

UHECR interactions

... and the production of astrophysical neutrinos

Credit: Steven Saffi

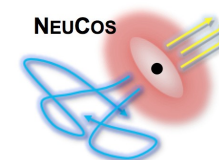
Daniel Biehl

DESY Zeuthen, Germany

20th ISVHECRI, Nagoya University

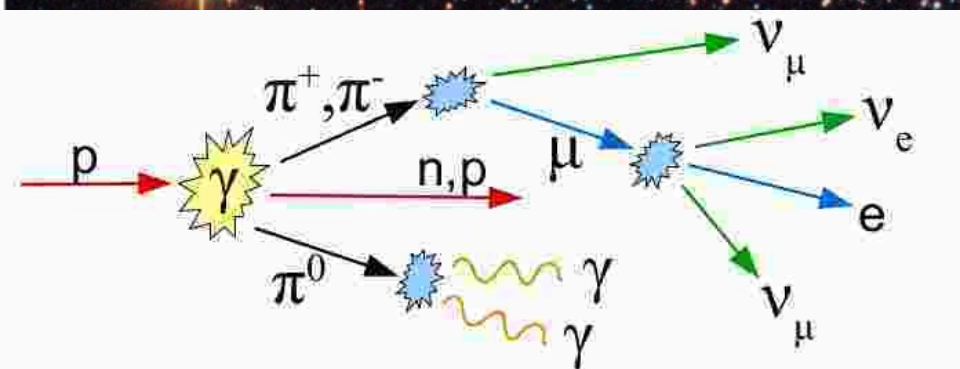
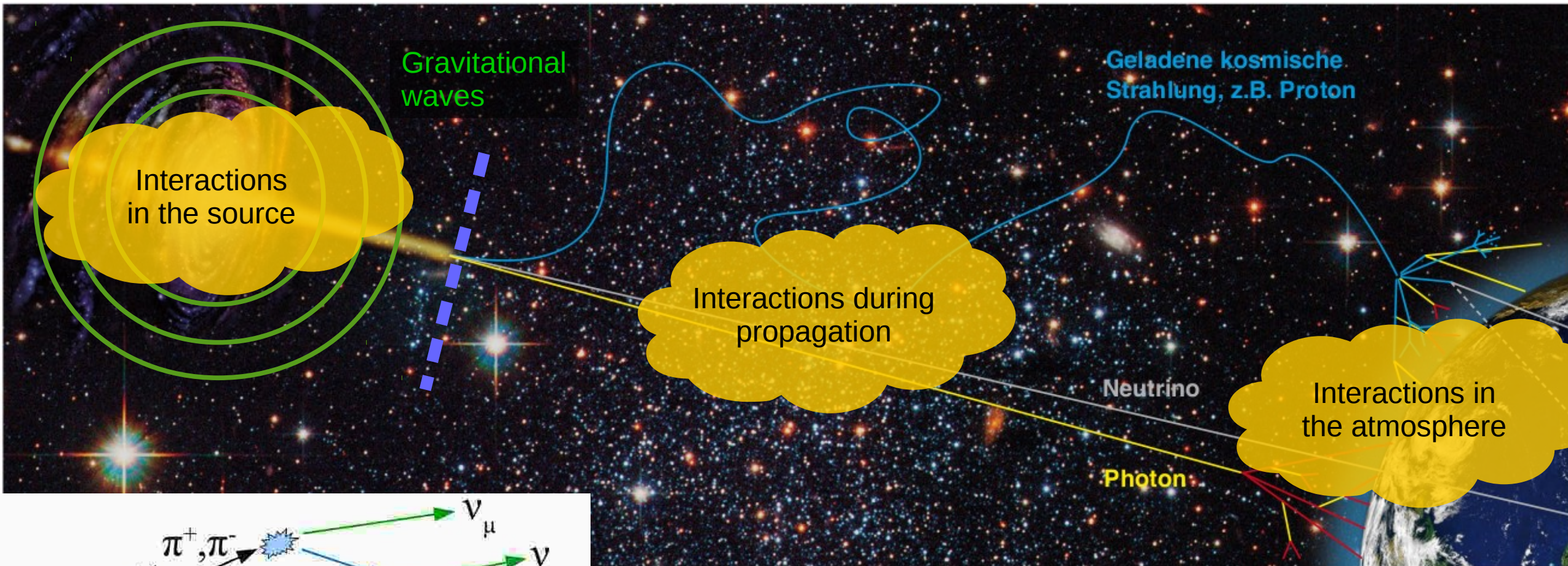
May 21, 2018

HELMHOLTZ RESEARCH FOR
GRAND CHALLENGES



Interactions of UHECRs in the multi-messenger context

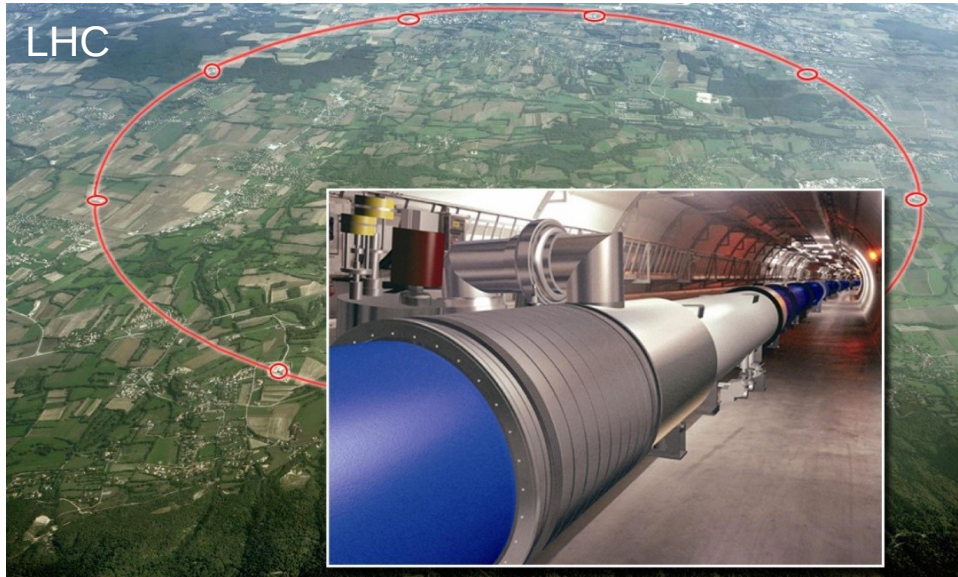
The connection between neutrinos and ultra-high energy cosmic rays



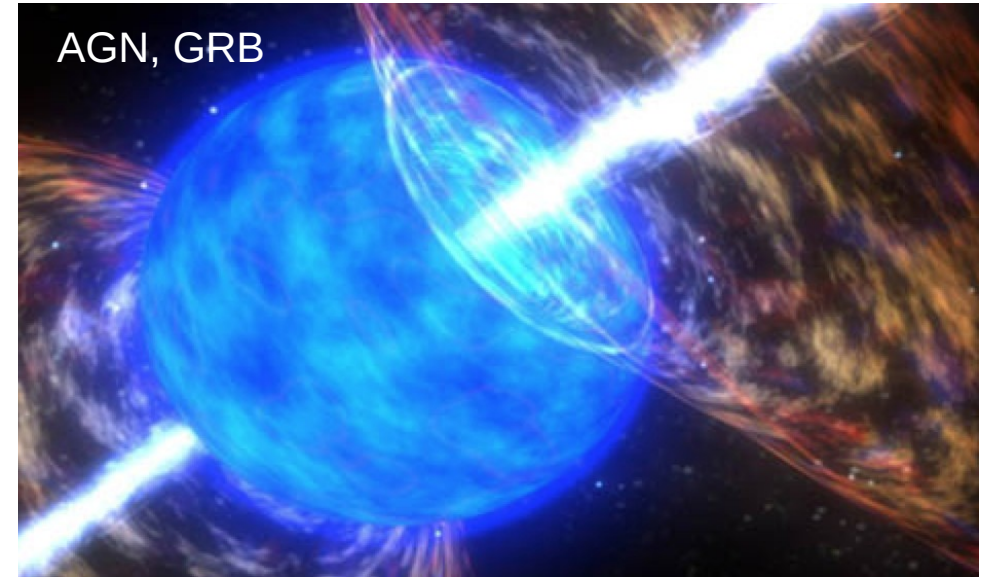
Multi-messenger interpretations rely on our understanding of the involved interactions!

Maximum reachable energy

Why size matters



$$E_{\max} \sim q B R$$



Terrestrial particle accelerators

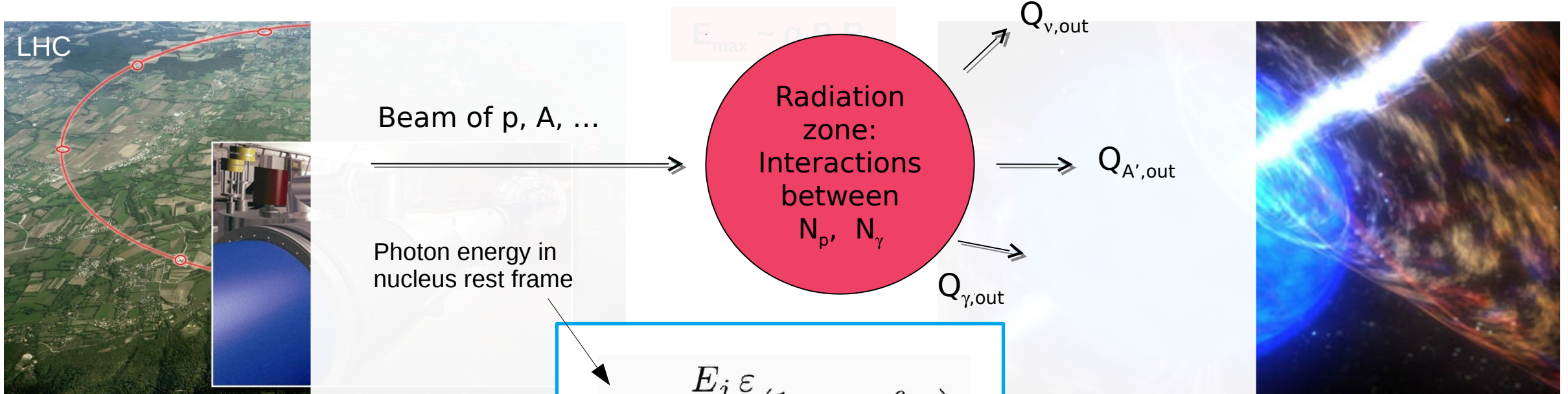
- $B > 8 \text{ T}$
- $R \sim 4.3 \text{ km}$
- $E_{\max} \sim 13 \text{ TeV}$

Cosmic particle accelerators

- $B \sim 1 \text{ mT} - 1 \text{ T}$
- $R \sim 100,000 - 10,000,000,000 \text{ km}$
- $E_{\max} \sim 300,000,000 \text{ TeV}$

Maximum reachable energy

Why size matters



$$\epsilon_r = \frac{E_j \epsilon}{m_j} (1 - \cos \theta_{j\gamma})$$

high energy nucleus
+
“low” energy photon
=
photo-nuclear physics in
MeV – GeV range

