

Probing the primordial power spectrum (and non-Gaussianity) with μ -distortions

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In collaboration with:

A. Melchiorri, E. Pajer: Phys. Rev. D 93, 083515 (2016)

E. Pajer and D. van der Woude, in preparation

Upcoming CMB experiments

Main “inflationary” target → primordial B -modes:

- primordial tensor modes: smoking gun of inflation
- Lyth bound: constraints on inflationary models (large field vs. small field)
- open possibility of testing $r = -n_t/8$

However:

- large field models (e.g. “natural” $V(\phi) \sim m^2\phi^2$) strongly disfavored by *Planck*
- lower bound $r > 10^{-53}$ from

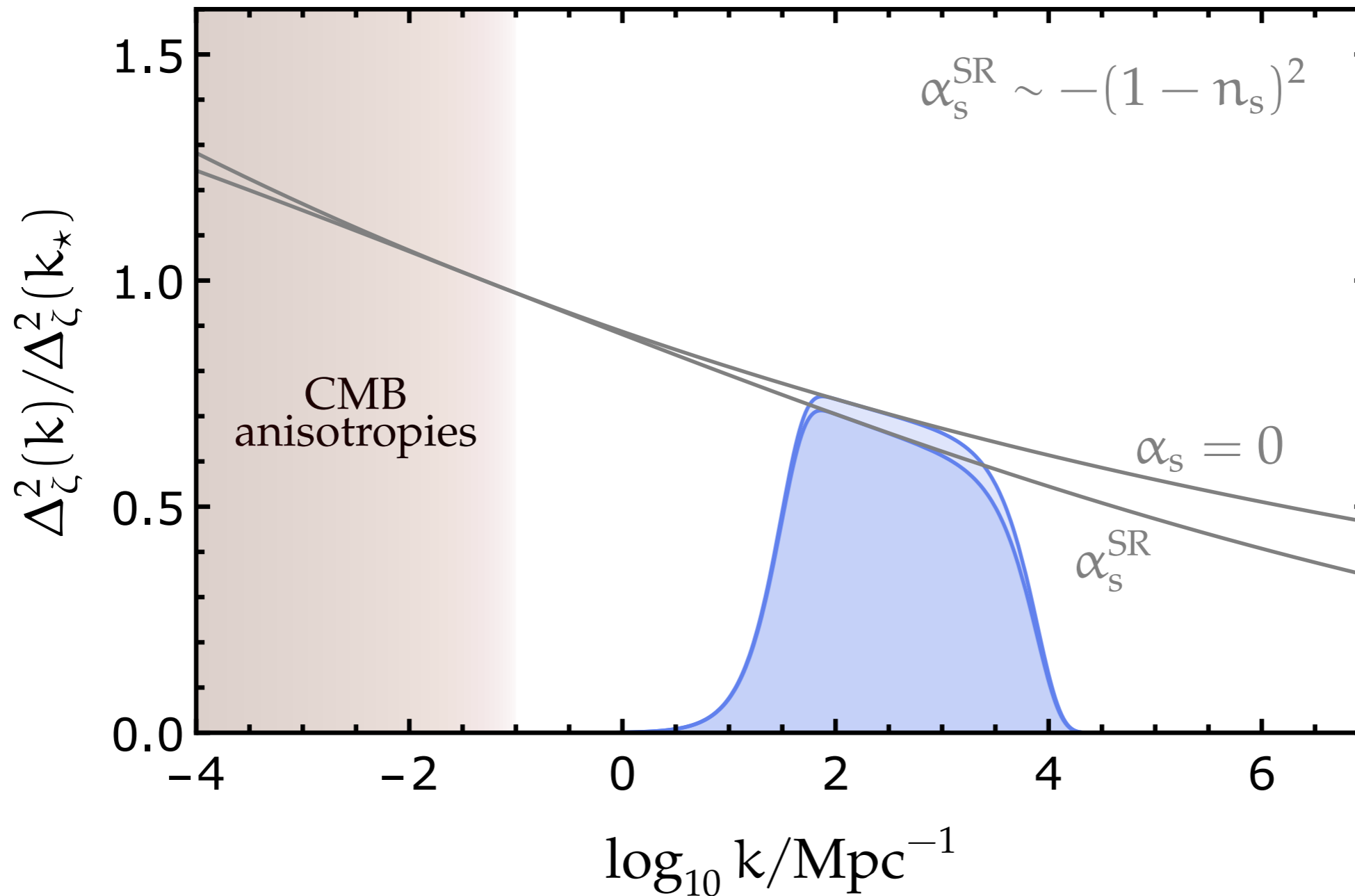
$$A_s \approx 10^{-1} \times \frac{H^2}{M_{\text{P}}^2} \frac{1}{r} \approx 10^{-9}, \quad M_{\text{P}}^2 H^2 > 10 \times \underbrace{T_{\text{reh}}^4}_{1 \text{ TeV}} \Rightarrow r > 10^{-53}$$

⇒ just scratching the surface of a wide theoretical prior

Scalar modes are still the safest bet? Power spectrum at small scales, primordial NG...

