Simulating NNLO QCD corrections for processes with giant K factors

Sebastian Sapeta

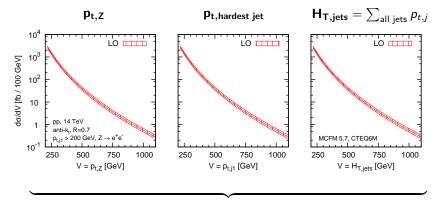
LPTHE, UPMC, CNRS, Paris

in collaboration with Gavin Salam and Mathieu Rubin 1

CERN TH Institute, Perturbative higher-order effects at work at the LHC, June 21 - July 9, 2010

The problem of giant K factors

► Z+j at the LHC

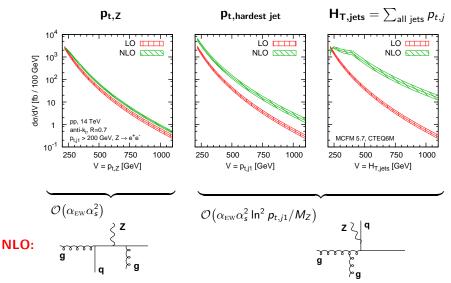


LO:



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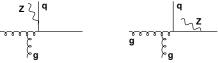


► The large K factor for the Z+jet comes from the new "dijet type" topologies that appear at NLO 72 | q



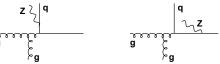


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- ▶ though formally NLO diagrams for Z+jet, these are in fact leading contributions to $p_{t,j1}$ and H_T spectra
- this raises doubts about the accuracy of these predictions
- need for subleading contributions for Z+jet, in this case NNLO

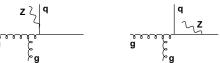
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 at NNLO $=$ $Z+3j$ tree $+$ $Z+2j$ 1-loop $+$ $Z+j$ 2-loop $Z+2j$ at NLO

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$$Z+j$$
 at NNLO = $Z+3j$ tree + $Z+2j$ 1-loop + $Z+j$ 2-loop $Z+2j$ at NLO

▶ 2-loop part

- ▶ we need it to cancel IR and collinear divergences from Z+2j at NLO result
- it will have the topology of Z+j at LO so it will not contribute much to the cross sections with giant K-factor

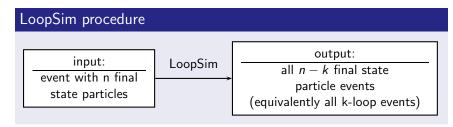
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use unitarity to simulate the divergent part of 2-loop diagrams

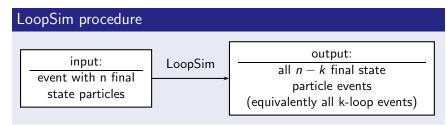
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notation:

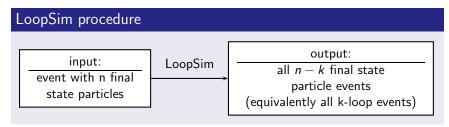
*n***LO** − simulated 1-loop

 $\bar{n}\bar{n}$ **LO** – simulated 2-loop and simulated 1-loop

 \bar{n} NLO − simulated 2-loop and exact 1-loop

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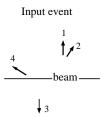
▶ notation: \bar{n} **LO** − simulated 1-loop

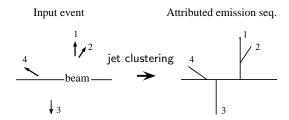
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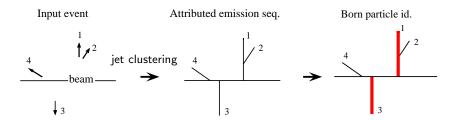
this will work very for well for the processes with large K factors e.g.

$$\sigma_{\bar{n}\mathsf{NLO}} = \sigma_{\mathsf{NNLO}} \left(1 + \mathcal{O}\left(\frac{\alpha_{\mathsf{s}}^2}{\mathsf{K}_{\mathsf{NNLO}}}\right) \right) \,, \quad \mathsf{K}_{\mathsf{NNLO}} \gtrsim \mathsf{K}_{\mathsf{NLO}} \gg 1$$

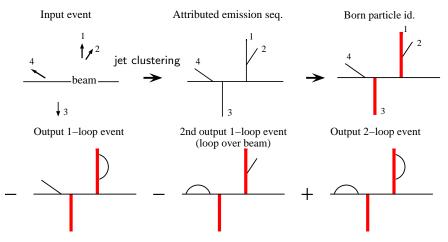




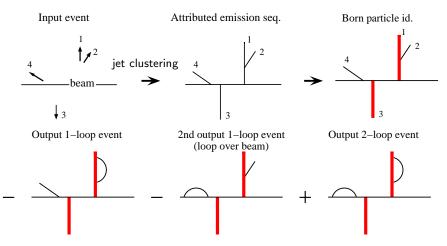
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- lacktriangle weight of an event $\sim (-1)^{\mathsf{nb.}}$ of loops and all weights sum up to zero (unitarity)
- beware: the loops above are just a shortcut notation!