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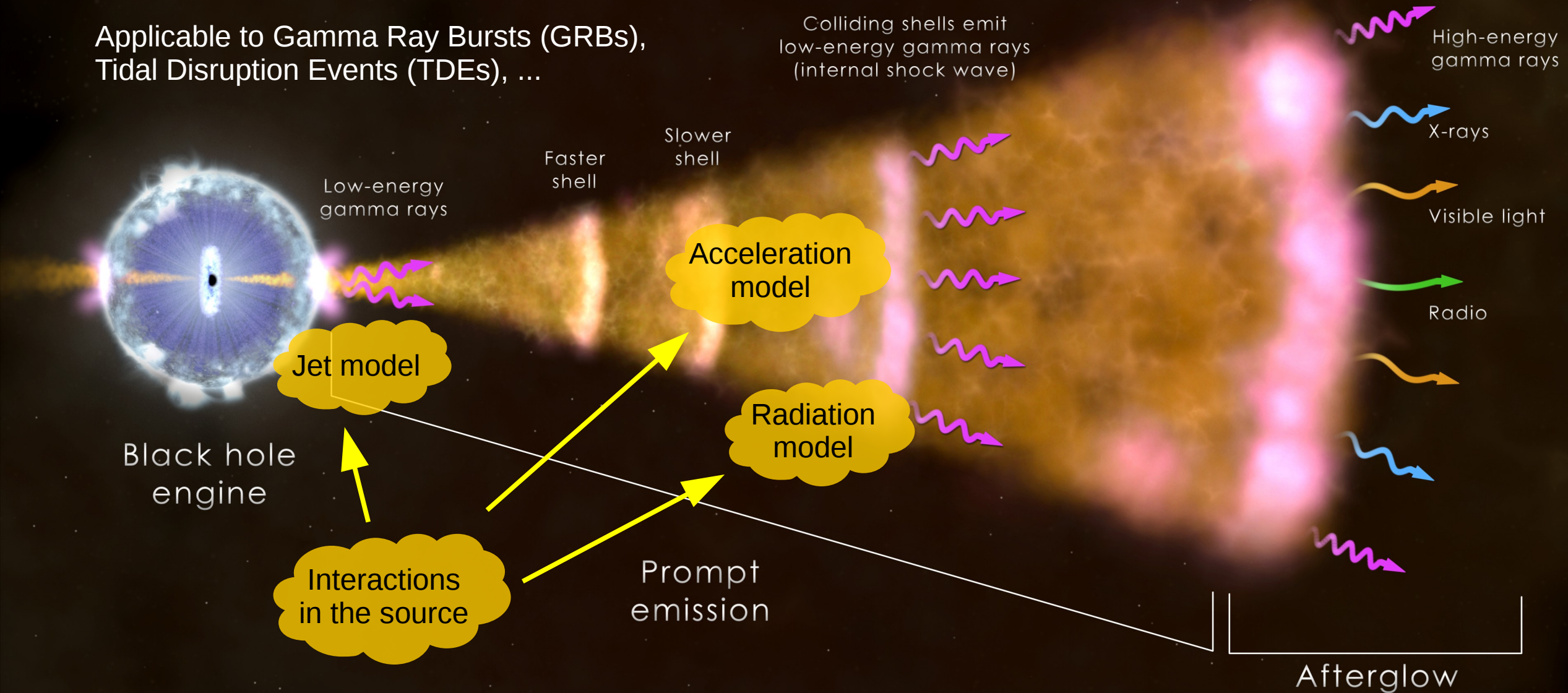
UHECR interactions with ambient photons

A possible scenario for a generic $\text{p}\gamma$ source

Applicable to Gamma Ray Bursts (GRBs),
Tidal Disruption Events (TDEs), ...

Jet collides with
ambient medium
(external shock wave)

Colliding shells emit
low-energy gamma rays
(internal shock wave)

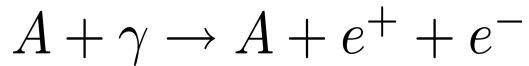


Relevant energy scales for UHECR interactions with photons

High-energy nucleus + 'low energy' photon = photo-nuclear physics in the MeV – GeV range

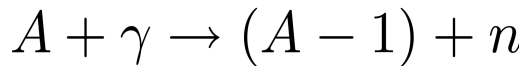
Photo-hadronic (A γ) interactions

- **QED scale:** e.g. pair-production



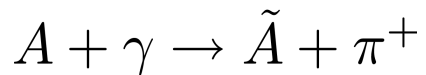
Photon energy in nucleus rest frame
 $\varepsilon_r > 1 \text{ MeV}$

- **Nuclear scale:** nuclear photo-disintegration, e.g.



$\varepsilon_r > 8 \text{ MeV}$

- **Mesonic scale:** baryonic resonances, photo-meson production (produces neutrinos), e.g.



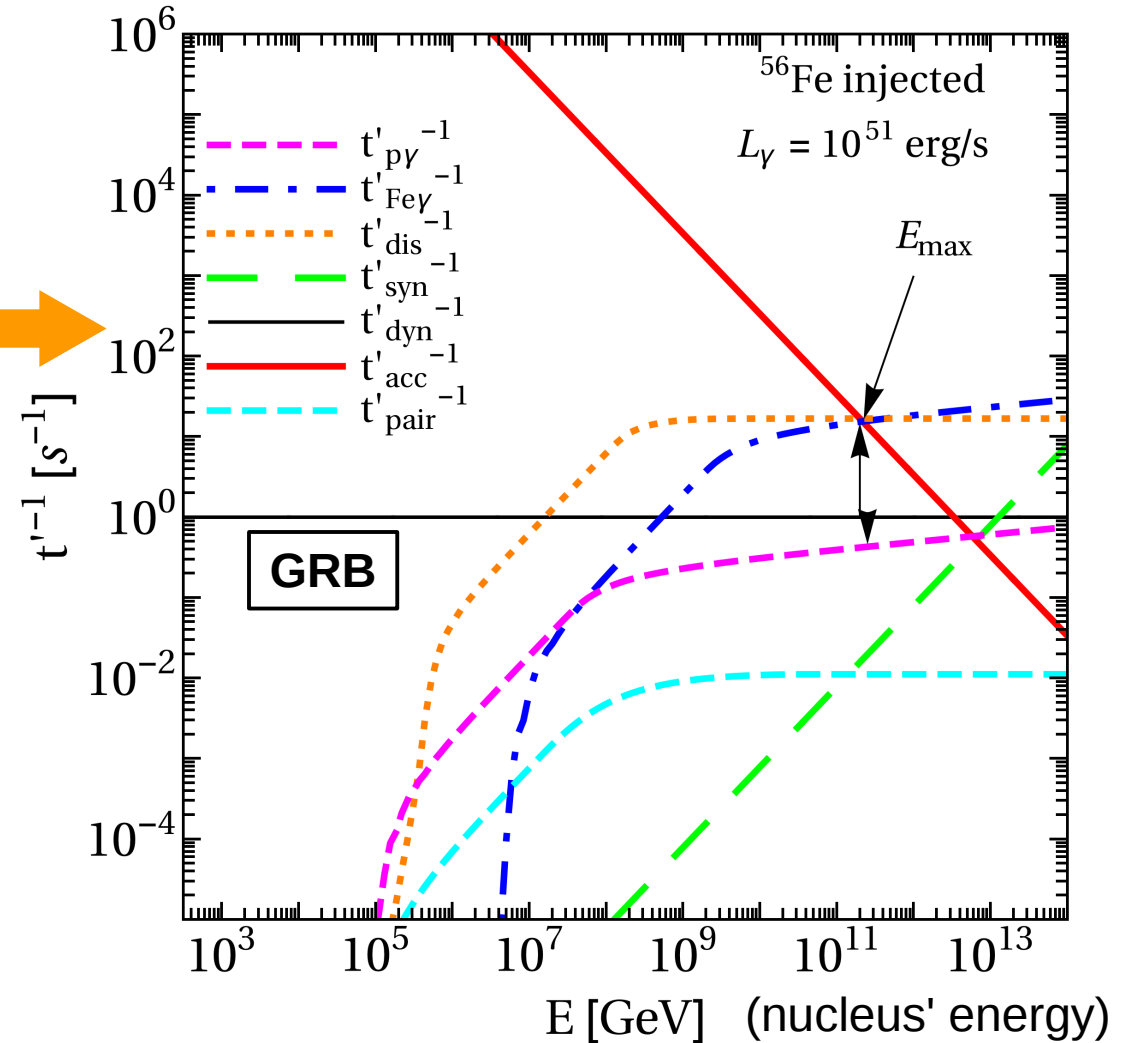
$\varepsilon_r > 140 \text{ MeV}$

- **Hadronic scale:** hadronic structure becomes relevant for the interaction

$\varepsilon_r > 1 \text{ GeV}$

Other processes

- Beta-decays, pp-interactions, spontaneous nucleon emission, spallation, de-excitation, ...



[DB, D. Boncioli, A. Fedynitch, W. Winter – A&A (2018)]

Relevant energy scales for UHECR interactions with photons

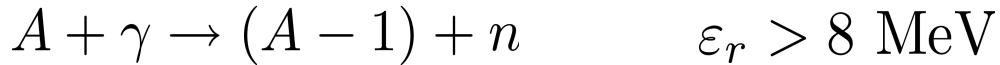
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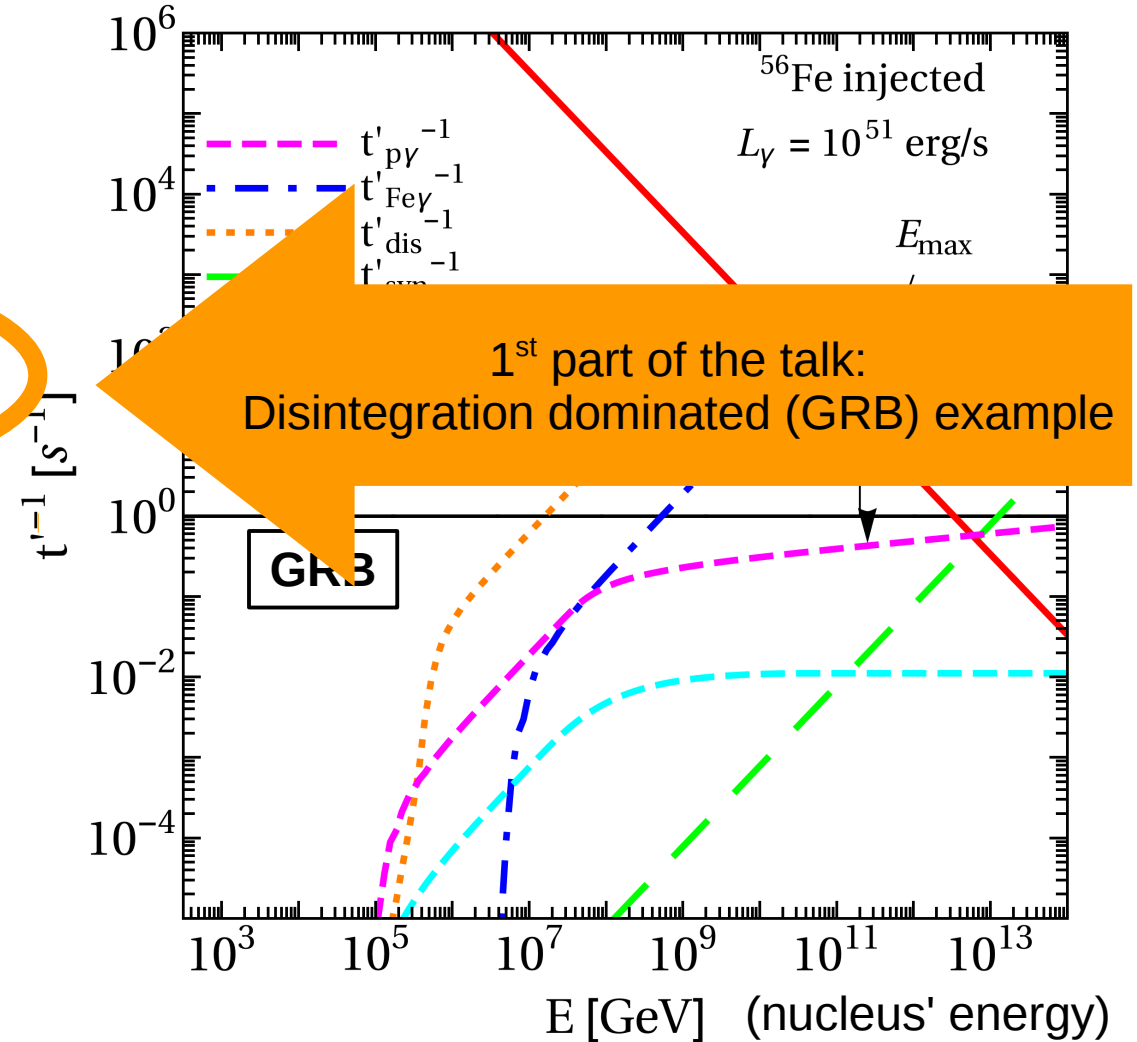
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$$\epsilon_r > 1 \text{ GeV}$$

