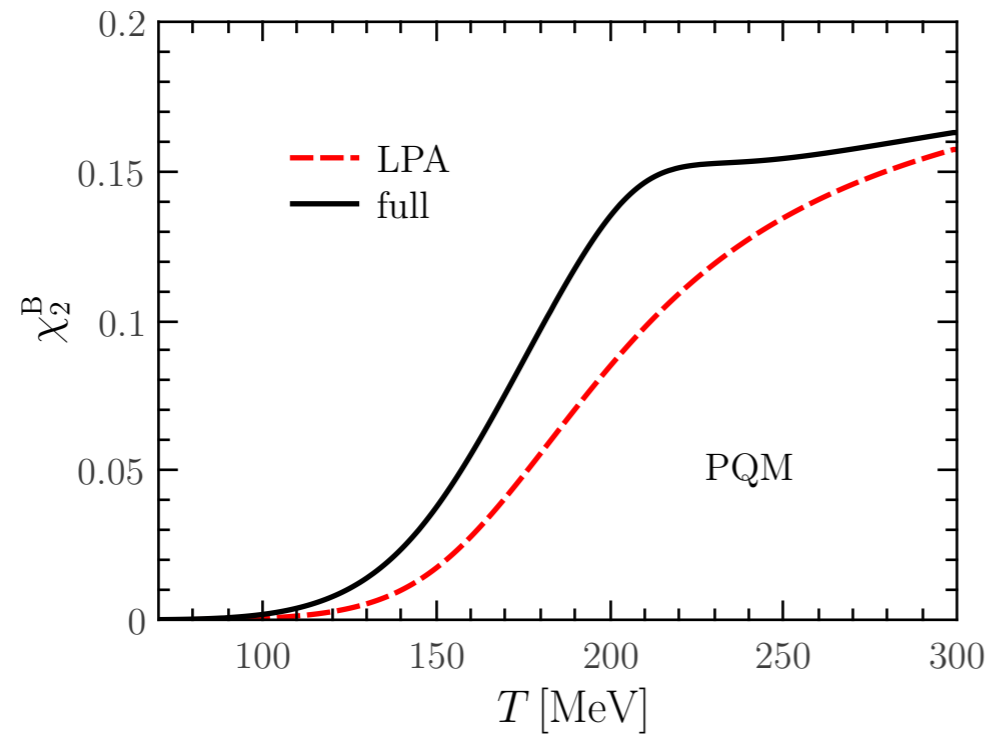
QM



PQM

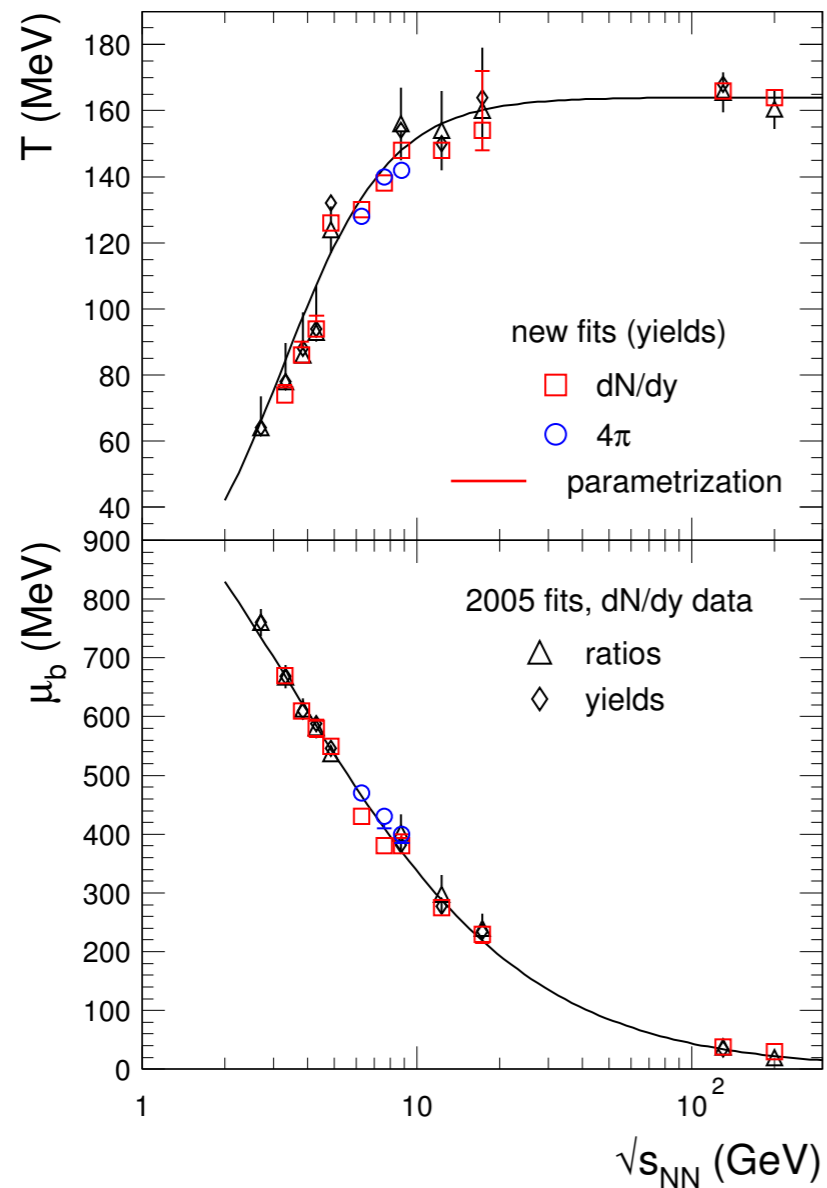


# Baryon number fluctuations

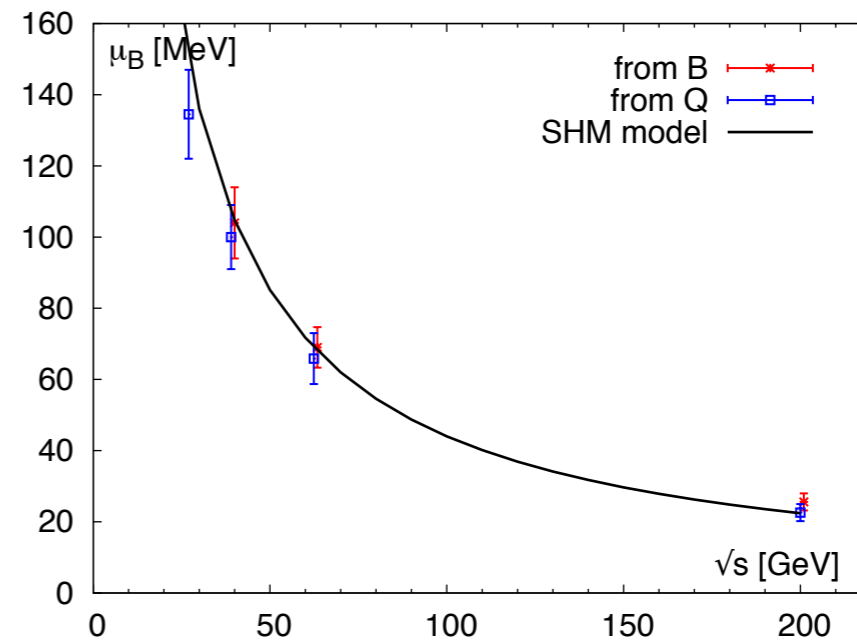


# Freeze-out line

see talk by Peter Braun-Munzinger



Freeze-out temperature and chemical potential obtained from the Statistical Hadronization Model