Monopole of μ : constraints on the power spectrum

μ -distortions and primordial scalar spectrum

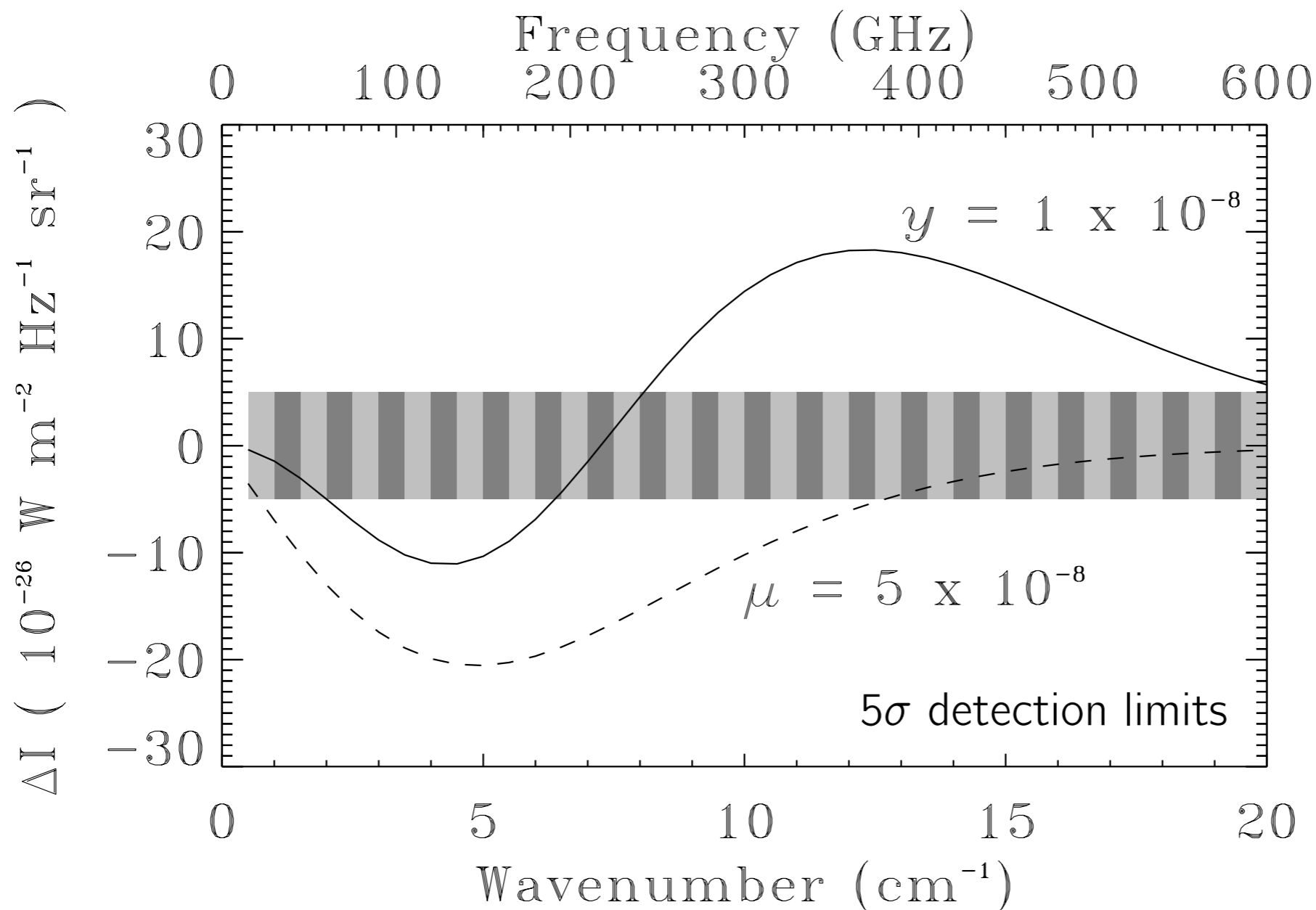


$$\Delta_{\zeta}^2 = A_{\zeta} \left(\frac{k}{k_{\star}} \right)^{n_s - 1 + \frac{\alpha_s}{2} \log \frac{k}{k_{\star}}}$$

$$\alpha_s^{\text{SR}} \sim -(1 - n_s)^2 \ll 1 - n_s \sim 0.04$$

PIXIE

- 400 channels from 30 GHz to 6 THz. $\sigma_\mu = 10^{-8}$ (very optimistic! Taking foregrounds into account: $\sigma_\mu \approx 10^{-7}$ Abitbol, Chluba, Hill, Johnson, 2017)
- \sim three (\sim two) orders of magnitude improvement over FIRAS