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Photon energy in nucleus rest frame



[DB, D. Boncioli, A. Fedynitch, W. Winter – A&A (2018)]

Development of the nuclear cascade

A qualitative and quantitative representation of interactions

Triggering the nuclear cascade

- Example: pure iron injected in a GRB shell, different luminosities
- Development of the nuclear cascade scales with the photon density

$$u'_\gamma \sim \frac{L_\gamma}{\Gamma^2 R^2}$$

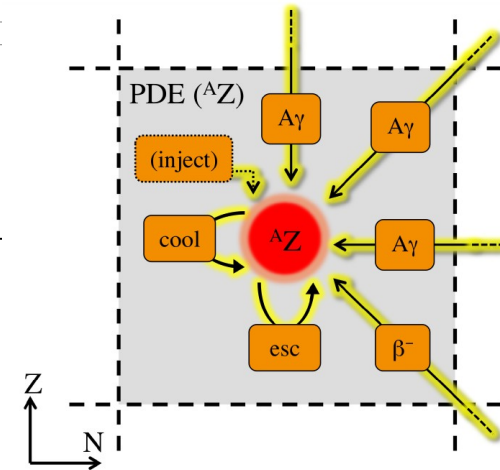
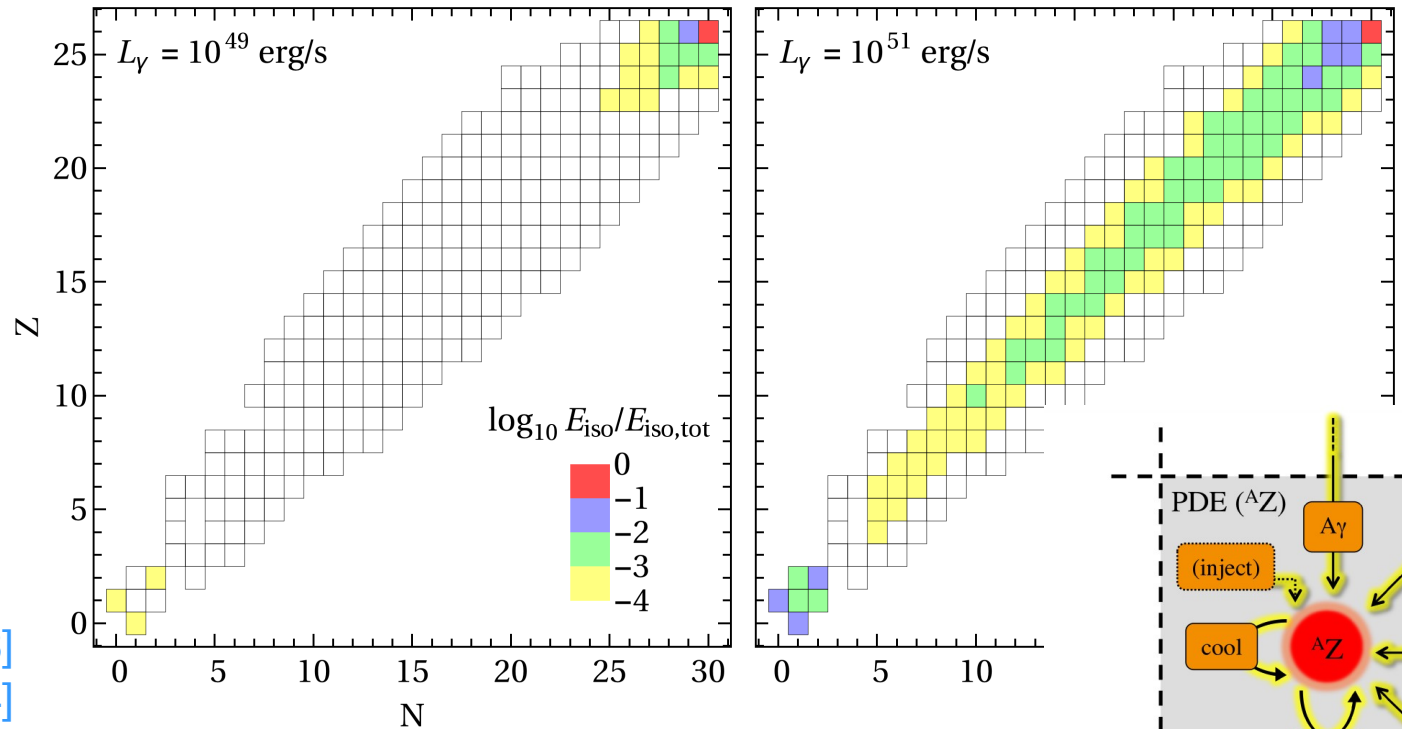
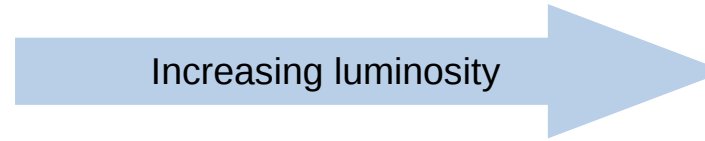
- Internal shock scenario: $R \simeq 2\Gamma^2 \frac{ct_v}{1+z}$

$$\longrightarrow f_{p\gamma} \propto L_\gamma / (\Gamma^4 t_v \epsilon_{\gamma,br})$$

[Waxman, Bahcall, 1998]

[Guetta et al., 2004]

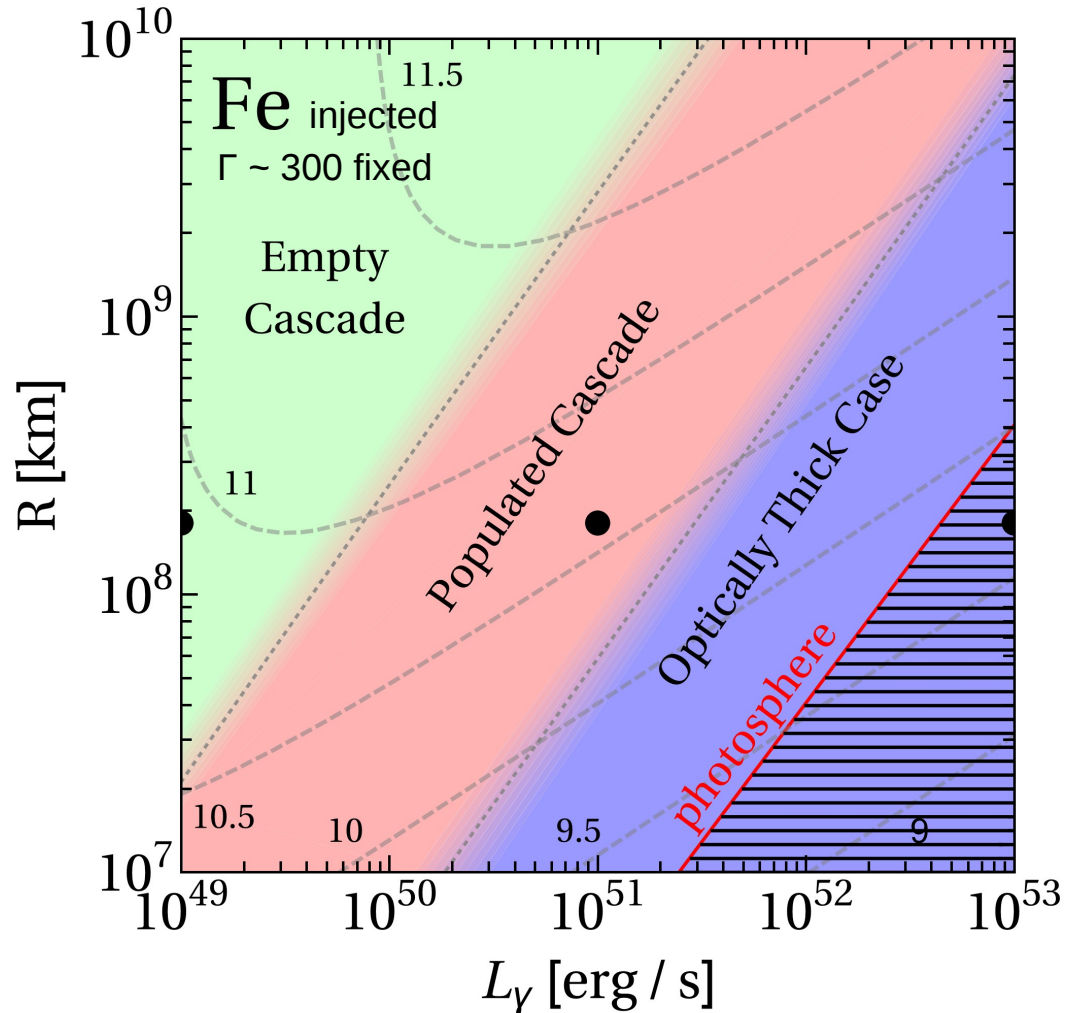
- Production radius R and luminosity L are the main control parameters for the nuclear cascade and neutrino production



[DB, D. Boncioli, A. Fedynitch, W. Winter – A&A (2018)]