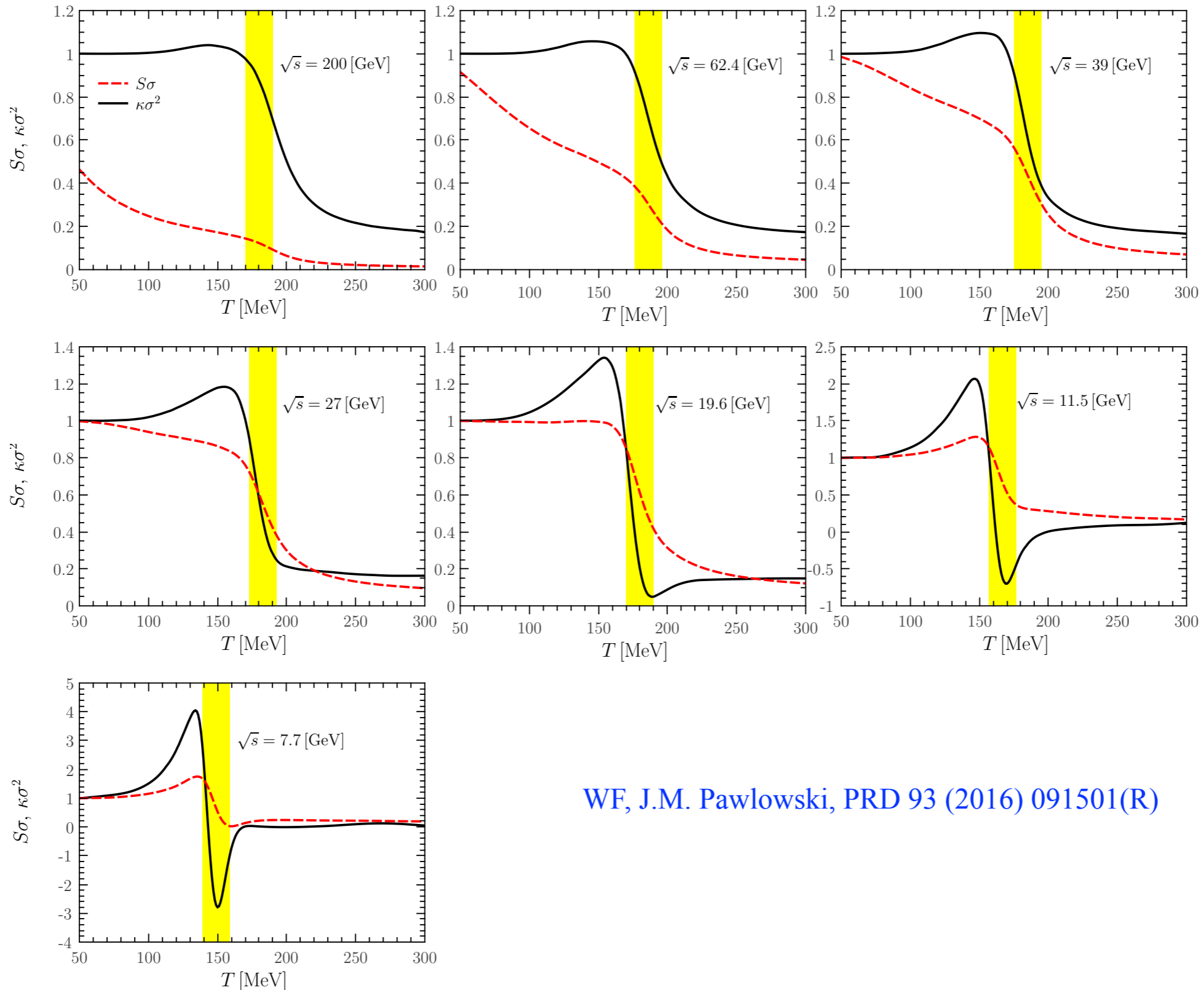
Freeze-out chemical potential obtained from lattice simulations

S. Borsanyi *et al.*, Phys.Rev.Lett. 113 (2014) 052301

A. Andronic, P. Braun-Munzinger, J. Stachel, Phys.Lett. B673 (2009) 142

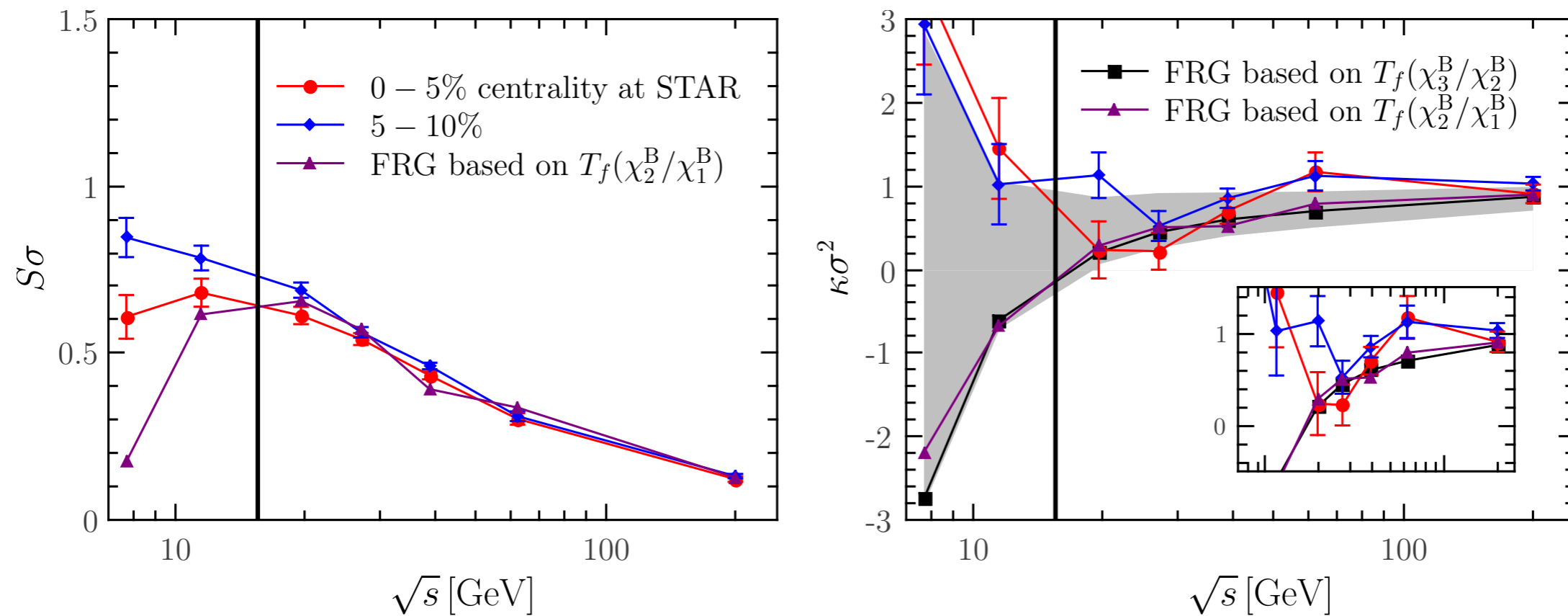


# Correlating the skewness and kurtosis of baryon number distributions



WF, J.M. Pawłowski, PRD 93 (2016) 091501(R)

# Comparison with experimental measurements



WF, J.M. Pawlowski, PRD 93 (2016) 091501(R)

# Silver Blaze Property and the Frequency Dependence

Frequency-dependent quark anomalous dimension:

$$\eta_{q,k}(p) = \frac{1}{Z_{q,k}(p)} \frac{1}{4N_c N_f} \frac{\partial^2}{\partial |\vec{p}|^2} \text{Tr} \left( i\vec{\gamma} \cdot \vec{p} \partial_t \tilde{\Gamma}_{q\bar{q},k}^{(2)}(p) \right)$$