 $E_{n,l}$ – input event with n final state particles and l loops

 U_l^b – operator producing event with b Born particles and l loops

 $U^b_orall$ — operator generating all necessary loop diagrams at given order

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How to introduce exact loop contributions?

$$U_{\forall}^{b}(E_{n,0})$$

generate all diagrams from the tree level event

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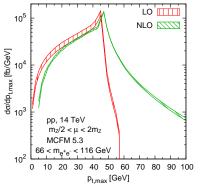
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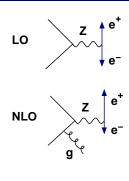
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- remove all approximate diagrams from $U_{\forall}^{b}(E_{n,0})$ that have exact counterparts provided by $U_{\forall}^{b}(E_{n-1,1})$
- ▶ inclusion of exact loops helps reducing scale uncertainties
- straightforward generalization to arbitrary number of exact loops

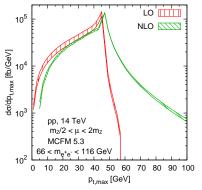


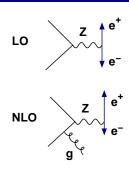
Validation



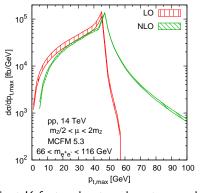


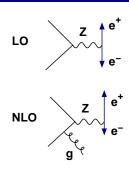
▶ giant K factor due to a boost caused by initial state radiation



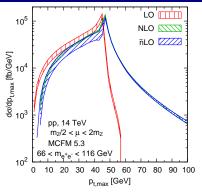


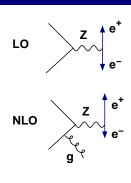
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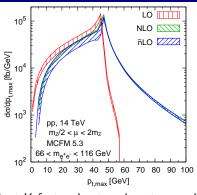
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- ▶ three regions of $p_{t,\text{max}}$: $\lesssim \frac{1}{2}M_Z$ $\left[\frac{1}{2}M_Z, 58\,\text{GeV}\right]$ > 58 GeV

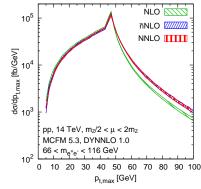




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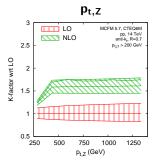
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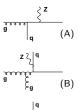




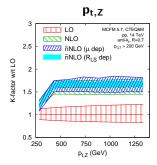
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*n***NLO** predictions for LHC





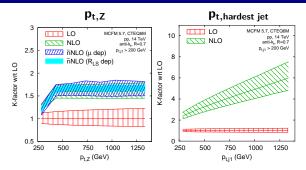




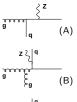
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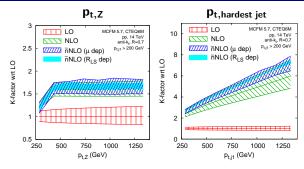






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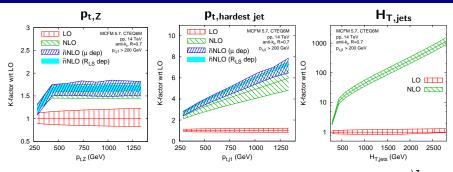
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Z+jet at $\bar{n}NLO = Z+j@NLO + LoopSim \circ (Z+2j@NLO_{only})$

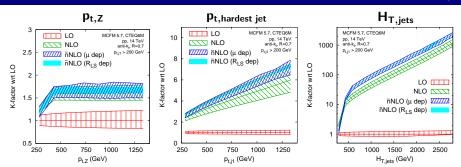


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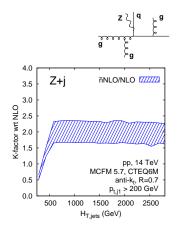
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- ▶ $H_{T,jets}$: significant correction; K factor \sim 2; given that its more like going from LO to NLO this may happen sometimes, especially for nontrivial observables like H_T ; can we understand it here?



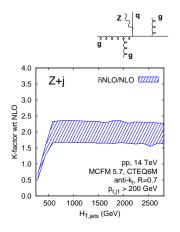


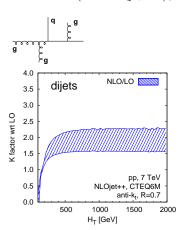






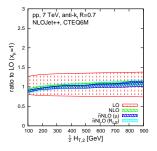






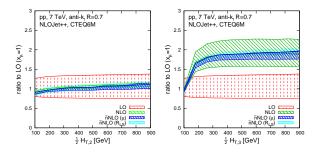
- \blacktriangleright H_T for dijets receives large contributions at NLO!
 - caused by appearance of the third jet from initial state radiation

Dijets at *n*NLO



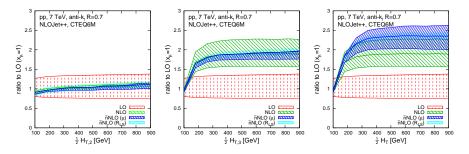
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- ▶ $H_{T,3}$ converges, significant reduction of scale uncertainty: the observable comes under control at $\bar{n}NLO$
- ▶ H_T does not converge: again caused by the initial state radiation, this time a second emission which shifts the distribution of H_T to higher values and causes no effect for the $H_{T,3}$ distribution