The nuclear cascade in the parameter space

Classification of interactions regions



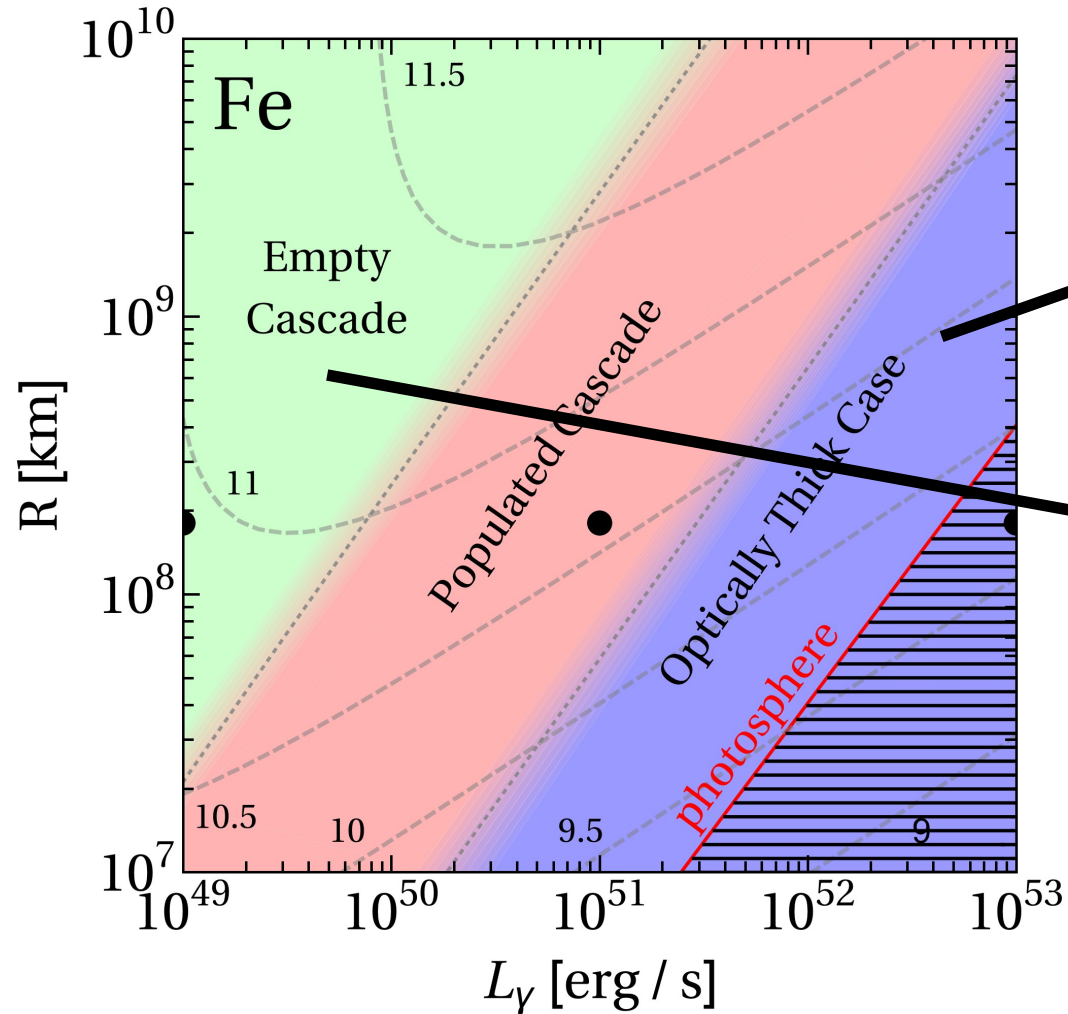
Parameter space regions for interactions

- **Empty cascade:**
 - Optically thin to photo-hadronic interactions of all species
 - Low neutrino production, nuclear cascade does not develop
- **Populated cascade:**
 - Optically thick to nuclei heavier than protons
 - Intermediate neutrino production, broad cascade develops
- **Optically thick case:**
 - Optically thick to all particles
 - High neutrino production, narrow cascade along the main diagonal

[DB, D. Boncioli, A. Fedynitch, W. Winter – A&A (2018)]

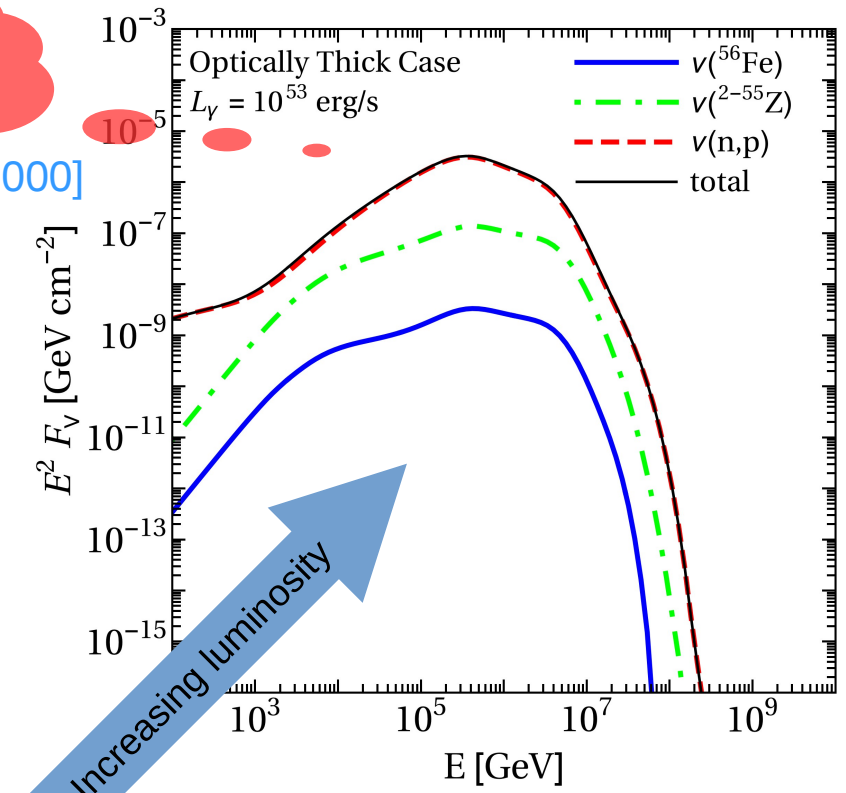
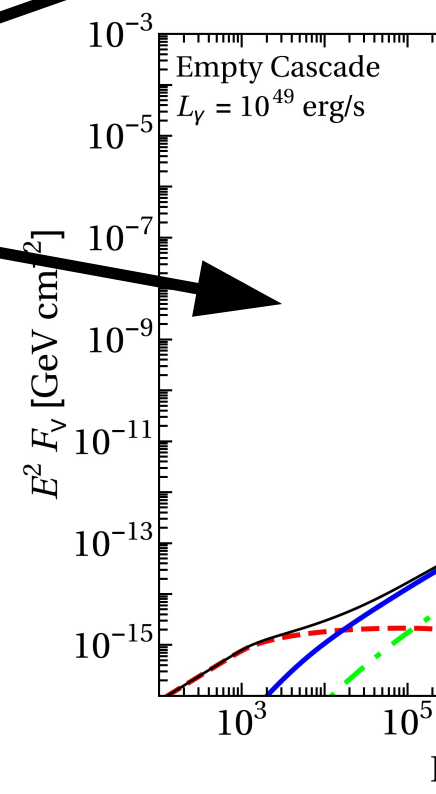
UHECR interactions as a measure for neutrino production

Meaning of the regions for neutrino production



[Mücke et al., SOPHIA, 2000]

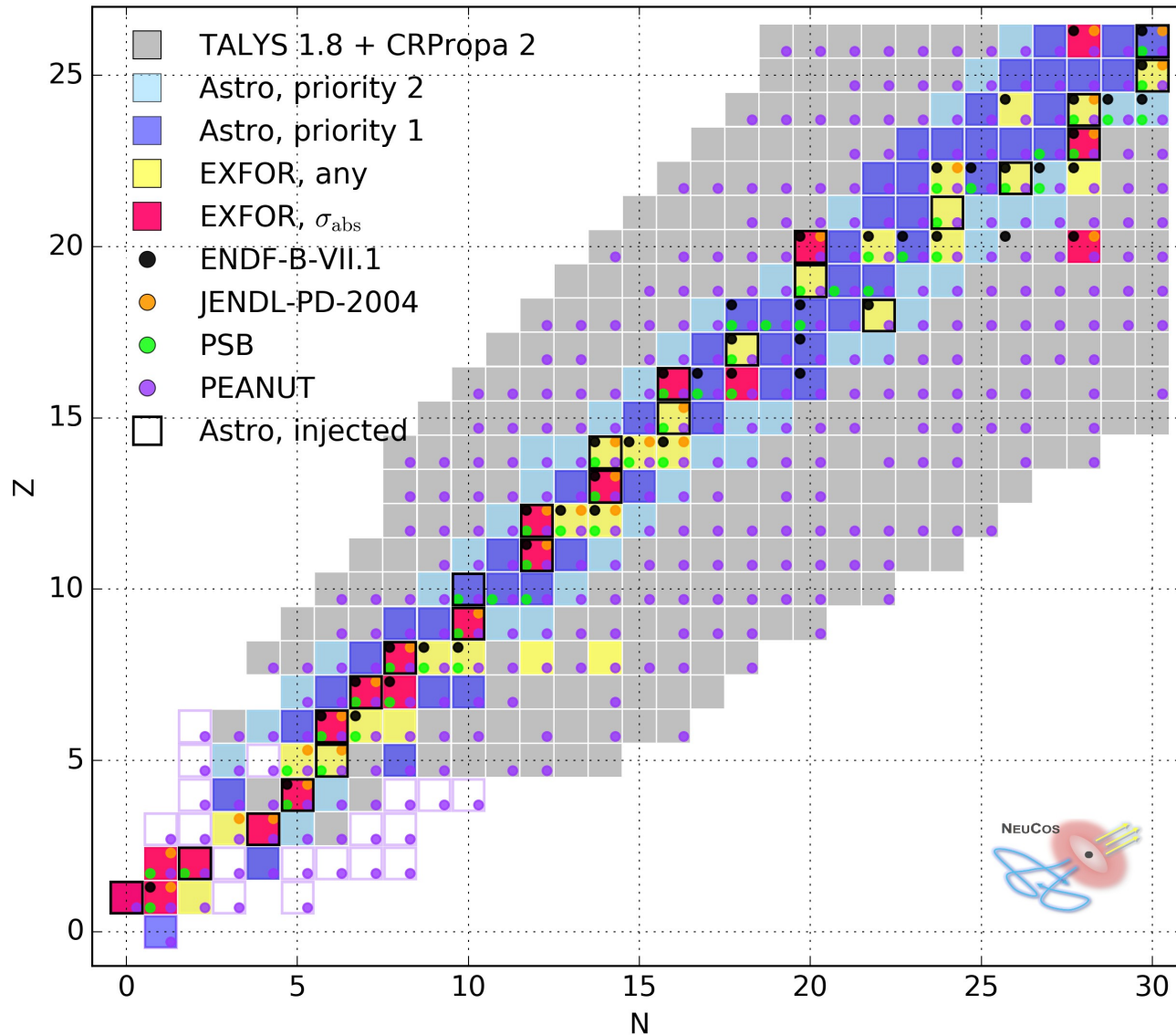
Well-understood cross-section



Poorly understood cross-section

Current situation on experimental data and theoretical models

Importance of future measurements and improved models



EXFOR data base cross-sections

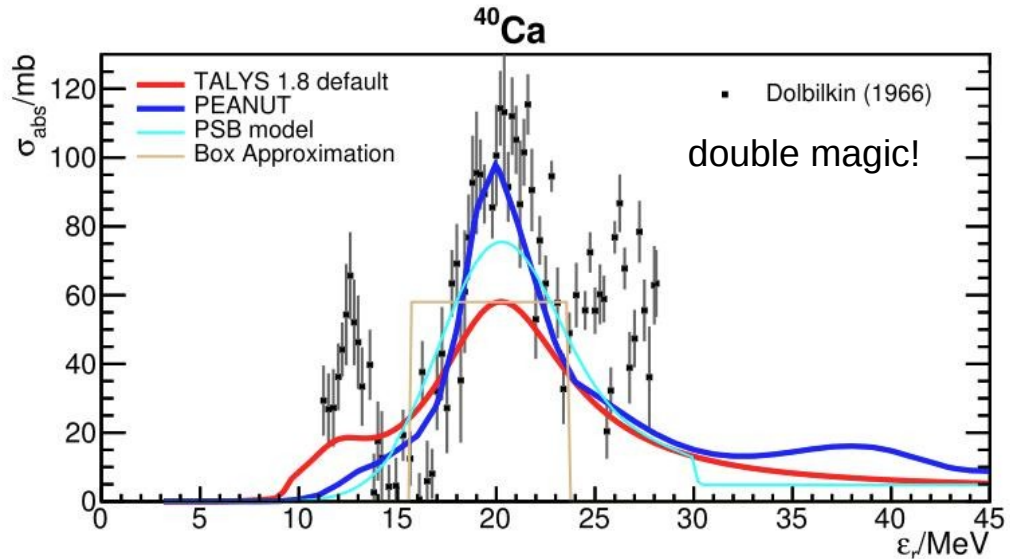
[\[https://www-nds.iaea.org/exfor/exfor.html\]](https://www-nds.iaea.org/exfor/exfor.html)