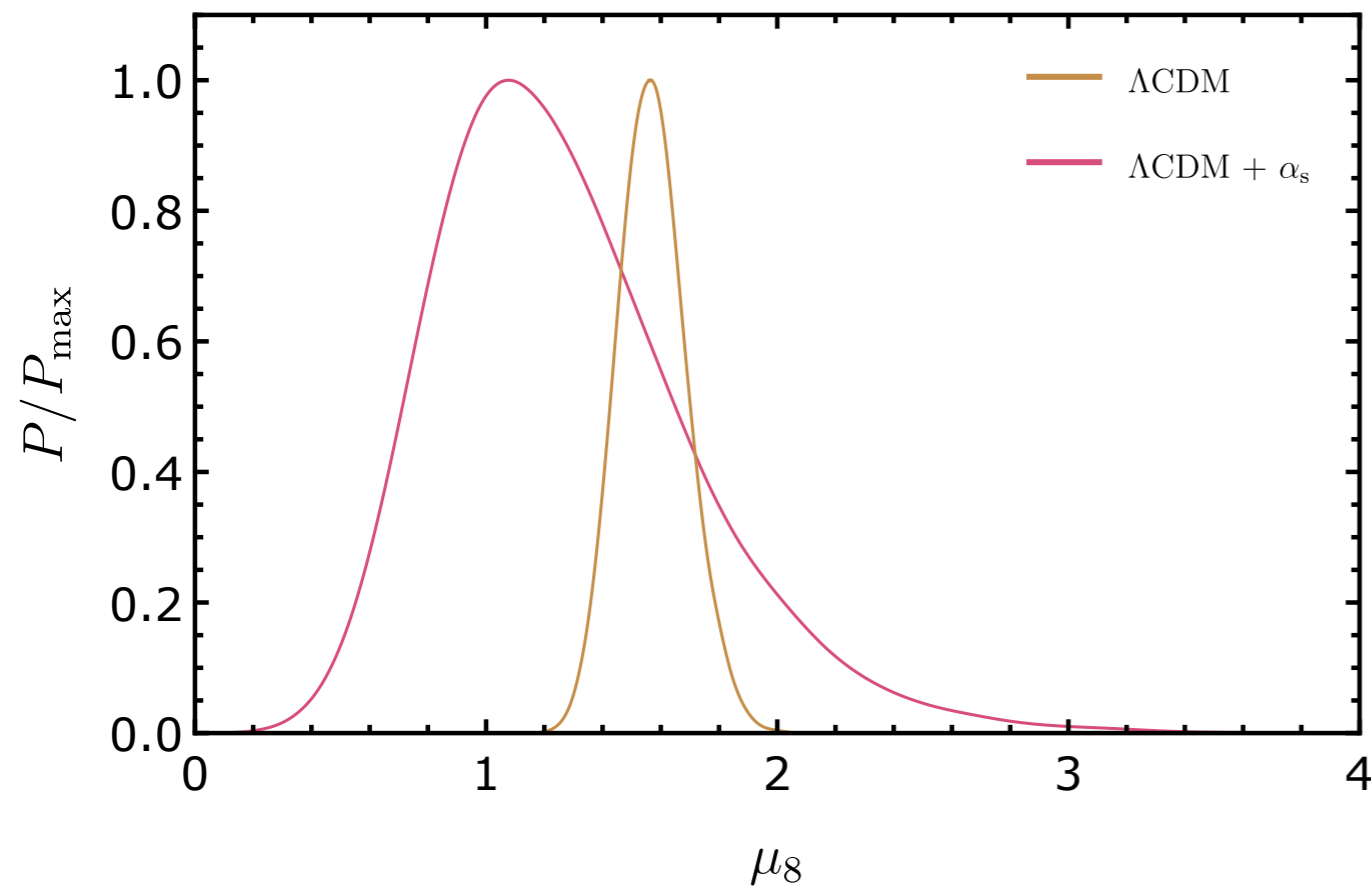


Kogut et al. (2011)

Predictions from CMB anisotropies

idistort code used
(Khatri and Sunyaev, 2013)

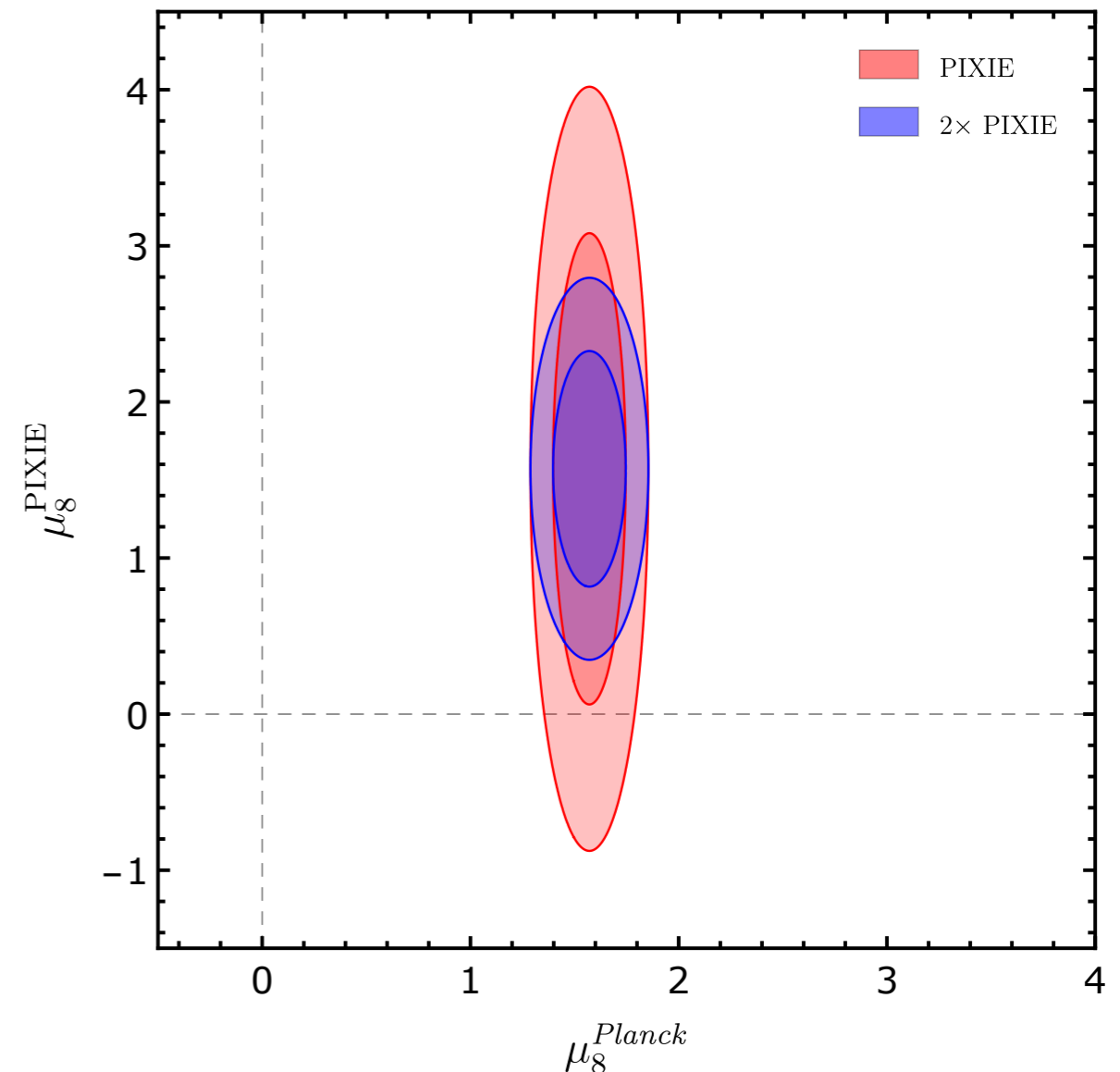
$$\sigma_\mu = (1/n) \times 10^{-8} \rightarrow \text{define } \mu_8 = \mu \times 10^8$$