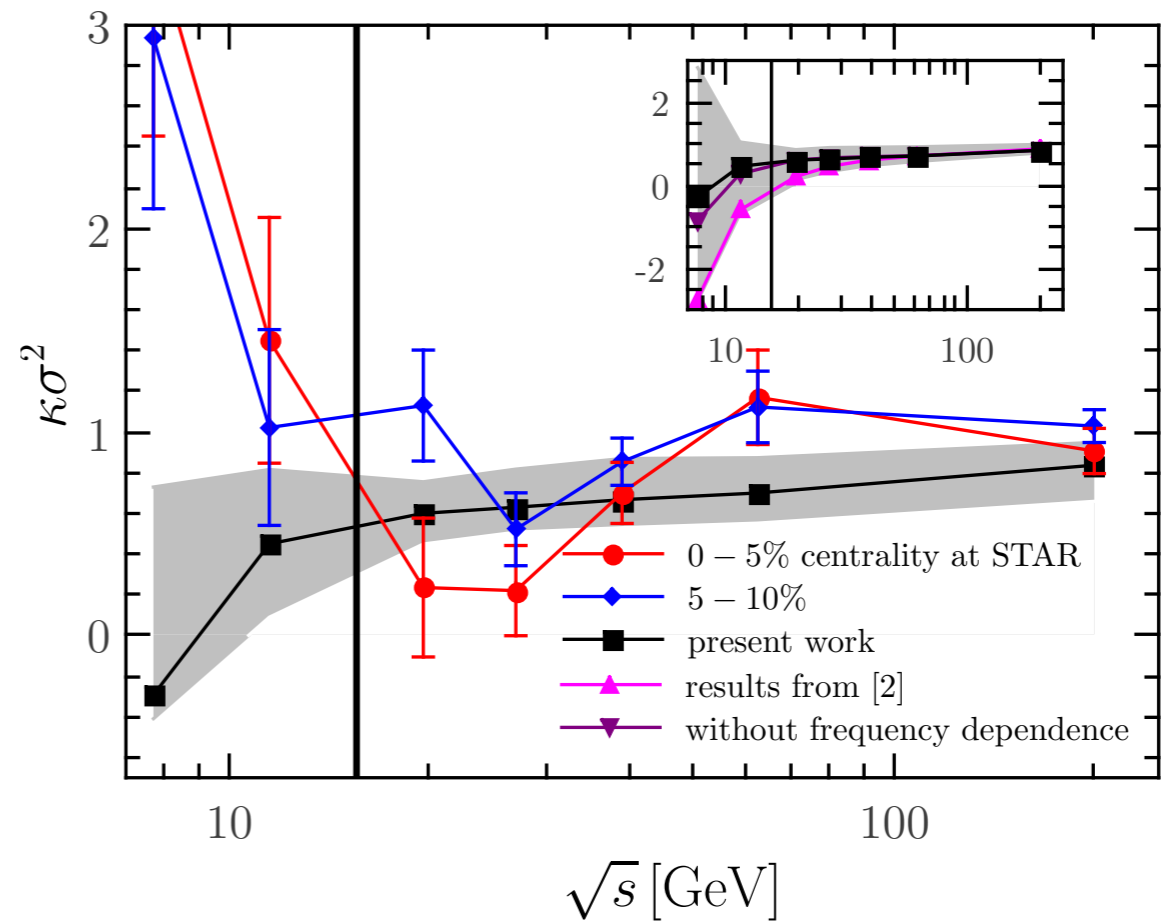
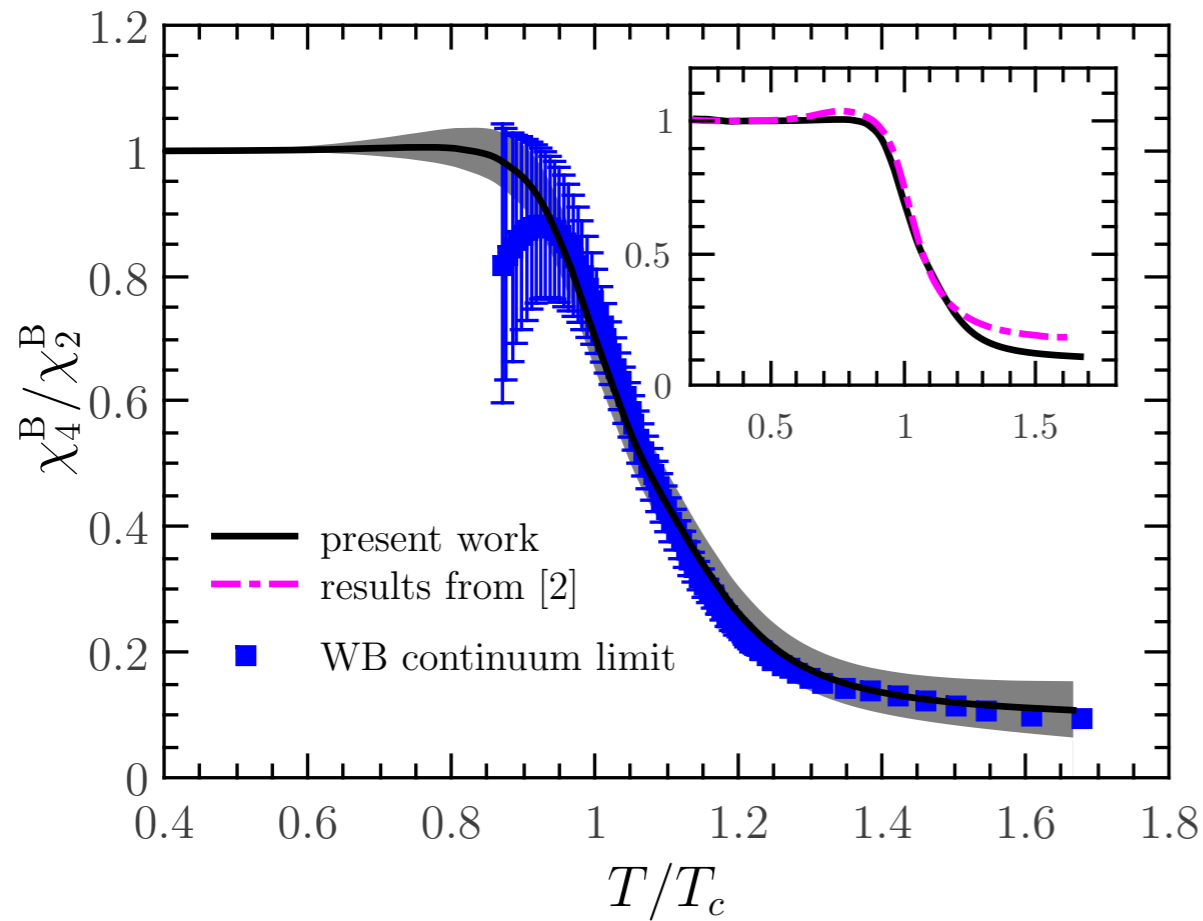
Insert this into the flow of effective potential, and perform the two-loop summation

$$\partial_t V_k^q = - \int_q \text{circle with } \otimes \text{ and } q \text{ arrows} \approx - \iint_{p,q} \text{circle with } \otimes \text{ and } p, q \text{ arrows}$$

One obtains

$$\begin{aligned} \partial_t V_k(\rho) = & \frac{k^4}{360\pi^2} \left\{ 12(5 - \eta_{\phi,k}) \left[ (N_f^2 - 1) \mathcal{B}_{(1)}(\bar{m}_{\pi,k}^2) \right. \right. \\ & + \left. \mathcal{B}_{(1)}(\bar{m}_{\sigma,k}^2) \right] - 5N_c \left( 48N_f \mathcal{F}_{(1)}(\bar{m}_{F,k}^2) \right. \\ & + \frac{1}{2\pi^2} (-4 + \eta_{\phi,k}) \bar{h}_k^2 \left[ \mathcal{F}\mathcal{F}\mathcal{B}_{(1,1,2)}(\bar{m}_{F,k}^2, \bar{m}_{\sigma,k}^2) \right. \\ & \left. \left. + (N_f^2 - 1) \mathcal{F}\mathcal{F}\mathcal{B}_{(1,1,2)}(\bar{m}_{F,k}^2, \bar{m}_{\pi,k}^2) \right] \right) \left. \right\}, \end{aligned}$$

# Two-loop Results



$$\partial_t V_k^q = - \int_q \text{[Diagram 1]} \approx - \iint_{p,q} \text{[Diagram 2]}$$

Diagram 1: A circle with a cross in the center and an arrow labeled  $q$  at the bottom.

Diagram 2: A circle with a cross in the center, an arrow labeled  $q$  at the bottom, and an arrow labeled  $p$  pointing right across the center.

# From effective models to QCD

## Summary on effective model including mesonic fluctuations

- Quantitative agreement with lattice results below  $\sim 1.2 T_c$
- No bump on the baryon number kurtosis
- Discrepancy observed at large temperature or density because of the UV cutoff effect
- UV cutoff should be pushed up higher, and glue quantum fluctuations should be included



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$$\partial_t \Gamma_k = \frac{1}{2} \left( \text{diagram 1} - \text{diagram 2} - \text{diagram 3} + \frac{1}{2} \text{diagram 4} \right)$$

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Come back!

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$$\Gamma_k = \int_x \left\{ Z_{q,k} \bar{q} (\gamma_\mu \partial_\mu - \gamma_0 \mu) q + \frac{1}{2} Z_{\phi,k} (\partial_\mu \phi)^2 + h_k \bar{q} \left( T^0 \sigma + i \gamma_5 \vec{T} \cdot \vec{\pi} \right) q + V_k(\rho) - c\sigma \right\} + \dots$$

## Rebosonized QCD:

$$\Gamma_k = \int_x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + Z_{c,k} (\partial_\mu \bar{c}^a) D_\mu^{ab} c^b + \frac{1}{2\xi} (\partial_\mu A_\mu^a)^2 + Z_{q,k} \bar{q} (\gamma_\mu D_\mu) q - \lambda_{q,k} [(\bar{q} T^0 q)^2 - (\bar{q} \gamma_5 \vec{T} q)^2] + h_k [\bar{q} (i \gamma_5 \vec{T} \vec{\pi} + T^0 \sigma) q] + \frac{1}{2} Z_{\phi,k} (\partial_\mu \phi)^2 + V_k(\rho) - c\sigma \right\} + \Delta \Gamma_{\text{glue}}$$

# Flow Equations

Wetterich equation:

$$\begin{aligned}\partial_t \Gamma_k[\Phi] &= \frac{1}{2} \text{STr} \left\{ \partial_t R_k (\Gamma_k^{(2)}[\Phi] + R_k)^{-1} \right\} \\ &= \frac{1}{2} \text{STr} \left\{ \tilde{\partial}_t \ln (\Gamma_k^{(2)}[\Phi] + R_k) \right\}\end{aligned}$$

with  $t = \ln(k/\Lambda)$  and

$$(\Gamma_k^{(2)}[\Phi])_{ij} := \frac{\overrightarrow{\delta}}{\delta \Phi_i} \Gamma_k[\Phi] \frac{\overleftarrow{\delta}}{\delta \Phi_j} \quad \Gamma_k^{(2)} + R_k = \mathcal{P} + \mathcal{F}$$