- Cross-sections only measured for very few isotopes (red)
- Located mostly on main diagonal (stable elements)
- All other isotopes need models prediction
→ not always well in reproducing the data
- **Need future measurements and improved models**

[\[D. Boncioli, A. Fedynitch, W. Winter – Sci. Rep. \(2017\)\]](#)

Impact of nuclear cross-section data and models

Large uncertainties originating from nuclear physics



Comparison between different models

- TALYS (CRpropa 2+ - style) predictions not / weakly depending on nuclear mass and element, e.g. ⁴⁰Ca is double magic, ⁴⁰Ar (no data) is not, so no reason for cross sections to be equal

[A. J. Koning et al., 2007]

[K.-H. Kampert et al., 2005]

- PEANUT (a module of FLUKA) predictions are different in the same isobar, if data available at least the central GDR peak is reproduced

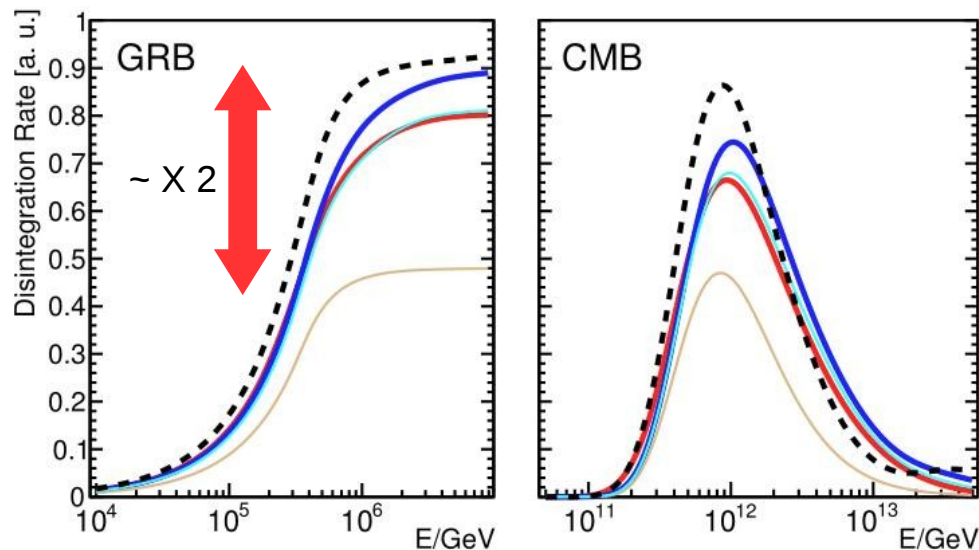
[A. Ferrari et al., 2005]

- Box approximation, e.g. used in [Murase, Beacom 2010] underestimates data and models, insufficient description

- Up to factor two differences in disintegration rates**

[D. Boncioli, A. Fedynitch, W. Winter – Sci. Rep. (2017)]

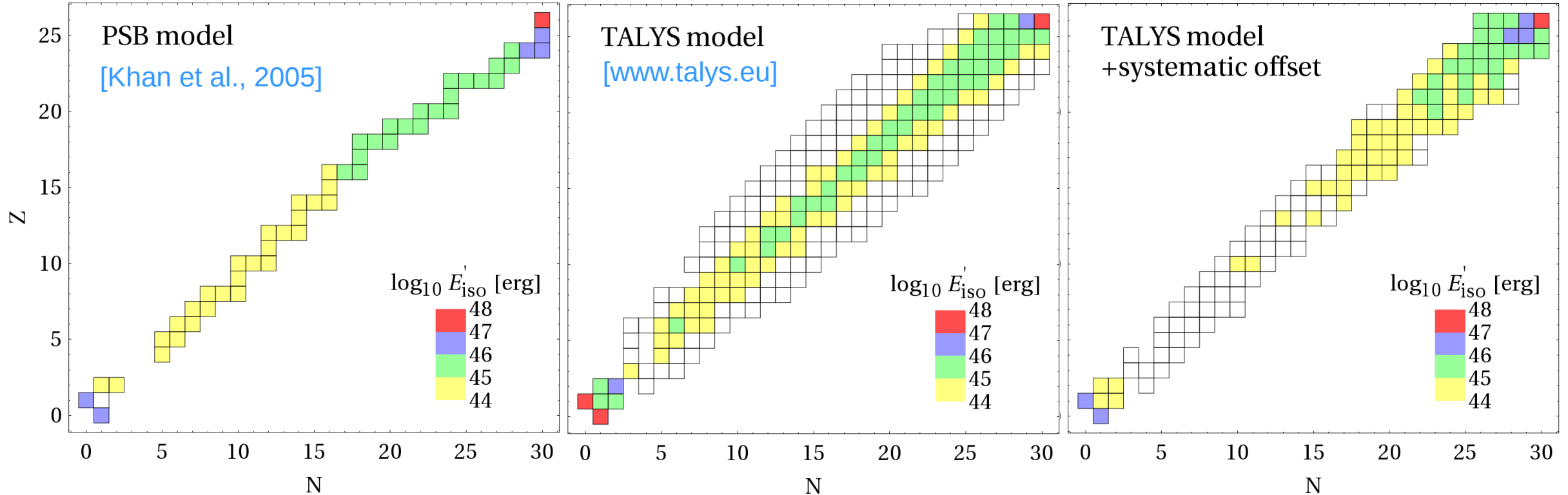
see also e.g. [Soriano et al., 1805.00409]



Impact of different models on the nuclear cascade

Disintegration strongly depending on interaction models

[D. Boncioli, A. Fedynitch, W. Winter – Sci. Rep. (2017)]



- PSB disintegration chain weakly describes multi-nucleon emission, only small fragments can be ejected
- TALYS provides much more channels, ejection of p, n, d, t, He-3, He-4
- Systematic offset = 'do not trust unmeasured cross-sections' → **cascade will not be populated** → **ejection composition!**

see also e.g. [Alves Batista et al., JCAP 2015]

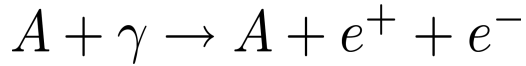
[Pierre Auger collaboration, JCAP 2017]

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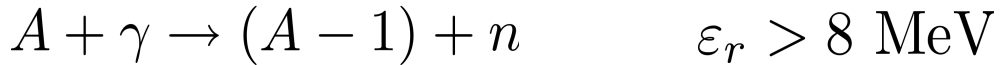
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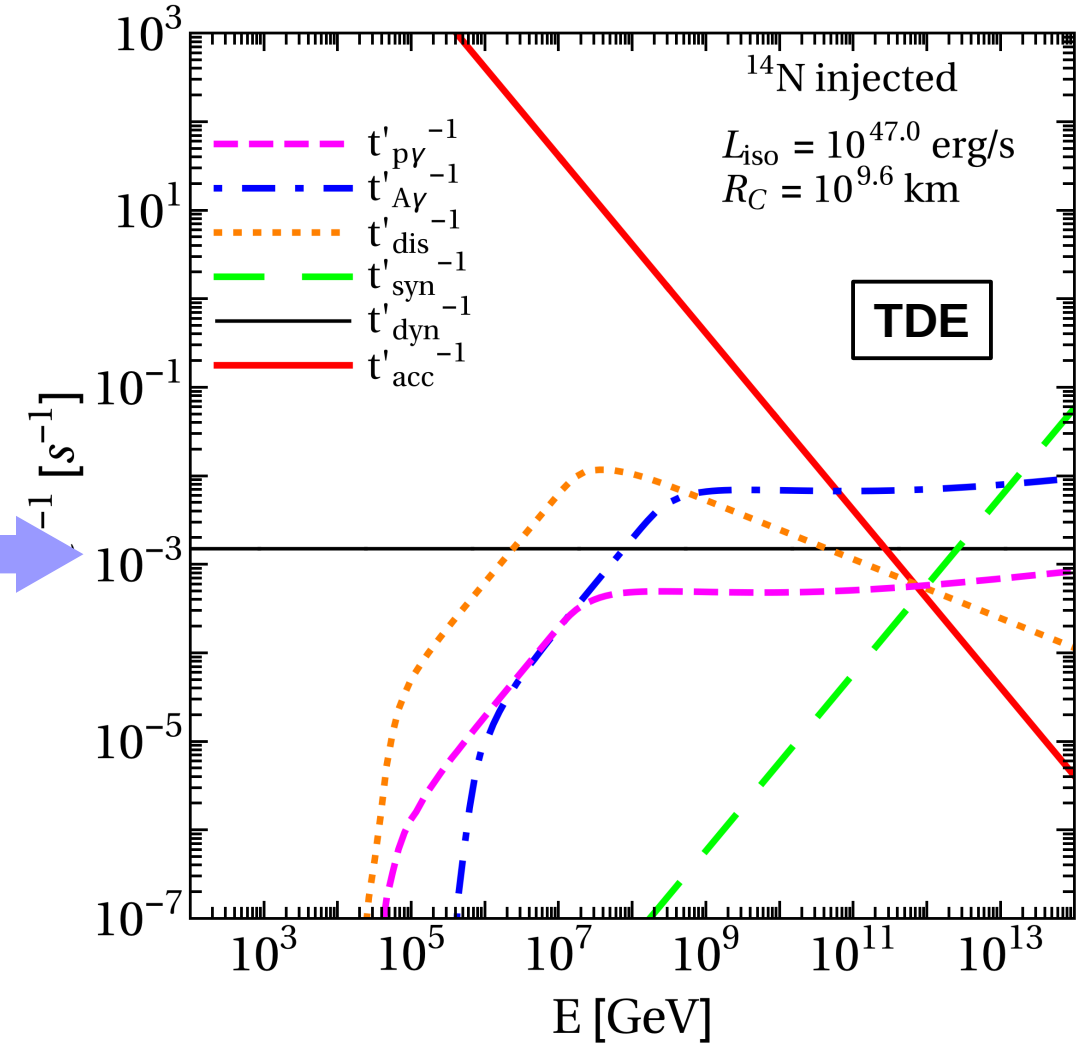
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Other processes

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TDEs as origin of UHECRs AND PeV neutrinos



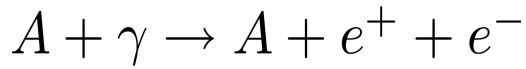
[DB, D. Boncioli, C. Lunardini, W. Winter – arXiv:1711.03555]