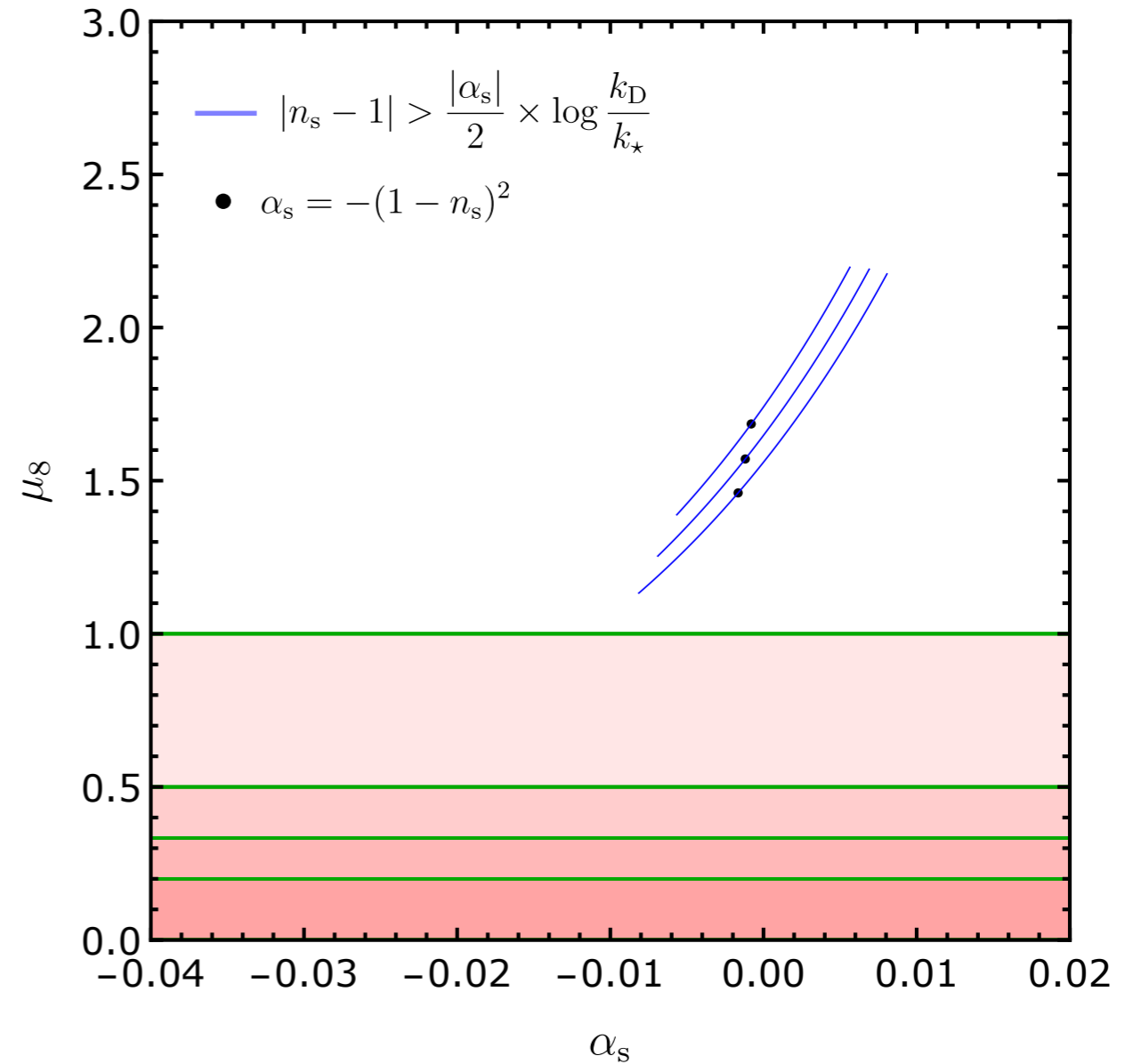
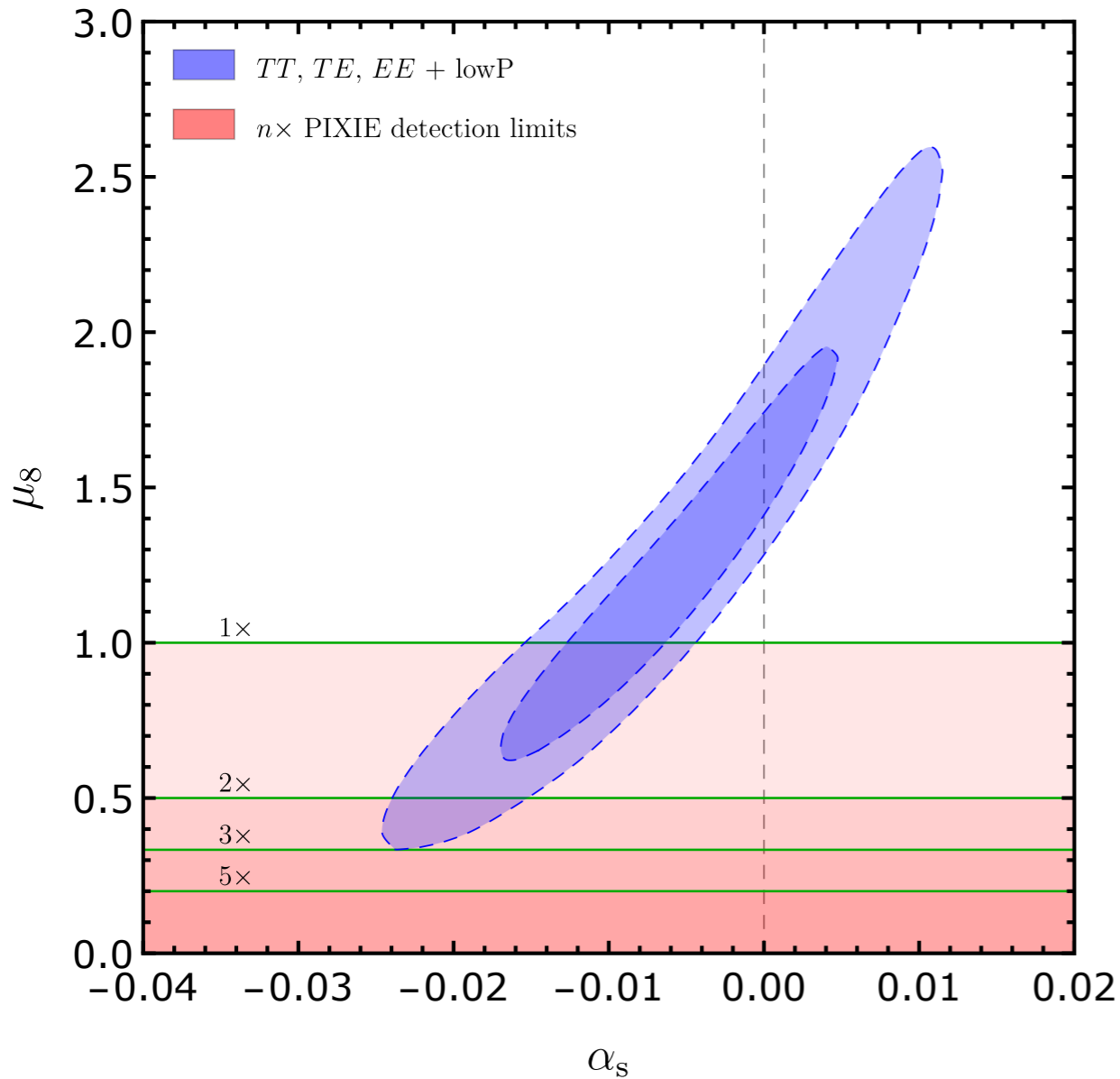
with $\sigma_\mu \sim 0.5 \times 10^{-8} \rightarrow$
 $\mu_8 \leq 0$ excluded at $\approx 3\sigma$