Vertex expansion:

$$\begin{aligned}\partial_t \Gamma_k &= \frac{1}{2} \text{STr} \{ \tilde{\partial}_t \ln(\mathcal{P} + \mathcal{F}) \} = \frac{1}{2} \text{STr} \tilde{\partial}_t \ln \mathcal{P} \\ &+ \frac{1}{2} \text{STr} \tilde{\partial}_t \left( \frac{1}{\mathcal{P}} \mathcal{F} \right) - \frac{1}{4} \text{STr} \tilde{\partial}_t \left( \frac{1}{\mathcal{P}} \mathcal{F} \right)^2 + \dots\end{aligned}$$

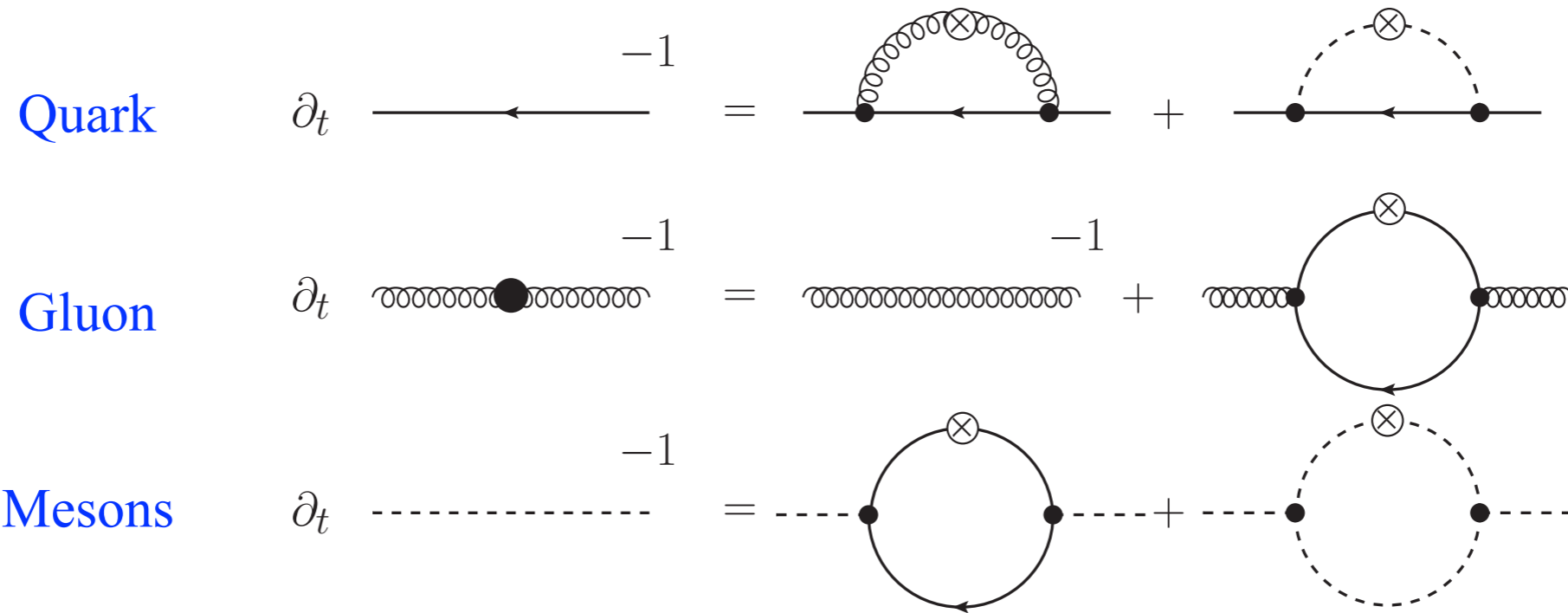
3d optimized regulator:

$$R_{\text{F},k}(q) = Z_{q,k} i \vec{q} \cdot \vec{\gamma} r_{\text{F}}\left(\frac{\vec{q}^2}{k^2}\right), \quad \text{with} \quad r_{\text{F}}(x) = \left(\frac{1}{\sqrt{x}} - 1\right) \Theta(1-x)$$

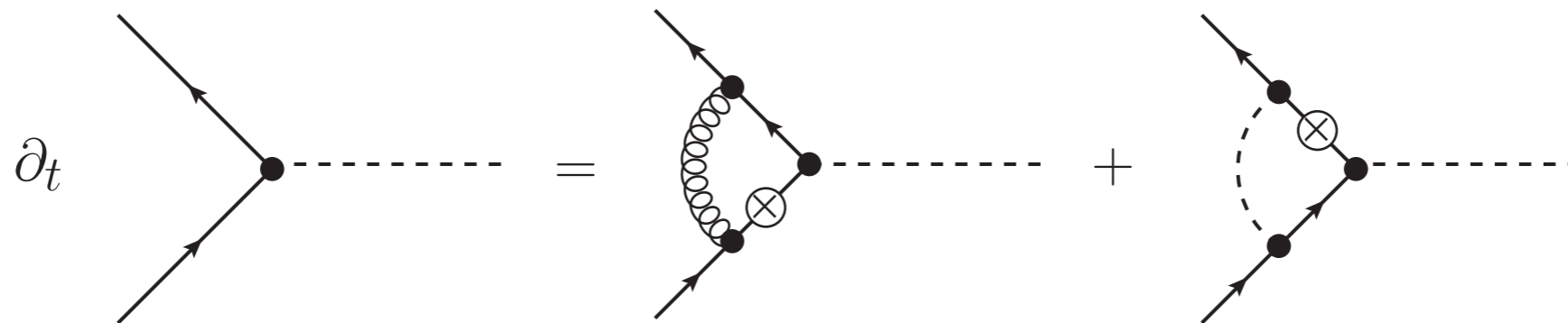
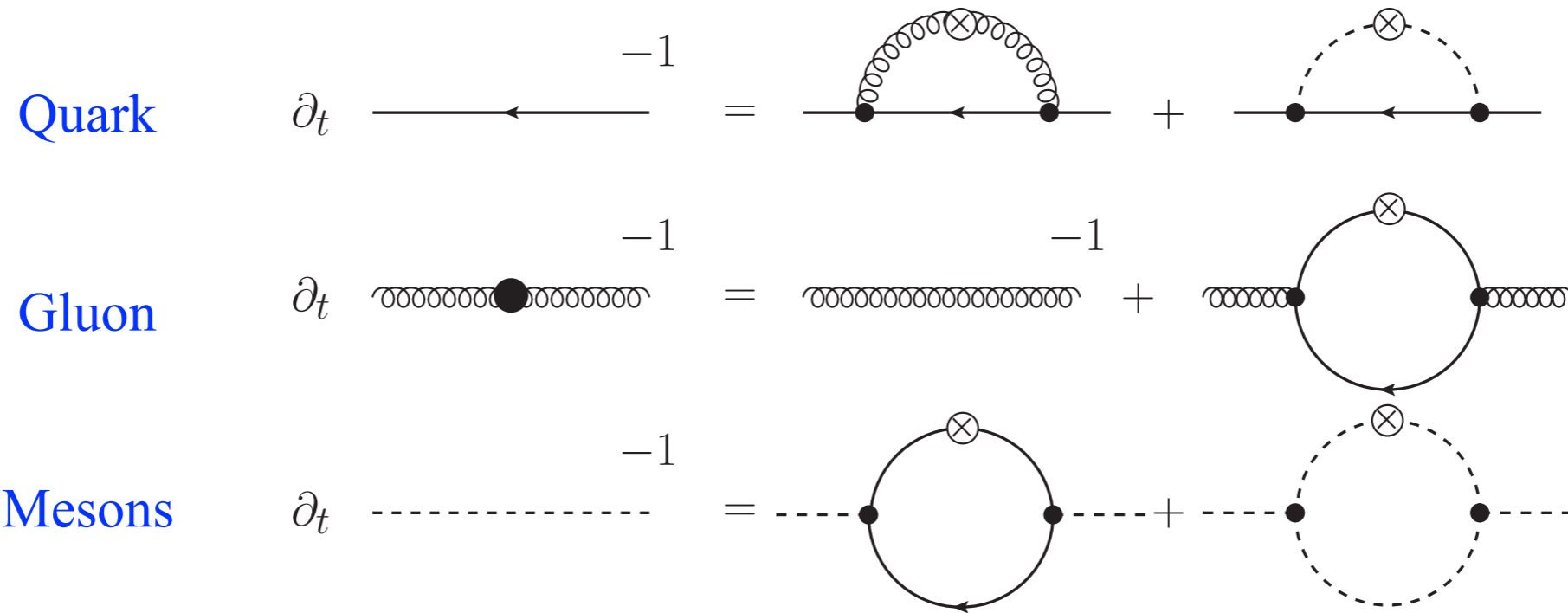
$$R_{\text{B},k}(q) = Z_{\phi,k} \vec{q}^2 r_{\text{B}}\left(\frac{\vec{q}^2}{k^2}\right), \quad \text{with} \quad r_{\text{B}}(x) = \left(\frac{1}{x} - 1\right) \Theta(1-x)$$