Photo-meson production dominated objects

A few examples

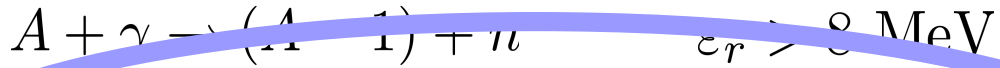
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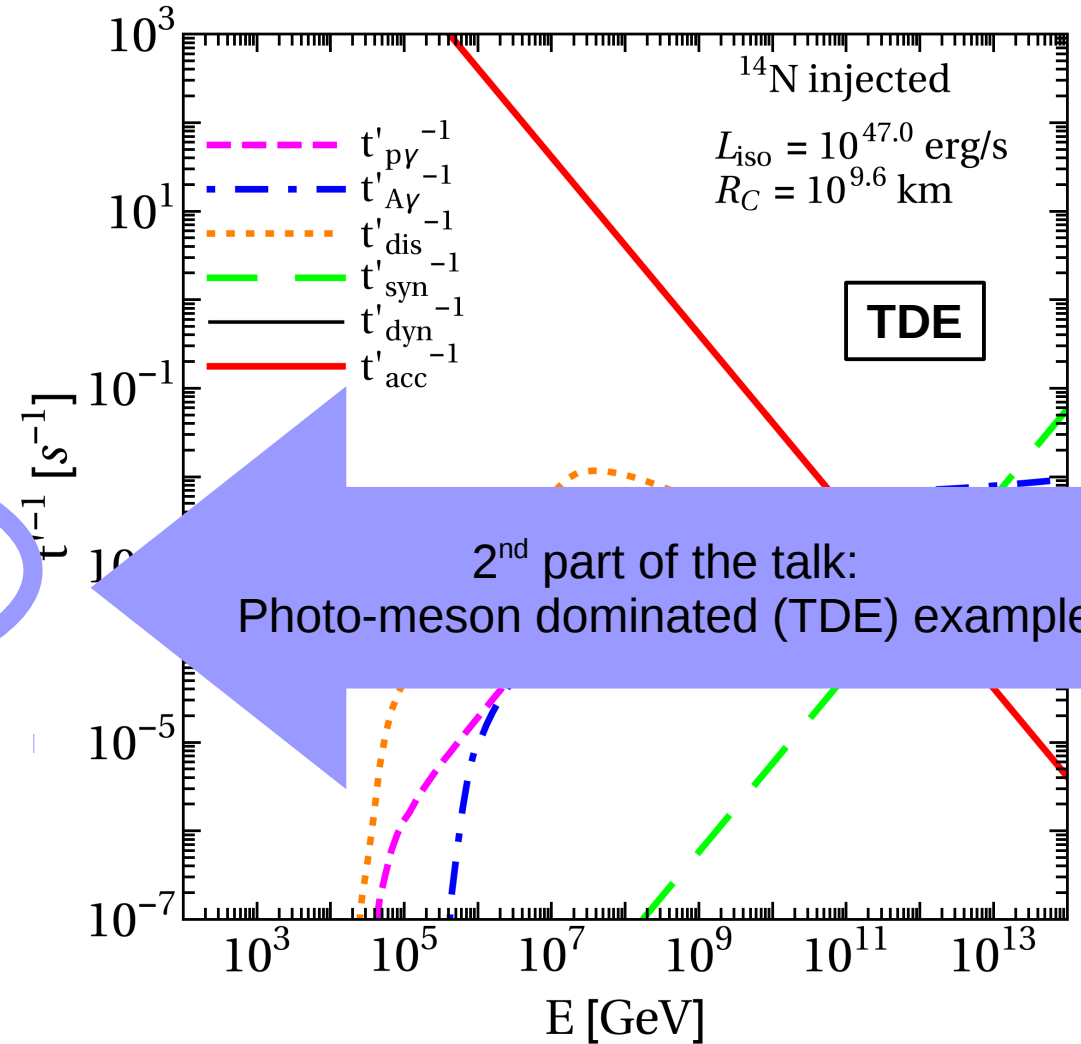
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Cross-section models for photo-meson production

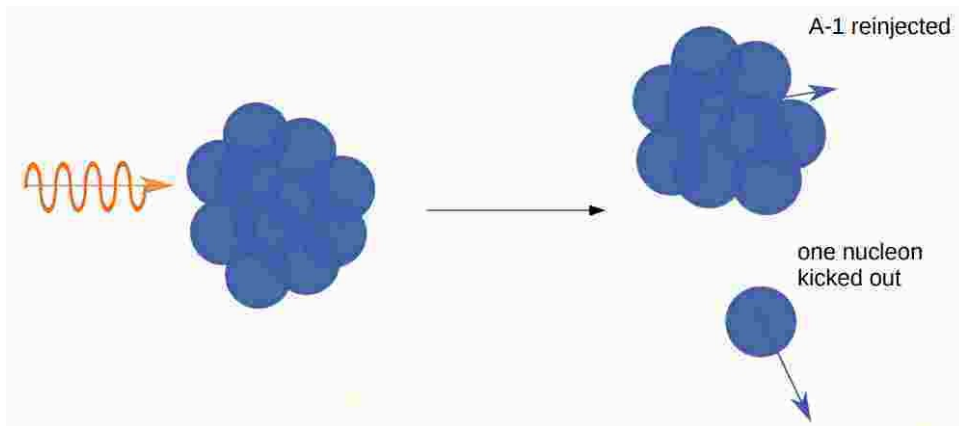
Comparison of different approaches

Superposition model vs. universal curve

- Current state-of-the-art photo-meson models: superposition model with individual nucleon interaction

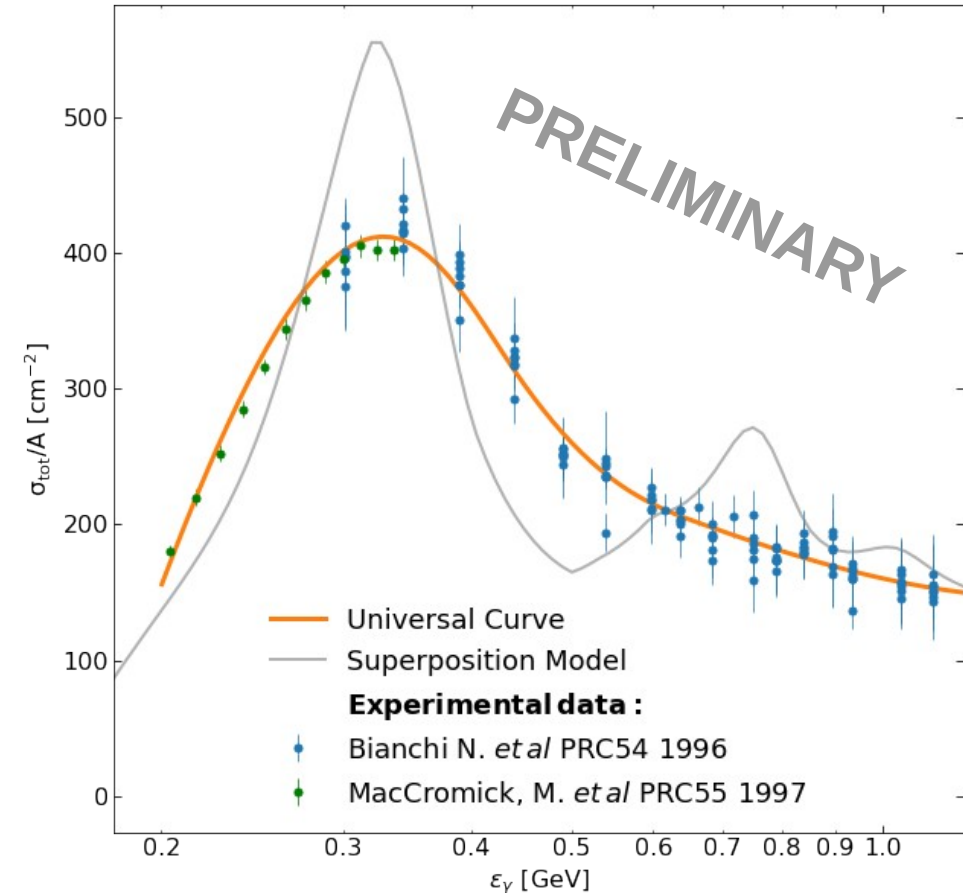
$$\sigma_A^{\text{tot}}(E) = \frac{Z}{A}\sigma_p(E) + \frac{N}{A}\sigma_n(E)$$

- Remaining nucleus (A-1) reinjected with no mediating de-excitations or decays



- CRpropa uses scaling $\sim A^{2/3}$ for the whole energy range

Instead: universal behaviour for nuclei observed, spline interpolation of data to obtain universal curve



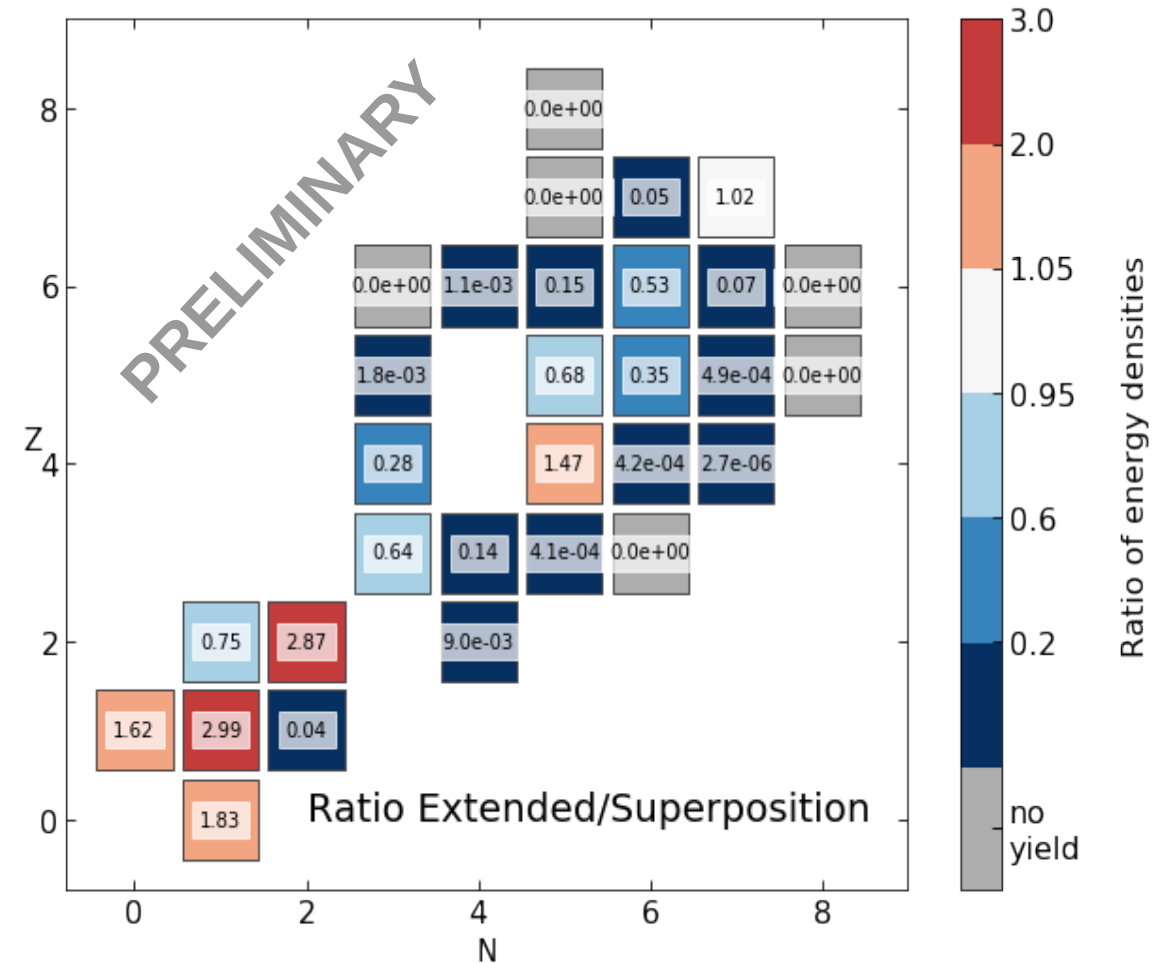
[L. Morejon et al., in preparation]