Summary

- several cases of observables with giant NLO K factor exist
- ▶ those large corrections arise due to the appearance of new topologies at NLO
- we developed a method, called LoopSim, which allows one to obtain approximate NNLO corrections for such processes
- the method is based on unitarity and makes use of combining NLO results for different multiplicities
- we gave arguments why the method should produce meaningful results and we validated it against NNLO Drell-Yan and also NLO Z+j and NLO dijets
- we computed approximated NNLO corrections to Z+j and dijets at the LHC finding, depending on observable, either indication of convergence of the perturbative series or further corrections
- ▶ the latter has been understood and attributed to the initial state radiation

Outlook

ightharpoonup processes with W, multibosons, heavy quarks, ...

BACKUP SLIDES

The LoopSim method: some more details

For a given input E_n event with n final state particles the weights of all diagrams generated by LoopSim sum up to zero

$$\sum_{\text{all diagrams}} w_n = \sum_{\ell=0}^{\upsilon} (-1)^\ell \binom{\upsilon}{\ell} = 0 \ , \qquad \ell - \text{number of loops, } \upsilon - \text{maximal } \ell$$

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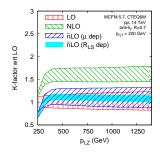
- infrared and collinear safety
- conservation of four-momentum
- choice of jet definition (algorithm, value of R)
- treatment of flavour (e.g. for processes with vector bosons)
 - Z boson can be emitted only from quarks and never emits itself
- extension to input events with exact loops; for example:

$$Z + j@\overline{n}NLO = Z + j@NLO + LoopSim \circ (Z + 2j@NLO_{only})$$



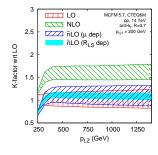
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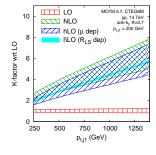
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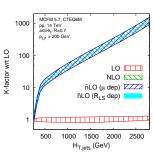


- \triangleright $p_{t,Z}$ (lack of large K-factor):
 - finite loop contributions matter
 - correctly reproduced dip towards p_t = 200 GeV

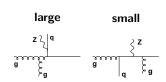
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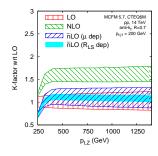


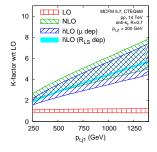


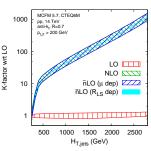
- \triangleright $p_{t,Z}$ (lack of large K-factor):
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 - correctly reproduced dip towards $p_t = 200 \text{ GeV}$
- ▶ $p_{t,j}$, $H_{T,jets}$ (giant K-factor):
 - ▶ very good agreement between n̄LO and NLO



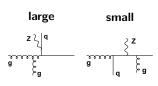
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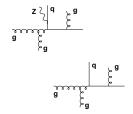


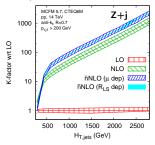


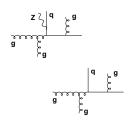


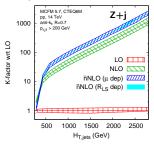
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 - ► finite loop contributions matter
 - correctly reproduced dip towards $p_t = 200 \text{ GeV}$
- \triangleright $p_{t,j}$, $H_{T,jets}$ (giant K-factor):
 - ▶ very good agreement between \bar{n} LO and NLO
- small R uncertainties driven only by subleading diagrams

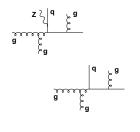


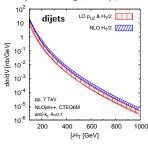




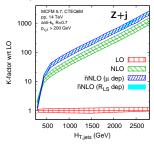


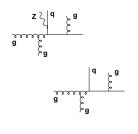


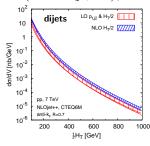




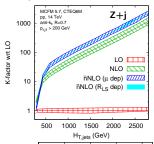
- \blacktriangleright H_T for dijets receives large contributions at NLO!
 - caused by appearance of the third jet from initial state radiation

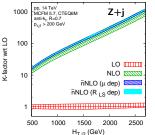


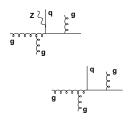


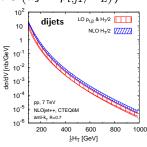


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- ▶ if the same is valid for Z + j we should see only small correction for $H_{T,j2} = \sum_{i=1}^{2} p_{t,j_i}$









- H_T for dijets receives large contributions at NLO!
 - caused by appearance of the third jet from initial state radiation
- ▶ if the same is valid for Z + j we should see only small correction for $H_{T,j2} = \sum_{i=1}^{2} p_{t,j_i}$
 - and indeed it is small!