# Feynman Diagrams

# Feynman Diagrams



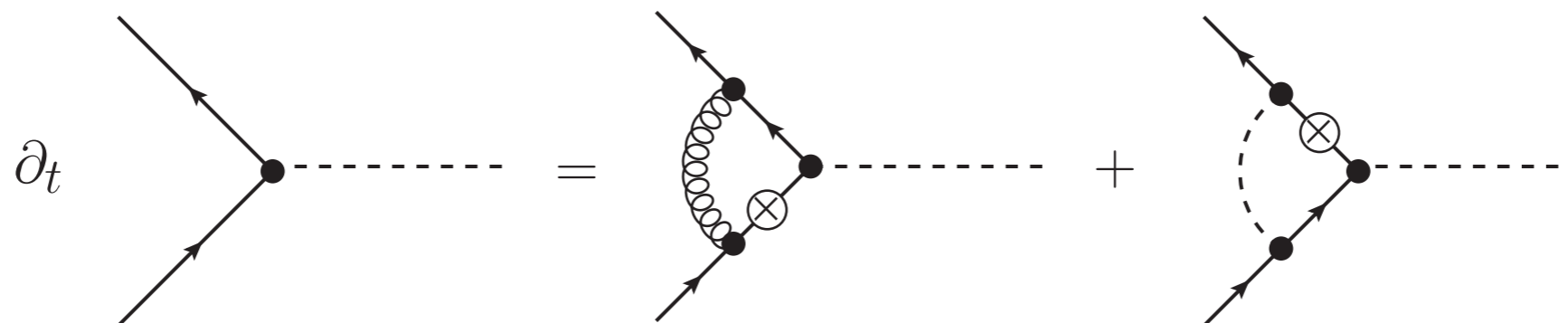
# Feynman Diagrams



Yukawa coupling



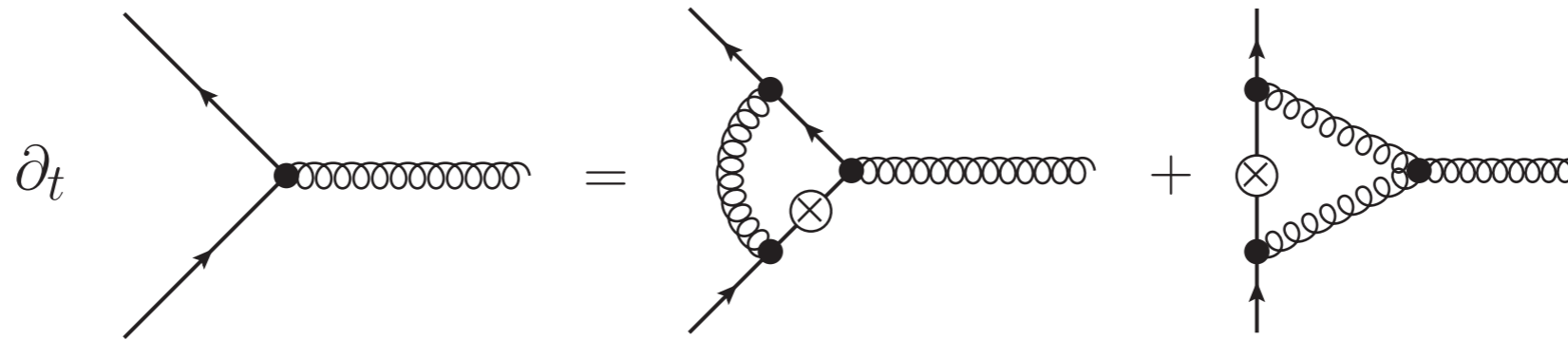
# Feynman Diagrams



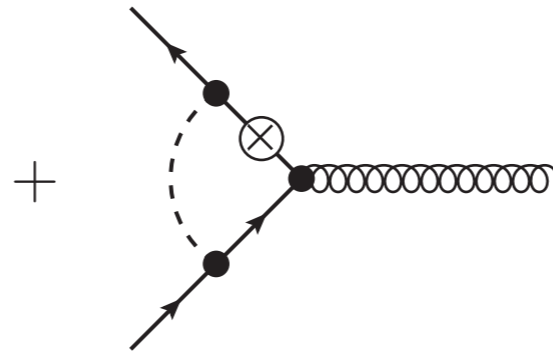
Yukawa coupling

# Feynman Diagrams

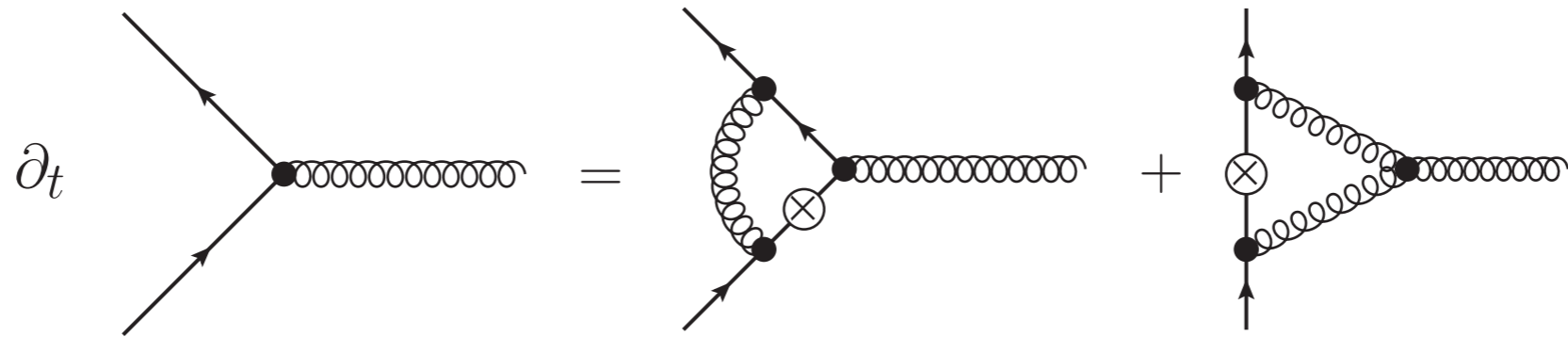
# Feynman Diagrams



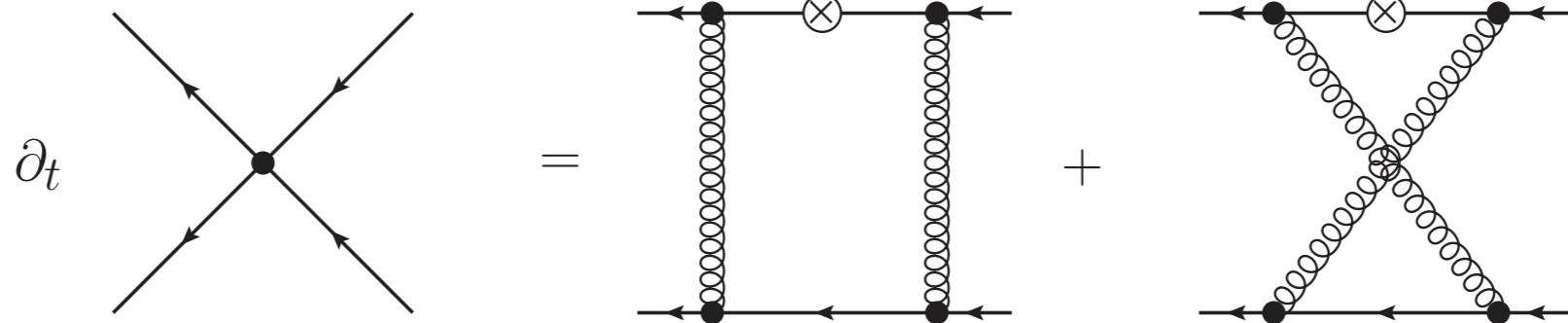
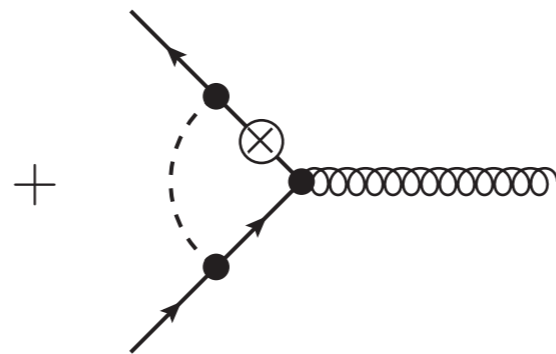
Quark gluon vertex