Impact of photo-meson models on nuclear cascade

Comparison between extended and superposition model

Disintegration chain and ejected composition

- Example: pure nitrogen injected in TDE (zoom in on the low-mass tail of the nuclear cascade shown before)
- Extended model allows for ejection of multi-nucleon fragments
 - **More energy** for specific channels along the main diagonal
 - **Less energy** off the main diagonal, some isotopes basically not populated anymore
- **Direct impact on neutrino production!**



[L. Morejon et al., in preparation]

Glashow resonance: $\bar{\nu}_e$ vs. pp interactions

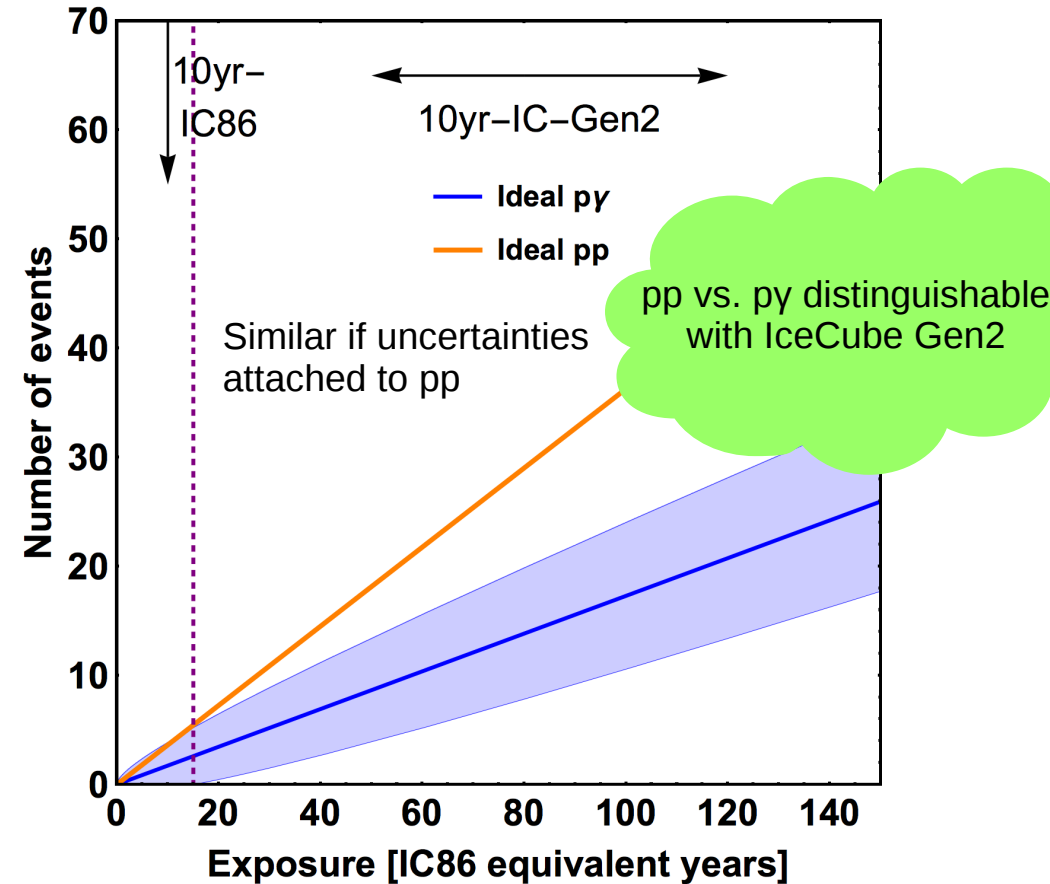
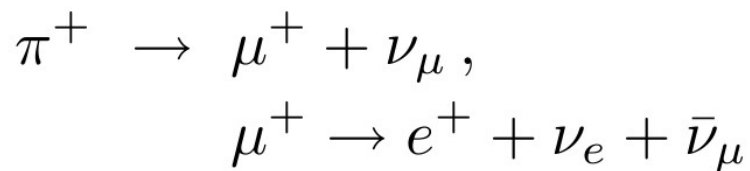
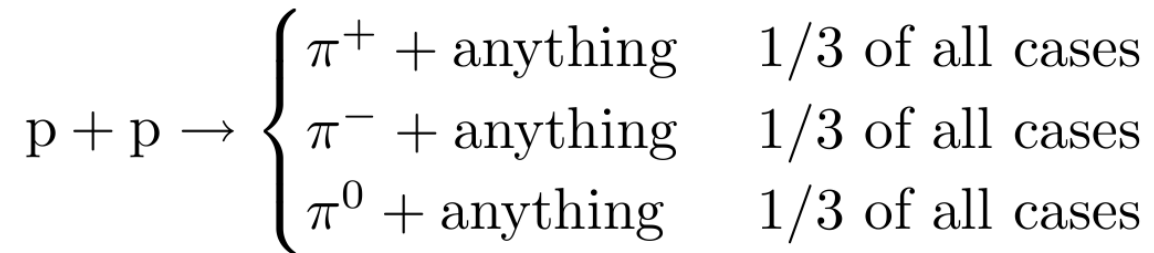
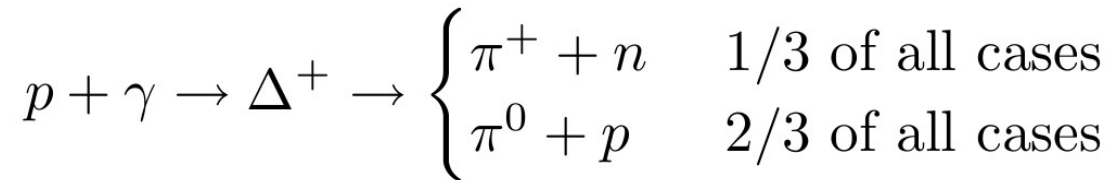
More realistic models making the difference

Triggering the resonance and flavor composition

- Resonant interaction of an astrophysical electron antineutrino with electron



- Origin of the incoming neutrino



[DB, A. Fedynitch, A. Palladino, T. Weiler, W. Winter – JCAP (2017)]

Glashow resonance: $p\gamma$ vs. pp interactions

More realistic models making the difference

More realistic description: pp source

- Pion charge ratio not exactly equal to one, reasonable estimate from hadronic interaction models: average of EPOS-LHC, QGSJet-II-04 and SIBYLL 2.3

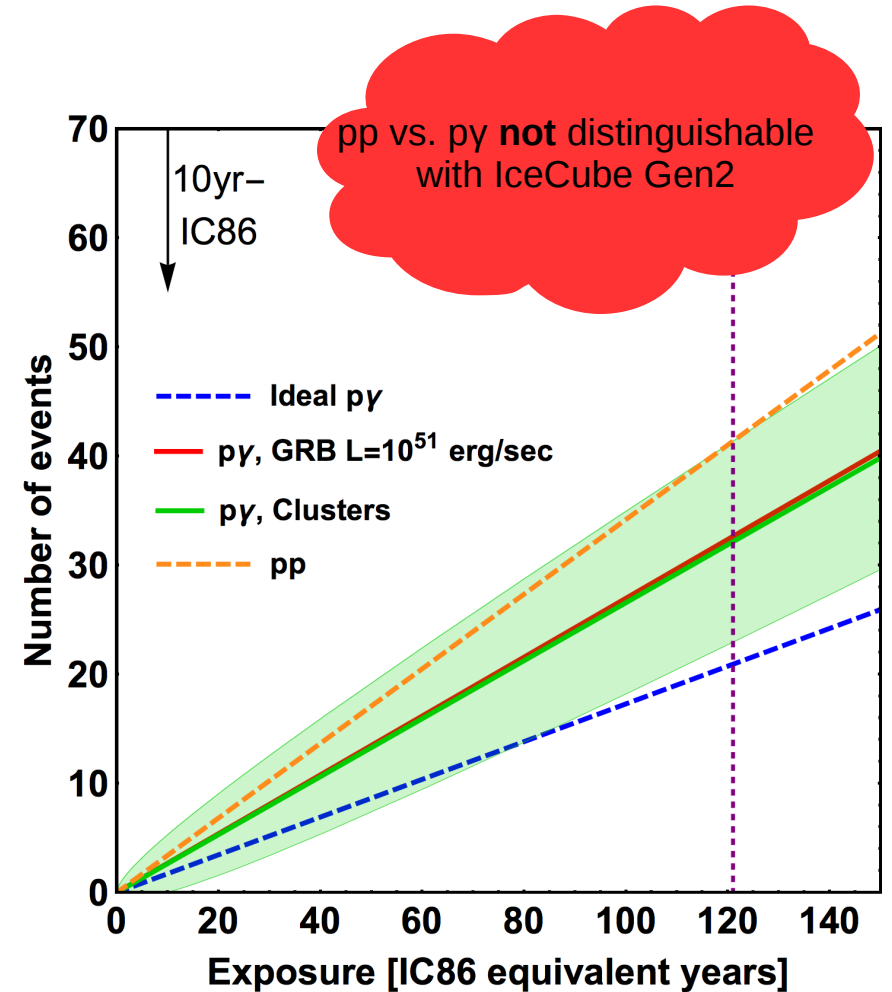
α	N_{π^+}/N_{π^-}
2.0	1.25 ± 0.20
2.3	1.34 ± 0.15
2.6	1.43 ± 0.24

Gets even worse with softer injection spectrum

More realistic description: $p\gamma$ source

- Multi-pion processes and neutrons from the disintegration of nuclei (if present) lead to contamination by π^-

$$n + \gamma \rightarrow \Delta^0 \rightarrow \begin{cases} \pi^- + p & 1/3 \text{ of all cases} \\ \pi^0 + n & 2/3 \text{ of all cases} \end{cases}$$



[DB, A. Fedynitch, A. Palladino, T. Weiler, W. Winter – JCAP (2017)]

... but $p\gamma$ vs. $A\gamma$ might work!