\rightarrow from *Planck* data

$$\mu_8 = 1.57^{+0.11}_{-0.13} \text{ (68\% CL)}$$



What about the running?



“Slow-roll prior”:

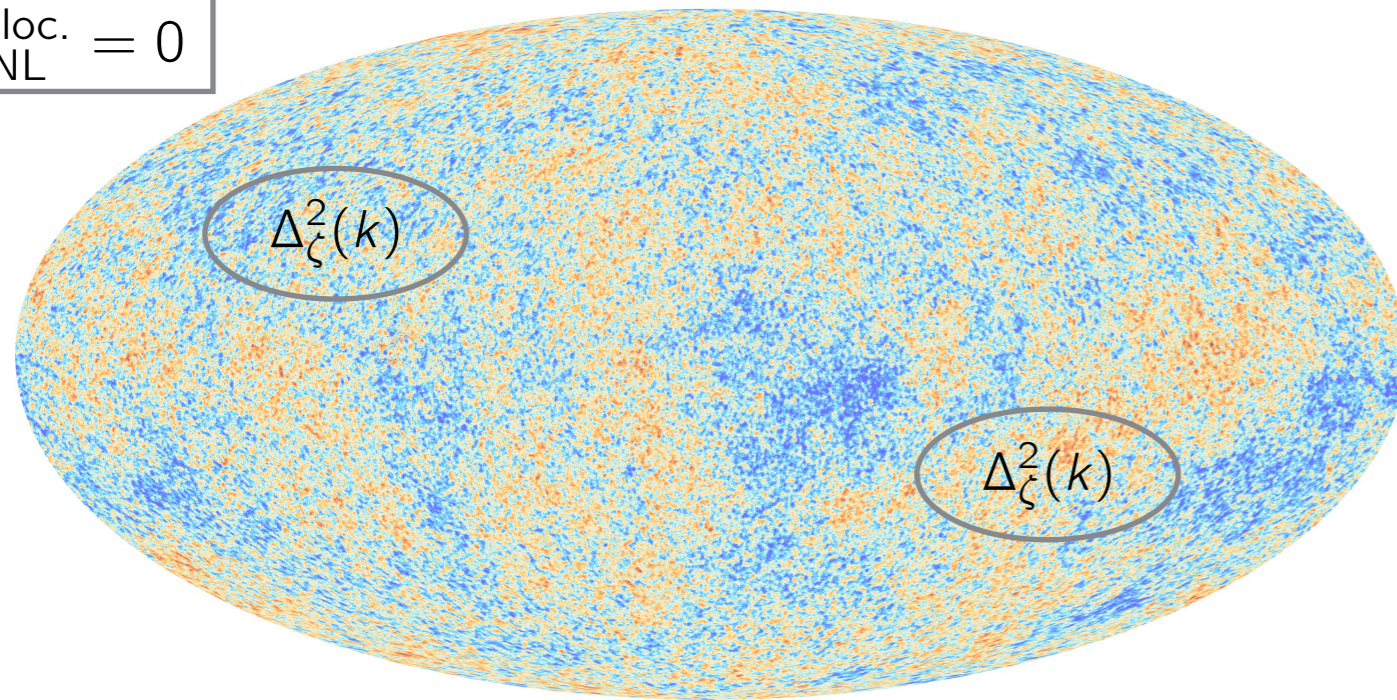
- *Planck*: $n_s = 0.9655 \pm 0.0062$ (68% CL)
- perturbative expansion holds down to k_D
- typical prediction of slow-roll models: $\alpha_s = -(1 - n_s)^2$