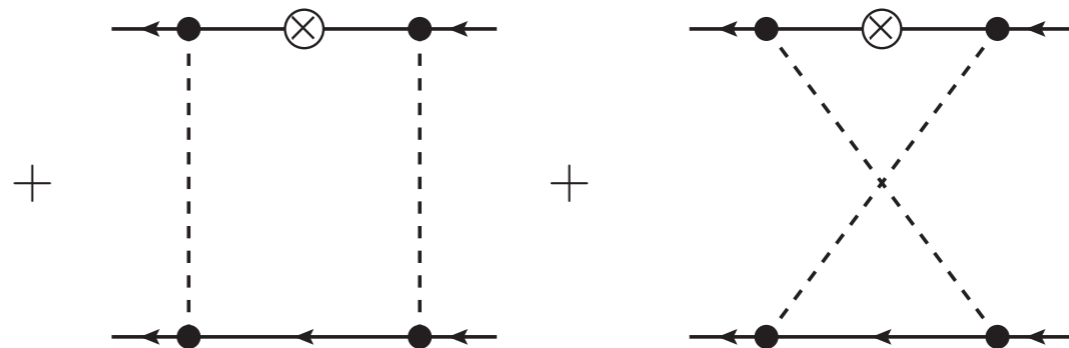
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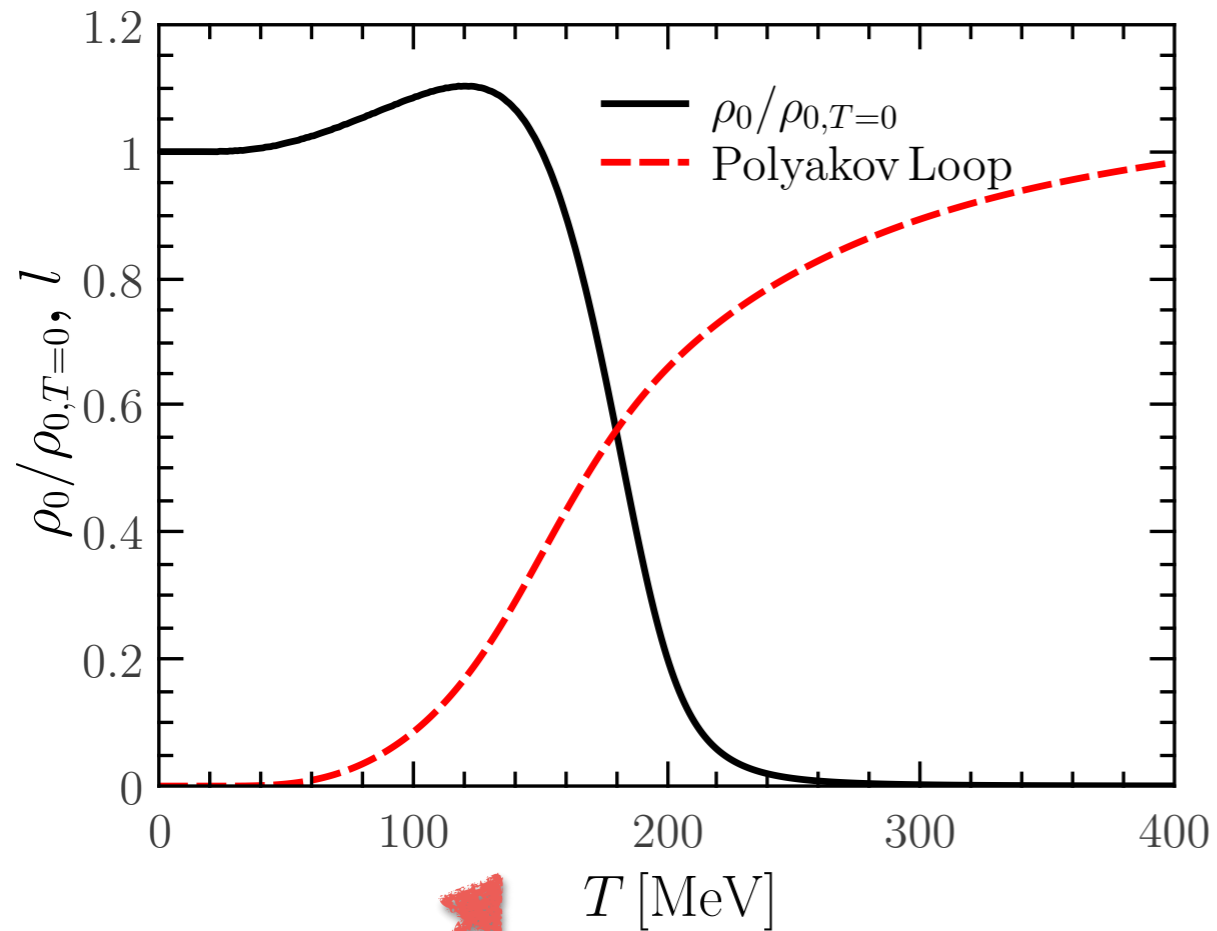
Quark gluon vertex