Summary and conclusions

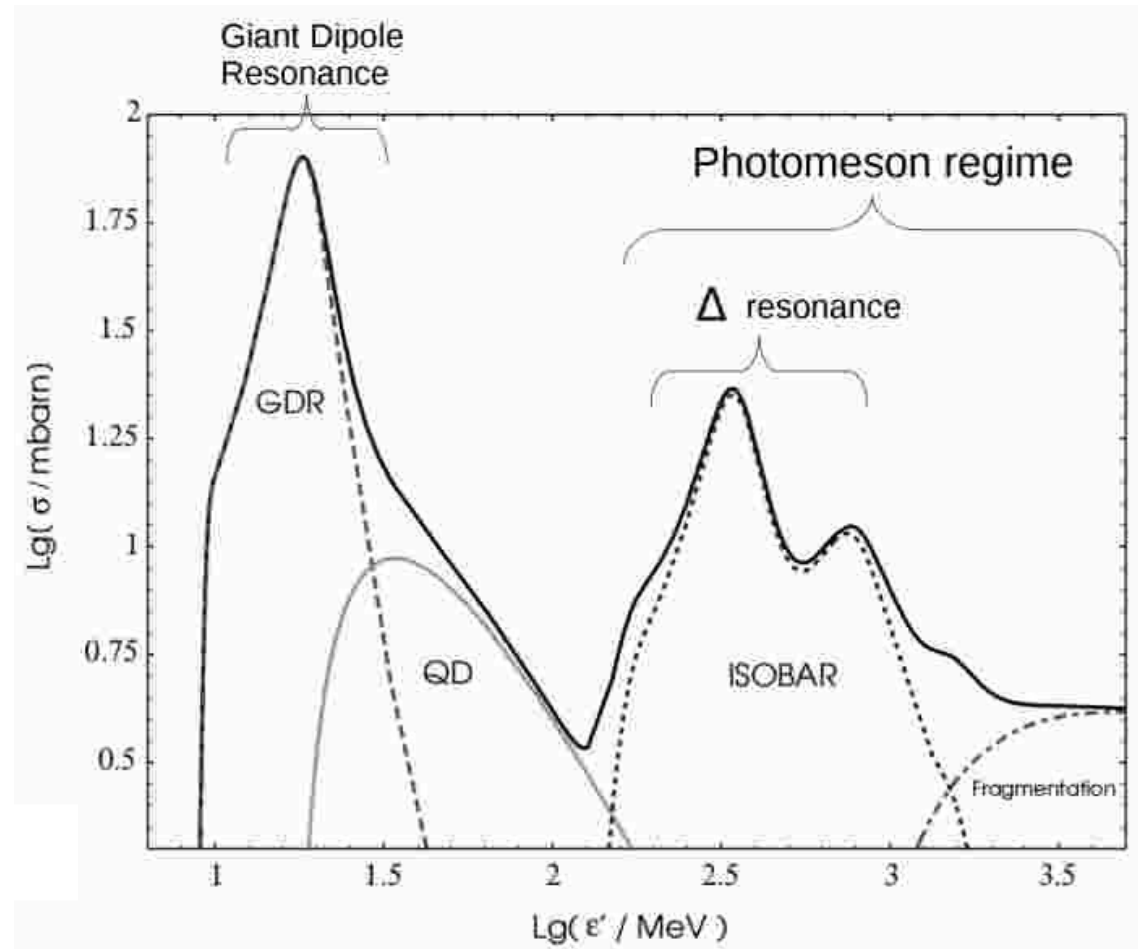
Neutrinos from UHECR interactions

- Efficient neutrino production requires high radiation densities, where on the other hand UHECRs efficiently disintegrate
 - Nuclear cascade as a measure of UHECR interactions and neutrino production, needs to be triggered for a combined treatment
 - Challenges the hypothesis of a common origin of neutrinos and cosmic rays at the highest energies
- Future measurements and improved theoretical models are essential for a better description of neutrinos from UHECR interactions, as the current data is sparse and the models do not always reproduce it well leading to large uncertainties
- Development of the nuclear cascade and neutrino production strongly depend on the model assumptions, different disintegration chains lead to different ejected compositions and neutrino yields
- Strong implications can be obtained when interactions producing neutrinos are well-understood, as in the case of the Glashow resonance, which with increasing exposure will help constrain the sources

BACKUP

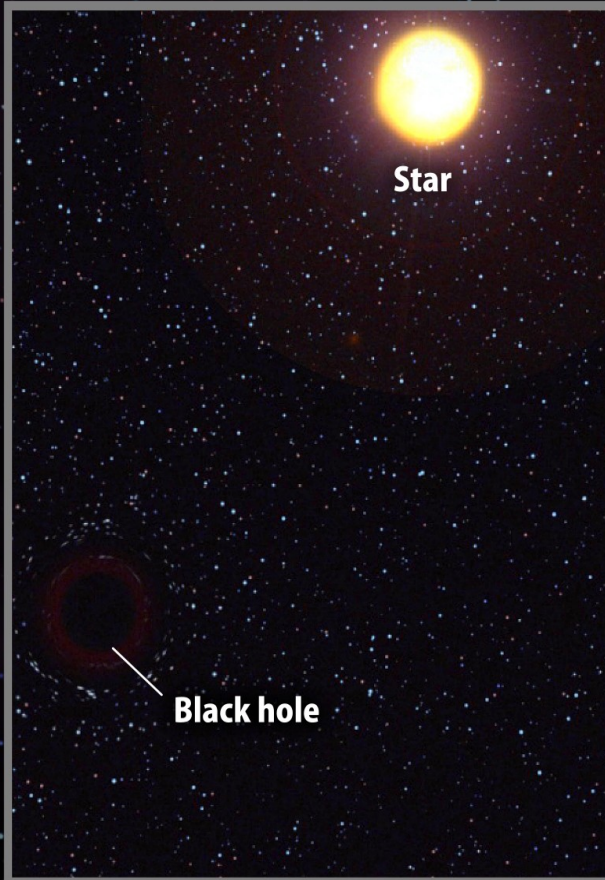
Cross section of different energy scales

And individual contributions to it

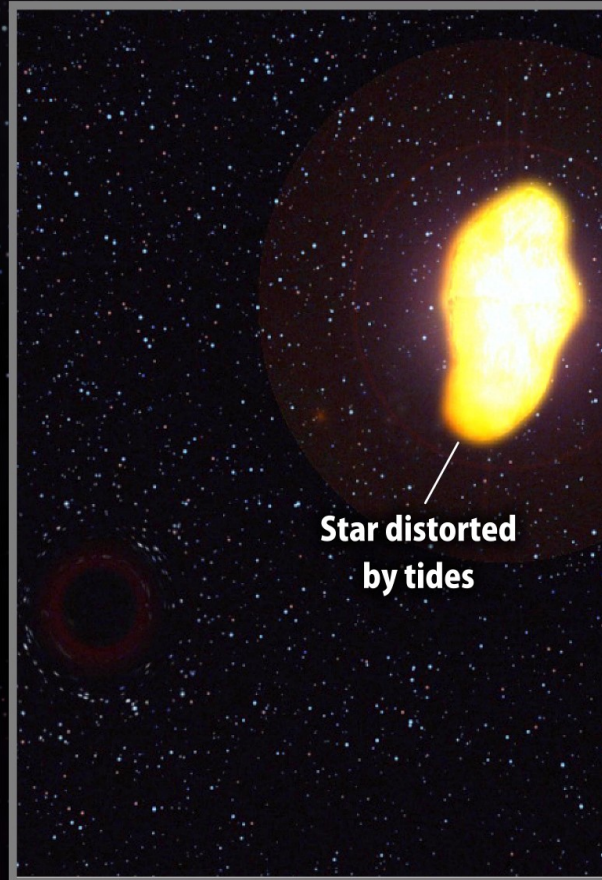


[Rachen J. P. 1996 PhD Thesis]

Swift J1644+57: Onset of a relativistic jet



1. A sun-like star on an eccentric orbit plunges toward the supermassive black hole in the heart of a distant galaxy.



2. Strong tidal forces near the black hole increasingly distort the star. If the star passes too close, it is ripped apart.



3. The part of the star facing the black hole streams toward it and forms an accretion disk. The remainder of the star just expands into space.



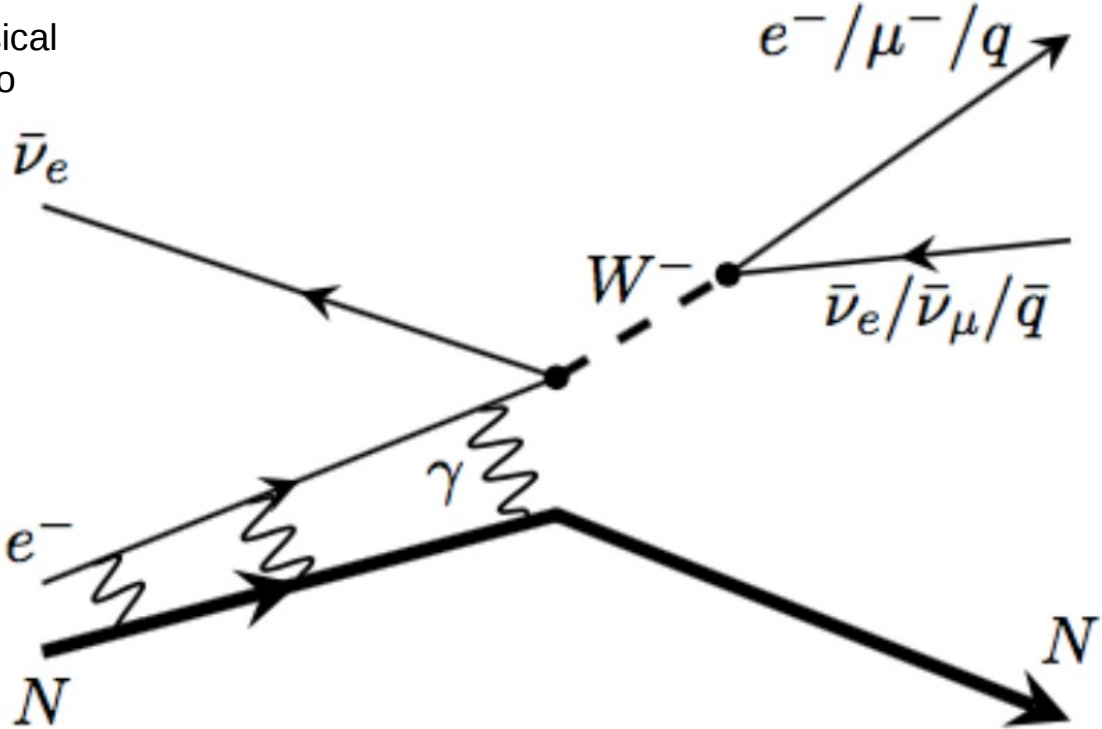
4. Near the black hole, magnetic fields power a narrow jet of particles moving near the speed of light. Viewed head-on, the jet is a brilliant X-ray and radio source.

Glashow resonance: $\bar{\nu}e$ vs. $\bar{\nu}p$ interactions

More realistic models making the difference

Incoming astrophysical
electron antineutrino
with $E \sim 6.3$ PeV

Electron from ambient
atom in the detector



Leptonic channel not distinguishable
from non-resonant events since
energy is carried away by neutrino

Description of models

Nuclear data libraries

What is TALYS?

www.talys.eu

TALYS is software for the simulation of nuclear reactions. Many state-of-the-art nuclear models are included to cover all main reaction mechanisms encountered in light particle-induced nuclear reactions. **TALYS** provides a complete description of all reaction channels and observables, and is user-friendly.

ENDF-B-VII.1 [18] is an evaluated nuclear data library based on calculations using the GNASH code system. Its photo-nuclear part contains absorption cross-sections and sometimes inclusive emission spectra of neutrons and protons, but no residual cross-sections. Comparisons with data reveal a very good agreement with the measurements.

JENDL/PD-2004 [19] is another evaluated library, based on Lorentz fits at GDR energies and quasi-deuteron emission above. Elements without σ_{abs} measurements are evaluated through branching ratios from pre-equilibrium and evaporation models, together with photo-neutron data. The description of σ_{abs} is good for all measured elements.