μ anisotropies and non-Gaussianity

With local non-Gaussianity

In presence of local NG, the heating rate becomes **spatially dependent!**

$$f_{\text{NL}}^{\text{loc.}} = 0$$

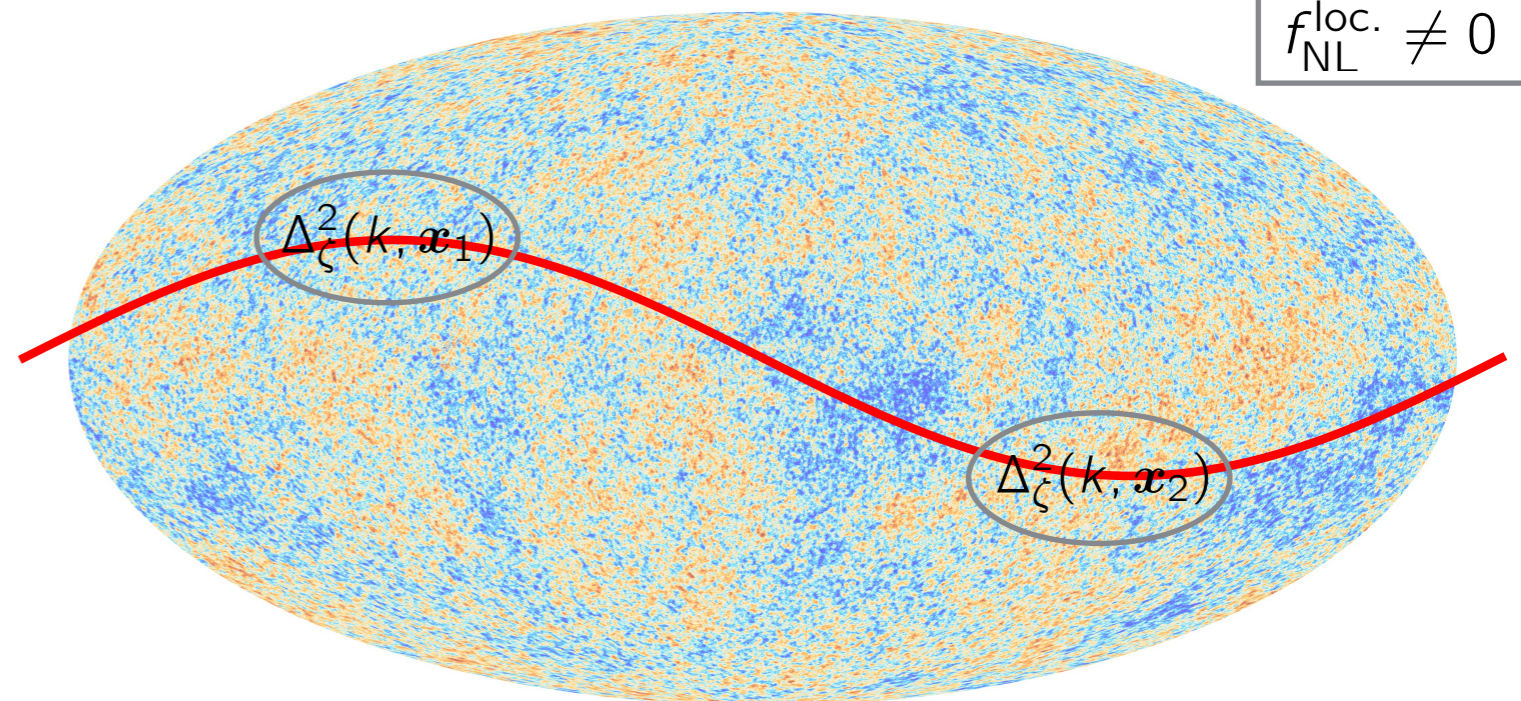


amplitude $\Delta_{\zeta}^2(k_{\text{small}})$
will be the same
everywhere



$$f_{\text{NL}}^{\text{loc.}} \neq 0$$

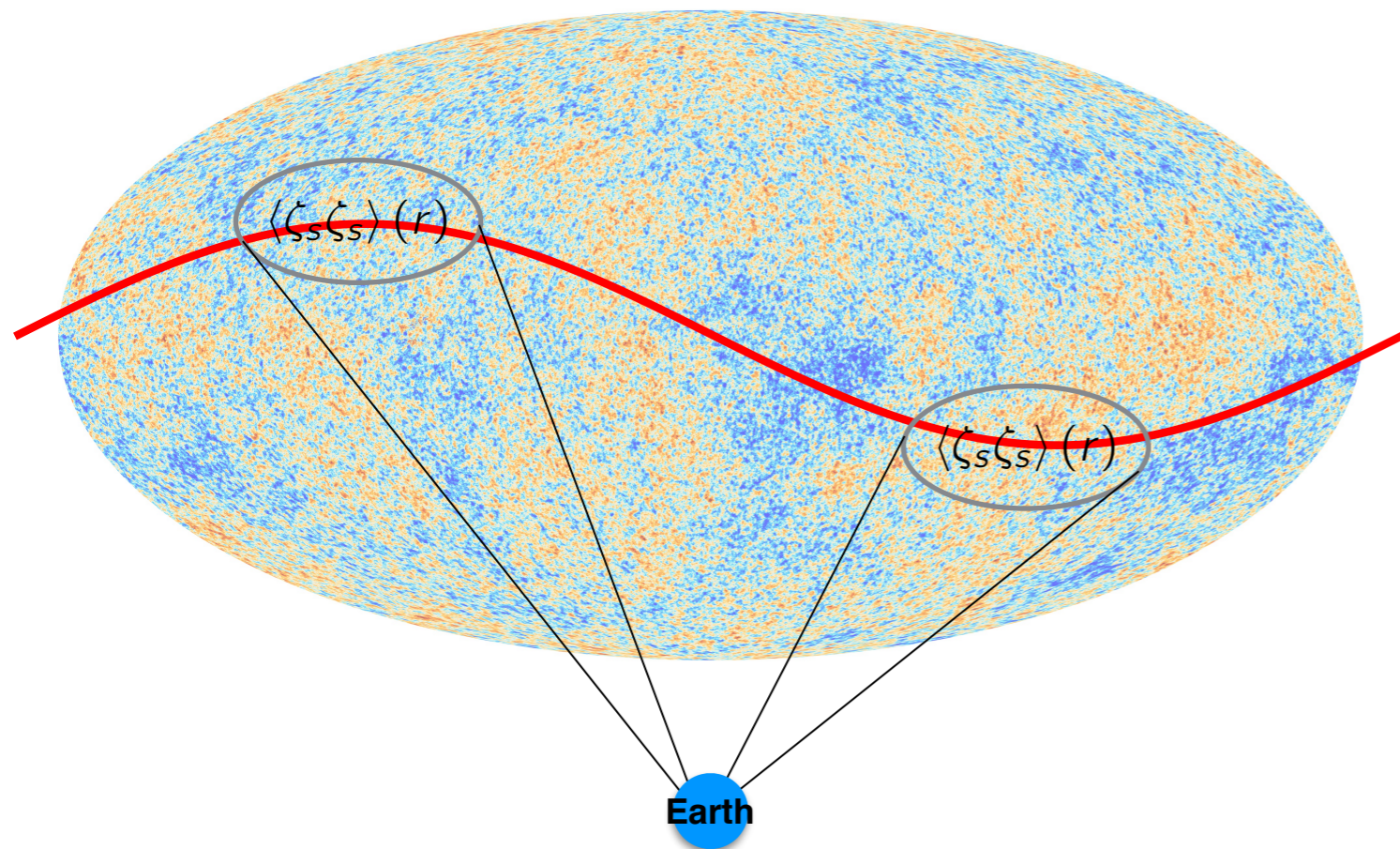
amplitude $\Delta_{\zeta}^2(k_{\text{small}})$
will differ from
one patch to another
→ correlated with ζ_{ℓ}



Short-scale power spectrum is correlated with the long mode \rightarrow **is correlated with the long-wavelength T perturbation!** Angular power spectrum $C_{\ell \lesssim 100}^{\mu T} \approx 12 f_{\text{NL}}^{\text{loc.}} C_{\ell}^{TT}$

However! In every patch a local observer cannot see the coupling with ζ_{ℓ} in single-field inflation! **No μ - T correlations in this case?**

But \rightarrow be careful about projection effects! E.g.: CMB squeezed bispectrum



$$B_{\ell_L, \ell_S}^{TTT} = C_{\ell_L}^{TT} C_{\ell_S}^{TT} \left(2 - \frac{d \log(\ell_S^2 C_{\ell_S}^{TT})}{d \log \ell_S} \right)$$

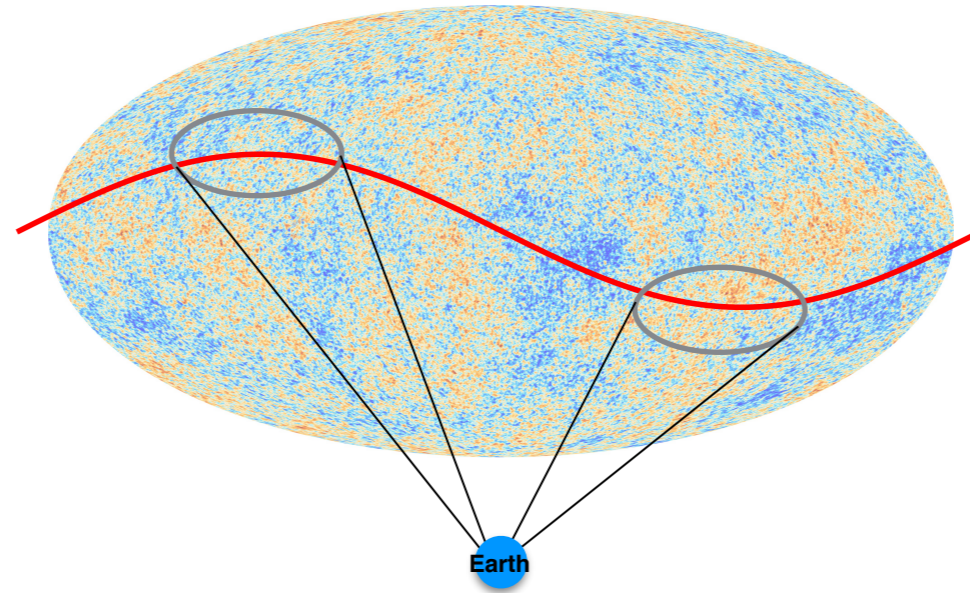