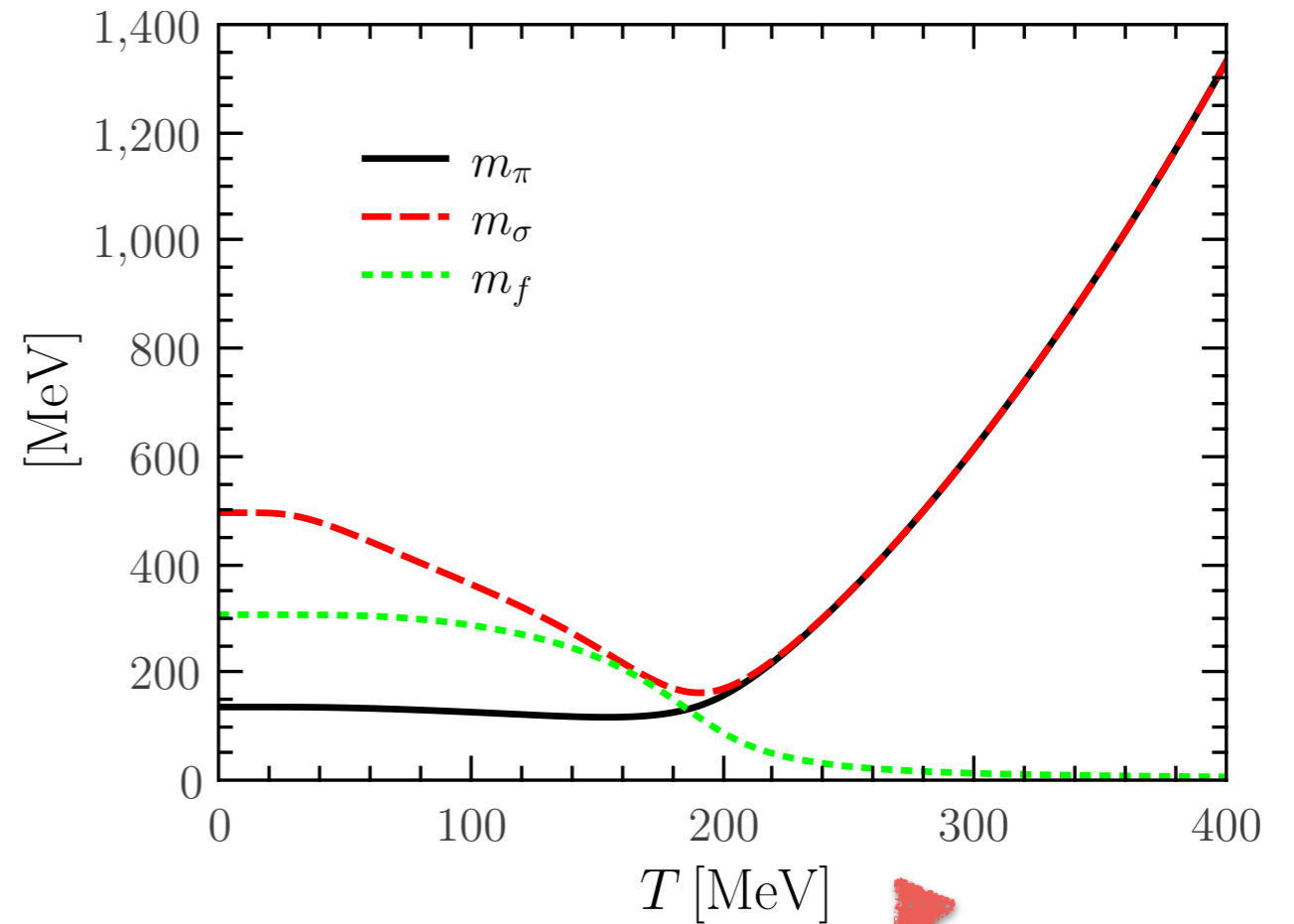
Four-fermion coupling



# QCD Phase Transition with T

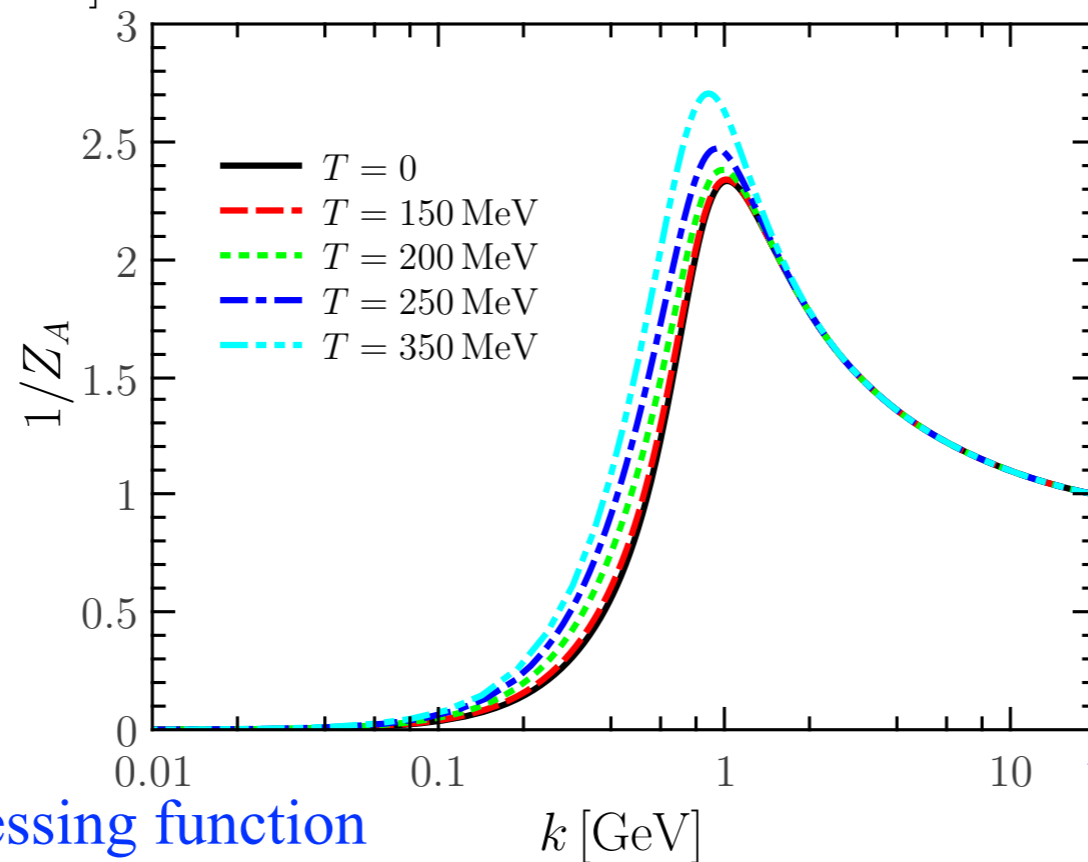
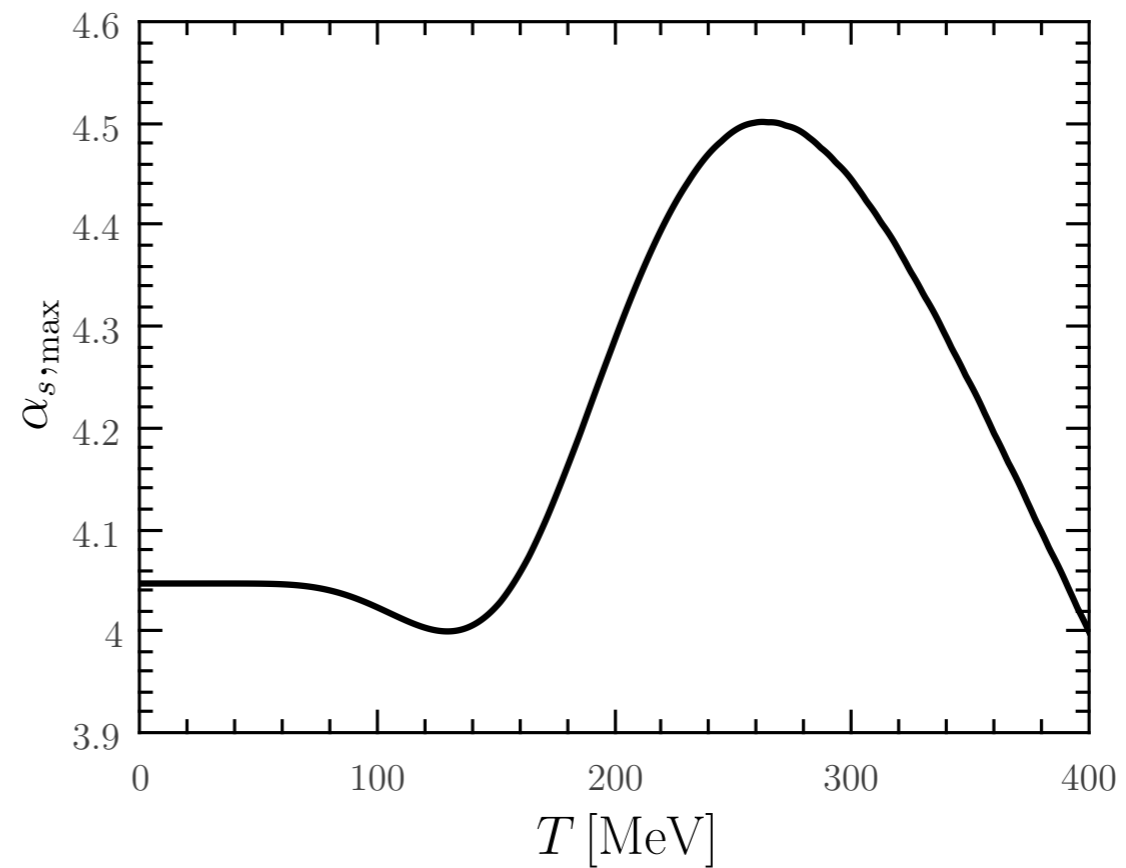
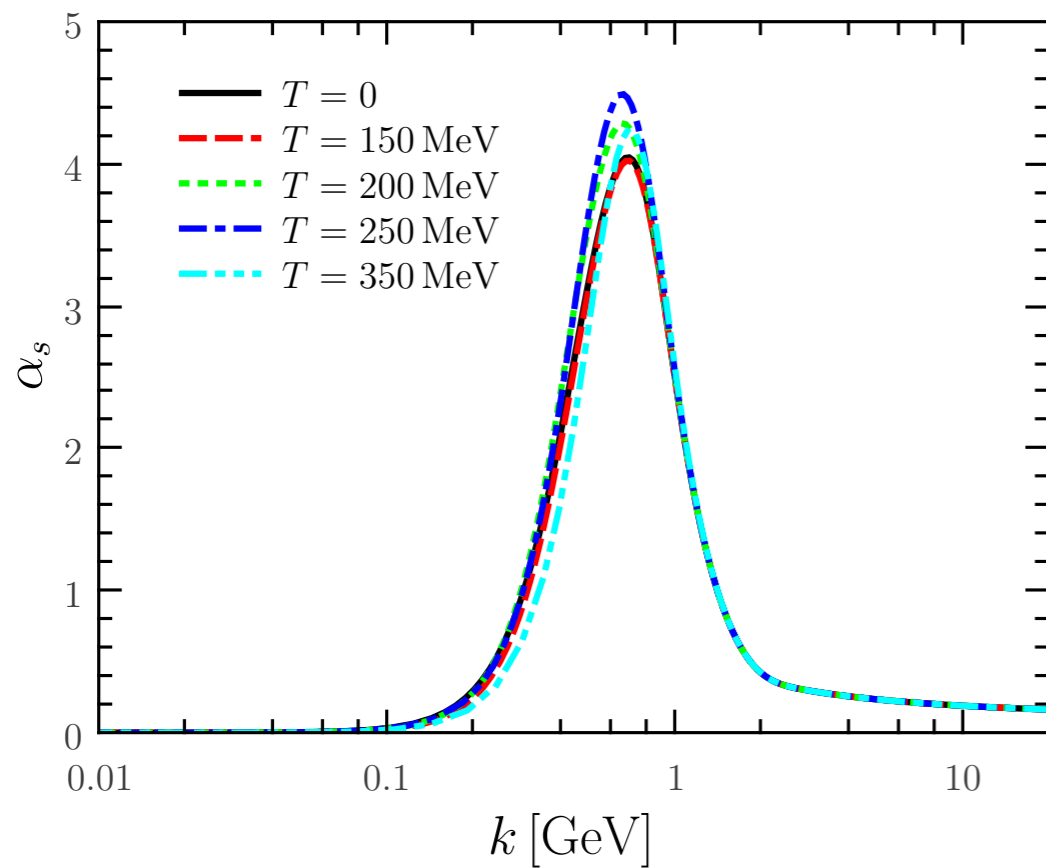