$$\ell_L \lesssim 100 \quad \text{Creminelli, Pitrou, Vernizzi (2011)} \\ \text{Bartolo, Matarrese, Riotto (2011)}$$

any deviation is
 $\propto f_{\text{NL}} - (1 - n_s)$

e.g.:

Pajer, Schmidt, Zaldarriaga (2013)
Mirbabayi and Zaldarriaga (2014)

Projection effects



1. same physical length in the local patch at recombination appears at different angular sizes to the observer
2. photons experience a different redshift in different directions due to the presence of the long mode



1. we are looking at the **average** μ in a patch \rightarrow no "ruler" whose length the long mode can perturb
2. μ does not redshift

$$\frac{Df_\gamma}{d\lambda} = 0 \Rightarrow \begin{cases} T_{\text{obs}} = \frac{E_{\text{obs}}}{E_{\text{rec}}} T_{\text{rec}} \\ \mu_{\text{obs}} = \mu_{\text{rec}} \end{cases}$$

No correlation from evolution from the LSS to the observer... **What is the leading effect?**

Evolution after the end of the μ -era \rightarrow damping of μ **inhomogeneities!**

Pajer and Zaldarriaga, 2012

 no effect on $\langle \mu \rangle \dots$

Then \rightarrow effect of long mode on μ creation! In principle, complicated non-linear dynamics...