Order parameter



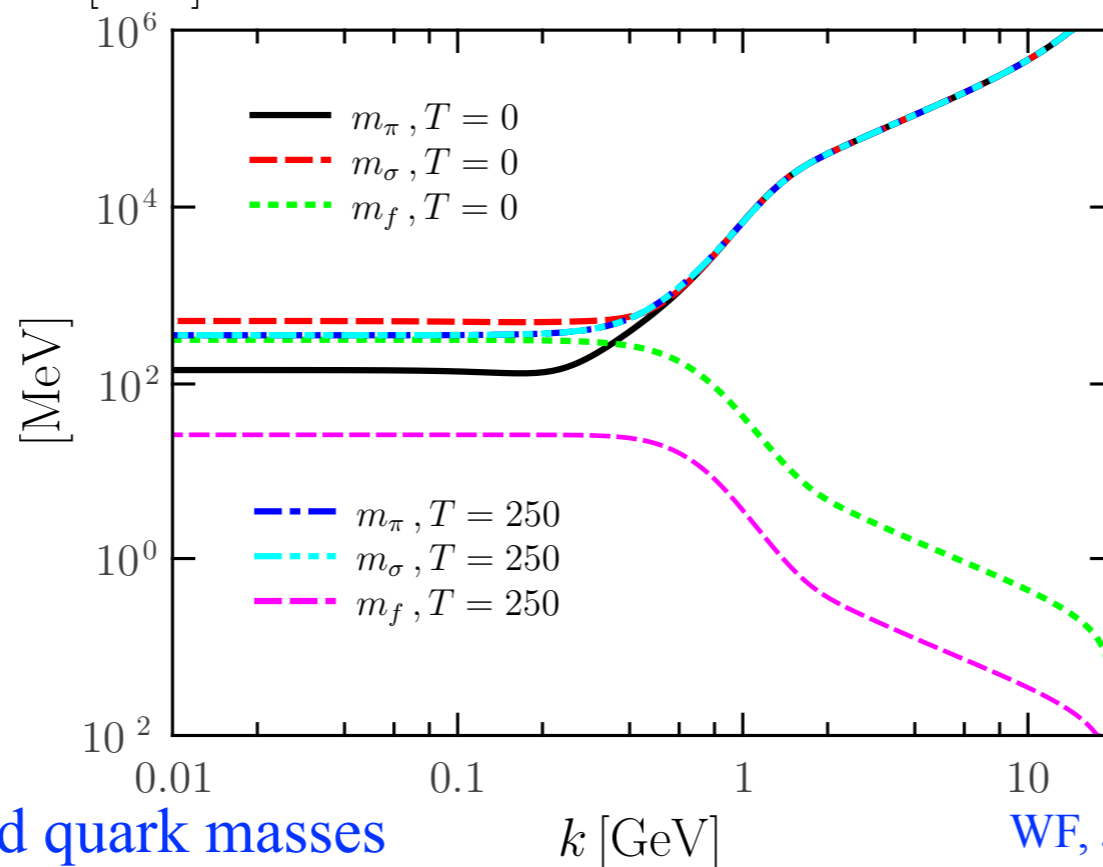
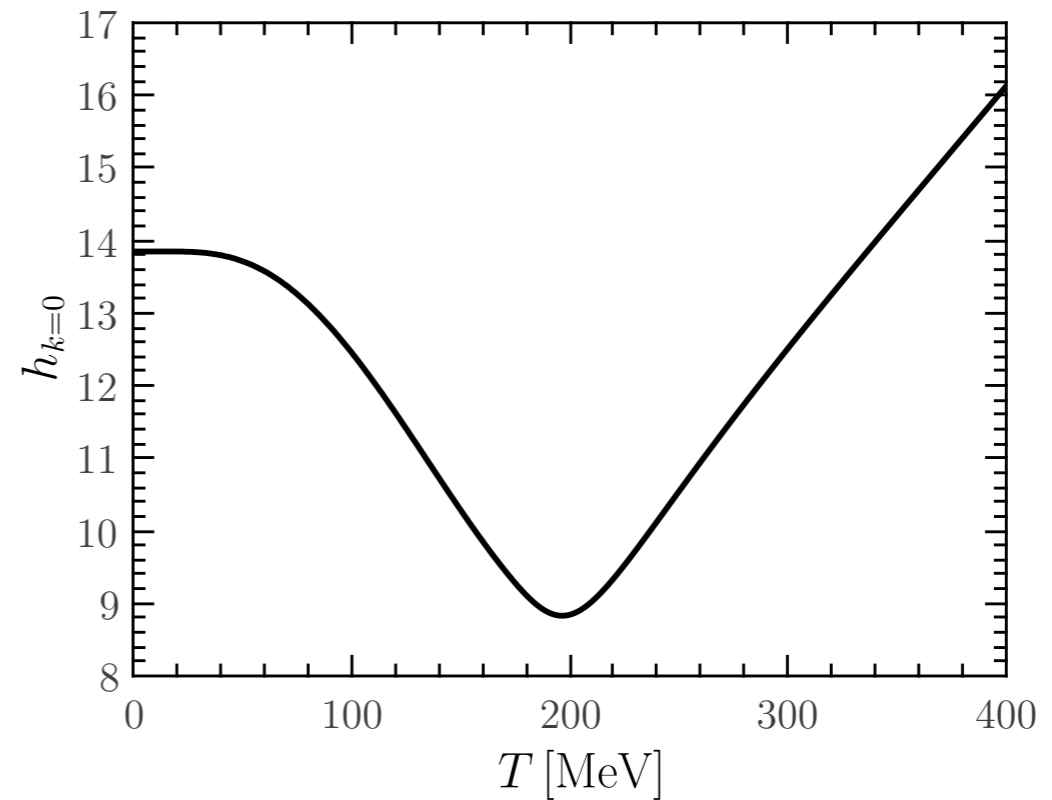
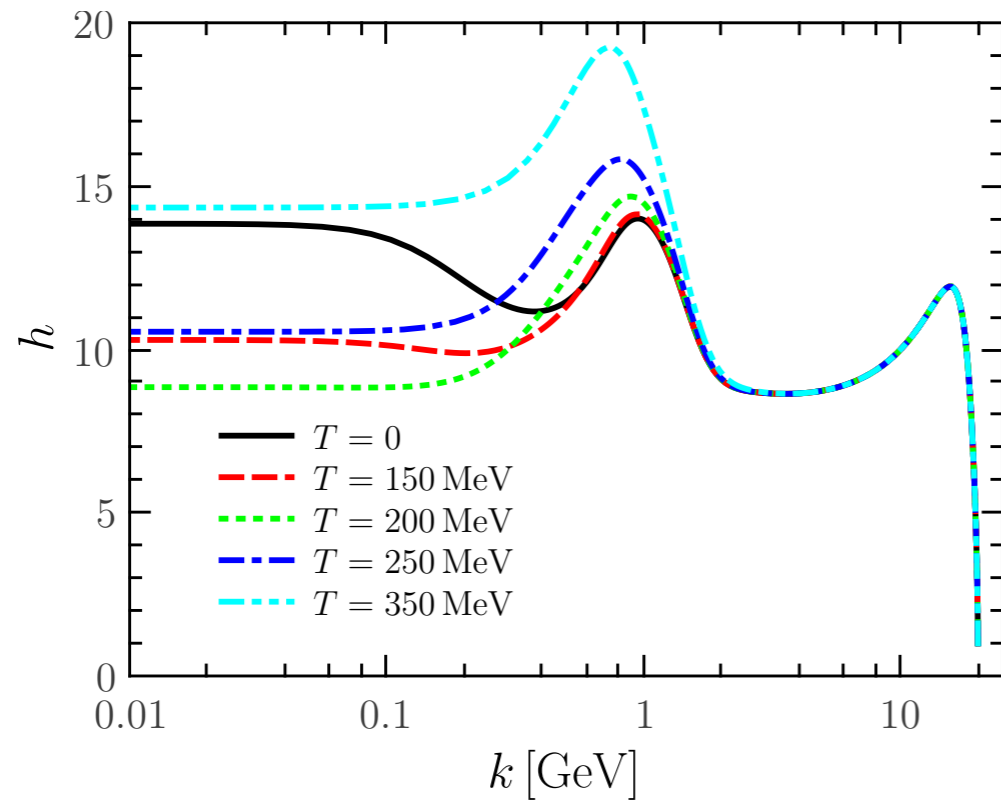
Meson and quark mass

# Running of the Strong Coupling



Gluon dressing function

# Yukawa coupling and Masses



Meson and quark masses

WF, J.M. Pawłowski, F. Rennecke, in preparation

# Summary and outlook



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- ★ Better agreement with lattice simulations and experiments is observed when more quantum fluctuations are included.
- ★ We have also performed FRG QCD calculations at finite temperature. The QCD phase transitions have been investigated.
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**Thank you very much for your attentions!**