Exploit another nice difference with CMB bispectrum: horizon at end of μ era is $\approx 10^{-1} \text{ Mpc}^{-1}$

Up to very high $\ell \gg 100$ we can treat the effect with the **separate universe approach!**

similarly to halo bias...
e.g.: Dai, Pajer, Schmidt, 2015

Long mode modifies the expansion history and adds curvature $K_F \propto \partial^2 \mathcal{R}$ to the FLRW:

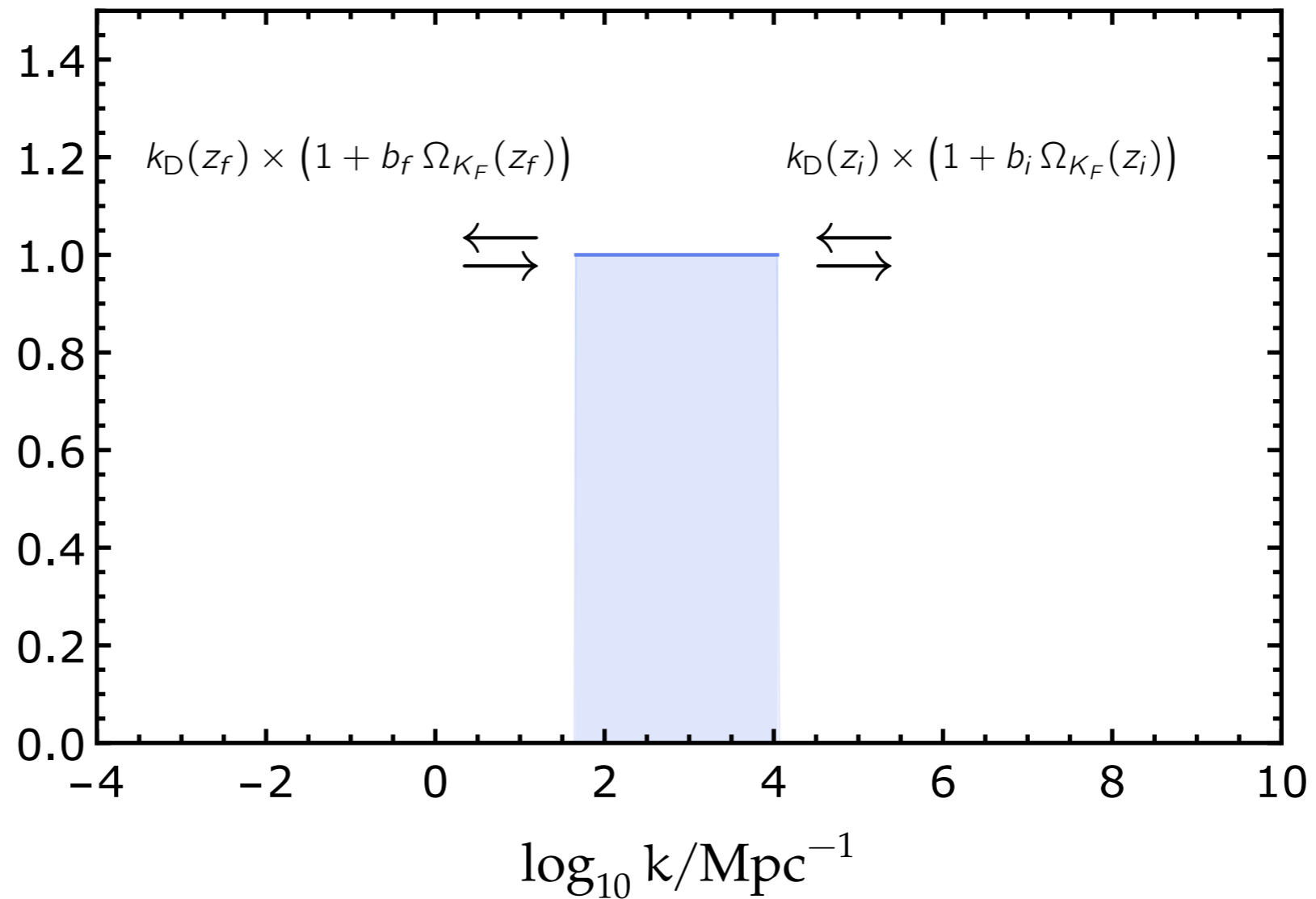
1. modification the duration of the μ era
2. direct effect on the evolution of short modes

GC, Pajer, Schmidt, 2016

3. additional coupling to $K_F \propto \partial^2 \mathcal{R}$ from inflation. It is of order $1 - n_s$: can be neglected unless there are (unlikely) cancellations of 1. and 2.

1.

$$\langle \mu \rangle \approx A_s \times$$



2.

Solve for the evolution of δ_γ in a curved FLRW...

A difference with δ_m : $c_s^2 \neq 0 \rightarrow$ not only corrections from modification to $a = a(K_F)$

Conclusions

- CMB spectral distortions probe small scales $50 \text{ Mpc}^{-1} < k < 10^4 \text{ Mpc}^{-1} \Rightarrow$ can be used to test scale-dependence of primordial power spectrum and the squeezed bispectrum
- good observable to rule out slow-roll inflation!
- in single-field inflation (practically) no “contamination” on μ - T correlations:
 - no projection effects from last-scattering to observer
 - no effects from tight-coupling evolution
 - only effect: “bias” on μ production. Suppressed by $k_L^2/\mathcal{H}(z_f)^2 \dots$